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STUDY OF INCOMPLETE MODELS
OF DYNAMIC STRUCTURES

Contract No.: NAS5-21007

Prepared by

Kaman Aerospace Corporation
Bloomfield, Connecticut

for

Goddard Space Flight Center
Greenbelt, Maryland

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SUMMARY

This report presents the theoretical development and a qualitative evaluation of a new concept in the mathematical modeling of dynamic structures. The method has several unique features. It uses both test data and analytical approximations to identify the parameters of what is termed an "incomplete model". The technique makes one of the first substantial applications of the matrix pseudo inverse to structural dynamics. The "incomplete model" itself is unusual in that, while the mass matrix and damping coefficient are physically meaningful, the stiffness matrix is, in a sense, not physically related to the true (but unknown) values. The model does, however, have the capability of predicting the response of the points of interest on the structure over the frequency range of interest and can be used to predict the changes in natural frequencies and normal modes due to structural changes. It is the ability to correctly predict the effects of structural changes which have been tested in the work presented in this report.

The theory was tested by running simulated tests on a relatively simple structure, identifying the parameters of the incomplete model, and using this model to predict

the effects on frequency and mode shapes of several mass and stiffness changes. The conditions of the tests were varied by selecting different numbers of points of measurement, varying the frequency range, and by including assumed measurement error. Several analytical approximations were used to formulate the incomplete model.

Based on the qualitative evaluation made, the following conclusions are indicated.

1. The method presented appears to be an excellent tool for the prediction of changes in normal modes and natural frequencies due to structural changes.
2. The predicted effects are quite insensitive to reasonable variations in the analytical mass distribution.
3. The procedure is relatively insensitive to measurement error.
4. Attempts to make the model more complete by using a greater frequency range for testing will not always be successful unless the number of points of measurement are also increased or off-diagonal masses are added. There is a point in frequency beyond which the model will tend to deteriorate.

Because of the success of this preliminary evaluation, it is recommended that the theoretical development be continued and that applications to more complex structures be carried out in order to develop a better understanding of the limitations and capabilities of the method. A successful, more definitive evaluation, could lead to immediate practical applications.

TABLE OF CONTENTS

	<u>PAGE</u>
SUMMARY.	ii
TABLE OF CONTENTS.	v
LIST OF FIGURES.	vii
LIST OF TABLES	viii
INTRODUCTION	1
I. DESCRIPTION OF THE PROBLEM	5
II. BASIC RELATIONSHIPS.	7
III. THE INCOMPLETE MODEL	14
IV. IDENTIFICATION OF THE MASS MATRIX.	18
V. MASS CHANGES	25
VI. STIFFNESS CHANGES.	27
VII. IDENTIFICATION OF DAMPING COEFFICIENT.	30
VIII. DESCRIPTION OF COMPUTER PROGRAM.	32
IX. THE SIMULATED TESTS.	37
X. THE IDENTIFIED MASSES.	43
XI. THE EFFECTS OF MASS CHANGES.	48
XII. THE EFFECTS OF STIFFNESS CHANGES	57
XIII. THE IDENTIFIED DAMPING COEFFICIENTS.	65
XIV. CONCLUSIONS.	67
NEW TECHNOLOGY	69

TABLE OF CONTENTS (Continued)

	<u>PAGE</u>
BIBLIOGRAPHY.	70
GLOSSARY OF SYMBOLS	72
APPENDIX I - PRINCIPAL IDEMPOTENTS	
APPENDIX II - PSEUDO INVERSES	
APPENDIX III - PROGRAM LISTING AND SAMPLE COMPUTATION	

LIST OF FIGURES

<u>FIGURE</u>		<u>PAGE</u>
1	Logic Diagram of Computer Program.	33
2	Schematic Representation of the Simulated Test Specimens	39
3	Exact Normal Modes of Test Specimens . . .	40
4a	Computed Mode Changes Due to Mass Change First Mode, Specimen I, Change B	53
4b	Computed Mode Changes Due to Mass Change Second Mode, Specimen I, Change B.	54
4c	Computed Mode Changes Due to Mass Change Third Mode, Specimen I, Change B	55
4d	Computed Mode Changes Due to Mass Change Third Mode, Specimen II, Change B.	56
5a	Computed Mode Changes Due to Stiffness Change First Mode, Specimen I, Change L	61
5b	Computed Mode Changes Due to Stiffness Change Second Mode, Specimen I, Change L	62
5c	Computed Mode Changes Due to Stiffness Change Third Mode, Specimen I, Change L	63
6	Stiffness Changes.	64

LIST OF TABLES

<u>TABLE</u>		<u>PAGE</u>
I	APPROXIMATE MASS DISTRIBUTIONS.	41
II	DESCRIPTION OF SIMULATED TESTS.	42
III	IDENTIFIED MASSES	44
IV	COMPUTED FREQUENCIES DUE TO MASS CHANGE	49
V(a)	COMPUTED FREQUENCIES DUE TO STIFFNESS CHANGE.	59
V(b)	COMPUTED FREQUENCIES DUE TO STIFFNESS CHANGE 'M'.	60
VI	IDENTIFIED GENERALIZED MASSES AND DAMPING COEFFICIENTS.	66

INTRODUCTION

The need for mathematical models of aerospace structures is obvious. No attempt to determine a structure's suitability for its intended mission could be made without a means of predicting its response to the expected loads. No rational means of improvement or optimization would be available without some ability to predict the effects of structural changes on the response of the structure.

Purely analytical modeling is necessary in the preliminary design stage of any structure. The evaluation of alternative basic designs, the selection of materials, the arrangement of components are some of the problems which must be solved prior to construction of a test specimen. The science of structural analysis has progressed to the point where such problems may be treated with reasonable confidence.

Prior to usage, any critical structure must be submitted to actual testing, since the analysis, after all, is only based on a hypothetical model of the system. An analytical model is as important at this stage as it was in the preliminary design. The actual mission force environment cannot economically be duplicated in test and the effects of possible changes cannot all be tested.

It would be reassuring if the results of dynamic testing could be used to give an analytical model of the actual structure. To date, however, there is no generally accepted method applicable to the dynamics of structures. In addition, it appears that the unique identification of an analytical model from test data alone must be limited to rather special conditions.

The number of degrees of freedom of a linear model of a dynamic system is equal to the number of independent spatial coordinates and also to the number of normal modes of the system. The response of any point on the structure is the superposition of the responses of each of the normal modes at that point. It is recognized that in linear systems, measurable excitation of each normal mode is only achieved by forces having frequencies near or above the natural frequency of the normal mode in question. Thus, in order that test data contain the information necessary for the unique identification of a linear model having a specified number of degrees of freedom, it is necessary that the applied forces must have had significant components at frequencies up through the same number of natural frequencies. Data obtained from smaller ranges of frequencies of excitation would not contain sufficient significant information for such a unique identification.

There are conditions where it is possible to perform a unique identification. One is where the number of points of interest on the structure is small and it is feasible to excite the structure over the appropriate frequency range. Such a situation is treated in Reference 1. Another condition is where the major portion of the motions are due to rigid body responses and there is in effect only six physical degrees of freedom. This is the primary area of past applications of identification techniques to such problems as determining aerodynamic stability derivatives.

The problem considered in this report is concerned with what is probably the more common situation when the application is made to elastic structures. It is taken to be impossible, uneconomical, or generally undesirable to excite the structure and measure responses over a sufficiently wide frequency range so as to provide the information necessary for the identification of a unique linear analytical model.

In this report, the concept of an "incomplete model" is developed. This is an analytical model which has fewer normal modes than coordinates and is valid only over a limited frequency range. It is shown how the parameters may be determined using test data together with analytical approximations. The attributes of the model are discussed

and certain of these are tested using computer simulated test data. It is shown that it is possible to determine a rational mass matrix and structural damping coefficient and that it is possible to use the incomplete model to predict the changes in frequencies and normal modes due to mass and stiffness changes. It is further shown that the procedures are relatively insensitive to measurement error.

This research project has demonstrated a new and useful analytical tool. It is suggested that further development of the theory and methods of application would result in procedures which will lead to better understanding of structures and increased capability to improve and optimize the design of aerospace vehicles.

I. DESCRIPTION OF THE PROBLEM

An aspect of the following problem is considered in this report. Consider a structure which is to be subjected to dynamic testing. There are certain deflections and rotations of points on the structure which are of interest. It is desired to be able to predict each of these motions under various loading conditions for the structure actually tested and for modified versions of the structure.

It is assumed that an analytical model consisting of P lumped masses interconnected by linear springs with scalar structural damping will be adequate to represent the deformation of the structure under consideration (the more general mass matrix with off-diagonal terms is not excluded). The motion of each of the P masses is considered to represent the motion of P points on the structure including the points of interest. It is also assumed that the loading conditions of interest will contain primarily force components at frequencies below some finite value including less than P natural frequencies of the model. It is required that the analytical model faithfully represent the dynamics of the "points of interest" over the "frequency range of interest". It is further required that the model have the capability of predicting the

changes in response due to structural changes including mass, stiffness and support changes.

The work reported here is limited in scope as follows. The analysis assumes knowledge of the normal mode shapes at each of P points and natural frequencies of the structure through the frequency range of interest (Knowledge of the modal masses is not required). It is considered that this data has been obtained from testing. The P points are distributed over the structure so as to represent the characteristics of the measured modes and include all the points of interest. In the rest of this report, all these points will be considered to be the "points of interest".

In the computer experiments described below, the dissipative component of the resonant response was taken as a good approximation to the normal mode. However, other more sophisticated procedures such as those given in References 1 or 2 might be justified, for example, if the resonances were not well separated. In addition, it is assumed that a "reasonable" analytical model of the mass matrix of the structure has been derived through analytical or intuitive means. Using these assumed data, a procedure is developed for identifying the parameters in the equations of motion such that the model has the capability of predicting the effects of changes in mass and stiffness on natural frequencies and modes.

II. BASIC RELATIONSHIPS

The material presented in this section is not new but is derived here in the form in which it will be referenced in following sections of this report. Definitions of all the symbols used are summarized in the Glossary of Symbols.

The matrix equation of the spring-mass-structural damping model discussed in the previous section may be written (see, for example, Reference 3)

$$M\ddot{y}(t) + (1 + ig)Ky(t) = f(t) \quad (1)$$

\ddot{y} , y , f are column matrices (vectors) representing the acceleration, displacement (or slope), applied force (or moment) as a function of time at each of the P points of interest. M and K are $P \times P$ symmetric matrices representing the mass and stiffness coefficients of the model. i is $\sqrt{-1}$ and g is the structural damping coefficient. The imaginary term is used here to indicate that the damping is in phase with the velocity but proportional to the displacement. Under steady state sinusoidal oscillation at a frequency, ω , the equation becomes

$$\{-\omega^2 M + (1 + ig)K\}y = f \quad (2)$$

where f may be thought of as the amplitude of the applied forces and y as the amplitude of the displacements. y will, in general, be complex because of the phase between the force and displacement.

Writing this equation in terms of velocity ($\dot{y} = i\omega y$)

$$\left\{ \frac{g}{\omega} K + i \left(\omega M - \frac{1}{\omega} K \right) \right\} \dot{y} = f \quad (3)$$

the (velocity) impedance matrix is given by

$$Z = \frac{g}{\omega} K + i \left(\omega M - \frac{1}{\omega} K \right) \quad (4)$$

and the mobility matrix by

$$Y = Z^{-1} = \left\{ \frac{g}{\omega} K + i \left(\omega M - \frac{1}{\omega} K \right) \right\}^{-1} \quad (5)$$

It is interesting to note that the quantity directly obtained through analysis is the impedance, Z , while the quantity measured in testing is the mobility, Y .

Consider now, the eigenvalue problem corresponding to Equation (2).

$$\{K - \Omega_i^2 M\} \phi_i = 0 \quad i = 1, 2, \dots, P \quad (6)$$

where Ω_i are the natural frequencies and ϕ_i are the normal modes of the system. This equation may be written in two ways:

$$M^{-1}K\phi_i = \Omega_i^2 \phi_i \quad i = 1, 2, \dots, P \quad (7)$$

or

$$CM\phi_i = \frac{1}{\Omega_i^2} \phi_i \quad i = 1, 2, \dots, P \quad (8)$$

where C, the influence coefficient matrix, equals K^{-1} . The orthogonality relationship is given by

$$\begin{aligned} \phi_i^T M \phi_j &= 0 & j \neq i \\ &= m_i & j = i \end{aligned} \quad (9)$$

It is important to note that the normal modes (eigenvectors) of $M^{-1}K$ and its inverse, CM , are the same and that the respective eigenvalues are reciprocals. The dominant mode of $M^{-1}K$ is the one having the highest frequency and the dominant mode of CM is the one having the lowest frequency.

It is convenient to write (7), (8), (9) using the ϕ matrix where

$$\phi = [\phi_1 \quad \phi_2 \quad \dots \quad \phi_p] \quad (10)$$

is a square $P \times P$ matrix. Then

$$M^{-1}K\phi = \phi [\Omega_i^2] \quad (11)$$

$$CM\phi = \phi \left[\frac{1}{\Omega_i^2} \right] \quad (12)$$

and

$$\phi^T M \phi = [m_i] \quad (13)$$

From each of the above three equations, one can write

$$K = M\phi \left[\frac{\Omega_i^2}{m_i} \right] \phi^{-1}$$

$$C = \phi \left[\frac{1}{\Omega_i^2} \right] \phi^{-1} M^{-1}$$

and

$$\phi^{-1} = \left[\frac{1}{m_i} \right] \phi^T M$$

Substituting the third equation into the first two

$$K = M\phi \left[\frac{\Omega_i^2}{m_i} \right] \phi^T M \quad (14)$$

$$C = \phi \left[\frac{1}{\Omega_i^2 m_i} \right] \phi^T \quad (15)$$

These two equations may be written in the equivalent form

$$K = \sum_{i=1}^p \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M \quad (16)$$

$$C = \sum_{i=1}^p \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T \quad (17)$$

In connection with the discussion in Section VI, these may be written in terms of principal idempotents (See Appendix I) as follows

$$M^{-1}K = \sum_{i=1}^p \Omega_i^2 \left(\frac{1}{m_i} \phi_i \phi_i^T M \right) \quad (18)$$

$$CM = \sum_{i=1}^P \frac{1}{\Omega_i} \left(\frac{1}{m_i} \phi_i \phi_i^T M \right) \quad (19)$$

Note that the product $\phi_i \phi_i^T$ appearing in Equations (16) - (19) are square matrices of order p but rank 1. When P of the matrices of this type are summed as indicated, there result $P \times P$ matrices of rank P which are thus nonsingular.

The impedances and the mobility may also be expressed in terms of the normal modes. Write Equation (4) as

$$Z = \frac{1}{\omega} (g - i)K + i\omega M$$

Use Equation (14) for K and factor M out on both sides and the equation becomes

$$Z = M \left\{ \frac{1}{\omega} (g - i) \phi \begin{bmatrix} \Omega_i^2 \\ \frac{1}{m_i} \end{bmatrix} \phi^T + i\omega M^{-1} \right\} M$$

From (13), one can write

$$M = \phi^{-T} \begin{bmatrix} 1 \\ m_i \end{bmatrix} \phi^{-1}$$

or

$$M^{-1} = \phi \begin{bmatrix} 1 \\ m_i \end{bmatrix} \phi^T$$

Substituting into the above equation for Z and factoring the ϕ and ϕ^T matrices, there results

$$Z = M \phi \left\{ \frac{1}{\omega} (g - i) \begin{bmatrix} \Omega_i^2 \\ \frac{1}{m_i} \end{bmatrix} + i \begin{bmatrix} 1 \\ m_i \end{bmatrix} \right\} \phi^T M$$

or

$$Z = \frac{1}{\omega} M\phi \left[\frac{g\Omega_i^2 + i(\omega^2 - \Omega_i^2)}{m_i} \right] \phi^T M \quad (20)$$

and in summation form

$$Z = \frac{1}{\omega} \sum_{i=1}^p \frac{\Omega_i^2 \{ g + i[(\frac{\omega}{\Omega_i})^2 - 1] \}}{m_i} M\phi_i \phi_i^T M \quad (21)$$

The mobility can be written as the inverse of the impedance, from Equation (20):

$$\begin{aligned} Y = Z^{-1} &= \omega \{ M\phi \left[\frac{g\Omega_i^2 + i(\omega^2 - \Omega_i^2)}{m_i} \right] \phi^T M \}^{-1} \\ &= \omega (\phi^T M)^{-1} \left[\frac{m_i}{g\Omega_i^2 + i(\omega^2 - \Omega_i^2)} \right] (M\phi)^{-1} \end{aligned}$$

but from (13)

$$(\phi^T M)^{-1} = \{ [m_i] \phi^{-1} \}^{-1} = \phi \left[\frac{1}{m_i} \right]$$

and

$$(M\phi)^{-1} = \{ \phi^{-T} [m_i] \}^{-1} = \left[\frac{1}{m_i} \right] \phi^T$$

then

$$Y = \omega \phi \left[\frac{1}{m_i [g\Omega_i^2 + i(\omega^2 - \Omega_i^2)]} \right] \phi^T \quad (22)$$

or in summation form

or

$$Y = \omega \sum_{i=1}^p \frac{g - i \left[\left(\frac{\omega}{\Omega_i} \right)^2 - 1 \right]}{i^2 \left[\left(\frac{\omega}{\Omega_i} \right)^2 - 1 \right]^2 + g^2} \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T \quad (23)$$

III. THE INCOMPLETE MODEL

The structure under consideration is represented analytically by g, M, K (see Equation (1)). Consider, for the time being, that g, M are known. It was seen that K and its inverse, C , could each be written as a sum of matrices containing the eigenvectors (Equations (16), (17)).

$$K = \sum_{i=1}^P \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M \quad (24)$$

$$C = K^{-1} = \sum_{i=1}^P \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T \quad (25)$$

and the impedance and mobility could be written in similar forms (Equations (21), (23))

$$Z = \frac{1}{\omega} \sum_{i=1}^P \frac{\Omega_i^2}{m_i} \{g - i[(\frac{\omega}{\Omega_i})^2 - 1]\} M \phi_i \phi_i^T M \quad (26)$$

$$Y = Z^{-1} = \omega \sum_{i=1}^P \frac{1}{\Omega_i^2 m_i} \frac{g - i[(\frac{\omega}{\Omega_i})^2 - 1]}{[(\frac{\omega}{\Omega_i})^2 - 1]^2 + g^2} \phi_i \phi_i^T \quad (27)$$

All the square matrices (K, C, Z, Y, M) are of order P , the number of degrees of freedom, and the ϕ vectors have P elements. The square matrices $\phi_i \phi_i^T$ are each $P \times P$ but

are of rank 1 and are thus individually singular. Since the ϕ_i 's are linearly independent, a linear combination of P of these simple products will be of rank P and thus the summations will be nonsingular (See Reference 4, p.6).

If g , M and all the normal modes of the system were known, the behavior of the system could be predicted by forming the complete equations of motion. The question posed is this: If incomplete information is available, i.e. only the first N normal modes, is it possible to generate a mathematical model that will give useful information about the behavior of the system and modifications of it?

The following postulate is advanced: The K, C, Z, Y matrices obtained by using less than P terms can be used to predict the behavior of the structure and certain modifications of it.

Part of the purpose of this report is to test aspects of the above postulate. The analytical model described by the incomplete summations will be called an incomplete model.

These matrices will be written

$$K_{inc} = \sum_{i=1}^N \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M \quad (28)$$

$$C_{inc} = \sum_{i=1}^N \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T \quad (29)$$

$$Z_{inc} = \frac{1}{\omega} \sum_{i=1}^N \frac{\Omega_i^2}{m_i} \{g + i[(\frac{\omega}{\Omega_i})^2 - 1]\} M \phi_i \phi_i^T M \quad (30)$$

$$Y_{inc} = \omega \sum_{i=1}^N \frac{1}{\Omega_i^2 m_i} \frac{g - i[(\frac{\omega}{\Omega_i})^2 - 1]}{[(\frac{\omega}{\Omega_i})^2 - 1]^2 + g^2} \phi_i \phi_i^T \quad (31)$$

Certain characteristics of the incomplete model are apparent:

(1) Since the terms containing the higher values of Ω_i are not included, the dominant terms of K and Z will be missing and thus K_{inc} and Z_{inc} will not resemble the true K and Z matrices.

(2) Conversely, the dominant terms of C and Y are included in C_{inc} and Y_{inc} . These are the matrices which represent the responses due to applied forces and for the model to have validity, it is necessary that they approach the true values for $\omega < \Omega_N$.

(3) The four matrices are of order P (and represent the P points of interest) but are of rank N . Thus, they are all singular and they must be all formed separately and not by inversion.

(4) The eigenvalue equation from (28) or (29) can be seen to be

$$\begin{aligned}
 M^{-1}K_{inc}\phi_j &= \sum_{i=1}^N \frac{\Omega_i^2}{m_i} \phi_i \phi_i^T M \phi_j = \Omega_j^2 \phi_j & j = 1, 2, \dots, N \\
 &= 0 & j > N
 \end{aligned}$$

and similarly for CM. Thus it may be said that the incomplete model contains only the first N modes of the corresponding complete model.

IV. IDENTIFICATION OF THE MASS MATRIX

The previous discussion has assumed knowledge of the normal modes. It is apparent that knowledge of the mass matrix is also required. Before proceeding further with the discussion of the incomplete model, a method will be described for identifying the mass matrix. As above, it is assumed that the first N natural frequencies and normal modes have been determined through testing. Each of these modes contains P elements representing the relative motion of all the points of interest.

The normal modes are orthogonal with respect to the mass as given in Equation (9).

$$\phi_i^T M \phi_j = 0 \quad j \neq i \quad (32)$$

Writing out the indicated multiplications in terms of the individual elements, this equation becomes

$$\sum_{k=1}^P \phi_{ki} \phi_{kj} m_{kk} + \sum_{k=1}^{p-1} \sum_{n=k+1}^p (\phi_{ki} \phi_{nj} + \phi_{ni} \phi_{kj}) m_{kn} = 0$$

$$i = 1, 2, \dots, N-1$$

$$j = i+1, \dots, N \quad (33)$$

where ϕ_{ki} indicates the kth element of ϕ_i and m_{kn} is an element of the M matrix. The mass matrix has been assumed symmetrical but not necessarily diagonal. Equation (33) is,

in reality, $N(N-1)/2$ linear equations having the mass elements as unknowns with products of the elements of the known normal modes as coefficients.

It is possible that certain mass elements may be known to be zero or to have some definite value. If they are to be zero, the corresponding terms are dropped from the equation. If they are to be restricted to a particular value, the corresponding terms are placed on the right-hand side of the equation.

In addition, if any of the generalized modal masses, m_i , are known, Equation (9) for $j=i$ can be used.

$$\phi_i^T M \phi_i = m_i \quad (34)$$

or

$$\sum_{k=1}^p \phi_{ki}^2 m_{kk} + 2 \sum_{k=1}^{p-1} \sum_{n=k+1}^p \phi_{ki} \phi_{ni} m_{kn} = m_i \quad (35)$$

There is another possible known condition. The total of the diagonal elements may be considered to be known (the total mass of the structure, for example). This leads to the equation

$$\sum_{k=1}^p m_{kk} = m_T \quad (36)$$

Equations (33), (35) and (36) may be written as

$$A\bar{m} = R \quad (37)$$

where \bar{m} is a column matrix made up of the unknown elements of M, A is a matrix formed by the coefficients of these unknowns from Equation (33), and, if used, from (35) and (36). R is a column matrix made up of the right-hand side terms corresponding to known masses, if any, known generalized masses, if any, and possibly the known total mass.

There are, then, at least, $N(N-1)/2$ equations and possibly as many as $N(N+1)/2+1$ if all the possible equations are used. A typical situation would use Equations (33) and (36) resulting in $N(N-1)/2+1$ equations. If none of the mass elements are considered to be known (other than zero), there are at least P unknown diagonal masses and as many as $P(P+1)/2$ if the matrix is taken to be completely filled. In many situations, it will be desired to obtain a purely diagonal mass matrix. In a typical situation, there will be, say, $N(N-1)/2+1$ equations in P unknowns.

When the number of equations is less than the number of unknowns, there are an infinite number of solutions. When the reverse is true, there will ordinarily be no solution.

This treatment will be limited to the first situation where there are an infinite number of solutions to the equations. This is not a severe restriction since the class of problems considered includes a relatively large number of points of interest and a relatively small number of known normal modes. If these were nearly equal, it may have been possible with a small amount of additional testing to have completed the information requirements for other techniques (Reference 1). A typical situation might include 4 normal modes plus the total mass giving 7 equations and thus requiring more than seven unknown masses. Similarly 6 normal modes would give 16 equations and 10 normal modes would give 46 equations. As more normal modes become known for the same number of points of interest, it may be necessary to introduce off-diagonal masses as unknowns.

Return to the consideration of Equation (37) where A is an $n_e \times n_v$ matrix (n_e being the number of equations and n_v the number of variables) and $n_e < n_v$. \bar{m} is $n_v \times 1$ and R is $n_e \times 1$. The equations have an infinite number of solutions, that is, there are an infinite number of mass distributions which will cause the modes to be orthogonal. In fact, it is quite possible for there to be an infinite number of

mass distributions which will also give the same generalized masses, m_i , and thus result in the same C_{inc} and Y_{inc} (see Equations (29),(31)). In other words, it is quite possible that different valid mass distributions used with the measured normal modes will predict identical responses of the system to sinusoidal forcing.

If, however, it is desired to use the model for making predictions under other conditions, especially to predict the effects of changes in parameters, then it is apparent that the masses used in the model should be as near to the "true" values as possible. The best information available as to what the "true" values are, is the approximation arrived at by the analyst. These analytical values will not, in general, satisfy the orthogonality condition of the normal modes, i.e. Equation (37).

The pseudo-inverse (see Appendix II) is an elegant mathematical tool which can be used to obtain the solution to Equation (37) which is the closest (in a least squares sense) to any specified analytical approximation. Another way of saying this is that the smallest possible changes in the approximation can be found so as to satisfy the conditions of orthogonality.

Define a column matrix, m_A , which is the approximation to \bar{m} and subtract Am_A from both sides of (37) giving

$$A(\bar{m} - m_A) = R - Am_A \quad (38)$$

At this point, a weighting function is introduced in the form of a diagonal matrix, W . Each element is a measure of the analysts confidence in the corresponding approximation. The result will be that masses having higher values of weighting functions will tend to vary least. Inserting the identity $W^{-1}W$ into the above equation results in

$$(AW^{-1})\{W(\bar{m} - m_A)\} = R - Am_A \quad (39)$$

Defining $(AW^{-1})^+$ as the pseudo-inverse of AW^{-1} , the solution given by

$$W(\bar{m} - m_A) = (AW^{-1})^+\{R - Am_A\} \quad (40)$$

is the one of the infinite number possible having the smallest weighted sum of squares of the differences of \bar{m} and m_A . Using ordinary matrix algebra from this point, there results

$$\bar{m} - m_A = W^{-1}(AW^{-1})^+\{R - Am_A\}$$

$$\bar{m} = W^{-1}(AW^{-1})^+R + \{I - W^{-1}(AW^{-1})^+A\}m_A$$

It is shown in Appendix II that

$$(AW^{-1})^+ = W^{-1}A^T\{A(W^{-1})^2A^T\}^{-1}$$

Then defining

$$B_R = W^{-1}(AW^{-1})^+ = (W^{-1})^2A^T\{A(W^{-1})^2A^T\}^{-1}$$

$$B_m = I - B_R A$$

the equation for the mass elements is

$$\bar{m} = B_R R + B_m m_A \quad (41)$$

V. MASS CHANGES

One of the criteria for evaluating the usefulness of an analytical model is its ability to predict the effects of changes. In this section, the use of an incomplete model to predict the effects of mass changes on the natural frequencies and normal modes is discussed.

The stiffness matrix and the influence coefficient matrix are independent of the mass of the system. The expression derived in terms of the normal modes do contain the mass, however. (See Equations (16), (17)). Thus, it must be concluded that when the mass is changed, the normal modes and frequencies must change in such a way that the summations remain invariant. In other words

$$K = \sum_{i=1}^p \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M = \sum_{i=1}^p \frac{(\Omega_i + \Delta\Omega_i)^2}{(m_i + \Delta m_i)} (M + \Delta M) (\phi_i + \Delta\phi_i) (\phi_i + \Delta\phi_i)^T (M + \Delta M) \quad (42)$$

$$C = \sum_{i=1}^p \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T = \sum_{i=1}^p \frac{1}{(\Omega_i + \Delta\Omega_i)^2 (m_i + \Delta m_i)} (\phi_i + \Delta\phi_i) (\phi_i + \Delta\phi_i)^T \quad (43)$$

where the Δ 's indicate the changes due to the change in mass, ΔM .

While the above expressions must be true when summed over all the modes, they will not be exact for incomplete summations, i.e. for the incomplete model. Of the two (K and C), it is to be expected that C_{inc} will be less sensitive to mass changes. The reason is that the dominant terms are included in C_{inc} and omitted in K_{inc} , thus, C_{inc} is much closer to the invariant matrix C than K_{inc} is to the invariant K .

This hypothesis has been tested by calculating the frequencies and modes of a modified system using matrix iteration on $C_{inc}(M + \Delta M)$. The changes predicted were in excellent agreement with the true values. These results are given in Section XI.

VI. STIFFNESS CHANGES

The effect of a change in the stiffness matrix cannot be handled as directly as the mass change discussed in the previous section. Both the K and C matrices must change when the stiffness is changed. Since the dominant terms of K are missing in K_{inc} it does not appear to be reasonable to hypothesize that $(K + \Delta K)_{inc} = K_{inc} + \Delta K$ since even small ΔK 's can easily be greater by orders of magnitude than the elements of K_{inc} .

As discussed in Appendix I (and also derived previously, Equation (16)), it is possible to write K in the following form

$$K = \sum_{i=1}^p \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M \quad (44)$$

if and only if the ϕ_i 's are eigenvectors of $M^{-1}K$. Otherwise there must be coupling terms of the form $a_{ij} M \phi_i \phi_j^T M$. Thus, if $K + \Delta K$ is expressed in terms of the eigenvectors of $M^{-1}K$, the equation must be of the form

$$K + \Delta K = \sum_{i=1}^p \sum_{j=1}^p a_{ij} M \phi_i \phi_j^T M \quad (45)$$

Now, pre- and post-multiply this equation by ϕ_k^T, ϕ_n

$$\phi_k^T (K + \Delta K) \phi_n = \sum_{i=1}^p \sum_{j=1}^p a_{ij} (\phi_k^T M \phi_i) (\phi_j^T M \phi_n) = a_{kn} m_k m_n \quad (46)$$

because of the orthogonality of the ϕ 's. And since $K\phi_n$

$$= \Omega_n^2 M\phi_n$$

$$\phi_k^T (K + \Delta K) \phi_n \equiv \phi_k^T K \phi_n + \phi_k^T \Delta K \phi_n \equiv \Omega_n^2 \phi_k^T M \phi_n + \phi_k^T \Delta K \phi_n = a_{kn} m_k m_n$$

or

$$a_{ij} = \frac{\Omega_j^2 \phi_i^T M \phi_j + \phi_i^T \Delta K \phi_j}{m_i m_j} = \frac{\Omega_i^2}{m_i} + \frac{\phi_i^T \Delta K \phi_i}{m_i^2} \quad j = i$$

$$= \frac{\phi_i^T \Delta K \phi_j}{m_i m_j} \quad j \neq i \quad (47)$$

Thus, substituting into (45), the expression for $K + \Delta K$ can be written

$$K + \Delta K = \sum_{i=1}^P \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M + \sum_{i=1}^P \sum_{j=1}^P \left(\frac{\phi_i^T \Delta K \phi_j}{m_i m_j} \right) M \phi_i \phi_j^T M \quad (48)$$

This expression when summed over all P modes is exact.

Note that the first summation is equal to K.

Now, truncating the series at the last known mode, the expression can be written

$$(K + \Delta K)_{inc} = K_{inc} + \sum_{i=1}^N \sum_{j=1}^N \frac{\phi_i^T \Delta K \phi_j}{m_i m_j} M \phi_i \phi_j^T M \quad (49)$$

This expression can now be evaluated and it is hypothesized that $M^{-1}(K + \Delta K)_{inc}$ can be used to obtain good approximations to the new natural frequencies and normal modes. This hypothesis has been tested and the results given later appear quite satisfactory.

VII. IDENTIFICATION OF DAMPING COEFFICIENT

If the identified model is required to predict responses, in addition to predicting the effects of changes on the natural frequencies, it is necessary to identify the damping coefficient. The responses to sinusoidal loading is, in effect, the mobility (See Equation (23)). The real component of the velocity response at the driving point is given by:

$$\dot{y}_R = \omega \sum_{i=1}^p \frac{g}{\left[\left(\frac{\omega}{\Omega_i} \right)^2 - 1 \right]^2 + g^2} \frac{1}{\Omega_i^2 m_i} \quad (50)$$

where the modes are normalized at this point. When the modes are reasonably well separated, \dot{y}_R at a resonance is nearly completely dependent on the term due to the one mode.

Thus

$$\dot{y}_R(\Omega_i) = \frac{1}{g \Omega_i m_i} \quad (51)$$

can be used to compute g when $\dot{y}_R(\Omega_k)$ has been measured.

If the modes are not separated, it is possible to use the incomplete version of (50) which also contains only the one unknown, g .

$$\dot{Y}_R(\Omega_j) = \Omega_j \sum_{i=1}^N \frac{g}{\left[\left(\frac{\Omega_j}{\Omega_i} \right)^2 - 1 \right]^2 + g^2} \frac{1}{\Omega_i^2 m_i} \quad (52)$$

Using this procedure, g's may be obtained at each resonance. Experiments discussed in following sections have shown fine consistency among these values and good agreement with the exact value.

Having obtained g and knowing the mass matrix, the natural frequencies, Ω_i , and the normal modes, ϕ_i , for $i = 1, N$ the dominant portion of the mobility matrix may be constructed at any frequency up to the vicinity of the N^{th} natural frequency. It is expected that this information allows the prediction of steady-state responses although this has not specifically been tested in the work reported in this report.

VIII. DESCRIPTION OF COMPUTER PROGRAM

In order to subject some of the hypotheses presented to a qualitative evaluation, a computer program has been developed. The program performs three functions: (1) it simulates a test; (2) it identifies the mass; (3) it finds the modes and frequencies after mass and stiffness changes. The program logic is outlined in Figure 1 and the program listing is given in Appendix III. Below is a general description of the program.

Part I - Simulated Test

In order to simulate a test, the program first accepts a complete description of the "actual system" consisting of a mass matrix, M , a stiffness matrix, K , (or optionally, an influence coefficient matrix, C) and a scalar damping coefficient, g . The "actual" number of degrees of freedom (the order of the matrices) is limited to 20.

The test to be simulated consists of measurements of the resonant response at P points on the structure due to sinusoidal excitation at one point. Thus, the data required for the simulation of the tests consists of the driving point, the points at which measurements are taken, the frequencies of excitation, and the measurement error.

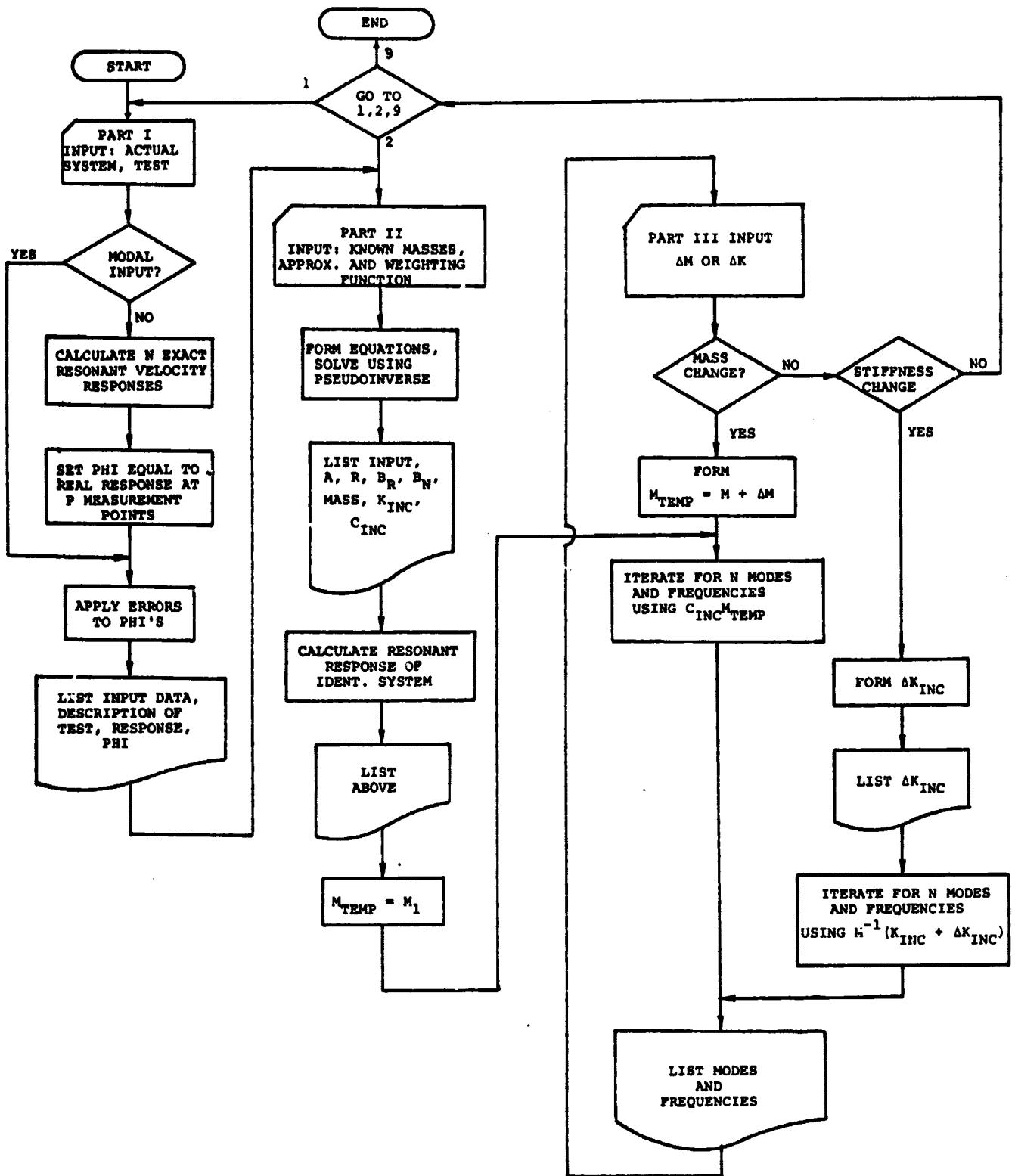


FIGURE 1. LOGIC DIAGRAM OF COMPUTER PROGRAM.

The frequencies of excitation, which should be the N resonant frequencies, are assumed to have been computed outside this program. The measurement error is taken to be a fixed bias error plus a uniform random distribution, each expressed as a percentage of the exact measurement. The procedure, then, is as follows: (1) calculate the impedance,

$\frac{q}{\omega} K + i(\omega M - \frac{1}{\omega} K)$, of the "actual" system at each of the N resonant frequencies; (2) invert these complex matrices to obtain the exact mobilities; (3) store the elements corresponding to the P measurement points of the column corresponding to the driving point; (4) apply the specified measurement errors to these values.

This data will be a good simulation of data which would be taken in an actual test. The test data required in the following parts of the program are the normal modes of the system. It is known that a column of the real component of the mobility near resonance will be a good approximation to the normal mode when the natural frequencies are separated. This can be seen from Equation (23). It is this information, then, which is passed on to Part II of the program as the N measured normal modes, each consisting of P points.

As an option to the above procedure, the program will accept as input the normal modes themselves. This procedure would simulate the use of a more sophisticated method of computing them. These modes are then polluted with errors as before, prior to being transferred to Part II of the program.

Part II - Mass Identification

The procedures given in Section IV of this report are carried out numerically. The input consists of any known generalized masses or mass elements and the masses which are to be treated as unknowns. For each unknown mass, an approximation and a weighting function is supplied. In addition, one may specify that the total of the diagonal masses must remain constant. Based on the input, the A matrix is formed as are A_{m_A} and W. The solution for the "best" mass distribution is obtained as given in Equation (41).

At this point, the identified mass matrix is formed, the generalized mass matrix $(\phi^T M \phi)$ is computed as are the incomplete stiffness and influence coefficient matrices (Equations (28), (29)). As a check, the resonant response and the natural frequencies and normal modes of the identified incomplete model are computed.

Part III - Mass or Stiffness Changes

There now is sufficient data to compute the changes in frequencies and modes due to mass or stiffness changes. The program will accept changes to either the mass matrix or the stiffness matrix. For mass changes the matrix $M + \Delta M$ is formed, where M is the identified mass matrix. The first N frequencies and modes are obtained by iteration on $C_{inc}(M + \Delta M)$ as discussed in Section V.

For stiffness changes, $(K + \Delta K)_{inc}$ is formed as given in Equation (37). The N frequencies and modes are computed by iteration from $M^{-1}(K + \Delta K)_{inc}$.

The program is written so that changes can be made on the identified model in sequence without rerunning the first two parts of the program.

IX. THE SIMULATED TESTS

In order to test the hypotheses presented regarding the characteristics of the incomplete model, it was necessary to select a structure on which to perform the simulated testing. It was decided not to select so simple a system that success would give no confidence that the methods would work in practice. On the other hand, it was not desirable to make the first tests using such a complex system that one would not have a good intuitive understanding of the system and modifications of it.

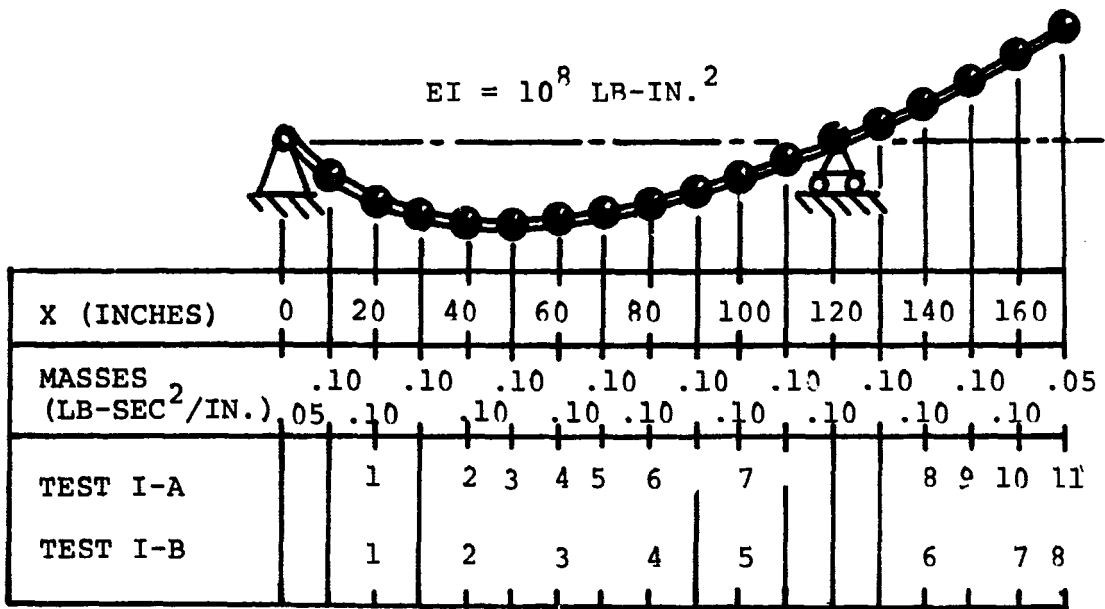
As a compromise, a structure was selected whose behavior is quite well understood yet is not so simple as to yield trivial results. The structure selected is a simple, thin, beam of constant EI having 18 lumped masses arranged so as to approximate a uniform mass distribution. The beam is allowed to deflect transversely only. The deflection is constrained at one end and at a point approximately 70 percent of the length of the beam, leaving an overhang of approximately 30 percent. A structural damping coefficient of .02 was used.

Most of the simulated testing was performed using this system. This system is designated Specimen I. Some tests were also carried out on a structure identical with the above except that two rotary inertias were added, one at the second support and one at the free tip of the beam. This system is called Specimen II. These beams are illustrated schematically in Figure 2. The natural frequencies and mode shapes of the two beams were computed using a standard computer program. These results are illustrated in Figure 3.

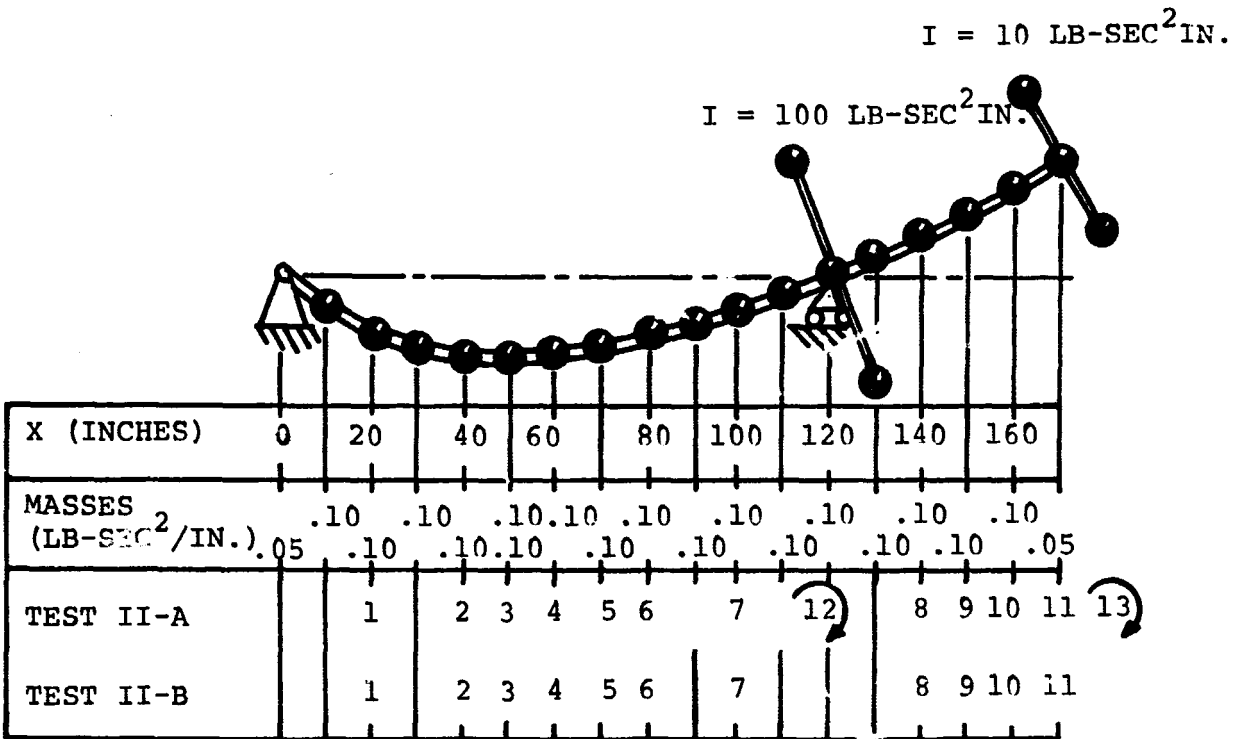
For each of the two specimens, two arrangements of test measurements were selected, referred to as tests A and B. These are also illustrated in Figure 2. In each test, the system was considered to be driven at station 60.

Associated with each specimen and each test is an approximate mass distribution. For Specimen I, test A, there are three approximate mass distributions (I, II, III). Each of these is somewhat reasonable. Table I lists the various mass approximations used. These are discussed further in Section X.

The simulated tests which were run are summarized and described in Table II.



SPECIMEN I



NOTE: INDICATES SLOPE MEASUREMENT

SPECIMEN II

Figure 2. Schematic Representation of the Simulated Test Specimens.

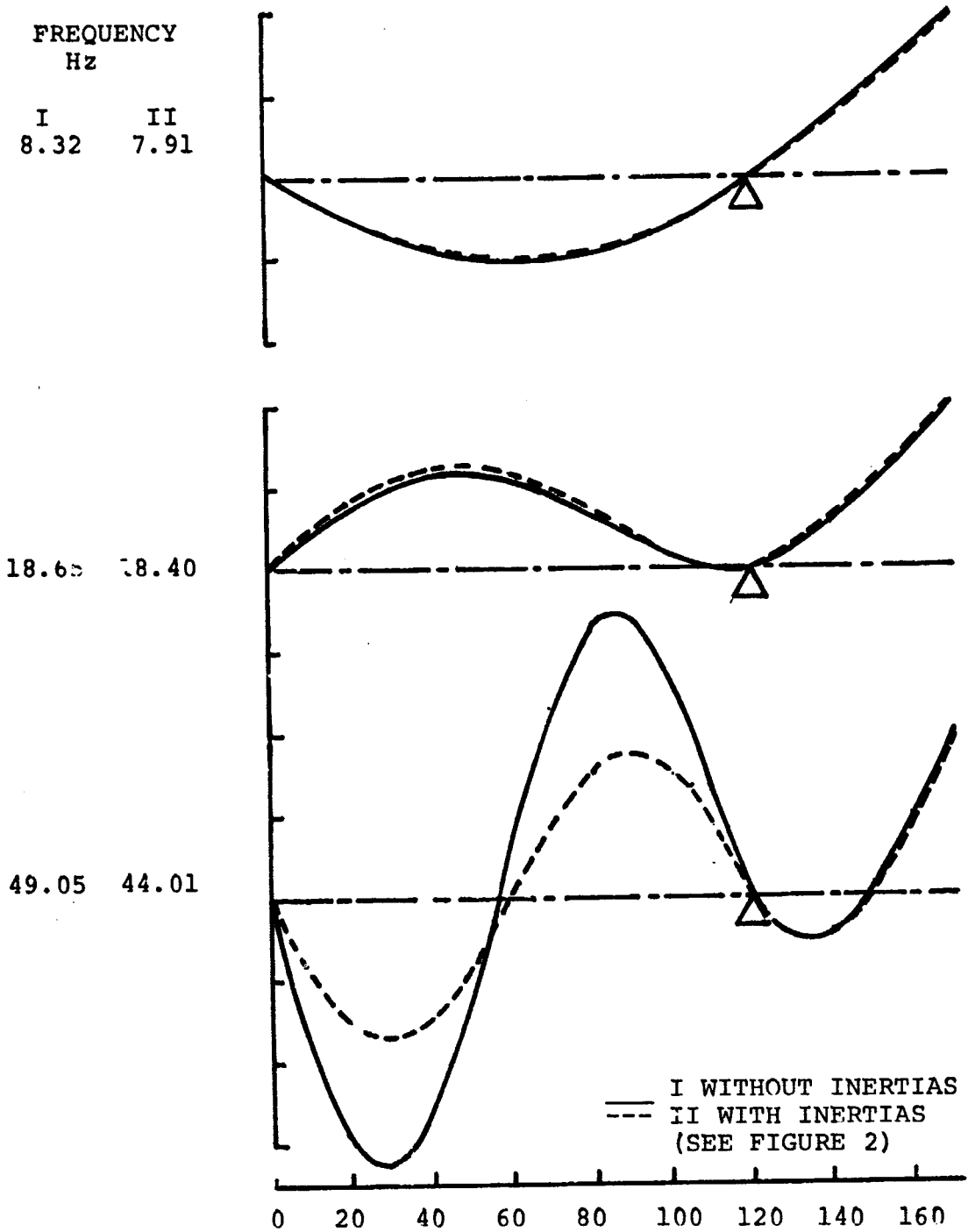


Figure 3. Exact Normal Modes of Test Specimens.

TABLE I. APPROXIMATE MASS DISTRIBUTIONS

Sta. In.	Specimen I						Specimen II		
	True Mass	Test A			W*	Test B	True Mass	Test A	Test B
		I	II	III					
0.Δ	.05						.05		
10.	.10						.10		
20.	.10	.30	.25	.25	0.1	.25	.10	.30	.30
30.	.10						.10		
40.	.10	.15	.15	.13	1.0	.20	.10	.15	.15
50.	.10	.10	.10	.12	0.5		.10	.10	.10
60.	.10	.10	.10	.10	1.0	.20	.10	.10	.10
70.	.10	.10	.10	.12	0.5		.10	.10	.10
80.	.10	.15	.15	.13	1.0	.20	.10	.15	.15
90.	.10						.10		
100.	.10	.30	.25	.25	0.2	.25	.10	.30	.30
110.	.10						.10		
120.Δ	.10						.10		
130.	.10						.10		
140.	.10	.25	.20	.15	0.5	.25	.10	.25	.25
150.	.10	.10	.10	.13	1.0		.10	.10	.10
160.	.10	.10	.10	.12	1.0	.15	.10	.10	.10
170.	.05	.05	.05	.05	1.0	.05	.05	.05	.05
I@120.	0						100.	100.	0
								(W=.01*)	
I@170.	0						10.	10.	0
								(W=0.1*)	

Δ Support points

* W = 1.0 except for test 6 when it was used as shown and for tests 18-20 when it was used to minimize the dominance of the large magnitude of the inertias.

TABLE II. DESCRIPTION OF SIMULATED TESTS

Test No.	Specimen (See Fig. 2)	Test	No. of Resonances or Modes	Mass Approx. (Table I)	Remarks*
1	I	A	4R**	I	
2	I	A	4M	I	
3	I	A	3R	I	
4	I	A	3M	I	
5	I	A	3R	I	Variable mass Weighting fun. (See Table I)
6	I	A	3R	I	
7	I	A	3R	II	
8	I	A	4R	III	
9	I	A	3R	III	Variable mass
10	I	A	3R	III	
11	I	A	3R	III	5% Bias and +5% Random Measurement Error
12	I	A	3R	III	
13	I	A	3R	III	
14	I	A	3R	III	
15	I	A	3R	III	
16	I	B	3R		Variable mass
17	I	B	3R		
18	II	A	4R		Inertias weighted (See Table I)
19	II	A	3R		
20	II	A	2R		
21	II	B	4R		
22	II	B	3R		

* Unless so indicated, total mass held constant, no weighting, and no errors assumed.

**R indicates resonant data used, M indicates actual mode shapes used.

X. THE IDENTIFIED MASSES

The masses identified by the program based on the simulated tests are given in Table III. The data is listed by test number. Refer to Table II for a description of the tests. The masses are given in the table to two significant figures for ease of reading except for tests 10-15 in order that the scatter may be observed. These are all identical except that tests 11-15 contain test measurement errors of 5 percent bias and +5 percent random on amplitude. In addition to the masses themselves, the table indicates the changes from the approximation in terms of the rms of the variance.

Certain characteristics of this data can be observed. It appears that "better" approximations have smaller variances associated with them. This is what would be expected since the better the approximation, the smaller the change required to make it satisfy the equations. In the cases of Specimen I, test A, approximations II and III are intuitively better than I since I includes the masses at the supports. It is apparent, by comparing similar conditions from tests 1-6 with 7-15 that the variances are significantly lower for approximations II and III. Comparison of 16, 17 with 9, 10 which use effectively the

TABLE III. IDENTIFIED MASSES

Test ⁺ No.	Point of Measurement (See Figure 1)											Var. RMS
	1	2	3	4	5	6	7	8	9	10	11	
Specimen I (See Figure 1)												
AI*	.30	.15	.10	.10	.10	.15	.30	.25	.10	.10	.05	
1	.28	.16	.12	.11	.12	.17	.26	.23	.088	.093	.074	.019
2	.25	.20	.12	.08	.11	.22	.21	.25	.11	.078	.092	.043
3	.29	.15	.11	.12	.12	.15	.29	.23	.077	.092	.076	.016
4	.29	.15	.11	.12	.12	.15	.29	.23	.077	.092	.075	.016
5	.29	.15	.12	.12	.12	.16	.29	.24	.081	.095	.077	.016
6	.27	.15	.13	.11	.13	.15	.25	.24	.097	.10	.066	.023
AII*	.25	.15	.10	.10	.10	.15	.25	.20	.10	.10	.05	
7	.25	.15	.11	.11	.11	.15	.24	.19	.091	.097	.061	.0065
AIII*	.25	.13	.12	.10	.12	.13	.25	.15	.13	.12	.05	
8	.25	.15	.12	.085	.12	.15	.26	.15	.11	.099	.059	.014
9	.25	.13	.13	.11	.12	.13	.25	.15	.12	.12	.047	.0040
10	.250	.134	.125	.106	.124	.131	.248	.146	.124	.115	.047	.0040
11	.256	.142	.130	.104	.117	.123	.245	.147	.125	.115	.044	.0066
12	.255	.138	.125	.101	.118	.126	.249	.150	.129	.116	.043	.0044
13	.251	.138	.130	.109	.125	.130	.244	.144	.120	.111	.048	.0069
14	.252	.138	.130	.109	.124	.126	.243	.143	.120	.113	.052	.0070
15	.257	.142	.127	.100	.114	.122	.247	.150	.129	.117	.044	.0059
B*	.25	.20	.20	.20	.25	.25	.15	.05				
16	.24	.20	.22	.21	.24	.24	.14	.055				.010
17	.24	.20	.22	.21	.24	.24	.14	.055				.010
Specimen II (See Figure 1)												
A**	.30	.15	.10	.10	.10	.15	.30	.25	.10	.10	.05	
18	.24	.16	.14	.12	.089	.14	.35	.26	.10	.071	.080	.034
19	.30	.15	.11	.11	.11	.16	.30	.24	.090	.098	.067	.015
20	.30	.15	.10	.10	.10	.15	.30	.25	.10	.10	.051	.0005
B*	.30	.15	.10	.10	.10	.15	.30	.25	.10	.10	.05	
21	.73	-.36	-.08	.41	.27	-.33	.27	1.11	-.31	-.19	.18	.409
22	.29	.12	.068	.082	.10	.17	.33	.28	.13	.11	.027	.023

* Mass approximation used for tests following (See Table 1)

** Inertias not tabulated. Identified inertias were virtually exact.

+ See Table II for description of test.

same distribution suggests that the smaller the number of points, the greater the variance will be. One reason for this is that the approximation having the greater number of masses is an intuitively better representation of the system, everything else being equal.

There is another, more mathematical reason for this effect. There are fewer equations than there are unknown masses (See Section IV). If these numbers were the same, the mass would be unique and have no necessary relation to the approximation. As the number of unknowns is increased, the set of solutions expands and the probability of finding one of these solutions closer to a given approximation increases. The same effect is true when the number of unknowns stays the same but the number of equations is decreased, as when the number of modes used is reduced. Notice tests 18, 19, 20, for example, where the variances associated with 4, 3, 2 modes are .034, .015, .0005 respectively.

Thus, it could be reasoned that, the more nearly complete the model (for a constant number of masses), the greater will be the variances in the identified masses from the approximation. The data presented here tends to bear out this argument.

The effect of the constraint on the sum of the masses has little effect as can be seen by comparing tests 3 and 5, 9 and 10, 16 and 17.

The use of exact modes instead of the resonant responses, in these examples, shows no significant improvement in terms of the variances. It is interesting to note that test 1 using four resonant responses has a variance of .019 while test 2 which was identical except that true normal modes were used has a variance of .043. In tests 3, 4, which are identical to 1,2 except that 3 modes were used, the results are virtually identical to each other.

For Specimen II, test B, an intentionally poor test and approximation was made by ignoring the large inertias and omitting the measurements of slope at these points. Test 21 using 4 resonances results in very poor masses including several negative values. However, when only three resonances were used, the identified masses show general agreement with the approximation. In actual testing, results such as from test 21 could point up the omission of a significant parameter in the analysis.

The effects of the measurement errors is rather small. In five simulated identical tests (11-15), errors of 5 percent bias and 5 percent random were applied to the response amplitudes. This is identical to a uniform random distribution between 0 and +10 percent. In each case, the variances are small. None of the individual mass elements vary by more than 8 percent over the five tests.

In general, the masses identified appear to be acceptable approximations (except in test 21). The test as to whether they are satisfactory or not will come when they are used to predict the effects of structural changes.

XI. THE EFFECTS OF MASS CHANGES

Two different lumped masses were considered to be added to the structure and the new frequencies and modes were calculated for the tests described in Table II. The method used has been described in Section V and its implementation in Section VIII.

The lumped masses were added at a point midway between the supports (change "A") and at the free end of the beam (change "B"). These masses were $1.0 \text{ lb-sec}^2/\text{in}$. This represents an increase in total mass of the beam of almost 60 percent.

Table IV summarizes the results of the frequency calculations. The table shows the frequencies at which the simulated testing was conducted (the frequencies of the original beam), the exact frequencies of the modified beams, and the predicted results for each of the simulated tests.

In general, the results are surprisingly good and relatively insensitive to the quantities that were varied. For Specimen I, change B (the more extreme of the effects), the first mode changes frequency by 54 percent and the poorest prediction is within 8 percent of the correct value.

TABLE IV. COMPUTED FREQUENCIES DUE TO MASS CHANGE

Mode	1		2		3		4	
f*	8.32		18.65		49.06		96.12	
+	A	B	A	B	A	B	A	B
f**	6.05	3.82	15.46	14.77	48.37	47.06	79.53	78.81
1	6.30	4.09	15.67	14.85	48.46	47.31	84.65	85.17
2	6.33	4.12	15.69	14.85	48.49	47.31	83.81	84.43
3	6.29	4.08	15.69	14.87	48.52	47.48		
4	6.29	4.08	15.68	14.87	48.55	47.48		
5	6.33	4.12	15.71	14.88	48.53	47.48		
6	6.32	4.11	15.70	14.88	48.50	47.41		
7	6.19	3.97	15.60	14.84	48.50	47.43		
8	6.20	3.98	15.59	14.82	48.44	47.28	83.48	84.22
9	6.21	3.99	15.62	14.84	48.48	47.37		
10	6.21	3.99	15.62	14.84	48.48	47.37		
11	6.23	3.83	15.51	14.96	48.47	47.46		
12	6.10	3.89	15.81	14.69	48.48	47.34		
13	6.26	4.03	15.54	14.73	48.46	47.42		
14	6.25	4.05	15.76	14.99	48.49	47.35		
15	6.12	3.91	15.71	15.00	48.45	47.29		
16	6.17	3.95	15.59	14.84	48.53	47.50		
17	6.18	3.95	15.60	14.84	48.53	47.50		
f*	7.91		18.40		44.01		65.94	
f**	5.91	3.78	14.92	14.58	44.00	40.23	60.33	62.72
18	6.14	4.02	15.15	14.68	44.00	40.40	61.42	62.82
19	6.19	4.01	15.15	14.69	44.00	40.69		
20	6.06	3.94	15.07	14.72				
21	5.01	2.92	14.30	14.37	43.99	40.36	63.46	64.48
22	5.88	3.73	14.91	14.56	43.99	40.01		

* Frequency before changes, i.e. frequency tested.
 ** Exact frequency after changes.
 + Addition of 1.0 lb-sec²/in. at Sta 60 ("A") or Sta 170 ("B").

For a more typical situation, Specimen I, change A, where the second mode changes frequency by about 17 percent, the poorest prediction is within about 2.5 percent.

The data bears out the expectation that the better mass approximations will result in better frequency predictions. The results of tests 7-10 (the better mass distributions) are seen to be slightly better than tests 1-6.

A better illustration of this effect is seen in test 21 which, as discussed in Section X, was given an intentionally bad mass distribution. Even though this model would predict the resonant responses and also the sinusoidal responses over the frequency range of interest (this has not been shown but is assumed to be true from theoretical considerations), its ability to predict the changes in frequency due to mass changes has been adversely affected by the poor approximation.

Test 22 which is the same as 21 except that only 3 resonances are used seems to contradict the above conclusions since the predicted frequencies are excellent. The fact that these predictions are better than those corresponding to better mass approximations is probably a coincidence. It is not surprising, however, that the 3 resonance data

gives better results than the 4 resonance data. The effect is discussed in Section X in connection with the mass identification. It may be surmised from this limited data that as the number of equations approaches the number of unknowns the solutions corresponding to poorer mass approximations will tend to deteriorate more rapidly.

The effect of measurement errors is seen to be slight from tests 11-15. The maximum scatter in predicted frequencies is about 2.5 percent even though the measurement error was randomly distributed between 0 to 10 percent.

The predicted normal modes behave in a manner similar to the predicted natural frequencies in that they are generally quite acceptable. Figure 4 illustrates the effect of change B on Specimen I for the first three modes. The figures show the original modes and the exact new modes. They are both normalized to be equal at station 60 which was the driving point in the simulated test. For modes 1 and 2, the results of tests 1-9 are all very close to the exact curve. The scatter obtained in tests 11-15 containing error compared to the exact curve and to test 10 (same conditions but without error) are quite satisfactory.

The results for the third mode show an interesting effect. The results based on data containing only 3 modes does not predict the proper shape as well as the tests containing 4 mode data. There is a slight tendency in this direction in the frequency data also (see tests 1, 2, 8, 18). The reason for this effect may be surmised. Assuming that the new modes contain components of the same and adjacent old modes, then one may expect that the predictions of the highest mode measured will be missing the significant effects of the next higher old mode.

Figure 4d illustrates the third mode of Specimen II, change B for the three tests: 18, 21, 22. Test 18 is the "good" approximation to Specimen II using 4 modes and shows good agreement with the exact mode. Tests 21, 22 represent the "bad" approximation for 4 and 3 modes, respectively. While both of these predict the frequency of this mode nearly equally well, the four mode data produces a considerably better third mode than the three mode data. This is in agreement with the effect previously noted.

While only sample modes are illustrated, these are quite typical of those obtained in all the computations made. The fourth modes are not shown because of the small amount of data available.

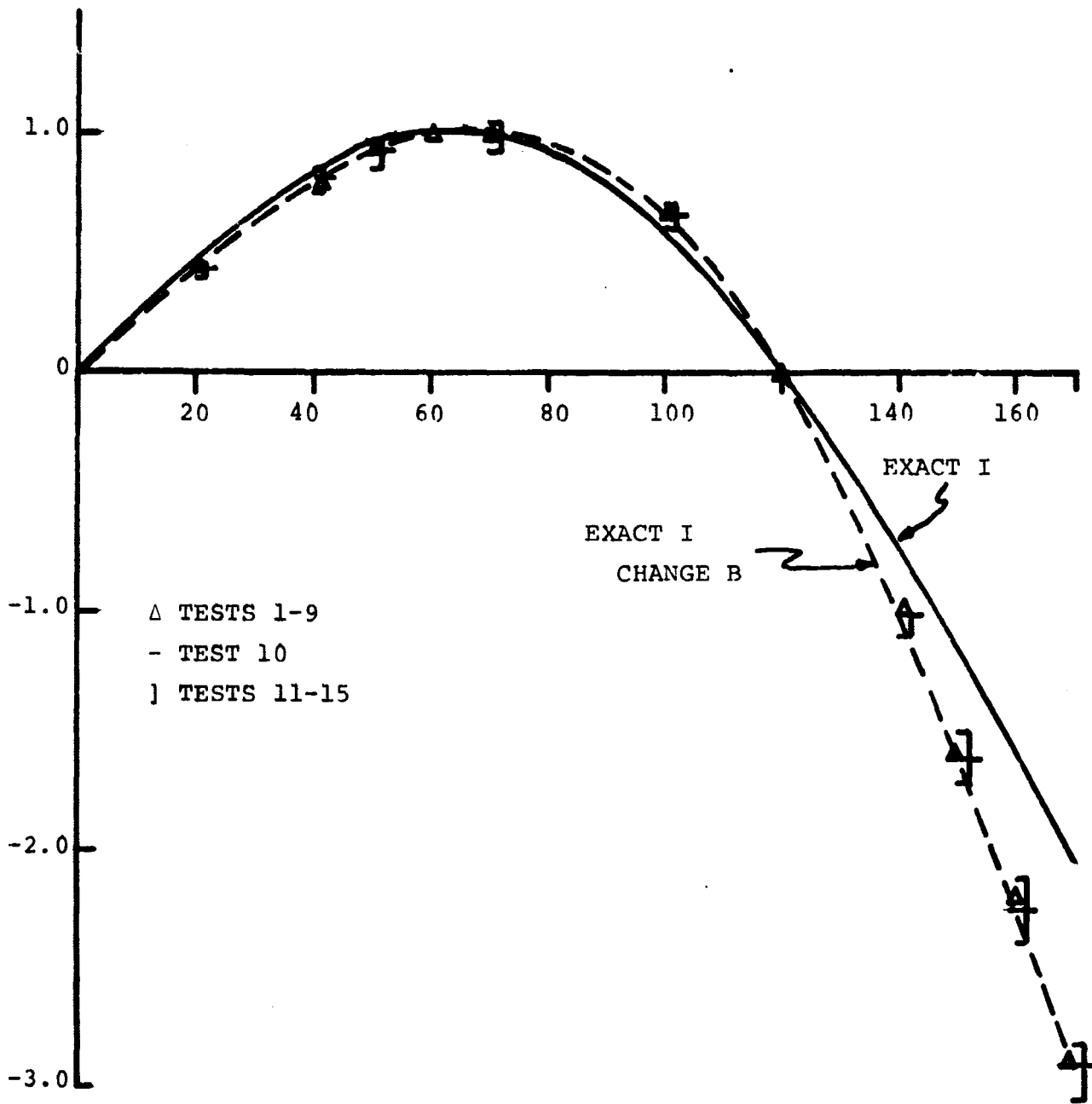


Figure 4a. Computed Mode Changes Due to Mass Change First Mode, Specimen I, Change B.

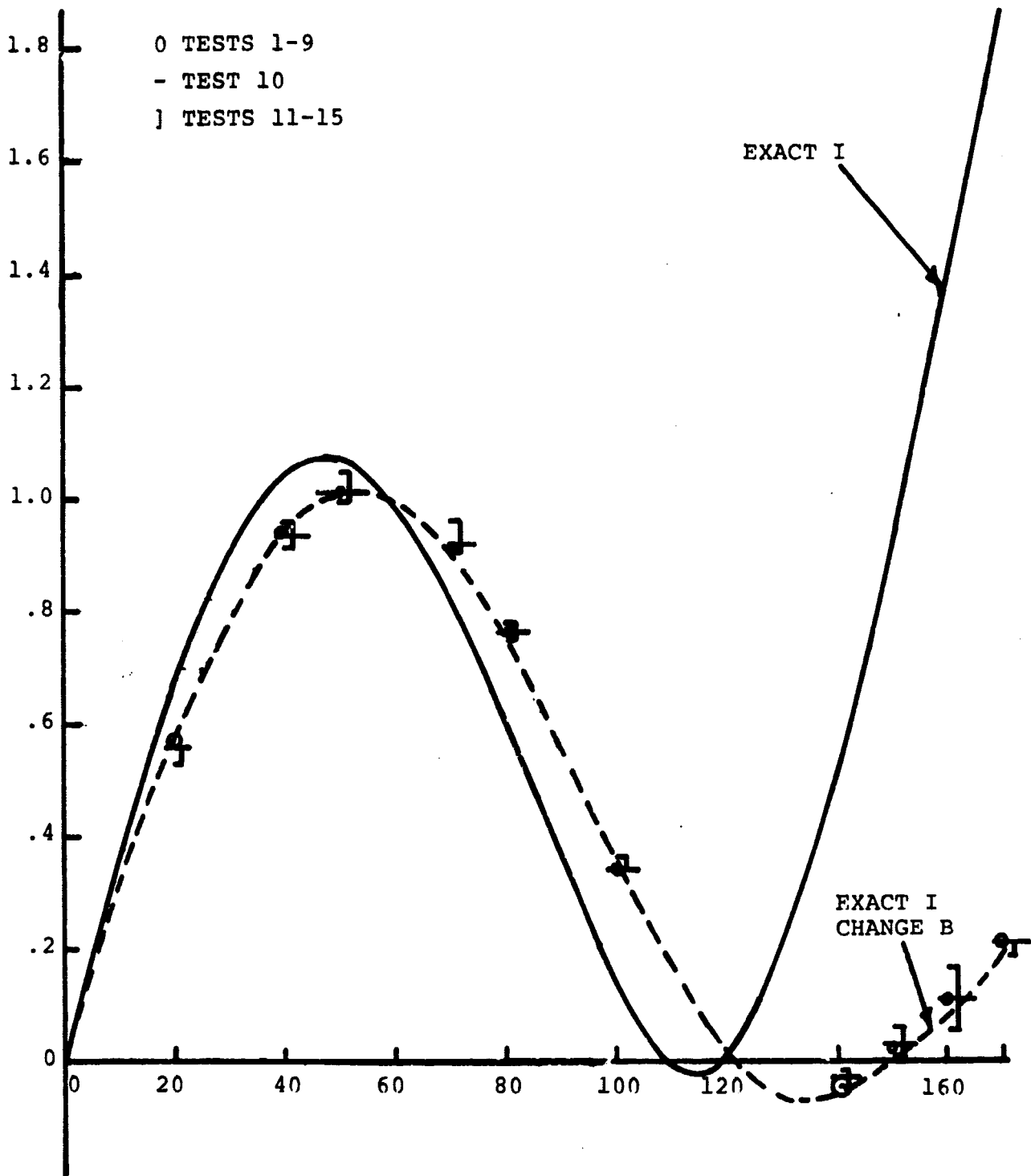


Figure 4b. Computed Mode Changes Due to Mass Change
 Second Mode, Specimen I, Change B.

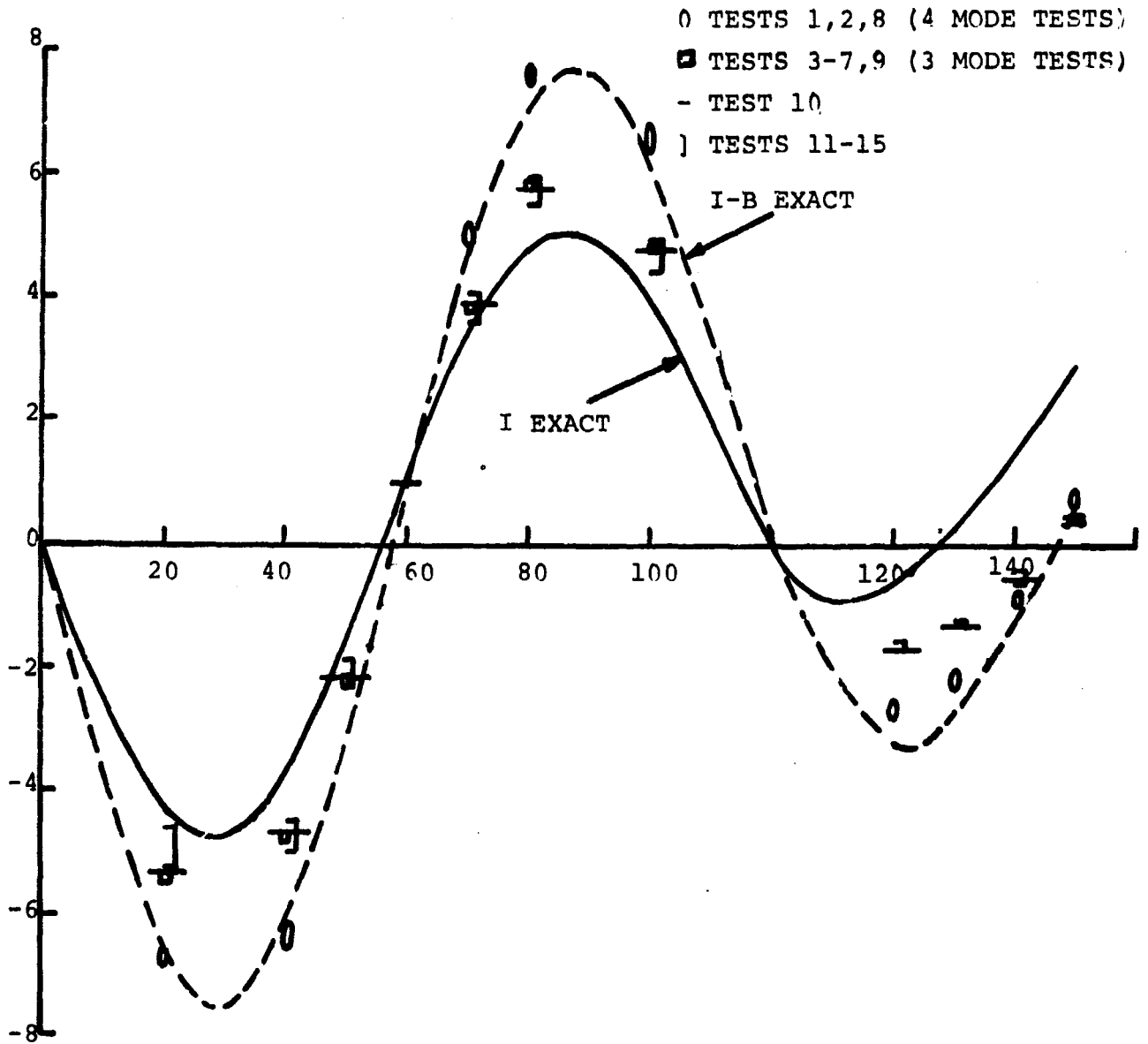


Figure 4c. Computed Mode Changes Due to Mass Change
Third Mode, Specimen I, Change B.

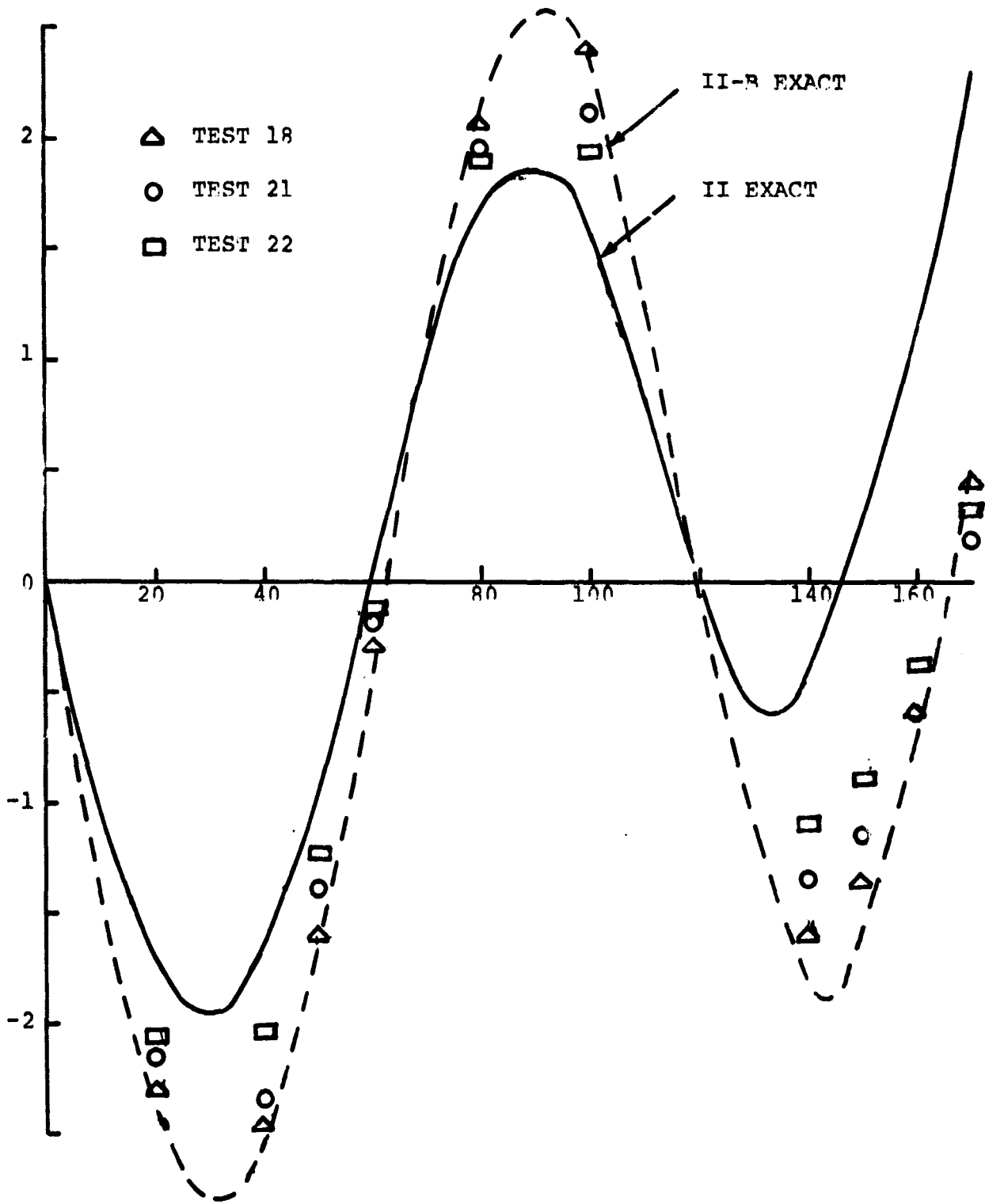


Figure 4d. Computed Mode Changes Due to Mass Change
Third Mode, Specimen II, Change B.

XII. THE EFFECTS OF STIFFNESS CHANGES

In general, the same conclusions may be drawn regarding stiffness changes as mass changes. The only possible exception being the conclusions regarding the prediction of the change in the highest mode shape measured. This is in doubt only because the stiffness changes made had very little effect on the third mode and no conclusions can be drawn.

Two kinds of stiffness changes were considered. The first type consisted of adding a spring to ground at each of the same points where the masses were added. These stations were the mid-point between the supports (change "K") and the tip of the beam (change "L"). The spring constants were 1000 lb/in. ΔK consisted of a change of the corresponding diagonal element of the K matrix (See Figure 6). As discussed in Section VI, the full $(K + \Delta K)_{inc}$ matrix was computed and the frequencies and modes were obtained by iteration on $M^{-1}(K + \Delta K)_{inc}$. The results of the frequency computations are given in Table IV and sample modes are illustrated in Figure 5.

The other stiffness change consisted of the addition of a uniform beam as illustrated in Figure 6, resulting in

a ΔK matrix consisting of nine terms. The results are given in Table V(b).

All the stiffness changes resulted in quite adequate predictions of the changes in frequency and normal mode with the exception of test 2) (as anticipated).

TABLE V(a). COMPUTED FREQUENCIES DUE TO STIFFNESS CHANGE

Mode	1		2		3		4	
f*	8.32		18.65		49.06		96.12	
+	K	L	K	L	K	L	K	L
f**	9.34	11.25	19.32	21.23	49.07	49.20	96.26	96.49
1	9.20	10.97	19.21	20.80	49.07	49.19	96.22	96.39
2	9.19	10.93	19.20	20.76		49.18	96.22	96.40
3	9.21	11.00	19.22	20.84		49.19		
4	9.21	11.00	19.22	20.85		49.19		
5	9.19	10.96	19.21	20.78		49.18		
6	9.20	10.97	19.21	20.80		49.19		
7	9.27	11.13	19.26	21.03	49.07	49.20		
8	9.26	11.10	19.26	20.98	49.07	49.20	96.24	96.44
9	9.26	11.10	19.26	20.99	49.07	49.20		
10	9.26	11.11	19.26	20.99	49.07	49.20		
11	9.24	11.34	19.28	21.13		49.21		
12	9.34	11.16	19.23	21.25		49.22		
13	9.22	11.02	19.26	21.02		49.20		
14	9.24	11.08	19.21	20.80		49.20		
15	9.32	11.26	19.25	20.99		49.22		
16	9.28	11.14	19.27	21.05	49.07	49.20		
17	9.28	11.14	19.27	21.05	49.07	49.19		
f*	7.91		18.40		44.01		65.94	
f**	8.86	10.80	19.11	20.66	44.01	44.41	66.03	66.15
18	8.73	10.53	19.00	20.29	44.01	44.36	66.02	66.12
19	8.75	10.57	19.01	20.34	44.01	44.37		
20	8.78	10.68	19.05	20.51				
21	9.47	11.89	19.75	22.77	44.01	44.82	66.01	66.11
22	8.89	10.87	19.14	20.77	44.01	44.52		
<p>* Frequency before changes, i.e. frequency tested. ** Exact frequency after changes. + Addition of 1000 lb/in. spring to ground at Sta 60 ("K") or Sta 170 ("L"). See Figure 6.</p>								

TABLE V(b). COMPUTED FREQUENCIES DUE TO STIFFNESS CHANGE 'M'

Mode	1	2	3	4
f*	8.32	18.65	49.06	96.12
f**	8.92	19.32	49.13	99.86
Test No.				
1	8.83	19.26	49.18	99.86
3	8.87	19.33	49.19	
6	8.86	19.32	49.20	
* Frequency before changes.				
** Exact frequency after change M. See Figure 6.				

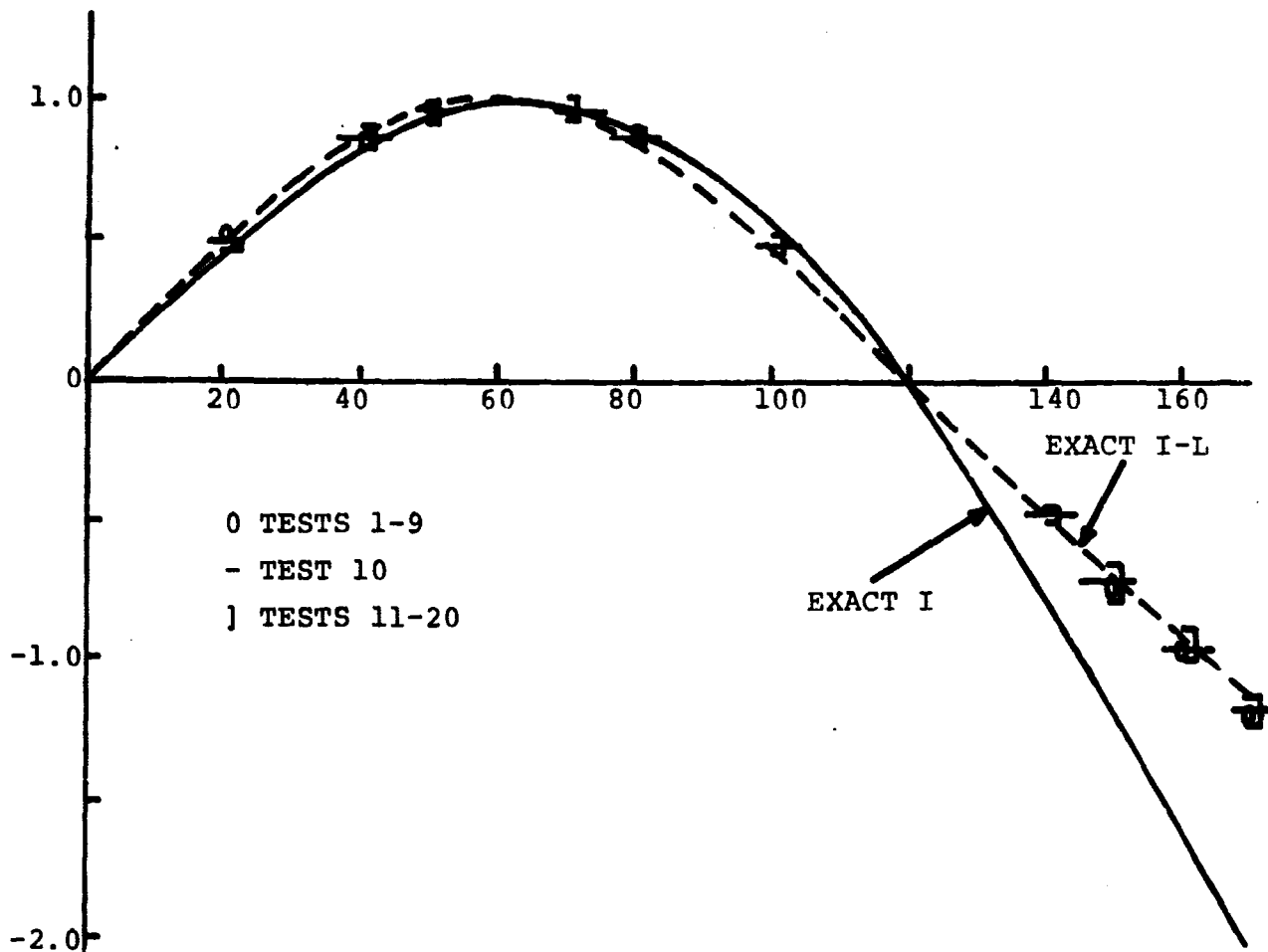


Figure 5a. Computed Mode Changes Due to Stiffness Change First Mode, Specimen I, Change L.

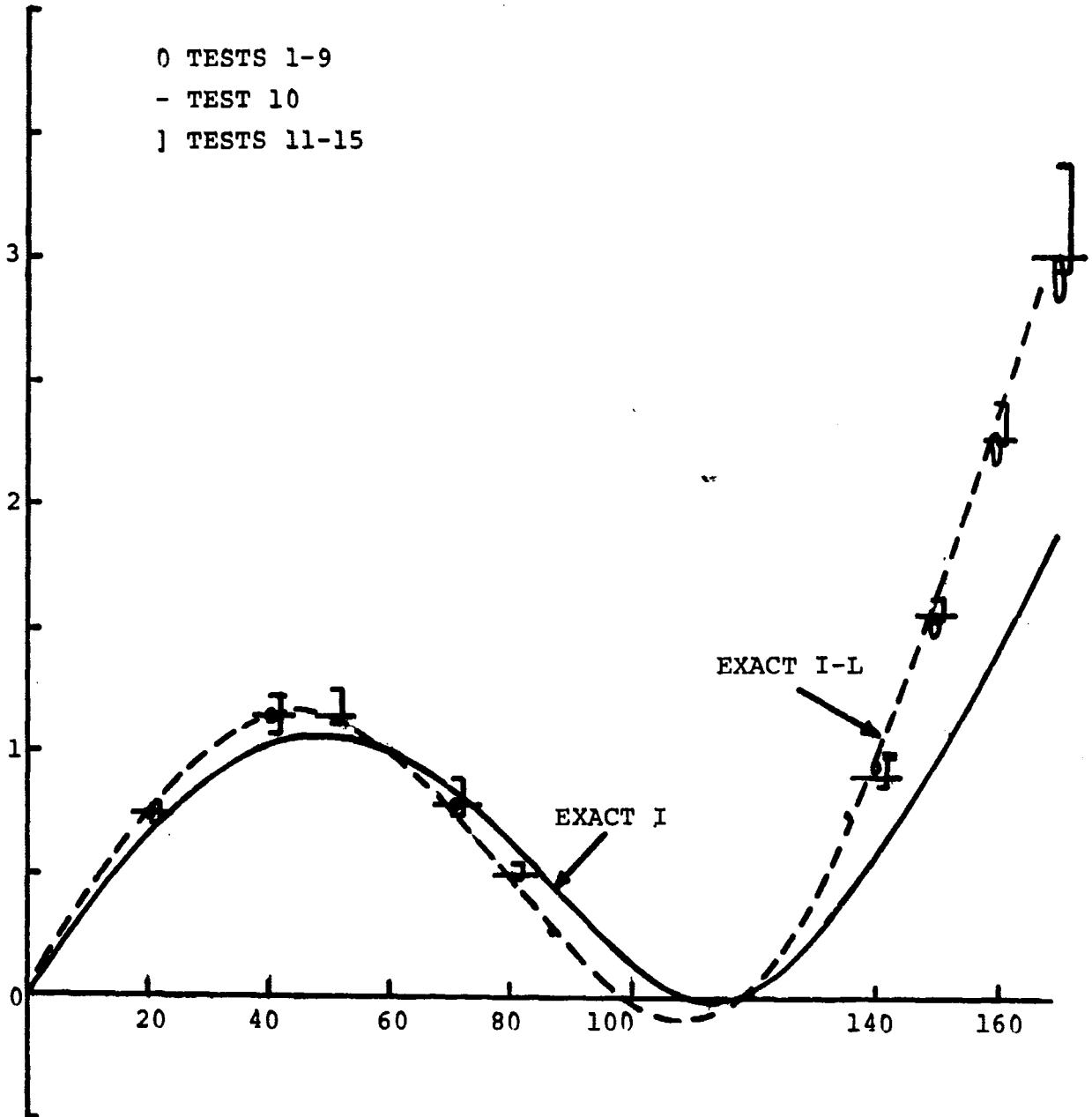


Figure 5b. Computed Mode Changes Due to Stiffness Change Second Mode, Specimen I, Change L.

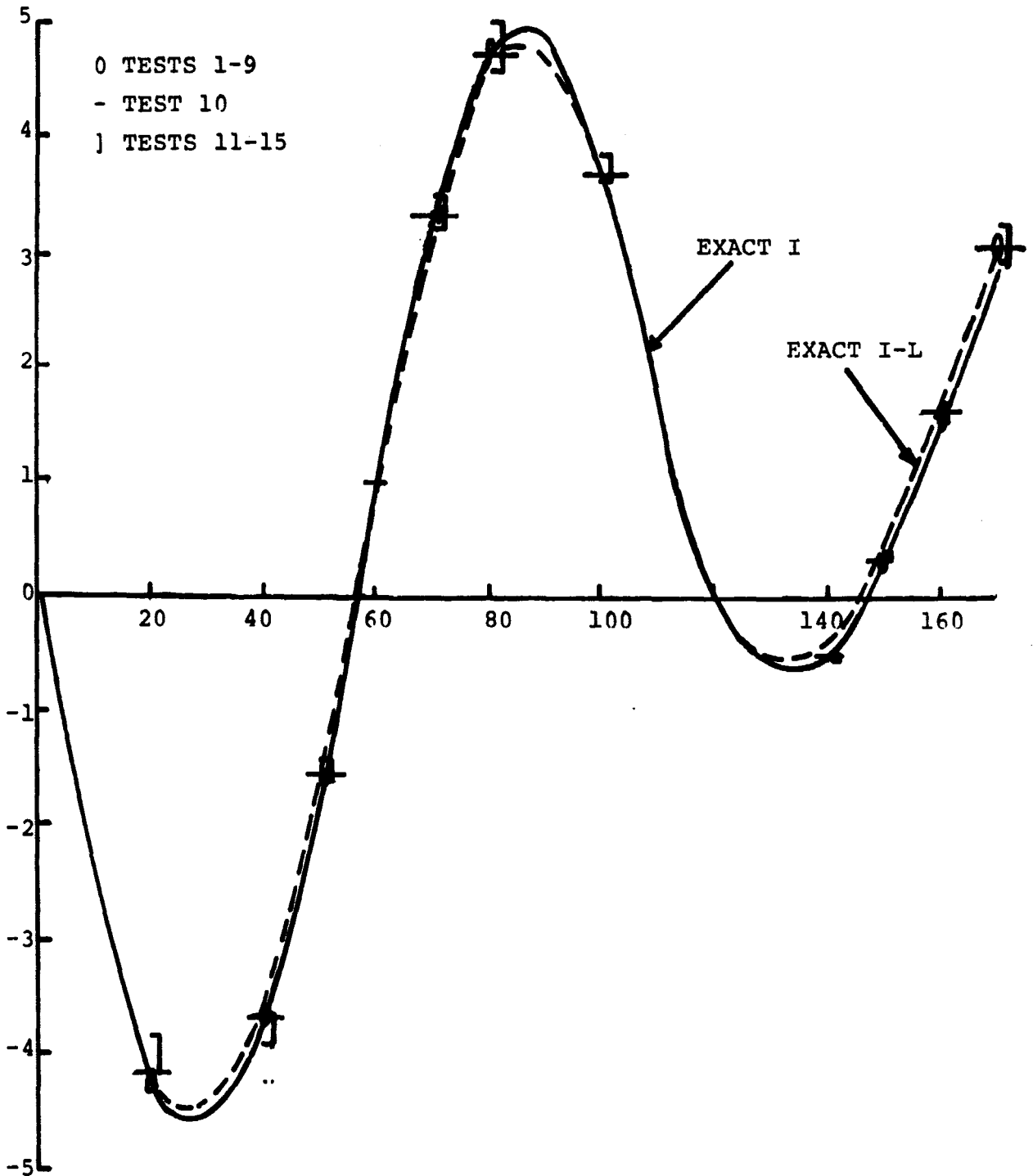


Figure 5c. Computed Mode Changes Due to Stiffness Change Third Mode, Specimen I, Change L.

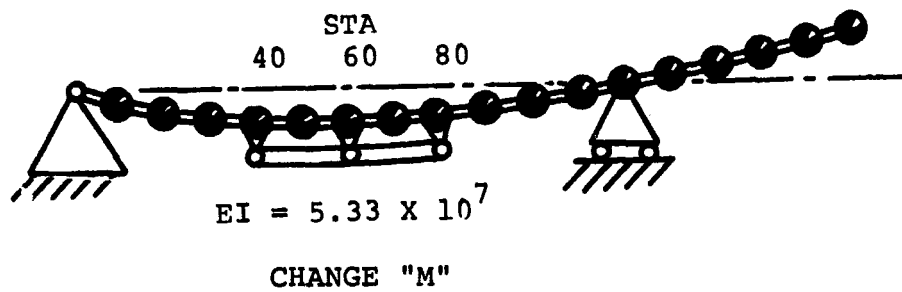
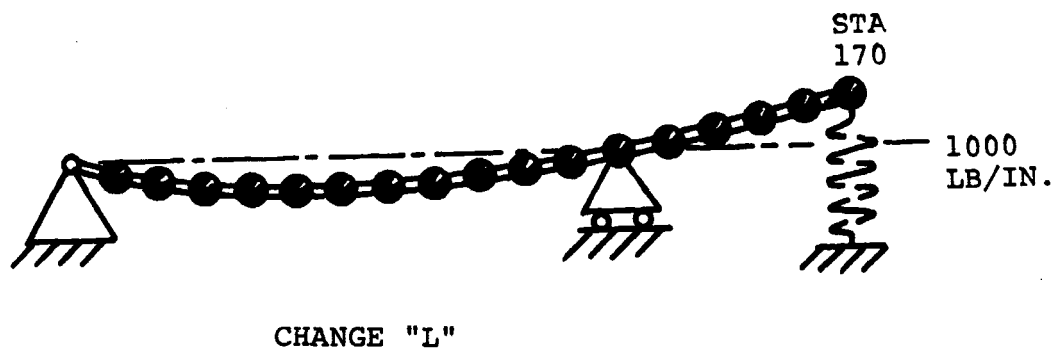
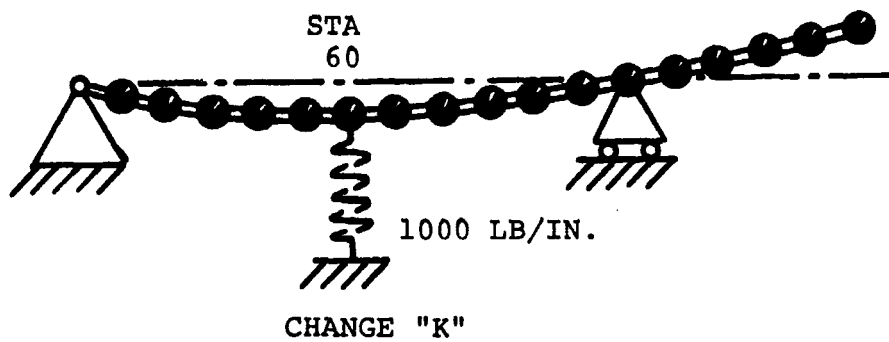


Figure 6. Stiffness Changes.

XIII. THE IDENTIFIED DAMPING COEFFICIENTS

The damping coefficient was obtained using Equation (51). This depends on the natural frequency, the real resonant velocity response, and the generalized mass. Thus, the damping coefficient may vary with the mode and with the identified mass. Table VI gives all the identified coefficients and it is seen that almost all the identifications fall within 15 percent. Tests 21, 22 which started with bad mass distributions behave as they did before. Test 21 gives very poor results while 22 gives quite good results for the first two modes. The poor results for the third mode of tests 18, 19 are not explained. It is noted, however, that the response is very small.

In the tests with error (11-15), the simulated erroneous response was used. The total scatter is of the order of 15 percent.

TABLE VI. IDENTIFIED GENERALIZED MASSES AND DAMPING COEFFICIENTS

Mode	1		2		3		4	
Ω_i	8.516		18.65		49.06		96.12	
\dot{Y}_R	.7487		.3985		.01097		.08263	
	m_1	g	m_2	g	m_3	g	m_4	g
1	1.512	.0175	1.268	.0169	17.32	.0170	1.338	.0150
2								
3	1.496	.0176	1.254	.0170	17.44	.0166		
4								
5	1.533	.0172	1.285	.0167	17.78	.0176		
6	1.521	.0174	1.247	.0172	16.83	.0188		
7	1.391	.0189	1.166	.0184	15.77	.0183		
8	1.404	.0188	1.178	.0182	16.18	.0192	1.134	.0176
9	1.411	.0187	1.182	.0181	15.40	.0192		
10	1.410	.0187	1.181	.0181	15.39	.0192		
11	1.436	.0178	1.139	.0188	15.04	.0196		
12	1.291	.0183	1.235	.0171	15.36	.0185		
13	1.473	.0170	1.168	.0173	14.95	.0187		
14	1.443	.0173	1.278	.0166	16.32	.0181		
15	1.317	.0181	1.189	.0175	14.35	.0189		
16	1.376	.0186	1.155	.0185	16.48	.0180		
17	1.379	.0186	1.157	.0185	16.51	.0179		
Ω_i	7.909		18.40		44.01		65.94	
\dot{Y}_R	.6938		.4301		.000178		.0510	
	m_1	g	m_2	g	m_3	g	m_4	g
18	1.705	.0170	1.108	.0170	587.4	.0345	2.585	.0183
19	1.674	.0173	1.159	.0173	587.5	.0346		
20	1.594	.0182	1.093	.0184				
21	.784	.0370	.542	.0371	277.8	.0732	2.894	.0164
22	1.401	.0207	.959	.0210	418.4	.0486		

XIV. CONCLUSIONS

It should be recognized that the following general conclusions are based on a small amount of simulated data on a relatively simple structure. While these conclusions are not proven, they are strongly suggested by the data obtained.

1. The concept of an incomplete model of a dynamic structure is valid and useful, at least, for predicting the effects of structural changes on the normal modes and frequencies.
2. The parameters of an incomplete model may be determined from measured modal data and a "reasonable" approximation to the mass matrix.
3. The identified mass matrix will be approximately equal to the assumed values when the assumption is reasonably valid. The identified structural damping coefficient will also be a reasonable approximation under this condition.
4. The identified stiffness matrix, while not approximating the true values, may be used to predict the effects of stiffness changes.

5. The identified influence coefficient matrix will be an approximation to the true values.

6. The effects of mass and stiffness changes on the normal modes and natural frequencies can be estimated using the techniques that have grown out of the incomplete model concept. These predictions are not overly-sensitive to the mass approximation and to errors of measurement.

7. The procedures may tend to deteriorate when fewer points or more modes are used, reducing the degeneracy of the equations. It is expected that this condition can be simply corrected by including off-diagonal masses as unknowns.

In general, the results of the qualitative study are quite encouraging. It is recommended that a more definitive evaluation of the concepts be carried out both from a theoretical viewpoint and by more detailed computer experimentation with more complex structures.

NEW TECHNOLOGY

This entire report is considered to be a "reportable item" in accordance with the New Technology Clause (NASA Form 1162).

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GLOSSARY OF SYMBOLS

A	Coefficient matrix in mass equation (See Section IV)
B_R	Matrix multiplying R in mass identification (See Section IV)
B_m	Matrix multiplying m_A in mass identification (See Section IV)
C	Influence coefficient matrix
f	Force vector
g	Structural damping coefficient
K	Stiffness matrix
M	Mass matrix
m_i	Generalized mass of i-th mode
m_{ij}	Element of M
m_A	Approximation to \bar{m}
m_T	Total of diagonal masses
\bar{m}	Vector consisting of unknown elements of M
N	Number of modes
n_e	Number of equations
n_v	Number of variables
P	Number of points of interest
R	Right-Hand side of mass equation (See Section IV)
W	Weighting matrix referring to confidence in m_A (diagonal)
Y	Mobility matrix
y	Displacement vector

GLOSSARY OF SYMBOLS (Continued)

\dot{Y}_R	Real component of velocity vector
Z	Impedance
Δ	Any change (as a prefix)
Φ	Matrix of modes
ϕ_i	Modal vector (i-th mode)
Ω_i	Natural frequency of i-th mode
ω	Forcing frequency
+	Superscript - pseudo inverse
$\uparrow \downarrow$	Indicates diagonal matrix

APPENDIX I
PRINCIPAL IDEMPOTENTS

The expression of positive definite matrices in terms of their principal idempotents has been discussed in Reference 5. It has been useful in the development of the incomplete model concept and some of the associated techniques. A brief discussion is given below.

Define ϕ_i as an arbitrary set of orthogonal vectors (P vectors each having P elements) and A as an arbitrary positive definite PxP matrix.

It is possible to express A as follows:

$$A = \sum_{i=1}^P \sum_{j=1}^P C_{ij} \phi_i \phi_j^T \quad (I-1)$$

since the ϕ 's are independent and there are as many C_{ij} 's as there are elements in A. Note that $\phi_i \phi_j^T$ is a non-zero square matrix for all i, j combinations but $\phi_j^T \phi_i$ is a scalar and equal to 0 unless $i=j$ since the vectors are orthogonal.

Now pre- and post-multiply the equations by ϕ_n^T and ϕ_k

$$\phi_n^T A \phi_k = \sum_{i=1}^p \sum_{j=1}^p C_{ij} (\phi_n^T \phi_i) (\phi_j^T \phi_k) = C_{nk} (\phi_n^T \phi_n) (\phi_k^T \phi_k)$$

since each term in the summation is zero except when $i=n$ and $j=k$. Thus (note that $\phi_n^T A \phi_k$ is a scalar)

$$C_{ij} = \frac{\phi_i^T A \phi_j}{(\phi_i^T \phi_i) (\phi_j^T \phi_j)} \quad (\text{I-2})$$

Thus for any set of orthogonal vectors it is possible to express the matrix A as in Equation (I-1) by evaluating the coefficients as in Equation (I-2).

However, if the ϕ 's are the eigenvectors of A, $A\phi_j = \lambda_j \phi_j$ and

$$C_{ij} = \frac{\lambda_j \phi_i^T \phi_j}{(\phi_i^T \phi_i) (\phi_j^T \phi_j)} = 0 \text{ for } i \neq j$$

since now $\phi_i^T \phi_j = 0$ for $i \neq j$, and

$$C_{ii} = \frac{\lambda_i}{\phi_i^T \phi_i}$$

Thus, the cross terms in Equation (I-1) become zero and

$$A = \sum_{i=1}^p \lambda_i \left(\frac{\phi_i \phi_i^T}{\phi_i^T \phi_i} \right) \quad (\text{I-3})$$

if and only if the ϕ 's are eigenvectors of A. Note that (for any orthogonal ϕ 's), if $V_i = \frac{\phi_i \phi_i^T}{\phi_i^T \phi_i}$, $V_i^2 = \frac{\phi_i (\phi_i^T \phi_i) \phi_i^T}{(\phi_i^T \phi_i)^2} = V_i$ thus V_i is idempotent. When the ϕ_i 's are the eigenvectors of A, V_i is called a principal idempotent.

By similar procedures it can be shown that the influence coefficient matrix and the stiffness matrix can be written in terms of their principal idempotents as follows (using the notation of the body of this report).

$$CM = \sum_{i=1}^p \frac{1}{\Omega_i^2} \left(\frac{\phi_i \phi_i^T M}{m_i} \right) \quad (I-4)$$

$$M^{-1}K = \sum_{i=1}^p \Omega_i^2 \left(\frac{\phi_i \phi_i^T M}{m_i} \right) \quad (I-5)$$

Where the terms in brackets are the principal idempotents of CM and $M^{-1}K$. It follows from the above two equations that

$$C = \sum_{i=1}^p \frac{1}{\Omega_i^2 m_i} \phi_i \phi_i^T \quad (I-6)$$

$$K = \sum_{i=1}^p \frac{\Omega_i^2}{m_i} M \phi_i \phi_i^T M \quad (I-7)$$

and that expressions of this form are only possible when the ϕ 's are the eigenvectors of CM or $M^{-1}K$.

APPENDIX II
PSEUDO INVERSES

The pseudo inverse, sometimes called the generalized inverse of a matrix is discussed in numerous publications (e.g. Reference 6). In this appendix, only the application made to the problem at hand will be discussed.

Consider the set of independent linear equations

$$Ax = b \quad (\text{II-1})$$

where there are more equations than unknowns. A , then, is a rectangular matrix, $n_e \times n_v$ where $n_e < n_v$. x then is $1 \times n_v$ and b is $1 \times n_e$. There are then an infinite number of solutions to these equations.

The pseudo inverse of A , written A^+ , is defined for this problem to be

$$A^+ = A^T(AA^T)^{-1} \quad (\text{II-2})$$

where the ordinary inverse of AA^T is used. This matrix, AA^T , is of order $n_e \times n_e$ and can be shown to be non-singular when the original equations are independent.

Now

$$\bar{x} = A^+ b$$

is a solution to Equation (II-1) as can be seen from

$$A\bar{x} \equiv AA^+ b \equiv AA^T (AA^T)^{-1} b \equiv IB = b$$

It is shown in the literature that this solution (of the infinity of solutions) is the one having the minimum sum of squares of the individual elements, i.e. the "smallest" in a least squares sense.

Using Equation (II-2) to obtain the pseudo inverse of AW^{-1} from Equation (39) in the main text (where W^{-1} is a diagonal matrix),

$$\begin{aligned} (AW^{-1})^+ &= (AW^{-1})^T \{ (AW^{-1}) (AW^{-1})^T \}^{-1} \\ &= W^{-1} A^T \{ AW^{-1} W^{-1} A^T \}^{-1} \end{aligned}$$

APPENDIX III

PROGRAM LISTING AND SAMPLE COMPUTATION

This appendix contains the Fortran listing of the computer program discussed in Section VIII. Included are the operating instructions and a sample computation of test 16.

The program was compiled and run on an IBM 360, Model 40, having 128K bytes of storage under DOS, release 17.

REPRODUCIBILITY OF THE ORIGINAL PAGE IS POOR

DISK OPERATING SYSTEM/360 FORTRAN 360N-FQ-451 31

```

***** INCMOD *****
C
C   INCOMPLETE MODEL THEORY - TEST PROGRAM - USES RESONANT RESP AS MIDE
C
C INPUT - PART I - SIMULATED TEST
C   1. HEAD1 COL 1 = IC, COL 2-80 HEADING
C       IC = 0, NORMAL (INPUT 2,(3),4,5,7,8FF)
C           1, SAME AS 0, BUT C INPUT INSTEAD OF K (2,(3),4,6,7,8FF)
C           2, NEW ACTUAL MASS AND FRFC ONLY (INPUT 4,7,8FF)
C           3, NEW FREQUENCIES ONLY (INPUT 7,8FF)
C           4, SAME ACTUAL SYSTEM, NEW TEST (INPUT 2A,(3),7,8FF)
C           5, NO ANALYTICAL MODEL, INPUT PHI'S DIRECTLY (2A,7,7A,8)
C
C   2. CONTROL CARD (PT. I)
C       COL 1,5 ND, NO OF ACTUAL DEG OF FREEDOM 15 20 MAX
C           6,10 NU, NO OF MEASUREMENT POINTS (USED) 15 20 MAX
C           IF NU .LT. NO, READS 3.
C           11,20 NROW, DRIVING POINT INDEX 110
C                   MUST BE A USED POINT
C           21,30 G, SCALAR STRUCT DAMPING COEFF F10
C           31,40 PCT, RANDOM ERROR CN AMPLITUDE, ERROR F10
C                   UNIFORMLY DISTRIBUTED BETWEEN -/+
C                   PCT*RESPONSE
C           41,50 PCTB, BIAS ERROR CN AMP = PCTB*RESP F10
C           71,80 IZ, SPEED FOR RANDOM NC GENERATOR 110
C
C   2A. SAME AS 2. EXCEPT DOES NOT READ ND
C
C   3. USED STA. CARD
C       USED(I),I=1,NU INDICES OF STATIONS
C       TO BE USED , IN NUMERICAL SEQUENCE 1914 19 MAX
C
C   4. ACTUAL MASS CARD(S)
C       FIRST CARD(S) M(I,I),I=1,ND DIAGONAL MASSES IN
C       SEQUENCE, AS MANY CARDS AS NECESSARY 8F10
C       NEXT CARD COL 1,10 NOD, NO OF OFF-DIAGONAL MASSES 110
C       NEXT CARDS (IF NOD .NE. 0) I,J,MASS 8(212,F6)
C       LOWER TRIANGLE ONLY (SYMMETRY ASSUMED)
C       B PER CARD, AS MANY CARDS AS NECESSARY
C
C   5. ACTUAL K CARDS K(I,J) I=1,ND ,J=1,I
C       STIFFNESS MATRIX, LOWER TRIANGLE INPUT ONLY, MATRIX
C       WILL BE SYMMETRIZED
C       START EACH ROW ON NEW CARD, END CN DIAGONAL
C       ELEMENT 8F10
C
C   6. ACTUAL C CARDS C(I,J)
C       INFLUENCE COEF MATRIX, SAME FORM AS K
C
C   7. ACTUAL NATURAL FREQ (USED IN SIMULATED TEST)
C       FIRST CARD NFRFC, NUMBER OF FREQUENCIES 110 10 MAX
C       NEXT CARD(S) FRFC(I),I=1,NFRFC 8F10
C
C   7A. ACTUAL PHI'S (PHI(I,J),J=1,NFRFC),I=1,NU

```

INPUT - PART II - IDENTIFICATION

8. HEAD2 COL 1 = IC2, COL 2-80 HEADING
 IC2 = 0, NORMAL (INPUT 9,10,11,12FF)
 1, NEW APPROX MASS - OTHER DATA UNCHANGED (11,12FF)
9. GENERALIZED MASS - KNOWN VALUES ONLY - IF NONE USE 1 BLANK CARD
 CD1 COL 1-10 NG, NO OF VALUES I10 10 MAX
 FOLLOWING CARD(S) I,G,M(I) 8(12,F9)
10. KNOWN MASSES - IF NONE USE 1 BLANK CARD
 CD1 COL 1-10 NK, NO OF VALUES I10 20 MAX
 FOLLOWING CARD(S) I,J,KP(I,J) (SEE *NOTE)
 ENTER LOWER TRIANGLE VALUES ONLY (I .GE. J) 8(212,F6)
 SYMMETRY ASSUMED
 ONLY NON ZERO MASSES
11. APPROX MASS CARDS - DEFINE UNKNOWN
 CD1 COL 1-10 NV, NO OF UNKNOWN MASSES I10 20 MAX
 COL 20 ICT = 0 NO CONTROL ON TOTAL MASS
 1 TOTAL UNKNOWN DIAGONAL MASSES CONSTANT
 FOLLOWING CARDS, ONE PER MASS - LOWER TRIANGLE ONLY
 COL 1-5 I (SEE *NOTE)
 6-10 J (SEE *NOTE)
 11-20 AM, APPROX MASS (MAY BE ZERO IF OFF DIAG)
 21-30 W, WEIGHTING FACTOR 215,2F10
 HIGHER W INDICATES GREATER CONFIDENCE IN ESTIMATE
 DO NOT USE W = 0

NOTE, ALL MASSES NOT DEFINED IN 10 OR 11 WILL BE ZERO
 SEE SUBROUTINE IDMASS FOR FRCP MESSAGES

INPUT - PART III - EFFECTS OF VARYING MASS CO STIFFNESS

12. HEAD3 COL 1 = IC3, COL 2-80 HEADING
 IC3 = 0, MASS CHANGE (INPUT 13,12FF)
 = 1, RETURN TO PART I (I)
 = 2, RETURN TO PART II (R)
 = 3, STIFFNESS (14,12FF)
 = 9, END OF RUN
13. CHANGED MASS CARDS
 FIRST CARD COL 1-10 NC, NO OF CHANGES I10
 NEXT CARD(S) I,J,MASS CHANGE (SEE *NOTE) 8(212,F6)
 LOWER TRIANGLE ONLY, SYMMETRY ASSUMED
14. CHANGED STIFFNESS CARDS
 SAME FORM AS 13. EXCEPT 4(212,F16)

NOTE, NEXT CARD READ IS 12. THESE MASS CHANGES ARE ONLY
 TEMPORARILY STORED. VARIOUS CHANGES CAN BE RUN IN SEQUENCE.

*NOTE I, J ARE INDICES OF TEST POINTS, NOT ACTUAL SYSTEM

01/03/70

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SUBROUTINES SYM, INVR5, MOUT2, PCB, CINV, MPPY, RANCU, ICMASS, PSEUDO, MITER, GEN

INTEGER HEAD(20), USED(20), HEAD2(20)
INTEGER HEAD3(20), ITN(10)
REAL M(20,21), K(20,21), C(20,21), FPEC(10), ZP(20,21), ZI(20,21),
1 YR(20,21), YI(20,21), RESR(10,20), RESI(10,20), GPP(10), PHI(20,10),
2 GPF(10), MINC(20,20), CINC(20,20), KINC(20,20)
REAL MTEMP(20,21), MOOE(20), PHIC(20,10), FRECC(10), CTMP(20,20)
REAL KTEMP(20,20), MINV(20,21), MOOM(20), CK(20,20), GMASS(10),
1 DCOF(20,10)
COMMON M, K, C, ZR, ZI, YR, YI, MINC, CINC, KINC, MTEMP, CTMP
COMMON KTEMP, MINV, DK, DCOF

C

PART I - SIMULATED TEST

1 READ (1,1000) IC, HEAD
1000 FORMAT(I1, A3, I9A4)
WRITE (3,1001) IC, HEAD
1001 FORMAT ('1'//T5, 'INCOMPLETE MCOFL THEORY - SIMULATED TEST'//
1 T5, 15(' '), I2, 3X, A3, I9A4, 5X, 15(' ')//)
ICC=IC+1
GO TO (10,10,20,70,11,11), ICC
10 READ (1,1002) ND, NU, NR0W, G, PCT, PCTA, IZ
1002 FORMAT (2I5, I10, 3F10.0, 20X, I10)
GO TO 13
11 READ (1,2002) NU, NR0W, G, PCT, PCTA, IZ
2002 FORMAT (5X, I5, I10, 3F10.0, 20X, I10)
13 IX=IZ*2+1
GO TO (14,14,14,14,14,18), ICC
14 IF(NU-ND)15,18,18
15 READ (1,1003) (USED(I), I=1, NU)
1003 FORMAT (I9I4)
DO 16 I=1, NU
IF(USED(I)-NR0W)16,17,16
16 CONTINUE
NR0WN=USED(1)
NR0WN=1
WRITE (3,999) NR0W
999 FORMAT (T5, '*** FORCING POINT INVALID. SET TC' I3//)
GO TO 20
17 NR0WN=I
GO TO 20
18 NR0WN=NR0W
DO 19 I=1, NU
19 USED(I)=I
20 GO TO (25,25,25,25,70,70), ICC
25 DO 30 I=1, ND
DO 30 J=1, ND
30 M(I,J)=0
READ (1,1004) (M(I,I), I=1, ND)
1004 FORMAT (3F10.0)
READ(1,1005) NDD
1005 FORMAT (I10)
IF (NDD) 40,40,35

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FORMATN

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35 READ (1,1006) (I,J,M(I,J),L=1,ND)
1006 FORMAT (8(2I2,F6.0))
    CALL SYM (M,ND)
40 GO TO (50,60,70,70,70,70),ICC
50 DO 55 I=1,ND
55 READ (1,1004) (K(I,J),J=1,I)
    CALL SYM(K,ND)
    CALL INVRS (K,ND,C)
    GO TO 70
60 DO 65 I=1,ND
65 READ (1,1004) (C(I,J),J=1,I)
    CALL SYM (C,ND)
    CALL INVRS (C,ND,K)
70 READ (1,1005) NFREQ
    READ (1,1004) (FREQ(I),I=1,NFREQ)
    GO TO (71,71,71,71,71,72),ICC
71 WRITE (3,1007) ND,G
1007 FORMAT (T10,'ACTUAL SYSTEM PARAMETERS'//10,' DEG OF FREEDOM, STRUCT
1 DAMPING COEF =F6.3/T10,24('-')//T50,'MASS MATRIX'//)
    GO TO 74
72 WRITE (3,2007) NFREQ,G
2007 FORMAT (T10,'ACTUAL MODES'//10,' MODES, STRUCTURAL DAMPING COEF ='
1 F6.3/T10,12('-')//)
    GO TO 90
74 CALL MOUT2 (M,ND,ND)
    GO TO (75,75,80,80,80,80),ICC
75 WRITE (3,1008)
1008 FORMAT ('1',T50,'STIFFNESS MATRIX'//)
    CALL MOUT2 (K,ND,ND)
    WRITE (3,1009)
1009 FORMAT ('1',T50,'INFLUENCE COEFFICIENT MATRIX'//)
    CALL MOUT2 (C,ND,ND)
    GO TO 90
80 WRITE (3,1010)
1010 FORMAT (//T10,'ACTUAL K AND C SAME AS PREVIOUS CASE')
90 WRITE (3,1011) NROW,NU
1011 FORMAT ('1',//T10,'DESCRIPTION OF SIMULATED TEST'//T10,29('-')//
1 T20,'POINT AT WHICH SYSTEM IS FORCED'//T20,'NO OF POINTS OF MEA
2SUREMENT'//8)
    GO TO (94,94,94,94,94,96),ICC
94 IF(ND-NU) 100,100,95
95 WRITE (3,1012) (USED(I),I=1,NU)
1012 FORMAT (T20,'POINTS AT WHICH MEASUREMENTS ARE TAKEN'//T30,20I4)
100 WRITE (3,1013) (FREQ(I),I=1,NFREQ)
1013 FORMAT (T20,'FREQUENCIES OF EXCITATION - HZ'//T30,1P5F15.4)
96 WRITE (3,1014) PCT,PCTB,IZ
1014 FORMAT (T20,'AVE RANDOM AMPLITUDE ERROR'//F10.3/T20,'BIAS AMPLITUDE
1 ERROR'//F16.3//T20,'SEED FOR RAND NO GENERATOR'//I2//)
    GO TO (109,109,109,109,109,101),ICC
101 DO 102 J=1,NU
102 READ (1,1004) (RFSR(I,J),I=1,NFREQ)
    GO TO 111
C
109 DO 110 I=1,NFREQ
    OM=FREQ(I)
    CALCULATE RESONANT RESPONSES

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CALL M08 (M,K,G,ND,CM,ZR,ZI,YR,YI)
DO 110 J=1,NU
JJ=USED(J)
C
C          COMPLEX FORM
RESR(I,J)=YR(NROW,JJ)
110 RESI(I,J)=YI(NROW,JJ)
C
C          CALC ASSUMED PHIS = REAL RESPONSE WITH ERROR
111 DO 120 I=1,NU
DO 120 J=1,NFREQ
CALL RANDU (IX,IY,YEL)
IX=IY
120 PHI(I,J)=RESR(J,I)*(1.0+2.0*PCT*(YEL-0.5)+PCTP)
C          NORMALIZE ON DRIVING POINT
DO 130 I=1,NFREQ
GPP(I)=RESR(I,NROWN)
GPPE(I)=PHI(NROWN,I)
DO 130 J=1,NU
PHI(J,I)=PHI(J,I)/GPPE(I)
IF(IC-5) 129,130,130
129 RESI(I,J)=RESI(I,J)/GPP(I)
130 RESR(I,J)=RESR(I,J)/GPP(I)
GO TO (135,135,135,135,135,131),ICC
131 WRITE (3,2015) (FREQ(I),I=1,NFREQ)
2015 FORMAT (//T10,'EXACT MODES OF ACTUAL SYSTEM, NORMALIZED ON DRIVING
1 POINT'//T3,'FREQ '1P10E12.4)
GO TO 139
135 WRITE (3,1015) (FREQ(I),I=1,NFREQ)
1015 FORMAT(//T10,'EXACT RESPONSE OF ACTUAL SYSTEM,IN/SEC/POUND, NORMAL
IZED ON REAL DRIVING POINT RESPONSE'//T30,'REAL PART'//T3,'FREQ
2*1P10E12.4)
WRITE (3,1016) (GPP(I),I=1,NFREQ)
1016 FORMAT(//T3,'D.P.'//T4,'RESP '1P10E12.4)
139 WRITE (3,1004)
DO 140 I=1,NU
140 WRITE (3,1017) USED(I),(RESR(J,I),J=1,NFREQ)
1017 FORMAT (16,3X,10F12.5)
GO TO (142,142,142,142,142,141),ICC
141 WRITE (3,2019)
2019 FORMAT ('1'//T10,'SIMULATED MEASURED MODES, ACTUAL PHIS WITH ERR
10P'//)
GO TO 154
142 WRITE (3,1018)
1018 FORMAT ('1'//T30,'IMAGINARY PART'//)
DO 145 I=1,NU
145 WRITE (3,1017) USED(I),(RESI(J,I),J=1,NFREQ)
WRITE (3,1019)
1019 FORMAT ('1'//T10,'SIMULATED TEST RESPONSE, REAL MOBILITY WITH ERR
10RS USED AS PHIS'//)
WRITE (3,1016) (GPPE(I),I=1,NFREQ)
WRITE (3,1004)
154 CALL MOUT2 (PHI,NU,NFREQ)
C          PART II - IDENTIFY MASSES
155 READ (1,1000) IC2,HEAD2
CALL IDMASS (IC2,HEAD2,FREQ,PHI,NU,NFREQ,MINC,CINC,KINC,GMASS)
C          CALCULATE RESONANT RESPONSE FROM KINC,MINC
```

01/03/70

FORTRAN

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DO 160 I=1,NRFREQ
  OM=FREQ(I)
  CALL MOR (MINC,KINC,G,NU,CM,ZR,7I,YR,YI)
  DO 160 J=1,NU
    RESR(I,J)=YR(J,NPCWN)
160 RESI(I,J)=YI(J,NPCWN)
  DO 170 I=1,NRFREQ
    GPP(I)=RESR(I,NPCWN)
  DO 170 J=1,NU
    RESI(I,J)=RESI(I,J)/GPP(I)
170 RESR(I,J)=RESR(I,J)/GPP(I)
  WRITE (3,1020) (FREQ(I),I=1,NRFREQ)
1020 FORMAT ('1'//T10,'RESONANT RESPONSE OF IDENTIFIED SYSTEM'//
  1 T30,'REAL PART'//
  2 T3,'FREQ '1P10F12.4)
  WRITE (3,1016) (GPP(I),I=1,NRFREQ)
  WRITE (3,1004)
  DO 180 I=1,NU
180 WRITE (3,1017) USED(I),(RESR(J,I),J=1,NRFREQ)
  WRITE (3,1018)
  DO 190 I=1,NU
190 WRITE (3,1017) USED(I),(RESI(J,I),J=1,NRFREQ)
  IFIRST = 0
  GO TO 2,5

```

PART III - CHANGED MASSES

```

200 READ (1,1000) IC3,HEAD3
  IF(IC3-9)220,210,210
210 CALL EXIT
220 ICC=IC3+1
  GO TO (250,1,155,250),ICC
250 WRITE (3,1030) IC3,HEAD3
1030 FORMAT ('1'//T5,'PART III - CHANGED MASSES OR STIFFNESSES'//T5,
  1 15('**'),I2,3X,A3,19A4,5X,15('**')//)
  READ (1,1005) NC
255 DO 260 I=1,NU
  DO 260 J=1,NU
    KTEMP(I,J)=0
260 MTEMP(I,J)=0
  IF(IFIRST)261,269,261
261 IF(IC3)265,265,300
265 READ (1,1006) (I,J,MTEMP(I,J),L=1,NC)
  CALL SYM (MTEMP,NU)
  WRITE (3,1031)
1031 FORMAT (T50,'MASS MATRIX CHANGES'//)
  CALL MOUT2 (MTEMP,NU,NU)
269 DO 270 I=1,NU
  DO 270 J=1,NU
    CTEMP(I,J)=CINC(I,J)
270 MTEMP(I,J)=MTEMP(I,J)+M*NC(I,J)
  ITERATE USING CINC * MTEMP
  DO 280 I=1,NRFREQ
  CALL MITER(CTEMP,MTEMP,NU,.00001,25,MODE,CMFG,IT)
  CON=M10E(NPCWN)
  DO 275 J=1,NU
  MODE(J)=MODE(J)/CON

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01/03/70

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275 PHIC(J,I)=MODE(J)
    ITN(I)=IT
    FREQC(I)=1.0/SQRT(OMEG)/6.2832
    GPPE(I)=GFN(MODE,MTEMP,NU)
    CON=OMEG/GPPF(I)
    DO 280 L=1,NU
    DO 280 J=1,NU
280 CTEMP(L,J)=CTEMP(L,J)-CON*MODE(L)*CCDF(J)
    IF(IFIRST)285,281,285
281 WRITE(3,2032) (FREQC(I),I=1,NFREQ)
2032 FORMAT ('1'//T10,'FREQUENCIES AND MODES OF IDENTIFIED SYSTEM'//
1 T3,'FREQ' '1P10F12.4)
    GO TO 289
285 WRITE (3,1032) (FREQC(I),I=1,NFREQ)
1032 FORMAT ('1'//T10,'NEW FREQUENCIES AND MODES OF MODIFIED SYSTEM'//
1 T3,'FREQ' '1P10F12.4)
289 WRITE (3,1033) (GPPF(I),I=1,NFREQ)
1033 FORMAT (/T2,'GFN MASS'1P10F12.4)
    WRITE (3,1034) (ITN(I),I=1,NFREQ)
1034 FORMAT (/T3,'ITER'10I12)
    WRITE (3,1004)
    DO 290 I=1,NU
290 WRITE (3,1017) USED(I),(PHIC(I,J),J=1,NFREQ)
    IFIRST=IFIRST+1
    GO TO 200
300 READ (1,3031) (I,J,KTEMP(I,J),L=1,NC)
2031 FORMAT (4(2I2,F16.0))
    CALL SYM (KTEMP,NU)
    WRITE (3,2031)
2031 FORMAT (T50,'K MATRIX CHANGES')
    CALL MOUT2 (KTEMP,NU,NU)
C
                                CALCULATE DELTA K INCOMPLETE
    CALL MPPY (KTEMP,PHI,NU,NU,NFREQ,DK)
    DO 310 I=1,NU
    DO 310 J=1,NFREQ
310 MINV(J,I)=PHI(I,J)
    CALL MPPY (MINV,DK,NFREQ,NU,NFREQ,DCDF)
    DO 320 I=1,NU
    DO 320 J=1,NU
320 DK(I,J)=0
    DO 330 IJ=1,NFREQ
    DO 330 JI=1,NFREQ
    DO 330 I=1,NU
    DO 330 J=1,NU
330 DK(I,J)=DK(I,J)+PHI(I,IJ)*PHI(J,JI)*CCDF(IJ,JI)/GMASS(IJ)/
1 GMASS(JI)
    CALL MPPY (MINC ,DK,NU,NU,NU,KTEMP)
    CALL MPPY (KTEMP,MINC ,NU,NU,NU,DK)
    WRITE (3,1040)
1040 FORMAT ('1'T50,'DELTA K INCOMPLETE')
    CALL MOUT2 (DK,NU,NU)
    DO 340 I=1,NU
    DO 340 J=1,NU
340 KTEMP(I,J)=KINC(I,J)+DK(I,J)
    CALL INVRS (MINC ,NU,MINV)

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C                                     IIFRATE USING MINV*KTEMP
      DO 380 II=1,NFREQ
      I=NFREQ-II+1
      CALL MITER(MINV ,KTEMP,NU,.00001,25,MODE,OMEG,IT)
      CON=MODE(NROWN)
      DO 375 J=1,NU
      MODE(J)=MODE(J)/CON
375  PHIC(J,I)=MODE(J)
      ITN(I)=IT
      CALL MOPY (MINC ,MODE,NU,NU,1,MODM)
      FREQC(I)=   SQRT(OMEG)/6.2832
      GPPE(I)=GFN(MODE,MINC ,NU)
      CON=OMEG/GPPE(I)
      DO 380 L=1,NU
      DO 380 J=1,NU
380  KTEMP(L,J)=KTEMP(L,J)-CON*MODM(L)*MODM(J)
      GJ TO 285
      END
```

```
SUBROUTINE MOUT2 (A,M,N)
REAL A(20,21)
ID=MIND(N,10)
WRITE (3,1000) (I,I=1,ID)
1000 FORMAT (/T5,10I12)
WRITE (3,1000)
DO 10 I=1,M
10 WRITE (3,1001) I,(A(I,J),J=1,ID)
1001 FORMAT (15,5X,1P10F12.4)
IF (ID-N) 20,50,50
20 WRITE (3,1000) (I,I=11,N)
WRITE (3,1000)
DO 30 I=1,M
30 WRITE (3,1001) I,(A(I,J),J=11,N)
50 RETURN
END
```

SUBROUTINE CINV (A,R,N,C,D)

C+I*D = INVERSE OF A+I*B I=SQRT(-1)

B ASSUMED NON SINGULAR

REAL A(20,21),B(20,21),C(20,21),D(20,21),E(20,21)
 CALL INVRS(B,N,C)
 CALL MPMY(C,A,N,N,N,E)
 CALL MPMY(A,E,N,N,N,C)
 DO 10 I=1,N
 DO 10 J=1,N
 10 C(I,J)=C(I,J)+R(I,J)
 CALL INVRS(C,N,D)
 CALL MPMY(E,D,N,N,N,C)
 DO 20 I=1,N
 DO 20 J=1,N
 20 D(I,J)=-D(I,J)
 RETURN
 END

```
SUBROUTINE SYM (A,N)
      FORMS SYMMETRIC MATRIX FROM LOWER TRIANGLE
      REAL A(20,21)
      N1=N-1
      DO 10 I=1,N1
      I1=I+1
      DO 10 J=I1,N
10  A(I,J)=A(J,I)
      RETURN
      END
```

DISK OPERATING SYSTEM/360 FCPTRAN 360N-FO-451 31

SUBROUTINE IDMASS (IC,HEAD,FR,PHI,P,N,MASS,CINC,KINC,GMASS)
 USES PSEUDO INVERSE TO OBTAIN UNKNOWN MASSES AND
 INCOMPLETE STIFFNESS AND INFLUENCE COEFFICIENT MATRICES
 SEE MAIN PROGRAM FOR INPUT DESCRIPTION

INPUT ERROR MESSAGES

MASS 'I' INDEX TOO LARGE
 MASS 'I' I LESS THAN J
 MASSES 'I' AND 'J' HAVE DUPLICATE INDICES
 WEIGHTING FUNCTION 'I' = 0
 APPROX MASS 'I' AND KNOWN MASS 'J' HAVE SAME INDICES

INTEGER HEAD(20),P,IGM(10),IKM(20),JKM(20),IAM(20),JAM(20)
 REAL PHI(20,10),GM(10),KM(20),AM(20),W(20),A(20,20),R(20),
 1 AW(20,20),RR(20,21),HM(20,20),MMIN(20),MPAR(20),DFL(20),
 2 FR(10),MASS(20,20),CCOFF(10),KCCFF(10),CINC(20,20),KINC(20,20)
 REAL GMASS(10)

INPUT

10 WRITE (3,1001) IC,HEAD
 1001 FORMAT ('I'//T10,'PART II - IDENTIFICATION'//T5,15(' '),12,3X,
 1 A3,19A4,5X,15(' ')//)
 IFR=0
 IERR=0
 ICC =IC+1
 1002 FORMAT (2I10)
 GO TO (41,42),ICC
 41 READ (1,1002) NG
 42 IF(NG) 50,50,60
 50 WRITE (3,1007)
 1007 FORMAT (//T20,'ALL GENERALIZED MASSES UNKNOWN')
 GO TO 70
 60 GO TO (61,62),ICC
 61 READ (1,1008) (IGM(I),GM(I),I=1,NG)
 1008 FORMAT (8(I2,F8.0))
 62 WRITE (3,1009) (IGM(I),GM(I),I=1,NG)
 1009 FORMAT (//T20,'KNOWN GENERALIZED MASSES'//T5, 'M('I2,') = '
 1 IPE10.4,T29,I2,') = 'E10.4,T27,'M('T51,I2,') = 'E10.4,T49,'M(''
 2 T73,I2,') = 'E10.4,T71,'M('T95,I2,') = 'E10.4,T93,'M(''
 70 GO TO (71,72),ICC
 71 READ (1,1002) NK
 72 IF (NK) 80,80,90
 80 WRITE (3,1010)
 1010 FORMAT (//T20,'ALL MASSES UNKNOWN')
 GO TO 100
 90 GO TO (91,92),ICC
 91 READ (1,1011) (IKM(I),JKM(I),KM(I),I=1,NK)
 1011 FORMAT (8(2I2,F6.0))
 92 WRITE (3,1012) (IKM(I),JKM(I),KM(I),I=1,NK)
 1012 FORMAT (//T20,'KNOWN MASSES'//T4, 'M('I2,','I2,') = 'IPE11.4,
 1T31,I2,','I2,') = 'E12.4,T29,'M('T57,I2,','I2,') = 'E12.4,T55,'M(''
 2T81,I2,','I2,') = 'E12.4,T79,'M('T106,I2,','I2,') = 'E12.4,T104,
 3'M(''
 INPUT ERROR TEST
 DO 920 I=1,NK

```

      L=IKM(I)
      K=JKM(I)
      IF(L-P) 901,901,902
901  IF(K-P) 905,905,902
902  IF(IER) 903,903,904
903  WRITE (3,2000)
2000 FORMAT (/T5,'INPUT ERRORS, KNOWN MASSES'/)
904  IER=IER+1
      WRITE (3,2001) I
2001 FORMAT (T10,'MASS'I4,' INDEX TOO LARGE')
905  IF(L-K) 906,910,910
906  IF(IER)907,907,908
907  WRITE (3,2000)
908  IER =IER+1
      WRITE (3,2002) I
2002 FORMAT (T10,'MASS'I4,' I LESS THAN J')
910  IF(I-NK) 911,920,920
911  J1=I+1
      DO 920 J=J1,NK
      IF(L-IKM(J)) 920,912,920
912  IF(K-JKM(J)) 920,913,920
913  IF(IER) 914,914,915
914  WRITE (3,2000)
915  IER=IER+1
      WRITE (3,2003) I,J
2003 FORMAT (T10,'MASSES'I4,' AND'I4,' HAVE DUPLICATE INDICES')
920  CONTINUE
100  READ (1,1013) NV,ICT,(IAM(I),JAM(I),AM(I),W(I),I=1,NV)
1013 FORMAT (I10,9X,I1/(2I5,2F10.0))
                                INPUT ERROR TEST
      DO 950 I=1,NV
      L=IAM(I)
      K=JAM(I)
      IF(L-P) 921,921,922
921  IF(K-P) 925,925,922
922  IF(IERR) 923,923,924
923  WRITE (3,2004)
2004 FORMAT (/T5,'INPUT ERRORS, APPROX MASSES'/)
924  IERR=IERR+1
      WRITE (3,2001) I
925  IF(L-K) 926,929,929
926  IF(IERR) 927,927,928
927  WRITE (3,2004)
928  IERR=IERR+1
      WRITE (3,2002) I
929  IF(W(I)) 935,930,935
930  IF(IERR) 931,931,932
931  WRITE (3,2004)
932  IERR=IERR+1
      WRITE (3,2005) I
2005 FORMAT (T10,'WEIGHTING FUNCTION'I4,' = 0')
935  IF(I-NV)936,942,942
936  J1=I+1
      DO 941 J=J1,NV
      IF(L-IAM(J)) 941,937,941

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01/03/70

IDMASS

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937 IF(K-JAM(J)) 941,938,941
938 IF(IERR)939,939,940
939 WRITE (3,2004)
940 IERR=IERR+1
    WRITE (3,2003) I,J
941 CONTINUE
942 IF(NK) 950,950,943
943 DO 950 J=1,NK
    IF(L-KM(J)) 950,944,950
944 IF(K-JKM(J))950,945,950
945 IF(IERR) 946,946,947
946 WRITE (3,2004)
947 IERR=IERR+1
    WRITE (3,2006) I,J
2006 FORMAT (T10,'APPROX MASS'14,' AND KNOWN MASS'14,' HAVE SAME INDICES
1')
950 CONTINUE
    IF(IER) 951,951,500
951 IF(IERR)101,101,500
101 NZE = N*(N-1)/2
    NE = NZE+NG+MINO(ICT,1)
    NR=NV-NE
    WRITE (3,1014) NF,NV,NR
1014 FORMAT (//T20,'PROBLEM DEFINITION NO OF EQUATIONS = '12/
1 T41,'NO OF VARIABLES = '12/T37,'ORDER OF DEGENERACY = '12)
    IF (NR) 110,120,149
110 WRITE (3,1015)
1015 FORMAT ('+'T65,'NOTE - NO EXACT SOLUTION POSSIBLE')
    GO TO 149
120 WRITE (3,1016)
1016 FORMAT ('+'T65,'NOTE - ONLY ONE SOLUTION POSSIBLE, APPROX MASSES N
NOT USED')
149 IF(ICT)130,151,130
130 WRITE(3,2016)
2016 FORMAT (//T41,'TOTAL OF DIAGONAL MASSES (CONSTANT')
151 IF(NE-20) 159,159,152
152 WRITE (3,3016)
3016 FORMAT (//T10,'*** TOO MANY EQUATIONS, LIMIT IS 20 ***')
    GO TO 500
159 GO TO (150,250),ICC
    FORM A MATRIX (NE X NV), 3 VECTOR (NF)
150 DO 160 I=1,NE
160 R(I)=0
    OFF-DIAGONAL EQUATIONS
    IE=0
    N1=N-1
    DO 200 I=1,N1
    J1=I+1
    DO 200 J=J1,N
    IF=IE+1
    DO 180 IV=1,NV
    L=IAM(IV)
    K=JAM(IV)
    IF(K-L) 170,165,170
165 A(IF,IV)=PHI(L,I)*PHI(L,J)

```

01/03/70

INMASC

0004

```

GO TO 190
170 A(IE,IV)=PHI(L,I)*PHI(K,J)+PHI(K,I)*PHI(L,J)
180 CONTINUE
    IF(NK)200,200,185
185 DO 199 IK=1,NK
    L=IKM(IK)
    K=JKM(IK)
    IF(K-L) 195,190,195
190 R(IE)=R(IE)-KM(IV)*PHI(L,I)*PHI(L,J)
    GO TO 199
195 R(IE)=R(IE)-KM(IK)*(PHI(L,I)*PHI(K,J)+PHI(K,I)*PHI(L,J))
199 CONTINUE
200 CONTINUE
    IF(NG) 250,250,210

```

EQS CORRESPONDING TO DIAG TERMS

```

210 DO 240 II=1,NG
    I=IGM(II)
    IE=IE+1
    DO 230 IV=1,NV
    L=IAM(IV)
    K=JAM(IV)
    IF (K-L) 225,220,225
220 A(IE,IV)=PHI(L,I)*PHI(L,I)
    GO TO 230
225 A(IE,IV)=2.0*PHI(L,I)*PHI(K,I)
230 CONTINUE
    R(IE)=GM(II)
    IF(NK) 240,240,231
231 DO 239 IK=1,NK
    L=IKM(IK)
    K=JKM(IK)
    IF(K-L) 235,232,235
232 R(IE)=R(IE)-KM(IK)*PHI(L,I)*PHI(L,I)
    GO TO 239
235 R(IE)=R(IE)-2.0*KM(IK)*PHI(L,I)*PHI(K,I)
239 CONTINUE
240 CONTINUE
250 IF(ICT)251,260,251

```

EQUATION FOR SUM OF DIAGONAL MASSES

```

251 IE=IE+1
    R(IE)=0
    DO 259 IV=1,NV
    L=IAM(IV)
    K=JAM(IV)
    IF(L<K) 258,252,258
252 A(IE,IV)=1.0
    R(IE)=R(IE)+AM(IV)
    GO TO 259
258 A(IE,IV)=0
259 CONTINUE

```

A MATRIX OUTPUT

```

260 WRITE (3,1017) NF,NV
1017 FORMAT ('1//T20, 'A MATRIX',4, ' X',13//)
    CALL MOUT2 (A,NF,NV)
    WRITE (3,1019)

```


01/03/70

IDMASS

0005

```

1019 FORMAT (///T20,' VECTOR'//)
WRITE (3,1020)(P(I),I=1,NF)
1020 FORMAT (/(T6,1P10E12.4))
C
FORM AW = A/W
DO 270 J=1,NV
W1=1.0/W(J)
DO 270 I=1,NF
270 AW(I,J)=A(I,J)*W1
C
FORM BR
CALL PSEUDO (AW,NE,NV,BR)
DO 280 I=1,NV
W1=1.0/W(I)
DO 280 J=1,NF
280 BR(I,J)=BR(I,J)*W1
C
FORM BM
CALL MPMY (BR,A,NV,NF,NV,BM)
DO 290 I=1,NV
DO 290 J=1,NV
BM(I,J)=-BM(I,J)
IF(I-J) 290,285,290
285 BM(I,J)=BM(I,J)+1.0
290 CONTINUE
WRITE (3,1021) NV,NF
1021 FORMAT ('1'//T20,' BR MATRIX'14,' X'13//)
CALL MOUT2 (BR,NV,NF)
WRITE (3,1022) NV,NV
1022 FORMAT ('1'//T20,' BM MATRIX'14,' X'13//)
CALL MOUT2 (BM,NV,NV)
C
FORM SOLUTION
CALL MPMY (BR,R,NV,NE,1,MMIN)
CALL MPMY (BM,AM,NV,NV,1,MMAR)
TM=0
TMB=0
DELS=0
DO 310 I=1,NV
MMAR(I)=MMAR(I)+MMIN(I)
TM=TM+AM(I)
TMB=TMB+MMAR(I)
DEL(I)=MMAR(I)-AM(I)
310 DELS=DELS+DEL(I)*DEL(I)
DELS=SQRT(DELS/NV)
C
MASS OUTPUT
WRITE (3,1023)
1023 FORMAT ('1'//T30,' REST MASS SOLUTION'//T5,' I J APPROX MASS W
1GT FACT MIN MASS REST MASS CHANGE'//)
WRITE (3,1024) (IAM(I),JAM(I),AM(I),W(I),MMIN(I),MMAR(I),DEL(I),
I I=1,NV)
1024 FORMAT (15,14,1PE15.4,0PF10.4,1P3E15.4)
WRITE (3,1025) TM,TMB,DELS
1025 FORMAT (//T4,' TOTALS'1PE15.4,T50,E15.4//T4,' RMS OF CHANGES'1E12.4)
C
FORM COMPLETE MASS MATRIX
DO 320 I=1,P
DO 320 J=1,P
320 MASS(I,J)=0
DO 330 I=1,NV

```

01/03/70

TDMASS

0004

```

L=IAM(I)
K=JAM(I)
MASS(L,K)=MRAP(I)
330 MASS(K,L)=MRAR(I)
IF(NK) 355,355,340
340 DO 350 I=1,NK
L=IKM(I)
K=JKM(I)
MASS(L,K)=KM(I)
350 MASS(K,L)=KM(I)
355 WRITE (3,1026)
1026 FORMAT ('1'//T30, 'FULL MASS MATRIX'//)
C          FORM GENERALIZED MASS MATRIX
C          (RM,BR DESTROYED)
DO 370 I=1,P
DO 370 J=1,N
370 BM(J,I)=PHI(I,J)
CALL MMPY (BM,MASS,N,P,P,BR)
CALL MMPY (BR,PHI,N,P,N,BM)
WRITE (3,1027)
1027 FORMAT ('1'//T30,'GENERALIZED MASS MATRIX'//)
CALL MOUT2 (BM,N,N)
DO 375 I=1,N
375 GMASS(I)=BM(I,I)
C          FORM C, K MATRICES
DO 390 I=1,N
OM=FR(I)*FP(I)*39.4784
CCOEF(I)=1.0/(OM*BM(I,I))
390 KCOEF(I)=OM/BM(I,I)
DO 395 I=1,P
DO 395 J=1,P
CINC(I,J)=0
395 KINC(I,J)=0
DO 400 IN=1,N
DO 400 I=1,P
DO 400 J=1,P
CINC(I,J)=CINC(I,J)+CCOEF(IN)*PHI(I,IN)*PHI(J,IN)
400 KINC(I,J)=KINC(I,J)+KCOEF(IN)*BR(IN,I)*BR(IN,J)
WRITE (3,1029) (KCOEF(I),I=1,N)
1029 FORMAT ('1'//T30,'INCOMPLETE STIFFNESS COEF'//T10,'MODAL COEFFICIE
INTS'/T6,1P10F12.4)
CALL MOUT2 (KINC,P,P)
WRITE (3,1028) (CCOEF(I),I=1,N)
1028 FORMAT ('1'//T30,'INCOMPLETE INFLUENCE COEF'//T10,'MODAL COEFFICIE
INTS'/T6,1P10F12.4)
CALL MOUT2 (CINC,P,P)
500 RETURN
END

```

SUBROUTINE MITER (A,B,N,TOL,ITMAX,FUN,VAL,IT)

ITERATES ON A*B FOR DOMINANT EIGENFUNCTION (FUN)
AND EIGENVALUE (VAL).

N IS ORDER

TOL IS DECIMAL (.01 PERCENT) TOLERANCE ON VAL.

ITMAX IS MAX NO OF ITERATIONS.

IT IS NUMBER OF ITERATIONS PERFORMED.

A,B ARE SQUARE OF ORDER N (DIMENSIONED (20,21)).

USES MPMY (A,B,N1,N2,N3,C)

REAL A(20,21),B(20,21),C(20,21),DUM(20),FUN(20)

CALL MPMY (A,B,N,N,N,C)

VAL0=100.

IT=1

DO 10 I=1,N

10 FUN(I)=1.0

11 CALL MPMY (C,FUN,N,N,1,DUM)

VAL=DUM(1)

DO 20 I=2,N

IF(ABS(VAL)-ABS(DUM(I)))15,20,20

15 VAL=DUM(I)

20 CONTINUE

DO 30 I=1,N

30 FUN(I)=DUM(I)/VAL

IF(ABS(VAL/VAL0-1.0)-TOL) 50,50,40

40 IT=IT+1

VAL0=VAL

IF(IT-ITMAX) 11,11,50

50 RETURN

END

FUNCTION GEN (FUN,A,N)

GEN = FUN(TRANS) * A * FIN

DIMENSION A(20,21),FUN(20)

GEN=0

DO 20 I=1,N

DUM=0

DO 10 J=1,N

10 DUM=DUM+A(I,J)*FUN(J)

20 GEN=GEN+DUM*FUN(I)

RETURN

END

SUBROUTINE PSFUDC (A, NR, NC, C)

C = PSEUDOINVERSE OF A A UNDISTURBED
 A IS A RECTANGULAR MATRIX OF MAXIMAL RANK (NR X NC)
 NR .GT. OR .LT. NC

$C = (A^T A)^{-1} A^T$ OR $A^T (A A^T)^{-1}$

NR, NC MAY NOT EXCEED 20

REAL A(20,20), R(20,21), C(20,21) R = A^T

DO 10 I=1, NR

DO 10 J=1, NC

10 R(J, I) = A(I, J)

IF (NR - NC) 20, 20, 30

NR .LE. NC

C = A A^T

20 CALL MPMY (A, R, NR, NC, NR, C)

A = INV OF C

CALL INVPS (C, NR, A)

C = PSEUDOINVERSE OF A (NC X NR)

CALL MPMY (R, A, NC, NR, NR, C)

GO TO 40

NC .LT. NR

C = A^T A

30 CALL MPMY (R, A, NC, NR, NC, C)

A = INV OF C

CALL INVPS (C, NC, A)

C = PSEUDOINVERSE OF A (NC X NR)

CALL MPMY (A, R, NC, NC, NR, C)

RESTORE A

40 DO 50 I=1, NR

DO 50 J=1, NC

50 A(I, J) = R(J, I)

RETURN

END

```

SUBROUTINE INVR5 (P,N,A)
A = INVERSE OF A      B UNDISTURBED
DIMENSION A(20,21),D(20,21),IROW(21),ICOL(21),B(20,21)
DO 1 I=1,N
DO 1 J=1,N
1 A(I,J)=B(I,J)
M=N+1
DO 7 I=1,N
IROW(I)=I
7 ICOL(I)=I
DO 20 K=1,N
AMAX=A(K,K)
DO 10 I=K,N
DO 10 J=K,N
IF(ABS(A(I,J))-ABS(AMAX))10,9,9
9 AMAX=A(I,J)
IC=I
JC=J
10 CONTINUE
KI=ICOL(K)
ICOL(K)=ICOL(IC)
ICOL(IC)=KI
KI=IROW(K)
IROW(K)=IROW(JC)
IROW(JC)=KI
IF(AMAX) 11,12,11
12 WRITE (3,13)
13 FORMAT(' SOLUTION OF EXISTING MATRIX NOT POSSIBLE')
GO TO 100
11 DO 14 J=1,N
E=A(K,J)
A(K,J)=A(IC,J)
14 A(IC,J)=E
DO 15 I=1,N
E=A(I,K)
A(I,K)=A(I,JC)
15 A(I,JC)=E
DO 16 I=1,N
IF(I-K) 18,17,18
17 A(I,M)=1.
GO TO 16
18 A(I,M)=0.
16 CONTINUE
PVT=A(K,K)
DO 8 J=1,M
8 A(K,J)=A(K,J)/PVT
DO 19 I=1,N
IF(I-K)21,19,21
21 AMULT=A(I,K)
DO 22 J=1,M
22 A(I,J)=A(I,J)-AMULT*A(K,J)
19 CONTINUE
DO 20 I=1,N
20 A(I,K)=A(I,M)

```

01/03/70

INVPS

0002

```
DO 25 I=1,N
DO 24 L=1,N
IF(IRJW(I)-L) 24,23,24
24 CONTINUE
23 DO 25 J=1,N
25 D(L,J)=A(I,J)
DO 26 J=1,N
DO 28 L=1,N
IF(ICDL(J)-L) 28,29,29
28 CONTINUE
29 DO 26 I=1,N
26 A(I,L)=D(I,J)
100 RETURN
END
```

SUBROUTINE MPMY (A,B,N1,N2,N3,C)

C = A * B
A (N1 X N2) B (N2 X N3) C (N1 X N3)

REAL A(20,21),B(20,21),C(20,21)
DO 20 I=1,N1
DO 20 J=1,N3
C(I,J)=0.
DO 20 K=1,N2
20 C(I,J)=C(I,J)+A(I,K)*B(K,J)
RETURN
END

SUBROUTINE RANDU (IX,IY,YFL)
THIS SUBROUTINE IS FROM SSP VERS. 11

C
IY=IX*65539
IF(IY)5,6,6
5 IY=IY+2147483647+1
6 YFL=IY
YFL=YFL*.4656613E-9
RETURN
END

```

SUBROUTINE MOB (M,K,G,N,OM,ZR,ZI,YR,YI)

```

```

CALCULATES COMPLEX IMPEDANCE AND MOBILITY
M IS SQUARE MASS MATRIX
K IS SQUARE STIFFNESS MATRIX
G IS SCALAR STRUCTURAL DAMPING
OM IS FREQUENCY IN HERTZ
N IS ORDER

```

```

IMPEDANCE IS ZR + I*ZI    (I = SQRT(-1))
MOBILITY = YR + I*YI

```

```

ALL SQUARE MATRICES ARE DIMENSIONED (20,21)

```

```

USES CINV, INVS, MNPY

```

```

REAL M(20,21),K(20,21),ZR(20,21),ZI(20,21),YR(20,21),YI(20,21)
OMR=OM*6.283185
CON=G/OMR
DO 10 I=1,N
DO 10 J=1,N
ZR(I,J)=CON*K(I,J)
10 ZI(I,J)=OMR*M(I,J)-K(I,J)/OMR
CALL CINV (ZR,ZI,N,YR,YI)
RETURN
END

```

INCOMPLETE MODEL THEORY - SIMULATED TEST

***** 1

(1.02) - R(3P) - TIC)

TEST 16

ACTUAL SYSTEM PARAMETERS

18 DEG OF FREEDOM, STRUCT DAMPING (COEF = 0.020)

MASS MATRIX

	1	2	3	4	5	6	7	8	9	10
1	5.0000E-02	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0
9	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0
10	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E-01
11	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
12	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
13	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
14	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
15	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
16	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
17	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
18	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0

	11	12	13	14	15	16	17	18
1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
9	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
10	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
11	1.0000E-01	0.0	0.0	0.0	0.0	0.0	0.0	0.0
12	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0	0.0
13	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0	0.0
14	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0	0.0
15	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0	0.0
16	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0	0.0
17	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E-01	0.0
18	0.0	0.0	0.0	0.0	0.0	0.0	0.0	5.0000E-02

REPRODUCIBILITY OF THE ORIGINAL PAGE IS POOR

STIFFNESS MATRIX

	1	2	3	4	5	6	7	8	9	10
1	1.1609E 04	-3.6482E 05	2.9824E 05	-6.8662E 04	1.8239E 04	-5.9123E 03	3.0578E 03	-4.8841E 03	1.0710E 04	-1.5171E 04
2	-3.6483E 05	9.8858E 04	-9.4075E 04	4.1430E 04	-1.1078E 05	3.2076E 04	-1.3473E 04	1.4254E 04	-2.9029E 04	4.2161E 04
3	2.9825E 05	-9.5077E 04	1.4019E 04	-1.0604E 04	4.4588E 05	-1.2378E 05	3.8338E 04	-1.8724E 04	2.8521E 04	-4.4296E 04
4	-6.8673E 04	4.1434E 05	-1.0604E 04	1.4340E 04	-1.0753E 06	4.5381E 05	-1.2184E 05	3.0628E 04	-1.8065E 04	2.6407E 04
5	1.8246E 04	-1.1081E 05	4.4593E 05	-1.0754E 06	1.4442E 06	-1.0756E 04	4.4196E 05	-1.0719E 05	2.5729E 04	-1.8699E 04
6	-5.9138E 03	3.2088E 04	-1.2381E 05	4.5384E 05	-1.0756E 06	1.4330E 04	-1.0583E 06	4.3051E 05	-1.0992E 05	3.7033E 04
7	3.0546E 03	-1.3416E 04	3.8338E 04	-1.2185E 05	4.4196E 05	-1.0583E 06	1.4205E 06	-1.0590E 06	4.4455E 05	-1.2676E 05
8	-4.8811E 03	1.4244E 04	-1.8713E 04	3.0623E 04	-1.0719E 05	4.3051E 05	-1.0590E 06	1.4351E 06	-1.0408E 06	4.6417E 05
9	1.0711E 04	-2.9030E 04	2.8523E 04	-1.6073E 04	2.9739E 04	-1.0993E 05	4.4457E 05	-1.0809E 06	1.4563E 06	-1.0917E 06
10	-1.5173E 04	4.2171E 04	-4.4314E 04	2.6426E 04	-1.8714E 04	3.7091E 04	-1.2679E 06	4.8121E 05	-1.0918E 06	1.4370E 06
11	1.2206E 04	-3.5194E 04	4.0561E 04	-2.9742E 04	2.2118E 04	-2.7963E 04	4.2690E 04	-1.7944E 05	4.5909E 05	-1.0832E 06
12	-6.0173E 03	1.8119E 04	-2.2854E 04	1.9700E 04	-1.6067E 04	1.3475E 04	-1.5627E 04	3.9827E 04	-1.2264E 05	4.5067E 05
13	3.4689E 03	-9.9553E 03	1.1940E 04	-1.0079E 04	8.0799E 03	-6.1097E 03	5.4219E 03	-9.4593E 03	3.1954E 04	-1.1935E 05
14	-4.8616E 03	1.3180E 04	-1.3139E 04	8.8194E 03	-7.2049E 03	4.3200E 03	-9.4684E 03	5.1721E 03	-9.1440E 03	3.1219E 04
15	6.9699E 03	-2.0459E 04	2.1889E 04	-1.5025E 04	1.2695E 04	-1.2311E 04	1.1564E 04	-8.9014E 03	6.4079E 03	-1.1119E 04
16	-6.7368E 03	2.2827E 04	-2.9414E 04	2.4238E 04	-1.9425E 04	1.6086E 04	-1.5127E 04	1.2789E 04	-8.9256E 03	9.8188E 03
17	3.5331E 03	-1.4867E 04	2.7263E 04	-2.2378E 04	1.7175E 04	-1.1511E 04	1.0786E 04	-9.8256E 03	7.3914E 03	-7.2288E 03
18	-7.3075E 02	4.1139E 03	-7.1322E 03	7.9279E 03	-5.9874E 03	3.3915E 03	-2.8271E 03	3.0017E 03	-2.3159E 03	2.0784E 03

	11	12	13	14	15	16	17	18
1	1.2206E 04	-6.0192E 03	3.4726E 03	-4.8691E 03	6.9790E 03	-6.7398E 03	3.5297E 03	-7.2839E 02
2	-3.5192E 04	1.8121E 04	-1.0004E 04	1.3200E 04	-7.0487E 04	2.2829E 04	-1.4850E 04	4.1042E 03
3	4.0548E 04	-2.2849E 04	1.1944E 04	-1.3157E 04	2.1404E 04	-7.8397E 04	2.2227E 04	-7.1152E 03
4	-2.9726E 04	1.9690E 04	-1.0021E 04	8.8205E 03	-1.5013E 04	2.4196E 04	-2.2282E 04	7.9105E 03
5	2.2103E 04	-1.4052E 04	8.0700E 03	-7.1633E 03	1.2659E 04	-1.9363E 04	1.7126E 04	-5.9720E 03
6	-2.2540E 04	1.3451E 04	-8.0955E 03	6.3038E 03	-1.2263E 04	1.6012E 04	-1.1458E 04	3.3765E 03
7	4.2652E 04	-1.5599E 04	5.4145E 03	-5.4619E 03	1.1526E 04	-1.5044E 04	1.0747E 04	-2.8151E 03
8	-1.2942E 05	3.5805E 04	-9.6672E 03	5.1761E 03	-8.8802E 03	1.2750E 04	-9.7980E 03	2.9943E 03
9	4.5905E 05	-1.2262E 05	3.1954E 04	-9.1361E 03	6.3819E 03	-8.8911E 03	7.3660E 03	-2.3305E 03
10	-1.0832E 06	4.5046E 05	-1.1935E 05	3.1200E 04	-1.1088E 04	9.7782E 03	-7.1935E 03	2.7648E 03
11	1.4430E 06	-1.0724E 06	4.4440E 05	-1.1930E 05	3.6708E 04	-1.8882E 04	1.0828E 04	-2.4700E 03
12	-1.0724E 06	1.4354E 06	-1.0704E 06	4.4909E 05	-1.2717E 05	4.5544E 04	-1.8998E 04	4.6229E 03
13	4.4689E 05	-1.0704E 06	2.4360E 06	-1.0748E 06	4.5830E 05	-1.3587E 05	4.2766E 04	-8.9398E 03
14	-1.1930E 05	4.4908E 05	-1.0748E 06	1.4431E 06	-1.0824E 06	4.5971E 05	-1.2125E 05	2.1445E 04
15	3.0217E 04	-1.2715E 05	4.5827E 05	-1.0824E 06	1.4399E 06	-1.0583E 06	4.8736E 05	-6.5872E 04
16	-1.8905E 04	4.5533E 04	-1.3579E 05	4.5452E 05	-1.0581E 06	1.3599E 06	-9.1220E 05	2.4403E 05
17	1.0861E 04	-1.4994E 04	6.7685E 04	-1.2104E 05	4.0709E 05	-9.1205E 05	9.4347E 05	-3.4784E 05
18	-2.8845E 03	4.6258E 03	-8.9110E 03	7.1365E 04	-6.5767E 04	2.4396E 05	-3.4747E 05	1.5429E 05

INFLUENCE COEFFICIENT MATRIX

Table with 10 columns and 20 rows of numerical data. The columns are labeled 1 through 10, and the rows are labeled 1 through 20. Each cell contains a number in scientific notation, such as 1.0000E-06, 9.1654E-07, etc.

DESCRIPTION OF SIMULATED TEST

POINT AT WHICH SYSTEM IS FORCED 7
 NO OF POINTS OF MEASUREMENT 9
 POINTS AT WHICH MEASUREMENTS ARE TAKEN
 3 5 7 9 11 15 17 18
 FREQUENCIES OF EXCITATION - HZ
 8.3160E 00 1.8654E 01 4.9055E 01

 AVE RANDOM AMPLITUDE ERROR 0.0
 BIAS AMPLITUDE ERROR 0.0

 SEED FOR RAND NO GENERATOR 0

EXACT RESPONSE OF ACTUAL SYSTEM, IN/SEC/POUND, NORMALIZED ON REAL DRIVING POINT RESPONSE

REAL PART

FREQ	8.3160E 00	1.8654E 01	4.9055E 01
D.P. RESP	7.4873E-01	3.9852E-01	1.7967E-02
3	0.47153	0.66158	-4.26195
5	0.83368	1.04127	-3.72077
7	1.00000	1.00000	1.00000
9	0.91745	0.60815	4.77066
11	0.57490	0.14263	3.75490
15	-0.75043	0.51650	-0.58588
17	-1.58366	1.39374	1.48294
18	-2.00865	1.87239	2.88484

SIMULATED TEST RESPONSE, REAL MOBILITY WITH ERRORS USED AS PHIS

D.P.	1	2	3
RESP	7.4873E-01	3.9852E-01	1.0967E-02
1	4.7153E-01	6.6158E-01	-4.2619E 00
2	8.3368E-01	1.0412E 00	-3.7208E 00
3	1.0000E 00	1.0000E 00	1.0000E 00
4	9.1745E-01	6.0815E-01	4.7707E 00
5	5.7490E-01	1.4263E-01	3.7549E 00
6	-7.5043E-01	5.1650E-01	-5.8588E-01
7	-1.5837E 00	1.3937E 00	1.4829E 00
8	-2.0087E 00	1.8724E 00	2.8848E 00

PART II - IDENTIFICATION

***** 0

1(02) - R(3R) - 1(C)

TFST 16

ALL GENERALIZED MASSES UNKNOWN

ALL MASSES UNKNOWN

PROBLEM DEFINITION NO OF EQUATIONS = 4
NO OF VARIABLES = 8
ORDER OF DEGENERACY = 4

TOTAL OF DIAGONAL MASSES CONSTANT

A MATRIX 4 X 8

1 2 3 4 5 6 7 8

1	3.1145E-01	8.6805E-01	1.0000E 00	5.5794E-01	8.1998E-02	-3.8760E-01	-2.2072E 00	-3.7610E 00
2	-2.0096E 00	-3.1019E 00	1.0000E 00	4.3768E 00	2.1587E 00	4.3966E-01	-2.3485E 00	-5.7946E 00
3	-2.8196E 00	-3.8741E 00	1.0000E 00	2.9013E 00	5.3556E-01	-3.0261E-01	2.0668E 00	5.4016E 00
4	1.0000E 00	1.0000E 00	1.0000E 00	1.0000E 00	1.0000E 00	1.0000E 00	1.0000E 00	1.0000E 00

R VECTOR

0.0 0.0 0.0 1.5500E 00

RR MATRIX 8 X 4

	1	2	3	4
1	-4.3193E-02	-4.2510F-03	-7.0886F-02	1.4666F-01
2	1.7264E-01	-9.0493F-02	-5.7130F-03	1.4509F-01
3	4.3394E-01	-1.2120E-01	1.7112F-01	1.3181F-01
4	8.4392F-02	4.0135F-02	6.8805E-02	1.4657F-01
5	-1.8300E-01	9.8495F-02	-7.0437E-02	1.5234E-01
6	-2.7448E-01	1.0553E-01	-1.1863F-01	1.4612F-01
7	-2.7324E-01	6.7839E-02	-8.0967E-02	9.8685F-02
8	8.2929E-02	-9.6052F-02	1.0671E-01	3.2784F-02

BM MATRIX 8 X 8

	1	2	3	4	5	6	7	8
1	6.5846E-01	-3.9691E-01	-2.8270E-02	1.0176E-01	-9.5919E-02	-1.8292E-01	-1.0541E-01	4.9210E-02
2	-3.9691E-01	4.0221E-01	-2.2153E-01	1.7123E-01	3.9159E-02	-4.0117E-02	3.5257E-02	1.0706E-02
3	-2.8270E-02	-2.2153E-01	3.8433E-01	-3.3990E-01	2.6029E-03	1.4145E-01	1.8768E-01	-1.2637E-01
4	1.0176E-01	1.7123E-01	-3.3990E-01	4.3106E-01	-2.7697E-01	-1.1068E-01	-8.2490E-03	3.1745E-02
5	-9.5919E-02	3.9159E-02	2.6029E-03	-2.7697E-01	6.8777E-01	-2.8789E-01	-1.7936E-01	1.1061E-01
6	-1.8292E-01	-4.0117E-02	1.4145E-01	-1.1068E-01	-2.8789E-01	6.6519E-01	-2.5892E-01	7.3882E-02
7	-1.0541E-01	3.5257E-02	1.8768E-01	-8.2482E-03	-1.7936E-01	-2.5892E-01	6.2489E-01	-2.9588E-01
8	4.9209E-02	1.0705E-02	-1.2637E-01	3.1745E-02	1.1061E-01	7.3882E-02	-2.9588E-01	1.4610E-01

BEST MASS SOLUTION

I	J	APPROX MASS	WGT FACT	MIN MASS	BEST MASS	CHANGE
1	1	2.5000E-01	1.0000	2.2723F-01	2.4410E-01	-5.9003E-03
2	2	2.0000E-01	1.0000	2.2489F-01	2.0163F-01	1.6285E-03
3	3	2.0000E-01	1.0000	2.0431F-01	2.1967F-01	1.9667E-02
4	4	2.0000E-01	1.0000	2.2718F-01	2.0853F-01	8.5347E-03
5	5	2.5000F-01	1.0000	2.3612F-01	2.4370F-01	-6.3010E-03
6	6	2.5000F-01	1.0000	2.2649E-01	2.1807E-01	-1.1926F-02
7	7	1.5000E-01	1.0000	1.5296F-01	1.3892F-01	-1.1085F-02
8	8	5.0000E-02	1.0000	5.0815E-02	5.5379E-02	5.3789E-03

TOTALS 1.5500E 00

1.5500E 00

RMS OF CHANGES 1.0191F-02

FULL MASS MATRIX

	1	2	3	4	5	6	7	8
1	2.4410E-01	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	2.0163F-01	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	2.1967F-01	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	2.0853F-01	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	2.4370E-01	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	2.3807E-01	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	1.3892E-01	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	5.5379E-02

GENERALIZED MASS MATRIX

	1	2	3
1	1.3760E 00	-5.9605F-08	2.563CE-06
2	5.9605E-08	1.1547F 00	-1.1921F-07
3	1.8477E-06	-5.9605F-08	1.6475F 01

INCOMPLETE STIFFNESS COEF

MODAL COEFFICIENTS

1.9841E 03 1.1897E 04 5.7663F 03

	1	2	3	4	5	6	7	8
1	6.5775E 03	4.9422F 03	-8.4557E 02	-5.6807F 03	-5.3906E 03	1.0322F 03	-9.1406E 02	-7.8457F 02
2	4.9422F 03	3.8258E 03	-3.2836E 02	-3.9231F 03	-3.8250E 03	8.5094E 02	-4.8096E 02	-4.6923E 02
3	-8.4557E 02	-3.2836E 02	9.4806F 02	1.6750F 03	1.3110E 03	6.6814E 01	6.7104E 02	4.2487E 02
4	-5.6807E 03	-3.9231F 03	1.6750E 03	5.9710E 01	5.3550E 03	-6.8244F 02	1.3904E 03	1.0307F 03
5	-5.3906E 03	-3.8250F 03	1.3110E 03	5.3550E 03	4.8817E 03	-7.3480E 02	1.1059E 03	8.5494E 02
6	1.0322E 03	8.5094E 02	6.6814E 01	-6.8244F 02	-7.3480E 02	3.5540E 02	1.9553E 02	6.2629E 01
7	-9.1406E 02	-4.8096F 02	6.7104F 02	1.3904F 03	1.1059E 03	1.9553E 02	7.8669E 02	4.7717E 02
8	-7.8457E 02	-4.6923F 02	4.2487F 02	1.0307F 03	8.5494E 02	6.2629E 01	4.7717F 02	2.9964E 02

INCOMPLETE INFLUENCE COEFF

MODAL COEFFICIENTS

2.6618E-04 6.3042E-05 6.3892E-07

1 2 3 4 5 6 7 8

1	9.8380E-05	1.5820E-04	1.6450E-04	1.2752E-04	6.7881E-05	-7.1050E-05	-1.4468E-04	-1.9187E-04
2	1.5820E-04	2.6720E-04	2.8517E-04	2.3217E-04	1.2801E-04	-1.3123E-04	-2.6347E-04	-3.2969E-04
3	1.6450E-04	2.8517E-04	3.2986E-04	2.8559E-04	1.6442E-04	-1.6756E-04	-3.3273E-04	-4.1478E-04
4	1.2752E-04	2.3217E-04	2.8559E-04	2.6190E-04	1.5731E-04	-1.6524E-04	-3.2878E-04	-4.0995E-04
5	6.7881E-05	1.2801E-04	1.6442E-04	1.5731E-04	9.8267E-05	-1.1160E-04	-2.2625E-04	-2.8362E-04
6	-7.1050E-05	-1.3123E-04	-1.6756E-04	-1.6524E-04	-1.1160E-04	1.6694E-04	3.6116E-04	4.6112E-04
7	-1.4468E-04	-2.6347E-04	-3.3273E-04	-3.2878E-04	-2.2625E-04	3.6116E-04	7.9144E-04	1.0140E-03
8	-1.9187E-04	-3.2969E-04	-4.1478E-04	-4.0995E-04	-2.8362E-04	4.6112E-04	1.0140E-03	1.3003E-03

RESONANT RESPONSE OF IDENTIFIED SYSTEM

REAL PART

FREQ	0.3160E 00	1.8654E 01	4.7055E 01
D.P. RESP	6.0333E-01	3.6940E-01	9.8470E-03
3	0.47377	0.66153	-4.26025
5	0.83608	1.04131	-3.71752
7	1.00000	1.00000	1.00000
9	0.92231	0.60820	4.77137
11	0.57813	0.14266	3.75450
15	-0.75294	0.51638	-0.58527
17	-1.58554	1.39374	1.48293
19	-2.01776	1.87224	2.88604

FREQUENCIES AND MODES OF IDENTIFIED SYSTEM

FRFO	3.3160E 00	1.8654E 01	4.9955E 01
GEN MASS	1.3760E 00	1.1547E 00	1.6475E 01
ITFR	10	7	3
3	0.47153	0.66158	-4.26195
5	0.83368	1.04122	-3.72076
7	1.00000	1.00000	1.00000
9	0.91744	0.60815	4.77061
11	0.57490	0.14263	3.75489
15	-0.75043	0.51650	-0.58589
17	-1.58366	1.39374	1.48291
18	-2.00865	1.87240	2.88485

PART III - CHANGED MASSES OR STIFFNESSES

***** 0

1(02) - R(30) - TIC) - A TEST 16

MASS MATRIX CHANGES

	1	2	3	4	5	6	7	8
1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	1.00000E+00	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0

NEW FREQUENCIES AND MODES OF MODIFIED SYSTEM

	6.1740E 00	1.5593E 01	4.8525E 01
FRFQ			
GEN MASS	2.1092E 00	7.0792E 00	1.2993E 02
ITER	9	6	4
3	0.48715	0.42100	-12.81041
5	0.85134	1.21949	-11.97098
7	1.00000	1.00000	1.00000
9	0.88426	0.32243	13.88652
11	0.53196	-0.27173	9.80743
15	-0.61529	1.79564	-1.58307
17	-1.26345	4.41410	3.86984
18	-1.59039	5.82078	7.57423

PART III - CHANGED MASSES OR STIFFNESSES

***** 0

1(,02) - 8(70) - 1(C) - 9 TEST 16

MASS MATRIX CHANGES

	1	2	3	4	5	6	7	8
1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E 00

NEW FREQUENCIES AND MODES OF MODIFIED SYSTEM

FREQ	3.9449E 00	1.4835E 01	4.7504E 01
GEN MASS	1.0368E 01	6.6600E-01	2.4805E 01
ITFR	6	5	3
3	0.44550	0.56205	-9.29938
5	0.80324	0.93473	-4.74020
7	1.00000	1.00000	1.00000
9	0.97245	0.75637	5.74816
11	0.65901	0.34192	4.71396
15	-1.02809	-0.03372	-1.64567
17	-2.24498	0.11025	-0.51286
1A	-2.87363	0.20256	0.44880

PART III - CHANGED MASSES OR STIFFNESSES

***** 3

(1.02) - B(1R) - I(C) - K TEST 16

K MATRIX CHANGES

	1	2	3	4	5	6	7	8
1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	1.0000E 03	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0

DELTA K INCOMPLETE

	1	2	3	4	5	6	7	8
1	2.4714E 01	4.1442F 01	5.4242F 01	4.9590F 01	3.0060F 01	-5.1006F 00	3.2559F 00	2.9921E 00
2	4.1442E 01	6.6789F 01	9.3866F 01	7.9922F 01	4.8446E 01	-8.2203F 00	5.2474E 00	4.8222F 00
3	5.8242E 01	9.3866F 01	1.3192F 02	1.1232F 02	6.8087F 01	-1.1553F 01	7.3748F 00	6.7772F 00
4	4.9590F 01	7.9922F 01	1.1232F 02	7.5636F 01	5.7972F 01	-9.8366F 00	6.2722E 00	5.7704F 00
5	7.0060F 01	4.8446E 01	6.8087F 01	5.7972F 01	3.5141E 01	-5.9627F 00	3.8063F 00	3.4979E 00
6	-5.1006F 00	-8.2203F 00	-1.1553F 01	-9.8366F 00	-5.9627F 00	1.0117F 00	-6.4586F-01	-5.9350F-01
7	3.2559E 00	5.2474F 00	7.3748F 00	6.2722F 00	3.8063F 00	-6.4586F-01	4.1228F-01	3.7887E-01
8	2.9921E 00	4.8222F 00	6.7772F 00	5.7704F 00	3.4979E 00	-5.9350F-01	3.7887F-01	3.4816E-01

NEW FREQUENCIES AND MODES OF MODIFIED SYSTEM

PARAMETER	4	9	9
FREQ	9.2806E 00	1.9269E 01	4.9971E 01
GEN MASS	1.6264E 00	1.0250E 00	1.5871E 01
ITER	4	9	9
1	0.45876	0.65373	-4.17724
5	0.81431	1.03218	-3.63470
7	1.00000	1.00000	1.00000
9	0.94080	0.62384	4.69599
11	0.60901	0.16623	1.69139
15	-0.85668	0.44015	-0.57543
17	-1.83513	1.21228	1.45784
18	-2.33705	1.63518	2.83554

PART III - CHANGED MASSES OR STIFFNESSES

***** 3

11.021 - R13R1 - TIC1 - L TEST 16

K MATRIX CHANGES

	1	2	3	4	5	6	7	8
1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
3	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
4	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
5	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
6	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
8	0.0	0.0	0.0	0.0	0.0	0.0	0.0	1.0000E 03

REPRODUCIBILITY OF THE ORIGINAL PAGE IS POOR

FFLTA K INCCPFLPYF

	1	2	3	4	5	6	7	8
1	7.7993E 00	3.2064E 00	-6.9364E 00	-8.8871E 00	-1.0671E 00	-3.8484E 01	-5.9277E 01	-3.1660E 01
2	3.2064E 00	1.3182E 00	-7.6872E 00	-3.6516E 00	-4.3869E-01	-1.9821E 01	-2.4368E 01	-1.3016E 01
3	-6.9364E 00	-2.6872E 00	5.4781E 00	7.4439E 00	8.9430E-01	3.2253E 01	4.9679E 01	2.6533E 01
4	-8.8871E 00	-3.6516E 00	7.4439E 00	1.0115E 01	1.2153E 00	4.3827E 01	6.7501E 01	3.6059E 01
5	-1.0671E 00	-4.3869E-01	8.9433E-01	1.2152E 00	1.4605E-01	5.2653E 00	8.1094E 00	4.3317E 00
6	-3.8484E 01	-1.9821E 01	3.2253E 01	4.3827E 01	5.2652E 00	1.8989E 02	2.9247E 02	1.5622E 02
7	-5.9277E 01	-2.4368E 01	4.9675E 01	6.7501E 01	8.1094E 00	2.9247E 02	4.9049E 02	2.4060E 02
8	-3.1660E 01	-1.3016E 01	2.6533E 01	3.6059E 01	4.3316E 00	1.5622E 02	2.4060E 02	1.2852E 02

NEW FREQUENCIES AND MODES OF MODIFIED SYSTEM

ITEM	3	10	9
FREQ	1.1144E 01	2.1051E 01	4.9195E 01
GEN MASS	9.9799E-01	2.0937E 00	1.6072E 01
3	0.50149	0.75213	-4.19111
5	0.86779	1.13590	-3.65111
7	1.00000	1.00000	1.00000
9	0.85995	0.48712	4.70362
11	0.48955	-0.01144	3.68883
15	-0.47878	0.90316	-0.51111
17	-0.93919	2.28498	1.67507
19	-1.16655	3.02693	3.05987