

# NATIONAL BUREAU OF STANDARDS REPORT

9008

*R-64*

AN ELECTRIC QUADRUPOLE TRANSITION IN THE  
A  $^1\Pi \leftarrow X^1\Sigma^+$  SYSTEM OF CO

by

S. G. Tilford and J. D. Simmons

Technical Report

to

National Aeronautics and Space Administration  
Washington, D. C.

GPO PRICE \$ \_\_\_\_\_

CFSTI PRICE(S) \$ \_\_\_\_\_

Hard copy (HC) *1.00* \_\_\_\_\_

Microfiche (MF) *.50* \_\_\_\_\_

**N66-16698**  
(ACCESSION NUMBER)  
*20*  
(PAGES)  
*CR 70045*  
(NASA CR OR TMX OR AD NUMBER)

\_\_\_\_\_  
(THRU)  
*1*  
(CODE)  
*24*  
(CATEGORY)



# 653 July 65

U.S. DEPARTMENT OF COMMERCE  
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NBS PROJECT

221-11-2210461

November 30, 1965

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## Preface

A manuscript based upon this report will be submitted for publication as a Letter to the Editor in The Journal of Chemical Physics.

AN ELECTRIC QUADRUPOLE TRANSITION IN THE

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N 66-16698

Abstract

In the fourth positive system, A  ${}^1\Pi \leftarrow X {}^1\Sigma^+$ , of CO electric quadrupole transitions ( $\Delta J = 2$ ) have been observed for the first time in bands of an electric dipole allowed transition. The quadrupole transition probability has been determined to be  $\sim 1.1 \times 10^3 \text{ sec}^{-1}$ . This value compares well with the value of  $\sim 2.0 \times 10^3 \text{ sec}^{-1}$  determined for the quadrupole transition probability of the electric dipole forbidden a  ${}^1\Pi_g \leftarrow X {}^1\Sigma_g^+$  transition in the isoelectronic  $N_2$  molecule.

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Work supported in part by the Office of Naval Research and the National Science Foundation.

\*\* Work supported in part by the National Aeronautics and Space Administration.

During the reinvestigation of the high-resolution absorption spectra of the forbidden transitions of carbon monoxide in the vacuum ultraviolet <sup>1)</sup> several weak lines were observed at the head of each of the strong fourth positive system,  $A^1\Pi \leftarrow X^1\Sigma^+$ , bands. These lines cannot be identified with any of the known systems of CO which occur in this region;  $a^3\Pi \leftarrow X^1\Sigma^+$ ,  $a'^3\Sigma^+ \leftarrow X^1\Sigma^+$ ,  $e^3\Sigma^- \leftarrow X^1\Sigma^+$ ,  $d^3\Delta \leftarrow X^1\Sigma^+$ , or  $I^1\Sigma^- \leftarrow X^1\Sigma^+$ . Similarly, the extra lines cannot be correlated with any isotopic analogues of the fourth positive system since appreciable quantities of stable species lighter than  $C^{12}$  and  $O^{16}$  do not exist.

A possible explanation of these lines is that they should be associated with bands of the fourth positive system. For this to be true, the selection rule  $\Delta J = +2$  would have to be imposed. That such is the case is illustrated in Table I in which lines of the S(J) branch ( $\Delta J = +2$ ) calculated from the observed P(J) and R(J) branches (which terminate on the  $\Pi_+$  levels of the  $A^1\Pi$  state) are compared with observed transitions. The observed (P(J) and R(J) branches, rather than the rotational constants, were used to calculate the S(J) branch since there is an indication of a weak perturbation at low J values in the  $v' = 4$  level of the  $A^1\Pi$  state as a result of interaction with the  $v = 14$  level of the  $a'^3\Sigma^+$  state <sup>2)</sup>.

The  $A^1\Pi \leftarrow X^1\Sigma^+$  transition in  $O_2$  is an electric dipole transition. Since  $\Delta J = 2$  transitions are rigorously forbidden for electric dipole radiation<sup>3)</sup>, the observed lines must be attributed to electric quadrupole radiation. This is, therefore, the first example in which electric quadrupole lines have been observed for an electric dipole allowed transition.

The selection rules for electric quadrupole radiation are  $+\leftrightarrow +$ ,  $-\leftrightarrow -$ , and  $+\nleftrightarrow -$ , just the opposite of those for electric dipole radiation, thus, the S(J) branch terminates on the  $\Pi_+$  levels as do the P(J) and R(J) branches. An O(J) branch ( $\Delta J = -2$ ) should also be observed. However, because of the strength of the electric dipole (P(J), Q(J), and R(J) branches and also branches of the isotopic  $C^{13}O^{16}$  and  $C^{12}O^{18}$  molecules which occur in the same region, we have not been able to assign any O(J) lines. From the relative orders of magnitude of transition probabilities for electric dipole to magnetic dipole to electric quadrupole radiation ( $1 : 1.3 \times 10^{-5} : 4.1 \times 10^{-7}$  at  $1418 \text{ \AA}$  taking into account the  $v^2$  term in the quadrupole contribution)<sup>4)</sup>, P(J), Q(J), and R(J) branches resulting from magnetic dipole transitions should also be observed in this region. The branches would be similar to those observed for the stronger electric dipole transition except the (P(J) and R(J) branches would terminate on the  $\Pi_-$  levels and the Q(J) branch would terminate on the  $\Pi_+$  levels, just the opposite of those for electric dipole radiation. Since the A - type

doubling is almost negligible for bands of the fourth positive system the magnetic dipole branches will be very difficult to observe except in regions where strong perturbations of only one  $\Lambda$  - type component of the  $\Pi$  state are found.

In Fig. 1 the (4 - 0) band of the  $A^1\Pi \leftarrow X^1\Sigma^+$  transition is illustrated. These spectrograms were photographed in the 4th order of the 21-foot vacuum spectrograph at the Naval Research Laboratory. A microwave excited krypton lamp was used as a source of background continuum<sup>5)</sup>. The lower spectrogram in Fig. 1 was obtained with a path length of approximately  $4 \times 10^{-2}$  meter-atmos. The upper photograph required only approximately  $5 \times 10^{-7}$  meter-atmos. The extra lines beyond the head of the (4 - 0) A - X band in the high pressure spectrum in Fig. 1 are high J transitions in the (7 - 0) band of the  $e^3\Sigma^- \leftarrow X^1\Sigma^+$  system (this corresponds to the (6 - 0) band reported earlier<sup>2,6)</sup>). The weaker lines in the lower pressure spectrum belong to the (14 - 0) band of the  $a'^3\Sigma^+ \leftarrow X^1\Sigma^+$  system and the  $C^1\Sigma^+ \leftarrow X^1\Sigma^+$  (4 - 0) band of  $A^1\Pi \leftarrow X^1\Sigma^+$ .

The spectra were first observed with the CO gas admitted directly into the body of the spectrograph. However, upon measurement of the proposed S(J) lines, we found that the lines were shifted from the positions calculated by use of the selection rule  $\Delta J = +2$ , by amounts up to approximately  $1 \text{ cm}^{-1}$  toward longer wavelengths. This shift was caused by the variation of the refractive index of the gas



as a function of wavelength.

Near an absorption band the refractive index may be written

$$n^2 \propto \left( 1 + \sum_i \frac{\rho_i}{(\nu_i^2 - \nu^2) + \frac{g_i \nu^2}{(\nu_i^2 - \nu^2)}} \right) \quad (1)$$

where  $\nu_i$  corresponds to the band origin,  $\rho_i$  is a constant and  $g_i$  is related to the width of the absorption band or line.<sup>7)</sup> The summation takes into account all absorption bands. From equation (1) it can be seen that as the wavelength approaches that of a strong absorption band, the refractive index can change by a large amount in a short wavelength region. When  $\nu$  is greater than  $\nu_i$  the refractive index becomes less than unity and the observed wavelength is greater than the corresponding vacuum wavelength (this is just the reverse for wavelengths normally observed in the visible and near ultraviolet regions) in agreement with our observations. Subsequent experiments with an absorption cell placed in front of the slit verified that we were observing a refractive shift very close to the band origin.

The appearance pressure necessary to observe the S(J) branch was approximately  $9 \times 10^4$  greater than that necessary to observe the R(J) branch with an identical path length.

This value is not too discordant with the order of magnitude of  $2.4 \times 10^6$  indicated above for the ratio of transition probabilities of an electric dipole to an electric quadrupole transition. From the absolute values of transition probabilities for the fourth positive system recently determined by phase-shift techniques<sup>8)</sup> a value of approximately  $1.1 \times 10^3 \text{ sec}^{-1}$  can be inferred for the absolute quadrupole transition probability observed here. The corresponding lifetime for the quadrupole transition would be approximately  $9 \times 10^{-8}$  sec.

From a rotational intensity analysis, the ratio of 0.33 between the relative transition probabilities of electric quadrupole to magnetic dipole radiation has recently been determined for the electric dipole forbidden  $a \ ^1\Pi_g \leftarrow X \ ^1\Sigma_g^+$  transition in the iso-electronic  $N_2$  molecule.<sup>9)</sup> From this value and the absolute transition probability of  $5.9 \times 10^3 \text{ sec}^{-1}$  measured for the magnetic dipole  $a \ ^1\Pi_g \leftarrow X \ ^1\Sigma_g^+$  transition in  $N_2$  by a molecular beam technique,<sup>10)</sup> the absolute quadrupole transition probability in  $N_2$  is calculated to be  $\sim 2.0 \times 10^3 \text{ sec}^{-1}$ .

The quadrupole transition probability for the  $a - X$  transition in  $N_2$  is approximately 1.8 times greater than that for the  $A - X$  transition in  $CO$ . This is in good agreement with the observation that to observe the  $CO$  quadrupole lines it is necessary to employ approximately 2.5 times the path length that is required to observe the  $N_2$  quadrupole lines.

It should be pointed out that in both  $N_2$  and  $CO$  no consideration has been taken into account for any possible line width variation for branches resulting from the different types of radiations. In the case of the quadrupole rotation-vibration spectrum of  $H_2$ , extremely small line widths have been observed, even at very high pressures<sup>11, 12</sup>).

The only other example of an electric quadrupole transition which has been reported is the  $b\ ^1\Sigma_g^+ - a\ ^1\Delta_g$  transition in molecular oxygen at 1.9 microns in the infrared<sup>13</sup>). The reported transition moment of  $2.5 \times 10^{-3} \text{ sec}^{-1}$  would correspond to  $4.5 \times 10^{-1} \text{ sec}^{-1}$  had this transition been observed at 1400 Å. The  $B\ ^3\Pi_g^- - X\ ^1\Sigma_g^+$  transition in  $N_2$  has also been observed, but no pure quadrupole lines in this system have yet been detected<sup>14</sup>).

The ratio of intensity contributions from electric quadrupole to electric dipole radiation is given by<sup>3, 15, 16</sup>)

$$\frac{I_q}{I_d} = \frac{3 \Pi^2 \nu^2 S_q}{10 S_d} \cdot X^2 \quad (2)$$

where  $\nu$  is the frequency,  $S_q$  and  $S_d$  are the appropriate line strength factors for an electric quadrupole and electric dipole transition, respectively, and  $X^2$  is the ratio between quadrupole matrix elements and dipole matrix elements.

Upon substitution of the observed appearance pressures which are inversely proportional to the intensities, the appropriate line strength factors for the R(J) branch<sup>3</sup>) and the S(J) branch<sup>16</sup>) and a

value of  $70500 \text{ cm}^{-1}$  for the frequency, a value of  $2.5 \times 10^{-16}$  is obtained for the observed ratio of quadrupole to dipole matrix elements. A crude theoretical prediction give an order of magnitude of  $3.6 \times 10^{-16}$  for the ratio of quadrupole to dipole matrix elements<sup>4</sup>).

#### ACKNOWLEDGEMENT

The authors wish to express their appreciation to Mr. Vincent Franklin and Mr. Rudolph Naber for setting up the equipment and obtaining much of the experimental data.

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CAPTIONS TO FIGURES

Fig. 1 The (4 - 0) absorption band of the  $A^1\Pi \leftarrow X^1\Sigma^+$  transition in  $\text{CO}$ . (a) Low pressure spectrogram with the R(J) electric dipole branch illustrated. (b) High pressure spectrogram with the S(J) electric quadrupole branch illustrated.

TABLE 1

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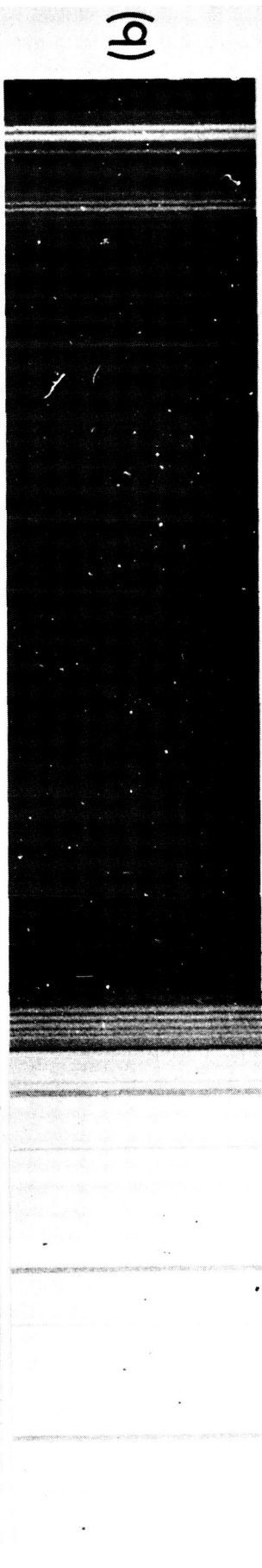
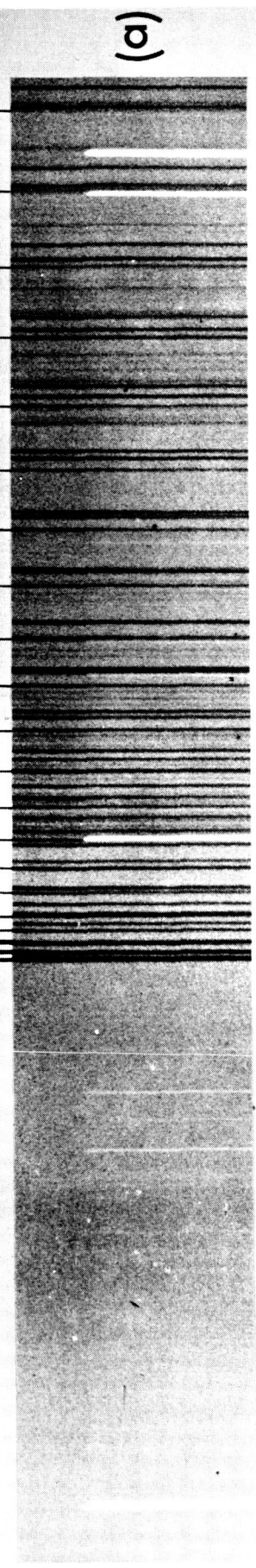
J	S(J) cm <sup>-1</sup>	
	calculated	observed
1	70481.00	70480.20*
2	485.25	485.24
3	488.77	488.78
4	491.37	491.52
5	493.30	493.41
6	494.31	494.26*
7	494.79	494.37*
8	494.01	493.87*
9	492.58	492.66
10	490.44	490.36
11	487.24	487.25
12	483.48	483.40

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\* blended

CO  $A^1\Pi - X^1\Sigma^+ (4-0)$

1418.969 Å



S

5 10

1 ← 1 Å → 1