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SYNTHESIS OF CALCULATIONAL METHODS FOR THE  
DESIGN AND ANALYSIS OF RADIATION SHIELDS FOR  
NUCLEAR ROCKET SYSTEMS. VOLUME 9:  
FASTER - A FORTRAN ANALYTIC SOLUTION OF THE  
TRANSPORT EQUATION BY RANDOM SAMPLING

T. M. Jordan

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VOLUME 9

**SYNTHESIS OF CALCULATIONAL METHODS FOR THE  
DESIGN AND ANALYSIS OF RADIATION SHIELDS  
FOR NUCLEAR ROCKET SYSTEMS**

Contract No. NAS-8-20414

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**FINAL PROGRESS REPORT**

**Volume 9**

**SYNTHESIS OF CALCULATIONAL METHODS  
FOR THE DESIGN AND ANALYSIS OF RADIATION  
SHIELDS FOR NUCLEAR ROCKET SYSTEMS**

**FASTER**

**A FORTRAN ANALYTIC SOLUTION OF THE  
TRANSPORT EQUATION BY RANDOM SAMPLING**

**by:**

T. M. Jordan

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## ABSTRACT

This report is Volume 9 of nine volumes of the final report on "Synthesis of Calculational Methods for the Design and Analysis of Radiation Shields for Nuclear Rocket Systems". Presented in this volume is a description of the FASTER program, a Fortran Analytic Solution of the Transport Equation by Random sampling.\*

FASTER is a Fortran IV Monte Carlo program which calculates energy dependent neutron or photon fluxes at points, surfaces and regions of complex geometries. This program contains all the data processing routines required for a wide variety of nuclear vehicle applications. The FASTER program is completely variable dimensioned, and hence, is capable of treating problems of varying complexity within the limitation of total computer storage. The program uses only the input and output tape units and is operational on the 32K WANL and MSFC IBM 7094 computers--using overlay--and on the 64K CDC 6600 computer at the Westinghouse Telecomputer Center in Pittsburgh.

FASTER utilizes the general quadric surface equation for describing the geometry. The more common equations for planes, cones, elliptical cylinders and ellipsoids can also be used for input description of the surfaces.

The FASTER program handles either neutron or photon sources. Each source is described in rectangular, cylindrical or spherical coordinates and the source geometry is superimposed on the problem geometry. The spatial, angular, and energy source distributions are assumed to be separable and are input as tabulated relative distribution data.

The FASTER program deals with the entire spectrum of particle energies simultaneously, thus eliminating costly repetition of geometric calculations which are usually required for treating individual mono-energetic particles. The scattered particle energy spectra include the effect of every possible scattering event at each scattering point. This eliminates the variance associated with the random selection of a single event. Neutron transport calculations utilize averaged multigroup cross sections which are available from several standard tabulations. Photon transport problems utilize point cross sections and the Klein-Nishina equation for Compton scattering.

Biased random sampling is used exclusively in the selection of source and scattering points. In performing flux calculations at a point, the FASTER program includes the singularity of the point flux estimator in the biasing functions. This permits more economical and more accurate point solutions in source and/or scattering volumes.

Analytic estimation is used to compute point fluxes and/or surface and volume averaged fluxes. These fluxes can be processed within the program to obtain various responses such as dose rates and heating rates. Contributions to the fluxes are obtained by source, number of collisions, and scattering region. Legendre angular moments and length-of-flight moments of the fluxes and responses can also be obtained.

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#### SECTION 1.0      INTRODUCTION AND SUMMARY

A major problem in the analysis of nuclear rockets is accuracy in predicting neutron and photon radiation levels both internal and external to the radiation sources. The most useful and fundamental expression of these radiation levels consists of energy dependent fluxes at specified points. These point fluxes can be utilized with a variety of energy dependent functions to yield integral quantities such as dose rate and heating rate. These point quantities can also be spatially integrated to obtain averages or totals for various surfaces and volumes, e.g., neutron and photon reactor leakage and volumetric heating. The accurate calculation of point fluxes requires a detailed treatment of source and material distributions and of basic particle cross sections. Accurate and economical calculations are particularly difficult for space vehicles and nuclear rocket test stands because of their complex geometries. The FASTER program was developed to handle these geometrically complex problems.

Energy dependent fluxes can be obtained, in principle, for any radiation analysis problem by computing the flux contributions by order of scatter. The geometric complexity of most realistic problems, however, limits a conventional numerical integration to the uncollided and, at most, the single scattered flux components. Therefore, the method of random sampling, i.e., the Monte Carlo method is used in the FASTER program. The Monte Carlo method permits a calculation of the flux components to an arbitrary order of scatter while simplifying the numerical procedures.

The FASTER program is an integral part of the "final" design method schematically shown in Fig. 1. This "final" design method is described more fully in Volume 1 of this report. As depicted in the figure, it begins with the POINT program (Volume 2) which prepares cross section and other basic data for use in the ODD-K two-dimensional transport program. The ODD-K two-dimensional transport program (Volume 6) provides neutron and photon energy fluxes throughout the reactor geometry. The NAGS data processing program (Volume 7) processes these fluxes and calculates neutron and photon radiation levels and neutron and photon source distributions within the reactor system. These sources can be employed in either the KAP-V point kernel program (Volume 4) or the FASTER Monte Carlo program for obtaining



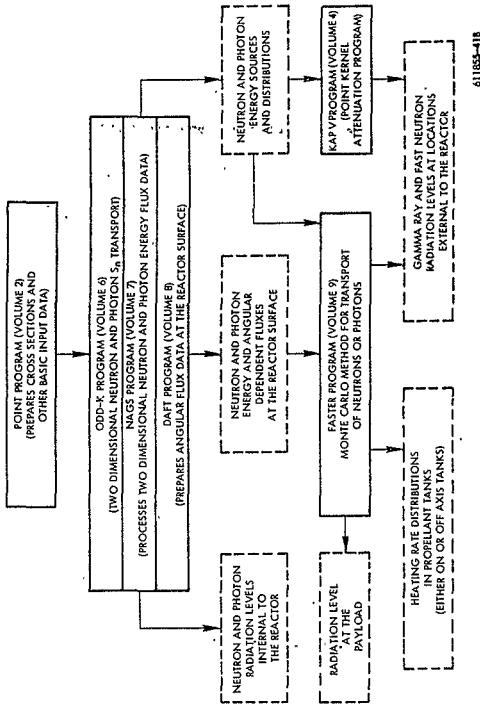


Figure 1. Schematic Diagram of the "Final" Design Method

radiation levels at locations internal and external to the reactor system. In addition, the FASTER program can compute heating rate distributions in the liquid hydrogen propellant (in either an on-axis or off-axis tank) and the radiation level of the payload. Alternately, the DAFT program (Volume 8) can prepare neutron and photon energy and angular dependent fluxes at the reactor surface from the ODD-K program for use in the FASTER Monte Carlo code.

Distinctive features of the Monte Carlo method as employed in the FASTER program, are described in Section 2.0 and include:

- 1) application of random sampling to the spatial and angular integrations only,
- 2) consistent use of energy-averaged sampling functions,
- 3) approximation of importance functions by point kernel techniques,
- 4) analytic treatment of the energy variable over its entire range, and
- 5) zero variance energy integration of the scattered source equations.

The Monte Carlo method presented in this report is based on techniques described in References 1, 2, and 3. Considerable improvement in the treatment of point detectors has been made since these early efforts were reported. In particular, the singularity in the point detector flux estimator has been included in the spatial sampling functions without sacrificing a detailed treatment of geometric effects.

The FASTER program utilizes this improvement of the Monte Carlo method to perform neutron or photon transport calculations in complex geometries. Computer oriented features of this program include:

- 1) coded in FORTRAN IV,
- 2) completely variable dimensioned,
- 3) completely internal—auxiliary tapes are not required, and
- 4) compatible with both the IBM 7094 and CDC 6600 computers.

The program logic is discussed in Section 3.0.

Subsequent sections of this report deal with the numerical techniques used in the FASTER program to implement the Monte Carlo method. First, is the geometric framework which utilizes the quadric surface equation. Basic features of the geometry described in Section 4.0 include:

- 1) separate description of surfaces,
- 2) simple input for the more common planes, cones, elliptical cylinders, and ellipsoids with an internal expansion to obtain the coefficients of the general quadric equation,
- 3) exclusive use of the general surface equation in all computations,
- 4) region description by listing bounding surfaces,
- 5) an internal calculation of the ( $\pm$ ) sign associated with region boundaries by using the coordinates of an arbitrary point in each region,
- 6) geometry consistency check using "point-in-region"
- 7) an internal calculation of "most-probable-next-regions" for boundary crossings,
- 8) an internal calculation of exterior boundaries, and
- 9) elimination of the ( $\epsilon, \delta$ ) boundary crossing search.

The FASTER program will treat multiple fixed sources where each source has separable spatial, angular and energy distributions. Other features of the fixed source description, described in Section 5.0, include:

- 1) rectangular, cylindrical and/or spherical coordinates,
- 2) each spatial variable can be continuous or discrete, permitting a variety of point, line, surface, and volume source geometries,
- 3) each angular variable can be continuous or discrete, running the range from monodirectional to angular,
- 4) spatial and angular distributions are specified by tabulating relative distributions at discrete points, in particular the distributions calculated by the NAGS and DAFT programs,
- 5) spectra are specified by tabulating either differential number or energy spectrum, or groupwise integrated number or energy spectrum, either one in an arbitrary group structure,
- 6) each source is normalized to total energy or particles.

Features of the treatment of cross sections for photons and neutrons are described in Section 6.0 and include:

- 1) compositions are accepted in  $10^{24}$  atoms/cm<sup>3</sup> or gm/cm<sup>3</sup>,
- 2) microscopic cross sections can be in barns/atom or cm<sup>2</sup>/gm,

- 3) photon total cross sections are defined at group boundaries and interpolated linearly in energy,
- 4) photon scattering uses the Klein-Nishina equation,
- 5) photon energy absorption coefficients are computed internally by element and composite material,
- 6) group averaged neutron cross sections are accepted from several standard tabulations, e.g., GAM-1,
- 7) neutron scattering cross sections are not limited in down scatter or order of the Legendre angular expansion,
- 8) neutron cross sections are manipulated internally to define transport corrected values or to remove this correction,
- 9) neutron kinetic heating responses are computed by element and composite material, and
- 10) hydrogen densities can be specified by region, eliminating the need of describing several composite materials which differ by hydrogen content only.

FASTER computes various flux moments in a collapsed set of energy groups for point, surface and/or volume detectors. All quantities are obtained by "analytic estimation", as shown in Section 7.0, and include:

- 1) groupwise number and energy flux, average energy, and variance,
- 2) groupwise differential and cumulative number and energy fluxes,
- 3) groupwise responses and total responses with limits on variance,
- 4) groupwise number flux and response function totals by source,
- 5) groupwise flux and response totals by scattering region,
- 6) groupwise flux and response totals by number-of-collisions,
- 7) Legendre angular moments of the groupwise flux and response totals, and
- 8) length-of-flight moments of the groupwise flux and response totals.

The sampling functions incorporated in FASTER, described in Section 8.0, are relatively easy to use. Typical random sampling input data or biasing data for point flux calculations consist of:

- 1) center and radius of a sphere enclosing all fixed sources,
- 2) group importance, e.g., flux-to-dose response,
- 3) linear buildup coefficient by group,
- 4) relative importance by group of forward-to-backward scattering for heavy elements,
- 5) relative importance by group of forward-to-backward scattering for hydrogen, and
- 6) a set of adjustment factors (ratios  $\approx 1.0$ ) applied to internally calculated sampling parameters.

Detailed input instructions are presented in Section 9.0 of this report. Section 10.0 contains a description of the input and output for a sample problem involving the calculation of photon and neutron fluxes at a point above a liquid hydrogen propellant tank. The FORTRAN IV listing of the FASTER program is given in Appendix C.

Computer times for the FASTER program depend on the problem complexity and the manner in which the program is used. Individual point detector flux calculations in and around a nuclear reactor require about 3 minutes per point for photons (23 groups) and 6 minutes per point for fast neutrons (13 groups). These times are typical for the IBM 7094 computer and yield computed variances in integral responses generally less than 10 percent.

Flux calculations for points in void regions and/or flux calculations for surfaces and volumes require more computer time per problem, but the fluxes are all obtained simultaneously. A typical photon problem involving volumetric heating rates for 50 regions of a reactor required about 1.2 hours on the IBM 7094 with computed variances generally less than 25 percent.

## SECTION

### 2.0 THE MONTE CARLO METHOD

This section describes the Monte Carlo Method as used in the FASTER program. The first topic is a general discussion of order-of-scatter solutions of the transport equation. The order-of-scatter equations are written explicitly starting with the equation for the uncollided scalar flux. The basic concepts of random sampling are then summarized and applied directly to the order-of-scatter equations. This is followed by a discussion of point kernel estimates of optimal sampling functions. Finally, the equations for angular flux estimation are derived for point, surface, and volume detectors.

#### 2.1 GENERAL COMMENTS

Fluxes at a point can be obtained in principle for any geometry by computing the order-of-scatter components, i.e., the uncollided flux, the single scattered flux, etc. This method utilizes a known fixed source distribution, which is numerically integrated with a simple attenuation kernel over the spatial extent of the source, to obtain the uncollided angular flux at all points in space. This uncollided flux is integrated with the scattering cross sections, over energy and solid angle, to yield the single scattered source distribution. Next, the procedure is repeated to an arbitrary order of scattering, finally yielding the total flux within an inherent error given by the uncomputed higher order-of-scattering flux components.

The FASTER program utilizes the Monte Carlo method in performing the numerical integrations. The Monte Carlo method involves an application of random sampling to the evaluation of definite integrals and it can be applied to all of the integrations involved in the order-of-scatter solution. Because this method is statistical in nature its application to each integration must be accompanied by an iteration procedure to reduce the associated variance. Therefore, it is desirable to perform as many of the integrations as possible by conventional techniques. For the order-of-scatter solution utilized in the FASTER program, the Monte Carlo method is applied only to the spatial and angular integrations, i.e., the method is applied to the variables that complicate direct numerical integration.

The application of the Monte Carlo method can proceed in many directions. The procedure used in the FASTER program involves the random selection of a single point characterizing the distributed fixed source. Then, the fixed source is evaluated at this point and the resulting point source is used to compute energy dependent uncollided fluxes throughout the geometry. Next, a first-scatter point is selected by random sampling. An energy-dependent angular flux (monodirectional) at this first scatter point is computed using the point representation of the fixed source. This energy dependent angular flux, when integrated with the scattering cross sections, yields a point representation of the single-scattered source for all scattered energies and directions. This single scattered point source is used to obtain the energy dependent single-scattered fluxes throughout the geometry. It is used to calculate the monodirectional flux at the next scattering point (obtained by random sampling). This procedure is continued to an arbitrary order-of-scatter and is repeated (starting with the fixed source) until the statistical error in computed fluxes is acceptable.

An implicit difficulty with this point-to-point procedure is the singularity in the attenuation kernel resulting from the inverse square law. This difficulty is present in the calculation of fluxes at both the next scatter point and at arbitrary detector points. It shall be seen later, that these difficulties are removed by considering the singularities in the random selection of the discrete position vectors.

An important consideration in every numerical integration is the selection of a procedure which will yield minimum error. Since, the Monte Carlo method is just another way of performing the spatial integrations, the same considerations are present. It is possible to devise many techniques for reducing the error in an integration performed by random sampling. The most fruitful technique involves the concept of spatial importance. Thus, in considering each possible source or scattering point, the following question is asked: How important will the selection of this point be to the final answer, i.e., how much flux, dose, etc. will eventually be derived from its selection?

Without solving the problem, there is one method which is readily available for approximating spatial importance. This is the point kernel method which approximates the

importance of future scattering events by simple alterations of the attenuation kernel, i.e., removal theory for neutrons and buildup factors for photons. This procedure will not yield a zero variance (or error) flux calculation, but it will yield information as to the approximate importance of space points. This importance function includes the singularities already mentioned and also approximates the importance of material and fixed source distributions.

## 2.2 THE ORDER-OF-SCATTER EQUATIONS

The solution of the time independent transport equation can be obtained by computing the order-of-scatter components of the flux. This method utilizes a known, differential source density where the uncollided flux is given by the familiar equation:

$$\phi_0(\vec{r}, E) = \iiint S_0(\vec{r}', \vec{\Omega}, E) \frac{\exp\left[-\int_0^s \Sigma^t(\vec{r} - s' \vec{\Omega}, E) ds'\right]}{s^2} dV \quad (2.1)$$

where

$$s = |\vec{r} - \vec{r}'|, \quad \vec{\Omega} = (\vec{r} - \vec{r}') / s$$

$\vec{r}, \vec{r}'$  are position vectors

$\vec{\Omega}$  is a unit direction vector

E is the particle energy (MeV)

s,  $s'$  are scalar distances from  $\vec{r}'$  along  $-\vec{\Omega}$  (cm)

dV is a differential volume element ( $\text{cm}^3$ )

$S_0(\vec{r}', \vec{\Omega}, E)$  is the fixed differential source density

$$\left( \frac{\text{particles}}{\text{cm}^3 \text{ steradian Mev} \cdot \text{sec}} \right)$$

$\phi_0(\vec{r}, E)$  is the scalar flux  $\left( \frac{\text{particles}}{\text{cm}^2 \cdot \text{Mev sec}} \right)$

$\Sigma^t(\vec{r}, E)$  is the total cross section ( $\text{cm}^{-1}$ )

This integration can be performed using spherical coordinates in a coordinate system centered at the detector point by using the transformation

$$\vec{r}' = \vec{r} - s\vec{\Omega}$$

$$dV = s^2 ds d\Omega$$

$$d\Omega = d\mu d\theta$$

$\mu$  = the cosine of the polar angle

$\phi$  = the azimuthal angle

$$\phi_0(\vec{r}, E) = \frac{4\pi}{4}\int_0^\infty S_0(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp\left[-\int_0^s \Sigma^t(\vec{r} - s'\vec{\Omega}, E) ds'\right] ds d\Omega \quad (2.2)$$

The most obvious reason for this transformation is that it removes the  $1/s^2$  singularity.

The angular flux at any point in the geometry can be obtained from the inner spatial integration of equation 2.2:

$$\phi_0(\vec{r}, \vec{\Omega}, E) = \int_0^\infty S_0(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp\left[-\int_0^s \Sigma^t(\vec{r} - s'\vec{\Omega}, E) ds'\right] ds \quad (2.3)$$

The single scattered source density is given by an integration of the product of this angular flux and the total differential scattering cross

$$S_1(\vec{r}, \vec{\Omega}, E) = \frac{4\pi}{4}\int_0^\infty \phi_0(\vec{r}, \vec{\Omega}', E') \frac{d^2\Sigma}{d\Omega dE}(\vec{r}, \vec{\Omega}', E' \rightarrow \vec{\Omega}, E) dE' d\Omega' \quad (2.4)$$

where  $\frac{d^2\Sigma}{d\Omega dE}(\vec{r}; \vec{\Omega}', E' \rightarrow \vec{\Omega}, E)$

is the total differential scattering cross section (Reference 4, pg. 265). The single scattered source density yields the single scattered flux and the process is then repeated for higher order scattered sources and fluxes:

$$\phi_k(\vec{r}, \vec{\Omega}, E) = \int_0^\infty S_k(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp\left[-\int_0^s \Sigma^t(\vec{r} - s'\vec{\Omega}, E) ds'\right] ds \quad (2.5)$$

$$S_{k+1}(\vec{r}, \vec{\Omega}, E) = \frac{4\pi}{4}\int_0^\infty \phi_k(\vec{r}, \vec{\Omega}', E') \frac{d^2\Sigma}{d\Omega dE}(\vec{r}, \vec{\Omega}', E' \rightarrow \vec{\Omega}, E) dE' d\Omega' \quad (2.6)$$

where  $k$  is the number of scattering events which the particles have experienced.

The FASTER program uses random sampling techniques to perform the integrations of the order-of-scatter equations, 2.5 and 2.6, thereby giving the uncollided and scattered components of the flux. The program uses a groupwise representation of the energy dependence in conjunction with other numerical techniques—described in Sections 3.0 through 8.0—in performing the integrations.

The remainder of this section is a general discussion of the application of random sampling to the order-of-scatter equations. In particular, consideration will be given to the development of sampling techniques which should minimize the error in the integration.

In devising optimal solutions of these order-of-scatter equations, it is noted that after computing the  $(k-1)$ th and lower order flux components, the unsolved portion of the problem, corresponding to the  $k$ th and higher order flux components, is given by a summation over the "future" components of the flux. This unsolved portion of the problem is denoted by

$\phi_{\geq k}(\vec{r}, \vec{\Omega}, E)$ —the angular flux from particles having  $k$  or more collisions—and is given by a summation of equation 2.5:

$$\begin{aligned} \phi_{\geq k}(\vec{r}, \vec{\Omega}, E) &= \sum_{k'=k}^{\infty} \phi_{k'}(\vec{r}, \vec{\Omega}, E) \\ &= \int_0^\infty S_{\geq k}(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp\left[\int_0^s \Sigma^t(\vec{r} - s'\vec{\Omega}, E) ds'\right] ds \end{aligned} \quad (2.7a)$$

where  $S_{\geq k}(\vec{r}, \vec{\Omega}, E)$  is the differential scattered source density of particles having  $k$  or more collisions.

$$S_{\geq k}(\vec{r}, \vec{\Omega}, E) = \sum_{k'=k}^{\infty} S_{k'}(\vec{r}, \vec{\Omega}, E)$$

For example, in initiating the order-of-scatter solution, the unsolved portion of the problem is:

$$\phi_{\geq 0}(\vec{r}, \vec{\Omega}, E) = \int_0^s S_{\geq 0}(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^t(\vec{r} - s\vec{\Omega}, E) ds' \right] ds \quad (2.7b)$$

where  $\phi_{\geq 0}(\vec{r}, \vec{\Omega}, E)$  is the total angular flux,  $\phi(\vec{r}, \vec{\Omega}, E)$ , from all orders-of-scatter and  $S_{\geq 0}(\vec{r}, \vec{\Omega}, E)$  is the total differential source density,  $S(\vec{r}, \vec{\Omega}, E)$ , including the scattered source.

In computing the  $k$ th order-of-scatter flux component, all higher order components should also be considered. If attention is given to a minimum error calculation of  $\phi_k(\vec{r}, \vec{\Omega}, E)$ , this will not be the calculation which minimizes the error in  $\phi_{\geq k}(\vec{r}, \vec{\Omega}, E)$ . However, it may be adequate since the two calculations are quite similar, i.e., it is theoretically possible to define kernels such that:

$$\begin{aligned} \phi_{\geq k}(\vec{r}, \vec{\Omega}, E) &= \int_0^\infty \left[ \iiint_0^\infty S_k(\vec{r} - s\vec{\Omega}, \vec{\Omega}', E') K(\vec{r} - s\vec{\Omega}, \vec{\Omega}', E' \rightarrow \vec{r}, \vec{\Omega}, E) dE' d\Omega' \right] \\ &\times \exp \left[ - \int_0^s \Sigma^t(\vec{r} - s\vec{\Omega}, E) ds' \right] ds \end{aligned} \quad (2.8)$$

where the kernel  $K(\vec{r} - s\vec{\Omega}, \vec{\Omega}', E' \rightarrow \vec{r}, \vec{\Omega}, E)$  usually varies more slowly than the material attenuation kernel.

This "resolvent" kernel will yield the solution with an integration over the fixed source. However, this kernel is as difficult to obtain as the order-of-scatter solution is and it involves a similar iterative process. Reference 5, pg. 522. The most well known examples of this kernel are the dose kernels derived from moments method calculations for point isotropic sources in infinite media.<sup>(6)</sup>

Use of the "resolvent" kernel is not practical for geometrically complex problems. However, the "success" of the approximate dose kernels in complicated geometries leads naturally to their use in estimating the importance of future scattering events. The use of approximate kernels is discussed in more detail in Section 2.5.

### 2.3 RANDOM SAMPLING CONCEPTS

The concepts involved in applying random sampling to the integration of equations 2.5 and 2.6 are simplified notationally by considering the evaluation of a simple definite integral with a non-negative integrand:

$$I = \int_{R(x)} f(x) dx = \int_{R(x)} \frac{f(x)}{p^*(x)} p^*(x) dx = \int_{R(x)} f^*(x) p^*(x) dx \quad (2.9)$$

where  $f^*(x) = f(x)/p^*(x)$

$R(x)$  is the range of  $x$

$p^*(x) \geq 0$

$p^*(x) > 0$  if  $f(x) > 0$

$$\int_{R(x)} p^*(x) dx = 1$$

$R(x)$

The conditions imposed on  $p^*(x)$  permit its use as a sampling function for obtaining  $x_i$  values  $x_i$  of  $x$ . This function is properly called a probability density function for the random variable  $x$ .

The mean or expected value of  $f^*(x)$  is simply:

$$E[f^*(x)] = \int_{R(x)} [f^*(x)] p^*(x) dx = I \quad (2.11)$$

The expected value of the mean square variation of  $f^*(x)$  from its expected value (the variance of  $f^*(x)$ ) is defined by:

$$\begin{aligned}\sigma^2 [f^*(x)] &= E [(f^*(x) - I)^2] \\ &= \int_{R(x)} [(f^*(x) - I)^2] p^*(x) dx = \int_{R(x)} [f^*(x)]^2 p^*(x) dx - I^2 \\ &= \frac{\int_{R(x)} f^2(x) dx}{R(x)} - I^2\end{aligned}\quad (2.12)$$

Both the mean and variance of  $f^*(x)$  involve analytic integrations. Of greater interest in numerical integrations are the sample mean:

$$I_n = \frac{1}{n} \sum_{i=1}^n f^*(x_i), \quad x_i \text{ is randomly selected from } p^*(x)\quad (2.13)$$

and the corresponding sample variance:

$$V_n^2 = \frac{1}{n-1} \sum_{i=1}^n [f^*(x_i) - I_n]^2 = \frac{1}{n-1} \left[ \sum_{i=1}^n [f^*(x_i)]^2 - n I_n^2 \right]\quad (2.14)$$

Simple manipulations (Reference 7, pg. 198) yield:

$$\begin{aligned}E[I_n] &= 1 \\ \sigma^2 [I_n] &= \frac{1}{n} \sigma^2 [f^*(x)] \\ E[V_n^2] &= \sigma^2 [f^*(x)].\end{aligned}$$

The last two equations imply that:

$$\sigma^2 [I_n] \approx \frac{1}{n} V_n^2\quad (2.15)$$

Finally, it is noted that since  $f(x) \geq 0$  for all  $x$  in  $R(x)$ , then the optimum (zero variance) integration is performed by sampling from:

$$p^*(x) = \frac{f(x)}{\int f(x) dx}\quad (2.16)$$

since any discrete point will give the correct answer. Subsequent sections deal with approximate relationships which utilize this obvious result.

It should be noted that a transformation to the integration variable  $u$ , where  $du = I^{-1} f(x) dx$ , also yields a zero error result in a conventional numerical integration. In fact, the Monte Carlo and conventional integrations would be the same except for the technique used to obtain discrete points  $x_i$  (random versus systematic).

## 2.4 INNER ITERATIONS

It was possible in developing the order-of-scatter equations to explicitly write an equation for the  $k^{th}$  order scalar flux component which involved a  $k$ -fold volume integration (and a  $k$ -fold energy integration), i.e., spatial integrations over the fixed source volume, the single scattered source volume, . . . , the  $k^{th}$  scattered source volume. The manner in which the order-of-scatter equations was developed obviated the need for explicitly displaying these volume integrations. However, in relating the discussion of Section 2.3 to the techniques used in solving the order-of-scatter equations, this  $k$ -fold volume integration must be recognized. In particular, the discrete random variable  $x_i$  used in the preceding section is equivalent to a series of discrete position vectors  $\vec{r}_i, 0^*, \vec{r}_i, 1^*, \vec{r}_i, 2^*, \dots, \vec{r}_i, k^*$ , obtained by random sampling.

Fortunately, the techniques used in the Monte Carlo integration of the order-of-scatter equations, i.e., the techniques used in selecting the discrete position vectors  $\vec{r}_{i,k}, k = 0, 1, \dots$ , can be discussed in an orderly fashion without displaying the  $k$ -fold volume integrations. The discussion of the integrations is given below with the "outer" iteration index  $i$  of Section 2.3 suppressed. The "inner" iteration index  $k$ , corresponding to the  $k^{th}$  order-of-scatter, will be retained.

The "order-of-scattering" inner iterants are obtained in a straight forward manner. A natural starting point is the equation for the  $k^{th}$  component of the scalar flux:

$$\begin{aligned}\phi_k^*(\vec{r}, E) &= \frac{1}{4\pi} \iint \phi_k(\vec{r}, \vec{\Omega}, E) d\Omega \\ &= \frac{1}{4\pi} \iiint S_k^{(0)}(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^t(\vec{r} - s^t \vec{\Omega}, E) ds^t \right] ds d\Omega \\ \phi_k(\vec{r}, E) &= \iiint S_k(\vec{r}, \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^t(\vec{r} - s^t \vec{\Omega}, E) ds^t \right] dV \quad (2.17)\end{aligned}$$

where  $dV$  is the general differential volume element, equivalent to  $s^2 ds d\Omega$  in a spherical coordinate system centered at  $\vec{r}$ . This equation has been transformed into a general volume integration to display the singularity ( $1/s^2$ ) associated with the flux calculation.

This equation has a definite value and is manipulated in the same manner as equation 2.9, i.e., the integrand is multiplied and divided by an arbitrary sampling function:

$$\phi_k(\vec{r}, E) = \iiint \left\{ \frac{S_k(\vec{r}', \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^t(\vec{r} - s^t \vec{\Omega}, E) ds^t \right]}{s^2 P_k^*(\vec{r}')} \right\} P_k^*(\vec{r}') dV \quad (2.18)$$

with the restrictions:

$$\left. \begin{aligned}P_k^*(\vec{r}') &\geq 0 \\ P_k^*(\vec{r}') &> 0 \text{ if } \iiint_0^{(0)} S_k(\vec{r}', \vec{\Omega}, E) dE d\Omega > 0 \\ \iiint P_k^*(\vec{r}') dV &= 1\end{aligned} \right\} \quad (2.19)$$

This equation can be evaluated in a manner analogous to that used in equation 2.13 for the sample mean, i.e.: select  $\vec{r}'_k$  at random from  $P_k^*(\vec{r}')$

Then the contribution to the  $k$ th component of the scalar flux is given by:

$$\Delta\phi_k^*(\vec{r}, E) = \frac{S_k(\vec{r}'_k, \vec{\Omega}, E)}{P_k^*(\vec{r}'_k)} \frac{\exp \left[ - \int_0^s \Sigma^t(\vec{r} - s^t \vec{\Omega}, E) ds^t \right]}{s^2} \quad (2.20)$$

where  $s = |\vec{r} - \vec{r}'_k|$  and  $\vec{\Omega} = (\vec{r} - \vec{r}'_k)/s$ .

Note that the energy varies over its entire range, i.e., random sampling has been limited to the spatial variables. This is equivalent to defining an energy dependent angular point source which represents the  $k$ th scattered differential source density,

$$W_k^*(\vec{\Omega}, E) = \frac{S_k(\vec{r}'_k, \vec{\Omega}, E)}{P_k^*(\vec{r}'_k)} \quad (2.21)$$

since the remaining factor in equation 2.20 is associated with the flux calculations and depends on the variable position vector  $\vec{r}$ . This point source is defined through the  $k$ th source distribution and can be evaluated for any desired direction  $\vec{\Omega}$ .

This procedure is also equivalent to representing the  $k$ th component of the differential source density by:

$$S_k^*(\vec{r}, \vec{\Omega}, E) = W_k^*(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}'_k) \quad (2.22)$$

where  $\delta(\vec{r} - \vec{r}'_k)$  is the Dirac delta function. This representation is particularly useful in later discussions where formal integrations over the  $k$ th scattered source volume are required.

In particular, the point representation of the fixed source ( $k = 0$ ) is:

$$W_0^*(\vec{\Omega}, E) = \frac{S_0(\vec{r}_0, \vec{\Omega}, E)}{P_0^*(\vec{r}_0)} \quad (2.23)$$

The expected value of  $S_o^*(\vec{r}, \vec{n}, E)$  is just the fixed source distribution:

$$\begin{aligned} E \left[ W_o^*(\vec{n}, -\delta(\vec{r} - \vec{r}_o)) \right] &= \iiint_{V_o} \left[ W_o^*(\vec{n}, E) \delta(\vec{r} - \vec{r}_o) \right] P_o^*(\vec{r}_o) d\vec{r}_o \\ &= \iiint_{V_o} S_o(\vec{r}_o, \vec{n}, E) \delta(\vec{r} - \vec{r}_o) P_o^*(\vec{r}_o) d\vec{r}_o \\ &= S_o(\vec{r}, \vec{n}, E) \end{aligned} \quad (2.24)$$

The definition of the energy dependent, angular, point sources for higher order scattered components is more involved. Assuming the inner iterations have progressed through the  $(k-1)$ th spatial integration, then the  $(k-1)$ th component of the source density is represented by the vector  $\vec{r}_{k-1}$  and the angular point source  $W_{k-1}(\vec{n}, E)$ . Then the  $k$ th scattered energy angular point source at  $\vec{r}_k$  is given by:

$$\begin{aligned} W_k^*(\vec{n}, E) &= \frac{1}{4\pi} \iiint_0^\infty \frac{\phi_{k-1}(\vec{r}_k, \vec{n}, E)}{P_k^*(\vec{r}_k)} \frac{d^2 \Sigma}{d\Omega dE} (\vec{r}_k, \vec{n}, E \rightarrow \vec{n}, E) dE' d\Omega' \\ &= \frac{1}{4\pi} \iiint_0^\infty \left[ \int_0^\infty S_{k-1}(\vec{r}_k - \vec{n}, \vec{n}, E') \exp \left[ \int_0^s \Sigma^i(\vec{r}_k - s, \vec{n}, E') ds \right] ds \right] \frac{d^2 \Sigma}{d\Omega dE} (\vec{r}_k, \vec{n}, E \rightarrow \vec{n}, E) dE' d\Omega' \\ &= \iiint \left[ \int_0^\infty W_{k-1}^*(\vec{n}', E') \exp \left[ - \int_0^s \Sigma^i(\vec{r}_k - s, \vec{n}', E') ds \right] \right] \frac{d^2 \Sigma}{d\Omega dE} (\vec{r}_k, \vec{n}, E \rightarrow \vec{n}, E) dE' d\Omega' \delta(\vec{r} - \vec{r}_{k-1}) dV \\ &= \int_0^\infty \frac{W_{k-1}^*(\vec{n}', E') \exp \left[ - \int_0^{s_k} \Sigma^i(\vec{r}_k - s, \vec{n}', E') ds \right]}{s_k^2 P_k^*(\vec{r}_k)} \cdot \frac{d^2 \Sigma}{d\Omega dE} (\vec{r}_k, \vec{n}, E \rightarrow \vec{n}, E) dE' \end{aligned} \quad (2.25)$$

where  $s_k = |\vec{r}_k - \vec{r}_{k-1}|$  and  $\vec{n}_k = (\vec{r}_k - \vec{r}_{k-1}) / s_k$

Note that the  $(k-1)$ th source density component was replaced by its equivalent point representation  $W_{k-1}^*(\vec{n}', E)$ . Thus, this equation is an equality in the sense that the expected value of the right hand side is the left hand side.

It now becomes expedient to define an energy dependent point monodirectional quantity at  $\vec{r}_k$  which characterizes the  $(k-1)$ th component of the flux:

$$W_{k-1}^*(E) = W_{k-1}^*(\vec{n}_k, E) \exp \left[ - \int_0^{s_k} \Sigma^i(\vec{r}_k - s, \vec{n}_k, E) ds \right] \frac{s_k^2}{s_k^2 P_k^*(\vec{r}_k)} \quad (2.26)$$

The factor associated with random selection of the discrete position vector  $\vec{r}_k$  has been included in this energy-dependent, point monodirectional representation of the flux to simplify the final equation for the  $k$ th scattered, angular point source:

$$W_k^*(\vec{n}, E) = \int_0^\infty W_{k-1}^*(E') \frac{d^2 \Sigma}{d\Omega dE'} (\vec{r}_k, \vec{n}_k, E' \rightarrow \vec{n}, E) dE' \quad (2.27)$$

The order of calculations may clarify this procedure:

- select  $\vec{r}_k$  from  $P_k^*(\vec{r}_k)$
- calculate  $s_k = |\vec{r}_k - \vec{r}_{k-1}|$
- calculate  $\vec{n}_k = \vec{r}_k - \vec{r}_{k-1} / s_k$
- calculate  $W_{k-1}^*(\vec{n}_k, E)$  for the discrete direction  $\vec{n}_k$   
 $k = 1, W_o^*(\vec{n}_k, E)$ : from equation 2.23 for  $\vec{n} = \vec{n}_k$   
 $k > 1, W_{k-1}^*(\vec{n}_k, E)$  from equation 2.27 for  $\vec{n} = \vec{n}_k$  using  $W_{k-2}^*(E)$ .
- calculate  $W_{k-1}^*(\vec{n}_k, E)$  using equation 2.26.
- then  $W_k^*(\vec{n}, E)$  can be calculated, as required for a given direction  $\vec{n}$ , by using equation 2.27.

Thus,  $W_k^s(\vec{r}, \vec{\Omega}, E)$  is not computed directly except as required for discrete directions. It is always available through the fixed source distribution for  $k = 1$ , or, for  $k > 1$ , through  $W_{k-1}^\phi(E)$  and the differential scattering cross sections. This is indicated in step d) above.

For  $k = 1$ , the verification of the representation of the source given by equation 2.27 is obtained by calculating expected values:

$$\begin{aligned}
 E \left[ W_1^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_1) \right] &= \iiint_0 V_1 \iiint \left[ W_1^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_1) \right] P_1^*(\vec{r}_1) P_0^*(\vec{r}_0) d\vec{r}_1 d\vec{r}_0 \\
 &= \iiint_0 \iiint_1 \left\{ \int_0^\infty \frac{S_0(\vec{r}_0, \vec{\Omega}, E') \exp \left[ - \int_0^{r_1} \sum s' (\vec{r}_1 - s' \vec{r}_1) E' ds' \right]}{s^2 P_0^*(\vec{r}_0) P_1^*(\vec{r}_1)} \right. \\
 &\quad \times \left. \frac{d^2 \sum}{d\vec{r} dE'} (\vec{r}_1, \vec{\Omega}', E' - \vec{\Omega}, E) dE' \right\} \delta(\vec{r} - \vec{r}_1) P_1^*(\vec{r}_1) P_0^*(\vec{r}_0) d\vec{r}_1 d\vec{r}_0 \\
 &= \iiint_0 \left\{ \int_0^\infty S_0(\vec{r}_0, \vec{\Omega}', E') \frac{\exp \left[ - \int_0^s (\vec{r}_1 - s \vec{\Omega}', E') ds' \right]}{s^2} \frac{d^2 \sum}{d\vec{r} dE'} (\vec{r}, \vec{\Omega}', E' - \vec{\Omega}, E) dE' \right\} d\vec{r}_0 \\
 &= S_1(\vec{r}, \vec{\Omega}, E) \tag{2.28}
 \end{aligned}$$

Verification for higher order components can be obtained by induction, i.e., by assuming:

$$\begin{aligned}
 E \left[ W_k^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_k) \right] &= \iiint_0 \dots \iiint_k W_k^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_k) P_k^*(\vec{r}_k) \cdots P_0^*(\vec{r}_0) d\vec{r}_k \cdots d\vec{r}_0 \\
 &= S_k(\vec{r}, \vec{\Omega}, E)
 \end{aligned}$$

and then showing that

$$E \left[ W_{k+1}^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_{k+1}) \right] = S_{k+1}(\vec{r}, \vec{\Omega}, E)$$

Similar verifications can also be obtained for arbitrary point kernels, e.g., those used for flux calculations (Section 2.6). The resulting integrations are similar to those above and are not shown.

The final set of procedures used in a single outer iteration is shown in Figure 2. The inner iterations cannot proceed indefinitely as is indicated by the order-of-scatter solution. Therefore, the inner iterations are terminated using various criteria such as total number of iterations (collision cutoff), all energies below some minimum (energy cutoff), or all flux contributions being negligible (weight cutoff).

## 2.5 SAMPLING FUNCTIONS

At this point, the definitions of the sampling or probability density functions are arbitrary except for the conditions imposed in Equation 2.19. The criterion selected for defining optimal functions is that the variance in the contributions from all future inner iterations be minimum.

After ( $k-1$ ) inner iterations, the scalar flux will have been estimated to the extent of:

$$\phi_{\leq k}(\vec{r}, E) = \sum_{k'=0}^{k-1} \phi_{k'}(\vec{r}, E) \tag{2.29}$$

The remainder associated with neglecting higher order components of the flux or, alternatively, the importance of these components is given by:

$$\phi_{\geq k}(\vec{r}, E) = \sum_{k'=k}^{\infty} \phi_{k'}(\vec{r}, E) \tag{2.30}$$

Since a detailed treatment of these "future" components is as difficult as solving the original problem--equation 2.7a would be solved--various approximations must be made. It is noted

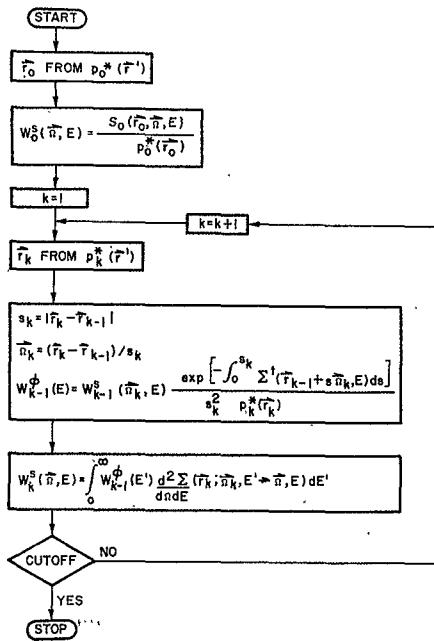


Figure 2. The Monte Carlo Method - Inner Iterations

that a single point criterion--independent of energy--is more easily applied, and much easier to estimate; e.g.,

$$D_{\leq k}(\vec{r}) = \int_0^\infty f(E) + k(\vec{r}, E) dE \quad (2.31)$$

where  $f(E)$  is a response function of particular importance to the problem being solved, e.g., the flux to dose conversion factor if dose rates are being calculated. The solution and remainder corresponding to this response function are then given by:

$$D_{\leq k}(\vec{r}) = \sum_{k'=0}^{k-1} D_{k'}(\vec{r}) \quad (\text{Total response already calculated}) \quad (2.32)$$

$$D_{\geq k}(\vec{r}) = \sum_{k'=k}^{\infty} D_{k'}(\vec{r}) \quad (\text{uncalculated contributions to total response}) \quad (2.33)$$

The calculation of the latter quantity is also as difficult as solving the original problem. However, for estimating purposes, the "point-kernel method" is available for approximating the remainder

$$\begin{aligned} D_{\geq k}(\vec{r}) &= \sum_{k'=k}^{\infty} \int_0^\infty f(E) \phi_{k'}(\vec{r}, E) dE \\ &= \sum_{k'=k}^{\infty} \int_0^\infty f(E) \iiint S_{k'}(\vec{r}', \vec{r}, E) \exp\left[-\int_0^s (\vec{r}' - \vec{s}') \cdot \vec{v}_s(E) ds'\right] dVdE \\ &= \iiint \left\{ \int_0^\infty \left[ \sum_{k'=k}^{\infty} S_{k'}(\vec{r}', \vec{r}, E) \right] \exp\left[-\int_0^s (\vec{r}' - \vec{s}') \cdot \vec{v}_s(E) ds'\right] f(E) dE \right\} dV \end{aligned} \quad (2.34)$$

$$\approx \tilde{D}_{\geq k}(\vec{r})$$

where

$$\tilde{D}_{\geq k}(\vec{r}) = \iiint \left[ \int_0^\infty \iint S_k(\vec{r}', \vec{\Omega}' E) K(\vec{\Omega}', E, \eta(E)) \exp \left[ -\eta(E) \right] f(E) dE d\Omega' \frac{dV}{s^2} \right] (2.35)$$

$$s = |\vec{r} - \vec{r}'|, \vec{\Omega} = (\vec{r} - \vec{r}')/s$$

$$\eta(E) = \int_0^\infty (\vec{r} - \vec{s}', \vec{\Omega}', E) ds'$$

$K(\vec{\Omega}', E, \eta(E))$  is some approximate representation of the response contributions by multiple scattering events, usually a simple function of the number of the mean-free-paths  $\eta(E)$  between a source (or scattering) point and the detector point.

Equation 2.35 not only estimates the importance of future scattering events but can also be used to estimate the importance of the "source", or scattering points. Since it is a feasible spatial importance estimator -- it is related in an approximate manner to the zero variance importance function -- and since it lends itself to a variety of further approximations, it plays an important role in the development of optimal sampling schemes.

Equation 2.35 can be manipulated in a manner similar to equation 2.18:

$$\tilde{D}_{\geq k}(\vec{r}) = \iint \left[ \int_0^\infty \iint S_k(\vec{r}', \vec{\Omega}', E) K(\vec{\Omega}', E, \eta(E)) \exp \left[ -\eta(E) \right] f(E) d\Omega' dE \right] \cdot p_k^*(\vec{r}') dV$$

Equation 2.36 then implies that the optimum sampling function, yielding zero variance for  $\hat{D}_{\geq k}(\vec{r})$ , is given by:

$$p_k^*(\vec{r}') = \frac{1}{\tilde{D}_{\geq k}(\vec{r}')} \int_0^\infty \iint \frac{S_k(\vec{r}', \vec{\Omega}' E) K(\vec{\Omega}', E, \eta(E)) \exp \left[ -\eta(E) \right] f(E) d\Omega' dE}{s^2} (2.36a)$$

This equation is still quite involved and its numerical implementation requires further simplification and approximation.

The form of this equation may be clarified by examining its equivalent for a one velocity problem. Assuming a uniform infinite medium and isotropic sources, either fixed or scattered, this equation can be approximated by:

$$p_k^*(\vec{r}') = \frac{S_k(\vec{r}') \exp \left[ -\Sigma s \right] / s^2}{\iiint S_k(\vec{r}') \frac{\exp \left[ -\Sigma s \right]}{s^2} dV} (2.36b)$$

where  $\Sigma$  is an effective cross section giving the attenuation characteristics of the total flux from the  $k$ th scattered source and  $S_k(\vec{r}')$  is the differential  $k$ th scattered source density.

For point detectors, random sampling of the fixed source distribution,  $k = 0$ , can utilize equation 2.36 directly. The more important aspects of this equation are:

- it includes the spatial divergence ( $1/s^2$ ) from the detector;
- it retains the exponential falloff of source point importance due to material attenuation, and
- it includes the fixed source distribution.

The techniques used in the FASTER program for sampling the fixed source are closely associated with equation 2.36b. An average source energy is used to define the necessary sampling parameters. The program includes, however, an explicit representation of the material distribution. The details of the sampling procedures are discussed in Section 8.3.

Applications of equation 2.36 to higher order flux components,  $k > 0$ , involves the scattered source definition using the  $(k-1)$ th point source representation:

$$p_k^*(\vec{r}') = \frac{1}{\tilde{D}_{\geq k}(\vec{r}')} \int_0^\infty \iint \int_0^\infty \sum_{k=1}^t \frac{S_{k-1}(\vec{r}', \vec{\Omega}', E') d^2 \sum_{k=1}^t (\vec{r}', \vec{\Omega}', E') \rightarrow \vec{\Omega}, E)}{d\Omega' dE'} (2.37a)$$

$$K(\vec{\Omega}', E, \eta(E)) \exp \left[ -\int_0^\infty \sum_{k=1}^t (\vec{r}' - t', \vec{\Omega}', E) dt' - \eta(E) \right] f(E) dE' d\Omega' dE$$

$$\text{where } t = |\vec{r}' - \vec{r}_{k-1}|, \vec{\Omega}' = (\vec{r}' - \vec{r}_{k-1})/s$$

Again, a large degree of approximation is required for numerical calculations.

A one-velocity approximation may also clarify this equation. Using the assumptions used in equation 2.36, the sampling function can be approximated by:

$$P_k^*(\vec{r}) = \frac{\exp \left[ -\Sigma^t t - \Sigma^r s \right]}{\int \int \int \frac{\exp \left[ -\Sigma^t t - \Sigma^r s \right]}{t^2 s^2} dV} \quad (2.37b)$$

where  $\Sigma^t$  is the total cross section.

For a point detector the following characteristics are noted for this scattering point sampling function:

- It includes the spatial divergence from both the source ( $1/t^2$ ) and the detector ( $1/s^2$ )
- It includes the exponential attenuation along both "legs" of the scattering triangle
- It includes the spatial and angular dependence of the scattering cross section, and
- It includes the angular dependence of the  $(k-1)th$  source component.

The FASTER program uses a one-velocity approximation of this sampling function with group averaged parameters being obtained at each scattering point. The material distribution and scattering angle effects are included by several alternate approximations. The details are discussed in Sections 8.4, 8.5, and 8.6.

Similar equations can be obtained for sampling functions which minimize the variance of volume and surface averaged flux calculations. However, these equations involve an integration over the spatial extent of these volumes and surfaces and the development is more complicated. These sampling functions have been approximated, therefore, by "solid angle" considerations as described in Section 8.6.

## 2.6 POINT ANGULAR FLUXES

The set of inner iterations yields a single estimate of the total differential source density:

$$S_i^*(\vec{r}, \vec{\Omega}, E) = \sum_{k=0} W_{ik}^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_{ik}) \quad (2.38)$$

where the index  $i$  corresponds to the  $i$ th repetition of this process, or a single outer iteration. Previous discussions verified that:

$$E[S_i^*(\vec{r}, \vec{\Omega}, E)] = S(\vec{r}, \vec{\Omega}, E) - R_i \quad (2.39)$$

where  $R_i$  is a remainder corresponding to the termination of the inner iterations after a finite value of  $k$ , i.e., the neglected higher order-of-scatter components.

Repetitive application of these techniques ( $n$  outer iterations) yields the final estimate of the total differential source density:

$$S(\vec{r}, \vec{\Omega}, E) = \frac{1}{n} \sum_{i=1}^n \left[ \sum_{k=0}^n W_{ik}^s(\vec{\Omega}, E) \delta(\vec{r} - \vec{r}_{ik}) \right] \quad (2.40)$$

This is not the end result, of course. The original intent was to obtain the flux at an arbitrary point. Equation 2.7b implies the total angular flux is obtained by the integration,

$$\phi(\vec{r}, \vec{\Omega}, E) = \int_0^\infty S(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \sum_t^s (\vec{r} - s'\vec{\Omega}, E) ds' \right] ds \quad (2.41)$$

This can be approximated using the above source density:

$$\phi^*(\vec{r}, \vec{\Omega}, E) = \int_0^\infty S^*(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \sum_t^s (\vec{r} - s'\vec{\Omega}, E) ds' \right] ds \quad (2.42)$$

It is a straightforward task to show that the expected value of this approximation is, indeed, the total angular flux within an error  $\bar{F}_U$  corresponding to the source errors  $R_i$  above:

$$E \left[ \phi^* (\vec{r}, \hat{\Omega}, E) \right] = \phi (\vec{r}, \hat{\Omega}, E) - \bar{F}_U \quad (2.43)$$

The verification involves integrations similar to those used to verify the representation of the differential source density, equations 2.24 and 2.28, and is not shown.

Substitution of equation 2.40 into equation 2.42 yields:

$$\begin{aligned} \phi^* (\vec{r}, \hat{\Omega}, E) &= \int_0^\infty \sum_{i=1}^n \left[ \sum_{k=0}^n W_{ik}^s (\hat{\Omega}, E) \delta (\vec{r} - \vec{r}_{ik}) \right] \exp \left[ - \int_0^s \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}, E) ds' \right] ds \\ &= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^n \int_0^\infty W_{ik}^s (\hat{\Omega}, E) \delta (\vec{r} - \vec{r}_{ik}) \exp \left[ - \int_0^s \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}, E) ds' \right] ds \end{aligned} \quad (2.44)$$

where the order of the summations and integration has been reversed. To eliminate some of the notational clutter, the individual elements in the summation will be examined with the outer iteration index suppressed:

$$\Delta \phi_k^* (\vec{r}, \hat{\Omega}, E) = \int_0^\infty W_k^s (\hat{\Omega}, E) \delta (\vec{r} - \vec{r}_k) \exp \left[ - \int_0^s \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}, E) ds' \right] ds \quad (2.45)$$

This is also the procedure used in numerical calculations, i.e., individual contributions are computed without considering the iteration index.

The spatial integration is performed using the relationship:

$$\delta (\vec{r} - s \hat{\Omega} - \vec{r}_k) = \delta (s - s_k) \frac{\delta (\hat{\Omega} - \hat{\Omega}_k)}{s_k^2} \quad (2.46)$$

where:

$$s_k = |\vec{r} - \vec{r}_k|, \quad \hat{\Omega}_k = (\vec{r} - \vec{r}_k)/s_k$$

Thus:

$$\begin{aligned} \Delta \phi_k^* (\vec{r}, \hat{\Omega}, E) &= \int_0^\infty W_k^s (\hat{\Omega}, E) \delta (s - s_k) \frac{\delta (\hat{\Omega} - \hat{\Omega}_k)}{s_k^2} \exp \left[ - \int_0^s \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}, E) ds' \right] ds \\ &= W_k^s (\hat{\Omega}, E) \delta (\hat{\Omega} - \hat{\Omega}_k) \frac{\exp \left[ - \int_0^{s_k} \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}, E) ds' \right]}{s_k^2} \end{aligned}$$

The delta function involving the direction vector complicates nothing. In particular, the corresponding contribution to the scalar flux is given by an integration over solid angle:

$$\begin{aligned} \Delta \phi_k^* (\vec{r}, E) &= \iint_{4\pi} \Delta \phi_k^* (\vec{r}, \hat{\Omega}, E) d\Omega \\ &= W_k^s (\hat{\Omega}_k, E) \cdot \frac{\exp \left[ - \int_0^{s_k} \sum_{i=1}^n (\vec{r} - s' \hat{\Omega}_k, E) ds' \right]}{s_k^2} \end{aligned} \quad (2.48)$$

as one would expect.

As indicated in Section 2.4, the sampling functions which minimize the variance in the flux at the point  $\vec{r}_k$  involve this point in a rather complicated manner. However, the techniques used in flux estimation are independent of these considerations and can be used for an arbitrary set of points. It should be noted that flux estimation for arbitrary points will yield a result with an infinite variance<sup>(8)</sup> unless these points are located in volumes which exclude source and scattering points. That is, the  $1/s_k^2$  factor is the trouble maker and its deleterious effect can only be removed by including it in the sampling function or by excluding small values of  $s_k^2$ .

## 2.7 SPATIALLY AVERAGED ANGULAR FLUXES

The problem of infinite variance flux estimates can also be removed by averaging the fluxes at arbitrary points over a specified surface or volume. While these averaged results are usually less desirable than a set of point results, there are instances when averages are the only requirement.

The contribution to a surface averaged flux is obtained by integrating the point result given by equation 2.47:

$$\Delta\phi_k^*(\vec{\Omega}, E)_A = \frac{1}{A} \iint_A W_k^s(\vec{\Omega}, E) \delta(\vec{\Omega} - \vec{\Omega}_k) \frac{\exp\left[-\int_0^{s_k} \Sigma^t(\vec{r} - s' \vec{\Omega}, E) ds'\right]}{s_k^2} dA \quad (2.49)$$

where  $A$  is the area of the specified surface,  $dA$  is a differential element of area, and  $\vec{r}$  is a point on the surface. The integration is then transformed to an integration over solid angle about the point  $\vec{r}_k$ :

$$\Delta\phi_k^*(\vec{\Omega}, E)_A = \frac{1}{A} \iint_A W_k^s(\vec{\Omega}, E) \delta(\vec{\Omega} - \vec{\Omega}_k) \frac{\exp\left[-\int_0^{s_k} \Sigma^t(\vec{r} - s' \vec{\Omega}, E) ds'\right]}{s_k^2} \frac{s_k^2 d\Omega_k}{|\vec{\Omega}_k \cdot \vec{n}|} \quad (2.50)$$

where  $dA = s_k^2 d\Omega_k / |\vec{\Omega}_k \cdot \vec{n}|$ ,  $\vec{n}$  is the unit normal to the surface at  $\vec{r}$ , and a summation over multiple points on the surface which yield the same direction vector  $\vec{\Omega}_k$  is implicit. Note that the integration is performed using a differential element of solid angle  $d\Omega_k$  since  $\vec{\Omega}_k$  is the direction vector determining the point ( $s$ )  $\vec{r}$  on the surface. Using the reciprocal nature of the Dirac delta function and the material attenuation kernel, the integration yields:

$$\Delta\phi_k^*(\vec{\Omega}, E)_A = \frac{1}{A} W_k^s(\vec{\Omega}, E) \frac{\exp\left[-\int_0^{s_k(\vec{\Omega})} \Sigma^t(\vec{r}_k + s' \vec{\Omega}, E) ds'\right]}{|\vec{\Omega} \cdot \vec{n}|} \quad (2.51)$$

where  $s(\vec{\Omega})$  is the distance to the surface.

The contribution to a volume averaged flux is also obtained by integrating the point result:

$$\Delta\phi_k^*(\vec{\Omega}, E)_V = \frac{1}{V} \iiint_V W_k^s(\vec{\Omega}, E) \delta(\vec{\Omega} - \vec{\Omega}_k) \frac{\exp\left[-\int_0^{s_k} \Sigma^t(\vec{r} - s' \vec{\Omega}, E) ds'\right]}{s_k^2} dV \quad (2.52)$$

where  $V$  is the volume over which the flux is averaged.

The integration is transformed using a spherical coordinate system centered at  $\vec{r}_k$ :

$$\Delta\phi_k^*(\vec{\Omega}, E)_V = \frac{1}{V} \iint_{4\pi} \int_{s \text{ in } V} W_k^s(\vec{\Omega}, E) \delta(\vec{\Omega} - \vec{\Omega}_k) \frac{\exp\left[-\int_0^{s_k} \Sigma^t(\vec{r}_k + s' \vec{\Omega}, E) ds'\right]}{s_k^2} s_k^2 ds_k d\Omega_k \quad (2.53)$$

where  $dV = s_k^2 ds_k d\Omega_k$  since  $\vec{\Omega}_k$  and  $s_k$  define the points  $\vec{r}$  in the volume.

Using the same arguments as before, the integration yields:

$$\Delta\phi_k^*(\vec{\Omega}, E) = \frac{1}{V} W_k^s(\vec{\Omega}, E) \int_{s \text{ in } V} \exp \left[ - \int_0^{s(\vec{\Omega})} \Sigma^t (\vec{r}_k + s^t \vec{\Omega}, E) ds^t \right] ds \quad (2.54)$$

The integration over distance  $s$ , is limited to points  $\vec{r}$  on  $\vec{r}_k + s\vec{\Omega}$  which lie in the volume. There may, of course, be several discrete intersections with the volume along  $\vec{\Omega}$ .

There are two interesting, special forms of this equation. Considering only one intersection with the volume, let  $s(\vec{\Omega})$  be the distance to the volume and  $\Delta s(\vec{\Omega})$  the distance in the volume. Then, if the volume is void:

$$\Delta\phi_k^*(\vec{\Omega}, E) = \frac{1}{V} W_k^s(\vec{\Omega}, E) \exp \left[ - \int_0^{s(\vec{\Omega})} \Sigma^t (\vec{r}_k + s^t \vec{\Omega}, E) ds^t \right] \Delta s(\vec{\Omega}) \quad (2.55)$$

and if the volume has constant material properties:

$$\Delta\phi_k^*(\vec{\Omega}, E)_V = \frac{1}{V} W_k^s(\vec{\Omega}, E) \exp \left[ - \int_0^{s(\vec{\Omega})} \Sigma^t (\vec{r}_k + s^t \vec{\Omega}, E) ds^t \right] \frac{[1 - \exp(-\Delta s(\vec{\Omega}) \Sigma^V(E))]}{\Sigma^V(E)} \quad (2.56)$$

where  $\Sigma^V(E)$  is the total cross section at any point in the volume.

There is a possible difficulty with both the surface and volume averaged flux contributions since they are still defined for all directions  $\vec{\Omega}$ ; i.e., a direct numerical integration to obtain various flux components could be prohibitive. If so, random sampling can be

used. A general equation for the integration of an angular kernel  $g(\vec{\Omega})$  is written first:

$$\begin{aligned} \Delta G_k^*(E) &= \iint_{4\pi} \Delta\phi_k^*(\vec{\Omega}, E) \{A\}_{V'} g(\vec{\Omega}) d\Omega \\ &= \iint_{4\pi} \left\{ \frac{\Delta\phi_k^*(\vec{\Omega}, E)}{q^*(\vec{\Omega})} \{A\}_{V'} g(\vec{\Omega}) \right\} q^*(\vec{\Omega}) d\Omega \\ &\approx \frac{1}{L} \sum_{j=1}^L \frac{\Delta\phi_k^*(\vec{\Omega}_j, E)}{q^*(\vec{\Omega}_j)} \{A\}_{V'} g(\vec{\Omega}_j) \end{aligned} \quad (2.57)$$

where  $\{A\}_{V'}$  denotes either a surface or volume,

$$\left. \begin{aligned} q^*(\vec{\Omega}) &\geq 0 \\ q^*(\vec{\Omega}) &> 0 \text{ if } \int_0^{\vec{\Omega}} \Delta\phi_k^*(\vec{\Omega}, E) \{A\}_{V'} dE > 0 \end{aligned} \right\} q^*(\vec{\Omega}) \text{ is a sampling function (2.59)}$$

$$\iint_{4\pi} q^*(\vec{\Omega}) d\Omega = 1$$

$L$  is the total number of discrete directions, and  $\vec{\Omega}_j$  is a discrete direction obtained by random sampling of  $q^*(\vec{\Omega})$ .

This is equivalent to representing the averaged angular fluxes by:

$$\Delta\phi_k^*(\vec{\Omega}, E)_A = \frac{1}{A} \frac{1}{L} \sum_{j=1}^L \frac{W_k^s(\vec{\Omega}, E)}{q^*(\vec{\Omega}_j)} \frac{\exp \left[ - \int_0^{s(\vec{\Omega})} \Sigma^t (\vec{r}_k + s^t \vec{\Omega}, E) ds^t \right]}{|\vec{\Omega} \cdot \vec{\Omega}_j|} \delta(\vec{\Omega} - \vec{\Omega}_j) \quad (2.60)$$

$$\Delta\phi_k^*(\vec{\Omega}, E)_V = \frac{1}{V} \frac{1}{L} \sum_{k=1}^L \frac{W_k^s(\vec{\Omega}, E)}{q^*(\vec{\Omega})} \int_{s \text{ in } V} \exp \left[ - \int_0^s \Sigma^t (\vec{\tau}_k + s^t \vec{\Omega}, E) ds^t \right] ds \delta(\vec{\Omega} - \vec{\Omega}_k) \quad (2.61)$$

for surfaces and volumes respectively since angular integrations with the arbitrary kernel  $g(\vec{\Omega})$  yield equation 2.58 above.

The definition of  $q^*(\vec{\Omega})$  which will minimize the error in the averaged flux angular integrations can be argued in a manner similar to the arguments used in defining the functions  $p_k^*(\vec{\tau})$  used in selecting the source and scattering points. The arguments are simplified since future scattering contributions need not be considered--this is just an angular integration. The arguments are complicated by the fact that the same set of random discrete directions will probably (not necessarily) be used for all the surfaces and/or volumes over which fluxes are being averaged and these surfaces and/or volumes may occupy widely varying spatial positions.

## 2.8 CONCLUDING REMARKS

The previous sections pertained to a development of the Monte Carlo method which utilized random sampling for all of the spatial integrations. This is the procedure used in the FASTER program. In particular, the major equations used in the program are the numerical equivalents of:

- a) equation 2.36 for selecting source points,
- b) equation 2.23 for the point representation of the source,
- c) equation 2.37 for selecting scattering points,
- d) equation 2.26 for the point representation of the flux of particles going into a collision,
- e) equation 2.27 for the point representation of scattered sources,
- f) equation 2.47 for angular point flux estimation,
- g) equation 2.60 for surface averaged angular flux estimation, and
- h) equations 2.61, in conjunction with equations 2.55 and 2.56, for volume averaged angular flux estimation.

Each integration involved in calculating the order-of-scatter fluxes could have been performed by either random sampling or by conventional numerical integrations. In particular, most Monte Carlo programs use random sampling for the energy integrations as well as the spatial integrations and the entire calculational procedure is equivalent to the simulation of individual particle histories.<sup>(9)</sup>

A variety of other combinations of integration techniques can be used. In particular, the variance associated with selecting discrete points from the fixed source in a Monte Carlo integration can be replaced by the systematic error involved in a conventional numerical integration over the spatial extent of the source volume. In fact, uncollided flux calculations can use various combinations of random sampling and direct numerical integration for the three spatial variables. One combined integration procedure is discussed in Appendix A.



### 3.0 PROGRAM LOGIC

The techniques described in Section 2.0 permit a large degree of separation of conventional numerical techniques from random sampling techniques. This separability is utilized in the structure of the FASTER program through a series of subprograms that perform conventional calculations, such as, source interpolation at a point and a single scattering calculation for a fixed scattering angle. Another series of subprograms are used in the random selection of the parameters, e.g., the source point or scattering angle, for these conventional calculations.

### 3.1 DATA REQUIREMENTS

Several major divisions in the data required by the FASTER program have been made. Detailed input instructions for these data are given in Section 9.0.

Section 1 data, i.e., data in the first section of input, include the limits and controls for the FASTER calculations. The requisite data are described in Section 9.2.

Section 2 data are used in describing the geometry of the problem. Details of the geometric calculations performed by FASTER are given in Section 4.0. Data input instructions for geometric parameters are given in Section 9.3.

Section 3 data involve the description of the distributed sources. Calculations associated with these sources are described in Section 5.0. Input instructions for the description of sources are given in Section 9.4.

Section 4 data include the microscopic cross sections used to obtain the macroscopic attenuating and scattering properties of the non-void regions of the geometry. The calculations requiring these cross sections are described in Section 6.0. Detailed input instructions for the cross sections are given in Section 9.5.

Section 5 data are used to specify the final form of the computed results. This includes the description of the various points, surfaces and/or volumes for which the FASTER

program will compute multigroup fluxes. Details of the various flux component estimates used in FASTER are given in Section 7.0. Input instructions for this section of data are given in Section 9.6.

Section 6 data pertain to the description of the random sampling functions. The sampling functions incorporated in FASTER are described in Section 8.0. Input instructions for the sampling parameters are given in Section 9.7.

### 3.2 PROGRAM FLOW

This section attempts to clarify the inter-relationships of data and calculational techniques by discussing the general flow within the FASTER program. The major subprograms of FASTER are also identified as to their function, i.e., which equations they contain.

The program is divided into two major parts. The first part involves data input and preparation and is controlled by the subroutine DEFINE. The second part involves the actual calculations and these are controlled by the subroutine SOLVIT.

#### Data Input and Preparation

DEFINE is the first routine called by FASTER. It calls other routines, in order, for the input of the data described in Section 9.0. The first subroutine entered is STORER. This routine reads the Section 1 data inputs and then allocates storage for all the dimensioned data. If multiple cases are run, it also manipulates the data arrays to account for dimension changes, etc.

Subroutine GEOMIN is then entered and all Section 2 data are input. After all geometric inputs have been processed, the  $\pm$  sign associated with the boundaries of the regions are calculated using the input coordinates of an arbitrary point in each region. Then the geometry is checked using the input point-in-region coordinates and the function subroutine LOCATE, which computes the region(s) occupied by arbitrary points.

Section 3 data are input next in subroutine SOURCE. This includes the energy group structure and the definition and normalization of all the sources. Subroutine INSECT accepts Section 4 cross section data and combines the microscopic data into the requisite macroscopic data.

Section 5 data are input in subroutine RESULT. This includes the flux groups, flux conversion factors, and detector definitions.

Finally, the input sampling parameters—Section 6 data—are input by subroutine RANDOM.

After all data are input, a check of an error count is made. If any errors were detected, the data for the next case is input. If no errors were detected, control is passed back to the FASTER program.

#### Calculations

With all requisite data well defined, control is passed from FASTER to SOLVIT. The SOLVIT routine passes control to one of two available calculational control subprograms. The first, SOBER, was written to compute surface and volume averaged fluxes and/or fluxes at multiple point detectors in void regions. The second control subprogram, SOLVER, was coded for the individual treatment of point detectors at arbitrary locations in the geometry.

As indicated, these routines perform almost identical functions. The following discussion attempts to describe both simultaneously. There may be minor variations from the actual order of some of the calculations, but they are unimportant in the overall picture. Some of the routines discussed below also require computations by other routines, however, this secondary control level will not be discussed.

The first step is the definition of a preferred point for use in the sampling procedures. In SOBER it is defined by input and is surrounded by a sphere with an input radius which encloses a volume in space where fluxes will be calculated. SOLVER contains an iteration over detector points and the preferred point is the detector point being treated. An average source group index, used in a one-velocity approximation of equation 2.36, is then computed by subroutine GROUP.

All other calculations are performed within the outer iteration loop. The first calculation within this loop is the random selection of a position vector  $\vec{r}_0$  from the fixed source. This is done by one of two subroutines; PSTAR, if the sampling is

performed in the source geometry coordinate system or SPHERE, which approximates equation 2.36.

The calculations enter the inner iteration loop where they remain until terminated by one of the cutoff criteria, i.e., maximum number of inner iterations (collision cutoff), weight cutoff, or energy cutoff. The first iteration in this inner loop differs from all subsequent inner iterations because the source point is actually a point in a source volume. For all inner iterations after the first, it is actually a scattering (scattered source) point.

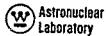
The first calculation in the inner loop involves the flux contribution to the point detectors. For each of these detectors, calculation of the distance and direction to the detector from the source point is performed by the function subroutine VECTOR. The calculation of the angular point source for the direction towards the detector is then performed by interpolation in the SZERO subroutine (equation 2.23) if this is the first inner iteration, or by SINGLE (equation 2.27) for subsequent inner iterations. If there is a non-zero source for this direction, the path lengths through the various regions lying between the source (scattering) point and the detector point are computed by subroutine PATH. The mean-free-paths (mfp) along the total path are then computed for each source group by subroutine KERNEL and the flux estimation performed in subroutine DETECT using equation 2.47.

The next step is performed only in the SOBER inner iterations. It involves the random selection of discrete directions for calculating surface and volume averaged fluxes. These directions are obtained from subroutine VSTAR. A possible exception is the first inner iteration where they can be obtained from subroutine QSTAR. The source for each fixed direction is computed using SZERO or SINGLE. If non-zero, the regions lying along the direction vector are computed by PATH. Each of the regions lying along the direction vector is checked to see if it is a volume detector and should receive a flux contribution. If so, the mfp's up to the region are computed by KERNEL and the flux is computed by DETECT using contribution equation 2.61 in conjunction with equations 2.55 and 2.56. Each boundary crossing between regions is also checked to see if it is a surface detector. If it is, the normal derivative at the

boundary is computed by subroutine NORMAL; the mfp's to the boundary are obtained from KERNEL and the flux contributions are then computed by DETECT using equation 2.60.

The final step in each inner iteration is the random selection of the next scattering point. The average flux contribution energy, as calculated by DETECT, is used to define an average energy group for the sampling procedures. The direction vector defining the scattering point is then obtained from VSTAR (or possibly QSTAR on the first inner iteration). The regions lying along the ray defined by this direction are computed by PATH. The distance to the collision point is obtained from USTAR, (equation 2.37) and the previous source point is then evaluated for the direction vector from this prior point to the new point, using SZERO or SINGLE. If non-zero, the mfp's to the scattering point are obtained from KERNEL, and the point monodirectional fluxes are computed at the scattering point using equation 2.26. The next inner iteration is initiated with the sources being obtained from SINGLE using the monodirectional fluxes. These inner iterations are continued until a cutoff is obtained.

The outer iterations are continued to a specified maximum with a printout of the flux edits being performed by subroutine ANSWER at specified intervals.



## SECTION

### 4.0 GEOMETRIC CONSIDERATIONS

The most important feature of the Monte Carlo method—in comparison with other "exact" solutions of the transport equation—is its applicability to complicated geometries. This feature is implemented in the FASTER program by utilizing the general quadric surface equation. The numerical analysis presented below follows that of Reference 10.

#### 4.1 QUADRIC SURFACES

The general quadric equation for a specified surface  $\vec{r}$  is:

$$u_i(\vec{r}) = a_{0,i} + a_{1,i}x + a_{2,i}y + a_{3,i}z \\ + a_{4,i}x^2 + a_{5,i}y^2 + a_{6,i}z^2 \\ + a_{7,i}xy + a_{8,i}yz + a_{9,i}zx \quad (4.1)$$

where  $a_{j,i}$ ,  $j = 0, 1, 2, \dots, 9$  are constants,

$$\vec{r} = x\hat{i} + y\hat{j} + z\hat{k}$$

$x, y, z$  are rectangular coordinates (cm), and

$\hat{i}, \hat{j}, \hat{k}$  are unit vectors parallel to the  $x$ -,  $y$ -, and  $z$ -axes, respectively. The value of this equation,  $u_i(\vec{r})$ , is zero for points  $\vec{r}$  on the surface.

These surfaces are described independently of the regions which define the material distributions to eliminate redundant input. Provision has been made in the FASTER program for recognizing more simple surfaces such as planes, cones, elliptical cylinders and ellipsoids. The equations for these simple surfaces are expanded by subroutine GEOMIN to obtain the coefficients of the general equation above. These special surfaces are shown in Figures 7, 8, and 9 and the equations are tabulated in Table 1 of the Section 9.3 Input instructions.

#### 4.2 SURFACE CALCULATIONS

Since surfaces are described independent of geometric regions, it is possible to define several quantities which are used in geometric calculations. For this discussion, the following are defined:

$$\vec{r} = (x, y, z), \quad \text{the position vector of any point in space,}$$

$\vec{\Omega} = (\alpha, \beta, \gamma)$ ,  
 a unit vector defining the direction of a straight line,  
 or ray, emanating from  $\vec{r}_i$ , where  $\alpha, \beta, \gamma$  are direction  
 cosines with respect to the  $x, y$ , and  $z$  axes, respectively,

$s, \quad 0 \leq s < \infty$  the scalar distance from  $\vec{r}_i$  along  $\vec{\Omega}$

$$\begin{aligned}\vec{r}' &= \vec{r}_i + s\vec{\Omega} \\ &\text{a point on the ray} \\ &= (x + \alpha s, y + \beta s, z + \gamma s)\end{aligned}$$

#### Intersection of a Line and a Surface

The value of the quadratic equation at  $\vec{r}'$  is given by:

$$u_i(\vec{r}') = u_i(\vec{r}_i) + 2s v_i(\vec{r}, \vec{\Omega}) + s^2 w_i(\vec{r}, \vec{\Omega}) \quad (4.2)$$

where  $v_i(\vec{r})$  is given by equation 4.1 above and  $v_i(\vec{r}, \vec{\Omega})$  and  $w_i(\vec{r}, \vec{\Omega})$  are obtained by expanding this equation for  $\vec{r}'$ , and collecting the coefficients of  $s$  and  $s^2$ :

$$\begin{aligned}v_i(\vec{r}, \vec{\Omega}) &= \frac{1}{2} [\alpha a_{1,i} + \beta a_{2,i} + \gamma a_{3,i}] \\ &+ [\alpha x a_{4,i} + \beta y a_{5,i} + \gamma z a_{6,i}] \\ &+ \frac{1}{2} [(\alpha y + \beta x)a_{7,i} + (\beta z + \gamma y)a_{8,i} + (\gamma x + \alpha z)a_{9,i}] \quad (4.3)\end{aligned}$$

$$\begin{aligned}w_i(\vec{r}, \vec{\Omega}) &= \frac{1}{2} [\alpha^2 a_{4,i} + \beta^2 a_{5,i} + \gamma^2 a_{6,i}] \\ &+ [\alpha \beta a_{7,i} + \beta \gamma a_{8,i} + \gamma \alpha a_{9,i}] \quad (4.4)\end{aligned}$$

Intersections of the ray with the surface are obtained by requiring:

$u_i(\vec{r}') = 0$ , i.e. this condition defines points on the surface using equation 4.2.

$$u_i(\vec{r}) + 2s_i v_i(\vec{r}, \vec{\Omega}) + s_i^2 w_i(\vec{r}, \vec{\Omega}) = 0 \quad (4.5)$$

a) one intersection if  $w_i(\vec{r}, \vec{\Omega}) = 0, v_i(\vec{r}, \vec{\Omega}) \neq 0$

$$s_i = -v_i(\vec{r}) / 2v_i(\vec{r}, \vec{\Omega}) \quad (4.6)$$

b) two intersections if  $v_i(\vec{r}, \vec{\Omega}) \neq 0, v_i^2(\vec{r}, \vec{\Omega}) > u_i(\vec{r}) w_i(\vec{r}, \vec{\Omega})$

$$s_i = \frac{-v_i(\vec{r}, \vec{\Omega}) \pm \sqrt{v_i^2(\vec{r}, \vec{\Omega}) - u_i(\vec{r}) w_i(\vec{r}, \vec{\Omega})}}{w_i(\vec{r}, \vec{\Omega})} \quad (4.7)$$

Note that the case of two equal intersections is not admitted since this is equivalent to no intersection. For all real intersections, the appropriate sign for multiple intersections is determined in the following manner:

a) The rate of change of  $u_i(\vec{r}')$  with respect to distance at the intersection is given by differentiation of equation 4.2:

$$\left. \frac{\partial u_i(\vec{r}')}{\partial s} \right|_{s=s_i} = 2 \left[ v_i(\vec{r}, \vec{\Omega}) + s_i w_i(\vec{r}, \vec{\Omega}) \right] \quad (4.8)$$

Using the intersection equation 4.7 above:

$$\left. \frac{\partial u_i}{\partial s}(\vec{r}') \right|_{s=s_i} = 2 \left[ v_i(\vec{r}, \vec{\Omega}) + w_i(\vec{r}, \vec{\Omega}) \right] \left\{ \frac{-v_i(\vec{r}, \vec{\Omega}) \pm \sqrt{v_i^2(\vec{r}, \vec{\Omega}) - u_i(\vec{r}) w_i(\vec{r}, \vec{\Omega})}}{w_i(\vec{r}, \vec{\Omega})} \right\}$$

$$= \pm 2 \sqrt{v_i^2(\vec{r}, \vec{\Omega}) - u_i(\vec{r}) w_i(\vec{r}, \vec{\Omega})} \quad (4.9)$$

Thus this derivative must have the sign (+) used in the intersection equation.

b) It is noted that each surface defines two disjoint volumes, which, for the sake of a convention, are described as:

$$\left. \begin{array}{l} \text{inner volume: } u_i(\vec{r}') < 0 \\ \text{outer volume: } u_i(\vec{r}') > 0 \end{array} \right\} \quad (4.10)$$

It follows that in crossing the surface from the inner volume,  $u_i(\vec{r}') < 0$ , to the outer volume,  $u_i(\vec{r}') > 0$ , along any straight line, that the rate of change of the value of the surface equation  $\partial u_i(\vec{r}') / \partial s$  is greater than zero at the intersection. Thus, if  $u_i(\vec{r}) < 0$  then the origin of the ray is inside the surface, and the first intersection with the surface is obtained from the quadratic equation 4.7 using the positive sign. A similar argument holds for crossing from the outside to the inside of the surface; i.e. if  $u_i(\vec{r}) > 0$ , then  $\partial u_i(\vec{r}) / \partial s < 0$  at the first intersection implying the negative sign.

#### Surface Normal

The normal vector to the surface at the intersection, used in surface averaged flux calculations, is calculated by subroutine NORMAL as:

$$\hat{n} = \frac{\nabla u_i(\vec{r}')}{\|\nabla u_i(\vec{r}')\|} = (c_1^n, c_2^n, c_3^n) \quad (4.11)$$

$$\text{where } \nabla = \vec{i} \cdot \frac{\partial}{\partial x'} + \vec{j} \cdot \frac{\partial}{\partial y'} + \vec{k} \cdot \frac{\partial}{\partial z'} \quad (4.12)$$

$$\begin{aligned} \hat{n} \cdot [\Delta u_i(\vec{r}')] &= [a_{1,i} + 2x' a_{4,i} + a_{7,i} y' + a_{9,i} z'] \vec{i} \\ &\quad + [a_{2,i} + 2y' a_{5,i} + a_{8,i} z' + a_{7,i} x'] \vec{j} \\ &\quad + [a_{3,i} + 2z' a_{6,i} + a_{9,i} x' + a_{8,i} y'] \vec{k} \\ &= c_1 \vec{i} + c_2 \vec{j} + c_3 \vec{k} \end{aligned} \quad (4.13)$$

$$c_i^n = \sqrt{\sum_{i=1}^3 c_i^2}, \quad i = 1, 2, 3 \quad (4.14)$$

#### 4.3 REGIONS

##### Material Properties

The surfaces referred to in the previous section are used to describe the extent of geometric regions or zones having constant material properties. These properties are specified for each region,  $i$ , by a composition indicator  $m_i$  ( $m_i < 0$  indicates that region  $i$  is void) and a separate hydrogen density  $\rho_i^h$ . The capability of specifying hydrogen densities by region simplifies the description of many problems, e.g., regions in a liquid hydrogen propellant tank. It is also helpful in describing hydrogen density variations in NERVA-type reactors where all other material properties are constant.

##### Source-in-Region

Additional regions may be required to correctly define the spatial extent of the fixed source volumes, i.e., there is an optional sampling technique in the FASTER program (described in Section 8.3) which requires that only one source be superimposed over a region and that the source cover the region. Alternatively, each source may cover more than one region.

This sampling technique is preferred since it requires much less *a priori* knowledge of the importance of various sources--point-kernel importance estimates are built in.

#### Region Boundaries

The geometric description of each region involves the listing of the surfaces which bound the region:

$$k_{j,i}, \quad i = 1, 2, \dots \quad \text{where } k_{j,i} \text{ is the index of the surface forming the } [i]^{th} \text{ boundary of the region.}$$

Also required are the components of an arbitrary point  $\vec{r}_i^g$  in the region:

$$\vec{r}_i^g = x_i^g \hat{i} + y_i^g \hat{j} + z_i^g \hat{k} \quad \text{where } (x_i^g, y_i^g, z_i^g) \text{ are specified.}$$

#### Ambiguity Indices

The description of each region is completed in subroutine GEOMIN by computing an "ambiguity index" for each boundary surface. This ambiguity index indicates whether the region is inside or outside each of its boundaries. It also yields the sign to be used in calculating distances to quadratic or quadric boundaries. The ambiguity index is computed using equation 4.1:

$$\delta_{j,i} = -\frac{v_k(\vec{r}_i^g)}{|v_k(\vec{r}_i^g)|}, \quad k = k_{j,i}, \quad i = 1, 2, \dots \quad (4.15)$$

#### Region Occupied by a Point

To ensure correct calculations, it is necessary that the ambiguity indices have the same sign for all points inside the region, i.e., a point  $\vec{r}$  is in region 1 if, and only if  $\delta_{j,i} v_{k,j}(\vec{r}) < 0$  for all boundaries. The region index calculation for an arbitrary point is performed by the function subprogram LOCATE.

#### Possible Region Description Errors

Restrictions must be imposed on possible region shapes to ensure that all points in a region are always on the same side of each region boundary. For example, the single region

indicated in Figure 3A is unacceptable since there are points in the region which are both inside and outside boundaries A and B. The obvious solution is to use two regions to describe such geometric shapes.

It is sometimes necessary to introduce fictitious boundaries. Figure 3B shows a typical situation requiring these boundaries. Examination of the shaded and cross hatched regions reveals that they form two sections of a single region since ambiguity indices of the boundaries have the same values for both sections. This condition can cause trouble if only one section of the region is desired, even if other regions occupy the second section. The specification of the fictitious boundary eliminates the problem without otherwise affecting the geometric calculations. These fictitious boundaries must be included in the initial surface descriptions.

#### Geometry Consistency Check

A rather simple check for correct geometric description involves a calculation of the region(s) occupied by each point  $\vec{r}_i^g$ . If the point  $\vec{r}_i^g$  is in any region other than region  $i$ , the geometry representation is incorrect. This geometry check is performed by the GEOMIN subroutine using the LOCATE function subprogram.

#### 4.4 RAY TRACING

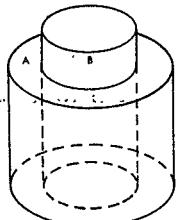
The procedure for ray tracing is similar to that discussed above for the intersection of a ray with individual surfaces. The cumulative path lengths through each region along a ray are computed at a single pass, in the order traversed, by subroutine PATH. As indicated before, the ray tracing calculations are related to the surfaces.

$$s_{l+1} = \min_{i=1, 2, \dots} \left\{ s_{j,i} \text{ such that } s_{j,i} \geq s_l \text{ and } \text{if } s_{j,i} = s_l, \quad \delta_{j,i} \partial v_k(\vec{r} + s_{j,i} \vec{n}) / \partial s > 0 \right\} \quad (4.16)$$

$s_l$  = distance up to the region, and

$s_{j,i}$  = is obtained from equation 4.6 or 4.7 for surface  $k = k_{j,i}$  using the sign of  $\delta_{j,i}$  for non-planar surfaces.

## 3A. AMBIGUOUS REGION BOUNDARIES



## 3B. DISJOINT REGIONS

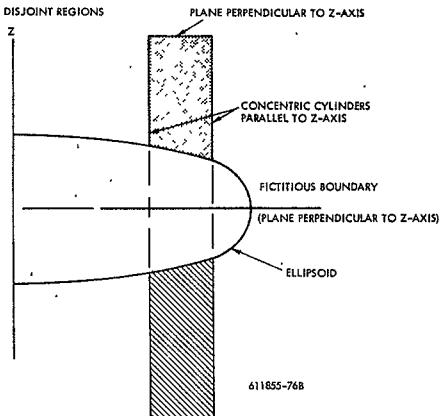


Figure 3. Problems in Region Descriptions

The surface crossed in leaving the region is that giving the minimum distance.

If the new total distance  $s_{j+1}$  exceeds a specified maximum, e.g., the distance between two points, the distance through the region is adjusted and the ray tracing terminates.

If not, the next region entered by the ray is computed by requiring that:

- the next region have the surface just crossed as a boundary,
- the next region be on the other side of this surface (opposite signs on ambiguity indices), and
- for all other boundaries

$$v_k(\vec{r} + s_{j+1}\vec{\Omega}) - \delta_{j+1}^{k+1} \leq 0, \quad k^i = k_{j+1}^i, \quad i = 1, 2, \dots \quad (4.17)$$

which is evaluated using equation 4.2.

The index of the region entered in crossing the above boundary is saved as the most probable next region for subsequent ray tracings. If no region is accepted, the exterior of the geometry is assumed and an indicator set. After several subsequent failures, testing for this boundary is never performed.

The constants  $u_k(\vec{r})$ ,  $v_k(\vec{r}, \vec{\Omega})$ , and  $w_k(\vec{r}, \vec{\Omega})$  and the intersections are computed only once during a given ray tracing. The current status of calculations for the  $k^i$  surface is indicated by  $n_k$  where

- $n_k = 0$  if the  $k^i$  surface has not been involved as yet, in the ray tracing,
- $n_k = 1$  if the constants  $u_k(\vec{r})$ ,  $v_k(\vec{r}, \vec{\Omega})$ ,  $w_k(\vec{r}, \vec{\Omega})$  have been computed,
- $n_k = 2$  if the constants and the intersections have both been calculated and at least one intersection is at a distance greater than the current cumulative distance  $s_j$ , and
- $n_k = 3$  if the constants and intersections have been computed but the intersections need no longer be considered e.g., if both are less than  $s_j$ , imaginary, etc.



## SECTION

### 5.0 FIXED SOURCES

A number of numerical techniques can be used for describing fixed source distributions. Those employed in the FASTER program incorporate the assumption of separable variables. These techniques are general enough to permit the description of a variety of real source distributions and the distributions used in generating basic data.

### 5.1 SPATIAL AND ANGULAR VARIABLES

The FASTER program will handle multiple sources in rectangular, cylindrical or spherical geometries. The geometry for each source is superimposed over the various geometric regions. The source geometries for each of the multiple sources need not be the same. In all geometries there are three spatial variables ( $v_1, v_2, v_3$ ) and two angular variables ( $v_4, v_5$ ). The relationships between these spatial and angular variables are shown in Figure 10 in the Section 9.4 input instructions.

#### Rectangular Geometry

The most simple geometry is that used for describing rectangular source volumes. The spatial variables ( $v_1, v_2, v_3$ ) are the rectangular coordinates ( $x, y, z$ ). The angular variables ( $v_4, v_5$ ) are the azimuthal angle  $\theta$  measured from the  $x$ -axis and the cosine of the polar angle,  $\mu$ , measured from the  $z$ -axis.

#### Cylindrical Geometry

The next allowed geometry, usually used in describing reactor sources, involves cylinders parallel to the  $z$ -axis. The spatial variables are:

$$v_1 = r = \sqrt{x^2 + y^2}, \text{ the radius}$$

$$v_2 = \theta = \tan^{-1}(y/x), \text{ the azimuthal angle measured from the } x\text{-axis}$$

$$v_3 = z, \text{ the axial coordinate}$$

The angular variables are measured in a coordinate system which rotates with the radius vector. The variables are:

$$v_4 = \theta^*, \text{ the azimuthal angle measured from } \theta.$$

$$v_5 = \mu^*, \text{ the cosine of the polar angle measured from the } z\text{-axis.}$$

This simplifies the description of angular sources on the surfaces of a cylindrical reactor.

#### Spherical Geometry

The final geometry, useful in describing sources such as capture gammas in the hemispherical bottom of a liquid hydrogen propellant tank, involves the spatial variables:

$$v_1 = \rho = \sqrt{x^2 + y^2 + z^2}, \text{ the spherical radius}$$

$$v_2 = \theta = \tan^{-1}(y/x) \text{ the azimuthal angle measured from the } x\text{-axis.}$$

$$v_3 = \mu = \frac{z}{\rho}, \text{ the cosine of the polar angle measured from the } z\text{-axis.}$$

The coordinate system used for the angular variables rotates with the spherical radius vector:

$$v_4 = \theta^*, \text{ the azimuthal angle measured as shown in Figure 10}$$

$$v_5 = \mu^*, \text{ the cosine of the polar angle measured from the spherical radius vector.}$$

#### Source Translations

In addition, each source is given a translation vector  $\vec{r}_i = (x_i, y_i, z_i)$  from the origin of the geometry coordinate system. Thus, the coordinates of the source points, expressed in the geometry coordinate system, are

$$\vec{r} = (x + x_i, y + y_i, z + z_i)$$

#### 5.2 SPATIAL AND ANGULAR DISTRIBUTIONS

Each of the distributions for the spatial and angular variables is described separately by tabulating relative distributions

$$(v_{k,i}, f_{k,i}), \quad k = 1, 2, \dots$$

where  $v_{k,i}$  is the  $k$ th value of the  $i$ th variable, and  $f_{k,i}$  is the relative distribution at  $v_{k,i}$ .

Each variable  $v_j$  may take on only one value  $v_j = v_{1,j}$  i.e.,  $f(v_j) = \delta(v_j - v_{1,j})$ . If more than one point is needed, then  $f(v_j)$  is assumed to be continuous.

The continuous distributions are normalized in subroutine SOURCE by integrating a linear interpolation formula and requiring that:

$$\sum_{k=1}^{v_{k+1,i}} \int_{v_{k,i}}^{v_{k+1,i}} \frac{(v_{k+1,i} - v_j) f_{k,i} + (v_j - v_{k,i}) f_{k+1,i}}{v_{k+1,i} - v_{k,i}} v_j^n dv_j = 1 \quad (5.1)$$

where  $n = 0$  except for the radial distributions of cylindrical and spherical sources where  $n = 1$  or  $2$ , respectively.

The final representation of the spatial and angular distributions is the product of the individual distributions for the five source variables:

$$p(\vec{r}, \vec{\Omega}) = \prod_{j=1}^5 f(v_j) \quad (5.2)$$

### 5.3 SOURCE SPECTRA

Particle energies generally decrease with the increase in the order of scattering and the FASTER program requires the same order in energies, i.e., a series of energy groups are defined with group 1 containing particles with the maximum energy. The same group structure is used for the source spectra and later for the cross sections. These energy groups are defined in subroutine SOURCE by:

$$\text{Group } i : E_i \geq E (\text{MeV}) \geq E_{i+1} \quad i = 1, 2, \dots \quad (5.3)$$

Some relaxation is allowed on describing the source spectra in that an arbitrary group structure, with energy group boundaries of decreasing energy, can be used. Various quantities are accepted as input. They are all reduced, however, to one form; a differential number spectrum:

$$\eta_k \left( \frac{\text{particles}}{\text{MeV} \cdot \text{sec}} \right) \text{ at energy } E_k \quad k = 1, 2, \dots$$

where the  $E_k$ 's define the input energy group boundaries.

This spectrum is then integrated into the group structure for the problem using a linear interpolation formula for the energy variation. The final spectrum is expressed as the number of particles in each group and the average energy of these particles:

$$n_i^o = \sum_{k=1}^{E_{ik}^h} \int_{E_{ik}^l}^{E_{ik}^h} \eta_k(E) dE \quad (\text{particles in group } i) \quad (5.4)$$

$$\overline{E}_i^o = \frac{1}{n_i^o} \sum_{k=1}^{E_{ik}^h} \int_{E_{ik}^l}^{E_{ik}^h} \eta_k(E) E dE \quad (\text{average energy of the particles in group } i) \quad (5.5)$$

$$\text{where } \eta_k(E) = \frac{(E - E_{k+1}) \eta_k + (E_k - E) \eta_{k+1}}{E_k - E_{k+1}} \quad (5.6)$$

$$\left. \begin{aligned} E_{ik}^l &= \max(E_{i+1}, E_{k+1}) \\ E_{ik}^h &= \min(E_i, E_k) \end{aligned} \right\} \quad (5.7)$$

The groupwise number spectrum  $n_i^o$ ,  $i = 1, 2, \dots$ , is then normalized to an input total source strength.

### 5.4 FIXED SOURCE ACQUISITION

#### Volume Sources

The definition of neutron and photon source distributions for reactor configurations can be a time-consuming task. For geometries where the discrete ordinates ( $S_n$ ) methods are applicable, source distributions can be obtained both efficiently and economically through their use. In particular, the coupled ODD-K - NAGS system (References 11, 12) can provide the relative distributions and spectra in the form required by the FASTER program. Included are separable radial and axial distributions and spectra for all reactor regions.

#### Angular Surface Fluxes

The design of nuclear rocket engines is such that given the internal reactor arrangements it is possible to define the reactor leakage within an error involving external reactivity

effects and that inherent in the calculational method. This same leakage can then be applied to a variety of external problems such as individual external components or a liquid hydrogen propellant tank.

Assuming a detailed internal calculation, it is possible to numerically integrate the equation for the unperturbed angular flux at an arbitrary point in space outside the reactor:

$$\begin{aligned}\phi^U(\vec{r}, \vec{\Omega}, E) &= \int_0^\infty S(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^I(\vec{r} - s'\vec{\Omega}, E) ds' \right] ds \\ &= \int_0^{s_1} S(\vec{r} - s\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^s \Sigma^I(\vec{r} - s'\vec{\Omega}, E) ds' \right] ds \\ \phi^U(\vec{r}, \vec{\Omega}, E) &= \int_0^{s_1 - s_0} S(\vec{r} - s\vec{\Omega} - t\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^t \Sigma^I(\vec{r} - s\vec{\Omega} - t\vec{\Omega}, E) dt' \right] dt \\ &= \phi(\vec{r} - s_0\vec{\Omega}, \vec{\Omega}, E)\end{aligned}\quad (5.8)$$

where  $S(\vec{r}, \vec{\Omega}, E)$  is the total differential source density in the reactor, i.e., including scattering.  
 $s_0$  is the distance to the reactor surface from  $\vec{r}$   
 $s_1$  is the distance through the reactor from  $\vec{r}$   
 $\phi(\vec{r} - s_0\vec{\Omega}, \vec{\Omega}, E)$  is the angular flux at the reactor surface, i.e., at  $\vec{r} = s_0\vec{\Omega}$

(For numerical integrations, the equation is sometimes transformed to an area integration over the reactor surface.)

A more general use of the angular leakage fluxes is as a surface source (or a "thin" volume source). The requisite data reduction of ODD-K angular leakage fluxes, both neutron and photon, is performed by the DAFT program.<sup>(13)</sup> Included in this reduction are averages over arbitrary mesh points, groups, etc.

The external environment is then given by a set of equations analogous to those in Section 2.2, where:

$$\phi_o(\vec{r}, \vec{\Omega}, E) = \phi(\vec{r} - s_o\vec{\Omega}, \vec{\Omega}, E) \exp \left[ - \int_0^{s_o} \Sigma^I(\vec{r} - s'\vec{\Omega}, E) ds' \right] \quad (5.9)$$

These reduced angular fluxes thus fit into the framework provided in the FASTER program for describing fixed sources, since the equation for the external uncollided flux is identical to that obtained for a surface source.

## 5.5 SOURCE EVALUATION

The source was defined as a function of the variables  $v_1, v_2, \dots, v_5$ . For the order of scatter flux calculations, the source must be evaluated for a specified point  $\vec{r}_o$  and direction  $\vec{\Omega}$ . This evaluation, when divided by the value  $p_o(\vec{r}_o)$  of the sampling function used in obtaining the position vector  $\vec{r}_o$ , yields the energy dependent angular point source  $W_o(\vec{\Omega}, E)$ . The source evaluation is performed by subroutine SZERO and involves several variable transformations to obtain the source variables equivalent to  $\vec{r}_o$  and  $\vec{\Omega}$ :

- define the source centered position vector,  
 $\vec{r} = \vec{r}_o - \vec{r}_f = (x, y, z)$
- calculate the spatial variables  $(v_1, v_2, v_3)$  from the rectangular coordinates using the transformation equations for the source geometry,
- calculate the direction vector rotation matrix,  $R_{ij}$ , (See Figure 4.),
- calculate the direction cosines in the rotated coordinate system for which the angular distribution of the source is defined

$$c_i^j = \sum_{i=1}^3 c_i R_{ij}, \quad j = 1, 2, 3 \quad \text{where } \vec{\Omega} = c_1 \vec{i} + c_2 \vec{j} + c_3 \vec{k}$$

e) calculate the angular variables  $v_4$  and  $v_5$ :

$$v_4 = \tan^{-1} (c'_2/c'_1) \quad v_5 = c'_3$$

A linear interpolation of the source variable distributions yields:

$$W_i^s = \left[ \frac{5}{\pi} \int_{\vec{r}_o}^{\vec{r}} f(v_i) d\vec{r} \right] n_i^o \quad i = 1, 2, \dots$$

$$\bar{E}_i^s = E_i^o$$

where  $W_i^s$  is the number of particles in group  $i$  and  $\bar{E}_i^s$  is the average energy of the particles in group  $i$ . The average group energy  $\bar{E}_i^s$  is just  $E_i^o$  as calculated during spectrum normalization.

Thus the general equation 2.23; for  $W_o^s(\vec{r}, E)$  is reduced to a groupwise representation:

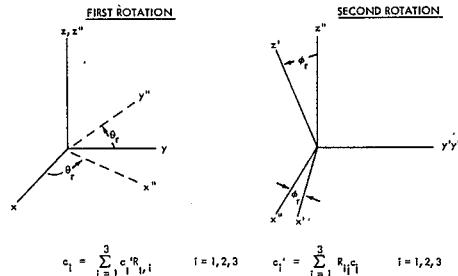
$$W_i^s = \int_{E_{i+1}}^{E_i} W_o^s(\vec{r}, E) dE$$

$$\bar{E}_i^s = \frac{1}{W_i^s} \int_{E_{i+1}}^{E_i} W_o^s(\vec{r}, E) E dE$$

The angular dependence has been suppressed since, in the numerical calculations, only one discrete direction is considered at a time. Moreover, the order-of-scatter subscript has also been suppressed since it is used only as a counter in the numerical calculations.

The same representation of the point angular sources is used for higher order-of-scatters i.e., for  $W_k^s(\vec{r}, E)$ . The equations are discussed in the next section.

$$\begin{aligned}\vec{n} &= \text{unit direction vector in geometry coordinate system} \\ &= c_1^{\perp} + c_2^{\perp} + c_3^{\perp} \\ \vec{n}' &= \text{unit direction vector in rotated coordinate system} \\ &= c_1^{\parallel} + c_2^{\parallel} + c_3^{\parallel}\end{aligned}$$



$$c_i = \sum_{j=1}^3 c_j^i R_{j,i} \quad i = 1, 2, 3 \quad c_i' = \sum_{j=1}^3 R_{j,i} c_j \quad i = 1, 2, 3$$

$$\begin{aligned}R_{1,1} &= \cos\theta_r \cos\phi_r & R_{1,2} &= \sin\theta_r \cos\phi_r & R_{1,3} &= -\sin\phi_r \\ R_{2,1} &= -\sin\theta_r & R_{2,2} &= \cos\theta_r & R_{2,3} &= 0 \\ R_{3,1} &= \cos\theta_r \sin\phi_r & R_{3,2} &= \sin\theta_r \sin\phi_r & R_{3,3} &= \cos\phi_r\end{aligned}$$

#### SOURCE GEOMETRY ROTATION ANGLES

ANGLE	RECTANGULAR	CYLINDRICAL	SPHERICAL
θ <sub>r</sub> (azimuthal)	0	$\tan^{-1}(y/x)$	$\tan^{-1}(y/r)$
φ <sub>r</sub> (polar)	0	0	$\tan^{-1}\left(z/\sqrt{x^2+y^2+z^2}\right)$

Figure 4. Direction Vector Rotations



## SECTION

### 6.0 TRANSPORT AND SCATTERING KERNELS

The equations used in the FASTEP program for defining the attenuating and scattering properties of homogeneous material compositions are described below. The general procedure involves a calculation of microscopic cross sections using input microscopic data. The two most common sets of units for microscopic data and compositions are both permitted in FASTER:

- $u_x = 0$  for microscopic cross sections in barns/atom
- $= 1$  for microscopic cross sections in  $\text{cm}^2/\text{gm}$
- $u_c = 0$  for compositions in  $10^{24}$  atoms/ $\text{cm}^3$
- $= 1$  for compositions in  $\text{gm}/\text{cm}^3$

i.e., the composition and cross section units can be mixed.

### 6.1 PHOTON CROSS SECTIONS

The equations used for photon cross sections are discussed in this section. Photon cross sections are defined at the boundaries of the energy groups defined by equation 5.3.

The requisite data for the  $i^{\text{th}}$  element is:

- $A_i$  the atomic mass ( $a, m, u_s$ ) of the element
- $Z_i$  the atomic number of the element
- $\rho_{m,i}$  the density of the element in composite material  $m$  with units according to  $u_c$  and
- $\sigma_{i,i}^t$  the microscopic total cross section for energy level  $i$  with units according to  $u_x$ .

The total cross section,  $\Sigma_{i,m}^t$ , is computed by subroutine INSECT for each energy level of each composite material by a summation over the element data:

$$\Sigma_{i,m}^t = \sum_{i=1}^f \left( \frac{0.6025}{A_i} \right)^{(u_c - u_x)} \rho_{m,i} \sigma_{i,i}^t \quad (\text{cm}^{-1}) \quad (6.1)$$

The total electron density is calculated by:

$$\rho_m^e = \sum_{i=1}^n \left( \frac{0.6025}{A_i} \right)^{u_c} \rho_{m,i} z_i \left( \frac{10^{24} \text{ electrons}}{\text{cm}^3} \right) \quad (6.2)$$

where  $(0.6025/A_i)^{u_c}$  converts compositions to  $10^{24} \text{ atoms/cm}^3$ , and  $(0.6025/A_i)^{u_x}$  converts cross sections to barns/atom.

The energy absorption coefficient by energy level is computed by element and material. The absorption coefficients can then be used as a flux-to-heating conversion factor under conditions noted in Section 9.5 of the input instructions. The equation for the microscopic energy absorption coefficients assumes all interactions, except Compton scattering, are absorptions (Reference 4, pg. 159):

$$\mu_{i,i}^a = \sigma_{i,i}^t - \frac{1.995}{8} \left( \frac{0.6025}{A_i} \right)^{u_x} \left[ \frac{\ln(1+2\eta)}{\eta^3} + \frac{2(1+\eta)(2\eta^2 - 2\eta - 1)}{\eta^2(1+2\eta)} + \frac{8\eta^2}{3(1+2\eta)^3} \right] \quad (6.3)$$

with units according to  $u_x$  and where  $\eta = E_1/0.511$ .

## 6.2 PHOTON TRANSPORT

The material attenuation kernel is written for the  $i$ th energy group as:

$$K_i^t(\vec{r}, \vec{r}') = \exp \left[ - \int_0^s \Sigma^t(\vec{r} + s\vec{\Omega}, \vec{E}_i^s) ds \right], \quad s = |\vec{r} - \vec{r}'|, \vec{\Omega} = (\vec{r} - \vec{r}')/s \quad (6.4)$$

This kernel is evaluated using  $\vec{E}_i^s$ , the average group energy of the multigroup representation of the angular point sources defined in equation 5.10. The assumption of transport without change in the average group energy simplifies this kernel.

The evaluation of this kernel uses a group averaged cross section obtained by a linear, energy interpolation of the cross sections at the group boundaries. Since each geometric region has constant material properties,

$$K_i^t(\vec{r}, \vec{r}') = \exp \left[ - \sum_{j=1}^n \Sigma_j^v \Delta s_j \right] \quad (6.5)$$

where  $\Delta s_j$  is the distance in the  $j$ th region traversed from

$\vec{r}$  to  $\vec{r}'$ , and

$\Sigma_j^v$  is the average total cross section for the region and group  $j$

This total cross section can be composed of two parts; one representing the composite material and another representing the hydrogen in the region as discussed in Section 4.3.

This photon attenuation kernel is used to calculate fluxes, and, in particular, to define the point monodirectional flux components which are used to represent the next order-of-scatter point source; i.e., equation 2.27

$$W_i^\Phi = \frac{w_i^s K_i^t(\vec{r}_{k-1}, \vec{r}_k)}{|\vec{r}_k - \vec{r}_{k-1}|^2 p_k^*(\vec{r}_k)}, \quad \bar{E}_i^\Phi = \bar{E}_i^s \quad (6.6)$$

where  $p_k^*(\vec{r}_k)$  is the value of the sampling function used to obtain  $\vec{r}_k$ .

This is merely a groupwise representation of equation 2.27, where the average energy of the particles in each group at the scattering point is assumed equal to the average energy of the previous source point:

$$W_i^\Phi = \int_{E_{i+1}}^{E_i} w_k^\Phi(E) dE$$

$$\begin{aligned} \bar{E}_i^\Phi &= \frac{1}{W_i^\Phi} \int_{E_{i+1}}^{E_i} W_k^\Phi(E) E dE \\ &\approx \bar{E}_i^s \end{aligned}$$

### 6.3 PHOTON SCATTERING

The photon scattering calculations performed in FASTER by subroutine SINGLE, use the Klein-Nishina equation for Compton scattering. All scattering is assumed to occur at the average group energy  $\bar{E}_i^\phi$  so that the ratio of the energies before and after scattering for group  $j$  is:

$$R_j = \frac{0.511}{0.511 + \bar{E}_j^\phi(1-\mu)} \quad (6.7)$$

where  $\mu = \hat{n}_k \cdot \hat{n}_i$ , the cosine of the scattering angle, and

$$\hat{n}_k = (\hat{r}_k - \hat{r}_{k-1}) / |\hat{r}_k - \hat{r}_{k-1}| \quad (6.8)$$

The scattered point source component—the scattered particles due to particles originally in group  $j$ —is then obtained from the Klein-Nishina equation:

$$\Delta W_j^s = W_j^\phi \frac{0.49875}{4\pi} N_e R_j^2 \left[ R_j + \frac{1}{R_j} - 1 + \mu^2 \right] \quad (6.9)$$

where  $N_e$  is the total electron density for the region in which the scattering occurs—including that due to the separate hydrogen.

The scattered contributions are grouped, according to the average scattered energies, to yield the final representation of the scattered point source:

$$\left. \begin{aligned} W_i^s &= \sum_{j \in J'} \Delta W_j^s \\ \bar{E}_i^\phi &= \frac{1}{W_i^s} \sum_{j \in J'} \Delta W_j^s R_j \bar{E}_j^\phi \end{aligned} \right\} \begin{array}{l} j \in J' \text{ if } E_j > R_j \bar{E}_j^\phi \geq E_{j+1} \\ \text{otherwise } \bar{E}_i^\phi \end{array} \quad (6.10)$$

where  $R_j \bar{E}_j^\phi$  is the average energy after scattering of particles originally in group  $j$ , and  $J'$  is the set of initial group indices  $j$ , for which this scattered energy is within the boundaries of group  $i$ .

The angular point source at the scattering point is evaluated only for discrete directions  $\hat{n}$ . Therefore, the angular dependence has been suppressed in this development.

### 6.4 NEUTRON CROSS SECTIONS

Neutron transport and scattering calculations utilize group averaged cross sections only. These can be obtained from various tabulations, e.g., Reference 14, which also discusses the averaging techniques.

The microscopic data for the  $i$ th element is supplied to the FASTER program by energy group  $j$  as:

$\sigma_{i,i}^t$  the average total cross section for element  $i$  and group  $j$

$\sigma_{j \rightarrow k, i}^l$  the  $l$ th Legendre expansion coefficient for element  $i$  of the differential elastic scattering cross section for transfer from group  $j$  to group  $k$ ,

$\sigma_{j \rightarrow k, i}^{ne}$  on isotropic, weighted, non-elastic transfer cross section

$$\sigma_{j \rightarrow k, i}^{na} = \sigma_{j \rightarrow k, i}^{inelastic} + 2 \cdot \sigma_{j \rightarrow k, i}^{(n-2n)} + \dots$$

The elastic scattering coefficients are assumed to contain the  $(2l + 1)$  factor associated with the Legendre series expansion.

The INSECT subroutine will transport correct these neutron cross sections, or it will remove the transport correction under conditions noted in the input instructions in Section 9.5. In these cross section manipulations, the first two Legendre expansion coefficients of the group-averaged elastic scattering cross sections are computed first:

$$\sigma_{j,i}^1 = \sum_k \sigma_{j-k,i}^1 \quad (6.12)$$

where  $\sigma_{j,i}^0$  is the average total elastic scattering cross section for the  $j$ th group of the  $i$ th element

$$\sigma_{j,i}^1 = 3\bar{\mu}_{j,i} \sigma_{j,i}^0$$

$\bar{\mu}_{j,i}$  is the average cosine of the scattering angle in the laboratory coordinate system for group  $j$  of the  $i$ th element.

The correction of the group averaged total cross section yields the group averaged transport cross section  $\sigma_{j,i}^{tr}$ :

$$\sigma_{j,i}^{tr} = \sigma_{j,i}^t - \bar{\mu}_{j,i} \sigma_{j,i}^0 = \sigma_{j,i}^t - \sigma_{j,i}^1/3 \quad (6.13)$$

The removal of this correction requires the equation:

$$\sigma_{j,i}^t = \sigma_{j,i}^{tr} + \sigma_{j,i}^1/3$$

If the total cross section is transport corrected, the corresponding correction to the group averaged elastic scattering cross section is applied to the in-group term of the isotropic Legendre expansion coefficient,  $\sigma_{j-k,i}^0$  only:

$$\sigma_{j-k,i}^{0,tr} = \sigma_{j-k,i}^0 - \sigma_{j,i}^1/3 \quad (6.14)$$

Special attention is given to hydrogen since its presence in large amounts is not adequately represented by the above group-to-group transfer cross sections. Hydrogen cross sections are supplied as group averaged totals  $\sigma_{j,i}^h$ . The scattering cross section is then assumed to equal the total cross section and the angular dependence of the scattering is treated correctly as described in Section 6.5.

Macroscopic cross sections, by composite material, are computed by energy group in a manner similar to that for photon cross sections. This includes the total cross section  $\Sigma_{i,m}^t$  and the scattering cross sections  $\Sigma_{j-k,m}^1$  and  $\Sigma_{j-k,m}^{ne}$  or the equivalent transport corrected cross sections.

Kinetic heating responses are computed by INSECT for each element and combined by composite material for use as heating conversion factors. Computed as average fractional energy loss cross sections, the groupwise equations are:

$$\sigma \left[ 1 - \left( \frac{E_{out}}{E_{in}} \right) \right]_{j,i} = \frac{2A_i}{(A_i+1)^2} \sigma_{j,i}^0 \left( 1 - \bar{\mu}_{j,i}^{c,m} \right) \quad (6.15)$$

where  $\frac{E_{in}}{E_{out}}$  is the energy of neutrons going into the elastic collision,  
 $\frac{E_{out}}{E_{in}}$  is the energy of neutrons coming out of this collision,

$$1 - \left( \frac{E_{out}}{E_{in}} \right)$$

is the average fractional energy loss  
 $\sigma_{j,i}^0$  is the total elastic scattering cross section as given by equation 6.12 for heavy elements ( $A_i > 1$ )  
=  $\sigma_{j,i}^h$  for hydrogen

$$\bar{\mu}_{j,i}^{c,m}$$

is the average scattering angle cosine in the center-of-mass coordinate system

$$\bar{\mu}_{j,i}^{c,m} = 0 \text{ for hydrogen}$$

$$\bar{\mu}_{j,i}^{c,m} \approx \frac{1}{A_i} \left[ \bar{\mu}_{j,i}^2 - 1 + \bar{\mu}_{j,i} \sqrt{\bar{\mu}_{j,i}^2 + \bar{\mu}_{j,i} - 1} \right] \text{ for heavy elements} \quad (6.16)$$

$$\bar{\mu}_{j,i}$$

is the average scattering angle cosine in the laboratory coordinate system.

## 6.5 NEUTRON ATTENUATION AND SCATTERING

The neutron attenuation kernel  $K_i^t(\vec{r}, \vec{r}')$  is handled in the same manner as the photon attenuation kernel. The only difference is that there is no energy interpolation required to define the average group cross sections.

The neutron source from scattering is developed in two parts. The first calculation yields the neutron source contribution from heavy element scattering:

$$\Delta W_{i'}^H = \sum_{j=1}^{m'} W_j^\phi \frac{1}{4\pi} \left[ \Sigma_{j \rightarrow i', m}^{ne} + \sum_{l=0} \Sigma_{j \rightarrow i', m}^l P_l(\mu) \right] \quad (6.17)$$

where  $m'$  denotes the material at the scattering point

$\mu$  is the cosine of the scattering angle,  $\mu = \vec{n}_k \cdot \vec{n}$

$P_l(\mu)$  is the  $l$ th Legendre polynomial, i.e., from Reference 16, pg. 308:

$$\left. \begin{aligned} P_0(\mu) &= 1 \\ P_1(\mu) &= \mu \\ P_l(\mu) &= \frac{1}{l!} [(2l-1)\mu P_{l-1}(\mu) - (l-1)P_{l-2}(\mu)] \end{aligned} \right\} \quad (6.18)$$

The neutron source for scattering from hydrogen is obtained only if  $\mu > 0$ , i.e., scattering in the laboratory coordinate system is restricted to angles  $< 90$  degrees. Then the scattered source contribution is computed as:

$$\Delta W_{i'}^h = 4\mu \rho^h \sum_{j \in J'} W_j^\phi \frac{\sigma_j^h}{4\pi} \quad (6.19)$$

where  $\mu^{c,m} = 2\mu^2 - 1$ ,  $d\mu^{c,m}/d\mu = 4\mu$ ,  $E_{out}/E_{in} = \frac{1+\mu^{c,m}}{2} = \mu^2$ ,

$j \in J'$  if  $E_{i'} > \mu^2 \overline{E}_{i'}^\phi \geq E_{i'+1}$  i.e., if the scattered energy  $\mu \overline{E}_{i'}^\phi$  is within the boundaries of group  $i'$ ; and  $\rho^h$  is the local hydrogen density.

These two components are then combined to yield the final value of the neutron scattered point source:

$$W_{i'}^s = \Delta W_{i'}^H + \Delta W_{i'}^h$$

$$\overline{E}_{i'}^s = \frac{1}{W_{i'}^s} \left[ \overline{E}_{i'}^n \Delta W_{i'}^H + 4\mu \rho^h \sum_{j \in J'} W_j^\phi \frac{\sigma_j^h}{4\pi} \mu^2 \overline{E}_{i'}^\phi \right] \quad (6.20)$$

where the  $\overline{E}_{i'}^n, i' = 1, 2, \dots$ , are input average scattered energies for neutron scattering from heavy elements.



## SECTION

### 7.0 FLUX ESTIMATION

The FASTER program computes multigroup neutron or photon flux components for arbitrarily located point, surface and/or volume detectors. The surface averaged fluxes are obtained for specified sections of surfaces which form boundaries of regions. Thus there may be two equivalent definitions of each surface detector, i.e., if two regions have the desired section of the surface as a common boundary. Volume averaged fluxes are obtained for specified regions of the geometry using the two special equations 2.55 and 2.56 for valid regions and constant material density regions respectively.

### 7.1 ANGULAR FLUX CONTRIBUTIONS

The random sampling technique discussed in Section 2.7 is used in integrating the angular dependence of the point sources to obtain surface and volume averaged flux components. Therefore, the individual contributions for all detector types have a similar form involving discrete directions. A full set of indices will be used in the following summary of flux estimation equations, i.e., the contribution to the  $i$ th energy group from the  $k$ th inner iteration of the  $l$ th outer iteration is given by:

a) point detector at  $\vec{r}$

$$\Delta\phi_{ijk}^*(\vec{\Omega}) = \left\{ \frac{w_{ijk}^s K_i^t (\vec{r}_{ik}, \vec{n})}{|\vec{r} - \vec{r}_{ik}|^2} \right\} \delta(\vec{\Omega} - \vec{\Omega}_o) \quad (7.1)$$

$$\vec{\Omega}_o = (\vec{r} - \vec{r}_{ik}) / |\vec{r} - \vec{r}_{ik}|$$

b) surface detector at a distance  $s$  from  $\vec{r}_{ik}$  along  $\vec{n}$

$$\Delta\phi_{ijk}^*(\vec{\Omega}) = \left\{ \frac{w_{ijk}^s K_i^t (\vec{r}_{ik}, \vec{r}_{ik} + s\vec{n}, \vec{n})}{A \cdot L \cdot |\vec{n}| \vec{\Omega}_i \cdot \vec{n} \cdot q^*(\vec{\Omega}_i)} \right\} \delta(\vec{\Omega} - \vec{\Omega}_o) \quad (7.2)$$

$$\vec{\Omega}_o = \vec{\Omega}_i$$

c) void volume detector at a distance  $s$  from  $\vec{r}_{ik}$  along  $\vec{\Omega}_i$

$$\Delta\phi_{ijk}^*(\vec{\Omega}) = \left\{ \frac{W_i^s K_i^t (\vec{r}_{ik}, \vec{r}_{ik} + s\vec{\Omega}_i) \Delta s}{V \cdot L \cdot q(\vec{\Omega}_i)} \right\} \delta(\vec{\Omega} - \vec{\Omega}_o). \quad (7.3)$$

$$\vec{\Omega}_o = \vec{\Omega}_i$$

d) non-void volume detector at a distance  $s$  from  $\vec{r}_{ik}$  along  $\vec{\Omega}_i$

$$\Delta\phi_{ijk}^*(\vec{\Omega}) = \left\{ \frac{W_i^s K_i^t (\vec{r}_{ik}, \vec{r}_{ik} + s\vec{\Omega}_i) \cdot [1 - \exp(-\Delta s \Sigma_i^v)]}{V \cdot L \cdot q(\vec{\Omega}_i)} \right\} \delta(\vec{\Omega} - \vec{\Omega}_o). \quad (7.4)$$

$$\vec{\Omega}_o = \vec{\Omega}_i$$

or in general:

$$\Delta\phi_{ijk}^*(\vec{\Omega}) = \Delta\phi_{ijk} \delta(\vec{\Omega} - \vec{\Omega}_o) \quad (7.5)$$

where  $\Delta\phi_{ijk}$  is the quantity in brackets  $\{ \}$  above. The angular dependence implied by the Dirac delta function is only a formality for use in defining angular moments of the flux as discussed in Appendix B. This multigroup representation of the point angular source is obtained, of course, for the discrete direction  $\vec{\Omega}_o$ . Surface and volume detectors may receive more than one flux contribution. This may occur if more than one discrete direction,  $\vec{\Omega}_i$ , ( $L > 1$ ) is used in the angular integration of the point angular sources. It will occur if the detectors are intersected more than once for each discrete direction.

## 7.2 SCALAR FLUXES AND COMPONENTS

The full set of indices, outer iteration  $i$ , inner iteration  $k$ , and energy group  $j$ , was introduced to simplify the discussion of the final equations for the various flux components. Unless a flux component is used in a subsequent equation, there will be no symbol introduced on the left side of the equation for this component.

The FASTER program actually includes a group collapse, if desired, to a coarser group structure for the flux edits. This will be ignored in the equations below since it involves nothing more than an inner summation. Thus, a summation over energy groups before any other summation is implicit. Also implicit is a summation over multiple intersections and/or discrete directions for surfaces and volumes.

### Scalar Number Flux (Sample Mean)

The total scalar flux in group  $j$  is obtained by an angular integration:

$$\begin{aligned} \Phi_j^o &= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^L \iint_{4\pi} \Delta\phi_{ijk} \delta(\vec{\Omega} - \vec{\Omega}_o) d\Omega \\ &= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^L \Delta\phi_{ijk} \end{aligned} \quad (7.6)$$

Other quantities are obtained in a similar manner. The index  $i$  is reserved for the energy dependence in all of the equations below.

### Sample Variance of Scalar Flux

$$V_j^2 = \frac{1}{n-1} \left[ \sum_{i=1}^n \left( \sum_{k=0}^L \Delta\phi_{ijk} \right)^2 - n(\Phi_j^o)^2 \right] \quad (7.7)$$

### Relative Error

$$E_j = \sqrt{\frac{V_j^2}{\Phi_j^o}} \quad (7.8)$$

Differential Number Flux

$$\phi_i^o / (E_i - E_{i+1}) \quad (7.9)$$

Cumulative Number Flux ( $E \geq E_{i+1}$ )

$$\sum_{j=1}^i \phi_j^o \quad (7.10)$$

Energy Flux

$$I_i^o := \frac{1}{n} \sum_{k=1}^n \overline{E_{ijk}} \Delta \phi_{ijk} \quad (7.11)$$

Average Group Energy

$$\bar{E}_i = I_i^o / \phi_i^o \quad (7.12)$$

Differential Energy Flux

$$I_i^o / (E_i - E_{i+1}) \quad (7.13)$$

Cumulative Energy Flux ( $E \geq E_{i+1}$ )

$$\sum_{j=1}^i I_j^o \quad (7.14)$$

Average Flux Between Outer Iterations  $n_1$  and  $n_2$ 

$$\frac{1}{n_2 - n_1} \sum_{i=n_1+1}^{n_2} \sum_{k=0}^n \Delta \phi_{ijk} \quad (7.15)$$

Flux Contribution from m<sup>th</sup> Fixed Source

$$\frac{1}{n} \sum_{i=1}^n \left\{ \left( \sum_{k=0}^n \Delta \phi_{ijk} \right) \text{ such that } \vec{r}_{i,o} \text{ is in the } m^{\text{th}} \text{ source} \right\} \quad (7)$$

Flux Contribution from i<sup>th</sup> Scattering Region

$$\frac{1}{n} \sum_{i=1}^n \sum_{k=1}^n \left\{ \Delta \phi_{ijk} \text{ such that } \vec{r}_{ijk} \text{ is in the } i^{\text{th}} \text{ region} \right\} \quad (7)$$

Flux Contribution from k<sup>th</sup> Order of Scatter

$$\frac{1}{n} \sum_{i=1}^n \Delta \phi_{ijk} \quad (7)$$

Azimuthally averaged Legendre moments of the angular flux can also be obtained from FASTER. In particular the zeroth moment is the scalar flux given by equation 7.6, first moment is the current. The equations used in obtaining these moments and the equations for externally manipulating these moments are summarized in Appendix B.

Length-of-flight moments of the scalar flux are also computed by FASTER. These moments are discussed in Appendix B.

### 7.3 FLUX CONVERSATIONS

Scalar fluxes are converted by group to more useful units using a linear interpolation in energy:

$$D_i = \phi_i^o \frac{f_i(\bar{E}_i - E_{i+1}) + f_{i+1}(E_i - \bar{E}_i)}{E_i - E_{i+1}} \quad (7)$$

where  $f_i$  is the point-wise energy dependent conversion factor or response function for  $i^{\text{th}}$  energy level.

The total response is also obtained with limits on its relative error:

$$D = \sum_{i=1}^D D_i \quad (7.20)$$

$$E_{\min} = \frac{1}{D} \left[ \sum_i E_i^2 D_i^2 \right]^{1/2} \quad (7.21)$$

$$E_{\max} = \frac{1}{D} \sum_i E_i D_i \quad (7.22)$$

The equation for  $E_{\max}$  is obtained by applying the Schwarz inequality to the covariance of the scalar fluxes in individual groups. Total responses are also obtained for each of the flux components and the angular and length-of-flight moments using equations like 7.19 and 7.20, above.

## SECTION

### 8.0 RANDOM SAMPLING TECHNIQUES

The preceding sections detailed the sampling functions in a formal manner only. This section describes the techniques used in the FASTER program for obtaining random discrete position vectors from the sampling functions and the considerations involved in defining these functions.

The development of sampling techniques for the FASTER program has proceeded historically from the approximation of the actual particle distributions to an approximation of optimal sampling functions. Both of these techniques are included in the FASTER program.

#### 8.1 GENERAL SAMPLING PROCEDURES

The technique used by FASTER in selecting discrete values of a random variable  $x$  involves the transformation:

$$p^*(x) dx = p(\xi) d\xi \quad (8.1)$$

where  $\xi$  is a random variable uniformly distributed on the open interval  $(0, 1)$

$$\begin{aligned} p(\xi) &= 1, \quad 0 < \xi < 1 \\ &= 0 \quad \text{otherwise} \end{aligned} \quad (8.2)$$

Various numerical techniques are available for randomly selecting discrete values of  $\xi$  from this distribution. A function subprogram RANNO is used in FASTER to obtain these discrete values of  $\xi$ .

The procedure for randomly selecting discrete values of  $x$  then reduces to solving the equation:

$$\int_{-\infty}^x p^*(x') dx' = \int_0^\xi p(\xi') d\xi' = \xi \quad (8.3)$$

Caution is exercised, therefore, to ensure that such a solution can be obtained with relative ease.

The form of  $p^*(x)$  frequently used in the FASTER program for one-dimensional sampling functions involves an exponential function. The probability density function for  $x$  is then written, for  $x$  on the interval  $(a, b)$ , as  $p^*(x; a, b, \alpha, x_0)$  with the following special cases:

- $p^*(x) = \delta(x - a)$  if  $a = b$
  - $p^*(x) = \frac{1}{b-a}$  if  $a < b$  and  $\alpha = 0$
  - $p^*(x) = \frac{\alpha}{c} \exp[-\alpha|x - x_0|]$  if  $a < b$  and  $\alpha \neq 0$
- (8.4)

where, for  $\alpha \neq 0$ ,

$x_0$  is the preferred value of  $x$ ,

$$a \leq x_0 \leq b$$

$$C = A + B$$

$$\begin{aligned} A &= \alpha \int_{x_0}^{x_0} \exp[-\alpha|x - x_0|] dx = \exp[-\alpha|x - x_0|] - 1 \\ B &= \alpha \int_{x_0}^b \exp[-\alpha|x - x_0|] dx = \exp[-\alpha|b - x_0|] - 1 \end{aligned} \quad (8.5)$$

The parameter  $\alpha$  is usually defined from a specification of the relative importance,  $P$ , of  $x_0$  as compared to the point,  $a$  or  $b$ , farthest away from  $x_0$ .

$$\begin{aligned} P &= 1/\exp[\alpha(x_0 - a)] \text{ or } P = 1/\exp[\alpha(b - x_0)] \text{, i.e.,} \\ \alpha &= -\frac{\ln P}{\max(x_0 - a, b - x_0)} \end{aligned} \quad (8.6)$$

The point  $x_0$  is "least preferred" if  $P$  is less than unity.

This probability density function is sampled using the function subprogram SAMPLE in the following manner:

- if  $a = b$   
set  $x = a$   
set  $p^*(x) = 1.0$  [actually equal to  $\delta(x - a)$ ]
- (8.7)

- if  $a < b$  and  $\alpha = 0$   
obtain  $\xi$  from RANNO  
solve equation 8.3,  
$$\xi = \int_a^x \frac{dx'}{b-a} = \frac{x-a}{b-a}$$
  
i.e.,  $x = a + \xi(b-a)$   
$$p^*(x) = 1/(b-a)$$
- (8.8)

- if  $a < b$  and  $\alpha \neq 0$   
obtain  $\xi$  from RANNO  
solve equation 8.3,  
i.e.,  $\xi = \frac{\alpha}{c} \int_a^x \exp[-\alpha|x - x_0|] dx$   
and  $\xi = \frac{1}{C} \left\{ A - \alpha \int_x^{x_0} \exp[-\alpha|x' - x_0|] dx' \right\}$   
or  $\xi C - A = -\left\{ \exp[-\alpha(x - x_0)] - 1 \right\}$
- (8.9)

- if  $\xi < A/C$ , then  $x < x_0$   
and  $\xi = \frac{1}{C} \left\{ A - \alpha \int_x^{x_0} \exp[-\alpha|x' - x_0|] dx' \right\}$   
or  $\xi C - A = -\left\{ \exp[-\alpha(x - x_0)] - 1 \right\}$
  - if  $\xi > A/C$ , then  $x > x_0$   
and  $\xi = \frac{1}{C} \left\{ A + \alpha \int_{x_0}^x \exp[-\alpha|x' - x_0|] dx' \right\}$   
or  $\xi C - A = \left\{ \exp[-\alpha(x - x_0)] - 1 \right\}$
- (8.10)

i.e., let  $\delta = \frac{\xi - A/C}{|\xi - A/C|}$   $\left\{ \begin{array}{l} +1 \text{ for case 1 above} \\ -1 \text{ for case 2 above} \end{array} \right.$

then  $\exp[\delta\alpha(x - x_0)] = 1 + \delta(\xi C - A)$

let  $D = 1 \pm \delta(\xi C - A)$ ,  $\cdot \cdot \cdot$

then  $x' = x_0 + \frac{\delta \ln D}{\alpha}$

and  $p^*(x) = \frac{\alpha D}{C}$

(8.9)

## 8.2 SPATIAL SAMPLING IN THE FIXED SOURCE COORDINATE SYSTEM

The FASTER program includes both the random selection of initial position vectors in a source centered coordinate system and a sampling procedure utilizing built-in importance functions. However, this latter sampling technique is limited to volume sources. Thus, the first technique is required for problems involving point or line sources. Surface sources can be treated by either technique since they are readily approximated by "thin" volumes.

Random selection of a point in the source coordinate system is performed by subroutine PSTAR. The first step in sampling in the fixed source coordinate system is the random selection of just one of the sources. The following equation defines the total sampling function:

$$p_o^*(\vec{r}) = \sum_{i=1}^{I-1} p_i^* q_i^*(\vec{r}) \quad (8.10)$$

where  $p_i^*$  is the input relative importance of the ith source

$$\sum_{i=1}^{I-1} p_i^* = 1 \quad (8.11)$$

$$\left. \begin{aligned} q_i^*(\vec{r}) &= \text{if } \vec{r} \text{ is in } V_i, \text{ the } i\text{th source volume} \\ q_i^*(\vec{r}) &= 0 \text{ if } \vec{r} \text{ is not in } V_i \\ \iiint q_i^*(\vec{r}) dV &= 1 \end{aligned} \right\} i = 1, 2, \dots \quad (8.12)$$

The appropriate volume is then selected using a random number  $\xi$  from RANNO

$$\xi = \iiint p_o^*(\vec{r}) dV = \sum_{i=1}^{I-1} p_i^* + p_I^* \iiint_{V_i} q_i^*(\vec{r}) dV \quad (8.13)$$

where  $f$  is some fraction, less than 1.0, of the volume  $V_i$ . Thus, the point  $\vec{r}$  is in the fixed source volume  $i$  for which

$$\sum_{i=1}^{I-1} p_i^* < \xi \leq \sum_{i=1}^I p_i^* \quad (8.14)$$

Having selected the fixed source, then each of the spatial variables  $v_1$ ,  $v_2$ ,  $v_3$  is obtained using the sampling function SAMPLE, described in Section 8.1

$$q_i^*(\vec{r}) dV = \prod_{j=1}^3 p_j^* \left( v_j; v_j^{\min}, v_j^{\max}, v_j^o, a_j \right) \frac{dv_j}{v_j^{\max} - v_j^{\min}} \quad (8.15)$$

where  $n = 1$  and 2 for cylindrical and spherical geometries, respectively

$= 0$  for rectangular geometries

$v_j^{\min}$  is the minimum value of the jth variable

$v_j^{\max}$  is the maximum value of the jth variable

$v_j^o$  is the preferred value of the jth variable

$$a_j = \frac{-\ln p_j}{\max(v_j^o - v_j^{\min}, v_j^{\max} - v_j^o)} \quad .$$

$p_j$  is the input relative importance of  $v_j^o$

The  $1/v_1^n$  factor is introduced to simplify the difference in differential volume elements for the different source geometries:

$$\frac{dV}{v_1^n} = \left( v_1^n dv_1 dv_2 dv_3 \right) / v_1^n = dv_1 dv_2 dv_3 \quad (8.16)$$

The rectangular components of  $\vec{r}_o$  are then obtained by the variable transformations discussed in Section 5.0. Since the point  $\vec{r}_o$  defined by  $(v_1, v_2, v_3)$  is in the  $i$ th volume, the value of  $p_o^*(\vec{r}_o)$  is:

$$p_o^*(\vec{r}_o) = p_i^* v_i^*(\vec{r}_o) \quad (8.17)$$

### 8.3 SOURCE SAMPLING USING A PSEUDO SPHERICAL SOURCE

Another sampling procedure was developed for fixed volume sources and attempts to minimize the variance associated with the source point selection. The function approximates equation 2.36 by a sampling function with a separable angular and spatial dependence.

$$p_o^*(\vec{r}) dV = \frac{u^*(\vec{\Omega}) v^*(\vec{s}; \vec{\Omega})}{s^2} (s^2 ds d\Omega) \quad (8.18)$$

where a transformation to a spherical coordinate system about a preferred point  $\vec{r}_p$ , e.g., a point detector, has been performed. The variables involved in this sampling function are shown in Figure 5. The random selection procedures are incorporated in subroutine SPHERE.

#### Angular Dependence

The  $z$ -axis of the spherical coordinate system is directed towards the center of the sources, i.e., towards  $\vec{r}_c$ , a point in the center of a sphere, of radius  $R$ , which encloses all the fixed sources. The random selection of the direction vector  $\vec{\Omega}^*$  in this rotated coordinate system uses the sampling functions:

$$v^*(\vec{\Omega}) d\Omega = p^*(\mu'; \mu^{\min}, 1, 1, a_{\mu'}) p^*(\theta; -\pi, \pi, 0, a_{\theta}) d\mu' d\theta' \quad (8.19)$$

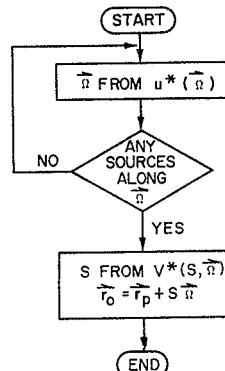
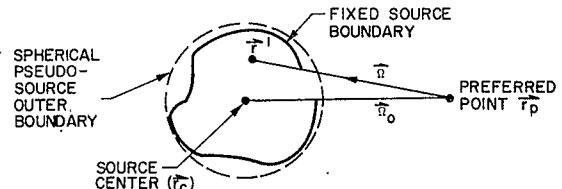


Figure 5. Optimal Fixed Source Sampling

i.e., discrete values of  $\mu'$  and  $\theta'$  are obtained from the function subprogram SAMPLE. The azimuthal angle varies over all possible values. The minimum value of the cosine of the polar angle is defined as:

$$\begin{aligned} \mu_{\min}' &= -1 \text{ if } R > R_0 = \left| \vec{r}_p - \vec{r}_c \right| \\ &= [R_0^2 - R^2]^{1/2} / R_0 \text{ if } R \leq R_0 \end{aligned} \quad (8.20)$$

i.e., the angular variables are limited to directions which intercept the sphere enclosing the fixed sources. The parameters  $a_{\mu'}$  and  $a_{\theta'}$  are defined from the input relative importances:

$$\begin{aligned} p_{\mu'} &= \frac{-\ln \rho_{\mu'}}{1 - \mu_{\min}'} \quad p_{\mu'} = \frac{\text{importance of } \mu' = 1}{\text{importance of } \mu' = \min'} \\ a_{\mu'} &= \frac{-\ln \rho_{\mu'}}{\pi} \quad p_{\theta'} = \frac{\text{importance of } \theta' = 0}{\text{importance of } \theta' = \pm \pi} \end{aligned} \quad (8.21)$$

The values of  $\mu'$  and  $\theta'$  obtained through this sampling procedure define the components of a direction vector in a rotated coordinate system

$$\begin{aligned} \vec{\Omega}' &= c'_1 \vec{i}' + c'_2 \vec{j}' + c'_3 \vec{k}' \quad c'_1 = \sqrt{1 - \mu'^2} \cos \theta' \\ &\quad c'_2 = \sqrt{1 - \mu'^2} \sin \theta' \\ &\quad c'_3 = \mu' \end{aligned} \quad (8.22)$$

These components are then transformed into equivalent values in the geometry coordinate system:

$$\vec{\Omega} = c'_1 \vec{i} + c'_2 \vec{j} + c'_3 \vec{k} \quad \text{where } c_i = \sum_{i=1}^3 c'_i R_{i,i} \quad i = 1, 2, 3 \quad (8.23)$$

The rotation matrix  $R_{i,j}$  is obtained as shown in Figure 4 for the rotation angles:

$$\begin{aligned} \phi_r &= \cos^{-1}(c_3^0) \\ \theta_r &= \tan^{-1}(c_2^0/c_1^0) \end{aligned} \quad (8.24)$$

where  $\vec{\Omega}_0 = c_1^0 \vec{i} + c_2^0 \vec{j} + c_3^0 \vec{k} = (\vec{r}_c - \vec{r}_p)/k$ , the unit vector directed towards the center of the sphere enclosing the fixed sources.

#### Spatial Dependence

The spatial sampling function incorporated in subroutine SPHERE is defined for all direction vectors  $\vec{\Omega}$  even if no sources lie along  $\vec{r}_p + s\vec{\Omega}$  for  $s \geq 0$ . If no source is intercepted along this ray, the outer iteration will yield a zero result. To ensure an adequate number of outer iterations, discrete directions are obtained by random sampling until one of the defined rays intercepts a fixed source. The misses are counted internally and factored into the flux averages; thus if  $n'$  is the requested number of iterations,  $n \geq n'$  iterations will be performed until  $n'$  non-zero contributions are obtained.

The spatial sampling function is defined as: for  $s$  in the  $i$ th region along  $\vec{\Omega}$  as:

$$v^*(s|\vec{\Omega}) = \frac{A_i \exp[-a_i(s - s_i)]}{\sum_i A_i [1 - \exp(-a_i(s_{i+1} - s_i))]} \quad (8.25)$$

where  $A_i = 0$ , if there is no source in the region

$$A_i = I_i P_i \left( \vec{r}_p + (s_i + \epsilon) \vec{\Omega} \right) \exp \left[ - \int_{s_i}^{s_{i+1}} \sum_{j=0}^i \left( \vec{r}_p + s' \vec{\Omega} \right) d s' \right] \quad (8.26)$$

if there is a source,  $i'$ , in the region

$$a_i = \beta_s \sum_{j=0}^i \left( \vec{r}_p + \vec{s}_j \vec{\Omega} \right) + \frac{2}{s_{i+1} - s_i} \quad \ln \left[ \frac{P_{i'} \left[ \vec{r}_p + (s_i + \epsilon) \vec{\Omega} \right]}{P_{i'} \left( \vec{r}_p + \vec{s}_{i+1} \vec{\Omega} \right)} \right] \quad (8.27)$$

$$\bar{s}_i = (s_i + s_{i+1})/2$$

$P_{i+1} (\vec{r}_p + (s_i + \epsilon) \hat{\Omega})$  is the normalized spatial source density just inside the region,  
i.e.  $\epsilon = \frac{s_{i+1} - s_i}{100}$

$P_p (\vec{r}_p + \bar{s}_i \hat{\Omega})$  is the normalized spatial source density half way through the region

$I_{i+1}$  is the total source intensity (Mev/sec),

$\hat{\tau}_o$  is an average energy group for the sources (described below),

$\beta_s$  is an input adjustment factor for the material attenuation importance, and a two point exponential curve fit of the spatial source density with distance has been used

The average source group index,  $\hat{\alpha}_o$ , corresponds to an average source energy  $\bar{E}_p$  computed for the preferred point  $\vec{r}_p$  by subroutine GROUP as:

$$\hat{\alpha}_o > \bar{E}_p \geq \hat{\alpha}_o$$

$$\text{where } \bar{E}_p = \frac{\sum_i G_i \bar{E}_i^o}{\sum_i G_i}, \quad G_i = \frac{g_i n_i^o \exp[-\eta_i]}{z_i^v} (1 + b_i \eta_i) \quad (8.28)$$

where  $g_i$  is an input group importance e.g., a flux-to-dose conversion factor

$b_i$  is an input linear buildup coefficient,

$n_i^o$ ,  $\bar{E}_i^o$  define the spectrum of the first source encountered on the line from  $\vec{r}_c$  to  $\vec{r}_p$  (the source closest to  $\vec{r}_c$ ),

$\eta_i$  is the number of mean free paths from the edge of this source nearest to the preferred point  $\vec{r}_p$ ,

$z_i^v$  is the cross section for the region over which the source is superimposed ( $\equiv 1.0$  if the region is void),

The spatial function is sampled by obtaining a random number  $\xi$  from RANNO and then solving the equivalent of equation 8.3.

$$\xi = \int_0^s v^*(s'; \hat{\Omega}) ds' \\ = \sum_{i=1}^{k-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp \left[ -\alpha_i (s_i + 1 - s_i) \right] \right] + \int_{s_k}^s A_k \exp \left[ -\alpha_k (s' - s_k) \right] ds' \\ \sum_{i=1}^{k-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp \left[ -\alpha_i (s_{i+1} - s_i) \right] \right] \quad (8.29)$$

$$s = s_k - \frac{1}{A_k} \ln \left[ 1 - \frac{\alpha_k}{A_k} (\xi A_{>o} - A_{}) \right] \\ v^*(s; \hat{\Omega}) = \left[ A_k - \alpha_k \left( \xi A_{>o} - A_{} \right) \right] / A_{>o} \quad (8.30)$$

where  $A_{>o} = \sum_{i=1}^{k-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp \left[ -\alpha_i (s_{i+1} - s_i) \right] \right]$ , the denominator of equation 8.29  
 $A_{} = \sum_{i=1}^{k-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp \left[ -\alpha_i (s_{i+1} - s_i) \right] \right]$ , that part of the numerator of equation 8.29 up to the region containing the source point

Thus the discrete position reactor is given by:

$$\vec{r}_o = \vec{r}_p + s \hat{\Omega}$$

$$\text{with } p_o^*(\vec{r}_o) = \frac{v^*(\hat{\Omega}') v^*(s; \hat{\Omega})}{s^2} \quad (8.31)$$

#### 8.4 SEPARATION OF VARIABLES FOR SCATTERING POINT SELECTIONS

Application of the procedures discussed in Sections 8.3 and 8.4 reduces the representation of the fixed source(s) to a single point source. The random sampling techniques for selecting the  $k$ th scattering point  $\vec{r}_k$  therefore involve a point angular source at  $\vec{r}_{k-1}$  and a preferred point at  $\vec{r}_p$  as shown in Figure 6. For point flux calculations in source and/or scattering volumes, the preferred point  $\vec{r}_p$  is the detector point  $\vec{r}$ .

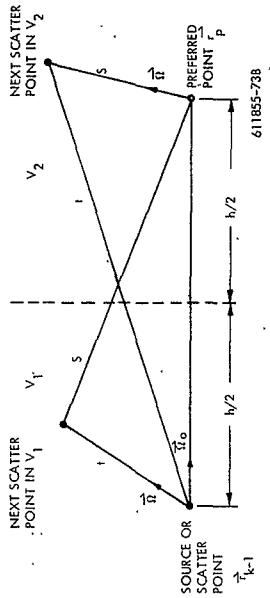


Figure 6. Optimal Scattered Source Sampling

Surface and volume averaged flux calculations utilize the preferred point  $\vec{r}_p$  to locate in a general manner the area in space where the detectors are located. For these calculations the point  $\vec{r}_p$  is assumed to be at the center of a sphere of radius  $R_p$  which contains all the detectors. It is introduced to permit a more comprehensive treatment of the spatial importance of scattering points.

The first technique used in selecting the scattering point is to reduce equation 2.37 to a more simple form by using a representative energy group and by defining the sampling function as a separable function of direction and distance:

$$p_k(\vec{r}') dV = \frac{q^*(\vec{\Omega}') u^*(x) dV}{x^2} \quad (8.32)$$

where both  $q^*(\vec{\Omega}')$  and  $u^*(x)$  depend on  $\hat{\gamma}$ , the index of the energy group corresponding to an average group energy determined from the flux estimation performed for the source point  $\vec{r}_{k-1}$ :

$$E_{\hat{\gamma}} > \bar{E} \geq E_{\hat{\gamma}+1}$$

$$\bar{E} = \frac{\sum_{\text{detectors}} \sum_{i=1}^n \Delta \phi_{i,j,k-1} g_i \eta_j b_i E_i^3}{\sum_{\text{detectors}} \sum_{i=1}^n \Delta \phi_{i,j,k-1} g_i \eta_j b_i} \quad (8.33)$$

The separation of the sampling function into two components permits the selection of a discrete direction vector  $\vec{\Omega}'$  with a cursory treatment of geometric effects. The distance is then selected from the second component using a more detailed treatment of the geometric effects.

It should be noted that the angular integration of the point sources used in obtaining surface and volume averaged fluxes involves a repetitive sampling of the angular part of this sampling function i.e., discrete directions  $\vec{\Omega}'$  are randomly selected from the angular sampling function  $q^*(\vec{\Omega}')$ .

For surface and volume flux calculations the discrete direction and distance defining  $\vec{r}_{k-1}$  are obtained in a spherical coordinate system which is always centered at the source point  $\vec{r}_{k-1}$ , i.e.,  $dV = r^2 d\Omega$ . A spherical coordinate system is also used for individual point detectors, except the origin can be located at the detector point  $\vec{r}_p$  with  $dV = r^2 d\Omega$ .

The considerations involved in selecting the origin for point detectors incorporate the fact that there are two singularities in the sampling function,  $1/r^2$  from the source point and  $1/r^2$  from the detector. Both singularities must be accounted for while also treating material distributions, etc. The singularity at the selected coordinate system origin is removed by the differential volume element. A very simple technique for the second singularity involves the definition:

$$u^*(\vec{r}; \Omega) dV = \left\{ p_1^* u_1^*(\vec{r}) + p_2^* u_2^*(\vec{r}) \right\} dV \quad (8.34)$$

where

$$u_1^*(\vec{r}) = 0 \text{ for points in } V_2$$

$$u_2^*(\vec{r}) = 0 \text{ for points in } V_1$$

$V_1$  is the volume on the source side of a plane perpendicular to and bisecting the line between the source and detector, and  $V_2$  is the volume on the other side of this plane.

The appropriate volume, and coordinate system origin, is randomly selected using the criterion that the relative importance of scattering in each of the volumes is proportional to the mean free paths encountered in the volumes along the line between the source and detector points:

$$p_1^* = \eta_1/\eta, \text{ the assumed importance of volume } V_1$$

$$p_2^* = \eta_2/\eta, \text{ the assumed importance of volume } V_2$$

$$\eta = \eta_1 + \eta_2$$

where  $\eta_1 = \int_0^{h/2} \sum_{\vec{r}} (\vec{r}_{k-1} + s^1 \vec{\Omega}_k) ds^1$  , an average number of mean free paths in volume  $V_1$

$$\eta_2 = \int_0^{h/2} \sum_{\vec{r}} (\vec{r}_p - s^1 \vec{\Omega}_k) ds^1 \text{ , an average number of mean free paths in volume } V_2 \quad (8.35)$$

The appropriate volume is selected using a random number  $\xi$  on  $(0, 1)$ .

$$V_1 \text{ if } \xi = \iiint p_k^*(\vec{r}) dV = f \cdot p_1^* \text{ i.e. } \xi < p_1^*$$

$$V_2 \text{ if } \xi = p_1^* + f' p_2^* \text{ i.e. } \xi > p_1^* \quad (8.36)$$

where  $f$  and  $f'$  are fractions. If  $p_1^*$  is zero, the point  $\vec{r}_p$  is automatically selected as the origin and the definition of the  $u_2^*(\vec{r})$  is extended over all space. Similarly if  $p_2^*$  is zero, the point  $\vec{r}_{k-1}$  is selected as an origin with the definition of  $u_1^*(\vec{r})$  being extended to all space points.

This procedure does not eliminate the consideration of the second singularity. It does, however, restrict it to values greater than  $4/(\vec{r}_{k-1} - \vec{r}_p)^2$ , which is sufficient to remove the major difficulty.

## 8.5 DIRECTION TO SCATTERING POINT

### Direction Vector Sampling-Source Coordinate System

There is a special form of  $q^*(\vec{\Omega})$  available for the case of  $k=1$ ; i.e., when  $\vec{r}_{k-1}$  is the original point in the fixed source. Since the capability of specifying angular sources is still permitted, these directions may require random sampling in the source coordinate system. In particular, a monodirectional source will always require this procedure.

The spatial part of the sampling function is then defined for all space, the source point is the origin, and the direction vector in the rotated source coordinate system is obtained by subroutine QSTAR using SAMPLE.

$$q^*(\hat{\Omega}) d\Omega^* = \prod_{i=4}^5 p^*(v_i; v_i^{\min}, v_i^{\max}, v_i^0, \alpha_i) dv_i \quad (8.37)$$

where  $v_4$  is the azimuthal angle,  $v_5$  is the cosine of the polar angle, and the other quantities have the same meaning as in equation 8.15. This is analogous to the procedure used for spatial sampling in the source coordinate system as discussed in Section 8.2. The selected variables  $v_4$  and  $v_5$  define the direction vector through a rotation procedure similar to that used in previous discussions.

#### Direction Vector Selection - General

The general procedure for selecting discrete directions attempts to include the importance of the scattered direction, as measured from the direction before scattering, and as measured from the unit vector towards the preferred point.

The unnormalized angular dependence for the scattering angle incorporates an assumed exponential variation with the cosine of the scattering angle:

$$f_a(\hat{\Omega} \cdot \hat{\Omega}_{k-1}) = f_a(\mu) = \exp(-\alpha_s \bar{a}_s \mu)$$

where  $\bar{a}_s = (\rho_h^h \alpha_1^h \alpha_2^h - \sum_{j=1}^4 \alpha_j^H) / (\rho_h^h + \sum_{j=1}^4)$  in subroutine SOBER

$$\bar{a}_s = (\eta^h \alpha_1^h + (\eta - \eta^h) \alpha_2^H) / \eta \quad \text{in subroutine SOLVER} \quad \left. \right\} \quad (8.38)$$

$\eta^h$  is the number of mean free paths in hydrogen between  $\hat{\Omega}_{k-1}$  and  $\hat{\Omega}_p$

$\rho_h^h$  is the local hydrogen density

$\sigma_t^h$  is the total microscopic hydrogen cross section for group  $i$

$\sigma_t^h$  is the total cross section-except hydrogen-for group  $i$

$$\alpha_i^h = -\frac{1}{2} \ln \rho_i^h$$

$\rho_i^h$  is an input relative importance of forward-to-backward scattering for the general group for hydrogen

$$\alpha_i^H = -\frac{1}{2} \ln \rho_i^H$$

$\rho_i^H$  is the forward-to-backward scatter importance for heavy elements

$\alpha_s$  is an input parameter used to increase or decrease this angular effect.

The unnormalized angular dependence with respect to the preferred point--the effect of the material attenuation--is also assumed to vary exponentially.

$$f_p(\hat{\Omega} \cdot \hat{\Omega}_p) = \exp(\alpha_p \bar{a}_p \hat{\Omega} \cdot \hat{\Omega}_p) \quad (8.39)$$

where-

$\alpha_p$  is an input adjustment factor,

$\bar{a}_p$  is an internally computed number

The latter number  $\bar{a}_p$  is obtained from a one velocity approximation of the scattering point importance which, by neglecting unimportant constants has the forms:

$$f_p(\mu) = \int_0^{s_{\max}} \frac{\exp[-\bar{s}(s+t)]}{s^2 t^2} s^2 ds \quad (8.40)$$

where

$$\Sigma = \eta/h \quad \eta = \eta_1 + \eta_2$$

$$s_{\max} = h/2\mu, \mu > 0$$

$$= \infty, \mu \leq 0$$

$$h = |\vec{r}_{k-1} - \vec{r}_p|$$

$$\text{Thus } f_p(\mu) \Big|_{\mu=1} = \int_0^{h/2} \frac{\exp[-\Sigma s]}{(h-s)^2} ds = \frac{e^{-\Sigma h}}{h} = \frac{e^{-\eta}}{h}.$$

$$\text{and } f_p(\mu) \Big|_{\mu=-1} = \int_0^{\infty} \frac{\exp[-\Sigma h - 2\Sigma s]}{(h+s)^2} ds \leq \frac{e^{-\eta}}{h} [12\eta]^{-1/2}$$

where Schwarz' inequality was used.

$$\text{Letting } f_p(1)/f_p(-1) = \exp[2\bar{\alpha}_p]$$

$$\text{Then } \exp[2\bar{\alpha}_p] \geq \sqrt{12\eta} \quad \text{or} \quad \bar{\alpha}_p \geq \frac{1}{4} \ln(12\eta) \quad (8.41)$$

The equality is used in the FASTER program.

The effect of this parameter has been examined for  $\mu = 1$  by computing:

$$\frac{\partial f(\mu)}{\partial \mu} \Big|_{\mu=1} = \left[ \frac{\partial}{\partial \mu} \int_0^{h/2\mu} \frac{\exp[-\Sigma(t+t)]}{t^2} dt \right]_{\mu=1} = \frac{e^{-\eta}}{h} \left( \frac{\eta}{2} - \frac{1}{3} \right)$$

where Leibnitz' rule was applied,  $\mu$  was then set equal to 1, and then the integration performed.

If the true behavior at  $\mu = 1$  was exponential, then

$$\bar{\alpha}_p = \frac{\eta}{2} - \frac{1}{3} \quad (8.42)$$

It is noted that for large values of  $\eta$ , this is a stronger angular dependence than that given by equation 8.41.

The total angular dependence of the sampling function is approximated by various combinations of the above functions. The angular variables are obtained using the function subprogram SAMPLE.

a) Towards Detector Most Important

$$q^*(\vec{\Omega}) = p^*(\mu'; -1, 1, -1, \alpha_s \bar{\alpha}_p) \\ p^*(\theta \cdot \theta_0; -\pi, \pi, 0, -\alpha_s \bar{\alpha}_s \sqrt{1-\mu'^2} \sqrt{1-\bar{\alpha}_{k-1} \cdot \bar{\alpha}_p^2}) \quad (8.43)$$

where

$$\vec{\Omega}_r = \vec{\Omega}_p$$

$$\mu' = \vec{\Omega}' \cdot \vec{\Omega}_r$$

and  $\theta_0$  is the azimuthal angle between the  $\vec{t}'$  axis and the projection of  $\vec{\Omega}_{k-1}$  in the  $x'-y'$  plane and where the rotated coordinate system has  $\vec{k}' = \vec{\Omega}_p$ .

b) Original Direction Most Important

$$q^*(\vec{\Omega}') = p^*(\mu'; -1, 1, -1, \alpha_s \bar{\alpha}_s) \\ p^*(\theta \cdot \theta_0; -\pi, \pi, 0, -\alpha_s \bar{\alpha}_p \sqrt{1-\mu'^2} \sqrt{1-\bar{\alpha}_{k-1} \cdot \bar{\alpha}_p^2}) \quad (8.44)$$

where

$$\vec{\Omega}_r = \vec{\Omega}_{k-1}$$

$$\mu' = \vec{\Omega}' \cdot \vec{\Omega}_r$$

$\theta_0$  is the azimuthal angle between the  $\vec{t}'$  axis and the projection of  $\vec{\Omega}_p$  in the  $x'-y'$  plane, and where the rotated coordinate system has  $\vec{k}' = \vec{\Omega}_{k-1}$ .

Only the latter technique, equation 8.44, is generally applicable to individual point detectors since the cosine of the polar angle  $\mu'$  used in the first technique, equation 8.43, applies only if the source point  $\vec{r}_{k-1}$  is selected as the coordinate system origin. As a matter of fact, both of these techniques neglect the effect of the scattering angle at  $\vec{r}'$ , the next scattering point being selected.

The effects of the scattering angle  $\cos \mu'$  at the next scattering point involves the use of the unnormalized approximation:

$$f(\mu') \Big|_{\mu=1} = \int_0^{h/2\mu} \frac{\exp [a_s \bar{a}_s \mu']}{h^2} ds \Bigg|_{\mu=1} = \frac{\exp [a_s \bar{a}_s]}{h}$$

$$f(\mu') \Big|_{\mu=-1} = \int_0^{h/2\mu} \frac{\exp [a_s \bar{a}_s \mu']}{h^2} ds \Bigg|_{\mu=-1} = \frac{\exp [-a_s \bar{a}_s]}{h}$$

Assuming an exponential variation  $\exp [\bar{a}\mu]$  for intermediate angles  $\mu$  implies

$$\bar{a} = \bar{a}_s \quad (8.45)$$

This parameter can be scaled through the input number  $a_s$ .

Combining this variation with that for material attenuation effects, equation 8.41 or 8.42 yields a third technique for selecting the direction vector with more or less equal applicability to either selected origin in point flux calculations:

c) Combination

$$q_i^*(\vec{\Omega}') = p^*(\mu'; -1, 1, -1, a_s \bar{a}_s - a_p \bar{a}_p) p^*(\theta', -\pi, \pi, 0, 0) \quad (8.46)$$

where  $\vec{\Omega}_r = \vec{\Omega}_p$

$$\mu' = \vec{\Omega}' \cdot \vec{\Omega}_p$$

$\theta'$  is a uniformly distributed azimuthal angle

All of the above angular sampling techniques, equations 8.43, 8.44, and 8.46, yield polar angle cosines,  $\mu'$ , and azimuthal angles  $\theta'$  in a rotated coordinate system. They require application of a rotation to be reduced to equivalent values in the geometry coordinate system. The  $z'$ -axis of the rotated coordinate system is  $\vec{\Omega}'$  as indicated above. The rotations are performed by previously described techniques.

### 8.6 DISTANCE-TO-SCATTERING POINT SELECTION

The spatial sampling function can contain a more detailed treatment of the effects of material distributions since it is defined for a fixed direction. For this discussion, the following notation will be adopted:

$\vec{r}_o$  is the position vector of the selected origin, either  $\vec{r}_{k-1}$  or  $\vec{r}_p$

$\vec{r}_b$  is the position vector of the other discrete point

$h$  is the distance from  $\vec{r}_o$  to  $\vec{r}_b = |\vec{r}_b - \vec{r}_o|$

$\vec{\Omega}$  is the unit direction vector from  $\vec{r}_o$  to  $\vec{r}_b$   
 $= (\vec{r}_b - \vec{r}_o)/h$

$\vec{\Omega}'$  is the direction vector selected from the angular sampling function.

$\mu_o$  is the cosine of the angle between  $\vec{\Omega}$  and  $\vec{\Omega}'$   
 $= \vec{\Omega} \cdot \vec{\Omega}'$

$s_f$  is the distance from  $\vec{r}_o$  along  $\vec{\Omega}'$  to the outer boundary of the geometry,

$s_m$  is the maximum distance along  $\vec{\Omega}'$  which can be considered for a scattering point.

$s_m = \min(s_f, h/2\mu_o)$  if sampling is restricted to the half space around  
 $\vec{r}_o$  and if  $\mu_o > 0$

$s_m = s_f$  in all other situations.

$\Sigma_{\hat{\gamma}}^V(s)$  is the total cross section at a distance  $s$  from  $\vec{r}_0$  for the energy group  $\hat{\gamma}$ , also denoted by  $\Sigma_{\hat{\gamma}}^V(s')$  where  $s'$  indicates the composite material located at this distance.

$\alpha$  is an adjustment factor for material attenuation importance along the first leg of the scattering triangle (from  $\vec{r}_0$ ).

$\beta$  is a similar adjustment factor for the second leg of this triangle (from  $\vec{r}_b$ ).

#### Exponential Transformation

There are two techniques available in subroutine USTAR for selecting the distance to the scattering point. The first is the "exponential transformation". There is no essential difficulty with using this technique for point detectors since the singularity is removed by the procedures used in selecting the half space.

The exponential transformation is written in the form:

$$\alpha^*(s) = \frac{-\frac{d}{ds} \exp \left[ -\int_0^s \Sigma_{\hat{\gamma}}^V(s') [\alpha - \beta \mu(s')] ds' \right]}{1 - \exp \left[ -\int_0^s \Sigma_{\hat{\gamma}}^V(s') [\alpha - \beta \mu(s')] ds' \right]} \quad (8.47)$$

where

$\mu(s)$  is the cosine of the scattering angle to  $\vec{r}_b$  at a distance  $s$  from

$\vec{r}_0$ , i.e., at  $\vec{r}' = \vec{r}_0 + s \hat{\gamma}$

$$\mu(s) = \frac{\mu_0 h - s}{t}$$

$t$  is the length of the second leg of the scattering triangle

$$= \sqrt{h^2 + s^2 - 2\mu_0 s h}$$

The extent to which this sampling function approximates equation 2.37 for the fixed direction  $\hat{\Omega}'$  can be determined by performing the differentiation

$$\alpha^*(s) = \frac{1}{C} \Sigma_{\hat{\gamma}}^V(s) [\alpha - \beta \mu(s)] \exp \left[ -\int_0^s \Sigma_{\hat{\gamma}}^V(s') [\alpha - \beta \mu(s')] ds' \right] \quad (8.48)$$

where  $C$  is the denominator of equation 8.47. The coefficient in front of the exponential should be approximating a term  $t^{-2} d \bar{\Sigma} [\epsilon, \mu(s)] / d\Omega$  except for a normalizing constant. Its actual form is:

$$\begin{aligned} \Sigma_{\hat{\gamma}}^V(s) [\alpha - \beta \mu(s)] &= \Sigma_{\hat{\gamma}}^V(s) \left( \alpha - \beta \frac{\mu_0 h - s}{t} \right) \\ &= \Sigma_{\hat{\gamma}}^V(s) \left[ \frac{(\alpha t + \beta s) - \beta \mu_0 h}{t} \right] \end{aligned} \quad (8.49)$$

which is not very close to the desired representation. It does include an approximation of the scattering cross section by the total cross section, however.

The exponential term can be examined best by performing the integration of the argument:

$$\int_0^s \Sigma_{\hat{\gamma}}^V(s) (\alpha - \beta \mu(s)) ds = \alpha \sum_i \Delta s_i \Sigma_{\hat{\gamma}, i}^V - \beta \sum_i \Sigma_{\hat{\gamma}, i}^V \int_{s_i}^{s_{i+1}} \mu(s) ds$$

$$= \alpha \sum_i \Delta s_i \Sigma_{\hat{\gamma}, i}^V + \beta \sum_i \Delta s_i \Sigma_{\hat{\gamma}, i}^V \quad (8.50)$$

where

$$\begin{aligned} \Delta s_i &\text{ is the incremental distance in the } i\text{th region along } \hat{\Omega}' \\ &= s_{i+1} - s_i \end{aligned}$$

$\Delta t_i$  is the change in distance on the other leg of the scattering triangle corresponding to  $\Delta s_i$ , i.e.,

$$\Delta t_i = - \int_{s_i}^{s_{i+1}} \mu(s) ds = t(s) \Big|_{s_i}^{s_{i+1}} = t_i - t_{i+1}$$

The first term in this expression properly accounts for attenuation along the first leg of the scattering triangle. The second term accounts for the relative effect of the change in attenuation along the second leg with an implicit assumption of material distributions being spherically symmetric about  $\frac{t}{b}$ . This assumption is as good as any since a detailed treatment of the second leg of the triangle is prohibitive.

Thus, this sampling function approximates to some degree the desired effects. It is sampled by obtaining a random number  $\xi$  from RANNO and solving the equivalent of equation 8.3:

$$\xi = \frac{1}{C} \left[ 1 - \exp \left( - \int_0^s \Sigma_{j=1}^X (s') (\alpha - \beta \mu(s')) ds' \right) \right]$$

$$\text{i.e. } \exp \left[ - \int_0^s \Sigma_{j=1}^X (s') (\alpha - \beta \mu(s')) ds' \right] = 1 - \xi C$$

$$\text{or } [\alpha(s - s_i) + \beta(t - t_i)] \Sigma_{j=i+1}^{i-1} + \sum_{j=1}^{i-1} \Sigma_{j=1}^X (\alpha \Delta s_j + \beta \Delta t_j) = -\ln(1 - \xi C)$$

This equation can be written in the form

$$t = As + B$$

where

$$A = -\frac{\alpha}{\beta}$$

$$B = -\frac{1}{\beta} \left\{ \alpha s_i + \beta t_i - \frac{1}{\Sigma_{j=1}^X} \left[ \ln(1 - \xi C) + \sum_{j=1}^{i-1} \Sigma_{j=1}^X (\alpha \Delta s_j + \beta \Delta t_j) \right] \right\}$$

$$\text{But } t^2 = s^2 + h^2 - 2\mu_0 sh = A^2 s + 2As + B^2$$

$$\text{or } A'^2 s^2 + 2B's + C' = 0 \quad A' = A^2 - 1$$

$$B' = AB + \mu_0 h$$

$$C' = B^2 - h^2$$

Thus

$$s = \frac{-B' + \sqrt{B'^2 - A'C'}}{A'}$$

(8.51)

$$p^*(s) = \Sigma_{j=1}^V (s) \left( \alpha - \beta \frac{\mu_0 h \cdot s}{A s + B} \right) \left( 1 - \xi C \right)$$

#### Curve Fit of Spatial Dependence ( $1/s^2$ )

The second technique used for selecting the distance to the scattering point is similar in some respects to the above function. In an unnormalized form, it attempts to approximate equation 2.37 by

$$u^*(s) = \Sigma_{j=1}^V \Delta \Omega(s) \exp \left[ - \int_0^s \Sigma_{j=1}^X (s') [\alpha - \beta \mu(s')] ds' \right] \quad (8.52)$$

where  $\Delta \Omega(s)$  is the  $1/s^2$  factor for point detector calculations, and is a fractional solid angle for surface and volume detector calculations.

This function is curve-fit from boundary to boundary of the regions traversed by:

$$u^*(s) = A_i \exp \left[ -a_i (s - s_i) \right] \quad \text{for } s_i < s < s_{i+1} \quad (8.53)$$

$$\text{where } u^*(s_i) = \Sigma_{j=1}^V \Delta \Omega(s_i) \exp \left[ \int_0^{s_i} \Sigma_{j=1}^X (s') [\alpha - \beta \mu(s')] ds' \right]$$

$$u^*(s_{i+1}) = \Sigma_{j=1}^V \Delta \Omega(s_{i+1}) \exp \left[ \int_0^{s_{i+1}} \Sigma_{j=1}^X (s') [\alpha - \beta \mu(s')] ds' \right]$$

i.e.,

$$A_i = u^*(s_i)$$

$$\alpha_i = -\frac{1}{s_{i+1}-s_i} \ln \left[ \frac{u^*(s_{i+1})}{u^*(s_i)} \right]$$

$$= \alpha \sum_{j=i}^V \frac{\Delta \Omega(s_{j+1})}{s_{j+1}-s_j} \left[ \ln \frac{\Delta \Omega(s_{i+1})}{\Delta \Omega(s_i)} - \beta \sum_{j=i}^V (t_{j+1} - t_j) \right] \quad (8.54)$$

The fractional solid angle is computed for the detector sphere in surface and volume flux calculations as:

$$\Delta \Omega(s_i) = 1 \text{ if } t_i < R_p$$

$$= \frac{1}{4} \left( \frac{R_p}{t_i} \right)^2 \left[ 1 + \frac{1}{4} \left( \frac{R_p}{t_i} \right)^2 \right] \text{ if } t_i > R_p \quad (8.55)$$

This is a second order Taylor series expansion and is sufficiently accurate considering the other approximations.

This function is sampled by the transformation equivalent to equation 8.3 using a random number  $\xi$  from RANNO:

$$\xi = \frac{\int_0^{s_i} u^*(s') ds'}{\int_0^{s_m} u^*(s') ds'} = \frac{1}{C} \left\{ \sum_{i=1}^{i-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp \left[ -\alpha_i (s_i - s_i) \right] \right] \right. \\ \left. + \int_{s_i}^s A_i \exp \left[ -\alpha_i (s - s_i) \right] ds' \right\}$$

$$\text{or } A_i \exp \left[ -\alpha_i (s - s_i) \right] = A_i - \alpha_i \left\{ \xi C - \sum_{i=1}^{i-1} \frac{A_i}{\alpha_i} \left[ 1 - \exp(-\alpha_i (s_i - s_i)) \right] \right\} = D \quad (8.56)$$

where D is the right side of the above equation and C is the denominator of the original.

Thus  $u^*(s) = D/C$

$$\text{and } s = s_i - \frac{1}{\alpha_i} \ln(D/A_i)$$

## SECTION

### 9.0 DATA INPUT INSTRUCTIONS

#### 9.1 GENERAL INPUT PROCEDURES

##### Input Formats

The FASTER program utilizes standard Fortran input statements, Reference 16 page 19. A variety of formats are used. Each format utilizes various combinations of the following data fields

hollerith information: A4 field (4 columns)

integer data: I3 field (3 columns)

floating point data: E9.0 field (9 columns)

Note that for floating point data entered without a decimal point, the decimal point is assumed to be to the right of the data field.

In preparing data, it should be remembered that all blanks in integer or floating point fields are interpreted as zeros. Therefore, all integers (including exponents of floating point numbers) must be right adjusted.

##### Card Input and Output

Each physical data card is written on the output file as soon as it is read from the input file. The resulting printout includes the information in card columns (cc) 73 through 80 of the data cards. Since the present version of FASTER does not print details of problem data except for the input cards, prolific use of card labeling is desirable. A note of warning: in obtaining the card identification from cc 73-80, all unused data fields in cc 1-72 are interpreted as data and these unused fields should be blank or contain valid data punches.

##### Input Data Sections

Data input to FASTER is divided into six sections to simplify the description of multiple change cases. The first data card of each input section is the minimum input requirements. The six sections of data input are:

- 1) title cards, limits, and options
- 2) surfaces and regions
- 3) fixed sources
- 4) cross sections
- 5) flux groups, response functions and detectors
- 6) sampling parameters

Detailed input instructions are given below. To simplify the setup of a problem, it is probably best to start with geometric input data and come back to the data input for section 1 later. This simplifies the necessity of specifying maximum array dimensions before the arrays can be used.

#### Input Control Procedure

The first data card of each input section contains integer constants controlling the input of the remainder of the data in the section. The general procedure is

- input control constant  $\leq 0$ , no input
- input control constant  $> 0$ , input

When input is present, the control constant may serve a dual purpose by also denoting the quantity of input. As an example, the first input-control constant (IN1) for each data section pertains to Hollerith or comment cards. If non-positive ( $IN1 \leq 0$ ), no comment cards should be supplied. If desired, any number of comment cards (up to 999) may be inserted immediately after the input control card for each section. The value of the first input-control constant is then set equal to the number of these comment cards, i.e.,  $IN1 = \text{total number of comment cards}$ .

#### 9.2 SECTION 1 DATA; TITLE CARDS, LIMITS AND OPTIONS

##### CARD 1-0, Input Controls for Section 1 Data

NOTE: This card is always required.

Column	Format	Symbol	Definition
1-3	I3	IN1	input control for card 1-1 (descriptive cards) omit card 1-1 if $IN1 \leq 0$ supply IN1 physical cards if $IN1 \geq 1$
4-6	I3	IN2	input control for card 1-2 (first title card) omit card 1-2 if $IN2 \leq 0$ supply card 1-2 if $IN2 \geq 1$
7-9	I3	IN3	input control for card 1-3 (second title card) omit card 1-3 if $IN3 \leq 0$ supply card 1-3 if $IN3 \geq 1$
10-12	I3	IN4	input control for card 1-4 (geometric limits) omit card 1-4 if $IN4 \leq 0$ supply card 1-4 if $IN4 \geq 1$
13-15	I3	IN5	input control for card 1-5 (fixed source limits) omit card 1-5 if $IN5 \leq 0$ supply card 1-5 if $IN5 \geq 1$
16-18	I3	IN6	input control for card 1-6 (cross section limits) omit card 1-6 if $IN6 \leq 0$ supply card 1-6 if $IN6 \geq 1$
19-21	I3	IN7	input control for card 1-7 (neutron scattering limits) omit card 1-7 if $IN7 \leq 0$ supply card 1-7 if $IN7 \geq 1$
22-24	I3	IN8	input control for card 1-8 (flux limits) omit card 1-8 if $IN8 \leq 0$ supply card 1-8 if $IN8 \geq 1$
25-27	I3	IN9	input control for card 1-9 (sampling options) omit card 1-9 if $IN9 \leq 0$ supply card 1-9 if $IN9 \geq 1$

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
28-30	I3	IN10	input control for card 1-10 (iteration limits) omit card 1-10 if $IN10 \leq 0$ supply card 1-10 if $IN10 \geq 1$
31-33	I3	IN11	printout control for storage map printout no printout of storage map if $IN11 \leq 0$ printout of storage map if $IN11 \geq 1$
34-72	I3I3	---	these columns are not used and should be left blank
73-80	2A4	---	any desired information for card identification

CARD 1-1, Descriptive Information for Section 1 Data

NOTE: a) Omit this card if  $INI \leq 0$   
b) Supply  $INI$  physical cards if  $INI \geq 1$

1-72    I8A4    ---    any desired information

73-80    2A4    ---    any desired information for card identification

CARD 1-2, First Title Card for Labeling the Printout

NOTE: a) Omit this card if  $IN2 \leq 0$   
b) Supply this card if  $IN2 \geq 1$

1-72    I8A4    ---    any desired information for identification of the problem; this will then appear on the first line of each printout page

73-80    2A4    ---    any desired information for card identification

CARD 1-3, Second Title Card for Labeling the Printout

NOTE: a) Omit this card if  $IN3 \leq 0$   
b) Supply this card if  $IN3 \geq 1$

1-72    I8A4    ---    any desired information for identification of the problem; this will appear on the second line of each printout page

73-80    2A4    ---    any desired information for card identification

CARD 1-4, Geometric Limits

NOTE: a) Omit this card if  $IN4 \leq 0$   
b) Supply this card if  $IN4 \geq 1$

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-3	I3	NSMAX	total number of surfaces required for the geometry description
4-6	I3	NAMAX	maximum number of coefficients required to define each surface in the expanded form, e.g. NAMAX = 6 for geometries involving any cones or spheres
7-9	I3	NRMAX	total number of regions required to describe the material distribution including voids
10-12	I3	NBMAX	maximum number of surfaces bounding a region
13-15	I3	NSRMAX	maximum number of regions which can be traversed by a single straight line or ray; the theoretical limit is $2-NSMAX - (\text{number of plane surfaces} + 1)$
16-72	I9I3	---	these columns are not used and should be left blank
73-80	2A4	---	any desired information for card identification

CARD 1-5, Fixed Source Limits

NOTE: a) Omit this card if  $IN5 \leq 0$   
b) Supply this card if  $IN5 \geq 1$

1-3    I3    NEMAX    number of energy groups used in this problem for source spectra and cross sections; source spectra may be described in a different group structure and are regrouped as noted later

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
4-6	I3	NVMAX	total number of fixed sources
7-9	I3	NXMAX	maximum number of points used to tabulate each spatial or angular source distribution
10-12	I3	NXEMAX	maximum number of energy points required to tabulate the differential source spectrum; if any source spectra are described by group integrated values, 2 energy points will be generated for each of these groups before the spectrum is integrated into the group structure for the problem.
13-72	20I3	---	these columns are not used and should be left blank
73-80	2A4	---	any desired information for card identification

CARD 1-6, Basic Cross Section Limits and Options

NOTE: a) Omit this card if IN6  $\leq$  0  
       b) Supply this card if IN6  $\geq$  1.

1-3	I3	NXSECT	particle type option. 0 for photons, 1 for neutrons
4-6	I3	NUNITD	composition units option. 0 for $10^{24}$ atoms/cm <sup>3</sup> 1 for gm/cm <sup>3</sup>
7-9	I3	NUNITX	microscopic cross sections units option. 0 for barns/atom, 1 for cm <sup>2</sup> /gm
10-12	I3	NIMAX	total number of different elements or isotopes
13-15	I3	NMMAX	total number of composite materials; hydrogen densities can be entered by region thus reducing the total number of different compositions, e.g. pure hydrogen compositions need not be defined

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
16-72	19I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

CARD 1-7, Neutron Scattering Limits

NOTE: a) Omit this card if IN7  $\leq$  0.  
       b) Supply this card if IN7  $\geq$  1.  
       c) For photon problems, define each variable on this card as zero.

1-3	I3	NORDER	Maximum number of elastic scattering transfer matrices for all elements. 1 for P <sub>0</sub> (transport approximation) (see notes after card 4-5 for internal cross section juggling). 2 for P <sub>1</sub> , 3 for P <sub>2</sub> , etc. Negative fluxes can, and will, occur for NORDER $>$ 1 and deep penetrations.
4-6	I3	NDOWN	Maximum group-to-group transfer for elastic scattering for all elements. 1 for in group only, 2 for down 1, etc.
7-9	I3	INELAS	Maximum number of groups for which non-elastic transfer can be initiated for all elements. 0 indicates none if these cross sections are included in the P <sub>0</sub> transfer.
10-12	I3	NDOWN	Maximum group-to-group transfer for non-elastic scattering for all elements. 1 for in group only, 2 for down 1, etc.
13-72	20I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

**CARD 1-8, Flux Limits**

NOTE: a) Omit this card if  $IN8 \leq 0$ .  
 b) Supply this card if  $IN8 \geq 1$ .

Column	Format	Symbol	Definition
1-3	I3	NGMAX	Number of flux groups; less than or equal to the number of groups used for source spectra and cross sections.
4-6	I3	NFMAX	Total number of response functions such as flux-to-dose, energy deposition. 0 indicates no response functions.
7-9	I3	NVMOD	Total number of fixed sources for which separate flux contributions are desired. 0 indicates none.
10-12	I3	NCMAX	Number of separate contributions to the flux by order-of-scatter 0 - none, 1 - uncollided flux, 2 - uncollided flux and single scattered flux, etc.
13-15	I3	NUMAX	Number of Legendre moments of the angular flux 0, $P_0$ moment (always obtained, $\equiv$ scalar flux). 1, $P_0$ and $P_1$ moments, etc.
16-18	I3	NTMAX	Number of length-of-flight moments of the flux. 0, zeroth moment (always obtained, $\equiv$ scalar flux), 1, zeroth and first moments, etc.
19-21	I3	NDMAX	Total number of detectors of all types.
22-24	I3	NSRMAX	Number of regions for which separate scattering contributions are desired.
25-72	16I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

**CARD 1-9, Random Sampling Options**

NOTE: a) Omit this card if  $IN9 \leq 0$ .  
 b) Supply this card if  $IN9 \geq 1$ .  
 c) The preferred value of each number on this card is 1.

Column	Format	Symbol	Definition
1-3	I3	NPOINT	0, calculate fluxes for all detectors simultaneously; the preferred point is defined by input, point detectors must be in void regions. 1, calculate fluxes for each point detector individually; the preferred point(s) is the point detector - surface and volume detectors are ignored.
4-6	I3	MODEL	0, randomly select the fixed source and then randomly select the spatial source variables in the source coordinate system. 1, randomly select the spatial source variables in a spherical coordinate system centered at the preferred point; all sources must be volumetric; source volumes must completely cover one or more regions.
7-9	I3	MODELQ	0, randomly select angular source variables in the source coordinate system. 1, randomly select angular source variables like MODELV below, with the direction-before-scattering defined as a unit vector from the center of all sources to the selected source point; all sources must be angular or isotropic.
10-12	I3	MODELU	Distance between scatters random selection option. 0, exponential transformation. 1, curve fit of optimum function (difficulties may be encountered for large volumes since the curve fit is from boundary to boundary).
13-15	I3	MODELV	Direction vector random selection option. 0, polar angle measured from direction before scattering; azimuthal angle measured from a unit direction vector towards the preferred point. 1, opposite of the above. 2 and 3, polar angle measured from unit vector towards preferred point with combined importance parameters from equations 7.42 and 7.45, or 7.41 and 7.45 respectively; azimuthal angles equiprobable.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
16-72	19I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.
<u>CARD 1-10, Iteration Limits</u>			
			NOTE: a) Omit this card if INT0 $\leq 0$ . b) Supply this card if INT0 $\geq 1$ .
1-3	I3	NPRINT	Total number of printouts during the flux calculations; yields a convergence history and protects against complete loss if the problem is terminated.
4-6	I3	NUNITS	Number of outer iterations between printouts, i.e., the number of packets of particles.
7-9	I3	NUMBER	Number of discrete directions obtained by random sampling used in integrating the angular point sources to obtain surface and volume averaged fluxes; not used if NPOINT = 1, or if all detectors are points.
10-12	I3	KALIDE	Maximum number of inner iterations per outer iteration, i.e., the number collisions per packet.
13-72	20I3	---	These columns are not used and should be left blank.
73-78	2A4	---	Any desired information for card identification.

### 9.3 SECTION 2 DATA; SURFACES AND REGIONS

#### CARD 2-0, Input Controls for Section 2 Data

NOTE: a) This card is always required.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-3	I3	IN1	Input control for Card 2-1 (descriptive cards). Omit Card 2-1 if IN1 $\leq 0$ . Supply IN1 physical cards if IN1 $\geq 1$ .
4-6	I3	IN2	Input control for Card 2-2 (surfaces). Omit Card 2-2 if IN2 $\leq 0$ . Supply IN2 physical cards if IN2 $\geq 1$ , i.e., IN2 surfaces will be described.
7-9	I3	IN3	Input control for Card 2-3 (regions) Omit Card 2-3 if IN3 $\leq 0$ . Supply IN3 physical cards if IN3 $\geq 1$ , i.e., IN3 regions will be described.
10-12	I3	IN4 <sup>(1)</sup>	Ambiguity index calculation option. 0, do not calculate ambiguity indices. 1, calculate ambiguity indices. Use IN4 = 1 on the first problem.
13-15	I3	IN5 <sup>(1)</sup>	Geometry consistency check option. 0, do not check geometry. 1, check geometry. Use IN5 = 1 on the first problem.
16-72	19I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

<sup>(1)</sup> IN4 = IN5 = 0 permits the redefinition of regions by surface description changes only. Extreme caution should be used to ensure that the ambiguity indices, calculated in previous problems, are still correct.

CARD 2-1, Description of Section 2 Data

NOTE: a) Omit this card if IN1  $\leq 0$ .  
       b) Supply IN1 physical cards if IN1  $\geq 1$ .

Column	Format	Symbol	Definition
1-72	18A5	---	Any desired information for description of the input data.
73-80	2A4	---	Any desired information for card identification.

CARD 2-2, Surface Description

NOTE: a) Omit this card if IN2  $\leq 0$ .  
       b) Supply IN2 physical cards if IN2  $\geq 1$ .

1-3	I3	I	Index of the surface being described.
4-6	I3	NTP(I)	Index (I) of the last non-zero coefficient if the surface is in the expanded form; calculated internally for all other surfaces.
7-9	I3	NEX	Form (n,) of the surface as input. 0, already in expanded form. 1, $\leq$ NEX $\leq$ 13, special form as indicated in figures 7, 8, and 9 and table 1.
10-18	E9.0	AA(1)	First parameter defining the surface.
19-27	E9.0	AA(2)	Second parameter defining the surface.
64-72	E9.0	AA(7)	Seventh parameter defining the surface.
73-80	2A4	---	Any desired information for card identification.

The requisite parameters are listed in the last column of table 1 and are input in the order shown. If the surface is in the expanded form and rotational terms are involved, supply these on Card 2-2' before supplying Card 2-2 for the next input surface.

TABLE 1  
SPECIAL SURFACE EQUATIONS

Surface Type	Surface Equation	Last Term	Type	Input
quadric	$a_0 + a_1(x - x_0)^2 + a_2(y - y_0)^2 + a_3(z - z_0)^2 + a_4(x - x_0)(y - y_0)^2 + a_5(x - x_0)(z - z_0)^2 + a_6(y - y_0)(z - z_0)^2 = 0$	$i(\text{input})$	$n_x$	$x_0 \ a_1 \ a_2 \ a_3 \ a_4 \ a_5 \ a_6$
plane I x-axis	$x = c$	1	$n_x$	$c$
plane I y-axis	$y = c$	2	$n_x$	$c$
plane I z-axis	$z = c$	3	$n_x$	$c$
plane II x-axis	$(x - x_0)/(z - z_0) = (y_1 - y_0)/(z_1 - z_0)$	3	$n_x$	$x_0 \ y_0 \ z_0 \ y_1 \ z_1$
plane II y-axis	$(x - x_0)/(z - z_0) = (x_1 - x_0)/(z_1 - z_0)$	3	$n_x$	$x_0 \ y_0 \ z_0 \ x_1 \ z_1$
plane II z-axis	$(x - x_0)/(y - y_0) = (x_1 - x_0)/(y_1 - y_0)$	2	$n_x$	$x_0 \ y_0 \ z_0 \ x_1 \ y_1$
cone II x-axis	$\left[ (y - y_0)^2 + (z - z_0)^2 \right]^{1/2} / (x - x_0) = (x_1 - x_0) / (x_1 - x_0)$	6	$n_x$	$x_0 \ y_0 \ z_0 \ x_1 \ y_1 \ z_1$
cone II y-axis	$\left[ (x - x_0)^2 + (z - z_0)^2 \right]^{1/2} / (y - y_0) = (x_1 - x_0) / (y_1 - y_0)$	6	$n_x$	$x_0 \ y_0 \ z_0 \ y_1 \ z_1$
cone II z-axis	$\left[ (x - x_0)^2 + (y - y_0)^2 \right]^{1/2} / (z - z_0) = (x_1 - x_0) / (z_1 - z_0)$	6	$n_x$	$x_0 \ y_0 \ z_0 \ x_1 \ y_1 \ z_1$
cylinder II x-axis	$(y - y_0)^2 / a^2 + (z - z_0)^2 / b^2 = 1$	6	$n_x$	$x_0 \ a \ z_0 \ b$
cylinder II y-axis	$(x - x_0)^2 / a^2 + (z - z_0)^2 / b^2 = 1$	6	$n_x$	$x_0 \ a \ z_0 \ b$
cylinder II z-axis	$(x - x_0)^2 / a^2 + (y - y_0)^2 / b^2 = 1$	5	$n_x$	$x_0 \ a \ y_0 \ b$
ellipsoid	$(x - x_0)^2 / a^2 + (y - y_0)^2 / b^2 + (z - z_0)^2 / c^2 = 1$	6	$n_x$	$x_0 \ a \ y_0 \ b \ z_0 \ c$

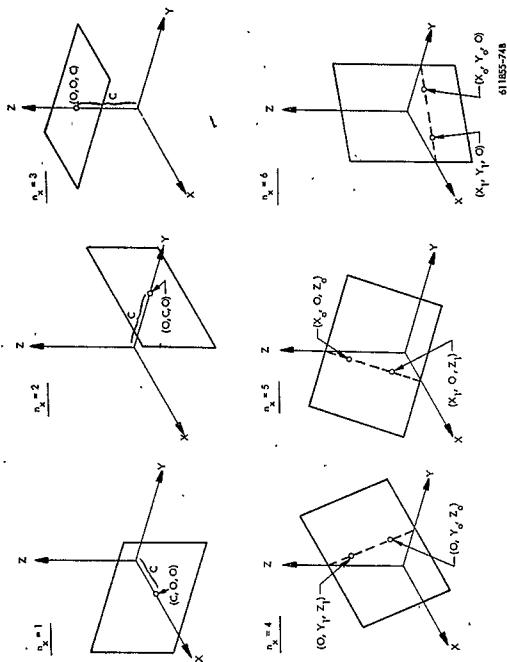


Figure 7. Plane Surfaces

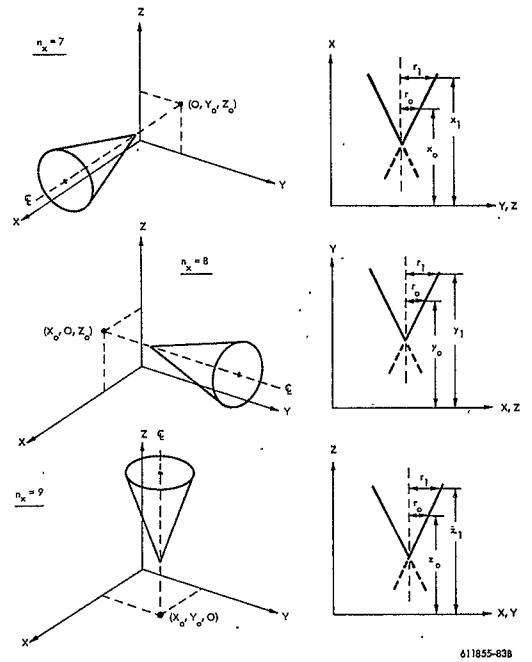


Figure 8. Conical Surfaces

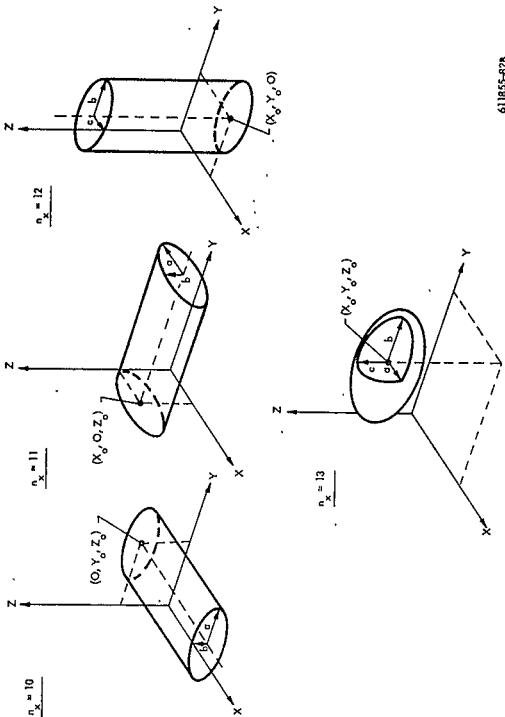


Figure 9. Elliptical Cylinders and Ellipsoids

611855-828

CARD 2-2', Rotated Surface Equation Terms

NOTE: a) Supply this card as required to finish the description of a surface, omit otherwise.

Column	Format	Symbol	Definition
1-9	E9.0	A(7, I)	Coefficient of $xy$ in the general surface equation.
10-18	E9.0	A(8, I)	Coefficient of $yz$ in the general surface equation.
19-27	E9.0	A(9, I)	Coefficient of $zx$ in the general surface equation.
28-72	5E9.0	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

CARD 2-3, Region Description

NOTE: a) Omit this card if  $IN3 \leq 0$ .  
 b) Supply  $IN3$  physical cards if  $IN3 \geq 1$ .

1-3	I3	I	Index of the region being described.
4-6	I3	ISV(I)	Index of the volume source superimposed over this region. 0, indicates none. Required if and only if MODELP = 1. May also be input or changed on Card 3-6.
7-9	I3	MTL(I)	$> 0$ , index of the composition for the region. $=0$ , the region contains hydrogen only. $< 0$ , the region is void. This index can be input or changed on Card 4-8.
10-12	I3	NS(1, I)	First boundary surface index.
13-15	I3	NS(2, I)	Second boundary surface index. 0 or blank if all boundaries have been listed.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
		•	
	i3	NS(j, i)	jth boundary surface index. 0 or blank if all boundaries have been listed.
		•	
34-36	i3	NS(9, i)	> 0, ninth boundary surface index if the region has exactly nine boundaries. 0, or blank if all boundaries have been listed (less than nine boundaries). -1, if the region has more than nine boundaries; the ninth and remaining boundaries are listed on Card 2-3'.
		•	
37-45	E9.0	RHO(i)	Hydrogen density in the region apart from that specified for the composite material in the region; units according to NUNITD, may be input or changed on Card 4-9.
		•	
46-54	E9.0	XR(1, i)	x-coordinate of any point in the region (cm).
55-63	E9.0	XR(2, i)	y-coordinate of the point in the region (cm).
64-72	E9.0	XR(3, i)	z-coordinate of the point in the region (cm).
73-80	2A4	---	Any desired information for card identification.
		•	
<b>CARD 2-3', Additional Region Boundaries</b>			
		•	
NOTE: a) Supply this card(s) for each region having more than nine boundaries, immediately behind Card 2-3 for the region; omit otherwise. b) This card contains data up to and including the maximum number of boundaries (more than 1 physical card if NBMAX > 32).			
		•	
1-72	24i3	NS(9, i)	Ninth boundary surface index.
		NS(10, i)	Tenth boundary surface index.
		NS(11, i)	Eleventh boundary surface index. 0 or blank if all boundaries are listed.
		•	

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
		NS(NBMAX, i)	Maximum boundary surface index. 0 or blank if all boundaries are listed.
73-80	2A4	---	Any desired information for card identification.
		•	
<b>SECTION 3 DATA; SOURCE DISTRIBUTIONS</b>			
		•	
<b>CARD 3-0, Input Controls for Section 3 Data</b>			
		•	
NOTE: a) This card is always required.			
		•	
1-3	i3	IN1	Input control for Card 3-1 (descriptive cards). Omit Card 3-1 if IN1 ≤ 0. Supply IN1 physical card if IN1 ≥ 1.
		•	
4-6	i3	IN2	Input control for Card 3-2 (energy levels). Omit Card 3-2 if IN2 ≤ 0. Supply Card 3-2 if IN2 ≥ 1.
		•	
7-9	i3	IN3	Input control for fixed sources, omit Cards 3-3 through 3-5 if IN3 ≤ 0. Supply Cards 3-3, 3-4 and 3-5 as required for IN3 fixed sources if IN3 ≥ 1.
		•	
10-12	i3	IN4	Input control for Card 3-6 (source-in-region). Omit Card 3-6 if IN4 ≤ 0. Supply source in region indices on Card 3-6 for IN4 regions if IN4 ≥ 1.
		•	
13-72	20i3	---	These columns are not used and should be left blank.
		•	
73-80	2A4	---	Any desired information for card identification.
		•	
<b>CARD 3-1, Description of Section 3 Data</b>			
		•	
NOTE: a) Omit this card if  IN1  ≤ 0. b) Supply IN1 physical cards if IN1 ≥ 1.			
		•	
1-72	18A4	---	Any desired information for describing the input data.
		•	
73-80	2A4	---	Any desired information for card description.
		•	

CARD 3-2, Energy Levels for Sources and Cross-Sections

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
---------------	---------------	---------------	-------------------

- NOTE: a) Omit this card if IN2  $\leq 0$ .  
 b) Supply this card(s) if IN2  $\geq 1$ .  
 c) Source spectra can be input in any desired group structure and will be regrouped to this set of groups.

1-72	8E9.0	ELL(1)	Upper energy boundary of the first energy group (Mev).
		ELL(2)	Lower energy boundary of the first energy group and upper boundary of the second group
		*	
		*	
		*	
		ELL(NEMAX+1)	Lower energy boundary of the last energy group; also defines the energy cutoff point.
72-80	2A4	----	Any desired information for card identification.

CARD 3-3, Fixed Source Constants and Input Options

- NOTE: a) Omit this card if IN3  $\leq 0$ .  
 b) Supply this card, and Cards 3-4 and 3-5 as required, for IN3 sources if IN3  $\geq 1$ .

1-3	I3	I	Index of source being described.
4-6	I3	NSG(I)	Source geometry type. 0, rectangular 1, cylindrical 2, spherical
7-9	I3	NPC(1, I)	First spatial variable distribution option. (See notes below.)
10-12	I3	NPC(2, I)	Second spatial variable distribution option. (See notes below.)
13-15	I3	NPC(3, I)	Third spatial variable distribution option. (See notes below.)

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
16-18	I3	NPC(4, I)	Azimuthal part of angular distribution option (fourth source variable). (See notes below.)
19-21	I3	NPC(5, I)	Polar part of the angular distribution, option (fifth source variable). (See notes below.)

NOTES: The source variables are shown in Figure 10 and are ordered as:

	<u>Rectangular</u>	<u>Cylindrical</u>	<u>Spherical</u>
J = 1	x (cm)	r (cm)	$\rho$ (cm)
J = 2	y (cm)	$\theta$ (radians)	$\theta$ (radians)
J = 3	z (cm)	z (cm)	$\mu$
J = 4	$\theta^1$ (radians)	$\theta^1$ (radians)	$\theta^1$ (radians)
J = 5	$\mu^1$	$\mu^1$	$\mu^1$

Azimuthal angles are in the range  $-\pi \leq \theta, \theta^1 \leq \pi$

Cosines of polar angles are in the range  $-1 \leq \mu, \mu^1 \leq 1$

If NPC(J, I)  $> 0$ , this is the number of tabulation points required to describe the Jth distribution using Card 3-4. If NPC(J, I)  $< 0$ , the distribution for variable J of source I<sup>1</sup> = -NPC(J, I) is used and Card 3-4 is not required.

22-24	I3	MAX	$> 0$ , number of energy points or energy groups required to describe the source spectrum $< 0$ , use the source spectrum for source number I <sup>1</sup> =MAX.
25-27	I3	NORM	Spectrum normalization option (the total source strength is carried in the spectrum). 0, normalize to total source in particles/sec; 1, normalize to total source in Mev/sec; 2, multiply spectrum by constant (if used with MAX $< 0$ , remember that the spectrum for source I <sup>1</sup> = -MAX has been normalized to the total source strength).

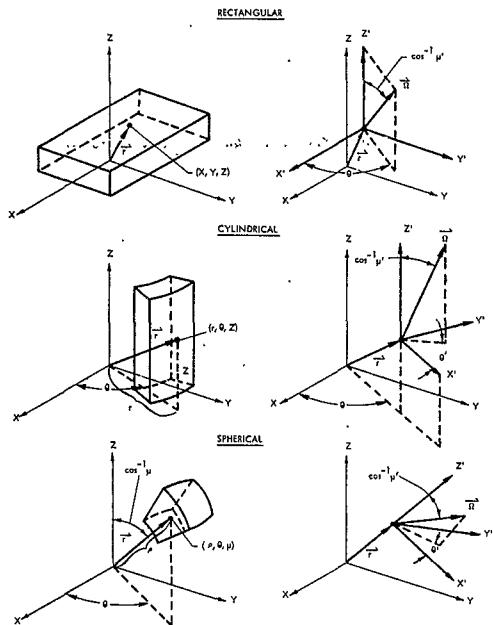


Figure 10. Source Distribution Variables

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
28-30	I3	ISP	Input spectrum units option if MAX > 0 0, differential number spectrum of energy points; 1, differential energy spectrum at energy points; 2, groupwise number spectrum (particles in group); 3, groupwise energy spectrum (energy in group).
31-33	I3	IAL	Spectrum format option if MAX > 0 0, alternating values of energy points and spectrum 1, spectrum input only using energy points previously input for this case 2, energy point input followed by spectrum input on separate card
34-36	I3	----	These columns are not used and should be left blank.
37-45	E9.0	TOT	Source normalization constant, particles/sec if NORM = 0, Mev/sec if NORM = 1, multiplying constant if NORM = 2.
46-54	E9.0	XTR(1, I)	x component of the translation of the source coordinate system origin from the geometry coordinate system origin (cm).
55-63	E9.0	XTR(2, I)	y component of the source translation (cm).
64-72	E9.0	XTR(3, I)	z component of the source translation (cm).
73-80	2A4	----	Any desired information for card identification.

Card 3-4. Source Variable Distributions

- NOTE:
- a) Supply these cards immediately behind Card 3-3 for the source to which they apply. Input for each variable in the order  $J = 1, 2, 3, 4, 5$ , starting a new physical card for each variable.
  - b) Omit for any variable for which  $NPC(J, I) < 0$ .
  - c) This distribution is normalized and interpolated linearly.
  - d) Histogram data can be used (points can coincide).

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	VEE(1, J, I)	Minimum value of the $J_{th}$ variable; if $NPC(J, I) = 1$ , this is the only value of the variable,
		VAL(1, J, I)	Relative value of the distribution function for the $J_{th}$ variable at its minimum value (not used if the variable is discrete, $NPC(J, I) = 1$ )
		*	
		*	
		*	
		VEE(K, J, I)	Kth value of the $J_{th}$ variable.
		VAL(K, J, I)	Relative value of the distribution function corresponding to $VEE(K, J, I)$ .
		*	
		*	
		*	
		VEE(L, J, I)	Maximum value of the $J_{th}$ variable where $L=NP(J, I)$ , the last tabulation point.
		VAL(L, J, I)	Relative value of the distribution function corresponding to $VEE(L, J, I)$ .
73-80	2A4	----	Any desired information for card identification.

NOTE: To avoid numerical difficulties, it is sometimes necessary to decrement the minimum value and increment the maximum value of these variables. In particular, for a uniform azimuthal distribution angular limits use  $3.1416 (\geq \pi)$  for  $\pi$  rather than  $3.14159 (< \pi)$  since the computer uses more significant figures than can be input.

#### CARD 3-5, Source Spectrum

NOTE: a) Supply this card(s) if and only if  $MAX > 0$ . It should be placed immediately after the source variable distributions, if any.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	----	See notes below.
73-80	2A4	----	Any desired information for card identification.
NOTES:			A. If $ISP \leq 1$ , the differential spectrum is tabulated at discrete energy points where $E(I)$ is the maximum spectral energy (Mev) • • • $E(MAX)$ is the minimum spectral energy $EN(K)$ is the differential spectrum corresponding to the Kth energy point $E(K)$ . The units of $EN(K)$ are particles/Mev if $ISP = 0$ or Mev/ Mev if $ISP = 1$ .
A. 1 If $ AL  = 0$ , the input on Card 3-5 consists of alternating values of energy and spectrum $E(1), EN(1), E(2), EN(2), \dots, E(MAX), EN(MAX)$			
A. 2 If $ AL  = 1$ , the energy points are already defined by prior input for this case (they will not be available from a previous case). The input consists of the relative spectrum at these points $EN(1), EN(2), \dots, EN(MAX)$			
A. 3 If $ AL  = 2$ , the energy points are defined first $E(1), E(2), \dots, E(MAX)$ and then followed by another card with the corresponding spectrum $EN(1), EN(2), \dots, EN(MAX)$			
B. If $ISP \geq 2$ , a groupwise integrated spectrum is tabulated by group where $EBG(1)$ is the upper energy boundary of group 1 • • • $EBG(MAX+1)$ is the lower energy boundary of the last spectrum group $ENG(K)$ is the integral spectrum for the Kth group with units of particles in group K if $ISP=2$ , or Mev in group K if $ISP = 3$ .			

**Column      Format      Symbol      Distribution**

B. 1 If IAL = 0, the input on Card 3-5 consists of alternating values of energy group boundaries and group spectrum.  
 EBG(1), ENG(1), EBG(2), ENG(2), . . . , EBG(MAX), ENG(MAX,  
 EBG(MAX+1))

B. 2 If IAL = 1, the energy group boundaries are already defined and the group-wise spectrum is supplied on Card 3-5  
 ENG(1), ENG(2), . . . , ENG(MAX)

B. 3 If IAL = 2, the energy group boundaries are defined first  
 EBG(1), EBG(2), . . . , EBG(MAX+1)  
 and then followed by another card with the groupwise spectrum  
 ENG(1), ENG(2), . . . , ENG(MAX)

**CARD 3-6, Source in Region Indices**

NOTE: a) Omit this card if IN4 ≤ 0.  
 b) Supply this card(s) if IN4 ≥ 1 with source indices for IN4 regions.  
 c) This data is required if and only if MODELP=1.  
 d) This data can also be input on Card 2-3.

1-72	24(3	I	First geometric region index	$\left[ \begin{array}{c} \text{ISV}(1) \\ \vdots \\ \text{ISV}(I) \end{array} \right]$	$\left\{ \begin{array}{l} K = 1 \\ \quad   \\ \quad 0 \text{ denotes none} \end{array} \right.$
		I	Second geometric region index		
		I	Index of source superimposed over Region I (source which completely covers the region), Region I source which completely covers the region)		
		I	Last geometric region index		
		I	Index of source superimposed over Region I (source which completely covers the region)		-K = IN4
73-80	***	2A4/ ***	Any desired information for card identification.		

**9.5 SECTION 4 DATA, MICROSCOPIC CROSS SECTIONS**
**CARD 4-0, Input Controls for Section 4 Data**

NOTE: a) This card is always required.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-3	I3	IN1	Input control for Card 4-1 (descriptive cards) Omit Card 4-1 if IN1 ≤ 0. Supply INT physical cards if INT ≥ 1
4-6	I3	IN2	Input control for Card 4-2 (neutron scattered energies). Omit Card 4-2 if IN2 ≤ 0. Supply Card 4-2 if IN2 ≥ 1.
7-9	I3	IN3	Input control for microscopic cross sections; Omit Cards 4-3 through 4-7 if IN3 ≤ 0; Supply Cards 4-3 through 4-7 as required if IN3 ≥ 1.
10-12	I3	IN4	Input control for Card 4-8 (material-in-region) Omit Card 4-8 if IN4 ≤ 0. Supply Card 4-8 with IN4 material-in-region indices if IN4 ≥ 1.
13-15	I3	IN5	Input control for Card 4-9 (hydrogen in region) Omit Card 4-9 if IN5 ≤ 0. Supply Card 4-9 with IN5 hydrogen-in-region densities if IN5 ≥ 1.
16-18	I3	IN6	Cross section output option (used only if IN6 ≥ 1), no output if IN6 ≤ 0. If IN6 ≥ 1, total cross sections are printed by group (neutrons) or level (photons) and by material. Heating responses are printed by energy level for each element in (MeV/atom/unit number flux) × 10 <sup>24</sup> and then for each composite material in (MeV/cm <sup>3</sup> /unit number flux).

Column    Format    Symbol    Definition

19-72    18I3    ----    These columns are not used and should be left blank.

73-80    2A4    ----    Any desired information for card identification.

CARD 4-1, Descriptive Information for Section 4 Data

- NOTE: a) Omit this card if IN1  $\leq 0$ .  
       b) Supply IN1 physical cards if IN1  $\geq 1$ .

1-72    18A4    ----    Any desired information for description of the input data.

73-80    2A4    ----    Any desired information for card identification.

CARD 4-2, Average Neutron Energies After Scatter

- NOTE: a) Omit this card if IN2  $\leq 0$ .  
       b) Supply this card (s) if IN2  $\geq 1$ .  
       c) This data is required if and only if NXSECT=1 (neutron problem)

1-72    8E9.0    ESB(1)    Average neutron energy for group 1 (Mev)  
                   ESB(2)    Average neutron energy for group 2  
                   .  
                   .  
                   .  
                   ESB(NMAX) Average neutron energy for the last energy group.

73-80    2A4    ----    Any desired information for card identification.

CARD 4-3, Composition Vector for Element

- NOTE: a) Omit Cards 4-3 through 4-7 if IN3  $\leq 0$ .  
       b) Supply Cards 4-3 through 4-7 as required in sets for each element, i.e., all data for the first element, all data for the second element, etc., through the data for element number NIMAX if IN3  $\geq 1$ .  
       c) The first element must always be hydrogen.

Column    Format    Symbol    Definition

1-72    8E9.0    ATW    Atomic weight of the element in atomic mass units.

73-80    2A4    ZEE    Atomic number of the element.

ATD(1)    ATD(1)    Density of the element in composite material 1, units according to NUNIT1.

ATD(2)    ATD(2)    Density of the element in composite material 2.

•  
   .  
   .

ATD(NMMAX)    Density of the element in the last composite material.

73-80    2A4    ----    Any desired information for card identification.

CARD 4-4, Microscopic Total Cross Sections

- NOTE: a) Supply this card immediately after Card 4-3 for each element.

1-72    8E9.0    XST(1)    Total microscopic cross section for energy level 1 (photons) or energy group 1 (neutrons), units according to NUNITX.

XST(2)    Total microscopic cross section for energy level 2 or energy group 2.

•  
   .  
   .

XST(NMAX) Total microscopic cross section for next-to-last photon energy level or last neutron energy group.

XST(NMAX+1) Total microscopic photon cross section for the last energy level. Do not input a number for neutrons.

73-80    2A4    ----    Any desired information for card identification.

CARD 4-5, Neutron Transfer Cross Section Array Limits

- NOTE: a) Omit this card for photon problems.  
 b) Omit this card for the first element (always hydrogen) of neutron problems.  
 c) Supply this card after Card 4-4, for all elements except the first of neutron problems.

Column      Format : " Symbol " " Definition "

1-72	2413	LMAX	[ LMAX ], number of elastic scattering transfer matrices.
			1 for $P_0$ only
			2 for $P_0$ and $P_1$ , etc.
			LMAX < 0, indicates total cross sections are transport corrected. (See notes below.)
		NDSM	Maximum group-to-group transfer for elastic scattering
			1 for in-group only
			2 for down 1, etc.
		JMAX	Maximum number of groups for which non-elastic transfer can be initiated. 0, none
		KMX(1)	Maximum non-elastic group-to-group transfer for initial group 1, = 1 for in group only, etc.
		KMX(2)	Maximum non-elastic group-to-group transfer for initial group 2, = 1 for in group only, etc.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
73-80	2A4	----	KMX(JMAX) Maximum non-elastic group-to-group transfer for the last possible initial group, = 1 for in group only, etc.
			Any desired information for card identification.
			NORDER = 1, LMAX = -2, Code transports corrects $P_0$ elastic transfer using $P_1$ transfer
			NORDER = 1, LMAX = +2, Code transports corrects $P_0$ elastic transfer and calculates transport cross sections using $P_1$ transfer.
			NORDER ≥ 2, LMAX = -2, Code calculates total cross section using $P_1$ transfer.
			NORDER = 1, LMAX = +1, Cross sections used as input.
			NORDER ≥ 2, LMAX ≥ 2, Cross sections used as input.

CARD 4-6, Neutron Elastic Transfer Coefficients

- NOTE: a) Supply this card for all elements except the first of neutron problems immediately after Card 4-5.

1-72	8E9.0	$\sigma_{i \rightarrow k}^0$	(1) $i$ th Legendre moment of the transfer
		XSE(J, K, L)	cross section from group i to k including (2L + 1) coefficient, e.g., GAM-1, GAM-2 printed output.
73-80	2A4	----	Any desired information for card identification.

(1) Start a card with  $P_0$  in-group transfer for all energy groups

$\sigma_{i \rightarrow j}^0, i = 1, 2, \dots, NEMAX$

Start a new card with  $P_0$  down 1 transfer for all groups except the last

$$\sigma_{i \rightarrow i+1}^{\text{ne}}, i = 1, 2, \dots, \text{NEMAX} - 1$$

•

•

•

Start a new card with  $P_0$  down (NDSM-1) transfer

$$\sigma_{i \rightarrow i + \text{NDSM} - 1}^{\text{ne}}, i = 1, 2, \dots, \text{NEMAX} - (\text{NDSM}-1)$$

Start a new card with  $P_1$  in group transfer

$$\sigma_{i \rightarrow i}^1, i = 1, 2, \dots, \text{NEMAX}$$

•

•

•

Start a new card with  $P_{|LMAX| - 1}$  down (NDSM-1) transfer

$$\sigma_{i \rightarrow i + \text{NDSM}-1}^{|LMAX|-1}, i = 1, 2, \dots, \text{NEMAX} - (\text{NDSM}-1)$$

#### CARD 4-7, Neutron Non-Elastic Transfer Coefficients

NOTE: a) Supply this card for all elements--except the first--of neutron problems immediately after Card 4-6.  
 b) Omit if  $JMAX = 0$ .

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	$\sigma_{i \rightarrow i+k-1}^{\text{ne}}$	(1) Non-elastic transfer coefficient from $XSI(K, J)$ group $j$ to group $i+k-1$ $= \sigma_{i \rightarrow i+K-1}^{\text{inelastic}} + 2 \sigma_{i \rightarrow i+k-1}^{n-2n} + \dots$
73-80	2A4	----	Any desired information for card identification.

#### (1) Start the first card with non-elastic transfer from group 1

$$\sigma_{i \rightarrow i+1+k-1}^{\text{ne}}, K = 1, 2, \dots, \text{KMX}(1)$$

Start a new card with non-elastic transfer from group 2

$$\sigma_{i \rightarrow i+2+k-1}^{\text{ne}}, K = 1, 2, \dots, \text{KMX}(2)$$

•

•

Start a new card with non-elastic transfer from group  $JMAX$

$$\sigma_{i \rightarrow i+JMAX+k-1}^{\text{ne}}, K = 1, 2, \dots, \text{KMX}(JMAX)$$

#### CARD 4-8, Material-In-Region Indices

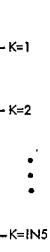
NOTE: a) Omit this card if  $|IN4| < 0$ .  
 b) Supply this card if  $|IN4| \geq 1$  with material indices for  $|IN4|$  regions.  
 c) These indices can also be input on Card 2-3.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	2413	I	First region index
		MTL(I)	$> 0$ , index of material in this region
			$= 0$ , region contains hydrogen only
			$< 0$ , region is void
		I	Second region index
		MTL(I)	Material index for region
			$\bullet$
			$\bullet$
			$\bullet$
		I	Last region index
		MTL(I)	Material Index for region

$K=1$   
 $K=2$   
 $\vdots$   
 $K=IN4$

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
73-80	2A4	----	Any desired information for card identification.

CARD 4-9, Hydrogen Density in Region

1-72	6(13,E9.0)	I	First Region Index	
	RHO(I)		Hydrogen density in region I units according to NUNJTD	
	I		Second region index	
	RHO(I)		Hydrogen density in region I	
	•		•	
	•		•	
	I		Last region index	
	RHO(I)		Hydrogen density in region I	

73-80	2A4	----	Any desired information for card identification.
-------	-----	------	--

## 9.6 SECTION 5 DATA; DETECTORS AND FLUX CONVERSIONS

CARD 5-0, Input Controls for Section 5 Data

NOTE: a) This card is always required.

1-3	I3	IN1	Input control for Card 5-1 (descriptive cards) Omit Card 5-1 if IN1 ≤ 0. Supply IN1 physical cards if IN1 ≥ 1.
4-6	**13**	IN2**	Input control for Card 5-2 (flux groups) Omit Card 5-2 if IN2 ≤ 0. Supply Card 5-2 if IN2 ≥ 1.

Column	Format	Symbol	Definition
7-9	I3	IN3	Input control for Card 5-3 (response functions) Omit Card 5-3 if IN3 ≤ 0. Supply IN3 response functions if IN3 ≥ 1.

10-12	I3	IN4	Input control for Card 5-4 (detectors) Omit Card 5-4 if IN4 ≤ 0. Supply IN4 detectors if IN4 ≥ 1.
13-15	I3	IN5	Input control for Card 5-5 (Δ total flux sources) Omit Card 5-5 if IN5 ≤ 0. Supply Card 5-5 if IN5 ≥ 1.
16-18	I3	IN6	Input control for Card 5-6 (scattered flux regions) Omit Card 5-6 if IN6 ≤ 0. Supply Card 5-6 if IN6 ≥ 1.
19-72	18I3	----	These columns are not used and should be left blank.
73-80	2A4	----	Any desired information for card identification.

CARD 5-1, Descriptive Information for Section 5 Data

NOTE: a) Omit this card if IN1 ≤ 0.  
b) Supply IN1 physical cards if IN1 ≥ 1.

1-72	18A4	----	Any desired information for describing the input data.
73-80	2A4	----	Any desired information for card identification.

CARD 5-2, Flux Groups

NOTE: a) Omit this card if IN2 ≤ 0.  
b) Supply this card if IN2 ≥ 1.  
c) This card is required if the number of flux groups (NGMAX) is less than the number of groups (NEMAX) used to run the problem.

1-72	24I3	NTG(1)	Index of flux group corresponding to first source and cross section group = 1.
------	------	--------	--

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
		NTG(2)	Index of flux group corresponding to the second source and cross section group. • • •
		NTG(NMAX)	Index of flux group corresponding to the last source and cross section group, = NGMAX.
73-80	2A4	---	Any desired information for card identification.
<u>CARD 5-3, Response Functions</u>			
			NOTE: a) Omit this card if IN3 $\leq$ 0. b) Supply this card for IN3 response functions if IN3 $\geq$ 1.
1-3	I3	I	Index of the response function.
4-6	I3	NTP	Response function type, 0, number flux response input by flux group boundary with units (response/particle cm <sup>-2</sup> sec <sup>-1</sup> ). 1, energy flux response input by flux group boundary with units (response/Mev cm <sup>-2</sup> sec <sup>-1</sup> ) <0, energy deposition response function for region I' = -NTP with units (Mev cm <sup>-3</sup> sec <sup>-1</sup> /particle cm <sup>-2</sup> sec <sup>-1</sup> ). NTP <0 requires input of microscopic cross sections in Section 4 data for this problem. (Requires no other input after Column 27 of this card. The response function will appear on the printout immediately below the data on this card.)
7-18	3A4	---	Any desired description of the response function used in labeling the output.
19-27	E9.0	FST	Response function scaling factor, this multiplicative factor can be used to convert the response units to more useful units. Do not use FST = 0, 0.
28-36	E9.0	RSP(I, I)	Response function for upper boundary of the first flux group. Leave blank if NTP < 0. • •

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
63-72	E9.0	RSP(5, I)	Response function for the upper boundary of the 5th flux group (lower boundary of the 4th flux group). Continue on card 5-3' if more than 4 flux groups and NTP $\geq$ 0. Leave blank if NTP < 0.
73-80	2A4	---	Any desired information for card identification.
<u>CARD 5-3', Response Function Continuation</u>			
			NOTE: Supply this card immediately behind Card 5-3 if there are more than 4 flux groups and if NTP $\geq$ 0.
1-72	8E9.0	RSP(6, I)	Response function for the lower boundary of the 5th flux group (upper boundary of the 6th group). • RSP(NGMAX+1, I) Response function for the lower boundary of the last flux group.
73-80	2A4	---	Any desired information for card identification.
<u>CARD 5-4, Detectors</u>			
			NOTE: a) Omit this card if IN4 $\leq$ 0. b) Supply IN4 physical cards if IN4 $\geq$ 1.
1-3	I3	I	Index of the detector being described.
4-6	I3	IDR(I)	0, for point detector >0, region index for a surface or volume detector.
7-9	I3	IDS(I)	0 for a point detector. 0 for a volume detector. >0, surface index for a surface detector (the detector is that part of surface IDS(I) which bounds region IDR(I)).
10-18	E9.0	VOL(I)	Not used for point detector, region volume (cm <sup>3</sup> ) for volume detector; detector area (cm <sup>2</sup> ) for surface detector (1.0 yields surface or volume integrated fluxes).

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
19-27	E9.0	CDT(1, I)	Relative direction cosine with respect to the x - axis of the unit direction vector used in obtaining Legendre moments of the angular flux (not used for surface detectors, angular moments are obtained with respect to the surface normal).
28-36	E9.0	CDT(2, I)	Relative direction cosine with respect to the y - axis.
37-45	E9.0	CDT(3, I)	Relative direction cosine with respect to the z - axis, the 3 direction cosines are normalized by the program.
46-54	E9.0	XDT(1, I)	x coordinate if a point detector (cm).
55-63	E9.0	XDT(2, I)	y coordinate if a point detector (cm).
64-72	E9.0	XDT(3, I)	z coordinate if a point detector (cm).
73-80	2A4	---	Any desired information for card identification.

CARD 5-5, Flux Contribution Sources

NOTE: a) Omit this card if  $IN5 \leq 0$ ,  
       b) Supply this card if  $IN5 \geq 1$ .

1-72	24I3	ISV(1)	Index of first source for which the individual scalar flux contribution is required.
		ISV(2)	Index of the second source for which the individual scalar flux contribution is required.
		ISV(NVMOD)	Index of the last source for which the individual scalar flux contribution is required.

73-80      2A4      ---      Any desired information for card identification.

CARD 5-6, Scattered Flux by Scattering Region

NOTE: a) Omit this card if  $IN6 \leq 0$   
       b) Supply this card if  $IN6 \geq 1$

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	24I3	ISR(1)	Index of the first non-void region from which the individual scattered scalar flux contribution is required.
		ISR(2)	Index of the second non-void region from which the individual scattered scalar flux contribution is required.
		ISR(NSRMAX)	Index of the last non-void region from which the scattered flux contribution is required.
73-80	2A4	---	Any desired information for card identification.
9.7	SECTION 6 DATA; RANDOM SAMPLING PARAMETERS		
			<u>CARD 6-0, Input Controls for Section 6 Data</u>
			NOTE: This card is always required.
1-3	I3	IN1	Input control for Card 6-1 (descriptive cards). Omit Card 6-1 if $IN1 \leq 0$ . Supply IN1 physical cards if $IN1 \geq 1$ .
4-6	I3	IN2	Input control for Card 6-2 (spherical source and detector). Omit Card 6-2 if $IN2 \leq 0$ . Supply Card 6-2 if $IN2 \geq 1$ .
7-9	I3	IN3	Input control for Card 6-3 (source importance). Omit Card 6-3 if $IN3 \leq 0$ . Supply Card 6-3 if $IN3 \geq 1$ .
10-12	I3	IN4	Input control for Card 6-4 (source variable importance). Omit Card 6-4 if $IN4 \leq 0$ . Supply Card 6-4 if $IN4 \geq 1$ .
13-15	I3	IN5	Input control for Card 6-5 (group importance). Omit Card 6-5 if $IN5 \leq 0$ . Supply Card 6-5 if $IN5 \geq 1$ .
16-18	I3	IN6	Input control for Card 6-6 (linear buildup). Omit Card 6-6 if $IN6 \leq 0$ . Supply Card 6-6 if $IN6 \geq 1$ .

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
19-21	I3	IN7	Input control for Card 6-7 (heavy element scatter). Omit Card 6-7 if IN7 $\leq 0$ . Supply Card 6-7 if IN7 $\geq 1$ .
22-24	I3	IN8	Input control for Card 6-8 (hydrogen scatter). Omit Card 6-8 if IN8 $\leq 0$ . Supply Card 6-8 if IN8 $\geq 1$ .
25-27	I3	IN9	Input control for Card 6-8 (scaling factors) Omit Card 6-9 if IN9 $\leq 0$ . Supply Card 6-9 if IN9 $\geq 1$ .
28-72	15I3	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

CARD 6-1, Descriptive Information for Section 6 Data

NOTE: a) Omit this card if IN1  $\leq 0$ .  
b) Supply IN1 physical cards if IN1  $\geq 1$ .

1-72	18A4	---	Any desired information for describing Section 6 data.
73-80	2A4	---	Any desired information for card identification.

CARD 6-2, Spherical Pseudo-Source and Detector

NOTE: a) Omit this card if IN2  $\leq 0$ .  
b) Supply this card if IN2  $\geq 1$ .

1-9	E9 0	RADIUS	Radius of a pseudo spherical source which encloses all the fixed sources (cm).
10-18	E9. 0	XCT(1)	x - coordinate of the center of the sphere (cm).
19-27	E9. 0	XCT(2)	y - coordinate of the center of the sphere (cm).
28-36	E9. 0	XCT(3)	z - coordinate of the center of the sphere (cm).
37-45	E9. 0	DELTA	Radius of a pseudo spherical detector which covers the area in space where fluxes are being calculated; the center of this sphere is the "preferred point" for surface and volume flux calculations (not used if NPOINT = 1).

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
46-54	E9. 0	BDC(1)	x - coordinate of the center of the detector sphere (cm).
55-63	E9. 0	BDC(2)	y - coordinate of the center of the detector sphere (cm).
64-72	E9. 0	BDC(3)	z - coordinate of the center of the detector sphere (cm).
73-80	2A4	---	Any desired information for card identification.

CARD 6-3, Relative Source Importances

NOTE: a) Omit this card if IN3  $\leq 0$ .  
b) Supply this card if IN3  $\geq 1$ .  
c) The data on this card is required if and only if, MODELP = 0 and more than one source is present.

1-72	8E9. 0	RSI(1)	Relative importance of fixed source number 1, (use intuitive knowledge or, better yet, a point kernel calculation of fractional contributions from each source).
...		RSI(NVMAX)	Relative importance of the last fixed source (these importances are normalized in the program).
73-80	2A4	---	Any desired information for card identification.

CARD 6-4, Source Variable Sampling (preferred values)

NOTE: a) Omit this card if IN4  $\leq 0$ .  
b) Supply this card and Card 6-4' for all sources if IN4  $\geq 1$ .  
c) The first three pieces of data on this card are required if MODELP = 0.  
d) The last two pieces of data are required if MODELQ = 0.

1-9	E9. 0	VMD(1, i)	Preferred value of the first source variable of the ith source, must be in the range of the variable including the minimum and maximum values.
-----	-------	-----------	--

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
46-72	3E9.0	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

CARD 6-4', Source Variable Sampling (relative importance)

NOTE: a) Omit this card if  $IN4 < 0$ ,  
       b) Supply this card immediately behind Card 6-4 for the same source if  $IN4 \geq 1$ , i.e.,

Card 6-4 }      Source 1

Card 6-4 }      Source 1

⋮

Card 6-4 }      Last Source

Card 6-4 }      Last Source

c) All numbers on this card must be greater than 0.0.

1-9	E9.0	ALP(1, I)	Relative importance of the preferred value of the first source variable for the $I^{th}$ source, expressed as a ratio to the importance of the value of the variable furthest away (either the minimum or maximum value of the variable).
			> 1.0, the preferred point is more important.
			$0.0 < ALP(1, I) < 1.0$ the preferred point is less important.
			= 1.0, all points are equally important (this number must be > 0.0 since its logarithm is computed).

⋮

37-45	E9.0	ALP(5, I)	Relative importance of the preferred value of the fifth source variable of the $I^{th}$ source.
-------	------	-----------	---

46-72	3E9.0	---	These columns are not used and should be left blank.
73-80	2A4	---	Any desired information for card identification.

CARD 6-5, Group Importance

NOTE: a) Omit this card if  $IN5 < 0$ ,  
       b) Supply this card if  $IN5 \geq 1$ .

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	GIM(I)	Relative importance of particles in the first source and cross section group; e.g., an average flux-to-dose conversion factor for the first group.

⋮

GIM(NEMAX)

Relative importance of particles in the last source and cross section group; e.g., an average flux-to-dose conversion factor for the last group.

73-80	2A4	---	Any desired information for card identification.
-------	-----	-----	--

CARD 6-6, Linear Building Coefficients

NOTE: a) Omit this card if  $IN6 < 0$ ,  
       b) Supply this card if  $IN6 \geq 1$ .

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	AIM(I)	Linear buildup coefficient for group 1 used to estimate the importance of future collisions, this number, when multiplied by the mean free paths to a detector, is used to approximate the future scattered contributions.

⋮

AIM(NEMAX)

Linear buildup coefficient for the last source and cross section group.

73-80	2A4	---	Any desired information for card identification.
-------	-----	-----	--

CARD 6-7, Heavy Element Scattering Importance

NOTE: a) Omit this card if  $IN7 < 0$ ,  
       b) Supply this card if  $IN7 \geq 0$ ,  
       c) All numbers on this card must be greater than 0.0.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-72	8E9.0	ALM(I)	Ratio of forward-to-backward scattering importance for heavy elements for the first energy group. (See note below.)

⋮

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
		ALM(NEMAX)	Similar ratio for the last energy group.
73-80	2A4	---	Any desired information for card identification.

NOTES: For neutrons, a ratio of 10.0 for each group has worked well; for photons, the ratio

$$\frac{d\Sigma}{d\Omega} (0^\circ \text{ scatter}) \times \text{energy after scatter } (0^\circ)$$

$$\frac{d\Sigma}{d\Omega} (180^\circ \text{ scatter}) \times \text{energy after scatter } (180^\circ)$$

using the Klein-Nishina formula for an average group energy has yielded good results.

#### CARD 6-8, Hydrogen Scattering Importance

NOTE: a) Omit this card if IN8  $\leq$  0.  
 b) Supply this card if IN8  $\geq$  1.  
 c) All numbers on this card must be greater than 0.0.

1-72	8E9.0	ALH(1)	Ratio of forward-to-backward scattering importance for hydrogen for the first energy group. (See notes below.)
		ALH(NEMAX)	Ratio of forward-to-backward scattering importance for hydrogen for the last energy group.

NOTES: For neutrons, there is no back scattering from hydrogen; large ratios, e.g.,  $10^3$  for each group have worked well. For photons, numbers identical to those on Card 6-7 have been used.

These numbers are applied only to the hydrogen density specified for the region, they are not applied to the hydrogen part, if any, of material compositions. Therefore, it is essential in neutron problems, that the hydrogen densities be specified by region to properly approximate the angular dependence of neutron scattering from hydrogen.

#### CARD 6-9, Sampling Parameter Scaling Factors

NOTE: a) Omit this card if IN9  $\leq$  0,  
 b) Supply this card if IN9  $\geq$  1.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
1-9	E9.0	ATA	Spherical pseudo source sampling, polar angle importance adjustment.  1.0, all angles equally important, >1.0, shifts importance towards small angles, <1.0, shifts importance towards large angles, Numbers in the range $1.0 \leq ATA \leq 10.0$ work fine. (This number must be greater than 0.0.)
10-18	E9.0	ATB	Spherical pseudo source sampling, azimuthal angle - important adjustment, 1.0, all angles equally important, >1.0, shifts importance towards $0^\circ$ , <1.0, shifts importance away from $0^\circ$ . This angle is measured in a rotated coordinate system and a little difficult to relate to the true coordinate system. The usual procedure is to use ATB = 1.0 (this number must be greater than 0.0).
19-27	E9.0	ATC	Spherical pseudo-source sampling, spatial importance adjustment. 1.0, uses built-in estimate of spatial importance >1.0, shifts importance to lower source energies (source points closer to the detector), <1.0, shifts importance to higher source energies (source points further away). General use of numbers $0.7 < ATC < 1.3$ yield good results. (The program can be tricked for leakage-type surface and volume detector calculations by putting the preferred detector in the center of the source and using ATC $\approx$ 1.0.)
28-36	E9.0	ATD	Flux contribution importance used in cutoff considerations; if all contributions to all detector groups on 2 successive inner iterations of a given outer iteration are less than ATD times the flux already obtained in this outer iteration, then the inner iterations are terminated.

<u>Column</u>	<u>Format</u>	<u>Symbol</u>	<u>Definition</u>
37-45	E9.0	AT	Scaling factor for the spatial importance on the first leg of the scattering triangle. 1.0 uses built-in parameters. <1.0, shifts importance to higher energies >1.0, shifts importance to lower energies General use of numbers $0.6 < AT < 1.2$ yields good results (must be greater than 0.0).
46-54	E9.0	BT	Scaling factor for the spatial importance on the second leg of the scattering triangle; should approximate higher order scattering effects, so it is generally less than AT. If MODELL = 0, it must be less than AT in absolute magnitude, i.e., $ BT  < AT$ . Numbers on the order of $0.4 < BT < 0.9$ yield good results. (If the trick mentioned in discussing ATC is used, it should also be used here; i.e., $-0.9 < BT < -0.4$ .) This number cannot = 0.0.
55-63	E9.0	AS	Scaling factor for preferred direction (towards detector) importance 1.0 uses built-in parameter >1.0 forces even more <1.0 forces less. 0.0 yields no effect. (<0.0 forces away and should be used when ATC and BT are <0.0.) Use of 1.0 yields good results for point detectors.
64-72	E9.0	BS	Scaling factor for scattered direction importance 1.0, uses group averaged parameter. >1.0 forces even more <1.0, forces less. 0.0, no effect from scattered direction. Use of 1.0 yields good results.
73-80	2A4	---	Any desired information for card identification.

## SECTION

### 10.0 SAMPLE PROBLEM

Numerical results obtained from the FASTER program for typical nuclear reactors are reported in References 17 and 18. The input data used in calculations for a large NERVA-type nuclear reactor are reproduced in the classified appendix, Reference 19.

#### Problem Description

A sample problem of moderate complexity is discussed below. Included in this discussion is a description of the printout obtained from running the FASTER program.

The sample problem involves the configuration used in a Lockheed study reported in Reference 20. Data taken directly from this study include: the geometrical model—shown in Figure 11, material compositions; spatial source distributions; and the differential photon spectrum. The specific problem is a dose rate calculation at a point detector above a partially empty liquid hydrogen propellant tank. The problem was run for both photons and neutrons in a single computer run using the "change case" capability of the FASTER program. This problem was run for 400 outer iterations for both photons and neutrons and required less than 4 minutes on the CDC 6600 computer. Approximately twice as much time would have been required on the IBM 7094 computer.

#### Input Data

The complete printout for this problem is shown in Table 2. The data cards for the problems are not shown since they appear on the printout in almost the same form. The major difference is that the card identification from columns 73-80 appears on the left side of the printout. The card identification (if any) is normally followed by 3 periods (...) after which the data from columns 1 - 72 is printed. However, the printout for this problem has had the card number entered above the 3 periods to simplify the examination of data on particular cards.

The listing of the data cards for each section is preceded by a line indicating the appropriate input section. Section 1 data are first and the printout of input data continues

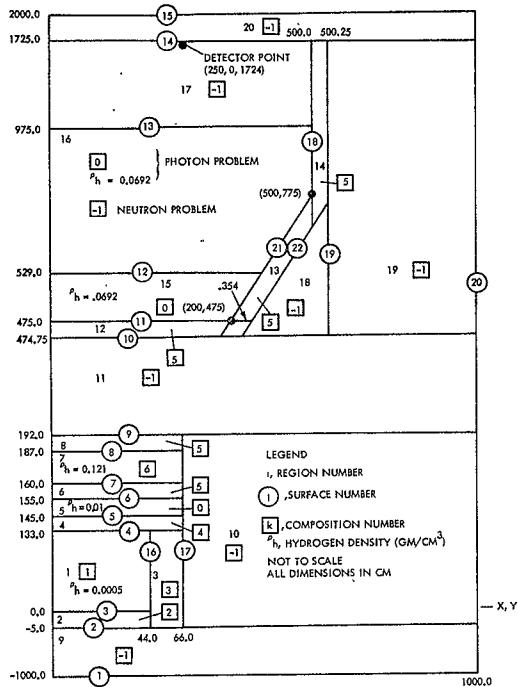


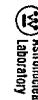
Figure 11. Sample Problem Geometry

611855-758

TABLE 2

THE SAMPLE PROBLEM PRINTOUT IS PRESENTED ON PAGES 154 THROUGH 169





SAMPLE PROBLEM 1: PHOTON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY EMPTY PROPELLANT TANK. REFERENCE ER-0236 (LOCKHEED STUDY FOR NASA-HSFC) \*\*\*\*THE FASTER CODE\*\*\*\*CASE 1



\*\*\*\*\*FLUXES FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 CALCULATED PRECISION ENERGY=REV THIN GROUP THIS GROUP DERIVATIVE DERIVATIVE CUMULATIVE CUMULATIVE  
 GROUP 1... 2.73937E+02 0. 1.2601E+07 0. 0. 0. 0. 0.  
 GROUP 2... 0. 0. 1.0291E+07 0. 0. 0. 0. 0.  
 GROUP 3... 1.2421E+01 4.7045E+00 6.4996E+07 3.0578E+08 3.2498E+08 1.5269E+08 7.7485E+07 9.1017E+08  
 GROUP 4... 1.1802E+01 3.1218E+00 1.2745E+08 3.9772E+08 5.4953E+08 2.6515E+08 2.0509E+08 7.4887E+08  
 GROUP 5... 1.0787E+01 2.7678E+00 1.1787E+08 3.1078E+08 4.8209E+08 2.4120E+08 1.8224E+08 6.9104E+08  
 GROUP 6... 2.4550E+01 1.4820E+00 1.4055E+08 2.1718E+08 4.9209E+08 4.3435E+08 3.1123E+08 1.1330E+09  
 GROUP 7... 2.6005E+01 9.6454E+00 2.8193E+08 2.7194E+08 6.2652E+08 6.0430E+08 7.9317E+08 1.6079E+09  
 GROUP 8... 3.1339E+01 5.5454E+00 3.1043E+08 1.3874E+08 7.7608E+08 1.3401E+09 1.1810E+09 1.8184E+09  
 GROUP 9... 3.1339E+01 5.5454E+00 3.1043E+08 1.3874E+08 7.7608E+08 1.3401E+09 1.1810E+09 1.8184E+09  
 GROUP 10... 6.4923E+01 1.1708E+00 1.1227E+08 1.3145E+08 9.3553E+08 1.0954E+09 2.6309E+09 2.0951E+09  
 GROUP 11... 6.4923E+01 6.3239E+02 7.5915E+07 4.6800E+06 1.5183E+06 9.6017E+07 2.7068E+09 2.0999E+09  
 \*\*\*\*\*NUMBER FLUX RESPONSES FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 REM/HOUR

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GROUP 1... 0.  
 GROUP 2... 6.9952E+01  
 GROUP 3... 1.4652E+02  
 GROUP 4... 4.5571E+02  
 GROUP 5... 5.2991E+02  
 GROUP 6... 3.6205E+02  
 GROUP 7... 2.3206E+02  
 GROUP 8... 6.2356E+02  
 GROUP 9... 6.8087E+01  
 GROUP 10... 2.2405E+02  
 GROUP 11... 6.8763E+01  
 TOTALS..... 3.4401E+03  
 HIRE ERROR... 9.8763E+02  
 MAX..... 2.3206E+02  
 \*\*\*\*\*NUMBER FLUX MOMENTS FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 ITERANT 1 REGION 1 REGION 2 REGION 3 REGION 4 REGION 5 REGION 6 REGION 7  
 GROUP 1... 0. 0. 0. 0. 0. 0. 0. 0.  
 GROUP 2... 1.2661E+07 0. 0. 0. 0. 0. 0. 0.  
 GROUP 3... 6.4998E+00 1.0890E+06 0. 1.3236E+05 3.8260E+05 0. 0. 0.  
 GROUP 4... 1.2743E+08 3.1946E+06 0. 1.3947E+06 1.4615E+06 0. 0. 0.  
 GROUP 5... 1.0787E+08 5.7365E+06 0. 4.8757E+06 1.7618E+06 0. 0. 0.  
 GROUP 6... 2.4550E+08 3.1043E+06 0. 5.2787E+06 1.8224E+06 0. 0. 0.  
 GROUP 7... 2.8193E+08 1.7082E+06 0. 6.3852E+06 3.4080E+06 0. 0. 0.  
 GROUP 8... 5.5454E+08 8.2677E+03 0. 3.6957E+06 2.5685E+06 0. 0. 0.  
 \*\*\*\*\*NUMBER FLUX MOMENTS FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 REGION 8 REGION 12 REGION 13 REGION 14 REGION 15 REGION 16 SCATTER 0 SCATTER 1  
 GROUP 1... 0. 0. 0. 0. 0. 0. 0. 0.  
 GROUP 2... 0. 0. 0. 0. 0. 0. 0. 0.  
 GROUP 3... 1.1729E+06 0. 0. 0. 0. 0. 0. 0.  
 GROUP 4... 9.3473E+06 0. 0. 0. 0. 0. 0. 0.  
 GROUP 5... 7.0084E+06 0. 0. 0. 0. 0. 0. 0.  
 GROUP 6... 1.7421E+06 0. 0. 0. 0. 0. 0. 0.  
 GROUP 7... 3.0564E+06 9.339HE+03 3.9501E+07 6.3105E+07 4.9934E+08 1.6105E+08 3.1205E+05 6.6505E+07  
 GROUP 8... 4.5719E+05 8.3341E+02 1.6128E+07 6.8349E+07 4.2217E+05 9.4273E+04 9.5295E+03 8.7282E+07  
 GROUP 9... 1.5765E+05 1.6128E+02 1.0050E+07 3.3620E+07 3.0076E+04 1.2621E+04 1.2023E+03 9.9542E+05  
 GROUP 10... 1.0787E+05 2.1657E+02 1.6202E+07 3.2047E+07 3.0151E+04 1.2023E+04 1.2023E+03 9.9542E+05  
 GROUP 11... 9.1598E+02 2.1657E+02 3.7032E+01 1.4046E+06 3.7175E+01 7.4563E+07 0. 0.  
 REM/HOUR 9.6452E+01 1.8694F+02 1.1675E+02 2.4323E+02 7.7941E+01 1.9497E+02 5.3457E+02 -1.3021E+03  
 \*\*\*\*\*SCATTERED NUMBER FLUXES ON THE SURFACE OF THE TANK\*\*\*\*\*  
 SCATTER 2 SCATTER 3 SCATTER 4 SCATTER 5 SCATTER 6 SCATTER 7 ANGLAR 1 SPATIAL 1  
 GROUP 1... 0. 0. 0. 0. 0. 0. 0. 0.  
 GROUP 2... 0. 0. 0. 0. 0. 0. 0. 0.  
 GROUP 3... 1.7328E+06 0. 0. 0. 0. 0. 0. 0.  
 GROUP 4... 2.1388E+07 0. 0. 0. 0. 0. 0. 0.  
 GROUP 5... 5.6404E+07 1.2551E+07 0. 0. 0. 0. 0. 0.  
 GROUP 6... 3.8343E+07 3.2024E+07 3.1213E+07 0. 0. 0. 0. 0.  
 GROUP 7... 1.0787E+07 2.1293E+07 6.0951E+07 5.2193E+07 0. 0. 0. 0.  
 GROUP 8... 1.3076E+08 1.9135E+08 8.7301E+07 5.8241E+07 0. 0. 0. 0.  
 GROUP 9... 8.0511E+06 2.9793E+07 6.0951E+07 5.1079E+07 9.1041E+06 0. 0. 0.  
 GROUP 10... 1.0787E+07 2.1293E+07 6.0951E+07 5.2193E+07 1.2722E+07 0. 0. 0.  
 GROUP 11... 9.5745E+03 6.3135E+02 1.6220E+06 1.1555E+05 9.2781E+05 3.4098E+07 0. 0.  
 REM/HOUR 7.9501E+02 4.2642E+02 2.2576E+02 1.0287E+02 2.7109E+01 3.2634E+01 3.2725E+03 3.5354E+03

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SAMPLE PROBLEM 1: PHOTON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY
EMPTY PROPELLANT TANK. REFERENCE EP-R239 (LOCKHEED STUDY FOR NASA-HSPC) ****THE FASTER CODE****CASE
*****N,JORDAN, VAN*****PAGE 1
*****BOUNDARY SEARCH PARAMETERS: (SURFACE=107 PROBABLE NEXT REGION)*****
REGION 1 (-3+ 21 l 4) ( 16- 3)
REGION 2 (-2+ 9) ( 3+ 1) ( 16- 3)
REGION 3 (-1+ 1) ( 5+ 1) ( 16- 3)
REGION 4 (+1+ 1) ( 5+ 1) ( 16- 3) (-1+ 10)
REGION 5 (+5+ 1) ( 6+ 6) ( 17+ 10)
REGION 6 (+6+ 5) ( 7+ 7) ( 17+ 10)
REGION 7 (+7+ 5) ( 8+ 7) ( 17+ 10)
REGION 8 (+8+ 7) ( 9+ 11) ( 17+ 10)
REGION 9 (-1+ 31) ( 20+ 0) ( 20+ 31)
REGION 10 (-1+ 8) ( 10+ 1) ( 20+ 31)
REGION 11 (-1+ 8) ( 10+ 1) ( 20+ 31) (-9+ 11)
REGION 12 (-10+ 11) ( 11+ 15) (-21+ 0)
REGION 13 (-10+ 0) ( 18+ 0) ( 21+ 16) (-22+ 18)
REGION 14 (-10+ 0) ( 18+ 0) ( 21+ 17) (-22+ 0) (-10+ 0)
REGION 15 (-11+ 12) ( 12+ 16) (-21+ 13)
REGION 16 (-12+ 15) ( 13+ 17) ( 18+ 14) (-21+ 13)
REGION 17 (-13+ 16) ( 14+ 20) ( 18+ 14)
REGION 18 (-14+ 16) ( 15+ 21) ( 18+ 14)
REGION 19 (-15+ 0) ( 20+ 31) (-10+ 11) (-1+ 20)
REGION 20 (-1+ 4) ( 15+ 31) ( 20+ 31)

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The logo consists of a stylized 'W' inside a circle, with the letters 'WPS' written vertically below it.



Astronucleon  
Laboratory

SAMPLE PROBLEM 2: NEUTRON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY  
 INERT PROPELLANT TANK. REFERENCE ER-0236 (LOCKHEED STUDY FOR NASA/NFSC) \*\*\*\*THE FASTER CODE\*\*\*\*\*CASE X  
 \*\*\*\*\*T.M.JORDAN,VAHIL\*\*\*\*\*PAGE 3

AS1	C	4	-6	+6.3065E+01	-1.5192E+00	+9.1332E-01	-1.3322E+00						
AS2	C	4	-7	+9.8860E+02	-1.0000E+00	+7.0302E+01	-6.9636E+01	0.					
AS2	C	4	-7	+9.8860E+02	-1.0000E+00	+7.0302E+01	-6.9636E+01	0.					
G1	C	4	-7	0	0.	1.2558E+01	1.5104E+01	7.7026E+02	2.0750E+03	0.3287E+04	5.4130E+04		
T 61	C	4	-7	0	0.	8.0900E+04	8.2336E+05	1.0000E+00	1.0000E+00				
T 62	C	4	-7	0	0.	8.0900E+04	8.2336E+05	1.0000E+00	1.0000E+00				
T 63	C	4	-7	0	0.	8.0900E+04	8.2336E+05	1.0000E+00	1.0000E+00				
T 03	C	4	-7	1.7810E+02	-0.5700E+00	2.6740E+02	-0.5700E+00	0.	2.0000E+00	7.0000E+00	1.2000E+00		
RD	FEN	4	-7	5.5874E+01	-2.6000E+01	1.1000E+01	-0.	0.	2.0000E+00	2.0395E+00	1.9117E+00		
SI	FEN	4	-7	1.7810E+02	-0.5700E+00	2.6740E+02	-0.5700E+00	0.	2.0000E+00	7.0000E+00	1.2000E+00		
LDI-FE	M	4	-7	2	13	11	9	0	-0	-0	-0		
55	FEN	4	-7	1.9941E+00	-2.2222E+00	2.2222E+00	-1.9941E+00	1.0000E+00	1.0000E+00	1.0000E+00	1.0000E+00		
01	FEN	4	-7	5.3240E+02	-2.7638E+02	9.3597E+02	-1.8622E+02	1.9199E+01	7.7770E+02	2.7935E+03	2.0285E+01		
DL	FEN	4	-7	3.6394E+02	-3.5347E+01	1.4220E+02	-2.1618E+01	0.	2.0000E+00	6.1697E+00	2.1524E+00		
A50	FEN	4	-7	5.0359E+02	-5.4773E+02	5.2339E+02	-5.2339E+02	0.	2.0000E+00	5.2339E+02	5.2339E+02		
A51	FEN	4	-7	4.8899E+02	-3.2910E+02	4.9720E+02	-1.1040E+01	-1.7556E+01	-6.6016E+02	-2.0460E+01	-2.5700E+01		
A51	FEN	4	-7	2.3173E+02	-9.9987E+01	-1.1980E+01	-2.1548E+01	-1.0000E+00	-1.0000E+00	-1.0000E+00	-1.0000E+00		
G1	FEN	4	-7	9.0303E+01	-4.1910E+01	-4.1910E+01	-1.0000E+00	1.0000E+00	1.0000E+00	1.0000E+00	1.0000E+00		
T 02	FEN	4	-7	2.0594E+02	-1.0000E+00	2.0594E+02	-1.0000E+00	5.00470E+03	5.00470E+03				
T 02	FEN	4	-7	1.4981E+02	-1.7455E+01	1.5570E+01	-1.6192E+01	0.	4.6860E+01	1.0593E+01	7.7610E+01	9.1646E+02	
T 02	FEN	4	-7	2.4610E+02	-2.6250E+02	1.0540E+02	-5.9736E+03	0.	4.6860E+01	5.3994E+01	1.2549E+01	9.3829E+02	
T 04	FEN	4	-7	8.7453E+01	-1.3012E+02	8.9202E+01	-1.0000E+00	0.	4.6860E+01	1.0593E+01	1.3134E+01	3.0911E+02	
H4	FEN	4	-7	2.4084E+01	-3.2113E+01	5.9709E+01	-4.5313E+02	0.	4.1717E+01	4.4865E+02	7.4787E+02		
T 06	FEN	4	-7	3.7903E+01	-3.6492E+02	8.1551E+01	-6.3833E+03	0.	5.0000E+00	7.1580E+03	6.3634E+03	7.7078E+02	
T 05	FEN	4	-7	1.8782E+03	-8.1551E+01	6.3833E+03	-5.0000E+00	0.	5.0000E+00	7.1580E+03	6.3634E+03	7.7078E+02	
T 07	FEN	4	-7	1.1129E+02	-6.8503E+02	9.7840E+02	-1.2615E+01	4.0076E+02	6.3890E+02	0.	6.4243E+02	9.5891E+02	6.7583E+02
ST	US	4	-7	2.5055E+00	-3.0485E+00	4.1320E+00	-4.7258E+00	0.	1.0000E+00	1.0000E+00	1.0000E+00	1.0000E+00	
ST	US	4	-5	5.9259E+00	-7.0168E+00	8.0903E+00	-9.7811E+00	1.4550E+01	5.1520E+00	-1.2171E+00	5.2229E+00		
LDI-US	U	4	-7	2	13	11	8	0	4	3	2		
SS	US	4	-7	4.9039E+00	-5.8900E+00	7.0164E+00	-8.4298E+00	1.2100E+01	4.3320E+00	3.0614E+00	3.4614E+00		
SS	US	4	-7	4.9039E+00	-5.8900E+00	7.0164E+00	-8.4298E+00	1.2100E+01	4.2221E+00	3.0614E+00	4.1074E+00		
D1	US	4	-7	1.3391E+02	-1.3390E+02	2.3101E+02	-2.7820E+02	6.2664E+02	3.2946E+02	0.8343E+02	9.5222E+02		

SAMPLE PROBLEM 2: NEUTRON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY EMPTY PROPELLANT TANK, REFERENCE ER-023B (LOCKHEED STUDY FOR NASA-HSFC) \*\*\*\*\*THE FASTER CODE\*\*\*\*\*CASE 2

ASD US4/-6	9.8067E+00	1.1762E+01	1.5815E+00	1.2722E+01	5.3946E+00	5.4058E+00
ASD US4/-6	5.313DE+00	5.2723E+00	5.7908E+00	6.1059E+00	1.8176E+01	6.3797E+00
ASD US4/-6	1.1250E+00	1.0800E+00	1.1314E+00	1.0800E+00	6.0846E+00	2.3244E+00
ASD US4/-6	1.3114E+00	1.2244E+00	1.3114E+00	1.2244E+00	6.0846E+00	2.0948E+00
T 61 US4/-7	1.5491E+00	2.1290E+00	6.079E+03	1.1623E+02	2.1765E+02	1.0805E+01
T 61 US4/-7	4.0598E+02	1.2193E+02	1.3882E+02	5.9407E+02	3.6946E+03	1.0805E+01
T 62 US4/-7	2.0245E+02	2.1290E+00	2.1290E+02	1.1623E+02	3.6946E+03	1.0805E+01
T 62 US4/-7	6.5247E+02	6.4250E+02	2.1290E+02	1.1623E+02	3.6946E+03	1.0805E+01
T 63 US4/-7	5.3668E+03	2.3332E+02	6.1445E+02	5.7677E+01	2.6983E+01	2.6965E+01
T 63 US4/-7	1.6088E+01	7.3110E+02	5.0291E+02	2.6983E+01	3.9787E+01	1.3272E+01
T 64 US4/-7	6.4657E+02	6.8486E+02	7.3110E+01	2.6244E+01	2.7070E+01	4.5686E+01
T 65 US4/-7	1.1725E+02	3.6751E+01	2.3332E+01	2.4307E+01	4.7937E+01	1.0533E+01
T 66 US4/-7	8.3496E+02	1.1725E+02	1.1725E+01	1.9240E+01	4.2205E+01	1.8341E+01
T 67 US4/-7	1.3287E+02	1.2050E+02	1.3287E+01	1.0805E+01	2.6752E+01	1.5589E+01
T 68 US4/-7	5.3007E+02	2.9781E+01	1.6733E+01	2.6244E+01	1.7477E+01	1.5489E+01
T 69 US4/-7	1.3608E+01	1.5374E+01	2.9133E+01	2.0505E+01	1.9409E+01	2.4802E+01
T 70 US4/-7	7.4242E+01	2.3332E+01	1.6733E+01	2.4307E+01	2.4802E+01	2.4429E+01
T 71 US4/-7	2.3494E+01	2.3494E+01	2.1622E+01	2.5482E+01	2.5482E+01	2.5482E+01
T 72 US4/-7	1.9649E+01	2.5482E+01	2.5482E+01	2.5482E+01	2.5482E+01	2.5482E+01
T 73 US4/-7	7.2141E+02	7.2141E+02	7.2141E+02	7.2141E+02	7.2141E+02	7.2141E+02
*****SECTION 5: SAMPLING CROSS SECTIONS REQUIRED 577 OF 10400 AVAILABLE LOCATIONS*****						
*****SECTION 6: SAMPLING CROSS SECTIONS REQUIRED 577 OF 10400 AVAILABLE LOCATIONS*****						
*****SECTION 7: SAMPLING CROSS SECTIONS REQUIRED 577 OF 10400 AVAILABLE LOCATIONS*****						
INPUTS 5-7	0 0 0 -0 -0 -0 -0	1.0805E+01	7.2141E+02	1.0805E+01	1.0805E+01	1.0805E+01
RESPONSES 5-7	1 0 0 RAD(TH)/HR	1.0805E+01	1.0805E+05	1.0805E+05	1.0805E+05	1.0805E+05
RESPONSES 5-7	1.1300E-05 8.6000E-06	7.2141E+02	6.2090E-05	5.2000E-06	4.0000E-06	3.9400E-06
RESPONSES 5-7	0 0 0 0 0 0 0	1.0805E+01	1.0805E+05	1.0805E+05	1.0805E+05	1.0805E+05
RESPONSES 5-7	1.4200E+05 1.1300E+05	1.0805E+01	2.7000E+05	2.7000E+05	2.3300E+05	2.1500E+05
RESPONSES 5-7	0 0 0 0 0 0 0	1.0805E+01	7.4000E+06	5.8000E+06	5.0000E+06	3.6000E+06
RESPONSES 5-7	0 0 0 REM/HOUR	1.0805E+01	1.1500E-04	1.1500E-04	1.1500E-04	9.8900E-05
RESPONSES 5-7	9.5000E-05 9.0000E-05	7.4000E+06	4.5000E+05	5.4000E+05	4.0000E+05	3.3000E+05
RESPONSES 5-7	0 0 0 0 0 0 0	1.0805E+01	1.0805E+05	1.0805E+05	1.0805E+05	2.1000E+05
INPUTS 6-7	0 0 0 0 1 1 1	1.0805E+01	-0 -0 -0 -0 -0 -0 -0	1.0805E+01	1.0805E+01	1.0805E+01
GROUP1H4-5	1.8000E+01 1.7200E+05	1.6200E+05	1.4400E+03	1.2000E+03	1.0000E+05	7.9000E+06
GROUP1H4-5	5.7000E+06 4.6000E+05	3.7000E+06	2.9500E+05	2.0000E+05	1.7000E+06	6.7000E+06
ATTEN-IM4-6	1.0000E+01 1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01
ATTEN-IM4-6	1.0000E+01 1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01	1.0000E+01
*****SECTION 8: SAMPLING INPUT AND PREPARATION COMPLETED*****						

\*\*\*\*\*FLUXES FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 CALCULATED AVERAGE NUMBER FLUX ENERGY FLUX NUMBER FLUX ENERGY FLUX NUMBER FLUX ENERGY FLUX  
 PREC1501 ENERGY-NEW THIS GROUP THIS GROUP DERIVATIVE DERIVATIVE CUMULATIVE CUMULATIVE  
 GROUP 1... 6.6185E+01 2.1524E+00 1.8556E+07 9.7388E+05 8.3932E+06 2.1932E+06 1.0566E+07  
 GROUP 2... 6.4742E+01 2.0712E+00 1.7132E+07 9.0102E+05 7.9007E+06 2.0414E+06 8.3901E+07  
 GROUP 3... 6.3851E+01 2.0712E+00 1.7132E+07 9.0102E+05 7.9007E+06 2.0414E+06 8.3901E+07  
 GROUP 4... 6.3927E+01 2.0740E+00 2.0397E+07 6.6760E+05 2.5057E+07 8.2039E+07 7.2224E+07 3.4074E+08  
 GROUP 5... 6.3942E+01 2.0740E+00 2.0397E+07 6.6760E+05 2.5057E+07 8.2039E+07 7.2224E+07 3.4074E+08  
 GROUP 6... 6.3942E+01 2.0740E+00 2.0397E+07 6.6760E+05 2.5057E+07 8.2039E+07 7.2224E+07 3.4074E+08  
 GROUP 7... 6.3804E+01 9.3843E+01 3.9292E+07 3.7119E+07 7.7008E+07 1.5592E+08 4.7684E+08 9.5851E+08  
 GROUP 8... 6.1616E+00 7.0919E+01 5.2727E+07 5.6279E+07 3.4559E+08 3.5923E+08 5.9558E+08 1.0521E+09  
 GROUP 9... 6.0201E+00 5.9253E+01 8.3204E+07 4.4393E+07 3.3149E+08 1.7511E+08 6.7875E+08 1.0560E+09  
 GROUP 10... 5.9253E+00 5.9253E+01 8.3204E+07 4.4393E+07 3.3149E+08 1.7511E+08 6.7875E+08 1.0560E+09  
 GROUP 11... 5.6227E+01 2.4982E+01 2.6943E+07 7.2303E+04 2.4321E+08 6.0769E+07 7.0505E+08 1.1527E+09  
 GROUP 12... 6.0946E+00 1.6553E+01 1.7525E+07 2.8557E+06 2.3961E+08 3.9662E+07 8.0768E+08 1.1372E+09  
 GROUP 13... 6.0258E+00 9.1025E+01 6.1625E+02 4.9271E+06 7.9608E+08 4.8783E+07 8.8920E+08 1.1212E+09  
 \*\*\*\*\*FLUXES FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 RAD(TIS)/HR RAD(ETH)/HR  
 REM  
 GROUP 1... 3.8775E+01 5.8153E+01 2.4773E+02  
 GROUP 2... 3.8775E+01 5.8153E+01 2.4773E+02  
 GROUP 3... 6.5356E+02 2.9276E+02 2.3406E+03  
 GROUP 4... 2.9398E+02 4.1316E+02 2.0914E+03  
 GROUP 5... 4.8044E+02 7.2017E+02 3.7897E+03  
 GROUP 6... 4.8044E+02 7.2017E+02 3.7897E+03  
 GROUP 7... 3.1282E+02 2.2252E+02 3.2520E+03  
 GROUP 8... 5.2240E+02 7.4438E+02 5.4310E+03  
 GROUP 9... 4.7915E+02 6.9010E+02 5.0044E+03  
 GROUP 10... 4.7915E+02 6.9010E+02 5.0044E+03  
 GROUP 11... 1.0816E+02 1.5771E+02 1.0489E+03  
 GROUP 12... 5.4889E+01 6.3399E+01 5.1908E+02  
 GROUP 13... 5.4889E+01 6.3399E+01 5.1908E+02  
 TOTALS 1... 7.4285E+03 1.6805E+03 6.4113E+04  
 MIN ERROR 3... 3.4951E+01 3.2971E+01  
 MAX ERROR 3... 5.8551E+01 5.9717E+01  
 \*\*\*\*\*FLUX MOMENTS FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 \*\*\*\*\*ITRANT 1 REGION 1 REGION 2 REGION 3 REGION 4 REGION 5 REGION 6 REGION 7\*\*\*\*\*  
 GROUP 1... 2.1542E+01 0.10701E+05 0... 6.7930E+04 3.2932E+04 8.0104E+03 3.4432E+04 3.1488E+05  
 GROUP 2... 9.3731E+00 4.9384E+05 0... 3.1241E+05 6.6202E+04 6.1888E+02 4.1745E+04 5.2598E+05  
 GROUP 3... 4.4022E+07 1.1550E+06 0... 1.6401E+06 1.4952E+05 1.4927E+03 2.7378E+06 4.4066E+05  
 GROUP 4... 2.0397E+07 1.5540E+04 0... 7.7371E+05 8.1111E+04 1.3031E+04 5.0814E+03 6.8499E+04

\*\*\*\*\*FLUX MOMENTS FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 REGION 1 REGION 2 REGION 3 REGION 4 REGION 5 REGION 6 REGION 7  
 GROUP 5... 3.9871E+07 5.8543E+03 0... 8.1361E+04 7.1213E+02 1.7872E+02 2.1450E+03 3.8481E+04  
 GROUP 6... 3.6454E+07 6.3833F+02 0... 7.4154E+05 7.0036E+04 1.3412E+02 3.7548E+02 1.3499E+04  
 GROUP 7... 3.9628E+07 3.7502E+00 0... 3.5534E+04 2.6404E+03 2.0046E+01 2.5147E+02 3.5216E+02  
 GROUP 8... 3.9628E+07 3.7502E+00 0... 3.5534E+04 2.6404E+03 2.0046E+01 2.5147E+02 3.5216E+02  
 GROUP 9... 6.3204E+07 4.5115E+04 0... 1.0214E+03 1.9553E+02 1.3592E+04 1.0996E+03 1.0399E+01  
 GROUP 10... 8.2904E+07 4.0476E+07 0... 4.7798E+01 9.3432E+00 8.2124E+02 1.0224E+05 0.0074E+01  
 GROUP 11... 2.0433E+07 5.2620E+10 0... 3.0208E+00 5.7751E+01 1.0335E+02 4.7224E+05 1.0524E+02  
 GROUP 12... 7.7237E+07 2.3824E+04 0... 6.1213E+00 8.1580E+01 5.8051E+05 3.5804E+10 1.0000E+00  
 GROUP 13... 8.0456E+07 1.9565E+27 0... 5.7393E+09 5.9632E+09 3.4940E+11 9.4635E+11 5.9498E+12  
 RAD(TIS)/HR 7.0524E+01 5.6557E+01 5.9717E+01  
 RAD(ETH)/HR 9.8050E+03 3.6268E+00 0... 6.1521E+01 7.0189E+00 4.5378E+02 1.8886E+00 2.3407E+01  
 RAD(ETH)/HR 9.8050E+03 1.1625E+01 0... 6.5194E+01 6.0944E+00 5.6392E+02 2.7224E+00 3.7706E+01  
 REGION 8 REGION 9 REGION 10 REGION 11 REGION 12 REGION 13 REGION 14 REGION 15 SCATTER 0 SCATTER 1 SCATTER 2  
 GROUP 1... 1.0743E+07 1.2740E+03 1.4855E+07 1.0245E+07 1.0245E+07 1.0245E+07 1.0245E+07 7.4322E+05 3.1222E+05  
 GROUP 2... 3.5377E+04 2.7820E+03 3.9191E+04 1.0707E+06 1.0548E+06 1.0232E+05 1.0495E+06 5.5616E+06  
 GROUP 3... 3.6268E+03 1.0525E+06 1.2195E+07 2.4484E+07 1.7171E+07 1.0568E+07 0.5107E+06  
 GROUP 4... 9.6821E+03 6.3579E+02 1.6573E+06 1.1542E+07 2.3248E+06 1.0200E+04 1.2089E+07 1.7964E+06  
 GROUP 5... 2.0202E+03 1.5444E+06 3.5215E+06 3.1816E+06 1.0596E+07 3.7495E+13 4.0482E+06 6.8539E+04  
 GROUP 6... 3.4499E+03 9.6592E+01 3.0649E+07 1.1204E+08 2.1102E+08 1.8581E+02 3.9937E+05 3.2742E+08  
 GROUP 7... 1.0734E+02 1.9387E+01 6.6751E+06 2.9411E+07 1.5037E+06 6.1613E+01 3.9155E+05 3.9441E+05  
 GROUP 8... 2.4638E+01 3.2720E+01 1.0141E+07 2.4395E+07 4.4708E+07 2.0595E+02 2.0570E+06 3.0062E+05  
 GROUP 9... 2.4638E+01 3.2720E+01 1.0141E+07 2.4395E+07 4.4708E+07 2.0595E+02 2.0570E+06 3.0062E+05  
 GROUP 10... 9.4901E+02 5.3207E+04 6.5236E+06 5.9885E+06 1.0392E+07 4.4951E+08 9.4461E+03 1.2820E+05  
 GROUP 11... 4.2501E+02 6.26677E+05 6.3433E+06 6.2469E+06 6.0352E+07 6.0304E+11 1.0569E+05 1.6457E+05  
 GROUP 12... 1.4156E+05 1.5444E+06 3.5215E+06 3.1816E+06 1.0596E+07 3.7495E+13 4.0482E+06 6.8539E+04  
 GROUP 13... 3.4499E+03 9.6592E+01 3.0649E+07 1.1204E+08 2.1102E+08 1.8581E+02 3.9937E+05 3.2742E+08  
 GROUP 14... 2.1448E+02 3.6268E+03 9.2217E+02 3.6095E+03 5.0844E+03 4.2290E+01 3.9418E+02 3.6337E+02  
 GROUP 15... 1.0734E+02 1.9387E+01 6.6751E+06 2.9411E+07 1.5037E+06 6.1613E+01 3.9155E+05 3.9441E+05  
 \*\*\*\*\*SCATTER FLUX MOMENTS FOR DETECTOR 1 AFTER 400 PACKETS\*\*\*\*\*  
 SCATTER 3 - SCATTER 4 - SCATTER 5 - SCATTER 6 - SCATTER 7 - SCATTER 8 ANGULAR 1 SPATIAL 1  
 GROUP 1... 1.4198E+03 8.1786E+04 2.1512E+03 5.5837E+02 1.9107E+01 6.9572E+00 2.1106E+02 2.1267E+01  
 GROUP 2... 9.4312E+03 6.3737E+05 1.7153E+04 5.2232E+06 1.4814E+05 1.4814E+05 3.7195E+07 9.4304E+06  
 GROUP 3... 6.4472E+03 1.0433E+06 2.1413E+06 5.3219E+06 1.6743E+06 5.0790E+04 1.7678E+07 2.3111E+07  
 GROUP 4... 1.7303E+07 1.2526E+05 8.1914E+05 6.1529E+05 8.6202E+05 2.8552E+05 3.3925E+07 4.0546E+07  
 GROUP 5... 1.0734E+02 3.6268E+03 9.2217E+02 3.6095E+03 5.0844E+03 4.2290E+01 3.9418E+02 3.6337E+02  
 GROUP 6... 1.0734E+02 3.6268E+03 9.2217E+02 3.6095E+03 5.0844E+03 4.2290E+01 3.9418E+02 3.6337E+02  
 GROUP 7... 9.4355E+03 1.0393E+06 2.1409E+06 5.3203E+06 1.7603E+06 1.9039E+05 3.4648E+07 4.6038E+07  
 GROUP 8... 6.3369E+06 9.0255E+05 2.9981E+05 2.3101E+07 2.1016E+06 2.5023E+05 7.4389E+07 8.9989E+07

SAMPLE PROBLEM 2: NEUTRON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY  
EMPTY PROPELLANT TANK, REFERENCE ER=0236 (LOCKHEED STUDY FOR NASA-HSFC) \*\*\*\*\*THE FASTER CODE\*\*\*\*\*CASE 2  
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GROUP 9...	6.1157E+06	1.1127E+06	5.2256E+05	2.1237E+07	2.7154E+04	3.9784E+05	7.4164E+07	9.4468E+07
GRDUP 10...	1.7787E+06	3.1560E+05	1.3075E+05	6.4893E+06	7.5750E+05	1.1898E+05	8.0415E+07	9.1955E+07
GROUP 11...	2.3382E+06	4.0064E+05	1.3888E+05	4.1991E+06	1.1307E+06	5.4589E+05	2.7220E+07	3.2730E+07
GROUP 12...	7.0102E+05	9.5089E+05	1.9700E+05	4.4313E+05	2.9125E+05	2.9125E+05	1.3487E+07	1.4930E+07
GROUP 13...	1.1848E+05	2.5035E+07	1.9705E+05	5.1564E+05	2.7220E+05	2.3000E+05	3.9924E+07	4.9830E+07
RAD(T15)/HR 1...	1.1822E+03	3.0003E+02	1.1302E+02	1.6404E+03	1.5000E+02	4.9203E+01	6.5119E+03	7.9730E+03
RAD(ETH)/HR 1...	1.6421E+03	4.2143E+02	1.5776E+02	2.2738E+03	2.0526E+02	6.7997E+01	9.0401E+03	1.1049E+04
REH/HOUR 1...	0.0392E+03	2.4640E+03	6.9632E+02	1.5576E+04	1.3635E+03	3.9158E+02	6.0395E+04	7.3919E+04

SAMPLE PROBLEM 2: NEUTRON DOSE RATE AT POINT DETECTOR ABOVE PARTIALLY \*\*\*\*\*THE FASTER CODE\*\*\*\*\*CASE 2  
EMPTY PROPELLANT TANK, REFERENCE ER=0236 (LOCKHEED STUDY FOR NASA-HSFC) \*\*\*\*\*T.M.JORDAN/WANL\*\*\*\*\*PAGE 9

\*\*\*\*\*PROBLEMS\*\*\*\*\*BOUNARY SEARCH PARAMETERS: (SURFACE,MOST PROBABLE NEXT REGION)\*\*\*\*\*

REGION 1	{ -3, 2)	{ 4, 4)	{ 1, 3)
REGION 2	{ -2, 9)	{ 3, 0)	{ 16, 3)
REGION 3	{ -2, 9)	{ 4, 6)	{ -10, 21) { 17, 10)
REGION 4	{ -5, 14)	{ 1, 14)	{ 17, 21)
REGION 5	{ -5, 4)	{ 6, 6)	{ 17, 10)
REGION 6	{ -6, 5)	{ 7, 7)	{ 17, 10)
REGION 7	{ -7, 6)	{ 8, 8)	{ 17, 10)
REGION 8	{ -8, 5)	{ 9, 11)	{ 17, 10)
REGION 9	{ -11, 31)	{ 2, 0)	{ 20, 31)
REGION 10	{ -17, 7)	{ 20, 31)	{ -2, 9) { 9, 11)
REGION 11	{ -19, 10)	{ 10, 12)	{ 20, 31)
REGION 12	{ -11, 16)	{ 11, 16)	{ -6, 9)
REGION 13	{ -10, 0)	{ 18, 14)	{ 21, 16) { -22, 18)
REGION 14	{ 14, 0)	{ -18, 17)	{ 19, 19) { -22, 0) { -10, 0)
REGION 15	{ -12, 15)	{ 14, 17)	{ -18, 14) { 21, 13)
REGION 16	{ -13, 16)	{ 14, 20)	{ 18, 14)
REGION 17	{ -13, 16)	{ 19, 19)	{ 22, 13)
REGION 18	{ -10, 11)	{ 19, 21)	{ 22, 13)
REGION 19	{ -19, 0)	{ 20, 31)	{ -10, 11) { 14, 20)
REGION 20	{ -14, 0)	{ 15, 31)	{ 20, 31)

by section, and terminates at page 4 of case 1 (upper right hand corner of the printout). Every page after the first has the title cards printed at the top of the page.

#### Flux Output

The next page, (5), of the printout starts the output of the computed fluxes by detector and energy group. The various calculated flux components are separated by lines containing asterisks.

The first set of output, labeled: \*\*FLUXES FOR DETECTOR XXX AFTER XXX PACKETS \*\*, is always printed first. (The term packet is just another way of describing an outer iteration; i.e., each energy dependent angular point source can be visualized as a packet of particles of different energies.) Each of the columns contains groupwise information as indicated on the left side of the printout. Note: there are less flux groups than were used in running the problem; i.e., 23 groups were collapsed to 11, as shown on Card 5-2. The first column is the coefficient of variation of the scalar flux, expressed as a fraction. The second column is the average energy in Mev. The third column is the number flux in particles/cm<sup>2</sup> sec unless the input data were juggled; e.g., using unit areas and volumes for surface and volume detectors. The fourth column contains the energy flux in Mev/cm<sup>2</sup> sec. The fifth and sixth columns contain the group averaged differential number and energy fluxes, i.e., columns 3 and 4, respectively, divided by the group width. The last two columns are running summations of the fluxes and give the total flux from particles with energies greater than the lower boundary of the group.

The next set of output is labeled: \*\* NUMBER FLUX RESPONSES FOR DETECTOR XXX AFTER XXX PACKETS \*\*. Only one response function was input and its title (entered on Card 5-3) appears over the first column. The response by group is then printed. The last three lines in this column are the total response, the coefficient of variation as if the group-wise responses were obtained independently (labeled MIN ERROR), and the coefficient of variation as if the group wise responses were strictly dependent (labeled MAX ERROR).

This set of response function output would be omitted if no response functions were input. It would be repeated, until all response functions were output, for problems with more than 8 input responses.

The final set of output is labeled \*\* NUMBER FLUX MOMENTS FOR DETECTOR XXX AFTER XXX PACKETS \*\*. Column 1 is always obtained in this set of output. It contains the average fluxes obtained between printouts. For this problem it is identical to the scalar flux. For problems with multiple printouts during the calculation, say every 100 outer iterations, this column would contain the average scalar flux as computed during the most recent 100 iterations. Subsequent columns contain other components of the total flux in the following order:

- a) total flux by source, labeled SOURCE XXX (not used in this problem since only 1 source was present)
- b) total scattered flux by scattering region, labeled REGION XXX (obtained in this problem for every nonvoid region)
- c) total flux by order-of-scatter, labeled SCATTER XXX (obtained through the 7th scatter for this problem)
- d) Legendre moments of the angular flux, labeled ANGULAR XXX (the current was obtained for this problem)
- e) length-of-flight moments of the flux, labeled SPATIAL XXX (first moment obtained for this problem)

After printing the contribution by group, each column contains the total contribution to each of the response functions, if any.

The final page of printout for a problem (page 7 of this printout) is sometimes helpful in correcting geometric errors. It contains a listing, by region, of the bounding surfaces (with the sign of the ambiguity index fixed) and the region entered the last time a ray crossed these boundaries (most-probable-next-regions). Most-probable-next-region indices greater than the total number of regions (greater than 20 for this problem) indicate that there was no next region and should correspond to the outer boundaries of the problem. A zero indicates the boundary was never crossed.

The remainder of this printout, labeled CASE 2, is the neutron problem which was run as a change case immediately behind the first case. Most of the input changes involve

the multigroup scattering cross sections. The flux printout (starting on page 6 of case 2) is interpreted in exactly the same manner as case 1.

The results contained in this printout are summarized in Table 3 along with appropriate results from Reference 20. Since the FASTER calculation was intended as a sample problem, comments on the relative accuracy and computer time requirements of the various calculations are not made.

TABLE 3  
SUMMARY OF SAMPLE PROBLEM RESULTS

	FASTER Monte Carlo	18-0 * Monte Carlo	Point Kernel
Photons (rad/hr)	$3.4 \times 10^3 \pm 0.34 \times 10^3$	$\sim 6.7 \times 10^3 \pm 3.0 \times 10^3$	$\sim 3.0 \times 10^3$
Neutrons (rad/hr)	$7.1 \times 10^3 \pm 2.7 \times 10^3$	$\sim 2.4 \times 10^4 \pm 1.5 \times 10^4$	$\sim 2.0 \times 10^2$ (14-0, A.W.) $\sim 5.0 \times 10^2$ (QAD, H <sub>2</sub> O) $\sim 6.6 \times 10^2$ (QAD, C)

\*Reference 20, Figures 30 and 55

## SECTION

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## APPENDIX A

## ALTERNATE MONTE CARLO PROCEDURES

Section 2.0 of this volume pertained to a development of the Monte Carlo method which utilized random sampling for all of the spatial integrations. This method is used in the FASTER program. It is not necessarily the most efficient procedure since there are many alternatives. An alternate technique for integrating the spatial dependence of the distributed fixed source is discussed below. The discussion is limited to 1) a technique which further reduces the integrations performed by random sampling and, 2) the consideration of a single point detector.

Alternate Uncollided Angular Flux Estimator

Equation for the uncollided scalar flux can be manipulated as:

$$\phi_o(\vec{r}, E) = \iint_{4\pi} \left\{ \frac{\phi_o(\vec{r}, \vec{\Omega}, E)}{q^*(\vec{\Omega})} \right\} q^*(\vec{\Omega}) d\Omega \quad (A.1)$$

$$\approx \frac{1}{n} \sum_{i=1}^n \phi_o^*(\vec{r}, \vec{\Omega}_i, E) \quad (A.2)$$

where  $\phi_o^*(\vec{r}, \vec{\Omega}_i, E) = \frac{1}{q^*(\vec{\Omega}_i)} \phi_o(\vec{r}, \vec{\Omega}_i, E)$  (A.3)

$\vec{\Omega}_i$  selected at random from  $q^*(\vec{\Omega})$  (A.4)

$$\left. \begin{aligned} q^*(\vec{\Omega}) &\geq 0 \\ q^*(\vec{\Omega}) &> 0 \quad \text{if} \quad \int_0^\infty \phi_o(\vec{r}, \vec{\Omega}, E) dE > 0 \end{aligned} \right\} \quad (A.5)$$

$$\iint_{4\pi} q^*(\vec{\Omega}) d\Omega = 1$$

Thus the angular flux is approximated for finite  $n$  by:

$$\phi_0^*(\vec{r}, \vec{\Omega}_1, E) = \frac{1}{n} \sum_{i=1}^n \phi_0^*(\vec{r}, \vec{\Omega}_1, E) \delta(\vec{\Omega} - \vec{\Omega}_i) \quad (\text{A.6})$$

But  $\phi_0^*(\vec{r}, \vec{\Omega}_1, E)$  is given by a simple line integral, i.e., equation 2.2:

$$\phi_0^*(\vec{r}, \vec{\Omega}_1, E) = \frac{1}{q(\vec{\Omega}_1)} \int_0^\infty S_0(\vec{r} - s\vec{\Omega}_1, \vec{\Omega}_1, E) \exp\left[-\int_0^s \Sigma^t(\vec{r} - s'\vec{\Omega}_1, E) ds'\right] ds \quad (\text{A.7})$$

This integration can be performed numerically with extreme accuracy for non-vacuum source regions with slowly varying source distributions by using the transformation:

$$du = \exp\left[-\int_0^s \Sigma^t(\vec{r} - s\vec{\Omega}_1, E) ds'\right] ds \quad (\text{A.8})$$

i.e., the discrete values of  $s$  corresponding to discrete values of  $u$  will be exponentially distributed with approximately equal contributions to the integration. Note: the transformation is performed for some average energy  $\bar{E}$  thus yielding the same "source" points for all energies. Solving A.8 for  $ds$  and substituting into A.7 gives the transformed equation:

$$\phi_0^*(\vec{r}, \vec{\Omega}_1, E) = \frac{1}{q(\vec{\Omega}_1)} \int_0^\infty S_0(\vec{r} - s(u)\vec{\Omega}_1, \vec{\Omega}_1, E) \exp\left[\int_0^{s(u)} \left\{ \Sigma^t(\vec{r} - s'\vec{\Omega}_1, \bar{E}) - \Sigma^t(\vec{r} - s'\vec{\Omega}_1, E) \right\} ds'\right] du \quad (\text{A.9})$$

where  $s(u)$  is the solution for  $s$  of

$$u = \int_0^{s(u)} \exp\left[-\int_0^{s'} \Sigma^t(\vec{r} - s\vec{\Omega}_1, \bar{E}) ds'\right] ds \quad (\text{A.10})$$

This procedure can also be generalized to include the effects of the source distribution in the transformation. All of these considerations have been incorporated into the sampling function described in Section 8.3, which approximates equation 2.36. The only difference is: random

discrete values of  $s$  (or  $u$ ) are obtained for performing the numerical integration.

Another procedure for computing the uncollided angular flux, is to perform all integrations numerically. When an angular integration is included in a rotated spherical coordinate system centered at the detector point, smaller computer times are expended than on most point kernel calculations. This is a result of the many "point sources" having the same "line-of-sight" to the detector; i.e., the geometric calculations are significantly reduced.

#### Alternate Single Scattered Flux Estimator

The use of the above technique does not yield a single point source which can be used to calculate the point representation of the scattered source at the first scattering point. However, the same technique can be applied at the first scattering point. The procedure, again, is to write the equation for the single scattered scalar flux. (This formalism isn't really necessary but it provides a convenient and consistent basis for the application of random sampling techniques. It is useful in inferring the form of optional sampling functions.)

$$\phi_1^*(\vec{r}, E) = \iiint \left\{ \frac{S_1(\vec{r}, \vec{\Omega}, E) \exp\left[-\int_0^s \Sigma^t(\vec{r} - s\vec{\Omega}, E) ds'\right]}{p_1(\vec{r}) s^2} \right\} p_1^*(\vec{r}') dV \quad (\text{A.11})$$

which yields a point, single-scattered source

$$W_1^*(\vec{\Omega}, E) = \frac{1}{p_1^*(\vec{r}')} S_1(\vec{r}', \vec{\Omega}, E) \quad (\text{A.12})$$

where  $\vec{r}'$  is selected at random from  $p_1^*(\vec{r}')$  as shown in Section 2.4.

In Section 2.4 it was assumed that there was an energy dependent angular point source,  $W_0(\vec{\Omega}, E)$ , at  $\vec{r}_0$  which finally yielded  $W_1^*(\vec{\Omega}, E)$ . The point source  $W_0(\vec{\Omega}, E)$  can be obtained at this point in the calculation although another technique is used for the uncollided flux. Nevertheless, other techniques can also be used. In particular, the technique discussed for the uncollided scalar flux can also be used for the single scattered source

$$W_1^*(\vec{r}, E) = \frac{1}{P_1(\vec{r}_1)} \int \left\{ \int_0^\infty \phi_0(\vec{r}_1, \vec{\Omega}', E) \frac{d^2 \Sigma^*}{d\Omega dE} (\vec{r}_1, \vec{\Omega}', E \leftarrow \vec{\Omega}, E) dE' \right\} q^*(\vec{\Omega}') d\Omega' \quad (A.13)$$

$$= \frac{1}{P_1(\vec{r}_1)} \frac{1}{q^*(\vec{\Omega}_1)} \int_0^\infty \phi_0(\vec{r}_1, \vec{\Omega}_1, E') \frac{d^2 \Sigma^*}{d\Omega dE} (\vec{r}_1, \vec{\Omega}_1, E' \leftarrow \vec{\Omega}, E) dE' \quad (A.14)$$

where the equality holds in the sense of the expected value. Multiple scattered flux contributions can be obtained using the remainder of the techniques described in Section 2.4.

There is an implicit assumption in the concept of using a line integral along a fixed direction vector that the angular variations can be approximated with facility in  $q^*(\vec{\Omega})$ . Intuitively, this appears correct since the angular flux at a point detector usually varies slowly within the solid angle cone subtended by the source. Especially, fluxes at a point detector in the most intense portion of a source could be computed with accuracy using only a few discrete directions.

The amount of detail which can be built into sampling functions internal to a computer program is limited and is usually based on line-of-sight observations between source or scattering points and the detector point. For more difficult problems, where line-of-sight estimates are not strictly correct, i.e., where short circuiting around a shield actually yields the major component of the flux, the effects can be approximated by using several alternate definitions of  $P_1^*(\vec{r})$ , where  $P_1^*(\vec{r})$  is used in randomly selecting the first-scatter points. This same capability is not usually warranted for multiple scattering events since the effect is harder to predict and control.

## APPENDIX B

## ANGULAR AND SPATIAL FLUX MOMENTS

Legendre moments of the angular flux are obtained by an integration of the individual flux contributions over solid angle. Using the notation of Section 7.0, the angular dependence of the flux contributions is denoted by:

$$\Delta \phi_{i,j,k}^*(\vec{\Omega}) = \Delta \phi_{i,j,k} \delta(\vec{\Omega} - \vec{\Omega}_o)$$

where  $\Delta \phi_{i,j,k}$  is the flux contribution to the  $j$ th energy group from the  $k$ th inner iteration of the  $i$ th outer iteration and is obtained for the fixed direction  $\vec{\Omega}_o$ .

The angular moments are averaged over the azimuthal angle and obtained with respect to a preferred direction  $\vec{\Omega}_p$ :

$$\phi_i^l = \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^{\infty} \int \int \Delta \phi_{i,j,k} \delta(\vec{\Omega} - \vec{\Omega}_o) P_j(\vec{\Omega} \cdot \vec{\Omega}_p) d\Omega \quad (B.1)$$

$$= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^{\infty} \Delta \phi_{i,j,k} P_j(\vec{\Omega}_o \cdot \vec{\Omega}_p) \quad (B.2)$$

$$\begin{aligned} \text{where } P_o(\psi) &= 1 \\ P_1(\psi) &= \\ P_l(\psi) &= \frac{1}{T} [(2l-1) P_{l-1}(\psi) - (l-1) P_{l-2}(\psi)] \end{aligned} \quad (B.3)$$

The preferred direction  $\vec{\Omega}_p$  is fixed by input for point and volume detectors. For surface detectors, it is the unit normal to the surface,  $\vec{n}_i$  as defined by equation 4.11. The zeroth moment,  $l=0$ , is the scalar flux as given by equation 7.6 and the first moment,  $l=1$ , is the particle current with respect to  $\vec{\Omega}_p$ .

The azimuthally averaged Legendre moments can be used to define azimuthally averaged differential angular fluxes and/or interval integrated values. This is not done by the FASTER program. However, the equations required for external manipulation of these moments are summarized below:

a) differential angular flux

$$\phi_i(\mu) = \sum_{l=0}^{\infty} \frac{2l+1}{4\pi} \phi_i^l P_l(\mu) \quad (B.4)$$

b) interval integrated flux,  $\mu_m \leq \mu \leq \mu_{m+1}$

$$\bar{\phi}_{i,m} = \sum_{l=0}^{\infty} \frac{2l+1}{2} \phi_i^l \int_{\mu_m}^{\mu_{m+1}} P_l(\mu) d\mu \quad (B.5)$$

Using the relationship (Reference 15, page 309):

$$P_l(\mu) = \frac{1}{2l+1} \left\{ \frac{dP}{d\mu} |_{l+1}(\mu) - \frac{dP}{d\mu} |_{l-1}(\mu) \right\} \quad (B.6)$$

Then

$$\bar{\phi}_{i,m} = \frac{1}{2} \sum_{l=0}^{\infty} \phi_i^l \left\{ P_{l+1}(\mu_{m+1}) - P_{l+1}(\mu_m) + P_{l-1}(\mu_{m+1}) - P_{l-1}(\mu_m) \right\} \quad (B.7)$$

#### Spatial Moments

Length-of-flight moments of the flux can also be obtained from FASTER. These are normalized to a reference distance  $t_o$ . Thus

$$\bar{\phi}_i^m = \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^{\infty} \Delta \phi_{i,k} \left( \frac{t_{i,k}}{t_o} \right)^m \quad (B.8)$$

$$t_o = \max(|\vec{r}_p - \vec{r}_c|, 100), \quad \vec{r}_p, \vec{r}_c \text{ are defined in Section 8.3} \quad (B.9)$$

$$t_{i,k} = \sum_{k'=0}^k |\vec{r}_{i,k'} - \vec{r}_{i,k'-1}| + \Delta t_{i,k} \quad (B.10)$$

where

$$\begin{aligned} \Delta t_{i,k} &= |\vec{r} - \vec{r}_{i,k}| \quad \text{for a point detector} \\ &= s, \text{ the distance to a surface detector} \end{aligned} \quad (B.11)$$

Moments of the average length-of-flight are obtained for volume detectors,

$$\begin{aligned} \Delta t_{i,k} &= s + \Delta s/2 \quad \text{if void} \\ &= s + 1/\Sigma_i^V - \Delta s / [\exp(\Sigma_i^V \Delta s) - 1] \quad \text{if not void} \end{aligned} \quad (B.12)$$

For photons and non-volume detectors, the length-of-flight moments can be related to the time of arrival of the flux

$$t_{i,k} = c \tau_{i,k}, \quad \tau_{i,k} = t_{i,k} / c \quad (B.13)$$

where  $c$  is the velocity of light and  $\tau$  is the time. If the source is assumed to have a time dependence  $\delta(\tau)$ , then, the temporal moments are given by

$$\begin{aligned} \hat{\phi}_i^m &= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^{\infty} \Delta \phi_{i,i,k} \tau_{i,k}^m \\ &= \frac{1}{n} \sum_{i=1}^n \sum_{k=0}^{\infty} \Delta \phi_{i,i,k} \left( \frac{t_{i,k}}{c} \right)^m \\ &= \left( \frac{t_o}{c} \right)^m \bar{\phi}_i^m \end{aligned} \quad \left. \right\} \quad (B.14)$$

These moments are not processed by the FASTER program but the obvious procedure is to use techniques of the type used for the spatial dependence in the moments<sup>(6)</sup>; e.g., assume a functional form and match moments.

$$\begin{aligned} \phi_i(\tau) &= \exp[-\alpha(\tau - \tau_o)] \sum_{m'=0} b_{m'} (\tau - \tau_o)^{m'}, \quad \tau \geq \tau_o \\ &= 0, \quad \tau < \tau_o \end{aligned} \quad \left. \right\} \quad (B.15)$$

where  $\tau_0$  is the earliest arrival time,  $\alpha$  is an exponential decay constant, and the  $b_m$ 's are constants in a polynomial representation of the departure from a true exponential decay.

Then

$$\begin{aligned}\hat{\Phi}_i^m &= \int_{\tau_0}^{\infty} \phi_i(\tau) \tau^m d\tau \\ &= \int_{\tau_0}^{\infty} \exp[-\alpha(\tau - \tau_0)] \sum_{m'=0} b_{m'} (\tau - \tau_0)^{m'} \tau^m d\tau \\ &= \sum_{m'=0} b_{m'} c_{m, m'}, \quad m = 0, 1, \dots\end{aligned}\quad (B.16)$$

where

$$c_{m, m'} = \int_{\tau_0}^{\infty} \exp[-\alpha(\tau - \tau_0)] \tau^m (\tau - \tau_0)^{m'} d\tau \quad (B.17)$$

Equation B.16 is solved for the constants  $b_{m'}$ ,  $m' = 0, 1, \dots$  and equation B.15 is used to compute the temporal dependence.

## APPENDIX C PROGRAM LISTING

The FORTRAN IV listing of the FASTER program is given in this appendix. Compatibility with other computer facilities can be obtained by the following changes in the control program on the first page of the listing:

- a) Change of input tape logical designation from 5 to 1  
replace M1 = 5 by M1 = i
- b) Change of output tape logical designation from 6 to 1  
replace M2 = 6 by M2 = i
- c) Change of maximum number of lines per printout page from 43 to k  
replace LINEX = 43 by LINEX = k
- d) Change of maximum number of locations for dimensioned arrays from 12000 to l.  
replace COMMON H(12000) by COMMON H(l)  
and replace NSTORE = 12001 by NSTORE = i + 1

The listing corresponds to an operational program for the IBM 7094 computer which uses a MAP random number generator as shown on the last page of the listing. For conversion to the CDC 6600 computer, the MAP routine is removed and a card in the control program (first page of listing) is changed:

replace IBMCDC = 0 by IBMCDC = 1

It is assumed that the random number generator, RANF, distributed by the Control Data Corporation is on the library tape and that the calling sequence is:

R = RANF(i), where i > 0, stores i as the generator

R = RANF(0) yields a random number R on (0, 1)

All calls to the random number generator are relayed through the function subprogram RANNO(n) where n is a do-nothing argument--is not used in the calling subprograms and can be defined arbitrarily in RANNO. Therefore, any other random number generator can be used by FASTER with appropriate changes in RANNO.

FORTRAN IV LISTING IS PRESENTED ON PAGES 184 THROUGH 259.



```

      NVMAX,NVOMCC,NPCINT,MODELPL,MODELQ,MODELU,MODELV, FAST0094
      NPRINT,NUNITS,NUMBER,KALIDE   FAST0095
COMMON/INDEXS/INTP ,IAZ ,ISV ,INTL ,IRHO ,IXR ,IELL , FAST0096
      IELB ,IAT ,IBE ,ISUV ,INTR ,ISCH ,ILSX ,FAST0097
      IXTR ,ISUV ,IATM ,IATW ,ISMH ,ICLN , FAST0098
      INTG ,IELF ,IFGW ,ITDS ,IIDIW ,IDOS ,IIDS , FAST0099
      IVOL ,ICDT ,IXCT ,IRS1 ,IALP ,IVMD ,IGIM , FAST0100
      IAIM ,ITALM ,IHAL ,IHS ,INS ,IVEE ,IVAL , FAST0101
      ISPN ,ISP8 ,IATD ,ISHP ,ISGT ,IIDE ,ITDI , FAST0102
      ISGS ,NEXT ,IXHP ,ISGR ,IWS ,IFS ,IWT ,FAST0103
      FEC ,IND ,IUE ,IV ,IWH ,ISI ,IST ,FAST0104
      INRG ,INSC ,LSP8 ,INRP ,INCP ,IFXP ,IFXS ,FAST0105
      IFXT ,IXE ,IFXA ,MM ,X(5)  FAST0106
COMMON/DUMLLC/NH   FAST0107
      COMMON/DUMLLC/NH ,MM ,X(5)  FAST0108
      D3 10 I=1,3  FAST0109
1) X111 = XXXX11  FAST0109
      NN = NNN  FAST0110
      CALL LOCUDM(NAHAX,NBMAX,1,H(INTP),H(IAZ),H(INS),H(IND),  FAST0111
      LOCATE = MM  H(TU))  FAST0112
      RETURN  FAST0113
      END  FAST0114
      FAST0115
$IRFTC LOCUDM_NH/2,XRT  FAST0116
LOCUDM//REGISTRY INDEX-CALCULATION FOR PROGRAM FASTER*T.H.JORDAN*WANL*196  FAST0117
      COMMON/INDEXS/INTP ,IAZ ,ISV ,INTL ,IRHO ,IXR ,IELL , FAST0118
      IELB ,IAT ,IBE ,ISUV ,INTR ,ISCH ,ILSX ,FAST0119
      IXTR ,ISUV ,IATM ,IATW ,ISMH ,ICLN , FAST0120
      COMMON/DUMLLC/NH ,NN ,X(5)  FAST0121
      COMMON/LIMITS/NSTCRE,NERRD,NSPA8 ,NAMAX ,NRMAX ,NMMAX ,NSTMAX ,  FAST0122
      1   NEMAX ,NVMAX ,NNMAX ,NEMAX ,NEMODD ,NXSECT ,NUNITS ,  FAST0123
      2   NUNITX ,NMAX ,NMAX ,NMAX ,NMAX ,NMAX ,NTRANS ,  FAST0124
      3   NVMAX ,NMAX ,NEMAX ,NEMODD ,NEMAX ,NMMAX ,NTMAX ,  FAST0125
      4   NDMAX ,NCMOD ,NDPENT ,NDMOD ,NPDMAX ,NPDMOD ,NSDMAX ,  FAST0126
      5   NVDMAX ,NVDCD ,NPPOINT ,MODELPL,MODELQ,MODELU,MODELV,  FAST0127
      NPRINT,NUNITS,NUMBER,KALIDE  FAST0128
      DO 100 I=1,3  FAST0129
      X111*3 = X(I1)*X(I2)  FAST0130
      J = X(I1)*X(I2)/3  FAST0131
100 X111*61 = X(I1)*X(J)  FAST0132
      DO 110 I=1,NMAX  FAST0133
110 ND(I) = 0  FAST0134
      DU 200 N1=1,2  FAST0135
      IF(N.GT.1)GL TO 130  FAST0136
      TMIN = NNN  FAST0137
      IF(110.LT.N1)*(NRMAX-MIN+1).GT.C)GO TO 120  FAST0138
      IMIN = NRMAX + 1  FAST0139
      GO TO 200  FAST0140
120 IMAX = NRMAX

```

```

      GO TO 140  FAST0141
130 IMAX = IMIN - 1  FAST0142
      IMIN = 1  FAST0143
140 IF(110.LT.N1)IMIN,IMAX  FAST0144
      DO 170 J=1,NMAX  FAST0145
      IF(NS(J,I).EQ.0)GC TO 180  FAST0146
      KP = NS(J,I)/1000000  FAST0147
      K = NS(J,I)/1000 - 1000*KP  FAST0148
      IF(ND(K).GT.0)GO TO 160  FAST0149
      ND(K) = 1  FAST0150
      MAX = NM(K)
      U(K) = AZ(K)
      DO 150 L=1,MAX  FAST0151
      U(L) = U(K) + X(L)*A(L,K)  FAST0152
150 IF(FLOAT(I*KP - 1)*U(K).GT.0.0)GO TO 190  FAST0153
160 IF(FLOAT(I*KP - 1)*U(K).GT.0.0)GO TO 190  FAST0154
170 CONTINUE  FAST0155
180 IF(110.LT.N1)GO TO 210  FAST0156
190 CONTINUE  FAST0157
200 CONTINUE  FAST0158
      NH = 0  FAST0159
210 RETURN  FAST0160
      END  FAST0161
      SORIGIN  ALPHA  FAST0162
$IRFTC DEDEFIN_NH/2,XRT  FAST0163
DEFINE//DATA INPUT AND PREPARATION FOR PROGRAM FASTER*T.H.JORDAN*WANL*66  FAST0164
      SUBROUTINE DEFINE .  FAST0165
      COMMON/V11  .  FAST0166
      COMMON/TAPEID/M1 ,M2  FAST0167
      COMMON/CASEID/SFILE ,HPAGE ,LINES ,LINEX ,TITLEA(18) ,TITLEB(18) ,  FAST0168
      COMMON/LIMITS/NSTCRE,NERRD,NSPA8 ,NAMAX ,NMMAX ,NEMAX ,NSTMAX ,  FAST0169
      NUNITX ,NMAX ,NNMAX ,NEMAX ,NEMODD ,NXSECT ,NUNITS ,  FAST0170
      1   NEMAX ,NVMAX ,NNMAX ,NEMAX ,NEMODD ,NXSECT ,NUNITS ,  FAST0171
      2   NUNITX ,NMAX ,NNMAX ,NEMAX ,NEMODD ,NXSECT ,NUNITS ,  FAST0172
      3   NMMAX ,NCMOD ,NDPENT ,NDMOD ,NPDMAX ,NPDMOD ,NSDMAX ,  FAST0173
      4   NVDMAX ,NVDCD ,NPPOINT ,MODELPL,MODELQ,MODELU,MODELV,  FAST0174
      5   NVDMAX ,NVDCD ,NPPOINT ,MODELPL,MODELQ,MODELU,MODELV,  FAST0175
      6   NVDMAX ,NVDCD ,NPPOINT ,MODELPL,MODELQ,MODELU,MODELV,  FAST0176
      COMMON/INDEXS/INTP ,IAZ ,ISV ,INTL ,IRHO ,IXR ,IELL , FAST0177
      IELB ,IAT ,IBE ,ISUV ,INTR ,ISCH ,ILSX ,FAST0178
      IXTR ,ISUV ,IATM ,IATW ,ISMH ,ICLN , FAST0179
      INTG ,IELF ,IFGW ,ITDS ,IIDIW ,IDOS ,IIDS , FAST0180
      IVOL ,ICDT ,IXCT ,IRS1 ,IALP ,IVMD ,IGIM , FAST0181
      IAIM ,ITALM ,IHAL ,IHS ,INS ,IVEE ,IVAL , FAST0182
      ISPN ,ISP8 ,IATD ,ISHP ,ISGT ,IIDE ,ITDI , FAST0183
      ISGS ,NEXT ,IXHP ,ISGR ,IWS ,IFS ,IWT ,FAST0184
      FEC ,IND ,IUE ,IV ,IWH ,ISI ,IST ,FAST0185
      INRG ,INSC ,LSP8 ,INRP ,INCP ,IFXP ,IFXS ,FAST0186
      IFXT ,IXE ,IFXA ,MM ,X(5)  FAST0187

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COMMON/INPUTS/INI ,IN1N(23)
DIMENSION TDUM(36),TYPE(5,6) ,ADUM(19)
DATA TDUM/
 1 4H(F4H),4HTRTA,4H(N,(A+4H)NAL,4H,4HYTIC,4H S1,4HOLUT,4HION,4H(1,4H)*,4H(TDUM
 2 4H(MAN),4HORT,4H(E0,4H)MATT,4HON,I,4H(JAN,4HDUM,4HSAMP,4HLING,4T0193
 3 4H(MAR),4HTRT,4HICAL,4H AND,4H(NUN,4HEKRYC,4HAL,A,4HMLAY,4HSLS,4H
 4 4H(CNC),4HUTER,4H PRG,4HROMA,4HTRG,4H AND,4H CHE,4HCKIU,4HTRT BY/FAST0195
 DATA TYPE/
 1 4H VAR,4HTABL,4HE 0,4HMEMS,4HION,4HION,
 2 4HSURF,4HSEAS,4H ACR,4H REG,4HIGNS,
 3 4H TPD,4HREP,4HIDENT,4H SOUIGNS,
 4 HBAUD,4HTRT,4HTRT,4H TPD,4HTRT,4H
 5 HDBTEB,4HTR,4H S,4HREP,4HSEAS,4HSEAS,
 6 4H SAM,4HPLX,4H GA,4H RAME,4HTERS/
10 KASE = KASE + 1
NPAGE = 0
LNINES = LINEX
1 IF (LNINES .EQ. 1) GO TO 30
DO 20 I=1,5
20 CONTINUE
30 TITLFA(J) = TDUM(J)
31 CONTINUE
DO 700 NNN=1,6
IF (LABEL11.EQ.01 WRITE(12,2000) NNN, (TYPE(I,NNN),I=1,5)
CALL READ12A,1N1
1 IF (INI.LE.01 GO TO 6C
DO 50 I=1,INI
50 CALL READA18,4DUM)
60 GOTO 1100,200,300,400,500,600), NNN
100 CALL SYSTEM
GO TO 700
200 CONTINUE
CALL GEOMIN(NAMAX,NBMAX,1,V(INTP),V(IAZ),V(IA),V((ISV),V(INTL),
1 V(INSL),V(RHO),V(XR)))
200 GO TO 700
300 IE = NXCRE - NMXENX
IEH = IE - NMEXAX
IENG = IE - NMEXMAX
IERG = IE - NMEXMAX
IF (IEBG.LT.NEXT) ERROR = NERRCH + 1
CALL SOURCEIRE(NAMAX,1,V(TEL),V(TELI),V(TELW),V(IAE),V(BF),V(INS),
1 V(INSL),V(LIS),V(LISL),V(LISX),V(LISU),V(LIXTR),V(LIVTE),
2 V(LIVAL),V(LSP),V(VISPE),V(VISL),V(VISV),V(IE, V(ICH)),
3 V(IEBG),V(IFIANG))
300 GO TO 700
400 IEAC = NSTPO - NEHGD*(NMMAX + NMHMAX)

```

```

IJHX = IEAC - NNESEC*|INELAS + 3|
IKSI = IJMX - NNESEC*|INELAS+NTRANS
IXSE = ISXI - NNESEC*|NDDWM*NEMAX*XAO(2,NORDER)
IXST = IXSE - NEMOD
INL = ISXI - ISGS
CALL INSECT(NEMAX,INDEL,NORDER,INELAS,NEMAX,1,V(|IE5B|),V(|ID0|),
1 V(|IS5H|),V(|IS5G|),V(|ID0E|),V(|ID0|),V(|IS5G|),V(|IS5H|),
2 V(|IM1L|),V(|IRHO|),V(|ISXI|),V(|ID0|),V(|IJHX|),V(|ISXI|),
3V(|ISXI|),V(|IEAC|),V(|IL1|),NTRANS)
NEXT = INL + NMAX
IF|NEXT .GT. IXMP| .NE|RRCR = NERRC + 1
GO TO 700
500 CALL RESULT(NGNDL+1,V(|IE1C|),V(|IE1F|),V(|IFGW|),V(|ID0|),V(|IRSP|),
1 V(|ID0|),V(|ID0|),V(|ID0|),V(|ID0|),V(|ID0|),V(|ID0|),
2V(|IM1L|),V(|IRHO|),V(|IEAC|),NEMOD)
GO TO 700
600 CALL RANDOM(V(|ISXI|),V(|ALP|),V(|IMHO|),V(|IGI|),V(|IAH|),V(|ALM|),
1V(|ALH|),V(|ICP|),V(|IE1C|),NMAX,5)
FAST0253
700 CALL RNDM(V(|ISXI|),V(|ALP|),V(|IMHO|),V(|IGI|),V(|IAH|),V(|ALM|),
1V(|ALH|),V(|ICP|),V(|IE1C|),NMAX,5)
FAST0254
IF|LABEL(1).GE.|0|WRITE(H2,10)
2010 FORMAT(1X,I3,35L1H*,36HDATA INPUT AND PREPARATION CMCPLETED,36(IH*))
IF|NERROR_E.C| GO TO 900
IF|LABEL(1).GE.|0|WRITE(H2,2020)
2020 FORMAT(1X,I3,35L1H*,23HPROCEEDING TO NEXT CASE,4Z(1H*))
GO TO 100
900 MAX = 10000
DO 910 I=NEXT,MAX
910 V(1) = 0.0
IF|LABEL(1).GE.|0|WRITE(H2,2040)
2040 FORMAT(1X,I3H*,35I3HG#0)
LINES = LINEX + 1
RETURN
END

$IBFTC STORE M942,XR7
C$TORM*LIMIT INPUT AND VARIABLE DIMENSIONING FOR PROGRAM FASTER*THJ=66
SUBROUTINE STORM
COMMON /N1/
COMMON /PROSE/NMAX,NTRANS
COMMON /TAPE01/M1
COMMON /CASED/KASE ,NPAGE ,LINES ,LINEX ,TITLEA(18) ,TITLEB(18)
COMMON /LIMITS/NSTORE,NERRR,NMSX,NAMAX ,NRMAX ,NBMAX ,NINSTX,
1 NEMAX ,NVMAX ,NMMAX ,NEMAX ,NEMOD ,NSECNT ,NUINTD,
2 NUNITS ,NMAX ,NMMAX ,NEMAX ,NEMOD ,NSECNT ,NUINTD,
3 NMMAX ,NEMAX ,NEMOD ,NSECNT ,NUINTD ,NMAX ,NIMAX ,
4 NDMAX ,NCHD ,NPOINT ,NMDOP ,NMDOP ,NMDOP ,NMDOP ,NMDOP ,
5 NVDMAX ,NHYD ,NPOINT ,NMDOP ,NMDOP ,NMDOP ,NMDOP ,NMDOP ,
6 NPRINT ,NUNITS ,NNUMBER ,KALIDE
FAST0255
FAST0256
FAST0257
FAST0258
FAST0259
FAST0260
FAST0261
FAST0262
FAST0263
FAST0264
FAST0265
FAST0266
FAST0267
FAST0268
FAST0269
FAST0270
FAST0271
FAST0272
FAST0273
FAST0274
FAST0275
FAST0276
FAST0277
FAST0278
FAST0279
FAST0280
FAST0281
FAST0282

```

```

COMMON/INDEXS/ I1SV ,IAZ ,I1SV ,IMTL ,IRHO ,IXR ,IELL ,FAST0283
1   ,IELV ,IAE ,I1SV ,INPC ,IJSN ,IJSK ,IJEN ,FAST0284
2   ,EXTR ,ISUV ,IATN ,IATM ,IJEN ,IJEN ,IJEN ,FAST0285
3   ,INTG ,IEFL ,IFCN ,ITOS ,IZDV ,IJDR ,IJEN ,FAST0286
4   ,IVOL ,ICDT ,IXDT ,IRSI ,IADP ,IWN0 ,IGIM ,FAST0287
5   ,IAIM ,IALM ,IAIH ,IA ,I1NS ,IVEE ,IVAL ,FAST0288
6   ,ISPW ,ISPE ,IATD ,ISPW ,ISGT ,IIDE ,IDDI ,FAST0288
7   ,I0SS ,I1SP ,I1XP ,I1SG ,I1WS ,I1ES ,I1MC ,FAST0289
8   ,IEC ,I1ND ,I1V ,I1W ,I1SI ,I1ST ,I1TC ,FAST0290
9   ,INRG ,INSC ,I1SP ,INRP ,INCP ,IPRP ,IFXS ,FAST0291
10  ,IFXT ,IFXE ,IFXA ,FAST0292
COMMON/INPUTS/ IN1 ,IN2 ,IN3 ,IN4 ,I1NS ,IN6 ,INT ,FAST0293
1  IN ,IN9 ,INIO ,INIL ,INH(13) ,FAST0294
DIMENSION LOPP(10),ARAL(10),LOC(1) ,FAST0295
EQUIVALENCE (LOC(1),INTP)
EQUIVALENCE (LOC(1),INTP)
1(LCP(40),L51),(LCP(44),L6),(LCP(45),L7),(LCP(46),L8),(LCP(47),L9),FAST0296
2(LCP(48),L10),(LCP(49),L11),(LCP(50),L12) ,FAST0297
DATA ARAL /0/ ,FAST0298
DATA ARR1 /0/ ,ARR2 /0/ ,ARR3 /0/ ,ARR4 /0/ ,ARR5 /0/ ,ARR6 /0/ ,ARR7 /0/ ,ARR8 /0/ ,ARR9 /0/ ,ARR10 /0/ ,ARR11 /0/ ,ARR12 /0/ ,ARR13 /0/ ,ARR14 /0/ ,ARR15 /0/ ,ARR16 /0/ ,ARR17 /0/ ,ARR18 /0/ ,ARR19 /0/ ,ARR20 /0/ ,ARR21 /0/ ,ARR22 /0/ ,ARR23 /0/ ,ARR24 /0/ ,ARR25 /0/ ,ARR26 /0/ ,ARR27 /0/ ,ARR28 /0/ ,ARR29 /0/ ,ARR30 /0/ ,ARR31 /0/ ,ARR32 /0/ ,ARR33 /0/ ,ARR34 /0/ ,ARR35 /0/ ,ARR36 /0/ ,ARR37 /0/ ,ARR38 /0/ ,ARR39 /0/ ,ARR40 /0/ ,ARR41 /0/ ,ARR42 /0/ ,ARR43 /0/ ,ARR44 /0/ ,ARR45 /0/ ,ARR46 /0/ ,ARR47 /0/ ,ARR48 /0/ ,ARR49 /0/ ,ARR50 /0/ ,ARR51 /0/ ,ARR52 /0/ ,ARR53 /0/ ,ARR54 /0/ ,ARR55 /0/ ,ARR56 /0/ ,ARR57 /0/ ,ARR58 /0/ ,ARR59 /0/ ,ARR60 /0/ ,ARR61 /0/ ,ARR62 /0/ ,ARR63 /0/ ,ARR64 /0/ ,ARR65 /0/ ,ARR66 /0/ ,ARR67 /0/ ,ARR68 /0/ ,ARR69 /0/ ,ARR70 /0/ ,ARR71 /0/ ,ARR72 /0/ ,ARR73 /0/ ,ARR74 /0/ ,ARR75 /0/ ,ARR76 /0/ ,ARR77 /0/ ,ARR78 /0/ ,ARR79 /0/ ,ARR80 /0/ ,ARR81 /0/ ,ARR82 /0/ ,ARR83 /0/ ,ARR84 /0/ ,ARR85 /0/ ,ARR86 /0/ ,ARR87 /0/ ,ARR88 /0/ ,ARR89 /0/ ,ARR90 /0/ ,ARR91 /0/ ,ARR92 /0/ ,ARR93 /0/ ,ARR94 /0/ ,ARR95 /0/ ,ARR96 /0/ ,ARR97 /0/ ,ARR98 /0/ ,ARR99 /0/ ,ARR100 /0/ ,ARR101 /0/ ,ARR102 /0/ ,ARR103 /0/ ,ARR104 /0/ ,ARR105 /0/ ,ARR106 /0/ ,ARR107 /0/ ,ARR108 /0/ ,ARR109 /0/ ,ARR110 /0/ ,ARR111 /0/ ,ARR112 /0/ ,ARR113 /0/ ,ARR114 /0/ ,ARR115 /0/ ,ARR116 /0/ ,ARR117 /0/ ,ARR118 /0/ ,ARR119 /0/ ,ARR120 /0/ ,ARR121 /0/ ,ARR122 /0/ ,ARR123 /0/ ,ARR124 /0/ ,ARR125 /0/ ,ARR126 /0/ ,ARR127 /0/ ,ARR128 /0/ ,ARR129 /0/ ,ARR130 /0/ ,ARR131 /0/ ,ARR132 /0/ ,ARR133 /0/ ,ARR134 /0/ ,ARR135 /0/ ,ARR136 /0/ ,ARR137 /0/ ,ARR138 /0/ ,ARR139 /0/ ,ARR140 /0/ ,ARR141 /0/ ,ARR142 /0/ ,ARR143 /0/ ,ARR144 /0/ ,ARR145 /0/ ,ARR146 /0/ ,ARR147 /0/ ,ARR148 /0/ ,ARR149 /0/ ,ARR150 /0/ ,ARR151 /0/ ,ARR152 /0/ ,ARR153 /0/ ,ARR154 /0/ ,ARR155 /0/ ,ARR156 /0/ ,ARR157 /0/ ,ARR158 /0/ ,ARR159 /0/ ,ARR160 /0/ ,ARR161 /0/ ,ARR162 /0/ ,ARR163 /0/ ,ARR164 /0/ ,ARR165 /0/ ,ARR166 /0/ ,ARR167 /0/ ,ARR168 /0/ ,ARR169 /0/ ,ARR170 /0/ ,ARR171 /0/ ,ARR172 /0/ ,ARR173 /0/ ,ARR174 /0/ ,ARR175 /0/ ,ARR176 /0/ ,ARR177 /0/ ,ARR178 /0/ ,ARR179 /0/ ,ARR180 /0/ ,ARR181 /0/ ,ARR182 /0/ ,ARR183 /0/ ,ARR184 /0/ ,ARR185 /0/ ,ARR186 /0/ ,ARR187 /0/ ,ARR188 /0/ ,ARR189 /0/ ,ARR190 /0/ ,ARR191 /0/ ,ARR192 /0/ ,ARR193 /0/ ,ARR194 /0/ ,ARR195 /0/ ,ARR196 /0/ ,ARR197 /0/ ,ARR198 /0/ ,ARR199 /0/ ,ARR200 /0/ ,ARR201 /0/ ,ARR202 /0/ ,ARR203 /0/ ,ARR204 /0/ ,ARR205 /0/ ,ARR206 /0/ ,ARR207 /0/ ,ARR208 /0/ ,ARR209 /0/ ,ARR210 /0/ ,ARR211 /0/ ,ARR212 /0/ ,ARR213 /0/ ,ARR214 /0/ ,ARR215 /0/ ,ARR216 /0/ ,ARR217 /0/ ,ARR218 /0/ ,ARR219 /0/ ,ARR220 /0/ ,ARR221 /0/ ,ARR222 /0/ ,ARR223 /0/ ,ARR224 /0/ ,ARR225 /0/ ,ARR226 /0/ ,ARR227 /0/ ,ARR228 /0/ ,ARR229 /0/ ,ARR230 /0/ ,ARR231 /0/ ,ARR232 /0/ ,ARR233 /0/ ,ARR234 /0/ ,ARR235 /0/ ,ARR236 /0/ ,ARR237 /0/ ,ARR238 /0/ ,ARR239 /0/ ,ARR240 /0/ ,ARR241 /0/ ,ARR242 /0/ ,ARR243 /0/ ,ARR244 /0/ ,ARR245 /0/ ,ARR246 /0/ ,ARR247 /0/ ,ARR248 /0/ ,ARR249 /0/ ,ARR250 /0/ ,ARR251 /0/ ,ARR252 /0/ ,ARR253 /0/ ,ARR254 /0/ ,ARR255 /0/ ,ARR256 /0/ ,ARR257 /0/ ,ARR258 /0/ ,ARR259 /0/ ,ARR260 /0/ ,ARR261 /0/ ,ARR262 /0/ ,ARR263 /0/ ,ARR264 /0/ ,ARR265 /0/ ,ARR266 /0/ ,ARR267 /0/ ,ARR268 /0/ ,ARR269 /0/ ,ARR270 /0/ ,ARR271 /0/ ,ARR272 /0/ ,ARR273 /0/ ,ARR274 /0/ ,ARR275 /0/ ,ARR276 /0/ ,ARR277 /0/ ,ARR278 /0/ ,ARR279 /0/ ,ARR280 /0/ ,ARR281 /0/ ,ARR282 /0/ ,ARR283 /0/ ,ARR284 /0/ ,ARR285 /0/ ,ARR286 /0/ ,ARR287 /0/ ,ARR288 /0/ ,ARR289 /0/ ,ARR290 /0/ ,ARR291 /0/ ,ARR292 /0/ ,ARR293 /0/ ,ARR294 /0/ ,ARR295 /0/ ,ARR296 /0/ ,ARR297 /0/ ,ARR298 /0/ ,ARR299 /0/ ,ARR300 /0/ ,ARR301 /0/ ,ARR302 /0/ ,ARR303 /0/ ,ARR304 /0/ ,ARR305 /0/ ,ARR306 /0/ ,ARR307 /0/ ,ARR308 /0/ ,ARR309 /0/ ,ARR310 /0/ ,ARR311 /0/ ,ARR312 /0/ ,ARR313 /0/ ,ARR314 /0/ ,ARR315 /0/ ,ARR316 /0/ ,ARR317 /0/ ,ARR318 /0/ ,ARR319 /0/ ,ARR320 /0/ ,ARR321 /0/ ,ARR322 /0/ ,ARR323 /0/ ,ARR324 /0/ ,ARR325 /0/ ,ARR326 /0/ ,ARR327 /0/ ,ARR328 /0/ ,FAST0329
FAST0330
FAST0331
FAST0332
FAST0333
FAST0334
FAST0335
FAST0336
FAST0337
FAST0338
FAST0339
FAST0340
FAST0341
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FAST0355
FAST0356
FAST0357
FAST0358
FAST0359
FAST0360
FAST0361
FAST0362
FAST0363
FAST0364
FAST0365
FAST0366
FAST0367
FAST0368
FAST0369
FAST0371
FAST0372
FAST0373
FAST0374
FAST0375
FAST0376

```

```

MOD1 = INELAS
MAXX = NNMAX
30 IF(I1SV .EQ. 0) CALL READA(10,TITLEA1)
IF(IN3 ,GT,0) CALL READA(10,TITLEB1)
IF(IN4 ,GT,0) CALL READA( 5,NSMAX )
IF(IN5 ,GT,0) CALL READA( 4,NEMAX )
IF(IN6 ,GT,0) CALL READA( 5,NXSECT)
IF(IN7 ,GT,0) CALL READA( 4,NORDER)
IF(IN8 ,GT,0) CALL READA( 4,NGPAK )
IF(IN9 ,GT,0) CALL READA( 5,NSRMAX )
IF(IN10,GT,0) CALL READA( 4,IPRINT)
IF(I1SV .EQ.0) GO TO 70
NARRAY = 51
IF(NXSECT,GT,MODX) GO TO 40
NNMOVE = 0
NNMAX = 0
40 NMOVE = LOCN(ARRAY)
NNMOVE = NSTORE - NMOVE
DO 50 I=1,NARRAY
50 LCP(I) = LOC(I) + NMOVE
1 NSTORE
J = NMOVE
NNMOVE = NMOVE - 1
DO 60 K=1,NNMOVE
I = I - 1
J = J + 1
60 NM = I + NM
70 NMODD = NEMAX + 1
NCODD = NNMAX + 1
NOMENT = NYMODD + NNMAX + NLMAX + NTMAX + 1
NOMENT = NOMENT + NSRMAX
NODOD = 1
IF(NPOINT .EQ. 0) NODOD = NMAX
INPC = INSC + 1
IAZ = I1PV + NNMAX
I1SV = IAZ + NNMAX
IMTL = I1SV + NRMAX
IRHO = IRHL + NRMAX
IXR = IRHO + NRMAX
I1ELW = I1EL + NMODD
IAE = I1EL + NEMAX
IBE = IAE + NEMAX
INSG = IBE + NEMAX
INPC = INSC + NVMAX
IJSN = INPC + NVMAX
FAST0329
FAST0330
FAST0331
FAST0332
FAST0333
FAST0334
FAST0335
FAST0336
FAST0337
FAST0338
FAST0339
FAST0340
FAST0341
FAST0342
FAST0343
FAST0344
FAST0345
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FAST0347
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FAST0356
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FAST0360
FAST0361
FAST0362
FAST0363
FAST0364
FAST0365
FAST0366
FAST0367
FAST0368
FAST0369
FAST0371
FAST0372
FAST0373
FAST0374
FAST0375
FAST0376

```

IJSX	=	IJSX + NVMAX	FAST0377
ISUV	=	IJSX + NVMAX*3	FAST0379
IATN	=	ISUV + NVMAX	FAST0380
IATW	=	IATN + NHMAX	FAST0381
IESW	=	IATW + NHMAX	FAST0382
IEWH	=	IESW + NEHMAX	FAST0383
IDEN	=	IEWH + NEHO	FAST0384
INTG	=	IDEN + NNMAX	FAST0385
IELF	=	INTG + NEFMAX	FAST0386
IFGN	=	IELF + NGPDO	FAST0387
ITDS	=	IFGN + NFMAX	FAST0388
ITFS	=	ITDS + NFMAX*3	FAST0389
IIODR	=	ITFS + RVHMAX	FAST0390
IIIDS	=	IIODR + NDMAX	FAST0391
IVOL	=	IIIDS + NDMAX	FAST0392
ICDT	=	IVOL + NDMAX	FAST0393
IXDT	=	ICDT + NDMAX*3	FAST0394
IALP	=	IXDT + NDMAX*3	FAST0395
IRSL	=	IALP + NVMAX	FAST0396
IVHD	=	IRSL + NVMAX*5	FAST0397
IGIM	=	IVHD + RVMAX*5	FAST0398
IAIM	=	IGIM + NEMAX	FAST0399
IALM	=	IAIM + NEMAX	FAST0400
IAHL	=	IALM + NEMAX	FAST0400
IA	=	IAHL + NEMAX	FAST0401
INS	=	IA + NSMAX+NMMAX	FAST0402
IVEE	=	INS + NRMAX+NMMAX	FAST0403
IVAL	=	IVEE + NVMAX*NNMAX*5	FAST0404
ISPW	=	IVAL + NVMAX*NNMAX*5	FAST0405
ISWT	=	ISPW + NVMAX*NNMAX	FAST0406
IATO	=	ISWT + NVMAX+NEHMAX	FAST0407
IRSP	=	IATO + NHMAX+NTHMAX	FAST0408
ISGT	=	IRSP + NFMAX+NGMD	FAST0409
ITOE	=	ISGT + ANHMAX+NEHMD	FAST0410
IID1	=	ITOE + NHMAX+NNDOM+*NDWN*NSECT	FAST0411
ISCG	=	IID1 + NHMAX+NNDOM+*NSECT	FAST0412
IFXT	=	ISCG + NHMAX+NMAX	FAST0413
IFXA	=	IFXT - NDMD+NGMAX+NDMONT	FAST0414
IFXE	=	IFXA - NDMD+NGMAX	FAST0415
IFXT	=	IFXE - NDMD+NGMAX	FAST0416
IFXS	=	IFXT - NDMD+NGMAX	FAST0417
IFXP	=	IFXS - NDMD+NGMAX	FAST0418
INCP	=	IFXP - NSTMAX	FAST0419
INRP	=	INCP - NSTMAX	FAST0420
ISTP	=	INRP - NSTMAX	FAST0421
INSC	=	ISTP - NSTMAX	FAST0422

```

NTP(I) = 10*I12
NEX = 10*K3
IF(NEX.GT.0)GO TO 230

C EXPANDED FORM
      AZ(I1) = AA(1)
      HAX = MIN(10,BMAX,6)
      DO 220 J=1,HAX
220   A(I1,J) = AA(J+1)
      IF(NTP(I).GT.6)CALL READE(NTP(I)-6,A(7,I))
      GO TO 230

230 NGT = (NEX*2)/3
      NEX = NEX - 3*(NGT-1)
      DO 231 J=1,NGT
231   A(I1,J) = AA(J)
      AZ(I1) = 0.0
      GO TO 232,233,235,260,2551,NGT

232 NTP(I) = NEX
      AZ(I1) = -AA(1)
      AT(NEX) = 1.0
      GO TO 230

233 FST = AA(3) - AA(1)
      D = -1.0
      K = 0
      DO 234 J=1,3
      IF(J.EQ.NEX)GO TO 234
      K = K + 1
      NTP(I) = J
      AA(j+1) = D*FST/(AA(K+2) - AA(K))
      AZ(I1) = AZ(I1) - AA(K)*AA(J+1)
234 CONTINUE
      GO TO 230

C CONVERGENCE TO HEX-AXIS
235 CONTINUE
      XX = (AA(16) - AA(4)*AA(5) - AA(3))
      YY = XXX*AA(13) - AA(4)
      NTP(I) = 6
      XX = XXX*#2
      AA(1,I1) = YY*#2 - XX*(AA(11)*#2 + AA(2)*#2)
      K = 0
      DO 240 J=1,3
      IF(J.EQ.NEX)GO TO 240
      AA(J+1) = 2.0*YY
      AA(J+3,I1) = 1.0
      GO TO 240
240 K = K + 1
      AA(1,I1) = 2.0*XXX*AA(13)
      GO TO 230

```

```

AI(J+3,I) = -X
250 CONTINUE
GO TO 290
CELLULAR CYLINDER AND ELLIPSOID
255 NEX = 0
260 CONTINUE
AZ(I,J) = -1.0
K = -1
DO 280 J=1,3
  IF(J.EQ.NEX) GO TO 280
  X = K + 2
  AI(J+3,I) = 1.0/A(K+1)*P2
  AI(J,1) = -2.0*A(K)*AI(J+3,I)
  AZ(I,J) = AZ(I,J) + AA(K)*P2*A(J+3,I)
280 CONTINUE
290 CONTINUE
C REGION DEFINITIONS
300 IF(INS.LE.0)GO TO 330
  FST = 1.0
  IF(NUNITD.EQ.1)FST = 0.6025/L.C0797
  DO 320 N=1,IN3
    CALL READR1D(N,ADM)
    I = IDM(N)
    ISV(I) = IDN(I)
    MTL(I) = IDN(3)
    RHO(I) = FST*ADM(I)
    DO 305 J=1,3
      305 XJ = RHO(I)*ADM(J+1)
      DO 312 J=1,NBMAX
        312 NS(J,I) = 0
        MOD = MIN(0,NBMAX),9
        DO 315 J=1,MOD
          315 NS(J,I) = 1000*IDM(J+3)
        IF(NBMAX.LT.9) GO TO 320
        IF(NBMAX.GT.9) GO TO 320
        CALL RGADIN(NBMAX-9,NS(1,I))
        DO 318 J=9,NBMAX
          318 NS(J,I) = 1000*NS(J,I)
320 CONTINUE
330 IF(LIN4.LE.0)GO TO 340
  DO 332 J=1,NRMAX
    332 J=1,3
    ADM(J,J) = XR(J,I)
    ADM(J+3) = ADM(J)*P2
    K = J + 1 - 3*(J/3)
331 ADM(J+6) = ADM(J)*XR(K,I)
  END

  DO 333 J=1,NRMAX
    333 NS(J,I) = 0
    K = NS(J,I)/1000
    K = K - 1000*(K/1000)
    MAX = NTPI(K)
    NS(J,I) = 1000*K
    FST = AZ(I,K)
    DO 332 J=1,NRMAX
      332 NS(J,I) = ADM(J)*P2*AL(J,K)
      IF(FST.LT.0.01NS(J,I)) NS(J,I) = NS(J,I) + 1000000
333 CONTINUE
340 CONTINUE
340 IF(INS.LE.0)GO TO 370
  DO 360 J=1,NRMAX
    360 J=1,3
    NS(J,I) = 0
    NGS = LOCATE(NGS1,XR(I,1,1))
    IF(NGS.EQ.1)GO TO 360
    IF(LABEL(1).GE.0)WRITEM(2,20001),NGS
  2000 FORMAT(1X,22H***1,35HGEOMETRY ERROR, THE POINT IN REGION,I5,18H
  1 ALSO IN REGION,I5,22H***1)
    NROR = NRORCR
    IF(NGS.GT.0)GO TO 350
  360 CONTINUE
  370 RETURN
END

$!FTFC SOURCE H94/2,XR7
CSOURCE#FIXED SOURCE INPUT AND NORMALIZATION FOR PROGRAM FASTER#T,JORDAN#FASTD3
SUBROUTINE SOURCE(L1,L2,L3,L4,E1,ELN,AE,BE,NSG,NPC,JSN,JSU,SUV,
  1 XTR,VEE,VAL,SPW,SPV,ISV,E,EN,EBG,ENG1
  2 ,NROR,NECR,CR,ELW1,ELW2,ELW3,ELW4,BELL1,BELL2,NSG1,I1,NPC1,L1,JSN1,I1,JSU1,I1,
  3 SUV11,XTR11,VEE11,L2,L3,L4,L1,VLL1,L2,L3,L4,SPW1,L1,L4,
  4 SPE11,ELW11,I1,JSV11,E11,EN11,EBG11,ENG11
  5 COMMON/TAPEID/1
  6 COMMON/LIMITS/NSTCRE,NROR,NSMAX,NMAX,NRMAX,NRMAX,NSTMAX,
  7 NSU,NSU1,NMAX,NMAX,NMAX,NECRD,NSSEC1,NUNITD,
  8 NSU2,NSU3,NMAX,NMAX,NMAX,NECRD,NSSEC2,NUNITD,
  9 NSU4,NSU5,NMAX,NMAX,NMAX,NECRD,NSSEC3,NUNITD,
  10 NSU6,NSU7,NMAX,NMAX,NMAX,NECRD,NSSEC4,NUNITD,
  11 NSU8,NSU9,NMAX,NMAX,NMAX,NECRD,NSSEC5,NUNITD,
  12 NSU10,NSU11,NMAX,NMAX,NMAX,NECRD,NSSEC6,NUNITD,
  13 NSU12,NSU13,NMAX,NMAX,NMAX,NECRD,NSSEC7,NUNITD,
  14 NSU14,NSU15,NMAX,NMAX,NMAX,NECRD,NSSEC8,NUNITD,
  15 NSU16,NSU17,NMAX,NMAX,NMAX,NECRD,NSSEC9,NUNITD,
  16 NSU18,NSU19,NMAX,NMAX,NMAX,NECRD,NSSEC10,NUNITD,
  17 NSU20,NSU21,NMAX,NMAX,NMAX,NECRD,NSSEC11,NUNITD,
  18 NSU22,NSU23,NMAX,NMAX,NMAX,NECRD,NSSEC12,NUNITD,
  19 NSU24,NSU25,NMAX,NMAX,NMAX,NECRD,NSSEC13,NUNITD,
  20 NSU26,NSU27,NMAX,NMAX,NMAX,NECRD,NSSEC14,NUNITD,
  21 NSU28,NSU29,NMAX,NMAX,NMAX,NECRD,NSSEC15,NUNITD,
  22 NSU30,NSU31,NMAX,NMAX,NMAX,NECRD,NSSEC16,NUNITD,
  23 NSU32,NSU33,NMAX,NMAX,NMAX,NECRD,NSSEC17,NUNITD,
  24 NSU34,NSU35,NMAX,NMAX,NMAX,NECRD,NSSEC18,NUNITD,
  25 NSU36,NSU37,NMAX,NMAX,NMAX,NECRD,NSSEC19,NUNITD,
  26 NSU38,NSU39,NMAX,NMAX,NMAX,NECRD,NSSEC20,NUNITD,
  27 NSU40,NSU41,NMAX,NMAX,NMAX,NECRD,NSSEC21,NUNITD,
  28 NSU42,NSU43,NMAX,NMAX,NMAX,NECRD,NSSEC22,NUNITD,
  29 NSU44,NSU45,NMAX,NMAX,NMAX,NECRD,NSSEC23,NUNITD,
  30 NSU46,NSU47,NMAX,NMAX,NMAX,NECRD,NSSEC24,NUNITD,
  31 NSU48,NSU49,NMAX,NMAX,NMAX,NECRD,NSSEC25,NUNITD,
  32 NSU50,NSU51,NMAX,NMAX,NMAX,NECRD,NSSEC26,NUNITD,
  33 NSU52,NSU53,NMAX,NMAX,NMAX,NECRD,NSSEC27,NUNITD,
  34 NSU54,NSU55,NMAX,NMAX,NMAX,NECRD,NSSEC28,NUNITD,
  35 NSU56,NSU57,NMAX,NMAX,NMAX,NECRD,NSSEC29,NUNITD,
  36 NSU58,NSU59,NMAX,NMAX,NMAX,NECRD,NSSEC30,NUNITD,
  37 NSU60,NSU61,NMAX,NMAX,NMAX,NECRD,NSSEC31,NUNITD,
  38 NSU62,NSU63,NMAX,NMAX,NMAX,NECRD,NSSEC32,NUNITD,
  39 NSU64,NSU65,NMAX,NMAX,NMAX,NECRD,NSSEC33,NUNITD,
  40 NSU66,NSU67,NMAX,NMAX,NMAX,NECRD,NSSEC34,NUNITD,
  41 NSU68,NSU69,NMAX,NMAX,NMAX,NECRD,NSSEC35,NUNITD,
  42 NSU70,NSU71,NMAX,NMAX,NMAX,NECRD,NSSEC36,NUNITD,
  43 NSU72,NSU73,NMAX,NMAX,NMAX,NECRD,NSSEC37,NUNITD,
  44 NSU74,NSU75,NMAX,NMAX,NMAX,NECRD,NSSEC38,NUNITD,
  45 NSU76,NSU77,NMAX,NMAX,NMAX,NECRD,NSSEC39,NUNITD,
  46 NSU78,NSU79,NMAX,NMAX,NMAX,NECRD,NSSEC40,NUNITD,
  47 NSU80,NSU81,NMAX,NMAX,NMAX,NECRD,NSSEC41,NUNITD,
  48 NSU82,NSU83,NMAX,NMAX,NMAX,NECRD,NSSEC42,NUNITD,
  49 NSU84,NSU85,NMAX,NMAX,NMAX,NECRD,NSSEC43,NUNITD,
  50 NSU86,NSU87,NMAX,NMAX,NMAX,NECRD,NSSEC44,NUNITD,
  51 NSU88,NSU89,NMAX,NMAX,NMAX,NECRD,NSSEC45,NUNITD,
  52 NSU90,NSU91,NMAX,NMAX,NMAX,NECRD,NSSEC46,NUNITD,
  53 NSU92,NSU93,NMAX,NMAX,NMAX,NECRD,NSSEC47,NUNITD,
  54 NSU94,NSU95,NMAX,NMAX,NMAX,NECRD,NSSEC48,NUNITD,
  55 NSU96,NSU97,NMAX,NMAX,NMAX,NECRD,NSSEC49,NUNITD,
  56 NSU98,NSU99,NMAX,NMAX,NMAX,NECRD,NSSEC50,NUNITD,
  57 NSU100,NSU101,NMAX,NMAX,NMAX,NECRD,NSSEC51,NUNITD,
  58 NSU102,NSU103,NMAX,NMAX,NMAX,NECRD,NSSEC52,NUNITD,
  59 NSU104,NSU105,NMAX,NMAX,NMAX,NECRD,NSSEC53,NUNITD,
  60 NSU106,NSU107,NMAX,NMAX,NMAX,NECRD,NSSEC54,NUNITD,
  61 NSU108,NSU109,NMAX,NMAX,NMAX,NECRD,NSSEC55,NUNITD,
  62 NSU110,NSU111,NMAX,NMAX,NMAX,NECRD,NSSEC56,NUNITD,
  63 NSU112,NSU113,NMAX,NMAX,NMAX,NECRD,NSSEC57,NUNITD,
  64 NSU114,NSU115,NMAX,NMAX,NMAX,NECRD,NSSEC58,NUNITD,
  65 NSU116,NSU117,NMAX,NMAX,NMAX,NECRD,NSSEC59,NUNITD,
  66 NSU118,NSU119,NMAX,NMAX,NMAX,NECRD,NSSEC60,NUNITD,
  67 NSU120,NSU121,NMAX,NMAX,NMAX,NECRD,NSSEC61,NUNITD,
  68 NSU122,NSU123,NMAX,NMAX,NMAX,NECRD,NSSEC62,NUNITD,
  69 NSU124,NSU125,NMAX,NMAX,NMAX,NECRD,NSSEC63,NUNITD,
  70 NSU126,NSU127,NMAX,NMAX,NMAX,NECRD,NSSEC64,NUNITD,
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  72 NSU130,NSU131,NMAX,NMAX,NMAX,NECRD,NSSEC66,NUNITD,
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  316 NSU618,NSU619,NMAX,NMAX,NMAX,NECRD,NSSEC310,NUNITD,
  317 NSU620,NSU621,NMAX,NMAX,NMAX,NECRD,NSSEC311,NUNITD,
  318 NSU622,NSU623,NMAX,NMAX,NMAX,NECRD,NSSEC312,NUNITD,
  319 NSU624,NSU625,NMAX,NMAX,NMAX,NECRD,NSSEC313,NUNITD,
  320 NSU626,NSU627,NMAX,NMAX,NMAX,NECRD,NSSEC314,NUNITD,
  321 NSU628,NSU629,NMAX,NMAX,NMAX,NECRD,NSSEC315,NUNITD,
  322 NSU630,NSU631,NMAX,NMAX,NMAX,NECRD,NSSEC316,NUNITD,
  323 NSU632,NSU633,NMAX,NMAX,NMAX,NECRD,NSSEC317,NUNITD,
  324 NSU634,NSU635,NMAX,NMAX,NMAX,NECRD,NSSEC318,NUNITD,
 
```

```

10 BE(I) = 0.0
C TOTAL SOURCE, ERROR CRITERION, TRANSLATION VECTOR
330 IF(IN3.EQ.01GO TO 505
      DO 335 IDN(I,I)=IN3
      CALL READ(EOM,ADM)
      VOLUME = 1.0
      I = IDM(1)
      NSG(I) = IDM(2)
      N=NSG(I)
      DO 405 J=1,5
      MAX = IDM(J+2)
      IF(MAX.GT.01 GO TO 299
      K = -MAX
      MAX = NPC(J,K)
      DO 200 L=1,MAX
      VEE(L,J,I) = VEE(L,J,K)
      200 VAL(L,J,I) = VAL(L,J,K)
      NPC(J,I,I) = MAX
      IF(MAX = 1) 405/405,21C
      299 NPC(J,I,I) = MAX
      CALL READ(E(MAX,VEE(1,J,I),VAL(1,J,I))
      IF(MAX.LE.-1GO TO 405
      FST = 0.0
      C  DO 987 K=1,MAX
      C  IF(VEE(K,J,I).EQ.0.0) VEE(K,J,I) = 1.0E-30
      C  987 CONTINUE
      DO 395 K=2,MAX
      BB = 0.0
      AB = 0.0
      IF(VEE(K,J,I).EQ.VEE(K-1,J,I)) GO TO 395
      BB = (VAL(K,J,I)-VAL(K-1,J,I))/(VEE(K,J,I)-VEE(K-1,J,I))
      AB = VAL(K,J,I)-BB*VEE(K,J,I)
      395 FST = FST+AB*(VEE(K,J,I)**(N+1))-VEE(K-1,J,I)**(N+1)/FLUAT(N+1)
      1          +BB*(VEE(K,J,I)**(N+2))-VEE(K-1,J,I)**(N+2)/FLUAT(N+2)
      DO 400 K=1,MAX
      400 VAL(K,J,I) = VAL(K,J,I)/FST
      210 CONTINUE
      IF(.LE.3)VOLUME = VOLUME*(VEE(MAX,J,I)**(N+1)-VEE(1,J,I)**(N+1))
      405 N = 0
      ISP = IDM(10)
      MAX = IDM(1)
      IF(MAX.GT.01 GO TO 299
      K = -MAX
      JSN(I) = JSN(K)
      JSX(I) = JSX(K)
      FST = 0.0
      EST = 0.0

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```

DO 220 J=1,NMAX
      SPM(J,I) = SPM(J,K)
      SPE(J,I) = SPE(J,K)
      FST = FST + SPM(J,I)
      220 EST = EST + SPM(J,I)*SPE(J,I)
      GO TO 299
299 CONTINUE
IF(ISP.GT.11 GO TO 100
C INPUT SPECTRUM IS DIFFERENTIAL IN PARTICLES OR INTENSITY
IF(IDM(11).GT.01 GO TO 600
      C 353 J=1,MAX
      EN = EMAX*EN
      GO TO 601
600 IF(TOM(11).EQ.2) CALL READ(E,EN)
      CALL READ(E(MAX,EN))
601 CONTINUE
IF(ISP.LE.01GO TO 354
C DIFFERENTIAL IN INTENSITY
      DO 354 J=1,MAX
      354 ENIJ = ENIJ/E(IJ)
      GO TO 354
C INPUT SPECTRUM IS GROUPWISE INTEGRATED, MAX = NC. GROUPS
100 IF(IDM(11).EQ.01 GO TO 351
      IF(DOM(11).EQ.2) CALL READ(E(MAX+1,EBG)
      CALL READ(E(MAX,ENG))
      GO TO 110
      351 CALL READ(E(MAX+1,ERG,ENG))
110 CONTINUE
      IF(ISP.EQ.2)GO TO 355
C TOTAL MEV IN GROUP, DIVIDE BY AVERAGE ENERGY
      DO 355 J=1,MAX
      355 ENIK = ENIJ/(EBG(J) + EBG(J+1))
      352 ENGJ = ENGJ/(ENG(J) + EBG(J+1))
C TOTAL PARTICLES IN GROUP, DIVIDE BY GROUP WIDTH
      355 K = 0
      DO 353 J=1,MAX
      FST = ENGJ/(EBG(J) + EBG(J+1))
      353 L=1,2
      K = L+1
      M = J + L - 1
      E(K) = EBG(M)
      353 ENIK = FST
      MAX = K
      354 FST = 0.0
      EST = 0.0
      JSN(I) = NMAX
      JSX(I) = 1
      DO 415 J=1,NMAX
      SPW(J,I) = 0.0

```

```

      SPC(J,I) = 0.0
      DO 410 K=2,MAX
      EMAX = AMINIEELL(JJ),E(K-1))
      EK = AMINIEELL(JJ),E(K))
      EKMIN= EK-EKMIN
      EKMAX= EK-EKMAX
      BB = (ENXK-1) - EN(XK)) / (C(K-1) - E(K))
      AA = EN(XK) - BB*EK
      DEL = (EMAX**2 - EMIN**2)/2.0
      SPW(J,I) = SPW(J,I) + AA*(EMAX - EMIN) + BB*DEL
      SPE(J,I) = SPE(J,I) + AA*DEL + BB*(EMAX**3 - EMIN**3)/3.0
      FAST0752
      FAST0753
      FAST0754
      FAST0755
      FAST0756
      FAST0757
      FAST0758
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      FAST0760
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      FAST0837
      FAST0838
      FAST0839
      FAST0840
      FAST0841
      FAST0842
      FAST0843
      FAST0844
      FAST0845
      410 CONTINUE
      IF(SPMW(J,I),ECL,0.0)GO TO 415
      JSH(I) = MIN0(JSH(I),J)
      JSX(I) = MAX0(JSX(I),J)
      FST = EST + SPW(J,I)
      EST = EST + SPE(J,I)
      SPE(J,I) = SPE(J,I)/SPW(J,I)
      415 CONTINUE
      297 CONTINUE
      DO 420 J=1,3
      420 XTR(J,I) = ADM(J+1)
      EAIVE = EST/FST
      ADM = ADM(J+1)*EAIVE
      421 FST = ADM(I)*FST
      SVU(I) = ADM(I)*EAIVE
      GO TO 422
      C NORMALIZE TO TOTAL PARTICLES
      422 FST = ADM(I)*EST
      SVU(I) = ADM(I)*EST
      GO TO 423
      C NORMALIZE TO TOTAL NEV
      423 FST = ADM(I)*EST
      SVU(I) = ADM(I)*EST
      GO TO 424
      C SCALE INPUT SPECTRUM
      424 FST = ADM(I)
      SVU(I) = ADM(I)*EST
      425 GO TO 425 J=1,NMAX
      425 SPW(J,I) = FST*SPW(J,I)
      430 CONTINUE
      505 IF(IN4.GT.0)CALL READIN(IN4,ISV)
      DO 510 I=1,NMAX
      510 IF(ISV(I).GT.NVPAX)ISV(I) = 0
      510 CONTINUE
      RETURN
      END
      $INFTC SEC1 M94/2,XR7
      CINSECT*CROSS SECTION INPUT, AND MIXING FOR PROGRAM FASTER*T.H,JORDAN*66*FAST0797
      SUBROUTINE INSECT(L1,L2,L3,L4,L5,L6,ESB,DEN,SSH,SGT,IOE,IOI,
      FAST0798
      FAST0799
      FAST0800
      FAST0801
      FAST0802
      FAST0803
      FAST0804
      FAST0805
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      FAST0840
      FAST0841
      FAST0842
      FAST0843
      FAST0844
      FAST0845
      200
  
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      DO 75 J=1,NMAX
 75  AT0(I,J) = FST*AT0(I,J)
 85  CALL READE(NEMIN,XST1)
  IF(NUUNITX.EQ.0)IGO TO 105
  FST = ATW/0.6025
  DO 95 J=1,NMIN
 95  XST(I,J) = XST1(J)
 105  IF(NKSECF.GT.0)GO TO 105
  COMPTON ENERGY ABSORPTION COEFFICIENT
  DO 106 J=1,NMOD
  XMA = ELL(I,J)/0.511
  XMB = L_0 + 2.0*XMA
 106  EAC(I,J) = XST1(J)-0.249375*ZEE*(ALDG(XMB)/XMA**3+2.0*(1.0-
  XMA**2*XMA**2-XMB)**2/(XMA*XMB)**2+8.0*XMA**2/(3.0*XMB**3))*
  109  1F11.GT.-1) GO TO 125
  DO 115 J=1,NMIN
 115  SSH(I,J) = XST1(J)
  IF(NKSECF.EC.0) GO TO 100
  DO 116 J=1,NMOD
  EAC(I,J,I) = 0.0
  IF(I,J,GT.1) EAC(I,J,I) = SSH(J-1)
  IF(I,J,LT.NMOD) EAC(I,J,I) = EAC(I,J,I) + SSH(I,J)
 116  EAC(I,J,I) = ELL(I,J)*EAC(I,J,I)/4.0
  GO TO 398
 125  IF(NUUNITX.EQ.0)IGO TO 100
  MOD = MIN0(L24,1,NELAS+3)
  CALL READI(MOD,JMX)
  NSIGT = 1
  IF(JMXK1(I).GT.0) NSIGT = 0
  LMAX = NMOD(JMXK1(I))
  NSDMOD = JMXK1(I)
  KMAX = JMXK1(I)
  IF(KMAX.GT.2)CALL READI(KMAX-2,JMXK1(25))
  DO 80 L=1,LMAX
  DO 80 K=1,NSDMOD
  80  CALL READI(KMAX+1-K,XSE(I,K,L))
  C TRANSPORT CORRECTION OF P-ZERO IN GROUP FROM P-ONE SCATTER
  XMA = ATW/1.0*(ATW + 1.0)**2
  IF(NUUNITX.GT.0) XMA = XMA*ATW/C.6025
  LMOD = MIN0(L,MAX)
  DO 82 J=1,NMAX
  FST = 1.0
  XMB = 0.0
  NSDMOD = MIN0(NSDM,NEMAX+1-J)
  DO 81 L=1,LMOD
  DO 81 K=1,NSDMOD
  81  XSET(L,I) = XSET(L,I) + XSE(I,J,K,L)

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  XMUL = FST/(3.0*XMB)
  XMUC = XMUL**2 - 1.0
  XMUC = (XMUC + XMUL*SQR(ATW**2 + XMUC))/ATW
  EAC(I,J,I) = XMUC*XMWB(I) - XC(I)
  IF(LMAX.GT.1)ORDEN(XST1(I,J,I),I) = XSE(I,J,I,I) - FST/3.0
  IF(NUUNITX.GT.0) FST = FST*ATW/0.6025
  82  XST(I,J) = XST1(J) + FST*FLOAT(NSIGT - 1/NORDER)/3.0
  83  CONTINUE
  LMAX = MIN0(NORDER,LMAX)
  XMA = 0.0
  DO 398 J=1,NMOD
  XMB = 0.0
  IF(I,J,LT.NMOD) XMB = EAC(I,J,I)
  EAC(I,J,I) = ELL(I,J)*(XMA + XMB)
  397  XMA = XMB
  IF(KMAX.LE.0)IGO TO 398
  C NELAS,NSDMOD,NSDM
  DO 90 K=1,XMAX
  90  CALL READE(JMXK1(31),XS1(I,K))
  398  DO 117 J=1,NMOD,NMAX
  117  EAC(I,J,I) = 2.0*EAC(I,J,I)
  C COMBINE MATERIAL TOTALS
  100  DO 118 I=1,NMAX
  118  IF(AT0(I,H).EQ.0.0)IGO TO 190
  DO 110 J=1,NMIN
  110  SGT(I,J,M) = SGT(I,J,M) + AT0(M)*XST(J)
  L = NMAX + M
  DO 111 L=1,LMAX
  111  EAC(I,L,I) = EAC(I,J,I) + AT0(M)*EAC(J,I)
  IF(NKSECF.GT.0)IGO TO 120
  DEN(M) = DEN(M) + 0.49875*ATD(M)*ZEE
  GO TO 190
  120  IF(I,GT.1)IGO TO 130
  DEN(M) = ATD(M)
  GO TO 130
  130  GST = ATD(M)
  IF(NUUNITX.NE.0)GST = GST*ATW/0.6025
  DO 165 L=1,LMAX
  DO 165 K=1,NSDM
  MM = 1-EK2(L,M)
  MAX = NEMOD - K
  IF(LMAX.GT.0) GO TO 140
  MM = NN
  NN = NN + MAX
  IDE(K,L,M) = MM
  140  IF((MM + MAX).LE.NNMOD) GO TO 150.
  NERROR = NERROR + 1

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      GO TO 165
150 DO 160 J=1,MAX
160 SGS(J,MN) = SGSI(J,MN) + GST*XSE(J,K,L)
165 CONTINUE
170 IF(KMAX.LE.0)GO TC 190
C INELASTIC NEXT
      DO 180 X=1,KMAX
      MN= IDI(K,MN)
      MAX = JM(X,K+3)
      IF(MAX.GT.0) GO TO 170
      MN = MN
      NN = NN + MIN0(NTRANS,NEMOD-K)
      MN = MN
      GO TO 173
170 MN = MN/1000
173 IDI(X,MN) = 1000*MN + MAX0(MAX,MN-1000*MN)
      MN = MN + MAX1.LE.NNNCD(MN) GO TO 175
      NERROR = NERROR + 1
      GO TO 185
175 DO 180 J=1,MAX
180 SGS(J,MN) = SGSI(J,MN) + GST*XSI(J,K)
185 CONTINUE
190 CONTINUE
200 CONTINUE
390 CONTINUE
      NNMAX = NN - 1
      IF(LABEL(1).GE.0) WRITE(M2,200)NNMAX,IN1
2000 FORMAT(1X,19I1#),34HSCATTERING CROSS SECTIONS REQUIRED,I6,3H OF,
     1 I6,20I4 AVAILABLE LOCATIONS,19(I#))
      IF(NLE.EQ.0) GO TO 210
      DO 210 I=1,NMAX
      DO 210 J=1,NMAX
      MAX = MIN0(J + 7,NMIN)
      MN = MIN0(J + 7,NMIN)
      IF(LABEL(1).GE.0) WRITE(M2,3000)I,(SGT(K,I),K=J,MAX)
3000 FORMAT(1X,6HSIGNAT,I5,1PBE12.4)
500 CONTINUE
510 CONTINUE
      MN = MN + NMAX * NMMAX
      DO 520 I=1,NMAX
      K = I
      IF(K.GT.NMAX) K = K - NMAX
      DO 520 J=1,NEMOD
      MN = MIN0(J + 7,NEMOD)
      IF(LABEL(1).GE.0) WRITE(M2,3010)K,(EAC(L,I),L=J,MAX)
3010 FORMAT(1X,6HE-DUMP,I5,1PBE12.4)
520 CONTINUE
530 CONTINUE
      }

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210 IF(IN4.GT.0)CALL READIN4,NTLG
      IF(IN5.LE.0)GO TO 250
      IF(NUNITD.EQ.0)GO TO 230
      FST = 0.6025/1.00797
      DO 220 I=1,NMAX
      RHO(I) = RHO(I)/FST
220 CALL READA(THS,RHO)
      IF(NUNITD.EC.0)GO TO 250
      DO 240 I=1,NMAX
      RHO(I) = FST*RHO(I)
240 RETURN
      END
      $IBTYPE ACKFOR MN/2,WRT
      CRESULT*RESPONSE FUNCTION AND DETECTOR INPUT FOR FASTER*T.M.JORDAN*WANL
      SUBROUTINE RESULT(L1,L2,NTG,ELF,FGW,TDS,RSP,IOR,IDS,VOL,CDT,XDT,
     1 ELL,MTL,RHO,EAC,L3)
COMMON/LIMITS/NSTORE,NERROR,NMAX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     1 NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     2 NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     3 NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     4 NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     5 NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,NMAXX,
     6 NPRINT,HVMDOM,HDPOINT,HDHELP,HDODEL,HDODELU,HDODELY
COMMON/HILL/NSRC,NSRCA,NSRCA,INSR
COMMON/INVS/NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN
COMMON/INVS/NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN,NSHUN
DIMENSION NTG(1),ELF(1),FW(1),TDS(1),RSP(1),VOL(1),CDT(1),XDT(1),
     1 IOR(1),IDS(1),ELL(1)
DIMENSION ADM(7),ODM(3),IDM(3)
COMMON/TAPEID/M1
COMMON/INVS/IN1,IN2,IN3,IN4,IN5,IN6,IN7,IN8,IN9,IN10,INN(14)
IF(NMAXX,LT,NEMAX)GO TO 20
DO 10 I=1,NEMAX
10 NTG(I) = 1
20 IF(IN2.GT.0) CALL READI(NEMAX,NTG)
30 ELF(1) = ELL(1)
      DO 40 I=1,NEMAX
      J = NGC(I)
      40 ELF(J+1) = ELL(I+1)
      DO 50 I=1,NCMAX
      50 NGM(I) = ELF(I)-ELF(I+1)
60 IF(IN3.LE.0)GO TO 120
      DO 110 I=1,IN3
      CALL READOF(IDH,BDM,ADM)
      I = TDM(I)
      DO 70 J=1,3

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70 TOS(I,J,I) = BDM(I,J)
IF(IOM(2),GE,0) GC TO 75
J = -IDM(2)
N = INTL(J) + NIMAX
RSPI(I,I) = RHO(I,J)*EAC(I,I)
IF((N,GT,NIMAX) RSPI(I,I) = RSP(I,I) + EAC(I,N)
DO 10 J=1,NMAX
L = N+I-K
RSPI(L+1,I) = RHO(I,J)*EAC(K+1,I)
IF((N,GT,NIMAX) RSPI(L+1,I) = RSP(L+1,I) + EAC(K+1,N)
11 CONTINUE
• (ILABEL(I+NGMOD/7).GE,0) WRITE(M2,2000)(RSPI(J,I),J=1,NGMOD)
2000 FORMAT(IX,1HHE ENERGY DUMP,1PE12.4)
GO TO 95
75 CONTINUE
MAX = MIN(5,NGMOD)
DO J=J1,MAX
80 RSPI(J,I) = RSP(J,I)+EAC(J,I)
IF((N,GT,5)ICAL READ(EHOD-5,RSPI(6,I))
IF(IDM(2)-FC,0)GO TO 100
DO 90 J=1,NGMOD
90 RSPI(J,I) = ELF(J)*RSPI(J,I)
95 CONTINUE
100 DO I=1,NGMOD
110 RSPI(I,I) = ADM(I)*RSPI(I,I)
120 IF((N4,LE,0)GO TO 552
DO I30 I=1,3
130 BDM(I) = 0.0
DO I50 N=1,IN4
CALL RDM(N,ADM,I,ADM)
I = IDM(I)
IDR(I) = IDM(2)
IDS(I) = IDM(3)
VOL(I) = ADM(I)
DO I40 I=1,3
CDT(I,J,I) = ADM(J+I)
140 XDT(I,J,I) = ADM(J+I)
150 FST = VECTOR(BDM,CDT(1,I),CDT(1,I))
552 NPDMAX = 0
NVDMAX = 0
DO I60 I=1,NMAX
I61(IDR(I),GT,0)GO TO 554
NPDMAX = NPDMAX + 1
GO TO 556
554 IF(IDS(I),LE,0)NVDMAX = NVDMAX + 1
556 CONTINUE
NVDMOD = NDMAX - NPDMAX
      
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NSDMAX = NVDMOD - NVCMAX
NPDMOD = NPCMAX
1F((N5,GT,0) CALL READIN(NVMOD,H1)           FASTI081
1F((N6,GT,0) CALL READIN(SRMAX,H1)             FASTI082
1F((N7,GT,0) CALL READIN(RMAX,H1)               FASTI083
1F((N8,GT,0) CALL READIN(RMAX,H1)               FASTI084
1F((N9,GT,0) RETURN                           FASTI085
1END                                         FASTI086
1$IBFTC RANDOM M94/2,XRT                   FASTI087
RANDOM$* SAMPLING PARAMETER INPUTS FOR PROGRAM FASTER*T,N,JORDAN*WANL*66*FASTI088
SUBROUTINE RANDOM(RHS1,ALP,VMD,GIM,AIM,ALM,ALH,NPC,VEE,L1,L2)          FASTI089
DIMENSION NPC(1),VEE(1),L2,L1),RAT(5)                                     FASTI090
DIREC(1),AIM(1),ALP(1),VMD(1),GIM(1),AIM(1),ALM(1),ALH(1)          FASTI091
COMMON/TAPERD/M1,M2
COMMON/OTHERS/RADIUS,XCT(3),DELTA ,BDC(3),ATA ,ATB ,ATC ,                FASTI092
1     ATD ,AT ,BT ,AS ,BS ,UMIN ,UMAX ,                           FASTI093
2     KHM ,KMX ,JBAR ,NZERO ,                           FASTI094
COMMON/LIMITS/NSTORE,NODR,NMAX ,NMMAX ,NMRAZ ,NPMAX ,NSTMAX ,          FASTI095
1     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI096
1     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI097
1     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI098
3     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI099
4     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI100
5     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI101
6     NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,NMRAZ ,          FASTI102
COMMON/INPUTS/INV,INV1,INV2,INV3,INV4,INV5,INV6,INV7,INV8,INV9,INV10,INV11,        FASTI103
1F((N2,GT,0)CALL READ(E0,RADIUS)           FASTI104
1F((N3,GT,0)CALL READ(INVMAX,RS1)           FASTI105
1F((MODEL,GT,0)GO TO 30                     FASTI106
1F((NVMAX ,EQ,1) GO TO 30
1F(PTM,GT,0)                                FASTI107
DO LO I=1,NVMAX                            FASTI108
10 FST = FST + RS1(I)
DO 20 I=1,NVMAX                            FASTI109
20 RS1(I) = RS1(I)/FST                      FASTI110
30 IF((N4,LE,0)GO TO 793
DO 40 I=1,NMAX
CALL READ(E5,VMD(I,I))
CALL READ(E5,VMD(I,I))
CALL READ(E5,VMD(I,I))
DO 40 J=1,5
ALP(I,J,I) = 0.0
K = -NPCI(J,I)
1F(K,GT,0) GO TO 40
ALP(I,J,I) = -ALOG(RAT(I,J))
1     /AMAX1(VMD(I,J,I)-VEE(I,J,I),VEE(K,J,I)-VMD(J,I))           FASTI112
40 CONTINUE
560 CONTINUE
793 IF((N5,GT,0)CALL READ(INEMAX,GIM)           FASTI113
1F((N6,GT,0)CALL READ(INEMAX,GIM)           FASTI114
1F((N7,GT,0)CALL READ(INEMAX,AIM)           FASTI115
1F((N8,GT,0)GO TO 60                         FASTI116
      
```

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      CALL READE(NMAX,ALK)
      DO 50 I=1,NMAX
      50 IF(I.NE.1)CALL ALG(ALM(I))
      60 IF((N9.LE.-0.02) TO 80)
      CALL READE(NMAX,ALH)
      DO 70 I=1,NMAX
      70 ALH(I) = 0.5*ALG(ALH(I))
      80 IF((N9.LE.-0.02) TO 90)
      CALL READE(NMAX,ATA)
      ATA = -ALG(ATA)
      ATB = -ALOG(ATB)/3.141596
      90 RETURN
      END

$1BFTC IREAD  M94/2,XRT
CREADE *FLOATING INPUT, 18A4,2413 FORMAT
      SUBROUTINE READE(MAX,ADM)
      COMMON//TAPEID/M1 ,M2
      DIMENSION IDM1(10000),H(2)
      DO 30 I=1,MAX,24
      MOD = I + 23
      IF(MOD.GT.MAX)GO TO 10
      READ(M1,1000)(IDM(J),J=I,MOD),H
      GO TO 20
      10 JMX = MOD - MAX
      MOD = MAX
      READ(M1,1000)(ADM(J),J=I,MOD),(K,J=1,JMX),H
      20 IF(LABEL(1).GE.0)WRITE(M2,2000H,(IDM(J),J=I,MOD))
      30 CONTINUE
      1000 FORMAT(124I3,24A4)
      2000 FORMAT(IX,2A4,3H...,IPBE12.4)
      RETURN
      END

$1BFTC IREAD  M94/2,XRT
CREADE *FLOATING INPUT, 8E9,0 FORMAT
      SUBROUTINE READE(MAX,ADM)
      COMMON//TAPEID/M1 ,M2
      DIMENSION ADM1(10000),H(2)
      DO 30 I=1,MAX,8
      MOD = I + 7
      IF(MOD.GT.MAX)GO TO 10
      READ(M1,1000)(ADM(J),J=I,MOD),H
      GO TO 20
      10 JMX = MOD - MAX
      MOD = MAX
      READ(M1,1000)(ADM(J),J=I,MOD),(X,J=1,JMX),H
      20 IF(LABEL(1).GE.0)WRITE(M2,2000H,(ADM(J),J=I,MOD))
      30 CONTINUE
      END

      1000 FORMAT(8E9.0,2A4)
      2000 FORMAT(IX,2A4,3H...,IPBE12.4)
      RETURN
      END

$1BFTC AREAD  M94/2,XRT
CREADE *HOLLERITH INPUT, 18A4 FORMAT
      SUBROUTINE READA(MAX,ADM)
      COMMON//TAPEID/M1 ,M2
      DIMENSION ADM1(10000),H(2)
      DO 30 I=1,MAX,10
      MOD = I + 10
      IF(MOD.GT.MAX)GO TO 10
      READ(M1,1000)(ADM(J),J=I,MOD),H
      GO TO 20
      10 JMX = MOD - MAX
      MOD = MAX
      READ(M1,1000)(ADM(J),J=I,MOD),(X,J=1,JMX),H
      20 IF(LABEL(1).GE.0)WRITE(M2,2000H,(ADM(J),J=I,MOD))
      30 CONTINUE
      1000 FORMAT(20A4)
      2000 FORMAT(IX,2A4,27I(H.),18A4)
      RETURN
      END

$1BFTC ISREAD M94/2,XRT
      SUBROUTINE READIS(MAX,ADM)
      DIMENSION IDM1(10000),JDM(2,12)
      DO 10 I=1,MAX,12
      MOD = I - (MAX+1-I)
      CALL READ(2*MOD,JDM)
      DO 10 J=1,MOD
      K = JDM(1,J)
      10 IDM(K) = JDM(2,J)
      RETURN
      END

$1BFTC AREAD  M94/2,XRT
      SUBROUTINE READAE(MAX,ADM,ADM)
      COMMON//TAPEID/M1 ,M2
      DIMENSION IDM1(10000),ADM1(1000),H(2)
      DO 30 I=1,MAX,6
      MOD = I + 5
      IF(MOD.GT.MAX)GO TO 10
      READ(M1,1000)(ADM(J),ADM(J),J=I,MOD),H
      GO TO 20
      10 JMX = MOD - MAX
      MOD = MAX
      READ(M1,1000)(ADM(J),ADM(J),J=I,MOD),(K,X,J=1,JMX),H
      20 IF(LABEL(1).GE.0)WRITE(M2,2000H,(IDM(J),ADM(J),J=I,MOD))
      END

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      30 CONTINUE
1000 FORMAT(6(I3,E9.0),2A4)
2000 FORMAT(1X,2A4,3H...,6(I4,1PE12.4))
      RETURN
    END
$1BFTC SREAD M94/2,XRT
      SUBROUTINE READE(MAX,ADM)
      DIMENSION ADM(100001),IDM(61),BDM(61)
      DO 10 I=1,MAX-K
      MOD = MIN0(I,MAX+I-1)
      CALL READE1(MOD,ADM,EDM)
      DO 10 J=1,MOD
      K = K+1
      ADM(K) = EDM(J)
      10 ADM(K) = BDM(J)
      RETURN
    END
$1BFTC EEREAD M94/2,XRT
      CREAD(EALTERNATING FLOATING INPUT, 8E9.0 FORMAT
      SUBROUTINE READEE(MAX,ADM,ADM1)
      DIMENSION ADM(100001),BDM(100001),CDM(2,4)
      K = 0
      DO 10 I=1,MAX,4
      MOD = MIN0(4,(MAX-K))
      CALL READE2(*#MOD,CDM)
      DO 10 J=1,MOD
      K = K+1
      ADM(K) = CDM(1,J)
      10 BDM(K) = CDM(2,J)
      RETURN.
    END
$1BFTC SREAD M94/2,XRT
      CREAD * SURFACE FORMAT, 3I3,7E9.0
      SUBROUTINE READS(IDM,ADM)
      COMMON/TAPE10/M1      ,M2
      DIMENSION IDM(3),ADM(7),H(2)
      READ(M1,1000)ADM,H
      IF(ILABEL(1).GE.0)WRITE(2,2000)H,ADM,ADM
1000 FORMAT(3I3,7E9.0,2A4)
2000 FORMAT(1X,2A4,3H...,3I4,1PE12.4)
      RETURN
    END
$1BFTC SREAD M94/2,XRT
      CREAD * VOLUME SOURCE FORMAT, 4E9.0,12I3 FORMAT
      SUBROUTINE READR(IDM,ADM)
      COMMON/TAPE10/M1      ,M2
      DIMENSION IDM(12),ADM(12),H(2)
      READ(M1,1000)ADM,H
      IF(ILABEL(1).GE.0)WRITE(2,2000)H,ADM,ADM
1000 FORMAT(12I3,4E9.0,2A4)
2000 FORMAT(1X,2A4,3H...,1I4,2X,12I3,1PE12.4)
      RETURN
    END
$1BFTC FREAD M94/2,XRT
      CREAD * RESPONSE FUNCTION FORMAT, 2I3,3A4,6E9.0
      SUBROUTINE READF(IDM,ADM,BDM)
      COMMON/TAPE10/M1      ,M2
      DIMENSION IDM(2),ADM(3),BDM(61),H(2)
      READ(M1,1000)ADM,BDM,H
      IF(ILABEL(1).GE.0)WRITE(2,2000)H,ADM,BDM
1000 FORMAT(2I3,3A4,6E9.0,2A4)
2000 FORMAT(1X,2A4,3H...,1I4,2X,3A4,2X,1PE12.4)
      RETURN
    END
$1BFTC RSVOLV M94/2,XRT
      SUBROUTINE SOLVIT
      COMMON/LIMITS/NSTCRE,NERROR,NSMAX ,NAMAX ,NRMAX ,NBMAX ,NSTMAX ,
1      NEMAX ,NVMAX ,NMAX ,NEMAX ,NEMOD ,NSECY ,NSECX ,NSECZ ,NTRANS ,
2      NMAX ,NMAX ,
3      NMAX ,NMAX ,
4      NMAX ,NMAX ,
5      NMAX ,NMAX ,
6      NMAX ,NMAX ,
      NPRINT,NUINTS,NUMBER,KALIDE
      COMMON/INDEXS/INTP   ,IAZ   ,IISV   ,IMTL   ,IRHO   ,IXR   ,IELL   ,
1      IELW   ,IAE   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
2      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
3      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
4      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
5      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
6      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
7      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
8      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
9      ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,ITOT   ,
1      IFXT   ,IFXA   ,IFXA
      COMMON/H(1)
      IF(NPOINT,G1,0)GO TO 20
      CALL SOLVER1(H(1),H(2),H(3),H(4),H(5),H(6),H(7),H(8),H(9),H(10),
1      H(11),H(12),H(13),H(14),H(15),H(16),H(17),H(18),H(19),H(20),
2      H(21),H(22),H(23),H(24),H(25),H(26),H(27),H(28),H(29),H(30),
3      H(31),H(32),H(33),H(34),H(35),H(36),H(37),H(38),H(39),H(40))
      GO TO 10
20 CALL SOLVER1(H(1),H(2),H(3),H(4),H(5),H(6),H(7),H(8),H(9),H(10),
1      H(11),H(12),H(13),H(14),H(15),H(16),H(17),H(18),H(19),H(20),
2      H(21),H(22),H(23),H(24),H(25),H(26),H(27),H(28),H(29),H(30),
      H(31),H(32),H(33),H(34),H(35),H(36),H(37),H(38),H(39),H(40))

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3H((ISTP),H(INPP),H(INCP),H(1DTR))
FAST1315
10 RETURN
FAST1316
END
FAST1317
$IBFIC SOBER2 M94/2,XR7
FAST1318
CDSDPR2 FOR THE SOLUTION OF THE BOLTZMANN EQUATION BY RANDOM SAMPLING
FAST1319
SUBROUTINE CDSDPR2(NST,NTOTAL,L,NM,NSC,MCCE,*,ST,NGK,*_ML,1DTR,IDS,VOL,
FAST1320
1D1,COT,FXP,FXS,XMP,NSC,AM,ALF,RHO,ELL,SSH,SGT,L3)
DIMENSION SSH(L1),SGT(L1)
FAST1321
COMMON/LIMITS/NSTCRE,NERROR,NSMAX,NMAXX,NMAXR,NMAXB,NSTMAX,
FAST1322
1 NEMAX,NVMAX,NXMAX,NXMAX,NEMOD,NXSEC,NUNITD,
FAST1323
2 NMUNITX,NRMAX,NRMAX,NRORDER,NNDOM,INLEAS,NTTRANS,
FAST1324
3 NMHD,NCODE,NHDF,NCODE,NHDF,NCODE,NHDF,NCODE,NHDF,
FAST1325
4 NOKAX,NCODE,NCODE,NNDOM,NHDF,NHDF,NHDF,NHDF,NHDF,
FAST1326
5 NVMAXX,NVNDOD,NPOINT,MODELP,MODELO,MODELU,MODELV,
FAST1327
6 NPRINT,NTOTAL,NUMBER,KALIDE
FAST1328
COMMON/FLUXES/KKK
,NTALLY,ERROR ,TCTALN,TOTALN,RHON ,SNDR ,TOTFLX1300
COMMON/POINTX/NTOTAL,ITI ,*M0N
FAST1329
COMMON/OTHERS/RADIUS,THICK,DELTA ,BDC(3)*ATA ,ATB ,ATC ,
FAST1330
1 KTH ,ATD ,AT ,BT ,AT ,AT ,BT ,AT ,BT ,AT ,BT ,AT ,BT ,AT ,
FAST1331
2 KRNH ,KMAX ,JBAR ,NZERO
FAST1332
DIMENSION ELL1()
FAST1335
DIMENSION NSR(NST),NSR1(NST),EC(1),XMP(1),ST(1),NRG(1),NSCL(1),
FAST1336
1 HTL(1),JDR(1),IDS(1),VOL(1),XTD(1),COT(1),1
FAST1337
2 EXP(L1,1),EXP(L1,2),FXS(L1,1),GIM(1),AIM(1),ALM(1),ALH(1),RHO(1),L1
FAST1338
3 DIMENSION C1(3),XC1(3),C1(3),CC1(3),CN(3),CP(3)
FAST1340
SNDR = AHA(1,100,0,VECTOR(RCT,BDC,C1)
FAST1341
NTOTAL = 0
FAST1342
TOMEGA = 1
FAST1343
IF(INVMDP,GT,0) TOMEGA = NUMBER + 1
FAST1344
CALL GRDP(BDC)
FAST1345
DO 250 NM=1,NPRINT
FAST1346
DO 240 NM=1,NRUNITS
FAST1347
NTALLY = 1
FAST1348
ERROR = ATD
FAST1349
RUNTD = 0.0
FAST1349
DO 210 K=1,NKALICE
FAST1350
1 IF(KALICE,GT,1)GO TO 10
FAST1350
2 IF(HODELP,GT,0)GO TO 5
FAST1350
5 PDT = PSTAR(NX,1)
FAST1350
GO TO 40
FAST1350
5 PDT = SPHERE(BDC,STC,CNST,ST,ARG,NSC)
FAST1350
1 IF(NST.EQ.0) GO TO 220
FAST1350
NN = NST - 1
FAST1350
DO 6 1=1,NN
FAST1350
6 X(1) = BDC(1) + C(1)*STC
FAST1350
GO TO 40
FAST1350
10 IF(INTALL+NTALLY)-LE,01GO TO 220
FAST1350

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      IF(TOTALN.EC.0)IGO TO 13
      EBAR = TOTALN/TOTALN
      DD 12 I=1,NMAX
      IF((EBAR.GE.ELL(I)+1))GO TO 14
12  CONTINUE
13  JMAX = (JMIN + JMAX)/2
14  JBAR = JMAX
      PDX = USTAR(STC,BDC,X,C,NST,ST,NRG)
      IF(INST.EQ.0)IGO TO 220
      RUNTOIT = RUNTOIT + STC
      NST = NRG+NST
      PDT = PDT/STC**2
      CALL KERNEL(I,NST,ST,NRG)
      DO 20 I=1,3
      X(I) = X(I) + C(I)*STC
20  C(I) = C(I)
      KMIN = JMIN
      KMAX = JMAX
      DO 30 I=1,NMAX,JMAX
      WC(I) = PDT*NS(I)*EXP(-XMP(I))
30  EC(I) = ES(I)
40  NTALLP = NTALLY
      NTALL = 0
      TOTALN = 0.0
      TOT = 0.0
      IF((NPDMAX.EC.0)) GO TO 90
      NPD = 0
      DO 80 I=1,NDMAX
      IF((DR(I)).GT.0)IGO TO 80
      NPD = NPD + 1
      STM = SATOR(X,KT(I), I,C)
      IF(KKKK.GT.1)IGO TO 50
      CALL SZENCMN,PCT,X,C)
      GO TO 60
50  CALL SINGLENN,COSINE(C,CC),I,0)
      DO 60 IJMN.GT.JMAX)IGO TO 90
      TOT = TOT+STM
      CALL PATHLEN(X,C,NST,ST,NRG,NSC)
      CALL KERNEL(I,NST,ST,NRG)
      CALL DETECT(I,-1,1,0/STM**2,COSINE(C,CDT(I,1),0))
      IF((NPDM.EQ.NPDMAX))IGO TO 90
70  CONTINUE
90  IF(KKKK.GT.1)GO TO 95
      IF((NPDM.EC.0))GO TO 210
95  IF(KKKK.GT.1)GO TO 102
      PDT = PDT
      IF((MODELQ.GT.0))GO TO 100

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      NGT = 0
      GO TO 108
100  FST = VECTORFXCT,X,CC1
      JBAR = JZERO
      GO TO 104
102  PDX = 1.0
      NGT = 1
104  FST = RHO(HH)+SSH(JBAR)
      GST = 0.0
      M = MTL(HH)
      IF(GT,GT,0) GST = SG(JBAR,R)
      ALPHA = 0.0
      HST = FST + GST
      IF(HST,GT,0.0) ALPHA = (FST+ALPHA(JBAR)) + GST+ALPHA(JBAR)/HST
      STM = VECTOR(X,BDC,CP)
      BETA = 0.0
      IF(FST,LE,DELTA) GO TO 108
      CALL PATH(HH,FST-DELTA,X,CP,NST,ST,HRG,MSI)
      XMPT = 0.0
      DO 106 I=1,NST
      J = NRG(I)
      K = MTL(J)
      SIG = RHO(J)*SSH(JBAR)
      IF(K,GT,0) SIG = SIG + SG(JBAR,K)
106  XMPT = XMPT + SIG*ST(I)
      IF(XMPT,GT,0.0) BETA = AS*(ALOG((2.0*XMPT))/4.0
      IF(MODEL,EG,2) BETA = AS*XMPT/2.0
108  DO 205 LLL=1,IOMEGA
      IF(MGT,GT,0) GO TO 110
      PDA = PDX*VSTAR(HH,X,C)
      GO TO 112
110  PDA = PDX*VSTAR(CP,CC,C,BETA,BS*ALPHA)
112  IF(LLL,LT,IOMEGA) PDA = PDA/FLOAT(NNUMBER)
      IF(KKK,GT,1) GO TO 114
      CALL SZERC(HH,PIA,X,C)
      GO TO 116
114  CALL SINGLE(HH,COSINE(C,CC),PDA)
116  IF(IJMIN,GT,JMAX) GO TO 202
      CALL PATH(HH,1.0E+30,X,C,NST,ST,HRG,MSI)
      IF(LLL,EQ,IOMEGA) GO TO 205
      NSTMIN = 1
      STT = 0
      TOT = RUMTOI
      DO 200 N=1,NST
      L = HRG(N)
      M = MTL(L)
      IF(HYDMAX,EQ,0)GO TO 150

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      FAST1409
      FAST1410
      FAST1411
      FAST1412
      FAST1413
      FAST1414
      FAST1415
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      FAST1449
      FAST1450
      FAST1451
      FAST1452
      FAST1453
      FAST1454
      FAST1455

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      DO 140 I=1,NMAX
      IF(IDR(),NE,1)GO TO 140
      IF(IDS(),NE,0)GO TO 140
      ANG = COSINE(C,CGT(),T)
      IF(M,LT,0)IC TO 120
      FST = 1.0/VCL()
      RHUN = RHO(L)
      GO TO 130
120  FST = ST(I)/VCL()
      TOT = TOT + ST(I)/2.0
130  CALL KERNEL(NSTMN,N-1,ST,HRG)
      CALL DETECT(I,M,FST,ANG,ST(I))
      IF(M,LT,0) TOT = TOT - ST(I)/2.0
      NSTMIN = N + 1
      GO TO 150
140  CONTINUE
150  STT = STT + ST(I)
      TOT = RUNTOT + STT
      K = NSCIN()
      IF(K,NE,NSDMAX,LE,0)GO TO 200
      MAX = 2 - N/NST
      NGT = 0
      DO 190 K=1,MAX
      DO 160 I=1,NMAX
      IF(IDR(),EQ,L,AND,IDS(),EQ,K)GO TO 170
160  CONTINUE
      GO TO 190
170  IF(MGT,GT,0)GO TO 180
      NGT = 1
      CALL NORMALIK(STT,X,C,CN)
      ANG = COSINE(C,CN)
      CALL KERNEL(NSTMN,N,ST,HRG)
      NSTMIN = N + 1
180  FST = 1.0/ABS(ANG)*VOL()
      CALL DETECT(I,-1,FST,ANG,0.0)
      L = HRG(N)
190  L = HRG(N+1)
200  CONTINUE
202  IF(LLL,EQ,IOMEGA) GO TO 220
205  CONTINUE
210  CONTINUE
220  CONTINUE
      DO 230 I=1,NMAX
      DO 230 J=1,AGMAX
      FXS(J,I) = FXS(J,I) + EXP(I,J)**2
230  FXP(J,I) = 0.0
235  CONTINUE
240  CONTINUE

```

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      FAST1456
      FAST1457
      FAST1458
      FAST1459
      FAST1460
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      FAST1462
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      FAST1469
      FAST1470
      FAST1471
      FAST1472
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      FAST1497
      FAST1498
      FAST1499
      FAST1500
      FAST1501
      FAST1502

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```

CALL ANSWER(NNN)
CONTINUE
RETURN
END

$IBFCT SOLVE #94/2,47
CSOLVER(SUBROUTINE SOLUTION OF THE BOLTZMANN EQUATION T.H.JORDAN#4NHL1966*
SUBROUTINE SOLVER(L1,L2,L3,XDT,CDF,SF,NRG,NSC,HTL,RHO,SSH,SGT,ELL,
LWS,ES,ESL,ALM,ALH,FXP,FKS,STP,MRP,NCP,LDI)
DIMENSION XT(3),CT(3),C(3),X(3),CC(3),CD(3),XC(3),XP(3),
XT(3,1),CT(3,1),C(3,1),X(3,1),CC(3,1),CD(3,1),XC(3,1),XP(3),
1 SGT(L2,L2),ELL(1),WSL(1),ESL(1),WC(1),EC(1),XPK(1),ALM(1),
2 ALH(1),XPX(1),FXS(L3,1),IRDI)
COMMON/TAPEDATA/WT,WT2/
COMMON/LIMITS/NSTORE,NERROR,NSMAX ,NMAX ,NBMAX ,NSTMAX ,
1 NEMAX ,NVMAX ,NPKMAX ,NXEMAX ,NEMOD ,NSSEC1,NUNITD,
2 NUNITX,NINMAX ,NMRD ,NDRD ,NDOWN ,INELAS ,NTRANS ,
3 NGMAX ,NMMAX ,NFMAX ,NVHOD ,NCMAX ,NLMAX ,NTMAX ,
4 NMODX ,NPDMOD ,NMDOM ,NPDMAX ,NPNDMOD ,NSDMAX ,
5 NUDELX ,NUDEL ,NUDELW ,NUDELW ,NUDELW ,NUDELV ,
6 NPRINT ,NUNITS ,NUMBER ,KALIDE /
COMMON/FLUXES/KXX ,NTALLY ,ERROR ,TCTALN ,TOTALE ,RHON ,SNORM ,JOTFT
COMMON/POINTX/NTOTAL ,III ,HOM
COMMON/OTHERS/RADIUS ,XCT(3) ,DELTA ,BCC(3) ,ATA ,BTA ,ATC ,
1 ATO ,AT ,BT ,AS ,BS ,JMIN ,JMAX ,
2 KMAX ,JBAR ,NZERO
3
4 DIMENSION BDC(3,3)
5 COMMON/ENDRUN/HPDRUN
6 NPDRUN = 0
7 DO 100 I=1,3
100 BDCDUM(I)= BDC(I)
27 IF(IEQ(1,1),NMAX ,NTRL)
1 IF(IEQ(1,1),GTL,GT .01 GO TO 27
NPDRUN = NPDRUN + 1
DO 1 I=1,3
1 XT(I)= XDT((I,(I,I))
1 CT(I)= CDT((I,(I,I))
1 CC(I)= CC((I,(I,I))
1 CD(I)= CD((I,(I,I))
1 XC(I)= XC((I,(I,I))
1 XP(I)= XP((I,(I,I))
1 SNORM = MAX((1.00 ,0.0 ,WEVCR(XCT,XT,C1))
DO 30 NNN=1,NPRINT
30 29 NMM=1,NUNITS
IF(MODELP,GT ,.01 GO TO 93
NHALF = 1
PDS = 1.0 ,PARIN,X1
1 IF(NNEQ,GT ,0) GO TO 26
GO TO 94
93 CONTINUE

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```

XMP1 = XMPT + SIG*ST(1)
IF(NGT,EQ,1)GO TO 9
STT = STT + ST(1)
IF(STT,LT,STH)GO TO 9
XMPH = XMPT - SIG*(STT - STH)
NOT H
9 CONTINUE
BETA = 0.0
IF(NHALF,EQ,2)XMPH = XMPT - XMPH
IF(XMPH,GT,0.01) BETA = AS*(ALOG(12.0*XMPH))/4.0
IF(MOD(LV,EG,2) BETA = AS*XMPH/2.0
NNC = NN
IF(IKKK,NEQ,0)GO TO 10
IF(XMPH,LT,XMPT)GO TO 11
10 PDF = 1.0
H = 0.0
GO TO 12
11 IF(XMPH,LEQ,0.0)GO TO 12
PHALF = XMPH/XMPT
IF(RANDN(HMB1,GT,PHALF))GO TO 15
PDF = 1.0/PHALF
H = 0.0
12 NHALF = 1
DO 13 I=1,3
CD(I) = C(I)
BD(C) = XT(I)
13 XC(I) = X(I)
GO TO 14
14 PDF = 1.0
NGT = 0
GO TO 16
15 PDF = 1.0/(L0 - PHALF)
H = 0.0
16 NHALF = 2
DO 17 I=1,3
XC(I) = X(I)
CD(I) = -C(I)
BD(C) = XT(I)
17 XT(I) = XT(I)
NN = NZERO
18 ALPHA = 0.0
IF(XMPT,GT,0.0)ALPHA=(SIGH*ALH(JBAR)+(XMPT-SIGH)*ALH(JBAR))/XMPT
IF(IKKK,NEQ,0)GO TO 300
PDA = QSTAR(NN,X,C)
GO TO 301
300 PDA = VSTAR(CD,CC,C,BETA,BS+ALPHA)

```

```

301 CONTINUE
STM = 1.0E+30
IF(NGT,EQ,0)GO TO 19
PST = 1.0E+30
IF(PST,GT,0.01)STH=STM/PSI
19 CALL PATHNN(STM,X,C,NST,ST,NRG,NSC)
PDT = USTAR(LST,BC,C,X,C,NST,ST,NRG)
IF(NST,EQ,0)GO TO 26
NN = NRG(NST)
DO 20 I=1,2
XP(I) = XC(I) + STC*C(I)
20 XC(I) = XC(I)
STM = VECTOR(X,XP,C)
NN = NNC
PDT = STM*PDA*PDT/STM**2
IF(IKKK,GT,1)GO TO 21
CALL ZEROINN(NSP,X,C)
GO TO 22
21 CALL SINGLE(NN,COSINE(C,CC),1,C)
22 IF(IJMIN,GT,JMAX)GO TO 26
IF(IJMIN,LT,2)GO TO 200
CALL KERNEL(L1,NST,ST,NRG)
GO TO 201
200 CALL PATHNN(STM,X,C,NSTP,STP,ARP,NCP)
CALL KERNEL(L1,NSTP,STP,NRP)
201 CONTINUE
DO 23 J=JMIN,JMAX
-EC(I) = ES(I)
23 WC(I) = PDT*WS(I)*EXP(-XMP(I))
-KMIN = JMIN
KMAX = JMAX
NN = NNP
RUNT = RUNTOT + STM
DO 24 I=1,2
X(I) = XP(I)
24 CC(I) = C(I)
25 CONTINUE
26 CONTINUE
DO 27 J=1,NDDO
DO 28 J=1,NMAX
FXS(J,I) = FXSJ,I) + EXP(J,I)**2
28 FXP(J,I) = 0.0
27 CONTINUE
29 CONTINUE
CALL ANSWER(NNN)
30 CONTINUE
27 CONTINUE

```

```

NZERO = LCATE(NGS,C)
STM = VECTTOP(XC,XCT,C)
CALL PATHNZERO(STM,X,C,NST,ST,NRG,NSC)
I = NST + 1
DO 1 K=1,NST
I = I+1
L = NRG(I)
J = LSV(L)
IF(J,L.GT.0) GO TO 2
1 CONTINUE
JZERO = NMHDOD/?
GO TO 100 .
2 NMHDOD = 0
C CALCULATE AVERAGE SOURCE GROUP INDEX
TOTN = 0.0
TOTE = 0.0
MIN = JSN(I,J)
MAX = JSX(I,J)
JMIN = MIN
JMAX = MAX
DO 10 I=MIN,MAX
WSI(I) = SPW(I,J)
10 ESI(I) = SPE(I,J)
IF(NTT.EQ.1)GO TO 30
CALL KERNEL(NI,NIT,NT,ST,NRG)
IF(XMP(1).LT.-0.010D0)GO TO 30
30 FST = GIMMI*WST
IF(FST.LT.0.0)FST = FST/XMP(1)
TOTN = TOTN + FST
40 TOTE = TOTE + FST*ESI()
EBAR = TOTE/TOTN
DO 40 I=MIN,MAX
IF(EBAR.GE.ELL(I+1))GO TO 60
50 CONTINUE
I = (MIN + MAX)/2
60 JZERO = I
100 RETURN . .
END

$IRFC SIZEOF MP4/XR7
CSZEROF RED SOURCE EVALUATION FOR PROGRAM FASTER@T.
SUBROUTINE SZEROF(NT,PDF,XXX,CCC)
DIMENSION XXX(1),CCC(1)

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3      NMMAX ,NMAX ,NEMAX ,NVMD ,NCMAX ,LMAX ,NTMAX ,
4      NDMAX ,NNDOD ,NNDENT ,NNDOD ,NPDDOD ,NNDMAX ,
5      NVDDC ,NVDC ,NPDOD ,NPDENT ,NPDOD ,NPDENT ,NPDOD ,NPDOD ,
6      NPRINT ,NUNITS ,NUMBER ,KALIDE
COMMON /OTHERS/ RADIUS ,XC(13) ,DELTA ,BC(13) ,ATA ,ATB ,ATC ,
1          ATD ,AT ,BT ,AS ,BS ,JMIN ,JMAX ,
2          KMT ,KMAX ,JBAR ,NZERO
COMMON /LLREG/ LREG
COMMON /SPASHL/ PSORS
IF(NNN.GT.0) N = ISVN(NN)
N = LLREG
MON = 0
IF(NVMD.EQ.01) GO TO 3
DO 10 I=1,NVMD
IF(I.GT.1) EG,H1 GO TO 2
1 CONTINUE
GO TO 3
2 MON = 1
3 CONTINUE
C UN-TRANSLATE
4 DO 10 I=1,3
10 Z(I) = X(I) - XTR(I,N)
C SOURCE GEOMETRY CHECK
IF(NSCG(N).GT.0)GO TO 30
C RECTANGULAR
5 DO 20 I=1,3
20 D(I) = C11
C SPATIAL COORDINATES
      V(I) = Z(I)
C DIRECTION COSINES
      20 D(I) = C11
      IF(NNN.LT.105)105,105,100
C CYLINDRICAL AND SPHERICAL, AZIMUTHAL ANGLE
      105 DO 110 I=1,3
      110 IF((ABS(Z1)) + ABS(Z2)).GT.0.01 VZ = ATAN2(Z2,Z1)
C RADIUS NEXT
      MAX = NSG(N) + 1
      VI = 0.0
      DO 40 I=1,MAX
      40 V(I) = V(I)*I**2
      VI = SORT(V1)
      IF(MAX.EQ.3)GO TO 50
C CYLINDRICAL, Z-COORDINATE
      V3 = Z3
      IF(NNN.LE.01) GO TO 105
C ROTATION SETUP, POLAR-
      R = V
      CPH = 1.0

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```

SPH = 0.0
GO TO 40
C SPHERICAL, COSINE OF POLAR ANGLE
50 V3 = Z3/V1
      IF(INN,LE,0) GO TO 105
C ROTATION SETUP, POLAR
      CPH = 0.0
      STH = 0.0
      R = V1*SPH
      R = V1*SPH
C ROTATION SETUP, AZIMUTHAL
      60 IF(R,GT,0.0)GO TO 7C
      CTH = 1.0
      STH = 0.0
      GO TO 10
      70 CTH = Z1/R
      STH = Z2/R
C CALCULATE ROTATION MATRIX
      80 CALL ROTATE(CPH,SPH,CTH,STH,ROT)
C ROTATED DIRECTION COSINES
      DO 90 I=1,3
      U(I) = 0.0
      DO 90 J=1,3
      90 D(IJ) = D(IJ) + ROT(I,J)*C(J)
C COSINE OF POLAR ANGLE, ANGULAR PDF
      100 VS = D3
C AZIMUTHAL ANGLE, ANGULAR PDF
      110 VV = 0.0
      120 IF((ABS(D(1,1)) + ABS(D(2,1)),GT,0.0) VV = ATAN2(D2,D1))
C ALL SOURCE VARIABLES DEFINED, SET ERROR INDICATOR, TOTAL PDF
      JMIN = JSKIN
      JMAX = 0
      105 CONTINUE
      130 CONTINUE
      140 I=1,5
C CHECK FOR DELTA FUNCTION, ASSUMED CORRECT
      IF(NPC(I,N).EQ.1)GO TO 130
C NOT DELTA FUNCTION, CHECK RANGE
      K = NPC(I,N)
      IF(VEE(I,1) - VEE(I,1,N)) * (VEE(K,I,N) - V(I)) .LT. 0.0)GO TO 150
C INSIDE RANGE, STORE IN MAX
      150 JMAX = 2,2,K
      IF(VEE(J,I,N).GE.V(I))GO TO 120
      110 CONTINUE
C LINEAR INTERPOLATION
      120 K = J - 1
      PDS=PDS*(VAL(J,I,N)*(V(I)-VEE(K,I,N))+VAL(K,I,N)*(VEE(J,I,N)-V(I)))
      1 / (VEE(J,I,N)-VEE(K,I,N))
      IF(INN,GT,0) GO TO 130
      FASTI1874
      FASTI1875
      FASTI1876
      FASTI1877
      FASTI1878
      FASTI1879
      FASTI1880
      FASTI1881
      FASTI1882
      FASTI1883
      FASTI1884
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      FASTI1886
      FASTI1887
      FASTI1888
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      FASTI1893
      FASTI1894
      FASTI1895
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      FASTI1898
      FASTI1899
      FASTI1900
      FASTI1901
      FASTI1902
      FASTI1903
      FASTI1904
      FASTI1905
      FASTI1906
      FASTI1907
      FASTI1908
      FASTI1909
      FASTI1910
      FASTI1911
      FASTI1912
      FASTI1913
      FASTI1914
      FASTI1915
      FASTI1916
      FASTI1917
      FASTI1918

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224

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      IF(I,LT,-3) GO TO 130
      PDSRS = PDS
      GO TO 150
      130 CONTINUE
C DEFINE NUMBER IN GROUP AND AVERAGE ENERGY
      JMAX = JSKIN
      DO 140 I=JMIN,JMAX
      140 I=JMIN,JMAX
      150 I=JSKIN
      160 I=1,N
      170 I=1,N
      180 I=1,N
      190 I=1,N
      200 I=1,N
      210 I=1,N
      220 I=1,N
      230 I=1,N
      240 I=1,N
      250 I=1,N
      260 I=1,N
      270 I=1,N
      280 I=1,N
      290 I=1,N
      300 I=1,N
      310 I=1,N
      320 I=1,N
      330 I=1,N
      340 I=1,N
      350 I=1,N
      360 I=1,N
      370 I=1,N
      380 I=1,N
      390 I=1,N
      400 I=1,N
      410 I=1,N
      420 I=1,N
      430 I=1,N
      440 I=1,N
      450 I=1,N
      460 I=1,N
      470 I=1,N
      480 I=1,N
      490 I=1,N
      500 I=1,N
      510 I=1,N
      520 I=1,N
      530 I=1,N
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      550 I=1,N
      560 I=1,N
      570 I=1,N
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      590 I=1,N
      600 I=1,N
      610 I=1,N
      620 I=1,N
      630 I=1,N
      640 I=1,N
      650 I=1,N
      660 I=1,N
      670 I=1,N
      680 I=1,N
      690 I=1,N
      700 I=1,N
      710 I=1,N
      720 I=1,N
      730 I=1,N
      740 I=1,N
      750 I=1,N
      760 I=1,N
      770 I=1,N
      780 I=1,N
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      810 I=1,N
      820 I=1,N
      830 I=1,N
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      870 I=1,N
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      1010 I=1,N
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      1120 I=1,N
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      1170 I=1,N
      1180 I=1,N
      1190 I=1,N
      1200 I=1,N
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      1470 I=1,N
      1480 I=1,N
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      1500 I=1,N
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      1890 I=1,N
      1900 I=1,N
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      1920 I=1,N
      1930 I=1,N
      1940 I=1,N
      1950 I=1,N
      1960 I=1,N
      1970 I=1,N
      1980 I=1,N
      1990 I=1,N
      2000 I=1,N
      2010 I=1,N
      2020 I=1,N
      2030 I=1,N
      2040 I=1,N
      2050 I=1,N
      2060 I=1,N
      2070 I=1,N
      2080 I=1,N
      2090 I=1,N
      2100 I=1,N
      2110 I=1,N
      2120 I=1,N
      2130 I=1,N
      2140 I=1,N
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      2190 I=1,N
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      2770 I=1,N
      2780 I=1,N
      2790 I=1,N
      2800 I=1,N
      2810 I=1,N
      2820 I=1,N
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      2860 I=1,N
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      2970 I=1,N
      2980 I=1,N
      2990 I=1,N
      3000 I=1,N
      3010 I=1,N
      3020 I=1,N
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      3070 I=1,N
      3080 I=1,N
      3090 I=1,N
      3100 I=1,N
      3110 I=1,N
      3120 I=1,N
      3130 I=1,N
      3140 I=1,N
      3150 I=1,N
      3160 I=1,N
      3170 I=1,N
      3180 I=1,N
      3190 I=1,N
      3200 I=1,N
      3210 I=1,N
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      3270 I=1,N
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      10330 I=1
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6   [ISPM ,ISPE ,IATD ,IRSP ,ISGT ,IIDE ,IIDI , * FAST2013
7   ISGS ,NEXT ,IXMP ,ISGR ,IMS ,IES ,IMC , * FAST2014
8   IEC ,INSC ,ISPR ,ITMP ,IMR ,IST ,ISI , * FAST2015
9   [INRG ,INSC ,ISPT ,INRP ,INCP ,IFXP ,IFKS , * FAST2016
1   IFXT ,IFXE ] * FAST2017
COMMON/REGCSA/NSRMAX,INSR
DIMENSION N11
EQUIVALENCE (N11,H11)
NN = N11
P1 = AAA
PDA = BBB
INSR = 0
IF(NSRMAX.LE.0) GO TO 10
DO 5 I=1,NSRMAX
  IF(NN.EQ.N11) INSR = I
  5 CONTINUE
10 CONTINUE
CALL SINDUM(NDOWN,NORDER,1,INELAS,NENAM,HL(MTL),H(IRHO),H(IIDE),
  1 H(11),H(1ISGS),H(1IES),H(1IRH),H(1IDEN),
  2 H(1INS),H(1IES),H(1ICL),H(1ECI),H(1ELL))
RETURN
END
SIBFCT SINDUM M94/2,XR7
CSINDUM SINGLE SCATTERING CALCULATION FOR PROGRAM FASTER+T.M.JORDAN+MANLF2036
SUBROUTINE SINDUM(L1,L2,L3,L4,LS,MTL,RHO,IDE,IOT ,SGS,ESB,SSH,
1  DEN,WS,E5,EC,ELL)
1  DIMENSION MTL(1),RHOD(1),DELL(1),L2(L1),L3(L1),L4(L4,L3),GS51(1,1),ESB1(1)
1  COMMON/DSUMS/INH,PSI,PSI1,PSI2,PSI3,PSI4,PSI5,PSI6,PSI7,PSI8,PSI9,PSI10,PSI11
COMMON/DUMS/INH, * PSI
COMMON/OTHERS/RADIUS,XCT3(1),DELTA ,BDC1(3),ATA ,ATB ,ATC ,
1  ATD ,AT ,BT ,AS ,BS ,JINR ,JMAX , * FAST2047
1  KNIN ,KMAX ,JBAR ,NZERO , * FAST2048
COMMON/LIMITS/NSTONE,NMAX,NMMAX,NMAXX,NMAXM,NEMAX ,NSTMAX,
1  NMINTX,NMMAX ,NMMAX ,NMAXX,NMAXM,NEMOD ,NSECST,NUNITD , * FAST2049
2  NMINTX,NMMAX ,NMMAX ,NMAXX,NMAXM,NMAX ,NMINTX,NTRANS , * FAST2050
3  NMMAX ,NGMAX ,NFMPC ,NVMD ,NCMAX ,NLMAX ,NITMAX , * FAST2051
4  NGMAX ,NGMOD ,NOPENT ,NMOD ,NPDMX ,NPDMOD ,NMODMAX , * FAST2052
5  NVDMAX ,NVONCD ,NPGINT ,MODELQ ,MODEL0 ,NUDELV , * FAST2053
6  NPRINT,NUNITS,NUMBER,KALIDE ] * FAST2054
PST = 0.0
H = MTL(1)
JMIN = NEMAX + 1
JMAX = 0
DO 10 I=1,NEMAX
  WS(I) = 0.0
10 ENDDO
15 IF(MT.LT.1) GO TO 820

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IF(NXNSECT.NE.1)GO TO 220
      PDA = PDA
100 IF(KM=0.0)GO TO 175
      PLM2 = 0.0
      PLM1 = 1.0
      DO 130 I=L1,NORDER
      DD = 1.0/NODDN
      IF(I>DEK,L1,M1-LE,0)GO TO 130
      KMIN = KNIN
      N = IDEK,L1,M1
      KMDN = MIN0(KMAX,NEMAX+1-K)
      IF(KMIN.GT.KMDN)GO TO 130
      I = KMIN + K - 2
      DO 120 J=KMIN,KMDN
      I = I + 1
110 W$1(J) = W$1(J) + WC(J)*SGS(J,N)*PSIN
      JMAX = MAX0(JMAX,I)
120 CONTINUE
130 PLMZ=PSI*PLN1*FLOAT(2*L-1)-PLM2*FLOAT(L-1)/FLOAT(L)
      NM1 = KM1
      NM2 = KNIN
      PSIN = PDA*PLMZ
139 CONTINUE
      KMDN = MIN0(KMAX,INELAS)
      IF(KMDN.LT.KMIN) GO TO 160
      DO 150 J=NM1,KMDN
      I = J - 1
      K = IDEK,J,N
      IF(K.EQ.0) GO TO 160
      N = K/1000
      MAX = MIN0(NEMOD-J,K-1000*N)
      I = J - 1
      DO 140 K=I,MAX
      I = I + 1
140 W$1(J) = W$1(J) + WC(J)*SGSIK,H)*PDA
      JMAX = MAX0(I,JMAX)
150 CONTINUE
160 DO 170 I=JMIN,JMAX
170 ES1(I) = ESB(I)*W$1(I)
      FS1 = DENIM
175 FS1 = FS1*PSI*PDA*(FS1 + RHO(NNN))
      IF(FS1.LE.0.0)GO TO 900
      RAT = PSI**2
      MIN = KNIN
      DO 200 J=KNIN,KMAX
      E = RAT*EC(J)
      DO 180 I=MIN,NEMAX
      IF(E.GE.ELL(I+1))GO TO 190

180 CONTINUE
      GO TO 210
190 MIN = I
      JMIN = MIN0(IJMIN,I)
      WSC = FS1*SSH(J1)*WC(J)
      W$1(I) = W$1(I) + WSC
200 ES1(I) = ES1(I) + WSC*E
210 JMAX = MAX0(JMIN,JMAX)
      GO TO 900
COMPTON SCATTERING
220 MIN = KNIN
      PSIN = 1.0 - PSI
      PSIS = PSI**2 - 1.0
      GST = 0.511/(1.0+0.49875
      IF(HG.GT.0.1)GST = GST + DER(H)
      GST = PDA*GST
      DO 250 J=KNIN,KMAX
      RAT = 0.511/(1.0+1.0*PSI*EC(J))
      E = RAT*EC(J)
      DO 230 I=MIN,NEMAX
      IF(E.GE.ELL(I+1))GO TO 240
230 CONTINUE
      GO TO 260
240 MIN = I
      JMIN = MIN0(IJMIN,I)
      FS1 = WC(J)*GST*RAT**2*(RAT + 1.0/RAT + PS(S)
      W$1(I) = W$1(I) + FS1
      ES1(I) = ES1(I) + FS1*E
250 CONTINUE
260 JMAX = MIN
900 IF((JMIN.GT.JMAX))GO TO 920
      DO 910 I=JMIN,JMAX
      IF(W$1(I).NE.0.0) ES1(I) = ES1(I)/W$1(I)
910 CONTINUE
920 RETURN
      END

$16FTC TRACE M94/2,XR7
CTRACER//PARTIAL PATH LENGTHS FOR PROGRAM FASTER+T,H,JORDAN+MANL+1966+
SUBROUTINE PATH (INN,SSS,XXX,CCC,MMN,ST,NG,NSC)
COMMON /UNIT1/ H1
COMMON /DUNTRA/ NSSTORE,NERORB,NSMAX,NMAX,HBMAX,NSMAX,N$TMX,
1      NMAX,NNMAX ,NMMAX ,NMMAX ,N$MAX,NCEDR,NDDN,INELAS,NTMAX,
2      NUNITX,NNMAX ,NMMAX ,NCEDR,NDDN,INELAS,NTMAX,
3      NNMAX ,NMMAX ,NMMAX ,NMMAX ,NCEDR,NDDN,INELAS,NTMAX,
4      NMMAX ,NCEDR,NCEDR,NDDN,INELAS,NTMAX
      END
```

```

6      NPRINT,NUNITS,NUMBER,KALIDE      FAST2154
COMMON/INDEXS/INTP ,IAE ,IISV ,IPT ,IEND ,IRR ,IELL ,FAST2155
1      IEND ,IRE ,IOP ,IHSG ,INPC ,IJSN ,ILISX ,FAST2156
2      ILSTR ,IATM ,IATM ,IATM ,IJSR ,ITSSH ,IDEN ,FAST2157
3      INTG ,IELF ,IFGW ,ITDS ,ILOV ,IIOR ,IIDS ,FAST2158
4      IVOL ,ICDT ,IXOT ,IRSI ,IALP ,IVMD ,IGIM ,FAST2159
5      IAIM ,ISPE ,IAHL ,IA ,INS ,IVEE ,IVAM ,FAST2160
6      ISPM ,ISPE ,ITDT ,IRSP ,ISCT ,IIDE ,IIDT ,FAST2161
7      IEGS ,IET ,IIMP ,ISGR ,IMS ,IRES ,IEST ,FAST2162
8      IEC ,IHO ,IU ,IIS ,IISI ,IST ,FAST2163
9      INRG ,INSC ,ISTP ,IHRP ,IACP ,IFXP ,IFXS ,FAST2164
1      IFXT ,IFXE ,IFXA ,DIMENSION XXX(1),CCC(1)      FAST2165
NN = NNN
STM = SSS
DO 10 I=1,3
  X(I) = XXX(I)
10 CCC(1) = CCC(I)
  CALL TRADM(NMAX,NBMAX,1,H(INTP),H(IAZ),H(A),H(INS),H(END),
  1           HEIU),H(CIV),HEIW),HEISI),ST,NRG,NSC)
  HMM = NST
  RETURN
  END
$IBFTC TRADM M94/2,XRT
CTRADUM PARTIAL PATH LENGTHS FOR PROGRAM FASTER@T.M.JORDAN/WANL*1966
SUBROUTINE TRADM(L1,L2,L3,HTP,AZ,A,N,ND,UD,V,W,S1,ST,NRG,NSC)
DIMENSION NSC(1),NSL(1),NSL1(1),NSL2(1),ST(1),NRG(1),NSC(L1),
1           ND(1),VLL(1),VLL1,WLL1,ST(1),NSL(1),NSL1(1),NSL2(1),
COMMON/DUNTRA/NN ,STM ,NST ,X191 ,C191 ,CX(91
COMMON/RNTRNS/NEMAX ,NVMAX ,NMAMAX ,NMAX ,NRMAX ,NSTMAX,
1           NUNITH ,NIMAX ,NMAMAX ,NMAX ,NRMAX ,NSTMAX,
2           NMAM ,NMAM ,NMAX ,NMAX ,NRMAX ,NSTMAX ,
3           NMAM ,NMAM ,NMAX ,NMAX ,NRMAX ,NSTMAX ,
4           NDMAX ,NCMOD ,NMDENT ,NDMOD ,NMDPNA ,NMDMOD ,NSDMAX ,
5           NVDMAX ,NVDMOD ,NPPOINT ,MODELPA ,MODELQ ,MODELW ,
6           NPRINT,NUNITS,NUMBER,KALIDE      FAST2166
  DO 110 I=1,3
    J = I + 1 - 3*(I/3)
    X(J) = X(I)*X(J)
    X(I+6) = X(I)*X(J)
    C(I+3) = 2.0*(C(I)*X(I)
    C(I+6) = C(I)*X(J) + C(J)*X(I)
    CX(I+3) = C(I)*X(2)
  110 CX(I+6) = C(I)*C(J)
  DO 120 I=1,NMAX
  120 ND(I) = 0
    II = NN

```

```

    STM = 0.0
    DO 410 I=1,NSTMAX
    DO 400 M=1,3
    GO TO(130,140,160),M
130 IMIN = II
    JJ = 0
    SP = 1.0E+38
    GO TO 140
140 IF(INXJ-EQ.0.0GO TO 400
    IMIN = NKT
    GO TO 170
160 IMIN = NKT
    IML = NMAX
    DO 390 I=IMIN,IMAX
    IF(WLLE-21GO TO 190
    IF((NXY - I)*IP - I)-EQ.0.0GO TO 390
    DO 180 J=1,NBMAX
    IF(NSL(J,I)-EQ.0.0GO TO 390
    IF(NSL(J,I)-NSL(J,I)/1000).EQ.1000GO TO 190
180 CONTINUE
    GO TO 390
190 DO 340 J=1,NBMAX
    IF(NSL(J,I)-EQ.0.0GO TO 350
    KK = NSL(J,I)/1000000
    K = NSL(J,I) - 1000*KK
    IF(K-EQ.0.0AND.(K-EQ.LL))GO TO 340
    KK = 2*KK - 1
    IF(ND(K))GT.0.0GO TO 210
    ND(K) = 1
    UIKJ = AZ(K)
    VIKJ = 0.0
    WIKJ = 0.0
    MAX = NIP(K).
    DO 200 L=1,MAX
    UIKJ = UIKJ + X(L)*A(L,K)
    VIKJ = VIKJ + C(L)*A(L,K)
    WIKJ = C(L)*W(L,K)
    200 CONTINUE
210 IF(WLLE-1)GO TO 220
    IF(IFLOAT(KK)-IUI(X) + STT*(V(K) + STT*N(K))) 340,34C,390
220 IF(ND(K))- 21 230,230,340
230 ND(K) = 2
    L = 100,250,240,260
240 IF(UIKJ-EQ.0.0GO TO 310
    SII(L,K) = -UIKJ/V(K)
    IF(SII(L,K).LT.0.0)GO TO 310

```

```

      SI(2,K) = SI(1,K)
      GO TO 290
  250 L = 2
      GO TO 270
  260 L = 1
  270 E = -0.5*V(XK)/W(K)
      F = U(K)/W(K)
      H = E*2 - F
      IF(F,LE,0.0)GO TO 280
      IF(E,LE,0.0)GO TO 310
      IF(H,LE,0.0)GO TO 310
  280 L=1,SORT(H)
      SI(L,K)=E-H
      L=L+1-2*(L/2)
      SI(L,K)=E+H .
  290 L=1
      IF(W(K),GT,0.01)=2
      SD=SI(L,K)-STT
      IF(SD) 300,320,330
  300 L=L+1-2*(L/2)
      IF(SI(L,K),GT,STT)GO TO 340
  310 ND(K)=3
      GO TO 340
  320 IF(PLOA(IKK)*(VKI)+2.0*STT*W(K)).LT.0.01GO TO 340
  330 IF(PB,GE,SP)GO TO 340
      SP=SP
      JJ=J
      LL=K
  340 CONTINUE
  350 IF(M,LE,1)GO TO 360
      NS(J,J,IP)=1000*NSP+I
      GO TO 410
  360 IF(IJJ,EQ,0)GO TO 415
  370 NRG(N)=I
      IF(STM,GT,STT+SP)GO TO 380
      NS(N)=N
      STM=STM-STT
      GO TO 420
  380 NSC(N)=LL
      STM>SP
      STM=STM+SP
      IF(INTP(LL,LE,3)*IND(LL)=3
      NS=N+INTP(JJ,LE,1)/1000
      NXT=NS(JJ,1)-1000*NSP
      NXK=0
      IF(NXT,LE,NRMAX) GO TO 385
      FAST2249
      FAST2250
      FAST2251
      FAST2252
      FAST2253
      FAST2254
      FAST2255
      FAST2256
      FAST2257
      FAST2258
      FAST2259
      FAST2260
      FAST2261
      FAST2262
      FAST2263
      FAST2264
      FAST2265
      FAST2266
      FAST2267
      FAST2268
      FAST2269
      FAST2270
      FAST2271
      FAST2272
      FAST2273
      FAST2274
      FAST2275
      FAST2276
      FAST2277
      FAST2278
      FAST2279
      FAST2280
      FAST2281
      FAST2282
      FAST2283
      FAST2284
      FAST2285
      FAST2286
      FAST2287
      FAST2288
      FAST2289
      FAST2290
      FAST2291
      FAST2292
      FAST2293
      FAST2294
      FAST2295
      FAST2296
      FAST2297
      FAST2298
      FAST2299
      FAST2300
      FAST2301
      FAST2302
      FAST2303
      FAST2304
      FAST2305
      FAST2306
      FAST2307
      FAST2308
      FAST2309
      FAST2310
      FAST2311
      FAST2312
      FAST2313
      FAST2314
      FAST2315
      FAST2316
      FAST2317
      FAST2318
      FAST2319
      FAST2320
      FAST2321
      FAST2322
      FAST2323
      FAST2324
      FAST2325
      FAST2326
      FAST2327
      FAST2328
      FAST2329
      FAST2330
      FAST2331
      FAST2332
      FAST2333
      FAST2334
      FAST2335
      FAST2336
      FAST2337
      FAST2338
      FAST2339
      FAST2340
      FAST2341
  385 IP=I
  390 CONTINUE
  400 CONTINUE
      NS(J,J,IP)=1000*NSP+MIN(999,NRMAX*NRXX+1)
      GO TO 420
  410 CONTINUE
      NST=NSTMAX+1
  415 N=N+1
  420 NST=N
  430 RETURN
      END
      $IBFT NORML M9/2,XRT
      CNDRL+SURFACE NORMAL CALCULATION FOR PROGRAM FASTER-T.M.JORDAN#WANL#66FAST2310
      SUBROUTINE NORML(NNN,SSS,XXX,CCC,DDD)
      COMMON/H1/
      COMMON/LIMITS/NSTORE,NERROR,NMAX ,NAMAX ,NRMAX ,NBMAX ,NSTMAX ,
      NEMAX ,NVMAX ,NMAX ,NEMAX ,NEMOD ,NSEC,NUNITD, /FAST2311
      NUNITX,NMAX ,NMAX ,NORDER ,NDOMA ,NELAS ,NTRANS , /FAST2312
      NMODA ,NMOD ,NEMOD ,NPOINT ,NMDA ,NPNDXA ,NPCHD ,NSDMAX , /FAST2313
      NVMAX ,NVDMOD ,NPPOINT ,NMDP ,NPQDGL ,MODELU ,MODELV , /FAST2314
      NPINT ,NUNITS ,NUMBER ,KALIDE , /FAST2315
      COMMON/INDEXS/INTP ,IAZ ,IISV ,INIL ,IRHO ,IXR ,IELL , /FAST2316
      IELW ,IAE ,IBE ,INSG ,INPC ,IJSN ,JSK , /FAST2317
      IXTR ,ISUV ,IATN ,IATG ,IESB ,ISSN ,IDEN , /FAST2318
      ITG ,ITF ,ITG ,ITF ,ITOS ,ITD ,ITD ,ITD , /FAST2319
      IVA ,ICOT ,ICOT ,ICOT ,ICOT ,ICOT ,ICOT ,ICOT , /FAST2320
      IATM ,IATM ,IATM ,IATM ,IATM ,IATM ,IATM ,IATM , /FAST2321
      ISP ,ISP ,ISP ,ISP ,ISP ,ISP ,ISP ,ISP , /FAST2322
      ISGS ,NEXT ,IXMP ,ISGR ,ISGT ,IIDE ,IID1 , /FAST2323
      TEE ,IND ,IU ,IV ,IW ,ISI ,IST , /FAST2324
      INRG ,INSC ,ISP ,INRP ,INCP ,IFXP ,IFXS , /FAST2325
      IFAE ,IFXA , /FAST2326
      CALL NORDUM(NMAX,1,HNIMP1,H1(A),NNN,SSS,XXX,CCC,DDD) /FAST2327
      RETURN
      END
      $IBFT NORDM1 M9/2,XRT
      CNDRL+SURFACE NORMAL CALCULATION FOR PROGRAM FASTER-T.M.JORDAN#WANL#66FAST2328
      SUBROUTINE NORDM1(L1,L2,NIP,A,NNN,STT,XP,CP)
      DIMENSION NIP(L1:L2),A(L1:L2),XP(3),CP(3),X(3)
      DO 10 I=1,3
      10 X(I)=XP(I)+STT*CP(I)
      I = NNN
      FST = 0.0
      
```

```

L = 2
MAX = NTP(1)
DO 20 J=1,3
  X(J) = 0.0
  K = L
  L = K + 1 - 3*(K/3)
  IF(L>6) LE ,MAX1(J) = X(L)*A(L+6,I)
  IF(J>6) LE ,MAX1(J) = C(J) + X(K)*A(J+6,I)
  IF(J>3) LE ,MAX1(J) = C(J) + 2.0*X(J)*A(J+3,I)
  IF(J>1) LE ,MAX1(J) = C(J) + A(J,I)
20  FST = FST + MAX1(J)
  IF(FST.NE.-1.0)FST = SORT(FST)
  DO 30 J=1,3
    C(JJ) = C(J)/FST
  30  RETURN
END

$IBFTC ROTAT M94/2,XR7
CROTATE*ROTATION MATRIX
  SUBROUTINE ROTATE(CPH,SPH,CTH,STH,ROT)
  DIMENSION ROT(3,3)
10  ROT(1,1) = CTH*CPH
  ROT(1,2) = STH*CPH
  ROT(1,3) = STH*SPH
  ROT(2,1) = STH*CPH
  ROT(2,2) = CTH
  ROT(2,3) = ROT(1,3)
  ROT(3,1) = -SPH
  ROT(3,2) = 0.0
  ROT(3,3) = CPH
  RETURN
END

$IBFTC CROT M94/2,XR7
CROTATE*ROTATION MATRIX GIVEN DIRECTION VECTOR
  SUBROUTINE ROTATC(C,ROT)
  DIMENSION C(3)
  CPH = C(3)
  SPH = SORT1(MAX1(0.0,1.0 - CPH**2))
  IF(SPH.GT.0.0)GO TO 10
  CTH = 1.0
  STH = 0.0
  GO TO 20
10  CTH = C(1)/SPH
  STH = C(2)/SPH
20  CALL ROTATE(CPH,SPH,CTH,STH,ROT)
  RETURN
END

$IBFTC DOT M94/2,XR7

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```

CCOSINE*DOT PRODUCT
  FUNCTION COSINE(X,Y)
  DIMENSION X(3),Y(3)
  A = 0.0
  DO 10 I=1,3
    A = A + X(I)*Y(I)
  10  COSINE = A
  RETURN
END

$IBFTC TALLY M94/2,XR7
CDETECT*FLUX CALCULATIONS FOR PROGRAM FASTER.T.M.JORDAN*WANL*1966
  SUBROUTINE DETECT(I,II,JJJ,AAA,BBB,CCC)
  COMMON/DOUTDET/   I,N,FST ,ANG ,STV
  COMMON/LIMITS/NSTORE,NERROR,NMAX ,NMAX ,NMAX ,NMMAX ,NSTMAX ,
1   NMAX ,NMMAX ,NNMAX ,NEMAX ,NEMOD ,NSECT,NUNITD,
2   NUNITX,NIMAX ,NMMAX ,NORDER,NNDOW ,NELAS,NTRANS ,
3   NMMAX ,NCHMAX ,NMMAX ,NVMD ,NCMAX ,NLMAX ,NTMAX ,
4   NDMAX ,NPDM ,NPDM ,NPDMOD ,NPDMOD ,NSDMAX ,
5   NMVMAX ,NVMD ,NPDM ,NPDMOD ,NDEGL ,NDEGL ,MODEL ,NDELV ,
6   NPRINT,NUNITS,NNUMBER,KALIDE
  COMMON/INDEXES/INTP ,ITAZ ,IISV ,IMIL ,IRHD ,IXR ,JELL ,
1   IELM ,IAE ,IBE ,INSG ,INPC ,IJSR ,IJSX ,FAST2389
2   ITRX ,ISUV ,IATM ,IATD ,IESB ,ISSR ,IDEN ,FAST2390
3   INTG ,IIP ,ITDS ,IDWD ,IDY ,IDZ ,FAST2391
4   IVE ,ICOT ,IXOT ,IINI ,IA ,IIMO ,ICST ,FAST2392
5   IAIM ,IALE ,IAHL ,IA ,ITNS ,IVEE ,IVAL ,FAST2393
6   ITSP ,ISPE ,IATO ,IRS ,ISGT ,IIDE ,IIDI ,FAST2394
7   ISGS ,INEXT ,IXMP ,IWS ,IES ,INC ,FAST2395
8   IEC ,IND ,IU ,IV ,IW ,ISI ,IST ,FAST2396
9   INNG ,INSC ,ISIP ,INRP ,INPC ,IFXP ,IFXS ,FAST2397
10  IFXT ,IFAC ,IPAA ,FAST2398
11  I = III'
M = JJJ
FST = AAA
ANG = BBB
STV = CCC
IP = NSXT
IT = IP+NMAX(DNLMAX,2)
CALL DEDOUT(I,NMAX,NEMOD,MOMENT,H(IWS),H(IES),H(INTG),H(IEXT),
1H(IFXP1),H(IFXE),H(IFXA),H(IRHD),H(IAE),H(IEI),H(ISSH),H(ISGT),
2H(IFXP),H(IGM),H(IP),H(CT),H(IAIM))
RETURN
END

$IBFTC TALLY M94/2,XR7
CDETECT*FLUX TALLY FOR PROGRAM FASTER
  SUBROUTINE DETDUM(L1,L2,L3,L4,M5,ES,HTG,FXT,FXP,FRE,FXA,RHO,AE,BE,FAST2435

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DUT3 = FXT(J,J)/XN
IF(DUT3.EQ.0.0)GO TO 100
OUT1 = OUT1//J//XN
OUT2 = OUT4//OUT3
DO 90 K=5,6
  OUT(K) = OUT(K-2)/FGW(J)
90 OUT(K+2) = OUT(K+2) + OUT(K-2)
  FST = (FXN*(IPXN-1)*XN*OUT3)/XN
  IF(OUT1.EQ.0.0)OUT1 = SQRT(FST)/OUT3
100 FORMATT(I1X,5HGROUP,I4,3H...,IPE11,4,7E12,4)
AEIJ = (OUT2 - ELF(J+1))/FGW(J)
BEIJ = (ELF(J) - OUT2)/FGW(J)
FLXIJ = CUT3
110 ENDIF
  IF(NFMAX.EQ.0)IGO TO 170
  DO 160 J=1,NFMAX
    MAX = MIN0(J,7,NFMAX)
    MOD = MAX - J + 1
    IF(LABEL(1).GE.0)WRITEM(2,2020)(IOUT,NPF,((TDS(L,K),L=L,3),K=J,MAX))FAST2644
2020 FORMAT(I1X,24I1H*) ,34HNUMBER FLUX RESPONSES FOR DETECTOR,I4,6H AFTER
  LINES = LINES - 1
  DO 120 K=1,MOD
    OUTS(K) = 0.0
    ENM(K) = 0.0
120 EH(XK) = 0.0
  DO 140 K=1,NMAX
    AE(K) = 0.0
    ENM(K) = 0.0
    DO 130 M=J,MAX
      L = L + 1
      FST = FL(XK)*(AE(K)*RSP(K,M) + BE(K)*RSP(K+1,M))
      OUT(L) = FST
      OUTS(L) = OUTS(L) + FST
      ENM(L) = ENM(L) + FST**2
130 EH(XL) = ENH(XL) + FST**2
    IF(LABEL(1).GE.0)WRITEM(2,2010)(K,(OUT(L),L=1,MOD))FAST2661
140 CONTINUE
    IF(LABEL(1).GE.0)WRITEM(2,203)(OUTS(L),L=1,MOD))FAST2663
2030 FORMATT(I1X,6I1H0TALS,6I1H...) ,IPE11,4,7E12,4)
  DO 150 I=1,MOD
    OUT(I) = 0.0
    FST = OUTS(I)
    OUTS(I) = 0.0
    IF(FST.EQ.0.0)IGO TO 150
    OUT(I) = SQRT(ENM(I))/FST
    FAST2665
  CONTINUE
  IF(NFMOD.EQ.0)IGO TO 240
  DO 230 J=1,NFMOD+8
    MAX = MIN0(J,7,NFMOD)
    MOD = MAX - J + 1
    IF(LABEL(1).GE.0)WRITEM(2,2050)(IOUT,NPF,((ITLE(K,L),K=L,2),I001L),I4,6H
    L=1,MOD))FAST2682
    LINES = LINES - 1
    DO 2050 FORMATT(I1X,25I1H*) ,32HNUMBER FLUX MOMENTS FOR DETECTOR,I4,6H AFTER
      LINES = LINES - 1
      DO 190 K=1,NGMAX
        L = 0
        DO 180 M=J,MAX
          L = L + 1
          OUT(L) = FXA(M,K,I)/XN
          IF(L.EQ.1) OUT1 = XTI*OUT1
          IF(LABEL(1).GE.0)WRITEM(2,2010)(K,(OUT(L),L=1,MOD))FAST2693
190 CONTINUE
      DO 200 K=NFMAX,50,1)GO TO 230
      DO 200 L=1,1)NFMAX
      DO 200 L=1,MOD
200 OUT(I) = 0.0
      DO 210 L=1,NGMAX
        M = 0
        DO 210 N=J,MAX
          M = M + 1
          OUT(M) = OUT(I) + FXA(N,L,I)*(AE(L)*RSP(L,K)+BE(L)*RSP(L+1,K))/XN
          IF(L.EQ.1) OUT1 = XTI*OUT1
          IF(LABEL(1).GE.0)WRITEM(2,2060)(TDS(L,K),L=L,3),(OUT(L),L=1,MOD))FAST2705
210 OUT(I) = OUT1
2060 FORMATT(I1X,3A4,IPE11,4,7E12,4)
220 CONTINUE
230 CONTINUE
240 CONTINUE
  DO 250 I=1,NDMD
    DO 250 J=1,NGMAX
      FXA(I,J,I) = 0.0
      IF(NNM.LT.NPRINT)GO TO 900
      IF(NPOINT.EQ.0) GO TO 300
      IF(NPOINT.EQ.1)GO TO 870
      IF(NPOINT.EQ.2)GO TO 870
      IF(LABEL(-1).GE.0)WRITEM(2,3000)
3000 FORMATT(I1X,22I1H*) ,63HNUMBER SEARCH PARAMETERS, (SURFACE,HOST PROFAT2717

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4      .EQLV ,.ICDT ,.IXMT ,.IRSI ,.IALP ,.IVND ,.IGIM ,.FAST2765
5      .IALW ,.IALM ,.IAH ,.IANS ,.IVEE ,.IVAL ,.FAST2766
6      .ISPW ,.ISPE ,.IATD ,.ISPB ,.ISGT ,.TIDE ,.IDI ,.FAST2767
7      .ISGS ,.NEXT ,.IXMP ,.SGNS ,.IWS ,.IES ,.INC ,.FAST2768
8      .IEC ,.IND ,.IU ,.IV ,.IW ,.ISI ,.IST ,.FAST2769
9      .INRG ,.INST ,.ISTP ,.INRC ,.INCP ,.IFXP ,.IFKS ,.FAST2770
1      IF(FX1>FX2)FXA=FX1
1      CALL PSTDMUNIX(L1,L2,I,NNSG,HI(NPC),HI(VEE),HI(WHO),HI(ALPI,
1          HI(NSG);HI(XTR))
1      NNN = NN
DO 10 I=1,3
10  XXX(I)= X(I)
PSTAR = PDT
RETURN
END
!$IBFT PSTOMI M94/2,XR7
CPSTMURANDOM SELECTION OF SOURCE VECTOR FOR PROGRAM FASTER*T.M.JORDAN!$IBFT2781
SUBROUTINE PSTDMUL1(L1,L2,I,RSI,NPC,VEE,VND,ALP,NSG,XTR)
COMMON/DUMPST/,,PDT,,XTR)
DIMENSION RSI(L1),NPC(1,I),VEE(L1,L2,L3),VND(5,1),ALP(5,1),NSG(1),
1           XTR(1,I)
COMMON/LIMITS/NSTORE,NERROR,NMAXM ,NMAXX ,NMAXM ,NSTMAX,
1           NMAXV ,NMAXX ,NMAXX ,NMAXEM ,NEMOD ,NSCECT ,NUINT,
2           NUNITX ,NMAXX ,NMAXX ,NORDER ,NCWD ,NLELAS ,NTRIM,
3           NMAXX ,NMAXX ,NINMAX ,NUNDM ,NCRAN ,NLMAXM ,NTRAM,
4           NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,
5           NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,NMAXX ,
6           NPRINT ,NUNITS ,NUMBER ,KALTOP
COMMON/POINTX/NTOTAL ,III ,MON
COMMON/LLSREG/LLREG
DIMENSION VB(15)
OLD METHOD, SELECT VOLUME
IF(NRMAX.GT.1)GO TO 10
PDT = 1.0
N = 1
GO TO 40
10 R = RANH(1NMB)
FST = 0.0
DO 20 N=1,NMAXX
FST = FST + RSI(N)
IF(FST-GE.R1)GO TO 30
20 CONTINUE
30 PDT = 1.0/RSI(N)
40 DO 50 I=1,3
J = NPC(I,N)
50  PRINT *,PDT*(SAMPLE(VEE(1,I,J),VEE(J,I,N),VND(1,N),ALP(1,I,N),VB(1,I)
     1)LLREG ,N)

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DO 150 J=1,3
150 C11 = C11 + D(J)*RCT(J,1)
CALL PATHNN,STM,X,_C,NST,ST,MRG,NSC
      STT = 0.0
      DO 160 I=1,NST
         J = NGC(I)
         IF((ISV(I).GT.0))GO TO 170
160 STT = STT + ST(I)
165 CONTINUE
      NST = 0
      GO TO 300
170 NTT = 0
      PTOT = 0.0
      XMPT = 0.0
      STX = STT
      DD 220 I=NST,NST
         J = I
         K = ISV(I)
         L = MT(I,J)
         SGR = ATC*RH0(I,J)+$SH(JZERC)
         IF(L.GT.0)SGR = SGR + ATC*SGL(JZERO,L)
         XMPP = SGR*ST(I)
         IF(K.GT.0)GG TO 180
         XMII = 0.0
         GO TO 210
180 GALL ARROWIX,STX + 0.01*ST(I),C,XS1
      LLREG = K
      CALL SZERC(0.1,0,XS,C)
      PDRS = 0.0
      CALL ARROW (X,STX + 0.5*ST(I),C,XS)
      CALL SZERO(0.1,0,XS,C)
      PDM = PSORS
      TRAI1 = PDRS*SU(V1)*EXP(-XMP)
      ARGARG < 0.0
      IF((V1.GT.0.0).AND.(GT.0.0)) ARGARG = ALDG(PCB/PDH)
      IF(ST(I).LT.C-.01) SGR = SGR + 2.0*ARGARS/ST(I)
      TRB1 = SGR
      IF(SGR.NE.0.0)GD TO 190
      XMII = ST(I)*TRA11
      GO TO 200
190 XMII = XMII - TRA11*(1.0 - EXP(-SGR*ST(I)))/SGR
200 PTOT = PTOT + XMII + XMII
210 XMPT = XMPT + XMPP
220 STX = STX + ST(I)
      PDS = PTOT*RAMC(NNP)
      DO 230 I=NST,NST
         IF(XMII(I).GT.PDS)GO TO 240

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8   IEC ,IND ,IU ,EV ,IW ,ISI ,IST : FAST299
9   INRG ,INSC ,ISP ,INRP ,INC ,IFXP ,IFXS : FAST299
1   IFXT ,IFXG ,IFXA : FAST299
QSTAR = QSTODUM(NN,X,C,MAXX,5,1,H((ISV),H((INPC),H((IVEE),H((IVMD),
1   H((ALP),H((XTR),H((NSG)))) : FAST299
RETURN
END

$IRFTC QSTOHL M94/2,XR7
QSTODUM#RANDOM SELECTION OF INITIAL DIRECTION VECTOR FOR FASTER*. JORDANFAST300
FUNCTION QSTODUM(NN,X,C,L1,L2,L3,I5V,INPC,VEE,VND,ALP+XTR,NSG)
DIMENSION X(3),C(1)
DIMENSION V(5,1),Z(1,5)
1   XTR(1,1),NSG(1,1),Z(1,1),ROT(3,3)
COMMON/OTHERS/RADIUS,XCT(3),DELTA ,BDC(3),ATA ,ATB ,ATC
1   ATD ,ATD ,BT ,AS ,BS ,JMIN ,JMAX :
COMMON/LLSREG/LLREQ
N = ISV(NNN)
N = LLREQ
PDA = 1.0
DO 10 I=4,5
J = NPC(1,I)
10 PDA = PDA*ROT(I-1,I)
SAMPLE=(VEE(I,1,N)*VEE(J,1,N)*VMDI(1,N),ALP(I,N)*VB(I))
D13 = VB(5)
SPH = SORT(1.0 - C(3)*#2)
D11 = SPH*COS(VB(4))
D12 = SPH*SIN(VB(4))
IF(NSGN(1).GT.-0.05) GO TO 30
DO 20 I=1,3
20 CTH = D11/I
GO TO 90
30 DO 40 I=1,3
40 Z(1) = X(1) - XTR(I,1)
MAX = NSGN(1) + 1
ARE = 0.0
DO 50 I=1,MAX
50 ARE = ARE + Z(I)**2
ARE = SORT(ARE)
IF(MAX.EQ.3)GO TO 60
CPH = 1.0
SPH = 0.0
GO TO 70
60 CPH = Z(3)/ARE
SPH = SORT(1.0 - CPH**2)
ARE = SPH*ARE
70 CTH = Z(1)/ARE

```

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DIMENSION ST(1),HRC(1),RHO(1),SSH(1),MTL(1),SGT(L1,L2),TRA(1),      FAST3089
     TRBL(1),XW(1)          FAST3090
COMMON/DUMUST/SIZ ,STS ,PSI ,SIC ,NST ,PDT          FAST3091
COMMON/OTHERS/RADIUS,XCT(3),DELTA ,BDC(3),ATA ,ATB ,ATC ,      FAST3092
1   KDN ,ATD ,BT ,AS ,IBS ,JMN ,JMAX ,      FAST3093
2   KDN ,KDN ,KMAX ,JBAR ,NZER ,NMAX ,NTPAX ,      FAST3094
COMMON/LIMITS/STDB ,NMAX ,NMMAX ,NMED ,NMMOD ,NMAX ,NTPAX ,      FAST3095
1   NMEX ,NMMAX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,      FAST3096
2   NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,      FAST3097
3   NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,      FAST3098
4   NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,      FAST3099
5   NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,NMEX ,      FAST3100
NVMAX ,NVMD ,NPMEN ,NMDD ,NPDMAX ,NPMD ,NSDMAX ,      FAST3101
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3102
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3103
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3104
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3105
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3106
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3107
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3108
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3109
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3110
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3111
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3112
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3113
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3114
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3115
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3116
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3117
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3118
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3119
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3120
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3121
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3122
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3123
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3124
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3125
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3126
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3127
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3128
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3129
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3130
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3131
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3132
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3133
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3134
NVMAX ,NVMD ,NPMGINT ,NPMDEL ,NPMDEL ,NPMDEL ,      FAST3135

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DDM = (0.5*DELTA/STDB)**2           FAST3136
DDM = DDM+1.0 + DOMI             FAST3137
0 TRAB = DDM*EXP(-XMTI) - XMTI/DOZ
XMT = 0.0                         FAST3138
1 IF((SGR*STN).EQ.0.01) GO TO 10
TRAIN = SGR*TRAAC                FAST3139
IF(TRAIN).EQ.0.01 GO TO 10
ARGU = 1.0                         FAST3140
IF(ARGU.GT.0.01 GO TO 9
TRBN = (XMAA - XBBB)/STN          FAST3141
XMAA = XBBB                         FAST3142
XBBB = XMTI + XMT                 FAST3143
TRAAC = TRAB                         FAST3144
DDM = DDM+1.0                      FAST3145
IF((STDB.LE.0.01) GO TO 8
DDM = 1.0                          FAST3146
GO TO 8
9 TRBN = ALOG(ARGU)/STN          FAST3147
XBBB = SGR*(TRAB - TRAA)/TRBN    FAST3148
11 TRD(N) = TRBN                  FAST3149
10 XH(N) = XHN                     FAST3150
20 FMP = FMP + XHN                FAST3151
21 FMB = FMP*ANNU(NMB)            FAST3152
30 FMB = FMB - XM(N)              FAST3153
31 STI = 0.0                         FAST3154
32 DO = 40 .NST                     FAST3155
33 IF(XK(N).GE.FMB) GO TO 50
34 STT = STT + STN                  FAST3156
35 FMB = FMB - XM(N)              FAST3157
36 NST = N                         FAST3158
37 ARGU = TRA(N)
38 TRAN = TRB(N)
39 VARG = FMB*TRBN + TRAN         FAST3159
40 STP = STN(Y)                   FAST3160
41 XYZ = VARG*TRAN                FAST3161
42 IF(XYZ.GT.0.01 STP = ALOG(XYZ)/TRBN
43 STI = STT + STP                  FAST3162
44 SGR = 1.0                         FAST3163
45 IN=N
46 GO TO 170
47 CONTINUE
48 PSIS=1.0
49 NGT = 1
50 TM11 = STZ
51 STT = 0.0
52 FBT = 0.0
53 FMB = 0.0
54 DO IZO = 1,NST
55 I = NRG(I)
56 J = MTL(I)
57 SGR = RHO(I)*SSH(JBAR)

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```

1 IF(I>GT-.01SGR = SCR + SGT(JBAR,J)
1 IF(SGR-GT-.01GO TO 100
1 IF(SGR-EQ.0-0IGO TO 110
1 XM(N) = STIN)*SGR*ANG
1 GO TO 120
100 STT = STT + STIN
100 XM(N) = STT*SGRTISTS + STT*(I2.0*PSI + STTPSIS))
1 TR0(N) = 0
1 IF(SGR-EQ.0-0IGO TO 110
1 TRB(N) = BT*SGR*(TRAINH+1) - 'TRA(NH)
1 FBT = FBT + TRB(N)
1 XM(N) = TRB(N) + AT*SGR*SY(N)
1 GO TO 120
110 XM(N) = 0.0
120 FMP = FMP + XM(N)
1 IF(FMP>GT.0-0IGO TO 125
95 NST = 0
GO TO 195
125 FMP = FMP + FMP
ARG = 1.0 - AAA
VARG = 1.0 - ARG*RANNG(NMB)
FMA = -ALOC(VARG)
FMB = FMA
FRC = 0.0
STC = 0
DD 140 I=1,NST
IF(FMB<LE.XM(I))GO TO 150
FMB = FMB - XM(I)
STC = STC + ST(I)
140 FRC = FRC - TRB(I)
150 NM = NHG(I)
MM = NTL(NNN)
NST=1
SGP = RHO(INN)*SSH(JBAR)
IF(MH>GT-.01SGP = SGP + SGT(JBAR,MH)
1 IF(NGT.0-0IGO TO 160
1 SGP = NGC*CP
DOX = IEFMAAA/ARG + FMA = 1.0)/ANG
STP = FBH/SGP
STC = STC + STP
GO TO 170
160 DOX = IEFMAAA/ARG + SCR + SGP
SGP = BT*SCP
CQ = (FMB * SGA*STC + SGB*TRA(I))/SGB
AQ = SGA/SGB
FST = AQ*#2 - PSIS
BS = (AQ*CQ + PSI)/FST

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HH = -(QC**2 - (QC**2 - STS)/FST)
HH=HH+LT(0,0) NH=0.0
HH=SGRT(HH)
ST0 = BQ - HH
IF(ST0.LT.STC)ST0 = BQ + HH
STP = ST0 - STC
STC = A0 - STC
ARE = GO - A0*STC
FNC = FNC + SGB*(ARE - IRA(T))
PSI = (PSI + PS1*STC)/ARE
ANG = AT * BT*PSI
SGP = ANG*SGP
170 (SGP+1) = STP
PDT = ARG*STC**2/(SGP*VARG*12.566371)
195 RETURN
END

$IBFTC VSTARL M94/2,XR7
CVSTAR *RANDOM SELECTION OF DIRECTION VECTOR FOR PROGRAM FASTER*T.JORDAN
FUNCTION VSTAR(CC,CP,C,BS,BS1)
COMMON//LIMITS,UNIT0,NUMBER,MAX,MAXX,NMAX ,NBMX ,NSTHMX
1      NMAXX ,NMMAX ,NMAXH ,NHED0 ,NSECNT,NUINT0,
2      NUNITX ,NMMAX ,NMAXP ,NCRDR,NDOMW ,INELAS,INTRANS
3      NMMAX ,NMAXP ,NFMAP ,NVMD ,NCMAX ,NLMAX ,INTMAX ,
4      NOHAX ,NCMCN ,HCMENT ,NDOMD ,NPFMD ,NPDMOD ,NSDOMD ,
5      NVDOHAX ,NWDPCD,NPCINT ,NDELPQ ,NDEQLU ,NDELYV ,
6      NPRTIN ,NUNITS ,NUMBER ,KALIDE
IF(MODELV = 1) P5,10,15
5 CONTINUE
VSTAR = VSTDUM(CC,CP,C,BS,BS1)
GO TO 20
10 VSTAR = VSTDUM(CC,CP,C,BS,BS1)
GO TO 20
15 VSTAR = VSTDUM(CC,CP,C,BS,BS1)
GO TO 20
20 RETURN
END

$IRFTC VSTDIM1 M94/2,XR7
FUNCTION VSTDIM1(CC,CP,C,BS,BS1)
DATA PI/3.14159265358979311600D0/
DIMENSION CP(3),CC(3),DC(3,1),ROT(3,31,D13)
EQUIVALENCE (D1(1),D1),(D2(1),D2),(D3(1),D3)
COLI RADTAC(CP,ROT)
DO 10 I=1,3
  D1(I) = 0.0
  DO 10 J=1,3
    D1(I,J) = 0.0
10 D1(I,J) = D1(I,J) * ROT(I,J)*CC(J)
    THZ = 0.0
    IF((ABS(D1(1)) .ABSD02)) .GT. 0.0 THZ = ATAN2(D2(1))
    FST3230
    FST3231
    FST3232
    FST3233
    FST3234
    FST3235
    FST3236
    FST3237
    FST3238
    FST3239
    FST3240
    FST3241
    FST3242
    FST3243
    FST3244
    FST3245
    FST3246
    FST3247
    FST3248
    FST3249
    FST3250
    FST3251
    FST3252
    FST3253
    FST3254
    FST3255
    FST3256
    FST3257
    FST3258
    FST3259
    FST3260
    FST3261
    FST3262
    FST3263
    FST3264
    FST3265
    FST3266
    FST3267
    FST3268
    FST3269
    FST3270
    FST3271
    FST3272
    FST3273
    FST3274
    FST3275
    FST3276
    FST3277
    FST3278
    FST3279

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```

SINZ = SQRT(1.0 - D3**2)
PDA = SAMPLE(-1.0,1.0,-1.0,AS,03)
SINE = SQRT(1.0 - D3**2)
PDE = PDA*SAMPLE(-PI,PI,0.,-BS*SINZ*SINE,THE1)
THE1 = PI
D1 = SINE*COS(THE1)
D2 = SINE*SIN(THE1)
DO 20 I=1,3
  C(I) = 0.0
  DO 20 J=1,3
    C(J,I) = C(I,J) + D1*J*ROT(J,I)
  VSTOUM = PDA
  RETURN
  END*
$IBFTC SAMPLR  M94/2,XR7
$SAMPLE RANDOM SELECTION OF VARIABLE
  FUNCTION SAMPLE(X0,X1,VMD,ALP,VEE)
  IF(X0.GT.VMD) GO TO 10
  C DELTA FUNCTION (OR ERROR)
  PDF = 1.0
  VEE = VMD
  GO TO 30
  10 IF(ALP.NE.0.0) GO TO 20
  C UNIFORM DISTRIBUTION
  PDF*= VMD - VMD
  VEE= VMD + R*PDF
  GO TO 30
  C EXPONENTIAL DISTRIBUTION
  20 PCT = EXP(ALP*(VMD - VMD)) - 1.0
  PGT = EXP(ALP*(VMD - VMD)) - 1.0 + PLT
  PMD = PLT/PGT
  T = -1.0
  IF(R.GT.PMD) T= 1.0
  R = 1.0 - T*(R*PCT - PLT)
  VEE = VMD + T*ALOG(T)/ALP
  PDF = PGT/(R*ALP)
  30 SAMPLE = PDF
  RETURN
  END
$IBFTC RANDUM  M94/2,XR7
$RANDOM NUMBER GENERATOR
  COMMON/DGCD3W/IBMDC
  DATA N,NNN/0,11111/
  IF(IBMDC.GT.0) GO TO 10
  R = RANF(NNN)
  GO TO 20

```

```

  10 CONTINUE
  IF(N.EQ.0) R = RANF(NNN)
  F = RANF(10)
  20 CONTINUE
  N = N + 1
  RANH = R *
  NMB = N
  RETURN
  END
$IBMAP RANI
*   RANDOM NUMBER GENERATOR
RANF  ENTRY  RANF
      LDO   L16
      MPY   L17
      LLS   4
      ALS   4
      LRS   4
      STO   L16
      ADD   L16
      STO   L16
      ARS   4
      ORA   L20
      FAD   L20
      TRA   1:4
L16   OCT   002312421637
L17   OCT   000000001737
L20   OCT   200000000000
      END
$DATA
      *
      :

```