

RESEARCH LETTER

10.1002/2016GL069154

Special Section:

First results from NASA's Magnetospheric Multiscale (MMS) Mission

Key Points:

- MHD models can reproduce well the dipolarizations seen at MMS and VAP. Space weather forecasting can predict K_p variations within 0.5 step
- Beams of O^+ flowing downstream appear to cross the separatrix and become a second energized population of the tail plasma sheet
- MHD models successfully reproduced the polar cap convection patterns and cross-polar cap potential drops for a range of IMF conditions

Supporting Information:

- Supporting Information S1
- Figure S1
- Movie S1
- Movie S2
- Movie S3
- Movie S4

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Multispacecraft observations and modeling of the 22/23 June 2015 geomagnetic storm

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Abstract The magnetic storm of 22–23 June 2015 was one of the largest in the current solar cycle. We present in situ observations from the Magnetospheric Multiscale Mission (MMS) and the Van Allen Probes (VAP) in the magnetotail, field-aligned currents from AMPERE (Active Magnetosphere and Planetary Electrodynamics Response), and ionospheric flow data from Defense Meteorological Satellite Program (DMSP). Our real-time space weather alert system sent out a “red alert,” correctly predicting K_p indices greater than 8. We show strong outflow of ionospheric oxygen, dipolarizations in the MMS magnetometer data, and dropouts in the particle fluxes seen by the MMS Fast Plasma Instrument suite. At ionospheric altitudes, the AMPERE data show highly variable currents exceeding 20 MA. We present numerical simulations with the Block Adaptive Tree-Solarwind - Roe - Upwind Scheme (BATS-R-US) global magnetohydrodynamic model linked with the Rice Convection Model. The model predicted the magnitude of the dipolarizations, and varying polar cap convection patterns, which were confirmed by DMSP measurements.

1. Introduction

On 22 June, one moderate and one giant coronal mass ejection (CME) passed the ACE spacecraft at 04:51 and 17:58 UT, respectively. The larger shock was observed by the magnetometer instrument [Smith *et al.*, 1998] as a jump in the interplanetary magnetic field (IMF) from about 10 to over 40 nT and by the Solar Wind Electron, Proton, and Alpha Monitor (SWEPAM) instrument [McComas *et al.*, 1998] as a jump in solar wind density from 20 to over 45 particles/cm³, with a corresponding increase in pressure to over 50 nPa. When propagated to the bow shock, it was forecast to impact at about 18:36 UT after a smaller shock at about 05:40 UT (Figure 1a). Coupled with a strong southward IMF (Figure 1d), the Boyle Index reached nearly 500 kV, prompting our forecast system (<http://mms.rice.edu/realtime/forecast.html>) to send out a “yellow alert” at 06:04 and a “red alert” at 18:34, even before the CME impacted the bow shock. The y component of the IMF was very strong, with the IMF clock angle (Figure 1e), rotating anticlockwise then clockwise nearly 360°. The neural network forecaster [Bala and Reiff, 2012] predicted K_p of over 8 and the K_p forecast status went to “red” at 19:02.

The magnetospheric flotilla of spacecraft included Magnetospheric Multiscale Mission (MMS), Van Allen Probes (VAP), Time History of Events and Macroscale Interactions during Substorms (THEMIS), and Cluster in the magnetosphere, plus AMPERE, Defense Meteorological Satellite Program (DMSP), and the International Space Station (ISS) at low altitudes. This paper will concentrate on the large-scale features of the activity. It will include selected MMS, VAP, AMPERE, and DMSP results and compare to the BATS-R-US model simulations. Other papers in this issue [e.g., Baker *et al.*, 2016; Nakamura *et al.*, 2016] will focus on other aspects of this event.

2. Space Weather Forecasting

The Rice University space forecast system (<http://mms.rice.edu/realtime/forecast.html>) predicts the K_p , AE , and Dst indices for 1 and 3 h ahead of real time, at a cadence of 15 m. The prediction is based on a neural network forecast that uses one or more base functions and lookback data of up to 24 h. Yellow alerts are sent out if the forecast K_p is greater than 4, and red alerts if the forecast K_p is 6 or greater [Bala and Reiff, 2012]. Statistically, the accuracy of the predicted K_p is approximately one unit. Three base formulas are used in the forecast:

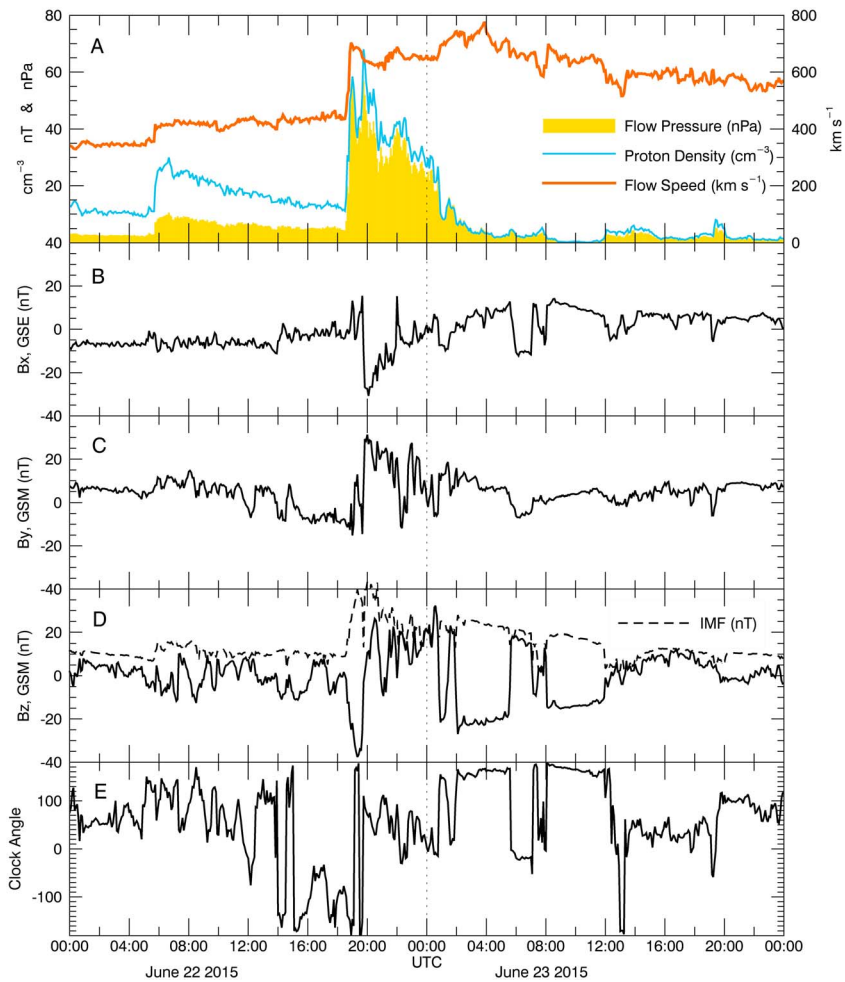


Figure 1. Interplanetary conditions during the event. (a) Solar wind density, velocity, and flow pressure; (b–d) IMF B_x , B_y , B_z (and $|B|$ shown dashed) components (GSM); (e) IMF clock angle ($0 = +B_z$; $90 = +B_y$ direction). The values are plotted from OMNI data, propagated to the bow shock.

the Boyle Index (marked as BI), the Boyle Index with a ram pressure term (RAM), and the Newell function (Newell) [Bala and Reiff, 2014]. All have similar prediction efficiencies, and the three predictions generally span the actual measured quantities. For this event, all three formulas gave similar predictions, with the AE index least well predicted ($r = 0.61$ to 0.72), whereas Kp and Dst were successful with $r = 0.79$ to 0.91 (Figure 2).

A comment on timing of “predictions” is in order. The forecast algorithm was trained on 1 h averages of AE , Dst , and Kp^* , with the Kp^* values being an overselection of Kp (because it is intrinsically a 3 h average). The algorithm predicts the 1 h AE , Dst , and Kp^* values for the hour following, based on the solar wind data of the previous hour. The calculation is done at the top of the hour after the real-time solar wind data are received, and a forecast typically posted at 6 min past the hour for the hour of the forecast (e.g. 18:06 for the hours 18–19). In this case, because the shock hit at the very end of hour 17, the AE predictions for hours 18–19 were not as high as the provisional AE turned out to be (Figure 2c), because the forecast only included 2 min of postshock data and because the high solar wind speed meant that the shock arrived at Earth well before the end of hour 18. The forecast 3 h Kp , based on 3 h averages of the 1 h predictions (Figure 2a), does show a very good fit ($r = 0.88$ to 0.91 , implying an accuracy of prediction from 0.36 to 0.47 step in Kp).

3. Magnetotail Observations

Since the MMS spacecraft suite [Burch et al., 2016] was in the “commissioning” phase of its mission, not all of the instruments were fully operational. All of the magnetometers [Russell et al., 2016] and many of the particle

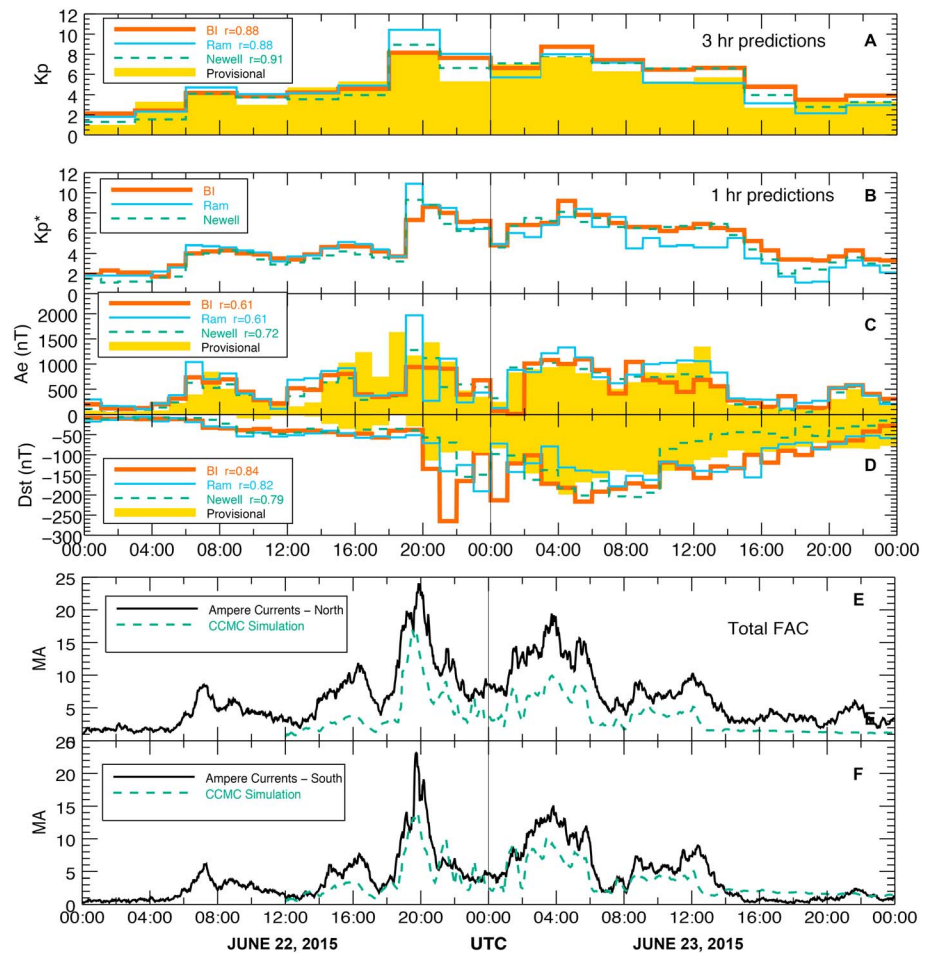


Figure 2. Comparison of predicted and modeled values compared to data. (a) The predicted 3 h K_p index for 22–23 June 2015, compared to the provisional K_p , with a 90% correlation coefficient. Predicted (histograms) and observed provisional (yellow) (b) K_p^* , (c) AE , and (d) Dst indices. All three base functions: BI (thick line), Ram (thin line), and Newell (dashed) performed well, some slightly better than others. Modeled integrated field-aligned currents (dashed) versus those calculated from AMPERE measurements for the (e) Northern and (f) Southern Hemispheres.

detectors were making measurements and saw the dynamic response of the magnetosphere. The Fast Plasma Instrument (FPI) on board MMS2 [Pollock *et al.*, 2016] saw a number of excursions between the plasma sheet and the lobe, shown in Figure 3 as particle flux dropouts at about 3:20–3:30 and 5:11–05:40 UT (Figures 3c and 3d). Clearly visible in the lobe are antisunward flowing ions, at a few hundred eV (Figures 3b and 3c). The Hot Plasma Composition Analyzer (HPCA) on board MMS1 [Young *et al.*, 2016] identified these beams as O^+ (Figure 3b) presumably from auroral outflow [Lu *et al.*, 1992] or possibly from the dayside cusp [Liao *et al.*, 2010] though cusp ions will more likely reach the neutral sheet farther down-tail [Liao *et al.*, 2012]. The HPCA instrument, because of a new AC sweep field, can reduce background H^+ fluxes by nearly 2 orders of magnitude, making the determination of the heavy ion species much more reliable. This cool lobe ion beam appears to be accelerated and heated as it crosses into the plasma sheet. In Figures 3e–3g, we show sample total ion particle distributions from FPI MMS2 as the spacecraft exits the plasma sheet into the lobe around 05:10:30. In the lobe around 05:11:19 (Figure 3g), we see two separated cold ion beams: one with a parallel velocity of around 50 km/s and another with a parallel velocity of ~ 200 km/s. The right image of that pair shows that those beams are also convecting at about the same $\mathbf{E} \times \mathbf{B}$ speed ($v_{\text{perp}1}$) as the parallel speeds, with no drift in the other direction ($v_{\text{perp}2}$) perpendicular to Figure 3b. From the HPCA data (Figures 3a and 3b), it is clear that the less energetic ions are H^+ and the ones with apparently higher velocity are just O^+ ions with the same parallel and perpendicular velocity as the H^+ ions. This is confirmed by the HPCA particle distribution plots (not shown). As the spacecraft exits

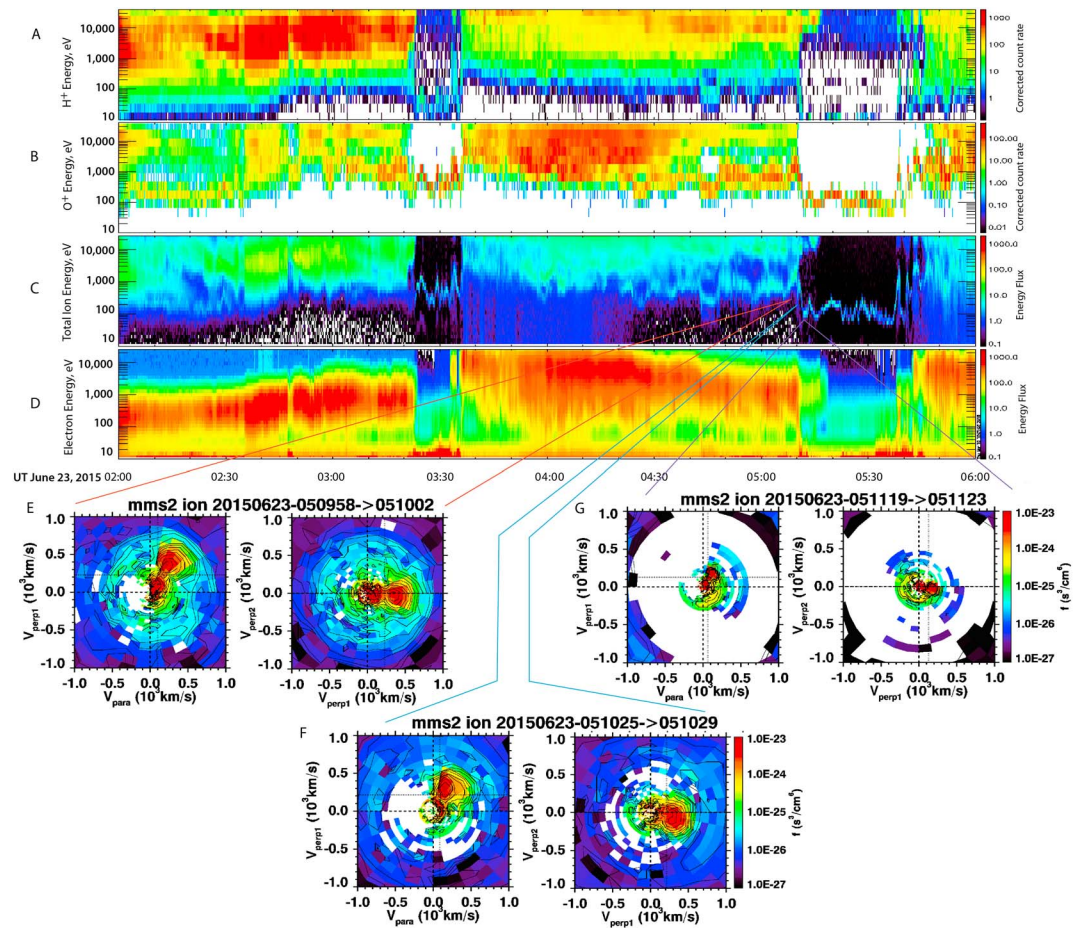


Figure 3. HPCA (MMS1) and FPI (MMS2) measurements in the magnetotail from 0200 to 0600 on 23 June. Energy spectrograms of the HPCA (a) H^+ and (b) O^+ corrected counts are shown with Figure 3b showing an O^+ beam from the ionosphere merging into the plasma sheet at each lobe/plasma sheet transition. (c) FPI ion and (d) electron spectrograms of differential energy flux are shown. (e–g) Three pairs of FPI distribution functions as the spacecraft exited the plasma sheet to the lobe. (left) The $v_{parallel}$ and v_{perp1} (along $\mathbf{E} \times \mathbf{B}$) components of the particle distribution functions and (right) the two perpendicular velocities, v_{perp1} and v_{perp2} . Similar distribution functions from HPCA are shown in the supporting information.

the plasma sheet boundary layer at around 05:10:30 (Figure 3f), we see the same two beams, but now with a $\mathbf{E} \times \mathbf{B}$ velocity double that in the lobe, This middle distribution does show some v_{perp2} , but it is not clear whether this is a time aliasing as the fields change direction during the 4 s measurement of the distribution function. The distribution function measured deepest in the plasma sheet (Figure 3e, at 05:09:58) shows the highest parallel and $\mathbf{E} \times \mathbf{B}$ drift speeds, and no significant v_{perp2} . For additional information on this event, see Nakamura *et al.* [2016].

At the same time, the magnetic fields observed at MMS in the tail and by VAP closer to the Earth showed dramatic dipolarizations as the magnetotail responded to the northward turnings of the IMF (Figure 4). At MMS-1 (Figure 4, top), the measured B_x went from -90 to -20 nT, and B_z increased from near 0 to 30 nT. The transitions from MMS from the plasma sheet to the lobe resulted primarily because of the thinning and expansion of the plasma sheet and partially because of the flapping of the magnetotail current sheet up and down.

4. Modeling Results

To put these observations into context, we ran the BATS-R-US model [Powell *et al.*, 1999; Tóth *et al.*, 2005, 2012] with Rice Convection Model in highest resolution available from the Coordinated Community Modeling Center (CCMC) (5 min cadence and $1/8 R_E$ (Earth radius) at the inner boundary), using the measured

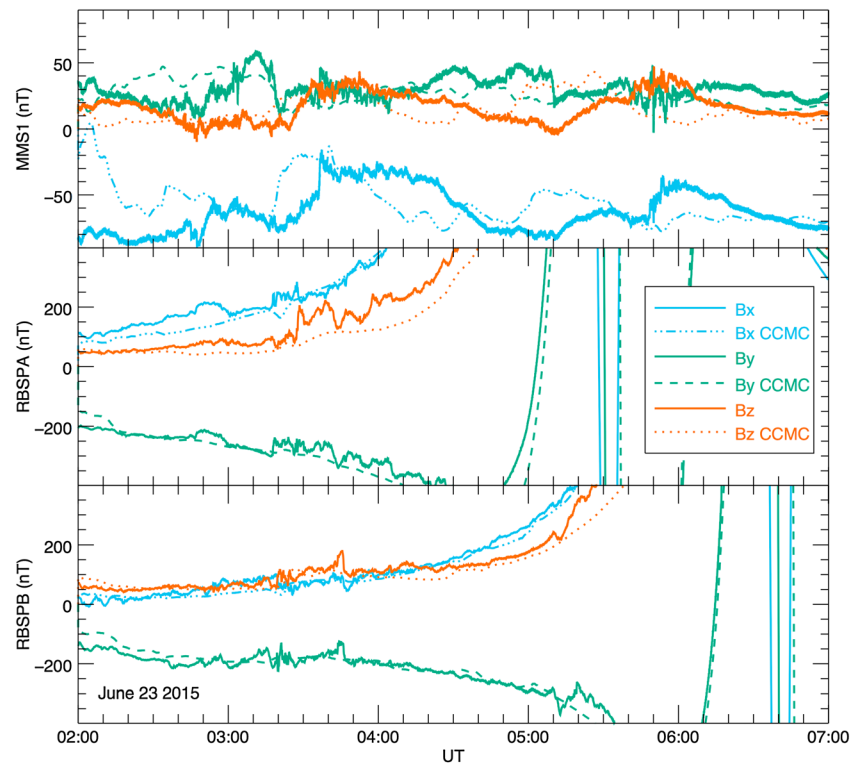


Figure 4. Tail magnetic fields observed by (top) MMS1 and Van Allen Probes (middle) A and (bottom) B (solid lines) and modeled (dashed lines) using a high-resolution run of BATS-R-US from the CCMC.

propagated solar wind and IMF from OMNIdata. The model showed many tail reconnection events in the time frame of 00:00 UT on 22 June to 24:00 UT on 23 June. Movies showing various cuts of the model for this event, including a cut that dynamically followed the MMS trajectory, can be found in the supporting information and at <http://mms.rice.edu/June22>. The model did especially well at capturing the several stretching and rapid dipolarizations observed in the 2–6 UT 23 June time frame (Figure 4, dashed). The first dipolarization in the model was between 3:15 and 3:20 (Figure 5a) and was observed at MMS at 3:16 [Nakamura *et al.*, 2016]. The second dipolarization in the model occurred between 04:55 and 05:00, seen at MMS at ~5:10 (see Movie S1 in the supporting information). The magnitude of the field changes predicted along the MMS path was quantitatively accurate, but with some variances in timing, reaching peak values somewhat earlier than measured. The model suggests that at the times of the dipolarizations, an X line is just downstream of MMS, which is very near the separatrix (Figure 5a). Since the tail was so stretched, a modest flapping of the tail affects the location of MMS relative to the model. The dipolarization signatures observed at VAP RBSP-A (Figure 4, middle) were accurate in location but low in amplitude. The dipolarization signatures observed at RBSP-B (Figure 4, bottom) showed good agreement both in magnitude and location. Thus, we argue that the O^+ ions seen flowing downstream in the lobe near the edge of the plasma sheet may be captured by the reconnection process and become the energetic O^+ seen in the plasma sheet.

When mapped to the ionosphere, the model showed extremely large field-aligned currents (FAC). As compared to the FAC's inferred from the AMPERE data [e.g., Coxon *et al.*, 2014], which reached an integrated value of nearly 25 MA, the predicted magnitudes were 20–50% smaller than that observed, although the timing structure was quite accurate (Figures 2e and 2f).

5. Polar Cap Convection Patterns

Because of the large magnetic field strengths (~ 40 nT) and the large and variable IMF y component, the model-predicted ionospheric convection varied from normal two-cell convection patterns during times of

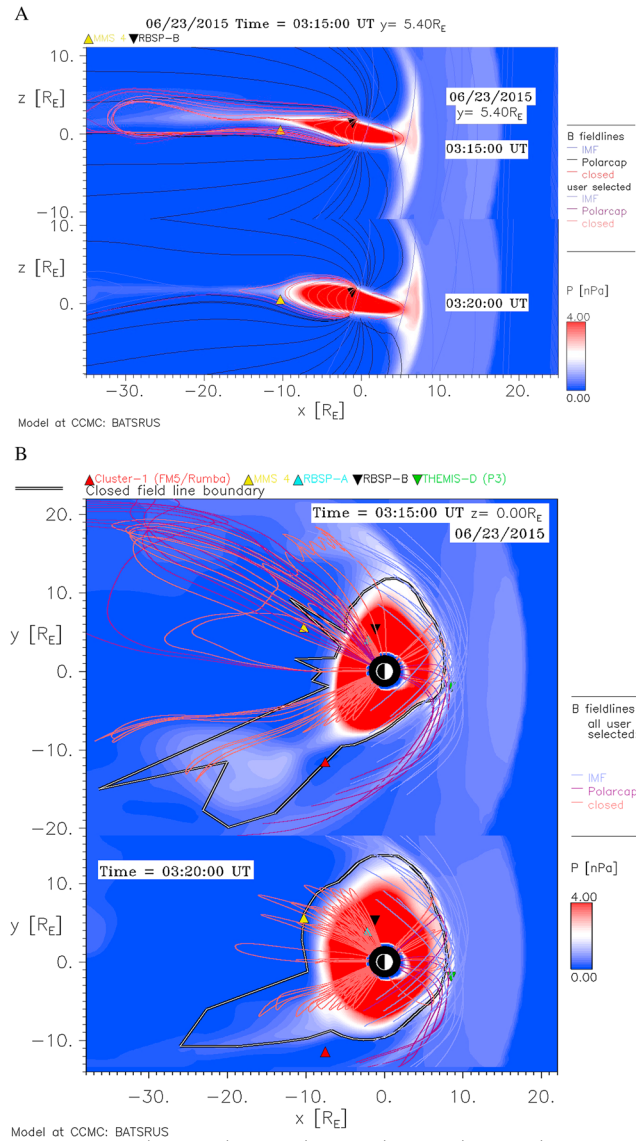


Figure 5. (a) Two frames from a movie of the CCMC run, showing MMS 1 spacecraft located very near the separatrix of an X line in the magnetotail. The dipolarization in the model is well observed between the stretched configuration of 03:15 and the more dipolar configuration of 03:20. The full movie including similar frames for the 05:00 dipolarization is in the supporting information and is available at <http://mms.rice.edu/June22/>. (b) Similar to Figure 5a, but showing the equatorial plane locations of the various magnetospheric spacecraft (MMS, VAP A and B, THEMIS, and Cluster) at the time of the dipolarization. The field lines in both times were started at the same locations on the equatorial plane. Note how less stretched each field line is after the dipolarization. A movie of this view is provided in the supporting information and on our website <http://mms.rice.edu/June22>.

large negative B_z (Figure 6a), to a four cell during high-positive IMF B_z intervals (Figure 6b) to a single convection cell rotating in opposite directions (clockwise in the North, counterclockwise in the South) in the two hemispheres after a long interval of strong $+B_y$ (Figure 6c). A true single-cell convection pattern is very rare: in this case the single cell was only predicted (and observed) in the North. A movie of the modeled convection patterns (Movie S2) is part of the supporting information and posted at <http://mms.rice.edu/June22>.

A comparison with DMSP plasma flow data from the same intervals confirms the basic features of the convection, but interesting differences are observed. For example, in the southward IMF case (Figure 6a), the polar cap is larger in the data than in the model by about 5°. The sunward flow channel in the data during northward IMF (Figure 6b) is much narrower than in the model, which predicts a very large area of sunward flow in

Ionospheric Convection Patterns June 23, 2015

Left: CCMC model (BATS-R-US)

Right: Crosstrack Plasma Drift DMSP (North)

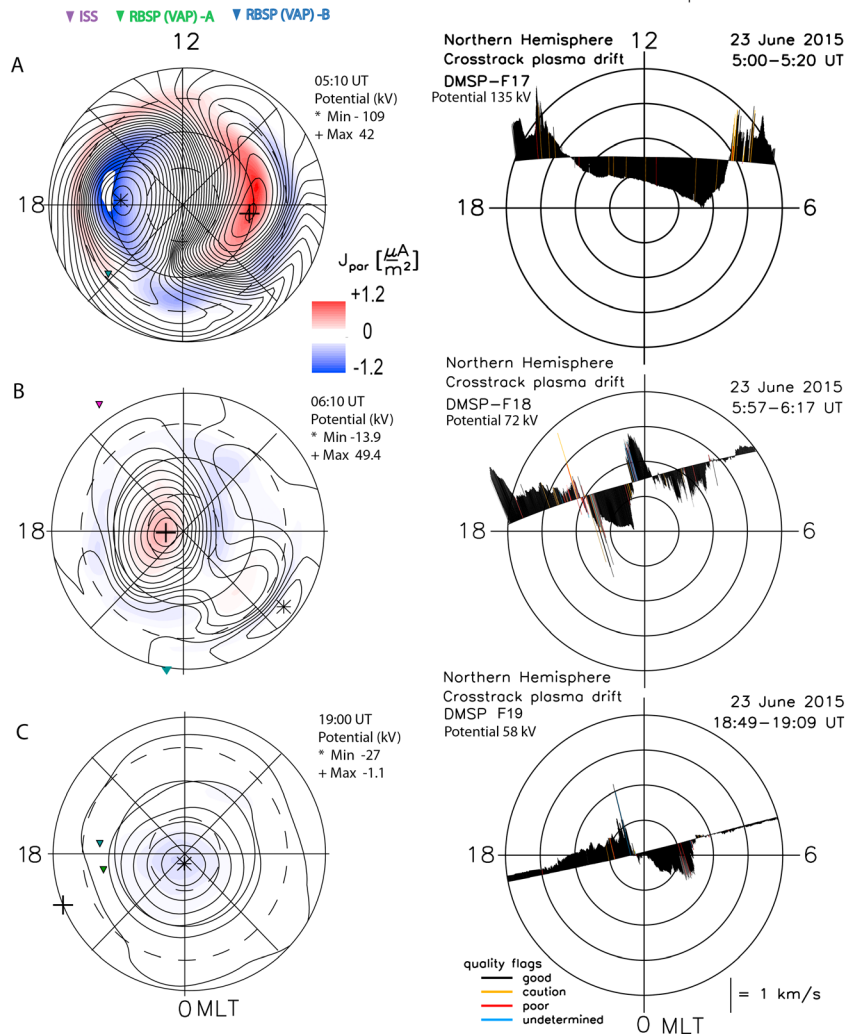


Figure 6. Northern polar cap convection equipotentials predicted from the BATS-R-US run (left image in each row). Colored areas are the field-aligned current densities, and equipotential lines are 4 kV apart in each plot. The movie is available at <http://mms.rice.edu/June22> and in the supporting information. (right column) The measured cross-track plasma drifts from DMSP at the same time as the predicted potentials. (a) A normal two-cell convection pattern during southward IMF, (b) a four-cell pattern with reversed flow in the central polar cap during strong northward IMF, (c) a single-cell clockwise convection observed during strong positive Y component of the IMF.

the polar cap. The model also missed the strong low-latitude sunward flows on the dusk side. The DMSP flow data confirm the presence of a single-cell convection pattern in the North (Figure 6c, right), with just a hint of a viscous cell at 7 MLT, 75 ILAT. The cross-polar cap potential drops estimated from the model and from the flow data are comparable. Both the model and the data show the center of the cell just duskward of the magnetic pole [Reiff and Burch, 1985]. The lobe cell convection is more common in times of high sunward dipole tilt [Crooker, 1992], and this event, which occurred on 22–23 June, had maximum sunward tilt in the North.

6. Conclusions

This event represents the first major storm of the new Heliospheric Flotilla era. With well-instrumented spacecraft strategically placed in the magnetosphere, and new computational models, our understanding of

magnetospheric dynamics, especially its response during dramatic events such as the one presented in this paper, is taking a leap forward. Despite this being an anomalously intense event with large magnetic fields, the BATS-R-US model did an admirable job of reproducing the amount of field change during the dipolarizations, estimating the polar cap convection and currents, and the approximate location of MMS near the separatrix during a very dynamic magnetotail sequence.

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