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# LIFETIME AND TRANSITION PROBABILITIES OF $np^4$ - (n + 1) p STATES OF Ne II, A II AND Kr II

S. H. Koozekanani and G. L. Trusty

The Ohio State University

## **ElectroScience Laboratory**

(formerly Antenna Laboratory) Department of Electrical Engineering Columbus, Ohio 43212

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Investigation of

Receiver Techniques and Detectors for Use at Millimeter and Submillimeter Wave Lengths

Subject of Report

Lifetime and Transition Probabilities of  $np^4$  - (n + 1) p States of Ne II, A II and Kr II

Submitted by

S. H. Koozekanani G. L. Trusty ElectroScience Laboratory Department of Electrical Engineering

Date

17 September 1968

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## ABSTRACT

Life time as well as transition probabilities of the first p-excited states of neon,<sup>+</sup> argon<sup>+</sup>, and krypton<sup>+</sup>have been calculated.

## LIFETIME AND TRANSITION PROBABILITIES OF $np^4$ - (n + 1) p STATES OF NeII, AII AND KrII

Lifetime and transition probabilities of some of the first excited "p" states of neon II, argon II and krypton  $\text{II}^{1,2}$ , <sup>3</sup> for the configuration of np<sup>4</sup> - (n + 1)p, have been calculated. This information is needed to determine the possible efficiencies of current and proposed laser systems and to aid in the identification of the energy levels involved.

The exact mixed wave functions  $\Psi_i$  ( $E_i$ ) corresponding to the experimentally found energy levels  $E_i$  can be expressed in terms of purely LS coupled wave functions  $\phi_i$ , with  $\phi_i$  corresponding to the LS coupled states of the form  $|p^4 LS; \ell_{ex}s; L_iS_iJ_iM_i>$ , i.e.,

(1) 
$$\Psi_i(\mathbf{E}_i) = \sum_j a_{ij} \phi_j.$$

where  $\overline{LS}$  are the orbital and spin angular momentum of the p<sup>4</sup> core electrons while  $L_j S_j$  are the total orbital and spin angular moments.  $\underline{J} = \underline{L} + \underline{S}$  and M is the projection of J on the z-axis.  $\ell_{ex}$  is the excited p electron with its spin s. The functions  $\Psi_i$  (E<sub>i</sub>) were found by treating the radial integrals as adjustable parameters so as to obtain the best fit with the experimentally found energy levels. This was done by choosing arbitrarily close test values for the parameters  $\begin{bmatrix} 1 \\ D \end{bmatrix}$ ,  $\begin{bmatrix} 3 \\ P \end{bmatrix}$  core energy states,  $F_2$ ,  $G_0$ ,  $G_2$  direct and exchange integrals and  $\zeta_{np}$  and  $\zeta_{\ell}$  ex (n + 1) p excited electron.

With the initially chosen test parameters  $\begin{bmatrix} 1 \\ D \end{bmatrix}_0 \begin{bmatrix} 3 \\ P \end{bmatrix}_0$ ... etc. and the Hamiltonian

(2) 
$$H = -\sum \frac{h^2}{2m} \nabla^2 - \sum_{i} \frac{e^2}{r_i} + \sum_{i>j} \frac{e^2}{r_{ij}} + \sum_{i} \xi(r) \ell_i, S_i,$$

the matrix  $[H_0]$  was constructed having elements  $\overline{H_0}_{ij}$  where

(3) 
$$[H_0]_{ij} = \langle \phi_i | H_0 | \phi_j \rangle_0$$

In Eq. (3), the subscript naught designated that in the calculation of the elements  $[H_{c]_{11}}$  the first set of test parameters are used.

A matrix  $M_0$  is found such that,

(4)  $[M_0]^{-1} [H_0] [M_0] = [D_0]$ 

with  $[D_0]$  being a diagonal matrix. If the choice of the parameters were appropriate  $[D_0]$  will differ little from the experimentally found energy levels. If not, a new matrix equation is constructed having the form,

(5)  $\left[M_0^{-1}\right]\left[H\right]\left[M_0\right] = \left[E\right]$ 

where [E] is again a diagonal matrix having for its diagonal elements the experimental energy levels  $E_i$ ,  $[M_0]$  is the matrix which was used to diagonalize  $[H_0]$ , and [H] is the Hamiltonian matrix with its elements expressed in terms of the radial parameters<sup>4</sup>, i.e.,

(6) 
$$[H]_{ij} = \langle \phi_i | H | \phi_j \rangle = a_{ij} [^1D] + b_{ij} [^3P] + C_{ij} F_2 + d_{ij} G_0 + \dots$$

Multiplying out  $[M_q^{-1} [H] [M]$  and equating it to [E] one obtains, taking the diagonal terms only, a linear set of nineteen simultaneous equations in terms of the seven parameters  $[^{1}D]$ ,  $[^{3}P]$ ,  $F_{2}$ ,  $G_{0}$ ,  $G_{2}$ ,  $\zeta_{np}$  and  $\zeta_{lex}$ . Solving the equations and least-square-fitting the results, we obtain a new set of parameters  $[^{1}D]_{1} [^{3}P]_{1} \dots$  etc. These parameters are now used to obtain the matrix  $[H_{1}]$ , with matrix elements  $[H_{1}]_{ij}$  in which the set of the newly found values of the radial parameters have been substituted, i.e.,

$$[H_1]_{ij} = a_{ij} [^1D]_1 + b_{ij} [^3\underline{P}]_1 + C_{ij} (F_2)_1 + d_{ij} (G_0)_1 + . .$$

Once again we find a matrix  $[M_1]$  which diagonalizes  $[H_1]$  giving the diagonal matrix  $[D_1]$  i.e.,

(7) 
$$\begin{bmatrix} M_1 \end{bmatrix}^{-1} \begin{bmatrix} H_1 \end{bmatrix} \begin{bmatrix} M_1 \end{bmatrix} = \begin{bmatrix} D_1 \end{bmatrix}$$

If again the elements  $[D_1]$  would differ much from the experimentally found energy levels  $E_1$  the process is repeated by solving the simultaneous equation obtained from the matrix equation,

(8)  $\begin{bmatrix} M_1 \end{bmatrix}^{-1} \begin{bmatrix} H \end{bmatrix} \begin{bmatrix} M_1 \end{bmatrix} = \begin{bmatrix} E \end{bmatrix}$ 

where again the elements of the matrix H are given by Eq. 6. A new set

of parameters  $\begin{bmatrix} {}^{1}D \end{bmatrix}_{2}$ ,  $\begin{bmatrix} {}^{3}P \end{bmatrix}_{2}$ , etc., are found. The process is repeated until the parameters  $\begin{bmatrix} {}^{1}D \end{bmatrix}_{n}$ ,  $\begin{bmatrix} {}^{3}P \end{bmatrix}_{n}$ , etc. converge and the matrix equation  $\begin{bmatrix} M_{n} \end{bmatrix}^{-1} \begin{bmatrix} H_{n} \end{bmatrix} \begin{bmatrix} M_{n} \end{bmatrix} = \begin{bmatrix} D_{n} \end{bmatrix} \simeq \begin{bmatrix} E \end{bmatrix}$ .

If the initial set of test parameters are close to the actual values, the process would terminate after a few cycles, and the parameters would converge to fixed values. If not, the convergence is not easily attained. In the above calculations the states arising from the <sup>1</sup>S core were not taken into consideration. However, their contribution to the states <sup>2</sup>  $P_{\frac{1}{2}}$ ,  $^{2}P_{\frac{3}{2}}$ ,  $^{2}P_{\frac{1}{2}}$  and  $^{2}P_{\frac{3}{2}}$  can be calculated by perturbation theory. The

shift in energy would be  $\triangle E = -\frac{H^2 ij}{Ei-Ej}$ . Table one gives the values of the parameters  ${}^{1}D$ ,  ${}^{3}P$ ,  $F_2$ ,  $G_0$ , ... etc. using Hartree Fock<sup>6</sup> wavefunctions and also the values which were obtained by the above fitting techniques. The two sets of values are more or less very close. Table II gives the theoretical and experimentally found energy levels and the mixing coefficients  $a_{ij}$ . Table III gives the transition probabilities for the states arising from the np<sup>4</sup> - (n + 1) p configuration of neon, argon, and krypton to the lower states of np<sup>4</sup> - (n + 1)s configuration. Finally Table IV gives the lifetimes of some of the higher states of interest. TABLE I

Ne

Kr

 $\mathbf{Ar}$ 

Fitting Cm <sup>-1</sup>	397	1802	171	2929	570	140669	152937	2044
Calculate Cm <sup>-1</sup>	415	2176	111	2593	251		1	2488
Fitting Cm <sup>-1</sup>	365	1752	86	1075	187	159918	173491	2262
Calculate Cm <sup>-1</sup>	424	2614	122.5	1198	107			2827
Fitting Cm <sup>-1</sup>	433	2452	167	062	-33	252853	278205	4225
Calculated	422.	4166.	160.	1104	63		1	4757
	F, (p	Go(P <sub>core</sub> P <sub>ex</sub> )	$G_2 (p_{core} - p_{ex})$	5 5 5 5 7 5 7 5 7 5 7 5 7 7 7 7 7	rcore 2 - 2			F2 (Pcore Pcore)

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TABLE II -a

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Experimental and theoretical energy levels of Ne II taken from Moore<sup>7</sup> , tables of atomic energy levels and this calculation. The mixing coefficients a<sub>ij</sub> are those given in Eq. (1).

2 F.	0.21956	-0.99976	0.00042	0.01307	0.01966	-0.99971	-0.00449												
<sup>2</sup> D'		<u></u>	0.01720	0.00271	-0.03385	-0.00511	0. 99926	0.00773	-0.00448	0.02494	-0. 02373	0.00375	-0. 10978	-0.99331		······································			
2 5							<u> </u>	0, 020 33	-0, 00908	-0.04265	-0.23011	0.00933	0.96654	-0. 10217	-0.01249	0.01401	-0.22804	-0. 10641	-0.96764
S S						<u></u>	<u> </u>	•			- <u></u>	<u> </u>			-0.04454	-0.00301	-0. 37847	0.92444	-0.01308
۲ ۲								-0.04273	0.04061	0.20859	0.94040	-0.11599	0.23087	-0.04370	-0.02210	-0.04947	0.89586	0.36197	-0. 25193
Q 2			0.03104	-0.13543	-0,98950	-0.02106	-0.03379	0,00625	0,09658	0.97118	-0,21197	0.04020	-0.00409	0.02967					
s.					<i></i>			0.06423	0.01201	0.01349	-0,11802	-0.99058	-0.01897	0.00196		- <u>-</u>			
с. *			0.99415	-0, 09696	0.04497	0.00010	-0.01525	0.99253	0.09102	-0.00654	0.05218	0.05933	-0.00729	0.00697	0.99729	-0.05422	0.00333	0.04932	0.00587
¢ D	0.99976	0.21956	0.10198	0.98594	-0.13166	0.01038	-0, 00883	0.09174	-0.99010	0.10309	0.02353	-0.00743	-0.00025	0.00721	0.05282	0.99720	0.04669	0.02493	0.00137
Etheor Cm <b>-1</b>	248652.5	274228.2	246314.4	249217.8	251158.4	274229.0	277407.1	246617.2	249553.3	251713.3	253922.7	253072.9	276396.3	277411.8	246861.7	249749.7	254197.1	252885.9	276598.9
E <sub>exp</sub> Cm <sup>-1</sup>	249110.8	274411.3	246194.8	249448.0	251013.3	274366.9	277346.1	246417.4	249697.7	251524.7	254167.0	252956.0	276278.6	277327.6	246599.9	249841.8	254294.0	252800.8	276514. 1
Ne II States	*D7	2 F.	4 1 1 2	4 ₽ ₽	2 D_5	2 1 2 2 2	2 D'5	<sup>4</sup> م 4	, <sup>4</sup>	2 D <sub>1</sub>	с Ч Д	ی م 4	2 P1 2 P1	<sup>2</sup> D <sup>1</sup>	, <sup>1</sup>	, d 1	ъ 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	s S T	2 P.

TABLE II-b

Experimental and theoretical energy levels of Ar II taken from Minnhagen<sup>8</sup> and from this calculation. The mixing coefficients of a<sub>ij</sub> are those given in Eq. (1).

2 Fr	0. 05680	-0. 99839	0. 00299	0.04919	0. 03689	0. 99759	-0. 03204		,									•	
2 Z			0. 03949	-0, 00525	-0.04014	-0. 03043	-0.99794	0. 02102	0.01261	0.00763	-0. 05244	0. 00052	0. 05873	-0. 99657					
2 Dr								0. 05959	0. 04422	-0.12847	-0.27314	0. 06355	-0.94747	-0.04059	0. 03795	-0, 04050	0. 31788	0.14933	0.93465
2 S															-0. 05442	0.01678	-0.49034	0.86912	0. 03085
с. 2								-0.11366	-0.17205	0.44766	0.81016	-0. 03572	-0.30918	-0.06202	-0. 06195	0. 09918	-0. 80018	-0.46980	0. 35401
D 7			0. 02312	-0.47202	-0.87867	0. 05688	0. 03701	-0, 1561	-0.31306	0.80472	-0. 50253	0. 01899	0. 0218	0. 02951				,	
*s								0.11854	-0.04193	0. 00421	0. 06872	0.98834	0. 05133	0. 00192					
ሲ •		•	0.98216	-0.15036	0.10820	0,00160	0. 03525	0, 96996	-0.18070	0.00919	0.09361	-0.13135	0.01493	0.01413	0.99105	0.10630	-0.07824	0.01627	-0.01162
а,	0.99839	0. 05680	0. 18291	0. 86726	-0.46181	-0. 02553	0. 02203	0.16724	0.91426	0. 36791	0.01598	0.01607	-0. 00095	0.01702	0. 09789	-0.98840	-0.11005	-0. 03676	-0, 00350
Etheor Cm - or	157176.6	170582.7	155128.2	157556.9	158710.3	170364.6	173432.7	155383.5	158007.6	159244.3	160354.7	161012.7	172213.8	173349.9	155781.9	158331.0	159869.7	161306.5	172726.4
Eexp Cm <sup>1</sup>	157233.93	170530.31	155043.07	157673. 32	158730.21	170400.94	173393.38	155351.03	158167.71	159393.31	160239.35	161048.64	172213.80	173347.83	155708.02	158428.03	159706.46	161089.31	172816.21
Ar II	<sup>4</sup> D <sub>T</sub>	2 E'7	4 P5	ب ف	<sup>2</sup> D5	2 F1 5	<sup>2</sup> D'5	P 4	10 <sup>4</sup> 1	2 D3	5 H	45, 2	2 P1 2 P1	<sup>2</sup> D <sup>1</sup> 2	ч Ч Ч	•D,	2 10 11	2 S <sub>1</sub>	2 P1 2

b.a

TABLE II-c

energy levels and this calculation. The mixing coefficients a<sub>ij</sub> are those given in Eq. (1). Experimental and theoretical energy levels of Kr II taken from Moore<sup>7</sup> , tables of atomic

2 F1	0, 15329	-0, 98818	0.02069	0.15873	0.08744	0.98066	-0,07092												
2 D'		-	0.10898	-0.04663	-0.08975	-0.05814	-0.98717	0.07257	-0.04741	0.14271	-0.04111	106:0.0-	0.15208	-0.97251				•	
ц г								0.15819	0.05687	0.25429	-0.21368	0.10487	-0.91809	-0. 08685	0.15203	0.17577	-0. 32072	0.24507	0.88491
S S		-									-				-0.19012	-0.40360	0.46584	0.76087	0.07095
Д, 7								-0.25507	-0. 39352	-0.43082	0.67428	0.06388	-0.32003	-0, 18255	-0.20833	-0.28601	0.56136	-0. 59032	0.45955
۵ ۲			-0.00191	-0.64380	-0.73946	0.17647	0.08704	-0.06918	-0.42128	0.80116	0.32338	-0.23728	0.07215	00660.0					
s S								0.24557	-0.23513	0.13765	0.03686	0.91550	0.16055	0.01389					<u> </u>
ф.			0.92956	-0.26376	0.24015	0.00875	0.09275	0.86841	-0. 33387	-0.21173	0.04557	-0.29324	0.02327	0.03093	0.92476	0,01265	0.37937	0,00795	-0.02610
<b>₽</b>	0.98818	0.15329	0.35157	0.69899	-0.61630	-0,06087	0.06542	0.29227	0.70387	0.15404	0.62449	0.05458	-0,00713	0.04903	0.20533	-0.85103	-0.47019	-0.11170	-0.00571
Etheor Cm <sup>-1</sup>	136046.4	149720.5	133855, 3	136913.7	139794.1	149096.5	152377.9	134032.7	138171.7	142424. 3	140117.4	141508.3	150519.5	152172.6	135638.4	139549.3	139042.2	142592.5	152094.4
Eexp Cm-1	135783.18	149704.55	133925.65	136071.00	140118.99	149173.42	152316, 26	134288.44	138381, 35	141995.68	140137.15	141722,72	150203, 48	152191.86	135783.03	140163.25	139103, 36	142363.55	152240.97
	<sup>4</sup> D <sub>7</sub>	2 F17	4- 4- 4-	<sup>4</sup> D <sub>5</sub>	2 D5	2 F) 5	<sup>2</sup> D <sup>1</sup> 5	4 C,42	⁺D <sub>2</sub>	<sup>2</sup> D <sub>2</sub>	2 P2	53 2	2 년 2 년	<sup>2</sup> D <sup>1</sup>	4 P1	<sup>4</sup> D <sub>1</sub>	<sup>2</sup> P <sub>1</sub>	<b>4</b> 21 21	2 Pt <u>1</u> 2

### TABLE III

Transition probabilities between 3p-3s levels of Ne II, 4p-4s levels of Ar II and 5p-5s levels of Kr II, each with a  $p^4 \begin{bmatrix} 3 \\ 2 \end{bmatrix}$  core. A in sec<sup>-1</sup>.

	Neor	1 II	Argon	n II	Krypton II				
	Pure L's	Mixed	Pure L's	Mixed	Pure L's	Mixed			
$D_{\frac{1}{2}} - P_{\frac{5}{2}}$	197.81 x 10 <sup>6</sup>	197.71 x 10°	146, 53 x 10 <sup>6</sup>	146.06 x 10 <sup>6</sup>	168.15 x 10 <sup>6</sup>	164 x 10 <sup>5</sup>			
${}^{2}D_{5} - {}^{4}P_{5}$	0	0.28	0	4.67	0	5.28			
- * P3	0.	4.12	0	31.65	0	91.27			
- <sup>2</sup> P <sub>3</sub>	143.30	140.31	103.65	80.03	140.97	77.08			
${}^{4}D_{5} - {}^{4}P_{5}$	61.37	43.08	46.53	18,91	52.37	4.59			
4 P. 2-	135.98	149.75	97.25	90.69	89, 85	68.27			
<sup>2</sup> P <sub>3</sub> 2	0	2.20	0	19.70	0	31, 39			
${}^{4}P_{\frac{5}{2}} - {}^{4}P_{\frac{5}{2}}$	101.86	114.65	79.28	96.24	91.38	122.90			
4 P3	41.20	28.96	30.06	14.84	27.86	4.29			
$4S_{\frac{1}{2}}$ $4P_{\frac{5}{2}}$	142.05	124.62	116.11	87.65	167.69	71.85			
*P3	90.42	93.83	70.32	77.64	87.52	93.20			
${}^{4}P_{\frac{1}{2}}$	44.00	52.87	33.11	49.39	32, 51	69.82			
<sup>2</sup> P <sub>3</sub>	0	1.41	0	0.088	0	0.26			
² P <u>1</u>	0	1.17 .	0	0.13	0	7.28			
$^{2}P_{3} + P_{5}$	0	0.59	0	2.80	0	10.91			
<sup>4</sup> P <sub>3</sub> 2	0	3.29	0	0.0073	0	44.51			
${}^{4}P_{\frac{1}{2}}$	0	1.31	0	0.088	0	24.04			
<sup>2</sup> P <sub>3</sub>	166.49	119.05	106.91	36.64	117,77	78.98			
${}^{2} P_{\frac{1}{2}}$	31.30	62.62	18.56	69.40	16.23	0.039			
$^{2}D_{3}$ $^{4}P_{5}$	0	0.20	0	1.60	o	0.036			
<sup>4</sup> P <sub>3</sub> 2	0	1.00	0	12,59	0	0.53			
<sup>4</sup> P <u>1</u>	0 5	1.35	0	8.83	0	17.24			
<sup>2</sup> P <sub>3</sub>	25.27	52.23	19.01	61.98	30.15	0.74			
<sup>2</sup> P <sub>1</sub> / <u>2</u>	118.08	91.00	82.01	29.97	107.19	105.87			
, ,		1							

	1	Neon II	con <sup>i</sup> t.	Argon II	con <sup>i</sup> t.	Krypton	I con <sup>†</sup> t.
TD.	4 P.	Pure L's	Mixed	Pure L's	Mixed	Pure L's	Mixed
D <del>]</del>	- <u>5</u>	10.40 x 10	4.77 2 10	0.20 x 10	0.47 x 10	11.60 x 10°	12,58 x 10°
	* P <sub>1</sub>	106.23	96.59	79.06	58.08	93.66	48.93
	<sup>4</sup> P <sub>1</sub> 2	80.54	93.00	57.72	65.89	51,96	41.17
	$\left  \begin{array}{c} 2 P_3 \\ -\frac{2}{2} \end{array} \right $	0	0,72	0	7.74	0	30,95
*	<sup>2</sup> P <sub>1</sub>	0	0.58	0	3.79	0	3.66
<sup>4</sup> P <sub>3</sub>	<sup>4</sup> P <sub>5</sub> 2	67.11	79.86	51.03	67.56	61.83	92, 73
	<sup>4</sup> P <sub>3</sub>	18.77	21.67	13.38	16.69	13,11	14.86
	<sup>4</sup> P <sub>1</sub> /2	56.72	41.97	38.69	20.49	26.90	4.76
	<sup>2</sup> P <sub>3</sub>	0	0.11	0	0.73	0	3.75
	<sup>2</sup> P <sub>1</sub> /2	0	0.040	0	0.052	0	0.055
<sup>2</sup> S <sub>1/2</sub>	* P1	0	0.82	o	0.0000053	0	0,41
	${}^{4}P_{\frac{1}{2}}$	0	0.0018	0	0.32	0	2, 34
	<sup>2</sup> P <sub>3</sub>	115.86	161,43	95.83	27.63	126.35	14.90
	<sup>2</sup> P <sub>1</sub>	. 54. 30	9.24	41.81	98.32	45,17	115.01
<sup>4</sup> D <sub>1/2</sub>	4 P3	33.68	25.84	25.55	14.40	36.51	24.72
	<sup>4</sup> P <sub>1</sub>	163.34	170.51	119.48	128.23	133.01	97.63
	<sup>2</sup> P <sub>3</sub>	0	0.12	0	0.50	0	34, 71
	$^{2}P_{\frac{1}{2}}$	0	0.17	o	0.43	0	0,000026
<sup>4</sup> P <sub>1/2</sub>	<sup>4</sup> P <sub>3</sub>	119.75	124.82	88.10	94.35	102.62	106.05
	${}^{4}P_{\frac{1}{2}}$	23.162	17.90	16.33	9.74	13.92	3.02
	<sup>2</sup> P <sub>3</sub>	0	0.20	0	0.41	0	5.47
	<sup>2</sup> P <sub>1</sub> 2	0	0.0045	0	0.020	0	0.16
	ł	ł		1			

## TABLE IV

	Ne II	A II	Kr II
<sup>4</sup> D <sub>7</sub> 2	5.0579	6.8457	6.0841
<sup>4</sup> D <sub>5</sub> 2	5.1276	7.7325	9.5710
<sup>2</sup> D <sub>5</sub> /2	6.9104	8.5926	5.7454
<sup>4</sup> P <sub>5</sub> 2	6.9598 <sup>°</sup>	8.9995	7.8594
<sup>4</sup> D <sub>3</sub>	5.1102	7.3537	7.2848
<sup>2</sup> D <sub>3</sub>	6.8593	8.6801	7.8240
<sup>4</sup> P <sub>3</sub>	6.9615	9.4748	8.6004
<sup>2</sup> P <sub>3</sub> 2	5.3463	9.0649	6.2623
<sup>4</sup> S <sub>3</sub>	3.6508	4.6504	4.1172
${}^{4}D_{\frac{1}{2}}$	5.0853	6.9650	6.3331
${}^{4}P_{\frac{1}{2}}$	6.9967	9.5673	8.7047
<sup>2</sup> S <sub>1/2</sub>	5.8307	7.8869	7.3880
$^{2}P_{\frac{1}{2}}$	5.5149	9.2892	8.5162
<sup>2</sup> F <sup>1</sup> 7/2	6.1745	8.0360	6.3348
<sup>2</sup> F <sup>1</sup> <u>5</u>	6.2084	7.9657	6.7005
² D' <u>5</u>	4.5647	5.5537	4.6621
<sup>2</sup> D' <sub>3</sub> 2	4.5505	5.4526	4.5695
<sup>2</sup> P' <sub>3</sub>	4.1451	5.3876	5.0771
² P' <u>1</u>	3.9419	4.7836	3.7494

Lifetimes of some of the 3p of Ne II, 4p of A II and 5p of Kr II (in nanoseconds)

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