Some Aspects of Difference Equations
by
James Hurt
B.S., University of Iowa, 1961
M.S., University of Iowa, 1963

Thesis
submitted in partial fulfillment of the requirements for the Degree of Doctor of Philosophy in the Division of Applied Mathematics at Brown University

June, 1967
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is.accepted in its present form by the Division of Applied Mathematics as satisfying the thesis requirement for the degree of Doctor of Philosophy

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Date:

Date:

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## Vita

The author of this thesis was born on He graduated from University High School in Iowa City in June of 1957. He attended the University of Iowa, Iowa City, Iowa, from 1957 until 1963, receiving a Bachelor of Science degree (with distinction) in Mechanical Engineering in June of 1961 and a Master of Science degree In Engineering Mechanics in August of 1963. While attending the University of Iowa, he was elected to membership in the Pi Tau Signa and Tau Beta Pi national honorary fraternities and in the Society of the Sigma Xi.

Durning the 1963-1964 school year, the author was a full time Instructor at the University of Iowa. During this period he co-authored with Royce Beckett the book Numerical Calculations and Algorithms, published by McGraw-Hill, New York, 1967.

In the fall of 1964, the author entered Brown University on a Brown University fellowship and received NASA Traineeships for the following two years while working for his Ph.D. degree. During this time he wrote a paper, "Some Stability Theorems for Difference Equations", which has been submitted to the SIAM Journal for Numerical Analysis for publication.

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## Abstract

In recent years there has been a great interest in the study of methods for computing the solution of an initial value problem numerically. Most methods used form a difference analog to the differential equation and use the solution of the difference analog as an approximation for the solution of the differential equation. Most of the work done to date is concerned with estimating the difference between the true solution and the computed solution on a finite interval. Such problems as determining stability properties of the differential equation or finding a periodic solution are intrinsically involved with unbounded intervals. When a difference analog is used in these cases, the question becomes not one of how accurate is the approximation on a finite interval but do the solutions of the difference analog display the same properties as do the solutions of the differential equation.

The questions concerning stability properties of the difference analog led to research which is reported in a paper by the author entitled "Some Stability Theorems for Difference Equations". This paper has been submitted to the SIAM Journal on Numerical Analysis for publication. The theorems in this paper have application to other areas of Numerical Analysis, notably iteration theory, and to sampled data control systems.

In this thesis, the problem of finding periodic, almost periodic, and bounded solutions of a difference analog to a differential equation are considered. In the introduction, several examples are given to show how the behavior of a difference analog can be different from that of the differential equation. The problem of periodic solutions of the difference analog is discussed and, in order to have the period depend continuously on the parameters.
in the difference analog, the solutions must be extended from the integers to the reals. Definitions of and theorems related to solutions on the reals are given.

The properties of solutions of Linear and Affine difference equations are discussed and a noncritical case is treated in some detail. The noncritica] case considered is when the linear part has a constant matrix which has no eigenvalues on the unit circle. A theorem is presented concerning the existence of a periodic, almost periodic, or bounded solution of a noncritical difference equation when the nonlinearity has a small parameter.

The simplest critical case, where the matrix of the linear part is the identity matrix, is discussed and a theorem is given concerning the existence of a $T$-periodic solution where $T$ is a rational number. Finally, an averaging method and an invariant manifold method are given, each with an application which proves that, under certain conditions, certain difference analogs do display the same behavior as the corresponding differential equation. The Crude Euler analog of $\dot{x}=\epsilon f(t, x)$ has an almost periodic solution for $\epsilon$ and $h$ small enough under the same conditions which insure the existence of an almost periodic solution of the differential equation. The central difference analog to the Van der Pol equation $\ddot{x}+2 \epsilon \dot{x}\left(x^{2}-1\right)+x=0$ has an invariant curve of radius approximately 2 for $\epsilon>0$ and small.

## INIRODUCTION.

Consider the initial value problem given in (1) where $x$ and

$$
\begin{equation*}
\dot{x}=f(t, x) \quad x\left(t_{0}\right)=x_{0} \tag{1}
\end{equation*}
$$

$f(t, x)$ are $n$-vectors (n-dimensional vectors), $t$ is a real number called time and $\dot{x}=d x / d t$. In recent years there has been a great interest in the study of methods for computing the solution to (1) numerically [1,2]. Most numerical methods divide time into increments with spacing $h$ such that $t_{k}=t_{0}+k h$ and form the difference analog (2).

$$
\begin{align*}
& x\left(t_{k+1}\right)=g\left(t_{k}, x\left(t_{k}\right), x\left(t_{k-1}\right), \ldots, x\left(t_{k-N}\right) ; h\right)  \tag{2}\\
& x\left(t_{k}\right)=x_{k} \quad, \quad 0 \leqq k \leqq N
\end{align*}
$$

The solution to the difference equation (2) is then taken as an approximation to the solution of (1). However, most of the work to date is concerned with computing the solution on a finite interval, say for $t_{0} \leqq t \leqq t_{1}$. Such problems as determining stability properties of (l) or finding periodic solutions of (1) are intrinsically involved with the unbounded intervals $t_{0} \leqq t<\infty$ or $-\infty<t<\infty$. When the difference analog (2) is used in these cases, the question becomes not one of how accurate is the approximation on .
a finite interval but do the solutions of (2) display the same properties as the solutions of (1)?

Several very simple cases will be covered here as a partial answer to this question. Consider first the difference equation (3).

$$
\begin{equation*}
x(k+1)=A x(k) \tag{3}
\end{equation*}
$$

where

$$
A=\left(\begin{array}{ll}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{array}\right) \quad \dot{x}(k)=\binom{x_{1}(k)}{x_{2}(k)}
$$

If the eigenvalues of the real matrix $A$ are $\lambda_{1}$ and $\lambda_{2}, \lambda_{1} \neq \lambda_{2}$, then all solutions of (3) are of the form

$$
x(k)=c \lambda_{1}^{k}+d \lambda_{2}^{k}
$$

where $c$ and $d$ are constant 2-vectors. The origin $x(k)=0$ is also a solution. The behavior of the solutions near the origin is completely determined by the values of $\lambda_{1}$ and $\lambda_{2}$. If $\left|\lambda_{1}\right|<1$ and $\left|\lambda_{2}\right|<1$, then all solutions approach the origin as $k \rightarrow \infty$. If $\lambda_{1}$ and $\lambda_{2}$ are real, then the origin is said to be a stable node (SN). If $\lambda_{1}$ and $\lambda_{2}$ are complex, then the origin is said to be a stable focus (SF). If $\left|\lambda_{1}\right|=\left|\lambda_{2}\right|=1$, then all solutions remain bounded and the origin is said to be a center (C). If $\left|\lambda_{1}\right|>1$ and $\left|\lambda_{2}\right|>1$, then all solutions become unbounded as $k \rightarrow \infty$. If $\lambda_{1}$ and $\lambda_{2}$ are real, the origin is said to be an unstable node (UN) and if
the eigenvalues are complex, the origin is said to be an unstable focus (UF). If $\left|\lambda_{1}\right|>1$ and $\left|\lambda_{2}\right|<1$, then the origin is said to be a saddle point (SP). The terms stable node, stable focus, center, etc., come from differential equations and are used here because the behavior of the solutions of the difference equation is very similar to the behavior of solutions of a differential equation.

For example, consider the difference equation (3) when the matrix
A is given by

$$
A=\left(\begin{array}{cc}
0 & 1  \tag{3a}\\
-1 & 2
\end{array} \cos \omega\right)
$$

The eigenvalues of $A$ are $\lambda_{1}=e^{i \omega}$ and $\lambda_{2}=e^{-i \omega}$, and so the origin is a center. Consider the functional $\mathrm{V}(\mathrm{x})$

$$
v(x)=x_{1}^{2}-2 \cos \omega x_{1} x_{2}+x_{2}^{2}
$$

We see immediately that $V(x(k+1))=V(x(k))$ for any solution $x(k)$ and hence that any solution remains on a level surface of $V$. A level surface of $\dot{V}$ is an ellipse and hence the name center for this type of behavior.

The eigenvalues of A depend entirely on two numbers, $\mathrm{B}=$
$+\frac{1}{2} \operatorname{trace}(A)=+\frac{1}{2}\left(a_{11}+a_{22}\right)$ and $C=\operatorname{det}(A)=a_{11} a_{22^{-a}}^{12^{a}} 21$. Figure 1 is a graph showing the various regions of behavior in the $B, C$ plane. It should be noted that, on the line $B^{2}=C, \lambda_{1}=\lambda_{2}$ and there is the possibility of A not having simple elementary divisors. However, since this line is always a transition line between two different types of behavior; this raises no serious question at this time.

Using the results shown in Figure 1 , several difference analogs of the differential equation (4) are studied.

$$
\begin{equation*}
\ddot{x}+2 d \dot{x}+x=0 \tag{4}
\end{equation*}
$$

In most cases, the differential equation (4) will be written in the vector form of (4').

$$
\begin{align*}
& \dot{x}_{1}=\bar{x}_{2} \\
& \dot{x}_{2}=-x_{1}-2 d x_{2}
\end{align*}
$$

The behavior of the solutions of (4) depend only on the parameter d. The origin is a center if $d=0$, a stable focus for $0<d<1$ and a stable node for $1<d<\infty$. The origin is an unstable focus for $-1<d<0$ and an unstable node for $-\infty<d<-1$. All the difference analogs will have an additional parameter, the spacing $h$, and the regions of the $d, h$ plane for each type of behavior for the difference analog should be compared to this ideal. The variations will be noted.

The first numerical method to be considered is the central difference analog where the derivatives of $x$ in (4) are replaced by their central difference approximations.

$$
\frac{1}{h^{2}}\left(x\left(t_{k+1}\right)-2 x\left(t_{k}\right)+x\left(t_{k-1}\right)\right)+\frac{2 d}{2 h}\left(x\left(t_{k+1}\right)-x\left(t_{k-1}\right)\right)+x\left(t_{k}\right)=0
$$

Letting $y_{1}(k)=x\left(t_{k-1}\right)$ and $y_{2}(k)=x\left(t_{k}\right)$, this becomes the difference. equation (5)

$$
\begin{equation*}
y(k+1)=A y(k) \tag{5}
\end{equation*}
$$

where

$$
y(k)=\binom{y_{1}(k)}{y_{2}(k)} \quad A=\left(\begin{array}{cc}
0 & 1  \tag{6}\\
-\frac{1-h d}{1+h d} & \frac{2-h^{2}}{1+h d}
\end{array}\right)
$$

The parameters are

$$
B=\frac{1}{2} \frac{2-h^{2}}{1+h d} \quad \text { and } \quad C=\frac{1-h d}{1+h d} .
$$

The regions of behavior are shown in Figure 2. Notice that for $h>2$, the behavior is always that of a saddle point, a behavior which does not at all resemble any of the behaviors of solutions of the differential equation. However, the main region of interest is for $h$ very small. Figure 2 shows the behavior of solutions of (5) for large $h$ because this gives a clearer picture of what happens for small $h$.

For small $h$, we notice that the origin is a center when $d=0$, is stable when $d>0$ and unstable when $d<0 \ldots$ in agreement with the origin of the differential equation. However, note that the transition from Focal behavior to Nodel behavior does not occur at $d= \pm 1$ but on the ellipse $4 d^{2}+h^{2}=4$. Still, from this point of view, the central difference analog is a good method for computing a solution to (4) numerically.

The next method considered is the Crude Euler method [3]. This method uses equation (4) and is described by the difference equation (5)


Figure 2
where

$$
y(k)=\binom{x_{1}\left(t_{k}\right)}{x_{2}\left(t_{k}\right)} \quad A=\left(\begin{array}{cc}
1 & h  \tag{7}\\
-h & 1-2 h d
\end{array}\right)
$$

The parameters are $B=1-h d$ and $C=1-2 h d+h^{2}$. The regions of behavior are shown in Figure 3. Notice that, for $h>2$, the behavior is always unstable. Also, for small $h$ and $d=0$, the behavior is not that of a center but that of an unstable focus. If $h$ is small and $0<d<\frac{1}{2} h$, the behavior of the difference equation is that of an unstable focus instead of a stable focus. For large $d$, it is possible to get unstable behavior of the difference equation by having $h$ large enough to be in the saddle point region even though it is still small. In order to get the desired behavior in this case, we must take $h$ much smaller than $2 d$ if $d$ is small and much smaller than $1 / d$ if $d$ is large. Notice that, for $d=0$, the Crude Euler method behaves like an unstable focus for any $h>0$ and not like the center for the differential equation. From this point of view, this method is not very acceptable. However, other considerations must be used in passing Judgement on any given method.

In the Corrected Euler method [3], the value of $y(k+1)$ obtained from the Crude Euler method is used as a first approximation to the solution at $t_{k+1}$. This first approximation is used along with the Trapezoidal rule of integration to obtain a second approximation which is used as the numeric solution. This difference analog of (4') is given by (5) where


Figure 3

Regions of Behavior - Crude Euler Method

$$
y(k)=\binom{x_{1}\left(t_{k}\right)}{x_{2}\left(t_{k}\right)} \quad A=\left(\begin{array}{ll}
1-\frac{1}{2} h^{2} & h(1-h d)  \tag{8}\\
-h(1-h d) & 1-\frac{1}{2} h^{2}-2 h d(1-h d)
\end{array}\right)
$$

The parameters are $B=1-\frac{1}{2} h^{2}-h c(1-h c)$ and $C=\left(1-\frac{1}{2} h^{2}\right)^{2}-2 h c(1-h c)\left(1-\frac{1}{2} h^{2}\right)+$

- $h^{2}(1-h c)^{2}$. The regions of behavior are shown in Figure 4. Notice that, while the line of centers is no longer a straight line and comes much closer to the h -axis, this figure is not much different from Figure 3 for the Crude Euler method and that the same comments apply.

The final method discussed here is the Iterated Euler method [3], sometimes called the Modified Euler method. In this method, the trapezoidal rule of integration is used to obtain second, third, etc., approximations to the solution at $t=t_{k+1}$ until these iterations converge. The difference analog of (4 $4^{\circ}$ ) is given by equation (5) where

$$
y(k)=\binom{x_{1}\left(t_{k}\right)}{x_{2}\left(t_{k}\right)} \quad A=\left(\begin{array}{ll}
\frac{4+4 h d-h^{2}}{4+4 h d+h^{2}} & \frac{4 h}{4+4 h d+h^{2}}  \tag{9}\\
-\frac{4 h}{4+4 h d+h^{2}} & \frac{4-4 h d-h^{2}}{4+4 h d+h^{2}}
\end{array}\right)
$$

The parameters are

$$
B=\frac{4-h^{2}}{4+4 h d+h^{2}} \quad \text { and } \quad C=\frac{\left(4+h^{2}\right)^{2}-16 h^{2} d^{2}}{\left(4+4 h d+h^{2}\right)^{2}}
$$

The regions of behavior are shown in Figure 5. The only discrepancy between this figure and the ideal figure is the saddle point region where there should be an unstable node. However, for $d<-1$, if we take $h$ much smaller than

CODE:

$$
\begin{aligned}
\mathrm{h} & =\text { Step size } \\
\mathrm{d} & =\text { Damping coefficient } \\
\mathrm{SN} & =\text { Stable Node } \\
\mathrm{SF} & =\text { Stable Focus } \\
\mathrm{Cen} & =\text { Center } \\
\mathrm{UN} & =\text { Unstable Node } \\
\mathrm{UF} & =\text { Unstable Focus } \\
\mathbf{S P} & =\text { Saddle Point }
\end{aligned}
$$



Figure 4
Regions of Behavior - Corrected Euler Method


Figure 5
$-1 / d$, then the behavior will be correctly that of an unstable node.
The significance of these examples is that, in most cases, the behavior of solutions of a difference analog to a differential equation is not the same as the behavior of solutions of the differential equation. In the case of a more general differential equation, this variation of behavior can be detected only by computing the solution twice, once with $h=h_{0}$ and once with a considerably different $h$, say $h=\frac{1}{2} h_{0}$ or $h=\frac{1}{10} h_{0}$. If the behavior of the solutions does not change drastically from one computed solution to another, then one may assume that the behavior of solutions of the difference analog and of the differential equation are the same. If the behavior does change drastically, even if the basic type of behavior doesn't change, then one must assume that the behavior of solutions of the difference analog and of the differential equation are different and the solution should be recomputed with another value of $h$. For example, if $d=0$ in equation (4) and we were using the Crude Euler method to compute a solution, we would note that the computed solution is spiralling outward like an unstable focus. When the solution is recomputed with half the original spacing, we still see the outward spiral but its rate of spiralling out will be reduced by about half, leading to the conclusion that the behavior of the differential equation is not that of an unstable focus.

A greater analysis of the center of a difference equation is in order. Consider the difference equation (3) where the matrix $A$ is as given in (3a). This happens to be the matrix $A$ given for the central difference analog in equation (6) with $d=0$ and $2 \cos \omega=2-h^{2}$, but we shall use it in the form given in (3) and (3a). As noted before, all solutions remain on a level surface of $V(x)=x_{1}^{2}-2 \cos \omega x_{1} x_{2}+x_{2}^{2}$. One such level surface, for
$\omega=2 \pi / 9$ and $x_{1}(0)=x_{2}(0)=1$, is the ellipse shown in Figure 6. The ticmarks on the ellipse show the points of the solution starting at $x_{1}(0)=x_{2}(0)$ $=1$. Since $x(9)=x(0)$, this solution is periodic with period 9. Any point on this ellipse can be used as the starting point $x(0)$ and will lead to a 9-periodic solution.

Suppose now that $\omega$ was changed slightly, say from $\omega=2 \pi / 9=$ $10 \pi / 45$ to $\omega=10 \pi / 46$. The ellipse $V(x)=$ constant will change very slightly but each solution will now take on 46 distinct points on the ellipse instead of 9 and we will have $x(46)=x(0)$, leading us to say that the solution is 46 -periodic. A change in the parameter $\omega$ of $\pi / 207 \doteq .0015$ changes the period by a factor of 5 . In fact, if $\omega$ is an irrational multiple of $\pi$, there will not be any integer $k$ such that $x(k)=x(0)$ and so the solution is not periodic. Yet, each point $x(k)$ of the solution will be on the level surface of $V(x)=V(x(0))$ and the points $x(k)$ will move around the ellipse in the same manner as shown in Figure 6.

Assume that the eigenvalues of $A$ are $e^{i \omega}$ and $e^{-i \omega}$ where $e^{i \omega} \neq e^{-i \omega}$. Then each solution of (3) can be written as in (10)

$$
\begin{equation*}
x(k)=a \cos \omega k+b \sin \omega k \tag{10}
\end{equation*}
$$

where $a$ and $b$ are constant 2 -vectors which are completely determined by $x(0)$. Letting $k=0$, we get $a=x(0)$. We determine $b$ from the equation. $x(1)=A x(0)$.

$$
b=\frac{(A-\cos \omega I)}{\sin \omega} x(0)
$$

Since $e^{i \omega} \neq e^{-i \omega}, \quad \sin \omega \neq 0$ and $b$ is uniquely determined.


Figure 6
A Level Surface of $V(x)=x_{1}^{2}-2 \cos (2 \pi / 9) x_{1} x_{2}+x_{2}^{2}$

Now, $x(k)=x(0)$ only when $\cos \omega k=1$, i.e., when $\omega k$ is some integer multiple of $2 \pi$. Then $\omega=2 \pi / 9$, then $9 \omega=2 \pi$ and the solution is 9-periodic. When $\omega=10 \pi / 46$, then $46 \omega=10 \pi$ and the solution is 46 -periodic. Notice that, in both cases, the period $k$ is the lowest integer such that $x(k)=x(0)$. When $\omega$ is an irrational multiple of 2 , then $\omega k$ is never equal to an integer multiple of $2 \pi$ and the solution is not periodic. Yet equation (10) gives the formula for a $2 \pi / \omega$-periodic function. Why can't we say that the solution (10) is $2 \pi / \omega$-periodic for any $\omega$ ? If we could, then the period would be continuous in the parameter $\omega$ (except when $e^{i \omega}=e^{-i \omega}$ ). The problem arises because the solution is defined only on the integers. What is meant by the statement, " $x$ is T-periodic" when $x$ is defined only on the integers and $T$ is not an integer? Since $x(k+T)$ is not defined whenever $k$ is an integer and $T$ is not, we cannot say $x$ is T-periodic whenever $x(k+T)=x(k)$. If we could define equation (10) as a solution of (3) for all real $k$, then we could say $x$ is $T$-periodic whenever $x(k+T)=x(k)$. However, (10) is a solution of (3) not only for $a$ and $b$ constant 2-vectors, but also for $a$ and $b$ any l-periodic functions, as in ( $10^{\prime}$ ).

$$
x(k)=a(k) \cos \omega k+b(k) \sin \omega k
$$

It is easily confirmed by substitution that $a$ and $b$ are no longer determined by $x(0)$ but that $x(k)$ is arbitrary for $0 \leqq k<1$. If $x(k)$ is defined (as initial value) for $0 \leqq k<1$, the $a(k)$ (hence $b(k)$ ) is completely determined and ( $10^{\prime}$ ) represents the solution with the initial values. Notice that, by defining the initial values for the solution on an interval
instead of at a point, we get the solution defined for all real $k$ instead of only on the integers and we can logically talk about T-periodic solutions for non-integer $T$. In (10') the vector $b(k)$ is given in terms of $a(k)$ by

$$
b(k)=\frac{(A-\cos \omega I)}{\sin \omega} a(k) \quad 0 \leqq k<1
$$

Just what does this "new" definition of a solution mean in terms of the simple difference equation (3)? Two different solutions of (3) for $\omega=2 \pi / 9$ are shown in Figure 7. The solid line is the solution which happens to have $a(k)$ and $b(k)$ constant and the dotted line has $a(k)$ and $b(k)$ non-constant. Notice that, on the integers (marked by the tics), both solutions look the same and that both solutions are 9-periodic.

This definition of a solution of a difference equation is formalized here. To emphasis that the solution is defined on some real interval and not on only the integers, $t$ is used as the independent variable instead of $k$. The difference equation under consideration is given by equation (11) where each $x$ and $f(t, x)$ is in $E^{n}$, an $n$-dimensional vector space.

$$
\begin{equation*}
x(t+1)=f(t, x(t)) \tag{11}
\end{equation*}
$$

DEFINITION: A function $x(t)=x\left(t ; t_{0}, x_{0}\right)$ where $t$ and $t_{0}$ are real numbers, $t_{0} \leq t<t_{0}+T$, some $T \geqq 1$, and $x_{0}$ is a $n$-vector valued function on $[0,1)$, is called a solution of the difference equation (11) if, for some $T \geqq 1$,

$$
\text { a) } x(t)=x_{0}\left(t-t_{0}\right) \text { for } t_{0} \leq t<t_{0}+1
$$



Figure 7
Two Solutions of $x(k+1)=\operatorname{Ax}(k)$

$$
\text { b) } x(t+1)=f(t, x(t)) \text { for } t_{0} \leqq t<t_{0}+T-1
$$

The difference between this definition of a solution and the usual definitions is that the initial value $x_{0}$ is now a function instead of a vector. The function $x_{0}$ defines the solution on the unit interval $\left[t_{0}, t_{0}+1\right.$ ) and the difference equation (11) is used to extend this definition to the right.

Notice that the solution is defined for all real $t$ in the interval $\left[t_{0}, t_{0}+T\right.$ ] instead of restricting $t-t_{0}$ to be an integer. Notice also that the solution space, which is finite dimensional when the solution is defined only on the integers, is now infinite dimensional. An initial function must be defined instead of an initial point.

For any real number $t$, the integer part of $t$, $[t]$, is the largest integer $n \leqq t$. The fractional part of $t,\langle t\rangle$, is given by the formula $\langle t\rangle=t-[t]$. Thus, as examples, $[3.14]=3,\langle 3.14\rangle=.14,[-1.62]=-2$, $\langle-1.62\rangle=.38,[0]=0$, and $\langle 0\rangle=0$. Let $R$ be the real line, $R^{+}$be the positive real half-line, and $E^{n}$ be an $n$-dimensional normed vector space. For each vector $x \in E^{n}$, denote the norm of $x$ by $|x|$. The most common norm used here will be the Euclidean norm $|x|=\left(x^{*} x\right)^{\frac{1}{2}}$ ( $x^{*}$ denotes the complex conjugate transpose of $x$ ) but discussion will not necessarily be limited to this norm.

For any function $x$ mapping $R$ (or $R^{+}$) into $E^{n}$, let $x_{t}$ be the function defined by $x_{t}(\theta)=x(t+\theta)$ for $0 \leqq \theta<1$. Let $X_{1}$ be a Banach space of $n$-vector valued functions defined on $[0,1)$. For each $x_{t} \in X_{1}$, denote the norm of $x_{t}$ by $\left\|x_{t}\right\|$. The usual space for $X_{1}$ will be $L_{2}[0,1)$, but discussion will not necessarily be limited to this space.

Let $E$ be the operator defined by $E x(t)=x(t+1)$. Then the difference equation (11) can be written in the somewhat shorter notation of
(11). The arguments on $x$ and Ex are understood to be $t$.

$$
\begin{equation*}
E x=f(t, x) \tag{119}
\end{equation*}
$$

THEOREM 1 On existence and uniqueness of solutions. Suppose $t_{o}$ and an $x_{t_{0}} \in X_{1}$ are given. If there are open, possibly unbounded, sets $G(t) \subset E^{n}$ for all $t \geqq t_{0}$ such that $f(t, x) \in G(t+1)$ for all $x \in G(t)$, all $t \geqq t_{0}$ and $x_{t_{0}}(\theta) \in G\left(t_{0}+\theta\right)$ for $0 \leqq \theta<1$, then the solution $x(t)=x\left(t ; t_{0}, x_{t_{0}}\right)$ exists for all $t \geqq t_{0^{\circ}}$ If $f(t, x)$ is single valued, the solution is unique.

PROOF: The proof is by induction. The solution exists and is unique for $t_{0} \leq t<t_{0}+1$ (defined by the initial function $x_{t_{0}}$ ). For each integer $N>0$, assume that the solution $x(t)$ exists, is unique and is in $G(t)$ for $t_{0} \leqq t<t_{0}+N$. Then the difference equation (11) gives the solution on $t_{0} \leqq t<t_{0}+N+1$. By assumptions, $x(t+1)=f(t, x(t)) \in G(t+1)$ since $x(t) \in G(t)$ for $t_{0}+N-1 \leqq t<t_{0}+N$ and $x(t+1)$ is defined and unique for $t_{0}+N-1 \leqq t<t_{0}+N_{.}$. . By induction, the solution exists and is unique for all $t \geqq t_{0}$ 。

It should be noted here that existence comes from the assumption that $f(t, x) \in G(t+1)$ for all $x \in G(t)$ and uniqueness comes from the (trivial?) assumption that $f(t, x)$ is single valued for all $x \in G(t)$. The sets $G(t)$ are open and possibly unbounded. The most common sets used in this work are $G(t)=E^{n}$ for all $t \geqq t_{0}$, but discussion is not necessarily limited to these sets. The fact that the sets are open play no vital role in this theorem. Note also that the theorem holds for each $\mathrm{x}_{\mathrm{t}_{0}}$ which satisfies the conditions and is not limited to a specific given $x_{t_{0}}$.

For the function $f(t, x)$, let $f_{t}(y)$, where $y$ is any function from $[0,1)$ into $E^{n}$, be the function defined by

$$
f_{t}(y)(\theta)=f(t+\theta, y(\theta)), \quad 0 \leqq \theta<1
$$

The following three theorems each show that $X_{t} \in X_{1}$ under proper conditions on $X_{t_{0}}$ and $f$ for three different choices of the space $X_{1}$. The statements of the theorems and their proofs are almost identical. The difference in these proofs, and the difficulty for a general Banach function space, is the proof that, if $x_{t_{0}} \in X_{1}$ and $x_{t_{0}+1} \in X_{1}$, then $X_{t} \in X_{1}$ for $t_{0} \leqq t \leqq t_{0}+1$. In Theorem $2 a$, for example, this is not true unless $x_{t_{0}}$ satisfies a certain boundary condition.

THEOREM 2a If, in addition to the assumptions of Theorem 1 , the space $X_{1}$ is the space of all continuous functions $C[0,1)$ with the uniform norm, $f_{t}(y) \in C[0,1)$ for each $y \in C[0,1)$ such that. $y(\theta) \in G(t+\theta), 0 \leqq \theta<1$, and each $t \geqq t_{0}$, and the initial function $x_{t_{0}}$ is in $C[0,1)$ and satisfies the boundary condition

$$
\lim _{\theta \rightarrow 1} x_{t_{0}}(\theta)=f\left(t_{0}, x_{t_{0}}(0)\right)
$$

then the solution $x(t)=x\left(t ; t_{0}, x_{t_{0}}\right)$ satisfies $x_{t} \in C[0,1)$ for each $t \geqq t_{0}$.

PROOF: The boundary condition on $x_{t_{0}}$ insures that the solution $x(t)$ is continuous at $t=t_{0}+1$. Since $x_{t_{0}+1}=f_{t_{0}}\left(x_{t_{0}}\right)$ is in $C[0,1)$, the solution is continuous for $t_{0} \leqq t<t_{0}+2$, i.e., $x_{t} \in C[0,1)$ for $t_{0} \leqq t<t_{0}+1$.

If $x_{t} \in \mathbb{C}[0,1)$ for $t_{0} \leqq t<t_{0}+N$, then, using the fact that $x_{t+1}=f_{t}\left(x_{t}\right)$ is in $C[0,1)$ whenever $x_{t} \in C[0,1)$, we get that $x_{t} \in C[0,1)$ for $t_{0} \leqq t<t_{0}+N+1$, and the theorem is proven by induction.

THEOREM Lb If, in addition to the assumptions of Theorem 1 , the space $X_{1}$ is the space of all $p$-th power integrable functions $L_{p}[0,1)$ with the usual norm, some $p \geqq 1, f_{t}(y) \in I_{p}[0,1)$ for each $y \in L_{p}[0,1)$ such that $y(\theta) \in G(t+\theta), 0 \leqq \theta<1$, and each $t \geqq t_{0}$, and $x_{t_{0}} \in L_{p}[0,1)$, then the solution $x(t)=x\left(t ; t_{0}, x_{t_{0}}\right)$ satisfies $x_{t} \in L_{p}[0,1)$ for each $t \geqq t_{0}$.

PROOF: We have $x_{t_{0}} \in L_{p}[0,1)$ and $x_{t_{0}+1}=f_{t_{o}}\left(x_{t_{o}}\right) \in L_{p}[0,1)$. Then, for any $t, t_{0} \leq t<t_{0}+1$,

$$
\begin{aligned}
\left\|x_{t}\right\|^{p} & =\int_{t}^{t+1}|x(s)|^{p_{d s}}=\int_{t}^{t_{0}+1}|x(s)|^{p} d s+\int_{t_{0}+1}^{t+1}|x(s)|^{p_{d s}} \\
& \leqq \int_{t_{0}}^{t_{0}+1}|x(s)|^{p_{d s}}+\int_{t_{0}+1}^{t_{0}+2}|x(s)|^{p} d s=\left\|x_{t}\right\|^{p}+\left\|x_{t_{0}+1}\right\|^{p}
\end{aligned}
$$

and hence $x_{t} \in L_{p}[0,1)$ for $t_{0} \leqq t<t_{0}+1$.
If $x_{t} \in L_{p}[0,1)$ for $t_{0} \leqq t<t_{0}+N$, then, using the fact that $x_{t+1}=f_{t}\left(x_{t}\right)$ is in $L_{p}[0,1)$ whenever $x_{t} \in L_{p}[0,1)$, we get that $x_{t} \in L_{p}[0,1)$ for $t_{o} \leqq t<t_{0}+N+1$ and the theorem is proven by induction.

Krasnosel'skii [8] gives a necessary and sufficient condition for $f_{t}(y)$ to be in $L_{p}[0,1)$ for all $y$ in $L_{p}[0,1)$ such that $y(\theta) \in G(t+\theta)$, $0 \leqq \theta<1$. This condition is given here without proof, $f_{t}(y)$ is in $L_{p}[0,1)$ if and only if there exists a non-negative valued function $a(\theta), 0 \leqq \theta<1$, with $\int_{0}^{l} a(\theta)^{p} d \theta<\infty$, and a non-negative constant $b$ such that

$$
|f(t+\theta, x)| \leqq a(\theta)+b|x|
$$

for all $x \in G(t+\theta), 0 \leqq \theta<1$. This condition is interesting for the case where $G(t+\theta)$ is unbounded because, while any linear function of $x$ satesflies it, there are large classes of non-linear functions of $x$ which do not. For the case where $G(t+\theta)$ is bounded for all $\theta, 0 \leqq \theta<1$, this condition reduces to the condition that $|f(t+\dot{\theta}, x)|$ be bounded for all $x \in G(t+\theta)$, $0 \leqq \theta<1$.

THEOREM Lc If, in addition to the assumptions of Theorem 1 , the space $X_{1}$ is the space of all functions of bounded variation $\operatorname{BV}[0,1)$ with the usual norm, $f_{t}(y) \in \operatorname{BV}[0,1)$ for each $y$ in $\operatorname{BV}[0,1)$ such. that $y(\theta) \in G(t+\theta), 0 \leqq \theta<1$, each $t \geqq t_{0}$, and $x_{t_{0}} \in \operatorname{BV}[0,1)$, then the solution $x(t)=x\left(t ; t_{0}, x_{t_{0}}\right)$ satisfies $x_{t} \in \operatorname{BV}[0,1)$ for each $t \geqq t_{0}$.

PROOF: Let $V\left(t_{1}, t_{2}\right)$ be the variation of the solution $x$ on the interval $\left[t_{1}, t_{2}\right)$ Then, if $t_{3} \geqq t_{2}$,

$$
v\left(t_{1}, t_{3}\right)=v\left(t_{1}, t_{2}\right)+v\left(t_{2}, t_{3}\right)
$$

Thus, for $t_{0} \leqq t<t_{0}+1$,

$$
\begin{aligned}
\left\|x_{t}\right\| & =|x(t)|+v(t, t+1) \\
& =|x(t)|+v\left(t, t_{0}+1\right)+v\left(t_{0}+1, t+1\right) \\
& \leqq|x(t)|+\left|x\left(t_{0}\right)\right|+v\left(t_{0}, t_{0}+1\right)+\left|x\left(t_{0}+1\right)\right|+v\left(t_{0}+1, t_{0}+2\right) \\
& \leqq|x(t)|+\left\|x_{t_{0}}\right\|+\left\|x_{t_{0}+1}\right\|
\end{aligned}
$$

and $x_{t} \in \operatorname{BV}[0,1)$ for $t_{0} \leqq t<t_{0}+1$ since $x_{t_{0}}$ and $x_{t_{0}+1}$ are in $\operatorname{BV}[0,1)$. If $x_{t} \in \operatorname{BV}[0,1)$ for $t_{0} \leqq t<t_{0}+N$, then, using the fact that $x_{t+1}=f_{t}\left(x_{t}\right)$ is in $\operatorname{BV}[0,1)$ whenever $x_{t} \in \operatorname{BV}[0,1)$, we get that $\dot{x}_{t} \in \operatorname{BV}[0,1)$ for $t_{0} \leqq t<t_{0}+N+1$ and the theorem is proven by induction.

The solution $x(t)$, while written as a function of $t, t_{0}, x_{t_{0}}$, really depends only on $t, t_{0}+\left\langle t-t_{0}\right\rangle$ and $x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$. Since $t=$ $t_{0}+\left\langle t-t_{0}\right\rangle+n$ where $n=\left[t-t_{0}\right]$ is an integer, the value of the solution at $t_{0}+\left\langle t-t_{0}\right\rangle$ determines the solution at $t_{0}+\left\langle t-t_{0}\right\rangle+k$ for $k=0,1,2$, $\ldots, n$, and thus determines the solution at $t$.

THEOREM 3 If $f(t, x)$ is a continuous function of $x$ for all $x \in G(t)$ and all $t \geqq t_{0}$ and the assumptions of Theorem 1 hold, then the solution $x(t)=x\left(t ; t_{0}, x_{t_{0}}\right)$ is a continuous function of $x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$.

PROOF: If $t_{0} \leqq t<t_{0}+1$, then $x(t)=x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$ and this is obviously continuous in $x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$.

Now $x(t+1)=f(t, x(t))$, a continuous function of a continuous function, is also continuous in $x_{t_{0}}\left(\left\langle t-t_{o}\right\rangle\right)$. Thus, the theorem is proven by induction for all $t \geqq t_{0}$.

THEOREM 4 Under the conditions of Theorem 1 and $2 b$, the solution $x_{t}$ considered as a function of $x_{t_{0}}$ is a continuous function of the initial data.

The proof of this theorem is very difficult and it is given on pages 20-26 of Krasnosel'skii [8].

In looking for periodic solutions of a difference equation, it would
be nice to be able to use only continuous functions, i.e., functions which are continuous in $t$. However, the following example illustrates that this is not always possible.

For $x(t)$ a real number, consider the difference equation (12).

$$
\begin{equation*}
E x=-x^{1 / 3} \tag{12}
\end{equation*}
$$

One solution of (12) is the function $x(t)$ given by (12')

$$
x(t)=\left\{\begin{array}{rl}
1 & 0 \leqq t<1  \tag{i}\\
-1 & 1 \leqq t<2
\end{array}\right.
$$

and $x(t+2)=x(t)$. This solution is 2-periodic so there exists a non-trivial 2-periodic solution of (12). However, this solution is not continuous in $t$.

Assume that there is a non-trivial 2-periodic solution of (12) which is continuous in $t$. Then there is a $t_{1}$ such that $x\left(t_{1}\right)>0$ and $x\left(t_{1}+1\right)<$ 0 . Since $x$ is continuous there is a $t_{2}$ such that $t_{1}<t_{2}<t_{1}+1, x\left(t_{2}\right)=0$, and $0<x\left(t_{2}-\epsilon\right)<l$. for some $\epsilon>0$. Then $x\left(t_{2}-\epsilon+2\right)=-x\left(t_{2}-\epsilon+1\right)^{1 / 3}=$ $\left(x\left(t_{2}-\epsilon\right)\right)^{1 / 9}>x\left(t_{2}-\epsilon\right)$. This is a contradiction since $x$ is 2-periodic, i.e., $x\left(t_{2}-\epsilon+2\right)=x\left(t_{2}-\epsilon\right)$. Thus, any non-trivial 2-periodic solution of (12) must be discontinuous in $t$.

LINEAR DIFFERENCE EQUATIONS
A linear difference equation is an equation of the form (13) where $A(t)$ is an $n$ by $n$ matrix (real or complex).

$$
\begin{equation*}
E x=A(t) x \tag{13}
\end{equation*}
$$

This equation is called linear because solutions of (13) obey the superposition rule.

LEMMA 1 Superposition rule. If $x_{1}(t)$ and $x_{2}(t)$ are any two solutions of (13), then $x(t)=\alpha_{1}(t) x_{1}(t)+\alpha_{2}(t) x_{2}(t)$ is also a solution for any scalars $\alpha_{1}, \alpha_{2}$ which are l-periodic.

PROOF: $\quad A(t) x(t)=\alpha_{1}(t) A(t) x_{1}(t)+\alpha_{2}(t) A(t) x_{2}(t)$

$$
\begin{aligned}
& =\alpha_{1}(t) x_{1}(t+1)+\alpha_{2}(t) x_{2}(t+1) \\
& =\alpha_{1}(t+1) x_{1}(t+1)+\alpha_{2}(t+1) x_{2}(t+1)=x(t+1) .
\end{aligned}
$$

In linear difference equations l-periodic scalars play a role very analogous to that played by constants in linear differential equations. Thus, the set of functions $x_{1}(t), x_{2}(t), \ldots, x_{N}(t)$ are said to be linearly dependent at $t$ over the integers if there exist scalars $\alpha_{1}, \alpha_{2}, \ldots, \alpha_{N}$, not all zero, such that

$$
\sum_{k=1}^{N} \alpha_{k} x_{k}(t+m)=0
$$

identically for all integers $m_{\text {: }}$ Clearly, these scalars $\alpha_{1}, \alpha_{2}, \ldots, \alpha_{N}$ are functions of $t$ and can be chosen to be l-periodic in $t$. If the functions $x_{1}(t), x_{2}(t), \ldots, x_{N}(t)$ are linearly dependent at $t$ over the integers for all
$t$, then they are said to be linearly dependent over the integers or merely linearly dependent. Repeating, the functions $x_{1}(t), x_{2}(t), \ldots, x_{N}(t)$ are linearly dependent (over the integers) if there exists l-periodic scalar functions $\alpha_{1}(t), \alpha_{2}(t), \ldots, \alpha_{N}(t)$ such that

$$
\sum_{k=1}^{N}\left|\alpha_{k}(t)\right|^{2}>0 \text { for all } t
$$

and

$$
\sum_{k=1}^{N} \alpha_{k}(t) x_{k}(t)=0 \text { for all } t
$$

The functions $x_{1}(t), x_{2}(t), \ldots, x_{N}(t)$ are linearly dependent if the only 1 periodic scalars satisfying this equation are all identically zero. The fact that the scalars $\alpha_{1}(t), \alpha_{2}(t), \ldots, \alpha_{k}(t)$ must be l-periodic is important. If each $X_{k}(t)$ is in $E^{n}$ and $N \geqq n+1$, then there always are scalars which satisfy this equation but they may not be l-periodic.

The following Lemmas are not new with this paper but are included here for completeness.

LEMMA 2 Linear dependence of solutions -- Any $n+1$ solutions $x_{1}(t), x_{2}(t)$, $\ldots, x_{n+1}$ ( $t$ ) defined for $t \geqq t_{0}$ of (13) are linearly dependent on $\left[t_{0}, \infty\right)$.

PROOF: For each $t, t_{0} \leqq t<t_{0}+1$, there exist scalars $\alpha_{1}(t), \alpha_{2}(t), \ldots, \alpha_{n+1}(t)$ such that

$$
\sum_{k=1}^{n+1}\left|\alpha_{k}(t)\right|^{2}>0
$$

and

$$
\sum_{k=1}^{n+1} \alpha_{k}(t) x_{k}(t)=0
$$

identically in $t$ on this interval ( $t_{0} \leqq t<t_{0}+1$ ). This is because, for each $t$, this equation is the relationship depicting the linear dependence of $n+1$ n-vectors. Operating on this equation with $A(t)$, we get

$$
\begin{aligned}
0=A(t) \sum_{k=1}^{n+1} \alpha_{k}(t) x_{k}(t) & =\sum_{k=1}^{n+1} \alpha_{k}(t) A(t) x_{k}(t) \\
& =\sum_{k=1}^{n+1} a_{k}(t) x_{k}(t+1) .
\end{aligned}
$$

If we extend $\alpha_{k}(t)$ to the right by letting $\alpha_{k}(t+1)=\alpha_{k}(t)$, we see that we have a set of l-periodic scalars which satisfy the condition for linear dependence.

LEMMA 3 If there are $n$ linearly independent solutions of (13) $x_{1}(t), x_{2}(t)$, $\ldots, x_{n}(t)$ defined for $t \geqq t_{0}$, then every solution $x(t)$ for $t \geqq t_{0}$ can be uniquely expressed as the sum

$$
\begin{equation*}
x(t)=\sum_{k=1}^{n} \alpha_{k}(t) x_{k}(t) \tag{14}
\end{equation*}
$$

where each $\alpha_{k}(t)$ is a 1-periodic scalar.

PROOF: For each $t$ in the interval $t_{0} \leqq t<t_{0}+1$, there exists scalars $\alpha_{1}(t), \alpha_{2}(t), \ldots, \alpha_{n}(t)$ which satisfy (14). These scalars are uniquely defined. If we extend these scalars to the right by using the relation $\alpha_{k}(t+1)=\alpha_{k}(t)$, then the right-hand side of (14) represents a solution to (13) which agrees with $x(t)$ on $t_{0} \leqq t<t_{0}+1$. Since the solution with this initial function is unique, we must have equation (14) holding for all $t \geqq t_{0}$.

If we let $X(t)$ be the matrix formed by the linearly independent
solutions $x_{1}(t), x_{2}(t), \ldots, x_{n}(t)\left(x_{k}(t)\right.$ is the $k-t h$ column of $\left.x(t)\right)$ and $y(t)$ be the vector with elements $\alpha_{k}(t)\left(\alpha_{k}(t)\right.$ is the k-th element of $\left.y(t)\right)$, then $y(t+1)=y(t)$ and equation (14) can be rewritten as (14').

$$
\begin{equation*}
x(t)=X(t) y(t) \tag{14'}
\end{equation*}
$$

The matrix $\mathrm{X}(\mathrm{t})$ is called a fundamental matrix solution of (13) and satisfies the matrix equation $X(t+1)=A(t) X(t)$. If $X(t)=I$, the unit matrix, for $t_{0} \leq t<t_{0}+1$, then $X(t)$ is the principal matrix solution of (13) and is written $X\left(t, t_{0}\right)$. We note that, with the principal matrix solution, the fundLion $y(t)$ in (14') becomes the initial function $x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$ and hence the solution $x(t)$ can be written as (14").

$$
\begin{equation*}
x(t)=x\left(t, t_{0}\right) x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right) \tag{14"}
\end{equation*}
$$

LEMMA 4. Existence of $n$ linearly independent solutions -- If $\operatorname{det}(A(t)) \neq 0$ for all $t \geqq t_{0}$, then there exists $n$ linearly independent solutions of (13) for $t \geqq t_{0}$.

PROOF: Let $X(t)$ be a matrix solution of (13) such that $\operatorname{det}(X(t)) \neq 0$ for $t_{0} \leqq t<t_{0}+1$. Then, since $\operatorname{det}(X(t+1))=\operatorname{det}(A(t) x(t))=\operatorname{det}(A(t)) \operatorname{det}(x(t))$, we get by induction that $\operatorname{det}(x(t)) \neq 0$ for all $t \geqq t_{0}$. The columns of $x(t)$ thus represent $n$ linearly independent solutions of (13).
. In everything that follows, I assume that $\operatorname{det}(A(t)) \nLeftarrow 0$ for all
$t \geqq t_{0}$ and hence that Lemma 4 holds.
Since $X(t)$ is non-singular for all $t \geqq t_{0}, X^{-1}(t)$ must exist for
all $t$. From the identity $X^{-1}(t) X(t)=I$, we get that

$$
x^{-1}(t+1) x(t+1)=x^{-1}(t+1) A(t) x(t)=X^{-1}(t) A(t)
$$

and, since $X(t)$ is non-singular, $X^{-1}(t)$ satisfies the difference equation (15).

$$
\begin{equation*}
x^{-1}(t)=X^{-1}(t+1) A(t) \tag{15}
\end{equation*}
$$

For $z(t)$ a row vector (a 1 by $n$ matrix), the difference equation (15') is called the adjoint to (13).

$$
\begin{equation*}
z(t)=z(t+1) A(t) \text { or } z(t+1)=z(t) A^{-1}(t) \tag{15'}
\end{equation*}
$$

A fundamental matrix solution of ( $15^{\prime}$ ) is the inverse of some fundamental matrix solution of (13). In particular, the principal matrix solution of (15') is the inverse of the principal matrix solution of (13).

For each vector norm $|x|$, define the norm of a matrix $M$ by $|M|$.

$$
|M|=\inf [b:|M x| \leqq b|x| \text { for } a l l x]
$$

This matrix norm satisfies all the usual properties of a norm [6].
(i) $|M| \geqq 0$ for all $M$ and $|M|=0$ only when $M=0$.
(ii) $\left|M_{1}+M_{2}\right| \leqq\left|M_{1}\right|+\left|M_{2}\right|$
(iii) $|\alpha M|=|\alpha||M|$ for any scalar $\alpha$ (iv) $\left|M_{1} M_{2}\right| \leqq\left|M_{1}\right|\left|M_{2}\right|$
(v) $|M x| \leqq|M||x|$ for any vector $x$

LEMMA 5 Boundedness of the Principal Matrix Solution -- Let $a(t)$ be such that $|A(t)| \leqq a(t)$ for all $t \geqq t_{0}$ and $\xi\left(t, t_{0}\right)$ be defined by

$$
\begin{aligned}
& \xi\left(t, t_{0}\right)=1, \quad t_{0} \leqq t<t_{0}+1 \\
& \xi\left(t, t_{0}\right)=\prod_{j=1} a(t-j) \text { for } t \geqq t_{0}+1
\end{aligned}
$$

Then $\left|X\left(t, t_{0}\right)\right| \leqq \xi\left(t, t_{0}\right)$ for all $t \geqq t_{0}$.

PROOF: The proof will proceed by induction. Since $X\left(t, t_{0}\right)=I$ for $t_{0} \leq t<t_{0}+1$, we get that $\left|X\left(t, t_{0}\right)\right|=1=\xi\left(t, t_{0}\right)$ on this interval. Assume that the inequality $\left|x\left(t, t_{0}\right)\right| \leqq \xi\left(t, t_{0}\right)$ holds for $t_{0} \leqq t<s$, some $s \geqq t_{0}+1$. Then, since $X\left(t, t_{0}\right)$ satisfies the difference equation $X\left(t+1, t_{0}\right)=$ $A(t) X\left(t, t_{0}\right)$, we get

$$
\left|x\left(t+1, t_{0}\right)\right| \leqq|A(t)|\left|x\left(t, t_{0}\right)\right| \leqq a(t) \xi\left(t, t_{0}\right) \equiv \xi\left(t+1, t_{0}\right)
$$

for all $t$, $t_{o} \leq t<. s$. Thus, if the inequality is true for $t_{0} \leq t<s$, then it is true for $t_{0} \leq t<s+1$ and, by induction, is true for all $t \geqq t_{0}$.

COROLLARY If $a(t) \leqq a$ for all $t \geqq t_{0}$, then $\xi\left(t, t_{0}\right) \leqq a^{\left[t-t_{0}\right]} \leqq k(a) a^{t-t_{0}}$ where $k(a)=1$ if $a \geqq I$ and $k(a)=a^{-1}$ if $a<1$.

PROOF:

$$
\xi\left(t, t_{0}\right)=\prod_{j=1}^{\left[t-t_{0}\right]} a(t-j) \leqq \prod_{j=1}^{\left[t-t_{0}\right]} a=a^{\left[t-t_{0}\right]}
$$

Since $\left[t-t_{0}\right]=t-t_{0}-\left\langle t-t_{0}\right\rangle$,

$$
\xi\left(t, t_{0}\right) \leqq a^{\left[t-t_{0}\right]}=a^{t-t_{0}-\left\langle t-t_{0}\right\rangle}=\left(a^{-\left\langle t-t_{0}\right\rangle}\right) a^{t-t_{0}} .
$$

If $a \geq 1, a^{-\left\langle t-t_{0}\right\rangle} \leq I=k(a)$. If $a<1, a^{-\left\langle t-t_{0}\right\rangle} \leq a^{-1}=k(a)$ and the corollary is proven.

IEMMA 5' On Boundedness of the inverse of the principal matrix solution -Let $a_{1}(t)$ be such that $\left|A^{-1}(t)\right| \leqq a_{1}(t)$ for all $t \geqq t_{0}$ and $\xi_{1}\left(t, t_{0}\right)$ be defined by

$$
\begin{aligned}
& s_{1}\left(t, t_{0}\right)=1 \quad \text { for } \quad t_{0} \leqq t<t_{0}+1 \\
& s_{1}\left(t, t_{0}\right)=\prod_{j=1}^{\left[t-t_{0}\right]} a_{1}(t-j) \text { for } t \geqq t_{0}+1
\end{aligned}
$$

then $\left|X^{-1}\left(t, t_{0}\right)\right| \leqq \xi_{1}\left(t, t_{0}\right)$ for all $t \geqq t_{0}$.
COROLLARY If $a_{1}(t) \leqq a_{1}$ for all $t \geqq t_{0}$, then

$$
\xi_{1}\left(t, t_{0}\right) \leq a_{1}^{\left[t-t_{0}\right]} \leq k\left(a_{1}\right) a_{1}^{t-t_{0}}
$$

where $k\left(a_{1}\right)=1$ if $a_{1} \geqq 1$ and $k\left(a_{1}\right)=a_{1}^{-1}$ if $a_{1}<1$.
The proofs are almost identical to the proofs for Lemma 5 and its corollary.

Notice that, under the conditions of the corollaries, the principal matrix solution (and hence any solution) and its inverse are both bounded above by an exponential. If, in the corollary to Lemma 5 , $a<1$, then all solutions approach zero exponentially as $t \rightarrow \infty$.

LINEAR AUTONOMOUS DIFFERENCE EQUATIONS
A linear autonomous difference equation is of the form (131) where A is some constant $n$ by $n$ matrix (real or complex).

$$
\begin{equation*}
E x=A x \tag{13'}
\end{equation*}
$$

In order to completely characterize the solutions of (13) we will need the following notation.

For any matrix $B$ and some integer $k \geqq 0$, Null $(B)^{k}$, is the set of all vectors $x$ such that $B^{k} x=0$. Null $(B)^{k}$ forms a linear subspace of X. The distinct complex numbers $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{r}$ such that $\operatorname{Null}\left(A-\lambda_{i} I\right)^{l}$ contains non-zero elements are called the eigenvalues of $A$. With each eigenvalue $\lambda_{i}$ of $A$ there is associated an integer $m_{i} \geqq 1$ which is the smallest integer such that $\operatorname{Null}\left(A-\lambda_{i} I\right)^{m_{i}}=\operatorname{Null}\left(A-\lambda_{i} I\right)^{m_{i}+1}$. The set of all non-zero vectors $x$ in $\operatorname{Null}\left(A-\lambda_{i} I\right)^{m_{i}}$ are called the generalized eigenvectors of $A$ corresponding to the eigenvalue $\lambda_{i}$. It can be proven [4, Theorem 2 on page 113-114] that any n-vector $x$ can be written uniquely as the sum

$$
x=\sum_{k=1}^{r} x_{k}
$$

where each $x_{k} \in \operatorname{Null}\left(A-\lambda_{k} I\right)^{m_{k}}$.
For each integer $n$, the factorial polynomial of degree $n$, written $t^{(n)}$, is defined by

$$
\begin{aligned}
& t^{(n)}=(t(t+1) \ldots(t(n+1)))^{-1} \text { if } n<0 \\
& t^{(0)}=1 \\
& t^{(n)}=t(t-1) \ldots(t-(n-1)) \text { if } n>0 .
\end{aligned}
$$

The factorial polynomial satisfies the very useful equality [5]

$$
(t+1)^{(n)}=t^{(n)}+n t^{(n-1)} .
$$

Any complex number $\lambda \neq 0$ can be written as $\lambda=r e^{i \omega}$ for some real $r>0$ and real $\omega$. For real $t$, we can define the exponential function $\lambda^{t}$ to be $\lambda^{t}=r^{t} e^{i \omega t}$. However, we can also write $\lambda=r e^{i(\omega+2 \pi k) t}$, and we have a different definition for $\lambda^{t}$ for each integer $k$. However, noting that $e^{i 2 \pi k t}$ is l-periodic for every integer $k$, we can chose $\omega$ such that $-\pi<\omega \leqq \pi$ and absorb all the other definitions into an arbitrary l-periodic coefficient. With the knowledge that there are many possible definitions of $\lambda^{t}$ but that their differences can be absorbed by a l-periodic coefficient, we will use the definition that $\lambda^{t}=r^{t} e^{i \alpha t}$ where $-\pi<\omega \leqq \pi$ in all of the following.

LEMMA 6 Linear independence of the factorial polynomials -- The factorial polynomials $t^{(k)}, k=0,1,2, \ldots$, are linearly independent over $\left[t_{0}, t_{0}+s\right)$ where $s$ is any number, $s>1$.

PROOF: The proof will proceed by induction, $t^{(0)}=1$ by definition, and hence $t^{(0)}$ by itself is linearly independent. Assume now that $t^{(0)}, t^{(1)}$, $\ldots, t^{(n-1)}$ are linearly independent but that $t^{(0)}, t^{(1)}, \ldots, t^{(n-1)}, t^{(n)}$ are linearly dependent for some integer $n \geqq 1$. Then we can write II

$$
t^{(n)}=\sum_{k=0}^{n-1} \alpha_{k}(t) t^{(k)}
$$

where each

$$
\begin{aligned}
(t+1)^{(n)} & =t^{(n)}+n t^{(n-1)} \\
& =\sum_{k=0}^{n-1} \alpha_{k}(t+1)(t+1)^{(k)} \\
& =\sum_{k=0}^{n-1} \alpha_{k}(t)\left(t^{(k)}+k t^{(k-1)}\right. \\
& =\sum_{k=0}^{n-1} \alpha_{k}(t) t^{(k)}+\sum_{k=0}^{n-1} \alpha_{k}(t){ }_{k t}(k-1) \\
& =t^{(n)}+\sum_{k=0}^{n-1} \alpha_{k}(t) k t^{(k-1)} .
\end{aligned}
$$

Thus, we get that

$$
n t^{(n-1)}=\sum_{k=0}^{n-1} \alpha_{k}(t) k t(k-1)
$$

This is possible only if $\alpha_{k}(t)=0, k=0,1,2, \ldots, n-2$, and $\alpha_{n-1}(t)(n-1)=n$, i.e.,

$$
\alpha_{n-1}(t)=\frac{n}{n-1}
$$

Then

$$
t^{(n)}=\frac{n}{n-1} t^{(n-1)}
$$

From the definition of the factorial polynomials, we get

$$
t^{(n)}=(t-(n-1)) t^{(n-1)} .
$$

Hence $\alpha_{n-1}(t)$ as above is not possible and we have a contradiction. Thus, the set $t^{(0)}, t^{(1)}, \ldots, t^{(n)}$ is linearly independent whenever $t^{(0)}, t^{(1)}, \ldots$, $t^{(n-1)}$ is. The lemma is proven by induction.

LEMMA 61 Linear independence of the exponential functions -- If $\lambda_{1}, \lambda_{2}, \ldots$, $\lambda_{r}$ are $r$ distinct scalars (complex numbers), none of which are zero, then the exponential functions $\lambda_{1}^{t}, \lambda_{2}^{t}, \ldots, \lambda_{r}^{t}$ are linearly independent.

PROOF: Let $\alpha_{1}(t), \alpha_{2}(t), \ldots, \alpha_{r}(t)$ be 1-periodic functions such that

$$
\sum_{k=1}^{r} \alpha_{k}(t) \lambda_{k}^{t}=0 \text { for all } t
$$

Then, substituting $t+j$ for $t$, we get

$$
\sum_{k=1}^{r} \alpha_{k}(t)_{k}^{t+j}=0 \text { for } j=0,1,2, \ldots, r-1
$$

The determinant of the coefficients on the $\alpha_{k}$ 's is $W(t)$.

$$
\begin{aligned}
W(t) & =\operatorname{det}\left(\begin{array}{cccc}
\lambda_{1}^{t} & \lambda_{2}^{t} & \ldots & \lambda_{r}^{t} \\
\lambda_{1}^{t+1} & \lambda_{2}^{t+1} & \ldots & \lambda_{r}^{t+1} \\
\vdots & & & \\
\lambda_{1}^{t+r-1} & \lambda_{2}^{t+r-1} & \ldots & \lambda^{t+r-1}
\end{array}\right) \\
& =\lambda_{1}^{t} \lambda_{2}^{t} \ldots \lambda_{r}^{t} \operatorname{det}\left(\begin{array}{llll}
1 & 1 & \ldots & 1 \\
\lambda_{1} & \lambda_{2} & \ldots & \lambda_{r} \\
\vdots & & & \\
\lambda_{1}^{r-1} & \lambda_{2}^{r-1} & \ldots & \lambda_{r}^{r-1}
\end{array}\right)
\end{aligned}
$$

Let there be constants $\beta_{o}, \beta_{1}, \ldots, \beta_{r-1}$ such that $p(\lambda)=\sum_{k=0}^{r-1} \beta_{k} \lambda^{k}$ satisfies $p\left(\lambda_{1}\right)=p\left(\lambda_{2}\right)=\ldots=p\left(\lambda_{r}\right)=0$. This $p(\lambda)$ is a polynomial in $\lambda$ of degree less than $r$ which has at least $r$ roots. The only possible polynomial with this property is the zero polynomial, i.e., $\beta_{0}=\beta_{1}=\ldots=\beta_{r-1}=0$. This in turn says that the rows of the last determinant above are linearly independent and hencethat determinant is non-zero. Since none of the $\lambda_{k}{ }^{\prime} s$ are zero, each factor $\lambda_{k}^{t}$ in $W(t)$ is non-zero, and so $W(t) \neq 0$ for all $t$. Since $W(t) \neq 0$ for all $t$, the only possible solution of the equations for the $\alpha_{k}^{\prime} s$ is $\alpha_{1}(t)=\alpha_{2}(t)=\ldots=\alpha_{r}(t)=0$ for all $t$, and the exponential functions $\lambda_{1}^{t}, \lambda_{2}^{t}, \ldots, \alpha_{r}^{t}$ are linearly independent.

Suppose now that $A$ has only one eigenvalue $\lambda \neq 0$ and that $m_{1}=m$ (or that we restrict our attention to the space $\operatorname{Null}(A-\lambda I)^{m}$ ). If $x_{o}$ is a non-trival solution of the equation $(A-\lambda I) x_{0}=0$, then one solution is of the form $x(t)=x_{0} \lambda^{t}$. However, this is only one of the $m$ linearly independent solutions which exist. Even if there were several linearly independent $x_{0}{ }^{\prime} s$ which satisfy the equation $(A-\lambda I) x_{0}=0$, in general there will not be $m$ of them. Since the factorial polynomials $t^{(0)}, t^{(1)}, t^{(2)}, \ldots$, form a linearly independent set on $\left[t_{0}, \infty\right)$, try the solution $x(t)$

$$
x(t)=\left(\sum_{k=0}^{\infty} t^{\left.(k)_{x_{k}}\right) \lambda^{t}}\right.
$$

where each $x_{k}$ is a constant vector. Then

$$
\begin{align*}
x(t+1)-A x(t) & =\left(\sum_{k=0}^{\infty}(t+1)^{(k)} x_{k}\right) \lambda^{t+1}-\left(\sum_{k=0} t^{(k)} A x_{k}\right) \lambda^{t} \\
& \left.=\sum_{k=0}^{\infty}\left(\lambda^{(t}(k)_{+k t}^{(k-1)}\right) x_{k}-t^{(k)} A x_{k}\right) \lambda^{t}
\end{align*}
$$

$$
=\sum_{k=0}^{\infty} t^{(k)}\left(\lambda x_{k}+(k+1) \lambda x_{k+1}-A x_{k}\right) \lambda^{t} .
$$

From the requirement that $x(t+1)-A x(t)=0$ for all $t$, we get that $x_{0}$ is arbitrary in $\operatorname{Null}(A-\lambda I)^{m}$ and

$$
x_{k+1}=\frac{1}{(k+1) \lambda} \cdot(A-\lambda I) x_{k}
$$

for $k=0,1,2, \ldots$. By repeatedly applying this identity, we get that

$$
x_{k}=\frac{1}{k!\lambda^{k}}(A-\lambda I)^{k} x_{0}
$$

Since $x_{0} \in \operatorname{Null}(A-\lambda I)^{m}$, we see that $x_{k}=0$ for all $k \geqq m$. Thus, the sum for $x(t)$ has only a finite number of non-zero terms and there is no worry about convergence. The solution is

$$
\begin{equation*}
x(t)=x\left(t ; \lambda, x_{0}\right)=\sum_{k=0}^{m-1} \frac{t^{(k)}}{k!}(A-\lambda I)^{k} x_{0} \tag{16}
\end{equation*}
$$

for any $x_{0} \in \operatorname{Null}(A-\lambda I)^{m}$. Note that this solution is a polynomial in $t$ times the exponential function $\lambda^{t}$ and is a linear function of $x_{0}$.

In the more general case where $A$ has the $r$ distinct eigenvalues $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{r}$, we can find $n$ linearly independent solutions in the following manner. Pick $n$ linearly independent vectors $x_{1}, x_{2}, \ldots, x_{n}$ such that each $x_{k}$ is in one of the spaces Null $\left(A-\lambda_{j} I\right)^{m_{j}}, j=1,2, \ldots, r$. This can always be done [4]. If $x_{k} \in \operatorname{Null}\left(A-\lambda_{j} I\right)^{m_{j}}$, form the solution $x_{k}(t)$ in (17) using the solution in (16).

$$
\begin{equation*}
x_{k}(t)=x\left(t ; \lambda_{j}, x_{k}\right) \text { for } k=1,2, \ldots, n \tag{17}
\end{equation*}
$$

IEMMA 7. The $n$ solutions (17) are linearly independent.
Let $x(t)$ be the solution $O f\left(13^{\prime}\right)$ with the initial condition $x_{0}(t)=x(t)$ for $0 \leqq t<1$. Write

$$
x_{0}(t)=\sum_{k=1}^{r} x_{k}(t) \quad 0 \leqq t<1
$$

where $x_{k}(t) \in \operatorname{Null}\left(A-\lambda_{k} I\right)^{m_{k}}$ for each $k$. Then the solution $x(t)$ is of the form (18) using the solutions in (16).

$$
\begin{equation*}
x(t)=\sum_{k=0}^{r} x\left([t] ; \lambda_{k}, x_{k}(\langle t\rangle)\right) \tag{18}
\end{equation*}
$$

Thus, any solution can be expressed as a linear combination of solutions of the form of (16), i.e., as a linear combination of exponential functions times factorial polynomials.

The behavior of each solution (16) depends on the eigenvalue $\lambda_{i}$ under consideration. If $\left|\lambda_{i}\right|>1$, then the solution (16) grows exponentially as $t \rightarrow \infty$. Equation (16) can be used as the definition of a solution for all $t$, and we see that the solution approaches zero exponentially as $t \rightarrow-\infty$. Let $X^{l}$ be the subspace of $X$ corresponding to all the eigenvalues $\lambda_{i}$ which satisfy $\left|\lambda_{1}\right|>1$. Then any solution which starts in $X^{l}$ (i.e., $x_{t_{0}}(\theta) \in X^{l}$ for $0 \leqq \theta<1$ ) remains in $X^{1}$ and grows exponentially as $t \rightarrow \infty$. The only solution in $X^{1}$ which is bounded for all $t$ in ( $-\infty, \infty$ ) is the trivial solution $x(t)=0$.

If $\left|\lambda_{i}\right|<1$, then the solution (16) approaches zero exponentially
as $t \rightarrow \infty$ and grows exponentially as $t \rightarrow-\infty$. Let $x^{2}$ be the subspace of $X$ corresponding to all the eigenvalues $\lambda_{i}$ which satisfy $\left|\lambda_{i}\right|<1$. Then any solution which starts in $X^{2}$ (i.e., $X_{t_{0}}(\theta) \in X^{2}$ for $0 \leqq \theta<1$ ) remains in $X^{2}$ and approaches zero exponentially as $t \rightarrow \infty$. The only solution in $x^{2}$ which is bounded for all $t$ in $(-\infty, \infty)$ is the trivial solution $x(t)=0$.

If $\left|\lambda_{i}\right|=1$, then the behavior of the solution depends on the values of $\left(A-\lambda_{i} I\right)^{k_{0}} x_{0}$ in equation (16). If $x_{0}$ is such that $\left(A-\lambda_{i} I\right) x_{0}=0$, then the solution is bounded for all $t$ and is, in fact, periodic. Suppose $\lambda_{i}=e^{i 2 \pi \omega}$ for some $\omega, 0<\omega \leqq 1$. Then the solution (16) is given by

$$
x(t)=e^{i 2 \pi \omega t} x_{0}
$$

which is ( $1 / \omega$ )-periodic. But, since $e^{i 2 \pi k t}, k$ an integer, is a l-periodic function, $e^{i 2 \pi k t} x(t)=e^{i 2 \pi(\omega+k) t_{0}}$ is also a solution of (13), and it is periodic with period $1 /(\omega+k)$. Since each T-periodic function is also mTperiodic ( $m$ integer, $m \neq 0$ ), the set $S$ of periods for solutions of (13) is $S_{\omega}=\{m /(\omega+k) ; m \neq 0$, all $k\}$. This set $S_{\omega}$ is dense in the real line, in contrast to the case in differential equations where $S_{\omega}$ (the set of periods of periodic solutions of $\dot{x}=A x$ ) consists of a countable set of distinct points.

If $\left|\lambda_{i}\right|=1$ and $x_{0}$ is such that $\left(A-\lambda_{i} I\right)^{m} x_{0}=0$ but $\left(A-\lambda_{i} I\right) x_{0} \neq$ 0 , then the solution (16) grows like a power of $t$ as $t \rightarrow \infty$. The only bounded solution is the trivial solution $x(t)=0$.

## AFFINE DIFFERENCE EQUATIONS

An affine difference equation is an equation of the form of (19)

$$
\begin{equation*}
E x=A(t) x+f(t) \tag{19}
\end{equation*}
$$

where $A(t)$ is an $n$ by $n$ matrix and $f(t)$ is an $n$-vector. Associated with each affine difference equation (19) is a corresponding linear difference aquaion (13).

$$
\begin{equation*}
E x=A(t) x \tag{13}
\end{equation*}
$$

Let $X\left(t, t_{0}\right)$ be the principal matrix solution of (13) and make the change of variables (20).

$$
\begin{equation*}
x(t)=X\left(t, t_{0}\right) y(t) \tag{20}
\end{equation*}
$$

Then

$$
\begin{aligned}
x(t+1) & =X\left(t+1, t_{0}\right) y(t+1) \\
& =A(t) x\left(t, t_{0}\right) y(t)+f(t) \\
& =X\left(t+1, t_{0}\right) y(t)+f(t)
\end{aligned}
$$

and so we get that $y(t)$ must satisfy the difference equation

$$
y(t+1)=y(t)+x^{-1}\left(t+1, t_{0}\right) f(t)
$$

For $t_{0} \leqq t<t_{0}+1$, we have $x\left(t, t_{0}\right)=I$ and so $y(t)=x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)$.

The solution for $y(t)$ is, for $t \geqq t_{0}+1$,

$$
y(t)=x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)+\sum_{j=1}^{\left[t-t_{0}\right]} x^{-1}\left(t+1-j, t_{0}\right) f(t-j)
$$

and the solution $x(t)$ is given by (21).

$$
x(t)=x\left(t, t_{0}\right) x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)+\sum_{j=1}^{\left[t-t_{0}\right]} x\left(t, t_{0}\right) x^{-1}\left(t+1-j, t_{0}\right) f(t-j)
$$

Equation (21) is called the variation of constants formula for (19).
Another form for the variation of constants formula can be obtained by a repeated application of equation (19).

$$
\begin{aligned}
x(t)= & A(t-1) x(t-1)+f(t-1) \\
= & A(t-1) A(t-2) x(t-2)+A(t-1) f(t-2)+f(t-1) \\
= & \ldots \\
= & A(t-1) \ldots A(t-n) x(t-n)+f(t-1)+A(t-1) f(t-2) \\
& +A(t-1) A(t-2) f(t-3)+\ldots+A(t-1) \ldots A(t-n+1) f(t-n) .
\end{aligned}
$$

If we let $n=\left[t-t_{0}\right]$, then $t-n=t_{0}+\left\langle t-t_{0}\right\rangle$ and

$$
x(t)=x\left(t_{,} t_{0}\right) x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)+\sum_{j=1}^{\left[t-t_{0}\right]} x\left(t, t_{0}+j\right) f\left(t_{0}+\left\langle t-t_{0}\right\rangle+j-1\right)
$$

This equation is more informative of the nature of the coefficients on $x_{t_{0}}$ and $f$ in the variation of constants formula.

Let $x^{0}(t)$ be the solution of (19) with zero initial function. Then $x_{t_{0}}\left(\left\langle t-t_{0}\right\rangle\right)=0$ for all $t$ and, for $t \geqq t_{0}+1$,

$$
x^{o}(t)=x^{o}(t ; f)=\sum_{j=1}^{\left[t-t_{0}\right]} x\left(t, t_{0}\right) x^{-1}\left(t+1-j, t_{0}\right) f(t-j)
$$

We see that $x^{0}$ is a linear function of the function $f$. Let $x^{1}(x)=$ $x^{1}\left(t ; x_{t_{0}}\right)$ be the solution of the homogeneous equation (13) with initial value $x_{t_{0}}$. Then $x^{1}$ is a linear function of the function $x_{t_{0}}$ and the solution $x(t)$ in (21) can be rewritten as in (21).

$$
\begin{equation*}
x(t)=x^{1}\left(t ; x_{t_{0}}\right)+x^{0}(t ; f) \tag{21'}
\end{equation*}
$$

Suppose now that $f \in C\left[t_{0}, t\right]$ and $\|f\|_{t}=\max \cdot\left(|f(s)|, t_{0} \leq s \leq t\right\}$.
Then

$$
\begin{aligned}
\left|x^{0}(t ; f)\right| & =1 \sum_{j=1}^{\left[t-t_{0}\right]} x\left(t, t_{0}+j\right) f\left(t_{0}+t-t_{0}+j-1\right) \mid \\
& \leqq \sum_{j=1}^{\left[t-t_{0}\right]}\left|x\left(t, t_{0}+j\right)\right|\left|f\left(t_{0}+t-t_{0}+j-1\right)\right|
\end{aligned}
$$

$$
\leqq \sum_{j=1}^{\left[t-t_{0}\right]} \xi\left(t, t_{0}+j\right)\|f\|_{t}=m\left(t, t_{0}\right)\|f\|_{t} .
$$

Hence, $x^{0}$ is a continuous linear function of $f$. Thus there exists a matrix valued function $\beta(t, s)$ such that each element, considered as a function of $s$ (holding $t$ fixed), is of bounded variation on $\left[t_{0}, t\right]$ and such that

$$
x^{0}(t ; f)=\int_{t_{0}}^{t} d_{s} \beta(t, s) f(s)
$$

See [7], for example.

BOUNDED SOLUTIONS -- NONCRITICAL CASE
Suppose that, in equation (19), where A is a constant matrix,

$$
\begin{equation*}
E x=A x+f(t) \tag{19}
\end{equation*}
$$

the function $f(t)$ is defined for all $t$ in $(-\infty, \infty)$ and is bounded for all $t$ in $(-\infty, \infty)$. Under what conditions is there a bounded solution $x(t)$ of (19) ?

Let $|x|$ be any vector norm for $x \in E^{n}$. A function $f:(-\infty, \infty) \rightarrow E^{n}$ is said to be bounded it there is some number $F \geqq 0$ such that $|f(t)| \leqq F$ for all $t$ in $(-\infty, \infty)$. Let $B$ be the set of all bounded functions and define addition and scalar multiplication on elements of $B$ in the usual sense. For each $f \in B$, let

$$
\|f\|=\sup \{|f(t)|: \text { all }, t\}
$$

With this norm, B becomes a Banach space.

LEMMA 8 Let $\left(u_{n}\right)$ be a sequence of bounded functions, $u_{n} \in B$, which converge to $u_{0} \in B$. If each $u_{n}$ is (uniformly) continuous in $t$, then $u_{0}$ is (uniformly) continuous in $t$. If each $u_{n}$ is almost periodic uniformly in $n$, then $u_{0}$ is almost periodic.

This lemma is a fairly standard one since convergence in $B$ means pointwise convergence uniformly in $t$. For the continuous and uniformly continuous proof, see any advanced calculus book, for example [16]. For the almost periodic proof, see, for example, [17].

The first case handled is the non-critical case where none of the eigenvalues of $A$ lie on the unit circle. The critical case is where one or more of the eigenvalues of $A$ lie on the unit circle. For the non-critical. case, we assume that $A$ has the following form

$$
A=\left(\begin{array}{ll}
A_{1} & 0 \\
0 & A_{2}
\end{array}\right)
$$

where $A_{1}$ and $A_{2}$ are square matrices such that all the eigenvalues of $A_{1}$ lie inside the unit circle and all the eigenvalues of $A_{2}$ lie outside the unit circle. If the matrix $A$ is not in this form, it can be arranged to this form by a transformation of coordinates. The assumption is that such a change of coordinates has already been performed. Let $x=\operatorname{col}\left(x_{1}, x_{2}\right)$ and $f(t)=\cdots$ $\operatorname{col}\left(f_{1}(t), f_{2}(t)\right)$. Then the difference equation (19) becomes the difference equation (22).

$$
\begin{align*}
& E x_{1}=A_{1} x_{1}+f_{1}(t)  \tag{22}\\
& E x_{2}=A_{2} x_{2}+f_{2}(t)
\end{align*}
$$

Let $x_{1}(t)$ be a solution of (22). Then

$$
\begin{aligned}
x_{1}(t) & =A_{1} x_{1}(t-1)+f(t-1)=A_{1}^{2} x_{1}(t-2)+A_{1} f_{1}(t-2)+f_{1}(t-1) \\
& =\ldots=A_{1}^{n} x_{1}(t-n)+\sum_{j=1}^{n} A_{1}^{j-1} f_{1}(t-j) .
\end{aligned}
$$

Since all the eigenvalues of $A_{1}$ are inside the unit circle, the matrices $A_{1}^{j-1}$ approach zero exponentially as $j \rightarrow \infty$. Since $f$ is bounded, this means that the sum converges as $n \rightarrow \infty$. If $x_{1}$ is bounded, then $A_{1}^{n} x_{1}(t-n) \rightarrow 0$ as
$n \rightarrow \infty$ and

$$
\begin{equation*}
x_{1}(t)=\sum_{j=1}^{\infty} A_{1}^{j-1} f_{1}(t-j) \tag{23}
\end{equation*}
$$

There are numbers $a_{1}>0$ and $b_{1}, 0 \leqq b_{1}<1$, such that

$$
\left|A_{1}\right| \leqq a_{1} b_{1}^{n} \quad \text { for all integers } n \geqq 0
$$

and so

$$
\begin{aligned}
\left|x_{1}(t)\right| & \leqq \sum_{j=1}^{\infty}\left|A_{1}^{j-1} f_{1}(t-j)\right| \leqq \sum_{j=1}^{\infty}\left|A_{1}^{j-1}\right|\left|f_{1}(t-j)\right| \\
& \leqq \sum_{j=1}^{\infty} a_{1} b_{1}^{j-1}\left\|f_{1}\right\|=\frac{a_{1}}{1-b_{1}}\left\|f_{1}\right\| .
\end{aligned}
$$

Hence the series for $X_{1}(t)$ converges absolutely. This inequality along with equation (23) shows that $x_{1}(t)$ does exist and that it is a continuous linear function of $f$.

If there were two bounded solutions to (22), say $x_{1}^{1}$ and $x_{1}^{2}$, then their difference $x_{1}^{1}-x_{1}^{2}$ would be bounded and satisfy the homogeneous difference equation $E x_{1}=A_{1} x_{1}$. The only bounded solution of this difference equation is the trivial solution $x_{1}=0$, hence $x_{1}^{1}=x_{1}^{2}$ and the bounded solution is unique. For future purposes, we write this solution as (23.')

$$
\begin{equation*}
x_{1}=L_{1}\left(f_{1}\right) \tag{23'}
\end{equation*}
$$

where $L_{1}: B \rightarrow B$ is the continuous linear operator defined by the right-hand side of equation (23).

In a similar manner, we get the unique bounded solution $x_{2}(t)$ as in equation (24).

$$
x_{2}(t)=L_{2}\left(f_{2}\right)(t)=-\sum_{j=0}^{\infty} A_{2}^{-j-1} f_{2}(t+j)
$$

Since all the eigenvalues of $A_{2}$ are outside the unit circle, there are numbers $a_{2}>0$ and $b_{2}, 0 \leqq b_{2}<1$, such that

$$
\left|A_{2}^{-n}\right| \leqq a_{2} b_{2}^{+n} \quad \text { for all integers } n \geqq 0
$$

and so

$$
\left|x_{2}(t)\right| \leqq+\sum_{j=0}^{\infty} a_{2} b_{2}^{j+1}\left\|f_{2}\right\|=\frac{a_{2} b_{2}}{1-b_{2}}\left\|f_{2}\right\|
$$

The series converges absolutely and $L_{2}: B \rightarrow B$ is a continuous linear operator.
We can write the unique bounded solution of (19) as in equation (25)

$$
\begin{equation*}
x(t)=L(f)(t) \tag{25}
\end{equation*}
$$

where. L: $B \rightarrow B$ is defined by

$$
L(f)=\binom{L_{1}\left(f_{1}\right)}{L_{2}\left(f_{2}\right)}
$$

and is a continuous linear operator.

THEOREM 5 If the matrix A does not have any eigenvalues on the unit circle, then there is a continuous linear operator $L: B \rightarrow B$ such that $x=L f$ is the unique bounded solution of (19).

$$
\begin{equation*}
E x=A x+f(t) \tag{19}
\end{equation*}
$$

This operator $I$ has the following properties.

1) If $f$ is uniformly continuous in $t$; then so is Lf.
2) If $f$ is T-periodic in $t$, then so is Lf.
3) If $f$ is almost periodic in $t$, then so is Lf.

PROOF: The existence and continuity of $L$ and the uniqueness of the solution $x=$ Lf are proven in the preceeding discussion. All that remains is to prove that $L$ possesses these three properties.

The proof of each depends on the following fact. Let $h$ be some real number and $f \in B$. Define $f_{1} \in B$ by $f_{1}(t)=f(t+h)$. Let $x=L f$ and defined $x_{1} \in B$ by $x_{1}(t)=x(t+h)$. Then $x_{1}=L f_{1}$ since, in equations (23.) and (24), $t$ can be replaced by $t+h$ on both sides of the equations without disturbing the equalities.

Since $L: B \rightarrow B$ is a continuous linear operator, there is a number $k, 0 \leqq k<\infty$, such that $\|L f\| \leqq k\|f\|$ for all $f \in B$.
1). If $f$ is uniformily continuous in $t$, then, for each $\in>0$, there is a $\delta=\delta(\epsilon)>0$ such that $|f(t+h)-f(t)| \leqq \epsilon$ for all $h,|h| \leqq \delta$. For any such $h$, let $f_{1}(t)=f(t+h)$. Then $f_{1} \in B,\left\|f_{1}\right\|=\|f\|$, and $\left\|f_{1}-f\right\|=$ $\sup \left\{\left|f_{1}(t)-f(t)\right|:\right.$ ail $\left.t\right\} \leqq \epsilon$. Let $x=\operatorname{Lf}$ and $x_{1}(t)=x(t+h)$. Then $x_{1}=L f_{1}$ and

$$
\left\|x_{1}-x\right\|=\left\|L f_{1}-L f\right\| \leqq k\left\|f_{1}-f\right\| \leqq k \in .
$$

Thus we have, for each $\epsilon>0,|x(t+h)-x(t)| \leqq\left\|x_{1}-x\right\| \leqq k \epsilon$ for ali $h,|h| \leqq \delta(\epsilon)$. Hence $x=L f$ is uniformly continuous in $t$.
2) If $f$ is T-periodic, then $f(t+T)=f(t)$. Let $f_{1}(t)=f(t+T)$.

Then $f_{1} \in B,\left\|f_{1}\right\|=\|f\|$, and $f_{1}-f=0$. Let $x=L f$ and $x_{1}(t)=x(t+T)$. Then $x_{1}=L \dot{f}_{1}$ and $x_{1}-x=L f_{1}-L f=L\left(f_{1}-f\right)=0$, implying that $x=L f$ is T-periodic.
3) If $f$ is almost periodic, then, for any $\epsilon>0$, there is an $l(\epsilon)>0$ such that, in any interval of length $l(\epsilon)$, there is a $T$ such that

$$
|f(t+T)-f(t)| \leqq \epsilon
$$

for all $t$. For this $\epsilon, l(\epsilon)$, and $T$, let $f_{1}$ be the function defined by $f_{1}(t)=f(t+T)$. Then $\left\|f_{1}-f\right\| \leqq \epsilon$. Let $x=L f$ and $x_{1}(t)=x(t+T)=L f_{1}(t)$. Then

$$
\left\|x_{1}-x\right\| \leqq k\left\|f_{1}-f\right\| \leqq k \epsilon
$$

and we see that $x=$ If is almost periodic.
As an example of the application of this theorem, consider the following interesting problem. Given the following differential equation where $f \in B$, find a solution $x \in B$

$$
\begin{aligned}
& \dot{x}_{1}=x_{2} \\
& \dot{x}_{2}=-x_{1}+g(t)
\end{aligned}
$$

numerically by using the Crude Euler method. For this differential equation and a given spacing $h>0$ in $t$, the Crude Euler method is given by the following difference equation.

$$
\begin{align*}
& E x_{1}=x_{1}+h \dot{x}_{1}=x_{1}+h x_{2} \\
& E x_{2}=x_{2}+h \dot{x}_{2}=-h x_{1}+x_{2}+h g(j s) \tag{26}
\end{align*}
$$

Here, the independent variable is changed to $s$ where $t=h s$ in order to bring this difference equation into the form of equation (19). In fact, equation (19) becomes equation (26) if we let

$$
A=\left(\begin{array}{cc}
1 & h \\
-h & 1
\end{array}\right) \quad f(s)=\binom{0}{h g(h s)}
$$

The eigenvalues of $A$ are $\lambda_{1}=1+i h$ and $\lambda_{2}=1$-ih. Neither $\lambda_{1}$ nor $\lambda_{2}$ are on the unit circle, so Theorem 5 holds. Indeed, since $\left|\lambda_{1}\right|>1$ and $\left|\lambda_{2}\right|>1$, equation (24) gives the bounded solution of (26) for any bounded function g. Since the differential equation has no bounded solution for $g(t)=\sin t$, we see that the behavior of the difference equation (26) in this case is considerably different from the differential equation.

Let $x$ be the bounded solution of (26). Then the bound on $x$ is given by

$$
|x(t)| \leqq \frac{\left(1+h^{2}\right)^{-\frac{1}{2}}}{1-\left(1+h^{2}\right)^{-\frac{1}{2}}}\|f\| \leqq \frac{1+\left(1+h^{2}\right)^{\frac{1}{2}}}{h^{2}}\|f\|=\frac{1+\left(1+h^{2}\right)^{\frac{1}{2}}}{h}\|\varepsilon\|
$$

and we see that, while $h>0$, if $h$ is small then the bound on $|x(t)|$ may be large indeed; in the limit as $h \rightarrow 0$, the bound becomes infinite (if $\|g\|>0$ ). For $g(t)=\sin \omega t$, we get the solution

$$
\begin{aligned}
& x_{1}(t)=\alpha_{1} \cos \omega t+\alpha_{2} \sin \omega t \\
& x_{2}(t)=\alpha_{3} \cos \omega t+\alpha_{4} \sin \omega t
\end{aligned}
$$

where

$$
\begin{aligned}
d & =\left((1-\cos \omega h)^{2}-(\sin \omega h)^{2}+h^{2}\right)^{2}+4(\sin \omega h)^{2}(1-\cos \omega h)^{2} \\
\alpha_{1} & =2 h^{2} \sin \omega h(1-\cos \omega h) / d \\
\alpha_{2} & =h^{2}\left((1-\cos \omega h)^{2}-(\sin \omega h)^{2}+h^{2}\right) / d \\
\alpha_{3} & =-h \sin \omega h\left((1-\cos \omega h)^{2}+(\sin \omega h)^{2}-h^{2}\right) / d \\
\alpha_{4} & =-h(1-\cos \alpha h)\left((1-\cos \omega h)^{2}+(\sin \omega h)^{2}+h^{2}\right) / d .
\end{aligned}
$$

We notice that $\left|x_{1}(t)\right|^{2} \leq \alpha_{1}^{2}+\alpha_{2}^{2}=h^{4} / \alpha$, that $x_{2}(t) \quad$ is approximately orthogonal to $x_{1}(t)$, and so the term $h^{4} / d$ will give us a good approximation to the upper bound of $|x(t)|^{2}$. Expanding $\sin \omega h$ and $\cos \omega h$ in their Taylor series, we get

$$
\begin{aligned}
1-\cos \omega h & =\frac{1}{2} \omega^{2} h^{2}+\ldots \\
\sin \omega h & =\omega h-\frac{1}{6} \omega^{3} h^{3}+\ldots
\end{aligned}
$$

and, after some simplification,

$$
d=\left(1-\omega^{2}\right)^{2} h^{4}+\frac{1}{6} \omega^{4}\left(7-\omega^{2}\right) h^{6}+\ldots
$$

The bound we desire is

$$
\frac{h^{4}}{d}=\frac{1}{\left(1-\omega^{2}\right)^{2}+\frac{1}{6} \omega^{4}\left(7-\omega^{2}\right) h^{2}+\cdots}
$$

For $\omega \neq 1$, this is close to the desired amplitude. The bounded solution of the differential equation satisfies

$$
|x(t)|^{2} \leqq 1 /\left(1-\omega^{2}\right)^{2}
$$

and we see that, if we ignore terms of order $h^{2}$, this bound is $h^{4} / a$. However, when $\omega=1$, there is no bounded solution of the differential equation and the bound on the bounded solution of the difference analog is

$$
|x(t)|^{2} \leqq \frac{h^{4}}{d}=\frac{1}{h^{2}+\ldots}
$$

This shows that the bound for $\|x\|$ can be reached. While the behavior of the bounded solution of the difference analog (26) is considerably different than the behavior of the bounded solution of the differential equation near $\omega=1$, this fact can be easily detected by computing the solution twice, once with $h=h_{0}$ and once with $h=\frac{1}{2} h_{0}$. If the two computed solutions vary greatly, then $h_{0}$ is in a region where the behavior of the difference analog differs greatly from the behavior of the differential equation. The solution is to use either a much smaller $h$ or a different numerical method.

THEOREM 6: If the matrix $A$ is such that none of its eigenvalues lie on the unit circle, then there is some $\epsilon_{I}>0$ such that there is a bounded solution $\dot{x} *(t)$ of (27) for each $\epsilon,|\epsilon|<\epsilon \epsilon_{1}$,

$$
\begin{equation*}
E x=A x+\epsilon f(t, x, \epsilon) \tag{27}
\end{equation*}
$$

where $f(t, x, \epsilon)$ is bounded for each fixed bounded $x$, each $\epsilon$, $|\epsilon|<\epsilon_{1}$, and all $t$, and $f$ is uniformly Lipschitz continuous with Lipschitz constant $F$ for all $t$ in ( $-\infty, \infty$ ) and all $x$
with $|x| \leqq M$, some $M>0$. This bounded solution $x^{*}(t)$ has the following properties.

1) If $f(t, x(t), \epsilon)$ is continuous in $t$ for fixed $\epsilon$, continuous $x$, then so is $x^{*}(t)$.
2) If $f(t, x(t), \epsilon)$ is $T$-periodic in $t$ for fixed $\epsilon$, T-periodic $x$, then so is $x^{*}(t)$.
3) If $f(t, x(t), \epsilon)$ is almost periodic for any almost periodic $x(t)$, then so is $x^{*}(t)$.

This result differs from that of Halanay [10] in that, in this paper, Halanay only considers the solution to be defined on the integers. Thus, when Halanay deals with T-periodicity, $T$ must be an integer. His work with almost periodic functions is for functions almost periodic on the integers. Thus, this result is a generalization of Halanay's result in [10]. If all the eigenvalues of $A$ are inside the unit circle, then this theorem is a specialization of the result of Halanay in [11]. At the same time, this result is more general than that of Halanay in [11] since this result allows $A$ to have eigenvalues outside the unit circle.

PROOF: Let $x$ be any bounded function with $\|x\| \leqq M$. Let $N$ be the positive number such that $|f(t, x, \epsilon)| \leqq N$ for all $t$, all $x$ with $|x| \leqq M$, and $|\epsilon|<\epsilon_{0}$, some $\epsilon_{0}>0$. For this $x$, define the operator $L^{*}$ by

$$
L^{*}(x)=\epsilon L(f(\cdot, x(\cdot), \epsilon))
$$

where $L$ is the linear operator of Theorem 5. Then

$$
\begin{aligned}
\left\|L^{*}(x)\right\| & =|\epsilon|\|L(f(\cdot, x(\cdot), \epsilon))\| \\
& \leqq|\epsilon| K\|f(\cdot, x(\cdot), \epsilon)\| \leqq|\epsilon| \mathrm{KN}
\end{aligned}
$$

where $K$ is the bound on $L$ from Theorem 5. Thus, $L^{*}$ maps bounded fundtions into bounded functions. Furthermore, if $|\epsilon| K N \leqq M$, we get that $\left\|L^{*}(x)\right\| \leqq M$.

If $x^{*}$ is a fixed point of $L^{*}$, i.e., $x^{*}=L^{*}\left(x^{*}\right)$, then $x^{*}$ is the bounded solution of

$$
E x=A x+\epsilon f\left(t, x^{*}(t), \epsilon\right)
$$

and hence is the desired solution. We will show that such a fixed point exists by showing that $L^{*}$ is a contradiction mapping for $|\epsilon|$ small enough.

Let $x_{1}, x_{2}$ be bounded solutions with $\left\|x_{1}\right\| \leqq M$ and $\left\|x_{2}\right\| \leqq M$.
Then

$$
\begin{aligned}
& \left\|L^{*}\left(x_{1}\right)-L^{*}\left(x_{2}\right)\right\|=|\epsilon|\left\|L\left(f\left(\cdot, x_{1}(\cdot), \epsilon\right)\right)-L\left(f\left(\cdot, x_{2}(\cdot), \epsilon\right)\right)\right\| \\
& \quad=|\epsilon|\left\|L\left(f\left(\cdot, x_{1}(\cdot), \epsilon\right)-f\left(\cdot, x_{2}(\cdot), \epsilon\right)\right)\right\| \\
& \quad \leq|\epsilon| K\left\|f\left(\cdot, x_{1}(\cdot), \epsilon\right)-f\left(\cdot, x_{2}(\cdot), \epsilon\right)\right\| \\
& \quad \leq|\epsilon| K F\left\|x_{1}(\cdot),-x_{2}(\cdot)\right\| \rightarrow\left\|x_{1}(\cdot)-x_{2}(\cdot)\right\|
\end{aligned}
$$

since $f$ is uniformly Lipschitz continuous with Lipschitz constant $F$.
If $\epsilon$ is such that $|\epsilon| K N \leqq M$ and $|\epsilon| K F<1$, then $L^{*}$ is a contraction mapping on the set $\|x\| \leqq M$, and $L^{*}$ has one and only one fixed point $x^{*}$ with $\left\|x^{*}\right\| \leqq M$. This is the case for all $\epsilon$ such that $|\epsilon|<\epsilon_{1}$ and

$$
\epsilon_{1}=\min \left(\epsilon_{0}, \frac{M}{K N}, \frac{1}{K F}\right)
$$

1) If $f(t, x, \epsilon)$ is continuous in $t$ for fixed $x, \epsilon$, then $f(t, x(t), \epsilon)$ is continuous in $t$ whenever $x$ is continuous in $t$, thus $L *(x)$ is continuous in $t$. In the iterations used to prove the existence and uniqueness of $x^{*}$, the initial function $x_{1}$ is chosen continuous in $t$, then each iterate $x_{K}$ will be continuous in $t$ and so the limit $x^{*}$ will be continuous in $t$.
2) If $f(t, x, \epsilon)$ is $T$-periodic in $t$ for fixed $x, \epsilon$, then $f(t, x(t), \epsilon)$ is T-periodic in $t$ for any T-periodic $x$ and $L^{*}(x)$ is Tperiodic will lead to the limiting function $\mathrm{x}^{*}$ being T -periodic.
3) If $f(t, x(t), \epsilon)$ is almost periodic in $t$ for any almost periodic $x$, then $L *(x)$ will be almost periodic and choosing an almost periodic initial function $x_{1}$ will lead to an almost periodic limiting function $x^{*}$.

## AFFINE DIFFERENCE EQUATIONS -- SIMPLEST CRITICAL CASE

The problem considered here is to find a T-periodic solution of the difference equation (28).

$$
\begin{equation*}
E x=x+f(t) \tag{28}
\end{equation*}
$$

where $f(t)$ is $T$-periodic for some $T \geqq I$ and $f$ is integrable.
This case is called critical because the homogeneous equation has non-trivial T-periodic solutions. It is the simplest critical case because the matrix $A$ is the simplest possible matrix, the unit matrix $I$.

We use the notation that [ $t$ ] is the largest integer less than or equation to $t$ and $t=[t]+\langle t\rangle$. We call [ $t$ ] the integer part of $t$ and $\langle t\rangle$ the fractional part of $t$ since $[t]$ is an integer and $0 \leqq\langle t\rangle<1$. For each integer $n>0$, the solution $x(t)$ is given by

$$
x(t)=x(t-n)+\sum_{j=1}^{n} f(t-j)
$$

Letting $n=[t]$, we get (29) for $t \geqq 1$

$$
\begin{equation*}
x(t)=x(\langle t\rangle)+\sum_{j=1}^{[t]} f(t-j) \tag{29}
\end{equation*}
$$

where $x(\langle t\rangle)=x_{0}(\langle t\rangle)$ is the initial function. The problem is to determine under what conditions of $f$ will there be an $x_{0}$ such that $x(t)$ given in (29) is T-periodic. From (29) we get, using $f(t+T)=f(t)$,

$$
x(t+T)-x(t)=x(\langle t+T\rangle)-x(\langle t\rangle)+\sum_{j=[t]+1}^{[t+T]} f(t+T-j)
$$

$$
\begin{equation*}
=x(\langle t+T\rangle)-x(\langle t\rangle)+\sum_{j=1}^{[t+T]-[t]} f(\langle t\rangle-j) \tag{30}
\end{equation*}
$$

While equation (30) was derived using the solution defined only for $t \geqq 1$, it can be easily shown that equation (30) hold for all $t$ in ( $-\infty, \infty$ ). Thus, if we can find an initial function $x_{0}$ such that the right hand side of (30) is zero identically in $t$, then that initial function will give a T-periodic solution.

The first, and simplest, case considered is when the period $T$ is an integer $\geqq 1$. Then $\langle t+T\rangle=\langle t\rangle$ and $[t+T]=T+[t]$ and (30) becomes (31)

$$
\begin{equation*}
x(t+T)-x(t)=\sum_{j=1}^{T} f(\langle t\rangle-j) \tag{31}
\end{equation*}
$$

and $x(t)$ is T-periodic if and only if $f(t)$ satisfies (32) for all $t$.

$$
\begin{equation*}
\sum_{j=1}^{T} f(\langle t\rangle-j)=0 \tag{32}
\end{equation*}
$$

Whenever the solution is defined only for $t$ integer, then (32) is the condition that is necessary and sufficient for a T-periodic solution of (28) to exist [10].

For any real $T$, let $x_{1}$ be a $T$-periodic solution of the homogeneous equation $E x=x$. Since all solutions of the homogeneous equation are l-periodic, $x_{1}$ is both l-periodic and $T$-periodic. If $T$ is irrational, then the only possibility is $x_{1}(t)=$ constant. If $T=m / n$ where $m$ and $n$ are relatively prime integers, then any $1 / n$-periodic function will do for $x_{1}(t)$.

If $T$ is an integer, then any l-periodic function will do for $x_{1}(t)$. Suppose that

$$
I=\int_{0}^{T} x_{1}(t+s)^{*} f(t+s) d s \neq 0
$$

where $x_{1}(t+s) *$ denotes the complex conjugate transpose of $x_{1}(t+s)$. This integral is independent of $t$. For any solution $x(t)$ of (28), consider

$$
r(t)=\int_{0}^{T} x_{1}(t+s) * x(t+s) d s
$$

Since $x_{1}(t+s+1)=x_{1}(t+s)$, we have

$$
\begin{aligned}
r(t+1) & =\int_{0}^{T} x_{1}(t+s+1) * x(t+s+1) d s \\
& =\int_{0}^{T} x_{1}(t+s) *(x(t+s)+f(t+s)) d s=r(t)+I
\end{aligned}
$$

and hence

$$
\gamma(t)=\gamma(t\rangle)+[t] I .
$$

Since $|\gamma(t)| \rightarrow \infty$ as $t \rightarrow \infty$, we have that $x(t)$ must be unbounded as $T \rightarrow \infty$ and hence cannot be periodic. This condition $I \neq 0$ is the condition of resonance and is exactly analogous to resonance in differential equations.

LEMMA 9 A necessary condition for a T -periodic solution of (28) to exist is that

$$
\begin{equation*}
\int_{0}^{T} x_{1}(s) * f(s) d s=0 . \tag{33}
\end{equation*}
$$

for every T-periodic solution $x_{1}$ of the homogeneous equation $E x=x$.

If $T$ is an irrational number, then condition (33) becomes simply

$$
\int_{0}^{T} f(s) d s=0
$$

If $T=m / n$ where $m$ and $n$ are relatively prime integers, then condition (33) becomes

$$
\int_{0}^{T} x_{1}^{*}(s) f(s) d s=0
$$

for all $1 /$ n-periodic functions $x_{1}(t)$.
In the case where $T$ is an integer, all l-periodic functions $x_{1}$ are permissible and this condition becomes

$$
\begin{aligned}
I= & \int_{0}^{T} x_{1}(s) * f(s) d s=\sum_{j=1}^{T} \int_{T-j}^{T-j+1} x_{1}(s) * f(s) d s \\
& =\sum_{j=1}^{T} \int_{0}^{1} x_{1}(s) * f(s-j) d s \\
& =\int_{0}^{1} x_{1}(s) * \sum_{j=1}^{T} f(s-j) d s=0
\end{aligned}
$$

One candidate for $x_{1}(s)$ is

$$
x_{1}(s)=\sum_{j=1}^{T} f(\langle s\rangle-j)
$$

which is l-periodic. Using this $x_{1}(s)$ and letting $|x|^{2}=x^{*} x$ for any n-vector $x$, we get equation (34).

$$
\begin{equation*}
I=\int_{0}^{1}\left|\sum_{j=1}^{T} f(s-j)\right|^{2} d s=0 \tag{34}
\end{equation*}
$$

Comparing this to condition (32), we see that condition (32) implies condition (34) and also that condition (34) implies condition (32) almost everywhere in $t$. If we identify all functions which are equal except on sets of measure zero as the same function, then conditions (32) and (34) are equivalent. In the work that follows, this identification of functions shall always be made. Thus, when we write $f(t)=g(t)$, we mean that they agree for all $t$ except possibly on a set of measure zero.

Let $X_{1}=L_{2}[0,1)$. For a given $T$, let $S$ be the set of all $T$ periodic.functions $f$ which are square integrable over one period. For each $f \in S$, define $\|f\|$ by

$$
\| f^{2}=\frac{1}{T} \int_{0}^{T}|f(s)|^{2} d s
$$

The functions $\phi_{n} \in S$,

$$
\Phi_{n}(t)=\exp (i \omega n t) \quad n=0, \pm 1, \pm 2, \ldots
$$

where $\omega=2 \pi / T$, form a complete orthogonal basis in this Hilbert space
That is, each $f$ in $S$ can be written as (35)

$$
\begin{equation*}
f(t)=\sum_{n} f_{n} \phi_{n}(t) \tag{35}
\end{equation*}
$$

where the sum on $n$ is over all the integers, positive and negative, and the $f_{n}$ are the Fourier coefficients of $f$.

$$
\begin{equation*}
f_{n}=\frac{1}{T} \int_{0}^{T} \phi_{n}(-s) f(s) d s \quad n=0, \pm 1, \pm 2, \ldots \tag{35'}
\end{equation*}
$$

Let $S_{1}$ be the set of all square integrable T-periodic solutions of the homogeneous equation $E x=x$. Then $S_{1}$ is a subspace of $S$ and consists of functions which are both l-periodic and $T$-periodic. If $T$ is an irrational number, then $S_{1}$ consists only of constant functions and, if $x \in S_{1}$, then $x(t)=x_{0} \phi_{0}(t)$ for some constant vector $x_{0}$. If $T$ is a rational number and $T=M / N$ where $M$ and $N$ are relatively prime integers, then $S_{1}$ consists of all square integrable $1 / N$-periodic functions and, if $\mathbf{x} \in \mathrm{S}_{1}$, then

$$
x(t)=\sum_{n} x_{M n}{ }^{\phi} M_{n}(t)
$$

where $X_{M n}$ is the Mn-th Fourier coefficient of $x$.
Let $P$ be the projection operator which maps $S$ onto $S_{1}$. That
is, if $x \in S$, then $P x \in S_{1}$ and, if $T$ is irrational

$$
\begin{equation*}
\operatorname{Px}(t)=x_{0} \phi_{0}(t)=\frac{1}{T} \int_{0}^{T} x(s) d s \tag{36}
\end{equation*}
$$

and, if $T=M / N$ as before, then

$$
\begin{equation*}
P x(t)=\sum_{n} x_{M n} \phi_{M n}(t) \tag{36'}
\end{equation*}
$$

Let $S_{2}$ be the orthogonal complement of $S_{1}$, i.e., all the functions $x$ in $S$ which satisfy $P x=0$.

If $\dot{x} \in S$, then $E x \in S$ and

$$
E x(t)=x(t+1)=\sum_{n} x_{n} \phi_{n}(t+1)=\sum_{n} x_{n} \phi_{n}(1) \Phi_{n}(t)
$$

and, if $T$ is irrational

$$
(P E x)(t)=\frac{1}{T} \int_{0}^{T} x(s+1) d s=\frac{1}{T} \int_{0}^{T} x(s) d s=(P x)(t)
$$

If $T=M / N$ as before

$$
(\operatorname{PEx})(t)=\sum_{\mathrm{n}} \mathrm{x}_{\mathrm{Mn}}{ }^{\phi} \mathrm{Mn}^{(I) \phi_{\mathrm{Mn}}}(\mathrm{t})=\sum_{\mathrm{n}} \mathrm{x}_{\mathrm{Mn}} \phi_{\mathrm{Mn}}(\mathrm{t})=(\mathrm{Px})(\mathrm{t})
$$

since $\phi_{M n} \in S_{1}$ (hence $\left.\Phi_{M n}(1)=1\right)$ for $n=0, \pm 1, \pm 2, \ldots$. In either case, we get the important fact that $P E x=P x$ for any $x \in S$.

IFMMM 10 The condition $P f=0$ is necessary for a T-periodic solution of the difference equation (28) to exist.

$$
\begin{equation*}
E x=x+f(t) \tag{28}
\end{equation*}
$$

This condition $P f=0$ is the same as conditions (32) and (33).

Lemma 10 and Lemma 9 are the same. Both are included since they are proven using different terminology. The condition $P f=0$ is the same as condition (33) but the former condition is easier to picture geometrically and
is easier to check. $P$ is a projection operator which projects any T-periodic $f$ in $S$ onto the subspace $S_{1}$ of $T$-periodic solutions of $E x=x$. Thus, the condition $P f=0$ is that $f$ must be orthogonal to the subspace $S_{1}$. $P$ has the explicit representation given in equation (36), so we can check if $P f=0$ merely by checking certain Fourier coefficients of $f$.

PROOF: Suppose $x$ is a T-periodic solution of (28). Using this $x$, operate with the projection $P$ on equation (28).

$$
P E x=P x+P f
$$

Since $P E x=P x$, we get $P f=0$.

LEMMA 11 If $x \in S$ is a T-periodic solution of (28), then there is a unique $T$-periodic solution $\hat{x}$ of (28) which satisfies $P \hat{x}=0$. This $\hat{x}$, considered as a function of $f \in S, P f=0$, is a linear function of f. If the T-periodic solution $x$ can be chosen in a manner such that it is a continuous function of $f$, then $\hat{x}: S_{2} \rightarrow S$ is a continuous linear function of $f$ and there is some $k, 0 \leqq k<\infty$, such that $\|\hat{x}(f)\| \leqq k\|f\|$.

PROOF: For any $x \in S$, the function $\hat{x}=(1-P) x$ is in $S$ and satisfies $P \hat{x}=0$.

Let $\hat{\mathbf{x}}$ be another $T$-periodic solution of (28) satisfying $p \hat{\hat{x}}=0$. Then $\hat{x}-\hat{\mathbf{x}}$ is a solution of the homogeneous equation $E x=x$ and thus is an element of $S_{1}$. So

$$
\hat{x}-\hat{\hat{x}}=P(\hat{x}-\hat{\hat{x}})=P \hat{x}-P \hat{x}=0
$$

and $\hat{x}=\hat{\hat{x}}$. Hence the $T$-periodic solution $\hat{x}$ satisfying $\hat{P} \hat{x}=0$ exists and is unique. Denote this as a function of $f$ by $\hat{x}=\hat{x}(f)$.

Since $\alpha \hat{x}(f)$ is a solution of (28) with $f$ replaced by $\alpha$ for any scalar $\alpha$, $\alpha \hat{x}(f)$ is T-periodic and satisfies $P(\alpha \hat{x}(f))=\alpha P \hat{x}(f)=0$, we see that $\alpha \hat{x}(f)=\hat{x}(\alpha f)$.

Since $\hat{x}\left(f_{1}\right)+\hat{x}\left(f_{2}\right)$ is a solution of (28) with $f$ replaced by $f_{1}+f_{2}$ for any $f_{1} \in S, f_{2} \in S, f_{1}=P f_{2}=0, \hat{x}\left(f_{1}\right)+\hat{x}\left(f_{2}\right)$ is $T$-periodic and satisfies $P\left(\hat{x}\left(f_{1}\right)+\hat{x}\left(f_{2}\right)\right)=P \hat{x}\left(f_{1}\right)+P \hat{x}\left(f_{2}\right)=0$, we see that $\hat{x}\left(f_{1}\right)+\hat{x}\left(f_{2}\right)=$ $\hat{x}\left(f_{l}+f_{2}\right)$. Hence, $\hat{x}(f)$ is linear in $f$.

If x is chosen in a manner such that it is a continuous function of $f$, then, since $\hat{x}=(I-P) x$ is a continuous function of $x, \hat{x}$ is a continuous function of $f$. The inequality $\|\hat{x}(f)\| \leqq k\|f\|$ follows from the fact that $\hat{x}$ is a continuous linear function of $f$ [7].

This preceding lemma shows that, if a T-periodic solution exists, then
there is a unique solution $\hat{x}$ which has all the desired properties, notably that $P_{0} \hat{x}=0$. The problem now is one of finding a $T$-periodic solution to (28). For the case when $T$ is a rational number, the solution $\hat{x}=\hat{x}(\dot{f})$ can be found by comparing Fourier series directly.

THEOREM 7 If $f \in S$ satisfies $P f=0$ and $T=M / N$ where $M$ and $N$ are relatively prime integers, then there is a continuous linear operator $L: S_{2} \rightarrow S_{2}$ such that $\hat{x}=L(f)$ is the unique $T$-periodic solution of (28) which

$$
\begin{equation*}
E x=x+f(t) \tag{28}
\end{equation*}
$$

PROOF: If there is such an $\hat{x} \in S_{2}$, then it can be represented by its Fourier Series.

$$
\hat{x}(t)=\sum_{n} x_{n} \phi_{n}(t)
$$

where the $x_{n}, n=0, \pm 1, \pm 2, \ldots$, are the Fourier coefficients of $\hat{x}$. Since $f \in S_{2}$, it has the Fourier series

$$
f(t)=\sum_{h} f_{n} \phi_{n}(t)
$$

The condition $P f=0$ is the condition that $f_{n}=0$ whenever $n / T$ is an integer. Then

$$
\hat{x}(t+1)-\hat{x}(t)-f(t)=\sum_{n}\left(\left(\phi_{n}(1)-1\right) x_{n}-f_{n}\right) \phi_{n}(t)
$$

If $\hat{\boldsymbol{x}}$ is to be a solution of (28), then each of the coefficients must be zero, i.e.,

$$
\left(\Phi_{n}(1)-1\right) x_{n}=f_{n} \quad n=0, \pm 1, \pm 2, \ldots
$$

If $n / T$ is an integer then $f_{n}=0, \phi_{n}(1)=1$, and $x_{n}$ is arbitrary. The condition $P \hat{x}=0$ specifies $x_{n}=0$ whenever $n / T$ is an integer. If $n / T$ is not an integer, then $\phi_{n}(1) \neq 1$ and $x_{n}$ is given by (37)

$$
\begin{equation*}
x_{n}=\frac{f_{n}}{\varphi_{n}(1)-1} \tag{37}
\end{equation*}
$$

Thus, $\hat{x}$ is a linear function of $f$. This preceding work is valid if we can show that this. $\hat{\mathrm{x}}$ is square integrable.

Since $T$ is a rational number and $\phi_{n}(t)=\exp ($ ian $), T=M / N$ where
$M$ and $N$ are relatively prime integers and

$$
\left|e^{i a n}-1\right|^{2}=4 \sin ^{2} \frac{1}{2} \text { an } \geqq 4 \sin ^{2} \frac{\pi}{M}=\epsilon^{2}>0
$$

for all $n$ such that $n / T \neq$ an integer. Then

$$
\sum_{n}\left|x_{n}\right|^{2}=\sum_{n}^{\prime} \frac{\left|f_{n}\right|^{2}}{\left|e^{i \omega n}-1\right|^{2}} \leqq \frac{1}{\epsilon} \sum_{n}\left|f_{n}\right|^{2}
$$

where $\sum_{n}^{\prime}$ means the sum over all the integers $n$ such that $n / T \neq$ an integer. This shows that $\hat{x}$ is in $S_{2}$ and is a continuous linear function of $f$. The linear operator $L$ is defined by

$$
\hat{x}(t)=(L f)(t)=\sum_{n}^{\prime} \frac{f_{n}}{\left(\Phi_{n}(1)-1\right)} \Phi_{n}(t)
$$

where $f_{n}$ is the n-th Fourier coefficient of $f$ and $\sum_{n}^{\prime}$ has the same meaning as before. By Lemma 11, we know that this is the only such solution.

$$
\text { If } T \text { is an irrational number, then the Fourier coefficients of } \hat{x}
$$ would still be defined by (37) (except for $x_{0}$ which is zero). The problem then is one of showing that

$$
\sum_{n}\left|x_{n}\right|^{2}=\|\hat{x}\|^{2}
$$

is bounded or under what conditions is $\hat{\mathbf{x}}$ a continuous linear function of $f$.

The permissible functions $f$ will probably have to be restricted to be some subset of $S_{2}$. For one solution of this problem, see Moser [11].

If there is a continuous linear operator $L: S_{2} \rightarrow S_{2}$ such that $\hat{x}=$ Lf is the unique solution of (28) which satisfies $\mathrm{P} \hat{\mathbf{x}}=0$, then a class of nonlinear problems could be studied in a manner similar to Hale [12]. However, there are problems in determining the meaning of the fact that the non-linear term must be bounded and Lipschitz continuous in the norm of $s$, i.e. the $L_{2}[0, T]$ norm.

As an example, consider the Crude Euler difference analog to the differential equation

$$
\dot{x}=f(t)
$$

where $f(t+T)=f(t)$. Letting $t=h s$, the analog is

$$
E x=x+h f(h s)=x+g(s)
$$

If we chose $h$ such that $T=h N$ for some integer $N$, then this difference analog' is like equation (28). In this case, the condition $\mathrm{Pg}=0$, which is necessary and sufficient for an N-periodic solution of the difference analog to exist, is that

$$
\int_{0}^{N} e^{i 2 \pi n s} g(s) d s=0 \quad n=0, \pm 1, \pm 2, \ldots
$$

or, in terms of $f$ and $t$

$$
\int_{0}^{T} e^{i \frac{2 \pi n t}{h}} f(t) d t=\int_{0}^{T} e^{i \frac{2 \pi}{T} N n t} f(t) d t=0 \quad n=0, \pm 1, \pm 2, \ldots
$$

If

$$
\int_{0}^{T} f(t) d t=0
$$

then a T-periodic solution to the differential equation does exist. However, we see that this is only one of the conditions needed to have a N-periodic solution of the difference analog to exist. In fact, every N-th Fourier coefficient of $f$ must be zero before there exists an $N$-periodic solution to the difference analog. If one or more of these Fourier coefficients are not zero, then there does not exist an $N$-periodic solution to the difference analog.

If $T / h=N$ is an irrational number, then the condition $\mathrm{Pg}=0$ reduces to the condition which is necessary and sufficient for a T-periodic solution of the differential equation to exist. However, I have not been able to show that the condition $\mathrm{Pg}=0$ is sufficient for an $N$-periodic solution of the difference analog to exist when $N$ is an irrational number.

The aim of all this work is to develop a bifurcation theory analogous to that in Hale [12]. The first step in this theory is to show that there exists a continuous linear mapping $L: S_{2} \rightarrow S_{2}$ such that $\hat{x}=L f$ is the unique $T-$ periodic solution of (28) which satisfies $P \hat{x}=0$. This has yet to be done for the case when $T$ is an irrational number. Yet, from the above example, it would appear that this would be the most interesting case.

AN AVERAGING METHOD WITH AN APPLICATION
Halanay [10] gives the following theorems for a system of almostperiodic difference equations. While all his work is concerned with the solution being defined only on the integers, his proofs are easily extended to include the solution being defined on the reals. In [10], Halanay first gives an approximation Lemma which is stated here.

LEMMA 12 If $f(t)$ is such that
a). $|f(t)| \leqq F$ for all $t$, some $F \geqq 0$
b) $\left|\frac{1}{N} \sum_{k=1}^{N} f(t-k)\right| \leqq \epsilon(N)$ for all $t$, some $E(N) \geqq 0$ and $\epsilon(N) \rightarrow 0$ as $N \rightarrow \infty$. For some $\eta, 0<\eta<1$, let $f_{\eta}$ be the function defined by

$$
f_{\eta}(t)=\sum_{k=1}^{\infty}(1-\eta)^{k-1} f(t-k) .
$$

Then there exists a continuous function $\delta(\eta)$ with $\delta(\eta) \rightarrow 0$ as $\eta \rightarrow 0$ such that

$$
\begin{aligned}
& \eta\left|f_{\eta}(t)\right| \leqq \delta(\eta) \text { for all } t \\
& \left|f_{\eta}(t+1)-f_{\eta}(t)-f(t)\right| \leqq \delta(\eta) \text { for all } t
\end{aligned}
$$

That is, $f_{\eta}$ is an approximate solution to the difference equation (28) which can be made as accurate as desired by taking $\eta$ small enough.

$$
\begin{equation*}
E x=x+f(t) \tag{28}
\end{equation*}
$$

Condition b) is the natural one to use for difference equations since the solution of (28) is a sum of $f(t-k)$. However, this condition is usually very hard to verify. When this lemma is applied to the numerical solution of differential equations, we will have $f(t)=g(h t)$ for some $h>0$ and

$$
\left|\frac{1}{T} \int_{0}^{T} g(t-s) d s\right| \leq \epsilon_{1}(T) \text { for all } t
$$

where $\epsilon_{1}(T) \rightarrow 0$ as $T \rightarrow \infty_{\text {. }}$
The following identity, a modification of the Euler-Maclaurin summation formula, is easily verified by integration by parts.

$$
\sum_{k=1}^{N} f(t-k)=\int_{0}^{N} f(t-s) d s-\int_{0}^{N} f^{\prime}(t-s)\langle s\rangle d s
$$

where $f^{\prime}(t)$ is the first derivative of $f$ with respect to $T$. Making the change of variables $f(t)=g(h t)$, we get

$$
\frac{1}{N} \sum_{k=1}^{N} f(t-k)=\frac{1}{T} \int_{0}^{T} g(h t-s) d s-\frac{h}{T} \int_{0}^{T} g^{\prime}(h t-s)\left\langle\frac{\dot{u}}{h}\right\rangle d u
$$

where $T \cdot=$ Nh. In general, this last integral will not go to zero as $T \rightarrow \infty$ even though the first one does. However, under fairly general conditions, e.g., $\left|g^{\prime}(t)\right|$ bounded for all $t$, we can bound this last integral by bG.

$$
\left|\frac{h}{T} \int_{0}^{T} g^{\prime}(h t-s)\left\langle\frac{s}{h}\right\rangle d s\right| \leqq h G \text { for all } t
$$

Then

$$
\left|\frac{1}{N} \sum_{k=1}^{N} f(t-k)\right| \leqq \epsilon_{1}(T)+h G=\epsilon_{1}(N h)+h G .
$$

In general, we get this sum going to zero only by letting $h$ go to zero. However, if we look at the proof of Lemma 12, we see that the critical inequality is

$$
\left|f_{\eta}(t)\right| \leqq \frac{N \epsilon(N)}{1-(1-\eta)^{N}}+N F
$$

and we get the desired $\delta(\eta)$ by letting $N=N(\eta)$ be chosen as a function of $\eta$ such that

$$
\epsilon(\mathbb{N}) \leqq 1-(1-\eta)^{\mathbb{N}}
$$

Then we get $\delta(\eta)=\eta N(\eta)(1+F)$.
If; instead of the sum of $f$ went to zero; we knew only that the
integral of $g$ went to zero where $f(t)=g(h t)$, then this crucial inequality becomes:

$$
\left|f_{\eta}(t)\right| \leqq \frac{N\left(\epsilon_{1}(N h)+h G\right)}{1-(1-\eta)^{N}}+N F
$$

We now let N and h be chosen as functions of $\eta$ such that

$$
\begin{array}{r}
2 \mathrm{hG} \leqq 1-(1-\eta)^{\mathrm{N}} \\
2 \epsilon_{1}(\mathrm{Nh}) \leqq 1-(1-\eta)^{\mathrm{N}}
\end{array}
$$

These functions $h=h(\eta)$ and $N=N(\eta)$ satisfy $h(\eta) \rightarrow 0$ as $\eta \rightarrow 0$, $h(\eta) N(\eta) \rightarrow \infty$ as $\eta \rightarrow 0$, and $\eta N(\eta) \rightarrow 0$ as $\eta \rightarrow 0$. Then we get

$$
\delta(\eta)=\eta \mathbb{N}(\eta)(1+F)
$$

Thus, Lemma 12 is still true. However, this solution has the undesirable property that we must let $h \rightarrow 0$ as $\eta \rightarrow 0$. This means that, in order to have $\delta(\eta)$ small, i.e., have $f_{\eta}(t)$ be an accurate solution of (28), then we must take both $N$ large and $h$ small.

There is one exception to this problem of $h(\eta) \rightarrow 0$ as $\eta \rightarrow 0$. If $g(t)$ is P-periodic and $\mathrm{P} / \mathrm{h}$ is an irrational number, then the theory of equidistributed sequences [ 14,15 ] gives us that

$$
\left|\frac{1}{N} \sum_{k=1}^{N} g(h t-h k)-\frac{1}{P} \int_{0}^{P} g(s) d s\right| \rightarrow 0 \text { as } N \rightarrow \infty .
$$

In this case, if

$$
\frac{1}{\bar{P}} \int_{0}^{P} g(s) d s=0
$$

then we have the existence of some $\epsilon_{3}(N)$ such that

$$
\left|\frac{1}{N} \sum_{k=1}^{N} g(h t-h k)\right| \leqq \epsilon_{3}(N) \text { for all } t
$$

and $\epsilon_{3}(N) \rightarrow 0$ as $N \rightarrow \infty$. Thus, in this case, the conditions of Lemma 12 are satisfied. If $\mathrm{P} / \mathrm{h}$ is rational, then we can construct a g such that this sume does not go to zero as $N \rightarrow \infty$, and we have to pick $h$ small in order to get the desired accuracy.

Consider the difference equation (38) where $A(t), A^{-1}(t)$, and $f(t, x, \epsilon)$ are almost periodic in $t$ uniformly in $x$ and $\epsilon$.

$$
\begin{equation*}
E x=A(t) x+\epsilon f(t, x, \epsilon) \tag{38}
\end{equation*}
$$

Let $X(t)$ be the principle matrix solution of the homogeneous equation (39) with $X(\theta)=I, 0 \leqq \theta<1$, and let $X(t)$ and $X^{-1}(t)$ be almost periodic in $t$.

$$
\begin{equation*}
E x=A(t) x \tag{39}
\end{equation*}
$$

If $A$ and $f$ are periodic in $t$, then they are also almost periodic. The problem in requiring $X(t)$ to be periodic is that this choice of the initial function $X(\theta)=I, 0 \leqq \theta<1$, will not in general lead to a periodic solution even if there exists a fundamental matrix solution of (39) which is periodic. Determining that the principle matrix solution is almost periodic is an easier task than finding a fundamental matrix solution which is periodic. For example, if $A(t)=A$, a constant 2 by 2 matrix, both the eigenvalues of $A$ have simple elementary divisors, and the eigenvalues of $A$ lie on the unit circle, then the principle matrix solution $X(t)$ is almost periodic. A very careful choice of $X(\theta), 0 \leqq \theta<1$, will have to be made before $X(t)$ will be periodic. In this sense, it seems more natural to look for almost periodic solutions of difference equations than to look for periodic solutions. By the change of variables $x=X(t) y$, the difference equation (38) becomes

$$
\begin{equation*}
E y=y+\epsilon g(t, y, \epsilon) \tag{40}
\end{equation*}
$$

where

$$
g(t, y, \epsilon)=X(t+1)^{-1} f(t, X(t) y, \epsilon)
$$

is almost periodic in $t$ uniformly in $y$ and $\epsilon$. Let $g_{0}(y, \epsilon)$ be

$$
\begin{equation*}
g_{0}(y, \epsilon)=\lim _{N \rightarrow \infty} \frac{1}{N} \sum_{k=1}^{N} g(t-k, y, \epsilon) \tag{41}
\end{equation*}
$$

The average value of $g(t, y, \epsilon)$ (since $g(t, y, \epsilon)$ is almost periodic, the limit exists uniformly in $y$ and $\epsilon$ and is independent of $t$ ). Let $y_{o}$ be a soluLion of $g_{0}(y, 0)=0$ and let the matrix $H$ be

$$
H=\frac{\partial g_{0}}{\partial y}\left(y_{0}, 0\right)
$$

The following theorem is a restatement of Theorem 5 in Halanay [10].

THEOREM 8 The principle of averaging: If the real parts of all the eigenvalues of $H$ are less than zero, then there exists an almost periodic solution $x(t, \epsilon)$ of (38) for each $\epsilon, 0<\epsilon<\epsilon_{0}$, some $\epsilon_{0}>0$, which reduces to $x(t, 0)=X(t) y_{0}$ as $\epsilon \rightarrow 0$.

Consider as an example of the application of Theorem 8 the Crude Euler difference analog for (42)

$$
\begin{equation*}
\dot{x}=\epsilon f(t, x, \epsilon) \tag{42}
\end{equation*}
$$

where $f(t, x, \varepsilon)$ is almost periodic in $t$ uniformly in $x$ and $\epsilon$. Let the average value of $f(t, x, \epsilon)$ be $f_{0}(x, \epsilon)$.

$$
f_{0}(x, \epsilon)=\lim _{T \rightarrow \infty} \frac{1}{T} \int_{0}^{T} f(t-s, x, \epsilon) d s
$$

The Crude Euler analog is given in (43)

$$
\begin{equation*}
E x=x+\epsilon g(t, x, \epsilon) \tag{43}
\end{equation*}
$$

where

$$
g(t, x, \epsilon)=h f(h t, x, \epsilon)
$$

This equation already is in the form of equation (40) so the change of variables $x=X(t) y$ does not have to be performed. The average value of $g(t, x, \epsilon)$, $\mathrm{g}_{1}(\mathrm{x}, \epsilon)$, is given by

$$
g_{1}(x, \epsilon)=\lim _{T \rightarrow \infty} \frac{1}{T} \int_{0}^{T} g(t-s, x, \epsilon) d s=h f_{0}(x, \epsilon)
$$

This is not the average value $g_{0}(x, \epsilon)$ used in Theorem 8 but, considering the discussion following Lemma 12, if we choose $h$ as a function of $\epsilon$ properly, then $h=h(\epsilon) \rightarrow 0$ as $\epsilon \rightarrow 0$ and the conclusions of Theorem 8 still hold when $\mathrm{g}_{1}(\mathrm{x}, \epsilon)$ replaces $\mathrm{g}_{\mathrm{o}}(\mathrm{x}, \epsilon)$. The matrix H

$$
H=\frac{\partial g_{1}}{\partial x}\left(x_{0}, 0\right)=h \frac{\partial f_{0}}{\partial x}\left(x_{0}, 0\right)
$$

goes to zero as $h \rightarrow 0$ but this does not create any problems since $g(t, x, \varepsilon) \rightarrow 0$ as $h \rightarrow 0$ and the effects of these two phenomena cancel each other.

If $x_{0}$ is a solution of $f_{0}(x, 0)=0$ and the real parts of all the
eigenvalues of $H$ are negative for $h>0$, then there exists an almost periodic solution $x(t, \epsilon)$ of (42) for each $\epsilon, 0<\epsilon<\epsilon_{0}$, some $\epsilon_{0}>0$, which reduces to $x(t, 0)=x_{0}$ as $\epsilon \rightarrow 0$. See, for example, [12] for a proof of this statement. The significance of this application of Theorem 8 is that, under the same conditions that insure the existence of an almost periodic solution of the differential equation (42), and the step size $h$ is chosen small enough, there is an almost periodic solution of the difference analog (43) and, furthermore, these two solutions reduce to the same solution $x(t, 0)=x_{0}$ as $\epsilon \rightarrow 0$. While this is a case where the differential equation and its difference analog display the same qualitative behavior, i.e., the existence of an almost periodic solution, the above analysis leaves several questions unanswered. For example, do the two almost periodic solutions display the same stability properties? How close is the almost periodic solution of the difference analog to the almost periodic solution of the differential equation?

AN INVARIANT MANIFOLD METHOD WITH AN APPLICATION
Consider the Van der pol equation (44) with $\in>0$.

$$
\begin{equation*}
\ddot{x}-2 \epsilon \dot{x}\left(1-x^{2}\right)+x=0 \tag{44}
\end{equation*}
$$

With a constant spacing $h>0$ and a change of independent variable from $t$ to $s=h t$, the central difference analog to (44) becomes (45)

$$
\begin{align*}
& E y=x \\
& E x=\left(2-h^{2}\right) x-y+E h f(x, y) \tag{45}
\end{align*}
$$

where

$$
f(x, y)=f(x, y, h, \epsilon)=\frac{\left(1-x^{2}\right)\left(\left(2-h^{2}\right) x-2 y\right)}{1-\epsilon h\left(1-x^{2}\right)}
$$

It is well known that the Van der Pol equation (44) has a stable limit cycle which is a periodic solution with amplitude approximately 2 for small $\boldsymbol{\epsilon}$ [12]. It would seem reasonable to expect the difference analog (45) to display this same type of behavior, at least for the spacing $h$ small enough. Yet the examples given at the beginning of this paper show that such blanket assumptions about the behavior of the difference analog reflecting the behavior of the differential equation cannot be made. The difference analog may display behavior completely different from that displayed by the differential equation. The purpose of this section is to show that the qualitative behavior of the difference analog (45) is the same as the qualitative behavior of the differential equation (44) for $\epsilon$ small.

The central difference analog is chosen for the following reasons. It
involves no more than the two variables $x$ and $y$ and so solutions can be displayed on a piece of paper without undue difficulty. Also, at $\epsilon=0$, both the differential equation (44) and its difference analog (45) have the behavior of a center. This desirable property is not shared by such other methods as the Crude Euler and the Corrected Euler. The Central Difference analog is chosen over the Iterated Euler analog because the difference equation for the Central Difference analog is much simpler.

$$
\text { For } \epsilon=0 \text {, we know that the solutions of (45) lie on a level curve }
$$ of $V(x, y)$ where

$$
v(x, y)=\frac{1}{4}(x+y)^{2}+\alpha^{2}(x-y)^{2}
$$

and

$$
\alpha^{2}=\left(4-h^{2}\right) /\left(4 h^{2}\right) \quad, \quad \alpha>0
$$

A level curve of $V(x, y)$ is an ellipse as shown in Figure 6 where $2 \cos \omega=2-h^{2}$. For $h$ small, these ellipses are very long and narrow. If we make a change of variables to $u$ and $v$ where

$$
\begin{array}{ll}
u=\frac{1}{2}(x+y) & x=u+(v / 2 \alpha)  \tag{46}\\
v=\alpha(x-y) & y=u-(v / 2 \alpha)
\end{array}
$$

then $v(x, y)=u^{2}+v^{2}$ and, in the $u, v$ plane, level surfaces of $V$ are circles The difference equation (45) becomes the difference equation (47)

$$
\begin{align*}
& E u=\cos \omega u+\sin \omega v+\epsilon h U(u, v)  \tag{47}\\
& E v=-\sin \omega u+\cos \omega v+\epsilon h 2 \alpha U(u, v)
\end{align*}
$$

where

$$
U(u, v)=U(u, v, h, \epsilon)=\frac{1}{2} f(u+(v / 2 \alpha), u-(v / 2 \alpha), h, \epsilon) .
$$

The following identities are used in deriving the following equations. $2 \cos \omega=2-h^{2}, \quad \sin \omega=h^{2} \alpha, \quad \sin \frac{1}{2} \omega=\frac{1}{2} h, \quad \cos \frac{1}{2} \omega=\alpha h$.

Since all solutions lie on circles in the $u, v$ plane when $\epsilon=0$,
it is natural to change coordinates again to polar coordinates.

$$
\begin{array}{ll}
u=r \cos \theta & x=r \cos \left(\theta-\frac{1}{2} \omega\right) /(\alpha h)  \tag{48}\\
v=r \cos \theta & y=r \cos \left(\theta+\frac{1}{2} \omega\right) /(\alpha h)
\end{array}
$$

If we let $w=r^{2}$, then the difference equation (47) becomes

$$
\begin{align*}
& E w=w+\epsilon h R(w, \theta)  \tag{49}\\
& E \theta=\theta-w+\epsilon h \Theta(w, \theta)
\end{align*}
$$

where $R$ and $\theta$ are $2 \pi$-periodic in $\theta$. For the Van der Pol equation (44), we have -

$$
\begin{align*}
& R\left(r^{2}, \theta\right)=R\left(r^{2}, \theta, h, \epsilon\right)=\frac{1}{h^{2}}\left\{\left(2-h^{2}\right) x-2 y+\epsilon h f(x, y, h, \epsilon)\right\} f(x, y, h, \epsilon)  \tag{50}\\
& \Theta\left(r^{2}, \theta\right)=\Theta\left(r^{2}, \theta, h, \epsilon\right)=\frac{1}{\epsilon h} \operatorname{Arctan}\left\{\frac{\epsilon f(x, y, h, \epsilon) \cos \left(\theta-\frac{1}{2} \omega\right)}{r+\epsilon f(x, y, h, \epsilon) \sin \left(\theta-\frac{1}{2} \omega\right)}\right\}
\end{align*}
$$

where the arctangent is taken such that $-\frac{1}{2} \pi \leqq \operatorname{Arctan}(x) \leqq \frac{1}{2} \pi$. Notice that $\theta$
is continuous and bounded in $\epsilon$ as $\epsilon \rightarrow 0$ even though there is a $1 / \epsilon$ coefficient on the arctangent. It is also easy to show that both $R$ and $\Theta$ are continuous and remain bounded as $h \rightarrow 0$ despite the $1 / h^{2}$ and $1 / h$ coefficients since the formulas for $x$ and $y$ in terms of $r, \theta$, and $h$ are given by (48).

In the difference equation (49) we can look for either of two things:

1) a solution for $w$ and $\theta$ which satisfies $w(t+T)=w(t)$ and $\theta(t+T)=\theta(t)+2 \pi$. This represents a T-periodic solution of the difference equations (45) and (47). The great difficulty with this method is that $T$ is not known a priori since it will be, in general, a function of $\epsilon$ and $h$.
2) a parametric representation for $w$ in terms of $\theta$ which is $2 \pi$-periodic and continuous in $\theta$ such that one solution of (49) is $\theta=\theta(t)$ and $w=w(\theta(t))$, i.e., an invariant manifold which is a closed curve in the $u, v$ or the $x, y$ planes. This is what is used in practice since, in computing solutions of Van der Pol's equation (44) numerically, a periodic solution is "found" when the numeric solution displays such a closed curve. The problem here is to show that such an invariant manifold exists and to compare it to the limit cycle of (44).

The difference equation (49) looks a lot like the difference equation studied by Moser [11] if we write $\Theta$ as the sum of its average value $\Theta_{0}$ and the remainder $\theta_{1}=\theta-\theta_{0}$.

$$
\Theta_{0}(w, h, \epsilon)=\frac{1}{2 \pi} \int_{0}^{2 \pi} \Theta(w, \theta, h, \epsilon) d \theta .
$$

Unfortunately, one condition Moser requires is that $\partial \Theta_{0} / \partial w \geqq 1$ and, since $f(x, y, h, \epsilon)$ is odd in $x$ and in $y$, we have $\Theta_{0}=0$ identically in $w, h, \epsilon$.

Thus, the work done by Moser does not apply here.
Halanay [9] developes a theorem which can be applied to this problem after a slight modification. This theorem, Theorem 4 in Halanay [9], is restated here without proof.

THEOREM 9: On Existence of an Invariant Manifold. Consider the difference equation (51) where $y \in E^{n}, \theta$ is a real number, and $Y_{1}$ and $\theta$ are $2 \pi$-periodic in $\theta$.

$$
\begin{align*}
& E y=Y_{0}(t, y)+\epsilon Y_{I}(t, y, \theta, \epsilon)  \tag{51}\\
& E \theta=\theta+\alpha(t)+\epsilon \Theta(t, y, \theta, \epsilon)
\end{align*}
$$

Assume there exists positive constants $H, K_{1}, \mu$, and $q$ with $0<q<1$ such that, for all $t \geqq t_{0},|y| \leqq H$, and $0 \leqq \theta \leqq 2, Y_{0}, Y_{1}$ and $\theta$ have continuous first partial derivatives with respect to $y$ and $\theta$ and
(1) $\left\|\frac{\partial Y_{o}}{\partial y}(t, y)\right\| \leqq K_{1},\left\|\frac{\partial \Theta}{\partial y}(t, y, \theta, \epsilon)\right\| \leqq K_{1},\left\|\frac{\partial \Theta^{\prime}}{\partial \theta}(t, y, \theta, \epsilon)\right\| \leqq K_{1}$
(2) $\cdot\left\|\frac{\partial Y_{o}}{\partial y}\left(t, y_{1}\right)-\frac{\partial Y_{0}}{\partial y}\left(t, y_{2}\right)\right\| \leqq K_{1}\left\|y_{1}-y_{2}\right\|^{\mu}$
(3) $\left\|\frac{\partial y_{1}}{\partial y}\left(\theta_{1}, y_{1}, \epsilon\right)-\frac{\partial y_{1}}{\partial y}\left(\theta_{2}, y_{2}, \epsilon\right)\right\| \leqq K_{1}\left(\left\|y_{1}-y_{2}\right\|^{\mu}+\left\|\theta_{1}-\theta_{2}\right\|^{\mu}\right)$
(4) $\left\|\frac{\partial \Theta}{\partial \theta}\left(\theta_{1}, y_{1}, \epsilon\right)-\frac{\partial \Theta_{2}}{\partial \theta}\left(\theta_{2}, y_{2}, \epsilon\right)\right\| \leqq K_{2}\left(\left\|y_{1}-y_{2}\right\|^{\mu}+\left\|\theta_{1}-\theta_{2}\right\|^{\mu}\right)$
(5) $\quad Y_{0}(t, 0)=0$
(6) Let $A(t)=\frac{\partial Y_{o}}{\partial y}(t, 0)$ and, if $z(t)$ is any solution of $E z=A(t) z$, then

$$
\|z(t)\| \leqq K_{1} q^{t-t_{0}}\left\|z\left(t_{0}\right)\right\|
$$

for all $t \geqq t_{0}$
Then, for $|\epsilon|$ small enough there exists a function $p(t, \theta)$, constans $\ell=\ell(\epsilon)$ and $L=L(\epsilon)$, continuous and monotone in $\epsilon$, such that $\ell(0)=L(0)$ and $l(\epsilon)>0, L(\epsilon)>0$ for $\epsilon \neq 0$ and
a) $\|\mathrm{p}(t, \theta)\| \leqq \ell(\epsilon)$
b) $\left\|p\left(t, \theta_{1}\right)-p\left(t, \theta_{2}\right)\right\| \leqq L(\epsilon)\left|\theta_{1}-\theta_{2}\right|$
c) If $\left\|y_{0}\right\| \leqq \ell(\epsilon)$ then

$$
\left\|y\left(t_{;} t_{0}, y_{0}, \theta_{0}\right)-p\left(t, \theta\left(t ; t_{0}, y_{0}, \theta_{0}\right)\right)\right\| \leqq K_{1}{ }^{t-t}{ }^{t}\left\|_{y_{0}}-p\left(t_{0}, \theta_{0}\right)\right\|
$$

d) If $y_{0}=p\left(t_{0}, \theta_{0}\right)$, then $y\left(t ; t_{0}, y_{0}, \theta_{0}\right)=p\left(t, \theta\left(t ; t_{0}, y_{0}, \theta_{0}\right)\right)$
e) This $p(t, \theta)$ is unique and $2 \pi$-periodic in $\theta$
f) If $Y_{0}, Y_{1}$, and $\Theta$ are T-periodic in $t$, then so is $p$.
g) If $Y_{0}, Y_{1}$, and $\Theta$ are independent of $t$, then so is $p$.

The set of points $(t, y, \theta)$ where $y=p(t, \theta)$ is an invariant manifold of the system (56) (result d), it is never very far from the set ( $t, 0, \theta$ ) (result a), it is Lipschitzian in $\theta$ (result b), and solutions near $y=0$ approach the manifold exponentially as $t \rightarrow \infty$ (result c).

For the difference equation (49), let

$$
R_{0}(w)=\frac{1}{2 \pi} \int_{0}^{2 \pi} R(w, \theta) d \theta
$$

and let $R_{1}(w, \theta)=R(w, \theta)-R_{0}(w)$. Let $w_{0}$ be a solution of $R_{0}(w)=0$ and let $w=w_{0}+y$. Let $R_{0}^{\prime}(w)$ denote the first derivative of $R_{0}$ with respect to $w$. Then the difference equation (49) becomes (52)

$$
\begin{align*}
& E y=\left(1+\epsilon h R_{0}^{!}\left(w_{0}\right)\right) y+\epsilon h R_{2}(y, \theta)  \tag{52}\\
& E \theta=\theta-\omega+\epsilon h \Theta\left(w_{o}+y, \theta\right)
\end{align*}
$$

where

$$
R_{2}(y, \theta)=R_{0}\left(w_{0}+y\right)-R_{0}^{\prime}\left(w_{0}\right) y+R_{1}\left(w_{0}+y, \theta\right) .
$$

This equation (52) is of the form of equation (51) with the matrix $A(t)=1+\epsilon h R_{0}^{\prime}\left(w_{0}\right)$. If $\epsilon h R_{0}^{\prime}\left(w_{0}\right)<0$, then the conditions of Theorem 9 seem to be satisfied. The problem is that the estimate $q$ is given by $q=1+\in h R_{0}^{\prime}(w)$ and $q \rightarrow 1$ as $\epsilon \rightarrow 0$ or as $h \rightarrow 0$. Since the proof of Theorem 9 depends on forming a contraction mapping where the contraction constant is $\alpha=\epsilon \mathrm{K} /(1-q)$ for some $K>0$, we see that this theorem cannot be really applied to this problem.

Try letting $w=p+\epsilon h u(p, \theta, \epsilon)$. Then the difference equation (49)
for $w$ becomes (53)

$$
\begin{gather*}
E p-p+\epsilon(u(E p, E \theta, \epsilon)-u(p, E \theta, \epsilon))=\epsilon h R_{0}(p, 0)-\epsilon h(u(p, E \theta, \epsilon)-u(p, \theta, \epsilon)- \\
\left.R_{1}(p ; \theta, \epsilon)\right) \tag{53}
\end{gather*}
$$

$$
+\epsilon h\left(R_{0}(p+\epsilon u, \epsilon)-R_{0}(p, 0)\right)+\epsilon h\left(R_{1}(p+\epsilon u, 8, \epsilon)-R_{1}(p, \theta, \epsilon)\right)
$$

where $E \theta$ is the function of $\theta, \epsilon$, and $w=p+\epsilon u$ given in the second equation of (49). Since $R_{0}$ and $R_{1}$ are continuous in all their arguments, there is some continuous monotone function $\delta(\epsilon)$ with $\delta(0)=0$ and $\delta(\epsilon)>0$ for $\epsilon>0$ such that

$$
\begin{aligned}
& \left|R_{0}(p+\epsilon u, \epsilon)-R_{0}(p, 0)\right| \leqq \delta(\epsilon) \\
& \left|R_{1}(p+\epsilon u, \theta, \epsilon)-R_{1}(p, \theta, \epsilon)\right| \leqq \delta(\epsilon)
\end{aligned}
$$

Furthermore, we can choose $h$ as a function of $p, \theta, \epsilon$ such that (see Lemma 12)

$$
\begin{equation*}
u(p, E \theta, \epsilon)-u(p, \theta, \epsilon)-R_{1}(p, \theta, \epsilon)=-\epsilon u(p, \theta, \epsilon) \tag{54}
\end{equation*}
$$

and

$$
\epsilon|u(p, \theta, \epsilon)| \leqq \delta(\epsilon)
$$

Thus, (53) becomes (55)

$$
\begin{align*}
E p-p+\epsilon & (u(E p, E \theta, \epsilon)-u(p, E \theta, \epsilon)  \tag{55}\\
& =\epsilon \epsilon_{R_{0}}(p, 0)+\epsilon \mathrm{hR}_{2}(p, \theta, \epsilon)
\end{align*}
$$

where

$$
\left|R_{2}(p, \theta, \epsilon)\right| \leqq 3 \delta(\epsilon) \rightarrow 0 \text { as } \epsilon \rightarrow 0
$$

Since $u$ satisfies (54), we get that

$$
\frac{\partial u}{\partial p}(p, E \theta, \epsilon)-\frac{\partial u}{\partial p}(p, \theta, \epsilon)-\frac{\partial R_{1}}{\partial p}(p, \theta, \epsilon)=-\epsilon \frac{\partial \dot{u}}{\partial p}(p, \theta, \epsilon)
$$

and that

$$
\epsilon\left|\frac{\partial u}{\partial p}(p, \theta, \epsilon)\right| \leqq \delta(\epsilon) \rightarrow 0 \text { as } \epsilon \rightarrow 0 .
$$

Thus, by the implicit function theorem, see [18] for example, there is a fundtron $H=H(p, y, \theta, \epsilon)$ such that $H(p, 0, \theta ; \epsilon)=0$ and

$$
\begin{equation*}
E p-p=\epsilon h R_{0}(p, 0)+\epsilon h R_{3}(p, \theta, \epsilon) \tag{56}
\end{equation*}
$$

is the solution of (55) for $\epsilon$ small enough where

$$
R_{3}(p, \theta, \epsilon)=R_{2}(p, \theta, \epsilon)+H\left(p, \epsilon h R_{0}(p, 0)+\epsilon h R_{0}(p, \theta, \epsilon), E \theta, \epsilon\right)
$$

and

$$
\begin{equation*}
\left|R_{3}(p, \theta, \epsilon)\right| \leqq \delta_{1}(\epsilon) \rightarrow 0 \text { as } \epsilon \rightarrow 0 \tag{56}
\end{equation*}
$$

Let $p_{0}$ be a root of $R_{0}(p, 0)=0$ and let $p=p_{0}+v$. Then the difference equation (56) becomes, in terms of $v$,

$$
\begin{equation*}
\mathrm{Ev}=\left(\mathrm{I}+\epsilon \mathrm{hR}_{\mathrm{o}}^{\mathrm{i}}\left(\mathrm{p}_{\mathrm{o}}, 0\right)\right) \mathrm{v}+\epsilon \mathrm{hR}_{4}(\mathrm{v})+\epsilon \mathrm{hR}_{3}\left(\mathrm{p}_{\mathrm{o}}+\mathrm{v}, \theta, \epsilon\right) \tag{57}
\end{equation*}
$$

where

$$
\begin{equation*}
R_{4}(v)=R_{0}\left(p_{0}+v, 0\right)-R_{o}^{\prime}\left(p_{0}, 0\right) v \tag{57'}
\end{equation*}
$$

and there is some $\delta_{2}(|v|)>0$ for $|v|>0, \delta_{2}(|v|) \rightarrow 0$ as $|v| \rightarrow 0$, and

$$
\left|R_{4}(v)\right| \leqq|v| \delta_{2}(|v|)
$$

If none of the eigenvalues of $R_{o}^{\prime}\left(p_{0}, 0\right)$ lie on the imaginary axis,
then there is some $\epsilon_{1} h_{1}>0$ such that, for $0<\epsilon h<\epsilon_{1} h_{1}$, none of the eigenvalues of ( $\mathrm{I}+\mathrm{ChR}_{\mathrm{o}}\left(\mathrm{p}_{\mathrm{o}}, 0\right)$ ) lie on the unit circle and, from Theorem 5, there is a continuous linear operator $L: B \rightarrow B$ such that $V=L f$ is the unique bounded solution of $E v=\left(I+\epsilon R_{o}^{\prime}\left(p_{o}, 0\right)\right) v+f$. There is some constant $K>0$ such that

$$
\|\operatorname{Lf}\| \leqq \frac{K}{h \epsilon}\|f\| .
$$

We are interested in a fixed point of the non-linear operator $N: B \rightarrow B$

$$
\mathrm{N}(v)=\mathrm{L}\left(\epsilon \mathrm{hR}_{4}(\mathrm{v})+\epsilon \mathrm{hR}_{3}\left(\mathrm{p}_{\mathrm{o}}+\mathrm{v}, \theta, \epsilon\right)\right)
$$

We see that, since we have the inequalities (56') and (57'), we can choose the set $S=\left\{v \in B:\|v\| \leqq \delta_{3}\right\}$, i.e., $\delta_{3}$, and $h$ as functions of $\epsilon$ for $\epsilon$ small. enough so that $N$ is a contraction mapping of $S$, hence $N$ has a fixed point $v$. Since $R_{3}$ is $2 \pi$-periodic in $\theta$, this fixed point will be $2 \pi$-periodic in $\theta$.

## These results are summarized in the following theorem.

THEOREM 10 Consider the difference equation (49).

$$
\begin{align*}
& E_{w}=w+\epsilon \mathrm{hR}_{0}(w, \epsilon)+\epsilon \mathrm{hR}_{1}(\mathrm{w}, \theta, \epsilon)  \tag{49}\\
& \mathrm{E} \theta=\theta-\omega+\epsilon \mathrm{h} \Theta(\mathrm{w}, \theta, \epsilon)
\end{align*}
$$

where $R_{0}$ and $R_{1}$ have continuous first partial derivatives with respect to $w$ and

$$
\frac{1}{2 \pi} \int_{0}^{2 \pi} R_{1}(w, \theta, \epsilon) d \theta=0 .
$$

If $w_{0}$ is a solution of $R_{0}\left(w_{0}, 0\right)=0$ and none of the eigenvalues of the matrix $R_{0}^{\prime}\left(w_{0}, 0\right)$ lie on the imaginary axis, then there is an $\epsilon_{1}>0$, and functions $h(\epsilon)$ and $w(\theta, \epsilon, h)$ such that, if $0<\epsilon<\epsilon_{1}, 0<h<h(\epsilon)$, then $\mathrm{w}(\theta, \epsilon, \mathrm{h})$ represents an invariant curve of (49). Also, $\mathrm{h}(\epsilon) \rightarrow 0$ as $\epsilon \rightarrow 0$ and $w(\theta, \epsilon, h) \rightarrow w_{0}$ as $\epsilon \rightarrow 0$.

This is a generalization of Halanay's theorem [9] which is given here as Theorem 9.

If the application to Van der Pol's equation (44), we have

$$
\begin{aligned}
& R_{0}(w, 0)=\frac{1}{2} w\left(4-h^{2}-w\right) \\
& R_{0}^{\prime}(w, 0)=\frac{1}{2}\left(4-h^{2}-2 w\right)
\end{aligned}
$$

The two roots of $R_{0}(w, 0)$ are $w_{1}=0$ and $w_{2}=4-h^{2}$. For $w=w_{1}, R_{0}^{\prime}\left(w_{1}, 0\right)=$ $\frac{1}{2}\left(4-h^{2}\right)>0$ and we have an invariant curve $w(\theta, \epsilon, h)=0$. For $w=w_{2}$, $R_{0}^{\prime}\left(w_{2}, 0\right)=-\frac{1}{2}\left(4-h^{2}\right)<0$, and we have an invariant curve $w(\theta, \epsilon, h) \doteq 4-h^{2}$. Thus we see that the central difference analog (45) to the Van der Fol equation (44) has the same qualitative behavior ( $w$ is the square of the amplitude).

SUGGESTIONS FOR FUTURE RESEARCH
The theory of difference equations as presented here is very incomplete. More work needs to be done with finding periodic solutions in critical cases, i.e., in cases when the homogeneous difference equation also has periodic solutions. This includes the case for $E x=x+f$, considered here but when the period $T$ is an irrational number. This also includes the case for $E x=A x+f$ where the matrix $A$ has eigenvalues on the unit circle which may or may not have simple elementary divisiors or where the matrix $A$ is a $T$ periodic function of $t$.

More work also needs to be done in studying the behavior of various difference analogs to differential equations which are not considered here. For example, what would a graph of the regions of behavior of the Kutta-Simpson analog for $\ddot{x}+2 c \dot{x}+x=0$ plotted as in Figure 3 look like? How do the discreprencies affect the usefullness of these methods for finding periodic solutions of a differential equation. If the differential equation has a periodic solution which is exponentially stable, can Theorem 10 be applied to the difference analog to show that it has an invariant manifold which is close to this periodic solution? If an invariant manifold of the difference analog is found which is exponentially stable, does this imply that the differential equation has a periodic or almost periodic solution near this manifold, These are all questions which need to be answered. There are many more questions not listed here which need to be answered.

BIBLIOGRAPHY
[1] Peter Henrici, Discrete Variable Methods in Ordinary Differential Equations, Wiley, New York, 1962. This book has a very good bibliography for the work done before 1962.
[2] Peter Henrici, Error Propagation for Difference Methods, Wiley, New York, 1963.
[3] Royce Beckett and James Hurt, Numerical Calculations and Algorithms; McGraw-Hill, New York, 1967.
[4] Paul R. Halmos, Finite-Dimensional Vector Spaces, D. Van Nostrand, New York, 1958.
[5] . Tomlinson Fort, Finite Differences and Difference Equations in the Real Domain, Clarendon Press, London, 1948.
[6] V.N. Faddeeva, Computational Methods of Linear Algebra (translated by Curtis D. Benster), Dover, New York, 1959.
[7] Riesz, F., and B. Sz. Nagy, Functional Analysis, Frederick Ungar Publ., New York, 1952.
[8] M.A، Krasnosel'skii, Topological Methods in the Theory of Nonlinear Integral Equations (translated by A.H. Armstrong and J. Burlak), Macmillan, New York, 1964.
[9] A. Halanay, "Invariant Manifolds for Discrete Systems", to appear in Proceedings of the Second Czeckoslovakian Conference on Differential Equar tions and their Applications (Equadiff II), Bratislava, Czeckoslavakia, 1966.
[10] A. Halanay, "Periodic and Almost Periodic Solutions of Systems of Finite Difference Equations", Archives of Rational Mechanics and Analysis, Vol. 12(1963), pp. 134-149.
[11] Jürgen Moser, "Perturbation Theory for Almost Periodic Solutions for Undamped Nonlinear Differential Equations", in Nonlinear Differential Equations and Nonlinear Mechanics, edited by J.P. LaSalle and S. Lefschetz, Academic Press, 1963, pp. 71-79.
[12] Jack K. Hale, Oscillations in Nonlinear Systems, McGraw-Hill, New York, 1963.
[13] N. Bogolinbov and Y.A. Mitropolski, "Asymptotic Methods in the Theory of Nonlinear Oscillations", translated by Gordon and Breach, 1962. Goz. Iz. Fiz. Mat. Lit.
[14] Philip J. Davis and New York, 1967.
[15] Philip J. Davis, Interpolation and Approximation, Blaisdell, New York, 1963.
[16] John M. H. Olmsted, Advanced Calculus, Appleton-Century-Crofts, 1956, 1959.
[17] A.S. Besicovitch, Almost Periodic Functions, Dover, New York, 1954.
[18] Wendell H. Fleming, Functions of Several Variables, Addision-Wesley, Reading, Massachusetts, 1966.

