

A CLASS OF LARGE GLOBAL SOLUTIONS FOR THE WAVE-MAP EQUATION

ELISABETTA CHIODAROLI AND JOACHIM KRIEGER

EPFL Lausanne
Station 8, CH-1015 Lausanne, Switzerland

ABSTRACT. In this paper we consider the equation for equivariant wave maps from \mathbb{R}^{3+1} to \mathbb{S}^3 and we prove global in forward time existence of certain C^∞ -smooth solutions which have infinite critical Sobolev norm $\dot{H}^{\frac{3}{2}}(\mathbb{R}^3) \times \dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$. Our construction provides solutions which can moreover satisfy the additional size condition $\|u(0, \cdot)\|_{L^\infty(|x| \geq 1)} > M$ for arbitrarily chosen $M > 0$. These solutions are also stable under suitable perturbations. Our method, strongly inspired by [9], is based on a perturbative approach around suitably constructed approximate self-similar solutions.

1. INTRODUCTION

1.1. Corotational wave maps. Let M be the Minkowski space $\mathbb{R}^{n,1}$, with coordinates $x = (x^0, x^1, \dots, x^n) = (t, \bar{x})$ and let N be a smooth, complete, rotationally symmetric k -dimensional Riemannian manifold without boundary. Following Tachikawa [14], we can then identify N , as a warped product, with a ball of radius $R \in \mathbb{R}^+ \cup \{\infty\}$ in \mathbb{R}^k equipped with a metric of the form

$$ds^2 = du^2 + g^2(u)d\theta^2, \quad (1.1)$$

where (u, θ) are polar coordinates on \mathbb{R}^k , $d\theta^2$ is the standard metric on the sphere \mathbb{S}^{k-1} , and $g : \mathbb{R} \rightarrow \mathbb{R}$ is smooth and odd,

$$g(0) = 0, \quad g'(0) = 1. \quad (1.2)$$

With these notations in hand, we can define a *wave map* $U : M \rightarrow N$ as a stationary point (with respect to compactly supported variations) of the functional

$$\mathcal{L}[U] = \frac{1}{2} \int_M \langle \partial_\mu U, \partial^\mu U \rangle = \frac{1}{2} \int_M \partial_\mu u \partial^\mu u + g^2(u) \Gamma_{A,B} \partial_\mu \theta^A \partial^\mu \theta^B. \quad (1.3)$$

Hence, they satisfy the Euler-Lagrange equations: if we denote the vector valued map U as $U := (u, \theta) \in \mathbb{R}^k$ then it satisfies

$$\begin{cases} \partial_\mu \partial^\mu u + g(u)g'(u)\Gamma_{A,B}\partial_\mu \theta^A \partial^\mu \theta^B = 0 \\ \partial^\mu (g^2(u)\Gamma_{A,B}\partial_\mu \theta^B) = 0. \end{cases} \quad (1.4)$$

We also introduce spatial polar coordinates $(t, r, \omega) \in \mathbb{R} \times \mathbb{R}^+ \times \mathbb{S}^{n-1}$ on M . In these coordinates the metric on M takes the form

$$ds^2 = -dt^2 + dr^2 + r^2 d\omega^2.$$

Since N is assumed to be rotationally symmetric, it becomes natural to consider *equivariant* wave maps by requiring that the orbit of any point in M under spatial rotations maps into the orbit of the image point in N . When this mapping of the orbits has degree 1, we call the map *corotational*. We suppose to be in the equivariant framework, thus we require that

$$u = u(t, r) \quad \text{and} \quad \theta = \theta(\omega). \quad (1.5)$$

It then follows that $\theta : \mathbb{S}^{n-1} \rightarrow \mathbb{S}^{k-1}$ has to be an *eigenmap*, i.e. a harmonic map of constant energy density

$$e = |\nabla_\omega \theta|^2.$$

Under this hypothesis, the wave map system (1.4) simplifies and reduces to the following simple scalar wave equation for the spatially radial function $u : M \rightarrow \mathbb{R}$

$$u_{tt} - u_{rr} - \frac{n-1}{r}u_r + \frac{f(u)}{r^2} = 0, \quad f(u) = eg(u)g'(u) \quad (1.6)$$

1.2. Wave maps into \mathbb{S}^3 . We are interested in equivariant corotational wave maps from $M = \mathbb{R}^{3+1} \rightarrow \mathbb{S}^3 = N$. In this case, it is immediate to represent the target manifold \mathbb{S}^3 , as a warped product. Let us show the standard way how the standard 3-dimensional sphere \mathbb{S}^3 can be written with a rotationally symmetric metric of the type $du^2 + g^2(u)ds_2^2$ where ds_2^2 is the canonical metric on $\mathbb{S}^2 \subset \mathbb{R}^3$. We consider the map

$$I : (0, \pi) \times \mathbb{S}^2 \rightarrow \mathbb{R} \times \mathbb{R}^3 \\ I(u, \theta) = (\cos u, \sin u \cdot \theta)$$

which maps into the unit sphere in \mathbb{R}^4 . In order to see that I is a Riemannian isometry we compute the canonical metric on $\mathbb{R} \times \mathbb{R}^3$ (“can”) using the coordinates $(\cos u, \sin u \cdot \theta)$. In order to carry out the computation we use that

$$1 = (\theta^1)^2 + (\theta^2)^2 + (\theta^3)^2 \\ 0 = 2(\theta^1 d\theta^1 + \theta^2 d\theta^2 + \theta^3 d\theta^3).$$

Thus we obtain

$$\begin{aligned} \text{can} &= (d \cos u)^2 + \sum \delta_{i,j} d(\sin u \theta^i) d(\sin u \theta^j) \\ &= \sin^2 u du^2 + \sum \delta_{i,j} (\theta^i \cos u du + \sin u d\theta^i) (\theta^j \cos u du + \sin u d\theta^j) \\ &= \sin^2 u du^2 + \cos^2 u du^2 + \sin^2 u \left(\sum (d\theta^i)^2 \right) \\ &= du^2 + \sin^2 u \left(\sum (d\theta^i)^2 \right), \end{aligned}$$

and the claim follows from the fact that $\sum (d\theta^i)^2$ is exactly the canonical metric ds_2^2 . According to our notation, the function g is here chosen to be $g(u) := \sin(u)$ and it satisfies $g(0) = 0$ and $g'(0) = 1$. Since we are dealing with $M = \mathbb{R}^{3+1}$, we have here $n = k = 3$: for $(t, r, \omega) \in \mathbb{R} \times \mathbb{R}^+ \times \mathbb{S}^2$ on M and $U := (u, \theta) \in \mathbb{R} \times \mathbb{S}^2$ we can choose the map $\theta = \theta(\omega)$ to be the identity map from \mathbb{S}^2 to \mathbb{S}^2 , whence $e = 2$. The same “ansatz” was also considered in [10] for harmonic maps into spheres. It is

now clear that the problem of looking for equivariant and corotational wave maps from \mathbb{R}^{3+1} to \mathbb{S}^3 reduces to solving the following equation (see (1.6))

$$u_{tt} - u_{rr} - \frac{2}{r}u_r + \frac{f(u)}{r^2} = 0, \quad f(u) = 2 \sin u \cos u \quad (1.7)$$

on \mathbb{R}^{1+1} .

2. BACKGROUND LITERATURE AND MAIN RESULT

The initial value problem for the wave map equation has drawn a lot of attention in the mathematical community in the last twenty years. In particular, wave maps have been studied extensively in the case of a flat background M . It is impossible here to account for the large amount of publications in the subject, hence we will quote some important contributions.

Let us make a short digression to explain the main problems related to the wave maps equation with particular attention to the case of flat backgrounds. As for other nonlinear evolutionary equations the first issues one is confronted with are the existence of global classical solutions and the development of singularities. The local-in-time existence is rather standard but it is challenging to identify classes of initial data for which global existence holds or on the contrary the identification of specific initial data that lead to a breakdown (blow-up) of the solution in finite time. If blow-up occurs it is interesting to get the idea behind its formation. For scaling invariant equations with a positive energy, such as the wave maps equation (see [6]), heuristically one would expect finite-time blow-up when the shrinking of the solution is energetically favorable, i.e. in the so called energy supercritical case, while global existence should occur when shrinking to smaller scales is energetically prohibited, i.e. in the energy subcritical space. The limit case when the energy itself is scaling invariant is called energy critical. For wave maps, the criticality classification depends on the spatial dimension of the base manifold M . The equation is energy subcritical, critical, or supercritical if $\dim M = 1 + 1$, $\dim M = 2 + 1$ or $\dim M \geq 3 + 1$, respectively.

As anticipated in the previous section, in this paper we study the simplest energy supercritical case: corotational wave maps from $(3 + 1)$ - Minkowski space to the three sphere which satisfy equation (1.7). Global well-posedness of the Cauchy problem for this equation when data are small in a sufficiently high Sobolev space follows from [13]. The existence of global weak solutions is shown in [11] for any initial data of finite energy, but otherwise arbitrary. Furthermore, a number of results concerning the Cauchy problem for equivariant wave maps are obtained in [12]; in particular, local well-posedness with minimal regularity requirements for the initial data is studied. On the other hand, in [11], Shatah has shown that in the case $M = \mathbb{R}^{3+1}$ and $N = \mathbb{S}^3$ the corotational wave map problem (1.7) admits self-similar solutions, thus making it possible to pose Cauchy problems with smooth data whose solutions develop singularities in finite time. This initiated the construction of blow-up solutions in the form of self-similar solutions: for further generalizations to more general targets and examples of non-uniqueness using self-similar blow-up profiles see [12] and [2]. As usual, by exploiting finite speed

of propagation, a self-similar solution can be used to construct a solution with compactly supported initial data that breaks down in finite time. In fact, (1.7) admits many self-similar solutions [1] and a particular one exhibiting blow-up was given in closed form in [15] and is known as *ground state* or *fundamental self-similar solution*. In [3], the blow-up of the ground state is shown to be stable indicating that blow-up is a generic situation, see also [4], [5] and [6]. It is clear that self-similar solutions play a crucial role in the blow-up theory for equation (1.7); they correspond to self-similar data at the time of blow-up. Of course these solutions leave the standard scaling critical Sobolev space $\dot{H}^{\frac{3}{2}}$ at the blow-up time. In [7], the author introduces a Besov space-based framework which includes the blowing up solutions of Shatah [11] and Bizoń [1]; in particular two types of solutions of (1.7) are constructed: on the one hand existence, uniqueness and scattering of solutions starting from data which possess infinite critical Sobolev norm, but are small in the sense of suitable Besov spaces is proven (i.e. in these Besov spaces blow-up does not occur), on the other hand large (in the sense of Besov spaces) data can be considered only in the strictly self-similar case where existence still holds but uniqueness is lost.

Moving from this background literature, where essentially global well-posedness for equation (1.7) is known only for finite critical norm and self-similar blow-up seems a generic situation, we present a completely new result on global existence of smooth solutions of (1.7) with infinite critical Sobolev norm $\dot{H}^{\frac{3}{2}} \times \dot{H}^{\frac{1}{2}}$. Moreover, we can construct such solutions so that they satisfy also the size condition $\|u(0, \cdot)\|_{L^\infty(r \geq 1)} > M$ for arbitrarily chosen $M > 0$. The building block for our construction are self-similar solutions which are smooth away from the light cone but singular across it. Our method, strongly inspired by [9], first consists in regularizing the self-similar solutions near the light cone, thus obtaining approximate self-similar solutions and then proceeds by solving a perturbative problem around the approximate self-similar solutions so to generate global solutions. In fact, the regularization destroys the scaling invariance and this turns out to be important for the ensuing perturbative argument. The smooth data thus constructed have infinite critical norm

$$\|u[0]\|_{\dot{H}^{\frac{3}{2}}(\mathbb{R}^3) \times \dot{H}^{\frac{1}{2}}(\mathbb{R}^3)} = \infty$$

only because of insufficient decay at infinity and not because of some singular behavior in finite space-time.

More precisely the main theorem of this paper is the following.

Theorem 2.1. *The equation (1.7) for corotational equivariant wave maps from \mathbb{R}^{3+1} to \mathbb{S}^3 admits smooth data $(f, g) \in C^\infty \times C^\infty$ decaying at infinity to zero, satisfying*

$$\|(f, g)\|_{\dot{H}^{\frac{3}{2}}(\mathbb{R}^3) \times \dot{H}^{\frac{1}{2}}(\mathbb{R}^3)} = \infty \quad \text{but} \quad \|(f, g)\|_{\dot{H}^s(\mathbb{R}^3) \times \dot{H}^{s-1}(\mathbb{R}^3)} < \infty$$

for any $s > \frac{3}{2}$ and such that the corresponding evolution of (1.7) exists globally in forward time as a C^∞ -smooth solution. In fact the initial data can be chosen such

that

$$\|f\|_{L^\infty(r \geq 1)} > M,$$

for arbitrary $M > 0$. The solutions thus obtained are stable with respect to a certain class of perturbations.

Moreover, the solutions of Theorems 2.1 exhibit a precise asymptotic description. Let us remark that the possibility of achieving $\|f\|_{L^\infty(r \geq 1)} > M$ leads to solutions which are not even small in Besov spaces (such as $\dot{B}_{2,\infty}^{\frac{3}{2}} \times \dot{B}_{2,\infty}^{\frac{1}{2}}$) as it is instead the case for the global solutions constructed in [7].

The paper is organized as follows. In Section 3 we construct smooth self-similar solutions of the form

$$Q_0(t, r) = Q\left(\frac{r}{t}\right)$$

for $r > t$ or $r < t$ by a reduction to a nonlinear Sturm-Liouville problem, see (3.1). We solve this ODE by use of fixed point arguments using smallness in L^∞ . As we will see, starting with small data at $a = 0$, Q_0 exhibits a singularity of the form $|a - 1| \log |a - 1|$ near $a = 1$ which precisely fails logarithmically to belong to the scaling critical Sobolev space $\dot{H}^{\frac{3}{2}}(\mathbb{R}^3)$. In the second part of Section 3 we glue together the two solutions residing inside and outside the light-cone, respectively, at $r = t$ to form a continuous function $Q_0(r, t)$ which decays at $a = \infty$ as a^{-1} (and thus fails to lie in $\dot{H}^{\frac{3}{2}}$ at ∞).

Section 4 shows that one of the parameters determining the self-similar solution of Section 3 near the singularity at $a = 1$ can be chosen arbitrarily large, leading to rapid growth and oscillation of the solution on the set $a > 1$ near $a = 1$; the nice behavior of non-linear terms allows to extend these solutions all the way to $a = \infty$, where they again decay asymptotically like a^{-1} . Thanks to this extension we can achieve the large solutions announced in Theorem 2.1.

Section 5 is devoted to the regularization of Q_0 near the light cone: we multiply the singular components of Q_0 by a smooth cut-off function localized at $a = 1$ which leaves untouched the solution far from the light cone. The smooth function thus obtained is no longer an exact solution of (1.7), but in the next section we show that it can be perturbed in a smooth way to obtain exact solutions. The concluding argument is the content of Section 6: it relies on the energy supercritical nature of the problem and imitates the ideas already present in [9]. The methods of Sections 5 and 6 apply to the case of small self-similar solutions constructed in Section 3 as well as to large-size self-similar solutions as constructed in Section 4.

3. SELF-SIMILAR SOLUTIONS

A *self-similar* solution of the Cauchy problem for a wave map u is a solution that depends only on the ratio r/t . Such solutions are thus constant along rays emanating from the origin in space-time, and consequently experience a gradient singularity at the origin (if non-trivial). The existence of non-trivial self-similar solutions was proven in the case of $M = \mathbb{R}^{3+1}$ (which we are concerned with) first by Shatah in [11] where $N = \mathbb{S}^3$, then it was extended to more general rotationally

symmetric, non-convex targets by Shatah and Tahvildar-Zadeh in [12] and also extended to higher dimensions for M by the same authors in a joined work with Cazenave [2]; a particular self-similar solution exhibiting blow-up was explicitly given in [15],

In the corotational setting, equation (1.7) admits a self-similar solution

$$u(t, r) = Q\left(\frac{r}{t}\right)$$

if Q satisfies the following ordinary differential equation in $a := r/t$

$$(1 - a^2)Q''(a) + \left(\frac{2}{a} - 2a\right)Q'(a) - \frac{1}{a^2}f(Q(a)) = 0, \quad (3.1)$$

with f defined as above to be $f(Q) = 2 \sin Q \cos Q$. Indeed we can compute

$$\begin{aligned} u_t &= -\frac{r}{t^2}Q'\left(\frac{r}{t}\right), & u_{tt} &= \frac{2r}{t^3}Q'\left(\frac{r}{t}\right) + \frac{r^2}{t^4}Q''\left(\frac{r}{t}\right) \\ u_r &= \frac{1}{t}Q'\left(\frac{r}{t}\right), & u_{rr} &= \frac{1}{t^2}Q''\left(\frac{r}{t}\right) \end{aligned}$$

whence equation (3.1). The natural initial conditions at $a = 0$ are

$$Q(0) = 0, \quad Q'(0) = d_0. \quad (3.2)$$

We immediately observe that the possible singularities for a solution of (3.1) on $[0, 1]$ can occur only at $a = 0$ and $a = 1$. In the following we will show existence of exact solutions to (3.1) by carefully analysing the behavior near $a = 0$ and $a = 1$. The strategies involved here are inspired by the analogous construction in [9].

Lemma 3.1. *There exists $\varepsilon > 0$ small enough such that, for any $0 \leq d_0 \leq \varepsilon$, the equation (3.1) admits a unique smooth solution on $[0, 1/2]$ with initial conditions (3.2). Furthermore*

$$\begin{aligned} Q(1/2) &= d_0\varphi_0(1/2) + O(d_0^3) \\ Q'(1/2) &= d_0\varphi'_0(1/2) + O(d_0^3), \end{aligned} \quad (3.3)$$

where

$$\varphi_0(a) = \frac{3}{4} \left[\frac{2}{a} + \frac{1 - a^2}{a^2} \log\left(\frac{1 - a}{1 + a}\right) \right].$$

Proof. We consider the associated linear equation

$$Q''(a) + \frac{2}{a}Q'(a) - \frac{2}{a^2(1 - a^2)}Q(a) = 0, \quad (3.4)$$

with fundamental system

$$\varphi_1(a) = \frac{1 - a^2}{a^2}, \quad \varphi_2(a) = \frac{2}{a} + \frac{1 - a^2}{a^2} \log\left(\frac{1 - a}{1 + a}\right). \quad (3.5)$$

We define the Green function $G(a, b)$ for $0 < b < a < 1$

$$G(a, b) := \frac{\varphi_1(a)\varphi_2(b) - \varphi_1(b)\varphi_2(a)}{W(b)} \quad (3.6)$$

where the Wronskian W is defined as

$$W(b) := \varphi_1(b)\varphi_2'(b) - \varphi_1'(b)\varphi_2(b). \quad (3.7)$$

By explicit computation we obtain

$$\begin{aligned} W(b) &= \frac{4}{b^2}, \\ G(a, b) &= \frac{1-a^2}{4a^2} \left(2b + (1-b^2) \log \left(\frac{1-b}{1+b} \right) \right) - \frac{1-b^2}{4} \left(\frac{2}{a} + \frac{1-a^2}{a^2} \log \left(\frac{1-a}{1+a} \right) \right). \end{aligned} \quad (3.8)$$

If we consider the inhomogeneous equation

$$\psi''(a) + \frac{2}{a}\psi'(a) - \frac{2}{a^2(1-a^2)}\psi(a) = H(a) \quad (3.9)$$

with initial conditions

$$\psi(0) = 0, \quad \psi'(0) = 0, \quad (3.10)$$

then a solution in integral form can be obtained via the Green function as

$$\psi(a) = \int_0^a G(a, b)H(b)db.$$

In view of this, we can make an ansatz for the self-similar solution of the non-linear equation (3.1) by choosing H as a non-linear function of Q itself, i.e. by thinking of H as the difference between the non-linear equation (3.1) and the linearized one (3.4): more precisely we define H as the following function of $Q(a)$ and a

$$H(Q(a)) := \frac{\sin(2Q(a)) - 2Q(a)}{a^2(1-a^2)} \quad (3.11)$$

and we seek for a solution Q of (3.1) on $0 \leq a \leq 1/2$ with initial conditions (3.2) of the form

$$Q(a) = \frac{3}{4}d_0\varphi_2(a) + \int_0^a G(a, b)H(Q(b))db. \quad (3.12)$$

We remark that φ_2 is analytic on $(-1, 1)$ with expansion

$$\varphi_2(a) = \frac{4}{3}a + \frac{4}{15}a^3 + O(a^5),$$

whence $\varphi(0) = 0$ and $\varphi'(0) = 4/3$. For the sake of brevity we introduce $\varphi_0(a) := \frac{3}{4}\varphi_2(a)$ so that we can simplify the ansatz (3.12) to

$$Q(a) = d_0\varphi_0(a) + \int_0^a G(a, b)H(Q(b))db, \quad (3.13)$$

so that the smallness parameter d_0 is highlighted. In order to solve (3.1) on $[0, 1/2]$ we set up a contraction argument by use of the ansatz (3.13). We define the map T as

$$(Tf)(a) := d_0\varphi_0(a) + \int_0^a G(a, b)H(f(b))db, \quad (3.14)$$

where f belongs to some suitable functional normed space X_{d_0} to be chosen such that $T : X_{d_0} \rightarrow X_{d_0}$ is a contraction. According to our ansatz, we assume that $0 \leq d_0 \leq \varepsilon$ and we introduce the space X_{d_0}

$$X_{d_0} := d_0\varphi_0 + \left\{ h(a) \mid h \in C^2([0, 1/2]), \|h\|_{C^2[0,1/2]} \leq d_0^2, |h(a)| \leq d_0^2 a^2 \right\}.$$

The set $\left\{ h(a) \mid h \in C^2([0, 1/2]), \|h\|_{C^2[0,1/2]} \leq d_0^2, |h(a)| \leq d_0^2 a^2 \right\}$ defines a closed convex subset of a linear space which we equip with the following norm

$$\|h\|_{X_{d_0}} := \|h\|_{C^2[0,1/2]} + \sup_{0 < a < \frac{1}{2}} \frac{|h(a)|}{a^2}.$$

Our claim can be precisely formulated: *there exists $\varepsilon > 0$ small enough such that for any $0 \leq d_0 \leq \varepsilon$ the equation (3.1) has a unique solution in X_{d_0} .*

We now proceed to the proof of the claim. As a first step we show that T maps X_{d_0} into itself. First of all, by (3.8) we have the following expansion for $0 < b < a$

$$G(a, b) = \frac{1-a^2}{4a^2} \left(\frac{4}{3}b^3 + O(b^5) \right) - \frac{1-b^2}{4} \left(\frac{2}{a} + \frac{1-a^2}{a^2} \log \left(\frac{1-a}{1+a} \right) \right).$$

Hence, if we consider $\frac{G(a,b)}{b^2(1-b^2)}$, we get for $0 < b < a$

$$\frac{G(a, b)}{b^2(1-b^2)} = \frac{1-a^2}{4a^2} \left(\frac{4}{3}b + O(b^3) \right) - \frac{1}{4b^2} \left(\frac{2}{a} + \frac{1-a^2}{a^2} \log \left(\frac{1-a}{1+a} \right) \right).$$

While the first term in $\frac{G(a,b)}{b^2(1-b^2)}$ is clearly analytic for $0 < b < a$, the second term exhibits a singularity of the type $-1/b^2$ near $b = 0$: nonetheless when dealing with T we can get rid of this singularity by integrating it against $[\sin(2f(b)) - 2f(b)]$ for $f \in X_{d_0}$.

Let us now consider a function $f \in X_{d_0}$ which we can write as $f(b) = d_0\varphi_0(b) + h_f(b)$: since f satisfies by definition the initial condition $f(0) = 0$ we are allowed to expand $H(f(b))(b^2(1-b^2))$ in a neighborhood of $b = 0 (= f(0))$ in the following standard way

$$\begin{aligned} H(f(b))(b^2(1-b^2)) &= \\ \sin(2f(b)) - 2f(b) &= -\frac{4}{3}f(b)^3 + \frac{4}{15}f(b)^5 + O(f(b)^7) \end{aligned} \quad (3.15)$$

Moreover, for $0 \leq b \leq 1/2$, any $f \in X_{d_0}$ satisfies $|f(b)| \leq cd_0b + d_0^2b^2 \leq Md_0b$ for an absolute constant $M > 0$ and if d_0 is small enough. This bound on $f(b)$ together with (3.15) gives the key estimate for killing the singularity in the second term of $\frac{G(a,b)}{b^2(1-b^2)}$. Indeed, thanks to all previous considerations, if $f \in X_{d_0}$ we can bound

$$h(a) := \int_0^a G(a, b)H(f(b))db = \int_0^a \frac{G(a, b)}{b^2(1-b^2)} [\sin(2f(b)) - 2f(b)]db$$

for $0 \leq a \leq 1/2$ as follows

$$\begin{aligned} |h(a)| &\leq M^3 d_0^3 a^2 \ll d_0^2 a^2, \\ |h'(a)| &\leq M^3 d_0^3 a \ll d_0^2 a, \\ |h''(a)| &\leq M^3 d_0^3 \ll d_0^2, \end{aligned} \quad (3.16)$$

provided d_0 is small enough, i.e. provided $\varepsilon > 0$ is chosen sufficiently small. This proves that $T : X_{d_0} \rightarrow X_{d_0}$.

To prove the claim it remains to show that T is indeed a contraction on X_{d_0} with respect to its norm. We consider $f, g \in X_{d_0}$ and using similar arguments as above we can estimate

$$\|Tf - Tg\|_{X_{d_0}} \leq C d_0^2 \|f - g\|_{X_{d_0}}$$

which implies that T is a contraction for d_0 sufficiently small, whence the claim.

In order to conclude for the regularity of the solution we can argue similarly as Krieger and Schlag did in [9, Lemma 2.1]. \square

The next step consists in solving equation (3.1) backwards starting from $a = 1$

Lemma 3.2. *There exists $\varepsilon > 0$ small enough such that, for any $d_1, d_2 \in (-\varepsilon, \varepsilon)$, the equation (3.1) admits a unique solution on $[1/2, 1)$ of the form*

$$Q(a) = d_1 \varphi_1(a) + d_2 \varphi_2(a) + (d_3 - d_2) \frac{2}{a} + Q_1(a) \quad (3.17)$$

where d_3 is given by $\sin(4d_3) = 4d_2$ and with

$$\begin{aligned} \varphi_1(a) &= \frac{1 - a^2}{a^2} = O(1 - a), \\ \varphi_2(a) &= \frac{2}{a} + \frac{1 - a^2}{a^2} \log\left(\frac{1 - a}{1 + a}\right) = 2(1 + O(1 - a)) + 2O((1 - a) \log(1 - a)), \\ Q_1(a) &= (|d_1|^3 + |d_2|^3) O((1 - a)^2 \log^2(1 - a)) \end{aligned} \quad (3.18)$$

where the expansions hold for $a \in [1/2, 1]$. Moreover for $a = 1/2$ we have

$$\begin{aligned} Q(1/2) &= d_1 \varphi_1(1/2) + d_2 \varphi_2(1/2) + 4(d_3 - d_2) + O(|d_1|^3 + |d_2|^3) \\ Q'(1/2) &= d_1 \varphi_1'(1/2) + d_2 \varphi_2'(1/2) - 8(d_3 - d_2) + O(|d_1|^3 + |d_2|^3). \end{aligned} \quad (3.19)$$

Proof. Similarly as we have done in the proof of Lemma 3.1, we argue here by contraction arguments. The ansatz for Q solution of (3.1) is given in integral form as

$$Q(a) = Q_0(a) + (d_3 - d_2) \frac{2}{a} - \int_a^1 G(a, b) \frac{\sin(2Q(b)) - 2Q_0(b) - 2Q_1(b)}{b^2(1 - b^2)} db \quad (3.20)$$

with Q_0 defined as

$$Q_0(a) := d_1 \varphi_1(a) + d_2 \varphi_2(a)$$

and where G is the Green function defined in (3.8). Assuming Q has the form (3.17), then the integral equation (3.20) can be viewed as an equation for Q_1 that can be solved by contraction. The asymptotics (3.18) are a consequence of the

behavior of $\frac{G(a,b)}{b^2(1-b^2)}$ for $1/2 \leq a < 1$ and $a < b < 1$. Indeed we can write $\frac{G(a,b)}{b^2(1-b^2)}$ as the sum of the following two terms

$$\begin{aligned} G_1(a,b) &:= \frac{1-a^2}{4a^2} \left(\frac{2}{b(1-b^2)} + \frac{1}{b^2} \log \left(\frac{1-b}{1+b} \right) \right) \\ G_2(a,b) &:= -\frac{1}{4b^2} \left(\frac{2}{a} + \frac{1-a^2}{a^2} \log \left(\frac{1-a}{1+a} \right) \right) \end{aligned} \quad (3.21)$$

so we clearly see that for $1/2 \leq a < 1$ and $a < b < 1$

$$\begin{aligned} G_1(a,b) &\approx (1-a) \left(\frac{1}{(1-b)} + \log(1-b) \right) \\ G_2(a,b) &\approx -(2 + 2(1-a) \log(1-a)). \end{aligned} \quad (3.22)$$

Now, integrating these two terms against $[\sin(2Q(b)) - 2Q_0(b) - 2Q_1(b)]$ which is of order $O((1-b) \log(1-b))$ and using the fact that $\sin(4d_3) = 4d_2$ allows to conclude that the integral in the ansatz (3.20) decays like $((1-a)^2 \log^2(1-a))$. \square

We remark that on the interval $(1/2, 1)$ we obtained a 2-parameter family (d_3 indeed is determined by d_2) of solutions: this allows us to solve the non-linear connection problem at $a = 1/2$, i.e. to “glue” smoothly at $a = 1/2$ the solutions of Lemma 3.1 and those of Lemma 3.2.

Corollary 3.3. *Given any d_0 small enough, the ordinary differential equation (3.1) admits a unique C^2 self-similar solution $Q(a)$ on the interval $[0, 1)$ with initial conditions (3.2). In a left neighborhood of $a = 1$ this solution Q has the form (3.17), i.e. it behaves as follows*

$$Q(a) = d_1 O(1-a) + 2d_2 O((1-a) \log(1-a)) + d_3 + Q_1(a) \quad (3.23)$$

with Q_1 as in (3.18) ($Q_1(1) = 0$).

Proof. In order to prove this corollary of Lemmas 3.1 and 3.2 we simply apply the inverse function theorem. More precisely, given any d_0 small enough, by Lemma 3.2 we can associate to it the unique solution Q . Now, we aim at finding d_1 and d_2 small such that (3.19) match the values of Q and Q' at $a = 1/2$, i.e. (3.3). Since the Jacobian determinant of (3.19) viewed as functions of d_1, d_2 in exactly the Wronskian of φ_1, φ_2 , we easily see that its value at $(d_1, d_2) = 0$ is 1. It being non-zero, we can invoke the inverse function theorem and find the desired d_1, d_2 small. \square

Let us finally note that the obtained solution just fails logarithmically to be in $\dot{H}^{\frac{3}{2}}$.

The next step consists in solving the self-similar ODE (3.1) in the exterior light cone, i.e. for $a = r/t > 1$. Similarly as for the interior light cone, we first solve the problem on the two intervals $(1, 2]$ and $[2, \infty)$ and then we will glue the two solutions at the connection point $a = 2$.

Lemma 3.4. *There exists $\varepsilon > 0$ small enough such that, for any $\tilde{d}_1, \tilde{d}_2 \in (-\varepsilon, \varepsilon)$, the equation (3.1) admits a unique solution on $(1, 2]$ of the form*

$$Q(a) = \tilde{d}_1 \tilde{\varphi}_1(a) + \tilde{d}_2 \tilde{\varphi}_2(a) - (\tilde{d}_3 - \tilde{d}_2) \frac{2}{a} + \tilde{Q}_1(a) \quad (3.24)$$

where \tilde{d}_3 is given by $\sin(4\tilde{d}_3) = 4\tilde{d}_2$ and with

$$\begin{aligned} \tilde{\varphi}_1(a) &= \frac{a^2 - 1}{a^2} = O(a - 1), \\ \tilde{\varphi}_2(a) &= -\frac{2}{a} + \frac{a^2 - 1}{a^2} \log\left(\frac{a - 1}{a + 1}\right) = 2(-1 + O(a - 1)) + 2O((a - 1) \log(a - 1)), \\ \tilde{Q}_1(a) &= (|\tilde{d}_1|^3 + |\tilde{d}_2|^3) O((a - 1)^2 \log^2(a - 1)) \end{aligned} \quad (3.25)$$

where the expansions hold for $a \in (1, 2]$. Moreover for $a = 2$ we have

$$\begin{aligned} Q(2) &= \tilde{d}_1 \tilde{\varphi}_1(2) + \tilde{d}_2 \tilde{\varphi}_2(2) - 4(\tilde{d}_3 - \tilde{d}_2) + O(|\tilde{d}_1|^3 + |\tilde{d}_2|^3) \\ Q'(2) &= \tilde{d}_1 \tilde{\varphi}_1'(2) + \tilde{d}_2 \tilde{\varphi}_2'(2) + 8(\tilde{d}_3 - \tilde{d}_2) + O(|\tilde{d}_1|^3 + |\tilde{d}_2|^3). \end{aligned} \quad (3.26)$$

We omit the proof of Lemma 3.4 since it can be carried out exactly along the line of the proof of Lemma 2.3.2

Lemma 3.5. *There exists $\varepsilon > 0$ small enough such that, for any $q_1, q_2 \in (-\varepsilon, \varepsilon)$, the equation (3.1) admits a unique solution on $[2, \infty)$ which, for $a \rightarrow \infty$, has the form*

$$Q(a) = q_1 \tilde{\varphi}_1(a) + q_2 \tilde{\varphi}_2(a) + O\left(\frac{1}{a^2}\right) \quad (3.27)$$

with

$$\begin{aligned} \tilde{\varphi}_1(a) &= \frac{a^2 - 1}{a^2}, \\ \tilde{\varphi}_2(a) &= -\frac{2}{a} + \frac{a^2 - 1}{a^2} \log\left(\frac{a - 1}{a + 1}\right) = -\frac{4}{a} + O\left(\frac{1}{a^2}\right), \end{aligned} \quad (3.28)$$

where the last expansion holds for $a \rightarrow \infty$. Moreover for $a = 2$ we have

$$\begin{aligned} Q(2) &= q_1 \tilde{\varphi}_1(2) + q_2 \tilde{\varphi}_2(2) + O(|q_1|^3 + |q_2|^3) \\ Q'(2) &= q_1 \tilde{\varphi}_1'(2) + q_2 \tilde{\varphi}_2'(2) + O(|q_1|^3 + |q_2|^3). \end{aligned} \quad (3.29)$$

Proof. As in the proof of Lemma 3.1 we make use of the Green function, now defined in terms of $\tilde{\varphi}_1$ and $\tilde{\varphi}_2$, i.e.

$$\tilde{G}(a, b) := \frac{\tilde{\varphi}_1(a)\tilde{\varphi}_2(b) - \tilde{\varphi}_1(b)\tilde{\varphi}_2(a)}{\tilde{W}(b)}, \quad (3.30)$$

where the Wronskian \tilde{W} is now defined as

$$\tilde{W}(b) := \tilde{\varphi}_1(b)\tilde{\varphi}_2'(b) - \tilde{\varphi}_1'(b)\tilde{\varphi}_2(b). \quad (3.31)$$

Similarly to the computation in the proof of Lemma 3.1 we obtain

$$\begin{aligned}\widetilde{W}(b) &= \frac{4}{b^2}, \\ \widetilde{G}(a, b) &= \frac{a^2 - 1}{4a^2} \left(-2b + (b^2 - 1) \log \left(\frac{b-1}{b+1} \right) \right) - \frac{b^2 - 1}{4} \left(-\frac{2}{a} + \frac{a^2 - 1}{a^2} \log \left(\frac{a-1}{a+1} \right) \right).\end{aligned}\tag{3.32}$$

In this case, the natural ansatz for Q is as follows

$$Q(a) = q_1 \widetilde{\varphi}_1(a) + q_2 \widetilde{\varphi}_2(a) + \int_a^\infty \widetilde{G}(a, b) \frac{\sin(2Q(b)) - 2Q(b)}{b^2(1-b^2)} db.\tag{3.33}$$

In order to solve (3.1) on $[2, \infty)$ we can set up a contraction argument for the ansatz (3.33) in analogy to the proof of Lemma 3.1. Let us note that we have the following

$$\begin{aligned}\frac{\widetilde{G}(a, b)}{b^2(1-b^2)} &= \frac{a^2 - 1}{4a^2} \left(\frac{2}{b(b^2 - 1)} - \frac{1}{b^2} \log \left(\frac{b-1}{b+1} \right) \right) \\ &\quad + \frac{1}{4b^2} \left(-\frac{2}{a} + \frac{a^2 - 1}{a^2} \log \left(\frac{a-1}{a+1} \right) \right)\end{aligned}$$

which to leading order (and up to constants) for $a, b \rightarrow \infty$ exhibits the following decay

$$\frac{\widetilde{G}(a, b)}{b^2(1-b^2)} \approx \frac{1}{b^3} - \frac{1}{ab^2}.$$

Integrating this against $[\sin(2Q(b)) - 2Q(b)]$ for $a < b < \infty$ gives the asymptotics (3.27). The values (3.29) are simply obtained by substituting $a = 2$. \square

In view of Lemmas 3.4 and 3.5 which both generate a two-parameter family of solutions we can glue those at $a = 2$ so to produce a one smooth solution for all $a > 1$. This is done in the following corollary.

Corollary 3.6. *Given any q_1, q_2 small enough, the ordinary differential equation (3.1) admits a unique C^2 self-similar solution $Q(a)$ on the interval $(1, \infty)$ with the asymptotics (3.27)-(3.28). In a right neighborhood of $a = 1$ the solution Q has the form (3.24).*

Proof. The corollary can be proven once more thanks to the inverse function theorem. Given q_1, q_2 small enough, Lemma 3.5 provides a unique solution Q for $a \in [2, \infty)$ with asymptotics (3.27)-(3.28). Now, we aim at finding \widetilde{d}_1 and \widetilde{d}_2 small (\widetilde{d}_3 is given as function of \widetilde{d}_2) such that (3.26) given by Lemma 3.4 match the values of Q and Q' at $a = 2$, i.e. (3.29). Since the Jacobian determinant of (3.26), i.e. the Wronskian, does not vanish and by smallness of q_1 and q_2 , we can apply the inverse function theorem in a neighborhood of $(\widetilde{d}_1, \widetilde{d}_2) = 0$ and solve the connection problem at $a = 2$. \square

Finally, to complete the construction of the self-similar solution we will connect just continuously the solution on $[0, 1)$ provided by Corollary 3.3 with the solution on $(1, \infty)$ provided by Corollary 3.6.

Corollary 3.7. *Given any small d_0 , there exists a unique C^2 solution $Q(a)$ of the ordinary differential equation (3.1) on $[0, 1)$ with initial conditions (3.2). By Corollary 3.6, this solution can be extended non-uniquely to a continuous function on $a \geq 1$ which solves (3.1) on $a > 1$ and behaves as*

$$Q(a) = q_1 - 4q_2a^{-1} + O(a^{-2}) \quad (3.34)$$

for $a \rightarrow \infty$. The global continuous solutions on $a \geq 0$ must satisfy the condition $2d_3 = -2\tilde{d}_3$ (see Lemmas 3.2 and 3.4). Let us denote these solutions on $a > 0$ by $Q_0(a)$, then we have the following global representation

$$Q_0(a) = C_1 \frac{|a^2 - 1|}{a^2} + C_2 \frac{|a^2 - 1|}{a^2} \log\left(\frac{|a - 1|}{a + 1}\right) + C_3 Q_3(a) + Q_4(a) \quad (3.35)$$

for $a \geq 1/2$, where $Q_3 = \frac{1}{a}$ is smooth at $a = 1$ and $Q_4(a) = O(|a - 1|^2 \log^2(|a - 1|))$ consists of higher order terms.

Proof. Given any small d_0 , Corollary 3.3 provides us with d_1, d_2, d_3 and with a C^2 solution Q on $[0, 1)$. On the other hand, thanks to Corollary 3.6 any q_1, q_2 small enough give \tilde{d}_1, \tilde{d}_2 and \tilde{d}_3 corresponding to a C^2 solution on $(1, \infty)$. The freedom of selecting q_1 and q_2 small allows us to make the choice in such a way that \tilde{d}_3 satisfies $-2\tilde{d}_3 = 2d_3$, i.e. such that we can extend continuously Q to $a > 1$. \square

4. LARGE SELF-SIMILAR SOLUTIONS

In this section we show that it is possible to choose a big constant \tilde{d}_1 in Lemma 3.4 thus allowing to construct solutions which are arbitrarily large in the exterior light-cone. Indeed, by Corollary 3.7, the continuity condition at the light cone involves only \tilde{d}_3 (and hence \tilde{d}_2) leaving freedom of choice for \tilde{d}_1 . Hence, the goal here is to prove the existence of a solution to the ODE (3.1) on $a > 1$ for \tilde{d}_1 large. The first step consists in proving the analogue of Lemma 3.4 on a right neighborhood of $a = 1$ for arbitrary \tilde{d}_1 .

Lemma 4.1. *There exists $\varepsilon > 0$ small enough such that, for any $\tilde{d}_2 \in (-\varepsilon, \varepsilon)$ and for any $\tilde{d}_1 \geq 1$ arbitrary, the equation (3.1) admits a unique solution Q on $(1, 1 + \ell]$ with $\ell = c\tilde{d}_1^{-\frac{1}{2}}$ for some absolute constant $c > 0$ small enough. This solution Q has the form*

$$Q(a) = \tilde{d}_1 \tilde{\varphi}_1(a) + \tilde{d}_2 \tilde{\varphi}_2(a) - (\tilde{d}_3 - \tilde{d}_2) \frac{2}{a} + \tilde{Q}_1(a) \quad (4.1)$$

where \tilde{d}_3 is given by $\sin(4\tilde{d}_3) = 4\tilde{d}_2$ (and hence \tilde{d}_3 is also small) and with

$$\begin{aligned} \tilde{\varphi}_1(a) &= \frac{a^2 - 1}{a^2} = O(a - 1), \\ \tilde{\varphi}_2(a) &= -\frac{2}{a} + \frac{a^2 - 1}{a^2} \log\left(\frac{a - 1}{a + 1}\right) = 2(-1 + O(a - 1)) + 2O((a - 1) \log(a - 1)), \\ \tilde{Q}_1(a) &= \tilde{d}_1 O((a - 1)^2 \log^2(a - 1)) \end{aligned} \quad (4.2)$$

where the expansions hold for $a \in (1, 1 + \ell]$. Moreover, there exists $a_* \in (1, 1 + \ell]$ such that

$$|Q(a_*)| \approx \tilde{d}_1^{\frac{1}{2}}. \quad (4.3)$$

As a consequence $Q(a_*)$ can be made arbitrarily large by choosing \tilde{d}_1 sufficiently large.

Proof. The key ideas are inherited from the proof of Lemma 3.2 except the fact that we have to deal here with lack of smallness for \tilde{q}_1 : this is overcome by a bootstrap argument.

The ansatz for Q is given in integral form as

$$Q(a) = \tilde{Q}_0(a) - (\tilde{d}_3 - \tilde{d}_2) \frac{2}{a} + \int_1^a \tilde{G}(a, b) \frac{\sin(2Q(b)) - 2\tilde{Q}_0(b) - 2\tilde{Q}_1(b)}{b^2(1-b^2)} db \quad (4.4)$$

with \tilde{Q}_0 defined as

$$Q_0(a) := \tilde{d}_1 \tilde{\varphi}_1(a) + \tilde{d}_2 \tilde{\varphi}_2(a)$$

and where \tilde{G} is the Green function defined in (3.30). Assuming Q has the form (4.1), the integral equation (4.4) can be solved by contraction for \tilde{Q}_1 . More precisely, for Q as in (4.1), we would like to obtain \tilde{Q}_1 as fixed point of the following equation

$$\tilde{Q}_1(a) = \int_1^a \tilde{G}(a, b) \frac{\sin(2Q(b)) - 2\tilde{Q}_0(b) - 2\tilde{Q}_1(b)}{b^2(1-b^2)} db. \quad (4.5)$$

For the sake of clarity let us define

$$\overline{Q}_1(a) := \frac{\tilde{Q}_1(a)}{(a-1)^2 \log^2(a-1)}$$

In order to run the fixed point argument for equation (4.5) we need to show that the bound

$$|\overline{Q}_1(a)| \leq C |\tilde{d}_1| \quad (4.6)$$

improves upon itself on the interval $a \in [1, 1 + \ell]$, with $\ell = c \tilde{d}_1^{-\frac{1}{2}}$ as given by the statement of the lemma, if inserted in the equation (4.5). More precisely, our goal is to prove that the bound (4.6) entails the following better one

$$|\overline{Q}_1(a)| \leq \frac{C}{2} |\tilde{d}_1|.$$

As a first step, we compute $\frac{\tilde{G}(a,b)}{b^2(1-b^2)}$ and we easily obtain

$$\begin{aligned} \frac{\tilde{G}(a,b)}{b^2(1-b^2)} &= \frac{a^2-1}{4a^2} \left(\frac{2}{b(b^2-1)} - \frac{1}{b^2} \log \left(\frac{b-1}{b+1} \right) \right) \\ &\quad + \frac{1}{4b^2} \left(-\frac{2}{a} + \frac{a^2-1}{a^2} \log \left(\frac{a-1}{a+1} \right) \right). \end{aligned}$$

Similarly to (3.21), we can write $\frac{\tilde{G}(a,b)}{b^2(1-b^2)}$ as the sum of two terms \tilde{G}_1 and \tilde{G}_2 which behave as follows for $1 < a \leq 1 + \ell$ and $1 < b < a$

$$\begin{aligned}\tilde{G}_1(a,b) &\approx (a-1) \left(\frac{1}{(b-1)} + \log(b-1) \right) \\ \tilde{G}_2(a,b) &\approx -(-2 + 2(a-1) \log(a-1)).\end{aligned}\quad (4.7)$$

The choice of l is such that on $[1, 1+\ell]$ the non-linearities involving \tilde{Q}_1 in $[\sin(2Q(b)) - 2\tilde{Q}_0(b) - 2\tilde{Q}_1(b)]$ are dominated by their first order (linear) approximation. We also introduce $\bar{Q}_0(a) := \tilde{Q}_0(a) - (\tilde{d}_3 - \tilde{d}_2)_a^2$ so that $Q(a) = \bar{Q}_0(a) + \tilde{Q}_1(a)$. Now, we can write $\sin(2Q(b)) - 2\tilde{Q}_0(b) - 2\tilde{Q}_1(b)$ as follows

$$\begin{aligned}\sin(2Q) - 2\tilde{Q}_0 - 2\tilde{Q}_1 &= \left(\sin(2\bar{Q}_0) \cos(2\tilde{Q}_1) - 2\tilde{Q}_1 \right) \\ &\quad + \left(\cos(2\bar{Q}_0) \sin(2\tilde{Q}_1) - 2\tilde{Q}_1 \right) \\ &=: A + B.\end{aligned}\quad (4.8)$$

The condition $\sin(4\tilde{d}_3) = 4\tilde{d}_2$ already ensures that $\tilde{Q}_1(1) = 0$. Now, we can expand \sin and \cos in (4.8) and integrate against \tilde{G}_1 and \tilde{G}_2 ,

At first we consider A . With the choice of $\ell = c\tilde{d}_1^{-\frac{1}{2}}$ for some small constant $c > 0$ and under the assumption (4.6) on \tilde{Q}_1 , $\cos(2\tilde{Q}_1)$ can be approximated by 1, being the higher order terms dominated by the first order one; as a consequence A can be further expanded as follows (remember that $\sin(4\tilde{d}_3) = 4\tilde{d}_2$)

$$\begin{aligned}A &\approx \left[\sin(-4\tilde{d}_3) \cos(\beta) + \cos(-4\tilde{d}_3) \sin(\beta) \right] \cos(2\tilde{Q}_1) - \beta + 4\tilde{d}_2 \\ &\approx \left[-4\tilde{d}_2 (1 + (\cos(\beta) - 1)) + (\beta + (\sin(\beta) - \beta)) \right] - \beta + 4\tilde{d}_2 \\ &\approx \left[-4\tilde{d}_2 (\cos(\beta) - 1) + (\sin(\beta) - \beta) \right]\end{aligned}\quad (4.9)$$

where $\beta \approx 4\tilde{d}_1(b-1) + 4\tilde{d}_2(b-1) \log(b-1)$ and where we also approximated $\cos(-4\tilde{d}_3)$ by 1 being \tilde{d}_3 small. If we integrate A against \tilde{G}_1 , in particular against the first and leading term of \tilde{G}_1 i.e. $(a-1)/(b-1)$ we obtain

$$\begin{aligned}&\tilde{d}^{-1}(a-1)^{-2} \left| \int_1^a \tilde{G}_1(b) A db \right| \\ &\lesssim \tilde{d}^{-1} \ell^{-2} (a-1) \int_1^a \frac{1}{b-1} \left| -4\tilde{d}_2 (\cos(\beta) - 1) + (\sin(\beta) - \beta) \right| db \\ &\lesssim c^{-2} (a-1) \int_1^a (4|\tilde{d}_2| + 2) \left| \frac{\beta}{b-1} \right| db \\ &\lesssim (4|\tilde{d}_2| + 2) \left((\tilde{d}_1 + |\tilde{d}_2|) \ell^2 + |\tilde{d}_2| \ell^2 |\log \ell| \right) \ll 1\end{aligned}\quad (4.10)$$

by choice of $\ell = c\tilde{d}_1^{-\frac{1}{2}}$. Similarly we can also estimate the term corresponding to the integral of A against \tilde{G}_2 . As far as the integral involving B is concerned we argue as follows. Using the fact that $\cos(2\bar{Q}_0)$ is bounded and that $\sin(2\tilde{Q}_1)$ behaves

to leading order as $(2\tilde{Q}_1)$ due to the choice of $\ell = c\tilde{d}_1^{-\frac{1}{2}}$ and to the assumption (4.6), we can estimate

$$\begin{aligned}
& \tilde{d}^{-1}(a-1)^{-2} \left| \int_1^a \tilde{G}_1(b) B db \right| \\
& \leq \tilde{d}^{-1} \ell^{-2} \left| \int_1^a \tilde{G}_1(b) (\cos(2\tilde{Q}_0(b)) \sin(2\tilde{Q}_1(b)) - 2\tilde{Q}_1(b)) db \right| \\
& \lesssim c^{-2} \left| \int_1^a \tilde{G}_1(b) [(\cos(2\tilde{Q}_0(b)) - 1) \sin(2\tilde{Q}_1(b)) + (\sin(2\tilde{Q}_1(b)) - 2\tilde{Q}_1(b))] db \right| \\
& \lesssim c\tilde{d}^{-\frac{1}{2}} \ll 1
\end{aligned} \tag{4.11}$$

and an analogous estimate can be obtained when integrating B against \tilde{G}_2 . The large value exhibited in (4.3) is achieved by choosing $a_* = 1 + \ell/2$. \square

Lemma 4.2. *The solutions to (3.1) provided by Lemma 4.1 on $(1, 1 + \ell]$ can be extended to $(1, \infty)$ as a smooth globally bounded solution $Q(a)$, which behaves as*

$$Q(a) = c_1 + c_2 a^{-1} + O(a^{-2}) \tag{4.12}$$

for $a \rightarrow \infty$ and for non-vanishing constants c_1 and c_2 .

Proof. Let us recall that we are dealing with the following ODE ((3.1))

$$Q''(a) + \frac{2}{a} Q'(a) - \frac{(2 \sin Q(a) \cos Q(a))}{a^2(1-a^2)} = 0, \tag{4.13}$$

where the non-linearity $(2 \sin Q(a) \cos Q(a))$ is clearly bounded. Hence, away from the singularities $a = 0$ and $a = 1$, standard elliptic estimates allow to prove L^∞ bounds independent of the existence time so that solutions can be extended for all time. \square

5. EXCISION OF SINGULARITY NEAR LIGHT CONE $a = 1$: APPROXIMATE SOLUTIONS

In the previous section we have shown the existence of self-similar solutions Q_0 to (3.1) which are smooth away from the light cone but are only continuous at $a = 1$ i.e. across the light cone. Our goal is to construct global smooth solutions to (1.7) which have infinite critical norm $\dot{H}^{\frac{3}{2}}$ departing from these self-similar solutions by excision of the singularity near $a = 1$. In order to achieve this, we introduce a smooth cut-off function $\chi(t-r)$ whose support lies at a distance C from the light cone: more precisely $\chi(x) = 1$ for $|x| \geq 2C$ and $\chi(x) = 0$ for $|x| \leq C$. In view of Corollary 3.7, we know that near $a = 1$, the function Q_0 is of the form

$$Q_0(a) = C_1 \frac{|a^2 - 1|}{a^2} + C_2 \frac{|a^2 - 1|}{a^2} \log \left(\frac{|a - 1|}{a + 1} \right) + C_3 Q_3(a) + Q_4(a) \tag{5.1}$$

where $Q_3 = \frac{2}{a}$ is smooth at $a = 1$ and $Q_4(a) = O(|a - 1|^2 \log^2(|a - 1|))$ consists of higher order terms. For the sake of simplicity we introduce

$$R(t, r) = R(a) = C_1 \frac{|a^2 - 1|}{a^2} + C_2 \frac{|a^2 - 1|}{a^2} \log \left(\frac{|a - 1|}{a + 1} \right).$$

Thus, we modify the exact singular solution near the light cone by introducing the following approximate solution

$$u_{approx}(t, r) := \chi(t-r) [R(a) + Q_4(a)] + C_3 Q_3(a) \quad (5.2)$$

Now, we will estimate the errors which arise when computing the expression

$$\partial_{tt} u_{approx} - \partial_{rr} u_{approx} - \frac{2}{r} \partial_r u_{approx} + \frac{f(u_{approx})}{r^2}.$$

Since for our estimates, the term Q_4 is of higher order we will neglect it in the following computations. First of all, we remark when applying derivatives only to χ we get $R(\partial_{tt}\chi - \partial_{rr}\chi)$ which vanishes since $\chi(t-r)$ solves the 1-dimensional wave equation. When only one time derivative falls on $\chi(t-r)$ we get

$$2\chi'(t-r)\partial_t R(t, r) = 2\chi'(t-r)\text{sign}(t-r) \left(2C_1 \frac{t}{r^2} - 2C_2 \frac{t}{r^2} \log\left(\frac{|r-t|}{r+t}\right) - 2C_2 \frac{1}{r} \right).$$

Similarly, when one derivative ∂_r falls on $\chi(t-r)$ we obtain

$$-2\chi'(t-r)\partial_r R(t, r) = 2\chi'(t-r)\text{sign}(t-r) \left(2C_1 \frac{t^2}{r^3} + 2C_2 \frac{t^2}{r^3} \log\left(\frac{|r-t|}{r+t}\right) + 2C_2 \frac{t}{r^2} \right).$$

Finally, we compute the contribution from $-\frac{2}{r}R\partial_r\chi$ which gives

$$-2\chi'(t-r) \left(C_1 \frac{|r^2-t^2|}{r^3} + C_2 \frac{|r^2-t^2|}{r^3} \log\left(\frac{|r-t|}{r+t}\right) \right).$$

Summing up all the terms involving χ' we have the following expression

$$\begin{aligned} & 2C_2\chi'(t-r) \log\left(\frac{|r-t|}{r+t}\right) \left(\frac{t}{r^2} \left| \frac{r}{t} - 1 \right| - \frac{t^2}{r^3} \left| \frac{r}{t} - 1 \right| \right) - 2C_2\chi'(t-r) \left(\frac{t}{r^2} \left| \frac{r}{t} - 1 \right| \right) \\ & + 2C_1\chi'(t-r) \left(\frac{t^2}{r^3} \left| \frac{r^2}{t^2} - 1 \right| - 2\frac{t^2}{r^3} \left| \frac{r}{t} - 1 \right| \right) \end{aligned}$$

which is of size t^{-3} , since χ' has support in the strip $C \leq |t-r| \leq 2C$ where $|r/t-1|$ behaves as $1/t$.

Moreover, the error from the nonlinear term is of the following form

$$\begin{aligned} & \frac{1}{r^2} f \left(\chi(t-r) \left(C_1 \frac{|r^2-t^2|}{r^2} + C_2 \frac{|r^2-t^2|}{r^2} \log\left(\frac{|r-t|}{r+t}\right) \right) + 2C_3 \text{sign}(t-r) \frac{t}{r} \right) \\ & - \frac{\chi(t-r)}{r^2} f \left(C_1 \frac{|r^2-t^2|}{r^2} + C_2 \frac{|r^2-t^2|}{r^2} \log\left(\frac{|r-t|}{r+t}\right) + 2C_3 \text{sign}(t-r) \frac{t}{r} \right) \\ & = O(t^{-3}) \end{aligned}$$

This implies that

$$\square u_{approx} + \frac{f(u_{approx})}{r^2} \in L^2(\mathbb{R}^3),$$

in light of the support properties of this expression, and is of order t^{-2} at fixed time t in this norm. Thus all the errors beat the scaling.

6. FROM AN APPROXIMATE SOLUTION TO AN EXACT SOLUTION

Here we construct exact solutions via the ansatz

$$u(t, r) = u_{approx}(t, r) + \epsilon(t, r).$$

We recall that we are considering the case of target \mathbb{S}^3 , in which case we have $g(u) = \sin u$. Then we obtain the following wave equation which is in fact on \mathbb{R}^{5+1} :

$$\begin{aligned} \left(\frac{\epsilon}{r}\right)_{tt} - \left(\frac{\epsilon}{r}\right)_{rr} - \frac{4}{r} \left(\frac{\epsilon}{r}\right)_r &= -\frac{1}{r} \frac{\sin(2\epsilon) - 2\epsilon}{r^2} \cos(2u_{approx}) \\ &\quad - \frac{\sin(2u_{approx})}{r^3} (\cos(2\epsilon) - 1) \\ &\quad - \frac{2\epsilon}{r^3} [\cos(2u_{approx}) - 1] \\ &\quad + \frac{e_0}{r}. \end{aligned} \tag{6.1}$$

Note that by introducing the new variable $v = \frac{\epsilon}{r}$, we get ‘‘essentially’’ the new wave equation

$$v_{tt} - \Delta_{\mathbb{R}^5} v = v^3 + \frac{vu_{approx}^2}{r^2} + \frac{v^2 u_{approx}}{r} + \frac{1}{r} e_0.$$

For the purely cubic term, we get the scaling $v(t, r) \rightarrow \lambda v(\lambda t, \lambda r)$, which on \mathbb{R}^{5+1} corresponds to $s_c = \frac{3}{2}$, as expected. Thus it is natural to try to run an iteration in the space $H_{\mathbb{R}^5}^{\frac{3}{2}}$. It is then natural to work with the Strichartz norms $L_t^\infty L_x^5 \cap L_t^2 L_x^{10} \cap L_t^2 (\nabla_x^{-\frac{1}{2}} L_x^5)$, with the same scaling as $L_t^\infty \dot{H}_{\mathbb{R}^5}^{\frac{3}{2}}$. Then the interaction terms $\frac{vu_{approx}^2}{r^2}$, $\frac{v^2 u_{approx}}{r}$, appear critical, since the first can be reduced to

$$\frac{vu_{approx}}{rt},$$

which fails logarithmically to belong to $L_t^1 \dot{H}^{\frac{1}{2}}$. Thus, as in the paper [9], we shall also be taking advantage of the Hamiltonian structure to handle this low frequency issue. We state

Proposition 6.1. *Let $C \geq 1$ a given constant, $T \geq 1$ sufficiently large, depending on \tilde{d}_1 in the approximate solution. Assume that the C^∞ -smooth data $v[T] = (v(T, \cdot), v_t(T, \cdot))$ are radial and supported in the annulus $r \in [T - C, T + C]$. Also, assume that \tilde{d}_2 for the approximate solution is sufficiently small (less than an absolute constant), and that for a $\delta_1 > 0 = \delta_1(C)$ sufficiently small, we have*

$$\|v[T]\|_{\dot{H}^{\frac{3}{2}}(\mathbb{R}^5) \cap \dot{H}^1(\mathbb{R}^5) \times \dot{H}^{\frac{1}{2}}(\mathbb{R}^5) \cap L^2(\mathbb{R}^5)} \leq \delta_1.$$

Then the problem (6.1) with initial data $v[T]$ at time $t = T$ admits a global-in-forward time solution $v(t, \cdot)$ of class C^∞ .

The existence of the solution v will follow from a standard local existence result as well as a more sophisticated bootstrap argument which is at the heart of the matter, as in [9]. The local existence result is as follows:

Proposition 6.2. *Given data $v[T]$ with the same support properties as above, satisfying*

$$\|v[T]\|_{\dot{H}^{\frac{3}{2}}(\mathbb{R}^5) \times \dot{H}^{\frac{1}{2}}(\mathbb{R}^5)} \ll 1,$$

then there exists a time $T_1 > T$ and a solution of (6.1) of class

$$v \in L_t^\infty \dot{H}^{\frac{3}{2}}([T, T_1] \times \mathbb{R}^5), v_t \in L_t^\infty \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)$$

with compact support on every time slice $t \in [T, T_1]$. If $v[T] \in \dot{H}^s \cap \dot{H}^{s-1}$, $s > \frac{3}{2}$, then so is the solution at all times $t \in [T, T_1]$.

The proof is standard and we refer to [9, Section 7] for a similar argument.

Before stating the key bootstrap proposition, we recall the following set of standard Strichartz estimates, in the \mathbb{R}^{5+1} -setting:

Lemma 6.3. *Let $\square_{5+1}u = 0$. Then for $\frac{1}{p} + \frac{2}{q} \leq 1$, $p \geq 2$, we have $(u[0] = (u(0, \cdot), u_t(0, \cdot)))$*

$$\|(-\Delta)^{\frac{\gamma}{2}}u\|_{L_t^p L_x^q} \leq C \|u[0]\|_{\dot{H}^{\frac{3}{2}} \times \dot{H}^{\frac{1}{2}}}, \gamma = -1 + \frac{1}{p} + \frac{5}{q}.$$

Next, the bootstrap proposition:

Proposition 6.4. *Let us assume all hypotheses of Proposition 6.1. Then there exists $C_1 > 1$ with $C_1\delta_1 \ll 1$, as well as a constant $\gamma = \gamma(\tilde{d}_{1,2}, \delta_1, T) > 0$, such that for any $T_1 > T$, the following conclusion holds:*

$$\|v\|_{L_t^2(L_x^{10} \cap \nabla_x^{-\frac{1}{2}} L_x^5)([T, T_1] \times \mathbb{R}^5)} + \sup_{t \in [T, T_1]} \|v[t, \cdot]\|_{\dot{H}^{\frac{3}{2}} \cap (\frac{1}{T})^\gamma \dot{H}^1(\mathbb{R}^5) \times \dot{H}^{\frac{1}{2}} \cap (\frac{1}{T})^\gamma L^2(\mathbb{R}^5)} \leq C_1 \delta_1$$

implies

$$\|v\|_{L_t^2(L_x^{10} \cap \nabla_x^{-\frac{1}{2}} L_x^5)([T, T_1] \times \mathbb{R}^5)} + \sup_{t \in [T, T_1]} \|v[t, \cdot]\|_{\dot{H}^{\frac{3}{2}} \cap (\frac{1}{T})^\gamma \dot{H}^1(\mathbb{R}^5) \times \dot{H}^{\frac{1}{2}} \cap (\frac{1}{T})^\gamma L^2(\mathbb{R}^5)} \leq \frac{C_1}{2} \delta_1.$$

Proof. We follow closely the procedure in [9]. We commence with the energy type norm, i. e.

$$\sup_{t \in [T, T_1]} \|v[t, \cdot]\|_{(\frac{1}{T})^\gamma \dot{H}^1(\mathbb{R}^5) \times (\frac{1}{T})^\gamma L^2(\mathbb{R}^5)}.$$

Multiplying (6.1) by v_t and integrating in space-time, we get

$$\begin{aligned} \int_{\mathbb{R}^5} [v_t^2 + |\nabla_x v|^2] dx|_{t=T_1} &= \int_{\mathbb{R}^5} [v_t^2 + |\nabla_x v|^2] dx|_{t=T} \\ &- \int_T^{T_1} \int_{\mathbb{R}^5} \frac{1}{r} \frac{\sin(2vr) - 2vr}{r^2} \cos(2u_{approx}) v_t dxdt \\ &- \int_T^{T_1} \int_{\mathbb{R}^5} \frac{\sin(2u_{approx})}{r^3} (\cos(2vr) - 1) v_t dxdt \quad (6.2) \\ &- \int_T^{T_1} \int_{\mathbb{R}^5} \frac{2v}{r^2} [\cos(2u_{approx}) - 1] v_t dxdt \\ &+ \int_T^{T_1} \int_{\mathbb{R}^5} \frac{e_0}{r} v_t dxdt. \end{aligned}$$

We treat the terms on the right via integration by parts. For the second term on the right, introducing

$$E(x) := \int_0^x (\sin y - y) dy, \quad |E(x)| \lesssim |x|^4,$$

we get

$$\begin{aligned} & \int_T^{T_1} \int_{\mathbb{R}^5} \frac{1}{r} \frac{\sin(2vr) - 2vr}{r^2} \cos(2u_{approx}) v_t dx dt \\ &= \int_{\mathbb{R}^5} r^{-4} E(vr) \cos(2u_{approx}) dx \Big|_T^{T_1} - \int_T^{T_1} \int_{\mathbb{R}^5} 2r^{-4} E(vr) \sin(2u_{approx}) u_{approx,t} dx dt. \end{aligned}$$

To bound these terms we use interpolation between $L_x^{\frac{10}{3}}$ and L_x^5 . With

$$\frac{1}{4} = \alpha \cdot \frac{3}{10} + (1 - \alpha) \cdot \frac{1}{5} = \frac{1}{5} + \frac{\alpha}{10},$$

we get $\alpha = \frac{1}{2}$, whence

$$\begin{aligned} \left| \int_{\mathbb{R}^5} r^{-4} E(vr) \cos(2u_{approx}) dx \Big|_T^{T_1} \right| &\lesssim \int_{\mathbb{R}^5} v^4 dx \lesssim \sum_{t=T, T_1} \|\nabla_x v(t, \cdot)\|_{L_x^2}^2 \left(\int_{\mathbb{R}^5} |v(t, \cdot)|^5 dx \right)^{\frac{2}{5}} \\ &\ll \sup_{t \in [T, T_1]} \|\nabla_x v(t, \cdot)\|_{L_x^2}^2, \end{aligned}$$

where in the last step we have used the bootstrap assumption. This can then be easily absorbed on the left hand side in (6.2).

On the other hand, for the space time integral, we have

$$\begin{aligned} & \left| \int_T^{T_1} \int_{\mathbb{R}^5} 2r^{-4} E(vr) \sin(2u_{approx}) u_{approx,t} dx dt \right| \\ &\lesssim \int_T^{T_1} \int_{\mathbb{R}^5} t^{-1} v^4 dx dt \leq \sup_{t \in [T, T_1]} \left(\int_{\mathbb{R}^5} |v(t, \cdot)|^5 dx \right)^{\frac{2}{5}} \int_T^{T_1} t^{-1} \|\nabla_x v(t, \cdot)\|_{L_x^2}^2 dt. \end{aligned}$$

To bound this last term, using the bootstrap assumption, we have (with an absolute implied constant independent of all other constants)

$$\begin{aligned} & \sup_{t \in [T, T_1]} \left(\int_{\mathbb{R}^5} |v(t, \cdot)|^5 dx \right)^{\frac{2}{5}} \int_T^{T_1} t^{-1} \|\nabla_x v(t, \cdot)\|_{L_x^2}^2 dt \\ &\lesssim (C_1 \delta_1)^2 \int_T^{T_1} t^{-1} (C_1 \delta_1)^2 \left(\frac{t}{T} \right)^{2\gamma} dt \leq \frac{(C_1 \delta_1)^2}{2\gamma} (C_1 \delta_1)^2 \left(\frac{T_1}{T} \right)^{2\gamma}, \end{aligned}$$

which suffices for the bootstrap, provided $\delta_1^2 \ll \gamma$. This deals with the second term on the right hand side of (6.2). To deal with the third term, write $F(x) :=$

$\int_0^x [\cos x - 1] dx$, whence $|F(x)| \lesssim |x|^3$. Then we obtain

$$\begin{aligned} & \int_T^{T_1} \int_{\mathbb{R}^5} \frac{\sin(2u_{approx})}{r^3} (\cos(2vr) - 1) v_t dx dt \\ &= \int_{\mathbb{R}^5} r^{-4} F(vr) \sin(2u_{approx}) dx \Big|_T^{T_1} \\ & - \int_T^{T_1} \int_{\mathbb{R}^5} r^{-4} F(vr) \cos(2u_{approx}) \cdot 2u_{approx,t} dx dt. \end{aligned}$$

Here we have

$$\begin{aligned} \left| \int_{\mathbb{R}^5} r^{-4} F(vr) \sin(2u_{approx}) dx \Big|_T^{T_1} \right| &\leq \sum_{t=T, T_1} \int_{\mathbb{R}^5} \frac{|v^3(t, \cdot)|}{r} dx \\ &\leq 2 \sup_{t \in [T, T_1]} \left\| \frac{v(t, \cdot)}{r} \right\|_{L_x^2} \left\| v^2(t, \cdot) \right\|_{L_x^2} \\ &\lesssim 2 \sup_{t \in [T, T_1]} \left\| \nabla_x v(t, \cdot) \right\|_{L_x^2}^2 \left\| v(t, \cdot) \right\|_{L_x^5} \\ &\ll \sup_{t \in [T, T_1]} \left\| \nabla_x v \right\|_{L_x^2}^2, \end{aligned}$$

where we have used Hardy's inequality and the bootstrap assumption. This can again be absorbed on the left hand side of (6.2). We similarly infer the bound

$$\begin{aligned} & \left| \int_T^{T_1} \int_{\mathbb{R}^5} r^{-4} F(vr) \cos(2u_{approx}) \cdot 2u_{approx,t} dx dt \right| \\ & \lesssim \frac{C\delta_1}{\gamma} (C_1\delta_1)^2 \left(\frac{T_1}{T} \right)^{2\gamma}, \end{aligned}$$

which suffices provided $\delta_1 \ll \gamma$.

For the fourth term on the right hand side of (6.2), we have to argue slightly differently, since now the smallness has to come from u_{approx} , which however is large immediately outside the light cone. Write

$$\begin{aligned} & \int_T^{T_1} \int_{\mathbb{R}^5} \frac{2v}{r^2} [\cos(2u_{approx}) - 1] v_t dx dt \\ &= \int_{\mathbb{R}^5} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx \Big|_T^{T_1} \\ &+ \int_T^{T_1} \int_{\mathbb{R}^5} \frac{v^2}{r^2} \sin(2u_{approx}) u_{approx,t} dx dt. \end{aligned}$$

For the first term on the right and evaluated at $t = T_1$, we get

$$\begin{aligned} \int_{\mathbb{R}^5} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx|_{T_1} &= \int_{r < T_1 - C} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx|_{T_1} \\ &\quad + \int_{r \in [T_1 - C, T_1 + C]} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx|_{T_1}. \end{aligned}$$

For the first term use that $|\cos(2u_{approx}) - 1| \ll_{\tilde{d}_2} 1$, whence by Hardy's inequality

$$\left| \int_{r < T_1 - C} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx|_{T_1} \right| \ll \|\nabla_x v(T_1, \cdot)\|_{L_x^2}^2,$$

which can be absorbed on the right hand side of (6.2). For the remainder term, smallness has to be a consequence of the additional r -localization. In fact, from Strauss' inequality for radial functions, we infer

$$|v(t, r)| \lesssim r^{-\frac{3}{2}} \|v(t, \cdot)\|_{\dot{H}^1},$$

and so

$$\begin{aligned} &\left| \int_{r \in [T_1 - C, T_1 + C]} \frac{v^2}{r^2} [\cos(2u_{approx}) - 1] dx|_{T_1} \right| \\ &\lesssim \|v(T_1, \cdot)\|_{\dot{H}^1}^2 \int_{r \in [T_1 - C, T_1 + C]} r^{-5} \cdot r^4 dr \ll \|v(T_1, \cdot)\|_{\dot{H}^1}^2 \end{aligned}$$

since $T_1 > T \gg 1$ by assumption. Hence this term can be absorbed on the right hand side of (6.2). For the space time integral above, we similarly divide it into an integral over $r < t - C$, $r \in [t - C, t + C]$, and by similar reasoning we obtain

$$\begin{aligned} &\left| \int_T^{T_1} \int_{\mathbb{R}^5} \frac{v^2}{r^2} \sin(2u_{approx}) u_{approx,t} dx dt \right| \\ &\ll \int_T^{T_1} \frac{\|\nabla_x v(t, \cdot)\|_{L_x^2}^2}{t} dt \end{aligned}$$

where the implied constant depends on \tilde{d}_1 as well as T (in particular, the latter needs to be large enough in relation to \tilde{d}_1 for this term to be small), and so we can again close provided the implied constant is small enough in relation to γ .

Finally, to control the last term on the right in (6.2), we use

$$\begin{aligned} \left| \int_T^{T_1} \int_{\mathbb{R}^5} v_t \frac{e_0}{r} dx dt \right| &\lesssim \int_T^{T_1} \|v_t(t, \cdot)\|_{L_x^2} \| \frac{e_0}{r}(t, \cdot) \|_{L_x^2} dt \\ &\lesssim \int_T^{T_1} \left(\frac{\|v_t(t, \cdot)\|_{L_x^2}^2}{t^2} + t^{-2} \right) dt \ll (C_1 \delta_1)^2 \left(\frac{T_1}{T} \right)^{2\gamma} \end{aligned}$$

provided we pick T sufficiently large. This completes the bootstrap for the norm

$$\sup_{t \in [T, T_1]} \|v[t, \cdot]\|_{(\frac{\gamma}{2})^\gamma \dot{H}^1(\mathbb{R}^5) \times (\frac{\gamma}{2})^\gamma L^2(\mathbb{R}^5)}.$$

We continue with the Strichartz type norms of critical scaling, given by

$$\|v\|_{L_t^2(L_x^{10} \cap \nabla_x^{-\frac{1}{2}} L_x^5)([T, T_1] \times \mathbb{R}^5)} + \|(v, v_t)\|_{L_t^\infty \dot{H}^{\frac{3}{2}} \times L_t^\infty \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)}.$$

Using the standard Strichartz estimates for free waves on \mathbb{R}^{5+1} , it suffices to prove the bound

$$\|F(v)\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \ll C_1 \delta_1$$

where $F(v)$ denotes the right hand side of (6.1). We estimate the individual components on the right:

The contribution of $-\frac{1}{r} \frac{\sin(2\epsilon) - 2\epsilon}{r^2} \cos(2u_{approx})$. We can bound this by

$$\begin{aligned} & \left\| \frac{1}{r} \frac{\sin(2\epsilon) - 2\epsilon}{r^2} \cos(2u_{approx}) \right\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \\ & \lesssim \|v^3\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} + \|v^3 (\nabla_x^{\frac{1}{2}} v) r\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} + \|v^4 r^{\frac{1}{2}}\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} \\ & \quad + \|v^3 t^{-\frac{1}{2}}\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)}, \end{aligned}$$

where we have taken advantage of writing

$$\frac{1}{r} \frac{\sin(2\epsilon) - 2\epsilon}{r^2} \cos(2u_{approx}) = v^3 \frac{\sin(2\epsilon) - 2\epsilon}{\epsilon^3} \cos(2u_{approx})$$

and also used the fractional derivative Leibniz rule. Then we estimate

$$\begin{aligned} \|v^3\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} & \lesssim \|\nabla_x^{\frac{1}{2}} v\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \\ & \lesssim \|v\|_{L_t^\infty \dot{H}^{\frac{3}{2}}([T, T_1] \times \mathbb{R}^5)} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \\ & \lesssim (C_1 \delta_1)^3 \ll C_1 \delta_1. \end{aligned}$$

Next, taking advantage of the Strauss' inequality $|v(t, r)| \lesssim r^{-1} \|v(t, \cdot)\|_{\dot{H}^{\frac{3}{2}}}$, we have

$$\begin{aligned} & \|v^3 (\nabla_x^{\frac{1}{2}} v) r\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} + \|v^4 r^{\frac{1}{2}}\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} \\ & \lesssim C_1 \delta_1 \|v^2 (\nabla_x^{\frac{1}{2}} v)\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} + (C_1 \delta_1)^{\frac{1}{2}} \|v^{\frac{7}{2}}\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} \\ & \lesssim C_1 \delta_1 \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \|\nabla_x^{\frac{1}{2}} v\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \\ & \quad + (C_1 \delta_1)^{\frac{1}{2}} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \|v^{\frac{3}{2}}\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \\ & \lesssim (C_1 \delta_1)^{\frac{1}{2}} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \|v\|_{L_t^\infty \dot{H}^{\frac{3}{2}}([T, T_1] \times \mathbb{R}^5)} \ll C_1 \delta_1. \end{aligned}$$

Finally, for the contribution of $v^3 t^{-\frac{1}{2}}$, we use the fact that by the Huyghen's principle, the support of the function $v(t, r)$ is contained in the set $r < t + C$, and so

$$\begin{aligned} \left\| v^3 t^{-\frac{1}{2}} \right\|_{L_t^1 L_x^2([T, T_1] \times \mathbb{R}^5)} &\lesssim \left\| v \right\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \left\| v \right\|_{L_t^\infty L_x^5([T, T_1] \times \mathbb{R}^5)} \left\| t^{-\frac{1}{2}} \right\|_{L_t^\infty L_x^{10}([T, T_1] \times \{r \leq t+C\})} \\ &\lesssim (C_1 \delta_1)^3. \end{aligned}$$

The contribution of $-\frac{2\epsilon}{r^3} [\cos(2u_{approx}) - 1]$. Here we distinguish between high and low frequency factors. Specifically, we write schematically

$$\frac{2\epsilon}{r^3} [\cos(2u_{approx}) - 1] = \frac{2\epsilon}{r} P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] + \frac{2\epsilon}{r} P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right].$$

For the second term on the right, we exploit that

$$P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right]$$

enjoys a special smallness property. In fact, by direct computation, we get

$$\begin{aligned} \nabla_x \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] &= \nabla_x \left[\frac{u_{approx}^2 \cos(2u_{approx}) - 1}{r^2 u_{approx}^2} \right] \\ &= \left[\nabla_x \left(\frac{u_{approx}}{r} \right) \frac{u_{approx}}{r} \frac{\cos(2u_{approx}) - 1}{u_{approx}^2} \right] \\ &\quad + \left[\frac{u_{approx}^2}{r^2} \nabla_x \left(\frac{\cos(2u_{approx}) - 1}{u_{approx}^2} \right) \right] \end{aligned}$$

and we bound these terms by $O\left(\frac{\log t}{t^3}\right)$. It follows that

$$\left\| P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\| \lesssim \frac{\log t}{t^{3-\delta}}.$$

This allows to bound the high frequency term by

$$\begin{aligned} \left\| \frac{2\epsilon}{r} P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{\dot{H}^{\frac{1}{2}}(\mathbb{R}^5)} &\lesssim \left\| v \right\|_{L_x^{10}} \left\| \chi_{r \leq t} \right\|_{L_x^5} \left\| P_{\geq t^{-\delta}} \nabla_x^{\frac{1}{2}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^5(r \leq t)} \\ &\quad + \left\| \nabla_x^{\frac{1}{2}} v \right\|_{L_x^5} \left\| \chi_{r \leq t} \right\|_{L_x^5} \left\| P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^{10}(r \leq t)}. \end{aligned}$$

We conclude that

$$\begin{aligned}
 & \left\| \frac{2\epsilon}{r} P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)} \left\| \left\| t P_{\geq t^{-\delta}} \nabla_x^{\frac{1}{2}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^5(r \leq t)} \right\|_{L_t^2[T, T_1]} \\
 & + \left\| \nabla_x^{\frac{1}{2}} v \right\|_{L_t^2 L_x^5([T, T_1] \times \mathbb{R}^5)} \left\| \left\| t P_{\geq t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^{10}(r \leq t)} \right\|_{L_t^2[T, T_1]} \\
 & \lesssim \left[\|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)} + \left\| \nabla_x^{\frac{1}{2}} v \right\|_{L_t^2 L_x^5([T, T_1] \times \mathbb{R}^5)} \right] \left\| \frac{\log t}{t^{1-\delta}} \right\|_{L_t^2[T, T_1]} \\
 & \ll C_1 \delta_1
 \end{aligned}$$

on account of $T \gg 1$.

Next, consider the low frequency term

$$\begin{aligned}
 & \frac{2\epsilon}{r} P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \\
 & = P_{< t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \\
 & + P_{\geq t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right].
 \end{aligned}$$

For the second term on the right, we get

$$\begin{aligned}
 & \left\| P_{\geq t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim \left\| t^{-\frac{3}{4}\delta} \langle \nabla_x \rangle^{\frac{1}{2}} P_{\geq t^{-\frac{\delta}{2}}} v \right\|_{L_t^\infty L_x^2([T, T_1] \times \mathbb{R}^5)} \left\| t^{\frac{3}{4}\delta} \langle \nabla_x \rangle^{\frac{1}{2}} P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 L_x^\infty}
 \end{aligned}$$

and then use that from our bootstrap hypothesis we have

$$\left\| t^{-\frac{3}{4}\delta} \langle \nabla_x \rangle^{\frac{1}{2}} P_{\geq t^{-\frac{\delta}{2}}} v \right\|_{L_t^\infty L_x^2([T, T_1] \times \mathbb{R}^5)} \lesssim C_1 \delta_1$$

while using Bernstein's inequality, we have

$$\begin{aligned}
 & \left\| t^{\frac{3}{4}\delta} \langle \nabla_x \rangle^{\frac{1}{2}} P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 L_x^\infty([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim \left\| \frac{u_{approx}}{r} \right\|_{L_t^2 L_x^{10+}([T, T_1] \times \mathbb{R}^5)}^2 \ll 1
 \end{aligned}$$

on account of $T \gg 1$. The conclusion is that

$$\left\| P_{\geq t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{< t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \ll C_1 \delta_1.$$

On the other hand, for the term where all factors have low frequency, i. e.

$$P_{<t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right],$$

we exploit the extra outer derivative and low frequency control (with a small loss): we have at fixed time t

$$\begin{aligned} & \left\| P_{<t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{\dot{H}^{\frac{1}{2}}(\mathbb{R}^5)} \\ & \lesssim \left\| \nabla_x^{\frac{1}{2}} P_{<t^{-\frac{\delta}{2}}} v \right\|_{L_x^{\frac{10}{3}}} \left\| P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^5(r \lesssim t)} \\ & + \left\| P_{<t^{-\frac{\delta}{2}}} v \right\|_{L_x^{\frac{10}{3}}} \left\| \nabla_x^{\frac{1}{2}} P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_x^5(r \lesssim t)} \lesssim t^{-\frac{\delta}{4} + \gamma - 1} T^{-\gamma} C_1 \delta_1 \end{aligned}$$

where we have taken advantage of

$$\left\| P_{<t^{-\frac{\delta}{2}}} v(t, \cdot) \right\|_{L_x^{\frac{10}{3}}} \lesssim \|v(t, \cdot)\|_{\dot{H}^1} \lesssim \left(\frac{t}{T} \right)^\gamma.$$

Also, the additional factors $t^{-\frac{\delta}{4}}$ which ensure integrability stem from the operator $\nabla_x^{\frac{1}{2}} P_{<t^{-\delta}}$. Thus if we arrange (as we may) that $\gamma \ll \delta$, we find

$$\begin{aligned} & \left\| P_{<t^{-\frac{\delta}{2}}} \left(\frac{2\epsilon}{r} \right) P_{<t^{-\delta}} \left[\frac{\cos(2u_{approx}) - 1}{r^2} \right] \right\|_{L_t^1 \dot{H}^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \lesssim C_1 \delta_1 \left\| t^{-\frac{\delta}{4} + \gamma - 1} T^{-\gamma} \right\|_{L_t^1 [T, T_1]} \\ & \ll C_1 \delta_1 \end{aligned}$$

provided T is sufficiently large (in relation to δ^{-1}). This finally concludes bounding the contribution from

$$-\frac{2\epsilon}{r^3} [\cos(2u_{approx}) - 1].$$

The contribution of $\frac{\sin(2u_{approx})}{r^3} (\cos(2\epsilon) - 1)$. We again use the high-low frequency method, which is somewhat simpler to implement here: write

$$\begin{aligned} & \frac{\sin(2u_{approx})}{r^3} (\cos(2\epsilon) - 1) \\ & = P_{<t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \left(\frac{\cos(2\epsilon) - 1}{r^2} \right) + P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \left(\frac{\cos(2\epsilon) - 1}{r^2} \right). \end{aligned}$$

For the second term on the right, use that

$$\begin{aligned} & \left\| \langle \nabla_x \rangle^{\frac{1}{2}} P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_x^{\frac{10}{3}}(r \lesssim t)} \lesssim \frac{\log t}{t^{\frac{1}{2} - \delta}}, \\ & \left\| P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_x^{10}(r \lesssim t)} \lesssim \frac{\log t}{t^{\frac{3}{2} - \delta}}, \end{aligned}$$

and so we find

$$\begin{aligned}
 & \left\| P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \left(\frac{\cos(2\epsilon) - 1}{r^2} \right) \right\|_{L_t^1 \dot{H}_x^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim \left\| \nabla_x^{\frac{1}{2}} P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)}^2 \\
 & + \left\| P_{\geq t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)} \left\| \nabla_x^{\frac{1}{2}} v \right\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim (C_1 \delta_1)^2 \ll C_1 \delta_1.
 \end{aligned}$$

For the low frequency term, we get

$$\begin{aligned}
 & \left\| P_{< t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \left(\frac{\cos(2\epsilon) - 1}{r^2} \right) \right\|_{L_t^1 \dot{H}_x^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \\
 & \lesssim \left\| \nabla_x^{\frac{1}{2}} v \right\|_{L_t^2 L_x^5([T, T_1] \times \mathbb{R}^5)} \|t^{-\gamma} v\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \left\| t^\gamma P_{< t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_t^2 L_x^\infty} \\
 & + \|v\|_{L_t^2 L_x^{10}([T, T_1] \times \mathbb{R}^5)} \|t^{-\gamma} v\|_{L_t^\infty L_x^{\frac{10}{3}}([T, T_1] \times \mathbb{R}^5)} \left\| t^\gamma \nabla_x^{\frac{1}{2}} P_{< t^{-\delta}} \left(\frac{\sin(2u_{approx})}{r} \right) \right\|_{L_t^2 L_x^{10}} \\
 & \lesssim (C_1 \delta_1)^2 \ll C_1 \delta_1.
 \end{aligned}$$

The contribution of $\frac{e_0}{r}$. Here we immediately check that $\left\| \frac{e_0}{r} \right\|_{L_t^1 \dot{H}_x^{\frac{1}{2}}([T, T_1] \times \mathbb{R}^5)} \lesssim T^{-1} \ll C_1 \delta_1$ if T is sufficiently large, which is as desired. The proof of Proposition 6.4 is thereby concluded.

7. PROOF OF THE MAIN RESULT

Here we shall show how to conclude Theorem 2.1 building on the previous sections. Indeed Theorem 2.1 follows from Proposition 6.1, the solution to (1.7) being given by $u := u_{approx} + \epsilon(t, r)$ where $\epsilon(t, r) := rv$ with v provided by Proposition 6.1. The initial data (f, g) are given by $(u(T, \cdot), u_t(T, \cdot))$ where $T > 0$ depending on \tilde{d}_1 is given by Proposition 6.1. Since (v, v_t) does have finite critical norm, but the approximate solution $(u_{approx}, \partial_t u_{approx})$ does not, we easily conclude that the initial data have infinite critical norm. Clearly the perturbation (ϵ, ϵ_t) lies in the spaces $\dot{H}^s \times \dot{H}^{s-1}$ for $s > \frac{3}{2}$ by construction (remind that (v, v_t) is compactly supported). Moreover, due to the asymptotics for $r \rightarrow \infty$ of the self-similar solutions given by formulas (3.34) and (4.12), in the small and large case respectively, we have that $(u_{approx} - q_1, \partial_t u_{approx})$, $(u_{approx} - c_1, \partial_t u_{approx})$ respectively, has finite norm in $\dot{H}^s \times \dot{H}^{s-1}$ for $s > \frac{3}{2}$: that is how we understand the finiteness in $\dot{H}^s \times \dot{H}^{s-1}$ for $s > \frac{3}{2}$ of the data (f, g) as claimed in Theorem 2.1. Of course the condition $\|f\|_{L^\infty(r \geq 1)} > M$, for arbitrary $M > 0$, can be achieved simply by choosing $\tilde{d}_1 > M^2$ in the context of large self-similar solutions as provided by Lemma 4.1. Finally the stability under a certain class of perturbations is a consequence of the fact that v belongs to an open set with respect to the norms of Proposition 6.2.

□

REFERENCES

- [1] Piotr Bizoń, *Equivariant self-similar wave maps from Minkowski spacetime into 3-sphere*, Comm. Math. Phys. **215** (2000), no. 1, 45–56.
- [2] Thierry Cazenave, Jalal Shatah, and A. Shadi Tahvildar-Zadeh, *Harmonic maps of the hyperbolic space and development of singularities in wave maps and Yang-Mills fields*, Ann. Inst. H. Poincaré Phys. Théor. **68** (1998), no. 3, 315–349.
- [3] Roland Donn timer, *On stable self-similar blowup for equivariant wave maps*, Comm. Pure Appl. Math. **64** (2011), no. 8, 1095–1147.
- [4] Roland Donn timer and Peter C. Aichelburg, *Spectral properties and linear stability of self-similar wave maps*, J. Hyperbolic Differ. Equ. **6** (2009), no. 2, 359–370.
- [5] ———, *A note on the eigenvalues for equivariant maps of the SU(2) sigma-model*, Appl. Math. Comput. Sci. **1** (2010), no. 1, 73–82.
- [6] Roland Donn timer, Birgit Schörkhuber, and Peter C. Aichelburg, *On stable self-similar blowup for equivariant wave maps: the linearized problem*, Ann. Henri Poincaré **13** (2012), no. 1, 103–144.
- [7] Pierre Germain, *Besov spaces and self-similar solutions for the wave-map equation*, Comm. Partial Differential Equations **33** (2008), no. 7-9, 1571–1596.
- [8] Joachim Krieger, *Global regularity of wave maps from \mathbf{R}^{3+1} to surfaces*, Comm. Math. Phys. **238** (2003), no. 1-2, 333–366.
- [9] Joachim Krieger and Wilhelm Schlag, *Large global solutions for energy supercritical nonlinear wave equations on \mathbf{R}^{3+1}* , J. Anal. Math. (accepted).
- [10] Richard Schoen and Karen Uhlenbeck, *Regularity of minimizing harmonic maps into the sphere*, Invent. Math. **78** (1984), no. 1, 89–100.
- [11] Jalal Shatah, *Weak solutions and development of singularities of the SU(2) σ -model*, Comm. Pure Appl. Math. **41** (1988), no. 4, 459–469.
- [12] Jalal Shatah and A. Shadi Tahvildar-Zadeh, *On the Cauchy problem for equivariant wave maps*, Comm. Pure Appl. Math. **47** (1994), no. 5, 719–754.
- [13] Thomas C. Sideris, *Global existence of harmonic maps in Minkowski space*, Comm. Pure Appl. Math. **42** (1989), no. 1, 1–13.
- [14] Atsushi Tachikawa, *Rotationally symmetric harmonic maps from a ball into a warped product manifold*, Manuscripta Math. **53** (1985), no. 3, 235–254.
- [15] Neil Turok and David Spergel, *Global Texture and the Microwave Background*, Phys.Rev.Lett. **64** (1990), no. 23, 2736–2739.