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#### SUMMARY REPORT

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#### GAMMA SPECTRAL DATA FOR SHIELDING AND HEATING CALCULATIONS

by

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#### ABSTRACT

Spectra of gamma rays following neutron absorption and inelastic-scattering events in H, Be, C, O, Al, Cr; Fe, Ni, Zr, natural W and the W isotopes, U<sup>235</sup> and  $U^{238}$  are tabulated. The gamma thermal capture spectra for Cd and Sm are presented. A detailed study of the prompt and delayed fission gammas (both intensities and time variations) is also given. Descriptions are given of the sources of information and the calculations performed. In addition, an evaluation of the re-Anthes liability of the data is given.

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#### 1. SUMMARY

The United Nuclear Corporation is presently modifying its UNC-SAM Monte Carlo radiation transport system of programs<sup>\*</sup> for use in the tungsten nuclear rocket program. The modified system and its main tracking routine will be called ATHENA (Attenuation, Tracking, and Heating for NASA). In association with this program, gamma heating studies in the core are being performed for which gamma spectral data are necessary. Therefore, a detailed study was undertaken to determine the gamma spectra and intensities following neutron absorption and inelastic-scattering events in the elements and nuclides of interest – H, Be, C, O, Al, Cr, Fe, Ni, Zr, natural W and the separate W isotopes, U<sup>235</sup>, and U<sup>238</sup>. The capture spectra for the strongly capturing elements Cd and Sm are given. Finally, a detailed study of the prompt and delayed fission gammas was made providing a detailed description of the post-fission sources from 0 to 10 hr.

The thermal capture gamma spectra are generally well-known. The energy and intensity of the discrete spectra are represented by their "best" experimental values. Continuous gamma spectra are represented by a sufficient number of discrete gamma energies to represent the spectrum adequately. Wherever discrete gamma lines of relatively high intensity are imposed on a continuum, the gammas of energy close to the discrete line have been lumped together with the discrete gamma. This procedure was also followed for the inelastic gamma spectra.

<sup>\*</sup>For a description of the UNC-SAM code see Reference 1.

It was assumed that the capture spectrum is independent of the neutron energy at which the capture occurs. This assumption is generally adequate for most problems. The assumption that no gammas follow the charged-particle reactions (except for oxygen) will tend to cause an underestimate of the total heating.

The gamma spectra following inelastic scattering, for neutron energies below  $\sim 4$  Mev, are believed to be adequate, since they are based on experimental levelexcitation cross sections. In the intermediate neutron energy range (4 to 8 Mev) the inelastic gamma spectra, based in part on statistical theory, are not uniformly reliable. For neutron energies above 8 Mev, the spectra are based on statistical theory which includes (n,2n) and (n,3n) processes. As the neutron energy increases above 8 Mev, the validity of the parameters used in the theory becomes increasingly questionable. In general, the shape of the spectra is given more reliably than the absolute values. For the problems in which it is planned to use the data presented in this report the inadequacy is not important because neutrons having energies > 8 Mev constitute only 1/2 of 1% of the fission source.<sup>2</sup>

Inasmuch as the tungsten being used in the reactor may be enriched in one or more isotopes, additional calculations were made to provide capture gamma spectra for the isotopes  $W^{182}$ ,  $W^{183}$ ,  $W^{184}$ , and  $W^{186}$ .

#### 2. GAMMA SPECTRAL DATA

#### 2.1 ABSORPTION AND INELASTIC GAMMA SPECTRA

#### 2.1.1 General Comments on the Absorption Spectra

This section describes a study made of the gamma spectra following absorption and inelastic-scatter events in elements of interest in the tungsten nuclear rocket program. The data are to be used in the modified United Nuclear Monte Carlo three-dimensional transport code to determine, among other quantities, gamma heating in and near the core.

The relevant data are presented in Tables 1 through 26, 32, and 33. Generally, there are two tables per element (each element being identified in the programs by a 5-digit integer). One table gives the absorption, and the other gives the inelastic spectrum. Each table consists of a matrix specifying, for a given neutron energy bin, the number of gammas produced per event (absorption or inelastic) for each of several discrete gamma energies.\* The neutron energies given represent the upper energy of the bin. The spectra are assumed to be constant within each neutron energy group.

The absorption spectrum for each element (except for Cd and Sm) includes both capture and charged-particle events, the spectrum of each being weighted by its

<sup>\*</sup>For several of the elements the gamma energies were combined so as to condense the tables; the complete, more detailed spectra are available in punchedcard format for use in the ATHENA system.

average cross section in the given neutron energy group. If  $P_c$  and  $P_{cp}$  are the capture and charged-particle number spectra (number of photons produced) with corresponding average cross sections  $\sigma_c$  and  $\sigma_{cp}$ , then the absorption number spectrum  $P_a$  for that neutron energy bin  $E_n$  is

$$\mathbf{P}_{a}(\mathbf{E}_{n}, \mathbf{E}_{\gamma}) = \frac{\sigma_{c}(\mathbf{E}_{n})\mathbf{P}_{c}(\mathbf{E}_{\gamma}) + \sigma_{cp}(\mathbf{E}_{n})\mathbf{P}_{cp}(\mathbf{E}_{n}, \mathbf{E}_{\gamma})}{\sigma_{c}(\mathbf{E}_{n}) + \sigma_{cp}(\mathbf{E}_{n})}$$
(1)

For most of the elements (the exception being oxygen) it was assumed that  $P_{cp}=0$ , i.e., no gammas are created following a charged-particle reaction. In general, charged-particle reactions produce low-energy (<1 Mev) gammas. Thus they can be neglected in shielding problems where the desired quantity is the gamma-ray dose at the outside of the shield. For problems in which one desires the total local gamma heating, neglect of gammas from charged-particle reactions will underestimate the total heating. A saving factor is that, in general, charged-particle reactions are important only at high (several Mev) neutron energies. Hence for a fission source where neutrons with E > 6 Mev represent only 2.5% of the source, the fraction becoming 0.5% for E > 8 Mev, the underestimate of the heating will be small. Moreover, at high neutron energies, in addition to the charged-particle reactions, inelastic-scattering events (with  $\sigma_{inel} \gtrsim \sigma_{cp}$  generally) yield gamma rays, the latter being of higher energies than those from the charged-particle reactions.

Eq. 1 implies that the capture spectrum is independent of the neutron energy at which the capture occurs. This is not generally valid for resonance regions in which different resonances may excite different levels. However, for a continuous neutron energy distribution the composite thermal spectrum is adequate. The possible variation of the capture spectrum with neutron energy, and its consequences for any given element, could be further investigated. The thermal capture spectrum represents high-energy (several Mev) captures even less adequately than in the resonance region. This does not represent any difficulty,

however, as the capture cross section is generally very small in the Mev range. Moreover, for a fission spectrum, there are relatively few high-energy neutrons.

#### 2.1.2 General Comments on the Inelastic Spectra

The inelastic gamma spectrum for each element includes the (n,2n) and (n,3n)processes in addition to the purely inelastic  $(n,n'\gamma)$  reaction. For neutron energies below ~4 Mev, the values tabulated are based on experimental level-excitation cross sections. In some cases these were supplemented by Hauser-Feshbach calculations. In the intermediate neutron energy range (4 to 8 Mev) the spectra are based on statistical-model calculations, supplemented by some meager experimental data. The data of Perkin<sup>3</sup> for Al, Fe, Ni, and W, though complete, are of dubious quality. He occasionally gives production cross sections for gammas of energies which are not physically possible for the particular element and neutron energy considered. Perkin also normalizes the cross section separately at each neutron energy for which data are presented. This introduces spurious structure in the gamma spectra, which we have attempted to remove.

As the neutron energy increases beyond 8 Mev, there is an increasing paucity of experimental data (except for some data at 14 Mev). To compound the difficulty, the validity of the parameters used in the statistical theory [which includes (n,2n) and (n,3n) processes] becomes increasingly questionable as the neutron energy extends further away from the energy range at which experimental data exist. In general, the spectral shapes are given more reliably than the absolute values. However, since the fission spectrum has very few neutrons above 8 Mev, inaccuracies in the associated gamma spectra are not too important.

#### 2.2 DISCUSSION OF THE REFERENCES USED AND CALCULATIONS PER-FORMED IN PREPARING THE SPECTRAL INFORMATION FOR EACH ELEMENT

#### 2.2.1 Hydrogen

The gamma spectrum following a neutron interaction with hydrogen is a model of simplicity. There are no inelastic scatterings, and an absorption at any energy produces a single 2.23-Mev gamma.<sup>4</sup> The gamma spectrum is shown in Table 1.

#### 2.2.2 Beryllium

The gamma spectrum following a neutron absorption in beryllium is based mainly on the work of Draper and Bostrom.<sup>5,6</sup> The spectrum is given in Table 2. There is no  $(n,n'\gamma)$  reaction in Be at the neutron energies of interest ( $\leq 18.0$  Mev). However, the (n,2n) reaction does lead to the creation of a 2.43-Mev gamma. The number of such gammas created per "inelastic" event as a function of neutron energy is given in Table 3. The data are based on the work of A. Krumbein.<sup>7</sup>

#### 2.2.3 Carbon

E. Troubetzkoy and H. Goldstein<sup>4</sup> give the spectrum of gamma rays following neutron capture in carbon. The spectrum was assumed constant for incident neutron energies below 223 ev, above which the  $(n,\gamma)$  cross section becomes zero. Above 223 ev no gammas are produced per neutron absorption as the gammas following the charged-particle reactions are neglected. The spectrum is shown in Table 4.

The gamma spectrum following inelastic scattering in carbon was based on statistical theory for neutron energies above 10 Mev. For  $E_n \leq 10$  Mev, experimental level-excitation cross sections were used. It was assumed that the 7.66-Mev level does not decay by gamma emission.<sup>8</sup> The data are presented in Table 5.

#### 2.2.4 Oxygen

There are no neutron capture events in oxygen. The absorption spectrum given in Table 6 represents the production of the 3.5-Mev gamma following the  $(n, \alpha)$  reaction. The gamma spectrum following inelastic scattering was based on the experimental level-excitation cross sections measured by the Rice group and quoted in BNL-325.<sup>9</sup> These were supplemented by statistical model calculations. The details are given in Reference 10. The gamma spectrum, as a function of neutron energy, is given in Table 7.

#### 2.2.5 Aluminum

The capture spectrum was taken from J. E. Draper and C. O. Bostrom<sup>6</sup> and the compilation of E. Troubetzkoy and H. Goldstein.<sup>4</sup> The gamma spectrum emitted per inelastic-scattering event in aluminum was deduced as follows. The work of Towle and Gilboy<sup>11</sup> and the compilations found in References 12 and 13 were supplemented by Hauser-Feshbach calculations to determine the excitation cross sections for the six lowest levels for incident neutron energies below 4.0 Mev. These agreed remarkably well (to within 15%) with the calculations of M. Leim-dörfer (personal communication).

For neutron energies above 4 Mev, statistical-model calculations were used. These agreed well with the discrete-level data of Reference 11 and the Hauser-Feshbach calculations near 4 Mev. The calculations were supplemented by the experimental data of Perkin<sup>3</sup> for neutron energies to 8.5 Mev and by the data of Thompson and Engesser at  $E_n = 14$  Mev (Reference 14).

The capture and inelastic gamma spectra are tabulated in Tables 8 and 9.

#### 2.2.6 Chromium

The spectral intensities of gamma rays resulting from neutron capture in chromium were taken from the compilation of Troubetzkoy and Goldstein.<sup>4</sup> The range of

gamma energies from 0 to 9.72 Mev has been divided into seven groups. Integrated intensities are given for each group. Since the charged-particle reactions [which far outweigh the  $(n,\gamma)$  reaction for  $E_n \ge 2.5$  Mev] are assumed to be accompanied by negligible gamma emission, the absorption gamma spectrum equals the capture spectrum for neutron energies up to 2.5 Mev. Thereafter it drops sharply to zero.

The gamma spectra following inelastic scattering in chromium were based on the level-excitation cross sections of Van Patter,<sup>15</sup> for incident neutron energies below 3 Mev. For incident neutron energies above 3 Mev the gamma emission was calculated from statistical theory.

The capture and inelastic gamma spectra following neutron interactions in chromium are given in Tables 10 and 11.

#### 2.2.7 Iron

The absorption gamma spectrum for iron was taken to be the capture spectrum, as given by Troubetzkoy and Goldstein,<sup>4</sup> for incident neutron energies below 4.5 Mev. For  $E_n > 4.5$  Mev the very small  $(n,\gamma)$  cross section is negligible compared with the (n,p) and  $(n,\alpha)$  reactions. It is assumed that the charged-particle reactions produce no gammas.

The production of gamma rays by inelastic neutron scattering from iron, for incident neutron energies below 4 Mev, were taken from the experimental data of Montague and Paul.<sup>16</sup> For neutron energies above 4 Mev, the data of Perkin<sup>3</sup> which extend to  $E_n = 8.5$  Mev, and the data of Caldwell<sup>17</sup> for  $E_n = 14$  Mev were used.

The capture and inelastic gamma spectra are tabulated in Tables 12 and 13.

#### 2.2.8 Nickel

The absorption gamma spectrum in nickel was assumed to be equal to the capture spectrum, which was taken from the compilation of Troubetzkoy and Goldstein.<sup>4</sup> The spectrum of gamma rays following inelastic scattering, for neutron energies below 4 Mev, was based mainly on the discrete-level excitation cross sections of Broder et al.,<sup>18</sup> supplemented by the work of Day<sup>19</sup> and Cranberg and Levin.<sup>20</sup> For neutron energies between 4 and 8.5 Mev the data of Perkin<sup>3</sup> were used. Above 8.5 Mev, the spectra were calculated by statistical theory with parameters adjusted to fit Perkin's data at  $E_n = 8.5$  Mev.

The spectra are tabulated in Tables 14 and 15.

#### 2.2.9 Zirconium

The capture gamma spectrum is taken from the compilation of Troubetzkoy and Goldstein<sup>4</sup> for gamma energies above 3 Mev. For  $E_{\gamma} < 3$  Mev the capture spectrum for zirconium was assumed to approximate that for molybdenum (which has a similar spectrum above  $E_{\gamma} = 3$  Mev). The values were taken from Reference 4. The gamma spectrum following inelastic scattering events in zirconium was obtained from the data of M. Fleishman,<sup>21</sup> using the total inelastic scattering cross sections of J. Ray (UNC Phys./Math Memo No. 1679, Dec. 1960). The spectra are tabulated in Tables 16a and 16b.

#### 2.2.10 Cadmium

The capture gamma spectrum given in Table 17 is based on the compilation of Troubetzkoy and Goldstein<sup>4</sup> and on the work of Smither.<sup>22</sup> Nothing is presented in this report on the inelastic gamma spectrum as cadmium will be present in relatively low concentrations for its effect as a strongly capturing medium.

#### 2.2.11 Samarium

The capture spectrum given in Table 18 is from Troubetzkoy and Goldstein<sup>4</sup> and Groshev.<sup>23</sup> No inelastic gamma spectrum is given.

#### 2.2.12 Tungsten

#### A. Natural Tungsten

The gamma spectrum following thermal neutron capture is based on the experimental data found in References 24 to 27. The spectrum, assumed to be independent of neutron energy, is given in Table 19.

For low neutron incident energies ( $E_n < 2$  Mev), the spectrum of gamma rays from inelastic scattering is based mainly on the discrete level-exitation cross sections of Smith.<sup>28</sup> At higher neutron energies the data from Perkin<sup>3</sup> were supplemented by statistical-model calculations. The values obtained are presented in Table 20.

#### B. The Tungsten Isotopes $-W^{182}$ , $W^{183}$ , $W^{184}$ , and $W^{186}$

Additional calculations have been made to provide capture gamma spectra for the separated tungsten isotopes  $W^{182}$ ,  $W^{183}$ ,  $W^{184}$ , and  $W^{186}$ . These spectra were obtained from that for natural tungsten, as described in Troubetzkoy and Goldstein<sup>4</sup> in conjunction with data on binding energies and on certain gamma lines assignable to particular isotopes, as given by Treado and Chagnon.<sup>29</sup>

Briefly, the procedure was as follows. The capture gammas for natural tungsten are represented by a 12-group spectrum  $(n_i E_i)$ , i = 1, 2, ..., 12; i.e., per absorption in natural tungsten there are emitted  $n_1$  gammas of energy  $E_1$ , etc. A "background" spectrum was constructed, assumed common to all of the isotopes, by subtracting (with proper weighting) contributions attributable to particular isotopes. This background is

$$n_{i}^{(B)} = n_{i} - \sum_{\substack{j=182,183\\184,186}} a_{j}n_{j,i}$$

where  $n_{j,i}$  is the number of gammas emitted, in the energy range containing  $E_i$ , by isotope j, and  $a_j$  is the abundance of isotope j in natural tungsten, multiplied by its thermal capture cross section and divided by the sum of the products of abundances and cross sections, i.e., aj represents the probability that a given capture in natural tungsten is a capture in the isotope j.

Following this, the background spectrum was renormalized separately for each isotope j, yielding the constants  $c_j$ , so as to set the total gamma emission per capture in isotope j equal to the binding energy  $U_j$ :

$$C_{j\sum_{i}} n_{i}^{(B)} E_{i} + \sum_{i} n_{j,i} E_{j,i}' = U_{j}$$
(3)

where  $E'_{j,i}$  is the actual energy of a line emitted by isotope j in the range represented by  $E_i$ , and  $U_j = 6.10$ , 7.48, 5.86, 5.34 for j = 182, 183, 184, 186, as reported by Treado and Chagnon. Finally, using the  $C_j$  found from Eq. 3, a 12-group capture gamma spectrum was computed for each isotope, following the formula

$$N_{j,i} = C_j n_i^{(B)} + \sum_k n_{j,k} \left( \frac{E_{j,k}}{E_i} \right), \qquad (4)$$

the summation being over lines k in the i<sup>th</sup> energy bin of the spectrum. The last factor ensures the energy balance satisfied by the spectra

$$\sum_{i=1}^{12} N_{j,i} E_i = U_j$$
 (5)

The capture spectra are presented in Tables 21 through 24.

11

(2)

#### 2.2.13 Uranium-235 and Uranium-238 - Nonfission Gammas

In the proposed Monte Carlo calculations, the fission-related gammas will be generated from a prescribed power pattern and operating history, without regard to the histories of individual neutrons after fission. On the other hand, neutron absorptions encountered during the neutron Monte Carlo are recorded on an interaction tape which is later processed by the GASP program to generate a source for the secondary gamma problem. Since no distinction is made in this recording between nonfission and fission captures, the element-dependent gamma-production input to GASP should define the nonfission gammas produced per absorption (i.e., per capture-or-fission), since the fission gammas are treated in a separate (primary gamma) calculation.

These nonfission gammas per absorption were computed as follows. Let  $n_{i,j}$  be the number of gammas of energy i produced as a result of the nonfission capture of a neutron in energy bin j. Then

$$N_{i,j} = \frac{\sigma_{\gamma}(E_j)}{\sigma_{\gamma}(E_j) + \sigma_f(E_j)} \cdot n_{i,j}$$

(where  $\sigma_{\gamma}$  and  $\sigma_{f}$  are the capture and fission cross sections) represents the number of nonfission gammas of energy i produced per fission or nonfission capture of a neutron in energy bin j. In the calculations for  $U^{235}$  the  $n_{i,j}$  were assumed to be independent of j (the neutron energy). The capture spectrum (for  $U^{235}$ ) was taken to be the same as that of the prompt fission gammas (to be discussed later). Their intensity corresponds to a total of 6.429 Mev/capture (see References 30, 31). To generate the  $N_{i,j}$  for  $U^{235}$  the following average cross section values were used (data from Reference 8).

Neutron	Lower Energy	
Bin	Limit, Mev	$\sigma_{\gamma}/(\sigma_{\gamma}+\sigma_{f})$
1	3.7 (-8)	0.15
2	1.0 (-7)	0.17
3	1.0 (-6)	0.25
4	4.0 (-6)	0.60
5	7.0 (-6)	0.30
6	1.0 (-5)	0.34
7	1.5 (-2)	0.28
8	1.0 (-1)	0.15
9	1.0 (+0)	0.035
10	4.0 (+0)	0.003
11	10.0 (+0)	0.0005
	1.81 (+1)	

#### AVERAGE CROSS SECTIONS $-\sigma_{\gamma}/(\sigma_{\gamma}+\sigma_f)$ vs ENERGY FOR U<sup>235</sup>

For  $U^{235}$  the N<sub>i,j</sub> (number of nonfission gammas per absorption) are given in Table 25.

The thermal capture gamma spectrum, the  $n_{i,j}$  for  $U^{238}$  is based on the work of Campion (see References 32 and 8). The spectrum was assumed to be independent of neutron energy; the total emission is 6.37 Mev/capture. Table 26, giving the number of nonfission gammas per absorption in  $U^{238}$ , is based on the following average cross section values (data from Reference 8).

Neutron Bin	Lower Energy Limit, Mev	$\sigma_{\gamma}/(\sigma_{\gamma}+\sigma_{\mathbf{f}})$
1	3.7 (-8)	1.00
2	1.0 (+0)	0.75
3	1.3 (+0)	0.35
4	1.5 (+0)	0.10
5	1.8 (+0)	0.08
6	2.3 (+0)	0.04
7	4.0 (+0)	0.01
8	7.0 (+0)	0.001
	2.0 (+1)	

#### AVERAGE CROSS SECTIONS $-\sigma_{\gamma}/(\sigma_{\gamma} + \sigma_f)$ vs ENERGY FOR U<sup>238</sup>

#### A. Prompt and Delayed Gammas from Fission in $U^{235}$

Presented below are the spectra, intensities, and time variations which have been compiled and calculated. This is followed by some discussion of the sources of information, and recommendations for future work.

#### **B.** Results

### Recommended Values for the Gamma Release Rates and Integrals Following Fission of $U^{235}$

- 1. The total gamma energy release rate and integral, 0 to 10 seconds, is shown in Table 27 and Fig. 1.
- 2. The spectrum of the prompt (0 to  $5.0 \times 10^{-8}$  second) radiation is shown in Table 28.
- 3. The integral data (gammas released during first 1.0 second after fission) indicated in the last column of Table 27 are summarized in Table 29.

At present the spectrum is assumed to be constant over the entire 1.0 second. Table 28 could be modified to represent the range 0 to 1.0 second by multiplying the entries for Mev/fission and photons/fission by 8.470/7.394 = 1.146.

- Photon release rates, in 12 groups (1.0 second to 10 hours) are shown in Table 30.
- 5. Integrated gamma output, Mev/fission, in 12 energy groups and three time bands from 1.0 second to infinity is shown in Table 31.

#### C. Discussion: Sources of Data – Calculations

#### 0 to $5 \times 10^{-8}$ Second

The principal source of data on the spectrum and time dependence of the "prompt" gammas (0 to  $5 \times 10^{-8}$  second) is Maienschein et al., References 33 and 34.



Fig. 1 — Gamma Energy Release Rate as Function of Time after Fission of  $U^{235}$ 

Fig. 4.1.1 of Reference 32 was integrated to provide gamma sources (for the first  $5 \times 10^{-8}$  second following fission) in each of the 12 gamma groups which are used in the tabulation of delayed gamma intensities, plus one group including the 5.5 to 7.5 Mev contributions. The energy range below 0.3 Mev, not reported in References 33 and 34, has been augmented by 0.24 Mev/fission, following Skliarev-skii<sup>35</sup> and Roos.<sup>36</sup> This led to a value of 7.39 Mev over the range of 0.02 to 7.5 Mev, 0 to  $5 \times 10^{-8}$  second.

#### $5\times10^{-8}$ to $1.0\times10^{-6}$ Second

Here (and over the whole first second) the spectrum is assumed to be the same as that for the very prompt radiation. The time dependence from  $5 \times 10^{-8}$  to  $10^{-6}$  second was estimated by combining the four intensity-vs-time curves of Fig. 4.2.2, Reference 34, each weighted by the average energy of the pertinent gammas. Small extrapolations of the 0.70-Mev and 1.30-Mev curves of that figure permitted a calculation of a plausible shape of the intensity-vs-time curve for the energy region of 0.15 to 1.42 Mev. This shape was assumed to describe the time variation of the entire gamma source in this time interval.

For the normalization of this portion of the curve, use was made of Maienschein's experimental result that, from  $5 \times 10^{-8}$  to  $10^{-6}$ , about 5.7% as many counts were observed over a fairly wide range (0.16 to 1.93 Mev) as were observed in the first  $5 \times 10^{-8}$  second for the same energy range. Hence the total energy emission in this time range is taken to be  $0.057 \times 7.39 = 0.421$  Mev.

#### $1.0\times10^{-6}$ to $6.0\times10^{-5}$ Second

This range was filled in by graphical interpolation between the earlier and later times. The possible error cannot be too large as the integrated energy release over this range is only about 0.1 Mev.

#### $6.0 \times 10^{-5}$ to 1.0 Second

The shape of the curve of total energy release rate vs time in this time range was taken to be the same as that given for  $E_{\gamma} \ge 0.51$  Mev in Reference 37, Fig. 9 and Table 1. The normalization from photons/fission-sec to Mev/fission-sec was made by matching the U<sup>235</sup> portion of that reference at 1.0 second to the absolute intensity (in Mev/fission-sec) implied by the table of intensities which we have given for times  $\ge 1.0$  second.<sup>38,39</sup> This normalization appears well-founded since the ratio of Walton's intensities to Zigman and Mackin's is very nearly constant over the range from 1 to 4 seconds.

#### 1.0 Second to 3.64 $\times$ 10<sup>4</sup> Seconds (10.14 Hours)

The data up to  $1.74 \times 10^4$  seconds are reproduced from the work of Zigman and Mackin<sup>38</sup> as reported in Watson.<sup>39</sup> The first nine columns (covering 0.02 to 4.0 Mev) were extrapolated graphically from  $1.74 \times 10^4$  to  $3.64 \times 10^4$  seconds. The last three columns (covering 4.0 to 5.5 Mev) were extrapolated from  $2.75 \times 10^3$  or  $4.03 \times 10^3$  seconds.

#### $3.64 \times 10^4$ Seconds to Infinity

The last time bands considered, extending the energy-release computations to  $10^8$  seconds and infinity, were treated by assuming that the exponents, c, in power fits to the various energy-rate curves, W' = at<sup>-C</sup> Mev/sec, were the same for  $t > 3.64 \times 10^4$  seconds as in the range  $1.74 \times 10^4 \le t \le 3.64 \times 10^4$  seconds.

#### 2.2.14 Uranium-235 Energy Distribution of Gamma Rays Following Neutron-Producing Reactions

The gamma-ray spectrum following inelastic-scattering events in U<sup>235</sup> was derived, for low incident neutron energies, from Hauser-Feshbach calculations, taking into account the known discrete levels.<sup>40</sup> At higher energies, statistical model calculations were performed. The spectra are given in Table 32.

#### 2.2.15 Uranium-238 Energy Distribution of Gamma Rays Following Neutron-Producing Reactions

The gamma-ray spectrum following inelastic scattering for  $E_n < 1$  Mev was derived from level-excitation measurements.<sup>41-44</sup> These measurements were extended by Hauser-Feshbach calculations, taking into account competition from fission and capture processes. The neutron penetrabilities were calculated using the nonlocal potential of Perey and Buck.<sup>45</sup> For  $E_n > 1$  Mev the gamma spectra are based on statistical theory, including (n,n'), (n,2n), and (n,3n) processes. The spectra are presented in Table 33.

# TABLE 1 — HYDROGEN –NUMBER OF GAMMA RAYSEMITTED PER ABSORPTION

	$E_{\gamma}$ , Mev
Energy	2.23
0.03 ev	1.0
18.02 Mev	1.0

### TABLE 2 --- BERYLLIUM -- NUMBER OF GAMMARAYS EMITTED PER ABSORPTION

		E <sub>y</sub> , Mev													
E, Mev	.8550	2.5900	3.3650	3.4410	5.9560	<u>6.8070</u>									
2,00000 <u>5</u> 01	0	-0	-0	- 0	- U	-0									
1,00000=-03	,2400	,2100	.2800	.1100	.0200	.6500									
1,000005#09	2400	,2100	<b>.28</b> 00	,1100	,0200	.6500									

#### TABLE 3 — BERYLLIUM – NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

	E	$\gamma$ , Mev
E, Mev		2.43
1.80200EE 1.63028EE 1.47514EE 1.33476EE 1.20774EE 9.88815EE 8.9947173EE 8.9955233EE 5.997673EE 5.997673EE 5.99103025 4.9103025	01 01 01 01 01 01 00 00 00 00 00 00 00 0	2610 2900 3200 4000 5100 5700 5700 5700 5700 5700 5700 5
3 63765E 3 29148E 2 97825E 2 69484E	00 00 00 00 00	4300 3000 1200
1,00000E-	10	Ũ

#### TABLE 4 — CARBON – NUMBER OF GAMMA RAYS EMITTED PER CAPTURE

	$E_{\gamma}$ , Mev	
1.27	3.68	4.95
.3000	.3000	.7000

TABLE 5 — CARBON – NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

	9.00	2704		1044	.0668	00000	<b>D610</b>	<b>10092</b>	500°	<b>6</b> 000	1000	0000	0000	0	0	0	0	0	0	c	a	o	0	0
	6.00	0196		0083	0057	.0041	.0030	6 T O O *	9000°	<b>0005</b>	0000	•	0	0	0	0	0	0	o	0	0	•	0	0
	5.25	6600 .		.0033	0022	0013	<b>1</b> 0000	1000	• • • • •	0000	0	•	0	9	0	0	•	0	0	0	0	0	0	0
	4.438	.0677		19397	4234	.4971	. 5778	. 6555	* 7032	.7811	.9447	, 8859	• 9170	.9458	0696	. 9836	. 9902	.9939	. 9974	.9979	9983	1,0000	1,0000	0
	3.75	0050		.0008	000	.0002	.000	<b>1000</b>	.0001	0000	0	0	0	0	0	0	0	0	0	0	0	0	<b>つ</b>	0
, Mev	3.25	6100.		9000	000	.000	.0001	0000.	0000	0000	.0000	0	0	0	0	•	0	0	0	0	0	0	0	0
Εγ	2.75	0100		000	0003	.0001	.0001	0000	0000	0000	0000	0	•	0	0	0	0	0	0	0	0	0	0	0
	2.25	9000.		0003	0005	.0001	.0001	0000	0000	0	o	0	0	0	0	э	Э	<b>ප</b>	э	Э	ð	.)	Э	0
	1.75	5000.	.0002	.0002	.0001	.0001	0000*	0000.	.0000	0000	0	•	0	0	•	0	ð	0	0	0	0	0	ð	Э
	1.25	.0001	1000	1000	,0000	• 0000	0000	0000	0000	0000	。	0	C	0	C	0	0	0	0	D	0	0	Э	0
	0.75	0000		0000	0000	0000	0000	0000.	0000	0000	0	0	c	0	0	0	0	0	c	Û	0	c	c	G
	0.25	0000	0000	0000	0000	0000	0000	0000	o	U	c	0	0	Э	n	0	2	0	0	n	c	2	ن	3
	E, Mev	1.80200E U1	1. 7-000F 01	1.55000E U1	1,47500E U1	1.40000E U1	1.33000E U1	1,27000E 01	1,21000E 01	1,15000E U1	1,09000E U1	1.04000E 01	9.890UUE UC	0.41000E 00	8.950U0E U0	8.65000E UD	8.51000E CO	A.35000E VU	A.20006 00	8.050UUE UU	7.0000E bU	7.83000E UC	4.0000E UL	4.82000E UI

#### TABLE 6 --- OXYGEN -- NUMBER OF GAMMA RAYS EMITTED PER ABSORPTION

	$\underline{\mathbf{E}}_{\boldsymbol{\gamma}}, \ \mathbf{Mev}$
E, Mev	3.5
2,00000E 01	4230
1,71000E 01	4330
1.63000E 01	4340
1,55000E 01	4310
1.47500E 01	4220
1.40000E 01	.4070
1.330005 01	,3930
1.2700UE 01	3990
1.210005 01	,3560
1.15000E 01	,367 <sub>0</sub>
1.090005 01	,462 <sub>0</sub>
1.04000F 01	,497 <sub>0</sub>
9.89000F 00	,5000
9.410005 00	<b>,</b> 5uoo
8.950005 00	.5000
8.510005 00	4500
A 100005 00	.4000
7.700005 00	.1000
7 33nnur All	0
	D D
<b>T<sup>6</sup> () n n n n n n 2 = = n</b>	-

#### .75 1.25 1.75 2.25 2.75 3.25 3.75 4.25 E, Mev 4 .0143 ,0099 ,0023 .0093 .0164 . v**17**8 .0060 1.81000E 01 .0010 ,0023 .0143 .0173 .0101 .0010 .0060 10099 .0161 1,71000E 01 • 0141 1.630002 01 ,0023 .0059 0158 . 1166 ,0104 10091 .0010 ,0022 0138 0152 , 9156 1,55000= 01 .0058 ,0106 0089 .0010 1,4750UE 01 .0107 .0135 .0146 ,0020 .0057 .0145 .0009 .0086 . .0112 0130 .0134 ,0055 .0083 .0136 1,40000 01 .0009 .0019 • .0125 .0116 .0016 .0124 ,0128 1.330002 01 .0052 jun74 .0008 1.27000E 01 1.21000E 01 ,0120 .0014 .0049 .0120 , 0123 .0000 .0073 . 0110 .0045 **j**12Ž .0005 .0069 .0117 ,0116 .0010 .0099 1.15000E 01 1.09000E 01 .0042 .0065 .0118 .0003 .0007 .0112 .0098 .0007 0062 .0005 .0041 .0113 .0061 .0094 .0031 +0001 . ( 1.040095 01 .0037 .0064 .0033 .0004 .0048 .0109 .0008 .0000 .0083 0023 9.89000E NJ .0027 .0034 .0003 .0008 Ũ 0 .0093 9.4100UE UU .0092 Ģ .0008 .0006 .0005 0 0 9.34009E 00 .0077 0 ŧ 0 U 0 Ű U .0062 9.200005 00 G Û 0 Û 0 11 D .0074 9,100002 00 Ü Ũ 0 Ð 0 U -U .0064 8,9500VE 00 Ũ D ü 0 0 Û U .0067 8,84003E 00 Û ٥ 1 ۵ 0 0 Ü ,0065 8.700095 00 G Ú G 0 Ũ J ŋ . Bu 73 Ō 8,510002 00 Ū, 0 0 0 υ ú .0000 8,35000E 00 ن ۵ 11 0 Û ü C .0070 8,1000uE au U U Ĵ, 0 0 0 0 .0046 7.8700uE uu 0 U. 6 n ۵ ۵ Ū. .0049 7,70000= 00 Ű. U J, Ð Û 0 Û 7.4000UE 00 . 3496 0 ü 3 Ð ΰ Û 0 7,33000E 00 0 ۵ IJ Ũ U 0 U Ũ 6.3n000± 00 ü a H $\mathbf{i}$ ΰ ΰ Ü 0 6.00000<u>2</u>00 Û J U Ú -8 Ü U п 1,000002-10 Û Û ť ų, 0 U Ü 9

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#### TABLE 7 — OXYGEN – NUMBER OF GAMMA

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RAYS EMITTED PER NEUTRON-PRODUCING REACTION	
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	Ε <sub>γ</sub> ,	Mev									
.75	5.25	5.75	6.25	6.75	7.25	7.75	9.0	<u>11.0</u>	13.0	15.0	<u>17.0</u>
)185	.0184	.0177	,1548	,1112	1029	.0112	2225	.2640	1429	1660	.0191
176	0172	0161	1559	,1115	1031	0080	2227	2672	1411	0561	0065
165	0158	0145	,1577	,1127	1644	0074	2233	2732	1361	U422	0000
153	.0142	0127	.1585	1131	1035	8052	2264	2797	1260	<u>u23</u> 0	٠ ١
139	0127	0112	,1645	1163	0997	0023	2344	2841	1089	10073	υ
126	.0114	8095	,1724	,1227	0968	0004	2469	2841	.0794	Ų	ΰ
115	0096	0067	,1803	1299	0941	0	2608	2756	e436	ü	Û
1098	0068	0034	,1927	1401	0917	n	2729	2555	0162	Û.	Ū
068	0033	0012	2045	1480	0889	Ő	2829	2135	0004	Ū	نا
032	0007	0	,2487	1630	.1844	n	3 <sub>0</sub> 42	1550	1 0	ü	0
007	Ű	Ű	3120	1824	.0766	Ó	3142	0634	0	Ū	Û
L U	ü	0	4356	1864	0669	n	2687	0220	Ū.	0	Ó
0	õ	Q	6976	1300	6389	ñ	1414	<b>0</b>	Ō	ŭ	Ű
0	Ð	Û	7738	1462	0421	ő	0424	ů .	Ď	ũ	ü
Ŭ Ŭ	Ŭ	Û	8421	1242	0361	n n	•	G	ŭ	Ű.	ŭ
Ú	Û	Ú	8705	0986	0279	n	0	n	6	Ğ	ů
0	Ŭ	Đ	8496	1166	0345	n	() ()	ŭ	(,	Č.	ů.
Ø	ŭ	0	8697	1019	0284	ti -	0	C C		ŭ	ŭ
D	Û	Û	8636	1051	.031 1	ñ	0	ő	1	Ū.	ŭ
0	6	Û	867 <sub>8</sub>	1039	. 129	n	â	6	11	- 6	
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0	ú	U	, A573	1087	0340	ú	Ŭ	ē	บ	ŭ	Ŭ
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#### TABLE 8 — ALUMINUM – NUMBER OF GAMMA RAYS EMITTED PER CAPTURE

#### E, Mev

. .

.20	2,5000
1.5	.9500
1.78	1.0000
2.5	.7000
4.0	.8000
6.0	.2000
7.73	.2500

E, Mev	.75	1.25	1.75	2.25	2.75	3.25	3
1.80200E 01 1.40000E 01 1.09000E 01 8.51000E 00 6.03000E 00 5.16000E 00 4.02000E 00 3.13600E 00 2.44000E 00 1.90000E 00 1.48000E 00 1.48000E 00 1.15000E 00 1.15000E 00 1.00000E-01	1000 1200 1300 1000 2000 2100 2200 1700 2900 3100 3600 1,0000 0	2100 2200 2500 3100 3500 4000 4700 7100 6900 6400 0	3100 3200 3100 2800 2400 1800 9400 90 00 00 00	1600 1700 1809 2000 2700 3700 3809 3309 3309 0 0 0 0 0 0	1100 1000 1300 1400 1300 0700 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	1100 1000 1200 1200 1100 0500 0100 0 0 0 0 0 0 0 0 0 0 0	

TABLE 9 --- ALUMINUM -- NUMBER OF GAMMA R

### 扣

	$E\gamma$ , Mev							
.75	4.25	4.75	5.25	5.75	6.25	6.75	7.5	8.5
100	.0900	.1100	.1000	<b>N9</b> 00	.n <b>9</b> 00	.0800	<b>,1</b> 500	<b>44</b> 00
000	.0900	<b>,1</b> 000	,0900	.080u	<b>.</b> 0 <b>7</b> 00	000	,1200	?3 <sub>00</sub>
900	,0900	<b>,</b> ១១០០	<b>,</b> ១៩០០	6703	a60a	0400	, 1600	,0686
700	.0990	,0800	<b>, 0600</b>	<u>∎</u> 0400	0300	.0100	0	ü
400	<u>,0600</u>	•0 <sup>5</sup> 00	<b>,</b> 0200	<b>005</b> 0	- n	U	Q	i,
100	.0100	0	0	J	Ο	0	Ú	J
0	Û	Ú	Û	U	0	9	Û	Ú
Ģ	Û	U	0	Û	0	D	D	Q
0	0	U	Ų	U	0	Ü	υ	3
0	Û	Ű	Û	u	Q	0	Ű	U
0	U	Ú	0	0	0	0	0	ü
0	0	U	Ú	ŋ	0	0	0	Û
0	Û	U	Q	9	0	0	0	t;
0	0	U	Ŭ	ΰ	Ō	0	0	Ŭ

AYS EMITTED PER NEUTRON-PRODUCING REACTION

 $\mathbf{E}_{\boldsymbol{\gamma}}, \mathbf{Mev}$ 

TABLE 10 — CHROMIUM –	NUMBER	OF	GAMMA	RAYS	EMITTED
PI	ER CAPTU	RE			

.

		E	y Mev			
0.5	1.5	2.5	4.0	6.0	8.0	9.716
.8500	.4100	,2100	,1200	,2300	.3900	.0640

TABLE 11 — CHROMIUM – NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

 $E_{\gamma}$ , Mev

10.5		•
8.0		•
5.5	4 4 808 9 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	;
4.75	4 8 6 8 6 7 7 6 7 7 7 7 7 7 7 7 7 7 7 7 7	•
3.75	$\begin{array}{c} \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$	
3.25	$\begin{array}{c} \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$	
2.75	NON004400000000000000000000000000000000	
2.25	$\begin{array}{c} 0 0 0 0 0 0 0 0$	
1.75	$\begin{array}{c} \textbf{H} \textbf{H} \textbf{H} \textbf{H} \textbf{H} \textbf{H} \textbf{H} H$	
1.25	$\begin{array}{c} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 $	
0.75	$\begin{array}{c} 0 \\ $	* * * •
0.25	「「「」」」」」」」」」」」」」」」」」」」」」」」」」」」」」」」」」」	
E, Mev		

TABLE 11 — CHROMIUM (CONTINUED)

			ti C	L I T		Ey Z	fev	1	4	1		
Mev	0.25	0.75	1.25	1.75	2.25	2.75	3.25	3.75	4.75	5.5	8.0	10.5
00 - 30	0	.1213	1,0400	1532	.0091	•	0	0	0	0	0	0
	•	1274	1,0192	7400	. 0161	0	0	0	0	0	•	0
	-					<b>ə</b> c	<b>&gt;</b> c		3 0		30	3 0
	<b>&gt;</b> c								50	<b>-</b> C	<b>-</b> -	<b>-</b> -
	• =	2023	0240	.111	.0051	• a	• =	• a	• -	• •	• ~	
	• 0	1521	0.01	0110		• •	• =	• •	• •		<b>,</b> a	•
006.00	- 0	1165	929	0115	) ca	• •	• •	• •	• •		0	0
00E: 00		.1162	91.6	• 6 0 0 •		0	• •	0	0	0	•	0
006-00	0	1209	. 9139	•	•	0	•	0	•	•	0	•
006 00	0	5555	. 8996	0	æ	0	•	0	0	0	0	0
00E 00	•	.1120	. 8880	•		0	0	0	0	•	•	0
00E: 00	0	. 1207	. 8792	0		0	0	0	0	0	0	0
005,00	0	.1239	.8762	•	3	0	0	•	0	0	0	0
00E-00	•	. 1930	. 8469	•	9	0	•	0	0	0	•	0
00E: 00	0	. 3156	. 6844	•	0	0	0	•	0	0	0	0
UOE: 00	0	. 9248	4752	•	0	•	•	0	0	0	0	0
00 - 300	0	6325	9694	•	0	•	0	0	0	0	0	•
00 200	0	. 6538	1940.	0	0	0	0	0	•	0	•	0
00E 00	•	7291	2708	0		•	0	0	•	0	0	a
00 E 00	0	. 7412	2568	•	•	0	•	•	0	o	0	0
00 300	0		2000	•	Ð	0	0	0	•	•	0	0
00 300	0	1,0000	0	0	0	•	0	0	0	0	0	a
00E=01	0	1,0000	•	a	0	•	0	0	•	0	0	0
006=01	0	1,0000	•	0	9	0	0	•	0	0	G	ð
005=01	0	1,0000	0	•	•	0	0	0	0	0	0	0
104300	0	1,0000	•	0	Ø	0	•	0	0	0	0	ø
006-01	0	1,0000	0	•	9	0	0	0	0	0	•	•
00E+01	0	1,0000	•	0	•	0	0	0	0	0	0	0
00e=01	0	1.0000	0	0		0	0	0	0	0	0	0
005=01	•	1,0000	•	•	Ð	•	0	0	0	0	0	0
005=01	0	1.0000	0	0		0	0	0	0	0	0	0
10100 10100	•	1,0000	0	•	0	0	0	0	•	0	•	0
	0	1,0000	Ð	-	-	0		Ð		6	•	0
006007	0	ð	•	0	•	•	•	0	0	0	•	0

#### TABLE 12 — IRON – NUMBER OF GAMMA RAYS EMITTED PER CAPTURE

$E_{\gamma}, Mev$	
.38	.7500
1.6	.6000
2.6	.2700
3.7	.2300
6.0	.2500
7.63	.3800
9.3	.0210

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	9.5	20000000000000000000000000000000000000	23
	8.5		
	7.5		00
	6.5		63
Mev	5.5	$\begin{array}{c} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} C$	
.ั ค	4.5		. 0 350 0 250
	3.5	$\begin{array}{c} ++++++++++++++++++++++++++++++++++++$	2530 2580 2680
	2.5	NNA 9 NNA 4 4 4 4 4 4 9 8 9 9 9 9 9 9 9 9 9 9 9 9	5550 5150
	1.5	4444444444 4444500000000000000000000000	6400 6100
	.85		1.1900
	E, Mev	$\begin{array}{c} 444444444444444444444444444444444444$	5.15000E 00 4.91000E 00

TABLE 13 — IRON (CONTINUED)

	9.5	0	<b>`</b>	3	0	0	0	0	0	0		0	0	3	0	0				9	0	0	0	ð	0	0	0	Э
	8.5	0	0	9	0	0	0	0	0	O	Ð	0	· د	0	Э	O I	D	0	0	D	0	0	0	0	0	0	0	0
	7.5	0	0	D	o	0	0	0	0	0	0	0	o	0	0	0	0	a	0	a	0	0	0	э	0	o	0	9
	6.5	0	0	a	0	0	0	0	0	0	0	0	0	0	a	•	a	0	Э	0	0	9	0	0	0	o	0	0
ev	5.5	0	c	0	C	0	0	0	0	0	0	0	0	Ð	0	0	0	c	0	Ð	0	0	0	0	0	0	G	Ð
$\mathbf{E}_{\gamma}, \mathbf{M}$	4.5	.0150	0 <b>0</b> 20	ŋ	0	0	O	n	n	Э	<b>つ</b>	Э	•	n	n	0	0	0	0	0	0	0	n	0	0	9	n	0
	3.5	,2200	,2000	1800	1550	125 <sup>0</sup>	.1000	,0560	Э	Ō	Э	0	Ð	Э	0	0	3	a	5	ø	c	0	0	0	Ö	a	0	0
	2.5	4950	<b>47</b> 00	4400	4000	,3600	.3200	.2000	,2000	.1001	00,00	3	<del>ر</del>	ŋ	Э	0	0	0	0	0	3	0	<b>つ</b>	0	0	O	0	0
	1.5	. 5800	<b>55</b> 00	5200	4900	4600	4300	3900	3400	3000	,2600	.2200	.1700	,1000	0980.	.0760	.0710	,0620	.0540	• 0 4 5 0	.0360	.0270	.0200	.0100	Э	0	a	a
	.85	1,1100	1.400	1,0300	1,0100	9800	.9600	.9400	. 0200	.9250	.9250	.0300	.9400	.000	.9140	,924 U	,9290	,9380	.9460	, 955 <sub>0</sub>	.9640	.9730	,9800	0066.	1,0000	1,0000	0	ŋ
	E, Mev	4.67000E 00	4 44000E 00	4 23000E 00	4.0200E 00	3,82000E 0J	3.64000E 00	3,46000E 0J	3,29000E 00	3,13000E UU	2.97004E 00	2,83000E 00	2,6900uE 0U	2,56000E 0U	2,44000E 00	2,3200UE DU	2,21000E 00	2,10000E 00	2,00000E 00	1,9000UE 00	1.81000E 00	1,72000E 00	1,63000E 00	1,5500UE 0U	1.4100JE 00	8,50000E=U1	8.4000E=01	1,000005-10

					ш	ι <sub>γ</sub> , Mev					
E, Mev	0.50	1.5	2.5	4.0	5.820	6.580	6.839	7.528	7.817	8.532	9.0
1,801705 01	. 6030	.0013	.000	.000	.000	.0002	.000	.0002	2000.	.000	012
1,40320E U1	. 8023	.0010	000	000	0005	0001	.000	1000	0003	4000	0000
1,09280E 01	0022	0100	9000	.000	0005	0001	5000	1000	0005	0003	0000
8.51080E 00	4008	.0011	.000	9000	.000	0001	.000	0007	.0002	4000	0000
6.6 <sup>2</sup> 8 <sup>2</sup> 05 UO	1000	.0014	.000	8000*	F000	0005	.000	.0002	0000	.0005	.0012
5,16210E 00	. 8050	.0020	.0012	.0012	+000°	0005	0011	2000	000	000	0010
4,02020E 00	0124	.0056	.0032	.0032	.0010	0001	.0030	.000	.0011	0020	0040
3,13100E 00	0010	.0081	.0047	0042	.0015	0100	0040	.0011	0010	0050	.0071
2,43640E UQ	.0526	.0234	.0135	5010.	<b>5400</b>	.0028	.0126	.0031	.0046	.0065	.0205
1,89900E UO	£730	.0764			.0148	000 °	.0420	.0102	.0152	0279	.0675
1.47900E 00	4570	.2030	.1170	.1170	.0378	0246	1110	.0268	0405	.0735	.1760
1,15180E UQ		0240	.2120	,2120	.0458	0495	. 2020	2640.	.0738	1350	.3270
8+97830Em01	.9000	. 4000	.2300	.2300	.0736	0484	2185	.0528	.0792	1450	3510

TABLE 14 --- NICKEL - NUMBER OF GAMMA RAYS EMITTED PER ABSORPTION

TABLE 15 --- NICKEL - NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

	10.5	9000000000 40 40
	<del>.</del> 0	10000000000000000000000000000000000000
	7.5	4000 900 900 900 900 900 900 900 90 90 90
	6.0	00000000000000000000000000000000000000
	4.5	0000N \$00000 90 N N 60 N 80 9 60 N 80 9 60 N 80 9 60 N 80 9 60
ſev	3.5	8004000 9004000 949490 949490 90008400000
$\mathbf{E}_{\gamma}, \mathbf{N}$	3.0	****** <b>****</b> <b>****</b> <b>****</b> <b>****</b> <b>****</b> <b>****</b> <b>****</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>***</b> <b>**</b> <b>**</b> <b>*</b> <b></b>
	2.5	8649000 8648000 8088000 9008000 900800000 900908000000
	2.0	
	1.5	
	<u>1.0</u>	00000000000000000000000000000000000000
	0.5	・・・・・・ ありまりのので あてのののので あてのののので あてのののでの。
	E, Mev	444 8 78 4 10 4 4 8 4 9 8 4 9 4 8 4 8 4 9 8 4 9 4 8 4 9 0 0 4 10 4 4 8 4 10 0 0 4 10 4 4 8 4 10 0 0 0 0 0 4 0 8 4 10 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0

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#### TABLE 16 — ZIRCONIUM – NUMBER OF GAMMA RAYS EMITTED PER CAPTURE AND PER NEUTRON-PRODUCING REACTION

(162	1) Numb Emitte	er of G ed per C	amma F apture	lays
		$E_{\gamma}, Me$	ev	
0.75	1.5	3.5	<u>6.0</u>	7.5
1.3	0.2	1.13	0.35	0.04

#### (16b) Number of Gamma Rays Emitted per Neutron-Producing Reaction

			$E_{\gamma}, M$	ev		
E, Mev	1.0	2.0	3.0	4.0	5.0	6.0
18.0	.39	.32	.25	.19	.10	.03
10.9	.50	.38	.31	.31	.19	.10
8.51	.53	.38	.32	.32	.21	.09
6.63	.58	.40	.30	.23	.12	1
5.16	.67	.42	.29	.12	0	
4.02	.80	.44	.15	0	1	
3.13	1.0	.50	0	1		
2.44	1.0	.30	1			
1.90	1.0	0				
1.48	1.0	1				
1.15	1.0					
0.90	0	ŧ	•	¥	•	t t

TABLE 17 --- CADMIUM - NUMBER OF GAMMA RAYS EMITTED PER CAPTURE

	8.0000	, 1100 , 1100
	5.5000	1700
	3.5000	.7300 .7300
<b>Aev</b>	2.5000	9600
۲_ ف	1.4000	. 9200 . 9200
	.6510	.1900 .1910
	.5580	. 8860 . 8810
	.5000	,3500 ,3500
	E, MeV	2,00005 01 1,00005-09

TABLE 18 — SAMARIUM – NUMBER OF GAMMA RAYS EMITTED PER CAPTURE

				$\mathbf{E}_{\mathcal{V}}, \mathbf{Mev}$				
E, Mev	.3340	.4390	.6500	1.3000	2.2000	3.5000	5.7000	7.2000
2.000005 U1 1.000005-09	. <u>5</u> 200	5400 5400	. 6500 . 6500	1.5000 1.5000	1,1600 1,1000	.4500 .4500	0050.	100.

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	6.60	,0720
	5.60	.0420
	5.25	• 0 • 0 0
	4.60	.1000
	3.80	1980
	3.25	.2300
E <sub>7</sub> , Mev 2.25 2.75 3.	.2600	
	<b>,33</b> 0 U	
	1.75	3950
	<u>1.25</u>	.4280
	0.75	.6536
	0.25	. 6000

TABLE 20 --- NATURAL TUNGSTEN - NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

						$\mathbf{E}_{\gamma}$ , Mev					
E, Mev	0.25	1.00	1.75	2.25	2.75	3.25	3.75	4.25	4.75	5.25	5.75
1.8000E U1	. 456 C	20000 20000 20000	.4100 .4100	.1800	0000.	0020	0300	.0200	100 00 00 00	0.000	0010
1.0900F 01	8000 8000	2,1000		3000	2040		00/0		0300	0200	0500
8 51000E 00	2400	2 1000	00 <b>4</b> 8	3600	2600	1400	1200	, UB00	0.200	0400	0200
6.63000E U0	3200	2,2000	1,0000	4000	2000	1600	4400	0200	0400	0200	0
5.16000 00	3600	2,1000	0046	3800	080	, N800			0	D	0
4.02000E UU	3600	1,9000	, A600	2300	۔ •	0	0	0	0	0	0
3.13000E 00	3800	1,7000	6900	0990	D	0	0	0	0	0	0
2,440U0E UC	4205	1 4000	4600	3	0	Ū	0	0	0	0	0
1.0000F UU	4200	1,0000	1500	0	0	0	0	0	0	0	0
1.48000E UJ	6500	.4000	υ	0	ņ	0	0	ø	0	0	0
1.15000E U	.9500	2000	2	0	2	0	0	0	0	0	0
0.0000F=01	1.0000	1000	ت	3	с.	3	3	o	0	C	0
7.00006=01	1,300.1	с	a	Э	Э	0	c	0	0	0	0
1.20000E-01	0	J	ŋ	<i>ت</i>	C	0	0	0	o	0	0

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		6,6000	0011.				6,6000	C470.
ANUT'		5.6000	.01UZ		TURE		5.6000	D
U PER CA		5.2500	.0633	1	PER CAF		5.2500	U Ó U 4
THITT		4.6000	6130		MITTED		4.6000	ĊAĈU.
CIAN AL		3.8000	1640		A RAYS E		3.8000	.2250
OF GAIMIN	Mev	3.2500	2100		F GAMM	Mev	3.2500	2910
VERTAIN	$\mathbf{E}_{\mathcal{Y}},$	2.7500	2600		UMBER C	Б <sub>Y</sub> ,	2.7500	3600
N - 70T - NAT TODATO T		2.2500	.33 <sub>00</sub>		N-183 — N		2.2500	4570
		1.7500	3950		UNGSTEI		1.7500	.5470
		1.2500	4280	1	LE 22 – T		1.2500	.5920
		.7500	<b>653</b> 0		TABI		.7500	<b>,</b> 9340
		.2500	• 5066				.2500	. 83UU

TABLE 21 — TUNGSTEN-182 – NUMBER OF GAMMA RAYS EMITTED DER CADTIBE

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		6,6000	14du.					6.6000	C140.
rur e		5.6000	<u>, 1516</u>			PTURE		5.6000	nca0.
PER CAP1		5.2500	, 0 4 3 y		ļ	PER CA		5.2500	<u>. 1530</u>
MITTED I		4.6000	9696			EMITTED		4.6000	.1221
RAYS EI		3.8000	.164,			IA RAYS		3.8000	.1720
F GAMMA	Mev	3.2500	2620			OF GAMN	Mev	3.2500	1620
IMBER OI	$\mathbf{E}_{\mathcal{Y}_{2}}$	2.7500	2510			NUMBER (	я Н	2.7500	2010
-184 – NU		2.2500	3329			N-186 – N		2.2500	,255 .
UNGSTEN		1.7500	3470			TUNGSTE		1.7500	305 و
.Е 23 – ТI		1.2500	4 3 <sub>13 U</sub>		ł	LE 24 '	ļ	1.2500	33 <sub>00</sub>
TABI		.7500	, 65 G J			TABI		.7500	5 _ 4 .
		.2500	, sc3,					.2500	4630

		6.4200	0000	000.	1200	. 7040 . 7048	. 0 8 3 0 8 5	. nu36	. 121
_		5.2440	. 4044		410c.	• 40 C 4	. J.01	4200 ·	4106. U-
		4.7430	0000	0000.	0024	.0055	0049 0198	0041	.0U24
		4.2430	0000	.000.	8400	. U104	.0190	0079	.0049
		3.7420	0000	• 10 n 2	00.76	1172	1304	.1127	0- 0-
		3.2400	.0000	. 1002	1125	1283	1220 1499	.9141	.125 u-
	E <sub>Y</sub> , Mev	2.7930	1000	0003	1174	1394	969II.	.0197	0
		2.3920	10002	0000 0000	1246	6491	1150	1325	1286
		1.9900	1001	. 0109	1432 1800	6260	173.	1640	1.432
		1.5590	5000.	c100.	142U	1730	1 5 U 5 U	1999 ()	n- 1-
		1.1020	C (1 (1 (1 )	. U368	1540 2940	3150 3150	6300	.1780	n
		6000	0, 14	. 1990 1991	7920	. 962 5 A 495	1,7000 7080	4610	
		.0894	134	2421	1.300 1.0300	2,4500 2,4700	4 1400 1.7200	1.(700	n
		E, Mev	1.82000= 01 1.60000 - 61	4.00000 00	1.00000 00 1.00000 00	1.56000=-62 2.00000=-62	7.00000=-00 4.00000=-00	1.00000=-00 1.00000=-0.2	3.7000u=-110

TABLE 25 --- URANIUM-235 - NUMBER OF NONFISSION GAMMA RAYS EMITTED PER ABSORPTION

TABLE 26 --- URANIUM-238 -- NUMBER OF NONFISSION (JAMMA RAYS EMITTED PER ABSORPTION

					$\mathbf{E}_{\gamma}$ , Mev				
E, Mev	.2500	.7500	1.2500	1.7500	2.2500	2.7500	3.2500	3.7500	4.5000
2.00000 01	4 n 7 - 4	1 10	2999	2006 -	.0003	0004	0000	0000-	. 001
		100 -1640 	. 1270	400. 1220	0031		5000. 8000	4000 4000	, 1007 1002
2.3n0nu= 8u 1.8n1nu= 8u	4650 4650	.1284 .16cu	0540. 	1432 1540	0246	n288	.1015	0032	1020
1.30000 00 1.30005 00	1 + 303 3 5012	1.2	512.	1903	.1070	1250	1007	0140	
1.00000 00	4 . 45 [1]	1.6.00	6839	5400	3070	1005	0190	0000	0020-
0.7000 E-08	n	0 1	<b>7</b>	- 1	-				

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#### TABLE 27 — TOTAL GAMMA ENERGY RELEASE RATE AND INTEGRAL, 0 TO 10 SECONDS

t, sec	Mev/sec-fission	Integral, <u>Mev/fission</u>
0-5.0(-8)	1.48(+8)(average)	7.394
5.0(-8) 1.0(-7) 2.0(-7)	3.51(+6)(instantaneous) 2.15(+6)(instantaneous) 8.93(+5)(instantaneous)	
4.0(-7) 6.0(-7) 8.0(-7)	2.72(+5) 1.02(+5) 4.26(+4)	0.421
1.0(-6)	1.95(+4)	) 0.102
6.0(-5) 8.0(-5)	6.20(+2) 5.14(+2)	
1.0(-4) 1.5(-4)	3.60(+2) 1.77(+2)	
2.0(-4) 3.0(-4)	8.99(+1) 2.66(+1)	0.0478
4.0(-4) 6.0(-4) 8.0(-4)	1.15(+1) 3.94(+0) 2.20(+0)	
1.0(-3)	1.58(+0)	J
1.5(-3) 2.0(-3)	1.21(+0) 1.12(+0)	
4.0(-3) 8.0(-3)	9.91(-1) 8.68(-1)	
1.0(-2) 2.0(-2)	8.06(-1) 7.60(-1)	0.0716
4.0(-2) 8.0(-2)	7.06(-1) 6.66(-1)	
1.0(-1) 2.0(-1)	6.50(-1) 6.05(-1)	
4.0(-1) 8.0(-1)	5.11(-1) 4.12(-1)	0.433
1.0(+0) 1.5(+0)	3.78(-1) 2.98(-1)	J
2.0(+0) 3.0(+0)	2.47(-1)	1.307
4.0(+0)	1.54(-1)	
ə.5(+0) 1.0(+1)	1.21(-1) 7.22(-2)	

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Energy,	Mev/fission, $0 \neq 5 \times 10^{-8}$ rec		Photons/fission,
Mev	0105×10 sec	$\mathbf{E} = \mathbf{V} \mathbf{E}_1 \mathbf{E}_2$	
0.02-0.4	0.710	0.08944	7.94
0.4 -0.9	1.950	0.6000	3.25
0.9 -1.35	1.327	1.102	1.204
1.35-1.8	0.912	1.559	0.585
1.8 -2.2	0.659	1,990	0.331
2.2 -2.6	0.526	2.392	0.220
2.6 -3.0	0.372	2.793	0.133
3.0 -3.5	0.310	3.240	0.0957
3.5 -4.0	0.218	3.742	0.0583
4.0 -4.5	0.155	4.243	0.0365
4.5 -5.0	0.089	4.743	0.0188
5.0 -5.5	0.061	5.244	0.0116
5.5 -7.5	0.105	6.423	0.0163

## TABLE 28 — SPECTRUM OF THE PROMPT (0 TO 5 $\times$ 10^{-8} SECOND) GAMMAS FROM U^{235} FISSION

Total 7.394

#### TABLE 29 — GAMMA ENERGY RELEASE, 0 TO 1 SECOND AFTER FISSION

Time Band, sec	Mev/fission
$0-5.0  imes 10^{-8}$	7.394
5.0(-8)-1.0(-6)	0.421
1.0(-6)-6.0(-5)	0.102
6.0(-5)-1.0(-3)	0.0478
1.0(-3)-0.1	0.0716
0.1-1.0	0.4332
Total	8.470

#### TABLE 30 --- PHOTONS PER SECOND PER FISSION AS A

		·····				-	
		.020	.400	.900	1.350	1.800	2
		.400	.900	1.350	1.800	2,200	2
Seconds		089	600	1 109	1 550	1 000	2
beconub				1.104	1.000	1.000	4
1.0000E	U D	1.6005-01	1.200E=01	7.200E-02	3.400E+02	2.500E-02	8.5
1.5000E	1:0	1,300E-01	9.60UE-02	5.300E-02	2.700E+02	1.800E-02	6.2
2.0000E	0.0	1.0006-01	8,100E-02	4.100E-02	2.300E-02	1.400E=02	7.ť
3.0000E	3 <b>0</b>	7.300E-02	6.100E-02	2.9005-02	1.700E=02	1.100E-02	6.7
4.0000E	9.0	5.500F-02	5.000E=02	2.300=-02	1.400E+02	A.500E-03	5.1
5. CUBUE	et 9	3.800E-02	3.7005-02	1.600E-02	1.600E=02	6.000E-03	4.6
9.0000E	÷ 0	2.400E-02	2.600E=02	1.1305-02	7.10003	4.200E-03	3.1
1,3000E	:1	1.7005-02	1.900E-02	7.600F-03	5.200E.03	3.000E=03	2.4
1.9000E	91	1.100E-02	1.300E-02	5.300E-03	3.700E-03	2.200E-03	1.1
2.8000E	11	8.800E-03	8.600E-03	3.8002-03	2.600E=03	1.500E-03	1.1
4.100UE	11	5.0005-03	5.900E=03	2.7005-03	1.800E+03	1.000E-03	8.0
5.0000E	° 1	3.300=-03	3.900E-03	1.9006-03	1.2005-03	6.700E=04	5.2
8.8000E	11	2.1005-03	2.600E-03	1.3005-03	8.500F-04	4.400E-04	3.5
1.2900E	.12	1.3005-03	1.790E-03	9.600E-04	5.700E-04	2.900E-04	2.3
1.89006	. 2	8.500E-04	1.2005-03	6.9105-04	3.800F=04	1.900E-04	1.5
2.7700E	2	5.600E-04	7.700E-04	4.800=-04	2.500F-94	1.300F-04	1.0
4.0600E	<b>2</b>	3.8005-04	5.200E=04	3.400E-04	1.700F-04	8-006E=05	7.0
5.950CE	12	2.9005-04	3.400E=04	2.300F-04	1.100F-04	5.200F-05	4.2
A. TURDE	2	2.3005-04	2.3005-04	1.5909-94	7.990F-05	3.400F-05	2.8
1.2800E	: 3	1.7008-04	1.500E-04	9.200E-05	4.600F-05	2.200F-05	1.7
1.8700E	<b>∋</b> 3	1.200F-04	1.000E-04	5.8000-05	3.100F-05	1.3005-05	1.1
2.7500E	13	8.300E-05	6.600E-05	3.000E-05	1.800F-05	8.200E=06	6.5
4.0300E	U <b>3</b>	4.9005-05	4.590E-05	1.70005	1.300F-05	5.000E-06	4.0
5.900UE	23	3.000F-05	2.900E-05	6.900E-06	7.690F+06	3.200F=06	2.5
8.6400E	03	1.700E-05	1.900E-05	4.700E-00	4.900E-06	2.1005-06	1.4
1.2700E	54	1.100E-05	1.200E-05	2.4005-06	2.700E-06	1.300E-06	6.8
1.7400E	<u></u> 4	7.800E-06	8.800E-06	1.600E-06	1.700E-06	9.60nE-07	4.2
3.6500E	04	3.100E-06	3.800E-06	5.600=-07	5.600E-07	4.2005-07	1.3

\*These data, as well as the spectra and intensities for earlier times  $(10^{-16})$  generator program written for the ATHENA system.

E1						
<b>E2</b>						
EG	Mev					
,		· · · · · · · · · · · · · · · · · · ·				
200	2.600	3.000	3.500	4.000	4,500	5,000
600	3,000	3,500	4.000	4.500	5,000	5,500
392	2,793	3.240	3.742	4,243	4 743	5 244
005-03	7 5005-07	6 3005-03	4 6886-03	3 7865-83		0 1005-04
005-03	4 500E-03	5 500E-03	4 000L-03	2 7885-8%		4 000E-04
005-03	5 0005-03	4.8865-83	3.60000-00	2 1845-83	1 2005-03	5.6105-04
1002-03		3.9565-63	2.088F-03	1.5005-03	8.7805-04	4.2005-04
005-03	4 00000-03	3.4005-03	2.4005-03	1 2005-03	6.8005-04	3.4005-04
005-03	3 2005-03	2.6000003	1.6006-03	8.7905-04	4.8005-04	2.4005-04
005-03	2 3005-03	1.8005-03	1.3001-03	6 3005-04	3.4005-04	1.7005-04
005-03	1 7000-03	1.3065-03	9.000F=04	4.4005-04	2.4005-04	1.2005-04
INDE-03	1.2005-03	9.400F=04	7.0005-04	3.2905-04	1.7008-04	8.400E=05
10E-03	8 2005-04	6.100F+04	4.600E-04	2.100FeU4	1.1005-04	5.400E-05
00F=04	5.8008-04	4.2008-04	3.100E-04	1.500E+04	7.5000-05	3.6005-05
00F-04	3.7095-04	2.700F=04	2.00UE-04	1.000F=04	4.700E-05	2.300E=05
90F=04	2.4005-04	1.700E-04	1.200E-04	0.700Em05	2.9005-05	1.400F=05
00E-04	1.5005-04	1.0008-04	7.600E-05	4.300E=05	1.8004-05	8.300F=06
100E-04	9.000=-05	5.0004-05	4.500E-05	2.500E=05	1.0005-05	4.400F=06
00E-04	5.300E-05	3.500E+05	2.7008-05	1.500E+05	5.200E-06	2.300E-06
105-05	3.300F-05	2.1008-05	1.600E-05	6.200E-06	2.800E-06	1.200E+06
106-05	1.8005-05	1.200E-05	8.700E-06	3.80000006	1.300E-06	4.600E-07
20E-05	1.2005-05	6.800E-06	5.000E-06	1.7006-06	5.4005-07	1.800E-07
00E-05	6.0005-06	3.700E-06	2.700E+06	6.200E+07	2.0006-07	4.900E=08
00E-05	3,400E-06	2.000E-00	1.400E-06	2 100E-07	6.400E-08	1.100E-08
10E+06	1,700E-06	9.800E-07	7.1008-07	4.500E=08	1.300E-08	1.000E-09
00E-06	8.300E-07	4.700E+07	3.400E-07	6,000E+09	2,000E-09	1,100E=10
00E-06	3,400E-07	1.900E-07	1.400E-07	7.500E+10	2.400E-10	1.200E-11
00E-06	1,300E=07	7.800E-06	5.700E+08	6.000E=11	2.200E-11	1.400E-12
10E=07	4,800E-08	2,800E=08	2.100E=08	7.500E=12	2,0005-12	1.100E=13
10E-17	2.500E-08	1,400E=08	1.000E=08	1.300Em12	3,300E-13	2.200E=14
D05-07	4.500E-09	2,200E+09	1.600E=09	1,500E=14	3.600E-15	3.100E-16

#### . FUNCTION OF ENERGY GROUP AND TIME AFTER FISSION\*

o 1.0 second), are included as internal data in the VANGEN source-

# TABLE 31 — INTEGRATED GAMMA OUTPUT, MEV/FISSION, IN12 GROUPS AND 3 TIME INTERVALS FROM1.0 SECOND TO INFINITY

		Tir	ne Interval, se	<u>c</u>	
		1.0 to	$3.64 imes10^4$	$1.0  imes 10^8$	Total,
Group	E, Mev	$3.64 \times 10^{4}$	to $1.0 \times 10^8$	<u>to</u> ∞	<u>1.0 to ∞</u>
1	0.089	1.77(-1)	3.42(-2)	5.90(-3)	0.217
2	0.600	1.25(+0)	4.04(-1)	2.38(-1)	1.89
3	1.102	1.07(+0)	5.15(-2)	2.09(-2)	1,12
4	1.559	9.29(-1)	6.24(-2)	1.25(-3)	0.993
5	1.990	6.31(-1)	1.53(-1)	1.01(-1)	0.885
6	2.392	5.53(-1)	1.96(-2)	2.20(-4)	0.573
7	2.793	3.61(-1)	3.50(-4)	0.	0.361
8	3.240	3.05(-1)	1.70(-4)	0.	0.305
9	3.742	2.59(-1)	1.50(-4)	0.	0.259
10	4.243	1.41(-1)	0.0	0.	0.141
11	4,743	7.71(-2)	0.0	0.	0.0771
12	5 <b>.24</b> 4	4.01(-2)	0.0	0.	0.0401
Total*		5.79	0.73	0.35	6.87
Total†					8.47
Total Fi	ssion Gam	mas			15.34

\*0.02-5.5 Mev.

.

†0.02-7.5 Mev, 0-1.0 sec.

TABLE 32 --- URANIUM-235 - NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

					$E_{\gamma}$ , Mev						
E, Mev	0.25	0.75	1.25	1.75	2.25	2.75	3.25	3.75	4.50	5.50	6.50
	.9777	1.4130	6998	.2072	.0547	.0167	0800.	• 0050	.0045	.010	.0001
.71390E U1	6006	1,1772	5261	1386	.0392	1610	.0143 .0043	0106	.0099 .0099	.0022	1000.
	7070 . 	1616.					. 1 2 2 1 . 7 5 5 5 1				
	1001	00004	/002. 0740	• • • • • •	• • • • • • • • • • • • • • • • • • •	- C C	100 140 140	. 721	673 	• 0 4 0 4	2000
	6372	4178	. 3025	2219	1779	1613	- 4 4 G	1104	1024	.0196	0008
334805 01	6429	4189	3558	2 6 8 4 4 8 0 C	2202	2064	1871	1446	1320	.0223	.0007
20070E U1	6556	4217	0000	. 2537	2185	.2178	.2033	1572	1368	.0174	.0000
1.20770E 01	. 6515	3773	.2754	.1980	.1879	.2049	.1960	,1477	.1123	.0070	0
1.14880E U1	5444	3341	.2161	.1545	.1655	.1932	.1812	.1243	.0682	.000	0
1.09280E 01	\$396	3050	1782	.1290	<b>151</b> 8	1769	1500	.0815	<b>,</b> 0221	0	0
10 3050E 01	6394	2935	.1623	.1163	1345	1416	.0945	.0292	.0014	0	0
.88820E UU	644 U	2976	.1598	.1029	.1009	<b>, ng36</b>	0321	,0020	0000.	0	0
0.45090E UU	, 6557	3127	1590	.0821	0568	.0278	.0028	.000°	1000.	0	0
.94720E UU	, 7202	3364	1577	.5634	,026u	, C Ü 62	.0015	<b>000</b>	1000.	0	0
3.51080E UU	, 7543	.3771	.1077	.0617	.0219	. nueo	<b>,</b> 0028	.0008	0005	0	0
.09570E U0	.8115	4558	<b>,</b> 2094	.0865	.0367	5440	0048	4100.4	400u •	0	0
70090E UU	9043	6043	3086	.1436	.0620	0243	.0083	• 0 0 <b>2</b> 4	.000.	0	9
.32530E 00	1.0489	.9638	4879	.2340	.1019	<b>0393</b>	.0132	.0038	.0011	0	0
5.96810E U0	1,2584	1,2635	.7452	.3562	1542	0591	.0197	.0056	.0015	Ð	0
5.62820E UU	1,5204	1.7482	1.0326	.4888	.2094	1670.	. 0260	0072	,020	0	0
5,30500E UV	1,7222	2,0681	1,2059	.5591	2332	.0854	.0220	.0072	.0019	0	0
5,99750E UU	1,7646	2,1.070	1,2024	,5389	, 2153	.0747	0220	• 0054	.0011	0	0
5,70500E UD	1,7002	2, J787	1,1578	.4988	,1895	.0617	.0167	.0036	2000.	0	0
5,42670E UU	1,7555	2,0493	1.1127	.4599	,1655	, <sup>n502</sup>	.0124	.0024	.0003	0	0
5.16210E 00	1,7504	2,0185	1,0568	.4217	1041.	• 0 <del>4</del> 0 1	.0089	.0015	.000 .	0	0
4,910305 U∪	1,7450	1,9865	1,0207	.3849	1225	.0315	.0062	- 0008 -	0001	0	0
1.67080E UU	1,7393	1,9532	, 9742	4645.	.1038	.0243	.0044		.0000	0	0
	1.7331	4 0 - 0 - 1	.9274	. 3152	0969	0193	- 200 ·	.0003	.0000	0	0
1.226305 US	1,7256	1,6625	100.		xt/0.	• 0 T 0 •	100.	1000 *	0000.	0	0
1.02020E 00	1,7196	1.8448	500H.	5162.	CBCU.	•00.00	0100.	1000.	0	0	0
5.82420E 00	1,7121	1.8061	.7873		- 040 	• 3006 •	.000.		0	• •	5
5.60/00E UD	1+0/1	H / 0.74						0000		- ·	2
5.40020E UD	1 • 6 4 5 5	1.7235	. 54 4 3	10/1.	.0250	2000			0	<b>.</b>	0 (
5+24120E UL	5020 T	1.0004	• • • • •	- 1 <del>-</del> 1	A T 2 0 •		.000	•	0	0	
3.131UGE 00	1.5769	1.6350	. 6052	.1250	. 0160	.0018	• 000 •	0	0	0	3
2,97830E 00	1,5000	1,5902	.5619	.1065	.0120	• 101 •	0	0	0	0	0
2.833U0E U0	1,6556	1,5427	-5194	• 0896	.0101	.000	0	0	0	0	0
2.69480E 00	1,5439	1,4943	.4783	.0750	.0080	,000 <b>4</b>	0	0	e	0	0
2.56340E Un	1,6314	1.4443	.4384	.no28	.0073	0000	0	0	0	0	0

TABL 3 22 --- URANIUM-235 (CONTINUED)

					Er, Mev						
E, Mev	0.25	0.75	1.25	1.75	2.25	2.75	3.25	3.75	4.50	5.50	6.50
2.4354DE UU	1,6182	1,3934	4003	+53n .	, 00 <sup>5</sup> y	9	3	0	9	0	0
2.31950E UU	1.6042	2 4 4 4 7	.3041	. 3468		0	0	0	o		
2.20630E UU	1,5892	1,2684	.3301	• 0 • 5 •	,0022	0	0	•	0	0	0
2,09870E 00	1,5736	1.2346	*562*	.0391	• 0 0 0 •	0	0	•	0	0	0
1.99640E UQ	1,5567	1.1789	.2724	1520.		0	0	0	0	0	0
1,89900E 00	1.5000	1,1229	. 2505	.0300	Э	•	•	0	•	3	0
1.80640E 00	1,9203	1.0647	2402.	• 0228	2	0	0	0	•	0	0
1.71630E UD	1,5004	1,0072	2237		3	0	0	0	•	D	•
1.64470E UU	1 4785	5046	2154	.0003	<b>D</b>	0	0	0	0	0	0
1,55480E UD	1,4569	8964	,2058	.0010	2	0	0	0	0	0	0
1.47900E UU	1,4337	.8468	.1916	0	5	•	0	•	Ð	0	0
1.40680E UU	1,4087		1735	D	7	0	•	9	0	0	0
1,33820E UU	1,3828	7713	1505	c	3	0		0	Ð	o	0
1.27300F UU		4 7 50 0 5	1217	•	э:	0	0:	0	0	0	•
		6202°			>	0	3	0	9	-	0
1, 308101, 1	9239 T	7485			<b>)</b> :	0	•	•	0	0	0
				<b>)</b>	с <b>э</b> -	<b>.</b>	<b>.</b>	0	•		0
			6 <b>4</b> 00 •		э ·	•	<b>-</b>	0	0	0	0
		0000	-		<b>.</b>			•	•	9	
			<b>D</b> :	<b>)</b> (	<b>)</b> :		2 0	0	•	0	-
			<b>)</b> (	<b>-</b> -	3 :	<b>-</b> c		<b>-</b> -	<b>-</b>	<b>-</b>	<b>&gt;</b>
	1 2000		<b>a</b> a		3 3	<b>&gt;</b> c	<b>)</b> c	<b>-</b> c	•	<b>-</b> c	<b>&gt;</b> c
	1 2540	3662	• =	• =			• =			<b>.</b>	<b>&gt;</b>
	1.2065	2000	, c	• •		• •	• •	• •			> c
6.066105-01	1,2808	2268					0	0			. 0
6 64540E+U7	1,2833	1596	0	0	0	•	0	0	0	0	0
6.32130E=01	1,2600	1031	0	o	<del></del>	0	0	0	0	o	•
010000mmc7	1.2682	0000	0	0	5	0	0	0	-	9	0
5, 71970E=01		0220	<b>.</b>	•	5	9	3	0	3	0	0
		0110	5	<b>.</b>	ר י י		<b>a</b> 1	•	-	•	
5+1/2-0EHU1 4-02300E-01	1 1 3 3 4		<b>&gt;</b> c	<b>ə</b> c	<u>ה</u> ב	<b>-</b> -	<b>ə</b> :	0 0		0 0	0 0
4 48295=U1	1,0073		> =	• =	-	<b>)</b> =	• -	<b>.</b> .	• •		<b>,</b> c
4.454905=01	1.0671	• •	) )					0 3		• =	• a
4.23730E+U1	1.0448	3	0	•		0			0	. 0	0
4.030606-01	1.0230	D	0	3	7	0	•	0	0	0	0
10=01+00=0	1,0103	0	0	-	8	0	0	0	0	0	o
10-2014-9-P		3	0	•	<b>n</b>	0	0	0	0	0	0
3.46920E=U1	1,0000	9			<b></b>	0	0	0	0	0	0
1,27000E+UZ	Ð	2	0	C	3	0	0	-	0	D	D

TABLE 33 --- URANIUM-238 - NUMBER OF GAMMA RAYS EMITTED PER NEUTRON-PRODUCING REACTION

E <sub>y</sub> , Mev	<u>5 1.75 2.25 2.75 3.25 3.75 4.50 5.50 6.5</u>	.34 .0955 .0584 .0424 0414 .0372 .0344 .000		73 000 1281 1280 0187 0181 0181 0181 0181 0181	.23 1014 0311 0247 0295 0366 0760 0516 009	08 1721 0485 0422 0513 0636 1332 0842 007		10 13455 "U914 LU854 LU71 1314 2601 1451 DU6	148 1276 0938 0920 1164 1396 2529 0879 002	11 1180 0893 0885 1091 1236 1834 0295 000	7. 1154 0447 0765 0863 0866 0858 0023 000					.80 ,0422 ,0114 ,0041 ,0043 ,0054 ,0126 ,0120	126 .0279 .0102 .0072 .0080 .0102 .0237 .0203 .003	705 .0274 .0167 .0128 .0144 .0183 .0429 .0363 .006	03 0426 0289 0222 0252 0320 0742 0629 011	150 .0686 .0465 .0361 .0412 .0528 .1217 .1021 .016		11 1303 0040 0113 1151 1151 0153		128 1367 1963 0897 1160 11518 3480 2149 (	148 11221 10908 1031 11235 11599 13555 1398	736 1149 0915 0979 1297 1653 3513 0628	866 1142 10959 1027 1334 1644 3097 10091 0	001 1184 1000 1046 1309 1533 2187 0	739 11249 11047 11021 11201 11362 11255 0	111 11314 11063 10947 11028 11142 10503 0 0	ian 1361 1040 0a3a 0a42 0a20 0111 0									
	.25 3.	0.414 .0			0295 0	0513 0		1071 .1	1164 1	1091	0863 .0	1404				0 0000	0800	0144 _0	0252 .0	0412 0	1640 0	1000	1006	1160 1	235	297 1	1334 1	1309	201 1	1028 1	0 842 0	0651 0	0473 0	0204	131					
Δê	2.75 3	0424		. 15763 .	0247	0422	. 696	.0854 .	0260	0885	0765					0041 .	0072	0128	0222	0361				0897	0.031	0979	1027	1046	1021	0.47	0,43,4	1040	057A				120	0058	0017	
E, Me	2.25	. 0584		1820	1150	0485	2040	0.014	.0 <b>93</b> 8	0,893	0447	0 7 3 1			9120	• 11 4	0102	.0167	0289	0465	4040	0040	1041	0063	0.008	0015	0959	1000	1047	1063	1040	1000	0050		4080			0477	0375	
	1.75	.0955		<b>.</b>	1014	0721		1015	1276	1180	1154	1122	000			0422	.0279	.0274	0426	0686		202	1516	1367	1221	1149	1142	1184	1249	4101	1361	1380	100	1380	1351		1100	1133	1055	
	1.25	1635	0/17	2422	1123	1100	1961	6191.	1948	1818	1776				00TT.	.1180	0926	0705	0603	0.050	1445	10.1	2182	2028	1848	1736	1686	1601	1739	1811	1AAD	1042	1005	2011	Mada	• • • • • • • • • • • • • • • • • • •	2000	1970	1040	
	0.75	2057	+ 6 8 2 4	.2688	2480	2404	2002	2980	3154	3076	3032	30.41			5062	2734	2465	2125	1767	1577	1754	2205			2467	2394	2357	2356	2384	2434	2480	2546	2505	9536	2667			- 10 - 10 - 10 - 10 - 10 - 10 - 10 - 10	2771	
	0.25	. 7385		4000.	.7265	7178	8011	0426	1931	7411	7380	7404			· 7 · 78	7542	.7630	. 7776	900e	101 1	0114		1.0075	1,1226	1,1207	1 1284	1,1360	1,1476	1,1602	1.1744	1.1.467	1 1084	1 2083	1 2163	1 2237	1 2304	1 2443	1,2514	1,2595	
	, Mev	020ne 01		2000E 01	5000E 01	7500c 01		2000E 01	7000E 01	1000E 01	5000c 01	0000F 01	40.001.01		17000E 00	5000E 00	1000E 00	0000E 00	00000 00	3000= 00	70005 00	30005 00		00000	0000E 00	3000E 00	6000F 00	1000E 00	7000E U0	4000 00	3000 00	20005 00	200rc 00	4000 C 00			7000E 00	3000 00	90005 00	

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TABLI: 33 --- URANIUM-238 (CONTINUED)

					E, 'I	Mev					
E, Mev	0.25	0.75	1.25	1.75	2.25	2.75	3.25	3.75	4.50	5.50	<u>6.50</u>
2,32000E 00	1,2870	.2804	.1812	.0770	n600*	0	0	Ð	0	0	•
2,21000E 00	1,2976	2806	1748		0037	00	-	•	0	•	0
			10/5		<b>~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ ~ </b>				<b>-</b> c		
	1.3371	2775			<b>,</b> 3			<b>.</b>	<b>,</b>	<b>,</b>	> c
1.81000E 00	1.3520	0440	1001	0153	• •	• •	• •		• •		• •
1.72000E 00	1,3691	2711	1200	.0077			0				. 0
1.63000E 00	1,3891	2657	1037	.0025	0	0	0	Ð	0	0	•
1,55000E 00	1,4099	25533	.0874	.0015	2	0	•	0	0	•	0
1,48040E U0	1,4299	, 2518	0694	0	Э	•	0	0	0	0	0
1.41000E U0	1040.	3290	.0527	0	9	0	Ð	0	0	0	0
1.34000E 00	1.035	2010 1	1420	0	3	0	0	0	•	0	0
					<b>ə</b> a	0 0	0	0 (	-	•	0
	001	6778 9273		<b>.</b>	<b>)</b>	5 0		51			•
	1 1 2 2 2 2	1000	<b>16</b> nn •			<b>-</b> c		5		•	•
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	+ + + 4000					- c	20	<b>-</b>	<b>.</b>	•	> <
	1 1000		• c		• =	• -		• c	) c	• c	
	1 2224	1101	20	0		• a	0	• •	• a	• •	• 0
8.12000E.01	1 2342	0260	0	0		0	0	0		0	•
7.720U0E-01	1.0454	0469	0	0	0	0	•	0	0	0	0
7.34000E=U1	1,2565	0220	0	0	0	0	0	0	o	0	0
6.99000E=01	1,2619	0	0	0	<b>.</b>	0	0	0	0	0	0
6.600UE=U1		0 0	0		3	0		•	0	0	0
6.34000E=01	1 1070	5 9	30	- 0		- c					90
5.7200F.01	1.6600	0			0	• •		• •		•	• •
5.440U0E=01	1,1527	0	0	0	5	0	0	0	0	0	0
5.18000E=01	1,1329	c	o	0	7	0	•	0	•	0	0
4.92000E-01	1.1980	0	0	0	9	0	0	0	0	0	•
4,660U0E+U1	1,1012	0	o ·	9	0	0	0	0	0	0	•
4.450006.01	1.0784	•	0		9	0			0	0	0
4.24000=01		0	0		э.	•	0	D	0		•
134U00090.4		э с	<b>.</b>		<b>ə</b> :		30	0	•	0 0	•
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1,*30,006*P	7 . UUUU	S	o	2	Þ	2	2	2	2	7	2

#### 3. CONCLUSIONS

The tables presented in this report represent an adequate representation of the gamma spectra following neutron absorptions and inelastic-scattering events from the elements of interest in the tungsten nuclear rocket program. These include H, Be, C, O, Al, Cr, Fe, Ni, Zr, natural W and four W isotopes,  $U^{235}$  and  $U^{238}$ .

The thermal capture gamma spectra are generally well-known. The assumption that the capture spectrum is independent of the neutron energy should be further investigated. The errors introduced in the gamma sources by neglecting gammas from charged particles, though probably small, should be investigated. Indeed, the charged particles themselves will, in general, deposit several Mev of energy locally at the point of interaction. Neglecting the latter contribution might be more important than neglecting the gammas produced in the charged-particle reactions.

The gamma spectra following inelastic scattering, for incident neutron energies below ~4 Mev, are reasonably accurate as they are based on experimental levelexcitation cross sections. The resultant gamma spectra become less reliable in proportion to the extent to which the experimental data are supplemented by Hauser-Feshbach calculations.

In the intermediate neutron energy range (4 to 8 Mev) the inelastic gamma spectra are not uniformly reliable. There is a real need for good experimental data in this energy range, for almost all the elements considered in this report. The

only relatively complete data, those of Perkin,<sup>3</sup> are of poor quality. The inelastic gamma spectra tabulated for this energy range are generally based (in varying degree) on statistical theory. The parameters used in the theoretical model are at times questionable, and the unreliability increases as the incident neutron energy increases.

The lack of good experimental data becomes more acute when one considers high (>8 Mev) incident neutron energies. There are practically no data except perhaps at  $E_n = 14$  Mev. For  $E_n > 8$  Mev the tabulated spectra are almost exclusively derived from statistical theory. In general, the spectral shapes are given more reliably than the absolute magnitudes. Fortunately the problems to be run in this program use a fission source, which has only one-half of one per cent of the source above 8 Mev.

As tungsten is an element of major importance in the program, it was decided to obtain capture spectra for several tungsten isotopes. The spectra are based on those for natural tungsten, modified for each isotope according to its binding energy and the gamma transitions appropriate to the particular isotope. It remains outside the scope of this report to obtain the inelastic gamma spectra for the various tungsten isotopes.

The gamma-ray spectra following neutron interactions with uranium-235 and uranium-238 were separated into two parts. The gammas associated with nonfission capture events were treated in the same manner as those of the other elements.

A very complete study was made of the fission gammas – both prompt and delayed (from 0 to 10 hr). These data are believed to be accurate to within the 15% uncertainty claimed for the prompt radiation, and as such should be adequate for most practical applications. However, some uncertainties remain which might be

investigated. One stems from the assumption of a constant gamma spectrum for the first second, which may not be entirely correct. For example, the average prompt photon energy implied by Table 28 is  $\overline{E} = 0.53$  Mev/photon, whereas Table 30 implies that  $\overline{E}$  varies from 0.85 Mev at 1 second to about 1.0 Mev over the range from 10 seconds to about 5 minutes, after which it decreases to 0.6 Mev at 10 hr.

Also, the extrapolations performed to extend Table 30 from about 5 hr to 10 hr, while introducing relatively small contributions to the total energy release, could be important in computing heating rates for long times after shutdown. Hence these extrapolations should be examined more carefully if problems in these time regions are contemplated.

In general, the accuracy of both the nonfission and the fission gamma spectra are believed to be adequate for the problems to be considered in the tungsten nuclear rocket program.

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