*c* **<sup>t</sup>**

provided by NASA Technical Reports Server

**COPED**<br>CONTEGOI

U.Alabdms

## .r.\* **PRELiMdNARY** . On Dielectric **Lenses**

In electromagnetic wave propagation in parollel-plane waveguides, one is confronted with the problem of launching the waves. When **holrrr are** used, they introduce 'fringing **due** to their side walls. **For** pxample, the electric field lines which are parallel to the walls (TE<sub>OI</sub> mode) are curved instead of being straight as they are **suppored** *to* be. A method of eliminating this difficulty is *to* use a dielectric lens.

The dielectric lens has been used *to* modify or correct the phase of + an electromagnetic wave (1, **2)** as shown in Figure 1. The **phase** front is usually spherical due to a point **source, or** circular due to a line source.

y spherical due to a point source, or circular due to a line source.<br>The shape-of the lens is determined by the fact that the electrical  $\frac{1}{2}$ <br>ce travelled through the air plus that travelled through the lens is<br>to th distance travelled through the air plus that travelled through the lens is **\r.g**  equal to the electrical equivalent of the constant distance R.

**If** the **source** is a horn as shown in Figure **2,** a different situation arises. The flux plot is no longer perfectly circular, and consequently the problem **of** designing the lens must be handled, graphically **as** well as analytically. **GPO PRICE** *8* 

**From Figure 2, one can formulate the following <b>CFSTI PRICE(S)** 

 $ds^{2} = dx^{2} + dy^{2}$  $s = \int_{0}^{L} [1 + g^2(x)]^{1/2} dx$  Microform

where s is the actual distance the wave travels, and  $g'(x) = \frac{dy}{dx}$ . The function  $g(x)$  is determined numerically at as many points as one likes for the desired accuracy. The shape of the lens is then determined by the imposed condition that the distance a wave front travels in an empty waveguide plus the distance the same front tmvels in the dielectric filled portion

S. Silver, editor, Microwave Antenna Theory and Design , McGraw-Hill **bok** Co., New **Yo&,** (1949). **1** 

Kraus, J. D., Antennas, McGraw-Hill Book Company, New York, (1950).



**SOD nnod NI-IIJW** 

.

Microfiche (MF).





of the guide is **eqwl** *to* the constant electrical disiunce **R.** Thiscodition is expressed by the following design equation.

$$
\frac{1}{x_{g}} \int_{0}^{x_{0}} [1 + f'^{2}(x)]^{1/2} dx + \frac{1}{x_{d}} (L - X_{o}) = \frac{1}{x_{g}} R
$$

where

**L** 

. **1** 

$$
\lambda_{g} = \frac{\lambda_{0}}{\sqrt{1 - (c/\rho)^{2}}}
$$

$$
\lambda_d = \frac{\lambda_o}{\sqrt{\epsilon_r - (f_c/\rho)^2}}
$$

**f** = **operating** frequency  $\lambda$  = free space wavelength at f  $f_c = \cot$ -off frequency of the guides = relative dielectric constant **r** 

If 
$$
n = \frac{9}{\lambda_d}
$$
, the above equation becomes

$$
\int_{0}^{x} [1 + {f'}^{2}(x)]^{1/2} dx + n (L - X_{0}) = R
$$

The position  $X_{\mathbf{O}}$  is determined by trial and error such that the design equation is satisfied. The equation may be **made** frequency independent if dielectrics with relatively high die1ectric~'constonts **are used.** 



**N. F.** Audeh H.Y.Yee University of **Alaboma**  Research Institute **Huntsvllle,** Alabama



