# LIFETIME AND TRANSITION PROBABILITIES OF $n p^{4}=(n+1) p$ StATES OF Ne II, A II AND Kr II 

S. H. Koozekanani<br>and

G. L. Trusty

# The Ohio State University ElectroScience Laboratory <br> (formerly A Atenna Laboratory) 

Department of Electrical Engineering Columbus, Ohio 43212

REPORT 1093-42
17 September 1968

Grant Number NsG-74-60


CSFTI PRICE(S) $\qquad$

Hard copy (HC).
Microfiche (MF).


## NOTICES

When Government drawings, specifications, or other data are used for any purpose other than in connection with a definitely related Government procurement operation, the United States Government thereby incurs no responsibility nor any obligation whatsoever, and the fact that the Government may have formulated, furnished, or in any way supplied the said drawings, specifications, or other data, is not to be regarded by implication or otherwise as in any manner licensing the holder or any other person or corporation, or conveying any rights or permission to manufacture, use, or sell any patented invention that may in any way be related thereto.

The Government has the right to reproduce, use, and distribute this report for governmental purposes in accordance with the contract under which the report was produced. To protect the proprietary interests of the contractor and to avoid jeopardy of its obligations to the Government, the report may not be released for non-governmental use such as might constitute general publication without the express prior consent of The Ohio State University Research Foundation.
REPORTby
The Ohio State University ElectroScience Laboratory
(Formerly Antenna Laboratory)
Columbus, Ohio 43212
Sponsor National Aeronautics and Space Administration Office of Grants and Research Contracts Washington, D.C. 20546
Grant Number ..... NsG-74-60
Investigation of Receiver Techniques and Detectors for Use at Millimeter and Submillimeter Wave Lengths
Subject of Report Lifetime and Transition Probabilities of $n p^{4}-(n+1) p$ States of Ne II, A II and Kr II
Submitted by S. H. Koozekanani G. L. Trusty
ElectroScience Laboratory
Department of Electrical Engineering
Date
17 September ..... 1968

## ABSTRACT

Life time as well as transition probabilities of the first p-excited states of neon, argon, and krypton ${ }^{\dagger}$ have been calculated.

```
LIFETIME AND TRANSITION PROBABILITIES OF
    np}\mp@subsup{}{}{4}-(n+1)p\mathrm{ STATES OF NeII, AII AND KrII
```

Lifetime and transition probabilities of some of the first excited " p " states of neon II, argon II and krypton II ${ }^{1,2,3}$ for the configuration of $n p^{4}-(n+1) p$, have been calculated. This information is needed to determine the possible efficiencies of current and proposed laser systems and to aid in the identification of the energy levels involved.

The exact mixed wave functions $\Psi_{i}\left(E_{i}\right)$ corresponding to the experimentally found energy levels $E_{i}$ can be expressed in terms of purely LS coupled wave functions $\phi_{i}$, with $\phi_{i}$ corresponding to the LS coupled states of the form $\mid p^{4} L S ; \ell e x s ; L_{i} S_{i} J_{i} M_{i}>$, i.e.,

$$
\begin{equation*}
\Psi_{i}\left(E_{i}\right)=\sum_{j} a_{i j} \phi_{j} . \tag{1}
\end{equation*}
$$

where $\bar{L} \bar{S}$ are the orbital and spin angular momentum of the $p^{4}$ core electrons while $L_{j} S_{j}$ are the total orbital and spin angular moments. $\underline{J}=\underline{L}+\underline{S}$ and $M$ is the projection of $J$ on the $z$-axis. $\ell_{\text {ex }}$ is the excited p electron with its spin s. The functions $\Psi_{i}\left(E_{i}\right)$ were found by treating the radial integrals as adjustable parameters so as to obtain the best fit with the experimentally found energy levels. This was done by choosing arbitrarily close test values for the parameters $\left[{ }^{1} \mathrm{D}\right],\left[{ }^{3} \mathrm{P}\right]$ core energy states, $\mathrm{F}_{2}, \mathrm{G}_{0}, \mathrm{G}_{2}$ direct and exchange integrals and $\zeta_{\mathrm{np}}$ and $\zeta_{\ell}$ ex
corresponding to the spin orbit interaction terms of the $n p^{4}$ core and ( $\mathrm{n}+1$ ) p excited electron.

With the inititially chosen test parameters $\left[{ }^{1} D\right]_{0}\left[{ }^{3} P\right]_{0}$. . etc. and the Hamiltonian

$$
\begin{equation*}
H=-\sum \frac{h^{2}}{2 m} \nabla^{2}-\sum_{i} \frac{e^{2}}{r_{i}}+\sum_{i>j} \frac{e^{2}}{r_{i j}}+\sum_{i} t(x) \ell_{i}, S_{i}, \tag{2}
\end{equation*}
$$

the matrix $\left[\mathrm{H}_{0}\right]$ was constructer having elements $\left[\mathrm{H}_{0}\right]_{i j}$ where

$$
\begin{equation*}
\left[H_{0}\right]_{i j}=\left\langle\phi_{i} \mid H^{\prime} \phi_{j}\right\rangle_{0} \tag{3}
\end{equation*}
$$

In Eq. (3), the subscript naught designated that in the calculation of the elements $\left[H_{C}\right]_{i j}$ the first set of test parameters are used.

A matrix $\left[M_{0}\right]$ is found such that,

$$
\begin{equation*}
\left[M_{0}\right]^{-1}\left[\mathrm{H}_{0}\right]\left[\mathrm{M}_{0}\right]=\left[\mathrm{D}_{0}\right] \tag{4}
\end{equation*}
$$

with $\left[D_{0}\right]$ being a diagonal matrix. If the choice of the parameters were appropriate $\left[D_{0}\right.$ ] will differ little from the experimentally found energy levels. If not, a new matrix equation is constructed having the form,

$$
\begin{equation*}
\left[\mathrm{M}_{0}^{-1}\right][\mathrm{H}]\left[\mathrm{M}_{0}\right]=[\mathrm{E}] \tag{5}
\end{equation*}
$$

where $[E]$ is again a diagonal matrix having for its diagonal elements the experimental energy levels $E_{i},\left[M_{0}\right]$ is the matrix which was used to diagonalixe $\left[\mathrm{H}_{0}\right]$, and $[\mathrm{H}]$ is the Hamiltonian matrix with its elements expressed in terms of the radial parameters ${ }^{4}$, $i_{\circ} e$. ,

$$
\begin{equation*}
[H]_{i j}=\left\langle\phi_{i}\right| H\left|\phi_{j}\right\rangle=a_{i j}\left[{ }^{1} D\right]+b_{i j}\left[{ }^{3} P\right]+C_{i j} F_{2}+d_{i j} G_{0}+\ldots \tag{6}
\end{equation*}
$$

Multiplying out $[M]^{-1}[H][M]$ and equating it to $[E]$ one obtains, taking the diagonal terms only, a linear set of nineteen simultaneous equations in terms of the seven parameters $\left[{ }^{1} \mathrm{D}\right],\left[{ }^{3} \mathrm{P}\right], \mathrm{F}_{2}, \mathrm{G}_{0}, \mathrm{G}_{2}, \xi_{\mathrm{np}}$ and $\zeta_{\ell \text { ex }}$. Solving the equations and least-square-fitting the results, we obtain a new set of parameters $\left[{ }^{1} D\right]_{1}\left[{ }^{3} P\right]_{1}$. . . etc. These parameters are now used to obtain the matrix $\left[\mathrm{H}_{1}\right]$, with matrix elements $\left[\mathrm{H}_{1}\right]_{\mathrm{ij}}{ }_{\mathrm{in}}$ which the set of the newly found values of the radial parameters have been substituted, i.e.,
$\left[H_{i}\right]_{i j}=a_{i j}\left[{ }^{1} D\right]_{1}+b_{i j}\left[{ }^{3} \underline{P}\right]_{1}+C_{i j}\left(F_{2}\right)_{1}+d_{i j}\left(G_{0}\right)_{1}+.$.
Once again we find a matrix $\left[M_{1}\right]$ which diagonalizes $\left[H_{1}\right]$ giving the diagonal matrix $\left[D_{1}\right]$ i.e.,

$$
\begin{equation*}
\left[\mathrm{M}_{1}\right]^{-1}\left[\mathrm{H}_{1}\right]\left[\mathrm{M}_{1}\right]=\left[\mathrm{D}_{1}\right] \tag{7}
\end{equation*}
$$

If again the elements $\left[D_{1}\right]$ would differ much from the experimentally found energy levels $F_{f}$ the process is repeated by solving the simultaneous equation obtained from the matrix equation,

$$
\begin{equation*}
\left[M_{1}\right]^{-1}[H]\left[M_{1}\right]=[E] \tag{8}
\end{equation*}
$$

where again the elements of the matrix $H$ are given by Eq. 6. A new set
of parameters $\left[{ }^{1} \mathrm{D}\right]_{2},\left[{ }^{3} \mathrm{P}\right]_{2}{ }^{\circ}{ }^{\circ}$. etc., are found. The process is repeated until the parameters $\left[{ }^{1} D\right]_{n} \quad\left[{ }^{3} P\right]_{n}$... etc. converge and the matrix equation $\left[M_{n}\right]^{-1}\left[H_{n}\right]\left[M_{n}\right]=\left[D_{n}\right] \simeq[E]$.

If the initial set of test parameters are close to the actual values, the process would terminate after a few cycles, and the parameters would converge to fixed values. If not, the convergence is not easily attained. In the above calculations the states arising from the ${ }^{1} S$ core were not taken into consideration However, their contribution to the states ${ }^{2} P_{\frac{1}{2}}$, ${ }^{2} P_{\frac{3}{2}},{ }^{2} P_{\frac{1}{2}}$ and ${ }^{2} P_{\frac{3}{2}}{ }^{2}$ can be calculated by perturbation theory. The
 wavefunctions and also the values which were obtained by the above fitting techniques. The two sets of values are more or less very close. Table II gives the theoretical and experimentally found energy levels and the mixing coefficients $\mathrm{a}_{\mathrm{ij}}$. Table III gives the transition probabilities for the states arising from the $n p^{4}-(n+1) p$ configuration of neon, argon, and krypton to the lower states of $n p^{4}-(n+1)$ s configuration. Finally Table IV gives the lifetimes of some of the higher states of interest.
TABLEI

e-II GTg

TABLE II-b
Experimental and theoretical energy levels of Ar II taken from Minnhagen ${ }^{8}$ and from this calculation. The mixing coefficients of $\mathrm{a}_{\mathrm{ij}}$ are those given in Eq. (1).

| Ar II States | $\begin{aligned} & E_{\exp }^{1} \\ & \mathrm{Cm}_{1} \end{aligned}$ | $\begin{aligned} & \mathrm{E}_{\text {thepr }} \\ & \mathrm{Cm}^{\prime}-9 \end{aligned}$ | ${ }^{4} \mathrm{D}$ |  | ${ }^{4} \mathrm{~S}$ | ${ }^{2} \mathrm{D}$ | ${ }^{2} \mathrm{P}$ | ${ }^{2} \mathrm{~S}$ | ${ }^{2} \mathrm{Pr}$ - | ${ }^{2} \mathrm{D}{ }^{\prime}$ | ${ }^{2} \mathrm{Fi}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{4} \mathrm{D}_{\frac{7}{2}}$ | 157233.93 | 157176.6 | 0.99839 |  |  |  |  |  |  |  | 0.05680 |
| ${ }^{2} \mathrm{~F}^{1} 7$ | 170530. 31 | 170582.7 | 0.05680 |  |  |  |  |  |  |  | -0.99839 |
|  |  |  | 0.18291 | 0.98216 |  | 0.02312 |  |  |  | 0.03949 | 0.00299 |
| ${ }^{4} \mathrm{P}_{5}$ | 155043.07 | 155128.2 | 0.18291 | 0.98216 |  |  |  |  |  |  |  |
| ${ }^{4} \mathrm{D}_{\frac{5}{2}}$ | 157673. 32 | 157556.9 | 0.86726 | -0.15036 |  | -0.47202 |  |  |  | -0.00525 | 0.04919 |
| ${ }^{2} \mathrm{D}_{5}$ | 158730.21 | 158710.3 | -0.46181 | 0.10820 |  | -0.87867 |  |  |  | -0.04014 | 0.03689 |
| ${ }^{2} \mathrm{~F}$ ' | 170400.94 | 170364.6 | -0.02553 | 0.00160 |  | 0.05688 |  |  |  | -0.03043 | 0.99759 |
|  | 173393.38 | 173432.7 | 0.02203 | 0.03525 |  | 0.03701 |  |  |  | -0.99794 | -0.03204 |
| ${ }^{\text {d }}$ | , |  |  |  |  |  |  |  |  |  |  |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 155351.03 | 155383.5 | 0.16724 | 0.96996 | 0.11854 | -0. 1561 | -0.11366 |  | 0.05959 | 0.021 |  |
| ${ }^{4} \mathrm{D}_{\frac{3}{2}}$ | 158167.71 | 158007.6 | 0.91426 | -0.18070 | -0.04193 | -0. 31306 | -0.17205 |  | 0.04422 | 0.01261 |  |
| ${ }^{2}{ }^{\text {d }}$ | 159393.31 | 159244.3 | 0.36791 | 0.00919 | 0.00421 | 0. 80472 | 0.44766 |  | -0.12847 | 0.00763 |  |
| ${ }^{2} \mathrm{P}_{3}$ | 160239.35 | 160354.7 | 0. 01598 | 0.09361 | 0.06872 | -0. 50253 | 0.81016 |  | -0. 27314 | -0.05244 |  |
| ${ }^{4} 5_{3}{ }^{2}$ | 161048.64 | 161012.7 | 0.01607 | -0.13135 | 0.98834 | 0.01899 | -0.03572 |  | 0.06355 | 0.00052 |  |
| ${ }^{2} \mathrm{P}$ | 172213.80 | 172213.8 | -0.00095 | 0.01493 | 0.05133 | 0.0218 | -0.30918 |  | -0.94747 | 0.05873 |  |
| ${ }^{2} \mathrm{D}^{1}$ | 173347.83 | 173349.9 | 0.01702 | 0.01413 | 0.00192 | 0.02951 | -0.06202 |  | -0.04059 | -0.99657 |  |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 155708.02 | 155781.9 | 0.09789 | 0.99105 |  |  | -0.06195 | -0.05442 | 0.03795 |  |  |
| ${ }^{4}$ | 158428.03 | 158331.0 | -0.98840 | 0. 10630 |  |  | 0.09918 | 0.01678 | -0.04050 |  |  |
| ${ }^{2}{ }^{2}$ | 159706.46 | 159869.7 | -0.11005 | -0.07824 |  |  | -0.80018 | -0.49034 | 0. 31788 |  |  |
|  | 161089. 31 | 161306.5 | -0.03676 | 0.01627 |  |  | -0. 46980 | 0.86912 | 0.14933 |  |  |
| 2 | 172816. 21 | 172726.4 | -0.00350 | -0.01162 |  |  | 0. 35401 | 0.03085 | 0.93465 |  |  |

TABLE II-c
Experimental and theoretical energy levels of Kr II taken from Moore ${ }^{7}$, tables of atomic
energy levels and this calculation. The mixing coefficients $\mathrm{a}_{\mathrm{ij}}$ are those given in Eq. (1).

|  | $\begin{aligned} & E_{\exp } \\ & C_{m}{ }^{-1} \end{aligned}$ | $E_{\text {theor }}$ $\mathrm{Cm}^{-1}$ | ${ }^{4} \mathrm{D}$ | ${ }^{4} \mathrm{p}$ | ${ }^{4}$ | ${ }^{2} \mathrm{D}$ | ${ }^{2} \mathrm{P}$ | ${ }^{2} \mathrm{~S}$ | ${ }^{2} \mathrm{PI}$ | ${ }^{2} \mathrm{D}{ }^{\text {P }}$ | ${ }^{2}{ }^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{4} \mathrm{D}_{\frac{7}{2}}$ | 135783. 18 | 136046.4 | 0.98818 |  |  |  |  |  |  |  | 0.15329 |
| ${ }^{2} \mathrm{~F}^{1} \frac{7}{8}$ | 149704.55 | 149720.5 | 0.15329 |  |  |  |  |  |  |  | -0.98818 |
| ${ }^{4}{ }^{\text {P }}$ | 133925.65 | 133855.3 | 0.35157 | 0.92956 |  | -0.00191 |  |  |  | 0.10898 | 0.02069 |
| ${ }^{4} \mathrm{D}_{\frac{5}{2}}$ | 136071.00 | 136913.7 | 0.69899 | -0.26376 |  | -0.64380 |  |  |  | -0.04663 | 0.15873 |
| ${ }^{2} \mathrm{D}_{\frac{5}{2}}$ | 140118.99 | 139794.1 | -0.61630 | 0.24015 |  | -0.73946 |  |  |  | -0.08975 | 0.08744 |
| ${ }^{2} \mathrm{~F}_{\frac{5}{2}}$ | 149173.42 | 149096.5 | -0.06087 | 0.00875 |  | 0.17647 |  |  |  | -0.05814 | 0.98066 |
| ${ }^{2} \mathrm{D}_{5}{ }^{\text {s }}$ | 152316. 26 | 152377.9 | 0.06542 | 0.09275 |  | 0.08704 |  |  |  | -0.98717 | -0.07092 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 134288. 44 | 134032.7 | 0.29227 | 0.86841 | 0.24557 | -0.06918 | -0. 25507 |  | 0.15819 | 0.07257 |  |
| ${ }^{4}{ }^{\frac{3}{2}}$ | 138381, 35 | 138171.7 | 0.70387 | -0.33387 | -0.23513 | -0.42128 | -0. 39352 |  | 0.05687 | -0.04741 |  |
| ${ }^{2} \mathrm{D}_{\frac{3}{2}}$ | 141995.68 | 142424.3 | 0.15404 | -0.21173 | 0.13765 | 0.80116 | -0.43082 |  | 0.25429 | 0.14271 |  |
| ${ }^{2} \mathrm{P}_{2}$ | 140137.15 | 140117.4 | 0.62449 | 0.04557 | 0.03686 | 0.32338 | 0.67428 |  | -0.21368 | -0.04111 |  |
| ${ }^{4} \mathrm{~S}_{\frac{2}{2}}$ | 141722.72 | 141508. 3 | 0.05458 | -0.29324 | 0.91550 | -0.23728 | 0.06388 |  | 0.10487 | -0.03901 |  |
| ${ }^{2}{ }^{\text {P }}{ }_{\frac{3}{2}}$ | 150203.48 | 150519.5 | -0.00713 | 0.02327 | 0. 16055 | 0.07215 | -0.32003 |  | -0.91809 | 0.15208 |  |
| ${ }^{2} \mathrm{D}_{\frac{3}{2}}$ | 152191.86 | 152172.6 | 0.04903 | 0.03093 | 0.01389 | 0.09900 | -0.18255 |  | -0.08685 | -0.97251 |  |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 135783.03 | 135638.4 | 0.20533 | 0.92476 |  |  | -0.20833 | -0.19012 | 0.15203 |  |  |
| ${ }^{4} \mathrm{D}_{\frac{1}{2}}$ | 140163.25 | 139549.3 | -0.85103 | 0.01265 |  |  | -0.28601 | -0.40360 | 0.17577 |  |  |
| ${ }^{2}{ }^{1} \frac{1}{2}$ | 139103.36 | 139042.2 | -0.47019 | 0.37937 |  |  | 0.56136 | 0.46584 | -0. 32072 |  |  |
| ${ }^{4} S_{\frac{1}{2}}$ | 142363. 55 | 142592.5 | -0.11170 | 0.00795 |  |  | -0.59032 | 0.76087 | 0.24507 |  |  |
| ${ }^{2}{ }^{\text {P1 }}{ }_{\frac{1}{2}}$ | 152240.97 | 152094.4 | -0.00571 | -0.02610 |  |  | 0.45955 | 0.07095 | 0.88491 |  |  |

TABLE III
Transition probabilities between $3 \mathrm{p}-3 \mathrm{~s}$ levels of Ne II, $4 \mathrm{p}-4 \mathrm{~s}$ levels of Ar II and $5 p-5 s$ levels of Kr II, each with a $\left.\mathrm{P}^{4}{ }^{3} \mathrm{P}\right]$ core. A in $\mathrm{sec}^{-1}$.

|  | Neon II |  | Argon II |  | Krypton II |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Pure L's | Mixed | Pure L's | Mixed | Pure L' ${ }_{\text {c }}$ | Mixed |
| ${ }^{4} \mathrm{D}_{\frac{7}{2}}-{ }^{3} \mathrm{P}_{\frac{5}{2}}$ | $197.81 \times 10^{6}$ | $197.71 \times 10^{6}$ | $146,53 \times 10^{6}$ | $146.06 \times 10^{6}$ | $168.15 \times 10^{6}$ | $164 \times 10^{6}$ |
| ${ }^{2} \mathrm{D}_{\frac{5}{2}}={ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 0 | 0.28 | 0 | 4.67 | 0 | 5.28 |
| $-{ }^{4} P_{\frac{3}{2}}$ | 0 | 4.12 | 0 | 31.65 | 0 | 91.27 |
| $-{ }^{2} \mathrm{P}_{\frac{3}{2}}$ | 143.30 | 140.31 | 103.65 | 80.03 | 140.97 | 77.08 |
| ${ }^{4} \mathrm{D}_{\frac{5}{4}}-{ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 61.37 | 43.08 | 46.53 | 18.91 | 52.37 | 4.59 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 135.98 | 149.75 | 97.25 | 90.69 | 89.85 | 68.27 |
| ${ }^{2} P_{\frac{3}{2}}$ | 0 | 2. 20 | 0 | 19.70 | 0 | 31. 39 |
| ${ }^{4} \mathrm{P}_{\frac{5}{2}}-{ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 101.86 | 114.65 | 79.28 | 96.24 | 91.38 | 122.90 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 41.20 | 28.96 | 30.06 | 14.84 | 27.86 | 4.29 |
| ${ }^{4} S_{\frac{7}{2}} \quad 4 P_{\frac{5}{2}}$ | 142.05 | 124.62 | 116.11 | 87.65 | 167.69 | 71.85 |
| ${ }^{4} P_{3}$ | 90.42 | 93.83 | 70.32 | 77.64 | 87.52 | 93.20 |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 44.00 | 52.87 | 33.11 | 49.39 | 32.51 | 69.82 |
| ${ }^{2} \mathrm{P}_{\frac{3}{2}}$ | 0 | 1.41 | 0 | 0.088 | 0 | 0.26 |
| ${ }^{2} \mathrm{P}_{\frac{1}{2}}$ | 0 | 1.17 | 0 | 0.13 | 0 | 7.28 |
| ${ }^{2} \mathrm{P}_{\frac{3}{2}} \quad{ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 0 | 0.59 | 0 | 2.80 | 0 | 10.91 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 0 | 3.29 | 0 | 0.0073 | 0 | 44.51 |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 0 | 1.31 | 0 | 0.088 | 0 | 24.04 |
| ${ }^{2} P_{\frac{3}{2}}$ | 166.49 | 119.05 | 106.91 | 36.64 | 117.77 | 78.98 |
| ${ }^{2} P_{\frac{1}{2}}$ | 31.30 | 62.62 | 18.56 | 69.40 | 16.23 | 0.039 |
| ${ }^{2} \mathrm{D}_{\frac{3}{2}} \quad{ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 0 | 0.20 | 0 | 1.60 | 0 | 0.036 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 0 | 1.00 | 0 | 12.59 | 0 | 0.53 |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 0 | 1.35 | 0 | 8.83 | 0 | 17.24 |
| ${ }^{2} \mathrm{P}_{3}$ | 25.27 | 52.23 | 19.01 | 61.98 | 30.15 | 0.74 |
| ${ }^{2} P_{\frac{1}{2}}$ | 118.08 | 91.00 | 82.01 | 29.97 | 107. 19 | 105.87 |


|  | Neon II <br> Pure L's | con't. <br> Mixed | Argon <br> Pure L's | on't. <br> Mixed | Krypton Pure L's | Mixed |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{r} D_{\frac{1}{2}} P_{\frac{5}{2}} \\ \end{array} P_{\frac{1}{2}}$ | $\begin{aligned} & 10.48 \times 10^{6} \\ & 106.23 \end{aligned}$ | $\begin{aligned} & 4.97 \times 10^{6} \\ & 96.59 \end{aligned}$ | $\begin{aligned} & \frac{\text { Pure L's }}{8.26 \times 10^{6}} \\ & 79.06 \end{aligned}$ | $\frac{\text { Mixed }}{0.47 \times 10^{6}}$ | $\frac{\text { Pure L's }}{11.60 \times 10^{6}}$ | $12.58 \times 10^{6}$ |
|  |  |  |  | 58,08 | 93.66 | 48.93 |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 80.54 | 93.00 | 57.72 | 65.89 | 51.96 | 41.17 |
| ${ }^{2} \mathrm{P}_{\frac{3}{2}}$ | 0 | 0.72 | 0 | 7.74 | 0 | 30.95 |
| ${ }^{2} \mathrm{P}_{\frac{1}{2}}$ | 0 | 0.58 | 0 | 3.79 | 0 | 3.66 |
| ${ }^{4} \mathrm{P}_{3}$ | 67.11 | 79.86 | 51.03 | 67.56 | 61.83 | 92.73 |
|  | 18.77 | 21.67 | 13.38 | 16.69 | 13.11 | 14.86 |
|  | 56.72 | 41.97 | 38.69 | 20. 49 | 26.90 | 4.76 |
|  | 0 | 0.11 | 0 | 0.73 | 0 | 3.75 |
|  | 0 | 0.040 | 0 | 0.052 | 0 | 0.055 |
| ${ }^{2} S_{\frac{1}{2}}$ | 0 | 0.82 | 0 | 0.0000053 | 0 | 0.41 |
|  | 0 | 0.0018 | 0 | 0.32 | 0 | 2. 34 |
|  | 115.86 | 161.43 | 95.83 | 27.63 | 126.35 | 14.90 |
|  | 54. 30 | 9.24 | 41.81 | 98.32 | 45.17 | 115.01 |
| ${ }^{4} \mathrm{D}_{\frac{1}{2}}$ | 33.68 | 25.84 | 25.55 | 14.40 | 36.51 | 24.72 |
|  | 163.34 | 170.51 | 119.48 | 128. 23 | 133.01 | 97.6. |
|  | 0 | 0.12 | 0 | 0.50 | 0 | 34. 71 |
|  | 0 | 0.17. | 0 | 0.43 | 0 | 0.000026 |
| ${ }^{4} P_{\frac{1}{2}}$ | 119.75 | 124.82 | 88. 10 | 94. 35 | 102.62 | 106.05 |
|  | 23.162 | 17.90 | 16.33 | 9.74 | 13.92 | 3.02 |
|  | 0 | 0.20 | 0 | 0.41 | 0 | 5.47 |
|  | 0 | 0.0045 | 0 | 0.020 | 0 | 0.16 |

TABLE IV
Lifetimes of some of the 3p of Ne II, 4p of A II and 5 p of Kr II (in nanoseconds)

|  | Ne II | A II | Kr II |
| :---: | :---: | :---: | :---: |
| ${ }^{4} \mathrm{D}_{\frac{7}{2}}$ | 5.0579 | 6.8457 | 6.0841 |
| ${ }^{4} \mathrm{D}_{\frac{5}{2}}$ | 5.1276 | 7.7325 | 9.5710 |
| ${ }^{2} \mathrm{D}_{\frac{5}{2}}$ | 6.9104 | 8.5926 | 5.7454 |
| ${ }^{4} \mathrm{P}_{\frac{5}{2}}$ | 6.9598 | 8.9995 | 7.8594 |
| ${ }^{4} \mathrm{D}_{\frac{3}{2}}$ | 5.1102 | 7.3537 | 7.2848 |
| ${ }^{2} \mathrm{D}_{\frac{3}{2}}$ | 6.8593 | 8.6801 | 7.8240 |
| ${ }^{4} \mathrm{P}_{\frac{3}{2}}$ | 6.9615 | 9.4748 | 8.6004 |
| ${ }^{2} \mathrm{P}_{\frac{3}{2}}$ | 5. 3463 | 9.0649 | 6.2623 |
| ${ }^{4} \mathrm{~S}_{\frac{3}{2}}$ | 3.6508 | 4.6504 | 4. 1172 |
| ${ }^{4} \mathrm{D}_{\frac{1}{2}}$ | 5.0853 | 6.9650 | 6.3331 |
| ${ }^{4} \mathrm{P}_{\frac{1}{2}}$ | 6.9967 | 9.5673 | 8.7047 |
| ${ }^{2} S_{\frac{1}{2}}$ | 5.8307 | 7.8869 | 7. 3880 |
| ${ }^{2} P^{\frac{1}{2}}$ | 5.5149 | 9.2892 | 8.5162 |
| ${ }^{2} F_{\frac{7}{2}}^{\prime}$ | 6.1745 | 8.0360 | 6.3348 |
| ${ }^{2} \mathrm{~F}_{\frac{1}{2}}$ | 6.2084 | 7.9657 | 6.7005 |
| ${ }^{2} \mathrm{D}_{\frac{5}{2}}^{\prime}$ | 4.5647 | 5.5537 | 4.6621 |
| ${ }^{2} D_{\frac{3}{2}}^{\prime}$ | 4. 5505 | 5.4526 | 4. 5695 |
| ${ }^{2} \mathrm{Pt}_{\frac{3}{2}}$ | 4.1451 | 5.3876 | 5.0771 |
| ${ }^{2} \mathrm{P}^{\mathrm{t}_{\frac{1}{2}}}$ | 3.9419 | 4.7836 | 3.7494 |

## REFERENCES

1. Garstang, R. H., Monthly Notices, Roy. Ast. Soc., 114, p. 118, 1954.
2. Koopman, D. W., "Line Strengths for Neutral and Singly Ionized Neon," J. of Op. Soc. of Am., Vol. 54, no. 11, November 1964.
3. Statz, H., F. A. Horrigan, S. H. Koozekanani, C. L. Tang and G. F. Koster, J. of Appl. Phys., Vol. 36, no. 7, p. 2278, July 1965.
4. See T. Yamanouchi and A. Amemiya, Proc. Phys. Soc. Japan, 1, p. 18, 1946, for the exchange and direct integral, and J. L. Tech, J. Res. Nat. Bur. Std., 67A, p. 555, 1963, for the spin orbit terms.
5. It can be shown that the equations arising from the off diagonal terms give no information.
6. Herman, F., and S. Steillman, "Atomic Structure Calculations," Prentice Hall Inc., Englewood Cliffs, N. Jo, 1963.
7. Moore, C. E., "Atomic Energy Levels," Circular of the Nat. Bur. of Std., U. S. Gov. Printing Office, Washington 25, D. C., 1949.
8. Minnhagen, L., "The Spectrum of Singly Ionized Argon, Ar II," Arkiv Foir Fysik, Band 25, nr. 19, Almquist and Wiksell, Stockholm, Sweden, 1963.
