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# Retrieval of Snow Properties for Ku- and Ka-band Dual-Frequency Radar

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## Abstract

The focus of this study is on the estimation of snow microphysical properties and the associated bulk parameters such as snow water content and water equivalent snowfall rate for Ku- and Ka-band dual-frequency radar. This is done by exploring a suitable scattering model and the proper particle size distribution (PSD) assumption that accurately represent, in the electromagnetic domain, the micro/macro-physical properties of snow. The scattering databases computed from simulated aggregates for small to moderate particle sizes are combined with a simple scattering model for large particle sizes to characterize snow scattering properties over the full range of particle sizes. With use of the single scattering results, the snow retrieval look-up tables can be formed in a way that directly links the Ku- and Ka-band radar reflectivities to snow water content and equivalent snowfall rate without use of the derived PSD parameters. A sensitivity study of the retrieval results to the PSD and scattering models is performed to better understand the dual-wavelength retrieval uncertainties. To aid in the development of the Ku- and Ka-band dual-wavelength radar technique and to further evaluate its performance, self-consistency tests are conducted using measurements of the snow PSD and fall velocity acquired from the Snow Video Imager/Particle Image Probe (SVI/PIP) during the winter of 2014 in the NASA Wallops flight facility site in Wallops Island, Virginia.

49 **1 Introduction**

50  
51 The Global Precipitation Measurement (GPM) core satellite, a joint Earth-observing mission  
52 between the National Aeronautics and Space Administration (NASA) and the Japan Aerospace  
53 Exploration Agency (JAXA), was successfully launched into orbit on 27 February of 2014 from  
54 Japan (Hou et al. 2008 and 2014). One of the goals of the Dual-frequency Precipitation Radar  
55 (DPR) aboard the GPM satellite is to provide measurements and estimates of snow precipitation  
56 rate and water content for mid- and high-latitude regions. This is usually done by estimating  
57 parameters of snow particle size distribution (PSD) that are often modeled by an analytical  
58 function, such as the exponential, gamma or lognormal distribution, with two or three unknown  
59 parameters (Gorgucci et al. 2000 and 2002; Bringi et al. 2002). The inability of the modeled PSD  
60 to represent actual snow spectra and to characterize their intrinsic variations in time and space  
61 can lead to errors in the estimates of precipitation rate obtained from the DPR. Additionally,  
62 uncertainties associated with scattering computations of snow aggregates also affect the accuracy  
63 of the dual-wavelength radar retrieval of snow arising from the complex shape and structure of  
64 snow aggregates and the corresponding variability in the scattering parameters. Therefore,  
65 understanding the uncertainties in snow precipitation estimation that depend on PSD  
66 parameterizations and scattering models of individual particles is important in evaluating the  
67 overall performance of DPR retrieval algorithms and in gaining insight into ways to improve the  
68 algorithms.

69  
70 Several studies have been carried out using dual-frequency radar for the retrieval of precipitating  
71 ice/snow parameters from the ground (Matrosov 1998; Szyrmer and Zawadzki 2014) and from  
72 airborne radar measurements (Liao et al. 2005, 2008; Heymsfield et al. 2005; Wang et al. 2005;

73 Matrosov et al. 2005). Although various combinations of frequencies were used in these studies,  
74 a common feature is that at least one of the radar frequencies operates in the non-Rayleigh  
75 regime to ensure a measurable difference in the reflectivities. It is this differential reflectivity  
76 that can be related to a characteristic size parameter of the snow particle distribution. Because  
77 of uncertainties in the snow microphysics arising from the natural variability of the particle  
78 density, shape, and orientation and also because of uncertainties associated with the particle  
79 backscattering cross section and terminal fall velocity as well as the natural variability in PSDs,  
80 it is important to assess the errors in the model and their impacts on the retrievals.

81  
82 The emphasis of this study is on the estimation of snow microphysical properties and the  
83 associated bulk parameters such as snow water content and water equivalent snowfall rate. As  
84 indicated earlier, one of the challenges in the radar retrieval of snow is to characterize the  
85 variability of the snow PSD and to efficiently compute scattering properties of the snowflakes  
86 over the full range of sizes. The aim of our study is to explore a suitable scattering model and an  
87 appropriate PSD that accurately represents, in the electromagnetic domain, the micro/macro-  
88 physical properties of snow.

89  
90 Although several scattering databases are available, which provide the scattering properties of  
91 snow aggregates (Liu 2004 and 2008; Nowell et al. 2013; Kuo et al. 2016), they are often limited  
92 to small-to-moderate particle sizes for a limited set of frequencies. These limitations arise not  
93 only because of the magnitude of the computational burden but also because of the convergence  
94 properties of the numerical solution. To develop an operational-type radar algorithm for the DPR  
95 snow retrieval, it is desirable to have a scattering model that provides efficient computation at an

96 arbitrary frequency over a large range of particle sizes. Comparisons of the scattering results  
97 between simple and more complicated snow models indicate that the scattering properties of  
98 aggregates at the DPR frequencies are fairly well reproduced by randomly-oriented ellipsoidal  
99 particles if the effective mass density of snow is constant with size (Liao et al. 2013). By taking  
100 advantage of both developed scattering databases and simple scattering models, we attempt to  
101 employ the scattering results of the aggregates from the scattering database for small to moderate  
102 particle sizes and use the results from the simple scattering models for large particles to cover the  
103 full range of particle sizes for characterizing snow scattering properties at Ku- and Ka-band.

104

105 One of the DPR algorithms for snow retrieval employs a fixed-snow-density spherical model for  
106 computations of the Ku- and Ka-band radar backscattering and extinction cross sections using  
107 the assumption of an exponential PSD (Seto et al. 2013). To improve snow retrieval accuracy,  
108 we will investigate retrieval uncertainties associated with the PSD and the particle scattering  
109 models, and introduce new forms of the retrieval look-up tables that directly link DPR Ku- and  
110 Ka-band radar reflectivities to snow water content and snowfall rate without the use of derived  
111 PSD parameters. Newly formed look-up tables provide an effective means not only for snow  
112 retrieval but for analysis of the retrieval uncertainties associated with the PSD model and the  
113 particle scattering models. In order to check the consistency of the snow retrievals,  
114 measurements of the snow PSD and fall velocity acquired from the Snow Video Imager/Particle  
115 Image Probe (SVI/PIP) are used (Newman et al. 2009). It is believed that a collection of long-  
116 term PSD data, fall velocities and information on particle mass spectra will provide a strong  
117 basis for evaluating the performance of the DPR Ku- and Ka-band techniques.

118

119 This article is organized as follows. Methods for snow retrieval are described in Section 2, and  
120 measurements of the PSD are discussed in Section 3, followed by remarks and a summary in  
121 Section 4.

122  
123 **2 Technical Approach and Methodology**  
124

125 Understanding the retrieval errors associated with the snow particle size/mass distributions  
126 models and particle scattering models employed by the DPR algorithms is important for the  
127 evaluation of algorithm performance. The study also provides insight into which of the models  
128 yield the most accurate DPR estimates of snow. Proper selection of the PSD and scattering  
129 models can improve the overall performance of the DPR profiling algorithm. What follows are  
130 discussions of snow scattering models and the parameterization of the particle size spectra and  
131 how these models affect the radar retrievals.

132  
133 **2.1 Single scattering model of snow aggregates**  
134

135 Several simulated aggregate models have been developed using the pristine ice crystal habits  
136 found in nature as the basic elements from which the aggregates are constructed (Draine and  
137 Flatau 1994; Liu 2004 and 2008; Weinman and Kim 2007; Petty and Huang 2010; Botta et al.  
138 2010 and 2011; Nowell et al. 2013; Kuo et al. 2016). For these particles, a numerical scheme is  
139 required to compute the scattered fields. Although these numerical computations are useful, they  
140 are time consuming, and are often limited to small-to-moderate particle sizes for a limited set of  
141 frequencies. A few scattering databases derived from simulated aggregates are available (Nowell  
142 et al. 2013; Kuo et al. 2016), but the maximum equivalent ice diameter is limited to around 2.5-3  
143 mm which is not large enough to cover entire particle size range.

144  
145 To account for the scattering contribution from the entire particle size range, the current DPR  
146 algorithms adopt a simple scattering model, namely, the ice-air mixed spheres with a fixed snow  
147 density of  $0.1 \text{ g/cm}^3$  for all particle sizes (Seto et al. 2013). To check the validity and accuracy of  
148 the simple geometric scattering model, a study was carried out by Liao et al. (2013) in which  
149 scattering results from aggregates comprised of 6-branch bullet rosette crystals were compared  
150 with those obtained from spherical or spheroidal ice-air mixed phase particles. Shown in Fig.1 is  
151 an example of these results, in which backscattering (left) and extinction (right) coefficients at  
152 35.6 GHz are plotted versus the equivalent ice diameter for 3 simulated snow aggregates and  
153 results from the sphere, oblate and prolate spheroids. A constant effective snow density of  $0.2$   
154  $\text{g/cm}^3$  is assumed for all spherical and spheroidal particle models. The spheroids are assumed to  
155 be randomly oriented, i.e., their symmetry axes are uniformly distributed in three-dimensional  
156 space. The aspect ratios ( $\gamma$ ) of the spheroids, defined as the ratio of polar to equatorial lengths,  
157 are taken to be either 0.5 or 2 to represent oblate and prolate spheroids respectively. The results  
158 of the study show that the scattering properties of the aggregates are fairly well reproduced by an  
159 equivalent-mass spheroidal particle when a constant snow density is assumed.

160  
161 Agreement between the spheroidal/fixed density model and the aggregates suggests the validity  
162 of the simple model and its utility for computing look-up tables for the DPR. An effective snow  
163 density of  $0.2 \text{ g/cm}^3$  is best able to reproduce the scattering parameters of the bullet-rosette  
164 aggregates at the DPR frequencies. It is important to note that the effective mass density, as  
165 defined in this paper, is that mass density of a spheroidal particle whose associated scattering  
166 parameters provide the best match to those of the simulated aggregates with the same mass. This  
167 definition is motivated by the desire to match the scattering properties of the spheroidal and

168 simulated aggregates and is distinct from the usual definition of snow density given by the ratio  
169 of the particle mass to the volume where the volume is taken to be that of a circumscribing  
170 sphere or spheroid. To distinguish our definition from others, the density we use for scattering  
171 computations is referred to as the effective snow density. It is also important to note that the  
172 snow mass is preserved in both definitions; that is, the product of the mass density or effective  
173 mass density and the particle volume yields the same mass.

174  
175 Recently, Kuo et al. (2016) have developed a comprehensive scattering database, which is  
176 computed using the discrete dipole approximation (DDA) from a collection of realistic  
177 aggregates simulated from a 3-D growth model with mass vs. size and fractal properties that are  
178 consistent with field observations (Gravner and Griffeath 2009). Because of its limited range of  
179 particle size mentioned earlier, the scattering tables of snow aggregates in our study will be taken  
180 as a hybrid form that combines the scattering results from the Kuo et al. scattering database for  
181 small to moderate particle sizes and the results from the simple scattering models for large  
182 particles. Illustrated in Figs.2 and 3 are these scattering results at Ku- and Ka-bands from the  
183 Kuo et al. database along with the results from an oblate spheroidal model for particle diameters  
184 up to 6 mm. The results from the simple models (red curves) are those obtained from the  
185 randomly-oriented oblate spheroids with an aspect ratio of 0.7 and a constant effective mass  
186 density of  $0.2 \text{ g/cm}^3$ . The results of the scattering tables, which are denoted by the term “fitted”  
187 and shown by the black curves, represent the mean values of the data from the scattering  
188 database, where the mean is taken from all types of aggregates having the same mass, and the  
189 results from the  $0.2 \text{ g/cm}^3$  oblate spheroid model in the size range where the scattering database  
190 is unavailable. The maximum liquid equivalent diameters of the current Kuo et al. database at



191 Ku- and Ka-bands are 3 and 2 mm, respectively. In other words, the scattering results from the  
 192 simple models are employed in the tables for the size ranges from 3-6 mm at Ku-band and 2-6  
 193 mm at Ka-band. The fairly good agreement of the scattering results in Figs. 2 and 3 between the  
 194 simple model and the scattering database (with mean differences less than 10% for the  
 195 backscattering and 7% for extinction cross sections) over the size range where the database is  
 196 valid suggests the validity of the simple models for the smaller particle sizes. As the scattering  
 197 database is updated to cover larger particle sizes, the validity of the simple scattering model will  
 198 be reassessed.

199

## 200 2.2 Particle size/mass distribution model

201

202 The three-parameter gamma distribution is one of the most common ways to mathematically  
 203 describe hydrometeor size/mass distributions (Gorgucci et al. 2000 and 2002; Bringi et al. 2002).

204 The form of the gamma distribution is expressed as

$$205 \quad N(D) = N_w f(\mu) \left( \frac{D}{D_m} \right)^\mu \exp(-\Lambda D), \quad (1)$$

206 where  $D_m$  is the mass-weighted diameter of the particle,  $N_w$  is a scale factor, and  $\mu$  is the shape  
 207 factor where

$$208 \quad f(\mu) = \frac{6(4 + \mu)^{\mu+4}}{4^4 \Gamma(\mu + 4)} \quad \text{and} \quad \Lambda = (4 + \mu)/D_m, \quad (2)$$

209 To describe snow particle size and mass spectra, the PSD in (1) and (2) is given as a function of  
 210 liquid equivalent or melted diameter  $D$ , which is also called particle melted-size distribution.  $D_m$   
 211 is the melted median mass diameter, defined by

212

213 
$$D_m = \frac{\int_0^\infty D^4 N(D) dD}{\int_0^\infty D^3 N(D) dD}. \quad (3)$$

214

215 In the inner swath, the DPR provides Ku- and Ka-band reflectivity factor measurements at each  
 216 range gate so only two parameters of the PSD can be determined. Typically, the shape factor ( $\mu$ )  
 217 is taken to be constant. Although  $\mu$  is often set to zero (exponential distribution) (Gunn and  
 218 Marshall 1958; Seto et al. 2013), the impact of this choice on the retrieval needs to be  
 219 investigated.

220

221 The differential frequency ratio (DFR), which is defined as the difference between the radar  
 222 reflectivity factors at Ku- and Ka-bands in decibels, is perhaps the most important quantity for  
 223 the dual-wavelength radar techniques in estimating hydrometeor micro/macro-physical  
 224 properties. As the DFR is independent of  $N_w$ ,  $D_m$  can be derived from the DFR relations once  $\mu$   
 225 has been fixed. However, the DFR- $D_m$  relation depends not only on  $\mu$  but on the particle shape,  
 226 orientation distribution and mass density. Fig.4 provides the results of DFR as a function of the  
 227 liquid equivalent median mass diameter using a randomly-oriented, fixed density spheroidal  
 228 particle model. The left plot shows the variations in the DFR- $D_m$  relation resulting from  
 229 different effective snow densities. Computations of the radar scattering parameters at different  
 230 effective snow densities are made in the same way as in the case of  $0.2 \text{ g/cm}^3$ . The particle sizes  
 231 (semi axes of spheroid) are solely determined by the density specified for a given particle mass.  
 232 The center plot shows the effects of particle shape where a  $\gamma$  value of 1 corresponds to a sphere  
 233 while  $\gamma$  values less than 1 correspond to an oblate spheroid. The plot on the right shows the  
 234 effect of changing  $\mu$ . Analysis of these results indicates that particle shape has a small effect on

235 the DFR- $D_m$  relation and changes the results by less than 4% for changes in  $\gamma$  from 0.5 to 1 while  
236 the shape factor leads to the change in the results of no more than 20%. On the other hand, the  
237 DFR- $D_m$  relation has a strong dependence on the effective snow density. In other words, the  
238 determination of  $D_m$  from the DFR is relatively insensitive to  $\mu$  and to particle shape, if the  
239 orientation is random, whereas the relationship is quite sensitive to the effective snow density  
240 used for computations of the scattering parameters. As can be seen in Figs.1-3, the extinction  
241 coefficients at both Ku- and Ka-bands, though in good agreement between the simple model and  
242 the scattering database, are small and can usually be neglected. It should be noted that the above  
243 conclusions are based on the assumption of random orientation of the aggregates. When this  
244 assumption is violated, then the orientation distribution as well as particle shape become  
245 important.

246  
247 Although an effective snow density of  $0.2 \text{ g/cm}^3$  is found to be suitable for the computations of  
248 the Ku- and Ka-band radar scattering parameters for equivalent ice diameters up to 2-3 mm,  
249 further testing will be necessary to assess this assumption when scattering results from larger  
250 aggregates become available. It is worth noting that the results of Liao et al. (2013) show  
251 reasonably good agreement between radar scattering parameters at higher frequencies (from 89-  
252 183.31 GHz) as derived from the simple models and the simulated aggregates for particle  
253 diameters up to 2.5 mm despite the fact that the simple models using spheres or nearly spherical  
254 particles produce backscattering results with more pronounced oscillations (resonance effects)  
255 than the aggregate results. This is encouraging in the sense that it shows that an effective density  
256 of  $0.2 \text{ g/cm}^3$  yields good agreement with the simulated aggregate results for electrically larger  
257 particles. On the other hand, it might be the case that the effective snow density may need to be

258 changed for larger particle types so as to ensure good agreement. In either case, the objective is  
259 to provide scattering tables at all relevant frequencies and particle sizes that incorporate the most  
260 recent scattering results.

261

### 262 2.3 Dual-wavelength retrieval algorithm

263

264 As discussed above,  $D_m$  can be derived from the DFR- $D_m$  relations for a given  $\mu$ . In principle,  
265 once  $D_m$  is determined,  $N_w$  is derived using the radar reflectivity at either Ku- or Ka-band.

266 Subsequently, snow water content (SWC) and equivalent snowfall rate ( $R_S$ ) can be computed  
267 from the derived PSD parameters. The fall velocity of snow is needed in order to estimate  $R_S$ .

268 For the development of an effective dual-wavelength radar retrieval technique, it is desirable to  
269 employ look-up tables (LUT) that are formed in such a way that the radar measurements are  
270 directly linked to the microphysical properties of snow ( $D_m$  and  $N_w$ ) and its associated bulk  
271 parameters (SWC and  $R_S$ ). With use of the LUTs different particle models and their scattering  
272 properties can be evaluated separately in the context of the same algorithm.

273

274 Illustrated in Fig.5 are such tables in which SWC (top-left),  $R_S$  (top-right),  $D_m$  (bottom-left) and  
275  $N_w$  (bottom-right) along the ordinate are given as a function of the DFR. A flowchart is provided  
276 in Fig.6 showing the procedures to compute the radar reflectivity factors and snow size and  
277 bulk hydrometeor parameters from an assumed mass spectrum model. In Fig.5, SWC,  $R_S$ ,  $N_w$   
278 have been normalized by the Ku-band radar reflectivity factor so that they can be expressed  
279 solely as a function of DFR for given PSD and scattering models. The way to normalize liquid  
280 water content by reflectivities has previously been adopted in the study of ice clouds (Hogan et

281 al. 2000; Botta et al. 2013). The results in Fig.5 are computed under the assumption that the  
282 snow particles are fixed-density, randomly-oriented oblate spheroids with an aspect ratio of 0.7  
283 that follow an exponential particle size distribution. As an example, and also for reference, the  
284 tables are plotted in Fig.5 for effective snow densities varying from 0.05 to 0.5 g/cm<sup>3</sup>. The  
285 terminal velocities of snowflakes used for the computations of  $R_S$  are based on the results of  
286 Magono and Nakamura (1965). It is important to note that the results from the LUTs shown in  
287 Fig.5 can be used to determine SWC and  $R_S$  as they directly link the DPR radar reflectivities to  
288 SWC and  $R_S$  without use of the derived PSD parameters.

289  
290 The procedure for the estimation of snow parameters is described as follows: given a pair of  
291 reflectivity factors ( $Z_{Ku}$ ,  $Z_{Ka}$ ), the DFR in dB is defined as  $10\text{Log}_{10}(Z_{Ku}/Z_{Ka})$ , from which we find  
292 the values of  $\text{SWC}/Z_{Ku}$  (left) and  $R_S/Z_{Ku}$  (right) for an assumed effective snow density. By  
293 multiplying by  $Z_{Ku}$ , the results of SWC and  $R_S$  are then obtained. Obviously the values of SWC  
294 and  $R_S$  depend on the effective snow density. The estimates of  $D_m$  and  $N_w$  can be achieved in a  
295 similar way. It is worth mentioning again that snow attenuations, though correctable, are  
296 typically negligibly small for most Ku- and Ka-band spaceborne radar measurements.

297  
298 As LUTs change with different scattering models and PSD parameterizations, a proper selection  
299 of the tables is critical to the accuracy of the retrieval. It is instructive to conduct a sensitivity  
300 study with respect to the model assumptions and to gain an understanding of the uncertainties  
301 associated with each of the models. Figure 7 provides such a sensitivity study in which the LUTs  
302 are checked against 3  $\mu$  values. Similar to the findings in Fig.4, a change in  $\mu$  leads to changes in  
303 the estimates of SWC and  $R_S$  of less than 20% so that the assumption of  $\mu$  equal to zero, as found

304 in many observations, yields a reasonable approximation for the estimates of snow. Although it  
305 is worth testing the scattering databases of the aggregates formed from various crystal habits, it is  
306 not the focus of the study to evaluate and validate these scattering databases. Because the mass  
307 of the aggregates is the dominant factor in the scattering parameters at Ku- and Ka-bands,  
308 significant differences among the various scattering databases are not expected. This is  
309 evidenced by the fact of that there is good agreement between the scattering databases derived by  
310 Nowell et al. (2013) and Kuo et al. (2016) and that the small variations of the scattering and  
311 extinction coefficients as computed from various aggregate models (Kuo et al. 2016) can be seen  
312 from the variations in the data (blue) shown in Figs.2-3.

313

### 314 **3 Assessment of snow retrieval: PSD model assumptions**

315

316 Because of the complexity of snowfall processes and the difficulties encountered in accurately  
317 measuring the microphysical properties, validating snow estimates is a challenging task. With the  
318 advent of more advanced digital cameras and image processing technologies, measurements of  
319 falling snow have been improved to the point where the snow particle size spectra and fall  
320 velocities can be obtained fairly accurately (Bohm 1989; Huang et al. 2010 and 2015; Garrett et  
321 al. 2012). An independent and direct measurement of the mass of individual snow particles is,  
322 however, still a difficult task, and therefore direct measurements of the snow mass spectrum are  
323 rarely available. Several investigations into deriving snow mass spectra are being pursued, which  
324 are in fact part of the effort in the GPM ground validation project. These methods are based on  
325 the principle that particle masses can be related to their fall velocities after accounting for air  
326 drag and other aerodynamics effects (Bohm 1989; Heymsfield et al. 2010). Understanding the  
327 microphysical properties of snow should further improve our ability to generate better scattering

328 representations and more accurate look-up tables for retrieving snow bulk properties from the  
329 DPR. It would be ideal to evaluate snow retrievals with co-located dual-wavelength radar  
330 measurements and in situ snow microphysical measurements. These data, though desirable, are  
331 not available. Our attention is therefore focused on the assessment of the PSD assumptions used  
332 in developing the retrieval algorithms using measured PSD.

333 To check the consistency of snow retrievals using the LUTs, measurements of snow spectra are  
334 used. The data were obtained from 8 snow events during winter of 2014 taken at the NASA  
335 Wallops Flight Facility in Wallops Island, Virginia using the Snow Video Imager/Particle Image  
336 Probe (SVI/PIP). Table 1 provides details of these events that include starting and ending times  
337 of snowfall, mean temperature as well as total accumulations of each event. In Wallops Island  
338 annual mean snowfall is about 200.66 mm, and in 2014 it was recorded to be 223.66 mm,  
339 slightly more than average. Although the PIP measures the dimensions or sizes of the  
340 snowflakes and their fall velocities, it does not provide measurements of particle mass. In order  
341 to compute the radar reflectivities and snow bulk parameters as in (4) and (5), the mass spectra or  
342 melted size spectra are needed. Conversion of the PSD measurements to the mass spectra,  
343 however, relies on the empirical mass-size relations. There are many such relationships available  
344 in the literature that can be used to derive  $m(D)$  (Nakaya 1954; Magono and Nakamura 1965;  
345 Zikmunda and Vali 1972; Locatelli and Hobbs 1974; Mitchell et al. 1990; Brandes et al. 2007;  
346 Heymsfield et al. 2010). These results show some variability depending on snow type, amount of  
347 riming and other conditions under which the measurements were made. In this study, two well-  
348 known mass-size relations, the results from Heymsfield et al. (2010) and Brandes et al. (2007),  
349 are used to test how the estimates of snow change with use of different mass-size relations when  
350 the same LUTs are used.

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An example of the PSD measurements is shown in Fig.8, in which  $N(D)$  of the PSD (top), shown along the vertical with the amplitude of the spectrum given by the color scale, is given as a function of time. The equivalent snowfall rate (middle) and median mass diameter ( $D_m$ ) (bottom) are also shown for the same time period. For the computations of snowfall rate and  $D_m$ , empirical snow mass-size relations are used in conjunction with the measured snow particle size spectra and fall velocities. The following equations are used for obtaining  $R_S$  and  $D_m$ :

$$D_m = \frac{\int_0^\infty D m(D)N(D)dD}{\int_0^\infty m(D)N(D)dD}, \tag{4}$$

$$R_S = \frac{36 \times 10^{-4}}{\rho_w} \int_{D_{min}}^{D_{max}} N(D) m(D) V(D) dD, \tag{5}$$

where  $D$  and  $D_m$  in (4) and (5) as well as in Fig.8 are, respectively, the actual particle diameter and median mass diameter rather than the melted sizes employed in the rest of the paper.  $m(D)$  is the particle mass-size relation, and  $\rho_w$  is the water mass density taken to be  $1 \text{ g/cm}^3$ . Note also that the data shown in Fig. 8 represent a measurement period of 1000 minutes of snow data (~three snow events) with a one-minute integration time. The mass-size relation of Heymsfield et al. (2010) is used for computations of  $R_S$  and  $D_m$ .

Illustrated in Fig.9 are the scatter plots (red dots) of SWC (top row) and  $R_S$  (bottom row) computed from the measured PSD with use of mass-size relations when the hybrid scattering tables are assumed. For reference, the LUTs derived from the constant effective density scattering models as shown in Fig.6 are superimposed in Fig.9. The SWC and  $R_S$  as well as the quantities associated with the Ku- and Ka-band radar reflectivities shown in Fig.9 are obtained from a total of about 8000 1-minute PSD measurements collected from the PIP during the winter



373 of 2014 at Wallops using the mass-size relations of Brandes et al. (left column) and Heymsfield  
374 et al. (right column). Mass-size relations are used to convert the measured particle size  
375 distribution (PSD) to the melted size distribution from which the scatterings and snow bulk  
376 parameters can be computed as in (2)-(5). The mass-size relations used in Fig.9 are those from  
377 Eq.(8) of Brandes et al. (2007) and Eq.(10) of Heymsfield et al. (2010). The procedures used in  
378 obtaining the radar parameters and SWC and  $R_s$  from the measured PSD are shown in the flow  
379 diagram of Fig.10. Analysis of the SWC results indicates that the snow water content derived  
380 from the measured PSD agree reasonably well with those from the tables when the effective  
381 snow density is taken to be  $0.2 \text{ g/cm}^3$ . Because the scattering table, which is a hybrid formed  
382 from the mean aggregate solution at small particle sizes with the spheroidal particle model at  
383 larger sizes, used for computing the reflectivities of the measured PSD and the scattering results  
384 with a density of  $0.2 \text{ g/cm}^3$  are nearly the same as shown in Figs.2-3, the differences in snow  
385 water content between the measured PSD and the table results from a density of  $0.2 \text{ g/cm}^3$  are  
386 mostly caused by the differences between measured and modeled melted particle size  
387 distributions. As noted earlier, the exponential melted-size distribution is assumed in the look-up  
388 tables while the measured melted-size distribution is derived from the measured particle diameter  
389 spectrum and the mass-size relation that generally will be different from an exponential  
390 distribution. Computation of the snowfall rate, on the other hand, depends not only on the  
391 particle mass (or melted-size) spectrum but also on the particle fall velocities. Most of the  
392 estimated snowfall rates, as derived from the measured PSD and the mass-size relations, and  
393 shown in the lower panels of Fig.9, lie between the table results with effective densities of 0.1  
394 and  $0.2 \text{ g/cm}^3$ . The mean differences of the SWC between the  $0.2 \text{ g/cm}^3$  snow density LUTs and  
395 the PSD-derived results are about 20% for Ku-band radar reflectivities greater than 15 dBZ (the

396 approximate minimum detectable signal of the DPR Ku-band channel). Larger differences  
397 between the  $R_S$  estimates are found, and can be attributed to differences in the snow fall velocity  
398 spectra between the measured and modeled mass distribution. The terminal velocities of  
399 snowflakes used for the computations of the LUTs are based on the results of Magono and  
400 Nakamura (1965) while the measured fall velocities are used for the computations of PSD  
401 snowfall rate.

402  
403 The overall agreement of the snow water content between the results from the measured PSD  
404 and the results from the LUTs suggests that the exponential particle distribution model assumed  
405 in the tables is reasonable. Different mass-size relations lead to different mass spectra for a given  
406 measured PSD. That the retrieval results from the Brandes et al. and Heymsfield et al. mass-size  
407 relations follow the trends of SWC and  $R_S$  similar to those derived from the LUTs further  
408 suggest that the Ku- and Ka-band dual-wavelength techniques adopted are relatively insensitive  
409 to the choice of either the Brandes or the Heymsfield mass-size relation. These results  
410 approximately yield the table values obtained from the  $0.2 \text{ g/cm}^3$  snow density. It is also worth  
411 mentioning that the results from the measured PSD are relatively insensitive to PSD integration time even  
412 though the scatter in the data is slightly reduced if a longer integration is used.

413

#### 414 **4 Summary**

415

416 The ultimate goal of this study is to better understand the estimation process in retrieving snow  
417 microphysical properties ( $N_w$  and  $D_m$ ) and the associated bulk parameters (SWC and  $R_S$ ) for  
418 improvement of the Ku- and Ka-band dual frequency radar retrieval. This is done by first finding  
419 suitable single scattering tables and PSD models and then using this information to construct

420 snow retrieval look-up tables. Presently available scattering databases, though accurate and  
421 useful, are limited to small and moderate particle sizes. To extend the results to larger sizes, a  
422 simple scattering model that agrees well with the scattering databases at small particle sizes is  
423 used. It is found that a snow particle model consisting of randomly oriented oblate spheroids  
424 with an effective mass density of  $0.2 \text{ g/cm}^3$  yields good agreement with the results from the  
425 scattering databases at Ku- and Ka-band. Thus the single-particle scattering database is a hybrid  
426 that uses the scattering database for small and moderate particles and a simple randomly oriented  
427 oblate with a constant effective mass density of  $0.2 \text{ g/cm}^3$  for large particles.

428

429 Using single scattering tables and an assumed PSD model, the Ku- and Ka-band radar  
430 reflectivity factors and snow bulk parameters are computed. Thus, the relationships between the  
431 results of DFR and SWC and  $R_S$  are established to form the dual-wavelength radar retrieval look-  
432 up tables. Retrievals of snow water content and snowfall rate, as the primary focus of this study,  
433 are therefore achieved by using newly introduced look-up tables that directly link Ku- and Ka-  
434 band radar reflectivities to hydrometeor parameters without the use of derived PSD parameters.  
435 The look-up tables are formed so that SWC and  $R_S$ , both of which are normalized by the Ku-  
436 band radar reflectivity factor, are expressed as a function of the differential frequency ratio of  
437 Ku- and Ka-bands. The look-up tables offer not only computational advantages but provide  
438 direct insight into how the model assumptions impact the retrieval results. The nature of one-to-  
439 one relations between the normalized hydrometeor parameters and DFR provides a means to  
440 obtain unique solutions of the snow parameters for a given PSD and single scattering model. To  
441 understand the uncertainties in the snow estimates associated with the PSD parameterizations  
442 and scattering models, a sensitivity study was done, finding that the choice of shape factor of the

443 gamma PSD has only a slight impact on the retrievals. As such, a value of  $\mu$  of zero, as  
444 supported by some observations, should yield reasonable estimates of snow parameters from the  
445 perspective of dual-wavelength radar retrieval.

446

447 Self-consistency of the snow retrievals has been checked using measurements of snow PSD and  
448 fall velocity acquired from the PIP during the winter of 2014 in Wallops. Among several  
449 assumptions that have been examined are conversions to particle mass spectra using different  
450 mass-size relations, scattering particle models and snow PSD. Analysis of nearly 8000 1-minute  
451 PSD measurements suggests that exponential PSD model ( $\mu=0$ ) is sufficiently accurate for the  
452 dual-wavelength radar retrieval of snow bulk parameters. It also indicates that the use of either  
453 the Heymsfield or the Brandes mass-size relation yields approximately the same snow estimates.  
454 However, these findings should be viewed as preliminary because of the limited data  
455 measurements at a single location. Collections of long-term PSD data, fall velocities and  
456 information on particle mass spectra at multiple sites will provide further evaluation of the  
457 performance of the Ku- and Ka-band radar techniques. Further tests of the scattering tables will  
458 be done by comparing the accuracy of the simple particle model against scattering results from  
459 larger simulated aggregates, as these results become available.

460

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464 providing and processing SVI/PIP data, and Dr. Kwo-Sen Kuo of University of Maryland for  
465 providing the scattering database.



467

## Appendix

468

### List of Symbols and Acronyms

469	GPM:	Global Precipitation Measurement
470	DPR:	Dual-frequency Precipitation Radar
471	NASA:	National Aeronautics and Space Administration
472	JAXA:	Japan Aerospace Exploration Agency
473	Ku-band:	Frequency of 13.6 GHz
474	Ka-band:	Frequency of 35.6 GHz
475	PSD:	Particle Size Distribution
476	SVI/PIP:	Snow Video Imager/Particle Image Probe
477	DDA:	Discrete Dipole Approximation
478	$N(D)$ :	Particle Size Distribution
479	$N_w$ :	Scale Factor of Particle Size Distribution
480	$\mu$ :	Shape Factor of Gamma Distribution
481	$D$ :	Particle Diameter
482	$D_m$ :	Mass-weighted Diameter
483	$\Lambda$ :	Slope Parameter
484	DFR:	Differential Frequency Ratio
485	SWC:	Snow Water Content
486	$R_S$ :	Equivalent Snowfall Rate
487	LUT:	Look-up Table
488	$Z_{Ku}$ :	Ku-band Radar Reflectivity Factor
489	$Z_{Ka}$ :	Ka-band Radar Reflectivity Factor
490	$m(D)$ :	Particle Mass as Function of Particle Diameter
491	$V(D)$ :	Particle Fall Velocity
492	$\rho_w$ :	Liquid Mass Density
493	$\rho_s$ :	Snow Mass Density
494	$D_{min}$ :	Minimum Diameter
495	$D_{max}$ :	Maximum Diameter
496	$f$ :	Frequency
497	$\gamma$ :	Aspect Ratio of Particle
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499

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639 Table 1 Snow events during winter of 2014 in Wallops Island, Virginia

640

<b>Events</b>	<b>Start Time (UTC)</b>	<b>End Time (UTC)</b>	<b>Accumulation (mm)</b>	<b>Mean Temperature (°C)</b>
1	JAN03 05:09	JAN03 11:30	46.41	-1.4
2	JAN21 22:05	JAN22 10:03	5.52	-5.4
3	JAN28 20:41	JAN29 12:40	62.12	-9.5
4	FEB14 01:58	FEB14 05:12	6.37	1.9
5	FEB15 20:41	FEB15 23:23	0.76	2.2
6	MAR03 14:40	MAR03 22:00	34.15	-4.4
7	MAR17 08:04	MAR17 20:53	20.93	0.8
8	MAR25 18:50	MAR26 06:13	47.40	1.0

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642 Figure captions:

643 Fig.1 Comparisons of backscattering (left) and extinction (right) coefficients of 3 snow  
644 aggregates with the results from the sphere, oblate and prolate spheroids at a frequency of 35.6  
645 GHz in terms of equivalent ice diameter. A constant snow density of  $0.2 \text{ g/cm}^3$  is assumed for all  
646 spherical and spheroidal particle models. The oblate and prolate spheroids are randomly oriented  
647 with aspect ratios ( $\gamma$ ) of 0.5 and 2, respectively.

648 Fig.2 Backscattering (left) and extinction (right) coefficients from the scattering database (blue  
649 dots), simple scattering model (red curves), which is the randomly-oriented oblate spheroid with  
650 a constant effective density of  $0.2 \text{ g/cm}^3$  and an aspect ratio of 0.7, and scattering-database-  
651 simple-model-combined results (black curves), also referred to as “fitted”, at Ku band. The mean  
652 values of the scattering results are used for the combined results over the data range.

653 Fig.3 Backscattering (left) and extinction (right) coefficients from the scattering database (blue  
654 dots), simple scattering model (red curves), which is the randomly-oriented oblate spheroid with  
655 a constant effective density of  $0.2 \text{ g/cm}^3$  and an aspect ratio of 0.7, and scattering-database-  
656 simple-model-combined results (black curves), also referred to as “fitted”, at Ka band. The mean  
657 values of the scattering results are used for the combined results over the data range.

658 Fig.4 The differential frequency ratio ( $\text{DFR}=10\text{LOG}_{10}(Z_{\text{Ku}}/Z_{\text{Ka}})$ ) as a function of equivalent-  
659 liquid median mass diameter  $D_m$ . (Left): DFR- $D_m$  relations are plotted with several effective  
660 snow densities ( $\rho_s$ ) from 0.1 to  $0.4 \text{ g/cm}^3$  as the shape factor ( $\mu$ ) of the gamma PSD is set to zero  
661 and the aspect ratio ( $\gamma$ ) of the oblate spheroid particles is set to 0.7. (Middle): DFR- $D_m$   
662 relations are plotted with the aspect ratios of 0.5, 0.7 and 1 at  $\rho_s=0.2 \text{ g/cm}^3$  and  $\mu=0$ . (Right): DFR- $D_m$   
663 relations are plotted with the values of  $\mu$  of 0, 3 and 6 at  $\rho_s=0.2 \text{ g/cm}^3$  and  $\gamma=0.7$ .

664 Fig.5 The retrieval look-up tables that show the snow water content (SWC) (top-left) and  
665 equivalent snowfall rate ( $R_s$ ) (top-right), both of which are normalized by the Ku-band radar  
666 reflectivity factor ( $Z_{\text{Ku}}$ ), as a function of the DFR, defined by  $10\text{LOG}_{10}(Z_{\text{Ku}}/Z_{\text{Ka}})$ , for several  
667 effective snow densities ( $\rho_s$ ) with the values from 0.05 to  $0.5 \text{ g/cm}^3$ . The liquid equivalent  
668 median mass diameter  $D_m$  (bottom-left) and the PSD scale parameter  $N_w$  normalized by  $Z_{\text{Ku}}$   
669 (bottom-right) are also plotted in terms of DFR.

670 Fig.6 Flowchart of computing radar parameters and snow size and bulk properties.

671 Fig.7 The look-up tables used for the retrieval of SWC (left) and  $R_s$  (right) with  $\mu$  of 0, 3 and 6,  
672 respectively, as computed from the single scattering tables depicted in Figs.2-3.

673 Fig.8 Example of a segment of the PSD measurements (1000 minutes) in time series taken from  
674 8 snow events during winter of 2014 at the NASA Wallops Flight Facility using the SVI/PIP.  
675 The particle size spectra ( $\text{mm}^{-1} \text{ m}^{-3}$ ), shown in the color scale, are given in the top panel while  
676 equivalent snow fall rate and actual median mass diameter are displayed in the middle and

677 bottom panels, respectively. The PSD data, obtained by averaging the measurements over 1-  
678 minute integration time, are merged from all the snow events into one data file with consecutive  
679 time.

680 Fig.9 The snow water content (SWC) (top row) and equivalent snowfall rate ( $R_s$ ) (bottom row),  
681 both of which are normalized by the Ku-band radar reflectivity factor ( $Z_{Ku}$ ), as a function of the  
682 DFR, defined by  $10\text{LOG}_{10}(Z_{Ku}/Z_{Ka})$ , for several effective snow densities ( $\rho_s$ ) with the values  
683 from 0.05 to 0.5 g/cm<sup>3</sup>. The scatter plots (red dots) are the results derived from the measured  
684 PSD that were collected by the SVP/PIP from 8 snow events in the winter of 2014 at the NASA  
685 Wallops Flight Facility. Two empirical density-size relationships reported by Brandes et al.  
686 (2007) and Heymsfield et al. (2010) are used in converting the measured PSD to the snow mass  
687 spectra, and the results from their respective relations are shown in the left and right panels.

688 Fig.10 Flowchart detailing the procedures of obtaining radar reflectivities and snow parameters.

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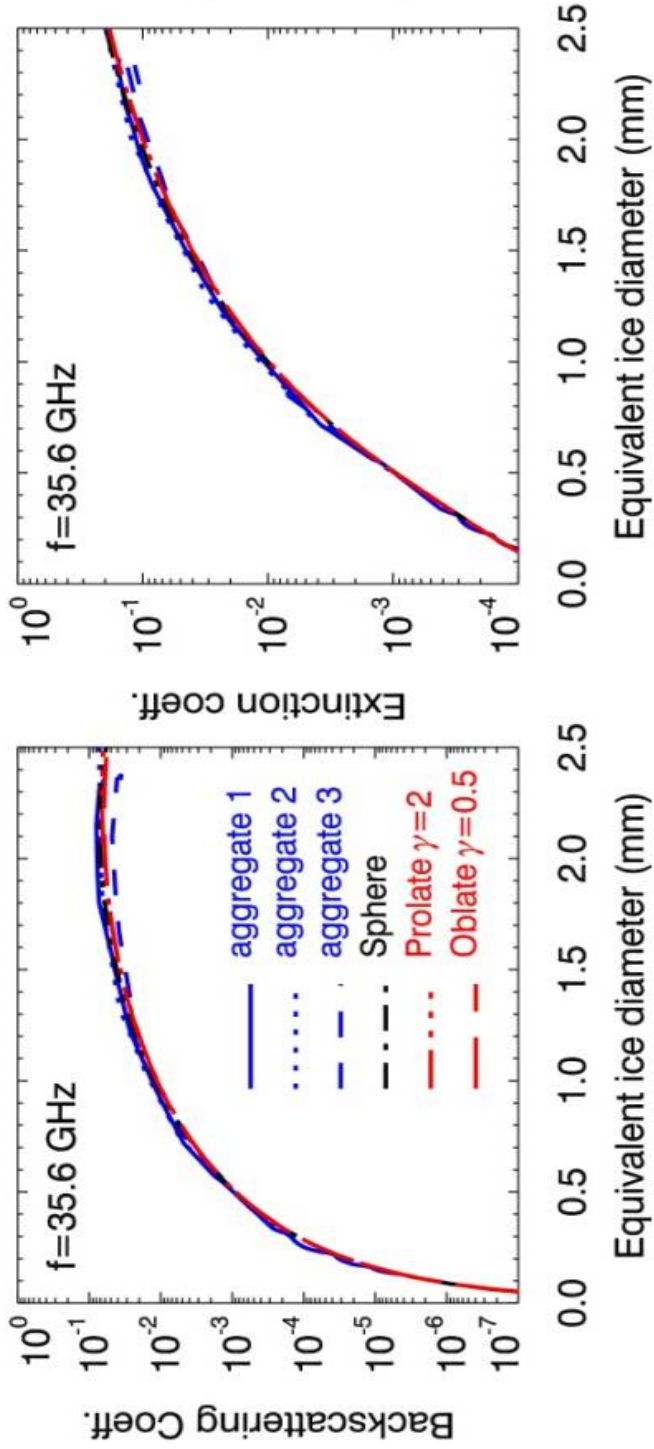
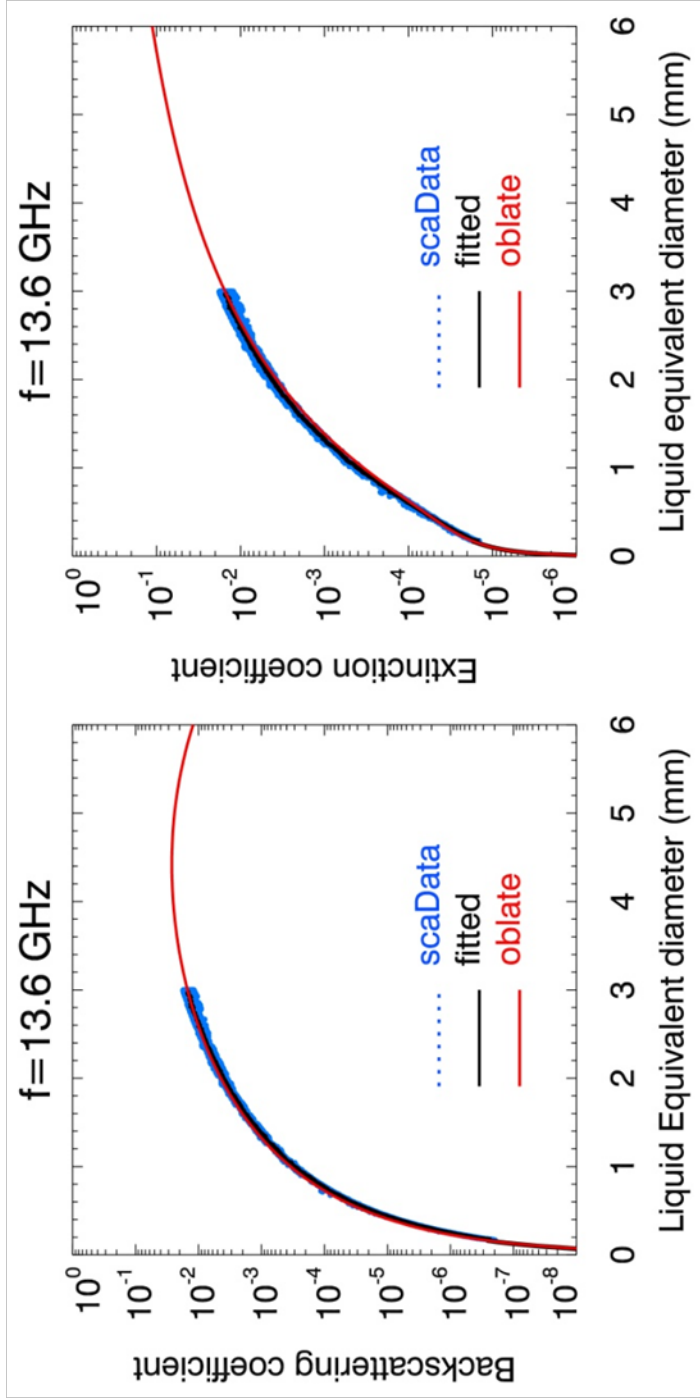


Fig.1 Comparisons of backscattering (left) and extinction (right) coefficients of 3 snow aggregates with the results from the sphere, oblate and prolate spheroids at a frequency of 35.6 GHz in terms of equivalent ice diameter. A constant snow density of  $0.2 \text{ g/cm}^3$  is assumed for all spherical and spheroidal particle models. The oblate and prolate spheroids are randomly oriented with aspect ratios ( $\gamma$ ) of 0.5 and 2, respectively.

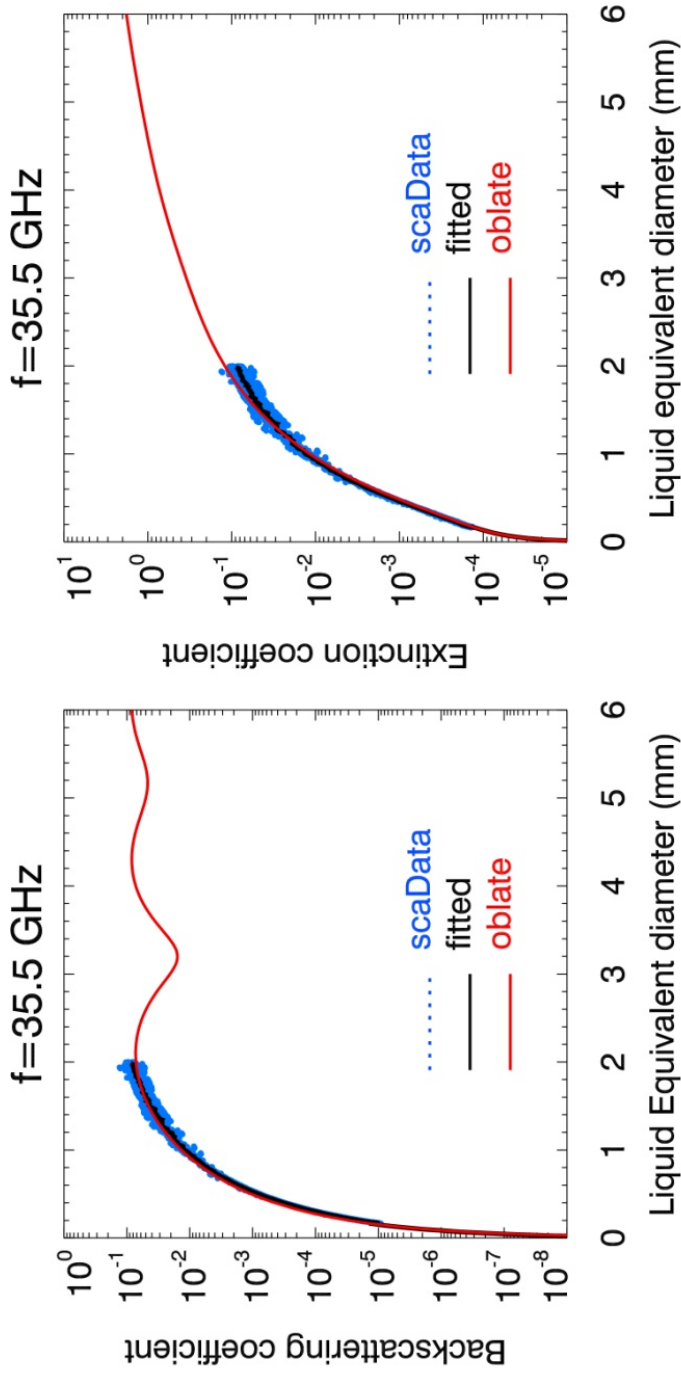
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Fig.2 Backscattering (left) and extinction (right) coefficients at Ku-band from the scattering database (blue dots), simple scattering model (red curves), which is the randomly-oriented oblate spheroid with a constant effective density of 0.2 g/cm<sup>3</sup> and an aspect ratio of 0.7, and scattering-database-simple-model-combined results (black curves), also referred to as "fitted". The mean values of the scattering results are used for the combined results over the data range.

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Fig.3 Backscattering (left) and extinction (right) coefficients from the scattering database (blue dots), simple scattering model (red curves), which is the randomly-oriented oblate spheroid with a constant effective density of  $0.2 \text{ g/cm}^3$  and an aspect ratio of 0.7, and scattering-database-



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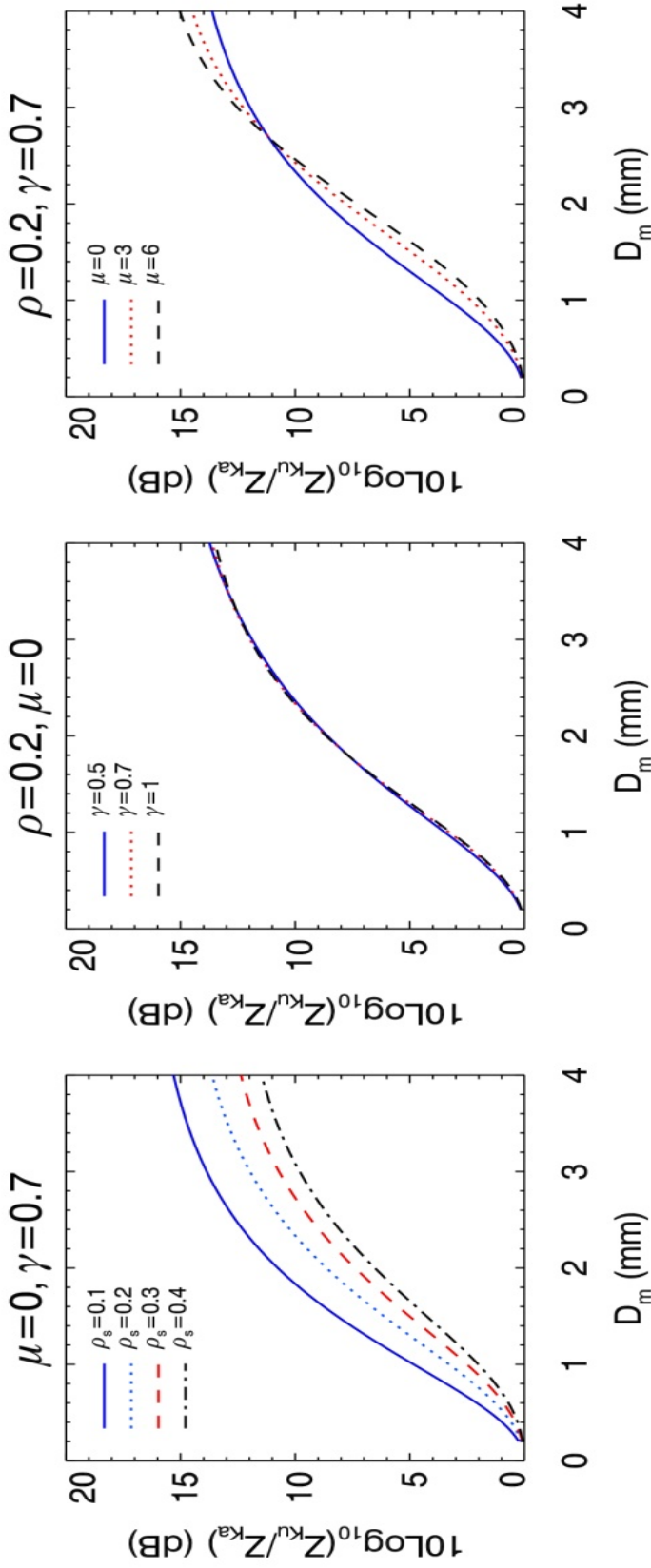
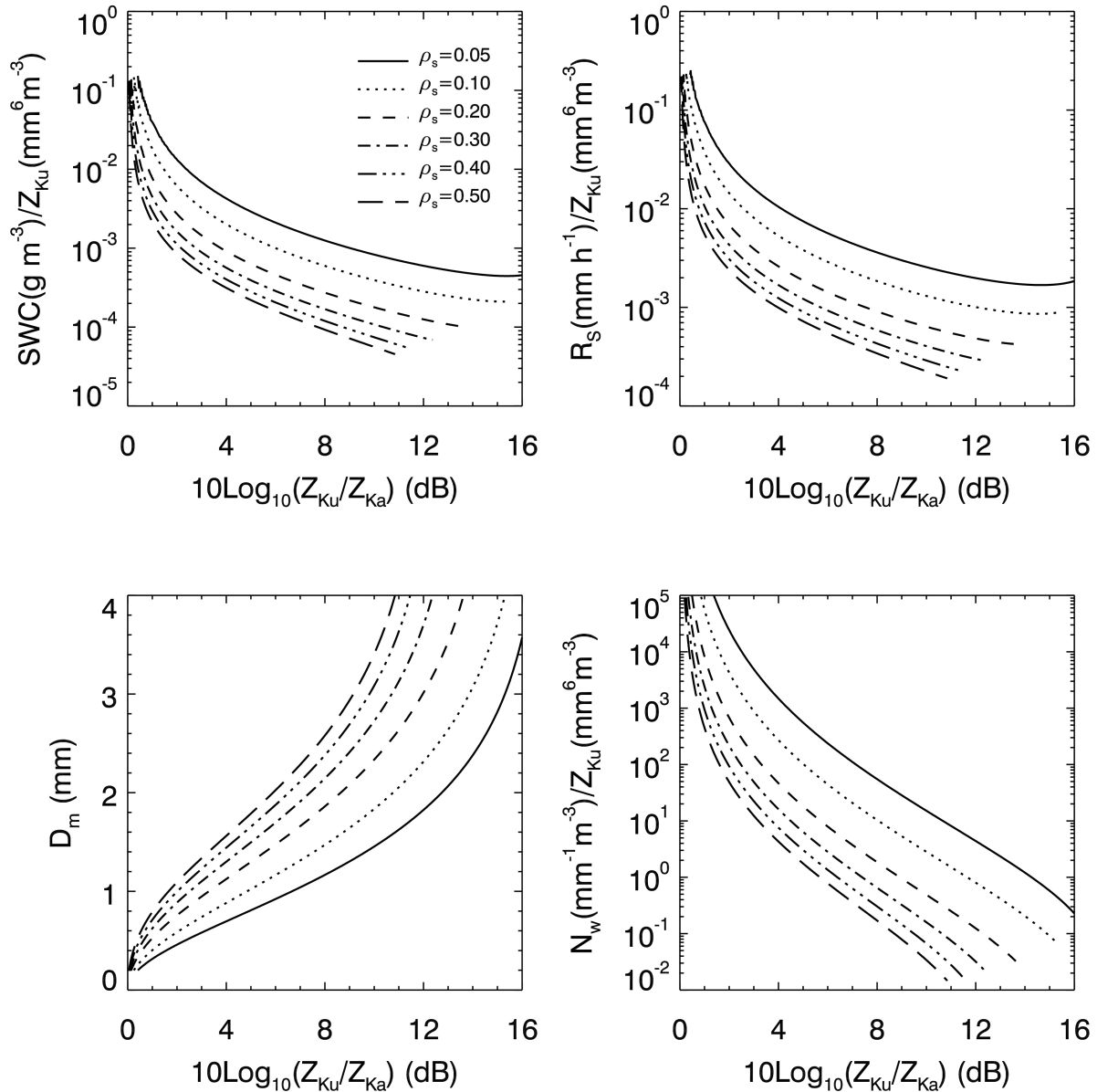


Fig.4 The differential frequency ratio (DFR= $10 \text{LOG}_{10}(Z_{Ku}/Z_{Ka})$ ) as a function of equivalent-liquid median mass diameter  $D_m$ . (Left): DFR- $D_m$  relations are plotted with several effective snow densities ( $\rho_s$ ) from 0.1 to 0.4 g/cm<sup>3</sup> as the shape factor ( $\mu$ ) of the gamma PSD is set to zero and the aspect ratio ( $\gamma$ ) of the oblate spheroid particles is set to 0.7. (Middle): DFR- $D_m$  relations are plotted with the aspect ratios of 0.5, 0.7 and 1 at  $\rho_s=0.2$  g/cm<sup>3</sup> and  $\mu=0$ . (Right): DFR- $D_m$  relations are plotted with the values of  $\mu$  of 0, 3 and 6 at  $\rho_s=0.2$  g/cm<sup>3</sup> and  $\gamma=0.7$ .



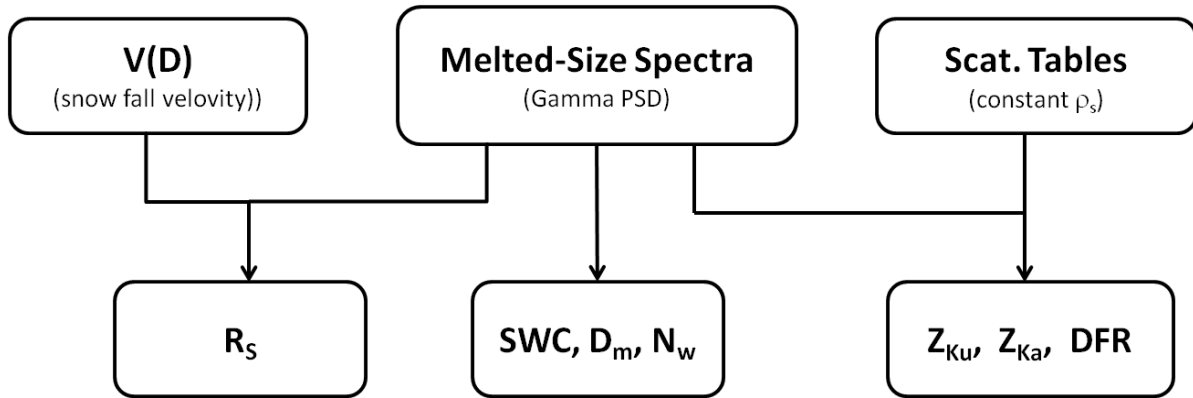
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726 Fig.5 The retrieval look-up tables that show the snow water content (SWC) (top-left) and  
 727 equivalent snowfall rate ( $R_s$ ) (top-right), both of which are normalized by the Ku-band radar  
 728 reflectivity factor ( $Z_{Ku}$ ), as a function of the DFR, defined by  $10\text{LOG}_{10}(Z_{Ku}/Z_{Ka})$ , for several  
 729 effective snow densities ( $\rho_s$ ) with the values from 0.05 to 0.5  $\text{g/cm}^3$ . The liquid equivalent  
 730 median mass diameter  $D_m$  (bottom-left) and the PSD scale parameter  $N_w$  normalized by  $Z_{Ku}$   
 731 (bottom-right) are also plotted in terms of DFR.

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735 Fig.6 Flowchart of computing radar parameters and snow size and bulk properties.

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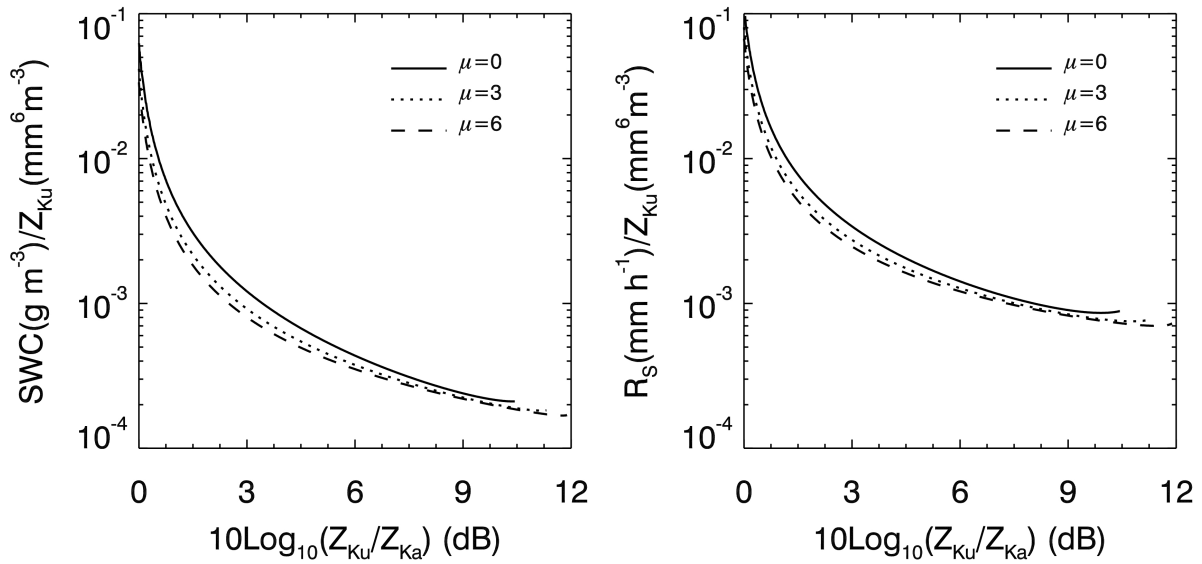
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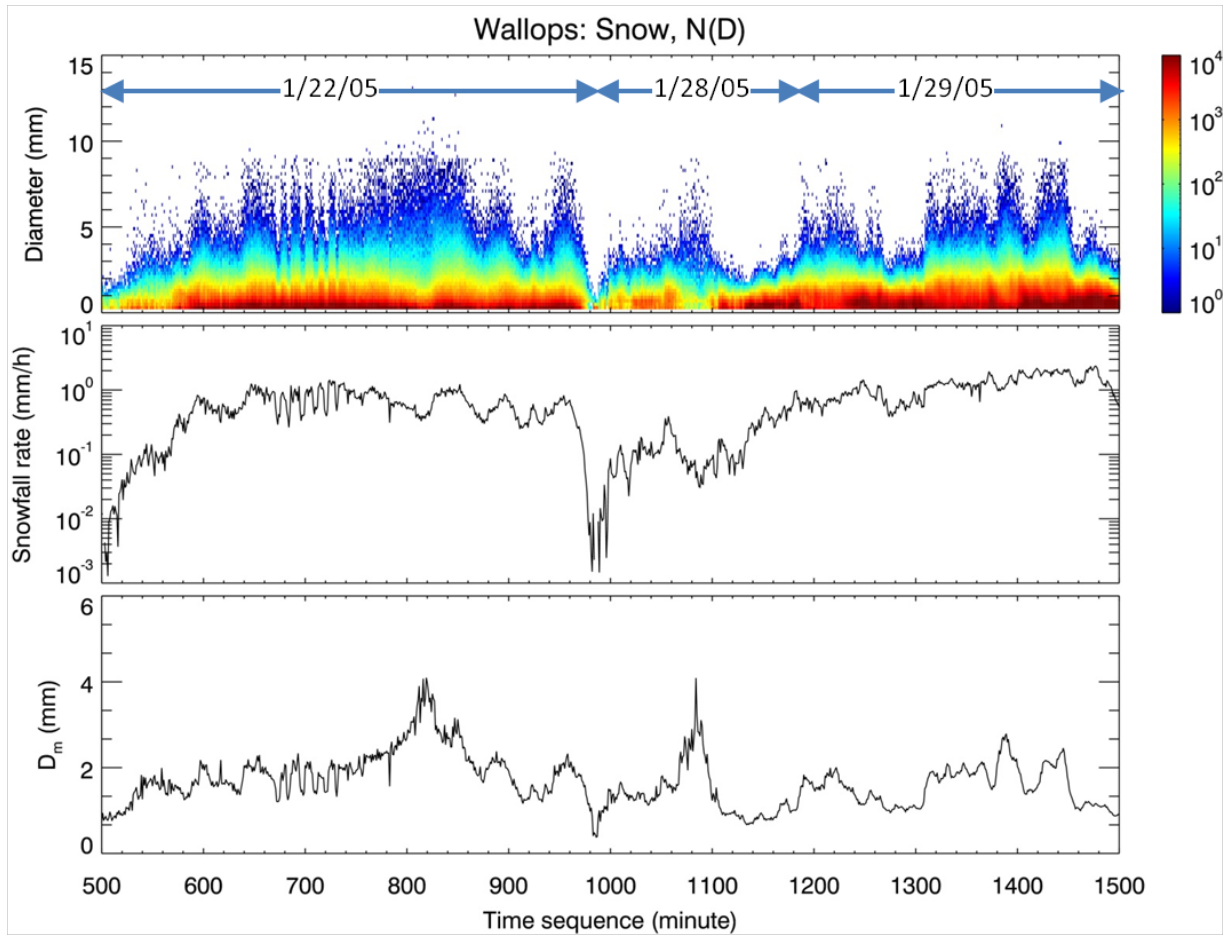
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746



747 Fig.7 The look-up tables used for the retrieval of SWC (left) and R<sub>S</sub> (right) with μ of 0, 3 and 6,  
748 respectively, as computed from the single scattering tables depicted in Figs.2-3.

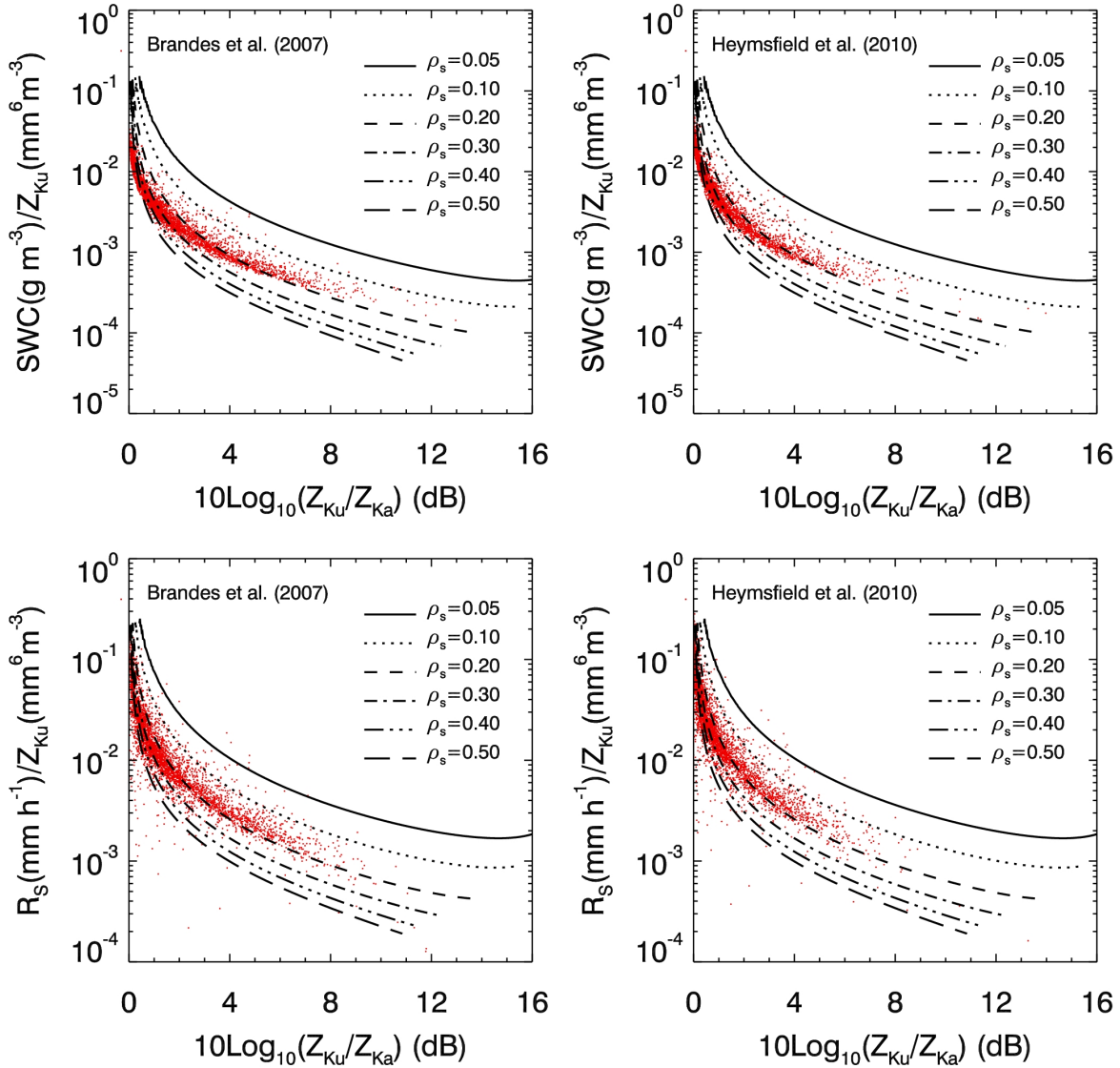
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751 Fig.8 Example of a segment of the PSD data, obtained by averaging the measurements over 1-  
 752 minute integration time, in time series taken from 8 snow events during winter of 2014 at the  
 753 NASA Wallops Flight Facility using the SVI/PIP. The particle size spectra ( $\text{mm}^{-1} \text{m}^{-3}$ ), shown in  
 754 the color scale, are given in the top panel while equivalent snow fall rate and actual median mass  
 755 diameter are displayed in the middle and bottom panels, respectively. The dates provided in the  
 756 top panel correspond to the snow events displayed.

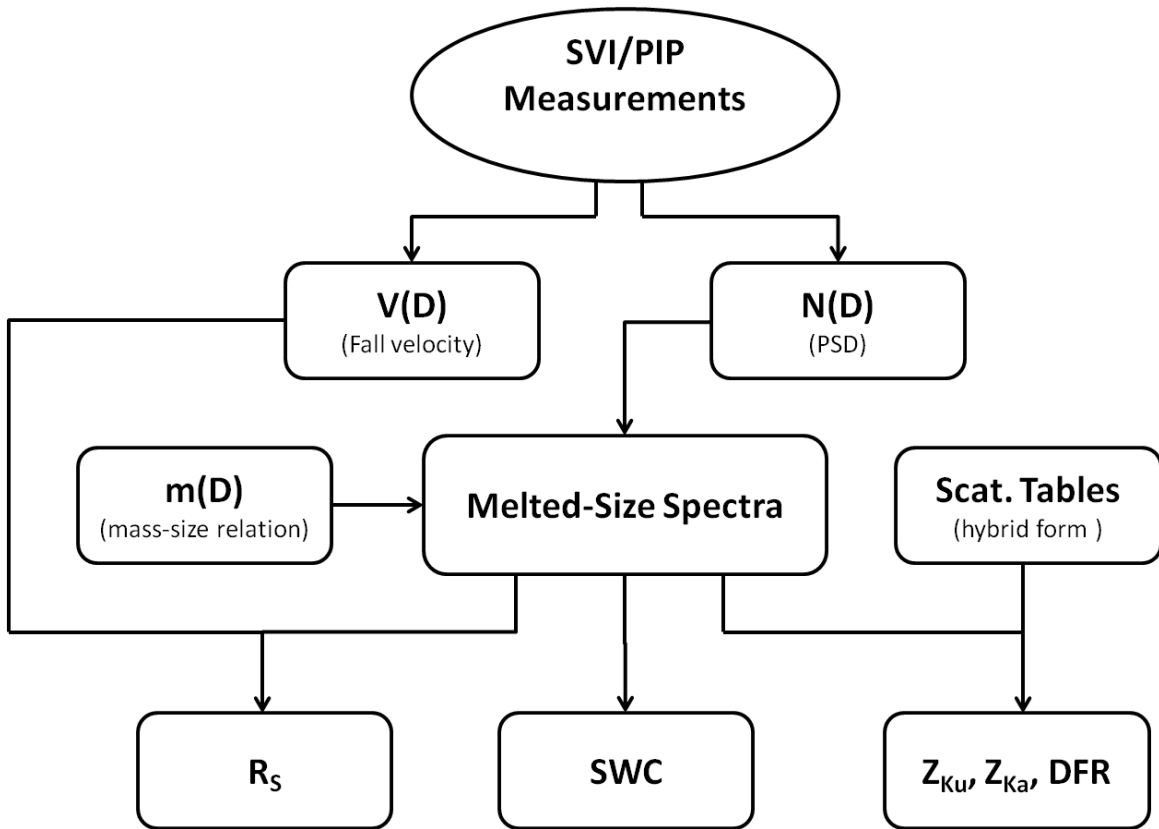
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759 Fig.9 The snow water content (SWC) (top row) and equivalent snowfall rate (R<sub>S</sub>) (bottom row),  
 760 both of which are normalized by the Ku-band radar reflectivity factor ( $Z_{Ku}$ ), as a function of the  
 761 DFR, defined by  $10\text{LOG}_{10}(Z_{Ku}/Z_{Ka})$ , for several effective snow densities ( $\rho_s$ ) with the values  
 762 from 0.05 to 0.5 g/cm<sup>3</sup>. The scatter plots (red dots) are the results derived from the measured  
 763 PSD that were collected by the SVP/PIP from 8 snow events in the winter of 2014 at the NASA  
 764 Wallops Flight Facility. Two empirical density-size relationships reported by Brandes et al.  
 765 (2007) and Heymsfield et al. (2010) are used in converting the measured PSD to the snow mass  
 766 spectra, and the results from their respective relations are shown in the left and right panels.

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770 Fig.10 Flowchart detailing the procedures of obtaining radar reflectivities and snow parameters.

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