

DOUGLAS REPORT SM-49213

NUMERICAL RESIDUAL PERTURBATION SOLUTIONS  
APPLIED TO THE PROBLEM OF A CLOSE SATELLITE  
OF THE SMALLER BODY  
IN THE RESTRICTED THREE-BODY PROBLEM

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ABSTRACT

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The purpose of this report is to demonstrate the new method of numerical residual perturbation solution as applied to the problem of a close satellite of the smaller body in the restricted three-body problem. Cowell demonstrated a method of numerically solving the total differential equations of motion of an orbiting object. The variation of parameters and Encke's methods take advantage of the known analytic solution to the two-body problem and numerically handle only the perturbations to the orbit. This report demonstrates the use of an analytic series perturbation solution as a reference orbit (rather than using conics as a reference) and the numerical solution of the residual (i.e., not accounted for by the analytic solution) perturbation equations of motion. The new method shows advantages over the full region of validity of the analytic reference orbit. The analytic theory of motion used is such as to limit the application of the results to orbits of small inclination and eccentricity with special initial conditions. The program is thus just for demonstration. This work was supported by contract NASw-901.

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Section 1  
INTRODUCTION

The purpose of this paper is to show the relative merits for prediction of future position of a satellite of the smaller body in the restricted three-body problem, of the analytic solution of reference 1, of numerical solution of the residual perturbation differential equations not accounted for by the theory of motion in reference 1, of the perturbation differential equations from a Keplerian ellipse, and of the total equations of motion.

In a numerical solution, the time interval is normally chosen so as to make the error due to series truncation in the approximating function approximately equal to the error due to truncation of numbers in the computation of the equations of motion and the approximating function. The balancing of these errors assures maximum accuracy from the method. (Actually, to determine mathematically the time interval, it was necessary in this program to make the series truncation error slightly larger than number truncation error.) In a numerical solution of the total equations, if  $\delta Y$  is added to  $Y_i$  at each point, then the error builds up as:

$$\begin{array}{r}
 Y_i = \text{xxxxxxxx} \\
 \quad \quad \quad \swarrow \quad \downarrow \\
 \quad \quad \quad \text{digit affected by series truncation} \\
 \quad \quad \quad \text{digit affected by number truncation} \\
 + \quad \delta Y = \text{xxxxxxxx} \\
 \hline
 = Y_{i+1} = \text{xxxxxxxx} \\
 \quad \quad \quad \swarrow \\
 \quad \quad \quad \text{digit where error accumulates}
 \end{array}$$

If it is possible to write  $Y$  as  $Y_a + \Delta Y$  where  $Y_a$  is obtained by an analytic approximation and only  $\Delta Y$  is obtained by numerical solution of the differential equations, then the error accumulates as:

$$\begin{array}{r}
\Delta Y_i = \text{xxxxxxxx} \\
+ \delta Y_i = \text{xxxxxxxx} \\
\hline
= \Delta Y_{i+1} = \text{xxxxxxxx} \\
+ Y_{ai+1} = \text{xxxxxxxx} \\
\hline
= Y_{i+1} = \text{xxxxxxxx} \text{ (x)}
\end{array}$$

digit affected by series truncation  
digit affected by number truncation  
digit where error accumulates  
digit where error accumulates  
digit where number truncation occurs  
but is not accumulated.

It is, of course, obvious that the better the approximation  $Y_a$ , the smaller  $\Delta Y$  will be and the slower the accumulation of number truncation.

This approach in this paper differs from the work of Encke primarily in that an approximate analytic solution of the perturbed motion is used as  $Y_a$  rather than  $Y_a$  being an unperturbed Kepler ellipse. The procedure assumes a differential equation of the form:

$$\ddot{Y} = f(X, Y, \dot{Y}) + g(X, Y, \dot{Y}) \tag{I-1}$$

It is approximated by:

$$\ddot{Y} = f(X, Y, \dot{Y}) \tag{I-2}$$

and equation (I-2) is solved analytically to yield the approximation  $Y_a$  in the form of a finite asymptotic series. The exact solution,  $Y$ , of equation (I-1) is written in the form:

$$Y = Y_a + Z \tag{I-3}$$

Equation (I-3) is differentiated twice and the result substituted for the left-hand side of equation (I-1). The value of  $Z$  is computed by a numerical

solution of this differential equation. The solution,  $Y$ , involves two approximations in this case. First, it is derived from equation (I-2), not equation (I-1); second, the truncation of the series for  $Y_a$  represents an error. It may be seen that the numerical procedure for  $Z$  corrects for both approximations.

The above discussion is independent of whether the machine computations are done single-precision, double-precision, or  $n$ -precision since numerical solution of the total equations of motion, of the perturbations from a conic, and of the residual perturbation from a theory may all be done to  $n$ -precision without changing the relative advantages.

Reference 1 was an early work in the series of papers on celestial mechanics described in the list of references. Besides being limited to small eccentricities and inclinations, it does not generalize initial conditions, but always begins with the perigee and node lined up along the Earth-Moon line. Therefore, rectification such as Encke used is not possible. In using it as a basis for the numerical solution, there results a demonstration for the new numerical techniques but not a useful analytic tool. For demonstration purposes, it was not deemed worth-while to attempt the cancellation of all oscillatory terms of the highest order carried by the approximation in order to provide the minimum residual perturbations. What has been done in this paper is to use the precessing elliptic orbit which provides the zero order portion of reference 1 as the reference orbit. This is an appreciable improvement over a fixed initial osculating ellipse and serves to demonstrate the usefulness of the technique.

In this example, the equations of motion are those of the restricted three-body problem. No other perturbations are considered. However, no real increase in complexity or decrease in accuracy result from adding other perturbations. This is being demonstrated in a companion report for a satellite of an oblate planet. In the companion report, an analytic approximation is used for  $J_2$  and  $J_4$  perturbations;  $J_3$  and higher-order Earth potential terms as well as Sun and Moon perturbations are treated numerically.

This report consists of a description of the program, a discussion of the results, and plotted results for a variety of mass ratios, semi-major axes, eccentricities, and inclinations. The general flow of the program is illustrated in figure 1. The subsection numbers on the top of each block refer to Section 4 where the equations are detailed.

In this report, the equal sign is often used in the FORTRAN sense, i.e.,

$$x = x + 1$$

means replace  $x$  by  $x + 1$ . Similarly, symbols such as  $A(N)$  are used to refer to an array of  $N$  numbers.

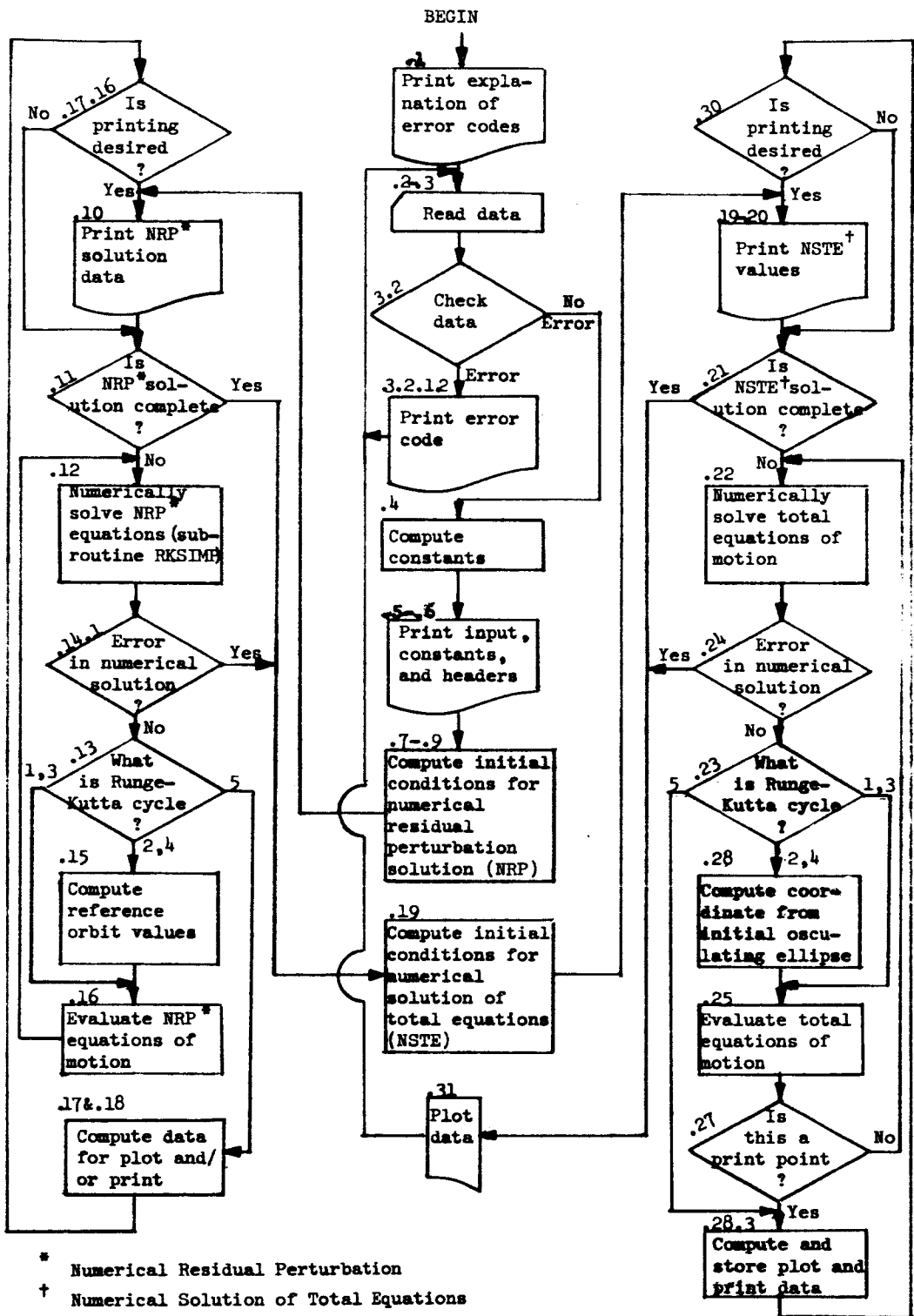


Figure 1

Section 2  
DEFINITION OF SYMBOLS

a	Semi-major axis																		
a	Subscript for analytic solution from reference 1																		
$a_1$	Semi-major axis in analytic solution = A1																		
$a_2$	$a^2$																		
$a_3$	$a^3 = A3$																		
$a_{12}$	$\sqrt{a}$																		
$a_{32}$	$a^{3/2}$																		
$a_{92}$	$a^{9/2} = A92$																		
$a_e$	$a(1-e^2) = AEC$																		
$a_{e1/2}$	$\sqrt{a(1-e^2)}$																		
$A_p$	Plotting array of dimension 13,000 with elements																		
	<table border="0" style="width: 100%;"> <tr> <td style="width: 50%;"><math>ap_{1+13i} = (\Delta C/C)_{ai}</math></td> <td style="width: 50%;"><math>ap_{7+13i} = (\Delta \rho/a)_{pki}</math></td> <td rowspan="6" style="font-size: 4em; vertical-align: middle; padding-left: 10px;">}</td> <td rowspan="6" style="vertical-align: middle; padding-left: 10px;">i = 0, 1, ..., 999</td> </tr> <tr> <td><math>ap_{2+13i} = (\Delta C/C)_{pki}</math></td> <td><math>ap_{8+13i} = (\Delta \rho/a)_{ci}</math></td> </tr> <tr> <td><math>ap_{3+13i} = (\Delta C/C)_{ci}</math></td> <td><math>ap_{9+13i} = (\Delta \rho/a)_{ki}</math></td> </tr> <tr> <td><math>ap_{4+13i} = (\Delta C/C)_{ki}</math></td> <td><math>ap_{10+13i} = (\Delta \dot{\rho} a^{1/2})_{ai}</math></td> </tr> <tr> <td><math>ap_{5+13i} = (\Delta C/C)_{pi}</math></td> <td><math>ap_{11+13i} = (\Delta \dot{\rho} a^{1/2})_{pki}</math></td> </tr> <tr> <td><math>ap_{6+13i} = (\Delta \rho/a)_{ai}</math></td> <td><math>ap_{12+13i} = (\Delta \dot{\rho} a^{1/2})_{ci}</math></td> </tr> <tr> <td></td> <td style="text-align: center;"><math>ap_{13+13i} = (\Delta \dot{\rho} a^{1/2})_{ki}</math></td> <td></td> <td></td> </tr> </table>	$ap_{1+13i} = (\Delta C/C)_{ai}$	$ap_{7+13i} = (\Delta \rho/a)_{pki}$	}	i = 0, 1, ..., 999	$ap_{2+13i} = (\Delta C/C)_{pki}$	$ap_{8+13i} = (\Delta \rho/a)_{ci}$	$ap_{3+13i} = (\Delta C/C)_{ci}$	$ap_{9+13i} = (\Delta \rho/a)_{ki}$	$ap_{4+13i} = (\Delta C/C)_{ki}$	$ap_{10+13i} = (\Delta \dot{\rho} a^{1/2})_{ai}$	$ap_{5+13i} = (\Delta C/C)_{pi}$	$ap_{11+13i} = (\Delta \dot{\rho} a^{1/2})_{pki}$	$ap_{6+13i} = (\Delta \rho/a)_{ai}$	$ap_{12+13i} = (\Delta \dot{\rho} a^{1/2})_{ci}$		$ap_{13+13i} = (\Delta \dot{\rho} a^{1/2})_{ki}$		
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	$ap_{13+13i} = (\Delta \dot{\rho} a^{1/2})_{ki}$																		
$a_s$	Semi-major axis of precessing ellipse																		



$a_{s32}$	$a_s^{3/2}$
A1	FORTRAN symbol for semi-major axis, $a$
A2	FORTRAN symbol for $a^2$
A3	FORTRAN symbol for $a_3 = a^3$
A12	FORTRAN symbol for $\sqrt{a}$
A32	FORTRAN symbol for $a^{3/2} = a_{32}$
A92	FORTRAN symbol for $a^{9/2} = a_{92}$
ACOB1	FORTRAN symbol for the Jacobi integral - $1/\epsilon^2 r$ (Computed at initial time using numerical residual perturbation solution values)
AEC	FORTRAN symbol for $a(1-e^2)$
AESQ	FORTRAN symbol for $\sqrt{a_s(1-e_s^2)}$
ALPHA1	FORTRAN symbol for $\alpha_1$
ALPHA2	FORTRAN symbol for $\alpha_2$
AP	FORTRAN symbol for $A_p$
AP(I)	FORTRAN symbol for $a_{pi}$
ARNODE	FORTRAN symbol for argument of node, $\Omega$
AS	FORTRAN symbol for $a_s$
AS32	FORTRAN symbol for $a_s^{3/2}$

B1	FORTTRAN symbol for $C_{15} + \rho_s C_{16}$
B2	FORTTRAN symbol for $C_{17} + \rho_s C_{18}$
BETA	FORTTRAN symbol for $\beta$
C	FORTTRAN symbol for the Jacobi constant
c	Subscript for values obtained from numerical solution of total equations of motion
$C_1$	$\frac{15}{4} \epsilon a^3 e$
$C_2$	$\frac{15}{4} \epsilon a^{1/2} e$
$C_3$	$\frac{\frac{3}{8} \epsilon i a^{5/2}}{1 - \epsilon a^{3/2}}$
$C_6$	$\dot{\Omega}$
$C_7$	$\epsilon^4 \lambda_1^2 = \dot{\Omega}^2$
$C_8$	$1 - \epsilon^4$
$C_9$	$a^{3/2} (1 - e^2)^{3/2}$
$C_{10}$	$\frac{\epsilon^2 \omega_1}{1 + \epsilon^2 \omega_1} + \epsilon_1 a^{3/2} (1 + \epsilon_1 m_1)$
$C_{11}$	$2P_1 - 1$
$C_{12}$	$2\dot{\Omega} / \sqrt{a(1-e^2)} = 2\dot{\Omega}/a_{e1/2}$
$C_{13}$	$2\dot{\omega} / \sqrt{a(1-e^2)} = 2\dot{\omega}/a_{e1/2}$

$C_{14}$	$2C_{\dot{\omega}\dot{\Omega}} = 2\dot{\omega}\dot{\Omega}$
$C_{15}$	$-\frac{2}{ae^{1/2}} (\dot{\Omega} + \dot{\omega} \cos i)$
$C_{16}$	$-2\dot{\omega}\dot{\Omega} - (\dot{\omega}^2 + \dot{\Omega}^2) \cos i$
$C_{17}$	$\frac{2}{ae^{1/2}} (\dot{\Omega} \cos i + \dot{\omega})$
$C_{18}$	$2\dot{\omega}\dot{\Omega} \cos i + \dot{\omega}^2 + \dot{\Omega}^2$
$C_{19}$	$\dot{\omega}^2 \sin i$
$C_{20}$	$\frac{2e \sin i}{ae^{1/2}} \dot{\omega}$
$C_{21}$	$\frac{2 \sin i}{ae^{1/2}} \dot{\omega}$
$C_{22}$	$-\frac{2e}{ae^{1/2}} (\dot{\Omega} + \dot{\omega} \cos i)$
$C_{23}$	$\frac{2e}{ae^{1/2}} (\dot{\Omega} \cos i + \dot{\omega})$
$C_{24}$	$\sqrt{1-e_s^2}$
$C_E$	$\cos E$
$C_i$	$\cos i$
$C_{init}$	Initial value of Jacobi constant

$C_{\max}$	Maximum of absolute values of the components of the velocity increment over the last two Runge-Kutta cycles = CMAX
$C_{P1\psi}$	$\cos (2p_1 - 1)\psi$
$C_{\text{part}}$	Part of Jacobi constant (see Section 3, equation (68))
$C_v$	$\cos v$
$C_{\epsilon t}$	$\cos \epsilon t$
$C_{\theta_1}$	$\cos \theta_1$
$C_{\phi_s}$	$\cos \phi_s = \cos (v + \omega)$
$C_\psi$	$\cos \psi$
$C_\Omega$	$\cos \Omega$
$C_\omega$	$\cos \omega$
$C_{\omega\Omega}$	$\dot{\omega}\dot{\Omega}$
$C_{\omega^2\Omega^2}$	$\dot{\omega}^2 + \dot{\Omega}^2$
C1 through C3	FORTRAN symbol for $C_1$ through $C_3$
C6 through C21	FORTRAN symbol for $C_6$ through $C_{21}$
CALPHA	FORTRAN symbol for $\cos \alpha_1$
CC11PS	FORTRAN symbol for $\cos (C_{11}\psi)$

CEC	FORTRAN symbol for $\cos E$
CEPT	FORTRAN symbol for $\cos \epsilon t$
CI	FORTRAN symbol for cosine inclination, $\cos i$
CIS	FORTRAN symbol for $\cos i_s$
CMAX	FORTRAN symbol for the maximum of absolute values of the components of the velocity increment over the last two Runge-Kutta cycles = C
CNO	FORTRAN symbol for $\cos \Omega$
COPER	FORTRAN symbol for $\cos \omega_s$
COV	FORTRAN symbol for $\cos v_s$
CP1	FORTRAN symbol for $\cos (2p_1 - 1)\psi$
CP <sub>3</sub>	FORTRAN symbol for $\cos p_3\psi$
CP1PS	FORTRAN symbol for $\cos (p_1\psi + \alpha_1)$
CP1PST	FORTRAN symbol for $\cos (p_1\psi)$
CPS	FORTRAN symbol for $\cos \psi$
CSNO	FORTRAN symbol for $\cos \Omega_s$
CTHETA	FORTRAN symbol for $\cos \theta_1$
CTRL	FORTRAN symbol for $\cos (v + \omega_s)$

$D_n$	$1 + e \cos E$
DEGI	FORTRAN symbol for $i$ (degrees)
DELM	FORTRAN symbol for the error in mean anomaly when solving Kepler's equation
DELXD, DELYD, DELZD	In the Runge-Kutta subroutine, FORTRAN symbols for the increment of the components of velocity accumulated for the time interval computation
DEN	FORTRAN symbol for $D_n$
DNOD	FORTRAN symbol for $\dot{\Omega} = C_6$
DPDN	FORTRAN symbol for $\ddot{\omega}\dot{\Omega} = C_{\omega\Omega}$
DPER	FORTRAN symbol for $\dot{\omega}$
D2PHDT	FORTRAN symbol for $d^2\phi/dt^2$
DPHIDT	FORTRAN symbol for $d\phi/dt$
DPNSQ	FORTRAN symbol for $\dot{\omega}^2 + \dot{\Omega}^2 = C_{\omega\Omega}^2$
D2PSDT	FORTRAN symbol for $d^2\psi/dt^2$
DPSIDT	FORTRAN symbol for $d\psi/dt$
DR	FORTRAN symbol for $(\Delta x^2 + \Delta y^2 + \Delta z)^2$
DRC	FORTRAN symbol for the magnitude of the vector position difference between the numerical residual perturbation solution and the fixed initial osculating ellipse solution at a given time

DRDC	FORTRAN symbol for the magnitude of the vector velocity difference between the numerical residual perturbation solution and the numerical solution of the total equation's of motion at a given time
DRDF	FORTRAN symbol for the magnitude of the vector velocity difference between the numerical residual perturbation solution and the initial osculating ellipse value at a given time
DRF	FORTRAN symbol for the magnitude of the vector position difference between the numerical residual perturbation solution and the initial osculating ellipse value at a given time
DSDPSI	FORTRAN symbol for $ds/d\psi$
DSDT	FORTRAN symbol for $ds/dt$
DT	FORTRAN symbol for $\Delta t$
DT2	FORTRAN symbol for $\Delta t/2$
DT3	FORTRAN symbol for $\Delta t/3$
DT4	FORTRAN symbol for $\Delta t/4$
DT6	FORTRAN symbol for $\Delta t/6$
DT1DPS	FORTRAN symbol for $\epsilon^2 \left( \frac{dt}{d\psi} \right)_1$
DTDPSI	FORTRAN symbol for $dt/d\psi$
DTM	FORTRAN symbol for multiplier on $\Delta t$ if error is less than the minimum allowable = 1.5 if input $W_9 \leq 1$
DTODPS	FORTRAN symbol for $\left( \frac{dt}{d\psi} \right)_0$

DX, DY, DZ,	FORTTRAN symbol for $\Delta x, \Delta y, \Delta z$
DXD, DYD, DZD,	$\Delta \dot{x}, \Delta \dot{y}, \Delta \dot{z}$
DX2D, DY2D, DZ2D	$\Delta \ddot{x}, \Delta \ddot{y}, \Delta \ddot{z}$

position, velocity and acceleration obtained from residual perturbation equations. Coordinate system references to plane of planetary motion.

**E** Eccentric anomaly in main program (measured from apogee). In subroutine RKSIMP, the maximum component of the difference between the Runge-Kutta and Simpson rule solutions for velocity

**e** Eccentricity (initial)

**e<sub>2</sub>**  $\sqrt{1-e^2}$

**E<sub>min</sub>** Minimum error in Runge-Kutta routine

**E<sub>n</sub>** Tentative guess at eccentric anomaly

**E<sub>r1</sub>**  $(1-\epsilon^4)\epsilon^2/r^3$

**E<sub>r2</sub>**  $(1-\epsilon^4)\epsilon^2[2(x' \cos \epsilon t + y' \sin \epsilon t) + \epsilon^2 \rho^2] \cdot \frac{[1+r(1+r)]}{r^3(1+r)}$

**EN** FORTTRAN symbol for **E<sub>n</sub>**

**e<sub>s</sub>** Eccentricity of precessing ellipse

**E<sub>2x</sub>**  $\epsilon^2 x$

**e<sub>2</sub>**  $\sqrt{1-e^2} = EC2$ ; (also temporary value of  $1-e^2$ )

**e<sub>3</sub>**  $1 + e$



EALL                   FORTRAN symbol for maximum allowable error in Runge-Kutta =  $|W_8|$   
                           set =  $10^{-7}$  if input  $W_8 = 0$

EC1                    FORTRAN symbol for e

EC2                    FORTRAN symbol for  $e_2 = \sqrt{1-e^2}$ ; (also temporary value of  $1-e^2$  in Sections 4.3.22 - 4.4)

EC3                    FORTRAN symbol for  $1+e$

ECAN                   FORTRAN symbol for eccentric anomaly (measured from apogee) in analytic solution

ELMB                   FORTRAN symbol for  $\lambda_1$

EM                     FORTRAN symbol; in subroutine PLOT, the time at the beginning of each series of  $\Delta C/C$ ,  $\Delta \rho/a$ ,  $\Delta \dot{\rho} a^{1/2}$  plots

EN                     FORTRAN symbol; in subroutine PLOT, the time at the end of each series of  $\Delta C/C$ ,  $\Delta \rho/a$ ,  $\Delta \dot{\rho} a^{1/2}$  plots

EP1                    FORTRAN symbol for  $\epsilon$

EP2                    FORTRAN symbol for  $\epsilon^2$

EP4                    FORTRAN symbol for  $\epsilon^4$

EPA                    FORTRAN symbol for  $\epsilon a^{3/2}$

EPR1                   FORTRAN symbol for  $(1-\epsilon^4)\epsilon^2/r^3$

EPR2                   FORTRAN symbol for  $(1-\epsilon^4)(1 - 1/r^3)$

EP2X                   FORTRAN symbol for  $\epsilon^2 x = E_{2X}$

ERMIN                   FORTRAN symbol for the minimum allowable error =  $|W_{12}|$

ERROR	FORTRAN symbol for $t(\psi) - t$ (should equal zero)
ES	FORTRAN symbol for $e_s$
ESA	FORTRAN symbol for the eccentric anomaly of two-body solution
ESTER	FORTRAN symbol for estimated error in Runge-Kutta subroutine
f	$\frac{1 - (1 + 2q)^{-3/2}}{q} = \text{Encke's } f \text{ series}$
FAC(40)	FORTRAN symbol for the factor in polynomial expression for f
FTD	FORTRAN symbol for a factor ( $< 1$ ) by which the calculated computing interval is multiplied before use to protect against round-off error
FPS(I)	FORTRAN symbol for $F_{\psi i}$
FQ	FORTRAN symbol for $f_q = 1 - (1 + 2q)^{-3/2}$ See Section 3, equation (102)
FRA(40)	FORTRAN symbol for the fraction by which a term in series for f must be multiplied to give succeeding term
G	$\frac{-\left(\frac{dt}{d\psi}\right)_0 \epsilon^2 \omega_1 + \epsilon^2 \frac{dt_1}{d\psi} (1 + \epsilon^2 \omega_1)}{\frac{dt}{d\psi} \left(\frac{dt}{d\psi}\right)_0 (1 + \epsilon^2 \omega_1)} - \epsilon^2 \omega_1 - \dot{\omega}$
GAMMA	FORTRAN symbol for $et$ modulo $2\pi$

GNU1

FORTRAN symbol for  $v_1$

HX, HY, HZ

FORTRAN symbols for dummy variables of position,

HXD, HYD, HZD

velocity, and acceleration in Runge-Kutta subroutine.

HX2D, HY2D, HZ2D

Equivalent to:

$\Delta \vec{x}$ ,  $\dot{\Delta \vec{x}}$ ,  $\ddot{\Delta \vec{x}}$  for numerical residual perturbation  
solution

$\vec{x}$ ,  $\dot{\vec{x}}$ ,  $\ddot{\vec{x}}$  for numerical solution of total equations  
of motion

HXDA, HYDA, HZDA

FORTRAN symbol for Simpson's rule; integrated values  
of  $\dot{x}$ ,  $\dot{y}$ ,  $\dot{z}$ , respectively

I

FORTRAN symbol, main program cards 220-230 -- index  
for zeroing input array W

Main program cards 563-564 -- index for zeroing first  
13 elements of plotting array AP

Main program card 1030 -- index for printing input  
array W

Main program cards 2510-2660 -- index for Encke's f  
series

Main program card 5380 -- index used to print plot-  
table data at computing point ITD

Runge-Kutta subroutine cards 6670-6680 -- index for  
restoring initial condition array S from array SP  
after print point, which is not a normal computing  
point while integrating total equations of motion

Runge-Kutta subroutine cards 7140-7150, 7670-7680,  
and 7850-7860 -- index for restoring initial condi-  
tion array S from array SS when computing interval  
fails

• i                   Inclination (initial) held constant in analytic solution (radians)

i°                    Inclination (degrees)

i<sub>s</sub>                  Inclination of precessing ellipse

ICH                   FORTRAN symbol for the flag to signal which equations are being numerically solved  
                       Residual perturbation equations if ICH = 1  
                       Total equations of motion if ICH = 2

IERR                  FORTRAN symbol; in subroutine PLOT for number of off-scale points

II                    FORTRAN symbol; in subroutine PLOT, an index giving the plot point number of the first point of each series of  $\Delta C/C$ ,  $\Delta \rho/a$ ,  $\Delta \dot{\rho} a^{1/2}$  plots

III                   FORTRAN symbol; in subroutine PLOT, an index giving the plot point number of the last point of each series of  $\Delta C/C$ ,  $\Delta \rho/a$ ,  $\Delta \dot{\rho} a^{1/2}$  plots

IIK                   FORTRAN symbol; in subroutine PLOT, an index for successive sets of  $\Delta C/C$ ,  $\Delta \rho/a$ ,  $\Delta \dot{\rho} a^{1/2}$  plots

IP                    FORTRAN symbol for the initial point flag  
                       IP = 1 on first point  
                       IP = 2 on all others  
                       Set = 1 by main program  
                       Set = 2 by Runge-Kutta subroutine



K In subroutine PLOT, subscript for initial osculating ellipse values

KC FORTRAN symbol for flag to signal Simpson's rule computations  
No if KC = 1  
Yes if KC = 2  
Under control of Runge-Kutta subroutine

KF FORTRAN symbol for intermediate failure counter in Runge-Kutta subroutine

KFAIL FORTRAN symbol for failure counter in Runge-Kutta subroutine

KHALT FORTRAN symbol for halt flag  
If = 1, run continued  
If = 2, run halted, next case called  
If = 3, perturbed trajectory completed, numerical solution of total equations of motion to be initiated

KR FORTRAN symbol for Runge-Kutta flag to signal which pass of Runge-Kutta. Initially set = 1 in main program; thereafter under control of Runge-Kutta subroutine

KT FORTRAN symbol for computing interval counter

L FORTRAN symbol for index used in computing coefficients of Encke's  $f$  series =  $I - 1$

L                   FORTRAN symbol; in subroutine PLOT, an index used to establish for which plot the grid and common labeling are being done

$$L = 1 \quad \Delta C/C$$

$$L = 2 \quad \Delta \rho/a$$

$$L = 3 \quad \Delta \rho a^{1/2}$$

M                   Mean anomaly

$$m_1 \quad \frac{-a_{3/2}(1 + 12e)}{6}$$

M1                  FORTRAN symbol for  $m_1$

MFAIL              FORTRAN symbol for the maximum number of failures allowed in Runge-Kutta subroutine =  $|W_{11}|$

MRKPT              FORTRAN symbol; in subroutine PLOT, an array to identify the plotting symbols to be used for subroutine APLOTV

<u>Term No.</u>	<u>Value</u>	<u>Symbol Indicated</u>	<u>Type of Solution</u>
1	38	o	Analytic Solution
2	55	x	Precessing Ellipse
3	63	□	Numerical Solution of Total Equations of Motion
4	44	*	Initial Osculating Ellipse
5	42	.	Numerical Residual Perturbation Solution

n                   Point number index in explanation of Runge-Kutta method in Section 3.4

NDIV

FORTTRAN symbol for the error code

<u>NDIV</u>	<u>Reason for Halt</u>
1	$a \leq 0$
2	$1 - e^2 \leq 0$
7	$S = 0$
8	$dt/d\psi = 0$

NX

FORTTRAN symbol; in subroutine PLOT, the number of characters in horizontal labels

NY

FORTTRAN symbol; in subroutine PLOT, the number of characters in vertical labels

OMEG1

FORTTRAN symbol for  $\omega_1$

P

Subscript for numerical residual perturbation solution results

$P_{ss}$

Period of mean precessing ellipse

$P_t$

A 700-element array of plot times

$P_{t(ITP)}$

ITPth element of array  $P_t$

$P_1$

$$\frac{1}{1 + \epsilon_2 \omega_1} - \epsilon_1 a^{3/2} (1 + \epsilon_1 m_1)$$

$P_3$

$$\frac{1}{1 + \epsilon_2 \omega_1} - \epsilon_a [2 + \epsilon(v_1 - 2\lambda_1)]$$

$P_\Omega$

Period of node

$P_\omega$

Period of perigee



P1	FORTRAN symbol for $p_1$
P3	FORTRAN symbol for $p_3$
PHI	FORTRAN symbol for $\phi$
pk	Subscript for precessing ellipse values
PNOD	FORTRAN symbol for the period of the node
PPER	FORTRAN symbol for the period of the perigee
P1PS	FORTRAN symbol for $2p_1\psi$ where $0 \leq \psi < 360$ approximately
PSI	FORTRAN symbol for $\psi$
PS(I)	ith guess on $\psi$ in solution of Kepler's equation = array of 25 elements
PSS	FORTRAN symbol for the two-body period of the precessing ellipse which is the reference orbit of the numerical residual perturbation solution
PTIME	FORTRAN symbol for the 700-element array $P_t$
Q	FORTRAN symbol for $q$
q	$= [(x'_s + \frac{\Delta x'}{2})\Delta x' + (y'_s + \frac{\Delta y'}{2})\Delta y' + (z'_s + \frac{\Delta z'}{2})\Delta z'] / \rho_s^2$
r	Distance from larger body
$r_a$	$\sqrt{1 + 2\epsilon^2(x' \cos \gamma + y' \sin \gamma) + \epsilon^2 \rho^2}$ where $\rho'^2 = x'^2 + y'^2 + z'^2$

$r_f$	Distance from larger body, according to fixed ellipse
$r_o$	Initial distance from larger body
$r$	$\sqrt{1 + \epsilon^4 \rho^2 + 2\epsilon^2(x \cos \epsilon t + y \sin \epsilon t)}$
RA	FORTRAN symbol for $r_a$
RADI	FORTRAN symbol for inclination, $i$ (radians)
RCU	FORTRAN symbol for $r^3$
RF	FORTRAN symbol for $r_f$
RHO	FORTRAN symbol for $\rho$
RHOASQ	FORTRAN symbol for $\rho_a^2$
RHOCU	FORTRAN symbol for $\rho^3$
RHOD	FORTRAN symbol for $\dot{\rho} = \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}$
RHOF	FORTRAN symbol for $\rho_f$
RHOFSQ	FORTRAN symbol for $\rho_f^2$
RHOS	FORTRAN symbol for $\rho_s$ = two-body orbit radius
RHOSCU	FORTRAN symbol for $\rho_s^3$
RHOSQ	FORTRAN symbol for $\rho_s^2$
RKSIMP	FORTRAN symbol for the name of Runge-Kutta subroutine

RK1X, RK1Y, RK1Z	FORTTRAN symbol for the Runge-Kutta values of $\Delta t$
RK2X, RK2Y, RK2Z	times the acceleration on successive passes of Runge-
RK3X, RK3Y, RK3Z	Kutta
RNAUT	FORTTRAN symbol for $r_0$
RSQ	FORTTRAN symbol for $r^2$
S	FORTTRAN symbol for $s$
s	Reciprocal of radius vector, $\rho$
$S_E$	$\sin E$
$S_i$	$\sin i$
$S_{P_1}$	$\sin (2p_1 - 1)\psi$
$S_{P_1\psi}$	$\sin (p_1\psi)$
$S_v$	$\sin v$
$S_{\theta_1}$	$\sin \theta_1$
$S_{\phi_s}$	$\sin (v + \omega)$
$S_{\epsilon t}$	$\sin (\epsilon t)$
$S_\psi$	$\sin \psi$
$S_\omega$	$\sin \omega$
$S_\Omega$	$\sin \Omega$

S(10)	FORTRAN symbol for the 10-element array in which values of time, position, velocity, and acceleration are saved for ordinary Runge-Kutta use
SALPHA	FORTRAN symbol for $\sin \alpha_1$
SCLIPS	FORTRAN symbol for $\sin (C_{11}\psi)$
SEC	FORTRAN symbol for $\sin E$
SEPT	FORTRAN symbol for $\sin \gamma = S_{\epsilon t}$
SI	FORTRAN symbol for $\sin i$
SIPER	FORTRAN symbol for $\sin \omega_s$
SIS	FORTRAN symbol for $\sin i_s$
SIV	FORTRAN symbol for $\sin v_s$
SNO	FORTRAN symbol for $\sin \Omega$
SNODE	FORTRAN symbol for $\Omega_s$
SP	FORTRAN symbol for the 10-element array in which values of time, position, velocity, and acceleration are saved when print or plot occurs at other than a regular compute point during the numerical solution of the total equations of motion
SP1	FORTRAN symbol for $\sin (2p_1 - 1)\psi$
SP1PS	FORTRAN symbol for $\sin (p_1\psi + \alpha)$

SP1PST                   FORTRAN symbol for  $\sin(p_1\psi)$

SP3                       FORTRAN symbol for  $\sin(p_3\psi)$

SPDT                     FORTRAN symbol for the saved value of computing interval when print or plot occurs at other than a regular compute point during the numerical solution of the total equations of motion

SPER                     FORTRAN symbol for  $\omega_s$

SPS                       FORTRAN symbol for  $\sin \psi$

SS                       FORTRAN symbol in subroutine RKSIMP for the 17-element array in which values of time, position, velocity, and acceleration are saved for

1. in case computing interval selection fails
2. Simpson's rule integration

$ss_1 = t$	$ss_7 = \dot{z}_h$	$ss_{13} = t_\Omega$
$ss_2 = x_h$	$ss_8 = \ddot{x}_h$	$ss_{14} = p_\omega$
$ss_3 = y_h$	$ss_9 = \ddot{y}_h$	$ss_{15} = \theta_1$
$ss_4 = z_h$	$ss_{10} = \ddot{z}_h$	$ss_{16} = \sin \theta_1$
$ss_5 = \dot{x}_h$	$ss_{11} = t_s$	$ss_{17} = \gamma$
$ss_6 = \dot{y}_h$	$ss_{12} = t_\psi$	

SSNO                     FORTRAN symbol for  $\sin \Omega_s$

STHET1                   FORTRAN symbol for  $\sin \theta_1$ . Old value not removed to avoid unnecessary change in code

STHETA	FORTTRAN symbol for $\sin \theta_1$
STRL	FORTTRAN symbol for $\sin (\nu + \omega_s)$
T	FORTTRAN symbol for time, t
t	Time
$t_{a32}$	$t \cdot a^{3/2}$
$T_{den i}$	ith term in Encke's series  $f = 3[1 - \frac{5}{2} q + \frac{5 \cdot 7}{2 \cdot 3} q^2 - \frac{5 \cdot 7 \cdot 9}{2 \cdot 3 \cdot 4} q^3 + \dots]$
$T_f$	Final time
$t_i$	Integer part of $t_R$
$T_{num i}$	Numerator by which ith term in series for f must be multiplied
$T_{num (i-1)}$	Numerator by which the (i-1)th term in series for f must be multiplied
$t_p$	Print and plot time during numerical residual perturbation solution
$t_{pT}$	Print time during numerical solution of total equations of motion
$t_R$	Time interval plotted on each page
$t_s$	Time, modulo the period of precessing ellipse

$t_0$	Initial time
$T_1$	Period of $t$ expression defined in $t(2\pi) = T_1 = a^{3/2} 2\pi + \frac{15}{4} \epsilon a^3 e \sin [(2p_1 - 1)2\pi] \\ - \epsilon^2 a^{9/2} \frac{11}{8} \sin 4p_1 \pi$
$t_\psi$	Time since last passage of apocenter by analytic solution
$t_\Omega$	Time since the last passage of the node through zero
$t_\omega$	Time since the last time the perigee was zero
TAS32	FORTRAN symbol for $t_s/a_s^{3/2}$
TDEN	FORTRAN symbol for the 40-element array of the denominators by which the terms in series for $f$ must be divided
TERM	FORTRAN symbol for the 40-element array of succeeding terms in series for $f$
TF	FORTRAN symbol for the total flight time; input as a run stop condition
THETA1	FORTRAN symbol for $\theta_1$
THETA2	FORTRAN symbol for $\theta_2$
TN	FORTRAN symbol for "next time" = $t + \Delta t$
TNOD	FORTRAN symbol for $t_\Omega$

TNUM	FORTRAN symbol for the 40-element array of numerators by which the terms in the series for $f$ must be multiplied
TPI	FORTRAN symbol for the print or plot interval
TPER	FORTRAN symbol for $t_{\omega}$
TPI	FORTRAN symbol for the print or plot time
TPRINT	FORTRAN symbol for $t_p$ and $t_{pt}$
TPSI	FORTRAN symbol for $t_{\psi}$
TR	FORTRAN symbol for $t_R$
TRUL	FORTRAN symbol for $v + \omega_s$
TS	FORTRAN symbol for $t_s$
V	FORTRAN symbol for true anomaly in two-body equations, $v$
v	true anomaly, $V$
VSQ	FORTRAN symbol for $\rho^2$
VSSQ	FORTRAN symbol for $V_s^2$ at initial point
W	FORTRAN symbol for the 30-element input array consisting of $W_1$ through $W_{30}$
$W_1$	Input $a$



$W_2$  Input  $e$

$W_3$  Input  $i$  (degrees)

$W_4$  Input  $\epsilon^4 (\equiv \mu)$

$W_5$  Input FTD modifier of  $\Delta t$

$W_6$  Not used

$W_7$  Input run stop time

$W_8$  Input maximum allowable error in Runge-Kutta

$W_9$  Input multiplier on computing interval if error  
 < minimum allowable. Must be > 1. If it is not, it  
 is set = 1.5

$W_{10}$  Input initial computing interval. If zero, it is  
 set = 0.005

$W_{11}$  Input maximum failures permitted by program in  
 selection of computing interval

$W_{12}$  Input minimum allowable error in Runge-Kutta

$W_{13}$  Input flag for secular rates  
 If  $W_{13} \leq 0$ , only initial and final points printed  
 If  $W_{13} > 0$ , approximately every  $1/4$  radian of  $\psi$   
 printed

WR FORTRAN symbol for the 30-element input reference  
 run array

X, Y, Z	FORTRAN symbol for x, y, z
XD, YD, ZD	$\dot{x}, \dot{y}, \dot{z}$
X2D, Y2D, Z2D	$\ddot{x}, \ddot{y}, \ddot{z}$

Total position, velocity, and acceleration components.  
 Coordinate system referenced to plane of planetary motion.

$x'$  Inertial coordinate in plane of massive bodies, from smaller body away from larger

$x'_a$  Inertial coordinate from analytic solution

$x''_a$  Coordinate in orbital plane direction to the node

$x'_f$  Inertial coordinate from fixed ellipse

$x_h, y_h, z_h$   
 $\dot{x}_h, \dot{y}_h, \dot{z}_h$   
 $\ddot{x}_h, \ddot{y}_h, \ddot{z}_h$  In subroutine RKSIMP, dummy variables of position, velocity, and acceleration, respectively

$X_L$  Time for the left edge of each plot (in subroutine PLOT) = EM

$x'_s$  Inertial coordinate from precessing ellipse

$X_T$  Time at right edge of each plot = EN

$x'_o$  Initial inertial coordinate

XD See X, Y, Z, for definition

X2D See X, Y, Z, for definition

XF	FORTRAN symbol for the $x'$ coordinate of position determined from the initial osculating ellipse
XL	FORTRAN symbol in subroutine PLOT for left-most value of abscissa
XNAUT	FORTRAN symbol for $x'_0$
XP, YP, ZP	FORTRAN symbol for $x', y', z'$
XDP, YDP, ZDP	$\dot{x}', \dot{y}', \dot{z}'$
X2DP, Y2DP, Z2DP	$\ddot{x}', \ddot{y}', \ddot{z}'$
	Approximate position, velocity and acceleration components. Coordinate system referenced to plane of planetary motion.
X2P, Y2P, Z2P	FORTRAN symbol for $x'', y'', z''$
XD2P, YD2P, ZD2P	$\dot{x}'', \dot{y}'', \dot{z}''$
X2D2P, Y2D2P, Z2D2P	$\ddot{x}'', \ddot{y}'', \ddot{z}''$
	Approximate position, velocity and acceleration components. Coordinate system referenced to orbit plane of satellite
XS, YS, ZS	FORTRAN symbol for $x'_s, y'_s, z'_s$
XSD, YSD, ZSD	FORTRAN symbol for $\dot{x}'_s, \dot{y}'_s, \dot{z}'_s$
$y'$	Inertial coordinate in plane of massive bodies, perpendicular to $x'$ in direction of motion
$y'_a$	Inertial position component from analytic solution
$y''_a$	Coordinate in satellite orbital plane from analytic solution

$y'_f$	Inertial coordinate from initial osculating ellipse
$y'_s$	Inertial coordinate from precessing ellipse
YACAF	FORTTRAN symbol for $\Delta C_f$ and $\Delta C_f/C_{init}$
YACAP	FORTTRAN symbol for $\Delta C_a$ and $\Delta C_a/C_{init}$
YACAS	FORTTRAN symbol for $\Delta C_s/C_{init}$
YACOBI	FORTTRAN symbol for Jacobi constant on first point $\Delta C_E$ . Thereafter equals difference between instantaneous value and initial value = $\Delta C_E/C_{init}$
YB	FORTTRAN symbol in subroutine PLOT for the bottom value of the scale
YF	FORTTRAN symbol for $y'$ coordinate of position as derived from the initial osculating ellipse
YFD	FORTTRAN symbol for $y'$ coordinate of the velocity derived from the initial osculating ellipse
YT	Top of ordinate in plots
$z'$	Inertial coordinate normal to plane of massive bodies, positive north
$z'_a$	Inertial coordinate from analytic solution
$z''_a$	Coordinate perpendicular to orbital plane from analytic solution
$z'_f$	Inertial coordinate from fixed ellipse

$z'_s$	Inertial coordinate from precessing ellipse
ZF	FORTRAN symbol for $z'$ coordinate of the position as derived from the initial osculating ellipse
ZFD	FORTRAN symbol for $z'$ coordinate of the velocity derived from the initial osculating ellipse
$\alpha_1$	Part of $2p_1 \psi_{tot}$ corresponding to an integral number of revolutions of $\psi = 2n\pi \cdot 2p_1$ , modulo $2\pi$
$\alpha_2$	$\alpha_1 + 2\pi \cdot 2p_1$ , modulo $2\pi = 2(n+1)\pi \cdot 2p_1$ , modulo $2\pi$
$\beta$	Part of $p_3 \cdot \psi_{tot}$ corresponding to an integral number of revolutions of $\psi = 2p_3 n\pi$ , modulo $2\pi$
$\gamma$	$\epsilon t$ , modulo $2\pi$
$\Delta C$	Error in Jacobi constant
$\Delta C_a$	Error in Jacobi constant calculated from analytic solution
$\Delta C_c$	Error in Jacobi constant calculated from numerical solution of total equations of motion
$\Delta C_E$	Error in Jacobi constant calculated from the residual perturbation numerical solution
$\Delta C_F$	Error in Jacobi constant calculated from initial osculating ellipse
$\Delta C_s$	Error in Jacobi constant calculated from precessing ellipse

$(\Delta C/C)_{ai}$	Fractional error of Jacobi constant of analytic solution
$(\Delta C/C)_{ci}$	Fractional error of Jacobi constant for numerical solution of total equations of motion
$(\Delta C/C)_{ki}$	Fractional error of Jacobi constant from initial osculating ellipse
$(\Delta C/C)_{pi}$	Fractional error in Jacobi constant from numerical residual perturbation solution
$(\Delta C/C)_{pki}$	Fractional error in Jacobi constant from precessing Kepler ellipse
$\Delta E$	Difference between successive guesses at $E$ in iterative solution of Kepler's equation
$\Delta E_p$	Difference between successive guesses at $E$ in iterative solution of perturbed Kepler's equation
$\Delta t$	Computing interval
$\Delta t_p$	Time interval for print
$\Delta t_{sp}$	In numerical solution of total differential equations, the saved normal computing interval when a special print point is being computed
$\Delta x', \Delta y', \Delta z'$	Increments in position, velocity, and acceleration obtained numerically from the residual perturbation equations. Coordinate system is referenced to plane of planetary motion

$\Delta \dot{x}, \Delta \dot{y}, \Delta \dot{z}$	In subroutine RKSIMP, the accumulated change in $\dot{x}$ , $\dot{y}$ , and $\dot{z}$ over two computing intervals
$\Delta \rho_a$	Distance between analytic solution and numerical residual perturbation solution
$\dot{\Delta \rho}_a$	Magnitude of vector velocity difference between the analytic solution and the numerical residual perturbation solution
$\Delta \rho_c$	Distance between numerical solution of total differential equations and the numerical residual perturbation solution
$\dot{\Delta \rho}_c$	Magnitude of vector velocity difference between the numerical solution of the total equations of motion and the numerical residual perturbation solution
$\Delta \rho_F$	Distance between initial osculating ellipse solution and the numerical residual perturbation solution
$\dot{\Delta \rho}_F$	Magnitude of vector velocity difference between the initial osculating ellipse solution and the numerical residual perturbation solution
$\Delta \rho_S$	Distance between the precessing ellipse solution and the numerical residual perturbation solution
$\dot{\Delta \rho}_S$	Magnitude of vector velocity difference between the precessing ellipse and the numerical residual perturbation solution
$\epsilon$	$\mu^{1/4}$

$\epsilon_a$	$\epsilon \cdot a^{3/2} = \epsilon a_{3/2}$
$\epsilon_1$	$\epsilon = \mu^{1/4}$
$\epsilon_2$	$\epsilon^2 = \mu^{1/2}$
$\epsilon_4$	$\epsilon^4 = \mu$
$\epsilon_4$	Mass ratio of sun to sun + planet = $\mu$
$\theta_1$	Part of $(2p_1-1) \psi_{tot}$ resulting from integer $n$ rotations of $\psi = 2\pi n (2p_1-1)$ , modulo $2\pi$
$\theta_2$	$\theta_1 + 2\pi (2p_1-1)$ , modulo $2\pi = 2(n+1)\pi(2p_1-1)$ , modulo $2\pi$
$\lambda_1$	Angular rate of change of $\Omega$
$\mu$	Mass ratio of smaller body to sum of smaller and larger bodies
$v_1$	$a^{3/2} (11 + 24e)/12 = \dot{\phi}/\epsilon^2$
$\rho$	Radial distance from smaller body
$\dot{\rho}$	Total velocity in blown up coordinates
$\rho_a$	Radial distance obtained from analytic solution
$\rho_s$	Radial distance obtained from precessing ellipse
$\phi$	Angle from node to satellite



$\phi_s$	Angular distance of the satellite from the node as obtained from precessing ellipse; $\phi_s = v + \omega$
$\phi_1$	Integer part of $\phi$ modulo $2\pi = \left[ \frac{2n\pi}{1+\epsilon^2 \omega_1} - \epsilon^2 v_1 t(2\pi n) \right]$ , modulo $2\pi$
$\psi$	True anomaly measured from apogee
$\Omega$	Argument of the node in analytic solution
$\omega$	Argument of pericenter
$\omega_1$	$\frac{a^3(24e - 7)}{12}$

### Section 3

#### DEVELOPMENT OF EQUATIONS

#### 3.0 GENERAL

The purpose of this section is to describe the sources of the equations used in the program and to outline the derivation of the exact forms programmed. When an equation from a reference is first introduced, the equation number in the source is shown to the left of the equation. The equation numbers of the present document are to the right of the equations. The equations are not derived in the exact order they are used; however, the numbers of the subsections in the Equations in Order of Solution, Section 4, where the equation is used, are stated after each derivation. Section 4 presents the equation numbers of this section so that the cross reference is complete.

#### 3.1 EQUATIONS OF MOTION

The motion considered is that of a satellite of negligible mass in close proximity to the smaller of the two massive bodies in the restricted three-body problem, which may serve as a simplified model for an actual case like a lunar satellite in the earth-moon system.

The coordinate system used is centered at the smaller mass and rotates with the constant angular velocity of the two massive bodies around their center of mass. With the notation as in reference 1, the  $x^*$ ,  $y^*$ , and  $z^*$  rectangular system is oriented such that the  $x^*$ ,  $y^*$  plane coincides with the plane of motion of the two massive bodies, the  $x^*$ -axis pointing away from the larger body, the  $y^*$ -axis  $90^\circ$  ahead in the direction of motion, and  $z^*$ -axis parallel to the rotation axis as to complete a right hand system.

In order to nondimensionalize the equations of motion, the following scales are used:

1. Unit of length = Distance between the two massive bodies
2. Unit of time = Period of rotation of the  $x^*$ ,  $y^*$ ,  $z^*$  system in inertial space, divided by  $2\pi$

3. Unit of mass = Total mass of the system.

The equations of motion of the satellite in these coordinates are given in reference 1 by:

$$(1) \frac{d^2 \vec{x}^*}{dt^{*2}} = \text{grad}^* \left[ \frac{1-\mu}{|\vec{x}^* + \hat{i}|} + \frac{\mu}{|\vec{x}^*|} \right] - 2\hat{k} \times \frac{d\vec{x}^*}{dt^*} - \hat{k} \times \{ \hat{k} \times [\vec{x}^* + (1-\mu)\hat{i}] \} \quad (1)$$

where  $\vec{x}^* = (x^*, y^*, z^*)$  and  $\hat{i}$ ,  $\hat{j}$ , and  $\hat{k}$  are the unit vectors in the  $x^*$ ,  $y^*$ , and  $z^*$  directions.

Since the motion of the satellite takes place in a close vicinity of the smaller body, the following "blown-up" coordinates are introduced:

$$x = \frac{x^*}{\mu^{1/2}}, \quad y = \frac{y^*}{\mu^{1/2}}, \quad z = \frac{z^*}{\mu^{1/2}}, \quad t = \frac{t^*}{\mu^{1/4}} \quad (2)$$

(Reference 1, page 205, equation above (6), there in vector notation) where  $\mu$  is the mass ratio of the smaller to the total mass.

In order to avoid terms of order  $\mu^{1/4} = \epsilon$  in the equations of motion, another transformation is made to inertially oriented coordinates by multiplying the vector  $\vec{x} = (x, y, z)$  by the matrix  $\{R\}$  with:

$$(7c) \quad \{R\} = \begin{bmatrix} \cos \epsilon t & -\sin \epsilon t & 0 \\ \sin \epsilon t & \cos \epsilon t & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (3)$$

[Reference 1, page 206, equation (7c)]

then

$$(7b) \quad \{\vec{x}\}' = \{R\}\{\vec{x}\} \quad (4)$$

The primed system is therefore inertially oriented, but centered at the smaller body.

When the transformations (2) and (4) are performed on the equations of motion (1), the following equations result:

$$\frac{d^2 \vec{x}'}{dt^2} = \frac{\vec{x}'}{\rho^3} - \frac{(1 - \epsilon^4) \epsilon^2 \vec{x}'}{r^3} + (1 - \epsilon^4) \left(1 - \frac{1}{r^3}\right) \begin{bmatrix} \cos \epsilon t \\ \sin \epsilon t \\ 0 \end{bmatrix} \quad (5)$$

where  $\vec{x}' = (x', y', z')$  (6)

$$\rho^2 = x'^2 + y'^2 + z'^2 \quad (7)$$

$$r^2 = 1 + \epsilon^4 \rho^2 + 2\epsilon^2 (x' \cos \epsilon t + y' \sin \epsilon t) \quad (8)$$

$$\epsilon = \mu^{1/4} \quad (9)$$

Equations (5), (6), and (7) are used in Section 4.25. Equations (8) and (9) are used in Sections 4.25, 4.19.2, 4.16.2, and 4.8.3.

Geometrically,  $\rho$  is the distance of the satellite from the smaller body and  $r$  the distance between satellite and larger body, both in blown-up coordinates.

### 3.2 THE JACOBI CONSTANT

In order to check the accuracy of the solutions, the Jacobian integral will be used. From reference 2, page 281, equation (7), one obtains:

$$\left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2 + \left(\frac{dz}{dt}\right)^2 = x^2 + y^2 + \frac{2(1 - \mu)}{r_1} + \frac{2\mu}{r_2} + C \quad (10)$$

This is in rotating coordinates, centered at the center of mass;  $r_1$  and  $r_2$  are the distances of the satellite from the larger and the smaller primary, respectively. To transform equation (10) to blown-up coordinates, centered at the smaller body, the above variables are to be replaced as follows:

$$\begin{aligned}
 x &\rightarrow 1 - \mu + \epsilon^2 x & \frac{dx}{dt} &\rightarrow \epsilon \dot{x} & r_1 &\rightarrow r \\
 y &\rightarrow \epsilon^2 y & \frac{dy}{dt} &\rightarrow \epsilon \dot{y} & r_2 &\rightarrow \epsilon^2 \rho \\
 z &\rightarrow \epsilon^2 z & \frac{dz}{dt} &\rightarrow \epsilon \dot{z} & \mu &\rightarrow \epsilon^4
 \end{aligned} \tag{11}$$

Then after applying the transformation (4), the Jacobi integral takes on the form:

$$\begin{aligned}
 \dot{x}'^2 + \dot{y}'^2 + \dot{z}'^2 + 2\epsilon(\dot{x}'y' - \dot{y}'x') - 2(1 - \epsilon^4)(x' \cos \epsilon t + y' \sin \epsilon t) \\
 - \frac{2}{\rho} - \frac{2(1 - \epsilon^4)}{\epsilon^2 r} + \frac{(1 - \epsilon^4)^2}{\epsilon^2} = C'
 \end{aligned} \tag{12}$$

The last terms on the left side are numerically large. To avoid the loss of significant figures in the computation of the error in the Jacobi constant, the initial value of the Jacobi constant is subtracted in the following way:

$$\begin{aligned}
 \dot{x}'^2 + \dot{y}'^2 + \dot{z}'^2 + 2\epsilon(\dot{x}'y' - \dot{y}'x') - 2(1 - \epsilon^4)(x' \cos \epsilon t + y' \sin \epsilon t) \\
 - \frac{2}{\rho} - \frac{2(1 - \epsilon^4)}{\epsilon^2 r} + \frac{(1 - \epsilon^4)^2}{\epsilon^2} - C_{\text{init}} = \Delta C
 \end{aligned} \tag{13a}$$

where  $C_{\text{init}} = -\dot{x}'^2(0) - \dot{y}'^2(0) - \dot{z}'^2(0) - 2\epsilon[\dot{x}'(0)y'(0) - \dot{y}'(0)x'(0)]$

$$+ 2(1 - \epsilon^4)x'(0) + \frac{2}{\rho(0)} + \frac{2(1 - \epsilon^4)}{\epsilon^2 r(0)} - \frac{(1 - \epsilon^4)^2}{\epsilon^2} \tag{13b}$$

The quantity  $\Delta C$  is zero if the solution is exact. For an approximate solution,  $\Delta C$  serves as a test quantity on the accuracy of the approximation. Define:

$$C_{\text{part}} = +\dot{x}'^2(0) + \dot{y}'^2(0) + \dot{z}'^2(0) + 2\epsilon[\dot{x}'(0)y'(0) - \dot{y}'(0)x'(0)] - 2(1 - \epsilon^4)x'(0) + \frac{2}{\rho(0)} \quad (14)$$

Using equation (8) and noting that with apocenter as the initial condition, then the difference of the  $\frac{1}{\epsilon^2}$  terms of (13a) is written as:

$$\begin{aligned} -\frac{2(1 - \epsilon^4)}{\epsilon^2 r} + \frac{2(1 - \epsilon^4)}{\epsilon^2 r_0} &= \frac{2(1 - \epsilon^4)}{\epsilon^2} \frac{(r^2 - r_0^2)}{rr_0(r + r_0)} \\ &= 2 \frac{1 - \epsilon^4}{\epsilon^2} \frac{1 + \epsilon^4 \rho^2 + 2\epsilon^2(x' \cos \epsilon t + y' \sin \epsilon t) - 1 - \epsilon^4 x_0'^2 - 2\epsilon^2 x'}{rr_0(r + r_0)} \\ &= 2(1 - \epsilon^4) \frac{2(x' \cos \epsilon t + y' \sin \epsilon t - x') + \epsilon^2(\rho^2 - x'^2)}{rr_0(r + r_0)} \quad (15) \end{aligned}$$

where  $r(0) = 1 + \epsilon^2 x_0' = r_0$  (\*)

Then the error in the Jacobi constant becomes:

$$\begin{aligned} \Delta C' &= (\dot{x}'^2 + \dot{y}'^2 + \dot{z}'^2) + 2[\epsilon(\dot{x}'y' - \dot{y}'x')] - \frac{1}{\rho} \\ &\quad - (1 - \epsilon^4)(x' \cos \epsilon t + y' \sin \epsilon t) - C_{\text{part}} \\ &\quad + 2(1 - \epsilon^4) \frac{[2(x' \cos \epsilon t + y' \sin \epsilon t - x_0) + \epsilon^2(\rho^2 - x_0^2)]}{rr_0(r + r_0)} \quad (16) \end{aligned}$$

Equation (16) is used in Sections 4.17.15, 4.28.1, and 4.28.2.

### 3.3 SOLUTIONS

In the following, three different methods for the solution of equation (5) will be outlined:

(\*) See equations (8), (42), and (61)

1. An approximate, analytic solution as given in reference 1.
2. A numerical residual perturbation method which uses the Keplerian part plus the secular perturbations to order  $\epsilon^2$  of the analytic solution as a reference orbit.
3. Direct application of numerical integration to the total exact equations of motion.

### 3.3.1 Analytic Solution

The developments as given in reference 1 start from an approximate set of equations that result from equation (5) when terms having  $\epsilon^3$  or a higher power of  $\epsilon$  as a factor are neglected.

The approximate equations of motion are:

$$\begin{aligned}
 \frac{d^2 x'}{dt^2} &= -\frac{x'}{\rho^3} + \frac{\epsilon^2 x'}{2} + \frac{3}{2} \epsilon^2 (x' \cos 2\epsilon t + y' \sin 2\epsilon t) \\
 (8a) \quad \frac{d^2 y'}{dt^2} &= -\frac{y'}{\rho^3} + \frac{\epsilon^2 y'}{2} + \frac{3}{2} \epsilon^2 (x' \sin 2\epsilon t - y' \cos 2\epsilon t) \quad (17) \\
 \frac{d^2 z'}{dt^2} &= -\frac{z'}{\rho^3} - \epsilon^2 z'
 \end{aligned}$$

These equations are identical with the vector equation (8a) of reference 1.

The solution of equation (17) is given in reference 1 in the following form:

$$(18a) \quad s(\psi) = s_0(\psi) + \epsilon^2 s_1(\psi, \epsilon) + O(\epsilon^3)$$

$$(18b) \quad t(\psi) = t_0(\psi) + \epsilon^2 t_1(\psi, \epsilon) + O(\epsilon^3) \quad (18)$$

$$(18c) \quad z_a''(\psi) = z_0(\psi) + \epsilon^2 z_1(\psi, \epsilon) + o(\epsilon^3)$$

where

$$(25a) \quad s_0(\psi) = \frac{1 - e \cos \psi}{a(1 - e^2)} \quad (19)$$

$$(25b) \quad t_0(\psi) = a^{3/2} \left[ \frac{e(1 - e^2)^{1/2} \sin \psi}{1 - e \cos \psi} + \cos^{-1} \frac{\cos \psi - e}{1 - e \cos \psi} \right] \quad (20)$$

$$z_0(\psi) = 0 \quad (21)$$

$$(38) \quad s_1/a^2 = \frac{2}{3} (1 + 5e) - \frac{1}{3} [5 + (95e/8)] \cos \psi + \cos 2p_1 \psi$$

$$+ (5/8)e \cos (2p_1 + 1)\psi - 15Qe \sin p_1 \psi$$

$$\times \sin (p_1 - 1)\psi + o(e^2)$$

$$(39) \quad t_1/a^{9/2} = [(10/3) + (83e/12)] \sin \psi - (11/8) \sin 2p_1 \psi$$

$$+ (3e/4) \sin (2p_1 - 1)\psi - 2e \sin (2p_1 + 1)\psi$$

$$+ (5e/2) \sin 2\psi - 30Qe \cos p_1 \psi$$

$$\times \sin (p_1 - 1)\psi + o(e^2)$$

$$(42) \quad z_1 = (3ia^{4/64}) [\epsilon a^{3/2}]^{-1} \cos \frac{1}{2} (p_3 + 1)\psi \sin \frac{1}{2} (p_3 - 1)\psi$$

$$+ \sin \frac{1}{2} (p_3 + 1)\psi \cos \frac{1}{2} (p_3 - 1)\psi]$$

$$Q = - \left[ \frac{1}{4\epsilon a^{3/2}} + \frac{67}{80} + o(\epsilon) \right]$$



Equation (19) is used in Section 4.17.3.

The geometric interpretation of the variables in equation (18) is as follows:

$s$  = reciprocal radius vector from the smaller body to the satellite, in the blown-up scale.

$t$  = time, in the blown-up scale.

$z''_a$  = distance of the satellite from a plane that passes through the smaller body and is oriented with an inclination  $i$  and a node  $\Omega$  against the  $x', y'$  plane. The node is varying with time according to:

$$(12) \quad \Omega = \Omega_0 + \epsilon^2 \lambda_1 t \quad (39)$$

In order to restrict  $\Omega$  to values smaller than  $2\pi$ , equation (39) is written as:

$$\Omega = \Omega_0 + \epsilon^2 \lambda_1 T_\Omega \quad (39a)$$

used in Section 4.17.5, where  $T_\Omega$  is reduced in Section 4.17.1 by:

$$P_\Omega = \frac{2\pi}{\epsilon^2 \lambda} \quad (39b)$$

whenever it exceeds  $P_\Omega$ .  $\Omega_0$  is zero by the special choice of initial conditions in reference 1.

The constants  $a$ ,  $e$ , and  $i$  are the semi-major axis, eccentricity, and inclination of a Keplerian ellipse represented by the leading terms  $s_0$ ,  $t_0$ , and  $z_0$ . The angular variable,  $\psi$ , represents the true anomaly, but counted from the apocenter. All the terms of order  $\epsilon^2 e$  were dropped from  $s_1$ ,  $t_1$ ,

and  $z_1$  in this program for consistency since  $\epsilon$ ,  $e$ , and  $i$  were all considered small variables in reference 1 and the paper omitted  $\epsilon^3$  and  $\epsilon ei$  terms. Therefore, the equations for the above quantities become:

$$s_1(\psi, \epsilon) = a^2 \left[ \frac{2}{3} - \frac{5}{3} \cos \psi + \cos 2p_1 \psi \right. \\ \left. + \frac{15}{8} \frac{e}{\epsilon a^{3/2}} (\cos \psi - \cos (2p_1 - 1)\psi) \right] \quad (22)$$

$$t_1(\psi, \epsilon) = a^{9/2} \left[ \frac{10}{3} \sin \psi - \frac{11}{8} \sin 2p_1 \psi \right. \\ \left. - \frac{15}{4} \frac{e}{\epsilon a^{3/2}} (\sin \psi - \sin (2p_1 - 1)\psi) \right] \quad (23)$$

$$z_1(\psi, \epsilon) = \frac{3}{8} i \frac{a^{5/2}}{\epsilon(1 - \epsilon a^{3/2})} [\sin p_3 \psi - p_3 \sin \psi] \quad (24)*$$

Equation (22) is used in Section 4.17.3; equation (24) in Section 4.17.4.  
The eccentric anomaly:

$$E = \cos^{-1} \left( \frac{\cos \psi - e}{1 - e \cos \psi} \right) = \sin^{-1} \left( \frac{\sqrt{1 - e^2} \sin \psi}{1 - e \cos \psi} \right) \quad (25)$$

is defined such that

$$2n\pi \leq E \leq 2(n+1)\pi$$

for

$$2n\pi \leq \psi \leq 2(n+1)\pi$$

The time  $t$  and the eccentric anomaly are therefore single valued and increasing function of  $\psi$ . However, in order to avoid large arguments of

\*Equation (24) was taken from equation (57) of reference 7 instead of being taken from equation (42) of reference 1.

the trigonometric functions for large times, a time,  $t_\psi$ , which is the time since the last passage of the apocenter, is introduced. If the total  $\psi$  and the total  $E$  are denoted by  $\psi_{\text{tot}}$  and  $E_{\text{tot}}$ , and the symbols  $\psi$  and  $E$  are used for the principal values only, such that:

$$\psi_{\text{tot}} = 2\pi n + \psi$$

$$E_{\text{tot}} = 2\pi n + E$$

Then the equation for the time is written [c.f. equations (20), (23), and (25)]:

$$t(\psi_{\text{tot}}) = a^{3/2} [E_{\text{tot}} + e \sin E_{\text{tot}}] - \frac{15}{4} \epsilon a^3 e [\sin \psi_{\text{tot}} - \sin (2p_1 - 1)\psi_{\text{tot}}] + \epsilon^2 a^{9/2} \left[ \frac{10}{3} \sin \psi_{\text{tot}} - \frac{11}{8} \sin 2p_1 \psi_{\text{tot}} \right] \quad (26)$$

The time after  $n$  revolutions is:

$$t(2n\pi) = a^{3/2} \cdot 2n\pi + \frac{15}{4} \epsilon a^3 e \sin [(2p_1 - 1) \cdot 2n\pi] - \epsilon^2 a^{9/2} \cdot \frac{11}{8} \sin (2p_1 \cdot 2n\pi) \quad (27)$$

Then, by subtraction of equation (27) from equation (26) and rearranging:

$$t(\psi_{\text{tot}}) - t(2n\pi) - t_\psi = a^{3/2} [E_{\text{tot}} - 2\pi n + e \sin E_{\text{tot}}] - \frac{15}{4} \epsilon a^3 e \{ \sin \psi_{\text{tot}} - \sin [(2p_1 - 1)(\psi_{\text{tot}} - 2\pi n) + 2\pi n(2p_1 - 1)] + \sin [(2p_1 - 1) \cdot 2n\pi] \} + \epsilon^2 a^{9/2} \left\{ \frac{10}{3} \sin \psi_{\text{tot}} - \frac{11}{8} \sin [2p_1 (\psi_{\text{tot}} - 2\pi n)] \right\}$$

$$+ 4p_1 n] + \frac{11}{8} \sin [4p_1 n\pi]$$

or

$$\begin{aligned} t_\psi = a^{3/2} [E + e \sin E] - \frac{15}{4} \epsilon a^3 e \{ \sin \psi - \sin (2p_1 - 1)\psi \cos \theta_1 \\ + [1 - \cos (2p_1 - 1)\psi] \sin \theta_1 + \epsilon^2 a^{9/2} \left\{ \frac{10}{3} \sin \psi \right. \\ \left. - \frac{11}{8} \sin 2p_1 \psi \cos \alpha_1 + \frac{11}{8} (1 - \cos 2p_1 \psi) \sin \alpha_1 \right\} \end{aligned} \quad (28)$$

where

$$\theta_1 = [2n\pi(2p_1 - 1)] \text{ modulo } 2\pi$$

and

$$\alpha_1 = [2n\pi p_1] \text{ modulo } 2\pi$$

When  $n$  becomes large, forming the angle and then computing the modulo with respect to  $2\pi$  does not improve the accuracy. It is necessary to form the  $\theta_1$  and  $\alpha_1$  without developing large angles. For this purpose, initially  $\theta_1$  and  $\alpha_1$  are set to zero. The period of the first revolution  $T_1$  is computed by substituting  $\psi = 2\pi$  in equation (26) yielding:

$$t(2\pi) = T_1 = a^{3/2} 2\pi + \frac{15}{4} \epsilon a^3 e \sin [(2p_1 - 1)2\pi] - \epsilon^2 a^{9/2} \frac{11}{8} \sin 4p_1 \pi \quad (28a)$$

When  $t_\psi$  exceeds  $T_1$ , we set:

$$t_\psi = t_\psi - T_1 \quad (28b)$$

$$\theta_1 = [\theta_1 + 2\pi(2p_1 - 1)] \text{ modulo } 2\pi \quad (28c)$$

$$\alpha_1 = \{[(\alpha_1 + 2\pi p_1) \text{ modulo } 2\pi] + 2\pi p_1\} \text{ modulo } 2\pi \quad (28d)$$

The period for the next revolution is computed by applying equation (26) to obtain:

$$T_1 = t(2(n + 1)\pi) - t(2n\pi)$$

or

$$T_1 = 2\pi a^{3/2} + \frac{15}{4} \epsilon a^3 e \{\sin \theta_2 - \sin \theta_1\} - \frac{11}{8} \epsilon^2 a^{9/2} \{\sin \alpha_2 - \sin \alpha_1\} \quad (28e)$$

where  $\theta_2$  is computed from the  $\theta_1$  just obtained by equation (28c) by:

$$\theta_2 = [\theta_1 + 2\pi(2p_1 - 1)] \text{ modulo } 2\pi \quad (28f)$$

and  $\alpha_2$  from the  $\alpha_1$  of equation (28d) by:

$$\alpha_2 = \{[(\alpha_1 + 2\pi p_1) \text{ modulo } 2\pi] + 2\pi p_1\} \text{ modulo } 2\pi \quad (28g)$$

Similarly the argument  $p_3\psi$  in  $z_a''$  [equation (24)] is taken as  $p_3\psi + \beta$ .  $\beta$  is set equal to zero initially and incremented each time  $t_\psi > T_1$  by:

$$\beta = (\beta + 2p_3\pi) \text{ modulo } 2\pi \quad (29)$$

Equations 28b through 29 are used in Section 4.17.1.2.

Now expression (28) may be solved for  $\psi$  when  $n$  and  $t_\psi$  are known. The resulting value is between 0 and  $2\pi$ . It is solved by Newton's method using  $M = \psi$  as a first guess in Section 4.17.2.4 with  $E$ , defined by equation (20), computed in Section 4.17.2.3.

The constants  $p_1$  and  $p_3$  are given by equations (27a) and (27c) of reference 1. However, these quantities can be obtained with more accuracy from (26a) and (26c) of reference 1,

by requiring:

$$\epsilon^2 = \sin 2 (\bar{\phi} - qt) = \epsilon^2 \sin 2 (p_1 \psi) + 0 (\epsilon^3)$$

$$\epsilon^2 \cos 2 (\bar{\phi} - qt) = \epsilon^2 \cos 2 (p_1 \psi) + 0 (\epsilon^3) \quad (29a)$$

$$\epsilon^2 \sin(\bar{\phi} - \ell t) = \epsilon^2 \sin (p_3 \psi) + 0 (\epsilon^3) \quad (29b)$$

$$\bar{\phi} - qt = p_1 \psi \quad (30)$$

$$\bar{\phi} - \ell t = p_3 \psi \quad (31)$$

where

$$(16a) \quad q = \epsilon[1 + \epsilon m_1] \quad (32)$$

$$(16b) \quad \ell = \epsilon[2 + \epsilon(v_1 - 2\lambda_1)] \quad (33)$$

$$(16c) \quad m_1 = v_1 - \lambda_1 \quad (34)$$

$$\text{past (42)} \quad v_1 = -\frac{1}{12} a^{3/2}(11 + 24e) \quad (35)$$

$$\text{past (41)} \quad \lambda_1 = -\frac{3}{4} a^{3/2} \quad (36)$$

With the relations between  $\psi$  and  $\bar{\phi}$  given in equation (40) below and observing only the secular part of the relation between  $t$  and  $\psi$  in

equations (20) and (23) (see also equation (71) below), which gives  $t = a^{3/2}$ , the following values for  $p_1$  and  $p_3$  are derived without expansion:

$$p_1 = \frac{1}{1 + \epsilon^2 \omega_1} - \epsilon a^{3/2} (1 + \epsilon m_1) \quad (37)$$

$$p_3 = \frac{1}{1 + \epsilon^2 \omega_1} - \epsilon a^{3/2} [2 + \epsilon (v_1 - 2\lambda_1)] \quad (38)$$

Equations (34) through (38) are used in Section 4.4. The relation between  $s$ ,  $t$ ,  $z''$  and the primed coordinates is given by the following transformations of variables:

$$(17) \quad \psi = (1 + \epsilon^2 \omega_1) \bar{\phi} \quad (40)$$

$$(14) \quad \bar{\phi} = \phi + \epsilon^2 v_1 t \quad (41)$$

$$(13a) \quad s = \frac{1}{\rho} \quad (42)$$

$$(13b) \quad x_a'' = \rho \cos \phi \quad (43)$$

$$(13c) \quad y_a'' = \rho \sin \phi \quad (44)$$

$$(11a-11c) \quad \begin{bmatrix} x_a'' \\ y_a'' \\ z_a'' \end{bmatrix} = \begin{bmatrix} \cos \Omega & \sin \Omega & 0 \\ -\sin \Omega \cos i & \cos \Omega \cos i & \sin i \\ \sin \Omega \sin i & -\cos \Omega \sin i & \cos i \end{bmatrix} \begin{bmatrix} x_a' \\ y_a' \\ z_a' \end{bmatrix} \quad (45)$$

$$\text{past (32)} \quad \omega = \omega_1 a^{-3} \quad (46a)$$

$$(36) \quad \omega = (-7 + 24e)/12 \quad (46b)$$

The initial conditions for any of the variables follow simply by putting  $\psi = 0$  for  $t = 0$ , which means that the motion starts at apocenter, which, in turn, coincides with the node for  $\Omega_0 = 0$ .

Equations (40) and (41) may be combined to give:

$$\phi = \frac{\psi}{1 + \epsilon^2 \omega_1} - \epsilon^2 v_1 t \psi \quad (47)$$

which is used in Section 4.17.1.2.

To avoid large arguments of trigonometric functions in equations (43) and (44), these equations are written as:

$$x_a'' = \frac{1}{s} \cos (\phi + \phi_1) \quad (48)$$

$$y_a'' = \frac{1}{s} \sin (\phi + \phi_1)$$

where  $\phi_1$  is the value of  $\phi$  at  $n$  revolutions, modulo  $2\pi$ :

$$\phi_1 = \left( \frac{2n}{1 + \epsilon^2 \omega_1} - \epsilon^2 v_1 t(2n\pi) \right) \text{ modulo } 2\pi \quad (49)$$

Equation (48) is used in Section 4.17.4. As in the case of  $\alpha_1$ ,  $\theta_1$ , and  $\beta$ ,  $\phi_1$  is formed in Section 4.17.1.2 at appropriate times by adding increments to the initial value  $\phi_1 = 0$  formed by evaluating equation (49) for  $n = 1$ . This avoids the angle ever becoming large, which would result in loss of accuracy.

Solving the linear system (45) for  $x'$ ,  $y'$ , and  $z'$  yields:



$$\left. \begin{aligned}
 x'_a &= x''_a \cos \Omega - y''_a \sin \Omega \cos i + z''_a \sin \Omega \sin i \\
 y'_a &= x''_a \sin \Omega + y''_a \cos \Omega \cos i - z''_a \cos \Omega \sin i \\
 z'_a &= y''_a \sin i + z''_a \cos i
 \end{aligned} \right\} \quad (50)$$

which are used in Section 4.17.6.

Differentiation of equations (18) through (24) with respect to  $\psi$  yields velocity components:

$$\begin{aligned}
 \frac{ds}{d\psi} &= \frac{e \sin \psi}{a(1 - e^2)} + \epsilon^2 a^2 \left( \frac{5}{3} \sin \psi - 2 \sin 2p_1 \psi \right. \\
 &\quad \left. - \frac{15}{8} \epsilon a^{1/2} e [\sin \psi - \sin (2p_1 - 1)\psi] \right) \quad (51)
 \end{aligned}$$

$$\begin{aligned}
 \frac{dt}{d\psi} &= \frac{a^{3/2} (1 - e^2)^{3/2}}{(1 - e \cos \psi)^2} - \frac{15}{4} \epsilon a^3 [\cos \psi - \cos (2p_1 - 1)\psi \cos \theta \\
 &\quad + \sin (2p_1 - 1)\psi \sin \theta] + \epsilon^2 a^{9/2} \left[ \frac{10}{3} \cos \psi - \frac{11}{4} \cos 2p_1 \psi \cos \alpha \right. \\
 &\quad \left. + \frac{11}{4} \sin 2p_1 \psi \sin \alpha \right] \quad (52)
 \end{aligned}$$

$$\frac{dz''_a}{d\psi} = \frac{3}{8} \frac{i \epsilon a^{5/2}}{(1 - \epsilon a)^{3/2}} [\cos p_3 \psi - \cos \psi] \quad (53)$$

Terms of order  $\epsilon^3$ ,  $\epsilon^2 e$ ,  $\epsilon^3 i$ , and  $\epsilon e i$ , are neglected in equation (51) through (53). Equation (51) is used in Section 4.17.9 and equation (52) is used in Sections 4.17.7 and 4.17.2.5. Differentiation of (48) with respect to  $t$  yields

$$\begin{aligned}
 \dot{x}''_a &= -\frac{1}{s} \frac{ds}{dt} \cos (\phi - \phi_1) - \frac{1}{s} \frac{d\phi}{dt} \sin (\phi - \phi_1) \quad (53a) \\
 \dot{y}''_a &= -\frac{1}{s} \frac{ds}{dt} \sin (\phi - \phi_1) + \frac{1}{s} \frac{d\phi}{dt} \cos (\phi - \phi_1)
 \end{aligned}$$

or

$$\dot{x}_a'' = -y_a'' \frac{d\phi}{dt} - \frac{x_a''}{s} \frac{ds}{dt} \quad (54a)$$

$$\dot{y}_a'' = +x_a'' \frac{d\phi}{dt} - \frac{y_a''}{s} \frac{ds}{dt} \quad (54b)$$

and from equation (53):

$$\dot{z}_a'' = \frac{3}{8} \frac{i\epsilon a^{5/2}}{(1-\epsilon a^{3/2})} [\cos p_3 \psi - \cos \psi] \frac{d\psi}{dt} \quad (54c)$$

where, from equation (47):

$$\frac{d\phi}{dt} = \frac{1}{1 + \epsilon^2 \omega_1} \frac{d\psi}{dt} - \epsilon^2 v_1 \quad (55)$$

and

$$\frac{d\psi}{dt} = \frac{1}{\frac{ds}{dt}} \cdot \frac{ds}{d\psi} \cdot \frac{d\psi}{dt} \quad (56)$$

Equations (54a), (54b) are used in Section 4.17.10. Equations (55) and (56) are used in Section 4.17.9.

Differentiation of (50) with respect to time yields the velocity components in the inertial frame:

$$\left. \begin{aligned} \dot{x}_a' &= (-x_a'' \sin \Omega - y_a'' \cos \Omega \cos i + z_a'' \cos \Omega \sin i) \dot{\Omega} \\ &\quad + \dot{x}_a'' \cos \Omega - \dot{y}_a'' \sin \Omega \cos i + \dot{z}_a'' \sin \Omega \sin i \\ \dot{y}_a' &= (x_a'' \cos \Omega - y_a'' \sin \Omega \cos i + z_a'' \sin \Omega \sin i) \dot{\Omega} \\ &\quad + \dot{x}_a'' \sin \Omega + \dot{y}_a'' \cos \Omega \cos i - \dot{z}_a'' \cos \Omega \sin i \\ \dot{z}_a' &= \dot{y}_a'' \sin i + \dot{z}_a'' \cos i \end{aligned} \right\} \quad (57)$$

which are computed in Section 4.17.12.

### 3.3.1.1 Initial Conditions for the Analytic Solution

The initial conditions are obtained by setting  $\psi = t = \Omega = 0$ . From equations (19) through (24):

$$\begin{aligned} s_0(0) &= \frac{1}{a(1+e)} & s(0) &= 0 \\ t_0(0) &= 0 & t_1(0) &= 0 \\ z_0(0) &= 0 & z_1(0) &= 0 \end{aligned} \tag{58}$$

From equations (51) through (53):

$$\frac{ds(0)}{d\psi} = 0 \tag{59a}$$

$$\frac{dt(0)}{d\psi} = \frac{a^{3/2}(1+e)^{3/2}}{(1-e^2)^{1/2}} + \frac{7}{12} \epsilon^2 a^{9/2} \tag{59b}$$

$$\frac{dz_a''(0)}{d\psi} = 0 \tag{59c}$$

Equation (59b) is used in Sections 4.8 and 4.19.2.

From equations (47), (48), and (98):

$$\phi(0) = 0 \tag{60}$$

$$x_a''(0) = \frac{1}{s(0)} = a(1+e) \tag{61}$$

$$y_a''(0) = 0 \tag{62}$$

$$z_a''(0) = 0 \quad (62a)$$

By (58) equations (61) and (62) are used in Section 4.8.3.

From equations (50) and (61):

$$\left. \begin{aligned} x_a'(0) &= x_a''(0) = a(1 + e) \\ y_a'(0) &= 0 \\ z_a'(0) &= 0 \end{aligned} \right\} \quad (63)$$

which are used in Section 4.8.3.

From (54a), (54b), (54c), (59a), (59c), and (62):

$$\left. \begin{aligned} \dot{x}_a''(0) &= 0 \\ \dot{y}_a''(0) &= x_a''(0) \cdot \frac{d\phi(0)}{dt} \\ \dot{z}_a''(0) &= 0 \end{aligned} \right\} \quad (64)$$

where by equation (55):

$$\frac{d\phi(0)}{dt} = [(1 + \epsilon^2 \omega_1) \frac{dt(0)}{d\psi}]^{-1} - \epsilon^2 v_1 \quad (65)$$

Equations (64) are computed in Sections 4.8.4 and equation (65) is used in Section 4.8.2 and 4.19.2.

And finally from (57), (61), (63), (62a), and (62):

$$\left. \begin{aligned}
 \dot{x}'_a(0) &= 0 \\
 \dot{y}'_a(0) &= x''_a(0) \cdot \dot{\Omega} + \dot{y}''_a(0) \cos i \\
 \dot{z}'_a(0) &= \dot{y}''_a(0) \sin i
 \end{aligned} \right\} \quad (66)$$

which are computed in Sections 4.8.4 and 4.19.2.

### 3.3.1.2 Solution for Perturbed Kepler's Equation

Since the original equations of motion (5) are written with  $t$  being the independent variable, equation (28) has to be inserted to give  $\psi$  as a function of  $t$ . To accomplish this, the following iteration scheme is used to solve the perturbed Kepler's equation for  $E$ .

First guess of  $E$  for a given  $t_\psi$ :

$$E_1 = a^{-3/2} t_\psi - e \sin(t_\psi / a^{3/2}) (1 - e \cos t_\psi / a^{3/2}) \quad (67)$$

i.e., true anomaly = mean anomaly

Then use the formula

$$E_i = E_{i-1} + \frac{t_\psi - t_\psi(E_{i-1})}{\frac{dt(E_{i-1})}{dE}} \quad i = 2, 3, \dots \quad (68)$$

to obtain a better value for  $E$ , and so on. On the right side of equation (68), the solution (28) with the argument  $E_{i-1}$  is used. From equation (25)  $\psi$  is obtained as a function of  $E$  as:

$$\sin \psi = \frac{\sqrt{1 - e^2} \sin E}{1 + e \cos E} \quad (68aa)$$

$$\cos \psi = \frac{\cos E + e}{1 + e \cos E} \quad (68ab)$$

In the limit, the second term of equation (68) is to go to zero. However,  $t_\psi$  as defined by equation (28) is a very complex function. Round off may limit the smallness of this quantity. Taking the derivative of equation (28) in the form:

$$\frac{dt}{dE} = a^{3/2} (1 + e \cos E) + \epsilon^2 \frac{dt_1}{d\psi} \frac{d\psi}{dE} \quad (68ac)$$

where

$$\begin{aligned} \epsilon^2 \frac{dt_1}{d\psi} = & -\frac{15}{4} \epsilon a^3 e \{ \cos \psi - \cos [(2p_1 - 1)\psi + \theta_1] \} \\ & + \epsilon^2 a^{9/2} \left\{ \frac{10}{3} \cos \psi - \frac{11}{4} \cos (2p_1 \psi + \alpha_1) \right\} \end{aligned} \quad (68ad)$$

and with equation (68ab):

$$\epsilon^2 \frac{dt_1}{dE} = \epsilon^2 \frac{dt_1}{d\psi} \frac{d\psi}{dE} = \frac{\sqrt{1 - e^2}}{1 + e \cos E} \epsilon^2 \frac{dt_1}{d\psi} \quad (68ae)$$

and substituting in equation (68) gives after cancellation and rearrangement:

$$E_i = \frac{a^{3/2} e (E_{i-1} \cos E_{i-1} - \sin E_{i-1}) + E_{i-1} (\epsilon^2 \frac{dt_1}{dE})_{i-1} + t_\psi - t_{1(i-1)}}{a^{3/2} (1 + e \cos E) + \epsilon^2 \frac{dt_1}{dE}} \quad (68af)$$

where

$$\begin{aligned} t_1 = & -\frac{15}{4} \epsilon a^3 e \{ \sin \psi - \sin [(2p_1 - 1)\psi + \theta_1] + \sin \theta_1 \} \\ & + \epsilon^2 a^{9/2} \left\{ \frac{10}{3} \sin \psi - \frac{11}{8} [\sin (2p_1 \psi + \alpha_1) - \sin \alpha_1] \right\} \end{aligned} \quad (68ag)$$

Equations (68aa) and (68ab) are computed in Section 4.17.2.2; equations (68ad), (68ae), (68af), and (68ag), in Section 4.17.2.4.

When  $\psi$  is determined by arc tangent function, it is placed in the correct one of the first four quadrants by comparing the signs of the sine and cosine. However, when  $E$  is near 0 or  $2\pi$ , an intermediate step of the iteration may result in a value of  $E$  in the fifth or minus first quadrant. To assure that  $\psi$  and  $E$  be in the same quadrant when  $\cos \psi > 0$  the equation is then:

$$\psi = \text{sign } E \left[ \psi + \pi(1 - \text{sign } \psi + 2 \text{ integer part of } \frac{E}{2\pi}) \right] \quad (68ah)$$

### 3.3.1.3 Jacobi Integral

The test quantity,  $\Delta C_a$ , for the analytic solution is simply obtained by subscripting the coordinates in equation (16) with the subscript  $a$  as shown in Section 4.17.15.

The initial value of the Jacobi constant is given by [see equations (13b) and (14)]:

$$C_{\text{init}} = C_{\text{part}} - \frac{2(1 - \epsilon^4)}{\epsilon^2 r_0} \quad (68a)$$

Substitution of the initial conditions (63) and (66) in equation (14) yields:

$$C_{\text{part}} = x_a''(0)^2 \dot{\Omega}^2 + 2x_a''(0) \dot{y}_a''(0) \dot{\Omega} \cos i + \dot{y}_a''(0)^2 \cos^2 i + \dot{y}_a''(0)^2 \sin^2 i \\ - 2x_a''(0) [x_a''(0) \dot{\Omega} + \dot{y}_a''(0) \cos i] - 2x_a''(0)(1 - \epsilon^4) - \frac{2}{x_a''(0)}$$

or

$$C_{\text{part}} = x'_a(0)^2 \left\{ \dot{\Omega}^2 + 2\dot{\Omega} \frac{d\phi(0)}{dt} \cos i + \left( \frac{d\phi(0)}{dt} \right)^2 - 2\epsilon \left[ \dot{\Omega} + \frac{d\phi(0)}{dt} \cos i \right] \right\} - 2(1 - \epsilon^4) x'_a(0) - \frac{2}{x'_a(0)} \quad (68b)$$

$C_{\text{init}}$  is computed in Section 4.8.6. The fractional error in the Jacobi constant is then given by  $\frac{\Delta C_a}{C_{\text{init}}}$

### 3.3.2 Numerical Residual Perturbation Method

The numerical residual perturbation method to be described below uses as a reference orbit the two-body ellipse as given by  $s_0, t_0$ , but accounts for the secular perturbations, which are reflected by the rotation of this ellipse with respect to the  $x', y', z'$  frame. This rotation is composed of the recession of the node as defined by equation (39) and the precession of the pericenter, which is comprised in equations (30) and (31). However, equations (30) and (31) contain some oscillatory terms due to the elliptic motion of the satellite. In order to separate the secular part, the argument of pericenter,  $\omega$ , is introduced as follows:

$$\phi = \psi + \omega \quad (69)$$

Substituting equations (52) and (41) in equation (40), one obtains:

$$\psi = \psi + \omega + \epsilon^2 \omega_1 \psi + \epsilon^2 \omega_1 \omega + \epsilon^2 v_1 t + \epsilon^4 \omega_1 v_1 t \quad (70)$$

Now from equation (26) or equation (27), it can be seen that the secular part of the relation between  $t$  and  $\psi$  is given by:

$$\psi = a^{-3/2} t \quad (71)$$



If equations (71), (46), and (35) are substituted in equation (70), one obtains for  $\omega$ :

$$\omega = \epsilon^2 \frac{\frac{3}{2} a^{3/2} - v_1 \epsilon^2 \omega_1}{1 + \epsilon^2 \omega_1} t_\omega = \dot{\omega} t_\omega \quad (72)$$

To keep  $\omega$  smaller than  $2\pi$ ,  $t$  is replaced by:

$$t_\omega = (t \text{ modulo } 2\pi) / \dot{\omega} \quad (72a)$$

The quantity  $\dot{\omega}$  is computed in Section 4.4,  $t_\omega$  is incremented in the RKSIMP subroutine (Section 7), and the modulo is formed in Section 4.17.1. The rotating ellipse is then described by:

$$\begin{aligned} \xi_s &= \rho_s \cos v \\ \eta_s &= \rho_s \sin v \\ \zeta_s &= 0 \end{aligned} \quad (73)$$

where the origin of  $\xi$ ,  $\eta$ ,  $\zeta$  is the smaller body,  $\xi$  points towards apocenter,  $\eta$  points  $90^\circ$  ahead of  $\xi$  in the direction of motion,  $\zeta$  is normal to the orbit plane, and

$$\begin{bmatrix} \xi_s \\ \eta_s \\ \zeta_s \end{bmatrix} = \{M\} \begin{bmatrix} x'_s \\ y'_s \\ z'_s \end{bmatrix} \quad (74)$$

The subscript  $s$  indicates approximate coordinates as derived from the reference orbit;  $\rho_s$  and  $v$  are radial distance and true anomaly, the latter measured from apocenter.

{M} is a matrix composed of three rotation matrices: A rotation around the z' axis by the angle  $\Omega$ , described by (see page 41 ff. reference 3):

$$\{\Omega\} = \begin{bmatrix} \cos \Omega & \sin \Omega & 0 \\ -\sin \Omega & \cos \Omega & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (75)$$

A rotation around the line of nodes by the angle  $i$ , described by:

$$\{i\} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos i & \sin i \\ 0 & -\sin i & \cos i \end{bmatrix} \quad (76)$$

A rotation around the polar axis of the orbital plane by the angle  $\omega$ , described by:

$$\{\omega\} = \begin{bmatrix} \cos \omega & \sin \omega & 0 \\ -\sin \omega & \cos \omega & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (77)$$

Then

$$\{M\} = \{\omega\} \cdot \{i\} \{\Omega\} \quad (78)$$

$$\{M\} = \begin{bmatrix} \cos \omega \cos \Omega & \cos \omega \sin \Omega & \sin \omega \sin i \\ -\sin \omega \sin \Omega \cos i & +\sin \omega \cos \Omega \cos i & \\ -\sin \omega \cos \Omega & -\sin \omega \sin \Omega & \cos \omega \sin i \\ -\cos \omega \sin \Omega \cos i & +\cos \omega \cos \Omega \cos i & \\ +\sin \Omega \sin i & -\cos \Omega \sin i & \cos i \end{bmatrix} \quad (79)$$

Now, using the vector notation:

$$\vec{\xi}_S = \begin{bmatrix} \xi_S \\ \eta_S \\ \zeta_S \end{bmatrix} \quad \vec{x}'_S = \begin{bmatrix} x'_S \\ y'_S \\ z'_S \end{bmatrix} \quad (80)$$

in inertial coordinates the two-body solution may be written as:

$$\vec{x}'_S = \{M\}^T \xi_S \quad (81)$$

or, using equation (73) and the transposed matrix of matrix (79)

$$\left. \begin{aligned} x'_S &= \rho_S [\cos(v + \omega) \cos \Omega - \sin(v + \omega) \sin \Omega \cos i] \\ y'_S &= \rho_S [\cos(v + \omega) \sin \Omega + \sin(v + \omega) \cos \Omega \cos i] \\ z'_S &= \rho_S \sin(v + \omega) \sin i \end{aligned} \right\} \quad (82)$$

Equation (82) is computed in Section 4.15.6.

The velocity components are obtained by differentiation of equation (81) with respect to time:

$$\dot{\vec{x}}'_S = \frac{d}{dt} \{M\}^T \cdot \vec{\xi}_S + \{M\}^T \dot{\xi}_S \quad (83)$$

Differentiating again yields the accelerations:

$$\ddot{\vec{x}}'_S = \{M\}^T \ddot{\xi}_S + 2 \frac{d\{M\}^T}{dt} \dot{\xi}_S + \frac{d^2\{M\}^T}{dt^2} \xi_S \quad (84)$$

Evaluating the matrix derivatives yields:

$$\begin{aligned}
 \frac{d\{M\}^T}{dt} = -\dot{\Omega} & \begin{bmatrix} \cos \omega \sin \Omega & -\sin \omega \sin \Omega & -\cos \Omega \sin i \\ +\sin \omega \cos \Omega \cos i & +\cos \omega \cos \Omega \cos i & \\ -\cos \omega \cos \Omega & \sin \omega \cos \Omega & -\sin \Omega \sin i \\ +\sin \omega \sin \Omega \cos i & +\cos \omega \sin \Omega \cos i & \\ 0 & 0 & 0 \end{bmatrix} \\
 -\dot{\omega} & \begin{bmatrix} \sin \omega \cos \Omega & \cos \omega \cos \Omega & 0 \\ +\cos \omega \sin \Omega \cos i & -\sin \omega \cos \Omega \cos i & \\ \sin \omega \sin \Omega & \cos \omega \sin \Omega & 0 \\ -\cos \omega \cos \Omega \cos i & +\sin \omega \cos \Omega \cos i & \\ -\cos \omega \sin i & \sin \omega \sin i & 0 \end{bmatrix} \quad (85) \\
 \frac{d^2\{M\}^T}{dt^2} = -\dot{\Omega}^2 & \begin{bmatrix} \cos \omega \cos \Omega & -\sin \omega \cos \Omega & +\sin \Omega \sin i \\ -\sin \omega \sin \Omega \cos i & -\cos \omega \sin \Omega \cos i & \\ \cos \omega \sin \Omega & -\sin \omega \sin \Omega & -\cos \Omega \sin i \\ +\sin \omega \cos \Omega \cos i & +\cos \omega \cos \Omega \cos i & \\ 0 & 0 & 0 \end{bmatrix} \\
 + 2\dot{\Omega}\dot{\omega} & \begin{bmatrix} \sin \omega \sin \Omega & \cos \omega \sin \Omega & 0 \\ -\cos \omega \cos \Omega \cos i & +\sin \omega \cos \Omega \cos i & \\ -\sin \omega \cos \Omega & -\cos \omega \cos \Omega & 0 \\ -\cos \omega \sin \Omega \cos i & +\sin \omega \sin \Omega \cos i & \\ 0 & 0 & 0 \end{bmatrix} \\
 -\dot{\omega}^2 & \begin{bmatrix} \cos \omega \cos \Omega & -\sin \omega \cos \Omega & 0 \\ -\sin \omega \sin \Omega \cos i & -\cos \omega \sin \Omega \cos i & \\ \cos \omega \sin \Omega & -\sin \omega \sin \Omega & 0 \\ +\sin \omega \cos \Omega \cos i & +\cos \omega \cos \Omega \cos i & \\ \sin \omega \sin i & \cos \omega \sin i & 0 \end{bmatrix} \quad (86)
 \end{aligned}$$

The rates of change of node and pericenter,  $\dot{\Omega}$  and  $\dot{\omega}$ , are obtained by differentiating equations (39) and (72) with respect to  $t$ :

$$\dot{\Omega} = -\epsilon^2 \lambda_1 \quad (87)$$

$$\dot{\omega} = \epsilon^2 \frac{\frac{3}{2} a^{3/2} - \epsilon^2 v_1 \omega_1}{1 + \epsilon^2 \omega_1} \quad (88)$$

Equations (87) and (88) are computed in Section 4.4.

The vector  $\dot{\xi}_s$  in Keplerian motion is given by:

$$\begin{aligned} \dot{\xi}_s &= -\frac{\sin v}{a^{1/2}(1-e^2)^{1/2}} \\ \dot{\eta}_s &= \frac{\cos v - e}{a^{1/2}(1-e^2)^{1/2}} \\ \dot{\zeta}_s &= 0 \end{aligned} \quad (89)$$

Equations (89) can be derived from equation (73) by differentiation with respect to time, making use of the area-integral  $\rho_s^2 \dot{v} = a^{1/2}(1-e^2)^{1/2}$  and the elliptic formula

$$\rho_s = \frac{a(1-e^2)}{1-e \cos v} \quad (90)$$

Substituting equations (73), (89), (79), and (85) in equation (83) and rearranging yields the following expression for the approximate velocity vector:

$$\begin{aligned} \dot{x}'_s &= \frac{-1}{\sqrt{a(1-e^2)}} \{ + \sin(v+\omega) \cos \Omega + \cos(v+\omega) \sin \Omega \cos i \\ &\quad - e(\sin \omega \cos \Omega + \cos \omega \sin \Omega \cos i) \} \\ &\quad - \rho_s \{ \cos(v+\omega) \sin \Omega [\dot{\Omega} + \dot{\omega} \cos i] \\ &\quad + \sin(v+\omega) \cos \Omega [\dot{\omega} + \dot{\Omega} \cos i] \} \end{aligned} \quad (91)$$

$$\begin{aligned}
\dot{y}'_s = & \frac{1}{\sqrt{a(1-e^2)}} \{-\sin \Omega \sin (v + \omega) + \cos \Omega \cos (v + \omega) \cos i \\
& + e[\sin \omega \sin \Omega - \cos \omega \cos \Omega \cos i]\} \\
& + \rho_s \{\cos (v + \omega) \cos \Omega (\dot{\Omega} + \dot{\omega} \cos i) \\
& - \sin (v + \omega) \sin \Omega (\dot{\omega} + \dot{\Omega} \cos i)\}
\end{aligned} \tag{92}$$

$$\dot{z}'_s = \frac{\sin i}{\sqrt{a(1-e^2)}} [\cos (v + \omega) - e \cos \omega] + \rho_s \dot{\omega} \sin i \cos (v + \omega) \tag{93}$$

Equations (91) through (93) are used in Section 4.17.13.

The residual perturbation equations of motion are obtained by substitution of the two-body equations of motion, i.e.:

$$\ddot{\xi}'_s = - \frac{\xi'_s}{\rho_s} \tag{94}$$

in equation (84), substitution of equation (74) in the result, and then subtraction of this from the exact restricted three-body equations (5), yielding:

$$\begin{aligned}
\ddot{\vec{x}}' - \ddot{\vec{x}}'_s = \ddot{\Delta \vec{x}}' = & - \frac{\vec{x}'}{\rho} + \frac{\vec{x}'_s}{\rho_s} - \frac{(1-\epsilon^4)\epsilon^2 \vec{x}'}{r^3} + (1-\epsilon^4)\left(1 - \frac{1}{r^3}\right) \begin{bmatrix} \cos \epsilon t \\ \sin \epsilon t \\ 0 \end{bmatrix} \\
& - 2 \frac{d\{M\}^T}{dt} \dot{\xi}'_s - \frac{d^2\{M\}^T}{dt^2} \xi'_s
\end{aligned} \tag{95}$$

The vector  $\Delta \vec{x}'$  is the correction vector which is to be added to the solution of the reference orbit in order to obtain the correct coordinates, i.e.,

$$\vec{x}' = \vec{x}'_s + \Delta\vec{x}' \quad (96)$$

(This addition is accomplished in Section 4.16.2.)

Equations (95) are to be integrated numerically. The first two terms on the right side of equation (95) represent a small difference of two large numbers. In order to avoid the loss of significant figures, the Encke transformation as described in detail in reference 3, starting at the bottom of page 176, will be applied.

Substitution of equations (85), (86), (89), and (73) in equation (95) and rearranging yields:

$$\begin{aligned} \ddot{\Delta x}' = & \frac{1}{\rho_s} \left[ f_q(x'_s + \Delta x') - x' \right] - \frac{(1 - \epsilon^4) \epsilon^2 x'}{r^3} + (1 - \epsilon^4) \left(1 - \frac{1}{r^3}\right) \cos \epsilon t \\ & + \frac{2\dot{\Omega}}{\sqrt{a(1 - e^2)}} \{- \sin(v + \omega) \sin \Omega + \cos(v + \omega) \cos \Omega \cos i \\ & + e (\sin \omega \sin \Omega - \cos \omega \cos \Omega \cos i)\} \\ & + \frac{2\dot{\omega}}{\sqrt{a(1 - e^2)}} \{+ \cos(v + \omega) \cos \Omega - \sin(v + \omega) \sin \Omega \cos i \\ & - e (\cos \omega \cos \Omega - \sin \omega \sin \Omega \cos i)\} \\ & - 2\rho_s \dot{\Omega} \dot{\omega} \{\sin(v + \omega) \sin \Omega - \cos(v + \omega) \cos \Omega \cos i\} \\ & + \rho_s (\dot{\omega}^2 + \dot{\Omega}^2) \{\cos(v + \omega) \cos \Omega - \sin(v + \omega) \sin \Omega \cos i\} \end{aligned} \quad (98)$$

$$\begin{aligned} \ddot{\Delta y}' = & \frac{1}{\rho_s} \left[ f_q(y'_s + \Delta y') - y' \right] - \frac{(1 - \epsilon^4) \epsilon^2 (y'_s + \Delta y')}{r^3} + (1 - \epsilon^4) \left(1 - \frac{1}{r^3}\right) \sin \epsilon t \\ & + \frac{2\dot{\Omega}}{\sqrt{a(1 - e^2)}} \{\sin(v + \omega) \cos \Omega + \cos(v + \omega) \sin \Omega \cos i \end{aligned}$$

$$\begin{aligned}
& - e (\sin \omega \cos \Omega + \cos \omega \sin \Omega \cos i) \} \\
& + \frac{2\dot{\omega}}{\sqrt{a(1-e^2)}} \{ \cos (v + \omega) \sin \Omega + \sin (v + \omega) \cos \Omega \cos i \\
& - e (\cos \omega \sin \Omega + \sin \omega \cos \Omega \cos i) \} \\
& - 2\rho_s \dot{\Omega} \dot{\omega} \{ - \sin (v + \omega) \cos \Omega - \cos (v + \omega) \sin \Omega \cos i \} \\
& + \rho_s (\dot{\Omega}^2 + \dot{\omega}^2) \{ \cos (v + \omega) \sin \Omega + \sin (v + \omega) \cos \Omega \cos i \} \quad (99)
\end{aligned}$$

$$\begin{aligned}
\Delta \ddot{z}' = \frac{1}{\rho_s^3} [f q (z'_s + \Delta z') - \frac{(1 - \epsilon^4) \epsilon^2 (z'_s + \Delta z')}{r^3} \\
+ \dot{\omega} \sin i \{ \left[ \frac{2}{\sqrt{a(1-e^2)}} + \rho_s \dot{\omega} \right] \sin (v + \omega) - \frac{2e}{\sqrt{a(1-e^2)}} \sin \omega \} \quad (100)
\end{aligned}$$

Equations (98) through (100) are used in Section 4.16.2 where the quantities  $f$  and  $q$  are defined by:

$$q = \frac{1}{\rho_s^2} \left[ (x'_s + \frac{1}{2} \Delta x') \Delta x' + (y'_s + \frac{1}{2} \Delta y') \Delta y' + (z'_s + \frac{1}{2} \Delta z') \Delta z' \right] \quad (101)$$

$$f = \frac{1 - (1 + 2q)^{-3/2}}{9} \quad (102)$$

$$\begin{aligned}
((7,6)) & = 3 \left[ 1 - \frac{5}{2} q + \frac{5 \cdot 7}{2 \cdot 3} q^2 - \frac{5 \cdot 7 \cdot 9}{2 \cdot 3 \cdot 4} q^3 + \dots \right] \\
& = 3 \left\{ 1 - \frac{5}{2} q \left( 1 - \frac{7}{3} q \left( 1 - \frac{9}{4} q \left( 1 - \frac{11}{5} q (\dots) \right) \right) \right) \right\}
\end{aligned}$$



The above expansion can be found in reference 4, page 98, equation ((7,6)). They are evaluated in Section 4.16.1.

Now it is still necessary to define  $\rho_s$  and  $v$  as functions of time. Introducing the eccentric anomaly  $E$ , the theory of elliptic motion (e.g., reference 2, page 164, equations (48) and (49), with  $k(\sqrt{1+m})$  being 1 in this case) gives:

$$t = a^{-3/2}[E + e \sin E] \quad (103)$$

$$\cos v = \frac{\cos E + e}{1 + e \cos E} \quad (104)$$

$$\sin v = \frac{\sqrt{1 - e^2} \sin E}{1 + e \cos E} \quad (105)$$

$$\rho_s = a(1 + e \cos E) \quad (106)$$

There are plus signs instead of minus signs in the above equations because all the anomalies are measured from apocenter. Equations (104) and (105) are evaluated in Section 4.15.2. Equations (106) is evaluated in Section 4.15.3. Kepler's equation (103) is solved for  $E$  in Section 4.15.1 by iteration as follows:

$$\begin{aligned} E_{i+1} &= E_i + \frac{a^{-3/2}[t_s - t_s(E_i)]}{\frac{dt_s(E_i)}{dE}} = E_i + \frac{a^{-3/2}t_s - E_i - e_s \sin E_i}{1 - e_s \cos E_i} \\ &= \frac{a^{-3/2}t_s - e_s(E_i \cos E_i - \sin E_i)}{1 - e_s \cos E_i} \end{aligned} \quad (107)$$

with a first estimate:

$$E_1 = a^{-3/2} t_s - e_s \sin(a^{-3/2} t_s) + \frac{1}{2} e_s^2 \sin 2(a^{-3/2} t_s) \quad (108)$$

Equation (108) is taken from reference 2, page 169, equation (59);  $M$  being  $a^{-3/2} t_s$ . The quantity  $t_s$  is the input time, modulo  $2\pi a^{3/2}$  (see Section 17.1).

### 3.3.2.1 Initial Conditions for the Residual Perturbation Solution

The initial conditions on the residual perturbation equations of motion (98) through (100) are chosen such that for  $t = 0$  the total values of  $x'$ ,  $y'$ , and  $z'$  agree with the values  $x'_a$ ,  $y'_a$ , and  $z'_a$  of the analytic solution.

$$\vec{x}'(0) = \vec{x}'_s(0) + \Delta \vec{x}'(0) = \vec{x}'_a(0) \quad (109)$$

They are set equal in Section 4.8.3. The initial conditions for the reference orbit are:

from equations (103) through (106), (73), (39), (72), and (82):

$$E(0) = v(0) = 0, \rho_s(0) = a(1 + e) \quad (110)$$

$$\xi_s(0) = \rho_s(0), \eta_s(0) = \zeta_s(0) = 0 \quad (111)$$

$$\Omega(0) = 0 \quad (112)$$

$$\omega(0) = 0$$

$$x'_s(0) = \rho_s(0)$$

$$y'_s(0) = z'_s(0) = 0 \quad (113)$$

The initial conditions from equations (113) are established in Section 4.8.3

From equations (91) through (93):

$$\dot{x}'_s(0) = 0 \quad (114)$$

$$\begin{aligned} \dot{y}'_s(0) &= [a(1 - e^2)]^{-1/2} \cos i (1 - e) + \rho_s(0)(\dot{\Omega} + \dot{\omega} \cos i) \\ &= x'_s(0) \left[ \dot{\Omega} + \left( \frac{1}{\frac{dt_o}{d\psi}(0)} + \dot{\omega} \right) \cos i \right] \end{aligned} \quad (115)$$

$$\begin{aligned} \dot{z}'_s(0) &= \{ [a(1 - e^2)]^{-1/2} (1 - e) + \rho_s(0) \dot{\omega} \} \sin i \\ &= x'_s(0) \left[ \frac{1}{\frac{dt_o}{d\psi}(0)} + \dot{\omega} \right] \sin i \end{aligned} \quad (116)$$

Equation (114) is computed in Section 4.8.6. Equations (115) and (116) are computed in Section 4.8.4.

Now from equation (109):

$$\begin{aligned} \Delta \vec{x}'(0) &= \vec{x}'_a(0) - \vec{x}'_s(0) \\ \Delta \dot{\vec{x}}'(0) &= \dot{\vec{x}}'_a(0) - \dot{\vec{x}}'_s(0) \end{aligned} \quad (117)$$

Substituting equations (113) through (116), (63), and (66) in equation (117) yields:

$$\left. \begin{aligned} \Delta x'(0) &= 0 \\ \Delta y'(0) &= 0 \\ \Delta z'(0) &= 0 \\ \Delta \dot{x}'(0) &= 0 \end{aligned} \right\} \quad (118)$$

$$\begin{aligned}\Delta \dot{y}'(0) &= x' \dot{G} \cos i \\ \Delta \dot{z}'(0) &= x' \dot{G} \sin i\end{aligned}\tag{119}$$

where

$$\begin{aligned}\dot{G} &= [(1 + \epsilon^2 \omega_1^2) \frac{dt(0)}{d\psi}]^{-1} - \epsilon^2 v_1 - \left(\frac{dt_o(0)}{d\psi}\right)^{-1} - \dot{\omega} \\ &= \frac{\frac{dt_o(0)}{d\psi} \left(\frac{dt_o(0)}{d\psi} + \epsilon^2 \frac{dt_1(0)}{d\psi}\right) (1 + \epsilon^2 \omega_1^2)}{\frac{dt(0)}{d\psi} (1 + \epsilon^2 \omega_1^2) \frac{dt_o(0)}{d\psi}} - \epsilon^2 v_1 - \dot{\omega} \\ &= \frac{-\epsilon^2 \omega_1 \frac{dt_o(0)}{d\psi} - \epsilon^2 \frac{dt_1(0)}{d\psi} (1 + \epsilon^2 \omega_1^2)}{\frac{dt(0)}{d\psi} \frac{dt_o(0)}{d\psi} (1 + \epsilon^2 \omega_1^2)} - \epsilon^2 v_1 - \dot{\omega}\end{aligned}\tag{120}$$

The  $\Delta$  velocities are not zero at the initial time because the precessing ellipse has been chosen to be a mean ellipse rather than an initial osculating ellipse. This was necessary in order that the period of the ellipse did not vary in a secular manner from the period of the actual motion. Equations (117) through (120) are computed in Section 4.8.4.

Substitution of the initial conditions into equations (98) through (100) yields the initial residual accelerations:

$$\begin{aligned}\Delta \ddot{x}'(0) &= \frac{1 - \epsilon^4}{r_o^3} [2\epsilon^2 x'(0) + 3\epsilon^4 x'(0)^2 + \epsilon^6 x'(0)^3] \\ &\quad + \frac{2(1 - e)}{\sqrt{a(1 - e^2)}} (\dot{\Omega} \cos i + \dot{\omega}) \\ &\quad + 2x'(0) \dot{\Omega} \dot{\omega} \cos i + x'(0) (\dot{\omega}^2 + \dot{\Omega}^2)\end{aligned}\tag{121}$$

$$\Delta \ddot{y}'(0) = \Delta \ddot{z}'(0) = 0\tag{122}$$

which are computed in Section 4.8.5.

### 3.3.3 Fixed Reference Orbit

For comparison, a fixed, osculating Kepler ellipse is also calculated from the initial conditions. The node and the argument of apocenter are zero at  $t = 0$ , and the inclination has the same value as in the previous solutions. However, the osculating values for  $a$  and  $e$  are different from those used previously because the perturbations in the velocities do not vanish initially in the analytic solution.

To obtain the osculating values for  $a$ , the formula in reference 3, page 48, equation (88) is used:

$$\frac{1}{a} = \frac{2}{r} - v^2 \quad (123)$$

or, in our notation:

$$a = \frac{1}{\frac{2}{\rho(0)} - \dot{\rho}(0)^2} \quad (124)$$

On the same page, further below, the following equation is found:

$$e \cos u = 1 - \frac{r}{a} \quad (125)$$

Since the eccentric anomaly,  $u$ , is  $180^\circ$  initially, this becomes in our notation:

$$e = \frac{\rho(0)}{a} - 1 = \frac{x'(0)}{a} - 1 \quad (126)$$

Equations (124) and (126) are used in Section 4.19.3. The coordinates and velocity components are then obtained as in equations (82) and (91) through (93), but setting  $\omega = \Omega = \dot{\omega} = \dot{\Omega} = 0$ :

$$\begin{aligned}
x'_f &= \rho_s \cos v \\
y'_f &= \rho_s \sin v \cos i \\
z'_f &= \rho_s \sin v \sin i \\
\dot{x}'_f &= - [a(1 - e^2)]^{-1/2} \sin v \\
\dot{y}'_f &= + [a(1 - e^2)]^{-1/2} (\cos v - e) \cos i \\
\dot{z}'_f &= + [a(1 - e^2)]^{-1/2} (\cos v - e) \sin i
\end{aligned} \tag{127}$$

Equations (127) are computed in Section 4.28.2.

Again, the Jacobi constant is calculated according to equation (16), subscripting all the variables with  $f$ , indicating that they are computed from the fixed orbit. The radius vector and the total velocity are also compared with the corresponding values obtained by the numerical residual perturbation method:

$$\begin{aligned}
\Delta\rho_f &= \frac{1}{a} [(x'_f - x')^2 + (y'_f - y')^2 + (z'_f - z')^2]^{1/2} \\
\Delta\dot{\rho}_f &= a^{1/2} [(\dot{x}'_f - \dot{x}')^2 + (\dot{y}'_f - \dot{y}')^2 + (\dot{z}'_f - \dot{z}')^2]^{1/2}
\end{aligned} \tag{128}$$

Equations (128) are used in Section 4.28.2

The quantities  $\Delta\rho_f$  and  $\Delta\dot{\rho}_f$  are normalized with the semi-major axis and the circular velocity  $a^{-1/2}$ .

### 3.3.4 Numerical Integration of Total Equations

For the integration of the total equations the same initial conditions as in the analytic solution are used, but the exact equations of motion (5)

are integrated numerically. The results are checked by means of the Jacobi integral equation (16). The radius vector and the total velocity are also compared with the values obtained from the numerical residual perturbation method by computing the differences, normalized with  $a$  and  $a^{-1/2}$ , respectively:

$$\begin{aligned}\Delta\rho_c &= \frac{1}{a}[(x'_c - x')^2 + (y'_c - y')^2 + (z'_c - z')^2]^{1/2} \\ \Delta\dot{\rho}_c &= a^{1/2}[(\dot{x}'_c - \dot{x}')^2 + (\dot{y}'_c - \dot{y}')^2 + (\dot{z}'_c - \dot{z}')^2]^{1/2}\end{aligned}\tag{129}$$

where  $x'_c, y'_c, z'_c, \dot{x}'_c, \dot{y}'_c,$  and  $\dot{z}'_c$  stand for the values obtained by numerical integration of the exact equations of motion (5). Equations (129) were used in Section 4.28.2.

### 3.4 METHOD FOR NUMERICAL SOLUTION OF DIFFERENTIAL EQUATIONS

The Runge-Kutta method is used for the numerical solution of the differential equations. The method is a simple extension of the methods for the second order and simultaneous equations given by Hildebrand (reference 5, page 237) which are:

1. Given the simultaneous first order equations:

$$\begin{aligned}(6.16.7) \quad \frac{dy}{dx} &= F(x,y,u), \\ \frac{du}{dx} &= G(x,y,u)\end{aligned}\tag{130}$$

the solution may be written as:

$$\begin{aligned}(6.16.8) \quad y_{n+1} &= y_n + \frac{1}{6} (k_0 + 2k_1 + 2k_2 + k_3) + O(h^5) \\ u_{n+1} &= u_n + \frac{1}{6} (m_0 + 2m_1 + 2m_2 + m_3) + O(h^5)\end{aligned}\tag{131}$$

where

$$\begin{aligned}(6.16.9) \quad k_0 &= hF(x_n, y_n, u_n), \\ k_1 &= hF(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_0, u_n + \frac{1}{2}m_0), \\ k_2 &= hF(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1, u_n + \frac{1}{2}m_1), \\ k_3 &= hF(x_n + h, y_n + k_2, u_n + m_2)\end{aligned}\tag{132}$$

and

$$\begin{aligned}(6.16.10) \quad m_0 &= hG(x_n, y_n, u_n), \\ m_1 &= hG(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_0, u_n + \frac{1}{2}m_0), \\ m_2 &= hG(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1, u_n + \frac{1}{2}m_1), \\ m_3 &= hG(x_n + h, y_n + k_2, u_n + m_2).\end{aligned}\tag{133}$$

2. The second order equation

$$(6.16.11) \quad \frac{d^2y}{dx^2} = G(x, y, y'),\tag{134}$$

can be written as the two simultaneous first order differential equations

$$\frac{dy}{dx} = u$$

and

$$\frac{du}{dx} = G(x, y, u)$$



Then equation (6.16.9) gives:

$$k_0 = hy'_n, \quad k_1 = hy'_n + \frac{h}{2} m_0, \quad k_2 = hy'_n + \frac{h}{2} m_1, \quad k_3 = hy'_n + hm_2,$$

and hence equations (6.16.8) and (6.16.10) give:

$$\begin{aligned} (6.16.12) \quad y_{n+1} &= y_n + hy'_n + \frac{h}{6} (m_0 + m_1 + m_2) + O(h^5), \\ y'_{n+1} &= y'_n + \frac{1}{6} (m_0 + 2m_1 + 2m_2 + m_3) + O(h^5), \end{aligned} \quad \left. \vphantom{\begin{aligned} y_{n+1} \\ y'_{n+1} \end{aligned}} \right\} (135)$$

where

$$\begin{aligned} (6.16.13) \quad m_0 &= hG(x_n, y_n, y'_n), \\ m_1 &= hG(x_n + \frac{1}{2}h, y_n + \frac{1}{2}hy'_n, y'_n + \frac{1}{2}m_0), \\ m_2 &= hG(x_n + \frac{1}{2}h, y_n + \frac{1}{2}hy'_n + \frac{1}{2}hm_0, y'_n + \frac{1}{2}m_1), \\ m_3 &= hG(x_n + h, y_n + hy'_n + \frac{1}{2}hm_1, y'_n + m_2). \end{aligned} \quad \left. \vphantom{\begin{aligned} m_0 \\ m_1 \\ m_2 \\ m_3 \end{aligned}} \right\} (136)$$

The extension of the above technique to three second-order differential equations:

$$\frac{d^2 x_i}{dt^2} = f_i(t, x_1, x_2, x_3) \quad \text{where } i = 1, 2, 3 \quad (137)$$

yields

$$\frac{dx_i}{dt} = g_i(t, x_1, x_2, x_3) \quad \text{for } i = 1, 2, 3 \quad (138)$$

$$\frac{dg_i}{dt} = f_i(t, x_1, x_2, x_3) \quad (139)$$

$$m_{oi} = hf_i(t_n, x_{1n}, x_{2n}, x_{3n}) \quad (140)$$

$$m_{1i} = hf_i(t_n + \frac{h}{2}, x_{1n} + \frac{h}{2} \dot{x}_{1n}, x_{2n} + \frac{h}{2} \dot{x}_{2n}, x_{3n} + \frac{h}{2} \dot{x}_{3n}) \quad (141)$$

$$m_{2i} = hf_i(t_n + \frac{h}{2}, x_{1n} + \frac{h}{2} \dot{x}_{1n} + \frac{h}{4} m_{01}, x_{2n} + \frac{h}{2} \dot{x}_{2n} + \frac{h}{4} m_{02}, \\ x_{3n} + \frac{h}{2} \dot{x}_{3n} + \frac{h}{4} m_{03}) \quad (142)$$

$$m_{3i} = hf_i(t_n + h, x_{1n} + h\dot{x}_{1n} + \frac{h}{2} m_{11}, x_{2n} + h\dot{x}_{2n} + \frac{h}{2} m_{12}, \\ x_{3n} + h\dot{x}_{3n} + \frac{h}{2} m_{13}) \quad (143)$$

$$x_{i(n+1)} = x_{in} + h\dot{x}_{in} + \frac{h}{6} (m_{0i} + m_{1i} + m_{2i}) + O(h^5) \quad (144)$$

$$\dot{x}_{i(n+1)} = \dot{x}_{in} + \frac{1}{6} [m_{0i} + 2(m_{1i} + m_{2i}) + m_{3i}] + O(h^5) \quad (145)$$

With the substitution of  $x, y, z$  for  $x_1, x_2, x_3$ , equation (140) and the parameter for equation (141) are computed in Section 7.8; equation (141) and the parameter for equation (142) are computed in Section 7.26; equation (142) and the parameters for equation (143) are computed in Section 7.27; and equation (143) through (145) are computed in Section 7.28.

After integration over two intervals of equal size, the results for the velocity components are compared with an integration over the same intervals using Simpson's rule which is also of fourth order accuracy. Simpson's rule is given on page 73 of reference 5 as:

$$\int_{x_0}^{x_2} f(x) dx = \frac{h}{3} (f_0 + 4f_1 + f_2) - \frac{h^5 f^{IV}(\xi)}{90} \quad (146)$$

$$\text{where } x_0 < \xi < x_2$$

or in the notation of this problem:

$$\left. \begin{aligned} \dot{x}_{n+2} &= \dot{x}_n + \frac{\Delta t}{3} [\ddot{x}_n + 4\ddot{x}_{n+1} + \ddot{x}_{n+2}] \\ \dot{y}_{n+2} &= \dot{y}_n + \frac{\Delta t}{3} [\ddot{y}_n + 4\ddot{y}_{n+1} + \ddot{y}_{n+2}] \\ \dot{z}_{n+2} &= \dot{z}_n + \frac{\Delta t}{3} [\ddot{z}_n + 4\ddot{z}_{n+1} + \ddot{z}_{n+2}] \end{aligned} \right\} \quad (147)$$

where the subscript  $n$  means evaluated at time  $t_n$ , the subscript  $n+1$  means evaluated at  $t_n + \Delta t$ , and the subscript  $n+2$  means evaluated at  $t_n + 2\Delta t$ .

By virtue of the comparison between the two integrated results, decisions, are made by the program concerning the accuracy of the integration, and the computing interval for the next two intervals is chosen. The logic underlying these program decisions will now be explained using one first order differential equation as an example.

Let the differential equation to be solved be of the form:

$$\dot{x} = \dot{x}(t, x) \quad (148)$$

If this equation is integrated over an interval,  $h$ , by Runge-Kutta methods of fourth order, then the numerical value of that function corresponds to a Taylor series expansion with an error term of  $O(h^5)$ , i.e.:

$$x_{n+1} = x_n + \dot{x}_n h + \frac{\ddot{x}_n h^2}{2!} + \frac{\dddot{x}_n h^3}{3!} + \frac{x^{IV} h^4}{4!} + O(h^5) \quad (149)$$

The complete functional form of the coefficient of the error term is unknown

but it is known to contain  $x^V$ . For the purposes of this program the coefficient of the fifth order term is assumed to be the next term in the Taylor series

$\frac{x^V}{5!}$  and  $x^V$  is assumed to be a slowly varying function. The coefficient of the

fifth order term in Simpson's rule is known to be  $-\frac{x^V}{90}$ . Thus if we let  $x_c$

be the correct value of  $x$  at the end of the two equal intervals, and let  $x_{RK}$  and  $x_{SR}$  be the Runge-Kutta and Simpson's rule integrated values respectively, we may write:

$$x_c = x_{RK} + 2\left(\frac{x^V h^5}{5!}\right) \quad (150)$$

$$x_c = x_{SR} - \frac{x^V h^5}{90} \quad (151)$$

Eliminating  $x_c$  between these two equations and solving for  $x^V$  results in:

$$x^V = \frac{36(x_{SR} - x_{RK})}{h^5} \quad (152)$$

From equations (150) through (152) the error in the Runge-Kutta solution is estimated to be:

$$\delta x = \frac{3}{5} (x_{SR} - x_{RK}) \quad (153)$$

A factor of  $\frac{3}{5}$  is dropped in the use of this equation in Section 7.14 because an arbitrary constant is introduced at this point.

Letting  $\Delta\dot{x}$ ,  $\Delta\dot{y}$ ,  $\Delta\dot{z}$  be the changes in the  $\dot{x}$ ,  $\dot{y}$ ,  $\dot{z}$  values over the double interval, then what is required in the program is that:

$$E = \text{maximum} (|\delta\dot{x}|, |\delta\dot{y}|, |\delta\dot{z}|) < E_{\text{all}} = \text{maximum} (|W_8|C_{\text{max}}), \quad (154)$$

$$10^{-9} \text{ maximum} (|\dot{x}|, |\dot{y}|, |\dot{z}|)$$

$$\text{where } C_{\text{max}} = \text{maximum} (|\Delta\dot{x}|, |\Delta\dot{y}|, |\Delta\dot{z}|) \quad (155)$$

and  $W_8$  is an input number designed to require a series truncation greater than number truncation but as small as possible. An error which is less than

$10^{-9}$  of the maximum of the absolute values of  $\dot{x}$ ,  $\dot{y}$ , and  $\dot{z}$  is always acceptable, since it will be lost in the first addition anyway because of the limits of machine word length.

If  $E \leq E_{all}$ , the computation proceeds. If  $E > E_{all}$ , the last two steps are done over.

If  $E$  is greater than an input minimum error  $E_{min} \cdot C_{max}$ , then  $\Delta t$  is computed by:

$$\Delta t_{new} = FDT * \Delta t_{old} \left( \frac{E_{all}}{E} \right)^{.25} \quad (156)$$

If it assumed that  $E_{all} = W_8 C_{max} = K \Delta t$ , where  $K$  is some constant (since  $x$  is roughly proportional to  $\Delta t$  and  $c_{max}$  is normally proportional to  $\Delta x$ ) and  $FDT = 1$ , then by equations (150), (153), and (154)  $\Delta t_{new}$  would result in an error of exactly  $E_{all}$ .  $FDT$  is an input number  $< 1$  to prevent  $\Delta t_{new}$  from resulting in an error  $E > E_{all}$  due to number truncation or change of  $x^V$  over the two new intervals as compared to the  $x^V$  of the previous two intervals.

If  $E < E_{min} \cdot C_{max}$ , then it is assumed that the error in  $x_{RK}$  is primarily due to number truncation in the computations. In this case equation (150) does not apply. The new computing interval is then computed by:

$$\Delta t_{new} = \Delta t_{min} \cdot \Delta t_{old} \quad (157)$$

where  $\Delta t_{min}$  is an input quantity  $> 1$ . The checks against  $E_{all}$  are made in Sections 7.14 and 7.15 and checks against  $E_{min}$  are made in Section 7.22. The computation of the time interval by formula (156) is accomplished in Section 7.19. the computation of  $\Delta t$  by equation (156) is done in Section 7.24.

Section 4  
EQUATIONS IN ORDER OF SOLUTION

4.1 PRINT EXPLANATION OF ERROR CODE

<u>Error Code</u>	<u>Reason for Halt of Run</u>
1	Semi-major axis = 0
2	Eccentricity equals or exceeds 1
7	S (Inverse of radius vector) = 0
8	Derivative of time with respect to $\psi = 0$

4.1.1 Set Input Arrays = 0

Set  $W_I = 0$  for  $I = 1, 2, \dots, 30$

4.2 PROGRAM INPUT

Up to 30 quantities can be input to this program. Those specified thus far are given below.

4.2.1 Osculating Orbital Elements

$W_1 = a$ , semi-major axis in units of  $\mu^{1/2} D$

$W_2 = e$ , eccentricity

$W_3 = i$ , inclination (degrees)

4.2.2 Planetary Mass Ratio

$W_4 = \mu \equiv \epsilon^4$

4.2.3 Time and Accuracy Specifications

$W_5 = \text{FTD}$ , modifier of estimated computing interval for the next pass

$W_6 = \text{Not used}$

$W_7$  = Run stop time

$W_8$  = Maximum allowable velocity error

$W_9$  = Modifier of previous computing interval if velocity error is less than the minimum allowable

$W_{10}$  = Initial value of computing interval

$W_{11}$  = Maximum number of failures permitted in computing interval determination

$W_{12}$  = Minimum allowable velocity error

$W_{13}$  = Flag to call for print

If  $W_{13} \leq 0$ , print will be suppressed.

If  $W_{13} > 0$ , print will be called for.

#### 4.3 MOVE INPUT TO WORKING STORAGE AND CHECK INPUT

##### 4.3.1 Move Input to Working Storage and Form Necessary Constants

$a = |\text{input } a|$

$a_2 = a^2$

$a_s = a$

$e = |\text{input } e|$

$e_s = e$

$i^0 = \text{input } i$

$\epsilon_4 = \text{input } \mu$

$T_f = W_7$

#### 4.3.2 Check on Constant Divisors, Halt Run if Any Are Zero

##### 4.3.2.1 Test Semi-Major Axis

If  $a \leq 0$ , go to Step 4.3.2.1.1.

If  $a > 0$ , go to Step 4.3.2.2.

##### 4.3.2.1.1 Set Error Code = 1

##### 4.3.2.1.2 Print Error Code and go to Step 4.2

##### 4.3.2.2 Compute and Test $e_2$

$$e_2 = 1 - e^2$$

If  $e_2 \leq 0$ , go to Step 4.3.2.2.1.

If  $e_2 > 0$ , go to Step 4.4.

##### 4.3.2.2.1 Set Error Code to 2 and go to Step 4.3.2.1.2

#### 4.4 COMPUTE CONSTANTS

$$a_e = ae_2$$

$$a_e^{1/2} = \sqrt{a_e}$$

$$e_2 = \sqrt{e_2}$$

$$a_3 = a^3$$

$$\epsilon_2 = \sqrt{\epsilon_4}$$

$$T_{\text{num } 1} = 3$$



$$T_{\text{den } 1} = 1$$

$$\left. \begin{aligned} T_{\text{num } i} &= T_{\text{num } i-1} + 2 \\ T_{\text{den } i} &= T_{\text{den } i-1} + 1 \end{aligned} \right\} \begin{array}{l} \text{Form for} \\ i = 2, 3, \dots, 40 \end{array}$$

$$\text{ITP} = 2$$

$$P_{t1} = 0$$

$$ap_i = 0 \quad i = 1, \dots, 13$$

$$a_{32} = \sqrt{a^3}$$

$$a_{92} = (a^{3/2})^3$$

$$a_{S32} = a_{32}$$

$$\epsilon = \sqrt{\epsilon_2}$$

$$C_1 = 3.75\epsilon a_3 e$$

$$v_1 = \frac{-a_{32} (11 + 24e)}{12} \quad (35)$$

$$\epsilon_a = \epsilon a_{32}$$

$$\lambda_1 = - .75 \cdot a_{32} \quad (36)$$

$$m_1 = \frac{-a_{32} (1 + 12e)}{6} \quad (34), (35), (36)$$

$$\omega_1 = \frac{a_3 (-7 + 24e)}{12} \quad (46)$$

$$p_3 = \frac{1}{1 + \epsilon_2 \omega_1} - \epsilon_a [2 + \epsilon(v_1 - 2\lambda_1)] \quad (38)$$

$$i = i^0/57.2957795$$

$$C_i = \cos i$$

$$S_i = \sin i$$

$$i_s = i$$

$$C_{is} = \cos i_s$$

$$S_{is} = \sin i_s$$

$$\dot{\omega} = \left( \frac{3}{2} \epsilon_2 a_{32} - \epsilon_4 v_1 \omega_1 \right) / (1 + \epsilon_2 \omega_1) \quad (88)$$

$$\dot{\Omega} = - \frac{3}{4} \epsilon \cdot \epsilon_a \quad (87), (36)$$

$$C_{\dot{\omega}\dot{\Omega}} = \dot{\omega}\dot{\Omega}$$

$$C_{\dot{\omega}^2\dot{\Omega}^2} = \dot{\omega}^2 + \dot{\Omega}^2$$

$$a_{12} = a^{1/2}$$

$$p_1 = \frac{1}{1 + \epsilon_2 \omega_1} - \epsilon_a (1 + \epsilon_1 m_1) \quad (37)$$

$$C_2 = 1.875 \epsilon_1 a_{12} e$$

$$C_3 = \frac{.375 \epsilon_1 i a a_{32}}{1 - \epsilon_1 a_{32}}$$

$$C_6 = \dot{\Omega}$$

$$C_7 = \dot{\Omega}^2$$

$$C_8 = 1 - \epsilon_4$$

$$C_9 = a_e a_{12} e_2$$

$$C_{10} = \frac{\epsilon_2 \omega_1}{1 + \epsilon_2 \omega_1} + \epsilon_1 a_{32} (1 + \epsilon_1 m_1)$$

$$C_{11} = 2p_1 - 1$$

$$C_{12} = \frac{2\dot{\Omega}}{a_e 1/2}$$

$$C_{13} = \frac{2\dot{\omega}}{a_e 1/2}$$

$$C_{14} = 2C_{\dot{\omega}\dot{\Omega}}$$

Compute  $-\frac{2}{a_e 1/2} (\dot{\Omega} + \dot{\omega} \cos i)$

$$C_{15} = -C_{12} - C_{13} C_i$$

Compute  $-2\dot{\omega}\dot{\Omega} - (\dot{\omega}^2 + \dot{\Omega}^2) \cos i$

$$C_{16} = -C_{14} - C_{\dot{\omega}\dot{\Omega}} \cdot C_i$$

Compute  $\frac{2}{a_e 1/2} (\dot{\Omega} \cos i + \dot{\omega})$

$$C_{17} = C_{12} C_i + C_{13}$$

Compute  $2\dot{\omega}\dot{\Omega} \cos i + \dot{\omega}^2 + \dot{\Omega}^2$

$$C_{18} = C_{14} C_i + C_{\dot{\omega}^2 \dot{\Omega}^2}$$

Compute  $\dot{\omega}^2 \sin i$

$$C_{19} = \dot{\omega}^2 S_i$$

Compute  $\frac{2e \sin i}{a_e^{1/2}} \dot{\omega}$

$$C_{20} = e S_i C_{13}$$

Compute  $\frac{2 \sin i}{a_e^{1/2}} \dot{\omega}$

$$C_{21} = S_i C_{13}$$

Compute  $-\frac{2e}{a_e^{1/2}} (\dot{\Omega} + \dot{\omega} \cos i)$

$$C_{22} = e C_{15}$$

Compute  $\frac{2e}{a_e^{1/2}} (\dot{\Omega} \cos i + \dot{\omega})$

$$C_{23} = e C_{17}$$

$$C_{24} = \sqrt{1 - e_s^2}$$

$$T_1 = 2\pi a_{32} + C_1 \sin 2\pi C_{11} \quad (28a)$$

$$- 1.375 \epsilon_2 a_{92} \sin 4\pi p_1$$

$$P_{ss} = 2\pi a_{32}$$

$$P_{\Omega} = 2\pi/\dot{\Omega} \quad (39b)$$

$$P_{\omega} = 2\pi/\dot{\omega} \quad (72a)$$

$$\Delta E_{\max} = 0$$

$$\Delta E_{p\max} = 0$$

$$\alpha_1 = 0$$

$$\theta_1 = 0$$

$$\phi_1 = 0$$

$$C_{\alpha} = 1$$

$$S_{\alpha} = 0$$

$$C_{\theta} = 1$$

$$S_{\theta} = 0$$

#### 4.5 PRINT CONSTANTS

$w_1$	$w_2$	$w_3$	$w_4$	$w_5$
$w_6$	$w_7$	$w_8$	$w_9$	$w_{10}$
$w_{11}$	$w_{12}$	$w_{13}$	$w_{14}$	$w_{15}$

$$\begin{array}{cccccc}
 \omega_1 & v_1 & \lambda_1 & p_1 & p_3 & \\
 c_1 & c_2 & c_3 & c_4 & c_5 & c_6 \\
 c_7 & c_8 & c_9 & c_{10} & c_{11} & \sqrt{1 - e^2} \\
 \dot{\omega} & \dot{\Omega} & \ddot{\omega} & \dot{\Omega}^2 + \ddot{\Omega} & & 
 \end{array}$$

4.6 PRINT HEADER FOR NUMERICAL RESIDUAL PERTURBATION SOLUTION LISTING THE INPUT VALUES OF  $a$ ,  $e$ ,  $i$ , and  $\mu$

4.7 SET PERTURBATION-TOTAL FLAG

ICH = 1; the signal that residual perturbation equations are to be numerically solved and that the analytic solution is to be computed.

4.8 COMPUTE INITIAL VALUES OF POSITION, VELOCITY, AND ACCELERATION

$$\text{Set } e_3 = 1 + e$$

$$\frac{dt_0}{d\psi} = \frac{a_{32}(e_3)^{1.5}}{(1 - e)^{1/2}} \quad (59)$$

$$\epsilon^2 \frac{dt_1}{d\psi} = \frac{7}{12} \epsilon_2 a_{92}$$

$$\frac{dt}{d\psi} = \frac{dt_0}{d\psi} + \epsilon_2 \frac{dt_1}{d\psi} \quad (59)$$

4.8.1 Test Value of  $\frac{dt}{d\psi}$

If  $\frac{dt}{d\psi} = 0$ , go to Step 4.17.8.1.

If  $\frac{dt}{d\psi} \neq 0$ , continue below.

4.8.2 Compute  $\frac{d\phi}{dt}$  and Set Initial Values of  $t$  and  $\psi$

$$\frac{d\psi}{dt} = 1/\frac{dt}{d\psi}$$

$$\frac{d\phi}{dt} = \frac{\frac{d\psi}{dt}}{(1 + \epsilon_2 \omega_1)} - \epsilon_2 v_1 \quad (65)$$

$$t = 0$$

$$\psi = 0$$

4.8.3 Position Coordinates

$$x''_a = a e_3, \text{ where } e_3 = 1 + e \quad (61)$$

$$y''_a = z''_a = 0 \quad (62)$$

$$x'_a = x''_a$$

$$y'_a = z'_a = 0 \quad (63)$$

$$x' = x'_a$$

$$y' = z' = 0 \quad (109)$$

$$\Delta x' = \Delta y' = \Delta z' = 0 \quad (118)$$

$$\rho = x'$$

$$r = 1 + \epsilon_2 x' \quad (t = 0) \quad (8)$$

$$r^3 = (1 + \epsilon_2 x')^3$$

$$x'_s = x'_a$$

$$y'_s = z'_s = 0 \quad (113)$$

#### 4.8.4 Velocity Components

$$\dot{x}''_a = 0 \quad (64)$$

$$\dot{y}''_a = x'_a \frac{d\phi}{dt}$$

$$\dot{z}''_a = 0$$

$$\dot{y}'_a = \dot{y}''_a C_i + x''_a \dot{\Omega} \quad (66)$$

$$\dot{z}'_a = \dot{y}''_a S_i$$

$$\dot{x}'_a = 0$$

$$\dot{y}'_a = \dot{y}'_a \quad (117)$$

$$\dot{z}'_a = \dot{z}'_a$$

$$\Delta \dot{x} = 0$$

$$\dot{G} = \frac{-\left(\frac{dt_0}{d\psi} \epsilon_2 \omega_1 + \epsilon_2 \frac{dt_1}{d\psi} (1 + \epsilon_2 \omega_1)\right)}{\frac{dt}{d\psi} \frac{dt_0}{d\psi} (1 + \epsilon_2 \omega_1)} - \epsilon_2 \omega_1 - \dot{\omega}$$



$$\Delta \dot{y}' = x' \dot{G} C_i$$

$$\Delta \dot{z}' = x' \dot{G} S_i \quad (119)$$

$$\dot{\rho} = (\dot{y}'^2 + \dot{z}'^2)^{1/2}$$

$$\dot{x}'_s = 0 \quad (114)$$

$$\dot{y}'_s = x' \left[ \dot{\Omega} + \left( \frac{1}{\frac{dt}{d\psi}} + \dot{\omega} \right) C_i \right] \quad (115)$$

$$\dot{z}'_s = x' \left( \frac{1}{\frac{dt}{d\psi}} + \dot{\omega} \right) S_i \quad (116)$$

#### 4.8.5 Acceleration Components

$$E_{2x} = \epsilon_2 x$$

$$\begin{aligned} \ddot{\Delta x}' &= C_8 E_{2x} [2 + E_{2x} (3 + E_{2x})] / r^3 \\ &+ 2(\dot{\omega} + \dot{\Omega} C_i) \sqrt{\frac{(1-e)}{a(1+e)}} \\ &+ x'_s (2 C_{\omega\Omega} \dot{\omega} \dot{\Omega} C_i + C_{\omega^2\Omega^2} \dot{\omega}^2 \dot{\Omega}^2) \end{aligned} \quad (121)$$

$$\ddot{\Delta y}' = \ddot{\Delta z}' = 0 \quad (122)$$

#### 4.8.6 Initial Value of Jacobi Constant

$$C_{\text{part}} = x'^2 \left\{ \left( \frac{d\phi}{dt} \right)^2 + \dot{\Omega}^2 - 2[\epsilon \dot{\Omega} - \frac{d\phi}{dt} (\dot{\Omega} - \epsilon) \cos i] \right\} - 2 \left[ \frac{1}{x'} + C_8 x' \right] \quad (68b)$$

$$C_{\text{init}} = C_{\text{part}} - \frac{2 C_8}{\epsilon_2 r}$$

(68a)

$$x'_0 = x'$$

$$r_0 = 1 + \epsilon_2 x'$$

$$(\Delta c/c)_{\text{approx}} = (\Delta c/c)_{\text{exact}} = 0$$

Print Jacobi integral,  $C_{\text{init}}$ .

#### 4.9 SET RUN START FLAGS

IP = 1

IPRINT = 2

KHALT = 1

KR = 1

##### 4.9.1 Test Perturbation-Total

If ICH = 1, go to Step 4.10.

If ICH = 2, go to Step 4.19.4.

#### 4.10 OUTPUT OF NUMERICAL RESIDUAL PERTURBATION SOLUTION

The format will be:

t	$\Delta x'$	$\Delta \dot{x}'$	$\Delta \ddot{x}'$	$\rho$
	$\Delta y'$	$\Delta \dot{y}'$	$\Delta \ddot{y}'$	$\dot{\rho}$
$\frac{\Delta C_a}{C_{\text{init}}}$	$\Delta z'$	$\Delta \dot{z}'$	$\Delta \ddot{z}'$	

$$\frac{\Delta C_E}{C_{init}} \begin{matrix} x' & \dot{x}' & x'_a & \dot{x}'_a \\ y' & \dot{y}' & y'_a & \dot{y}'_a \\ z' & \dot{z}' & z'_a & \dot{z}'_a \end{matrix}$$

If this is the initial point (IP = 1), then go to Step 4.12; otherwise, go to Step 4.18.1.

#### 4.11 TEST HALT FLAG, KHALT

If KHALT = 1, continue numerical residual perturbation at Step 4.12.

If KHALT = 2, go to Section 4.14.1 to discontinue the solution.

If KHALT = 3, initiate numerical solution of total equations of motion at Step 4.19.

#### 4.12 CALL NUMERICAL SOLUTION SUBROUTINE FOR RESIDUAL PERTURBATION EQUATIONS

The arguments for subroutine RKSIMP are:

KR	Runge-Kutta flag
IP	Initial point flag
KC	Simpson's rule flag
t	Time
$\Delta t$	Computing interval
ICH	Perturbation-total flag
$\Delta x', \Delta y', \Delta z'$	} Position, velocity, acceleration components
$\Delta \dot{x}', \Delta \dot{y}', \Delta \dot{z}'$	
$\Delta \ddot{x}', \Delta \ddot{y}', \Delta \ddot{z}'$	

KHALT

Halt flag

$t_p$

Print time

$T_f$

Run stop time

#### 4.13 TEST RUNGE-KUTTA FLAG, KR

If KR = 1, go to Step 4.14.

If KR = 2 or 4, go to Step 4.15.

If KR = 3, go to Step 4.16.

If KR = 5, go to Step 4.17.

#### 4.14 TEST HALT FLAG, KHALT

If KHALT = 1 or 3, go to Step 4.16.

If KHALT = 2, go to Step 4.14.1.

##### 4.14.1 Print Computing Interval Selection Failed

Set  $T_f$  = last print time.

If ICH = 1, initiate numerical solution of total equations of motion at Step 4.19.

If ICH = 2, begin plotting results (Step 4.31).

#### 4.15 COMPUTE TWO-BODY EQUATIONS

##### 4.15.1 Solve Kepler's Equations

###### 4.15.1.1 Form Constant

$$t_{a32} = t_s / a_{32}$$

4.15.1.2 Compute a First Estimate for Eccentric Anomaly

$$E = t_{a32} - e_s \sin(t_{a32}) [1 + e_s \cos(t_{a32})] \quad (108)$$

$$\Delta E = 10,000$$

4.15.1.3 Compute New Guess at E

$$S_E = \sin E$$

$$C_E = \cos E$$

$$D_n = 1 + e_s C_E$$

$$E_n = \frac{t_{a32} + e_s (EC_E - S_E)}{D_n}$$

If  $E_n = E$ , go to Step 4.15.2.

If  $|E_n - E| \geq \Delta E$ , go to Step 4.15.1.4; otherwise set  $E = E_n$ ,  
 $\Delta E = |E_n - E|$  and repeat Step 4.15.1.3.

4.15.1.4 Store Maximum  $\Delta E$

$$\Delta E_{\max} = \max(|\Delta E|, \Delta E_{\max})$$

4.15.2 Compute Sine and Cosine of the True Anomaly

$$\sin v = \frac{\sqrt{1 - e_s^2} \sin E}{1 + e_s \cos E} = \frac{C_{24} S_E}{D_n} \quad (105)$$

$$\cos v = \frac{\cos E + e_s}{1 + e_s \cos E} = \frac{C_E + e_s}{D_n} \quad (104)$$

4.15.3 Compute  $\rho_s = a_s (1 + e C_E)$

(106)

4.15.4 Test Perturbation-Total Flag

If ICH = 1, continue below.

If ICH = 2, go to Step 4.28.2.

4.15.5 Compute Arguments of Node and Perigee and Their Sines and Cosines

$$\Omega = \dot{\Omega} t_{\Omega} ; \quad C_{\Omega} = \cos \Omega ; \quad S_{\Omega} = \sin \Omega$$

$$\omega = \dot{\omega} t_{\omega} ; \quad C_{\omega} = \cos \omega ; \quad S_{\omega} = \sin \omega$$

$$C_{\phi_s} = \cos v \cos \omega - \sin v \sin \omega$$

$$S_{\phi_s} = \sin v \cos \omega + \cos v \sin \omega$$

4.15.6 Compute the Precessing Two-Body Position Coordinates in the Inertial Coordinate System

$$x'_s = \rho_s (C_{\phi_s} C_{\Omega} - S_{\phi_s} S_{\Omega} C_i)$$

$$y'_s = \rho_s (C_{\phi_s} S_{\Omega} + S_{\phi_s} C_{\Omega} C_i)$$

$$z'_s = \rho_s S_{\phi_s} S_i$$

$$C_{\epsilon t} = \cos \gamma$$

$$S_{\epsilon t} = \sin \gamma$$

#### 4.16 ACCELERATION COMPONENTS FOR NUMERICAL RESIDUAL PERTURBATION EQUATIONS OF MOTION

##### 4.16.1 Computation of the Small Differences Between the Two-Body Terms

$$\text{Set } q = \frac{(x'_s + \frac{\Delta x'}{2})\Delta x' + (y'_s + \frac{\Delta y'}{2})\Delta y' + (z'_s + \frac{\Delta z'}{2})\Delta z'}{\rho_s^2} \quad (101)$$

Determine the term of lowest order in the following series which is less than  $2 \times 10^{-9}$

$$- \frac{5}{2!} q + \frac{5 \cdot 7}{3!} q^2 - \frac{5 \cdot 7 \cdot 9}{4!} q^3 + \dots$$

up to a limit of 40 terms. If the fortieth term still exceeds  $2 \times 10^{-9}$ , print its value, set final time  $t_f = P_t(\text{ITP} - 1)$ , and begin the solution of the total equations of motion at Step 4.19. If the absolute value of the nth term is  $< 2 \times 10^{-9}$  and  $n \leq 40$ , then compute

$$fq = 3q \{1 - \frac{5}{2} q (1 - \frac{7}{3} q (1 - \frac{9}{4} q (1 - \frac{11}{5} q \dots)))\} \quad (102)$$

etc., to the nth term

##### 4.16.2 Compute Numerical Residual Perturbation Acceleration Components

$$x' = x'_s + \Delta x'$$

$$y' = y'_s + \Delta y' \quad (96)$$

$$z' = z'_s + \Delta z'$$

$$\rho^2 = x'^2 + y'^2 + z'^2$$

$$r^2 = 1 + \epsilon_2 [2(x'_s C_{\epsilon t} + y'_s S_{\epsilon t}) + \epsilon_2 \rho^2] \quad (8)$$

$$r = \sqrt{r^2}$$

$$r^3 = r^2 r$$

$$E_{r_1} = C_8 \epsilon_2 / r^3$$

$$E_{r_2} = C_8 \epsilon_2 [2(x' C_{\epsilon t} + y' S_{\epsilon t}) + \epsilon_2 \rho^2] [1 + r(1 + r)] / r^3 (1 + r)$$

$$\rho_s^3 = (\rho_s)^3$$

$$\begin{aligned} \text{Compute } B_1 &= - \left\{ \frac{2(\dot{\Omega} + \dot{\omega} \cos i)}{\sqrt{a(1 - e^2)}} + \rho_s [2\dot{\omega} \dot{\Omega} + (\dot{\omega}^2 + \dot{\Omega}^2) \cos i] \right\} \\ &= C_{15} + \rho_s C_{16} \end{aligned}$$

$$\begin{aligned} \text{Compute } B_2 &= \frac{2(\dot{\Omega} \cos i + \dot{\omega})}{\sqrt{a(1 - e^2)}} + \rho_s [2\dot{\omega} \dot{\Omega} \cos i + (\dot{\omega}^2 + \dot{\Omega}^2)] \\ &= C_{17} + \rho_s C_{18} \end{aligned}$$

Following the preliminary computations, the acceleration components may be computed as:

$$\begin{aligned} \ddot{\Delta x'} &= (f_q x' - \Delta x') / \rho_s^3 - E_{r_1} x' + E_{r_2} C_{\epsilon t} \\ &\quad + S_{\phi_s} S_{\Omega} B_1 + C_{\phi_s} C_{\Omega} B_2 - (S_{\omega} S_{\Omega} C_{22} + C_{\omega} C_{\Omega} C_{23}) \quad (98) \end{aligned}$$

$$\begin{aligned} \ddot{\Delta y'} &= (f_q y' - \Delta y') / \rho_s^3 - E_{r_1} y' + E_{r_2} S_{\epsilon t} \\ &\quad - S_{\phi_s} C_{\Omega} B_1 + C_{\phi_s} S_{\Omega} B_2 + (S_{\omega} C_{\Omega} C_{22} + C_{\omega} S_{\Omega} C_{23}) \quad (99) \end{aligned}$$



$$\begin{aligned} \ddot{\Delta z'} &= (f_q z' - \Delta z')/\rho_s^3 - E_{r_1} z' + S_{\phi_s} S_i (C_{13} + \rho_s \omega^2) - e s_i S_\omega C_{13} \\ &= (f_q z' - \Delta z')/\rho_s^3 - E_{r_1} z' + S_{\phi_s} (C_{21} + \rho_s C_{19}) - S_\omega C_{20} \end{aligned} \quad (100)$$

where  $C_\Omega$ ,  $S_\Omega$ ,  $C_\omega$ ,  $S_\omega$ ,  $C_{\phi_s}$ , and  $S_{\phi_s}$  are defined in Section 4.15.5,  $C_{\epsilon t}$  and  $S_{\epsilon t}$  in Section 4.15.6, and  $C_{13}$  and  $C_{19}$  through  $C_{23}$  in Section 4.4.

#### 4.16.3 Test Runge-Kutta Flag, KR

If KR = 1, go to Step 4.16.4.

If KR > 1, go to Step 4.12.

#### 4.16.4 Test Print Flag, IPRINT

If IPRINT = 2, go to Step 4.12.

If IPRINT  $\neq$  2, go to Step 4.17.

### 4.17 COMPUTATIONS FOR PRINT ONLY

#### 4.17.1 Reduce All Angles that Increase with Time by Integer Multiples of $2\pi$

$$T_s = T_s, \text{ modulo } P_{ss} \quad (108)$$

$$T_\Omega = T_\Omega, \text{ modulo } P_\Omega \quad (39a, b)$$

$$T_\omega = T_\omega, \text{ modulo } P_\omega \quad (72a)$$

$$\gamma = \gamma, \text{ modulo } 2\pi$$

4.17.1.1 Test Value of  $T_\psi$

If  $t_\psi \geq T_1$ , go to Step 4.17.1.2.

If  $t_\psi > T_1$ , go to Step 4.17.2.1.

(28b)

4.17.1.2 Set  $T_\psi = T_\psi - T_1$

Recompute  $T_1$

$$\phi_1 = [\phi_1 + \frac{2\pi}{1 + \epsilon_2 \omega_1} - \epsilon_2 v_1 T_1] \text{ modulo } 2\pi \quad (49)$$

$$\alpha_1 = \{ [(\alpha_1 + 2\pi p_1) \text{ modulo } 2\pi] + 2\pi p_1 \} \text{ modulo } 2\pi \quad (28a)$$

$$S_{\alpha_1} = \sin \alpha_1$$

$$C_{\alpha_1} = \cos \alpha_1$$

$$\alpha_2 = \{ [(\alpha_1 + 2\pi p_1) \text{ modulo } 2\pi] + 2\pi p_1 \} \text{ modulo } 2\pi \quad (28g)$$

$$\theta_1 = (\theta_1 + 2\pi C_{11}) \text{ modulo } 2\pi \quad (28c)$$

where  $C_{11}$  is defined in Section 4.4.

$$S_{\theta_1} = \sin \theta_1$$

$$C_{\theta_1} = \cos \theta_1$$

$$\theta_2 = (\theta_1 + 2\pi C_{11}) \text{ modulo } 2\pi \quad (28f)$$

$$T_1 = 2\pi a_{32} + C_1 (\sin \theta_2 - S_{\theta_1}) - 1.375 \epsilon_2 a_{92} (\sin \alpha_2 - S_{\alpha_1}) \quad (28e)$$

where  $C_1$  is defined in Section 4.4.

$$\beta = (\beta + 2\pi p_3) \text{ modulo } 2\pi \quad (29)$$

#### 4.17.2 Solution of the Perturbed Kepler Equation

As in the unperturbed case,  $\psi$  (true anomaly) is an implicit function of time and  $\psi(t)$  must be found by iteration.

##### 4.17.2.1 Compute Initial Guess at E

$$E = t_{\psi}/a_{32} - e \sin (t_{\psi}/a_{32})(1 - e \cos t_{\psi}/a_{32})$$

$$\Delta E_p = 10,000$$

##### 4.17.2.2 Compute Sine and Cosine of $\psi$

$$S_E = \sin E$$

$$C_E = \cos E$$

$$D_n = 1 + e C_E$$

$$S_{\psi} = \sin \psi = \frac{e_2 S_E}{D_n} \quad (68aa)$$

where  $e_2 = \sqrt{1 - e^2}$  is computed in Section 4.4

$$C_{\psi} = \cos \psi = \frac{C_E + e}{D_n} \quad (68ab)$$

#### 4.17.2.3 Compute $\psi$

##### 4.17.2.3.1 Test $\sin \psi$

If  $\sin \psi = 0$ , go to Step 4.17.2.3.2; otherwise go to Step 4.17.2.3.3.

##### 4.17.2.3.2 Set $\psi = 0$ or $\pi$

$$\psi = \frac{\pi}{2} (1 - 1 \text{ sign } C_{\psi})$$

Go to Step 4.17.3.

##### 4.17.2.3.3 Compute principal value of $\psi$ ( $-\pi/2 < \psi < \pi/2$ )

$$\psi = \tan^{-1} \left( \frac{\sin \psi}{\cos \psi} \right)$$

##### 4.17.2.3.4 Test $\cos \psi$

If  $\cos \psi < 0$ , then set  $\psi = \psi + \pi$  and go to Step 4.17.2.4.

If  $\cos \psi = 0$ , go to Step 4.17.2.3.5.

If  $\cos \psi > 0$ , go to Step 4.17.2.3.6.

##### 4.17.2.3.5 Set $\psi = \pi/2$ or $3\pi/2$

$$\psi = \frac{\pi}{2} (2 - 1 \text{ sign } S_{\psi})$$

Go to Step 4.17.2.4.

##### 4.17.2.3.6 Place $\psi$ in first or third quadrant

$$\psi = \text{sign } E \left[ \psi + \pi(1 - 1 \text{ sign } S_{\psi} + 2 \text{ integer part of } \frac{E}{2\pi}) \right] \quad (68ah)$$

##### 4.17.2.4 Compute New Guess at E

Compute  $(2p_1 - 1)\psi$

$$C_{11\psi} = C_{11} \cdot \psi$$

Compute  $\sin (2p_1 - 1)\psi$

$$S_{C11\psi} = \sin C_{11\psi}$$

Compute  $\cos (2p_1 - 1)\psi$

$$C_{C11\psi} = \cos C_{11\psi}$$

Compute  $\sin (2p_1 - 1)(\psi + \theta_1)$

$$S_{P1} = S_{C11\psi}C_{\theta_1} + C_{C11\psi}S_{\theta_1}$$

Compute  $\cos (2p_1 - 1)(\psi + \theta_1)$

$$C_{P1} = C_{C11\psi}C_{\theta_1} - S_{C11\psi}S_{\theta_1}$$

Compute  $\cos 2p_1\psi$

$$C_{2P1\psi} = C_{C11\psi}C_{\psi} - S_{C11\psi}S_{\psi}$$

Compute  $\sin 2p_1\psi$

$$S_{2P1\psi} = S_{C11\psi}C_{\psi} + C_{C11\psi}S_{\psi}$$

Compute  $\cos 2p_1(\psi + \alpha_1)$

$$C_{P1\psi} = C_{2P1\psi}C_{\alpha_1} - S_{2P1\psi}S_{\alpha_1}$$

Compute  $\sin 2p_1(\psi + \alpha_1)$

$$S_{P1\psi} = S_{2P1\psi}C_{\alpha_1} + C_{2P1\psi}S_{\alpha_1}$$

$$\epsilon^2 \frac{dt_1}{d\psi} = -C_1(C_\psi - C_{P1}) + \epsilon_2 a^{9/2} \left[ \frac{10}{3} C_\psi - 2.75 C_{P1\psi} \right] \quad (68ad)$$

$$\epsilon^2 \frac{dt_1}{dE} = \frac{e_2}{D_n} \epsilon^2 \frac{dt_1}{d\psi} \quad (68ae)$$

$$\epsilon^2 t_1 = -C_1(S_\psi - S_{P1} + S_{\theta_1}) + \epsilon_2 a^{9/2} \left[ \frac{10}{3} S_\psi - 1.375(S_{P1\psi} - S_{\alpha_1}) \right] \quad (68af)$$

$$E_n = \frac{a_{32} e(E \cos E - \sin E) + E \frac{dt_1}{dE} + t_\psi - \epsilon^2 t_1}{a_{32}(1 + e \cos E) + \epsilon^2 \frac{dt_1}{dE}} \quad (68ag)$$

If  $E_n = E$ , go to Step 4.17.3.

If  $|E_n - E| \geq \Delta E_p$ , go to Step 4.17.2.5.

Otherwise set

$$\Delta E_p = |E_n - E|$$

and

$$E = E_n$$

and go to Step 4.17.2.2.

4.17.2.5 Store Maximum Final Value of  $\Delta E_p$

$$\Delta E_{Pmax} = \max(\Delta E_{Pmax}, |\Delta E_p|)$$

4.17.3 Compute Auxiliary Functions of  $\psi$  and  $t$

$$C_{P_3} = \cos p_3 \psi \cos \beta - \sin p_3 \psi \sin \beta \quad (29)$$

$$S_{P_3} = \sin p_3 \psi \cos \beta + \cos p_3 \psi \sin \beta$$

$$s = \frac{1 - eC_\psi}{a_e} + C_2(C_\psi - C_{P1}) + \epsilon_2 a_2 [(2 - 5C_\psi)/3 + C_{P1\psi}] \quad (18, 19, 22)$$

where  $C_2$  is defined in Section 4.4,  $C_\psi$  and  $S_\psi$  in Section 4.17.2.2, and  $C_{P1}$  and  $C_{P1\psi}$  in Section 4.17.2.4.

Test the value of  $s$ .

If  $s \leq 0$ , set and print the error code = 7 and go to Step 4.3.2.1.2.

If  $s > 0$ , continue below.

#### 4.17.4 Compute Double Prime Coordinates

$$\phi = \frac{\psi}{1 + \epsilon_2 \omega_1} - \epsilon_2 v_1 t_\psi \quad (49)$$

$$x''_a = (\cos \phi \cos \phi_1 - \sin \phi \sin \phi_1)/s \quad (48)$$

$$y''_a = (\sin \phi \cos \phi_1 + \cos \phi \sin \phi_1)/s$$

$$z''_a = C_3(S_{P_3} - p_3 S_\psi) \quad (24)$$

where  $C_3$  is defined in Section 4.4,  $S_{P_3}$  in Section 4.17.3, and  $S_\psi$  in Section 4.17.2.2.

#### 4.17.5 Compute Node

$$\Omega = C_6^T \dot{\Omega} \quad \text{where } C_6 = \dot{\Omega} \quad (39)$$

$$C_\Omega = \cos \Omega$$

$$S_{\Omega} = \sin \Omega$$

4.17.6 Compute Approximate Coordinates in Plane of Planets

$$\begin{aligned} x'_a &= x''_a C_{\Omega} - y''_a S_{\Omega} C_i + z''_a S_{\Omega} S_i \\ y'_a &= x''_a S_{\Omega} + y''_a C_{\Omega} C_i - z''_a C_{\Omega} S_i \\ z'_a &= y''_a S_i + z''_a C_i \end{aligned} \quad (50)$$

where  $S_i = \sin i$  and  $C_i = \cos i$

4.17.7 Compute Time Derivatives of  $\psi$ ,  $\phi$ , and  $s$

$$\frac{dt}{d\psi} = C_9 / (1 - eC_{\psi})^2 - C_1(C_{\psi} - C_{Pl}) + \epsilon_2 a_{92} \left( \frac{10}{3} C_{\psi} - 2.75 C_{Pl\psi} \right) \quad (52)$$

where  $C_9$  is defined in Section 4.4 and  $C_{\psi}$ ,  $C_{Pl}$ , and  $C_{Pl\psi}$  are defined in Section 4.17.2.

4.17.8 Test Value of  $\frac{dt}{d\psi}$

If  $\frac{dt}{d\psi} = 0$ , go to Step 4.17.8.1.

If  $\frac{dt}{d\psi} \neq 0$ , go to Step 4.17.9.

4.17.8.1 Set the Error Code = 8 and Go to Step 4.3.2.1.2

4.17.9 Compute Time Derivatives of  $\phi$  and  $s$

$$\frac{d\psi}{dt} = \frac{1}{\frac{dt}{d\psi}} \quad (56)$$



$$\frac{d\phi}{dt} = \frac{d\psi}{dt (1 + \epsilon_2 \omega_1)} - \epsilon_2 v_1 \quad (55)$$

$$\frac{ds}{d\psi} = \frac{eS_\psi}{a_e} - C_2 (S_{P1} - S_\psi) + \epsilon_2 a_2 \left( \frac{5}{3} S_\psi - 2 S_{P1\psi} \right) \quad (51)$$

where  $S_{P1}$ ,  $S_\psi$ , and  $S_{P1\psi}$  are defined in Section 4.17.2 and  $C_2$  is defined in Section 4.4.

$$\frac{ds}{dt} = \frac{ds}{d\psi} \frac{d\psi}{dt}$$

#### 4.17.10 Compute Approximate Velocity Components in Orbit Plane

$$\begin{aligned} \dot{x}_a'' &= y_a'' \frac{d\phi}{dt} - \frac{x_a''}{s} \frac{ds}{dt} \\ \dot{y}_a'' &= x_a'' \frac{d\phi}{dt} - \frac{y_a''}{s} \frac{ds}{dt} \end{aligned} \quad (54)$$

$$\dot{z}_a'' = C_3 (C_{P3} - C_\psi) \frac{d\psi}{dt}$$

where  $C_3$  is defined in Section 4.4,  $C_{P3}$  in Section 4.17.3, and  $C_\psi$  in Section 4.17.2.2.

#### 4.17.11 Compute $\rho$

$$\rho = \sqrt{\rho^2}$$

##### 4.17.11.1 Test Perturbation-Total Flag

If ICH = 1, go to Step 4.17.12.

If ICH = 2, go to Step 4.17.14.1.

4.17.12 Compute Approximate Velocity Components in Inertial System

$$\dot{x}'_a = C_6 (-x''_a S_\Omega - y''_a C_\Omega C_i + z''_a C_\Omega S_i) + \dot{x}''_a C_\Omega - \dot{y}''_a S_\Omega C_i + \dot{z}''_a S_\Omega S_i$$

$$\dot{y}'_a = C_6 (x''_a C_\Omega - y''_a S_\Omega C_i + z''_a S_\Omega S_i) + \dot{x}''_a S_\Omega + \dot{y}''_a C_\Omega C_i - \dot{z}''_a C_\Omega S_i$$

$$\dot{z}'_a = \dot{y}''_a S_i + \dot{z}''_a C_i \quad (57)$$

where  $C_6 = \dot{\Omega}$ ,  $C_\Omega = \cos \Omega$ ,  $S_\Omega = \sin \Omega$ ,  $C_i = \cos i$ , and  $S_i = \sin i$  as computed previously.

4.17.13 Compute Precessing Keplerian Velocity Components in Inertial System

$$\dot{x}'_s = - [S_{\phi_s} C_\Omega + C_{\phi_s} S_\Omega C_i - e(S_\omega C_\Omega + C_\omega S_\Omega C_i)] / a_e^{1/2} \quad (91)$$

$$- \rho_s [\dot{\Omega}(C_{\phi_s} S_\Omega + S_{\phi_s} C_\Omega C_i) + \dot{\omega}(S_{\phi_s} C_\Omega + C_{\phi_s} S_\Omega C_i)]$$

$$\dot{y}'_s = - [S_{\phi_s} S_\Omega - C_{\phi_s} C_\Omega C_i - e(S_\omega S_\Omega - C_\omega C_\Omega C_i)] / a_e^{1/2} \quad (92)$$

$$+ \rho_s [\dot{\Omega}(C_{\phi_s} C_\Omega - S_{\phi_s} S_\Omega C_i) - \dot{\omega}(S_{\phi_s} S_\Omega - C_{\phi_s} C_\Omega C_i)]$$

$$\dot{z}'_s = [(C_{\phi_s} - eC_\omega) / a_e^{1/2} + \rho_s \dot{\omega} C_{\phi_s}] S_i \quad (93)$$

where  $C_\Omega$ ,  $S_\Omega$ ,  $C_\omega$ ,  $S_\omega$ , and  $C_{\phi_s}$  are defined in Section 4.15.5 and  $a_e^{1/2}$ ,

$C_i$ , and  $S_i$  are computed in Section 4.4.

4.17.14 Compute Total Inertial Velocity Components and Vector Magnitudes

$$\dot{x}' = \dot{x}'_s + \Delta\dot{x}'$$

$$\dot{y}' = \dot{y}'_s + \Delta\dot{y}' \quad (117)$$

$$\dot{z}' = \dot{z}'_s + \Delta\dot{z}'$$

4.17.14.1 Compute Inertial Velocity

$$v^2 = \dot{z}'^2 + \dot{y}'^2 + \dot{x}'^2$$

$$\dot{\rho} = \sqrt{v^2}$$

4.17.15 Compute Jacobi Constant Fractional Errors

$$r = \sqrt{r^2}$$

$$\Delta C_E = v^2 + 2[\epsilon_1(\dot{x}'y' - \dot{y}'x') - 1/\rho - C_8(x'C_{\epsilon t} + y'S_{\epsilon t})] \quad (16)$$

$$-C_{\text{part}} + 2C_8 \frac{[2(x'C_{\epsilon t} + y'S_{\epsilon t} - x_o) + \epsilon_2(\rho^2 - x_o^2)]}{rr_o(r_o + r)}$$

where  $v$  is defined in Section 4.17.14.1,  $C_8$  in Section 4.4, and  $C_{\epsilon t}$  and  $S_{\epsilon t}$  in Section 4.15.6.

$$\frac{\Delta C_E}{C} = \Delta C_E / C_{\text{init}}$$

$$\rho_a^2 = x_a'^2 + y_a'^2 + z_a'^2$$

$$r_a = \{1 + \epsilon_2 [2(x'_a C_{\epsilon t} + y'_a S_{\epsilon t}) + \epsilon_2 \rho_a^2]\}^{1/2}$$

$$\Delta C_a = \dot{x}'_a{}^2 + \dot{y}'_a{}^2 + \dot{z}'_a{}^2 + 2[\epsilon_1 (\dot{x}'_a y'_a - \dot{y}'_a x'_a) - \frac{1}{\sqrt{\rho_a^2}}$$

$$- C_8 (x'_a C_{\epsilon t} + y'_a S_{\epsilon t})] - C_{\text{part}} \quad (16)$$

$$+ 2C_8 \frac{[2(x'_a C_{\epsilon t} + y'_a S_{\epsilon t} - x_o) + \epsilon_2 (\rho_a^2 - x_o^2)]}{r_a r_o (r_o + r_a)}$$

$$\frac{\Delta C_a}{C} = \frac{\Delta C_a}{C_{\text{init}}}$$

$$\rho_s^2 = x_s^2 + y_s^2 + z_s^2$$

$$\rho_s = \sqrt{\rho_s^2}$$

$$r_s = \{1 + \epsilon_2 [2(x'_s C_{\epsilon t} + y'_s S_{\epsilon t}) + \epsilon_2 \rho_s^2]\}^{1/2}$$

$$\frac{\Delta C_s}{C} = \{\dot{x}'_s{}^2 + \dot{y}'_s{}^2 + \dot{z}'_s{}^2 + 2[\epsilon_1 (\dot{x}'_s y'_s - \dot{y}'_s x'_s) - 1/\rho_s$$

$$- C_8 (x'_s C_{\epsilon t} + y'_s S_{\epsilon t})] - C_{\text{part}}$$

$$+ 2C_8 \frac{[2(x'_s C_{\epsilon t} + y'_s S_{\epsilon t} - x_o) + \epsilon_2 (\rho_s^2 - x_o^2)]}{r_s r_o (r_o + r_s)} \} / C_{\text{init}}$$

If KHALT = 1, go to Step 4.17.16.

If KHALT = 2, go to Step 4.14.1

If KHALT = 3, go to Step 4.10.

4.17.16 Test if Printing is Desired

If  $W_{13} \leq 0$ , go to Step 4.18.1.

If  $W_{13} > 0$ , go to Step 4.10.

4.18 SAVE PLOT VALUES AND TEST IF PLOTTING ARRAY IS FULL

4.18.1 Save Values from Numerical Solution of Total Equations of Motion

$$P_t(\text{ITP}) = t$$

$$P_t(\text{ITP} + 1) = T_f$$

$$\text{IT} = 13 \text{ITP} - 12$$

$$a_p(\text{IT}) = \left| \frac{\Delta C_a}{C} \right|$$

$$a_p(\text{IT} + 1) = \left| \frac{\Delta C_s}{C} \right|$$

$$a_p(\text{IT} + 2) = x'$$

$$a_p(\text{IT} + 3) = y'$$

$$a_p(\text{IT} + 4) = \left| \frac{\Delta C_E}{C} \right|$$

$$a_p(\text{IT} + 5) = [(x'_a - x')^2 + (y'_a - y')^2 + (z'_a - z')^2]^{1/2} / a$$

$$a_p(\text{IT} + 6) = [(x'_s - x')^2 + (y'_s - y')^2 + (z'_s - z')^2]^{1/2} / a$$

$$a_p(\text{IT} + 7) = z'$$

$$a_p(\text{IT} + 8) = \dot{x}'$$

$$a_{p(IT + 9)} = a_{12} [(\dot{x}'_a - \dot{x}')^2 + (\dot{y}'_a - \dot{y}')^2 + (\dot{z}'_a - \dot{z}')^2]^{1/2}$$

$$a_{p(IT + 10)} = a_{12} [(\dot{x}'_s - \dot{x}')^2 + (\dot{y}'_s - \dot{y}')^2 + (\dot{z}'_s - \dot{z}')^2]^{1/2}$$

$$a_{p(IT + 11)} = \dot{y}'$$

$$a_{p(IT + 12)} = \dot{z}'$$

Increment index of plot arrays.

#### 4.18.2 Test Value of ITP

If  $ITP > 700$ , set  $T_f = P_{t(ITP)}$ , and go to Step 4.19.

If  $ITP \leq 700$ , set  $ITP = ITP + 1$ , set  $P_{t(ITP)} = T_f$ , and to to Step 4.11.

### 4.19 INITIATE NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTION

#### 4.19.1 Set Perturbation-Total Flag

$$ICH = 2$$

$$KT = 0$$

#### 4.19.2 Compute Initial Conditions for Numerical Solution of Total Equations of Motion

$x' = ae_3$  where  $e_3$  is computed in Section 4.8.

$$y' = 0$$

$$z' = 0$$

$$\gamma = 0$$

$$\rho = x' \quad (59)$$

$$t = 0$$

$$\frac{dt}{d\psi} = \frac{a_{32} e^{1.5}}{\sqrt{1-e}} + 7 \epsilon_2 a_{92} / 12$$

where  $a_{32}$ ,  $a_{92}$ , and  $\epsilon_2$  are computed in Section 4.4.

$$\frac{d\psi}{dt} = \frac{1}{\frac{dt}{d\psi}}$$

$$\frac{d\phi}{dt} = \frac{d\psi}{dt(1 + \epsilon^2 \omega_1)} = \epsilon_2 v_1 \quad (65)$$

$$\dot{x}' = 0$$

$$\dot{y}' = x' \left( \frac{d\phi}{dt} C_i + C_6 \right) \quad (66), (64)$$

$$\dot{z}' = x' \frac{d\phi}{dt} S_i \quad (66), (64)$$

$$\dot{\rho} = (\dot{y}'^2 + \dot{z}'^2)^{1/2}$$

$$r^3 = (1 + \epsilon_2 x')^3$$

$$E_{2x} = \epsilon_2 x'$$

$$\ddot{x}' = -\frac{1}{x'^2} + C_8 E_{2x} (2 + E_{2x} (3 + E_{2x})) / r^3$$

where  $C_8$  is computed in Section 4.4.

$$\ddot{y}' = \ddot{z}' = \frac{\Delta C}{C} = 0 \quad (5)$$

4.19.3 Compute Semi-Major Axis and Eccentricity of Initial Osculating Ellipse

$$a_s = \frac{1}{\frac{2}{\rho} - \dot{\rho}^2} \quad (124)$$

$$e_s = \frac{x'}{a_s} - 1 \quad (125)$$

$$a_{s32} = (a_s)^{1.5}$$

$$P_{ss} = 2\pi a_{s32}$$

$$a_{e1/2} = \sqrt{a_s(1 - e_s^2)}$$

Go to Step 4.9.

4.19.4 Print Header for Numerical Solution of Total Equations of Motion

4.20 PRINT NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTIONS

$$t' \quad x' \quad \dot{x}' \quad \ddot{x}' \quad \rho$$

$$y' \quad \dot{y}' \quad \ddot{y}' \quad \dot{\rho}$$

$$\frac{\Delta C}{C} \quad z' \quad \dot{z}' \quad \ddot{z}'$$

$$\frac{\Delta C_a}{C} \quad \frac{\Delta C_s}{C} \quad \frac{\Delta C_c}{C} \quad \frac{\Delta C_F}{C} \quad \frac{\Delta C_E}{C}$$

$$\Delta \rho_a \quad \Delta \rho_s \quad \Delta \rho_c \quad \Delta \rho_F$$

$$\dot{\Delta \rho}_a \quad \dot{\Delta \rho}_s \quad \dot{\Delta \rho}_c \quad \dot{\Delta \rho}_F$$



#### 4.21 TEST HALT FLAG, KHALT

If KHALT = 1, continue below.

If KHALT = 2, go to Step 4.14.1.

If KHALT = 3, go to Step 4.31.

#### 4.22 CALL NUMERICAL SOLUTION SUBROUTINE FOR TOTAL EQUATIONS OF MOTIONS

The arguments for subroutine RKSIMP are:

KR

IP

KC

t

$\Delta t$

ICH

$x', y', z'$

$\dot{x}', \dot{y}', \dot{z}'$

$x'', y'', z''$

KHALT

$t_{PT}$

$T_p$

4.23 TEST RUNGE-KUTTA FLAG, KR

If KR = 1, continue below.

If KR = 2, 3, or 4, go to Step 4.25.

If KR = 5, go to Step 4.28.1.

4.24 TEST HALT FLAG, KHALT

If KHALT = 2, set  $t_f = P_t(\text{ITP} - 1)$  and go to Step 4.14.1.

If KHALT = 1 or 3, continue below.

4.25 COMPUTE TOTAL EQUATIONS OF MOTION

$$\rho^2 = x'^2 + y'^2 + z'^2 \quad (7)$$

$$\rho^3 = (\rho^2)^{1.5}$$

$\gamma = \gamma$  modulo  $2\pi$   $\gamma = \epsilon t$  is developed in subroutine RKSIMP

$$C_{\epsilon t} = \cos \gamma$$

$$S_{\epsilon t} = \sin \gamma$$

$$r^2 = 1 + \epsilon_4 \rho^2 + 2\epsilon_2 (x' C_{\epsilon t} + y' S_{\epsilon t}) \quad (8)$$

$$r = (r^2)^{.5}$$

$$r^3 = r^2 r$$

$E_{r_1} = C_8 \epsilon_2 / r^3$   $C_8$  is defined in Section 4.4.

$$E_{r_2} = C_8 \epsilon_2 [2(x' C_{\epsilon t} + y' S_{\epsilon t}) + \epsilon_2 \rho^2] \frac{[1 + r(1 + r)]}{r^3(1 + r)} \quad (5)$$

$$\left. \begin{aligned}
 \ddot{x}' &= -\frac{x'}{\rho^3} - E_{r_1} x' + E_{r_2} C_{\epsilon t} \\
 \ddot{y}' &= -\frac{y'}{\rho^3} - E_{r_1} y' + E_{r_2} S_{\epsilon t} \\
 \ddot{z}' &= -\frac{z'}{\rho^3} - E_{r_1} z'
 \end{aligned} \right\} (5)$$

When the expression for  $E_{r_1}$  and  $E_{r_2}$  are substituted in the equations above, one arrives at a set of differential equations identical to (5) except that the last term on the right side of (5) has been put over the common demoninator  $r^3(1+r)$  in order to cancel the large terms in the numerator.

#### 4.26 TEST RUNGE-KUTTA FLAG, KR

If KR = 1, continue below.

If KR = 2, 3, 4, or 5, go to Step 4.22.

#### 4.27 TEST PRINT FLAG, IPRINT

If IPRINT = 1, continue below.

If IPRINT = 2 or 3, go to Step 4.22.

#### 4.28 COMPUTE DATA FOR PRINT AND FOR PLOT ROUTINE

##### 4.28.1 Compute Error in Jacobi Integral and Limit $t_s$ for Numerical Solution of Total Equations

$$\rho = (\rho^2)$$

$$r = (r^2)^{.5}$$

$$v^2 = \dot{x}'^2 + \dot{y}'^2 + \dot{z}'^2$$

$$\dot{\rho} = (v^2)^{.5}$$

$$\begin{aligned} \Delta C_c = v^2 + 2[\epsilon_1(\dot{x}'y' - \dot{y}'x') - \frac{1}{\rho} - C_8(x'C_{et} + y'S_{et}) - C_{part} \\ + 2C_8 \frac{[2(x'C_{et} + y'S_{et} - x_o) + \epsilon_2(\rho^2 - x_o^2)]}{rr_o(r_o + r)}] \end{aligned} \quad (16)$$

where  $C_8$  is computed in Section 4.4 and  $C_{et}$  and  $S_{et}$  are computed in Section 4.25.

$$\frac{\Delta C_c}{C} = \frac{\Delta C_c}{C_{init}}$$

$$T_s = T_s \text{ modulo } P_{ss}$$

Go to Step 4.15 for solution of Kepler's equation.

#### 4.28.2 Compute the Fixed Orbit Values

$$x'_f = \rho_s C_v$$

$$y'_f = \rho_s S_v C_i$$

$$z'_f = \rho_s S_v S_i$$

$$\dot{x}'_f = -S_v/a_e 1/2 \quad (127)$$

$$\dot{y}'_f = \frac{(C_v - e_s)}{a_{e_s} 1/2} C_i$$

$$\dot{z}'_f = \frac{(C_v - e_s)}{a_{e_s} 1/2} S_i$$

$$\rho_f^2 = x_f'^2 + y_f'^2 + z_f'^2$$

$$\rho_f = (\rho_f^2)^{.5}$$

$$r_f = \{ 1 + \epsilon_2[2(x_f' C_{\epsilon t} + y_f' S_{\epsilon t}) + \epsilon_2 \rho_f^2] \}^{.5}$$

$$\begin{aligned} \Delta C_f = & \dot{x}_f'^2 + \dot{y}_f'^2 + \dot{z}_f'^2 + 2[\epsilon_1(\dot{x}_f' y_f' - \dot{y}_f' x_f')] - \frac{1}{\rho_f} \\ & - C_8(x_f' C_{\epsilon t} + y_f' S_{\epsilon t}) - C_{\text{part}} \\ & + 2C_8 \frac{[2(x_f' C_{\epsilon t} + y_f' S_{\epsilon t} - x_o) + \epsilon_2(\rho_f^2 - x_o^2)]}{r_f r_o (r_o + r_f)} \end{aligned} \quad (16)$$

where  $C_8$  is computed in Section 4.4 and  $C_{\epsilon t}$  and  $S_{\epsilon t}$  are computed in Section 4.25.

$$IT = 13 \text{ ITP} - 12$$

$$\frac{\Delta C_f}{C} = \frac{\Delta C_f}{C_{\text{init}}}$$

$$\Delta \rho_f = \frac{1}{a} [(x_f' - a_{p(IT+2)})^2 + (y_f' - a_{p(IT+3)})^2 + (z_f' - a_{p(IT+7)})^2]^{1/2}$$

$$\begin{aligned} \Delta \dot{\rho}_f = & a^{1/2} [(\dot{x}_f' - a_{p(IT+8)})^2 + (\dot{y}_f' - a_{p(IT+11)})^2 \\ & + (\dot{z}_f' - a_{p(IT+12)})^2]^{1/2} \end{aligned}$$

$$\begin{aligned} \Delta \rho_c = & [(x' - a_{p(IT+2)})^2 + (y' - a_{p(IT+3)})^2 \\ & + (z' - a_{p(IT+7)})^2]^{1/2} / a \end{aligned}$$

$$\begin{aligned} \Delta \dot{\rho}_c = & a^{1/2} [(\dot{x}' - a_{p(IT+8)})^2 + (\dot{y}' - a_{p(IT+11)})^2 \\ & + (\dot{z}' - a_{p(IT+12)})^2]^{1/2} \end{aligned}$$

#### 4.28.3 Set Quantities Just Computed Equal to Values in Plot Array

$$a_p(IT + 2) = |\Delta C_c / C|$$

$$a_p(IT + 3) = |\Delta C_f / C|$$

$$a_p(IT + 7) = \Delta \rho_c$$

$$a_p(IT + 8) = \Delta \rho_f$$

$$a_p(IT + 11) = \dot{\Delta \rho}_c$$

$$a_p(IT + 12) = \dot{\Delta \rho}_f$$

#### 4.29 TEST VALUE OF HALT FLAG, KHALT

If KHALT = 3, go to Step 4.20.

If KHALT = 2, go to Step 4.14.1.

If KHALT = 1, continue below.

#### 4.30 TEST, IF PRINTING IS DESIRED

If  $W_{13} \leq 0$ , go to Step 4.22.

If  $W_{13} > 0$ , go to Step 4.20.

#### 4.31 CALL THE PLOT SUBROUTINE

The arguments for subroutine PLOT are:

$a^{3/2}$

$T_f$

a

e

i<sup>o</sup>

μ

C<sub>init</sub>

P<sub>ss</sub>

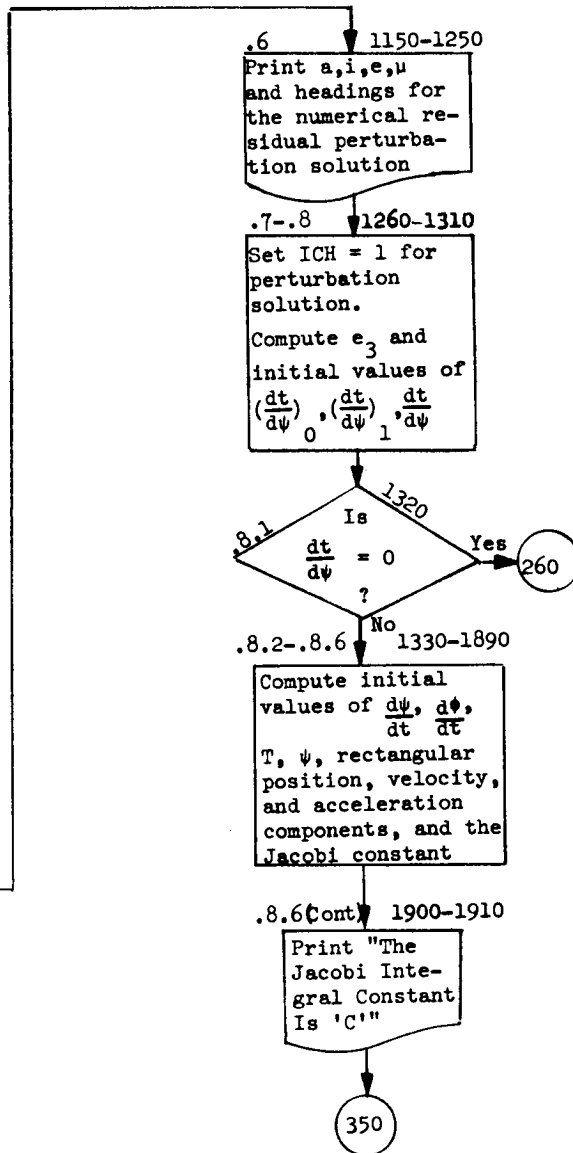
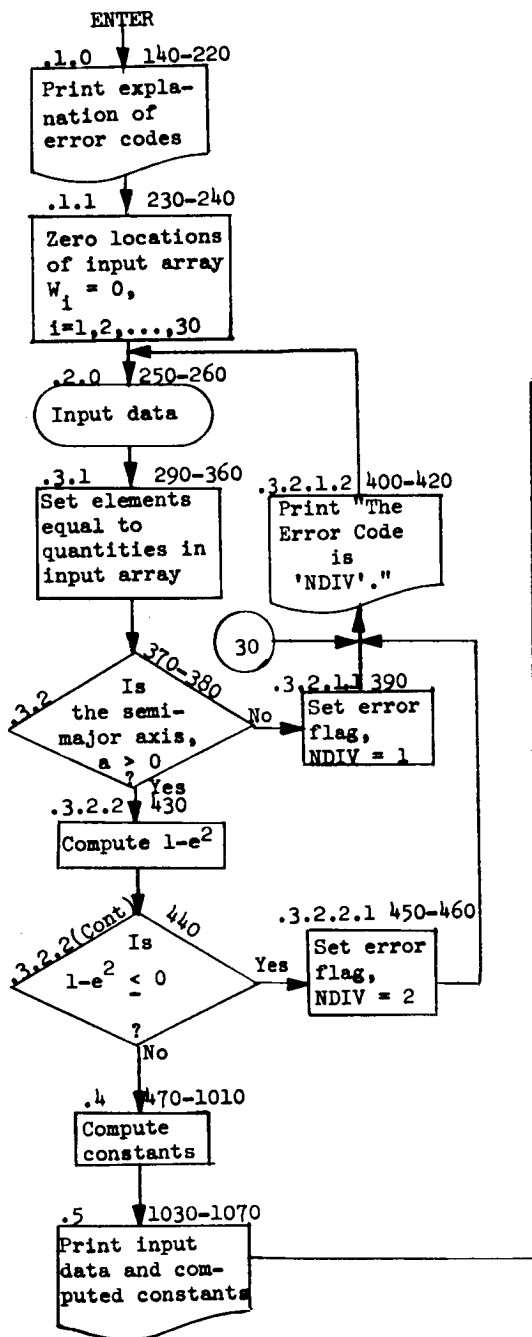
4.32 RUN IS COMPLETED, CALL THE NEXT CASE AT STEP 4.2

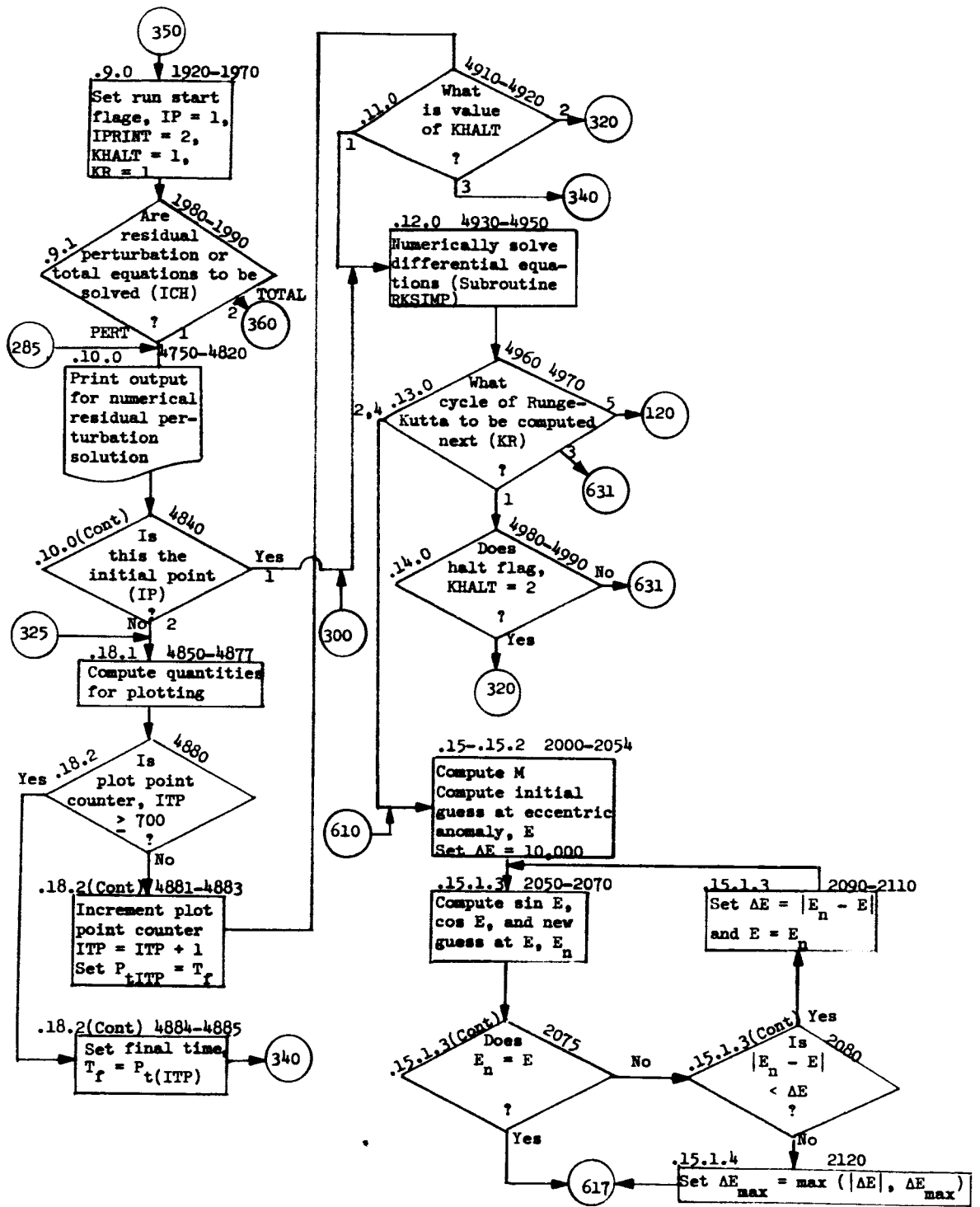
Section 5  
MAIN PROGRAM DETAIL FLOW CHARTS

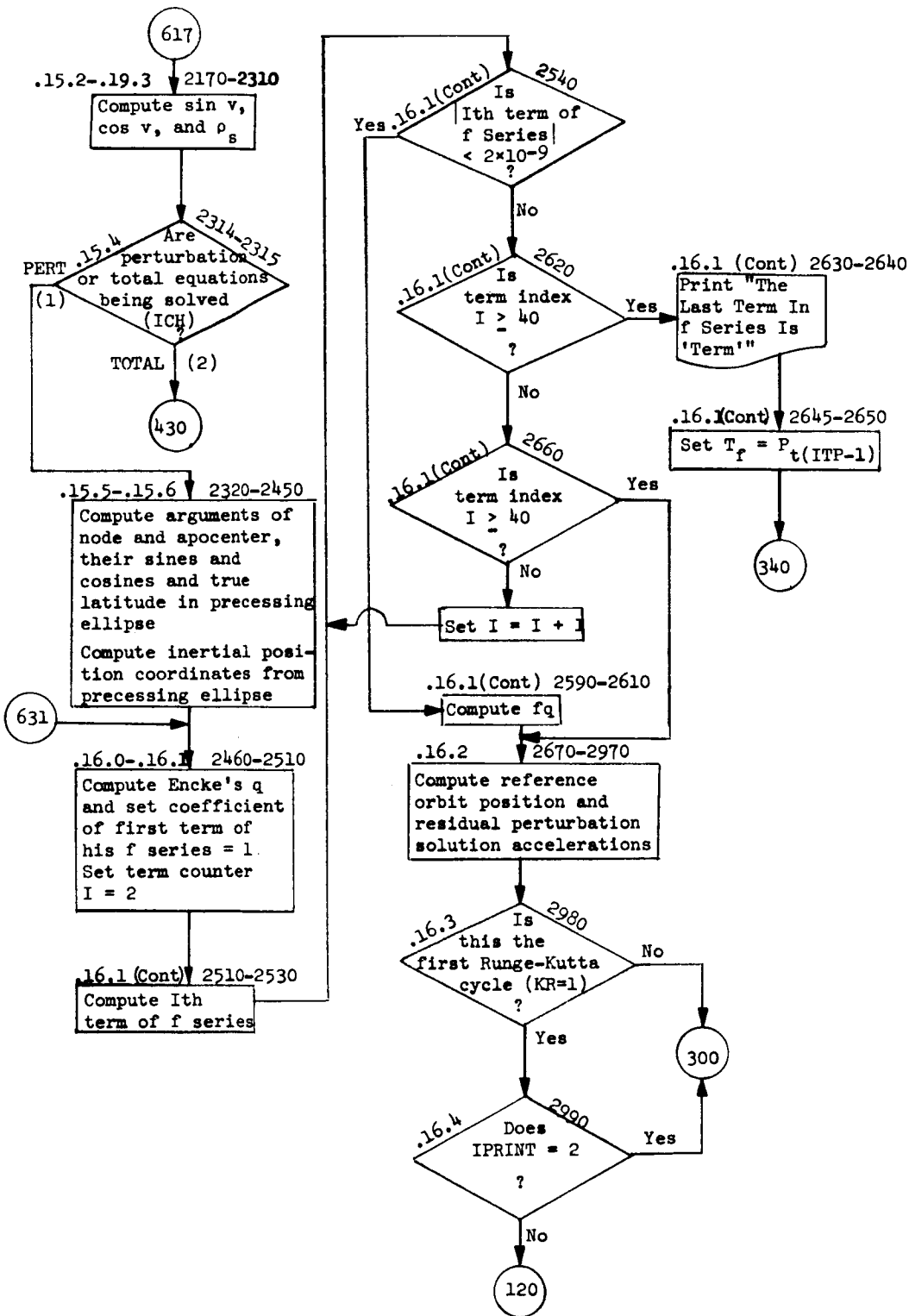
5.1 EXPLANATION

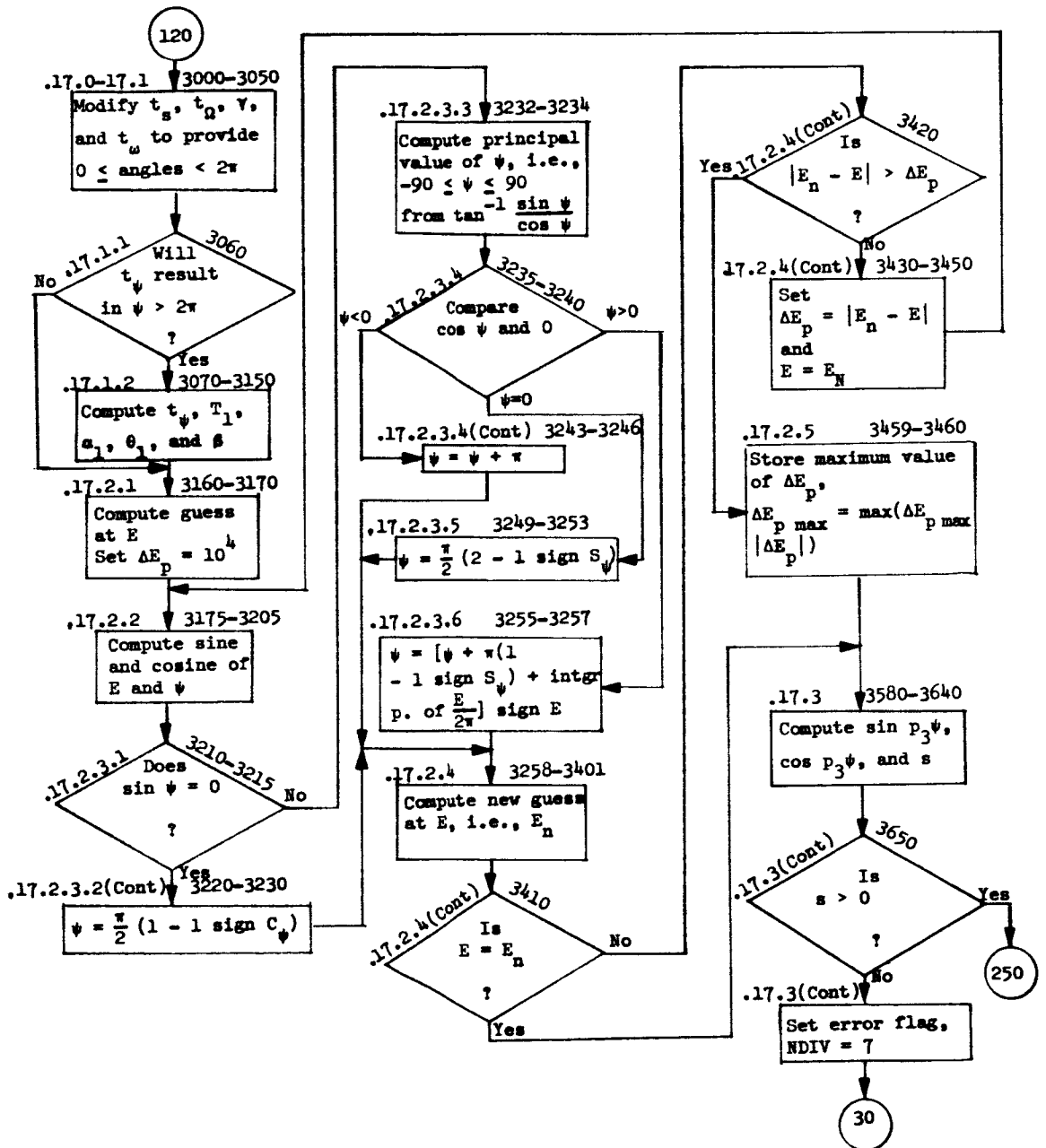
The flow charts of this section describe the logic of the main program completely, i.e., all conditional transfers are shown. The numbers to the top left at each box are the subsection numbers of Section 4 where the operations mentioned in the box are detailed. The numbers on the top right are the card numbers of the program listing, Section 6.

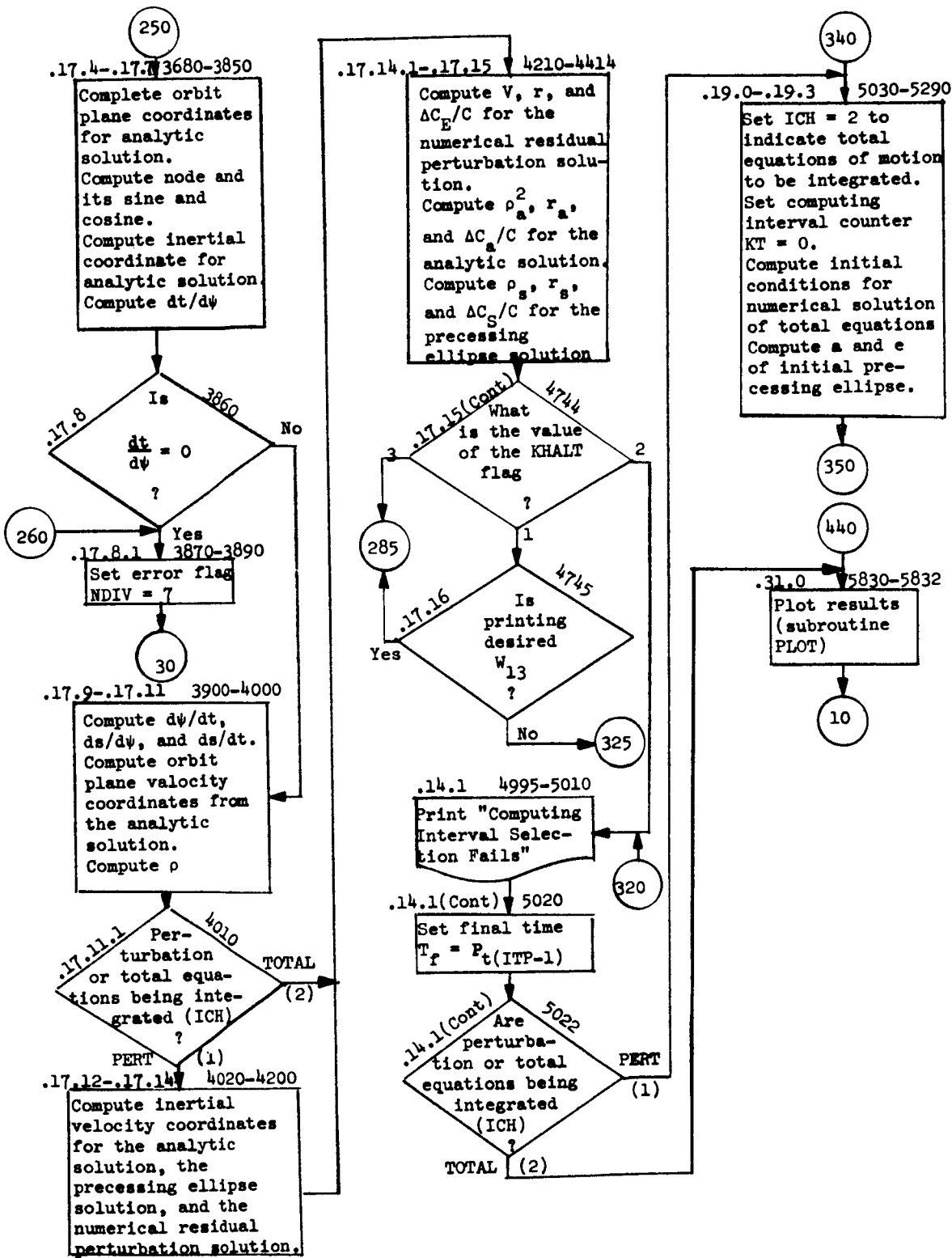


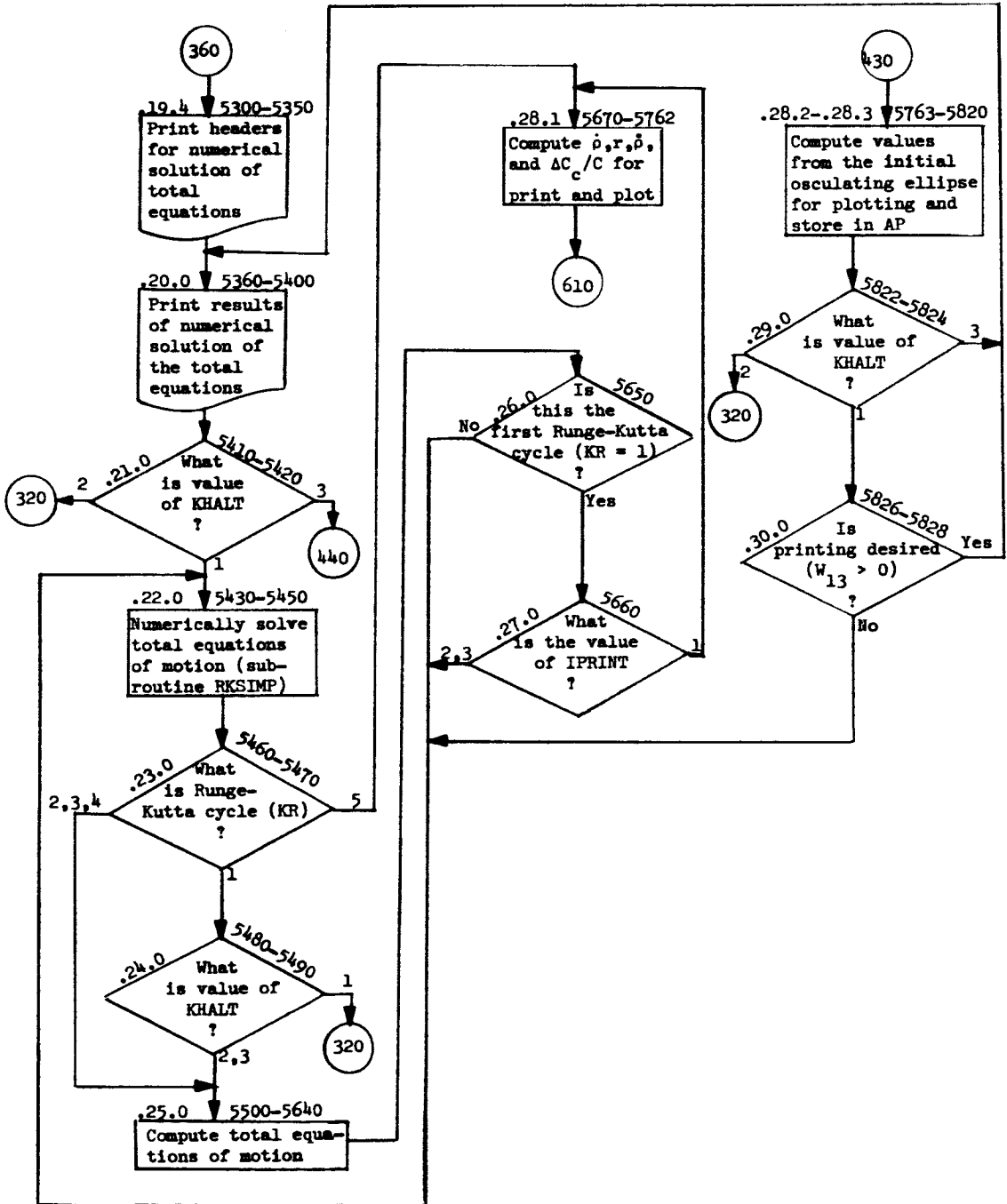












## Section 6

### MAIN PROGRAM LISTING

This section contains the listing of the main program. The subsection numbers on the comment cards refer to the subsections of Section 4, Equations in Order of Solution, wherein the code is explained.

M144 JFRY 09/01/65  
EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

```

C NUMERICAL RESIDUAL PERTURBATION SOLUTIONS APPLIED TO THE PROBLEM 0070H144
C OF A CLOSE SATELLITE OF THE SMALLER BODY IN THE RESTRICTED THREE- 0080H144
C BODY PROBLEM 0085H144
C DIMENSION M(30), WR(30), PS(25), FPS(25) 0090H144
C TIMEASION EGS(25), TERM(40), TNUM(40), IDEN(40), 0100H144
C I FRAI(40), FAC(40), PTIME(700), AP( 9100) 0110H144
C COMPN M, IPRINT, .5, TPSI, TNOD, TPER, GAMMA, EPI, 0120H144
C I JTP, PTIME, AP 0130H144
C1-0 PRINT EXPLANATION OF ERROR CODE 0140H144
      WRITE (6, 901) 0150H144
C1-1 901 FORMAT(1H1.3CH ERROR CODE REASON FOR HALT//4X2H 1.8X20H SEMI-MA0160H144
      1 JOR AXIS = 0/4X2H 2.8X33H ECCENTRICITY EQUALS OR EXCEEDS 1/ 0170H144
      5 4X2H 7.8X33H S (THE APPROX RADIUS VECTOR) = 0/4X2H 8.8X31H0210H144
      6 DERIVATIVE OF TIME WRT PSI = 0) 0220H144
C1-1 SFT INPUT ARRAYS = 0 0225H144
      DO 7 I = 1, 30 0230H144
      7 W(I) = 0. 0240H144
      KKKK = 1 0245H144
C 2-0 PROGRAM INPUT 0250H144
      10 CALL INPUT1 (W(1), W(30), WR(1)) 0260H144
      KFAIL = 0 0280H144
C3-0 MOVE INPUT TO WORKING STORAGE AND CHECK INPUT 0285H144
C3-1 MOVE INPUT TO WORKING STORAGE AND FORM NECESSARY CONSTANTS 0290H144
      A1 = ABS (W(1)) 0300H144
      A2 = A1 ** 2 0310H144
      A3 = A1 0320H144
      FC1 = ABS (W(2)) 0330H144
      ES = FC1 0340H144
      DFGJ = W(3) 0350H144
      YF = W (7) 0355H144
      FP4 = ABS (W(4)) 0360H144
C3-2 CHCK ON CONSTANT DIVISORS, HALT RUN IF ANY ARE ZERO 0370H144
C3-2-1 TEST SEMI MAJOR AXIS 0375H144
      IF (A1) 20, 20, 50 0380H144
C3-2-1-1 SET ERROR CODE = 1 0385H144
      20 NDIV = 1 0390H144
C3-2-1-2 PRINT ERRRC CODE AND GO TO STEP 2-0 0395H144
      30 WRITE (6, 40) NDIV 0400H144
      40 FORMAT(1H 18H THE ERROR CODE IS 15) 0410H144
      GO TO 10 0420H144
C3-2-2 COMPUTE ANC TEST E SUB 2 0425H144
      50 FC7 = 1. - EC1 ** 2 0430H144
      IF (FC2) 60, 60, 70 0440H144
C3-2-2-1 SET ERROR CODE TO 2 AND GO TO STEP 3-2-1-2 0445H144
      60 NDIV = 2 0450H144
      GO TO 30 0460H144
C4-0 COMPUTE CONSTANTS 0465H144
      70 AEC = A1 * EC2 0470H144
      AESO = SORT (AEC) 0480H144
      FC2 = SORT (FC2) 0490H144
      A3 = A1 * A1 * A1 0500H144
      EP2 = SORT (EP4) 0510H144
      TNUM(1) = 3. 0520H144
      IDEN(1) = 1. 0530H144
      DO AC 1 = 2, 40 0540H144

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I-144 JERRY EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

```

C24 = SQRT (1. - E5 ** 2)
P55 = 6.2831853 * A32
PNDD = 6.2831853 / DNOD
PPFR = 6.2831853 / DPER
DEM = 0.
DEPM = 0.
ALPHA1 = 0.
T-ETA1 = 0.
PH11 = 0
CALPHA = 1.
SALPHA = 0.
ST-ETA = 0.
CI-ETA = 1.
C5-0 PRINT CONSTANTS
WRITE (6, 1) (M(1), I = 1, 15), OMEG1, GNUL, ELMB, PI,
1 P1, C1, C2, C3, C4, C5, C6, C7, C8, C9, C10, C11, EC2
2 . DPER, DNOD, DPNB, DPNSO

1 FORMAT(1H0.5E18.8/5E18.8/5E18.8/5E18.8/6E18.8/6E18.8/
1 74F20.8)
C6-0 PRINT HEADER FOR NUMERICAL RESIDUAL PERTURBATION SOLUTION LISTING
C
117 WRITE (6, 502) A1, DEG I, ECL, EP4
902 FORMAT(1H1.49H THE INITIAL CONDITIONS ARE SEMI-MAJOR AXIS E17.1160H144
18.4417H INCLINATION+DEG E17.8/31X14H ECCENTRICITY E20.84X12H MASS1170H144
2 RATIO E22.8)
WRITE (6, 903)
903 FORMAT(1H0/61H TIME DELTA X DELTA X DOT DELTA X DBL DELTA X DBL
1 T RHO/65H DELTA Y DELTA Y DOT DELTA Y DBL DELTA Y DBL
2 RHO DOT/53H C APPROX DELTA Z DELTA Z DOT DELTA Z DBL
3 DOT/51H X X DOT X APPROX X DOT APPROX/51H
4 Y DOT Y DOT APPROX Y DOT APPROX/51H C EXACT Z
5 Z DOT Z APPROX Z DOT APPROX
C7-0 SET PERTURBATION-TOTAL FLAG
ICM = 1
C8-0 COMPUTE INITIAL VALUES OF POSITION, VELOCITY, AND ACCELERATION
EC1 = 1. * EC1 EC3 ** 1.5 / SORT (1. - EC1)
DTODPS = A32 * EP2 * A92
DTIDPS = DTODPS * DTIDPS
C8-1 IF (DTODPS) 119, 260, 119
C8-2 COMPUTE THE DERIVATIVE OF PHI WITH RESPECT TO T AND SET INITIAL
C VALUES OF T AND PSI
119 DPSTDT = 1. / DTODPS
DPHDT = DPSTDT / (1.+EP2 * OMEG1) - EP2 * GNUL
T = C.
PSI = 0.
C 8-3 POSITION COORDINATES
X2P = A1 * EC3
Y2P = 0.
Z2P = 0.
XP = X2P
YP = 0.
ZP = 0.
X = XP

```

```

0964H144 *87
0970H144 *88
0980H144 *89
0990H144 *90
0993H144 *91
0996H144 *92
1000H144 *93
1010H144 *94
1020H144 *95
1021H144 *96
1022H144 *97
1023H144 *98
1025H144 *99
1030H144
1040H144
1050H144
1060H144
1070H144
1070H144
1145H144
1146H144
1150H144
1150H144
1180H144
1190H144
1260H144
1270H144
1280H144
1290H144
1300H144
1310H144
1315H144
1320H144
1325H144
1326H144
1330H144
1340H144
1350H144
1360H144
1370H144
1380H144
1390H144
1400H144
1410H144
1420H144
1430H144
1440H144

```

```

*100 *101 *102 *103 *104 *105
*106 *107 *108
*109 *110
*111
*112
*113
*114
*115
*116
*117
*118
*119
*120
*121
*122
*123
*124
*125
*126
*127

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M144 JERRY EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

```

Y = C.
Z = C.
DX = 0.
DY = 0.
DZ = 0.
RND = X
R = 1. + FP2 * X
RCU = (1. + EP2 * X) ** 3
XS = XP
YS = 0.
ZS = 0.
XSD = 0.
VELOCITY COMPONENTS
XD2P = 0.
YD2P = XP * DPHIDT
ZD2P = 0.
XDP = 0
YDP = YD2P * CI + X2P * C6
ZDP = YD2P * SI
XD = 0.
YD = YDP
ZD = ZDP
DXD = 0.
GDOT = -(DTODPS * FP2 * OMEG1 + DTIDPS * (1. + EP2 * OMEG1))
1 / (DTODPS * (1. + EP2 * OMEG1) + DTODPS) - EP2 * GNUM1-DPER
DVD = X * GDOT * CI
D7D = X * GDOT * SI
R-ND = SORT (YD ** 2 + ZD ** 2)
XSD = 0.
YSD = X * (DNDD + (1. / DTODPS + DPER) * CI)
ZSD = X * (1. / DTODPS + DPER) * SI
C R.5 ACCELERATION COMPONENTS
FP2X = FP2 * X
DX2D = CB * EP2X * (2. + EP2X * (3. + EP2X)) / RCU
1 + 2. * (DPER + DNDD * CI) * SORT((1. - ES) / AS + (1.
2 + FS)) + XS * (2. + DPN * CI + DPNsq)
DY2D = 0.
DZ2D = 0.
C8.6 INITIAL VALUE OF JACOBI CONSTANT
ACORI = X ** 2 + (DPHIDT ** 2 + C7 - 2. * (C6 + EP1 -
1 / DPHIDT * (C6 - EP1) * CI)) - 2. * (1. / X + C8 * X)
RFAL JACOBI
JACOBI = ACORI - 2. * C8 / EP2 / (1. + EP2 * X)
XNAU = X
ANAU = 1. + EP2 * X
YACORI = 0.
YACAPP = 0.
WRITE (6, 391) JACOBI
C9.0 SET RUN START FLAGS
391 FORMAT(1H0.32H THE JACOBI INTEGRAL CONSTANT IS E17.8)
390 IP = 1
IPRINT = 2
KHALT = 1
KR = 1
C9.1 TFST PERTURBATION-TOTAL FLAG
GO TO (285, 360), ICH

```

1450H144	.128
1460H144	.129
1470H144	.130
1480H144	.131
1490H144	.132
1500H144	.133
1510H144	.134
1520H144	.135
1530H144	.136
1540H144	.137
1550H144	.138
1560H144	.139
1570H144	.140
1580H144	.141
1590H144	.142
1600H144	.143
1610H144	.144
1620H144	.145
1630H144	.146
1640H144	.147
1650H144	.148
1660H144	.149
1670H144	.150
1680H144	.151
1690H144	.152
1700H144	.153
1710H144	.154
1720H144	.155
1730H144	.156
1740H144	.157
1750H144	.158
1760H144	.159
1770H144	.160
1780H144	.161
1790H144	.162
1800H144	.163
1810H144	.164
1820H144	.165
1830H144	.166
1840H144	.167
1850H144	.168
1860H144	.169
1870H144	.170
1880H144	.171
1890H144	.172
1900H144	.173
1910H144	.174
1920H144	.175
1930H144	.176
1940H144	.177
1950H144	.178
1960H144	.179
1970H144	.180
1980H144	.181
1990H144	.182

C15.0	COMPUTE TWO-BODY EQUATIONS	2000H144	
C15.1	SOLVE KEPLER'S EQUATION	2010H144	.174
C15.1.1	FORM CONSTANTS	2015H144	
610	TAS32 = TS / AS32	2020H144	
C15.1.2	COMPUTE A FIRST ESTIMATE FOR ECCENTRIC ANOMALY	2025H144	.175
ESA = TAS32 - ES * SIN(TAS32)*11. + ES*CCOS (TAS32)		2030H144	.176
DI TF = 10000.		2040H144	
C15.1.3	CCMPUTF NEW GUESS AT E	2050H144	.177
616	SFSA = SIN (FSA)	2057H144	.178
CFSA = COS (ESA)		2058H144	.179
DEN = 1. + ES * CESA		2060H144	.180
FSAN = (TAS32 + FS * (ESA * CESA - SESAL)) / DEN		2070H144	.181
IF (FSA - ESAN) 612.617.612		2075H144	.182
612	IF ( ARS(ESA - ESAN) - DLTE) 613.220.220	2080H144	.183
613	DLTE = ARS (ESA - ESAN)	2090H144	.184
ESA = ESAN		2100H144	.185
GO TO 616		2110H144	.186
720	DFM = AMAX1 ( DEM, ABS (DLTE))	2120H144	.187
C15.2	COMPUTE SINE AND COSINE OF TRUE ANOMOLY	2170H144	.188
617	SIV = C24 * SESA / DEN	2190H144	
COV = (CESA * ES) / DEN		2200H144	.189
C15.3	COMPUTE RHO SUB S	2305H144	
623	RHOS = AS * DEN	2310H144	.190
C15.4	TEST PERTURBATION-TOTAL FLAG	2314H144	
GO TO 1431. 4301. ICM		2315H144	.191
C15.5	COMPUTE ARGUMENTS OF NODE AND PERIGEE AND THEIR SINES AND COSINES	2320H144	.192
431	SNODE = DNOD * TNOD	2330H144	.193
SPFR = DPER * TPER		2340H144	.194
CSNO = COS (SNODE)		2350H144	.195
SSNO = SIN (SNODE)		2360H144	.196
COPER = COS (SPFR)		2363H144	.197
SIPER = SIN (SPFR)		2373H144	.198
STRL = SIV * COPER + COV * SIPER		2380H144	
CTRL = COV * COPER - SIV * SIPER		2390H144	
C15.6	COMPUTE THE PRECESSING TWO BODY POSITION COORDINATE IN THE	2420H144	.199
C	INFRTIAL COORDINATE SYSTEM	2421H144	.200
XS = RHOS * (CTRL * CSNO - STRL * SSNO * CIS)		2430H144	.201
YS = RHOS * (CTRL * SSNO + STRL * CSNO * CIS)		2440H144	.202
ZS = RHOS * STRL * SIS		2450H144	.203
SFPT = SIN (GAMMA)		2453H144	
CFPT = COS (GAMMA)		2456H144	
C16.0	ACCELERATION COMPONENTS FOR NUMERICAL RESIDUAL PERTURBATION	2460H144	.204
C	FOULATIONS OF MOTION	2465H144	
C16.1	COMPUTATION OF THE SMALL DIFFERENCES BETWEEN THE TWO BODY TERMS	2470H144	
631	O = ( XS + DX / 2.) * DX + (YS + DY / 2.) * DY +	2480H144	.205
1	(ZS + DZ / 2.) * DZ ) / RHOS ** 2	2490H144	.206
TFRM(1) = 1.		2500H144	.207
OO 625 I = 2. 40		2510H144	.208
FRA (I) = O * TNUM(I) / TDEN(I)		2520H144	.209
TFRM(I) = TERM(I - 1) * FRA(I)		2530H144	.210
IF (ARS(TFRM(I))-2.E-9)624.624. 626		2540H144	.211
624	FAC(I) = 1. - FRA(I)	2550H144	.212
L = I - 1		2560H144	.213
OO 625 J = 2. L		2570H144	.214
K=I-J		2580H144	.215
625	FAC(J)=1.-FRA(K+1)*FAC(J-1)	2590H144	.216

EXTERNAL FORMULA NUMBER	SOURCE STATEMENT	INTERNAL FORMULA NUMBER(S)
H144	GO TO 340	2600H144
	FO = 3.0 * 0 + FAC(L)	217
	GO TO 630	2610H144
	626 IF (I - 40) 629, 627, 627	218
	627 WRITE (6, 628) TERM(I)	219
	628 FORMAT(1H0,29H THE LAST TERM IN F SERIES IS E17.6)	220, 221, 222
	IF = PTIME (TTP - I)	223
	TTP = I	224
	GO TO 340	225
629	CONTINUE	226, 227
C16.2	COMPUTE NUMERICAL RESIDUAL PERTURBATION ACCELERATION COMPONENTS	228
630	X = XS + DX	229
	Y = YS + DY	230
	Z = ZS + DZ	231
	RHOSC = X*Y + Y*Y + Z*Z	232
	RSO = 1. + EP2 * (2. + IX * CEPT + Y * SEPT) + EP2 * RHOSO	233
	R = SQRT (RSO)	234
	RCU = RSO * R	235
	FR1 = CA * EP2 / RCU	236
	FR2 = CB * EP2 * (2. + IX * CEPT + Y * SEPT) + EP2 * RHOSO	237
	1 RHOSO * (1. + R * (1. + R)) / (1. + R)	238
	RHOSCU = RHOS * 3	239
	R1 = C15 + C16 * RHOS	240
	R2 = C17 + C18 * RHOS	241
	DX20 = (FO * X - DX) / RHOSCU - EPRI * X + EPR2 * CEPT	242
	1 * S1RL * SSNO * B1 + CTRL * CSNO * B2 - S1PER * SSNO * C22 - COPER	243
	2 * CSNO * C23	244
	DY20 = (FC * Y - DY) / RHOSCU - EPRI * Y + EPR2 * SEPT	245
	1 - S1RL * CSNO * B1 + CTRL * SSNO * B2 + S1PER * CSNO * C22 - COPER	246
	2 * SSNO * C23	247
	DZ20 = (FO * Z - DZ) / RHOSCU - EPRI * Z	248
	1 + S1RL * (C21 + RHOS * C19) - S1PER * C20	249
C16.3	TEST THE RUNGE-KUTTA FLAG, KR	250
	GO TO (121, 300, 300, 300, 300), KR	251
C16.4	TEST THE PRINT FLAG, IPRINT	252
	GO TO (120, 300, 120), IPRINT	253
C17.0	COMPUTATIONS FOR PRINT ONLY	254
C17.1	REDUCE ALL ANGLES THAT INCREASE WITH TIME BY INTEGER MULTIPLES OF 2 PI	255
	120 IS = AMOD (TS, PSS)	256
	TNOD = AMOD (TNOD, PNOO)	257
	TPFR = AMOD (TPFR, PPER)	258
	GAMMA = AMOD (GAMMA, 6.2831853)	259
C17.1.1	TEST VALUE OF T SUB PSI	260
	IF (TPSI - T1)1007, 632, 632	261
C17.1.2	SFT T SUB PSI	262
632	TPSI = T PSI - T1	263
	PHI1=AMOD(PHI1+6.2831853 / (1.+ EP2 * OMEG1)) - EP2 * GNU1 * T1,	264
	1 6.2831853)	265
	ALPHA1 = AMOD (AMOD (ALPHA1 + 6.2831853 * P1, 6.2831853)	266
	1 + 6.2831853 * P1, 6.2831853)	267
	CALPHA = CCS (ALPHA1)	268
	SALPHA = SIN (ALPHA1)	269
	ALPHA2 = AMOD (AMOD (ALPHA1 + 6.2831853 * P1, 6.2831853)	270
	1 + 6.2831853 * P1, 6.2831853)	271
	TMETAI = AMOD (TMETAI + 6.2831853 * C11, 6.2831853)	272
	STMETA = SIN (TMETAI)	273



09/01/65

JFRY - INTERNAL FORMULA NUMBERS

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1144 JFRY - SOURCE STATEMENT - INTERNAL FORMULA NUMBERS
FATERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBERS
C17-2.5 STEME MAXIMUM FINAL VALUE OF DELTA E PERTURBED .296
6 DEPM = AMAX1 ( DEPM, ABS (DLTEP)) 3459H144
C17-3 COMPUTE AUXILIARY FUNCTIONS OF PSI AND T .297
235 CP3 = COS (P3 * PSI) * COS (BETA) - SIN (P3 * PSI) * 3580H144
1 SIN (BETA) 3590H144
SP3 = SIN (P3 * PSI) * COS (BETA) + COS (P3 * PSI) * 3600H144
1 SIN (BETA) 3610H144
S = (1. - ECI * CPS) / AEC * C2 * (CPS - CPI) 3620H144
1 * A2 * (1. - 5. * CPS) / 3. * CPIPS) * EP2 3630H144
IF (S) 240. 240. 250 3640H144
240 NDIV = 7 3650H144
GO TO 30 3660H144
C17-4 COMPUTE DOUBLE PRIME COORDINATES 3670H144
290 P41 = PSI / (1. + EP2 * OMEG1) - EP2 * GNUL * TPST 3680H144
22P = (COS (PH1) * COS (PH1)) - SIN (PH1) * SIN (PH1) / S 3690H144
22P = (SIN (PH1) * COS (PH1)) + COS (PH1) * SIN (PH1) / S 3700H144
22P = C3 * (SP3 - P3 * SP3) 3710H144
C17-5 COMPUTE NODE 3720H144
ANOCF = C6 * TNOD 3730H144
CNO = COS (ARNODE) 3740H144
SNO = SIN (ARNODE) 3750H144
C17-6 COMPUTE APPROXIMATE COORDINATES IN PLANE OF PLANETS 3760H144
C SINGIF PRIME SYSTEM 6.1.6.24.3
XP = X2P * CNO - Y2P * SNO * CI + Z2P * SNO * SI 3770H144
YP = X2P * SNO + Y2P * CNO * CI - Z2P * CNO * SI 3780H144
ZP = 1. - Y2P * SNO * CI + Z2P * CNO * SI 3790H144
C17-7 COMPUTE TIME DERIVATIVES OF PSI, PHI AND S 3800H144
DTPSI = C9 / (1. - ECI * CPS) * 2) - CI * 3810H144
1 (CPS - CPI) 3820H144
2 * A52 * (10. / 3. * CPS - 2.75 * CPIPS) * EP2 3830H144
C17-8 TEST THE VALUE OF THE DERIVATIVE OF TIME WITH RESPECT TO PSI 3840H144
IF (DTPSI) 270. 260. 270 3850H144
260 NDIV = 8 3860H144
GO TO 30 3870H144
C17-9 COMPUTE TIME DERIVATIVES OF PHI AND S 3880H144
270 DPSIDT = 1. / DTPSI 3890H144
DPHICT = DPSIDT / (1. + EP2 * OMEG1) - EP2 * GNUL 3900H144
GSDPSI = FC1 * SPS / AEC * C2 * (SPL - SPS) 3910H144
1 * A2 * (5. / 3. * SPS - 2. * SPLPS) * EP2 3920H144
DSDI = DSDPSI * DPSIDT 3930H144
C17-10 COMPUTE APPROXIMATE VELOCITY COMPONENTS IN ORBIT PLANE 3940H144
X22P = - Y2P * DPHICT - X2P * DSDT / S 3950H144
Y22P = X2P * DPHICT - Y2P * DSDT / S 3960H144
Z22P = C3 * (CP3 - CPS) * DPSIDT 3970H144
C17-11 COMPUTE RHO 3980H144
RHO = SORT (RHUSQ) 3990H144
C17-11.1 IF ICH = 1 GO TO 17.12. IF ICH = 2 GO TO 17.14.1 4000H144
GO TO (800. 801). ICH 4005H144
C17-12 COMPUTE APPROXIMATE VELOCITY COMPONENTS IN INERTIAL SYSTEM 4010H144
800 ADP = C6 * (1 - X2P * SNO - Y2P * CNO * CI + Z2P * CNO * SI) 4015H144
1 * X22P * CNO - Y22P * SNO * CI + Z22P * SNO * SI 4020H144
YDP = C6 * (X2P * CNO - Y2P * SNO * CI + Z2P * CNO * SI) 4025H144
1 * X22P * SNO + Y22P * CNO * CI - Z22P * CNO * SI 4030H144
ZDP = Y22P * SNO * CI 4035H144
C17-13 COMPUTE PRECESSING KEPLERIAN VELOCITY COMPONENTS IN INERTIAL 4040H144

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H144 JERRY EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

```

C
  SYSTEM
  XSD = (STRL * CSNO + CTRL * SSNO * CIS
  1 - FS * (SEPER * CSNO + COPER * SSNO * CIS) / AESO
  2 - RHOS * (DOND * (CTRL * SSNO + STRL * CSNO * CIS)
  3 * DPFR * (STRL * CSNO + CTRL * SSNO * CIS))
  YSD = (STRL * SSNO - CTRL * CSNO * CIS
  1 - FS * (SEPER * SSNO + COPER * CSNO * CIS) / AESO
  2 + RHOS * (DOND * (CTRL * CSNO - STRL * SSNO * CIS)
  3 - DPFR * (STRL * SSNO - CTRL * CSNO * CIS))
  ZSD = ((CTRL * SSNO + COPER) / AESO
  1 + R-OS * DPFR * CTRL) * SIS
C17.14 COMPUTE TOTAL INERTIAL VELOCITY COMPONENTS AND VECTOR MAGNITUDES
  AD = XSD + D4D
  YD = YSD + D7D
  ZD = ZSD + D7D
C17.14.1 COMPUTE INERTIAL VELOCITY
  801 VSO = AD * XD + YD * YD + ZD * ZD
  RHOD = SQR (VSO)
  R = SQR (RSO)
  YACOR1 = VSO + 2 * (EPI * (XD * Y - YD * X) - 1. / RHO
  1 - C8 * (X * CEPT + Y * SEPT)) - ACOR1 + 2 * C8 * (2.
  2 * (Y * CEPT + X * SEPT - XNAUT) + EP2 * (RHOSO - XNAUT
  3 ** 2)) / R / RNAUT / (RNAUT + R)
  YACOR2 = YACOR1 / JACOR1
  RHODASO = XP ** 2 + YP ** 2 + ZP ** 2
  RA = SQR (1. + EP2 * (2. * (XP * CEPT + YP * SEPT) +
  1 EP2 * RHODASO))
  YACAPP = XDP ** 2 + YDP ** 2 + ZDP ** 2 + 2 * (EPI *
  1 (XDP * YP - YDP * XP) - 1. / SORT (RHODASO) - C8 * (XP
  2 * CEPT + YP * SEPT)) - ACOR2 + 2 * C8 * (2. * (XP *
  3 CEPT + YP * SEPT - XNAUT) + EP2 * (RHODASO - XNAUT ** 2
  4)) / RA / RNAUT / (RNAUT + RA)
  YACAPP = YACAPP / JACOR1
  RHORSO = XS * XS + YS * YS + ZS * ZS
  RHORS = SQR (RHORSO)
  RS = SQR (1. + EP2 * (2. * (XS * CEPT + YS * SEPT) +
  1 + EP2 * RHORSO))
  YACAS = (XSC * XSD + YSD * YSD + ZSD * ZSD + 2 * (EPI * (XSD
  1 * YS - YSD * XS) - 1. / RHOS - C8 * (XS * CEPT + YS * SEPT)
  2 - ACOR1 + 2 * C8 * (2. * (XS * CEPT + YS * SEPT - XNAUT)
  3 + EP2 * (RHOSU - XNAUT * XNAUT)) / RS / RNAUT / (RNAUT + RS))
  4 / JACOR1
  GO IC (810, 320, 285), KHALI
C17.16 TEST IF PRINTING IS DESIRED
  810 IF (N(13)) 325, 325, 285
C18.0 OUTPUT OF NUMERICAL RESIDUAL PERTURBATION SOLUTION
  785 MATIF (6, 904) T, DX, DXD, DX2D, RHO, DY, DYD, DY2D,
  1 RHCG, YACAPP, DZ, DZO, DZ2D
  7, X, XD, XP, YD, YP, YD, YP, YOP, YACOR1, Z, ZO, ZP, ZDP
  904 FORMAT(IHO, /5F17-8/E3+.8, 3E17-8 /4E17-8 /E3+.8, 3E17-8 /
  1E3+.8, 3E17-8 /5F17-8)
  GO IC (300, 325), IP
C18.0 SAVE PLOT VALUES AND TEST IF PLOTTING ARRAY IS FULL
C18.1 SAVE VALUES FOR PLOT ROUTINE AND FOR LATER COMPARISON WITH VALUES
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  5791H144
  5792H144
  5793H144
  5794H144
  5795H144
  5796H144
  57
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INTERNAL FORMULA NUMBER(S)

EXTERNAL FORMULA NUMBER	SOURCE STATEMENT	INTERNAL FORMULA NUMBER(S)
*144	JFRMY	
325	PTIME (ITP) = T	4859H144
	IT = 13 * ITP - 12	4860H144
	AP (IT) = ARS (YACAPP)	4861H144
	AP (IT + 1) = ABS (YACAS)	4862H144
	AP (IT + 2) = X	4863H144
	AP (IT + 3) = Y	4864H144
	AP (IT + 4) = ABS ( YACOBI)	4865H144
	AP (IT + 5) = SORT ((XP - X) * 2 + (YP - Y) * 2	4866H144
	1 + (ZP - Z) * 2) / A1	4867H144
	AP (IT + 6) = SORT ((XS - X) * 2 + (YS - Y) * 2	4868H144
	1 + (ZS - Z) * 2) / A1	4869H144
	AP (IT + 7) = Z	4870H144
	AP (IT + 8) = XU	4871H144
	AP (IT + 9) = A12 * SORT((XDP - XD) * 2 + (YDP -	4872H144
	1 YD) * 2 + (ZDP - ZD) * 2)	4873H144
	AP (IT + 10) = A12 * SORT((XSD - XD) * 2 + (YSD -	4874H144
	1 YD) * 2 + (ZSD - ZD) * 2)	4875H144
	AP (IT + 11) = YD	4876H144
	AP (IT + 12) = ZD	4877H144
	AP (IT + 13) = 7D	4878H144
	C18.2 TEST THE VALUE OF ITP	4879H144
	IF (ITP - 700) 326, 327, 327	4880H144
	326 ITP = ITP + 1	4881H144
	PTIME (ITP) = TF	4882H144
	GO TC 290	4883H144
	327 TF = PTIME (ITP)	4884H144
	GO TC 340	4885H144
	C11.0 TEST HALT FLAG, KHALT	4910H144
	GO TC (300, 320, 340), KHALT	4920H144
	C12.0 CALL NUMERICAL SOLUTION SUBROUTINE FOR RESIDUAL PERTURBATION	4930H144
	C EQUATIONS	4935H144
	300 CALL RKSTMP (KR, IP, KC, T, DT, ICH, DX, DY, DZ, DXD,	4940H144
	1 DYD, DZD, DXZD, DYZD, DZZD, KHALT, TPRINT, TF)	4950H144
	C 13.0 TEST RUNGF-KALTA FLAG	4960H144
	GO TO (310, 610, 631, 610, 120), KR	4970H144
	C 14.0 TEST HALT FLAG FOR INTEGRATION ACCURACY ERROR	4980H144
	310 GO TC (631, 320, 631), KHALT	4990H144
	C14.1 PRINT COMPUTING INTERVAL SELECTION FAILED	4995H144
	320 WRITE (6, 330)	5000H144
	330 FORMAT(1H0.35H COMPUTING INTERVAL SELECTION FAILED)	5010H144
	TF = PTIME (ITP - 1)	5020H144
	C19.0 INITIATE NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTION	5022H144
	GO TC (340, 44C), ICH	5030H144
	C19.1 SET PERTURBATION-TOTAL FLAG	5035H144
	340 ICH = 2	5040H144
	KT = 0	5050H144
	C19.2 COMPUTE INITIAL CONDITIONS FOR NUMERICAL SOLUTION OF TOTAL	5060H144
	C EQUATIONS OF MOTION	5065H144
	X = A1 * EC3	5070H144
	Y = C.	5080H144
	Z = C.	5090H144
	GAMMA = 0.	5100H144
	RHO = X	5110H144
	T = C.	5120H144
	DTDPSI = A12 * EC3 * 1.5 / SORT (1. - EC1)	5130H144
	1 * A52 * EP2 * 7.	5140H144
	DPSIOT = 1. / DTDPSI	5150H144
		5160H144
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M144 JFRYR          09/01/65
EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

DPHJDT = DPHJDT / ( 1. + EP2 * OMEG1 ) - EP2 * GNU1
X0 = 0.
Y0 = X * (DPHJDT * C1 + C6)
Z0 = X * DPHJDT * S1
RHOD = SORT (Y0 ** 2 + Z0 ** 2)
RCU = (1. + FP2 * X) ** 3
EP2X = EP2 * X
XZD = - 1. / X ** 2 + C8 * EP2X * (2. + EP2X * (3. +
1 FP2X)) / RCU
YZD = 0.
ZZD = 0.
YACOBI = 0.
C19.3 COMPUTE SEMI-MAJOR AXIS AND ECCENTRICITY OF INIT. OSC. ELLIPSE
AS = 1. / (7. / RHOD - RHOD ** 2)
FS = X / AS - 1.
AS32 = AS ** 1.5
PSS = 6.7831853 * AS32
AESQ = SORT (AS * (1. - ES ** 2))
GO TO 350
C19.4 PRINT HEADER FOR NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTION
C
360 WRITE (6, 370)
370 FORMAT(1H1.37H INTEGRATION OF COMWELL EONS OF MOTION//6X5H TIME,16X53Z0H144
12H X,17X6H X DOT,13X10H X DBL DOT,14X5H RHO/27X2H Y,17X6H Y DOT,135330H144
2X10H Y DBL DOT,12X8H RHO DOT/17H JACOBI CONSTANT,10X2H Z,17X6H Z
30T,13X10H Z DBL DOT)
C20.0 PRINT NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTION DATA
380 IT = 13 * ITP - 12
IIO = IT + 12
WRITE (6, 390) T, X, XD, XZD, RHO, Y, YD, YZD, RHOD, YACOBI, Z, ZD,
1 ZD0, (AP1), I = 11,170)
390 FORMAT(1H0./E16.8,E19.8,3E21.8/E35.8,3E21.8/E16.8,
1 F19.8, 7E21.8 /5F20.8/4E20.8/4E20.8)
C21.0 TEST THE HALT FLAG, KHALT
400 GO TO (610, 320, 440), KHALT
C22.0 CALL NUMERICAL SOLUTION SUBROUTINE FOR TOTAL EQUATIONS OF MOTION
410 CALL RKSTMP (KR, IP, KC, T, DT, ICH, X, Y, Z, XD, YD,
1 ZD, XZD, YZD, 7ZD, KHALT, IPRINT, TF)
C23.0 TEST THE RUNGE-KUTTA FLAG, KR
420 GO TO (620, 275, 275, 275, 283), KR
C24.0 TEST THE HALT FLAG, KHALT
430 GO TO (275, 320, 275), KHALT
C25.0 COMPUTE TOTAL EQUATIONS OF MOTION
275 RHOSO = X * X + Y * Y + Z * Z
RHOCU = RHOSO ** 1.5
GAMMA = AMOD (GAMMA, 6.2831853)
CEPT = COS (GAMMA)
SFPT = SIN (GAMMA)
R50 = 1. + EP4 * RHOSO + 2. * EP2 * (X * CEPT + Y * SEPT)
R = SORT (R50)
RCU = R50 * R
FPRI = C8 * EP2 / RCU
EPR2 = C8 * EP2 * (2. * (X * CEPT + Y * SEPT) + EP2 *
1 RHOSO) * (1. + R * (1. + R)) / RCU / (1. + R)
XZD = - X / RHOCU - EPR1 * X + EPR2 * CEPT

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INTERNAL FORMULA NUMBER(S)	INTERNAL FORMULA NUMBER(S)
.394	5170H144
.395	5180H144
.396	5190H144
.397	5200H144
.398	5210H144
.399	5220H144
.400	5230H144
.401	5240H144
.402	5250H144
.403	5260H144
.404	5270H144
.405	5280H144
.406	5290H144
.407	5300H144
.408	5310H144
.409	5320H144
.410	5330H144
.411	5340H144
.412	5350H144
.413	5360H144
.414	5370H144
.415	5380H144
.416	5390H144
.417	5400H144
.418	5410H144
.419	5420H144
.420	5430H144
.421	5440H144
.422	5450H144
.423	5460H144
.424	5470H144
.425	5480H144
.426	5490H144
.427	5500H144
.428	5510H144
.429	5520H144
.430	5530H144
.431	5540H144
.432	5550H144
.433	5560H144
.434	5570H144
.435	5580H144
.436	5590H144
.437	5600H144
.438	5610H144
.439	5620H144

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INTERNAL FORMULA NUMBER(S)

M144 JFRYR - SOURCE STATEMENT -

Y2D = - Y / RHOCU - EPRI \* Y + EPR2 \* SEPT  
 Z2D = - Z / RHOCU - EPRI \* Z  
 C26.0 TEST THE RUNGE-KUTTA FLAG, KR  
 GO TO (411, 410, 410, 410, 410), KR  
 C27.0 TEST THE PRINT FLAG, IPRINT  
 411 GO TO (283, 410, 410), IPRINT  
 C28.0 COMPLETE DATA FOR PRINT AND FOR PLOT ROUTINE  
 C28.1 COMPUTE ERROR IN JACOBI INTEGRAL AND LIMIT T SUB S FOR NUMERICAL  
 SOLUTION OF TOTAL EQUATIONS  
 C 283 RHO = SORT (RHOSO)  
 R = SORT (RSO)  
 VSO = XD \*\* 2 + YD \*\* 2 + ZD \*\* 2  
 RHOD = SORT (VSO)  
 YACORI = VSO \* 2. \* (LEPI \* (XD \* Y - YD \* X) - 1. / RHO  
 1 - CR \* (X \* CEPT + Y \* SEPT)) - ACOR1 \* 2. \* C8 \* (2.  
 2 \* (X \* LEPT + Y \* SEPT - XNAUT)) + EP2 \* (RHOSO - XNAUT  
 3 \*\* 2) / R / RNAUT / (RNAUT + R)  
 YACORI = YACORI / JACOBI  
 TS = AMOD (TS, PSS)  
 GO TO 610  
 C28.2 COMPUTE THE FIXED ORBIT VALUES  
 430 XF = RHOS \* COV  
 YF = RHOS \* SIV \* CI  
 ZF = RHOS \* SIV \* SI  
 XFD = -SIV / AESO  
 YFD = (COV - ES) / AESO \* CI  
 ZFD = (COV - FS) / AESO \* SI  
 RHOFSO = XF \*\* 2 + YF \*\* 2 + ZF \*\* 2  
 RHOF = SORT (RHOFSO)  
 RF = SORT ((LEPT \* 2. \* (XF \* CEPT + YF \* SEPT)) +  
 1 \* FP) \* RHOFSO)  
 YACAF = XFD \*\* 2 + YFD \*\* 2 + ZFD \*\* 2 + 2. \* (LEPI \*  
 1 (XFC \* YF - YFD \* XF) - 1. / RHOF - C8 \* (XF \* CEPT +  
 2 YF \* SEPT)) - ACOR1 \* 2. \* C8 \* (2. \* (XF \* CEPT + YF  
 3 \* SEPT - XNAUT) + EP2 \* (RHOFSO - XNAUT \*\* 2)) / RF /  
 4 RNAUT / (RNAUT + RF)  
 IT = 13 \* ITP - 12  
 YACAF = YACAF / JACOBI  
 DRF = SORT ((XF - AP(11 + 2)) \*\* 2 + (YF - AP(11 +  
 1 3)) \*\* 2 + (ZF - AP(11 + 7)) \*\* 2) / AI  
 DRDF = AI \* 2 \* (ZF - AP(11 + 8)) \*\* 2 + (YFD  
 1 - AP(11 + 11)) \*\* 2 + (ZFD - AP(11 + 12)) \*\* 2)  
 DRDC = SORT ((X - AP(11 + 2)) \*\* 2 + (Y - AP(11 +  
 1 3)) \*\* 2 + (Z - AP(11 + 7)) \*\* 2) / AI  
 DRDC = AI \* 2 \* (Z - AP(11 + 8)) \*\* 2 + (YD -  
 1 AP(11 + 11)) \*\* 2 + (ZD - AP(11 + 12)) \*\* 2)  
 C28.3 SFT QUANTITIES JUST COMPUTED EQUAL TO VALUES IN PLOT ARRAY  
 AP(11 + 2) = ABS (YACORI)  
 AP(11 + 3) = ABS (YACAF)  
 AP(11 + 7) = ORC  
 AP(11 + 8) = DRF  
 AP(11 + 11) = DRDC  
 AP(11 + 12) = DRDF  
 C29.0 TEST VALUES OF HALT FLAG, KNALT  
 GO TO (815, 320, 380), KNALT  
 C30.0 TEST IF PRINTING IS DESIRED  
 5630H144 .436  
 5640H144 .437  
 5645H144 .438  
 5650H144 .439  
 5660H144 .440  
 5670H144 .441  
 5675H144 .442  
 5680H144 .443  
 5690H144 .444  
 5700H144 .445  
 5710H144 .446  
 5720H144 .447  
 5730H144 .448  
 5740H144 .449  
 5750H144 .450  
 5760H144 .451  
 5770H144 .452  
 5780H144 .453  
 5790H144 .454  
 5795H144 .455  
 5800H144 .456  
 5810H144 .457  
 5820H144 .458  
 5825H144 .459  
 5830H144 .460  
 5840H144 .461  
 5850H144 .462  
 5860H144 .463  
 5870H144 .464  
 5880H144 .465  
 5890H144 .466  
 5900H144 .467  
 5910H144 .468  
 5920H144 .469  
 5930H144 .470

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EXTERNAL FORMULA NUMBER - INTERNAL FORMULA NUMBER(S)

```

815 IF (M(13)) 410, 410, 380
C31.0 CALL THE PLOT SUBROUTINE
440 CALL PLOT (A32,TF,A1,EC1,DEGI,EP4,JACOBI,PSS)
WRITE (6,230) DEM,DEPM
230 FORMAT(1H01GMAXIMUM DELTA E= E17.8,27HMAXIMUM DELTA E PERTURBED
1 F17.8
C32.0 RUN IS COMPLETED, CALL THE NEXT CASE AT STEP 4.2
GO TO 10
END
5828H144 *471
5829H144 *472
5831H144 *473 *474 *475
5832H144
5833H144
5834H144
5835H144 *476
5836H144 *477
5838H144

```

## Section 7

### SUBROUTINE FOR NUMERICAL SOLUTION OF DIFFERENTIAL EQUATIONS, RKSIMP

#### 7.1 EQUATIONS IN ORDER OF SOLUTION

##### 7.1.1 Subroutine Input

##### 7.1.1.1 Calling Sequence

KR	Runge-Kutta flag; takes on five values from 1 to 5 indicating the cycle of Runge-Kutta. Initially set = 1 by the main program and thereafter under the control of the subroutine.
IP	Initial point flag; signals first point calculations by the subroutine. Initially set to 1 by the main program and advanced to 2 by the subroutine.
KC	Simpson's rule computation flag. The computations are bypassed if $KC = 1$ , computed if $KC = 2$ . Wholly under the control of the subroutine.
t	Time; initial value set by main program, thereafter under control of the subroutine.
$\Delta t$	Computing interval; initial value is input; if input is left blank, it is set equal to 0.005 by the subroutine.
ICH	A flag that signals which equations are being integrated.  ICH = 1 corresponds to numerical residual perturbation solution

ICH = 2 corresponds to numerical solution of total equations of motion

Wholly under control of the main program.

$x_h, y_h, z_h$   
 $\dot{x}_h, \dot{y}_h, \dot{z}_h$   
 $\ddot{x}_h, \ddot{y}_h, \ddot{z}_h$

Dummy variables of position, velocity, and acceleration, respectively.

KHALT Halt = 2, call next case.

$t_p$  Print time

$T_f$  Run stop time

#### 7.1.1.2 Common

W The input array

IPRINT Print flag; print if equal to 1.  
Initially set = 1 by main program for numerical residual perturbation solution  
Initially set = 2 by main program for numerical solution of total equations of motion  
Under control of subroutine during numerical solution of total equations of motion

$t_s$  Time modulo the period of the precessing ellipse

$t_\psi$  Time since last passage of the apocenter by the analytic solution

$t_{\Omega}$	Time modulo the period of the node
$t_{\omega}$	Time modulo the period of the apocenter
$\gamma$	$= \epsilon t$
$\epsilon$	$\mu^{1/4}$
ITP	Plot point index

#### 7.1.2 Test Runge-Kutta Flag, KR

If KR = 1, continue below (Step 7.1.3).

If KR = 2, go to Step 7.1.26.

If KR = 3, go to Step 7.1.27.

If KR = 4, go to Step 7.1.28.

If KR = 5, go to Step 7.1.30.

#### 7.1.3 Test Initial Flag, IP

If IP = 1, first point computations must be executed; continue below (Step 7.1.4).

If IP = 2, go to Step 7.1.9.

#### 7.1.4 First Point Computations

##### 7.1.4.1 Set Flags and Initial Conditions

Set IP = 2      Advance initial point flag.

LL = 0            Obsolete.

KT = -2           Set computing interval counter = -2.

ITP = 1           Set initial value of index for plotting arrays  $A_p$  and  $P_t$ .

KC = 1            Indicates next two Runge-Kutta cycles to be at the same  $\Delta_t$ .

$E_{\min} = |W_{12}|$     Set value for the minimum allowable error equal to input value.

FDT =  $|W_7|$         Modifier of allowable computing interval.

KF = 0            Set intermediate and total failure  
KFAIL = 0           counters to zero.

MFAIL =  $|W_{11}|$     Set the maximum number of failures allowed equal to an input value.

$t_{pT} = P_{t1}$        Set first print time for numerical solution of total equations of motion.

(Note: Unnecessary on numerical residual perturbation solution, but saves coding if included here.)



$\Delta\dot{X} = 0$	}	Sets velocity increments for use in computing interval tests initially equal to zero.
$\Delta\dot{Y} = 0$		
$\Delta\dot{Z} = 0$		
$\gamma = 0$	}	Sets initial values of angle and time variables to zero.
$t_{\psi} = 0$		
$\beta = 0$		
$t_{\Omega} = 0$		
$t_{\omega} = 0$		
$t_s = 0$		
$\Delta t_p = a^{3/2}/4$	}	Sets print interval to 1/4 radian of M
$t_p = \Delta t_p$		

7.1.4.2 Set the Maximum Allowable Error Equal to the Input Value or Equal to  $10^{-7}$  if the Input Value is Zero

$$E_{\text{all}} = |W_8| \text{ or } E_{\text{all}} = 10^{-7} \text{ if } W_8 = 0$$

7.1.4.3 Set the Computing Interval Multiplier (Used if Error is Less Than the Minimum Allowed) Equal to Input Value or Equal to 1.5 if Input Value  $\leq 1$

$$\Delta t_{\text{min}} = W_9 \text{ or } \Delta t_{\text{min}} = 1.5 \text{ if } W_9 \leq 1$$

7.1.4.4 Set the Initial Computing Interval Equal to Input Value or Equal to 0.005 if the Input Value is Zero

$$\Delta t = W_{10} \text{ or } \Delta t = 0.005 \text{ if } W_{10} = 0$$

7.1.4.5 Test Perturbation - Total Flag

If ICH = 1, go to Step 7.1.5 (numerical residual perturbation solution).

If ICH = 2, go to Step 7.1.4.6 (numerical solution of total equations of motion).

7.1.4.6 Test the First Print Time in Numerical Solution of Total Equations of Motion Versus Twice the Computing Interval

If  $t_{pT} < 2\Delta t$ , go to Step 7.1.4.7.

If  $t_{pT} = 2\Delta t$ , go to Step 7.1.4.8.

If  $t_{pT} > 2\Delta t$ , go to Step 7.1.5.

7.1.4.7 Set  $\Delta t = \frac{1}{2} t_{pT}$

7.1.4.8 Set IPRINT = 3 to Signal Both Print and Testing of Computed Values

7.1.5 Save Quantities at Start of Span for Computing Interval Calculations

Save positions, velocities, accelerations, time,  $t_s$ ,  $t_\psi$ ,  $t_\Omega$ ,  $t_\omega$ ,  $\theta_1$ , and  $\sin \theta_1$  for restart in case computing interval selection fails and for use in Simpson's rule calculations.

The array SS(17) is used for this purpose with:

$ss_1 = \text{time, } t$

$ss_2, ss_3, ss_4 = \text{position coordinates, } x_h, y_h, z_h$

$ss_5, ss_6, ss_7 = \text{velocity coordinates, } \dot{x}_h, \dot{y}_h, \dot{z}_h$

$ss_8, ss_9, ss_{10} = \text{acceleration components, } \ddot{x}_h, \ddot{y}_h, \ddot{z}_h$

$ss_{11} = t_s$

$ss_{12} = t_\psi$

$ss_{13} = t_\Omega$

$ss_{14} = t_\omega$

$ss_{15} = \theta_1$

$ss_{16} = \sin \theta_1$

$ss_{17} = \gamma$

#### 7.1.5.1 Increment Computing Interval Counter

$$KT = KT + 2$$

#### 7.1.6 Save Positions, Velocities, Accelerations, and Time for Ordinary Runge-Kutta Use

The array  $S(17)$  is used for this purpose and the indices are ordered as is the  $SS$  array in Step 7.1.5.

#### 7.1.7 Compute the Next Value of Time (i.e., At the End of the Next Computing Interval) and Make Special Time Tests

$$TN = s_1 + \Delta t$$

##### 7.1.7.1 Test Perturbation - Total Flag

If  $ICH = 1$  (numerical residual perturbation solutions), go to Step 7.1.7.2.

If  $ICH = 2$  (numerical solutions of total equations of motion), continue below (Step 7.1.7.2).

##### 7.1.7.2 Test Run Stop Time Against the Next Print Time

If  $T_f > t_p$ , go to Step 7.1.7.6.

If  $T_f \leq t_p$ , continue below (Step 7.1.7.3).

##### 7.1.7.3 Test the Next Value of Time Against Run Stop Time

If  $TN > T_f$ , continue below (Step 7.1.7.4).

If  $TN = T_f$ , go to Step 7.1.7.5.

If  $TN < T_f$ , go to Step 7.1.8.

##### 7.1.7.4 Set Computing Interval So That Time Will Equal the Run Stop Time at the End of the Next Interval

$$\Delta t = T_f - t$$

#### 7.1.7.5 Set Print and Halt Flags

IPRINT = 1

KHALT = 3

Go to Step 7.1.8.

#### 7.1.7.6 Test the Next Value of Time Against the Next Print Time

(Enter here from Step 7.1.7.2 if  $T_f > t_p$ .)

If  $TN > t_p$ , continue below (Step 7.1.7.7).

If  $TN = t_p$ , go to Step 7.1.7.8.

If  $TN < t_p$ , go to Step 7.1.8.

#### 7.1.7.7 Save Initial Conditions and Compute Interval for Print Only

Save the value of computing interval for use after print, and then set the computing interval so that the "next time" will equal the print time.

$$\Delta t_{sp} = \Delta t$$

$$\Delta t = t_p - t$$

Set the print flag.

IPRINT = 1

Save the values of positions, velocities, accelerations, and time in the array SP(10) for use after print. The indices correspond to those in Steps 7.1.5 and 7.1.6.

Set  $sp_I = s_I$  for  $I = 1, 2, \dots, 10$ .

Go to Step 7.1.8.

#### 7.1.7.8 Set IPRINT = 3

### 7.1.8 Complete First Runge-Kutta Pass

Compute intermediate value of time.

$$t = s_1 + \frac{1}{2} \Delta t$$

$$t_\psi = t_\psi + \frac{1}{2} \Delta t$$

$$t_s = t_s + \frac{1}{2} \Delta t$$

$$t_\Omega = t_\Omega + \frac{1}{2} \Delta t$$

$$t_\omega = t_\omega + \frac{1}{2} \Delta t$$

$$\gamma = \gamma + \frac{\varepsilon \Delta t}{2}$$

Compute the intermediate position values.

$$x_h = s_2 + \frac{1}{2} \Delta t \times s_5$$

$$y_h = s_3 + \frac{1}{2} \Delta t \times s_6$$

parameters for (141)

$$z_h = s_4 + \frac{1}{2} \Delta t \times s_7$$

Compute the Runge-Kutta parameters.

$$RK1X = \Delta t \times s_8$$

$$RK1Y = \Delta t \times s_9 \text{ computing intervals, etc. (140)}$$

$$RK1Z = \Delta t \times s_{10}$$

Set Runge-Kutta flag, KR = 2 and return to main program.

### 7.1.9 Test Whether Residual Perturbation of Total Equations are Being Solved

(Enter here from Step 7.1.3 on a first pass of Runge-Kutta which is not the first point of the trajectory.)

If ICH = 1 (residual perturbation), go to Step 7.1.12.

If ICH = 2 (total), continue below (Step 7.1.10).

### 7.1.10 Test Print Flag

If IPRINT = 1 (print), continue below (Step 7.1.11).

If IPRINT = 2 or 3 (don't print), go to Step 7.1.12.

### 7.1.11 Advance Print Index and Set the Next Print Time

$$ITP = ITP + 1$$

$$t_{PT} = P_t(ITP)$$

If IPRINT = 1, go to Step 7.1.11.1.

If IPRINT = 2, go to Step 7.1.5.

If IPRINT = 3, go to Step 7.1.11.2.

#### 7.1.11.1 Restore Values Saved Previously in Step 7.1.7.7

$$\Delta t = \Delta t_{sp}$$

$$s_I = sp_I \text{ for } I = 1, 2, \dots, 10$$

Set IPRINT = 2

Return to Step 7.1.7.

#### 7.1.11.2 Set IPRINT = 2 and Go to Step 7.1.5

#### 7.1.12 Test Simpson's Rule Flag

(Enter here from Step 7.1.9 when numerical residual perturbation solution is being computed and from Step 7.1.10 on a nonprint point during the numerical solution of the total equations of motion.)

If  $KC = 1$ , continued below (Step 7.1.13).

If  $KC = 2$ , go to Step 7.1.14.

#### 7.1.13 Advance Simpson's Rule Flag

$KC = 1$  signals that the first phase of the two-cycle Runge-Kutta is in progress.  $KC = 2$ , and return to Step 7.1.6.

#### 7.1.14 Reset Simpson's Rule Flag and Compute Actual and Allowable Errors

(Enter here from Step 7.1.12 if  $KC = 2$ . Signals that two Runge-Kutta cycles have been completed and accuracy tests can now be made.)

$$KC = 1$$

Compute the velocity increments by Simpson's rule:

$$\dot{x}_h = 1/3 \Delta t (ss_8 + 4s_8 + \ddot{x}_h)$$

$$\dot{y}_h = 1/3 \Delta t (ss_9 + 4s_9 + \ddot{y}_h)$$

$$\dot{z}_h = 1/3 \Delta t (ss_{10} + 4s_{10} + \ddot{z}_h)$$

Compute the estimated error (E) as follows:

Set  $C_{\max}$  = the maximum absolute value of the Runge-Kutta velocity increments over the last two cycles of Runge-Kutta  $\Delta X$ ,  $\Delta Y$ ,  $\Delta Z$ .

Determine E as the maximum absolute value of the respective differences between the Simpson's rule and Runge-Kutta velocity increments.

Reset the Runge-Kutta increments to zero for proper computation during the next two cycles:

$$\Delta\dot{X} = \Delta\dot{Y} = \Delta\dot{Z} = 0$$

Compute maximum and minimum allowable errors:

$$E_{\text{all}} = \max (|W_8| C_{\text{max}}, 10^{-9} \max (|\dot{x}_h|, |\dot{y}_h|, |\dot{z}_h|))$$

$$E_{\text{min}} = W_{12} C_{\text{max}}$$

7.1.15 Test the Estimated Error Against the Maximum Allowable Error

If  $E > E_{\text{all}}$ , continue below (Step 7.1.16).

If  $E < E_{\text{all}}$ , go to Step 7.1.22.

7.1.16 Increment and Test the Total Failure Counter

$$\text{KFAIL} = \text{KFAIL} + 1$$

Test the number of failures against the maximum allowed.

If  $\text{KFAIL} > \text{MFAIL}$ , continue below (Step 7.1.17).

If  $\text{KFAIL} < \text{MFAIL}$ , go to Step 7.1.18.

7.1.17 Set Halt Flag to Stop Trajectory Computations and Compute Final Time

$$\text{KHALT} = 3$$

$$T_f = t$$

Go to step 7.1.29.



### 7.1.18 Increment the Intermediate Failure Counter

(Enter here from Step 7.1.16 if the number of failures is less than the maximum allowed.)

$$KF = KF + 1$$

$$KHALT = 1$$

### 7.1.19 Compute a New $\Delta t$ Based on the Value of Estimated Error, E

$$\Delta t = FDT \cdot \Delta t (E_{\text{all}}/E)^{1/4}$$

If  $\Delta t/t > 10^{-8}$ , go to Step 7.1.20; otherwise print the computer interval and go to Step 7.1.17.

### 7.1.20 Test Intermediate Failure Counter

If  $KF \leq 0$ , continue below (Step 7.1.20.1).

If  $KF > 0$ , go to Step 7.1.21.

#### 7.1.20.1 Test Perturbation - Total Flag

If  $ICH = 1$ , go to Step 7.1.20.2.

If  $ICH = 2$ , go to Step 7.1.20.3.

#### 7.1.20.2 Perturbation Solution Calculations

$$KR = 5$$

##### 7.1.20.2.1 Check time against print time

If  $t \geq t_p$ , go to Step 7.1.20.2.2.

If  $t < t_p$ , go to Step 7.1.5.

7.1.20.2.2 Increment print time

$$t_p = t_p + \Delta t_p$$

Go to Step 7.1.29.

7.1.20.3 Test if This is Both a Print and a Regular Compute Point (IPRINT=3)

If IPRINT = 1 or 2, go to Step 7.1.5.

If IPRINT = 3, go to Step 7.1.20.2.

7.1.21 Restore Values Saved in Step 7.1.5 to Those in Step 7.1.6

Set  $s_I = ss_I$  for  $I = 1, 2, \dots, 17$

IPRINT = 2

Return to Step 7.1.8.

7.1.22 Test the Estimated Error Against the Minimum Allowable Error,  $E_{min}$

(Enter from Step 7.1.15 if the estimated error < the maximum allowed.)

If  $E \leq E_{min}$ , continue below (Step 7.1.23).

If  $E > E_{min}$ , go to Step 7.1.25.

7.1.23 Increment Total Failure Counter

$$KFAIL = KFAIL + 1$$

Test the number of failures against the maximum allowed.

If  $KFAIL > MFAIL$ , return to Step 7.1.17.

If  $KFAIL \leq MFAIL$ , continued below (Step 7.1.24).

7.1.24 Increment Intermediate Failure Counter, Compute New  $\Delta t$ , and Restore Initial Conditions

$$KF = KF + 1$$

$$KHALT = 1$$

Compute new  $\Delta t$  by multiplying the old by an input factor ( $>1$ ):

$$\Delta t = \Delta t_{\min} \Delta t$$

Restore the values saved in Step 7.1.5 to those in Step 7.1.6.

Set  $s_I = ss_I$  for  $I = 1, 2, \dots, 17$

Return to Step 7.1.17.

7.1.25 Set the Intermediate Failure Counter to Zero

(Enter here from Step 7.1.22 with  $E_{\min} < E < E_{\text{all}}$ .)

$$KF = 0$$

7.1.26 Second Pass Computations

(Enter here on second pass of Runge-Kutta ( $KR = 2$ ).)

Advance Runge-Kutta:

$$KR = 3$$

Compute new intermediate position values:

$$x_h = x_h + \frac{1}{4} (\Delta t) \quad (\text{RKIX})$$

$$y_h = y_h + \frac{1}{4} (\Delta t) \quad (\text{RKIY}) \quad \text{parameters for (142)}$$

$$z_h = z_h + \frac{1}{4} (\Delta t) \quad (\text{RKIZ})$$

Compute Runge-Kutta parameters:

$$RK2X = \Delta t \ddot{x}_h$$

$$RK2Y = \Delta t \ddot{y}_h \quad (141)$$

$$RK2Z = \Delta t \ddot{z}_h$$

Return to main program.

### 7.1.27 Third Pass Computations

(Enter here on third pass of Runge-Kutta (KR = 3).)

Advance Runge-Kutta flag:

$$KR = 4$$

Compute new time and intermediate position values:

$$t_\psi = t_\psi + \frac{1}{2} \Delta t$$

$$t_s = t_s + \frac{1}{2} \Delta t$$

$$t_\Omega = t_\Omega + \frac{1}{2} \Delta t$$

$$t_\omega = t_\omega + \frac{1}{2} \Delta t$$

$$\gamma = \gamma + \epsilon \Delta t / 2$$

$$T = s_1 + \Delta t$$

Compute Runge-Kutta parameters:

$$x_h = s_2 + \Delta t \times s_5 + \frac{1}{2} \Delta t \times \text{RK2X}$$

$$y_h = s_3 + \Delta t \times s_6 + \frac{1}{2} \Delta t \times \text{RK2Y} \quad \text{parameters for (143)}$$

$$z_h = s_4 + \Delta t \times s_7 + \frac{1}{2} \Delta t \times \text{RK2Z}$$

$$\text{RK3X} = \Delta t \ddot{x}_h$$

$$\text{RK3Y} = \Delta t \ddot{y}_h \quad (142)$$

$$\text{RK3Z} = \Delta t \ddot{z}_h$$

Return to main program.

#### 7.1.28 Fourth Pass Computations

(Enter here on fourth pass of Runge-Kutta (KR = 4).)

Reset Runge-Kutta flag:

$$\text{KR} = 1$$

Compute integrated position, velocities, and velocity increments.

$$x_h = s_2 + \Delta t \times s_5 + \frac{1}{6} \Delta t \quad (\text{RK1X} + \text{RK2X} + \text{RK3X})$$

$$y_h = s_3 + \Delta t \times s_6 + \frac{1}{6} \Delta t \quad (\text{RK1Y} + \text{RK2Y} + \text{RK3Y}) \quad (144)$$

$$z_h = s_4 + \Delta t \times s_7 + \frac{1}{6} \Delta t \quad (\text{RK1Z} + \text{RK2Z} + \text{RK3Z})$$

$$\dot{x}_h = s_5 + \frac{1}{6} (RK1X + 2 (RK2X + RK3X) + \Delta t \ddot{x}_h) \quad (143,$$

$$\dot{y}_h = s_6 + \frac{1}{6} (RK1Y + 2 (RK2Y + RK3Y) + \Delta t \ddot{y}_h) \quad 145)$$

$$\dot{z}_h = s_7 + \frac{1}{6} (RK1Z + 2 (RK2Z + RK3Z) + \Delta t \ddot{z}_h)$$

7.1.28.1 Test Whether Residual Perturbation Equations or Total Equations of Motion are Being Integrated

If ICH = 1 (perturbation), go to Step 7.1.28.3.

If ICH = 2 (total), go to Step 7.1.28.2.

7.1.28.2 Test Value of Print Flag

If IPRINT = 1, go to Step 7.1.29.

If IPRINT = 2 or 3, go to Step 7.1.28.3.

7.1.28.3 Compute Two Interval Increments in Velocity

$$\Delta \dot{X} = \Delta \dot{X} + \dot{x}_h - s_5$$

$$\Delta \dot{Y} = \Delta \dot{Y} + \dot{y}_h - s_6$$

$$\Delta \dot{Z} = \Delta \dot{Z} + \dot{z}_h - s_7$$

7.1.29 Return to Main Program

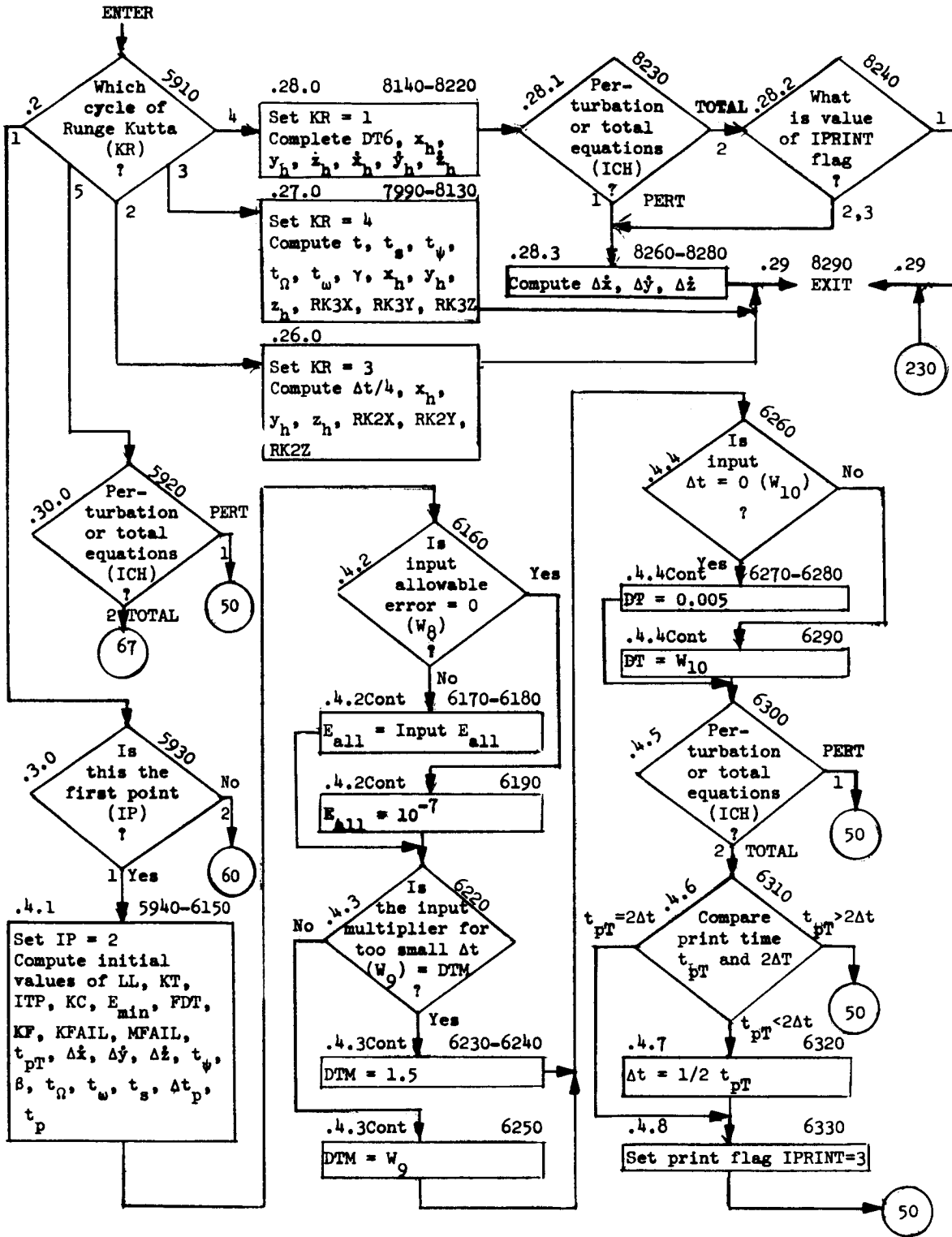
7.1.30 Test Whether Residual Perturbation or Total Equations of Motion are Being Integrated

If ICH = 1 (perturbation), go to Step 7.1.5.

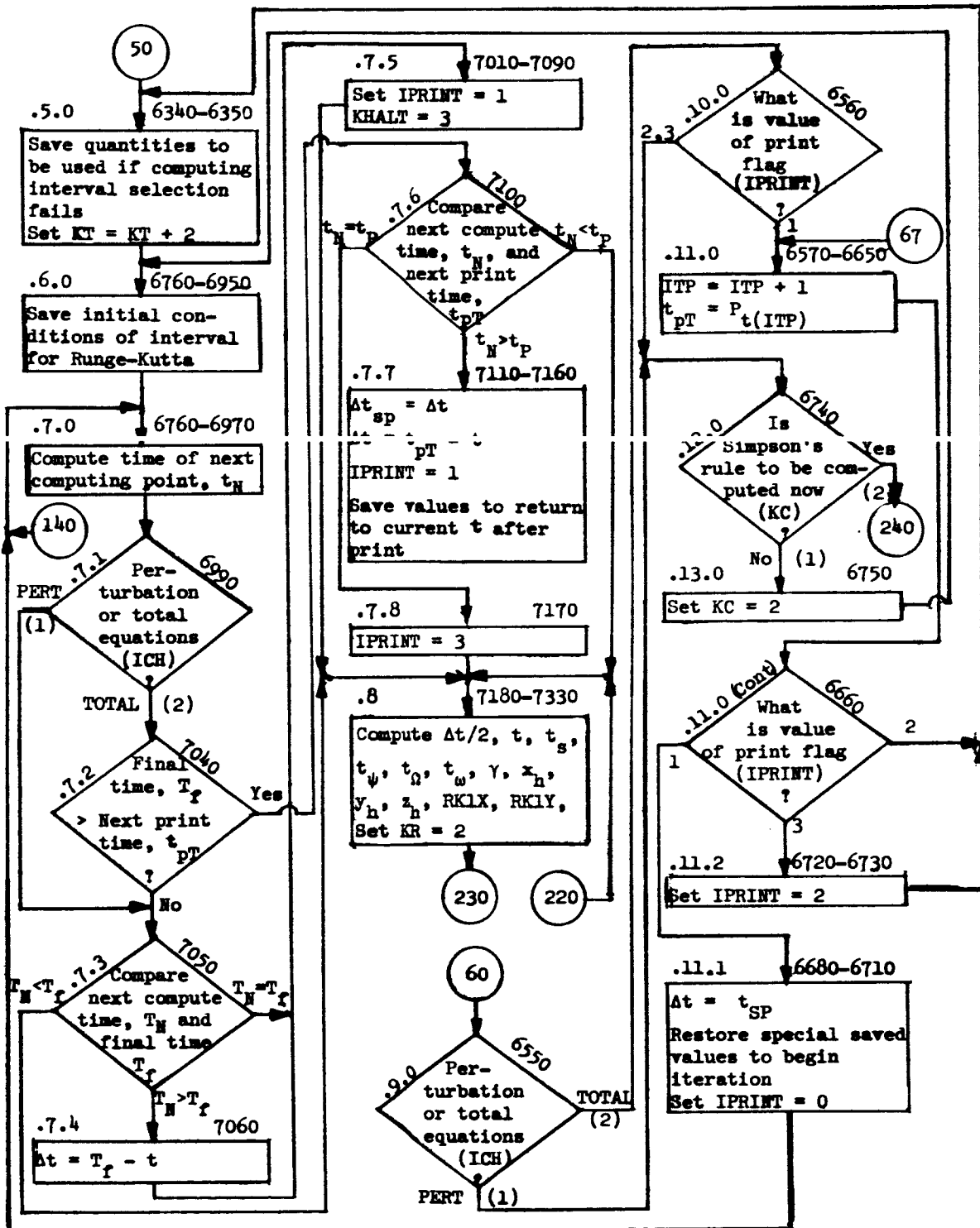
If ICH = 2 (total), go to Step 7.1.11.

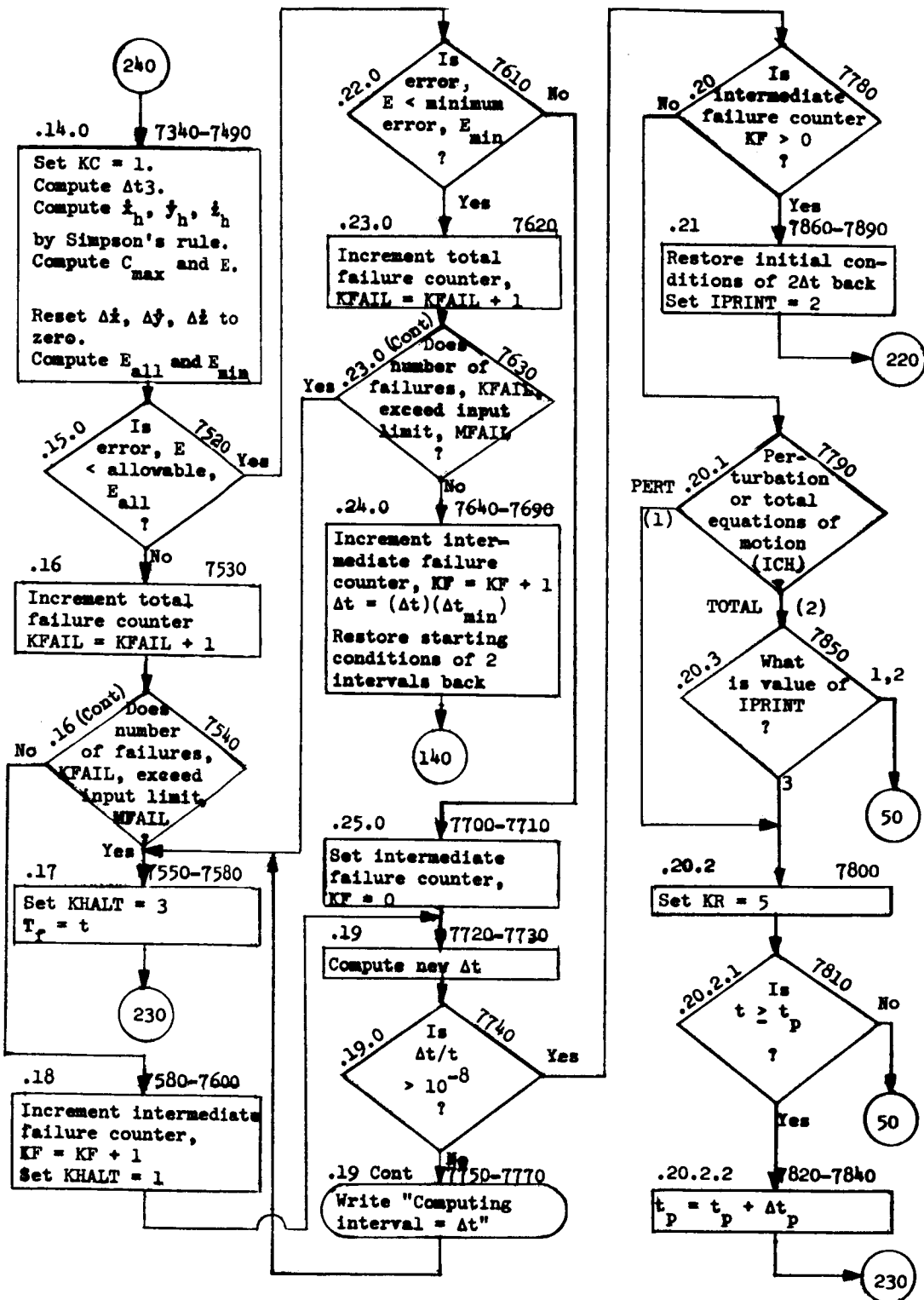
## 7.2 SUBROUTINE RKSIMP DETAIL FLOW CHARTS

The flow charts of this section describe the logic of subroutine RKSIMP completely, i.e., all conditional transfers are shown. The numbers to the top left of each box are the subsection numbers of Section 7.1 where the operations mentioned in the box are detailed. The numbers at the top right are the card numbers of the subroutine listing, Section 7.3.









### 7.3 RKSIMP LISTING

This section contains the listing for subroutine RKSIMP. The subsection numbers on the comment cards refer to the subsections in Section 7.1, wherein the code is explained.

M144 RKSIMP 09/01/65  
 EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

```

SUBROUTINE RKSIMP (KR, IP, KC, T, DT, ICH, HX, HY, HZ,
1 HXD, HYD, HXD, HX2D, HY2D, HZ2D, KHALT, TPRINT, TF)
COMMON M, IPRINT, TS, TFSI, TNOD, TPER, GAMMA, EPL,
1 TTP, PTIME, AP
DIMENSION W(30), S(17), SS(17), SP(17),
1 PTIME( 700),
2 AP( 9100)
C2.0 TEST RUNGE-KUTTA FLAG, KR
GO TO (10, 360, 370, 380, 51, KR
C3.0 TEST WHETHER RESIDUAL PERTURBATION OR TOTAL EQUATIONS OF MOTION
C ARF BEING INTEGRATED
5 GO TO (50, 67), ICH
C3.0 TEST THE INITIAL POINT FLAG, IP
10 GO TO (20, 60), IP
C4.0 FIRST POINT COMPUTATIONS
C4.1 SET FLAGS AND INITIAL CONDITIONS
20 IP = 2
LL = 0
KT = - 2
ITP = 1
KC = 1
FR MIN = ABS (W(12))
FDT = ABS (W(5))
KF = 0
KFAIL = 0
MFAIL = ABS (W(11))
TPRINT = PTIME(1)
DELXD = 0.
DELYD = 0.
DELZD = 0.
GAMMA = 0.
TPSI = 0.
BETA = 0.
TNOD = 0.
TPER = 0.
TS = 0.
TPI = .25 * (ABS (W(11))) ** 1.5
TPI = TPI
C4.2 SFT THE MAXIMUM ALLOWABLE ERROR EQUAL TO THE INPUT VALUE OR EQUAL
C TO 10 TO THE -7 IF THE INPUT VALUE IS ZERO
IF (W(8)) 21, 22, 21
21 FALL = ABS (W(8))
GO TO 30
22 FALL = 1. E-7
C4.3 SFT THE COMPUTING INTERVAL MULTIPLIER (USED IF ERROR IS LESS THAN
C THE MINIMUM ALLOWED) EQUAL TO INPUT VALUE OR EQUAL TO 1.5 IF INPUT
C VALUE IS LESS THAN OR EQUAL TO 1
30 IF (W(9) - 1.) 32, 32, 33
32 DTM = 1.5
GO TO 40
33 DTM = W(9)
C4.4 SET THE INITIAL COMPUTING INTERVAL EQUAL TO INPUT VALUE OR EQUAL
C TO 0.005 IF INPUT VALUE = 0
40 IF (W(10)) 42, 41, 42
41 DT = 0.005
GO TO 45
    
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5850H144  
 5860H144  
 5870H144  
 5880H144  
 5890H144  
 5900H144  
 5905H144  
 5910H144  
 5915H144  
 5916H144 \*1  
 5920H144 \*2  
 5925H144  
 5930H144  
 5933H144  
 5937H144 \*3  
 5940H144 \*4  
 5950H144 \*5  
 5960H144 \*6  
 5970H144 \*7  
 5980H144 \*8  
 5990H144 \*9  
 6000H144 \*10  
 6010H144 \*11  
 6020H144 \*12  
 6030H144 \*13  
 6040H144 \*14  
 6050H144 \*15  
 6060H144 \*16  
 6070H144 \*17  
 6080H144 \*18  
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 6100H144 \*20  
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 6156H144 \*26  
 6160H144 \*27  
 6170H144 \*28  
 6180H144  
 6190H144  
 6215H144 \*29  
 6220H144 \*30  
 6230H144 \*31  
 6240H144 \*32  
 6250H144  
 6255H144 \*33  
 6256H144 \*34  
 6260H144 \*35  
 6270H144  
 6280H144 \*36

H144 RKXIMP EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

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42 DT = W(10) 6290H144 .37
43 TEST PERTURBATION-TOTAL FLAG 6295H144
44 GO TO (50, 43), ICH 6300H144
C4.6 TEST FIRST PRINT TIME IN NUMERICAL SOLUTION OF TOTAL EQUATIONS IF 6305H144
C MOTION VS. TWICE THE COMPUTING INTERVAL 6306H144
43 IF (IPRINT - 2 * DT) < 46, 50 6310H144
C4.7 SET DT = 1/2 T SUB PT 6315H144
44 DT = 0.5 * IPRINT 6320H144
C4.8 SET IPRINT = 3 TO SIGNAL BOTH PRINT AND TESTING OF COMPUTED VAL 6325H144
C9.0 SAVE QUANTITIES AT START OF SPAN FOR COMPUTING INTERVAL 6330H144
C CALCULATIONS
90 S(1) = T 6340H144
S(2) = HX 6350H144
S(3) = HY 6360H144
S(4) = HZ 6370H144
S(5) = HXD 6380H144
S(6) = HYD 6390H144
S(7) = HZD 6400H144
S(8) = HX2D 6410H144
S(9) = HY2D 6420H144
S(10) = HZ2D 6430H144
S(11) = TS 6440H144
S(12) = TPI 6450H144
S(13) = TNGD 6460H144
S(14) = TPER 6470H144
S(15) = TMTA1 6480H144
S(16) = STMET1 6490H144
S(17) = GAMMA 6500H144
C5.1 INCREMENT COMPUTING INTERVAL COUNTER 6510H144
KI = KI + 2 6520H144
GO TO 110 6530H144
C9.0 TEST WHETHER RESIDUAL PERTURBATION OR TOTAL EQUATIONS ARE BEING 6540H144
C SOLVFD 6545H144
60 GO TO (90, 63), ICH 6550H144
C10.0 TEST THE PRINT FLAG 6555H144
63 GO TO (67, 90, 90), IPRINT 6560H144
C11.0 ADVANCE PRINT INDEX AND SET NEXT PRINT TIME 6565H144
67 ITP = ITP + 1 6570H144
IPRINT = PTIME (ITP) 6580H144
IF (IPRINT - 2) > 50, 85 6590H144
C11.1 RESTORE VALUES SAVED PREVIOUSLY IN STEP 7.1.7.7 6600H144
75 DT = SPDT 6605H144
GO RC I = 1, 17 6610H144
80 S(1) = SP(1) 6615H144
IPRINT = 2 6620H144
GO TO 140 6625H144
C11.2 SET IPRINT = 2 AND GO TO STEP 7.1.5 6630H144
85 IPRINT = 2 6635H144
GO TO 50 6640H144
C17.0 TEST THE SIMPSON'S RULE FLAG 6645H144
90 GO TO (100, 240), KC 6650H144
C13.0 ADVANCE THE SIMPSON'S RULE FLAG 6655H144
100 KC = 2 6660H144
C6.0 SAVE POSITIONS, VELOCITIES, ACCELERATIONS AND TIME FOR 6665H144
C ORDINARY RUNGE-KUTTA USE 6670H144
6675H144 .69
6680H144 .70
6685H144 .71
6690H144 .72
6695H144 .73
6700H144 .74
6705H144 .75

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F144 RKXIMP 09/01/65  
 EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

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110 S (1) = T
    S(2) = HX
    S(3) = HY
    S(4) = HZ
    S(5) = HXD
    S(6) = HYD
    S(7) = HZD
    S(8) = HXD2
    S(9) = HYD2
    S(10) = HZD2
    S(11) = TS
    S(12) = TPSI
    S(13) = INOD
    S(14) = TPER
    S(15) = THETA1
    S(16) = STHET1
    S(17) = GAMMA
    GO TO 140
C7.0 COMPUTE THE NEXT VALUE OF TIME AND MAKE SPECIAL PRINT TIME TESTS
    140 TN = S(1) + DT
C7.1 TEST PERTURBATION-TOTAL FLAG
    139 GO TO (150, 145), ICH
C7.2 TEST RUN STOP TIME AGAINST THE NEXT PRINT TIME
    145 IF (TF - TPRINT) 150, 150, 180
C7.3 TEST THE NEXT VALUE OF TIME AGAINST RUN STOP TIME
    150 IF (TN - TF) 220, 170, 160
C7.4 SET COMPUTING INTERVAL SO THAT TIME WILL EQUAL THE RUN STOP TIME
    C AT THE END OF THE NEXT INTERVAL
    160 CI = TF - T
C7.5 SET THE PRINT AND HALT FLAGS
    170 IPRINT = 1
        KHALT = 3
        GO TO 220
C7.6 TEST THE NEXT VALUE OF TIME AGAINST THE NEXT PRINT TIME
    180 IF (TN - TPRINT) 220, 200, 190
C7.7 SAVE INITIAL CONDITIONS AND COMPUTE INTERVAL FOR PRINT ONLY
    190 SPDT = DT
        DT = TPRINT - T
        IPRINT = 1
        DO 210 I = 1, 17
            S(I) = S(I)
        GO TO 220
C7.8 INDICATE COMPUTE POINT = PRINT POINT
    200 IPRINT = 3
C8.0 COMPLETE 1ST R-K PASS, COMPUTE NEW TIME AND POSITIONS
    220 DT2 = DT / 2.
        T = S(1) + DT2
        TS = S(11) + DT2
        TPSI = S(12) + DT2
        TMOD = S(13) + DT2
        TPER = S(14) + DT2
        GAMMA = S(17) + EP1 * DT2
        HX = S(2) + DT2 * S(5)
        HY = S(3) + DT2 * S(6)
        HZ = S(4) + DT2 * S(7)
        RKIX = DT * S(8)
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09/01/55  
INTERNAL FORMULA NUMBER(S)

M144 RKSTMP  
EXTERNAL FORMULA NUMBER - SOURCE STATEMENT -

901	GO TO 230	7300H144	.122
		7310H144	.123
		7320H144	.124
		7330H144	.125
		7340H144	.126
		7350H144	.127
		7360H144	.128
		7370H144	.129
		7380H144	.130
		7390H144	.131
		7400H144	.132
		7410H144	.133
		7420H144	.134
		7430H144	.135
		7440H144	.136
		7450H144	.137
		7460H144	.138
		7470H144	.139
		7480H144	.140
		7490H144	.141
		7500H144	.142
		7510H144	.143
		7520H144	.144
		7530H144	.145
		7540H144	.146
		7550H144	.147
		7560H144	.148
		7570H144	.149
		7580H144	.150
		7590H144	.151
		7600H144	.152
		7610H144	.153
		7620H144	.154
		7630H144	.155
		7640H144	.156
		7650H144	.157
		7660H144	.158
		7670H144	.159
		7680H144	.160
		7690H144	.161
		7700H144	.162
		7710H144	.163
		7720H144	
		7730H144	
		7740H144	
		7750H144	

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RK1V = DT * S19)
RK1Z = DT * S110)
KR = 2
901 GO TO 230
C14.0 RSFT SIMPSONS RULE FLAG AND COMPUTE ACTUAL AND ALLOWABLE ERRORS
240 KC = 1
DT3 = DT / 3.
HXDA = DT3 * (SS(8) + 4. * S18) + HXZD)
HYDA = DT3 * (SS(9) + 4. * S19) + HYZD)
HZDA = DT3 * (SS(10) + 4. * S110) + HZZD)
C COMPUTE MAXIMUM ESTIMATED ERROR
C MAX = AMAX1 (ABS(DELXD), ABS(DELXD), ABS(DELZD))
ESTPR = AMAX1 (ABS(DELXD) - HXDA), ABS(DELXD - HYDA).
DELXC = 0.
DELYD = 0.
DELZD = 0.
FALL = AMAX1 (ABS (W18)) * CMAX, 1.E-9 * AMAX1 (ABS (HXD
1 J. ABS (HYD), ABS (HZD)))
ERMIN = ABS (W12)) * CMAX
C15.0 TEST THE ESTIMATED ERROR AGAINST THE MAXIMUM ALLOWABLE ERROR
C16.0 INCREMENT AND TEST THE TOTAL FAILURE COUNTER
250 KFAIL = KFAIL + 1
C17.0 SFT HALT FLAG TO STOP TRAJECTORY COMPUTATIONS AND COMPUTE FINAL
C TIME
260 KHALT = 3
TF = T
GO TO 230
C18.0 INCREMENT THE INTERMEDIATE FAILURE COUNTER
270 KF = KF + 1
KHALT = 1
GO TO 330
C22.0 TEST THE ESTIMATED ERROR AGAINST THE MINIMUM ALLOWABLE ERROR
C (FRMIN)
280 IF (FSTER - ERMIN) 290, 290, 320
C23.0 INCREMENT THE TOTAL FAILURE COUNTER AND TEST THE NUMBER OF
C FAILURES AGAINST THE MAXIMUM ALLOWED
290 KFAIL = KFAIL + 1
C24.0 IF (KFAIL - MFAIL) 300, 300, 260
C AND RESTORE INITIAL CONDITIONS
300 KF = KF + 1
KHALT = 1
DT = DTM * DT
DO 310 I = 1, 17
310 S (I) = SS (I)
GO TO 140
C25.0 SFT THE INTERMEDIATE FAILURE COUNTER TO ZERO
320 KF = 0
GO TO 330
C19.0 COMPUTE NEW DT BASED ON THE VALUE OF THE ESTIMATED ERROR, ESTER
330 DT = FDT * DT * (EALL / ESTER) ** 0.25
IF (DT / T - 1.E-8) 334, 334, 336
334 WRITE (6,335) DT

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H144  RKSTMP
INTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)
09/01/65
335 FORMAT (1H0.16H COMP INTERVAL = E17.8)
GO TO 260
C20.0 TEST INTERMEDIATE FAILURE COUNTER
336 IF (KF) 331, 331, 340
C20.1 TEST PERTURBATION-TOTAL FLAG
331 GO TO (332, 333), ICH
C20.2 PERTURBATION SOLUTION CALCULATIONS
332 KR = 5
C20.2.1 CHECK TIME AGAINST PRINT TIME
IF (T - TPI) 50, 337, 337
C20.2.2 INCREMENT PRINT TIME
337 TPI = TPI + TPI
GO TO 230
C20.3 TEST IF THIS IS BOTH A PRINT AND A REGULAR COMPUTE POINT
C (IPRINT = 3)
333 GO TO ( 50, 50, 332), IPRINT
C21.0 RSTORE VALUES SAVED IN STEP 7.1.5 TO THOSE IN STEP 7.1.6
340 DO 350 I = 1, 17
350 S (I) = SS (I)
IPRINT = 2
GO TO 220
C26.0 SECOND PASS OF RUNGE-KUTTA
360 KR = 3
DT4 = DT / 4.
HX = HX + DT4 * RK1X
HY = HY + DT4 * RK1Y
HZ = HZ + DT4 * RK1Z
RK2X = DT * HX20
RK2Y = DT * HY20
RK2Z = DT * HZ20
GO TO 230
C27.0 THIRD PASS OF RUNGE-KUTTA
370 KR = 4
T = S(1) + DT
TS = S(11) + DT
TPS1 = S(12) + DT
TNDD = S(13) + DT
TPFR = S(14) + DT
GAMMA = S(17) + FPI * DT
HX = S(2) + DT * S(5) + DT2 * RK2X
HY = S(3) + DT * S(6) + DT2 * RK2Y
HZ = S(4) + DT * S(7) + DT2 * RK2Z
RK3X = DT * HX20
RK3Y = DT * HY20
RK3Z = DT * HZ20
GO TO 230
C28.0 FOURTH PASS OF RUNGE-KUTTA
380 KR = 1
DT6 = DT / 6.
HX = S(2) + DT * S(5) + DT6 * (RK1X + RK2X + RK3X)
HY = S(3) + DT * S(6) + DT6 * (RK1Y + RK2Y + RK3Y)
HZ = S(4) + DT * S(7) + DT6 * (RK1Z + RK2Z + RK3Z)
HXD = S(5) + (RK1X + 2 * (RK2X + RK3X) + DT * HX20) / 6.
HYD = S(6) + (RK1Y + 2 * (RK2Y + RK3Y) + DT * HY20) / 6.
HZD = S(7) + (RK1Z + 2 * (RK2Z + RK3Z) + DT * HZ20) / 6.
C28.1 TEST WHETHER RESIDUAL PERTURBATION EQUATIONS OR TOTAL EQUATIONS

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7760H144
7770H144
7775H144
7780H144
7785H144
7790H144
7795H144
7800H144
7805H144
7810H144
7815H144
7820H144
7830H144
7835H144
7836H144
7840H144
7845H144
7850H144
7860H144
7870H144
7880H144
7890H144
7900H144
7910H144
7920H144
7930H144
7940H144
7950H144
7960H144
7970H144
7980H144
7990H144
8000H144
8010H144
8020H144
8030H144
8040H144
8050H144
8060H144
8070H144
8080H144
8090H144
8100H144
8110H144
8120H144
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8140H144
8150H144
8160H144
8170H144
8180H144
8190H144
8200H144
8210H144
8220H144
8225H144

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M144	RKSI MP	EXTERNAL FORMULA NUMBER	SOURCE STATEMENT	INTERNAL FORMULA NUMBER(S)
			OF MOTION ARE BEING INTEGRATED	8226H144 .207
		GO TO 1387, 381), ICH		8230H144
		C28.7 TEST VALUE OF PRINT FLAG		8235H144 .208
		381 GO TO (230, 382, 382), IPRINT		8240H144
		C29.3 COMPUTE TWO INTERVAL INCREMENT IN VELOCITY		8245H144 .209
		382 DFLXD = DELXD + (RK1X + 2. * (RK2X + AK3X) + DT * HX2D) / 6.		8250H144 .210
		DFLYD = DELYD + (RK1Y + 2. * (RK2Y + AK3Y) + DT * HY2D) / 6.		8260H144 .211
		DFLZD = DELZD + (RK1Z + 2. * (RK2Z + AK3Z) + DT * HZ2D) / 6.		8270H144
		C29-0 RETURN TO MAIN PROGRAM		8275H144 .212
		230 RETURN		8280H144 .213
		END		8290H144 .214

Section 8  
SUBROUTINE PLOT

8.1 DEVELOPMENT OF EQUATIONS

The purpose of this subroutine is to plot the error in the Jacobi constant,  $\Delta C/C$ , for the analytic solution, for the numerical residual perturbation solution, for the precessing ellipse which forms the basis of the numerical perturbation solution, for a straightforward numerical solution of the total equations of motion, and for the initial osculating ellipse which would form the basis of an ordinary Encke method; and to plot the vector position and velocity differences of each of the other solutions from the numerical residual perturbation solution. The numerical residual perturbation solution is taken as the standard of comparison because it normally holds the Jacobi integral most constant. The routines are taken from reference 8.

The abscissa of the plots is time. The plotting symbols are 6 rasters wide and the length of the scales is 1,024 rasters less space required for labeling and scales. This allows about 152 points to a plot. Hand plots indicate a point about every quarter radian of true anomaly is necessary to describe the motion adequately. In the dimensionless variables, the mean motion is  $a^{-3/2}$  (ref 1, page 208). Thus, the time to be plotted per page to produce 152 points (corresponding to  $152/4 = 38$  radians) is given by:

$$t_R = (38a^{3/2})$$

To make the plot readable,  $t_R$  is rounded to have only one significant figure. If

$1 \times 10^n \leq t_R < 2 \times 10^n$ , then grid is drawn at intervals of  $0.1 \times 10^n$ ; if  $2 \times 10^n \leq t_R < 5 \times 10^n$ , then the grid is drawn at intervals of  $0.2 \times 10^n$ ; and if  $5 \times 10^n \leq t_R < 10^{n+1}$ , then the grid is drawn at intervals of  $0.5 \times 10^n$ .

## 8.2 EQUATIONS IN ORDER OF SOLUTIONS

### 8.2.1 Input

The input to subroutine PLOT is:

- $a_{32}$  = The semi-major axis to the three-halves power
- $T_f$  = Final time
- $a$  = The semi-major axis
- $e$  = The eccentricity
- $i^\circ$  = The inclination in degrees
- $\mu$  = The ratio of the mass of the smaller body to the sum of the masses of the two bodies
- $C_{init}$  = The Jacobi constant
- $P_{ss}$  = The anomalistic period of the precessing ellipse
- PTIME = Up to 700 values of time at which the data are to be plotted
- AP = A 9,100-element array defined as follows, (where  $i = 0, 1, 2, \dots, 699$  is the plot point count index):

$ap_{1+13i} = (\Delta C/C)_{ai}$       The subscripts are defined by:

$ap_{2+13i} = (\Delta C/C)_{pki}$        $a$  = analytic solution from reference 1

$ap_{3+13i} = (\Delta C/C)_{ci}$        $pk$  = precessing ellipse value

$ap_{4+13i} = (\Delta C/C)_{ki}$        $c$  = value obtained from numerical solution of total equations of motion

$$ap_{5+13i} = (\Delta C/C)_{pi} \quad k = \text{initial osculating ellipse value}$$

$$ap_{6+13i} = (\Delta \rho/a)_{ai} \quad p = \text{value from numerical residual perturbation solution}$$

$$ap_{7+13i} = (\Delta \rho/a)_{pki}$$

$$ap_{8+13i} = (\Delta \rho/a)_{ci}$$

$$ap_{9+13i} = (\Delta \rho/a)_{ki}$$

$$ap_{10+13i} = (\Delta \dot{\rho} a^{1/2})_{ai}$$

$$ap_{11+13i} = (\Delta \dot{\rho} a^{1/2})_{pki}$$

$$ap_{12+13i} = (\Delta \dot{\rho} a^{1/2})_{ci}$$

$$ap_{13+13i} = (\Delta \dot{\rho} a^{1/2})_{ki}$$

### 8.2.2 Initial Setup

1. The machine is instructed to provide both 9 in. by 9 in. transparencies and 35mm slides.
2. The plotting symbols are chosen as:

o = Analytic solution

x = Precessing ellipse

□ = Numerical solution of total equations of motion

\* = Initial osculating ellipse

• = Numerical residual perturbation solution

3. Compute abscissa scale.

Set initial time for initial plot = 0.

A. Compute unrounded  $t_R$ :

$$t_R = 38. + a_{32}$$

B. Set number of characters to be displayed in horizontal labels:

$$NX = 3$$

C. Develop  $A = 10^n$  such that  $10^n \leq t_R \leq 10^{n+1}$

(1) Set  $A = 1$

(2) Compare  $t_R$  and 1:

If  $t_R < 1$ , go to Step 8.2.2-3 C(3).

If  $t_R = 1$ , go to Step 8.2.2-3 C(6).

If  $t_R > 1$ , go to Step 8.2.2-3 C(5).

(3) Develop  $A < 1$ :

Set  $A = A/10$  and increment digit count  $NX$ :

$$NX = NX+1$$

Compare  $t_R$  and  $A$ :

If  $t_R < A$ , repeat Step 8.2.2-3 C(3).

If  $t_R \geq A$ , go to Step 8.2.2-3 C(6).

(4) Set  $A = B$  and increment  $NX$ :

$$A = B$$

$$NX = NX + 1$$

(5) Develop  $A > 1$ :

$$B = 10 A$$

Compare  $t_R$  and  $B$ :

If  $t_R \leq B$ , go to Step 8.2.2-3 C(6).

If  $t_R > B$ , go to Step 8.2.2-3 C(4).

(6) Compute integer portion of  $t_R/A = t_i$ :

Compare  $t_i$  and 2:

If  $t_i < 2$  (i.e., = 1), go to Step 8.2.2-3D.

If  $t_i = 2$ , go to Step 8.2.2-3E.

If  $t_i > 2$ , go to Step 8.2.2-3F.

D. Compute  $t_R$  and vertical line spacing for  $t_i = 1$ :

$$\delta X = A/10$$

$$t_R = 10\delta X t_i$$

Go to Step 8.2.2-3H.

E. Compute  $t_R$  and vertical line spacing for  $t_i = 2, 3, \text{ or } 4$ :

$$\delta X = A/5$$

$$t_R = 5\delta X t$$

Go to Step 8.2.2-3H.

F. Compare  $t_i$  and 4:

If  $t_i \leq 4$  (i.e., = 3 or 4), go to Step 8.2.2-3E.

If  $t_i > 4$ , go to Step 8.2.2-3G.

G. Compute  $t_R$  and vertical line spacing for  $4 < t_i < 10$ :

$$\delta X = A/2$$

$$t_R = 2\delta X t_i$$

H. Print scale values:

$$t_R, A, \delta X, t_i, NX$$

4. The machine is instructed to use semi-log scales, log on vertical, and linear on horizontal.
5. The point count index II for beginning of plot point number is set  $II = 1$ . The page count IIK is set = 1.

### 8.2.3 Computation for Each Set of Three Plots

1. Set point count index III for maximum end of plot point number  $III = II + 152$ .

2. Compute initial time EM and final time EN for each set of plots:

$$EM = EN$$

$$EN = EM + t_R$$

8.2.4 Set L = 1 to Indicate Plot for  $\Delta C/C$

8.2.5 Draw Grid and Do Common Labeling

1. Draw grid (subroutine GRIDV), calling sequence: general (L, XL, XR, YB, YT, DX, DY, N, M, I, J, NX, NY)

L = 1 indicates film to be advanced, job number, and frame count to be displayed

XL = EM is first time

XR = EN is last time

YB =  $10^{-10}$  is bottom of ordinate

YT = 1 is top of ordinate

DX =  $\delta X$  indicates vertical lines drawn at increments of  $\delta X$

DY = 1 indicates horizontal lines drawn at increments of 1

N = 0 causes all vertical lines to be of equal intensity

M = 2 with logarithmic ordinate, this is a dummy

I = -1 causes each vertical line to be labeled and labels to be below the scale

J = -4 with logarithmic ordinate, this is a dummy



NX = NX indicates NX significant places on X scale

NY = 1 indicates 1 significant place in Y scale

2. List and define symbols.

<u>Symbol</u>	<u>Location</u> (Raster Counts)	<u>Definition</u>	<u>Location</u> (Raster Counts)
o	125, 939	ANALYTIC SOLUTION	145-289, 939
x	125, 919	PRECESSING ELLIPSE (REFERENCE ORBIT FOR NUMERICAL RESIDUAL PRETURBATION SOLUTION)	145-408, 919 145-432, 899
□	125, 879	NUMERICAL SOLUTION OF TOTAL EQUATIONS	145-416, 879
*	125, 859	INITIAL OSCULATING ELLIPSE	145-336, 859

3. Label horizontal scale.

The work "TIME" goes in raster counts 475-507, 0.

4. List run parameters and time scale.

"SEMI-MAJOR AXIS = 'a' ECCENTRICITY = 'e' INCLINATION (DEGREES)  
= 'i' MASS RATIO = 'μ'" is printed in raster counts 162-942, 989.

"TIME RANGE FROM 'EN' TO 'EM'" is printed in raster counts 434-658,  
969.

5. If L = 1, go to Step 8.2.6.

If L = 2, go to Step 8.2.7.

If L = 3, go to Step 8.2.8.

### 8.2.6 Label and Plot Error in Jacobi Integral, $\Delta C/C$

1. Label vertical scale and title.

"ABSOLUTE VALUE OF  $\Delta C/C$ " is entered in raster counts 0, 800-348.

Label plot "ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL" in raster counts 302-727, 1012.

2. Add additional symbol definition and parameter values. The symbol '.' is printed at raster locations 125, 839.

"NUMERICAL RESIDUAL PERTURBATION SOLUTION" is printed in raster positions 145-465, 839.

"JACOBI CONSTANT,  $C = 'C_{init}'$ " is printed at raster positions 132-348, 969.

"ANOMALISTIC PERIOD = ' $P_{ss}$ '" is printed at raster positions 718-950, 969.

3. Plot  $\Delta C/C$ .

A. Set point count index  $i = II$ .

B. If  $t_i \leq t_f$ , go to Step 8.2.6-3C; otherwise go to Step 8.2.7.

C. Set index  $K$  to find location of  $(\Delta C/C)_{ai}$  in array AP:

$$K = -12 + 13i$$

D. If  $t_i \leq EN$ , go to Step 8.2.6-3E; otherwise go to Step 8.2.7.

E. Plot  $(\Delta C/C)_{ai}$ ,  $(\Delta C/C)_{pki}$ ,  $(\Delta C/C)_{ci}$ ,  $(\Delta C/C)_{ki}$ , and  $(\Delta C/C)_{pi}$ .

F. If  $i < III$ , set  $i = i+1$  and return to Step 8.2.6-3B; otherwise go to Step 8.2.7.

### 8.2.7 Label and Plot Position Differences from Numerical Residual

#### Perturbation Value, $\Delta\rho/a$

1. Draw grid and do common labeling.

Set  $L = 2$  and go to Step 8.2.5.

2. Label vertical scale and title.

"DELTA RHO/A" is printed in raster locations 0, 602-446.

"NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION" is printed in raster locations 302-822, 1012.

3. Plot  $\Delta\rho/a$

A. Set point count index  $i = II$ .

B. If  $t_i \leq t_f$ , go to Step 8.2.7-3C; otherwise go to Step 8.2.8.

C. Set index  $K$  to find location of  $(\Delta\rho/a)_{ai}$ :

$$K = -7 + 13i$$

D. If  $t_i \leq EN$ , go to Step 8.2.7-3E; otherwise go to Step 8.2.8.

E. Plot  $(\Delta\rho/a)_{ai}$ ,  $(\Delta\rho/a)_{pki}$ ,  $(\Delta\rho/a)_{ci}$ , and  $(\Delta\rho/a)_{ki}$ .

F. If  $i < III$ , set  $i = i+1$  and return to Step 8.2.7-3B; otherwise go to Step 8.2.8.

8.2.8 Label and Plot Velocity Differences from Numerical Residual Perturbation Value,  $\Delta \dot{\rho} \sqrt{a}$

1. Draw grid and do common labeling. Set  $L = 3$  and go to Step 8.2.5.
2. Label vertical scale and title.

"DELTA RHO DOT TIMES SQ RT A" is printed in rasters 0, 620-324.

"NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION" is printed in raster count 258-866, 1012.

3. Plot  $(\Delta \dot{\rho} a^{1/2})_{ai}$ ,  $(\Delta \dot{\rho} a^{1/2})_{pki}$ ,  $(\Delta \dot{\rho} a^{1/2})_{ci}$ , and  $(\Delta \dot{\rho} a^{1/2})_{ki}$ .
  - A. Set point count index  $i = II$ .
  - B. If  $t_i \leq t_I$  go to Step 8.2.8-3C; otherwise go to Step 8.2.10.
  - C. Set index  $K$  to find location of  $(\Delta \dot{\rho} a^{1/2})_{ai}$  in array AP:

$$K = -3 + 13i$$

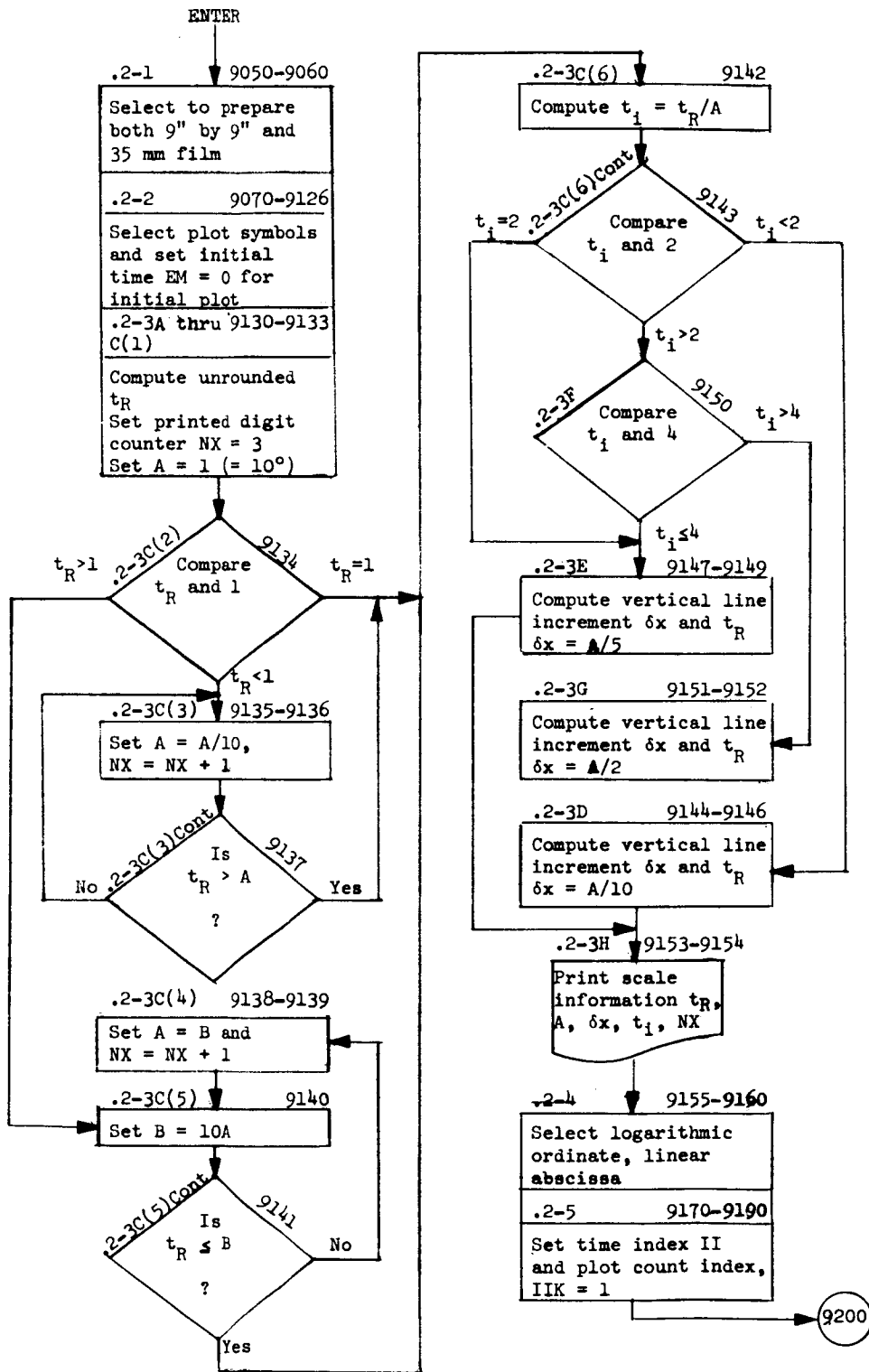
- D. If  $t_i < EN$ , go to Step 8.2.8-3E.  
If  $t_i = EN$ , go to Step 8.2.8-3F.  
If  $t_i > EN$ , go to Step 8.2.8-3G.
    - E. Plot  $(\Delta \dot{\rho} a^{1/2})_{ai}$ ,  $(\Delta \dot{\rho} a^{1/2})_{pki}$ ,  $(\Delta \dot{\rho} a^{1/2})_{ci}$ , and  $(\Delta \dot{\rho} a^{1/2})_{ki}$ .  
If  $i < III$ , set  $i = i+1$  and return to Step 8.2.8-3B; otherwise set  $II = III$  and go to Step 8.2.9.
    - F. Plot  $(\Delta \dot{\rho} a^{1/2})_{ai}$ ,  $(\Delta \dot{\rho} a^{1/2})_{pki}$ ,  $(\Delta \dot{\rho} a^{1/2})_{ci}$ , and  $(\Delta \dot{\rho} a^{1/2})_{ki}$ .
    - G. Set  $II = i$ .

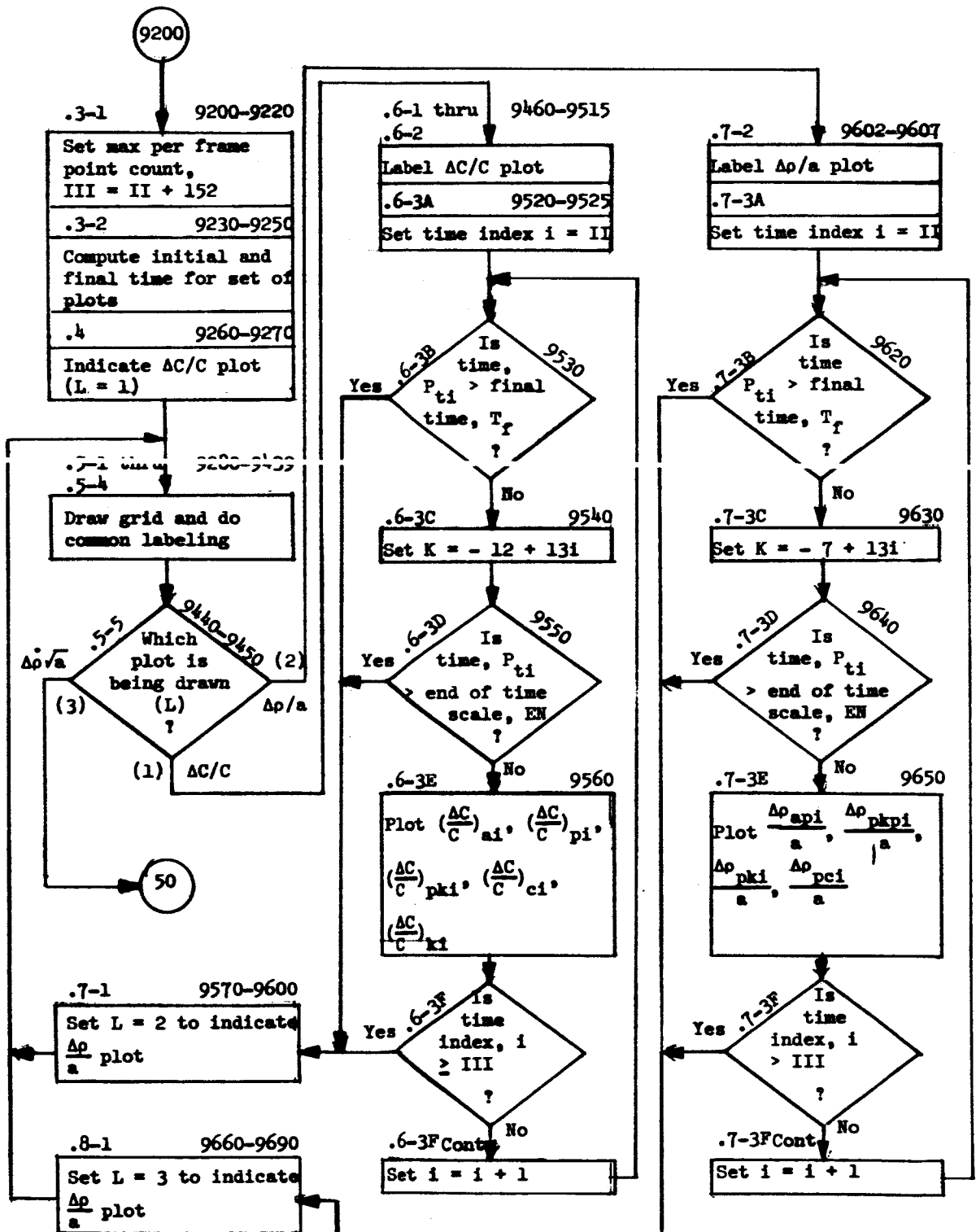
8.2.9 If  $IIK < 6$ , set  $IIK = IIK + 1$  and Return to Step 8.2.3-1; Otherwise  
Go To Step 8.2.10.

8.2.10 Return to Main Program

### 8.3 SUBROUTINE PLOT DETAIL FLOW CHARTS

The flow charts of this section describe the logic of subroutine PLOT completely, i.e., all conditional transfers are shown. The numbers to the top left of each box are the subsection numbers of Section 8.2 where the operations mentioned in the box are detailed. The numbers at the top right are the card numbers of the subroutine listing, Section 8.4.









#### 8.4 PLOT LISTING

This section contains the listing for subroutine PLOT. The subsection numbers on the comment cards refer to the subsections in Section 8.2, wherein the code is explained.

H144 PLOT 09/01/65  
EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

SUBROUTINE PLOT (A32,IF,AL,ECL,DEGI,EP4,JACOBI,PSS)  
DIMENSION W(30), PTIME( 700),API( 9100), MRKPT(5)  
COMMON W,IPRINT,IS,TPSI,TNOD,TPER,GAMMA,EPI,IIP,PTIME  
I,AP

C8.2.2 INITIAL SETUP

C8.2.2-1 SET FOR ROTH 9 IN BY 9 IN AND 35 MM OUTPUT

CALL CAMRAV (935)

C8.2.2-2 SELECTION OF PLOTTING SYMBOLS

MRKPT (1) = 38

MRKPT (2) = 55

MRKPT (3) = 63

MRKPT (4) = 44

MRKPT (5) = 42

C8.2.2-3 COMPUTE ABSCISSA SCALE

EN = 0.

TR = 38. \* A32

NX = 3

A = 1.

210 IF (TR - A)220,230,240

270 A = A /10.

NX = NX + 1

IF (TR - A) 220,230,230

250 A = R

NX = NX + 1

240 R = A \* 10.

IF (TR - R) 230,230,250

230 TI = AINT (TR / A)

IF (TI - 2.) 260,270,280

260 DX = A /10.

TR = DX \* 10. \* TI

GO TO 300

270 CX = A/5.

TR = DX \* 5. \* TI

GO TO 300

280 IF (TI - 4.) 270,270,290

290 DX = A/2.

TR = DX \* 2. \* TI

300 WRITE (6,310) TR, A, DX, TI, NX

310 FORMAT (1H 4E18.8, 18)

C8.2.2-4 SET FOR LOGARITHMIC ORDINATE, LINEAR ABSCISSA

CALL SMXYV ( 0, 1)

IF (I = 1

DO 150 I,1,1.6

C8.2.3 COMPUTATION FOR EACH SET OF 3 PLOTS

C8.2.3-1 SET MAXIMUM PER FRAME POINT COUNT

III = II + 152

IF (I = 1

C8.2.3-2 COMPUTE INITIAL AND FINAL TIME FOR EACH SET OF PLOTS

EM = FN

EN = FN + TR

C8.2.4 SET L = 1 TO INDICATE PLOT FOR DELTA C/C

L=1

C8.2.5 DRAW GRID AND DO COMMON LABELING

C8.2.5-1 DRAW GRID (SUBROUTINE GRIDV)

20 CALL GRIDIV (1,EM,EN,1,E-10,1,0,DX,1,0,0,2, -1,-4,NX,-1)

9010H144 \*1  
9020H144 \*2  
9030H144 \*3  
9031H144 \*4  
9040H144 \*5  
9050H144 \*6  
9060H144 \*7  
9070H144 \*8  
9080H144 \*9  
9090H144 \*10  
9100H144 \*11  
9110H144 \*12  
9120H144 \*13  
9123H144 \*14  
9126H144 \*15  
9131H144 \*16  
9132H144 \*17  
9133H144 \*18  
9134H144 \*19  
9135H144 \*20  
9136H144 \*21  
9137H144 \*22  
9138H144 \*23  
9139H144 \*24  
9140H144 \*25  
9141H144 \*26  
9142H144 \*27  
9143H144 \*28  
9144H144 \*29  
9145H144 \*30  
9146H144 \*31  
9147H144 \*32  
9148H144 \*33  
9149H144 \*34  
9150H144 \*35  
9151H144 \*36  
9152H144 \*37  
9153H144 \*38  
9154H144 \*39  
9155H144  
9160H144  
9170H144  
9180H144  
9190H144  
9200H144  
9210H144  
9220H144  
9230H144  
9240H144  
9250H144  
9260H144  
9270H144  
9280H144  
9290H144  
9300H144

M144 PLOT FATFARNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S) 09/01/65

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CB.2.5-2 LIST AND DEFINE SYMBOLS
CALL PLOTV(125,939,38)
CALL PRINTV (18,18M ANALYTIC SOLUTION,145,939)
CALL PLOTV(125,919,55)
CALL PRINTV(19,39PRECESSING ELLIPSE (REFERENCE ORBIT FOR,137,919)935OH144
CALL PRINTV(41,41NUMERICAL RESIDUAL PERTURBATION SOLUTION),145,89536OH144
19)
CALL PLOTV(125,879,63)
CALL PRINTV (36,36M NUMERICAL SOLUTION OF TOTAL EQUATIONS,137,879)938OH144
CALL PRINTV (125,859,44)
CALL PRINTV (27,27M INITIAL OSCILLATING ELLIPSE,137,859)
CB.2.5-3 LABEL HORIZONTAL SCALE
CALL PRINTV ( 4, 4TIME, 475, 0 )
CB.2.5-4 LIST RUN PARAMETERS AND TIME SCALE
CALL PRINTV ( 17,17SEMI-MAJOR AXIS =, 132, 989 )
CALL LABLV (A1,268,989,7,1,1)
CALL PRINTV ( 14,14ECCENTRICITY =, 340,989)
CALL LABLV (EC1,452,989,7,1,0)
CALL PRINTV ( 23,23INCLINATION (DEGREES) =, 524,989)
CALL LABLV (DEU,718,989,7,1,3)
CALL PRINTV ( 12,12MASS RATIO =, 790,989)
CALL LABLV (EP4,886,989,-3,1,1)
CALL PRINTV(15,15TIME RANGE FROM,434,969)
CALL LABLV (FM,562,969,4,1,3)
CALL PRINTV ( 2, 2TIME RANGE TO,402,969)
CALL LABLV (EM,650,969,4,1,3)
CB.2.5-5 DETERMINE WHICH PLOT IS BEING DRAWN
GO TC 130,46,501,1
CB.2.6 LABEL AND PLOT ERROR IN JACOBI INTEGRAL, DELTA C/C
CB.2.6-1 LABEL VERTICAL SCALE AND TITLE
30 CALL PRINTV(13,13ABSOLUTE VALUE OF DELTA C / C, 0,800)
CALL PRINTV(53,53ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI
INTEGRAL, 302,1012)
CB.2.6-2 ADD ADDITIONAL SYMBOL DEFINITION AND PARAMETER VALUES
CALL PLOTV (125,839,42)
CALL PRINTV (19,19M JACOBI CONSTANT,C =, 132,969)
CALL LABLV (JACOBI,292,969, 7, 1,6)
CALL PRINTV (21,21M ANOMOLISTIC PERIOD =,718,969)
CALL LABLV (PSS,894,969, 7,1, 2)
CALL PRINTV(40,40NUMERICAL RESIDUAL PERTURBATION SOLUTION,145,8399510H144
1)
CB.2.6-3 PLOT DELTA C/C
DO9C1 (1,1)
IF(PTIME(1)-TF) 10,10,100
10 K =-12+139 I
90 CALL PLOTV ( 5,PTIME(1), APIX), 0,1,5,MKPT,(ERR)
CB.2.7 LABEL AND PLOT POSITION DIFFERENCES FROM NUMERICAL RESIDUAL
C. PERTURBATION VALUE, DELTA RHO/A
CB.2.7-1 CHAN GRID AND DO COMMON LABELING
100 L=7
GO TC 20
CB.2.7-2 LABEL VERTICAL SCALE AND TITLE
40 CALL PRINTV (0,-12, 13,13DELTA RHO / A,0,602)
CALL PRINTV (45,45NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL
SATURATION POSITION, 302,1012)

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\* 77

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M144 PLOT                                09/01/65
EXTERNAL FORMULA NUMBER - SOURCE STATEMENT - INTERNAL FORMULA NUMBER(S)

C8.2-7-3 PLOT DELTA RHO/A                9608H144 .81
DO 60 I = 11,111                        9610H144 .82
IF (PTIME(I) - TF)130,130 .110          9620H144 .83
130 K = 7 * 13 * I                       9630H144 .84
      (EPTIME - EN)160,60-110            9640H144 .85
60 CALL APLQTV (4,PTIME(I),APIK) ,0.1,4,MRKPT,IERR)
C8.2-8 LABEL AND PLOT VELOCITY DIFFERENCES FROM NUMERICAL RESIDUAL
C PERTURBATION VALUE, DELTA RHO DOT, SO RT A
C8.2-8-1 DRAW GRID AND DO COMMON LABELING
110 L=9
GO TC 20
C8.2-8-2 LABEL VERTICAL SCALE AND TITLE
50 CALL APRNTV (0,-12.27,27)DELTA RHO DOT TIMES SO RT A, 0, 620)
CALL PRINTV (76.76)NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL
RESIDUAL PERTURBATION SOLUTION, 258,1012)
C8.2-8-3 PLOT (DELTA RHO DOT * SO RT A) SUB A1, (DELTA RHO DOT * SO
C RT A) SUB PK1, (DELTA RHO DOT * SO RT A) SUB C1, AND (DELTA RHO DOT * SO
C RT A) SUB KI
DO 70 I = 11,111
IF (PTIME(I) - TF)120,120,80
70 K = 3 * 13 * I
IF (PTIME(I) - EN)170, 160,160
70 CALL APLQTV (4,PTIME(I),APIK),0.1,4,MRKPT,IERR)
II = III
GO TC 150
160 CALL APLQTV (4,PTIME(I),APIK),0.1,4,MRKPT,IERR)
140 II = I
C8.2-9 IF IIK IS LESS THAN 6, SET IIK = IIK + 1 AND RETURN TO STEP
8.2.3-1 OTHERWISE GO TO STEP 8.2.10
150 CONTINUE
C8.2.10 RETURN TO MAIN PROGRAM
80 RETURN
END

```

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9608H144 .81
9610H144 .82
9620H144 .83
9630H144 .84
9640H144 .85
9650H144
9660H144
9665H144
9670H144
9680H144
9680H144 .86
9680H144 .88
9690H144
9700H144 .89
9710H144
9715H144
9716H144
9720H144
9724H144 .91
9730H144 .92
9740H144 .93
9750H144 .94
9760H144 .95
9770H144 .96
9780H144 .98
9790H144 .99
9800H144
9810H144
9820H144
9825H144
9830H144
9840H144
9850H144
9860H144

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.87
.87
.97
.100
.101
.102
.103
.104
.105

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Section 9  
INPUT AND OUTPUT

9.1 INPUT

The input to this program is by the Douglas subroutine INPUT 1. A load sheet for this program is shown as figure 2. The card format is:

column 1 - a one  
columns 2 through 6 - location number of piece of data  
columns 7 through 15 - input number  
columns 16 and 17 - location of decimal place from beginning of field  
positive if to the right

Three other pieces of data may be entered on the card. The location numbers are punched in columns 18-22, 34-38, and 50-54. The data are punched in columns 23-31, 39-47, and 55-63. The exponents, as explained above, are punched in columns 32-33, 48-49, and 64-65, respectively. The remaining information required is:

columns 66-68, zeros  
columns 69-70, reference run number  
columns 71-73, case number.

This routine allows identification on the card of each piece of input data by relative location number; only non-zero numbers need be entered. It has a "Reference Run," "Case" setup. If the case number (card columns 71 to 73) is non-zero, but the reference run number (card columns 69, 70) is zero, then the data on the load sheet are assumed to be sufficient and the case is computed. If the case number is zero and the reference run number is non-zero, the data are stored in array WR and no case is attempted. If the following load sheets with non-zero case numbers have also the reference run number of the stored array, then a case is run using the input of array WR as modified by the new load sheet.



The order of stacking cases is then:

1. All cases with zero reference run number
2. First reference run (zero case number)
3. All cases with first reference run number and non-zero case number
4. Second reference run (zero case number)

It is one of the peculiarities of this routine that cases must be run in numerical order.

## 9.2 OUTPUT

The first information that is printed each time the program is run consists of an explanation of the error codes. The program makes certain checks on the input to make sure that the input is reasonable and consistent with the theory being used. Throughout the program there are additional checks made to ascertain if division by zero has occurred which would cause the program to produce nonsense if the case continued. In these cases, an error code is printed and the case terminated. The print is as follows:

ERROR CODE	REASON FOR HALT
1	SEMI-MAJOR AXIS = 0
2	ECCENTRICITY EQUALS OR EXCEEDS 1
7	S (THE APPROX RADIUS VECTOR) = 0
8	DERIVATIVE OF TIME WRT PSI = 0 (i.e., $\frac{dt}{d\psi}$ )

The error code is set to 1 in Section 4.3.2.1.1. The error code is set to 2 in Section 4.3.2.2.1. The error code is set to 7 in Section 4.17.3. The error code is set to 8 in Section 4.17.8.1. The error code is printed in Section 4.3.2.1.2.



The program then reads in a case and INPUT 1 prints:

REFERENCE RUN NO. (RR#) CASE NO. (CASE #)

There follows a floating point print of the input array W and of constants computed by the program as described in Section 4.4.1.

Section 4.6 then prints the headings for the print during the numerical residual perturbation solution as follows:

THE INITIAL CONDITIONS ARE	SEMI-MAJOR AXIS	a	INCLINATION	i
	ECCENTRICITY	e	MASS RATIO	$\mu$
TIME	DELTA X	DELTA X DOT	DELTA X DBL DOT	RHO
	DELTA Y	DELTA Y DOT	DELTA Y DBL DOT	RHO DOT
C APPROX	DELTA Z	DELTA Z DOT	DELTA Z DBL DOT	
	X X DOT	X APPROX	X DOT APPROX	
	Y Y DOT	Y APPROX	Y DOT APPROX	
C EXACT	Z Z DOT	Z APPROX	Z DOT APPROX	

The algebraic interpretation of these headings is given in Section 4.10.

The initial value of the Jacobi constant,  $C_{init}$ , is printed in Section 4.8.6 as:

THE JACOBI INTEGRAL CONSTANT IS ( $C_{init}$ )

During computation of the numerical residual perturbation solution, the printing according to the above heading is accomplished in Section 4.10. If the input quantity  $W_{13} > 0$ , then printing takes place about every  $1/4$  radian of  $\psi$ . Otherwise printing takes place at the first and last point only.

The headers for the identification of the print during the numerical solution of the total equations of motion are printed in Section 4.19.4 and are as follows:

## NUMERICAL SOLUTION OF TOTAL EQUATIONS OF MOTION DATA

TIME	X	X DOT	X DBL DOT	RHO
	Y	Y DOT	Y DBL DOT	RHO DOT
JACOBI CONSTANT,	Z	Z DOT	Z DBL DOT	

The printing for the numerical solution of the total equations of motion is done in Section 4.20 and the explanation of the above headers is shown in that section. Additional print for which there are no headers is also shown there. Printing during the solution of the total equations of motion has the same rules with regard to the input quantity  $W_{13}$  as does the print during the numerical residual perturbation solution.

The plotted output is described in subroutine PLOT, Section 8.

SAMPLE PRINTED OUTPUT

ERROR CODE	REASON FOR HALT
1	SEMI-MAJOR AXIS = 0
2	ECCENTRICITY EQUALS OR EXCEEDS 1
7	S (THE APPROX RADIUS VECTOR) = 0
8	DERIVATIVE OF TIME MRT PSI = 0

REFERENCE RUN NO.	3	CASE NO.	90
0.1478000E 01	0.		0.9000000E 00
0.2660000E-04	0.9000000E 03	0.0999999E-01	0.3001999E-05
0.3000000E 03	0.0999999E-08	0.	0.1500000E 01
			0.
-0.1883389E 01	-0.1447111E 01	-0.1347638E 01	0.9294126E 00
0.	0.	0.	0.8504231E 00
0.9492679E-05	0.9999700E 00	0.1796849E 01	-0.0000000E-19
			-0.2334933E-02
0.4679927E-02	-0.2334933E-02	-0.1091719E-04	0.2731960E-04
			0.8588231E 00
			0.0999999E 01





Section 10  
DISCUSSION OF RESULTS

A summary of numerical results is included in graphical plot form in the Appendix. For the numerical studies, the mass ratios were  $\mu = 3 \times 10^{-6}$ , which corresponds to the sun-earth (or quite closely to the sun-Venus) systems and  $\mu = 0.0123$ , which corresponds to the earth-moon system. In reference 1 it was stated that the solutions should be valid for a region of order  $\mu^{1/2}$  around the origin in the natural variables, i.e., semi-major axes of order unity in the blown-up variable could be tolerated. This would represent the limit to which the series could be taken. Further approximation is involved in that the actual series are only computed to order  $\epsilon^2$ . All  $\epsilon^2 e$  terms were dropped in this program for consistency since the switch-back terms from  $\epsilon^3$  (presumed to be of order  $\epsilon^2 e$ ) and terms of forms  $\epsilon e i$  and  $\epsilon^3$  are not included in reference 1. For the earth-moon mass ratio, semi-major axes of 0.041 (corresponding to a near grazing lunar satellite), 0.1, and 0.4 were used. For the earth-sun mass ratio, the semi-major axes were chosen to be 0.15 (corresponding to about the synchronous altitude), 0.5, and 1.478 (corresponding to the moon). In reference 1 the values of eccentricity and inclination that could be tolerated consistently with the accuracy with respect to  $\mu$  were not specified explicitly, but it was merely stated that these elements must be small. We have used eccentricities of 0 and 0.05 and inclinations of  $0^\circ$  and  $5^\circ$ . The plots show the time history of the error in the Jacobi constant, the vector distance from the numerical residual perturbation value, and the vector velocity difference from the numerical residual perturbation value. The summary of the cases computed is given in Table I.

The error in the Jacobi constant is represented as the variation with time along the solution, divided by the initial value. It is presented for the analytic solution of reference 1, for a precessing mean Kepler ellipse which has essentially the same precession rate as the analytic solution but no oscillatory perturbations, for a numerical integration of the total equations of motion, for the fixed initial osculating ellipse, and for the numerical residual perturbation solution. The precessing ellipse is the reference orbit of the numerical residual perturbation solution. The fixed initial osculating

ellipse is provided for comparison since it would form the basis of a normal Encke's method solution. The variation of the Jacobi constant should, of course, be zero in the exact solution of the restricted three-body problem. The other two variables were chosen to represent position and velocity errors, and were referenced to the numerical residual perturbation solution since it was anticipated that this solution would, in general, maintain the Jacobi integral most accurately. For each set of parameters, the progress of these results was monitored for approximately 28 revolutions. This limit is set by the program which, by the way it is coded, imposes a storage limit at this level. Since it was desired to show both the short- and long-term performance of the technique, an expanded time scale was used. This causes a plot for 28 revolutions to extend over five to six pages. For the 24-hour circular earth orbit and the moon's orbit, the entire time span is exhibited. For the other cases, only the beginning and ending pages are included.

It will be seen from the error in Jacobi integral plots that for the 24-hour circular earth orbit (or 13,000-mile Venus orbit) and the grazing lunar orbit, the numerical residual perturbation solution approximates the sun's perturbations no better than the analytic solution, which is already at the round off level of the machine throughout the 28 orbits. It will also be noted that for these cases, the initial osculating ellipse maintains the Jacobi integral better than the particular numerical solution of the total equations of motion used. It will be noted that in most cases the initial osculating ellipse holds the Jacobi integral better than the precessing mean-ellipse. However, the position and velocity errors of the initial osculating ellipse are greater than (sometimes orders of magnitude greater) or equal to the precessing ellipse errors. This may be explained as follows. For motion of a close satellite around the smaller primary, the Jacobi integral is sensitive to errors in radius and its derivative. It is almost totally insensitive to errors along the direction of motion. The  $s$ , or one-over-radius, expression in the analytic solution contains a constant term in the perturbations of order  $\epsilon^2$  showing that the radius vector oscillates about a value which is not the Keplerian value of the precessing ellipse but is actually less. The  $dt/d\psi$  expression contains a perturbation term whose value is not zero at



initial time. This term is used in the computation of the elements of the osculating ellipse but not in the computation of the elements of the mean ellipse. The result is that the period of the osculating ellipse is in error, but the perigee is lowered, partially compensating for the neglected term in  $s$ . It is thus better in radius but poorer in angle.

At the largest semi-major axes considered (the case of the moon about the earth, and the 10 lunar radii lunar satellite), the analytic solution is very poor. This would be anticipated since de Pontecoulant demonstrated in the last century that this series was, at best, only barely convergent for the moon and only second order terms have been carried in this study. Over half of the precession of perigee and over half the "evection" were lost when the higher order terms were dropped (ref 3, pages 322 and 326). At these semi-major axes, the advantage of the numerical residual perturbation solution without rectification over the numerical solution of the total equations of motion in holding the Jacobi integral drops to a factor of 6 to 7. However, the sharp difference in slope between the two values of error in the Jacobi constant for small times suggests that with rectification, the numerical residual perturbation solution will be much better than the numerical solution of the total equations of motion. But since the  $\Delta p/a$  and  $\Delta \dot{p}\sqrt{a}$  plots show that the initial osculating ellipse is nearly as good as the precessing mean ellipse at small times, the advantage of the numerical residual perturbation solution over a numerical perturbation solution is not so clear. This suggests that when frequent rectification is required, as it seems to be with the large orbits, the use of a precessing ellipse whose initial orbital parameters are equal to those of an initial osculating ellipse might be better than using a precessing mean ellipse.

Comparing the results for  $e = 0.05$  with the results for  $e = 0$ , it seems that  $e = 0.05$  is a large  $e$  for this expansion. However, this program was constructed merely for demonstration of the residual perturbation method and is not deemed to be a practical analysis tool. In reference 6, the theory is developed without the restriction of small eccentricity or small inclination. The accuracy shown here for the case of zero inclination

and zero eccentricity then would be the only one applicable to a program based on that theory, and it would be expected that it would hold to the same order of accuracy at other eccentricities and inclinations. From the cases with zero-eccentricity, it appears that the assumption that the precessing ellipse would be nearly as useful a reference orbit as the analytic solution is not a good one. In hind sight it is seen that this would follow from the fact that in an accumulation, as is done in the numerical solution of differential equations, the decimal place where error accumulates is proportional to the maximum value of the solution. Thus, when an oscillation is being accumulated numerically, the error accumulates at a rate roughly proportional to the full magnitude of the oscillation. If the method of reference 6 were programmed, then it is anticipated that the additional algebra required to algebraically cancel all  $\epsilon$  terms and  $\epsilon^2$  in the residual perturbation equations of motion would result in appreciably improved numerical results.

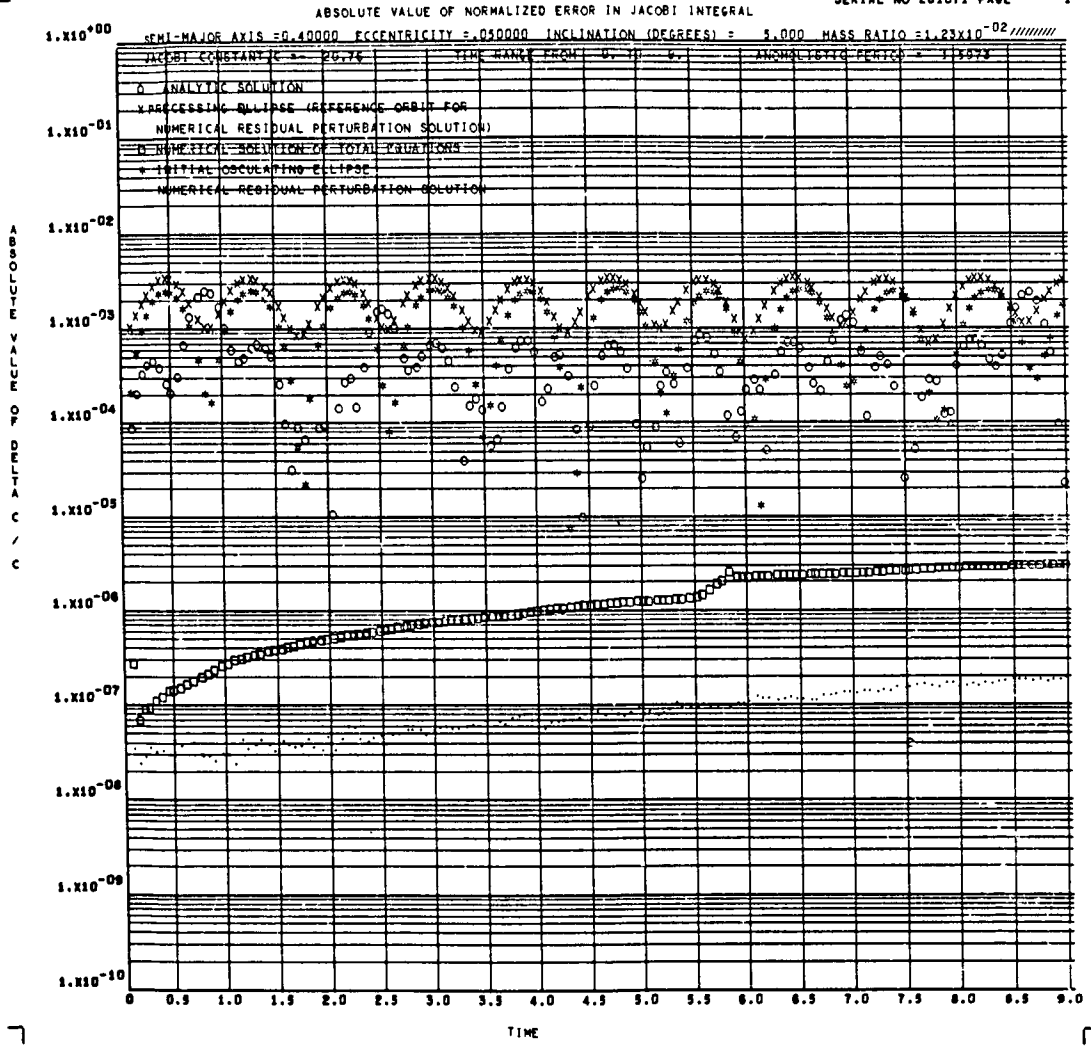
Table I. - SUMMARY OF COMPUTED CASES

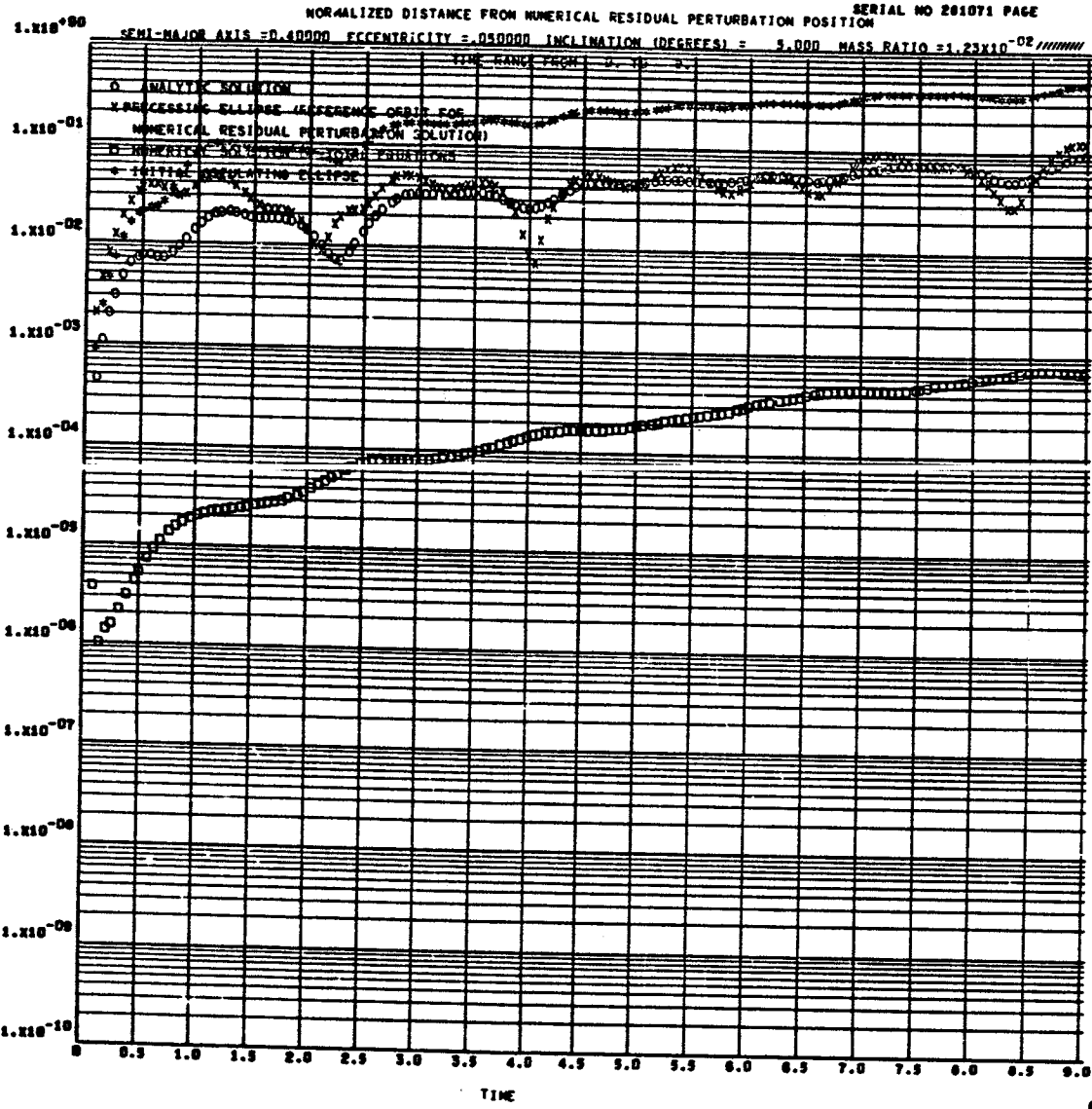
Mass Ratio $\mu$	Semi-major Axis $a$	Eccentricity $e$	Inclination $i^\circ$	Notes	Frame Identification	Serial Number	Page Numbers
$3 \times 10^{-6}$	0.15	0	0	Synchronous earth or 13,000-mile Venus satellite	281071	281071	145-162
	0.15	0.05	5		281071	281071	127-129, 139-144
	0.5	0.05	5		281071	281071	109-111, 124-126
	1.478	0	0		302074	302074	1-3, 13-18
$1.23 \times 10^{-2}$	1.478	0.0549	5.15	Moon's orbit or 120,000-mile Venus satellite	281071	281071	91-108
	0.041	0	0	Grazing lunar satellite	281071	281071	55-57, 67-69
	0.041	0.05	5		281071	281071	37-39, 49-51
	0.1	0.05	5		281071	281071	19-21, 34-36
	0.4	0.05	5	10 lunar radii lunar satellite	281071	281071	1-3, 13-15

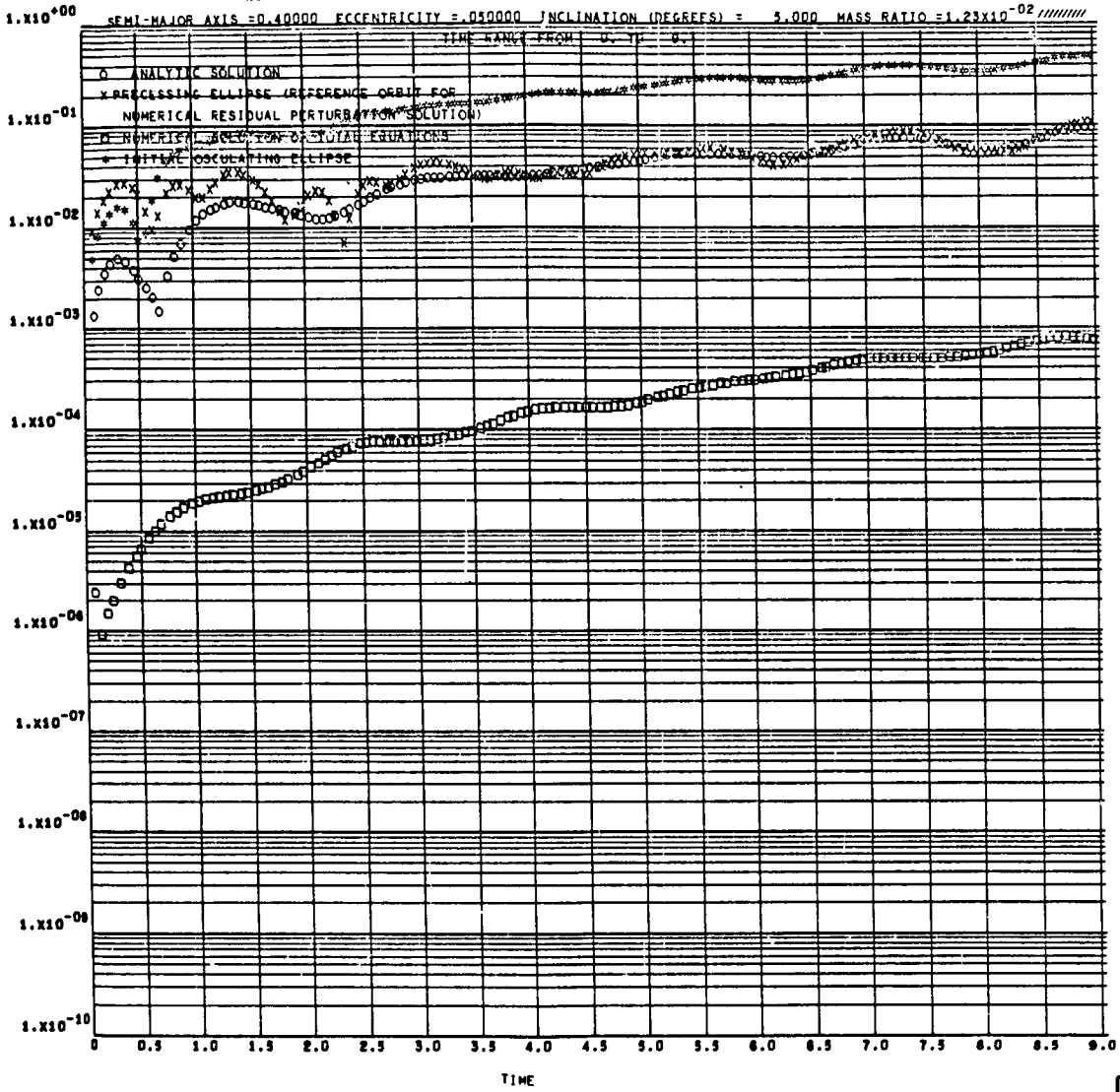
## REFERENCES

1. Kevorkian, J.: A Uniformly Valid Asymptotic Representation for All Times of a Satellite in the Vicinity of the Smaller Body in the Restricted Three-Body Problem. *Astron. J.*, vol. 67, no. 4, no. 1299, May 1962 (Revised version of reference 7.)
2. Moulton, Forest Ray: *An Introduction to Celestial Mechanics*. The Macmillian Company, 1914.
3. Brouwer, Dirk; and Clemence, Gerald M.: *Methods of Celestial Mechanics*. Academic Press, 1961.
4. Herget, Paul: *The Computation of Orbits*. 1948.
5. Hildebrand, F. B.: *Introduction to Numerical Analysis*. McGraw-Hill Book Co., Inc., 1956.
6. Eckstein, M.; Shi, Y; and Kevorkian, J.: Satellite Motion for Arbitrary Eccentricity and Inclination Around the Smaller Body in the Restricted Three-Body Problem. *Douglas Paper No. 3489*, April 1965.
7. Kevorkian, J.: A Uniformly Valid Asymptotic Representation for All Times of a Satellite in the Vicinity of the Smaller Body in the Restricted Three-Body Problem. *Douglas Paper No. 1163*, May 1961.
8. Anon.: *North American Aviation, Inc., Engineers Computing Manual, Region 74, S-C 4020 Subprograms*, Oct. 1963.

**APPENDIX**  
**SAMPLE GRAPHICAL RESULTS**

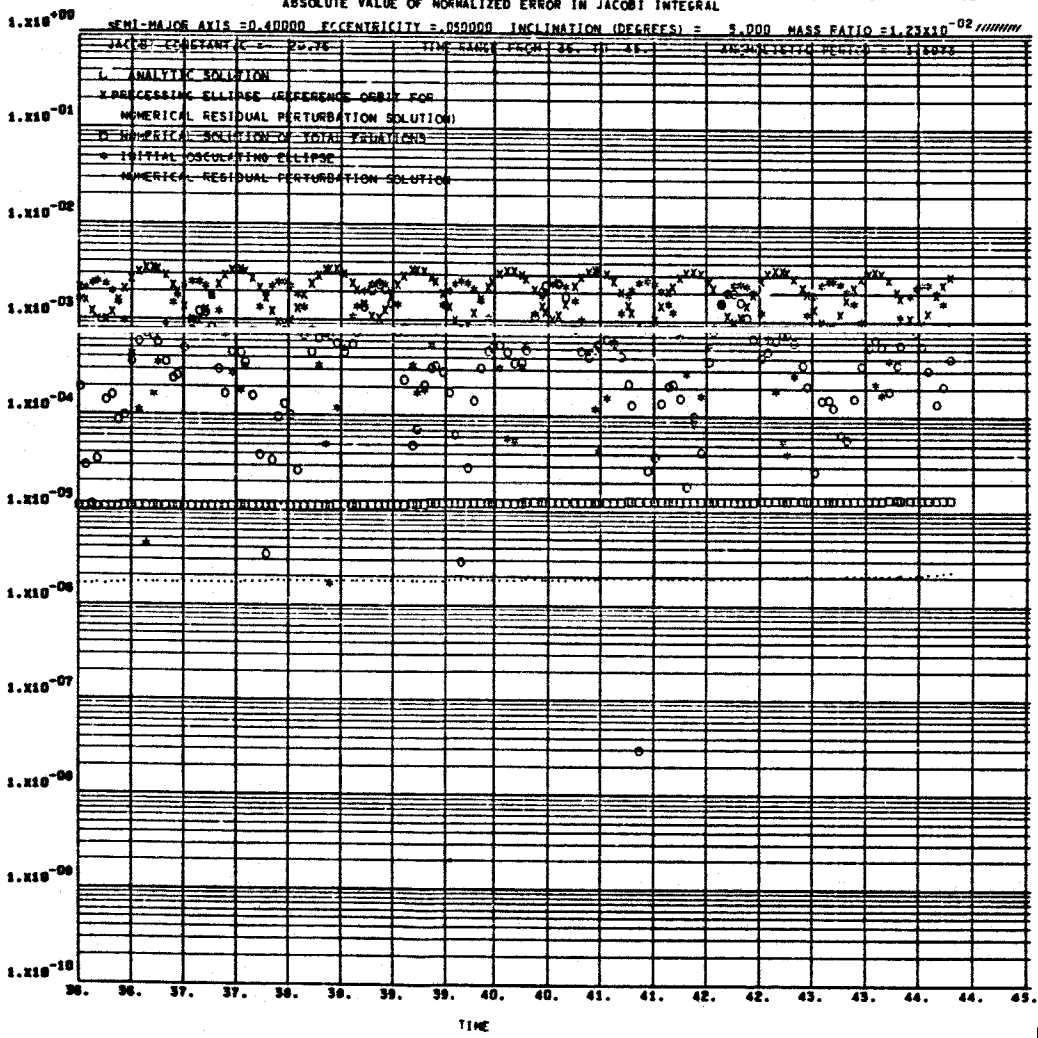


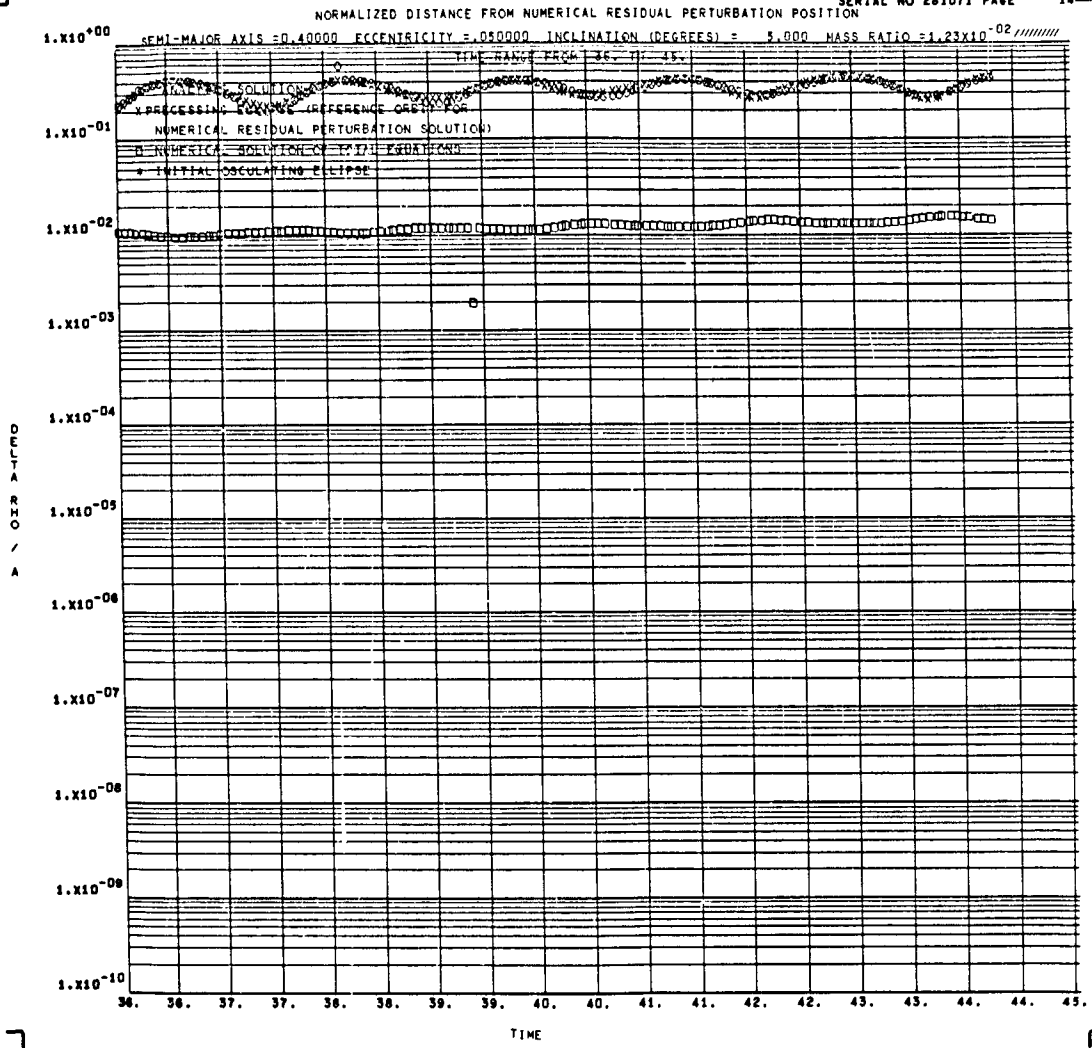


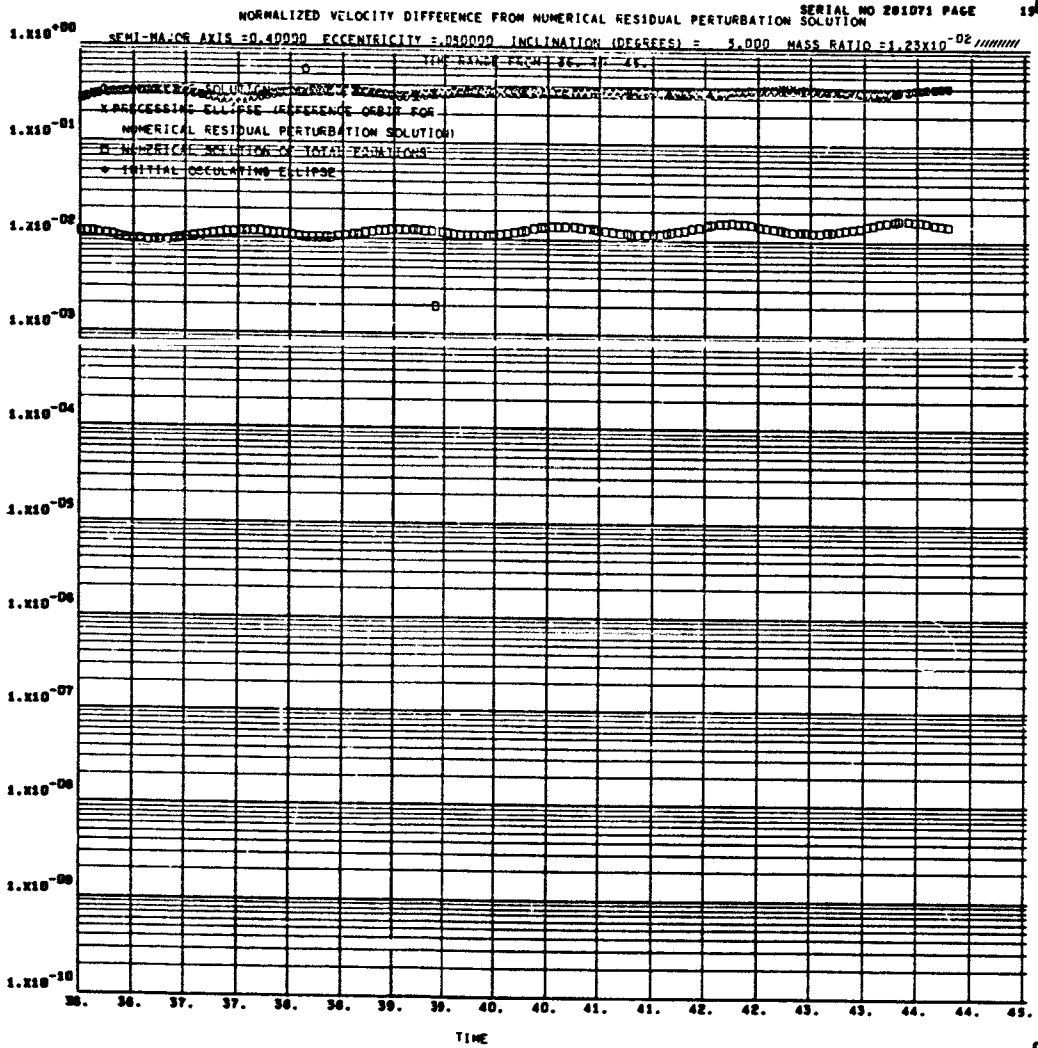


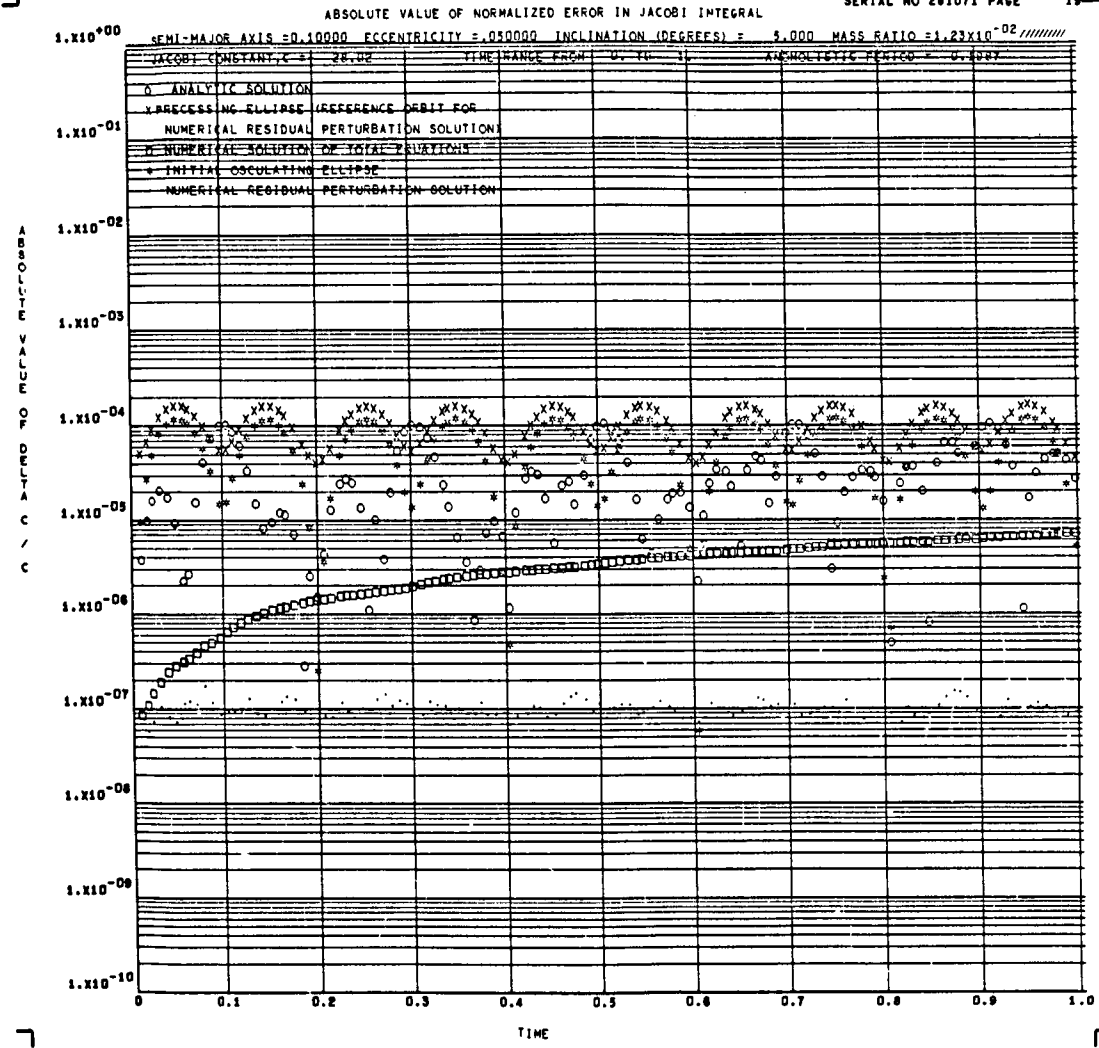


ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

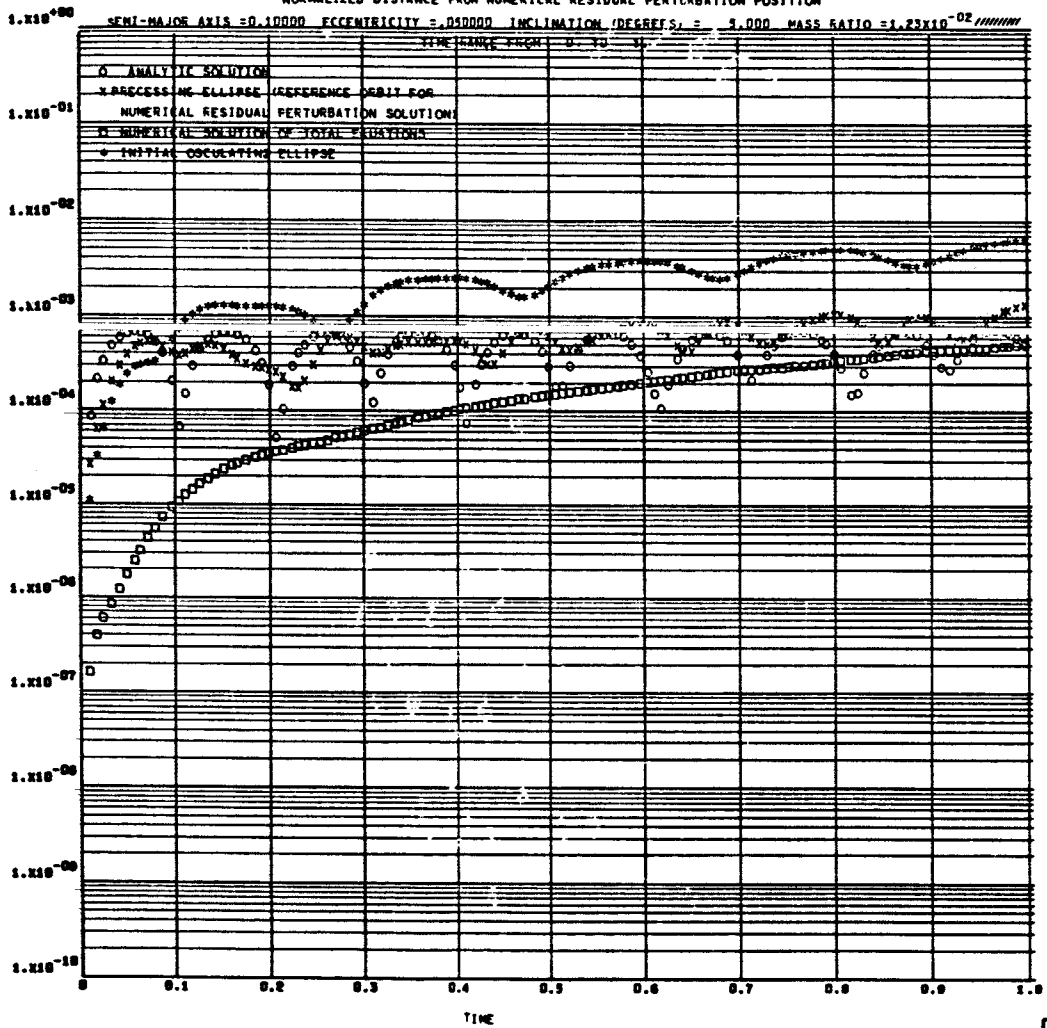




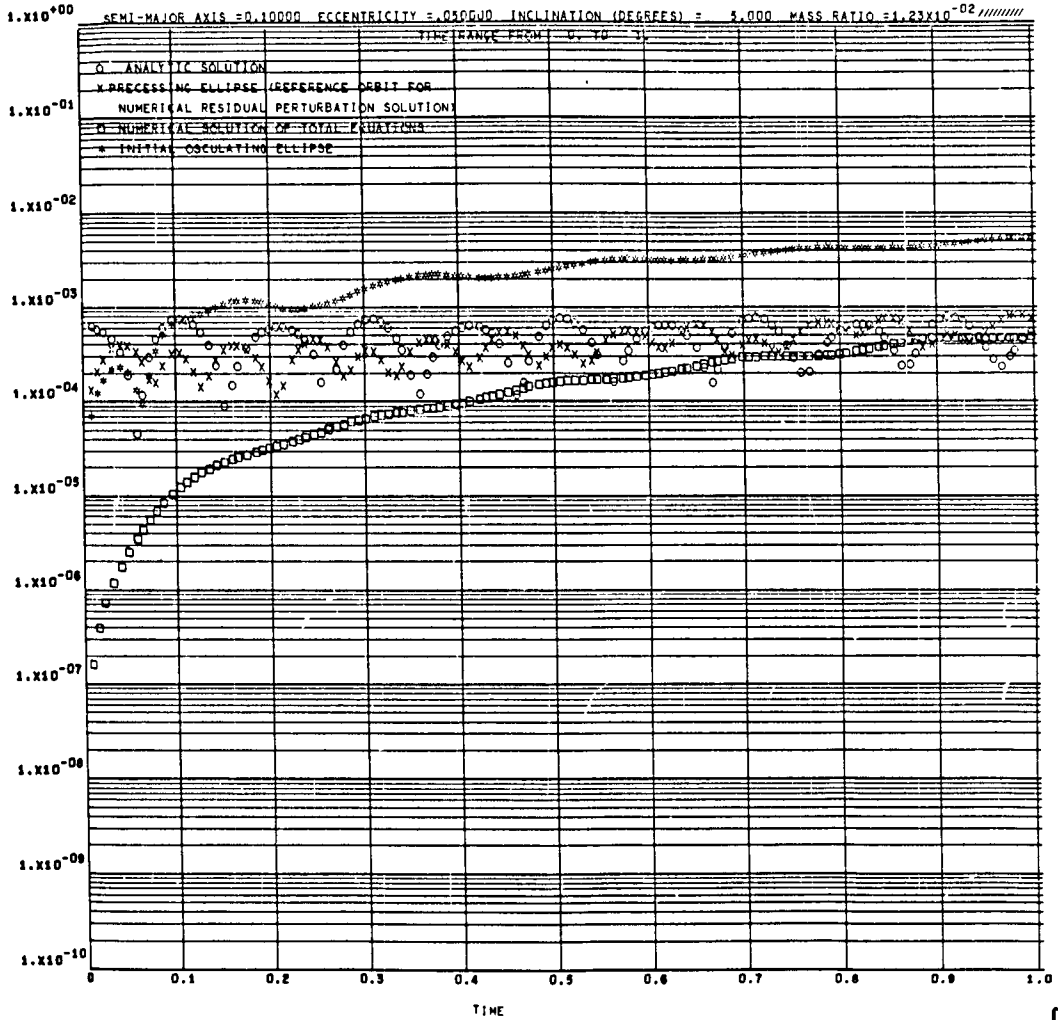




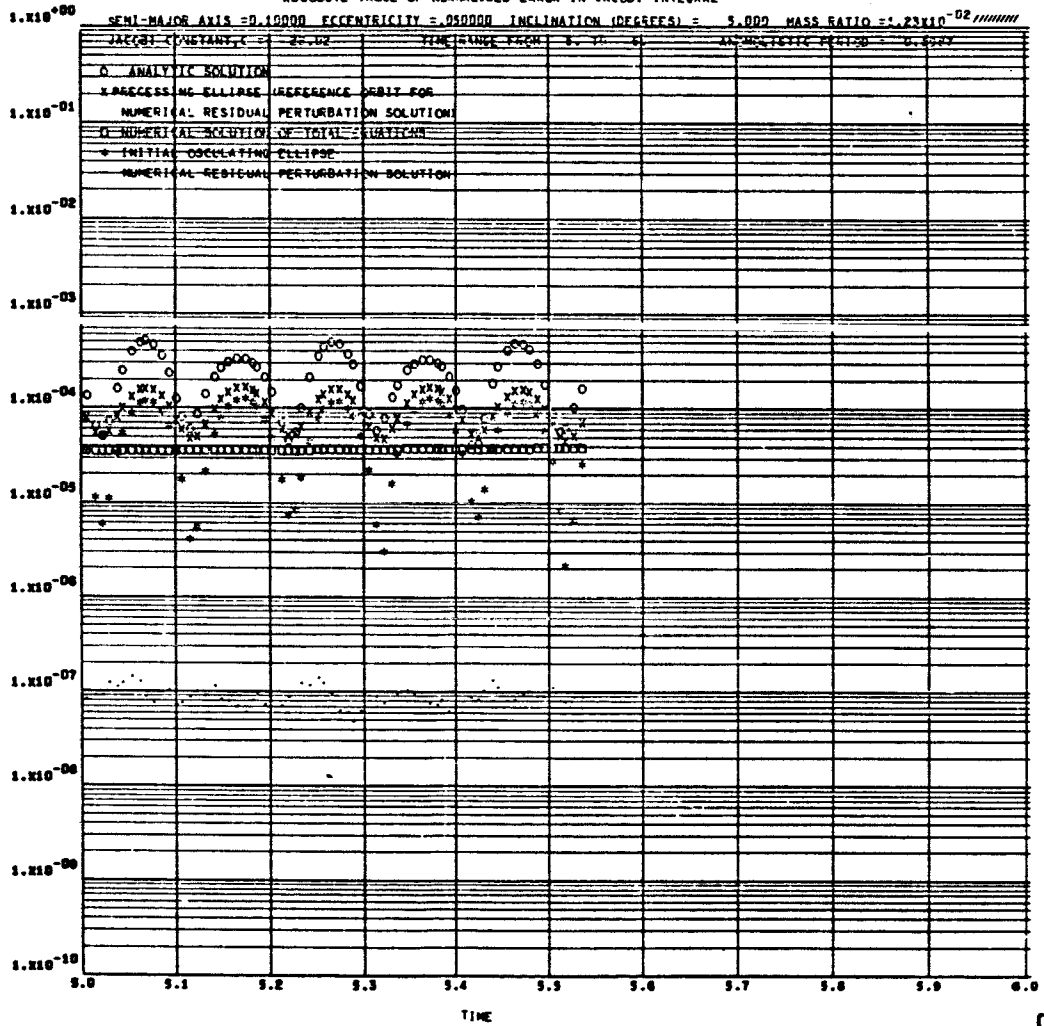
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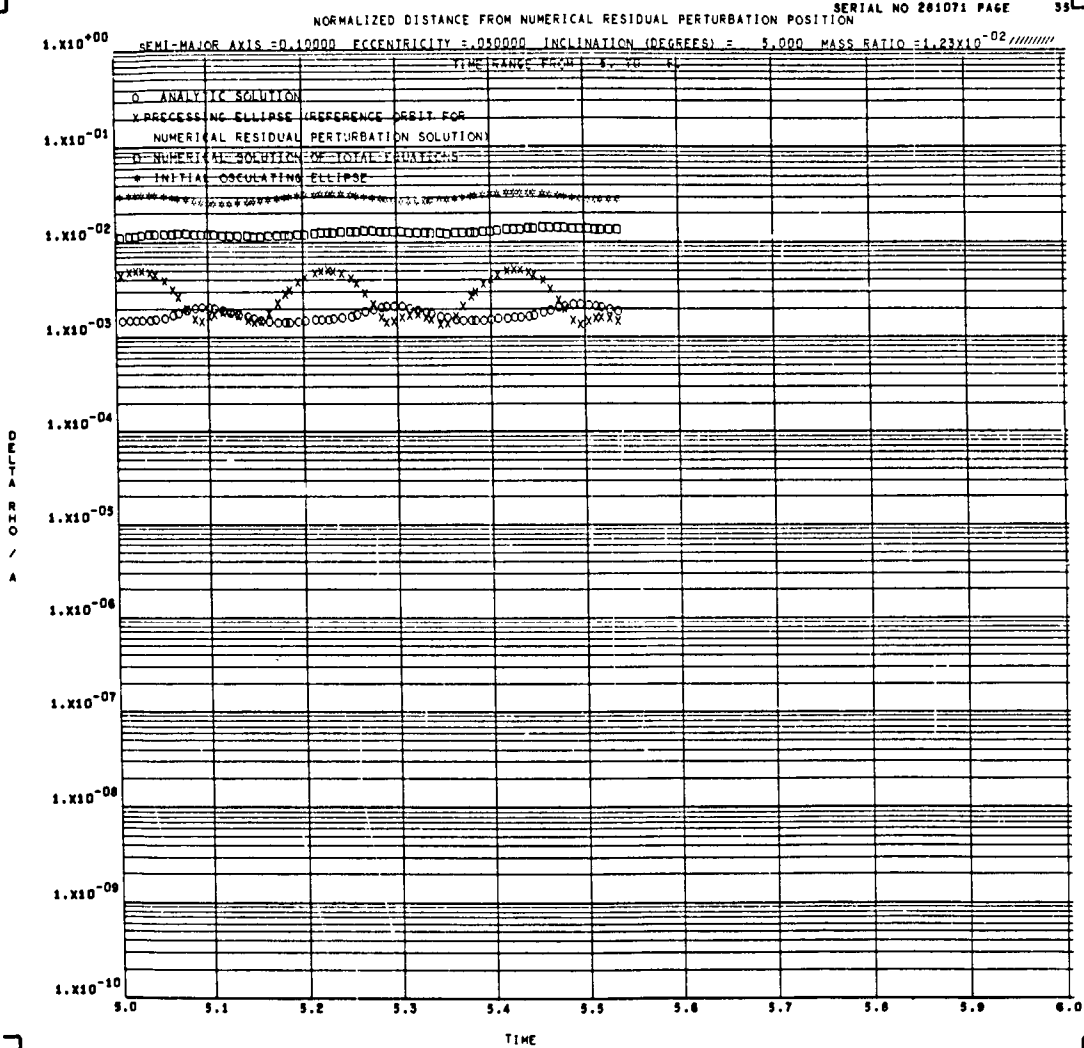


NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



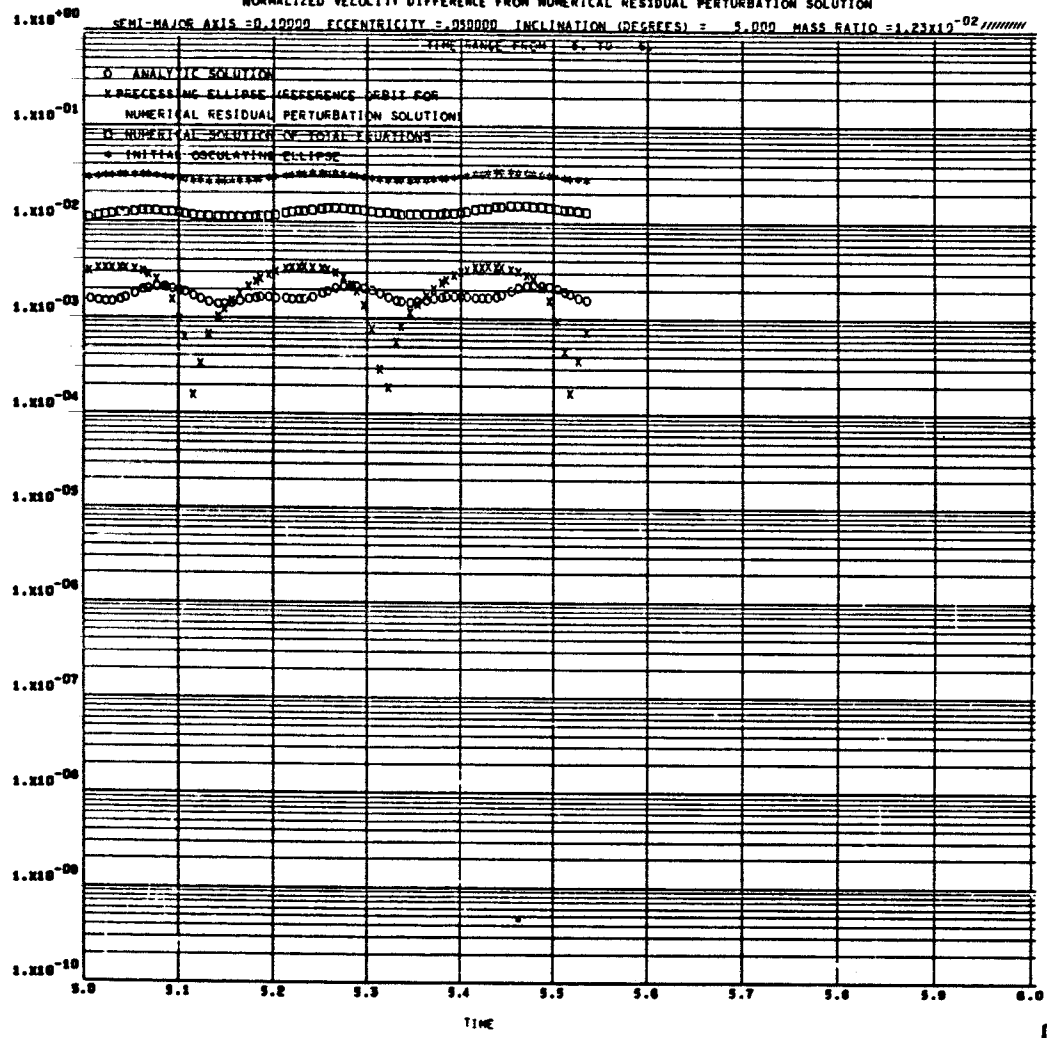
ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL





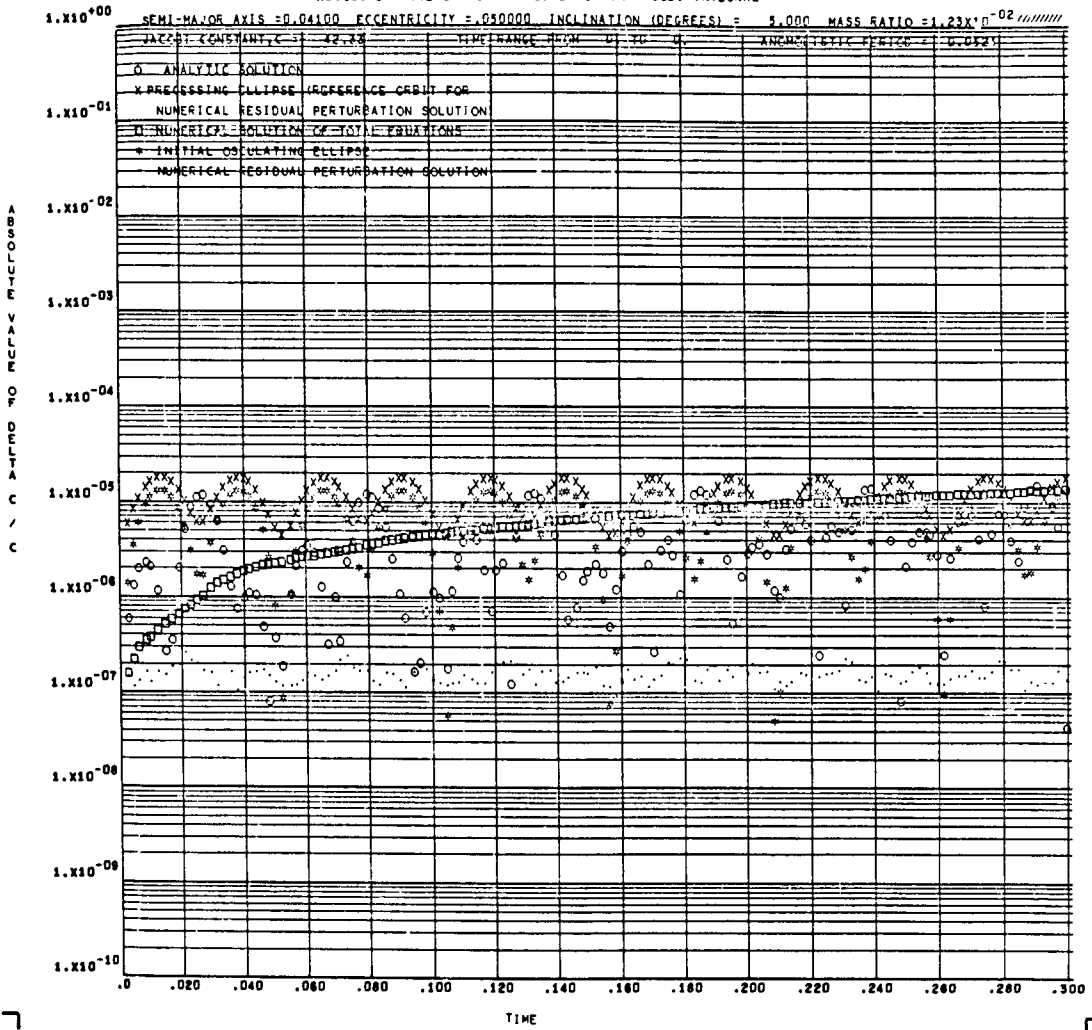


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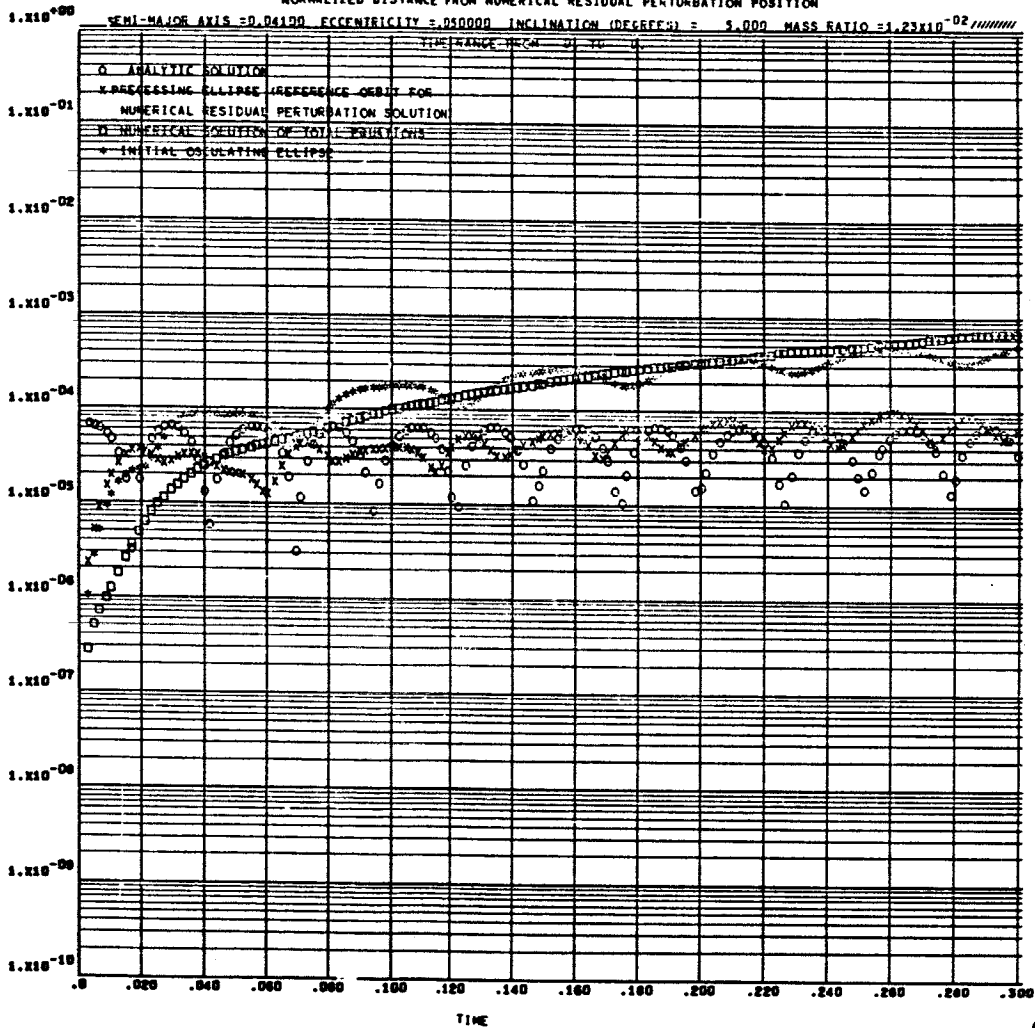


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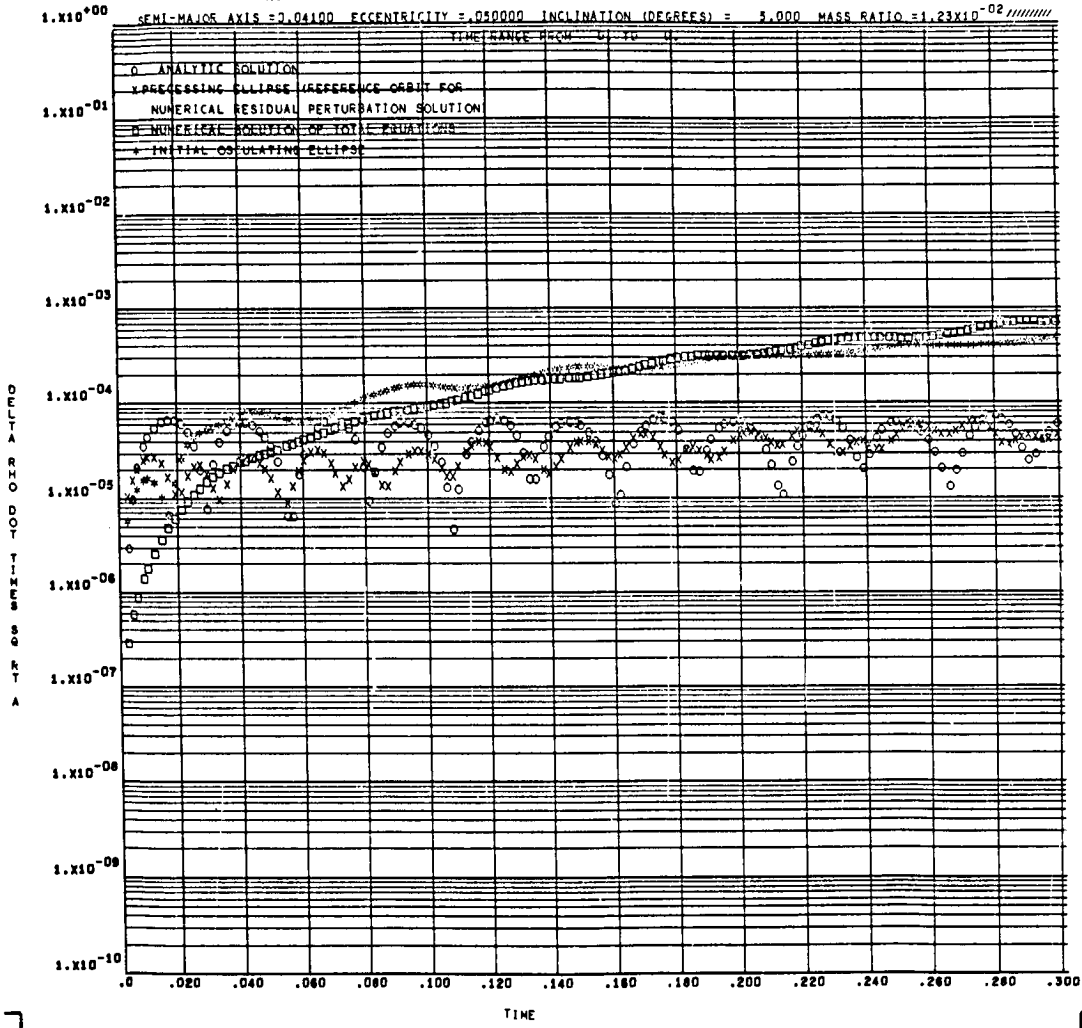
SERIAL NO 281071 PAGE 37



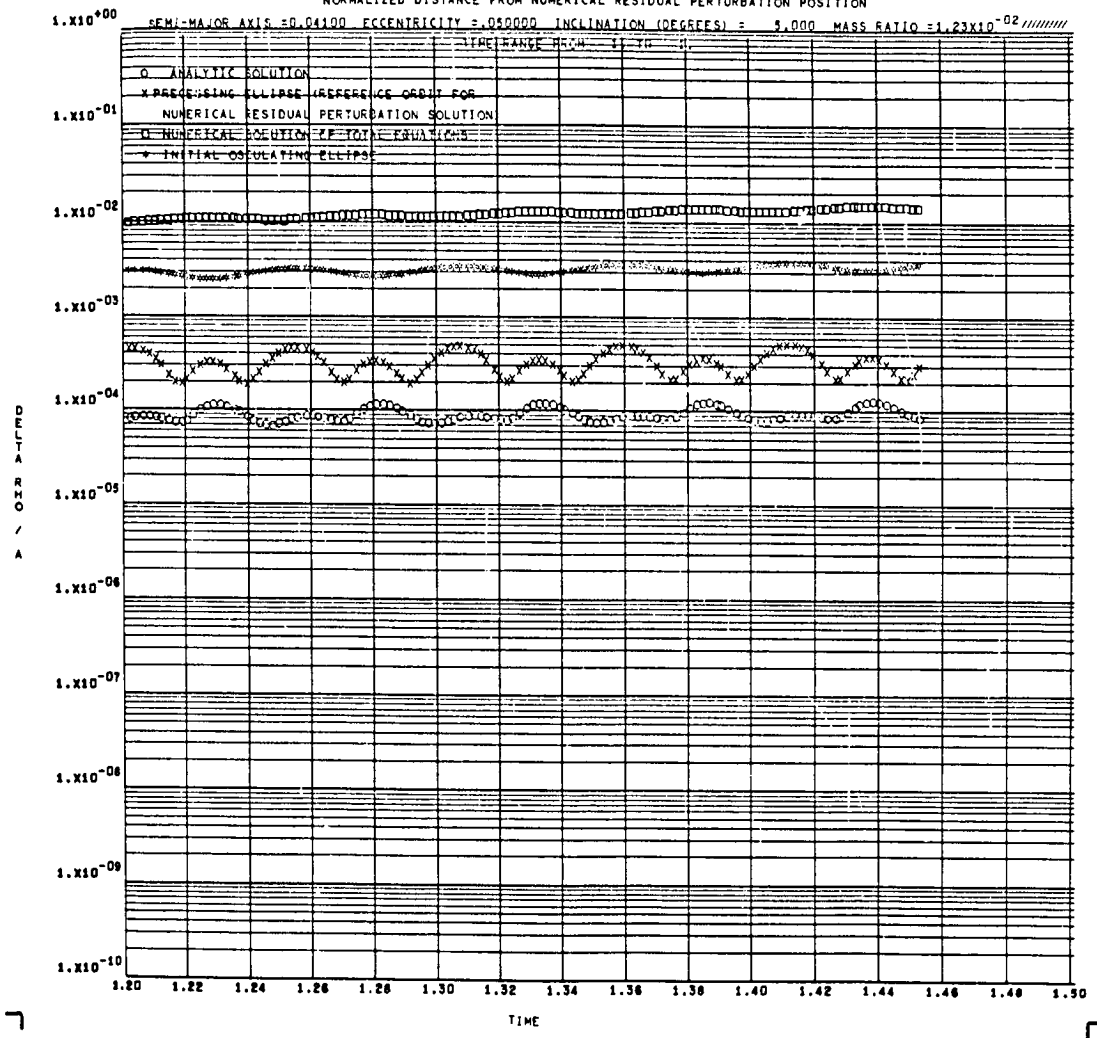
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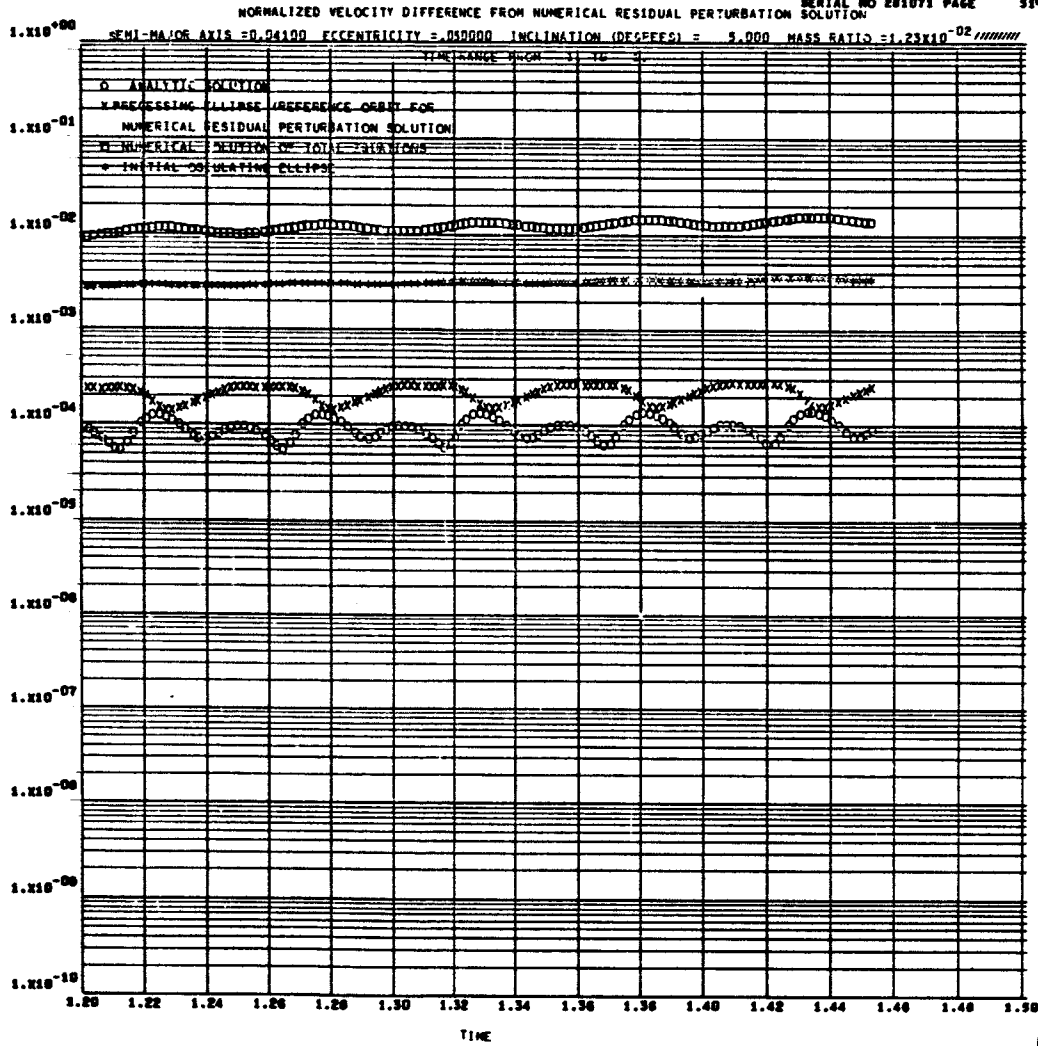


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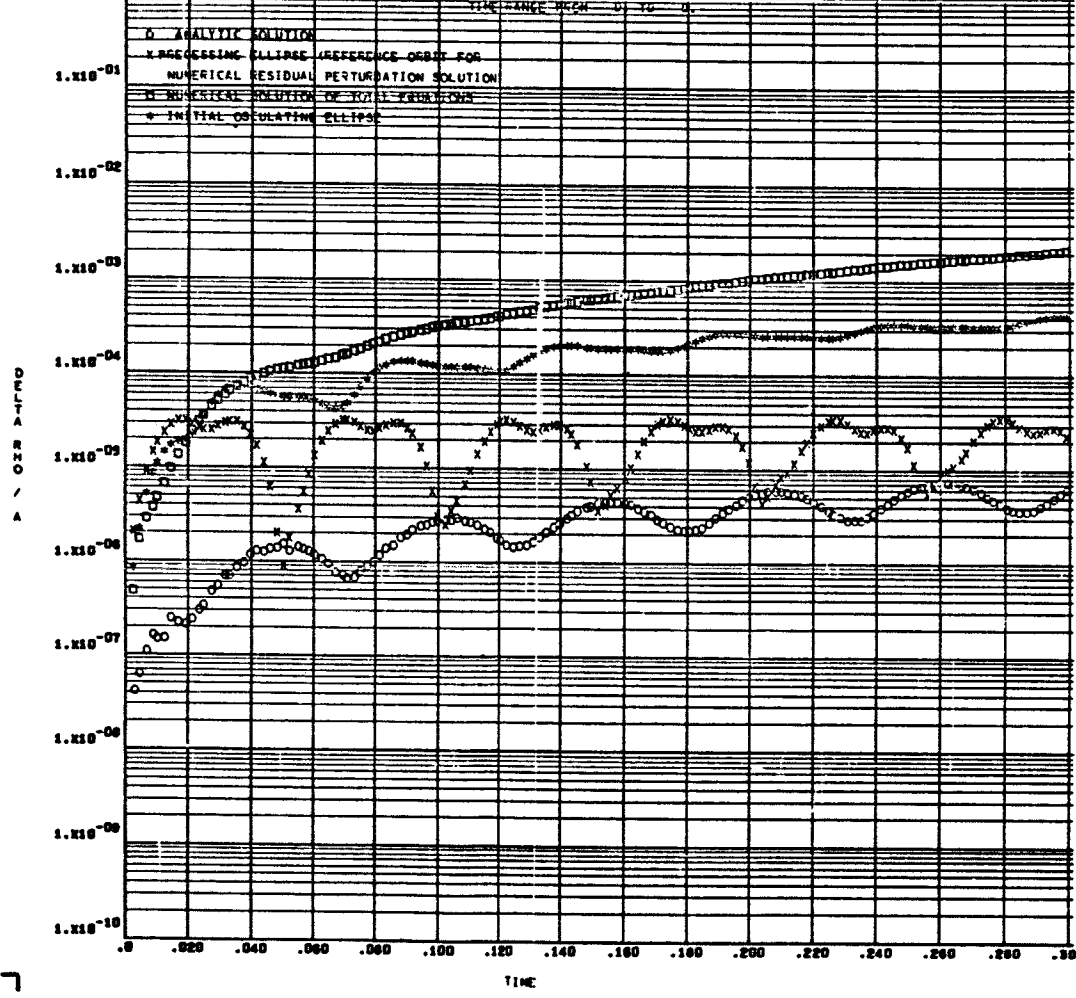


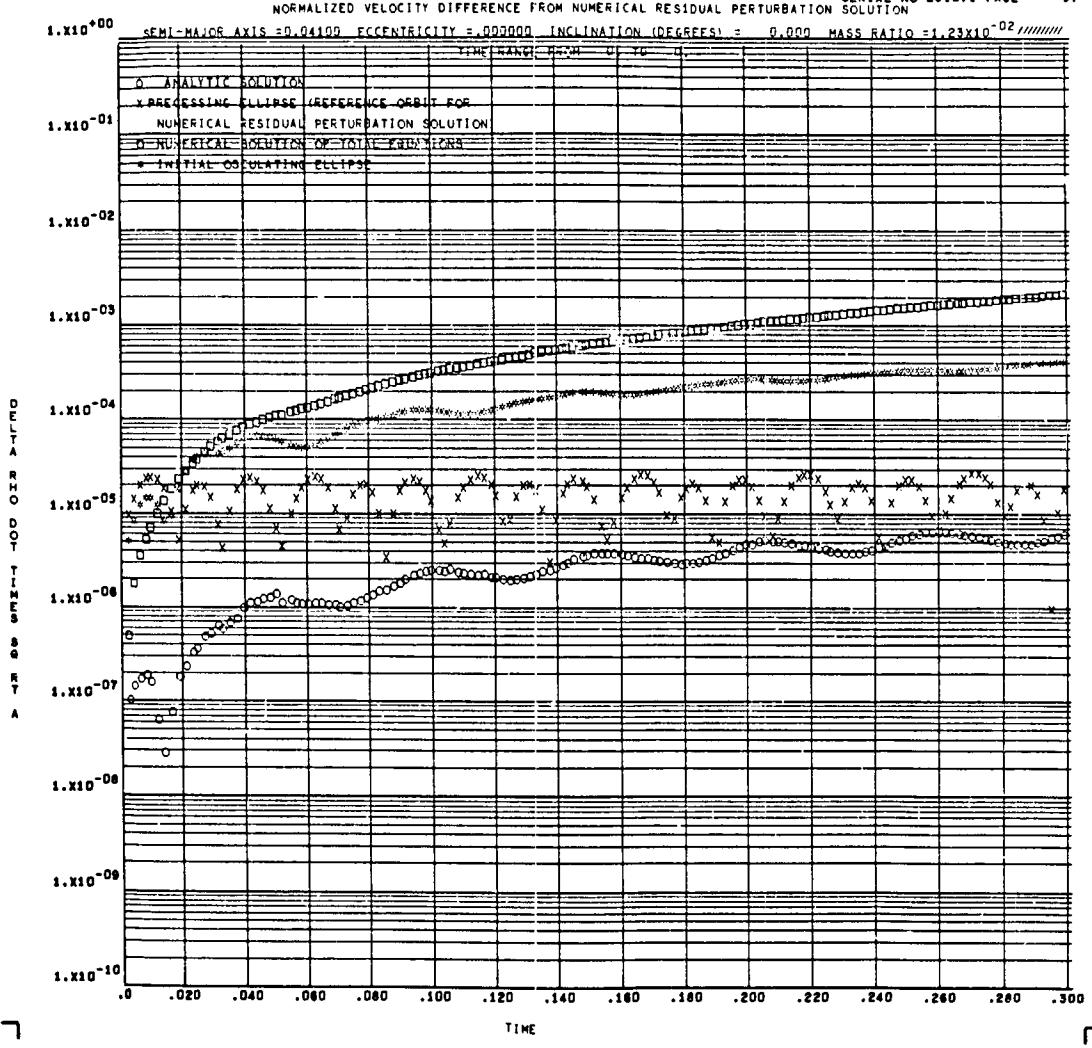




NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

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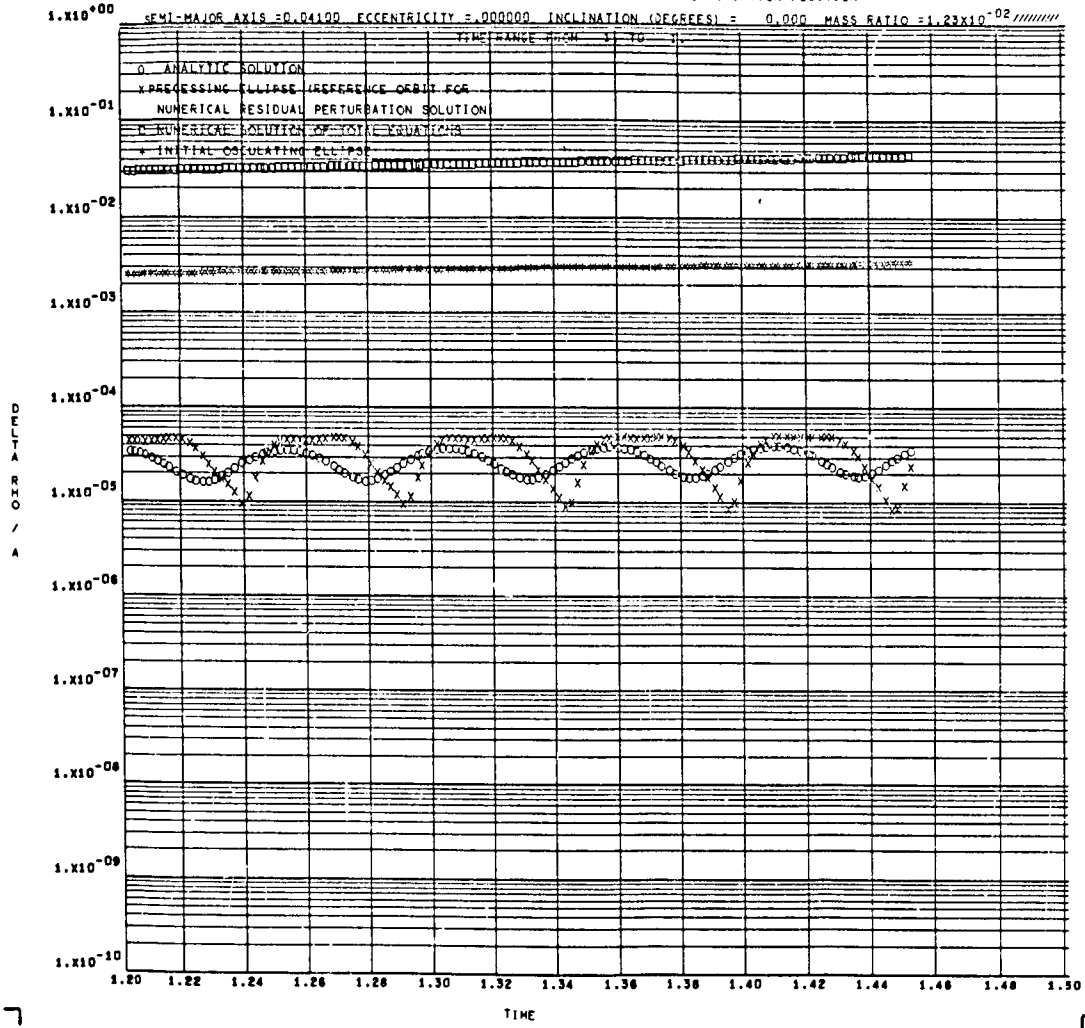




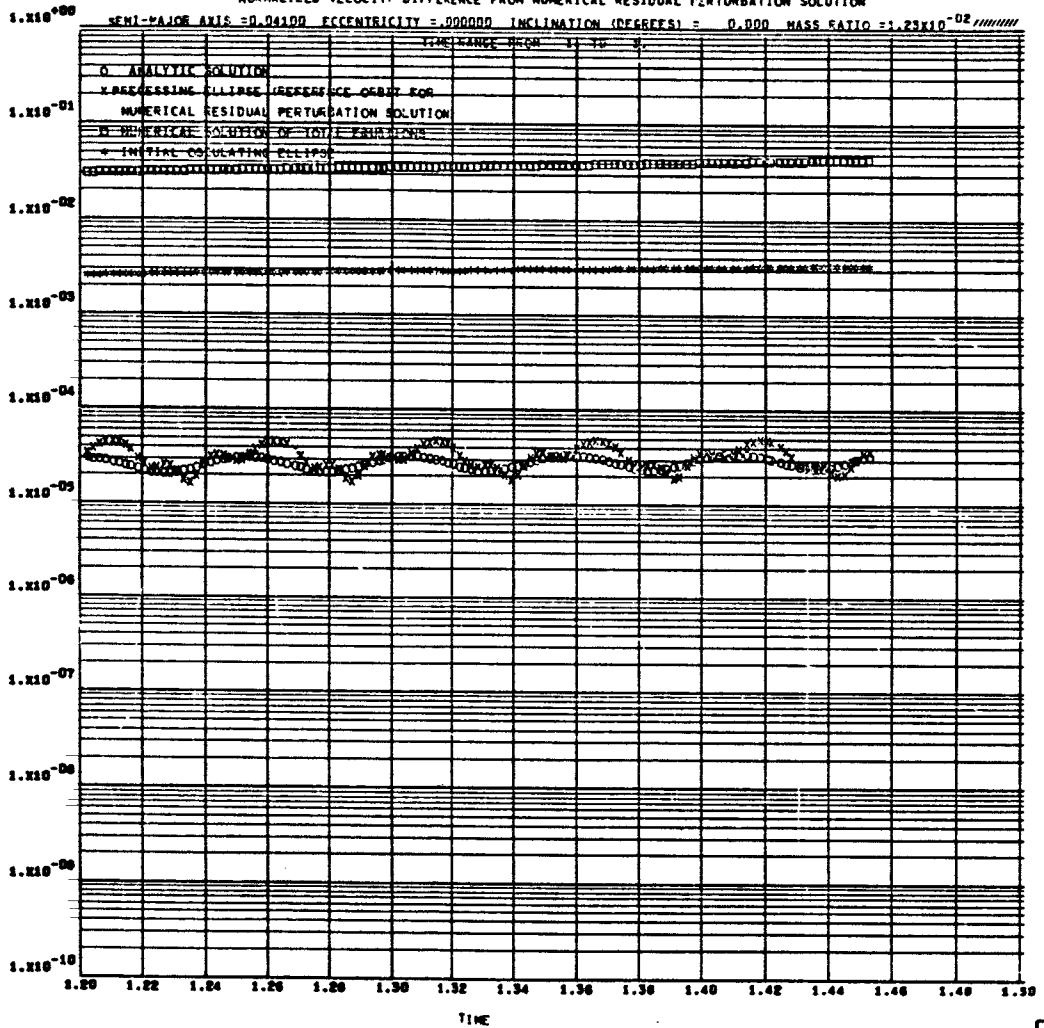


NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

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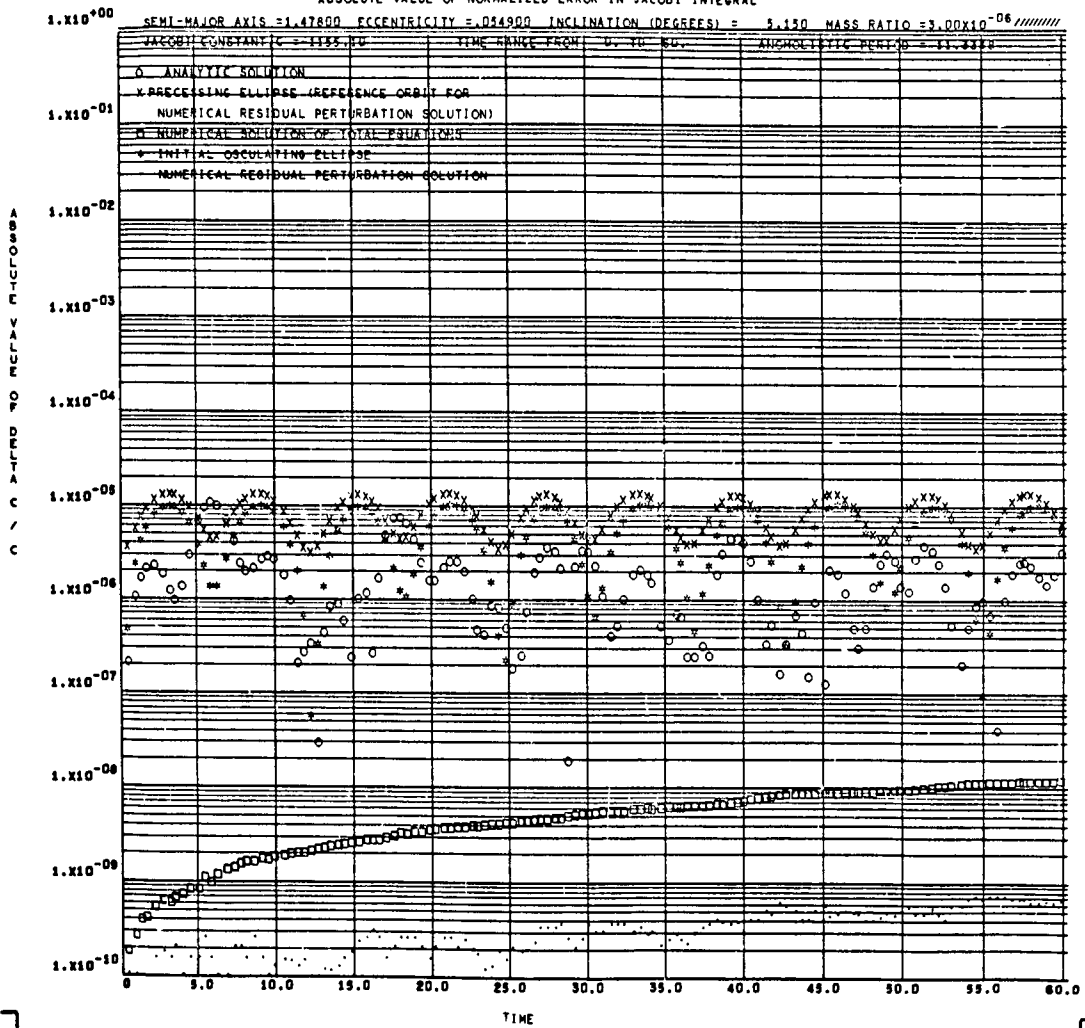


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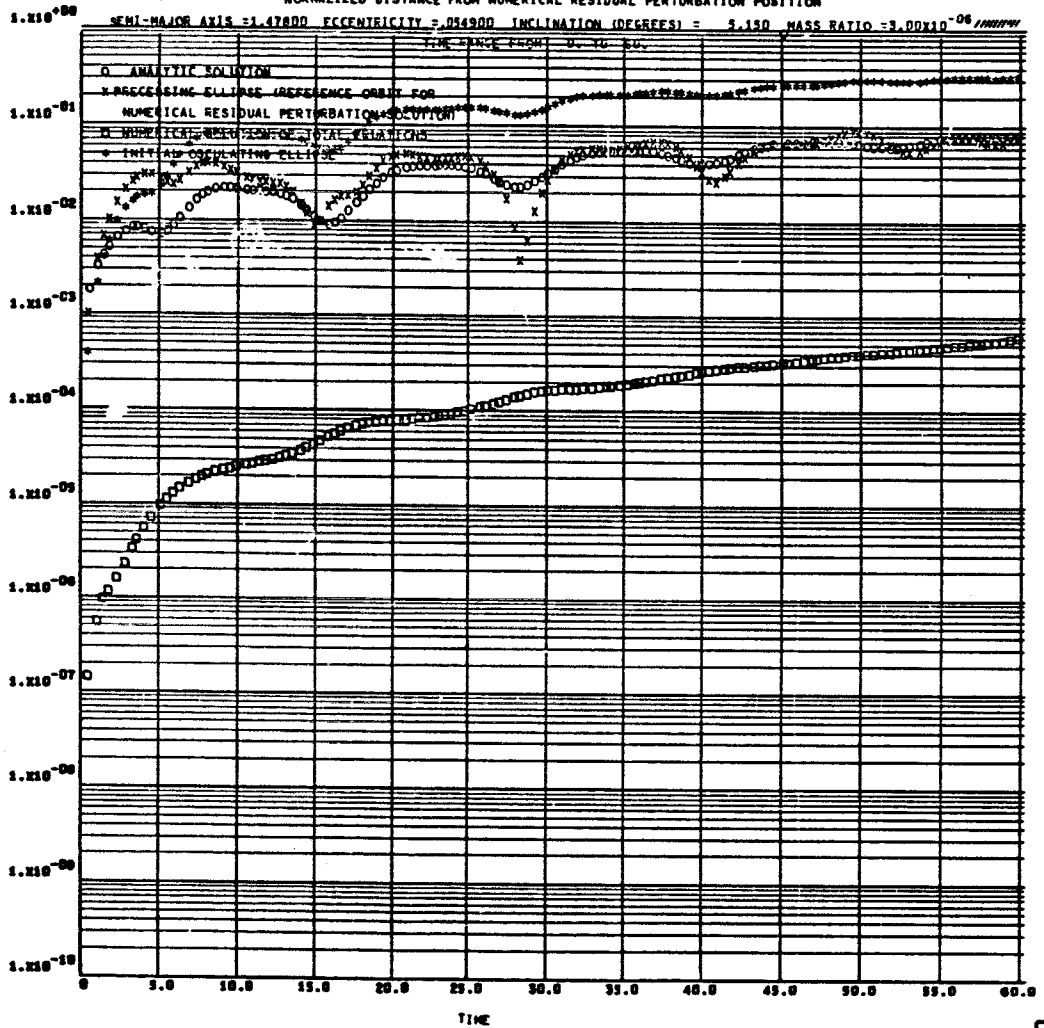


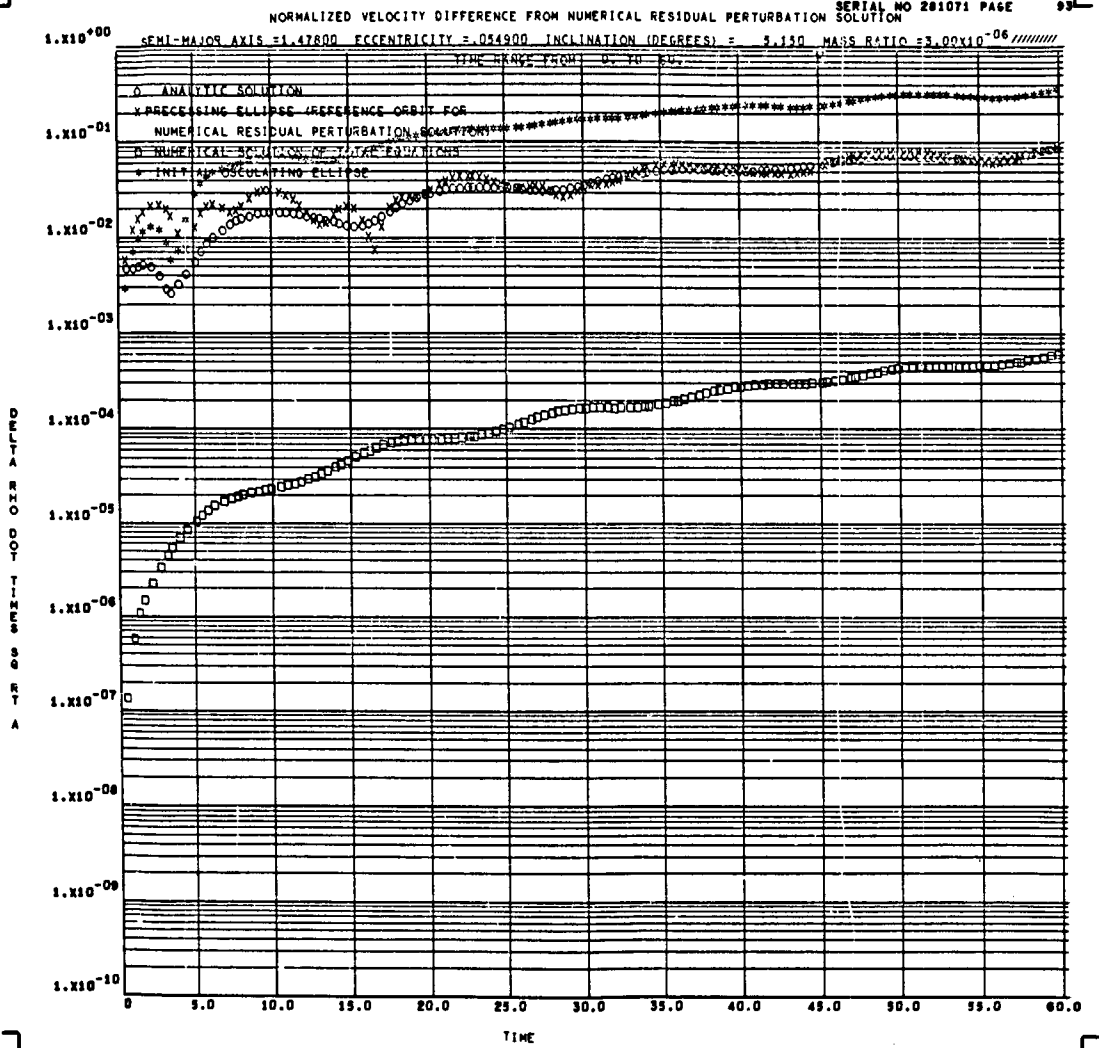
ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

SERIAL NO 281071 PAGE 91



NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

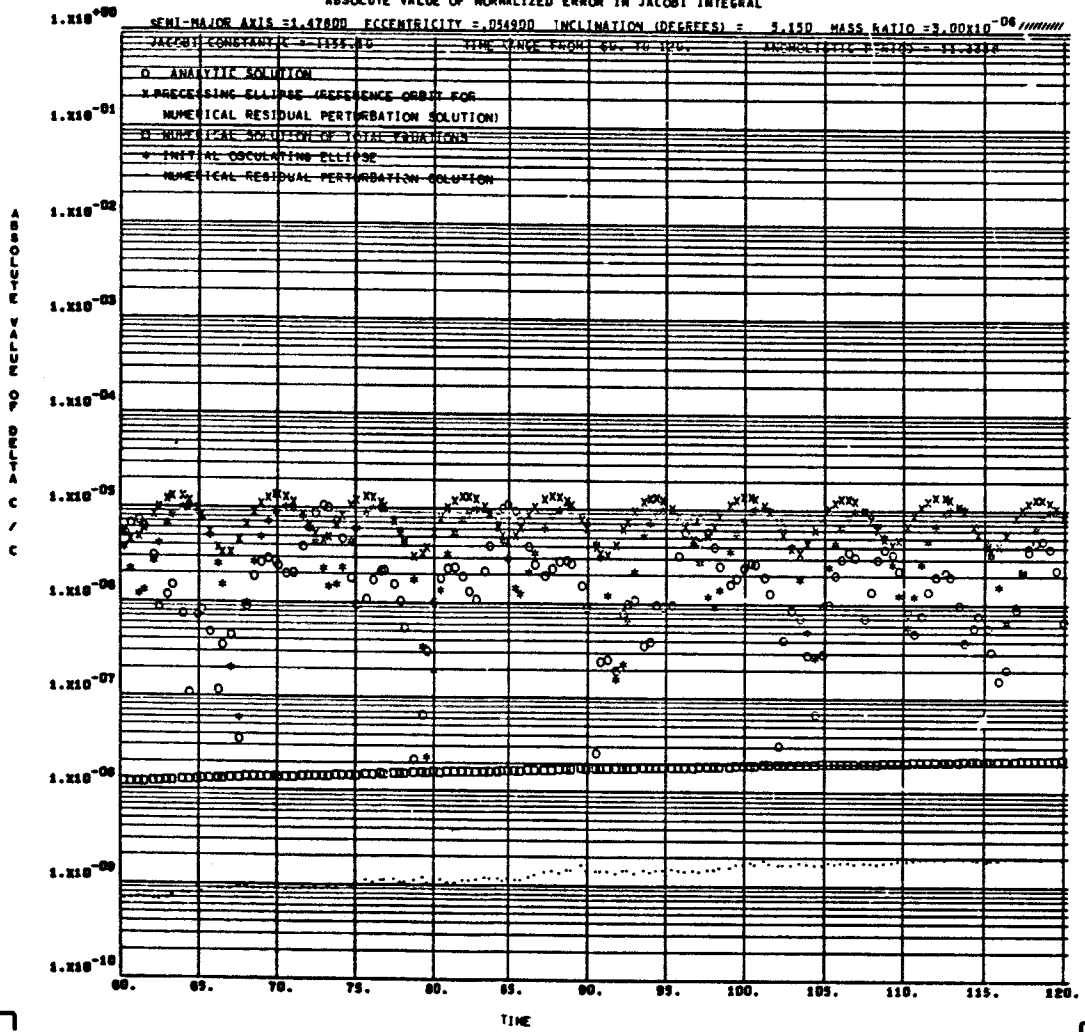


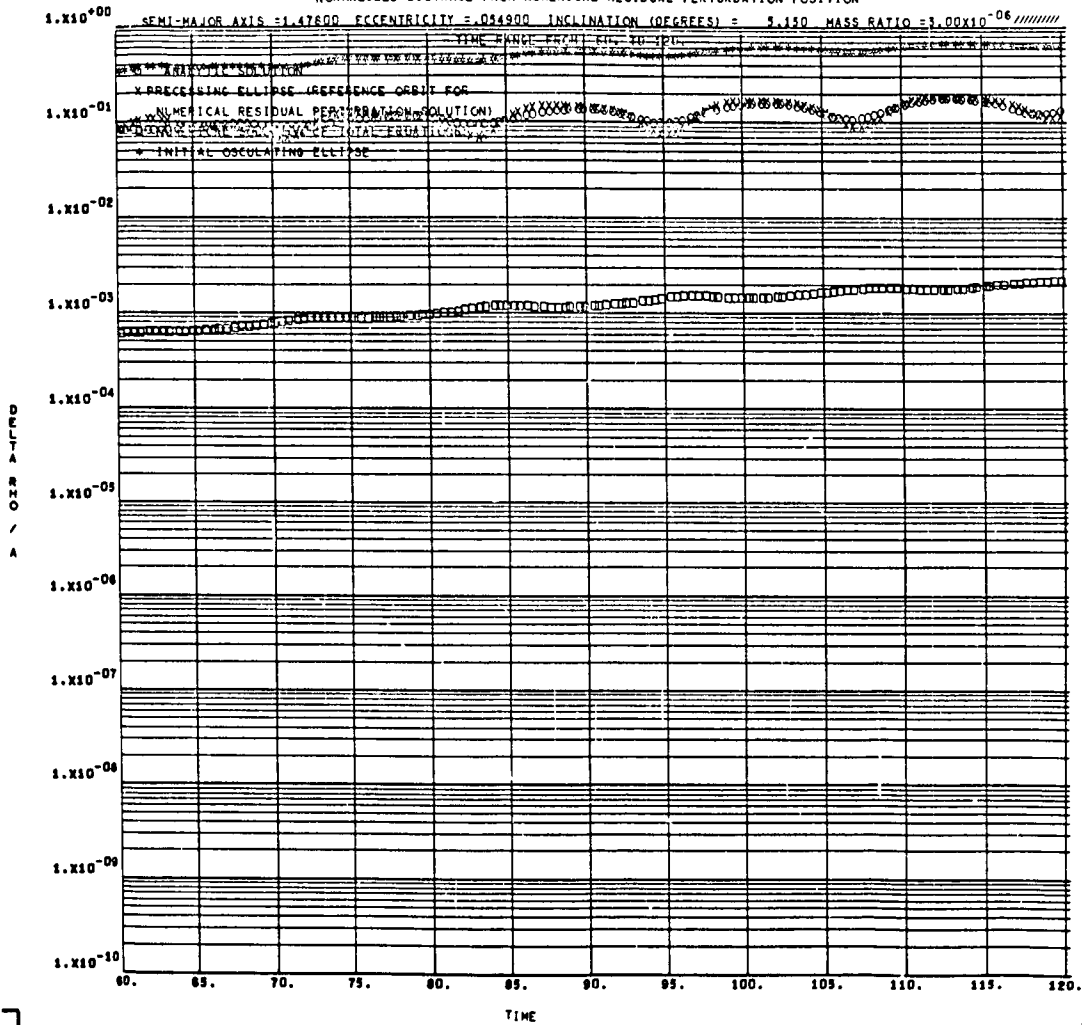




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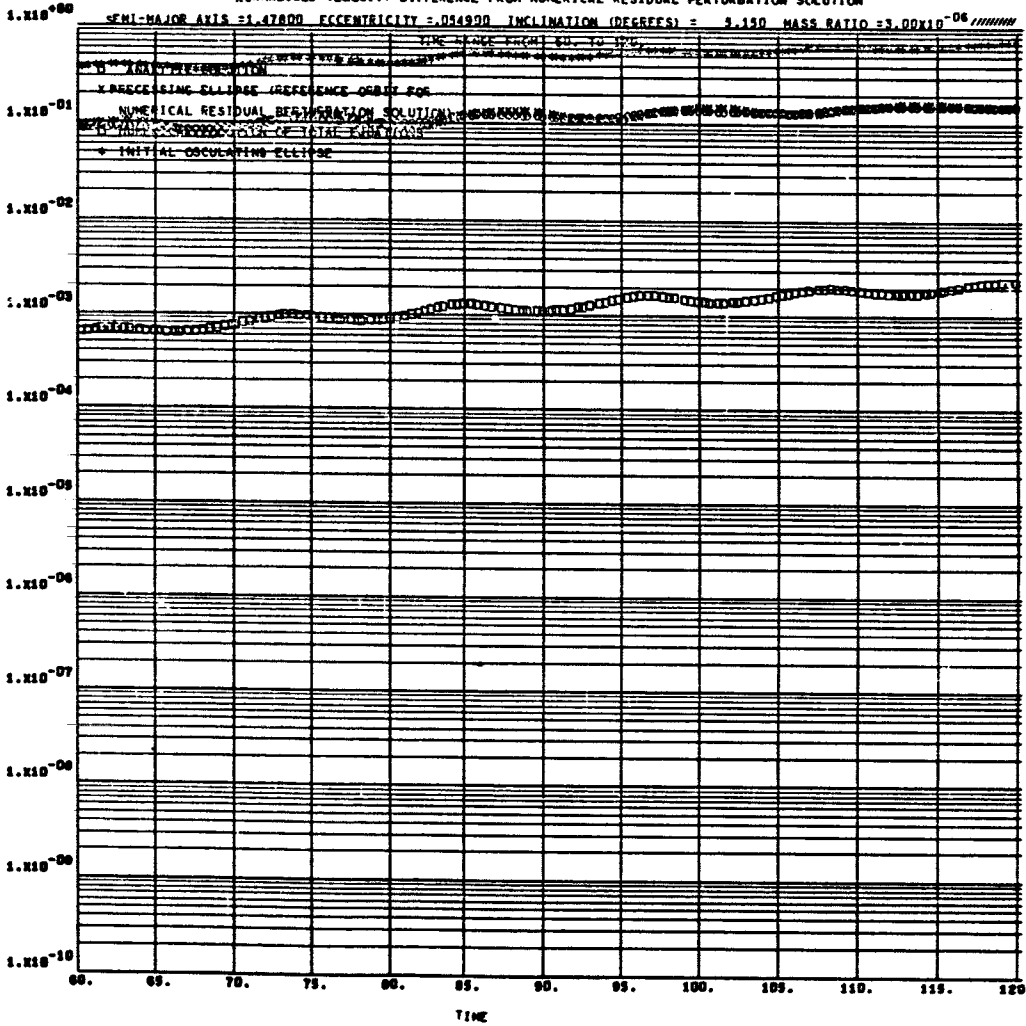
SERIAL NO 201071 PAGE



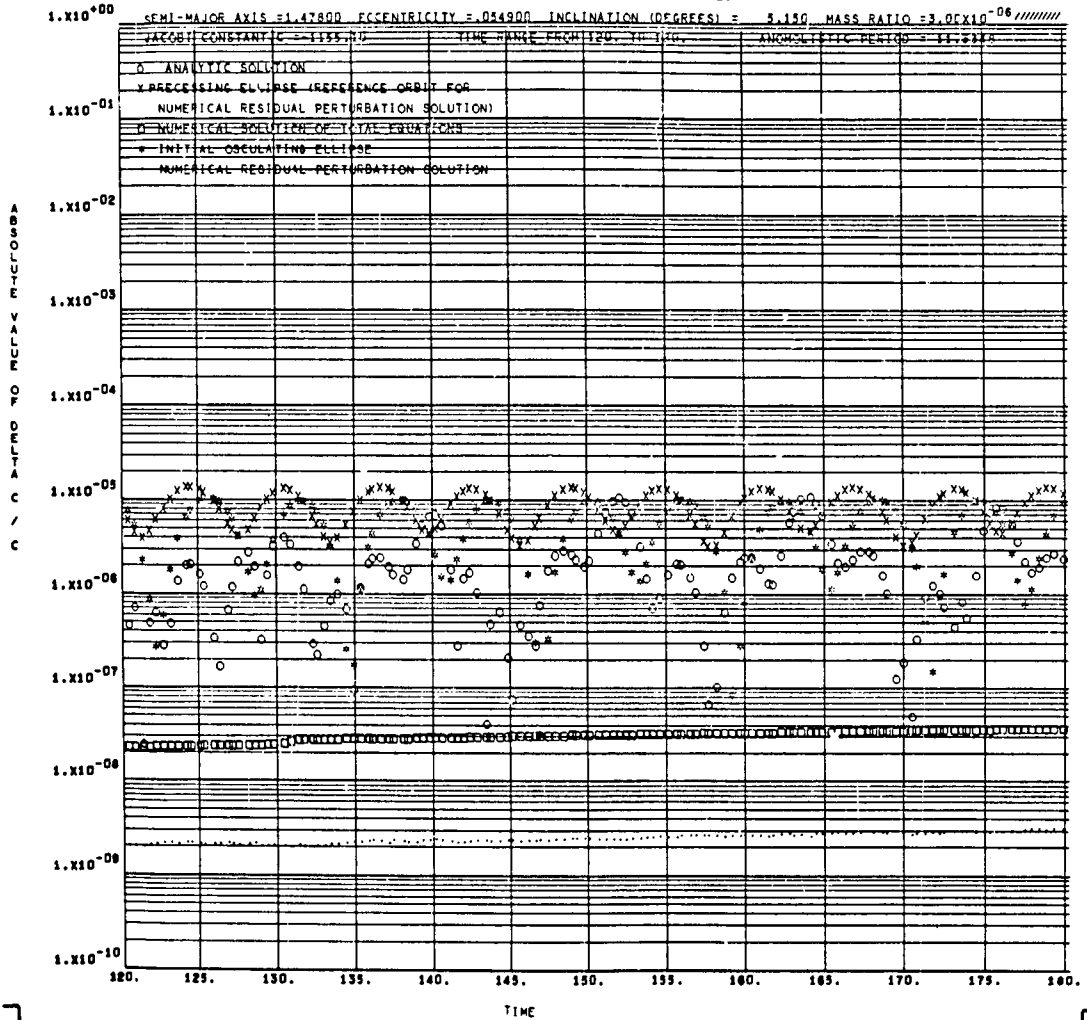


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SERIAL NO 261671 PAGE 2

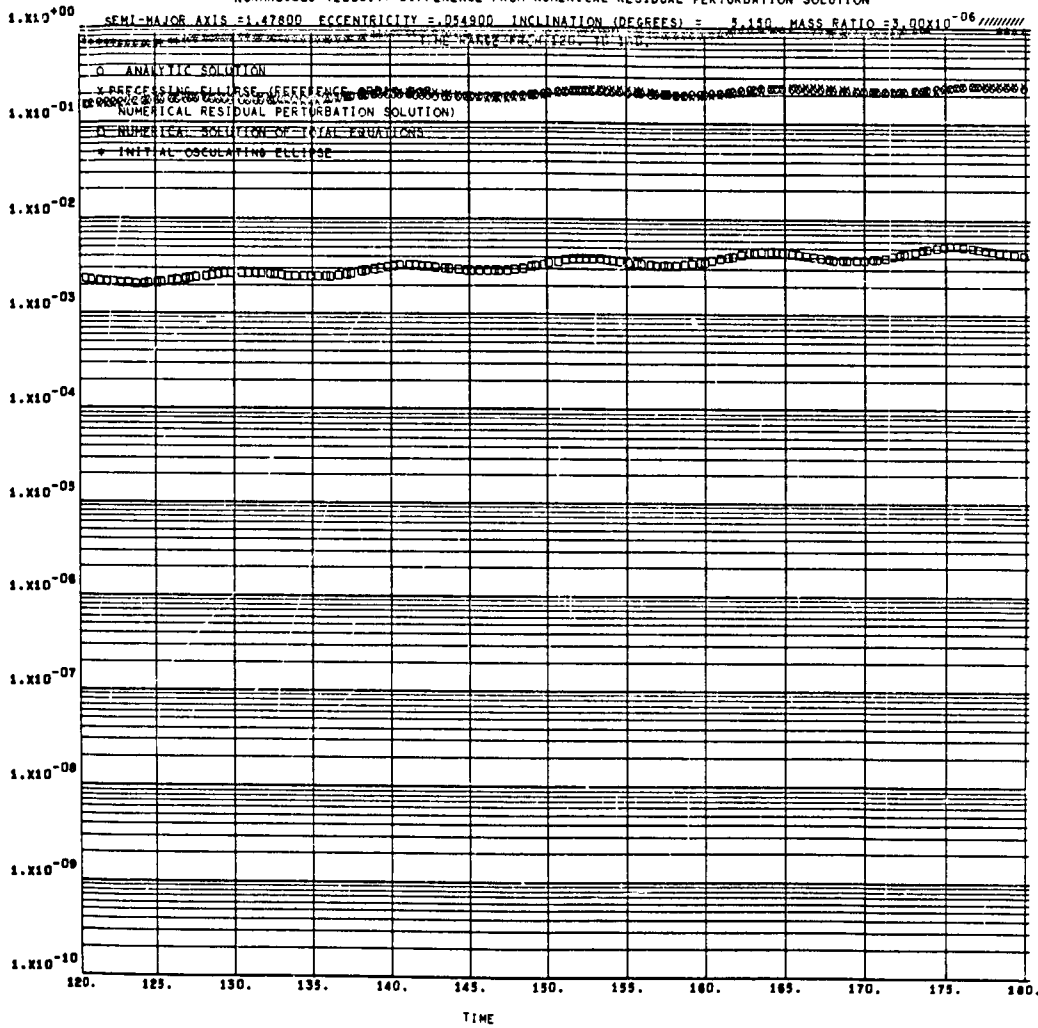


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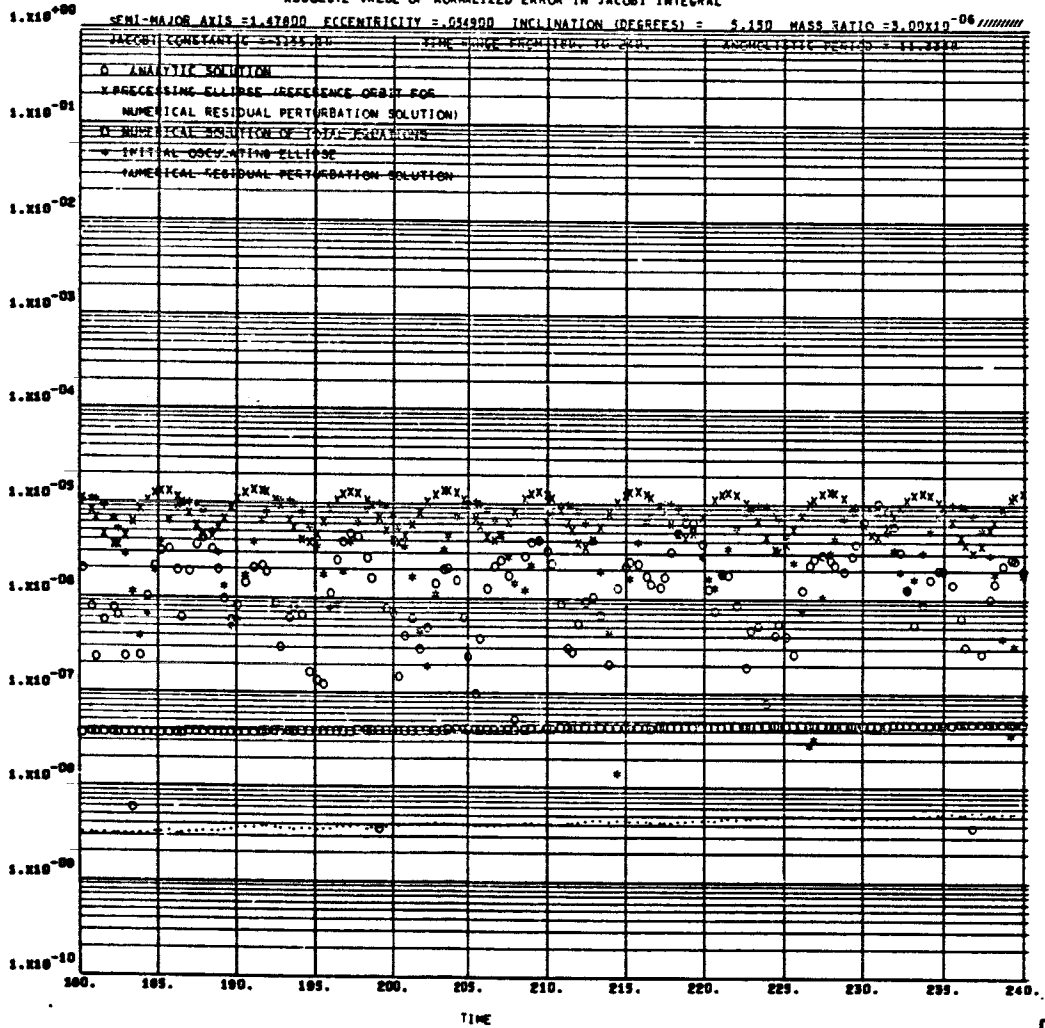




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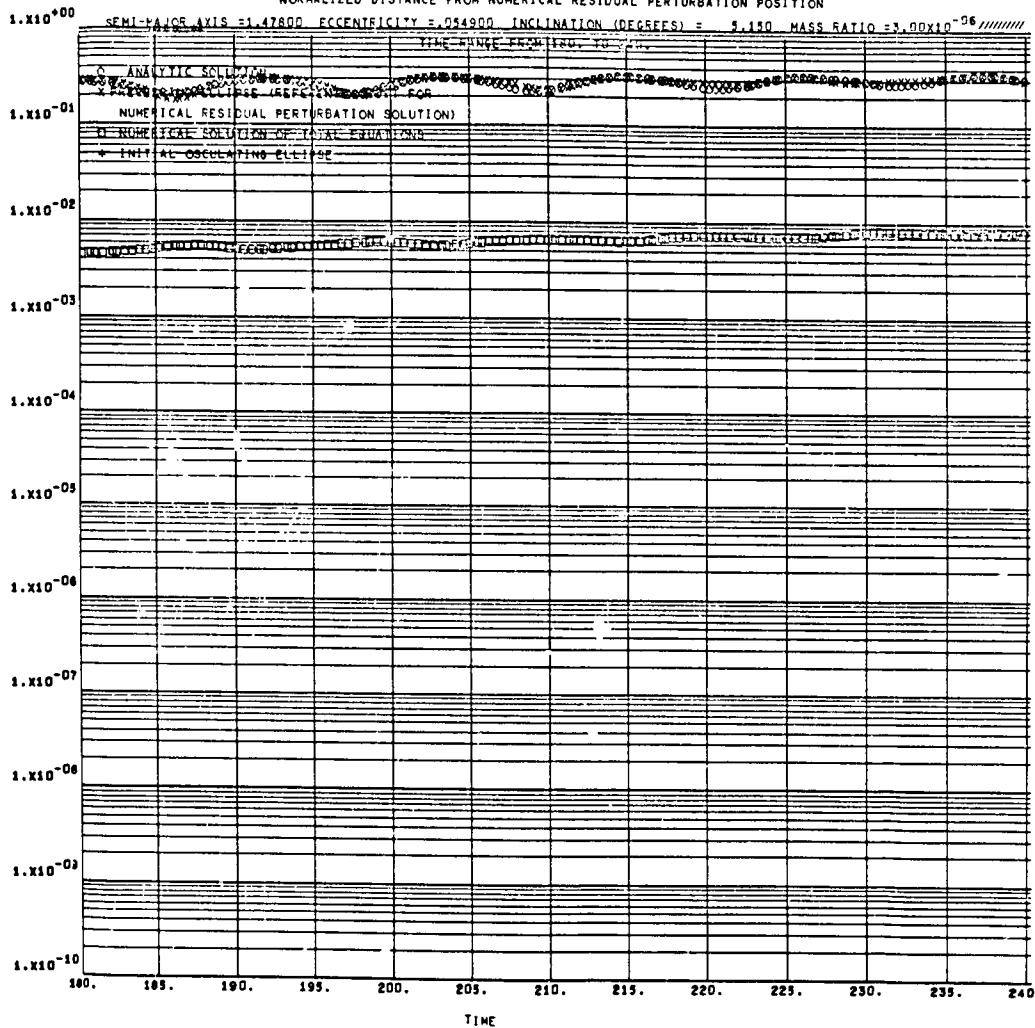


ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL



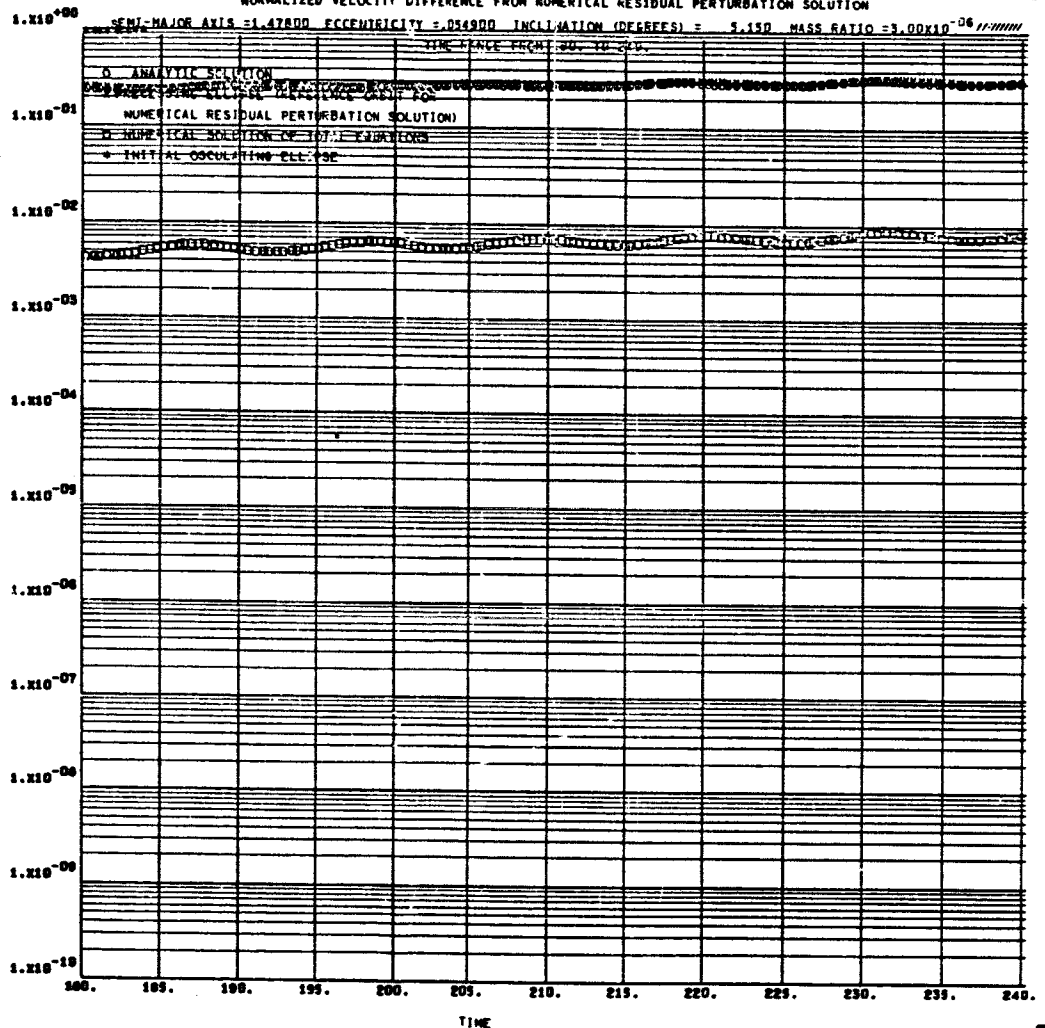
NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

SERIAL NO 281071 PAGE 101

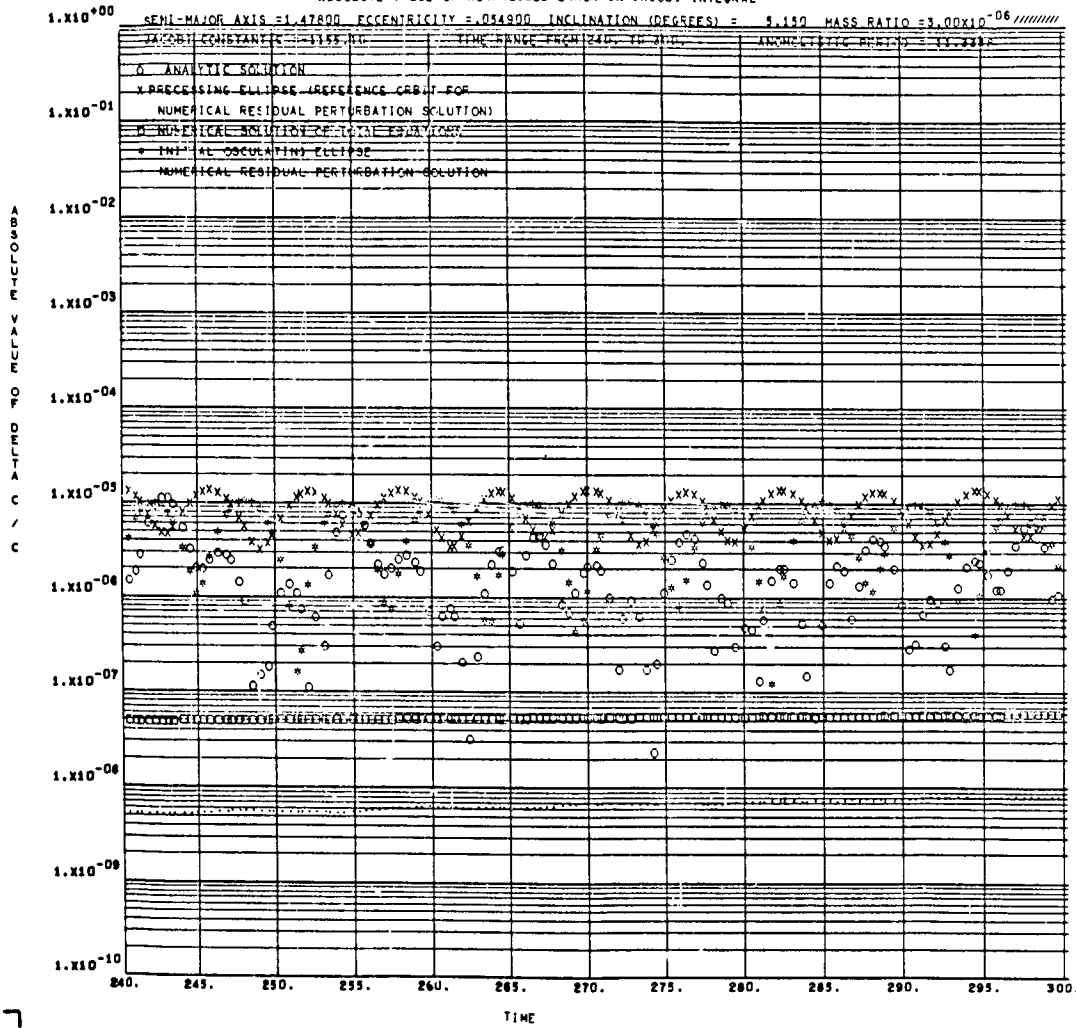


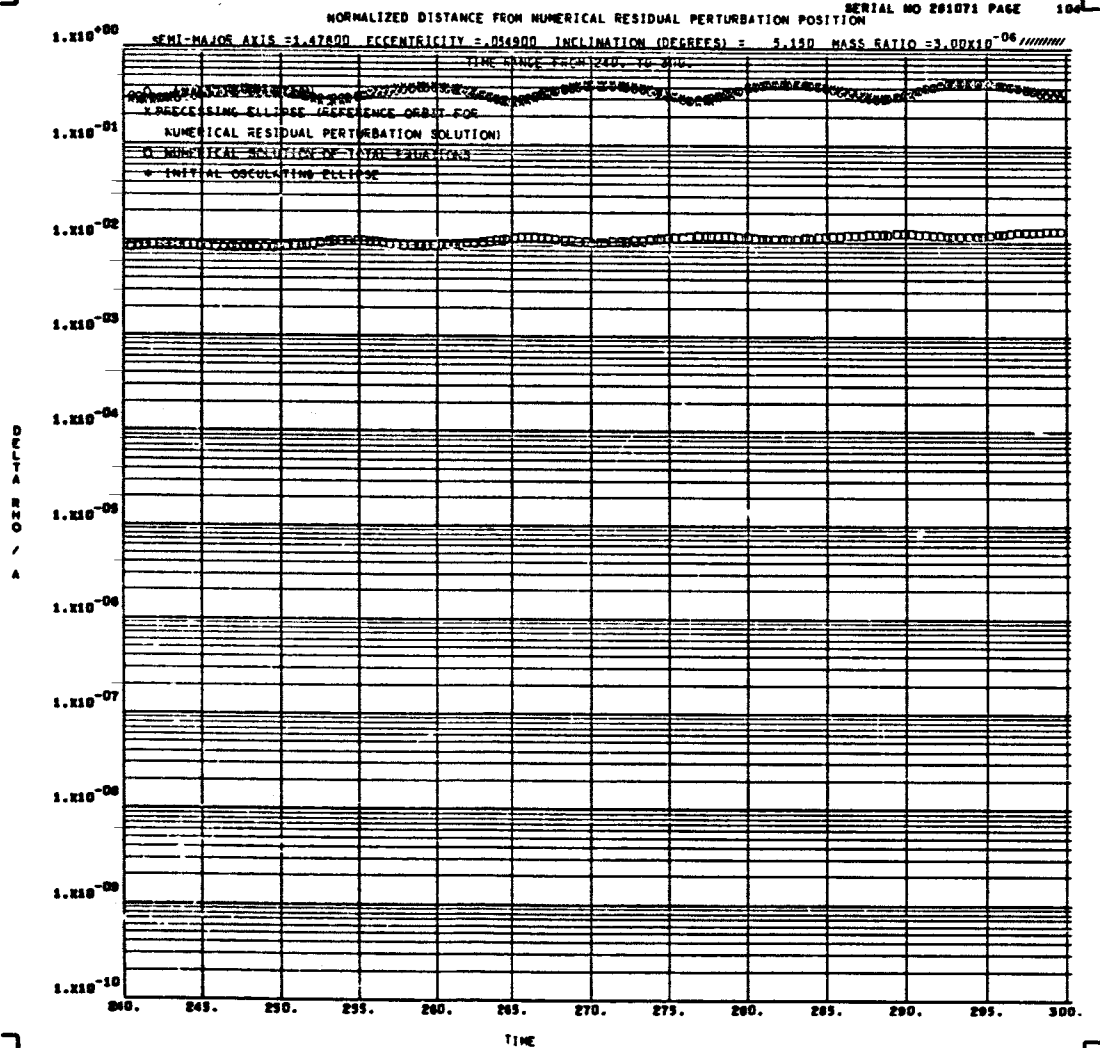


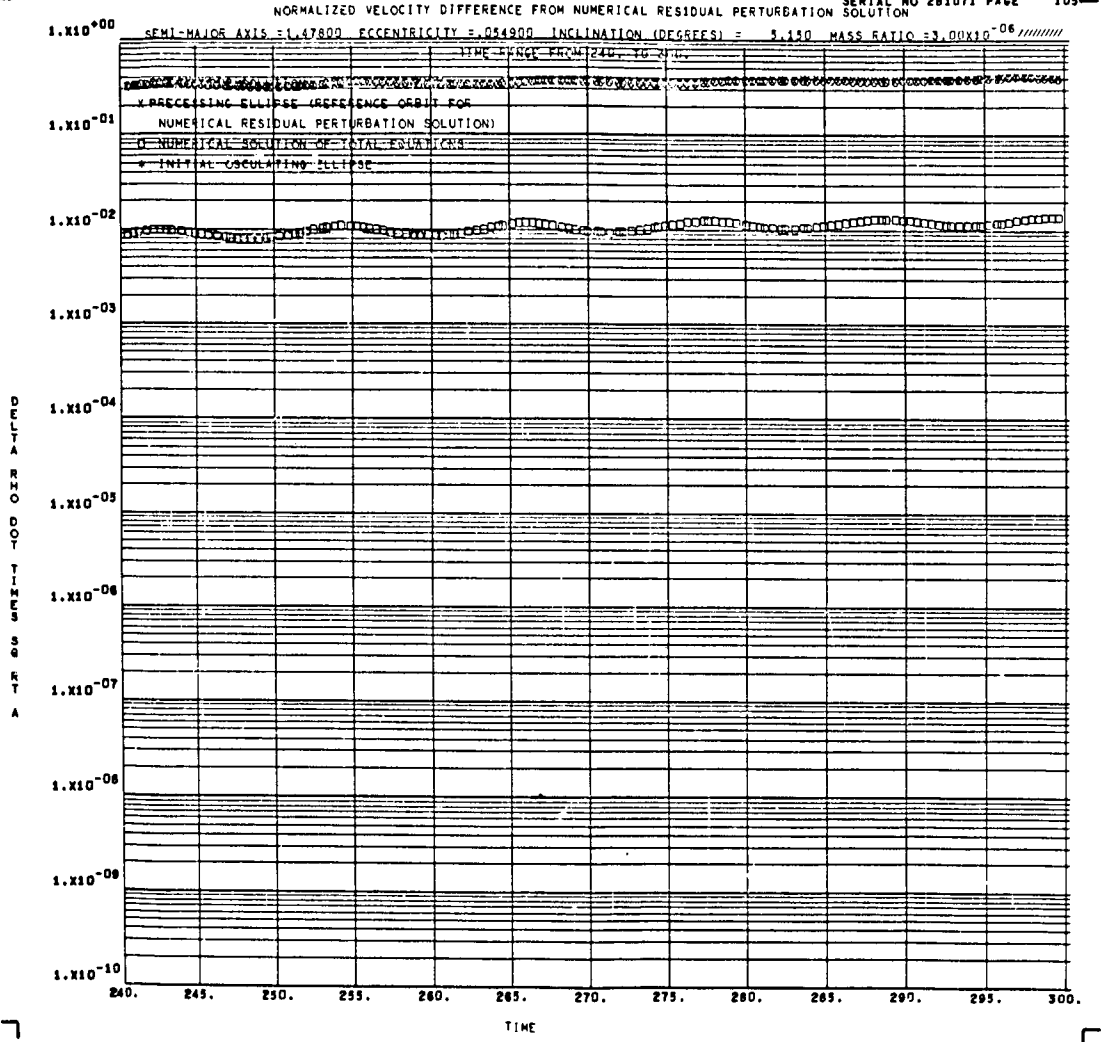
NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

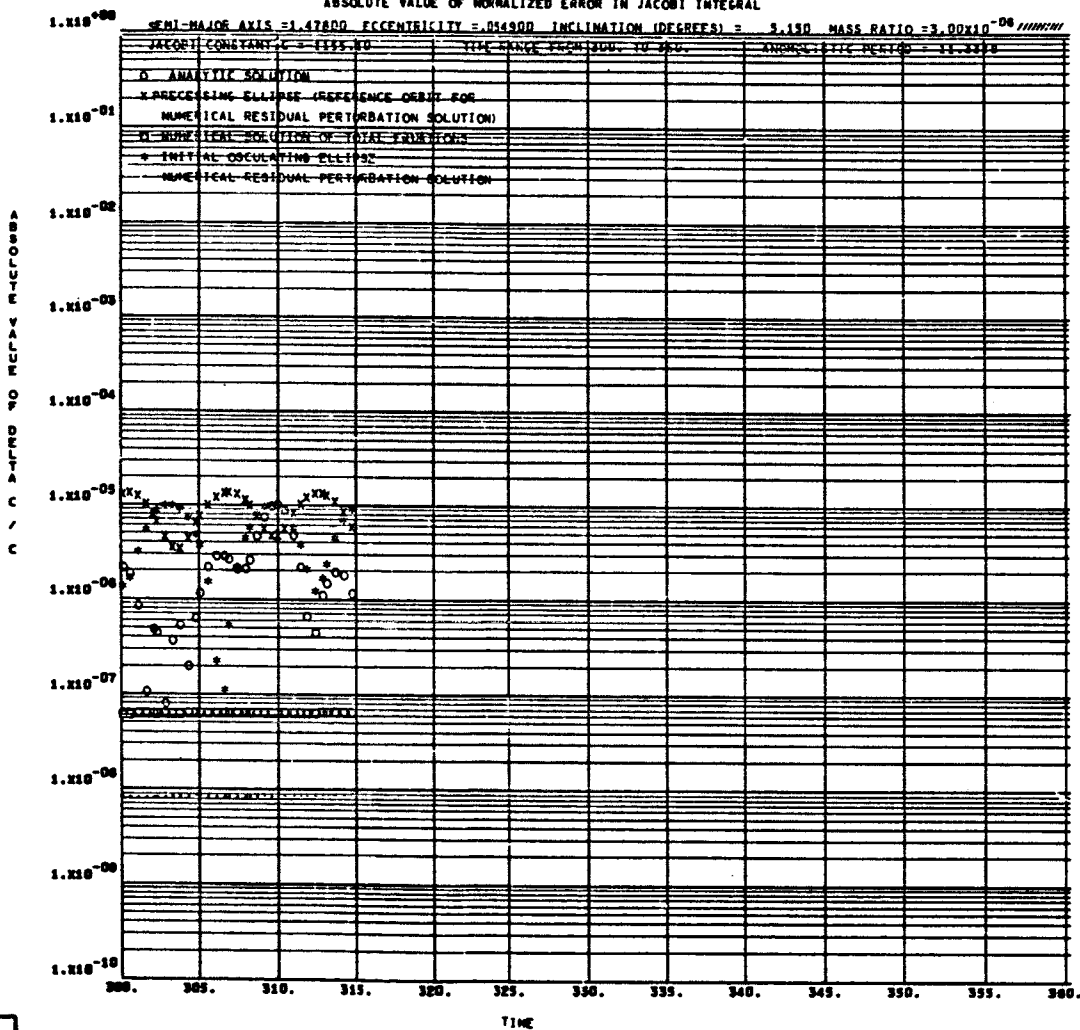






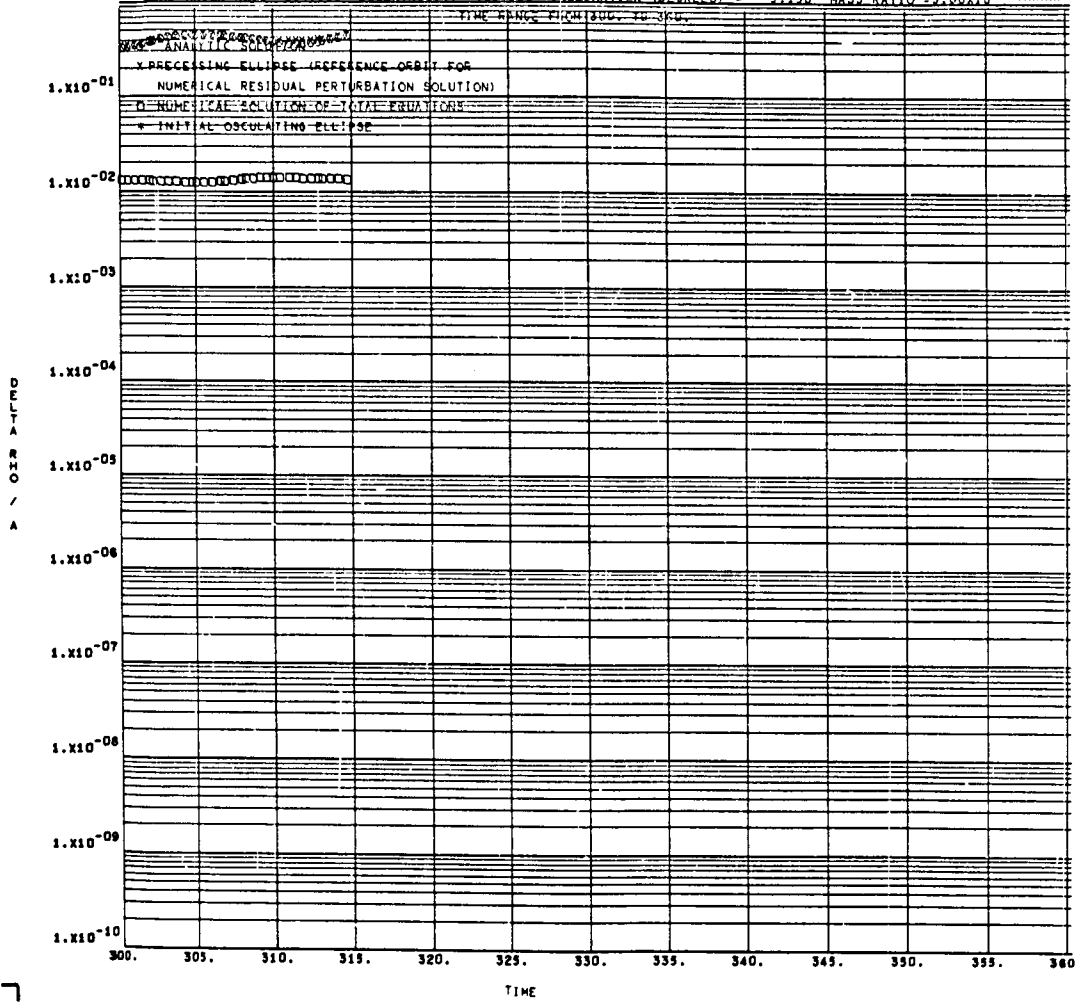
ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

SERIAL NO 201071 PAGE 10



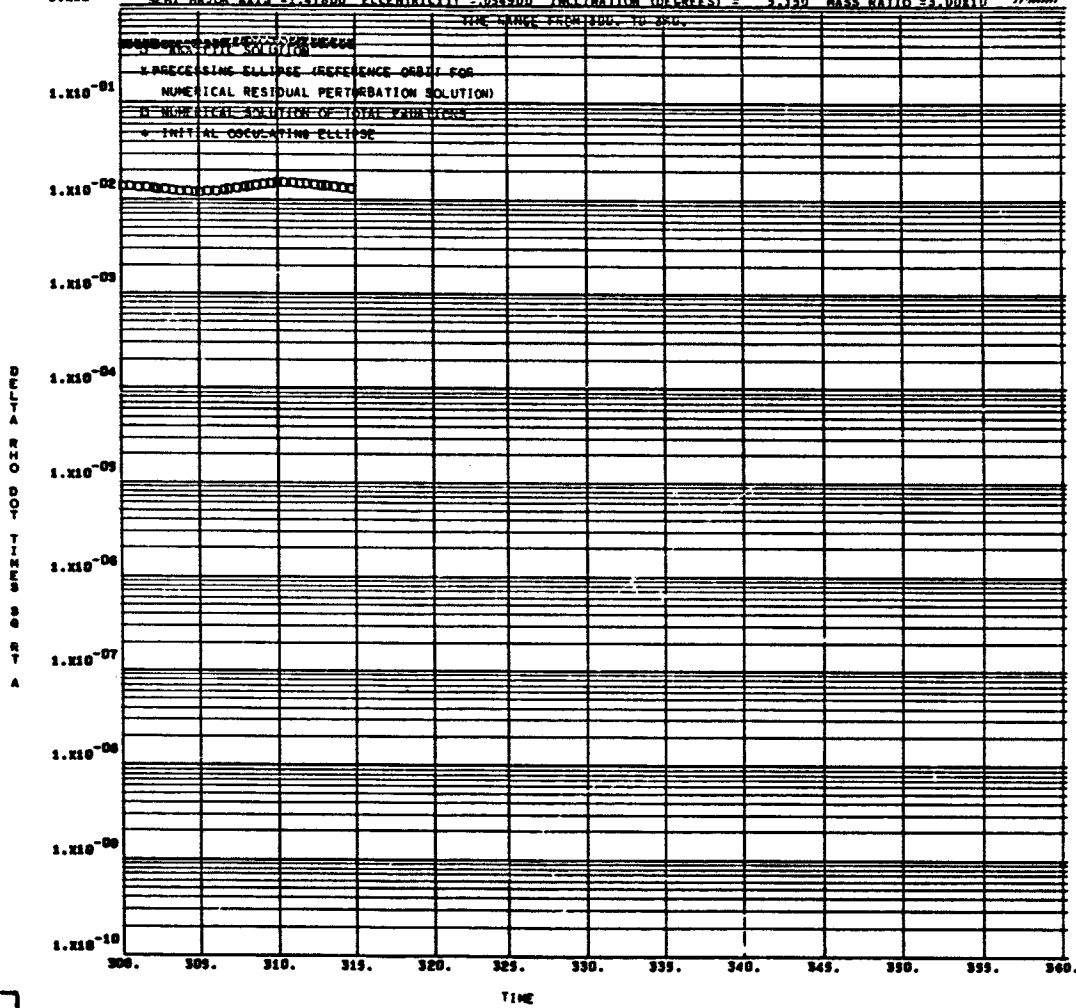
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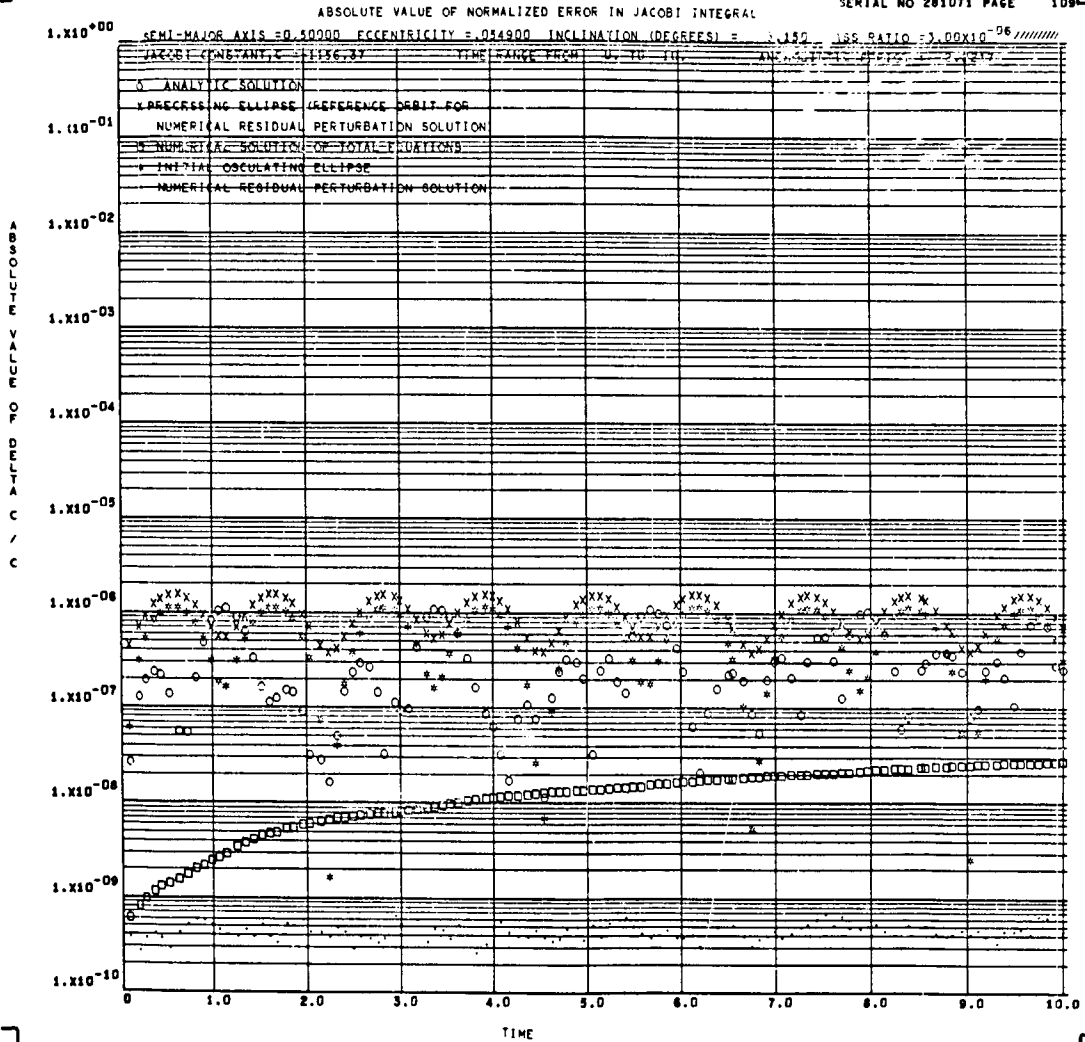
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NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

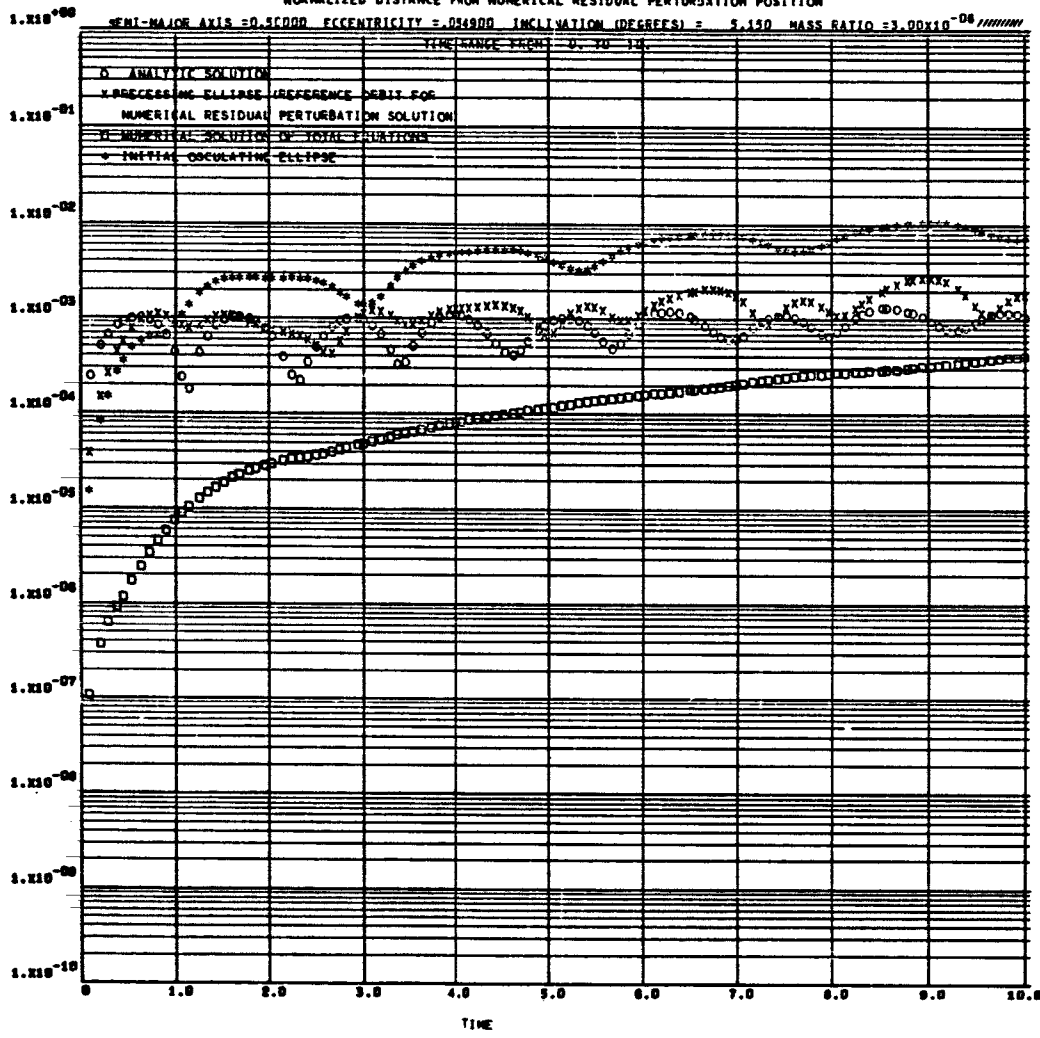
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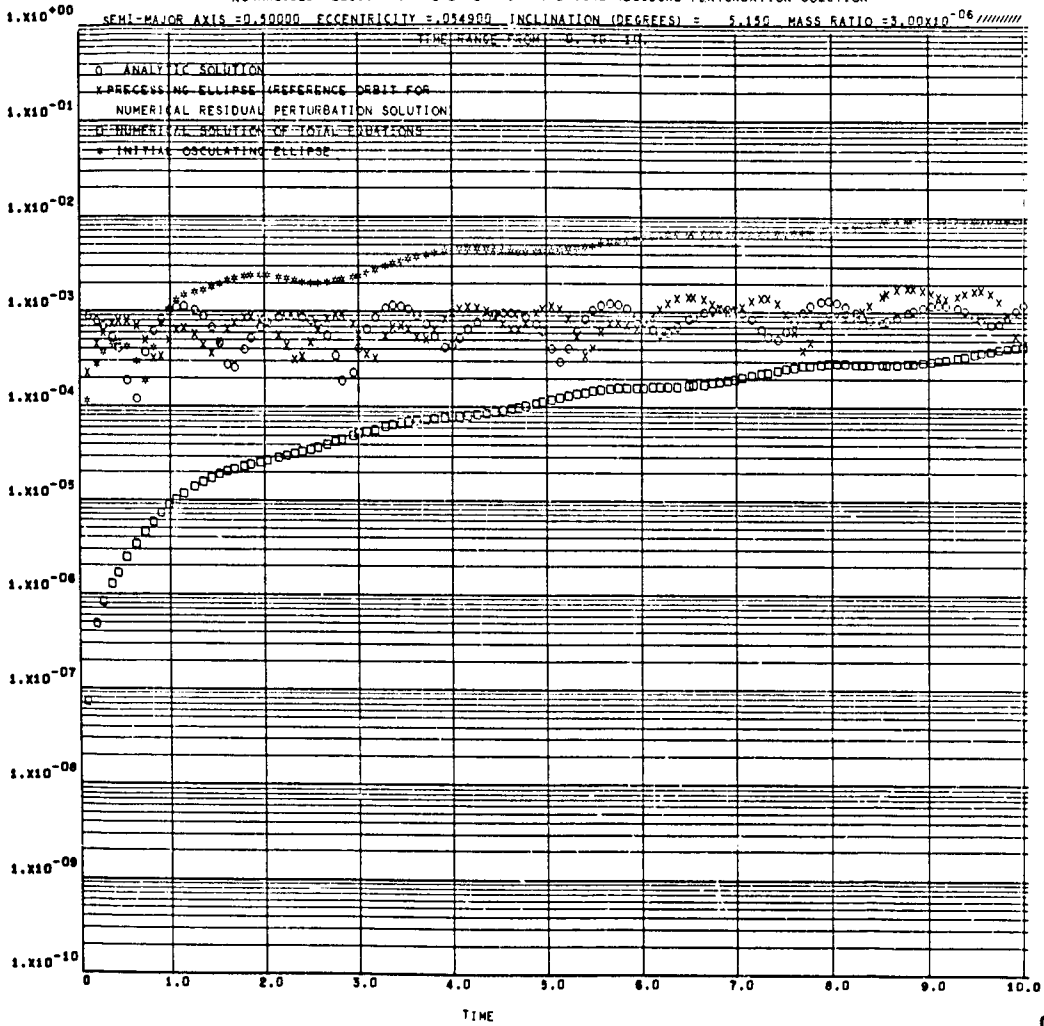


NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

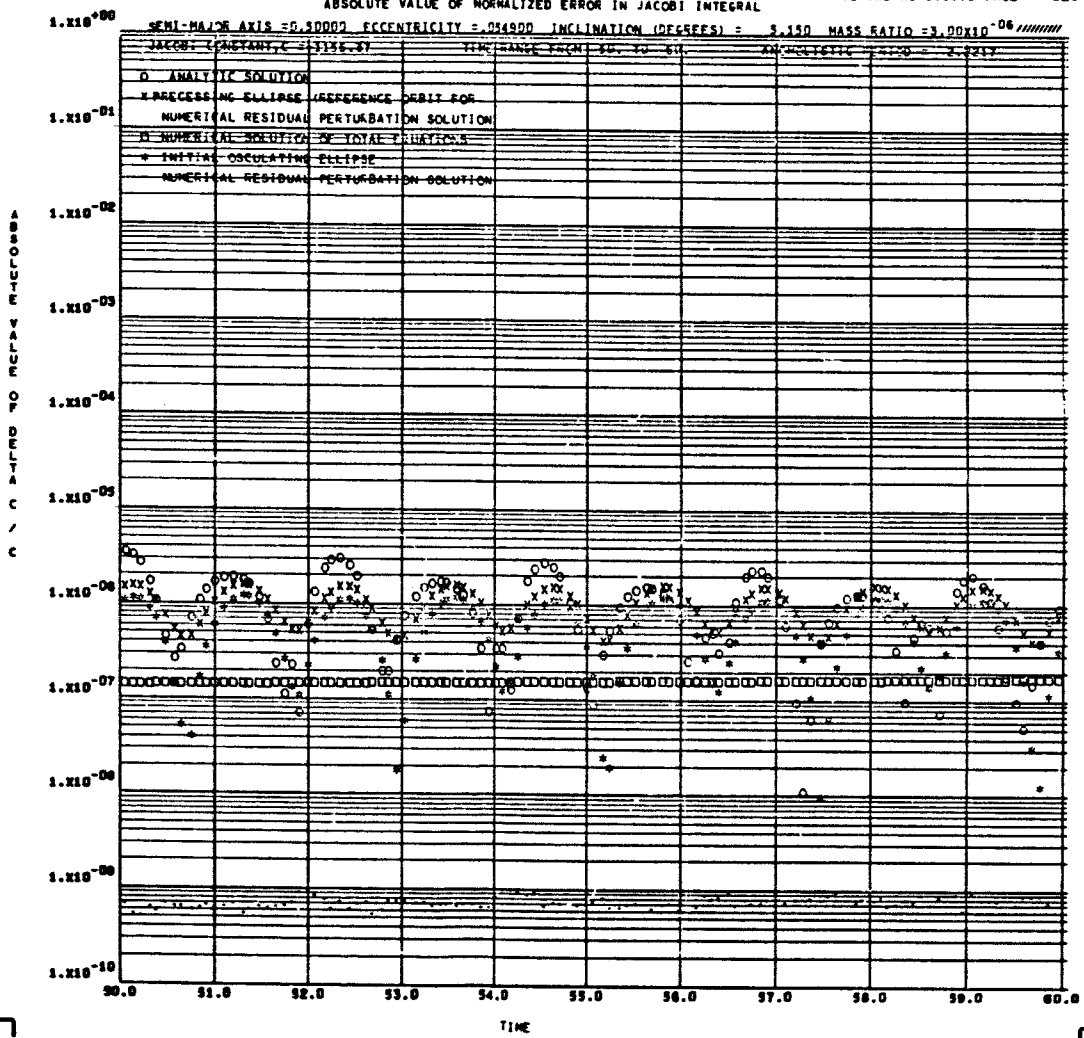


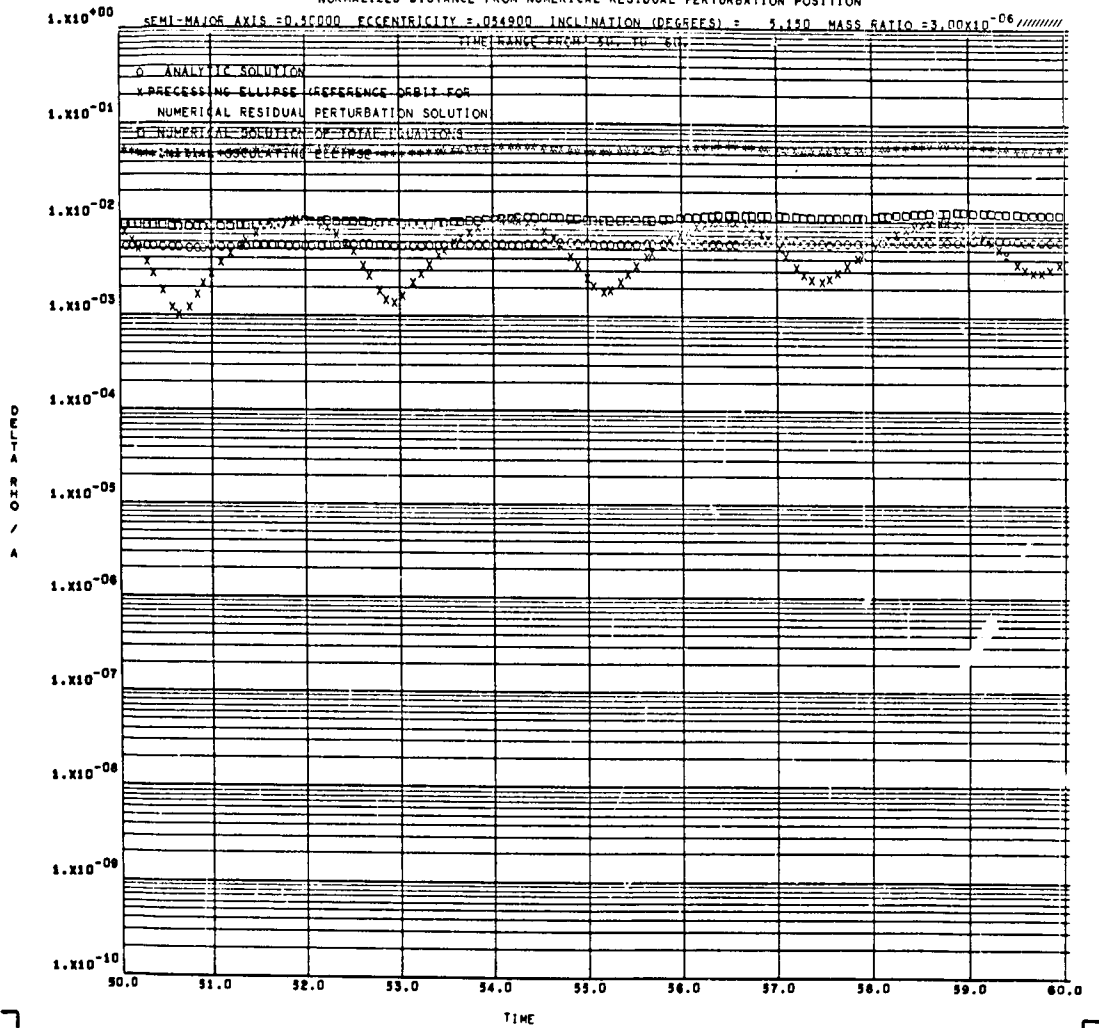
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NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



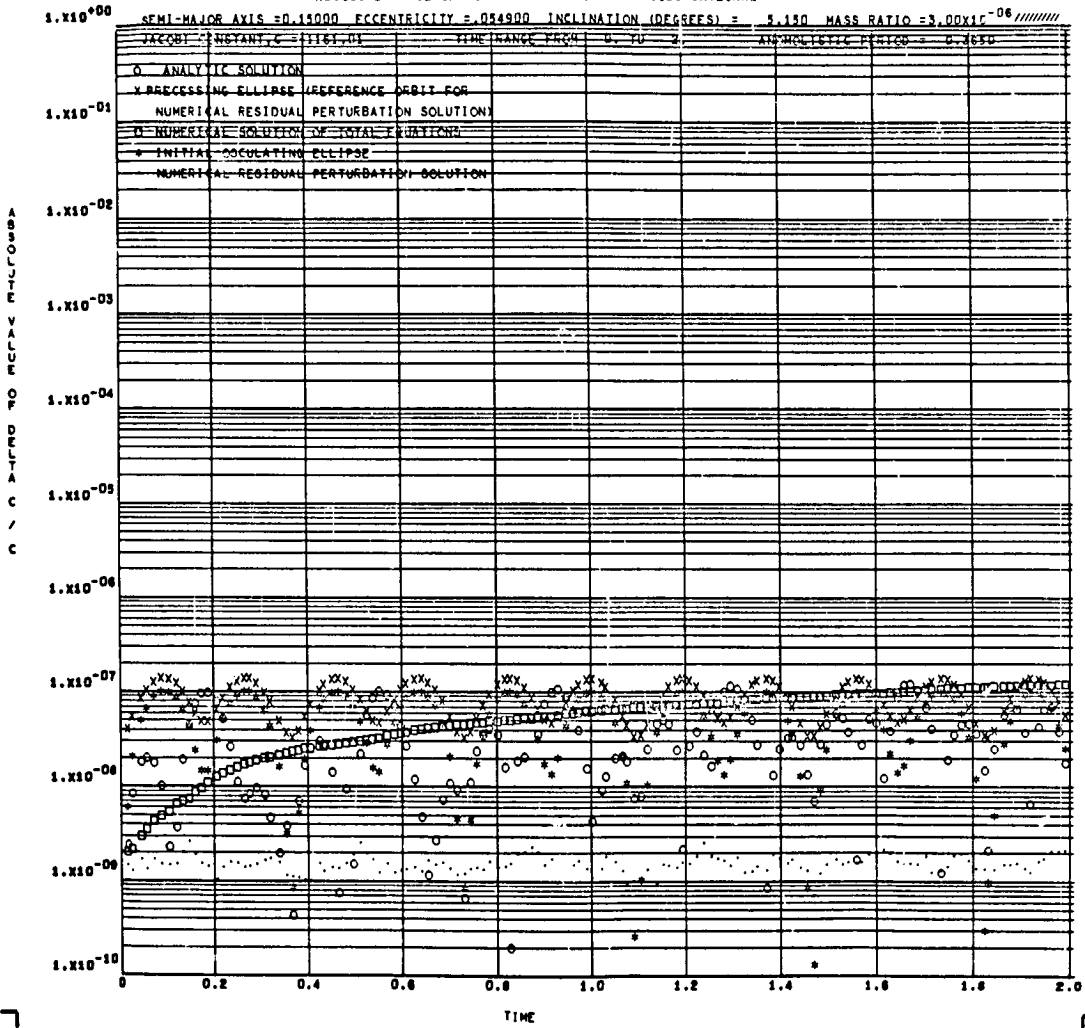
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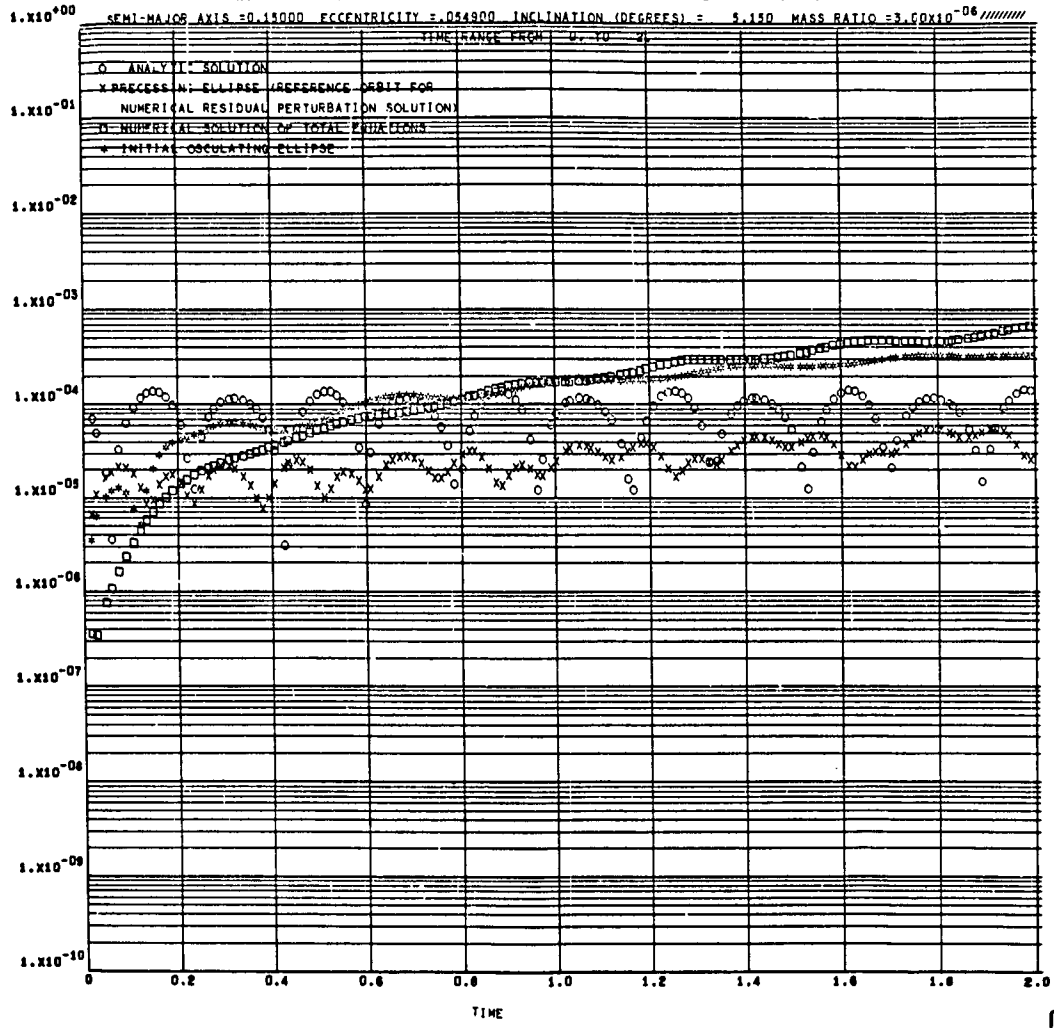


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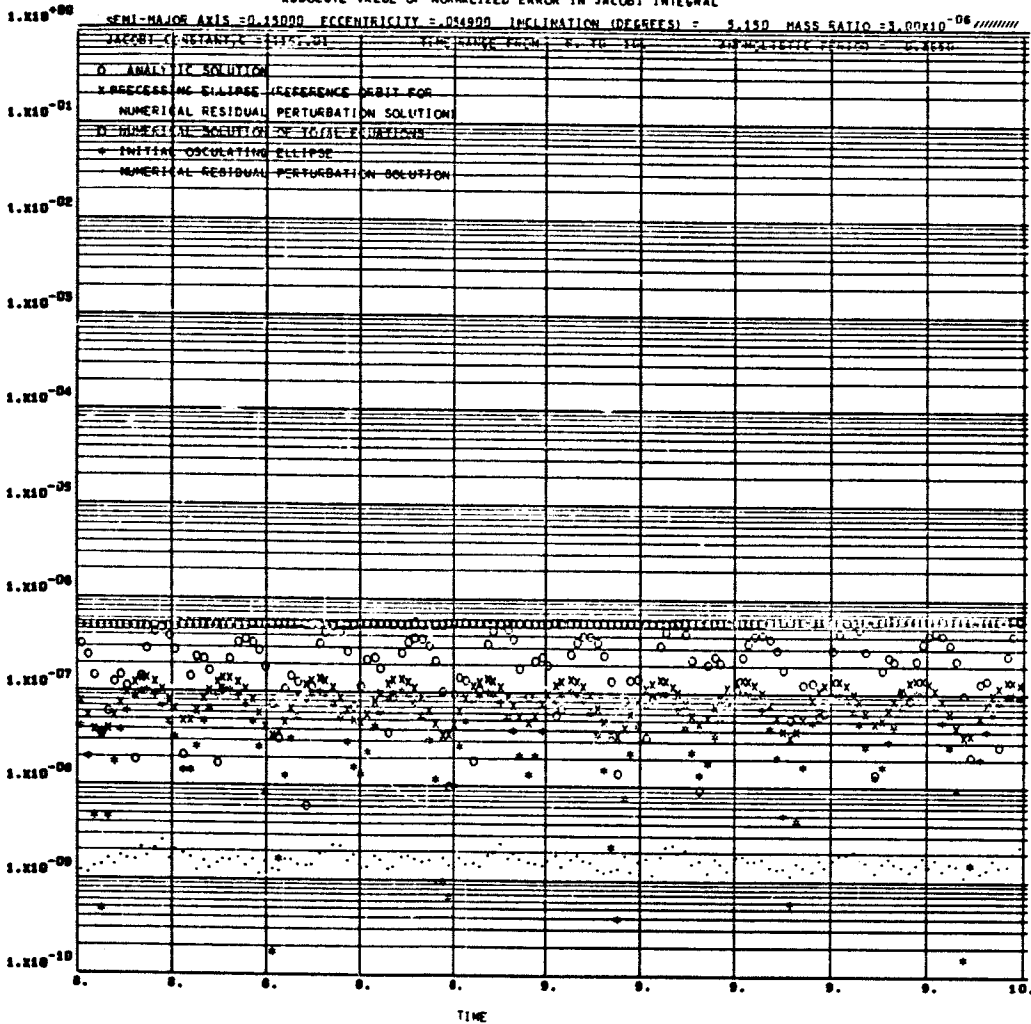


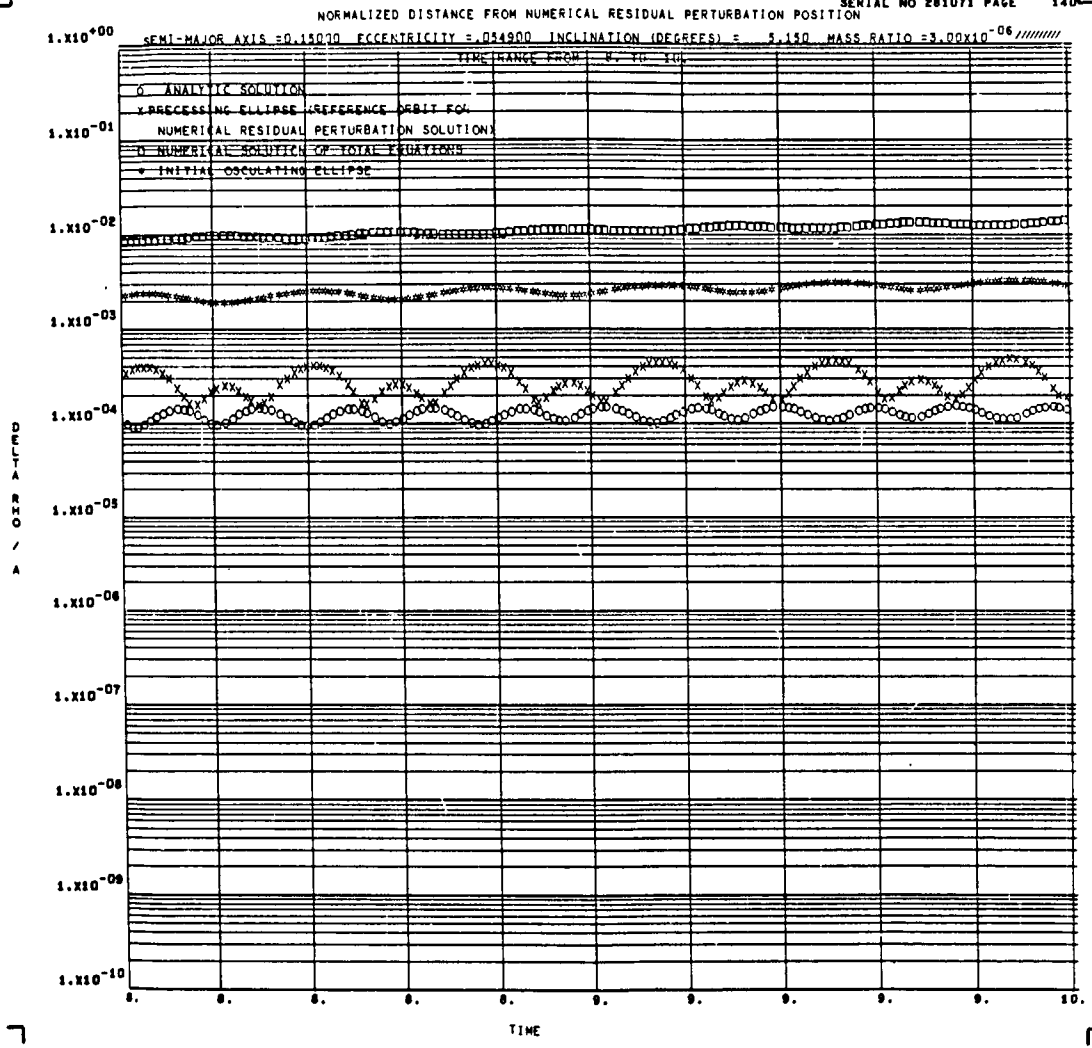
NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION





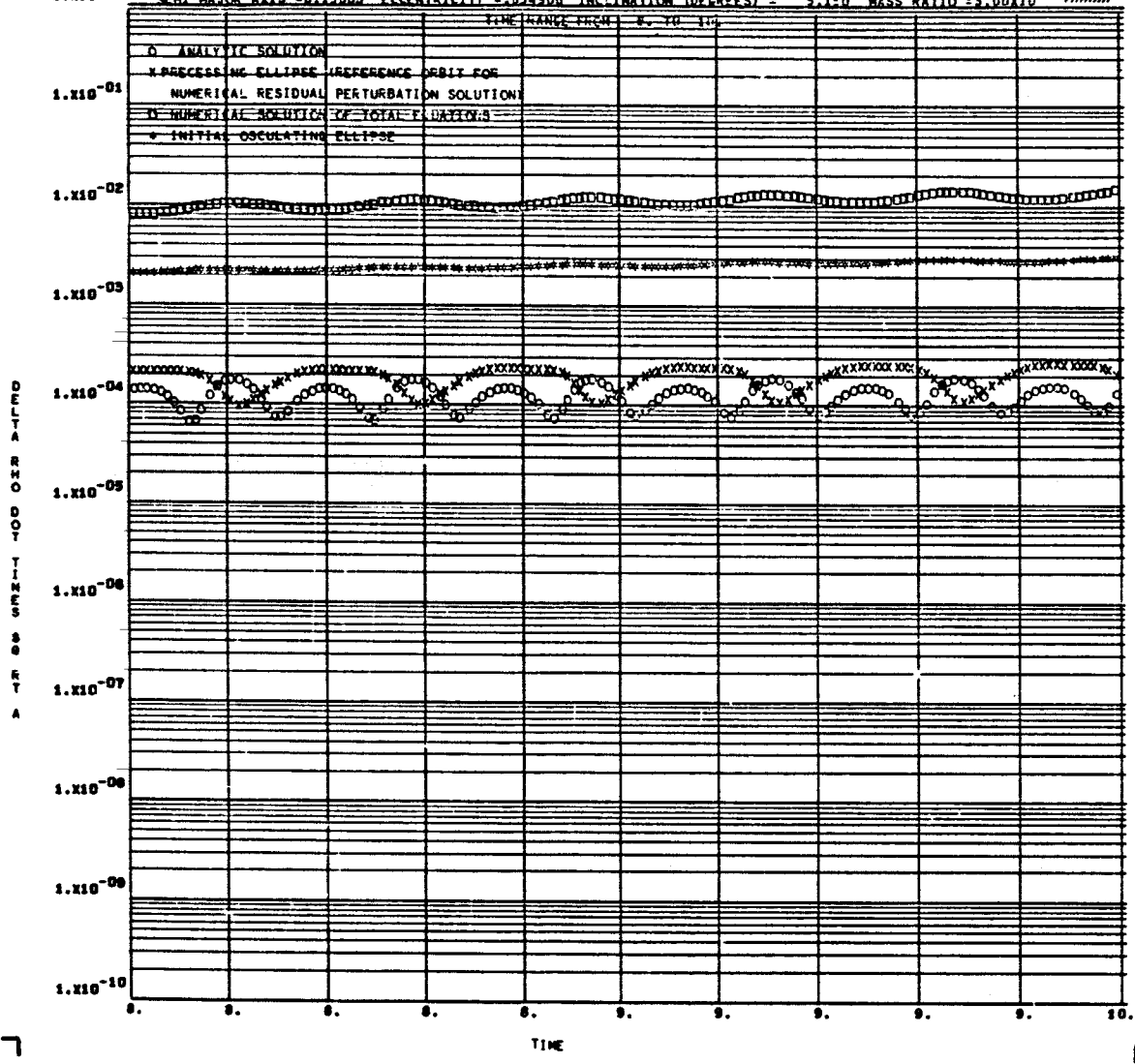
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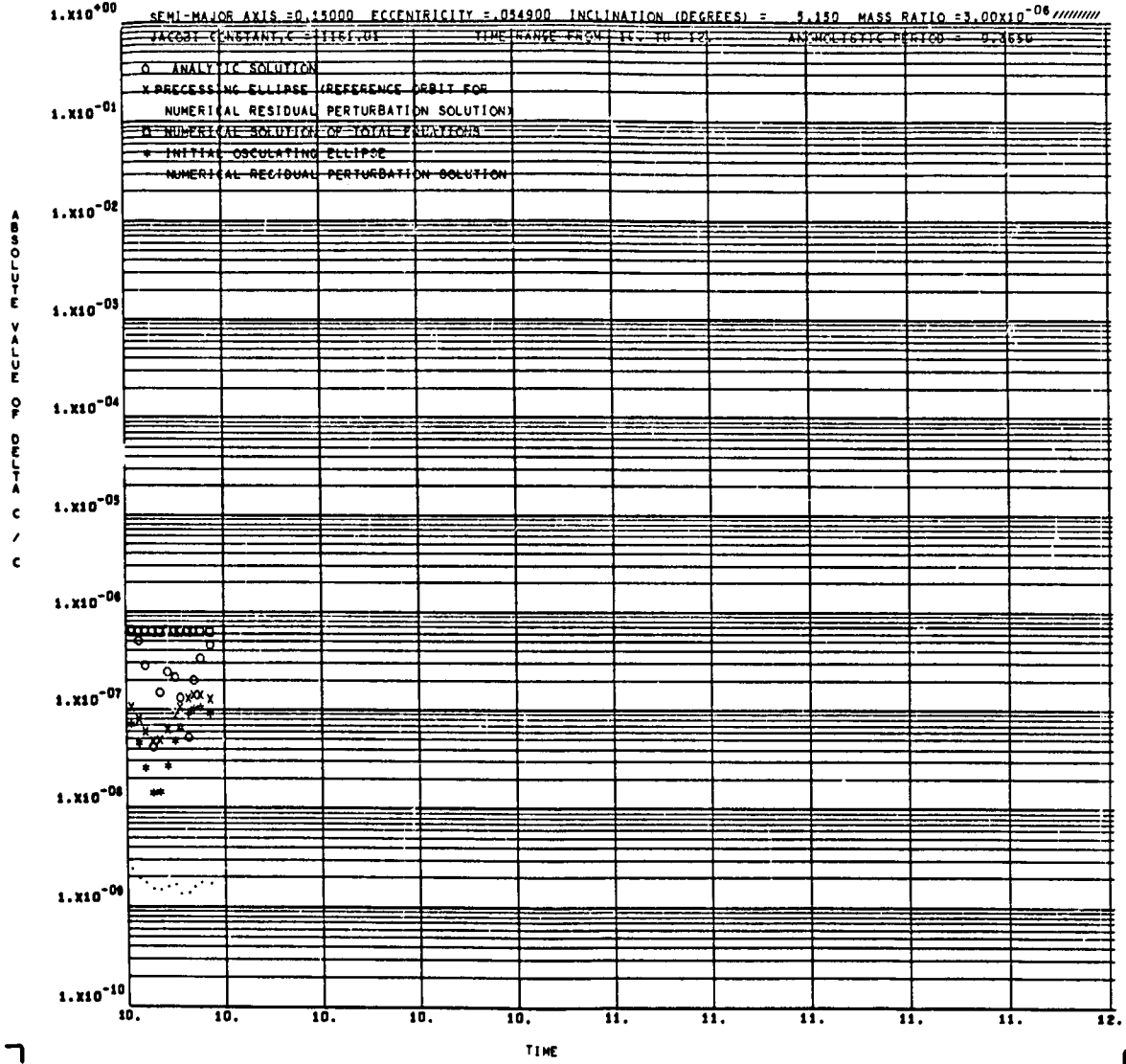


NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

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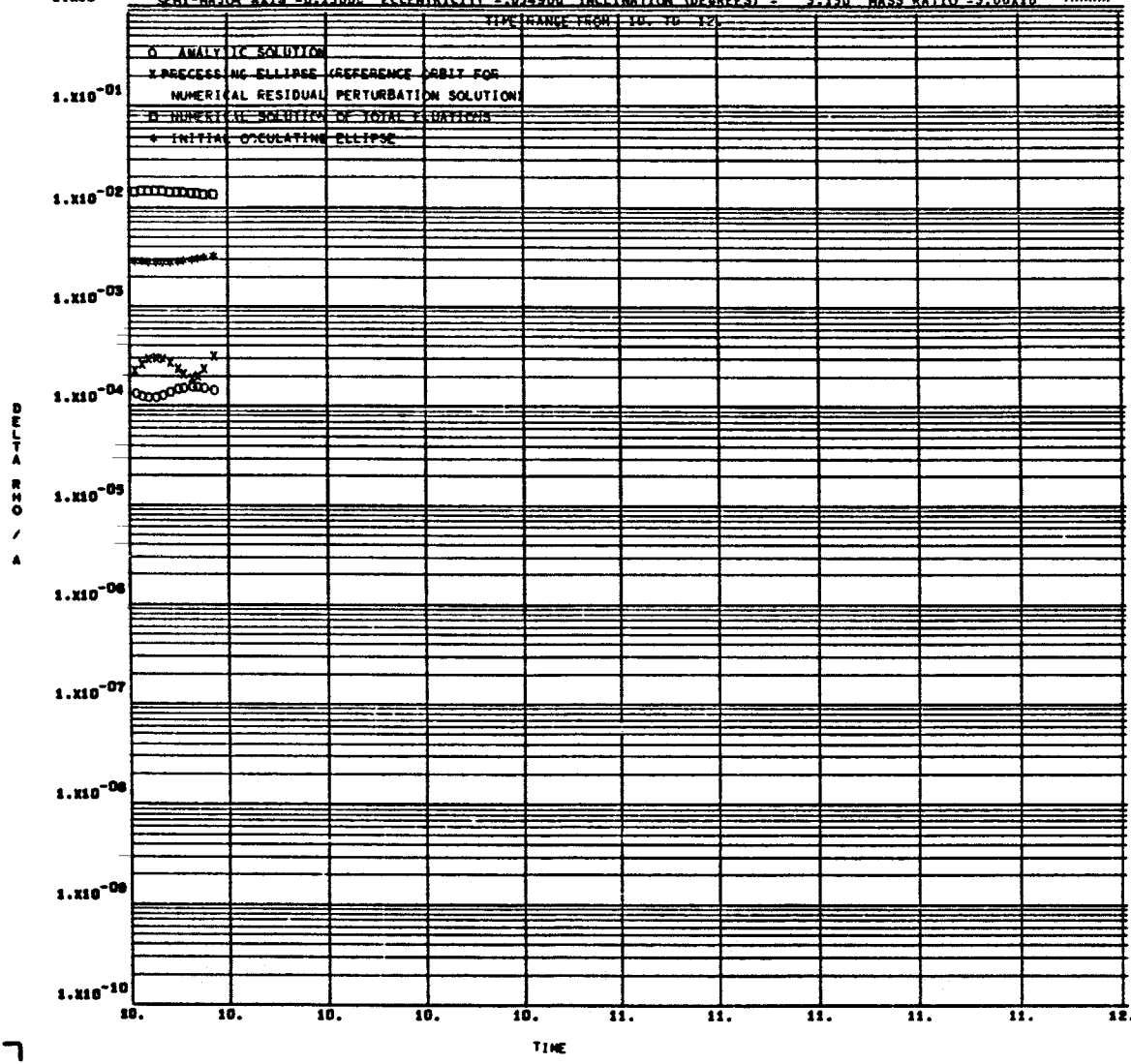


ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

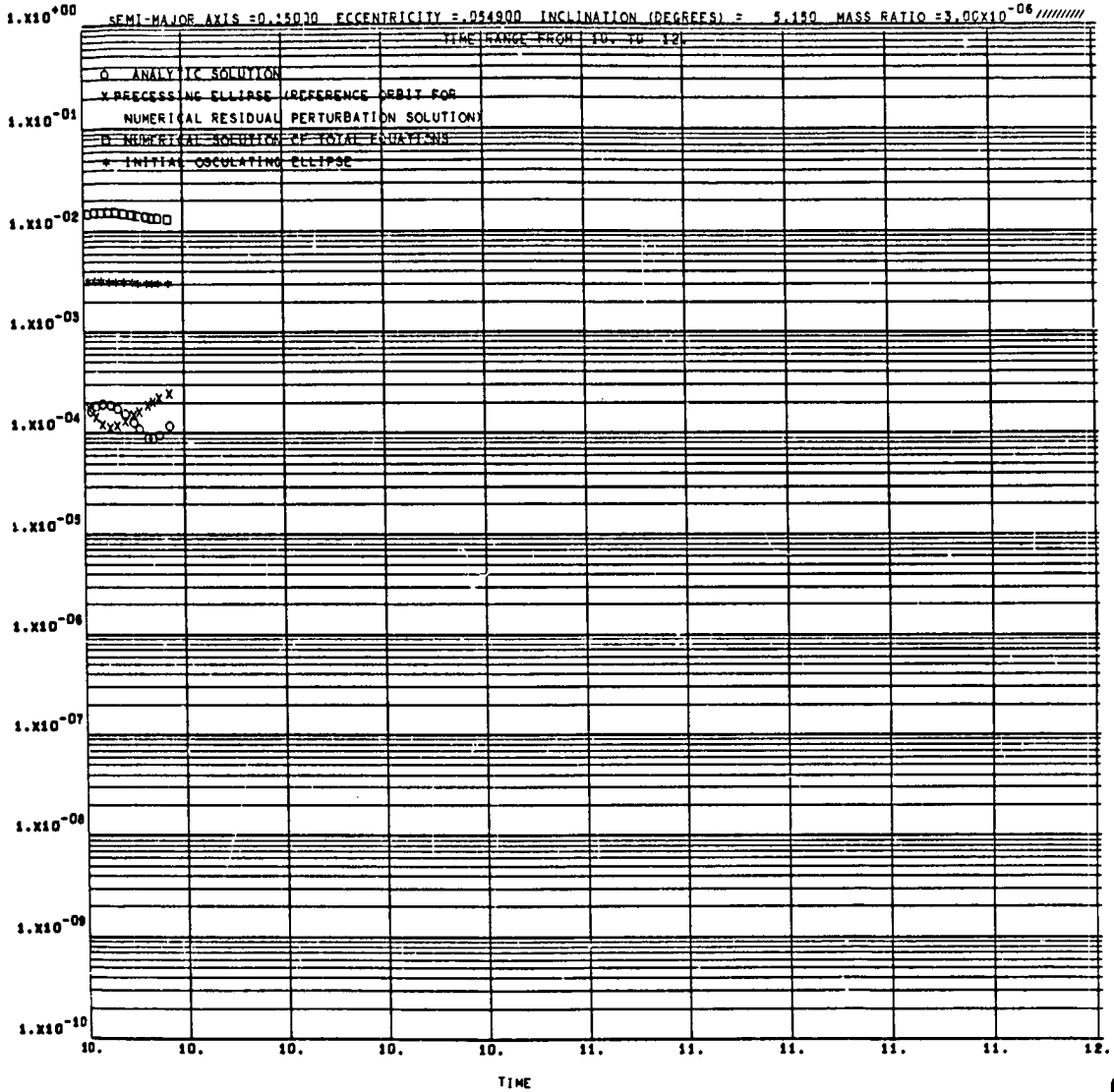


NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

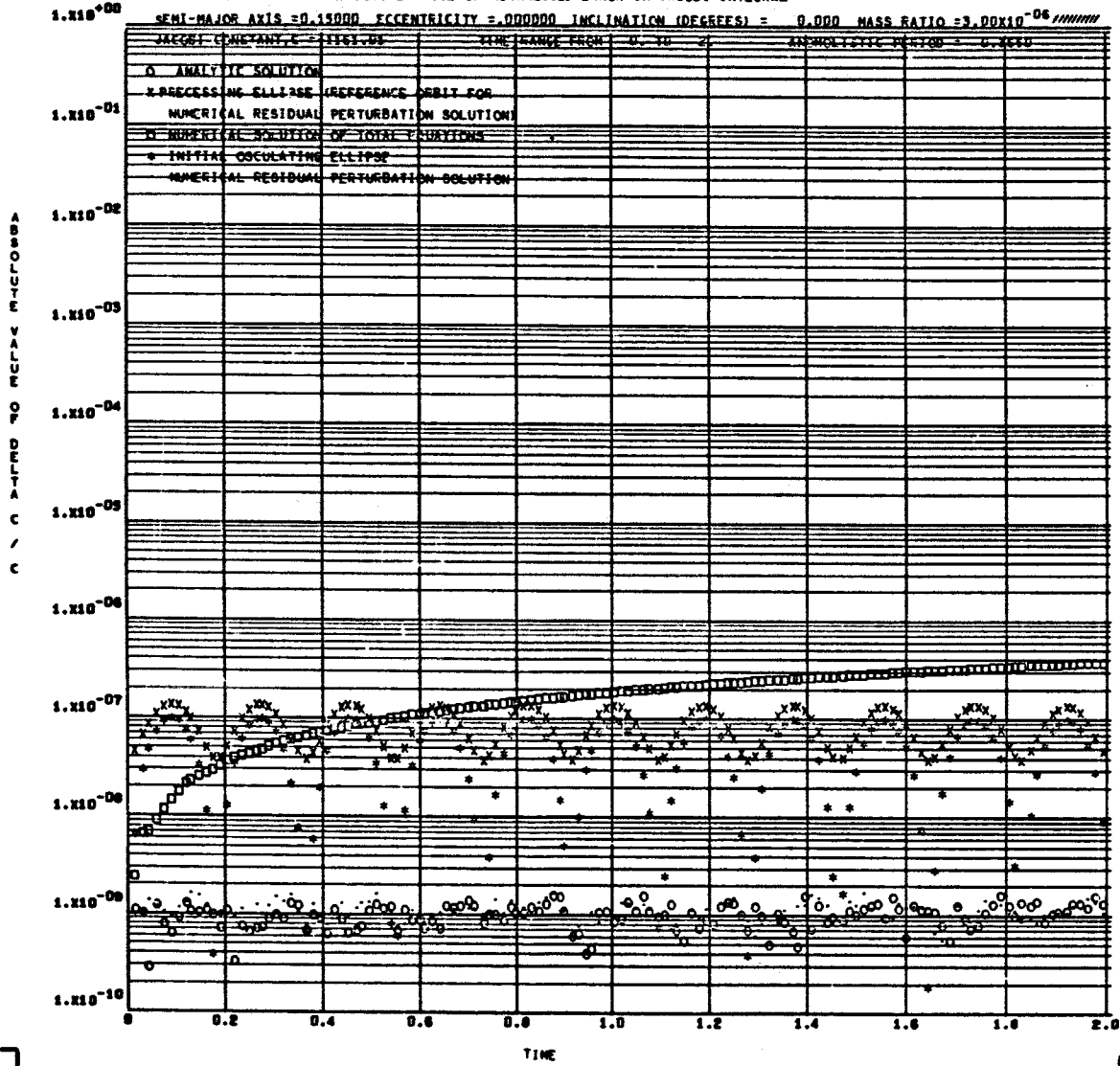
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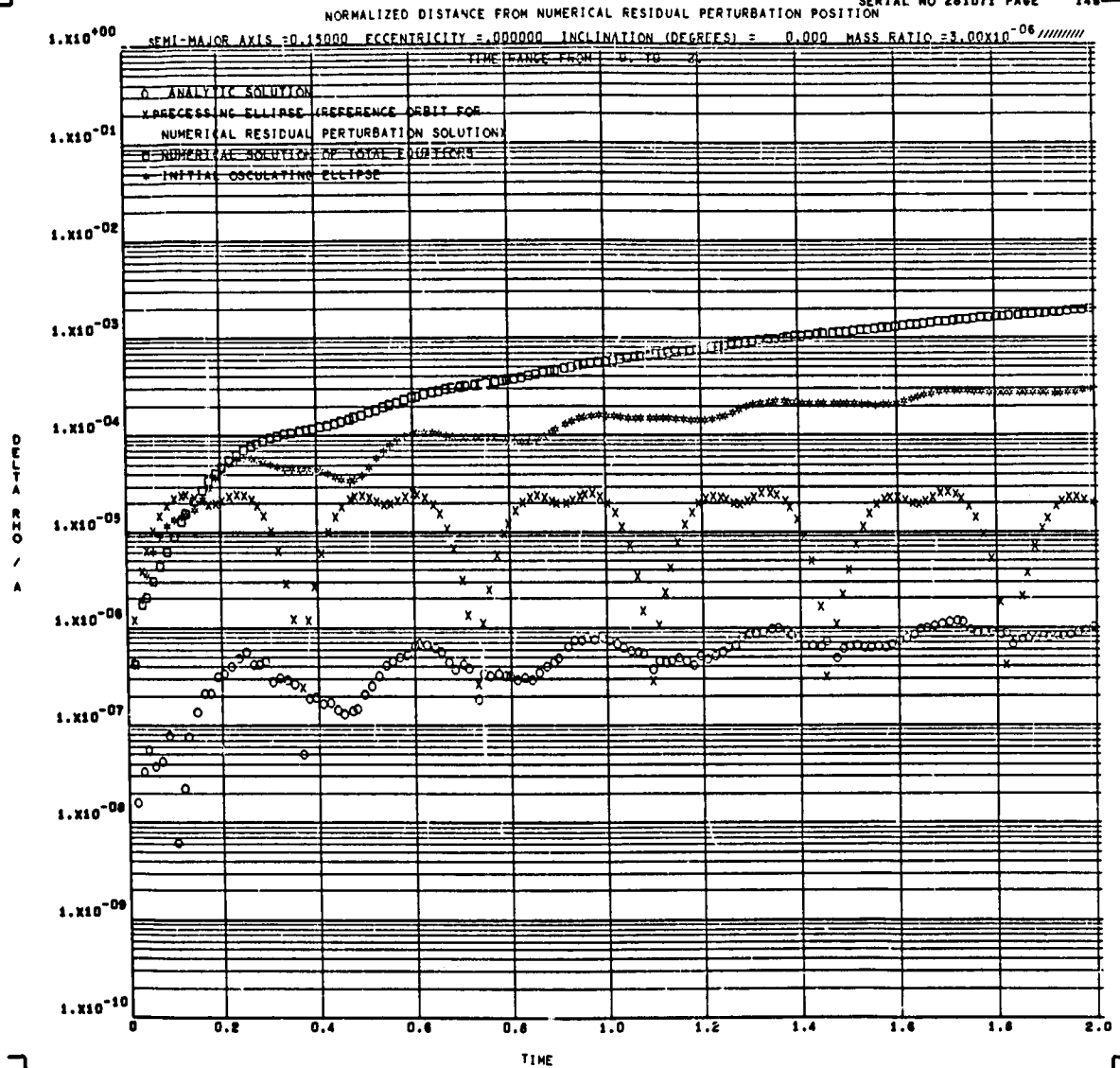


NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL





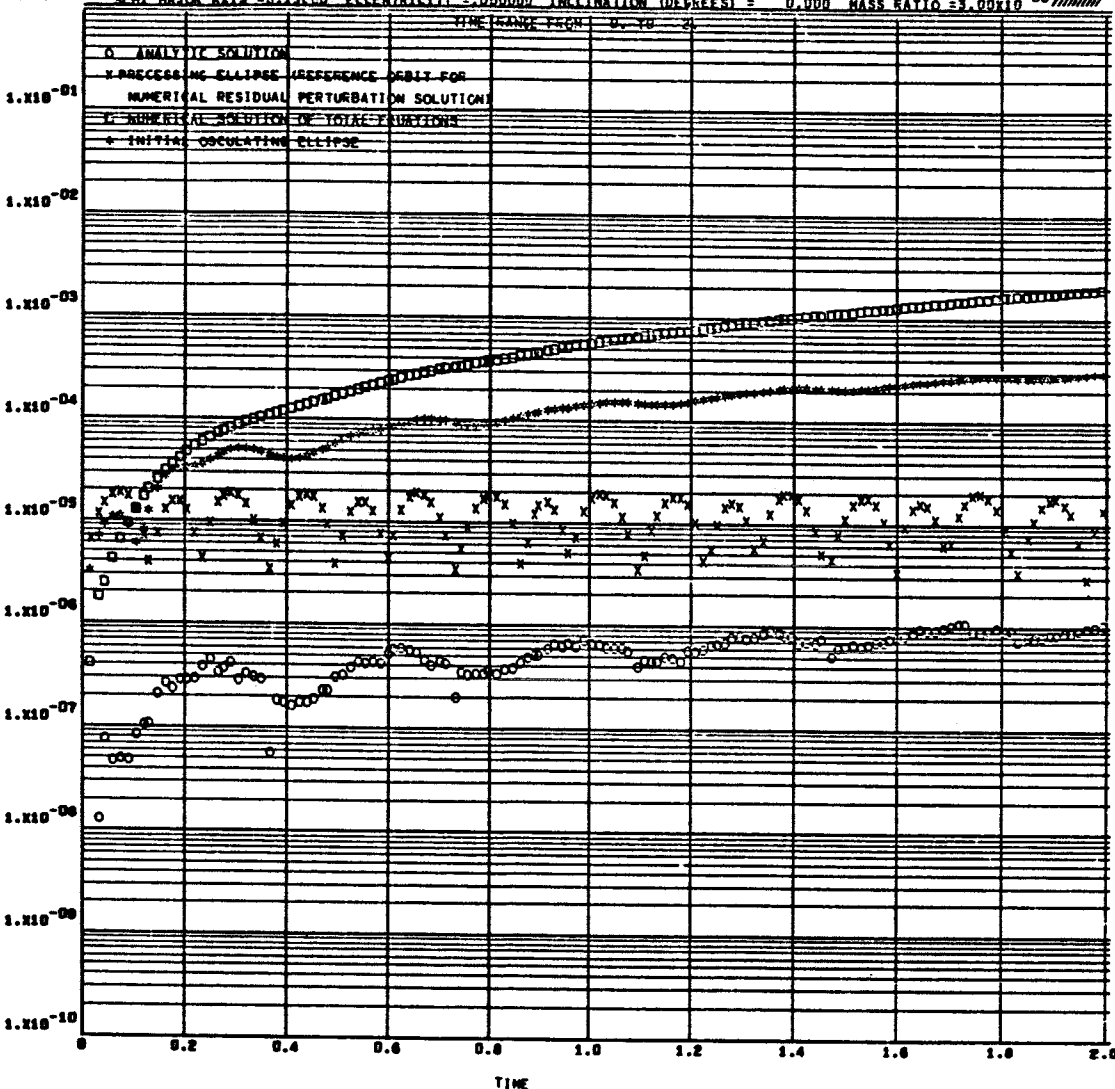


NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

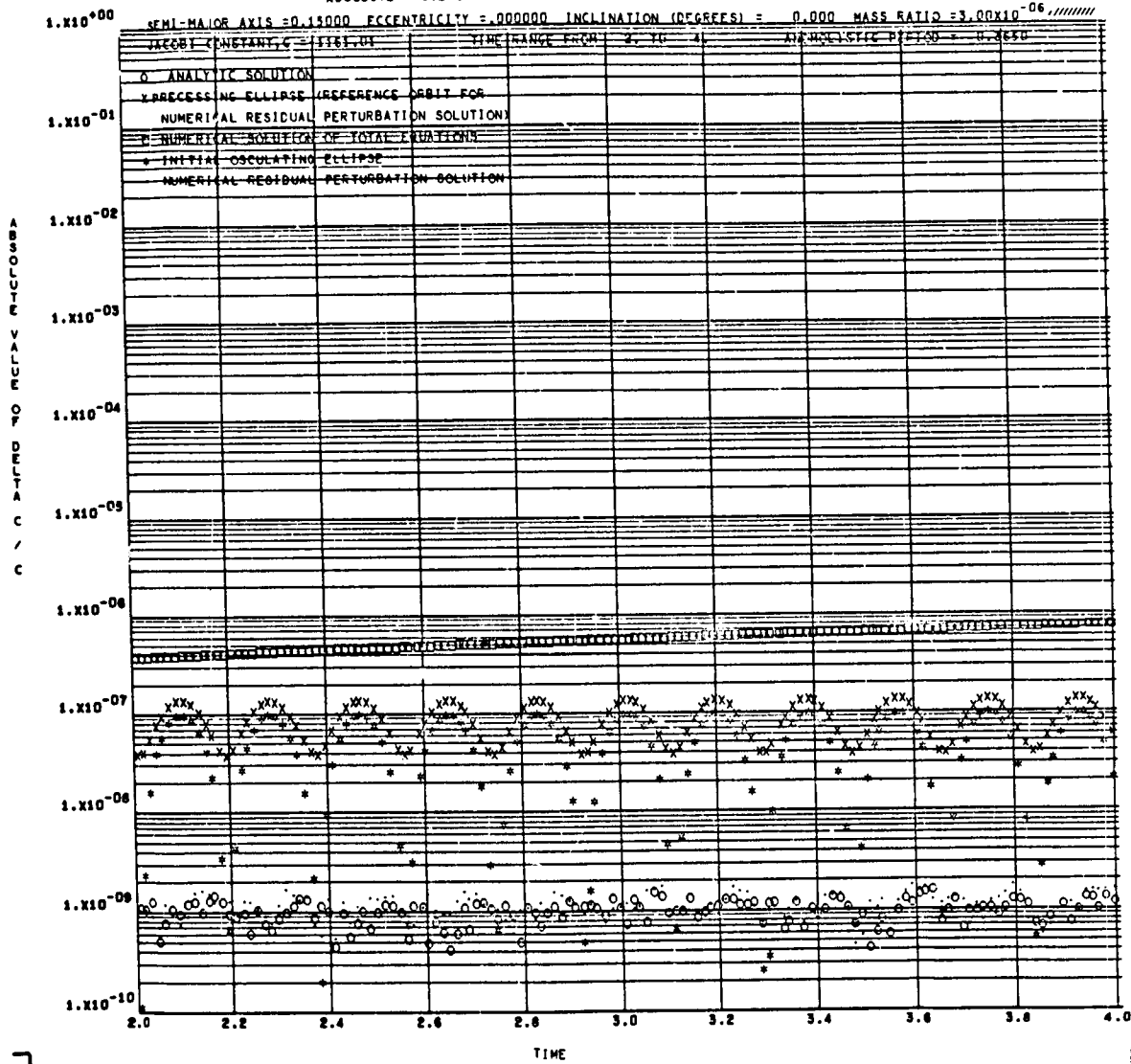
SERIAL NO 201071 PAGE 147

SEMI-MAJOR AXIS = 6.15E00 ECCENTRICITY = 0.00000 INCLINATION (DEGREES) = 0.000 MASS RATIO = 3.00E-06

DELTA RHO DOT TIMES S E T A

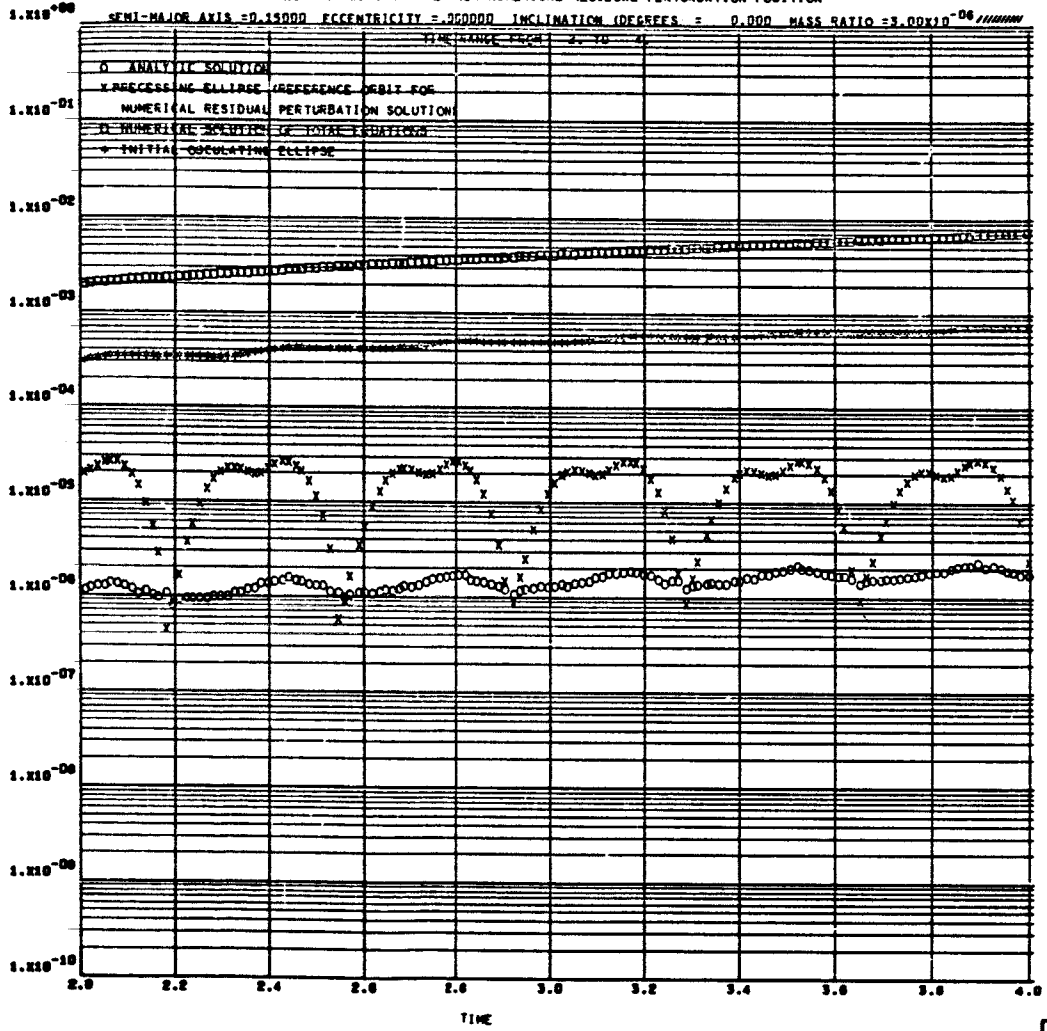


ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

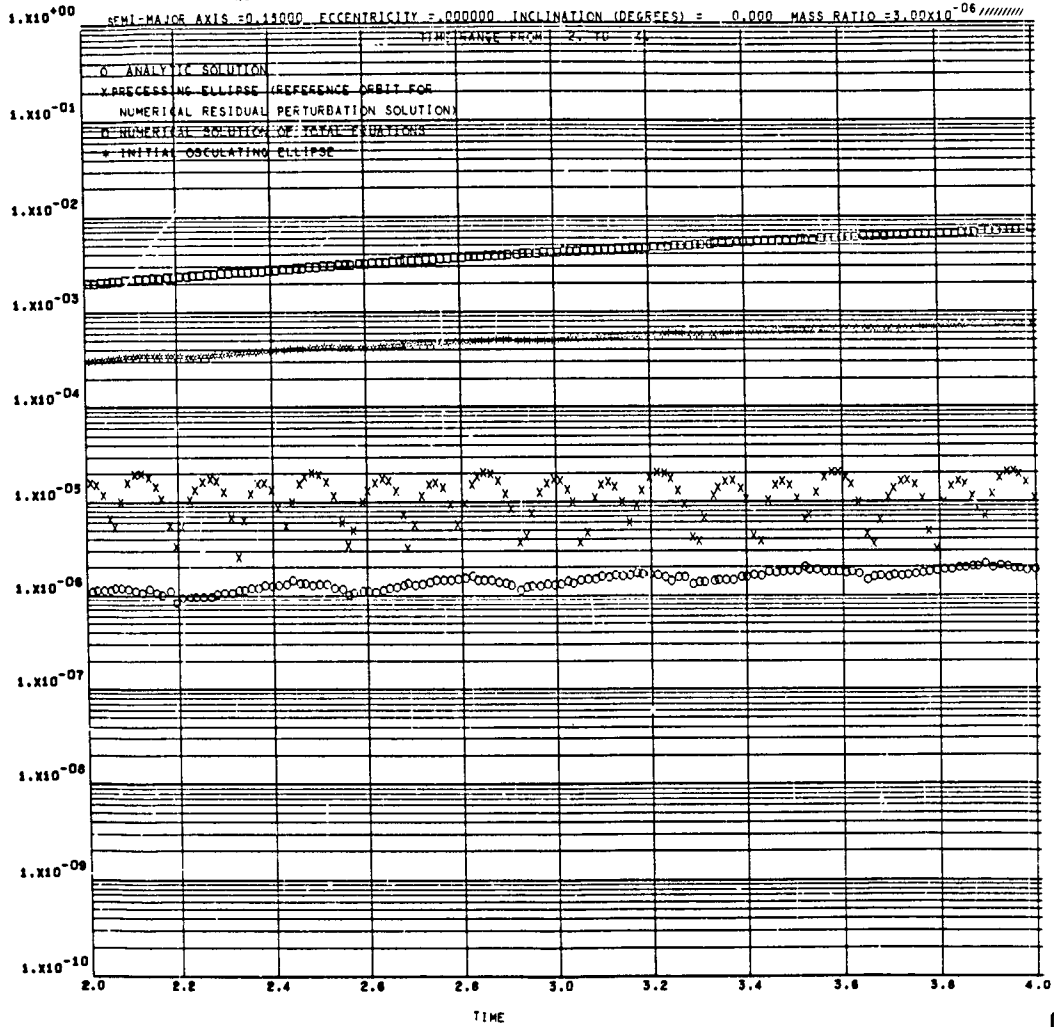


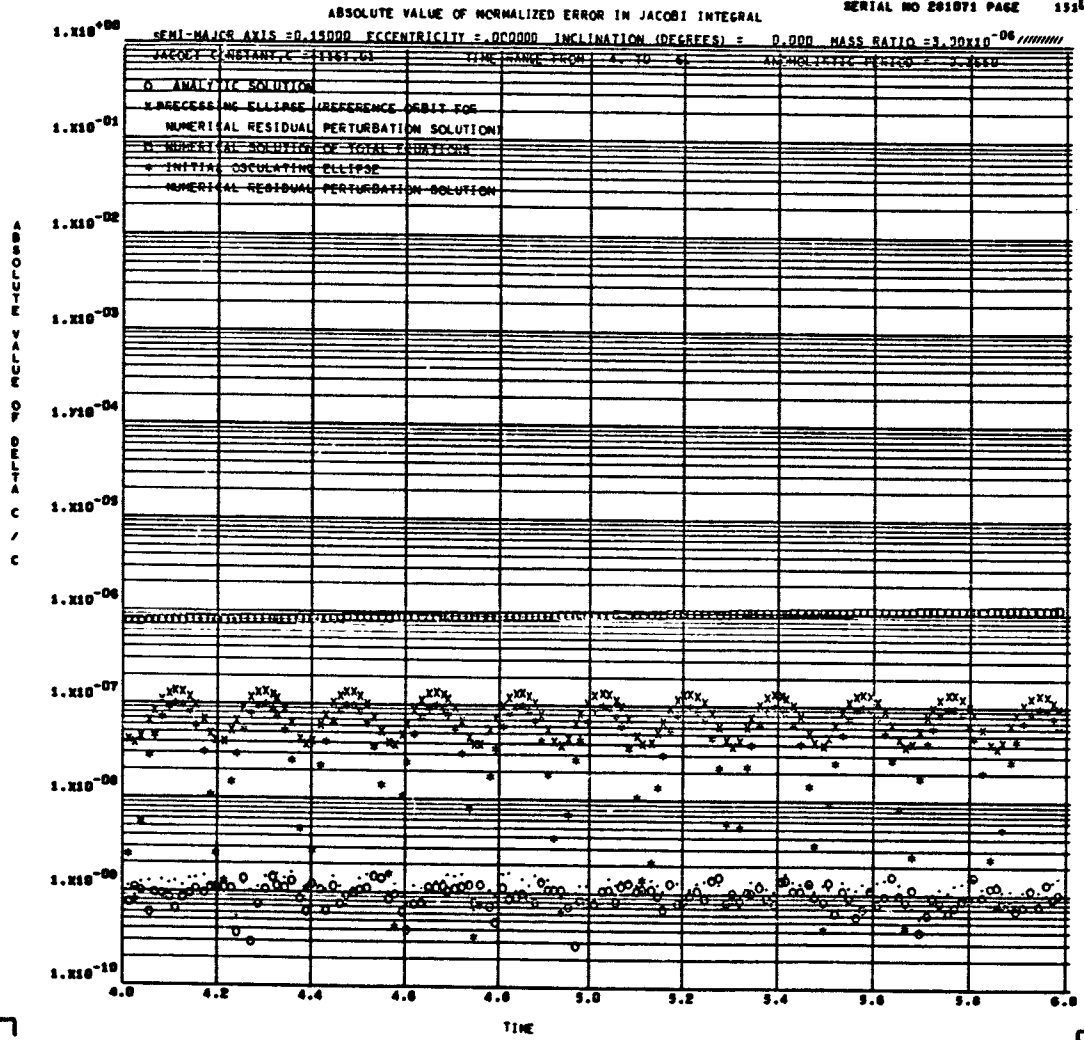
NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

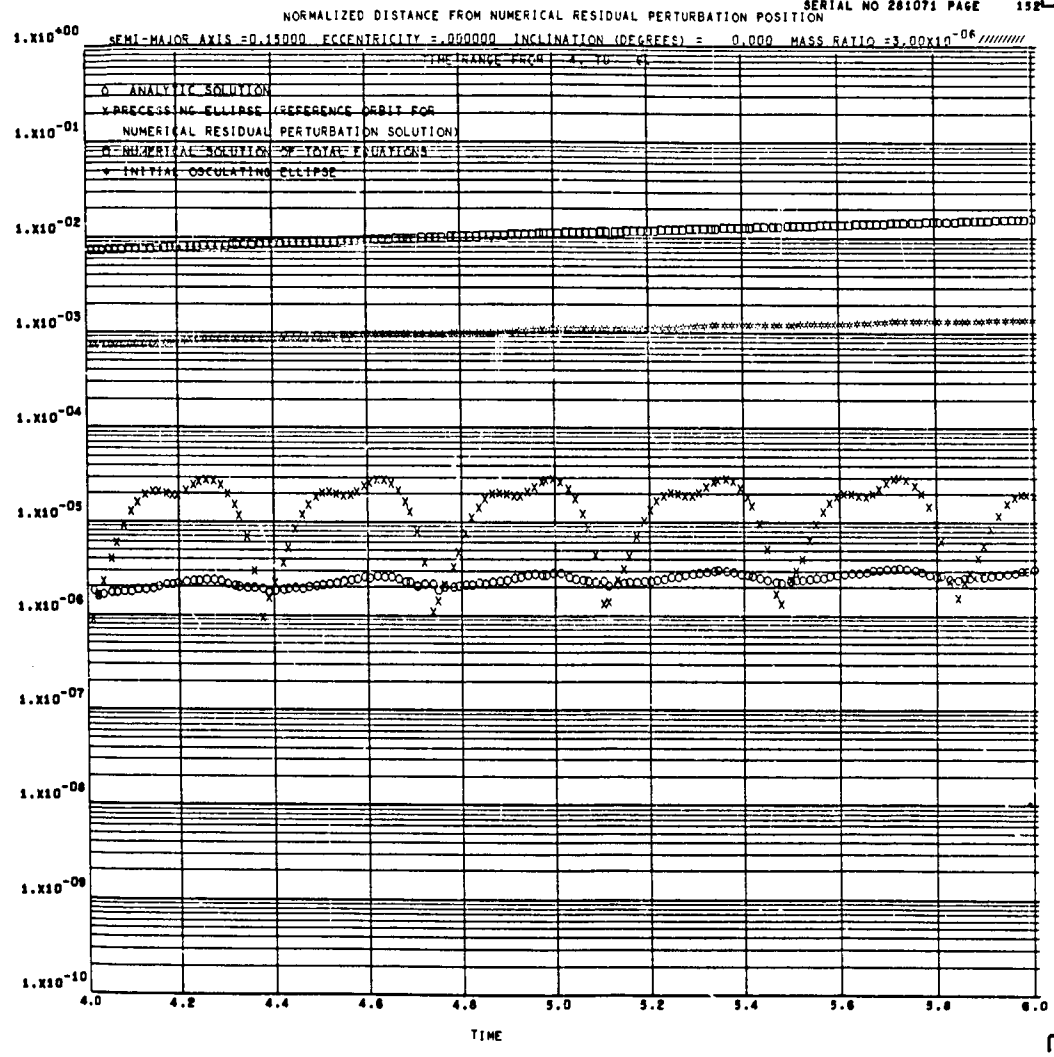
SERIAL NO 201071 PAGE 100



NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

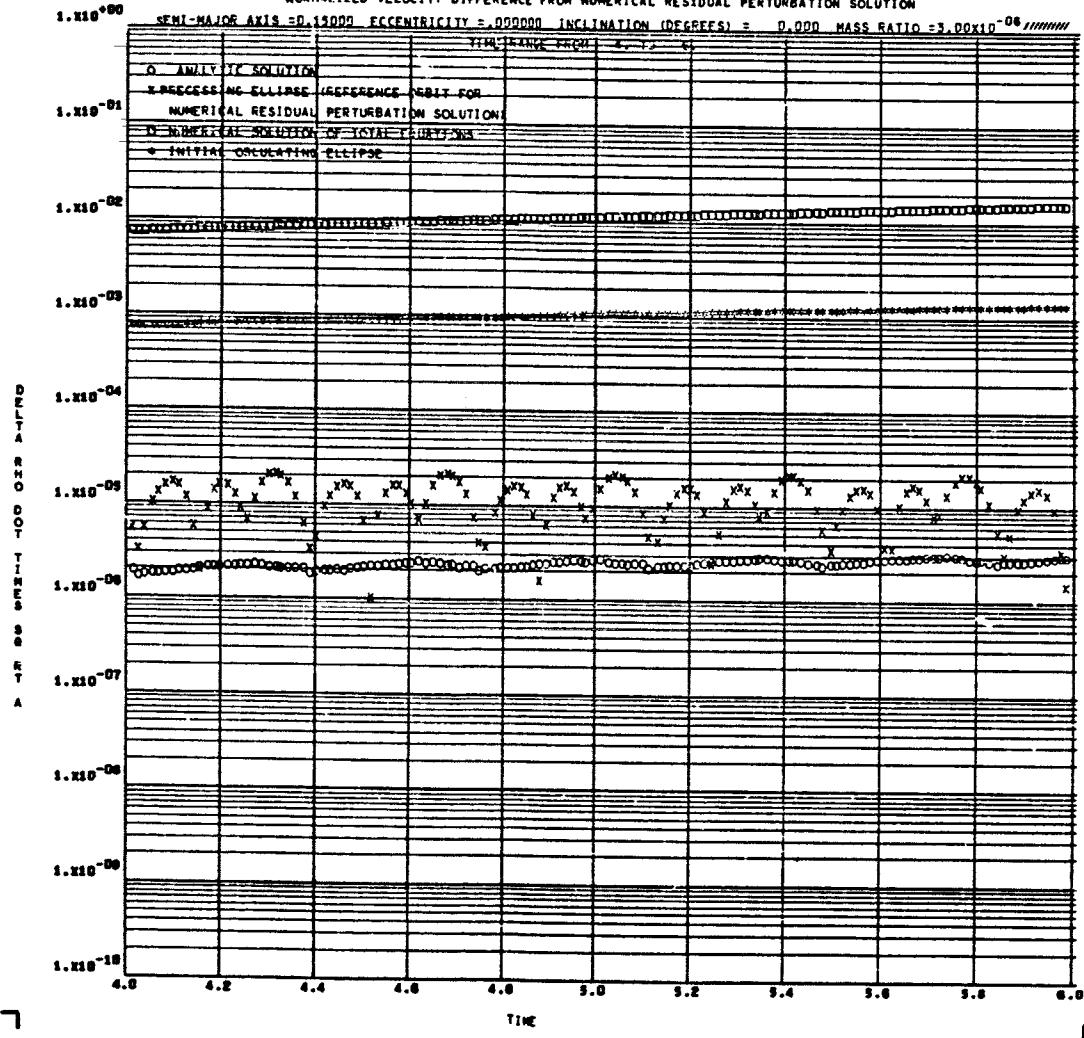


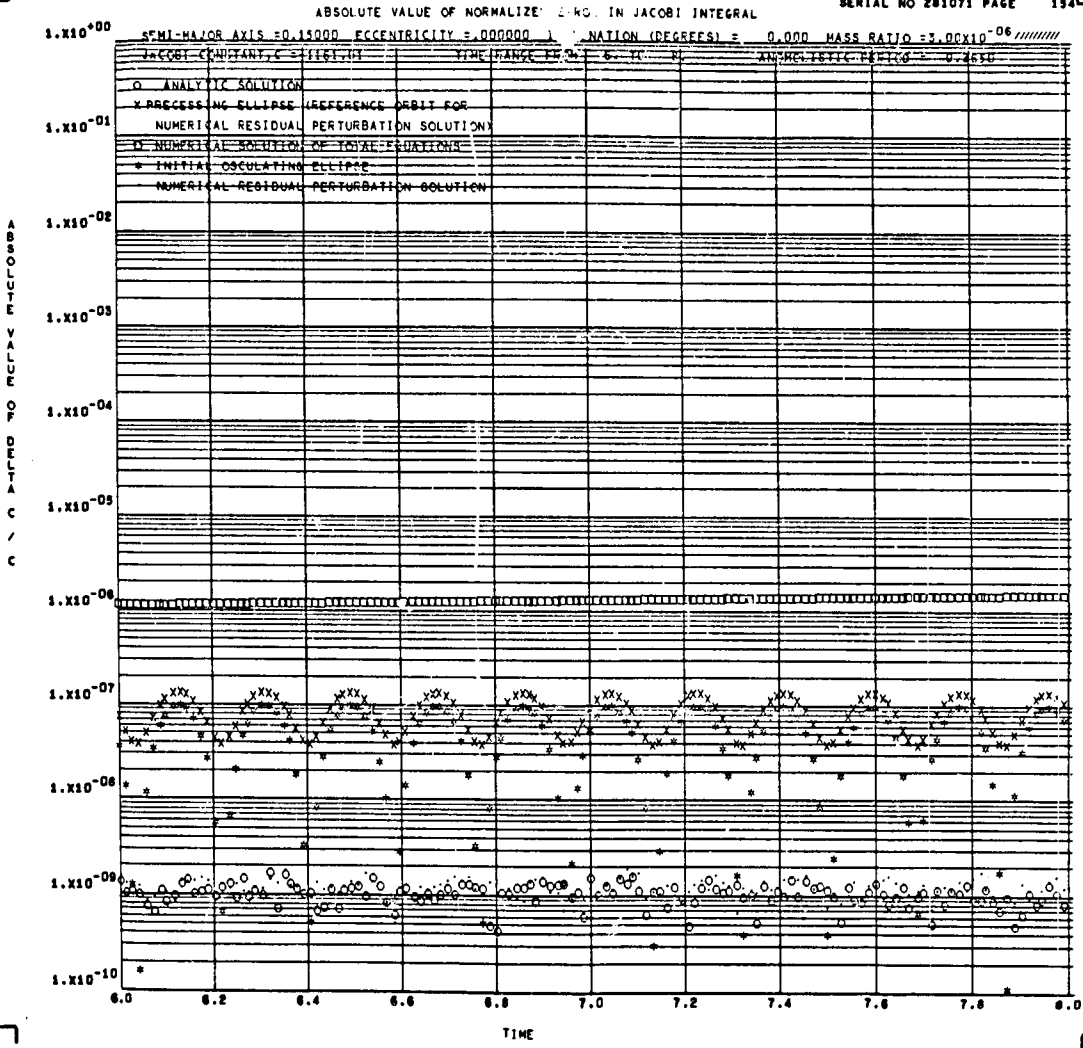




NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

SEMI-MAJOR AXIS = 0.15000 ECCENTRICITY = 0.00000 INCLINATION (DEGREES) = 0.000 MASS RATIO = 5.00X10<sup>-06</sup> //

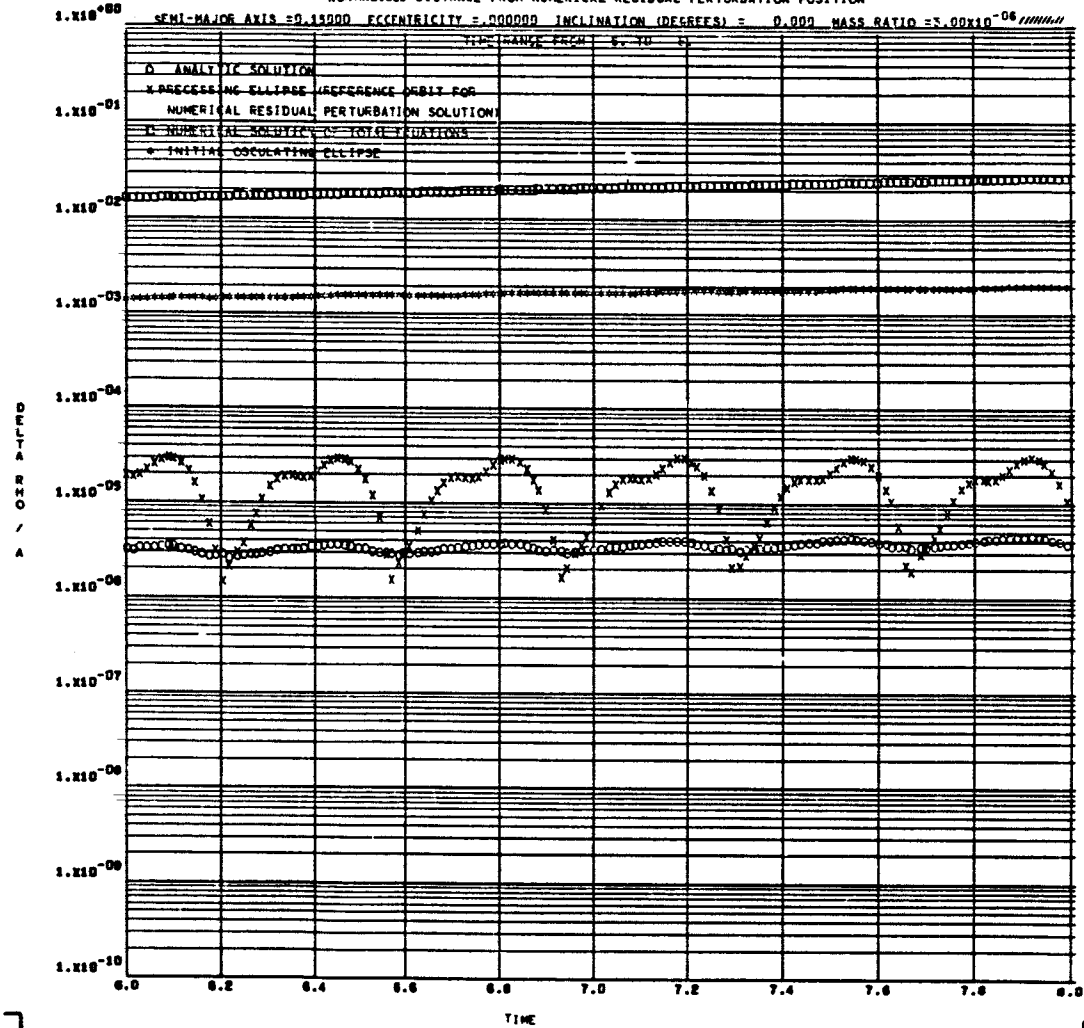




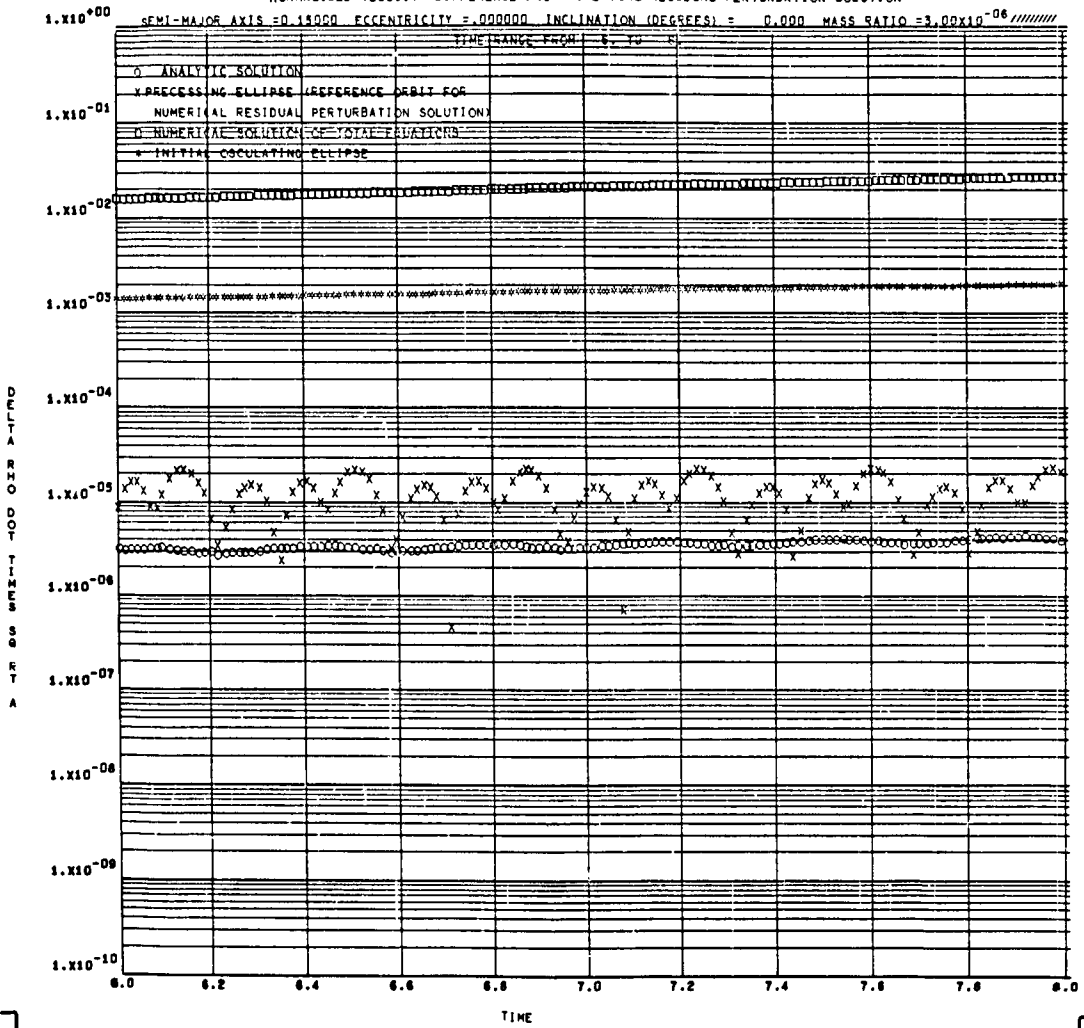


NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

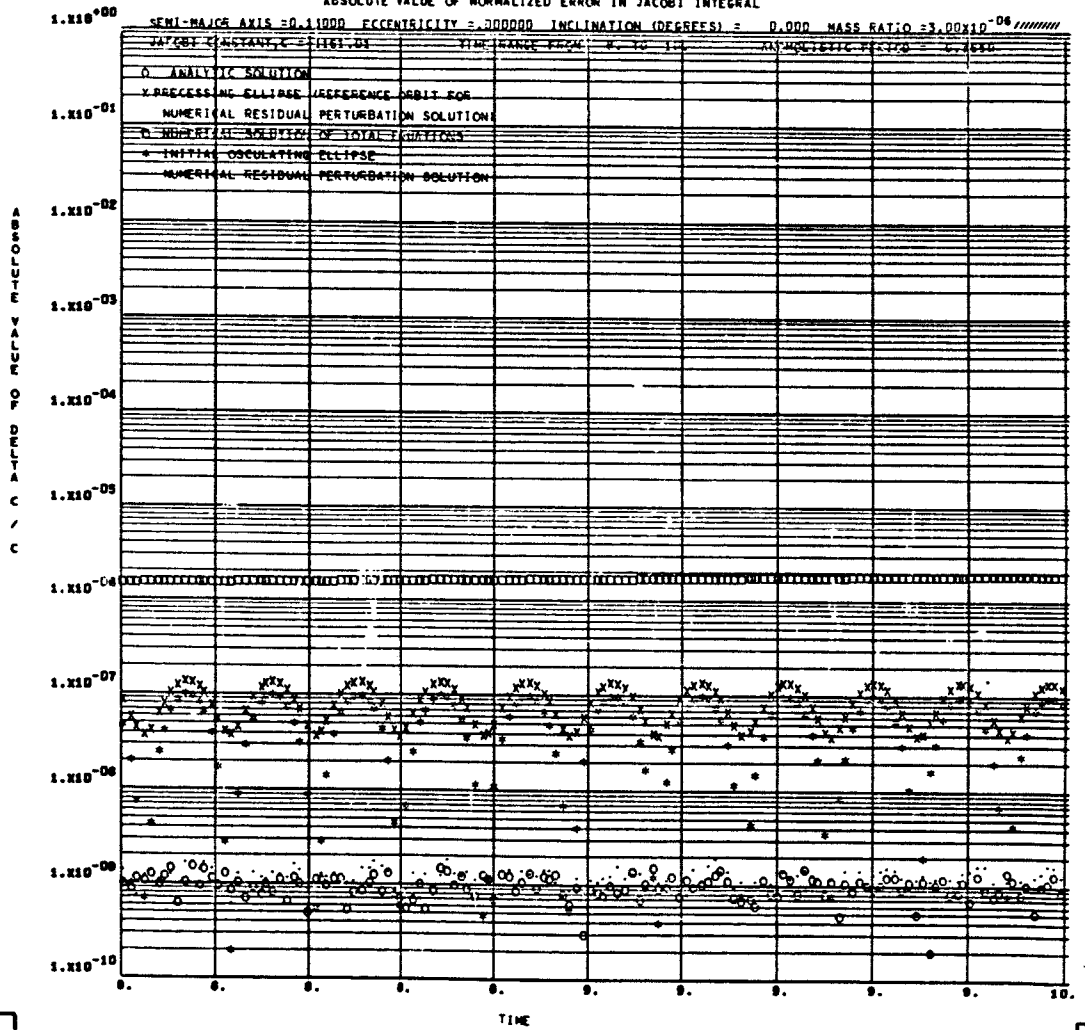
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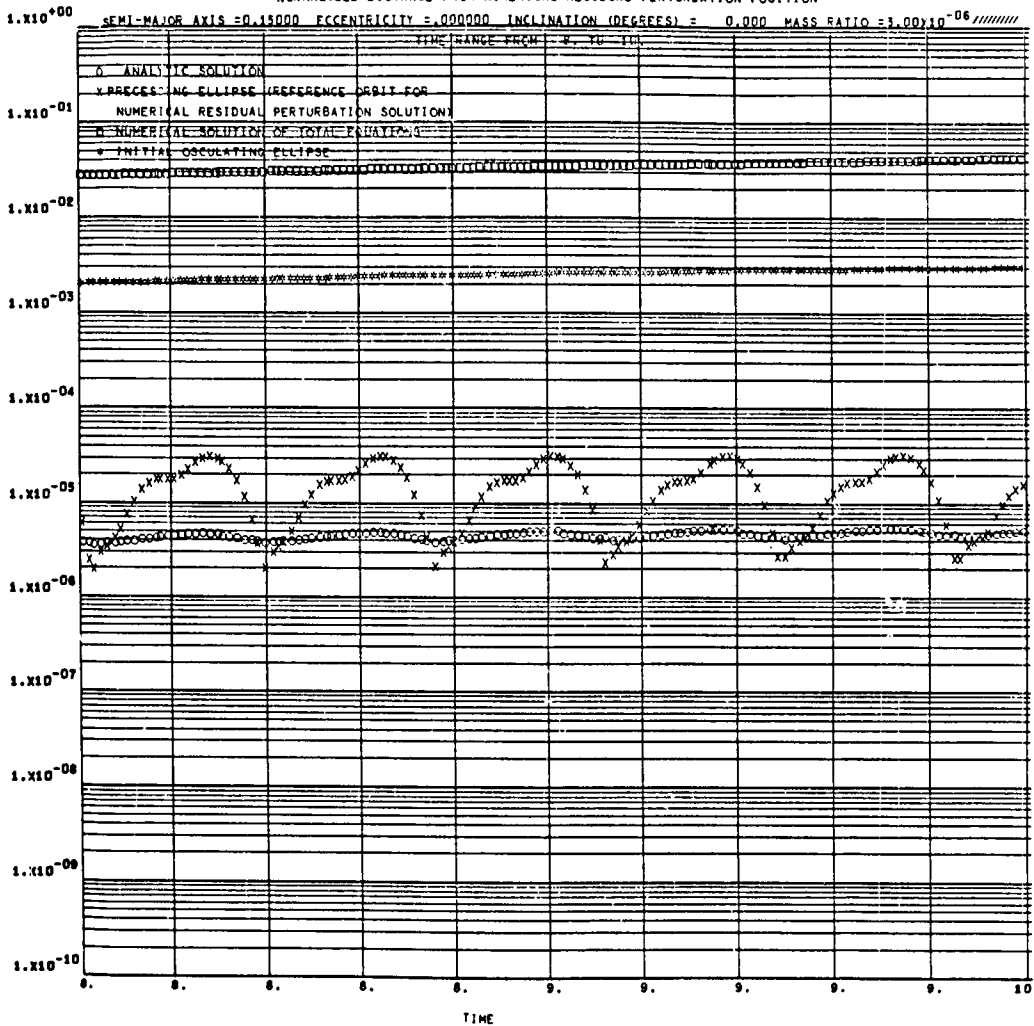
NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



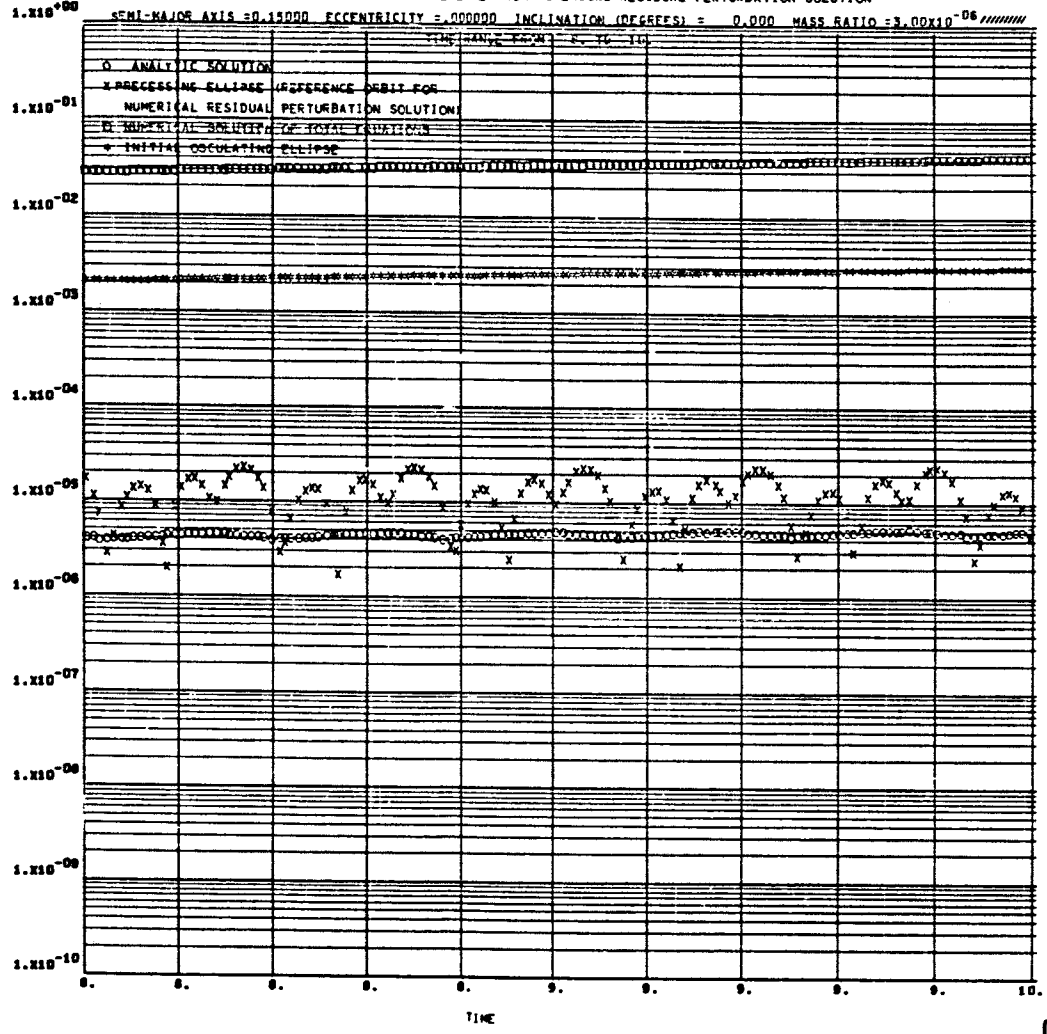
ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL



NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

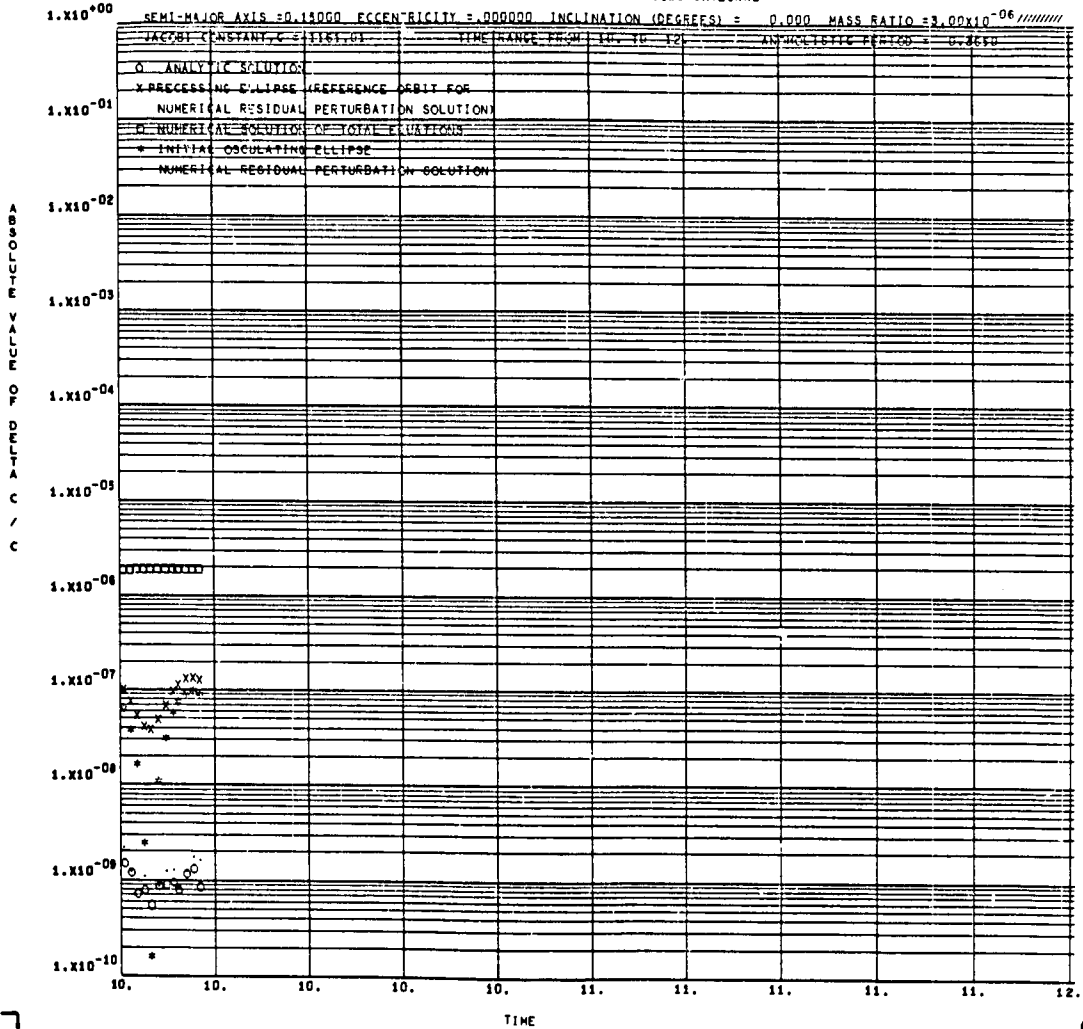


NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

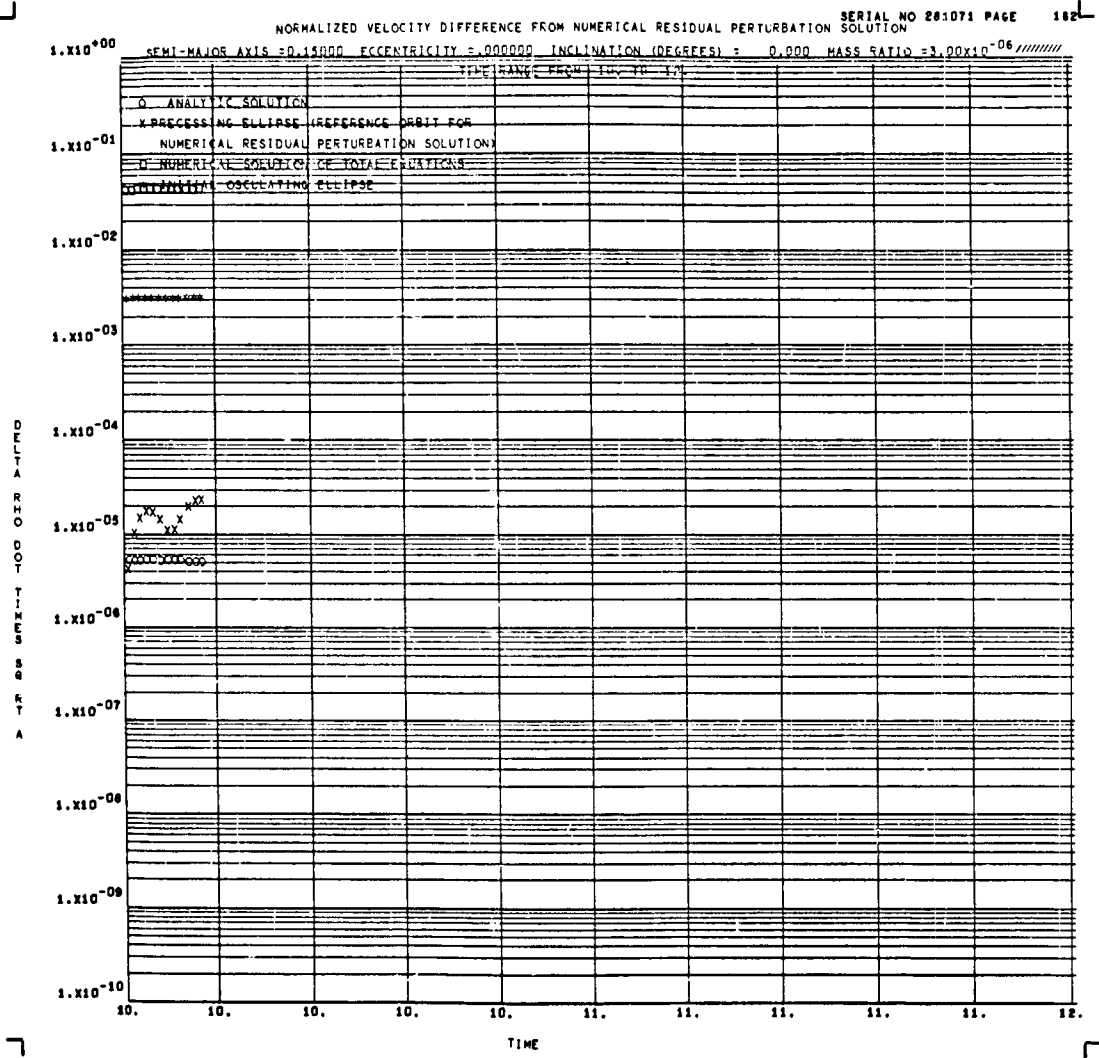


ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

SERIAL NO 281071 PAGE 160





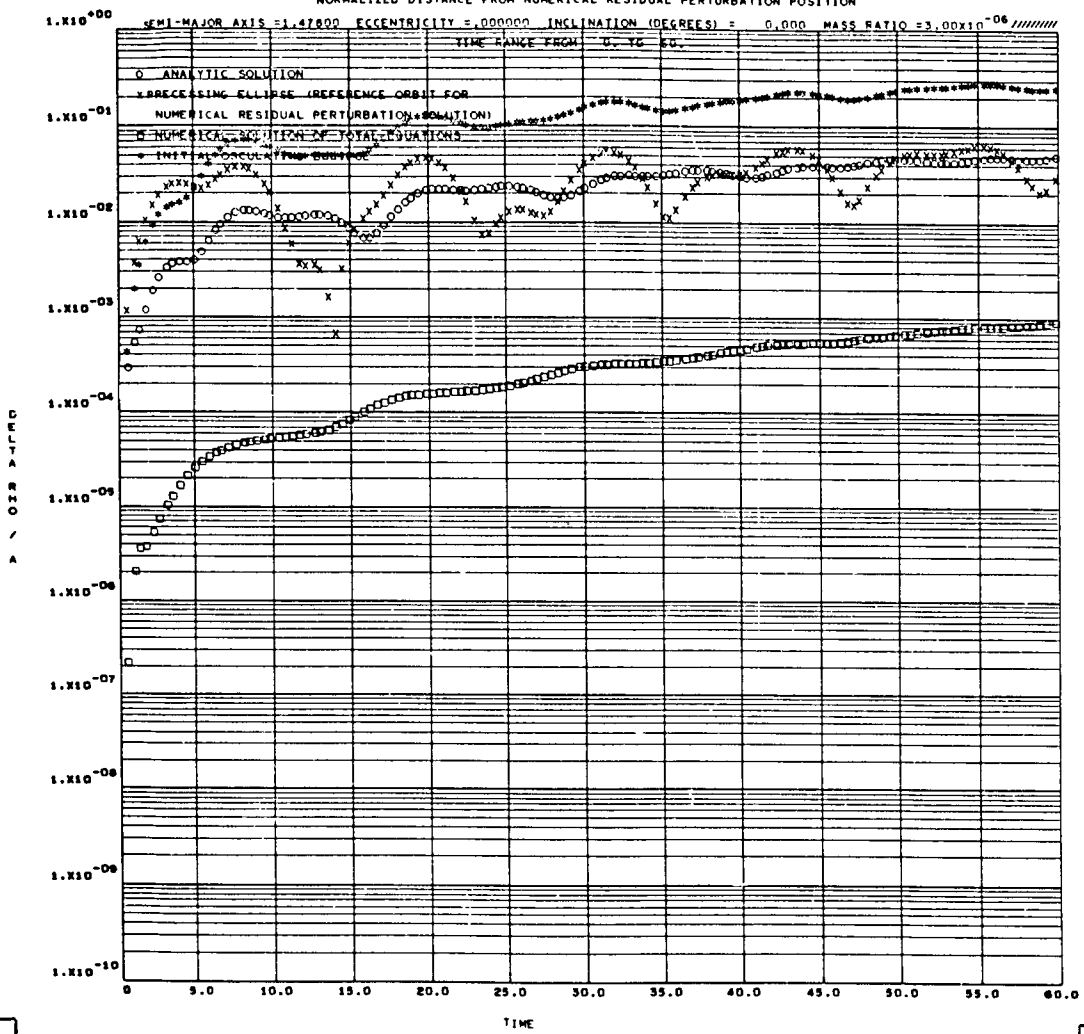






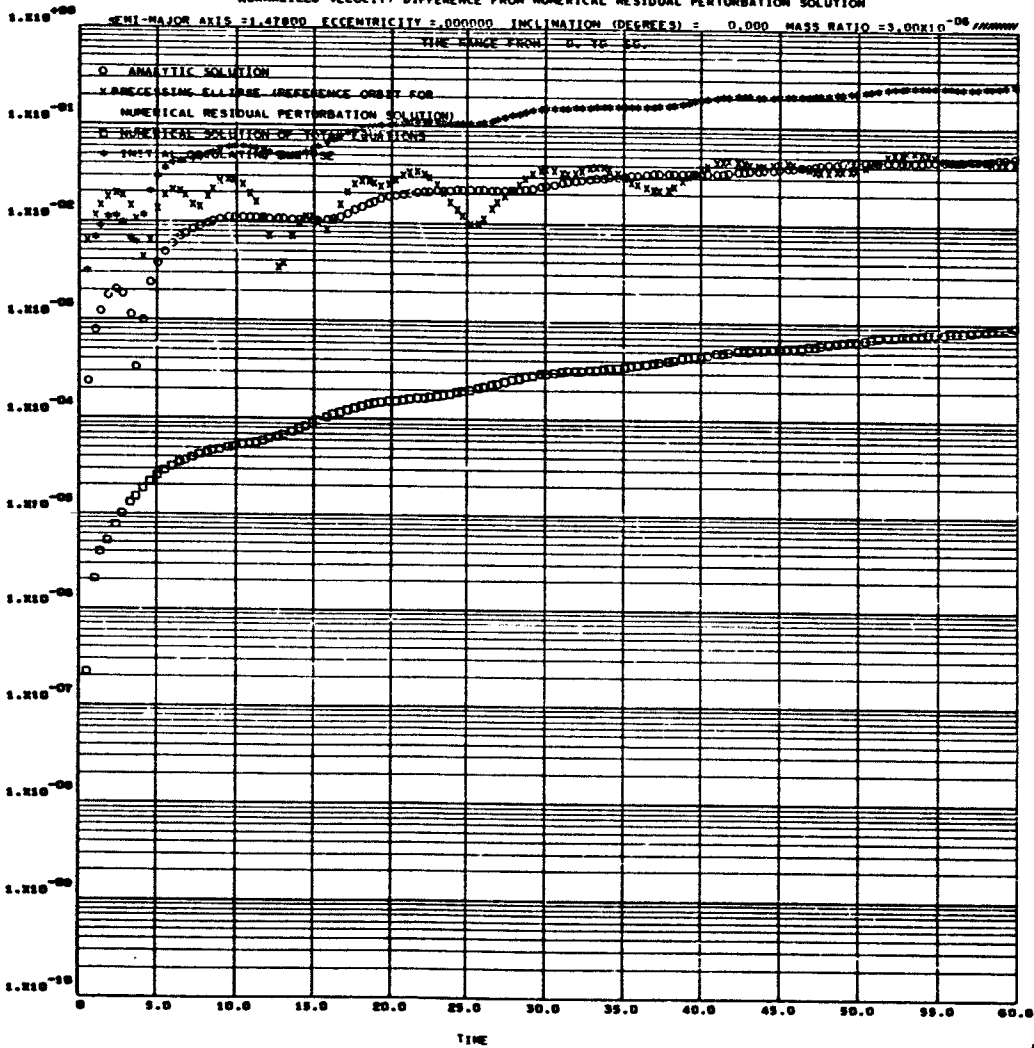
NORMALIZED DISTANCE FROM NUMERICAL RESIDUAL PERTURBATION POSITION

SERIAL NO 302074 PAGE



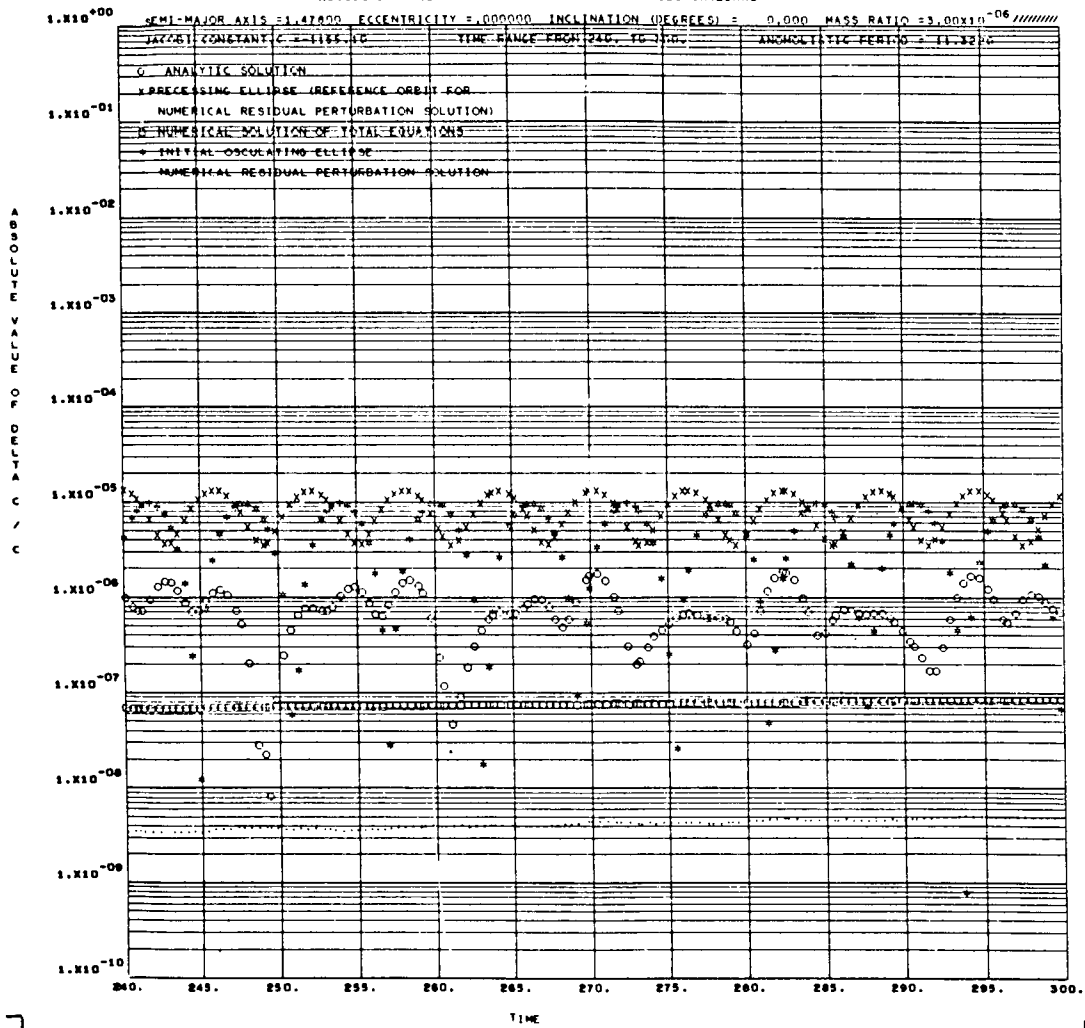
NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

SERIAL NO 302074 PAGE



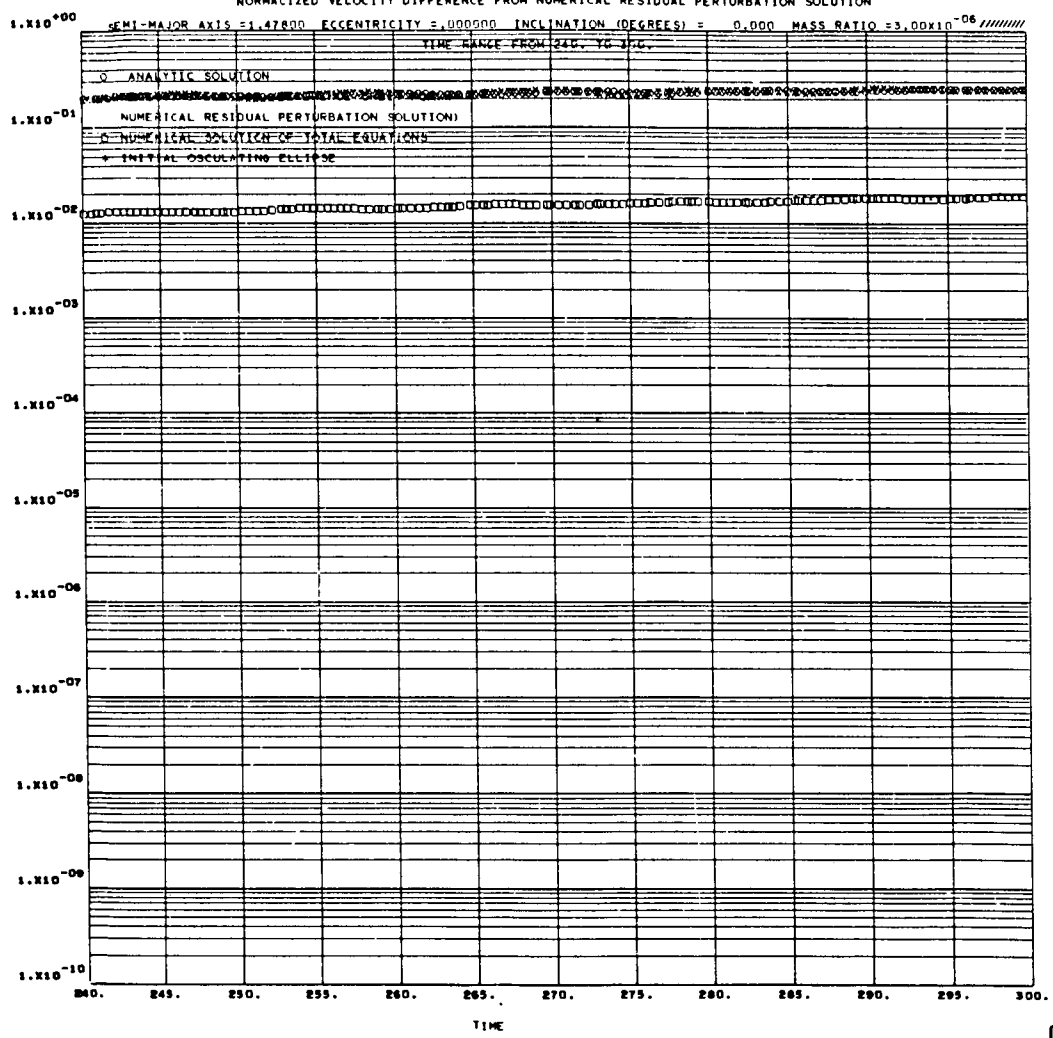
ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

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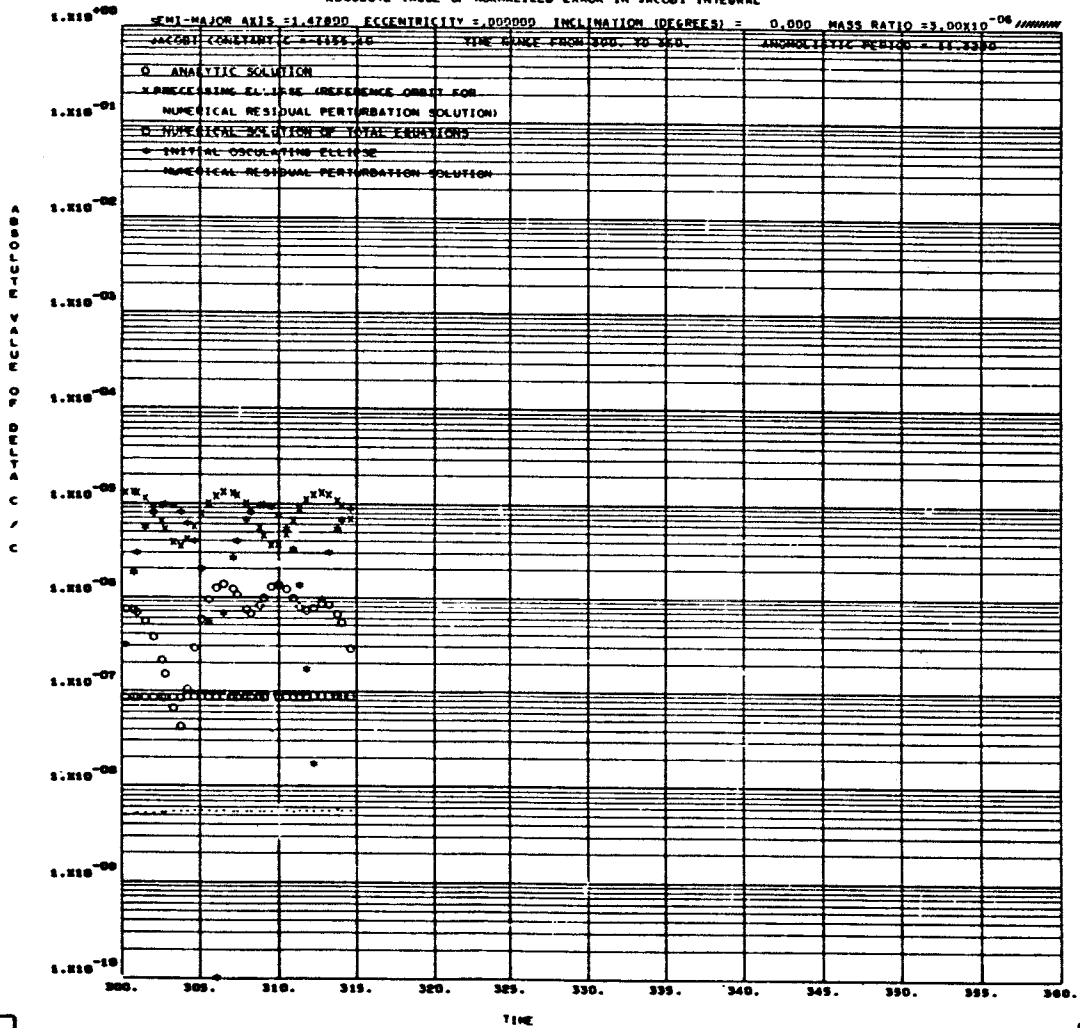


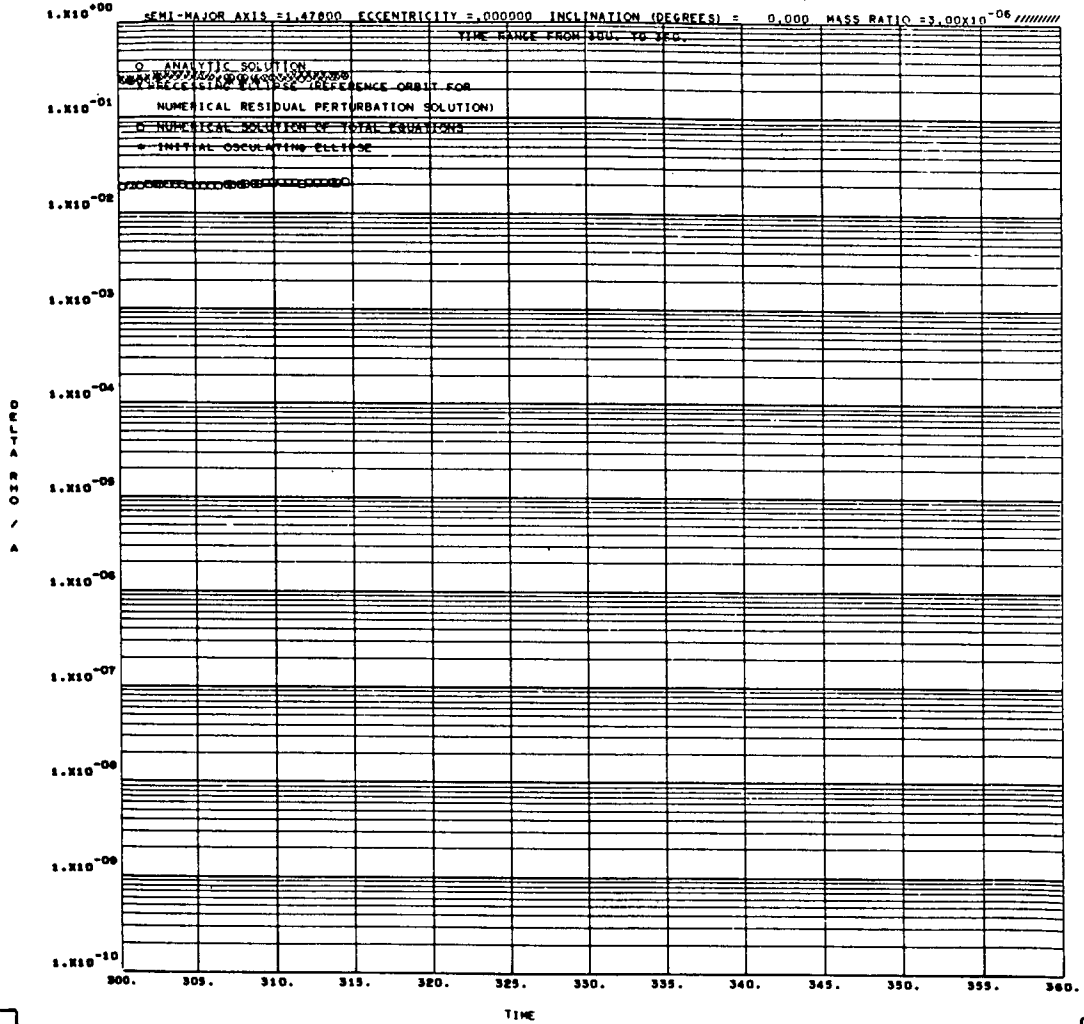
NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION



ABSOLUTE VALUE OF NORMALIZED ERROR IN JACOBI INTEGRAL

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NORMALIZED VELOCITY DIFFERENCE FROM NUMERICAL RESIDUAL PERTURBATION SOLUTION

