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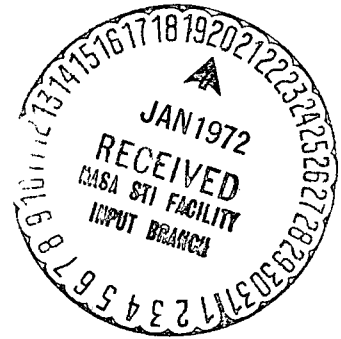
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Lu Ting and Fanny Kung

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DIFFRACTION OF A PLANE WAVE BY A THREE-DIMENSIONAL CORNER^{*}

by

Lu Ting[†] and Fanny Kung^{††}

ABSTRACT

By the superposition of the conical solution for the diffraction of a plane pulse by a three-dimensional corner, the solution for a general incident plane wave is constructed. A numerical program is presented for the computation of the pressure distribution on the surface due to an incident plane wave of any wave form and at any incident angle. Numerical examples are presented to show the pressure signature at several points on the surface due to incident wave with a front shock wave, two shock waves in succession or a compression wave with the same peak pressure. The examples show that when the distance of a point on the surface from the edges or the vertex is comparable to the distance for the front pressure raise to reach the maximum, the peak pressure at that point can be much less than that given by a regular reflection, because the diffracted wave front arrives at that point prior to the arrival of the peak incident wave.

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1. INTRODUCTION

Although the pressure wave created by a supersonic airplane is three-dimensional in nature, the radius of curvature of the wave front is usually much larger than the length scale of a structure. Therefore, the incident waves can be approximated by progressing plane waves composed of compression, expansion and shock waves. By the decomposition of a plane wave to a succession of plane pulses, the basic problem is therefore the diffraction of a plane pulse by a three-dimensional corner.

After the incidence of a plane pulse on a three-dimensional corner at the instant $t = 0$, the disturbed regions behind the incident plane wave are either a simple reflection from the surface of the corner, a two-dimensional diffraction by an edge or the three-dimensional diffraction by the vertex as shown in Fig. 1. The last region is confined by a sonic sphere $r = Ct$, centered at the vertex. The solution for the diffraction by an edge is a two-dimensional conical solution obtained by Keller and Blank [1]. Once the appropriated two-dimensional solutions corresponding to the incident angles are constructed, the boundary data on the sonic sphere about the vertex are obtained, and the three-dimensional conical solution inside the sonic sphere is constructed by the determination of the eigenfunctions and their coefficients [2]. In the next section, the essential procedure and the equations required for the numerical program I are presented for the computation of the following items: i) the two-dimensional conical solution for each edge corresponding to the direction cosines of the incident wave, ii) the boundary data on the sonic sphere around the vertex and iii) the coefficients

in the eigenfunction expansions of the solution inside the sonic sphere.

In section 3, the superposition of the solution of the diffraction of a plane pulse to that of a plane wave of a given wave form is described. The superposition is carried out by numerical program II with the coefficients of eigenfunctions expansions from program I and the incident wave form as input data. Numerical results are presented showing the pressure signatures received by several points on the surface of the corner corresponding to various types of incident wave forms.

Both the numerical program I and II are described and listed in the appendix.

2. INCIDENCE BY A PLANE PULSE

Fig. 1 shows a unit plane pulse incident on a corner of the cube. The three edges are chosen as the three coordinate axes x_j with $j = 2, 3, 4$ *. The direction cosines of the normal to the incident pulse are designated as n_j with $n_2^2 + n_3^2 + n_4^2 = 1$. The equation for the plane of the incident pulse is

$$H_0 = n_2 \bar{x}_2 + n_3 \bar{x}_3 + n_4 \bar{x}_4 = 1 \quad (1)$$

$$\text{with } \bar{x}_j = x_j / (Ct)$$

where C is the speed of sound and t is the time after the passing of the plane pulse over the vertex.

If the plane pulse hit the j - th edge first before hitting the vertex, n_j will be negative. This happens when the incident pulse diffracted by the first corner of the cube is subsequently diffracted by the adjacent corners.

* They begin with 2 so that the index $j-1$ will never be zero which will not be accepted by the computing machine.

For the incidence of a plane wave with the first corner of a building,
 n_j 's are nonnegative, i.e.,

$$n_j \geq 0 \quad j = 2, 3, 4 \quad (2)$$

In order to avoid the complicated discussion for various cases when one,
two, or all of the n_j 's are negative, the discussions in this report
will be restricted to the case of Eq. (2).

The plane pulse is intercepted by the j -axis at $X_j = 1/n_j$, and inter-
sects the x_j - x_{j-1} plane along the line

$$n_{j-1} \bar{x}_{j-1} + n_j \bar{x}_j = 1 \quad (3)$$

For the convenience of programming, quantities with subscript 1 and 5 are
identified with those with subscript 4 and 2 respectively.

The diffraction due to the j -th edge is confined inside the Mach cone
 G_j with vertex at X_j on the \bar{x}_j axis,

$$G_j: (1 - n_j \bar{x}_j) > [(\bar{x}_{j-1})^2 + (\bar{x}_{j+1})^2]^{1/2} (1 - n_j^2)^{1/2} \quad (4)$$

The diffraction by the vertex is confined inside the sonic sphere

$$S: \bar{x}_2^{-2} + \bar{x}_3^{-2} + \bar{x}_4^{-2} < 1 \quad (5)$$

The plane pulse will be reflected by the face, $\bar{x}_j = 0$, (the \bar{x}_{j-1} - \bar{x}_{j+1}
plane) since $n_j > 0$ and the plane of the reflected wave is

$$P_j = n_{j-1} \bar{x}_{j-1} - n_j \bar{x}_j + n_{j+1} \bar{x}_{j+1} = 1 \quad (6)$$

Across the reflected wave i.e., from $P_j > 1$ to $P_j < 1$, the pressure rises

from unity to 2.

Inside the cone G_j but outside and ahead of the sonic sphere S , the solution is a function of two conical coordinates ξ_j and η_j with

$$\xi_j = \frac{\bar{x}_{j-1} (1-n_j^2)^{\frac{1}{2}}}{1-n_j \bar{x}_j} \quad \text{and} \quad \eta_j = \frac{-x_{j+1} (1-n_j^2)^{\frac{1}{2}}}{1-n_j \bar{x}_j} \quad (7)$$

The cone G_j becomes the domain inside a unit circle,

$$\xi_j^2 + \eta_j^2 = 1 \quad .$$

The reflected wave P_{j-1} in ξ, η variables becomes

$$-n_{j-1} \xi_j - n_{j+1} \eta_j = (1-n_j^2)^{\frac{1}{2}} \quad (8)$$

and it is tangential to the unit circle at the point

$$A_j^-: \quad \xi_j^- = -n_{j-1}/(1-n_j^2)^{\frac{1}{2}} \quad \text{and} \quad \eta_j^- = -n_{j+1}/(1-n_j^2)^{\frac{1}{2}} \quad (9)$$

Similarly the reflected wave P_{j+1} becomes

$$n_{j-1} \xi_j + n_{j+1} \eta_j = (1-n_j^2)^{\frac{1}{2}}$$

and the point of contact to the unit circle is

$$A_j^+: \quad \xi_j^+ = n_{j-1}/(1-n_j^2)^{\frac{1}{2}} \quad \text{and} \quad \eta_j^+ = n_{j+1}/(1-n_j^2)^{\frac{1}{2}} \quad (10)$$

as shown in Fig. 2. The boundary condition on the unit circle, $\rho_j=1$, for the disturbance pressure p is

$$\begin{array}{ll} p = 2 & \omega_j^+ > \theta_j > 0 \\ p = 1 & \omega_j^- > \theta_j > \omega_j^+ \\ p = 2 & 3\pi/2 > \theta_j > \omega_j^- \end{array}$$

where $\omega_j^+ = \arcsin \eta_j^+$

$$\omega_j^- = \pi + \omega_j^+$$

and ρ_j, θ_j are the polar coordinates in ξ_j, η_j plane.

The disturbance pressure which lies inside the sonic cone G_j but ahead of the sonic sphere, $x_{j+1}^2 + x_j^2 + x_{j-1}^2 > 1$ and $x_j > n_j$, is given by the two-dimensional conical solution [1],

$$p = 1 + p_j(\rho_j, \theta_j) \text{ for } \rho_j \geq 0 \text{ and } 3\pi/2 > \theta_j > 0 \quad (11)$$

with

$$p_j = p_j^-(\rho_j, \theta_j, \omega_j^-) - p_j^+(\rho_j, \theta_j, \omega_j^+) + 1$$

and $p_j^+(\rho_j, \theta_j, \omega_j) = \frac{1}{\pi} \arctan \frac{(1-\tilde{\rho}^2) \sin(2\omega_j/3)}{2\tilde{\rho} \cos(2\theta_j/3) - (1+\tilde{\rho}^2) \cos(2\omega_j/3)}$ where the arc-tangent lies in the first and second quadrants and $\tilde{\rho} = \left\{ \rho_j / [1 + (1-\rho_j^2)^{1/2}] \right\}^{2/3}$.

In the domain common to the cones G_j and G_{j+1} from two edges j and $j+1$ but outside the sonic sphere S , the disturbance pressure is

$$p = 1 + p_j(\rho_j, \theta_j) + p_{j+1}(\rho_{j+1}, \theta_{j+1}) \quad (12)$$

In the domain ahead the two cones G_j and G_{j+1} , and behind the reflected shock P_{j-1} , the disturbance pressure is

$$p = 2 \quad (13)$$

In the remaining domain outside the sonic sphere and behind the incident shock, the disturbance pressure is

$$p = 1 \quad (14)$$

Ahead of the incident shock, the pressure is of course undisturbed,

$$p = 0 \quad (15)$$

Equations (11) to (15) define the disturbance pressure outside the sonic sphere S.

For the pressure distribution inside the sonic sphere S, the pressure distribution is represented by the eigenfunction expansions,

$$p(\zeta, \mu, \varphi) = \sum_{\lambda} K_{\lambda} Z_{\lambda}(\zeta) G_{\lambda}(\mu, \varphi) \quad (16)$$

where $\zeta = r/(Ct) = (\bar{x}_2^2 + \bar{x}_3^2 + \bar{x}_4^2)^{1/2}$

$$\bar{x}_4 = -\zeta\mu$$

$$\bar{x}_2 = \zeta (1-\mu^2)^{1/2} \sin(\varphi-3\pi/2),$$

and
$$\bar{x}_3 = \zeta (1-\mu^2)^{1/2} \cos(\varphi-3\pi/2).$$

The eigenvalues λ 's, the eigenfunctions $G_{\lambda}(\mu, \varphi)$ and the associated function $Z_{\lambda}(\zeta)$ are defined and are determined by the first numerical program in Ref. 2. The constants $A_{m\lambda}$, $B_{m\lambda}$, $C_{j\lambda}$, $D_{j\lambda}$ characterizing the function $G_{\lambda}(\mu, \varphi)$ and the eigenvalues λ 's are now introduced as the input data for the programs in this report.

For a given set of direction cosines (n_j) of the incident pulse, program I computes the following items: i) the pressure distributions in the various regions outside of the sonic sphere by the appropriate equation of Eqs. (11 to 15), ii) the boundary data, $F(\mu, \varphi)$ on the sonic sphere and iii) the coefficients K_{λ} in Eq. (16) for the solution inside the sonic sphere. The coefficients K_{λ} are related to the boundary data by the equation (see section 5 of [2])

$$K_{\lambda} = \frac{1}{2} \left\{ \int_0^1 d\mu \int_{-\pi}^{\pi} d\varphi F(\mu, \varphi) G(\mu, \varphi) + \int_{-1}^0 d\mu \int_{-3\pi/2}^{3\pi/2} d\varphi F(\mu, \varphi) G(\mu, \varphi) \right\} \quad (17)$$

Program I in this report is a generalization and an extension of the second program in Ref. 2 to compute the coefficients K_{λ} for any incident angle and to compute the pressure distribution outside the sonic sphere. With the knowledge of K_{λ} , the pressure distribution inside the sphere is given by Eq. (16) and are computed by the third program in Ref. 2. Figures 3, 4, and 5 show the pressure distribution on the surface due to the incidence of a unit plane pulse with direction cosines, 0.3, 0.4 and 0.83333. The discontinuities in the slope of the pressure distribution occur at the crossing of the sonic cones around an edge and that of the sonic sphere.

3. INCIDENCE OF A PLANE WAVE

The incident plane wave p_i can be represented in general as

$$p_i = \Psi(s)$$

with $s = Ct - (n_2x_2 + n_3x_3 + n_4x_4)$ where the wave form Ψ is a given function of its phase s and n_j 's are the direction cosines with respect to the axes x_j .

When the wave form is a Heaviside function, the diffraction due to the three-dimensional corner is given by the conical solution described in the preceding section. It will be designated as $p^*(r/(Ct), \mu, \varphi)$. The solution

corresponding to a plane wave of wave form $\Psi(\eta)$ is given by the Stieltjes integral

$$\begin{aligned}
 p(r, \theta, \varphi, t) &= \int p^* \left(\frac{r}{Ct - \eta}, \mu, \varphi \right) d\Psi(\eta) \\
 &= \sum_i p^* \left[\frac{r}{Ct - \eta_j} \right] \left[\Psi(\eta_{j+1}) - \Psi(\eta_j) \right]
 \end{aligned}
 \tag{18}$$

The second form is employed in the numerical program II for the computation of the pressure signature for points in the surface of the corner.

Numerical examples have been carried out for various wave forms with $n_j \equiv 1/\sqrt{3}$. Fig. 6 shows the pressure signature received at points $r = 0, \frac{1}{2}$ and 1 along the line dividing the top surface of the corner ($\mu = 0, \varphi = \pi$). The incident wave is a simple N-wave in sonic boom problems with front shock strength ϵ . The length of the N-wave, which is 4, is nearly the length of an airplane. The unit length scale in the numerical examples is therefore of the order of hundreds of feet. As shown in the figure, the pressure signature at the vertex is $8/7$ times the incident wave form in agreement with the theorem stated in [2]. At point $r = 0.5$, the front part of the pressure signature is equal to twice the incident wave form, i.e. the same as a regular reflection and then decreases from the value of a regular reflection after the arrival of the diffracted waves from the edges and corner. Similar phenomenon is observed for the pressure signature at $r = 1$ with a relative delay in the arrival of the diffracted waves.

Fig. 7 shows the pressure signature at $r = 0.5, 1.0$ and 2.0 along the same line $\mu = 0, \varphi = \pi$, while the pressure raise in the incident wave in Fig. 6 is now spread over a thickness of 0.3 . The peak pressure received at $r = 0.5$ and at $r = 1.0$ is less than twice the total pressure raise, i.e., the

value of regular reflection while, at $r = 2.0$ it is equal to the value of regular reflection. Fig. 8 shows the pressure signatures at the same set of points with the pressure raise in the incident wave spread over a thickness of 0.5. The peak pressure at all these points are less than the value given by a regular reflection.

The differences in the value of the peak pressure at a point on the surface from that of a regular reflection is due to the arrival at the point of the diffracted wave front prior to that of the peak incidence wave front. This is the case for all the points in Fig. 8. In Fig. 7 the pressure raise is faster, so that at $r = 2$, the peak incident wave arrives prior to the diffracted waves and the peak pressure at $r = 2$ is the same as the value of a regular reflection. In Fig. 6, the shock thickness is zero, therefore, the peak value at any surface point is the same as in a regular reflection with the exception of points along the edges and at the vertex which always lies inside the diffracted region.

Fig. 9 shows the pressure signatures at $r = 0.5, 1.0,$ and 2.0 along the same line $\mu = 0, \varphi = \pi$ when the front shock in Fig. 6 is split to two shock waves joined by the expansion wave of thickness 0.3. Although the peak pressure is the same as that in the single shock, the peak pressure received at points $r = 0.5$ and 1.0 are nearly 25% less than that in the case of regular reflection. The peak pressure at point $r = 2.0$ does reach the value of regular reflection i.e. 2ϵ . Fig. 10 shows the pressure signature at the same three points when the front shock is split into two shocks separated by an expansion wave of thickness 0.6. The peak pressure for all three points are now nearly 25% less than the value 2ϵ in a regular reflection. Figs. 9 and 10 again demonstrate that when the diffracted wave front arrives prior to the arrival

of the peak incident wave, the local peak pressure can be much less than the value given by a regular reflection.

The time interval between the arrival at point r along the line ($\mu = 0, \varphi = \pi$) of the leading incident wave and that of the leading diffracted wave from the edge (both edges in the present example with $n_2 = n_3 = n_4 = 1/\sqrt{3}$) is

$$\Delta T_1 = (d/C) [\sqrt{2}-1]/\sqrt{6} \quad (19)$$

The time interval between the former and the arrival of the diffracted wave from the vertex is

$$\Delta T_2 = (r/C) [\sqrt{3}-\sqrt{2}]/\sqrt{3} \quad (20)$$

When the distance between the two split front shock waves is d , the time interval between the arrival of the two shock waves at the point r is

$$\Delta T^* = (d/C) \quad (21)$$

The condition for the point r to receive a peak pressure less than that of a regular reflection is $\Delta T_1 < \Delta T^*$, that is,

$$r/d < \sqrt{6}/(\sqrt{2}-1) = 5.92 \quad (22)$$

In the problem of sonic boom, the length scale of the incident wave is about one quarter the length of an airplane. When the distance between two front shocks is 0.3, the maximum radial distance r allowed by the criterion (22) is nearly 1.8, therefore it is of the order of half an airplane length or hundreds of feet. In other words, when the front shock is split into two shock waves (Figs. 9, 10) or spread out to finite thickness, the area on the surface of the corner within a significant distance from the vertex (of hundreds of feet) will receive a peak pressure much less than the value given by a

regular reflection due to the relief from diffraction by edges and the vertex.

CONCLUSION

In this report, numerical programs are presented for the computation of pressure distribution on the surface of a three-dimensional corner due to an incident plane wave of any wave form at any incident angle. Numerical examples show that the area on the surface of the corner within a significant distance from the vertex which can be (of the order of hundreds of feet) will receive a peak pressure much less than that of a regular reflection when the front shock is split to two shock waves or when the total pressure raise is spread to a finite thickness.

REFERENCES

1. Keller, J.B. and Blank, A., "Diffraction and Reflection of Pulses by Wedges and Corners," *Communication on Pure and Applied Mechanics*, Vol. 4, No. 1, pp. 75-94, June 1951.
2. Ting, L. and Kung, F., "Diffraction of a Pulse by a Three-Dimensional Corner," NASA CR-1728, March 1971.

APPENDIX
 NUMERICAL PROGRAMS

Program I : determination of coefficients K_λ of the eigenfunction from the given eigenvalues of λ 's at a given set of direction cosines (n_j) of the incident pulse

Input Definition

ETA(J) : n_2, n_3, n_4 , the direction cosines of the normal to the incident pulse

II : control constant

II=1 ; for the calculation of K_λ for odd function

II equal to any integer other than one; for the calculation of K_λ for even function

NMAX: }
 LMAX: } number of terms in the eigenfunction

XLAM: λ , eigenvalue

BMIN (J): } associated constants for odd function
 DMIN (L): } associated constants for even function

INPUT FORMAT

```

II = 2,      NMAX,  LMAX                (3I5)
             Set{ XLAM                  (F15.0)
                  1 { BMIN(J)           (5F15.0)
                     DMIN(L)           (5F15.0)
                  }
             }

             Set{ }
             2 { }
             .
             .
             .
             END FILE

II = 1,      NMAX  LMAX
             { }
             { }

             END FILE
  
```


Note that the input data of XLAM, BMIN(J), DMIN(L) are determined by the first program in Ref. [2], and the input data of NMAX, LMAX must be the same as that program. The calculation of K_λ at $\lambda = 0$ for the even function must be calculated first, and the end file cards are used to separate even function and odd function.

Output and Definition

XLAM	: λ	These are input data listing { for identification and verification
BMIN(J)	: coefficients	
DMIN(L)	:	
ETA(2) , ETA(3) , ETA(4)	: n_2, n_3, n_4 , a set of direction cosines (n_j)	
K(LAMDA)	:	K_λ , the coefficients of the eigenfunction of incident pulse.

Program II: to determine the pressure distribution due to plane wave
of a given wave form

Input Definition

R:	r	} Spherical coordinates of pts.
THE:	θ	
PSI:	φ	
ETA(2), ETA(3), ETA(4):	n_2, n_3, n_4 a set of direction cosines (n_j)	
RSTART:	the first value of r for pressure calculation	
TSTART:	the initial value of T for pressure calculation at each value of r	
RMAX:	maximum value of r for the pressure calculation	
TMAX:	maximum value of t for the calculation of pressure distribution at each r	

DELTT: Δt , time increment $T = T_i - n \cdot \Delta t$
 NR: number of points r between RSTART and RMAX
 NT: number of points T between TSTART and TMAX
 NUMBOG: total number of λ for even or odd function
 NFMAX: total numbers of $F(I)$
 F(I): the increment of incident wave form function Ψ between phases S_i and S_{i-1}

INPUT FORMAT

II = 2, NMAX, LMAX (3I5)

NUMBOE (I5)

Set₁ { XLAM
 BMIN(J)
 DMIN(L) } (5F15.0)

Set₂ { }

.
 .
 .

END FILE

II = 1, NMAX LMAX

NUMBOE Set₁ { }

Set₂ { }

.
 .
 .

END FILE

Output

The input data are printed in the first part of the output, and there are NR numbers of tables in the second part. Each table is for each value of r , and the pressure distributions are printed in the first column, and the second column are values of T from TSTART to TMAX.

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C PROGRAM I
C DETERMINATION OF THE COEFFICIENTS OF THE EIGENFUNCTION EXPANSION FROM
C THE INITIAL DATA

DIMENSION F(70),FUP(70,70),FBT(70,70),SFL(70,70)
DIMENSION SUMJ(70),SUML(70),BMIN(70),DMIN(70)
DIMENSION ETA(10),WF1(10),WF2(10),BWF1(10),BWF2(10)
COMMON ETA,WF1,WF2,BWF1,BWF2,XXI
PI=3.1415926

XXI=1.

READ(5,7000) (ETA(I),I=2,4)

7000 FORMAT(3F15.8)

ETA(1)=ETA(4)

ETA(5)=ETA(2)

DO 402 J=2,4

SIGNEP=SIGN(1.,ETA(J+1))

SIGNEM=SIGN(1.,ETA(J-1))

ABSEM=ABS(ETA(J-1))

ABSEP=ABS(ETA(J+1))

EJPM=ETA(J-1)**2+ETA(J+1)**2

IF(SQRT(EJPM)-(1.E-08)) 444,444,445

444 W0=PI/4.

BW0=(5.*PI)/4.

GO TO 446

445 W0=ASIN(ABSEM/SQRT(EJPM))

BW0=ASIN(ABSEP/SQRT(EJPM))

446 CONTINUE

WF1(J)=(PI/2.-W0*SIGNEM)*(2./3.)

WF2(J)=2.*PI-WF1(J)

BWF1(J)=(PI+BW0*SIGNEP)*(2./3.)

BWF2(J)=2.*PI-BWF1(J)

402 CONTINUE

LLMAX=50

IDIM=70

MAXU1=60

MAXTOP=60

MAXU2=60

MAXBOT=60

1000 READ(5,1001) II,NMAX,LMAX

1001 FORMAT(3I5)

IF(ENDFILE 5) 9999,1002

1002 CONTINUE

IF(II .EQ. 1) 101,103

C

C

C ODD FUNCTION

```

101 IIMAX=LMAX+NMAX
    WRITE(6,1003)
1003 FORMAT(1H1*   ODD FUNCTION*)
    NCALL=NMAX
    LCALL=LMAX
    GO TO 104
C
C   EVEN FUNCTION
103 IIMAX=LMAX+NMAX+2
    WRITE(6,1004)
1004 FORMAT(1H1,*   EVEN FUNCTION*)
    NCALL=NMAX+1
    LCALL=LMAX+1
104 CONTINUE
    WRITE(6,7532) ETA(2),ETA(3),ETA(4)
7532 FORMAT(* ETA(2)=*,E12.5,3X,*ETA(3)=*,E12.5,3X,*ETA(4)=*,E12.5)
    PI=3.1415926
    ALPHA=PI/2.
    XMU=0.
    PPHI=2.*(PI-ALPHA/2.)
    XNU1=PI/PPHI
    HALPHI=PPHI/2.
    EPS  =0.000001
    WRITE(6,502) NMAX,LMAX,IIMAX
502 FORMAT(   * NMAX=*,I5,5X,*LMAX=*,I5,5X,*IIMAX=*,I5)
400 READ(5,100) XLAM
    IF(ENDFILE 5) 1000,1111
1111 CONTINUE
    100 FORMAT(5F15.0)
    READ(5,100) (BMIN(I),I=1,NCALL)
    READ(5,100) (DMIN(L),L=1,LCALL)
    U0=0.
    MAXPUS=MAXTOP+1
    INDEX=1
    WRITE(6,32) XLAM
32 FORMAT(   * XLAM=*,E15.8)
    WRITE(6,503)
503 FORMAT(//*   COEFFICIENTS OF EIGENFUNCTION*)
    WRITE(6,304) (BMIN(I),I=1,NCALL)
    WRITE(6,304) (DMIN(L),L=1,LCALL)
304 FORMAT(5E20.8/(5X,5E20.8))
    DELU1=(1.-U0-2.*EPS )/MAXU1
    MAXUST=MAXU1+1
    UK=U0+EPS
    DO 5 I=1,MAXUST
    TOTARG=0.

```

```
CALL FFF(UK,INDEX,MAXTOP,F,II,PPHI,IDIM)
DO 4 J=1,MAXPUS
FUP(I,J)=F(J)
4 CONTINUE
IROW=I
DO 13 J=1,NCALL
IF(II .EQ. 1) 111,113
```

C

C ODD FUNCTION

```
111 CONTINUE
XVAL=J
CALL SININT(FUP,XVAL,IROW,MAXTOP,PI,0,IDIM ,TRIINT)
GO TO 114
```

C

C EVEN FUNCTION

```
113 CONTINUE
XVAL=J-1
CALL COSINT(FUP,XVAL,IROW,MAXTOP,PI,0,IDIM ,TRIINT)
```

```
114 CONTINUE
```

```
SS=XVAL
CALL PPDD(SS,XLAM,UK,LLMAX,PP,DP)
TOTARG=TOTARG+BMIN(J)*PP*TRIINT
```

```
13 CONTINUE
```

```
SUMJ(I)=TOTARG
UK=UK+DELUI
```

```
5 CONTINUE
```

```
TOTEM1=SUMJ(1)
DO 19 I=2,MAXU1
COEF=3.+(-1.)*I
TOTEM1=TOTEM1+COEF*SUMJ(I)
```

```
19 CONTINUE
```

```
TOTEM1=(TOTEM1+SUMJ(MAXUST))*DELUI/3.
MAXPUS=MAXBOT+1
```

```
INDEX=2
```

```
DELUI=(1.+U0-2.*EPS )/MAXU2
```

```
MAXUSM=MAXU2+1
```

```
UK=-1.+EPS
```

```
DO 9 I=1,MAXUSM
```

```
TOTARG=0.
```

```
CALL FFF(UK,INDEX,MAXBOT,F,II,PPHI,IDIM)
```

```
DO 8 J=1,MAXPUS
```

```
FBI(I,J)=F(J)
```

```
8 CONTINUE
```

```
IROW=I
```

```
DO 23 L=1,LCALL
```

```
IF(II .EQ. 1) 121,123
```

C

C ODD FUNCTION

121 CONTINUE

XVAL=(2*L-1)*XNU1

CALL SININT(FBT,XVAL,IROW,MAXBOT,HALPHI,0,IDIM,TRIINT)

GO TO 124

C

C EVEN FUNCTION

123 CONTINUE

XVAL=2*(L-1)*XNU1

CALL COSINT(FBT,XVAL,IROW,MAXBOT,HALPHI,0,IDIM,TRIINT)

124 CONTINUE

SS=XVAL

UF=-UK

CALL PPDD(SS,XLAM,UF,LLMAX,PP,DP)

TOTARG=TOTARG+DMIN(L)*PP*TRIINT

23 CONTINUE

SUML(I)=TOTARG

UK=UK+DELU2

9 CONTINUE

TOTEM2=SUML(1)

DO 29 I=2,MAXU2

COEF=3.+(-1.)**I

TOTEM2=TOTEM2+COEF*SUML(I)

29 CONTINUE

TOTEM2=(TOTEM2+SUML(MAXUSM))*DELU2/3.

EEEM=TOTEM1+TOTEM2

WRITE(6,11)EEEM

11 FORMAT(/20X,*K(LAMDA)=*,E15.8)

WRITE(6,401)

401 FORMAT(////)

GO TO 400

9999 STOP

END

```

SUBROUTINE SININT(FP,XVAL,IROW,MAXST,UPLIM,BOTLIM,IDIM,TRIINT)
C  FILON'S METHOD FOR THE NUMERICAL EVALUATION OF TRIGONOMETRICAL
--C  INTEGRALS--(INTEGRAND=FP(P)*SIN(X*P))
  DIMENSION FP(IDIM,IDIM)
  HH=(UPLIM-BOTLIM)/MAXST
  S2S=0.5*FP(IROW,1)*SIN(XVAL*BOTLIM)
  DO 14 J=3,MAXST,2
    P=BOTLIM+(J-1)*HH
    S2S=S2S+FP(IROW,J)*SIN(XVAL*P)
  14 CONTINUE
  J=MAXST+1
  S2S=S2S+0.5*FP(IROW,MAXST+1)*SIN(XVAL*UPLIM)
  S2SM=0.
  DO 16 J=2,MAXST,2
    P=BOTLIM+(J-1)*HH
    S2SM=S2SM+FP(IROW,J)*SIN(XVAL*P)
  16 CONTINUE
  THE=XVAL*HH
  IF (THE-0.2) 25,21,21
  21 ALPHA=(THE**2+THE*SIN(THE)*COS(THE)-2.*SIN(THE)**2)/THE**3
  BETA=2.*(THE*(1.+COS(THE)**2)-2.*SIN(THE)*COS(THE))/THE**3
  GARM=4.*(SIN(THE)-THE*COS(THE))/THE**3
  GO TO 31
  25 ALPHA=2.*THE**3/45.-2.*THE**5/315.+2.*THE**7/4725.
  BETA=2./3.+2.*THE**2/15.-4.*THE**4/105.+2.*THE**6/567.
  GARM=4./3.-2.*THE**2/15.+THE**4/210.-THE**6/11340.
  31 FA=FP(IROW,1)
  FB=FP(IROW,MAXST+1)
  TRIINT=HH*(-ALPHA*(FB*COS(XVAL*UPLIM)-FA*COS(XVAL*BOTLIM))+BETA*
  I S2S+GARM*S2SM)
  RETURN
  END

```

```

SUBROUTINE COSINT(FP,XVAL,IROW,MAXST,UPLIM,BOTLIM,IDIM,TRIINT)
C   FILON'S METHOD FOR THE NUMERICAL EVALUATION OF TRIGONOMETRICAL
C   INTEGRALS--(INTEGRAND=FP(P)*COS(X*P))
      DIMENSION FP(IDIM,IDIM)
      HH=(UPLIM-BOTLIM)/MAXST
      S2S=0.5*FP(IROW,1)*COS(XVAL*BOTLIM)
      DO 14 J=3,MAXST,2
        P=BOTLIM+(J-1)*HH
        S2S=S2S+FP(IROW,J)*COS(XVAL*P)
14  CONTINUE
      J=MAXST+1
      S2S=S2S+0.5*FP(IROW,MAXST+1)*COS(XVAL*UPLIM)
      S2SM=0.
      DO 16 J=2,MAXST,2
        P=BOTLIM+(J-1)*HH
        S2SM=S2SM+FP(IROW,J)*COS(XVAL*P)
16  CONTINUE
      THE=XVAL*HH
      IF (THE-0.2) 25,21,21
21  ALPHA=(THE**2+THE*SIN(THE)*COS(THE)-2.*SIN(THE)**2)/THE**3
      BETA=2.*(THE*(1.+COS(THE)**2)-2.*SIN(THE)*COS(THE))/THE**3
      GARM=4.*(SIN(THE)-THE*COS(THE))/THE**3
      GO TO 31
25  ALPHA=2.*THE**3/45.-2.*THE**5/315.+2.*THE**7/4725.
      BETA=2./3.+2.*THE**2/15.-4.*THE**4/105.+2.*THE**6/567.
      GARM=4./3.-2.*THE**2/15.+THE**4/210.-THE**6/11340.
31  FA=FP(IROW,1)
      FB=FP(IROW,MAXST+1)
      TRIINT=HH*( ALPHA*(FB*SIN(XVAL*UPLIM)-FA*SIN(XVAL*BOTLIM))+BETA*
IS2S+GARM*S2SM)
      RETURN
      END

```



```
SUBROUTINE FFF(UU,INDEX,MAX,F,II,PPHI,IDIM)
DIMENSION F(IDIM)
PI=3.1415926
IF(INDEX.EQ.1) 4,6
4 DELPSI=PI/FLOAT(MAX)
GO TO 7
6 DELPSI=PPHI/FLOAT(2*MAX)
7 CONTINUE
PSI=0.
MAXPUS=MAX+1
DO 19 N=1,MAXPUS
BPSI=PSI-PPHI/2.
CALL FSF(BPSI,UU,FS,PPHI)
FA=FS
BPSI=2.*PI-PSI-PPHI/2.
CALL FSF(BPSI,UU,FS,PPHI)
FB=FS
IF(II.EQ.1) 14,16
14 F(N)=(FA-FB)/2.
GO TO 17
16 F(N)=(FA+FB)/2.
17 CONTINUE
PSI=PSI+DELPSI
19 CONTINUE
RETURN
END
```

```

SUBROUTINE CON2D(RHO,TAU,W1,W2,SS,PC)
PI=3.1415926
EPSIL=0.000001
CON1=1.-RHO**2
IF(CON1+EPSIL) 10,10,14
10 WRITE(6,11)
11 FORMAT(* 1-RHO**2+EPSIL .LE. ZERO, CHECK THE PROGRAM *)
STOP
14 IF(CON1 )15,15,17
15 GG=1.
GO TO 21
17 GG=(RHO/(1.+SQRT(CON1)))**SS
21 BTAU=SS*TAU
DD=(1.-GG**2)*SIN(0.5*(W2-W1))
CC=(1.+GG**2)*COS(0.5*(W2-W1)) -2.*GG*COS(BTAU-(W2+W1)*0.5)
SQCD=SQRT(CC**2+DD**2)
IF(SQCD-EPSIL) 23,25,25
23 PC=0.5
GO TO 30
25 IF(CC) 28,26,26
26 PC=ASIN(DD/SQCD)/PI
GO TO 30
28 PC=ASIN(DD/SQCD)/PI*(-1.)+1.
30 CONTINUE
RETURN
END

```

```

SUBROUTINE FSF(BPSI,UU,FS,PPH)
DIMENSION ETA(10),X(10),F(10),WF1(10),WF2(10),BWF1(10),BWF2(10)
COMMON ETA,WF1,WF2,BWF1,BWF2,XXI
EPS=1.E-06
PI=3.1415926
CONS23=2./3.
EEPS=1.E-05
199 CONTINUE
X(2)=XXI*SQRT(1.-UU**2)*SIN(BPSI)
X(3)=XXI*SQRT(1.-UU**2)*COS(BPSI)
X(4)=-XXI*UU
TESTX2=ABS(X(2))
TESTX3=ABS(X(3))
TESTX4=ABS(X(4))
IF(TESTX2 .LE. EPS) 191,192
191 X(2)=0.
192 IF(TESTX3 .LE. EPS) 193,194
193 X(3)=0.
194 CONTINUE
198 CONTINUE
X(1)=X(4)
X(5)=X(2)
IF(XXI-1.) 201,203,203
201 WRITE(6,202)
202 FORMAT(* ERROR XXI .LT. 1---CASE A*)
STOP
203 IF(X(2) .GT. 0. .AND. X(3) .GT. 0. .AND. X(4) .GT. 0. )204,206
204 THALPI=ABS(BPSI-1.57079632)
IF(THALPI-EEPS) 208,208,209
208 BPSI=BPSI+EEPS
GO TO 199
209 WRITE(6,205)
205 FORMAT(* ERROR X(2),X(3), AND X(4) .GT. 1E-06---CASE B*)
STOP
206 SUMXN=0.
DO 207 J=2,4
207 SUMXN=SUMXN+ETA(J)*X(J)
IF(SUMXN-1.) 215,213,211
211 FS=0.
GO TO 300
213 FS=0.5
GO TO 300
215 DO 258 J=2,4
IF(X(J)-ETA(J)) 221,223,223
221 F(J)=100.
GO TO 259

```

```

223 VAL1=SQRT(ETA(J-1)**2+ETA(J+1)**2)
    VAL2=1.-ETA(J)*X(J)
    Y=(X(J-1)*VAL1)/VAL2
    Z=(-X(J+1)*VAL1)/VAL2
    RHO=SQRT(Y**2+Z**2)
    IF(RHO-1.) 225,225,251
225 CONTINUE
    IF(RHO-EPS)226,226,228
226 TAU=PI/2.
    GO TO 240
228 CONTINUE
    TAU0=ASIN(ABS(Z/RHO))
    IF(Z) 235,231,231
231 IF(Y) 233,232,232
232 TAU=TAU0
    GO TO 240
233 TAU=PI-TAU0
    GO TO 240
235 IF(Y) 236,236,237
236 TAU=PI+TAU0
    GO TO 240
237 IF(X(J)) 1237,1237,1238
1237 F(J)=100.
    GO TO 259
1238 CONTINUE
    WRITE(6,238) TAU0,TAU,Y,Z
238 FORMAT(* ERROR FOR Y .GT. AND Z .LT. 0----CASE C*/4E20.8)
    WRITE(6,253) J,VAL1,VAL2,RHO,X(2),X(3),X(4),XXI,BPSI ,UU
    STOP
240 CONTINUE
    W1=WF1(J)
    W2=WF2(J)
    BW1=BWF1(J)
    BW2=BWF2(J)
    CALL CON2D(RHO,TAU,W1,W2,CONS23,FIF)
    CALL CON2D(RHO,TAU,BW1,BW2,CONS23,FIIF)
    SIGNEP=SIGN(1.,ETA(J+1))
    SIGNEM=SIGN(1.,ETA(J-1))
243 F(J)= ((1.-FIF)*SIGNEP+FIIF*SIGNEM)
    GO TO 259
251 F(J)=100.
259 CONTINUE
258 CONTINUE
253 FORMAT(I3,9E13.5)
    SUMF=0.
    DO 261 J=2,4

```

```
SUMF=SUMF+F(J)
261 CONTINUE
IF(SUMF-EPS) 260,260,262
260 FS=SUMF+1.
GO TO 300
262 CONTINUE
IF(SUMF-90.) 263,263,265
```

C

C-----INSIDE ALL THREE CONES

```
263 WRITE(6,264) (F(J),J=1,4),FIF,FIIF
264 FORMAT(* ERROR SUMF .LE. 90---CASE D*/6E20.8)
WRITE(6,200) VAL1,VAL2,Y,Z,RHO,TAU0,W1,W2,FIF,BW1,BW2,FIIF
STOP
265 IF(SUMF-190.) 267,267,269
```

C

C-----INSIDE TWO CONES AND OUTSIDE THE THIRD CONE

```
267 FS=SUMF-100.
GO TO 300
269 IF(SUMF-290.)271,271,272
```

C

C-----INSIDE ONE CONE ONLY

```
271 FS=SUMF-199.
GO TO 300
272 CONTINUE
```

C

C-----OUTSIDE ALL THREE CONES

```
273 DO 279 J=2,4
TEST=ETA(J-1)*X(J-1)-ETA(J)*X(J)+ETA(J+1)*X(J+1)
DJ=ETA(J+1)
DM=ETA(J-1)
IF(X(J) .LE. 0. .AND. TEST .LT. 1.) 275,279
275 IF(X(J+1) .GE. DJ .OR. X(J-1) .GE. DM)281,279
279 CONTINUE
FS=1.
GO TO 300
281 FS=2.
300 CONTINUE
301 FORMAT(* SUB FSF(NEW)* 5X, 5E20.8)
200 FORMAT(6E20.8)
RETURN
END
```

SUBROUTINE PPDD(SS,XLAM,XMU,LLMAX,PP,DP)

C SUBROUTINE PPDD--WITHOUT DP

XNU=SS

EPS=1.E-06

DD=1.

TEM1=1.

ZX=(1.-XMU)/2.

IF(ABS(ZX) .LE. EPS.AND. ABS(XNU) .LE. EPS) 35,31

31 CONTINUE

DO 40 L=2,LLMAX

ZL=L-1

DD=DD*(ZL-XLAM -1.)*(ZL+XLAM)/(ZL**2+ZL*XNU)

TEM1=TEM1+DD*ZX**(L-1)

40 CONTINUE

ARTFL=((1.-XMU)/(1.+XMU))**(XNU/2.)

PP=ARTFL*TEM1

GO TO 41

35 PP=1.

41 CONTINUE

DP=0.

RETURN

END

```

C PROGRAM II
  DIMENSION EDIM(20),RMIN(20),DMIN(20)
  DIMENSION ETA(10),WF1(10),WF2(10),BWF1(10),BW2(10)
  DIMENSION ELAM(40),ZZZ(40,20),PP(400),F(400),GG(40)
  COMMON /BLOCK1/ETA,WF1,WF2,BWF1,BWF2,XXI
  READ(5,7000) (ETA(I),I=2,4)
  READ(5,10) RMAX,TMAX,DELTT ,TSTART,RSTART
7000 FORMAT(3F15.8)
  READ(5,7001) NFMAX,NT,NR
7001 FORMAT(3I5)
  WRITE(6,7002)
7002 FORMAT(1H1)
  ETA(1)=ETA(4)
  ETA(5)=ETA(2)
  READ(5,10) (F(I),I=1,NFMAX)
 10 FORMAT(5F10.0)
  READ(5,10)THE,PSI
  UU=COS(THE)
  LLMAX=50
  PI=3.1415926
  ALPHA=PI/2.
  PPHI=2.*(PI-ALPHA/2.)
  XNUI=PI/PPHI
  BPSI=PSI-PPHI/2.
  DELR=RMAX/NR
  DELT=TMAX/NT
  WRITE(6,1) (ETA(I),I=1,3),RMAX,TMAX,DELTT,DELR,DELT
 1 FORMAT(* ETA(1)=*,E12.5,5X,*ETA(2)=*,E12.5,5X,*ETA(3)=*,E12.5/
 1* RMAX=*,E12.5,5X,*TMAX=*,E12.5,5X,*DELTT=*,E12.5,5X,*DELR=*,
 2E12.5,5X,*DELT=*,E12.5)
  WRITE(6,7005) UU,PSI
 7005 FORMAT(2X,*UU=*,E15.8,5X,*PSI=*,E15.8)
  WRITE(6,7003)
 7003 FORMAT(/3X,*F(I)*)
  WRITE(6,440) (F(I),I=1,NFMAX)
 440 FORMAT(5F20.5)
  DO 442 J=2,4
 5 SIGNEP=SIGN(1.,ETA(J+1))
  SIGNEM=SIGN(1.,ETA(J-1))
  ABSEM=ABS(ETA(J-1))
  ABSEP=ABS(ETA(J+1))
  FJPM=ETA(J-1)**2+ETA(J+1)**2
 3 IF(SQRT(EJPM)-(1.E=08)) 444,444,445
 444 W0=PI/4.
 2 BW0=(5.*PI)/4.
  GO TO 446
 1

```

```

445 W0=ASIN(ABSEM/SQRT(EJPJM))
    BW0=ASIN(ABSEP/SQRT(EJPJM))
446 CONTINUE
    WF1(J)=(PI/2.-W0*SIGNEM)*(2./3.)
    WF2(J)=2.*PI-WF1(J)
    BWF1(J)=(PI+BW0*SIGNEP)*(2./3.)
    BWF2(J)=2.*PI-BWF1(J)

```

```

442 CONTINUE
    NUMB=0.
1000 READ(5,1001) II,NMAX,LMAX
1001 FORMAT(3I5)
    IF(ENDFILE 5) 9999,1002
1002 CONTINUE
    IF(II .EQ. 1) 101,103

```

C ODD FUNCTION

```

101 IIMAX=LMAX+NMAX
    READ(5,1001) NUMBOE
    NBBB=0
    WRITE(6,2)
    2 FORMAT(1H1,//* ODD FUNCTION*)
    NCALL=NMAX
    LCALL=LMAX
    GO TO 104

```

C EVEN FUNCTION

```

103 IIVAX=LMAX+NMAX+2
    READ(5,1001) NUMBOE
    NBBB=0
    WRITE(6,3)
    3 FORMAT(1H1,//* EVEN FUNCTION *)
    NCALL=NMAX+1
    LCALL=LMAX+1

```

```

104 CONTINUE
    WRITE(6,502) NMAX,LMAX,IIMAX,ALPHA,PPHI,XNUI
502 FORMAT( * NMAX=*,I5,5X,*LMAX=*,I5,5X,*IIMAX=*,I5/5X, *
    1ALPHA=*,E12.5,5X,*PPHI=*,E12.5,5X,*XNUI=*,E12.5)
110 READ(5,100) XLAM,EEEM
100 FORMAT(5E15.8)
    IF(ENDFILE 5) 1000,1111
1111 CONTINUE
2000 NUMB=NUMB+1
    EOLD=EEEM
    NBBB=NBBB+1
    EEEM=(NUMBOE-(NBBB-1))*EEEM/NUMBOE
    READ(5,100) (BMIN(I),I=1,NCALL)

```



```

READ(5,100) (DMIN(L),L=1,LCALL)
CALL T102(XLAM,1.,EDIM,BBBBEB)
AAAAAA=BBBBBB
WRITE(6,32) XLAM,AAAAAA,EOLD
32 FORMAT( * XLAM=*,F14.9,5X,*AA0=*,F14.9,5X,*K(LAMDA)=*,F14.9)
WRITE(6,503)
503 FORMAT(// * COEFFICIENTS OF EIGEN-FUNCTIONS*)
WRITE(6,304) (BMIN(I),I=1,NCALL)
WRITE(6,304) (DMIN(L),L=1,LCALL)
304 FORMAT(5E20.8/(5X,5E20.8))
WRITE(6,401)
401 FORMAT(////)
ELAM(NUMB)=EEEM
CALL T102(XLAM,AAAAAA,EDIM,BBBBBB)
DO 2001 I=1,10
2001 ZZZ(NUMB,I)=EDIM(I)
IF(UU) 2100,2002,2002
2002 CONTINUE
IF(II .EQ. 1) 200,203
C
C ODD FUNCTION
200 XMU=UU
PHI=PSI
PPUS=0.
DO 201 J=1,NCALL
SS=J
CALL PPDD(SS,XLAM,XMU,LLMAX,PP,DP)
SINFU=SIN(J*PHI)
PPUS=PPUS+BMIN(J)*PP*SINFU
201 CONTINUE
GG(NUMB)=PPUS
GO TO 110
C
C EVEN FUNCTION
203 CONTINUE
XMU=UU
PHI=PSI
PPUS=0.
DO 206 J=1,NCALL
SS=J-1
CALL PPDD(SS,XLAM,XMU,LLMAX,PP,DP)
COSFU=COS(SS*PHI)
PPUS=PPUS+BMIN(J)*PP*COSFU
206 CONTINUE
GG(NUMB)=PPUS
GO TO 110

```

2100 IF(II .EQ. 1) 301,400

C

C ODD FUNCTION

301 XMU=UU

PHI=PSI

PMIN=0.

DO 302 L=1,LCALL

SS=(2*L-1)*XNU1

CALL PPDD(SS,XLAM,-XMU,LLMAX,PP,DP)

SINFU=SIN(SS*PHI)

PMIN=PMIN+DMIN(L)*PP*SINFU

302 CONTINUE

GG(NUMB)=PMIN

GO TO 110

C

C EVEN FUNCTION

400 XMU=UU

PHI=PSI

PMIN=0.

DO 402 L=1,LCALL

SS=2*(L-1)*XNU1

CALL PPDD(SS,XLAM,-XMU,LLMAX,PP,DP)

COSFU=COS(SS*PHI)

PMIN=PMIN+DMIN(L)*PP*COSFU

402 CONTINUE

GG(NUMB)=PMIN

GO TO 110

9999 MAXNUM =NUMB

DO 2006 I=1,10

2006 CONTINUE

R=-DELR +RSTART

2200 R=R+DELR

IF(R-RMAX) 2202,2202,2201

2201 STOP

2202 CONTINUE

WRITE(6,7004) R

7004 FORMAT(/75X,*R=*,E12.5,11X,*P*,17X,*T*)

T=-DELT+TSTART

2525 T=T+DELT

IF(T-TMAX) 2602,2602,2200

2602 CONTINUE

DO 2400 N=1,NFMAX

ST=T-N*DELTT +0.0001

IF(ST) 2203,2203,2205

2203 BIGP=0.

GO TO 2400

```
2205 XXI=R/ST  
IF(XXI-1.) 2300,2207,2207  
2207 CONTINUE  
CALL FSF(BPSI,UU,FS,PPHI)  
BIGP=FS  
GO TO 2400  
2300 BIGP=0.  
DO 2309 I=1,MAXNUM  
CALL INTER(XXI,I,ZZZ,ZXXI)  
BIGP=BIGP+ELAM(I)*ZXXI*GG(I)  
2309 CONTINUE  
2400 PP(N)=BIGP  
SUMFP=0.  
DO 2405 M=1,NFMAX  
2405 SUMFP=SUMFP+ PP(M)*F(M)  
WRITE(6,2407) SUMFP,T  
2407 FORMAT(19X,2E20.8)  
GO TO 2525  
2501 STOP  
END
```

```
SUBROUTINE INTER(XXI,ICOL,ZZZ,ZXXI)
DIMENSION ZZZ(40,20)
```

```
N=0
```

```
X=0.
```

```
10 N=N+1
```

```
X=X+0.1
```

```
IF(X-XXI) 10,20,30
```

```
20 ZXXI=ZZZ(ICOL,N)
```

```
GO TO 50
```

```
30 IF(N .GT. 10) GO TO 40
```

```
IF(N .LE. 1) GO TO 60
```

```
ZXXI=ZZZ(ICOL,N-1)+(ZZZ(ICOL,N)-ZZZ(ICOL,N-1))*(XXI-X+0.1)/0.1
```

```
GO TO 50
```

```
40 WRITE(6,41) XXI
```

```
41 FORMAT(//* ERROR IN SUB. INTER --FOR XXI=*,E12.5,* IS OUT OF RANGE*)
```

```
STOP
```

```
50 CONTINUE
```

```
C WRITE(6,51) XXI,ZXXI,ZZZ(ICOL,N-1),ZZZ(ICOL,N)
```

```
51 FORMAT(* SUB INTER*, 5E20.8)
```

```
GO TO 65
```

```
60 IF(ICOL .GT. 1) GO TO 61
```

```
ZZZZZZ=1.
```

```
GO TO 64
```

```
61 ZZZZZZ=0.
```

```
64 CONTINUE
```

```
ZXXI=(ZZZ(ICOL,1)-ZZZZZZ)*XXI+ZZZZZZ
```

```
65 CONTINUE
```

```
RETURN
```

```
END
```

```

SUBROUTINE FSF(BPSI,UU,FS,PPhi)
DIMENSION ETA(10),X(10),F(10),WF1(10),WF2(10),BWF1(10),BWF2(10)
COMMON /BLOCK1/ETA,WF1,WF2,BWF1,BWF2,XXI
EPS=1.E-06
PI=3.1415926
CONS23=2./3.
EEPS=1.E-05
199 CONTINUE
X(2)=XXI*SQRT(1.-UU**2)*SIN(BPSI)
X(3)=XXI*SQRT(1.-UU**2)*COS(BPSI)
X(4)=-XXI*UU
TESTX2=ABS(X(2))
TESTX3=ABS(X(3))
TESTX4=ABS(X(4))
IF(TESTX2 .LE. EPS) 191,192
191 X(2)=0.
192 IF(TESTX3 .LE. EPS) 193,194
193 X(3)=0.
194 CONTINUE
198 CONTINUE
X(1)=X(4)
X(5)=X(2)
IF(XXI-1.) 201,203,203
201 WRITE(6,202) XXI,X(2),X(3),X(4),UU,BPSI
202 FORMAT(* ERROR XXI .LT. 1---CASE A* 6E14.5)
STOP
203 IF(X(2) .GT. 0. .AND. X(3) .GT. 0. .AND. X(4) .GT. 0. )204,206
204 THALPI=ABS(BPSI-1.57079632)
IF(THALPI-EEPS) 208,208,209
208 BPSI=BPSI+EEPS
GO TO 199
209 WRITE(6,205)
205 FORMAT(* ERROR X(2),X(3), AND X(4) .GT. 1E-06---CASE B*)
STOP
206 SUMXN=0.
DO 207 J=2,4
207 SUMXN=SUMXN+ETA(J)*X(J)
IF(SUMXN-1.) 215,213,211
211 FS=0.
GO TO 300
213 FS=0.5
GO TO 300
215 DO 258 J=2,4
IF(X(J)-ETA(J)) 221,223,223
221 F(J)=100.
GO TO 259

```

```

223 VAL1=SQRT(ETA(J-1)**2+ETA(J+1)**2)
    VAL2=1.-ETA(J)*X(J)
    Y=(X(J-1)*VAL1)/VAL2
    Z=(-X(J+1)*VAL1)/VAL2
    RHO=SQRT(Y**2+Z**2)
    IF(RHO-1.) 225,225,251
225 CONTINUE
    IF(RHO-EPS)226,226,228
226 TAU=PI/2.
    GO TO 240
228 CONTINUE
    TAU0=ASIN(ABS(Z/RHO))
    IF(Z) 235,231,231
231 IF(Y) 233,232,232
232 TAU=TAU0
    GO TO 240
233 TAU=PI-TAU0
    GO TO 240
235 IF(Y) 236,236,237
236 TAU=PI+TAU0
    GO TO 240
237 IF(X(J)) 1237,1237,1238
1237 F(J)=100.
    GO TO 259
1238 CONTINUE
    WRITE(6,238) TAU0,TAU,Y,Z
238 FORMAT(* ERROR FOR Y .GT. AND Z .LT. C----CASE C*/4E20.8)
    WRITE(6,253) J,VAL1,VAL2,RHO,X(2),X(3),X(4),XXI,BPSI ,UU
    STOP
240 CONTINUE
    W1=WF1(J)
    W2=WF2(J)
    BW1=BWF1(J)
    BW2=BWF2(J)
    CALL CON2D(RHO,TAU,W1,W2,CONS23,FIF)
    CALL CON2D(RHO,TAU,BW1,BW2,CONS23,FIIF)
    SIGNEP=SIGN(1.,ETA(J+1))
    SIGNEM=SIGN(1.,ETA(J-1))
243 F(J)= ((1.-FIF)*SIGNEP+FIIF*SIGNEM)
    GO TO 259
251 F(J)=100.
259 CONTINUE
258 CONTINUE
253 FORMAT(I3,9E13.5)
    TIMEF=0.
    DO 261 J=2,4

```

```

SUMF=SUMF+F(J)
261 CONTINUE
IF(SUMF-EPS) 260,260,262
260 FS=SUMF+1.
GO TO 300
262 CONTINUE
IF(SUMF-90.) 263,263,265
C
C-----INSIDE ALL THREE CONES
263 WRITE(6,264) (F(J),J=1,4),FIF,FIIF
264 FORMAT(* ERROR SUMF .LE. 90---CASE D*/6E20.8)
WRITE(6,200) VAL1,VAL2,Y,Z,RHO,TAU0,W1,W2,FIF,BW1,BW2,FIIF
STOP
265 IF(SUMF-190.) 267,267,269
C
C-----INSIDE TWO CONES AND OUTSIDE THE THIRD CONE
267 FS=SUMF-100.
GO TO 300
269 IF(SUMF-290.)271,271,272
C
C-----INSIDE ONE CONE ONLY
271 FS=SUMF-199.
GO TO 300
272 CONTINUE
C
C-----OUTSIDE ALL THREE CONES
273 DO 279 J=2,4
TEST=ETA(J-1)*X(J-1)-ETA(J)*X(J)+ETA(J+1)*X(J+1)
DJ=ETA(J+1)
DM=ETA(J-1)
IF(X(J) .LE. 0. .AND. TEST .LT. 1.) 275,279
275 IF(X(J+1) .GE. DJ .OR. X(J-1) .GE. DM)281,279
279 CONTINUE
FS=1.
GO TO 300
281 FS=2.
300 CONTINUE
301 FORMAT(* SUB FSF(NEW)* 5X, 5E20.8)
200 FORMAT(6E20.8)
RETURN
END

```

```

SUBROUTINE CON2D(RHO,TAU,W1,W2,SS,PC)
EPSIL=0.000001
CON1=1.-RHO**2
IF(CON1+EPSIL) 10,10,14
10 WRITE(6,11)
11 FORMAT(* 1-RHO**2+EPSIL .LE. ZERO, CHECK THE PROGRAM *)
STOP
14 IF(CON1) 15,15,17
15 GG=1.
GO TO 21
17 GG=(RHO/(1.+SQRT(CON1)))**SS
21 BTAU=SS*TAU
DD=(1.-GG**2)*SIN(0.5*(W2-W1))
CC=(1.+GG**2)*COS(0.5*(W2-W1)) -2.*GG*COS(BTAU-(W2+W1)*0.5)
SQCD=SQRT(CC**2+DD**2)
IF(SQCD-EPSIL) 23,25,25
23 PC=0.5
GO TO 30
25 IF(CC) 28,26,26
26 PC=ASIN(DD/SQCD)/PI
GO TO 30
28 PC=ASIN(DD/SQCD)/PI*(-1.)+1.
30 CONTINUE
RETURN
END

```



```

SUBROUTINE T102(XLAM,AAAAA,EDIM,BBBBBB)
DIMENSION EDIM(20)
COMMON XXXXL
XXXXL=XLAM
XMAX=0.9989
X0=0.001
X0XL=F3(XLAM,X0)
X0XLP=F3(XLAM+1.,X0)
X0XLM=F3(XLAM-1.,X0)
X0XLM2=F3(XLAM+2.,X0)
10 F0=(X0XL+XLAM*(XLAM+1.)*X0XLM2/(4.*XLAM+6.))*AAAAA
G0=(XLAM*X0XLM+(XLAM+1.)*(XLAM+2.)*XLAM*X0XLP/(4.*XLAM+6.))*AAAAA
DELX=0.001
X=X0
F=F0
G=G0
I=1
XTEST=0.0999
20 XK1=F1(X,F,G)*DELX
XL1=F2(X,F,G)*DELX
XK2=F1(X+DELX/2., F+XK1/2., G+XL1/2.)*DELX
XL2=F2(X+DELX/2., F+XK1/2., G+XL1/2.)*DELX
XK3=F1(X+DELX/2., F+XK2/2., G+XL2/2.)*DELX
XL3=F2(X+DELX/2., F+XK2/2., G+XL2/2.)*DELX
XK4=F1(X+DELX, F+XK3, G+XL3)*DELX
XL4=F2(X+DELX, F+XK3, G+XL3)*DELX
DELF=1./6.*(XK1+2.*XK2+2.*XK3+XK4)
DELG=1./6.*(XL1+2.*XL2+2.*XL3+XL4)
X=X+DELX
F=F+DELF
G=G+DELG
IF(X .GE. XTEST) 201,203
201 EDIM(I)=F
XTEST=XTEST+0.1
I=I+1
203 CONTINUE
IF(X-0.0099999) 20,31,31
31 IF(X-0.98999) 32,33,33
32 DELX=0.01
GO TO 20
33 IF(X-0.9989) 34,40,40
34 DELX=0.001
GO TO 20
40 ZETA=1.-X
C1=-XLAM*(XLAM+1.)/2.
GEND=(F+ZETA*G)/(1.+C1*ZETA)

```

```
41 BBBBBB=1./GEND  
FDIM(I)=GEND  
RETURN  
END
```

```
FUNCTION F1(A,B,C)
COMMON XLAM
F1=C
RETURN
END
```

```
FUNCTION F2(A,B,C)
```

```
COMMON XLAM
```

```
F2=(-2.*A*(A**2-1.)*C-XLAM*(XLAM+1.)*B)/(A**2*(A**2-1.))
```

```
RETURN
```

```
END
```

```
FUNCTION F3(D,E)
F3=EXP(D*ALOG(E))
RETURN
END
```

```

SUBROUTINE PPDD(SS,XLAM,XMU,LLMAX,PP,DP)
C SUBROUTINE PPDD--WITHOUT DP
XNU=SS
EPS=1.E-06
DD=1.
TEM1=1.
ZX=(1.-XMU)/2.
IF(ABS(ZX) .LE. EPS.AND. ABS(XNU) .LE. EPS) 35,31
31 CONTINUE
DO 40 L=2,LLMAX
ZL=L-1
DD=DD*(ZL-XLAM -1.)*(ZL+XLAM )/(ZL**2+ZL*XNU)
TEM1=TEM1+DD*ZX**(L-1)
40 CONTINUE
ARTFL=((1.-XMU)/(1.+XMU))**(XNU/2.)
PP=ARTFL*TEM1
GO TO 41
35 PP=1.
41 CONTINUE
DP=0.
RETURN
END

```

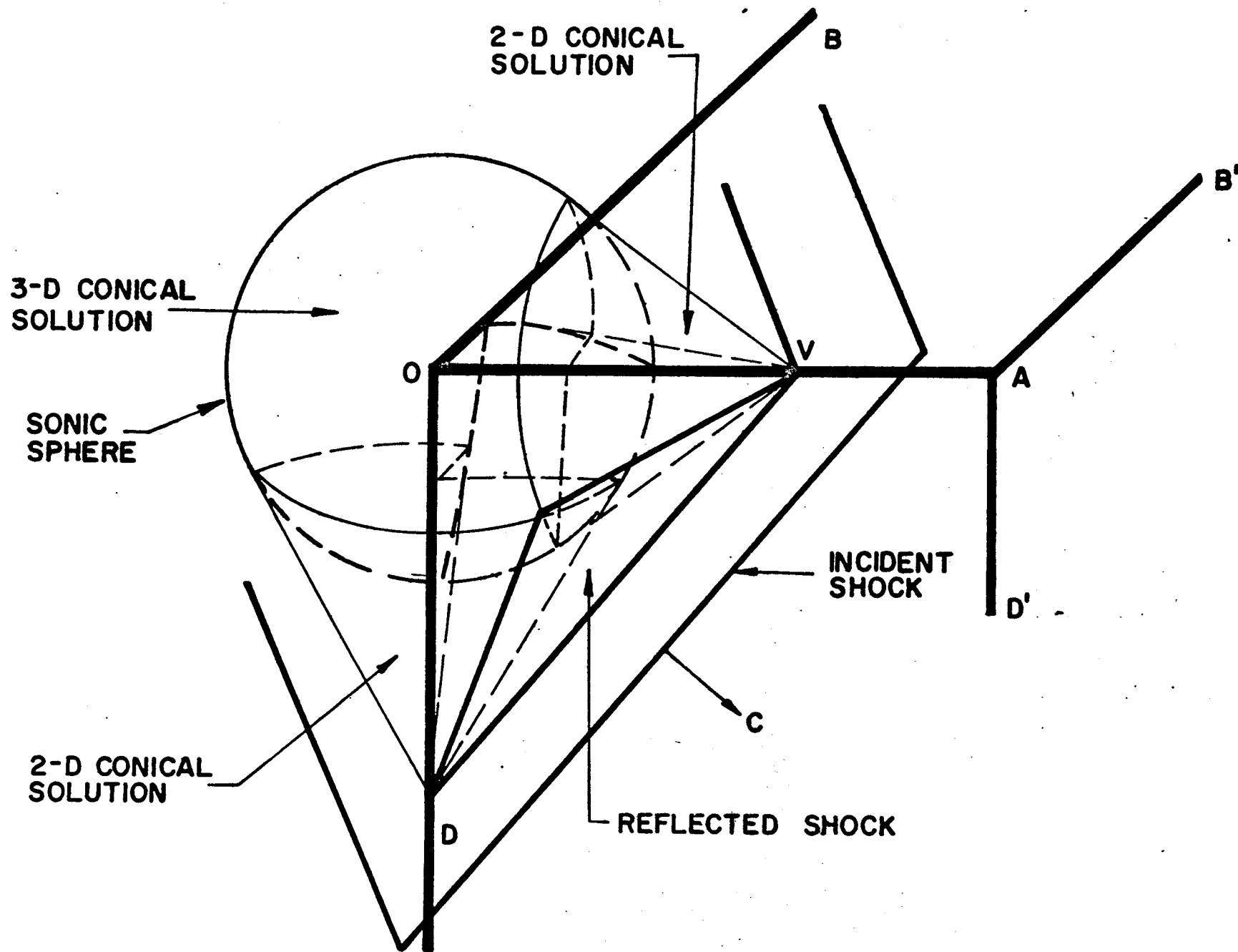


Fig. 1 Oblique incidence of a plane pulse on a three-dimensional corner

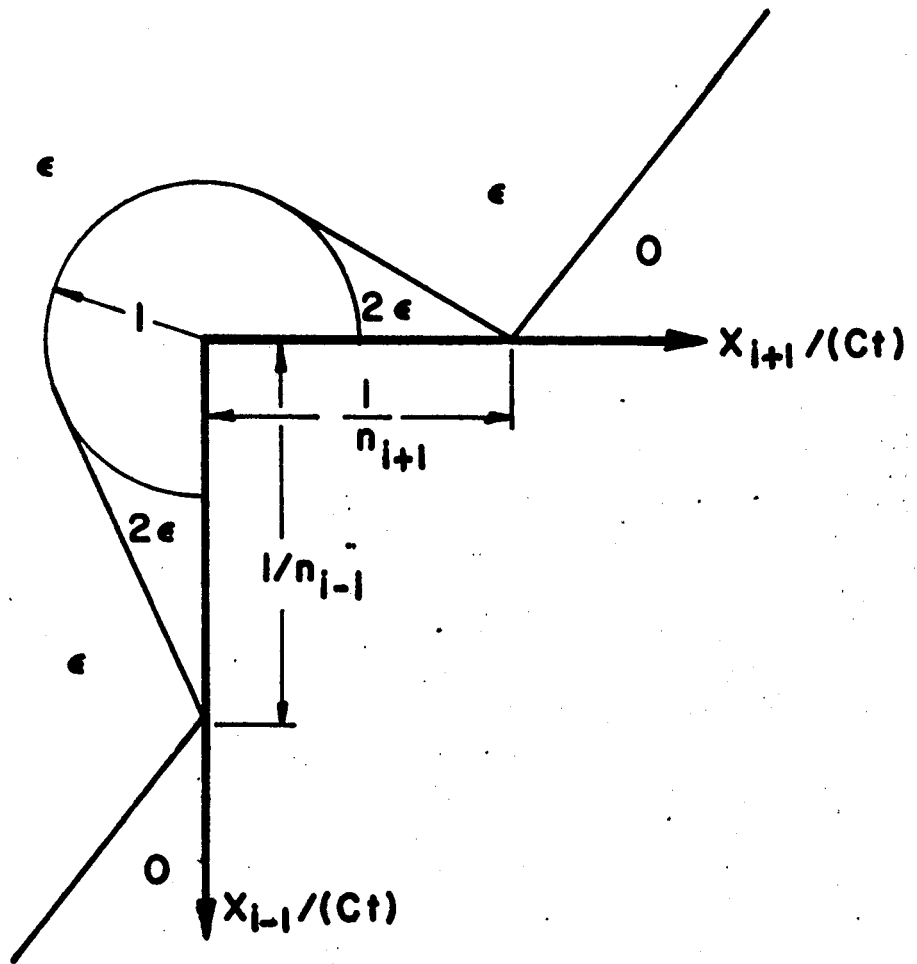


Fig. 2 Section normal to the i -th axis

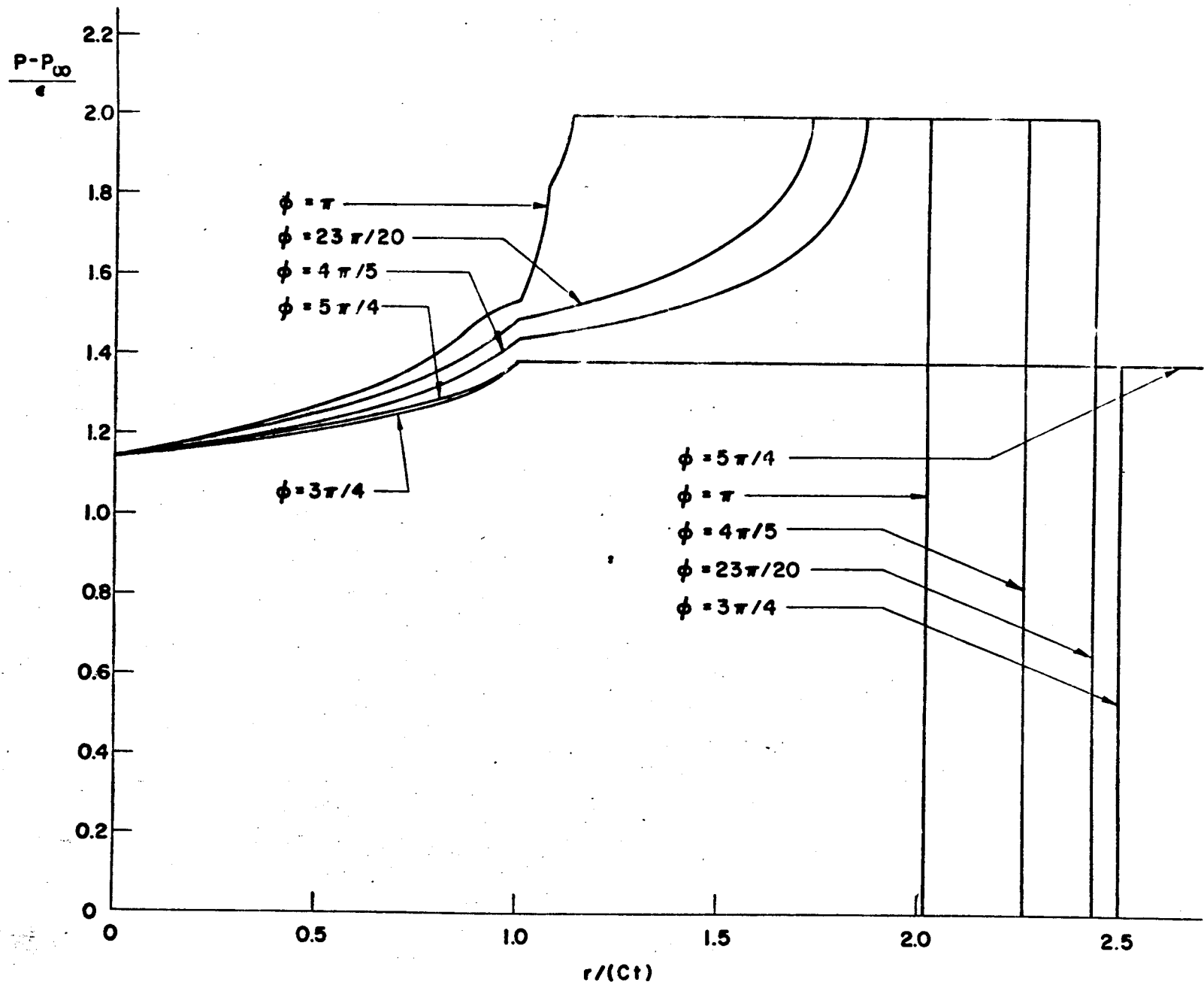


Fig. 3 Pressure distribution on the face $\theta = 0$ due to the incidence of a plane pulse with direction cosines 0.3, 0.4, 0.866

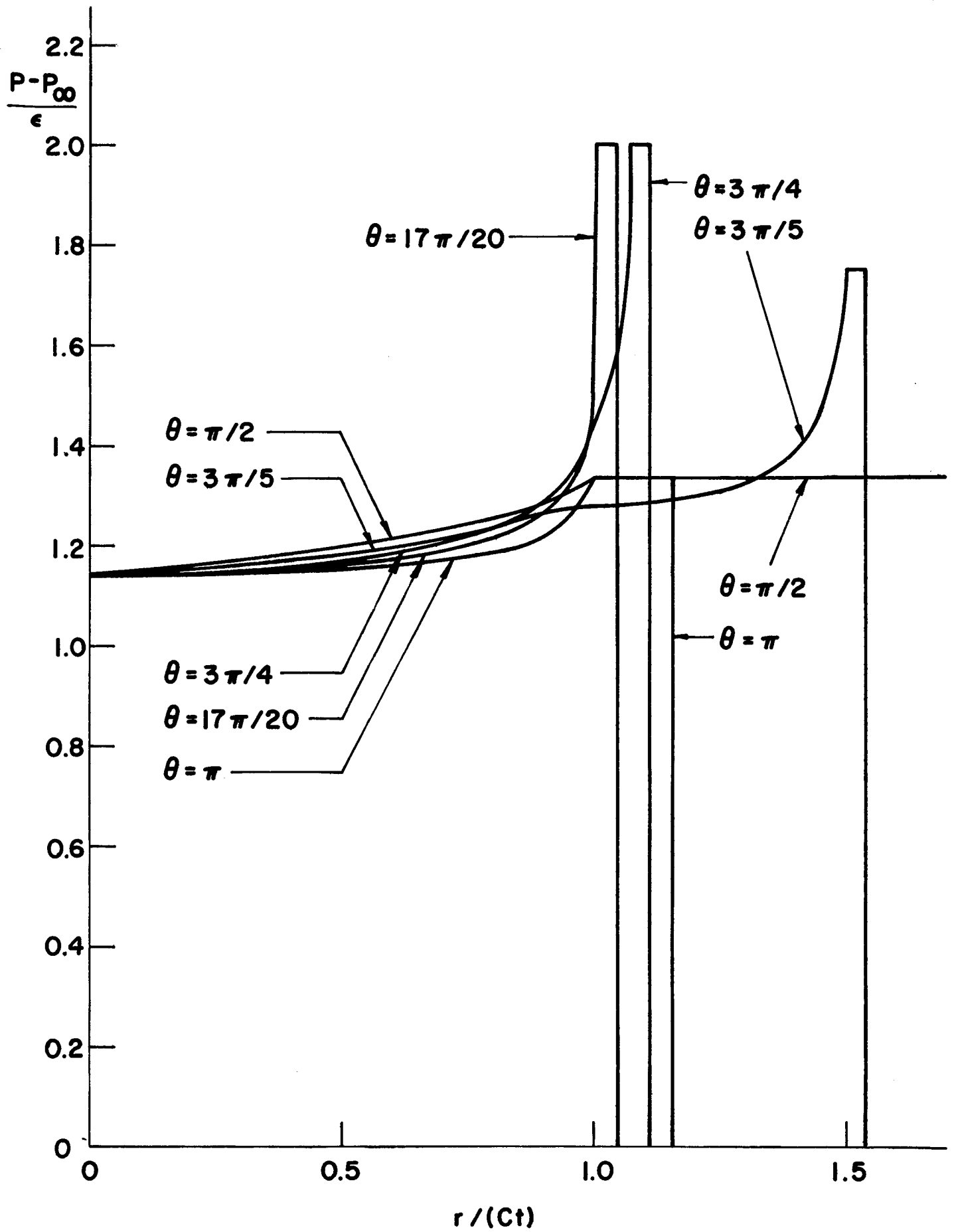


Fig. 4 Pressure distribution on the face $\phi = -3\pi/4$ due to the incidence of a plane pulse with direction cosines 0.3, 0.4, 0.866

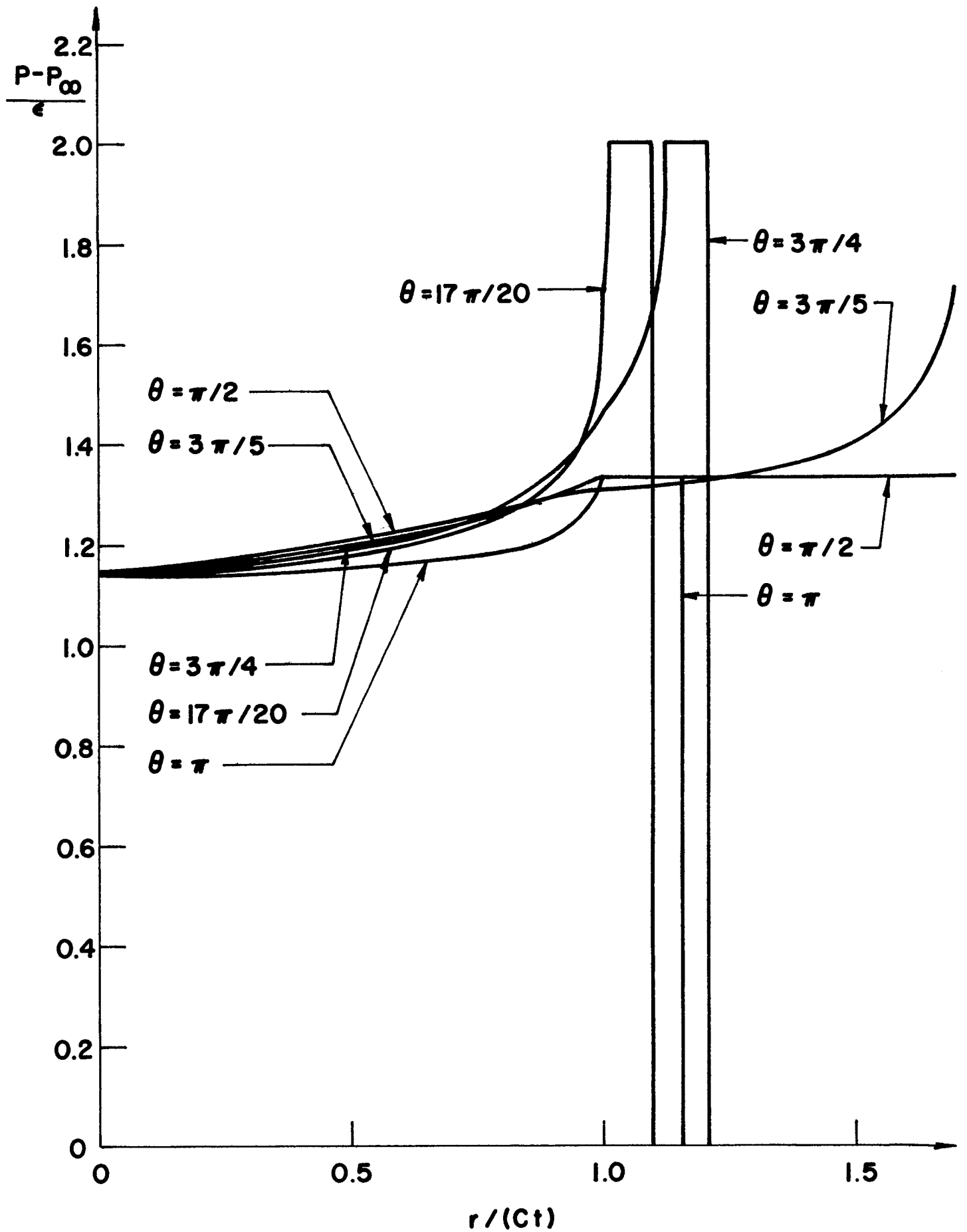


Fig. 5 Pressure distribution on the face $\varphi = 3\pi/4$ due to the incidence of a plane pulse with direction cosines 0.3, 0.4, 0.866

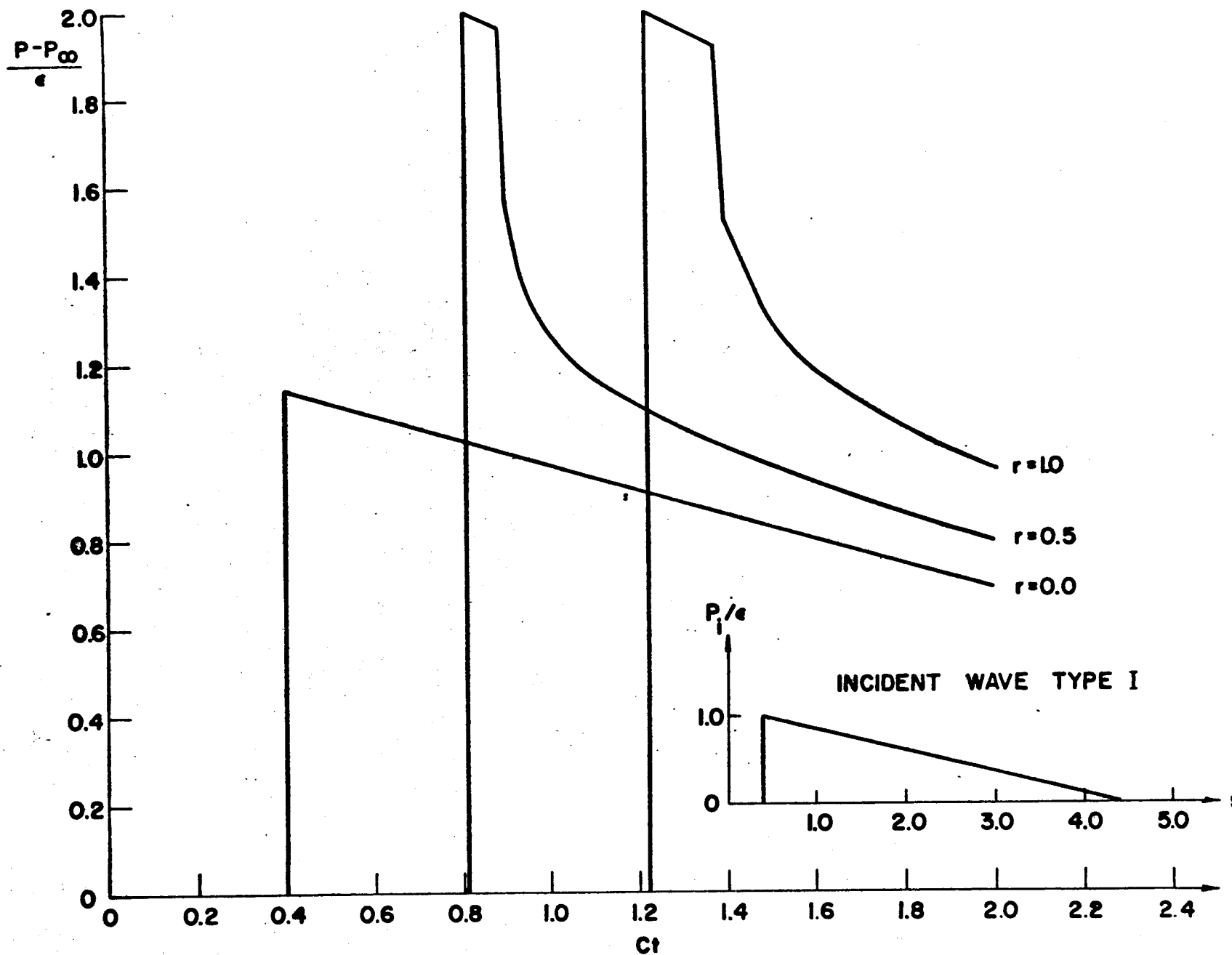


Fig. 6 Pressure signature received at points along the line $\theta = \pi/2$, $\varphi = \pi$ for plane wave incident at equal angles with the edges and with wave form of type I

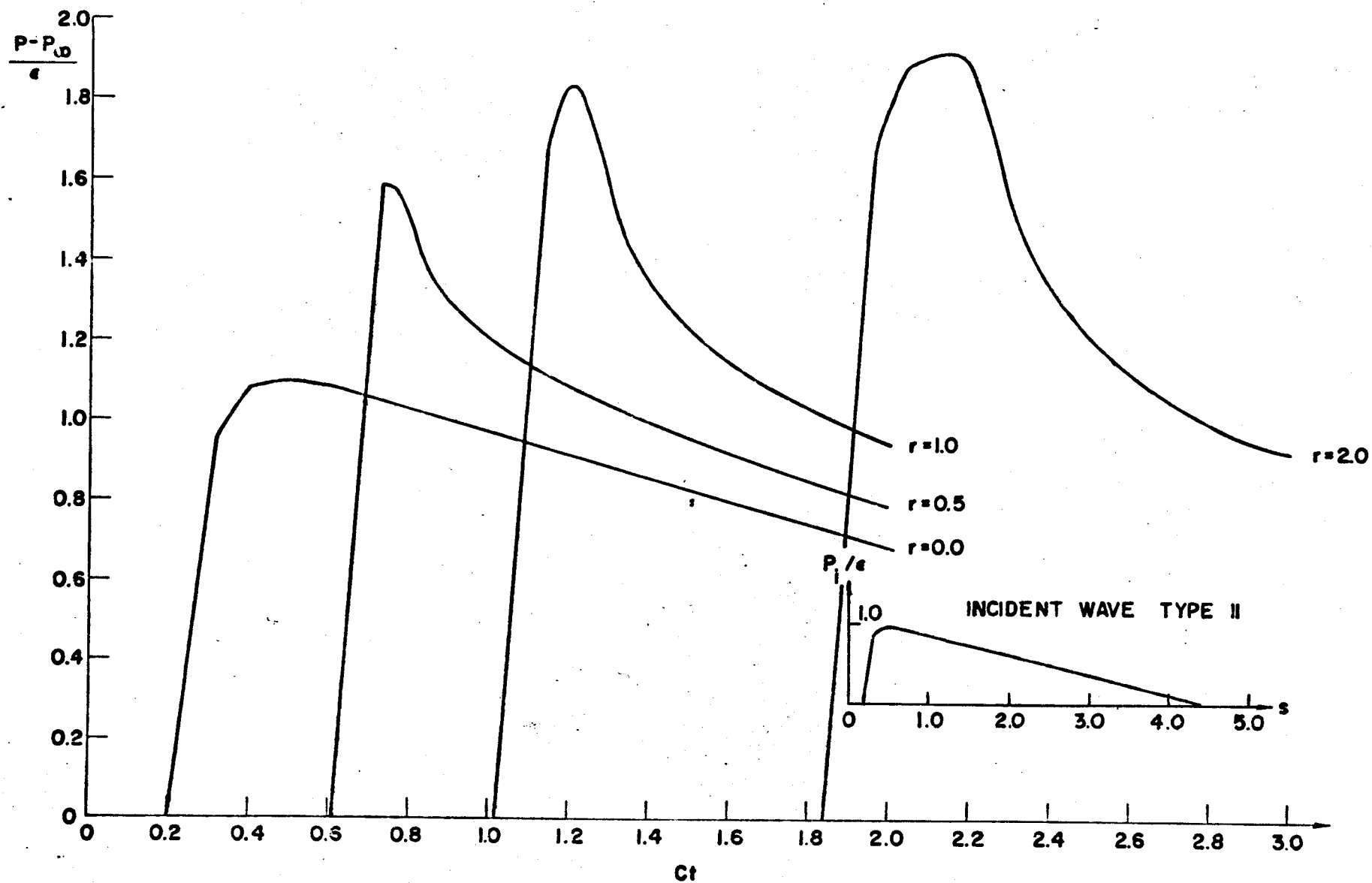


Fig. 7 Pressure signature received at points along the line $\theta = \pi/2$ $\phi = \pi$ for plane wave incident at equal angles with the edges and with wave form of type II

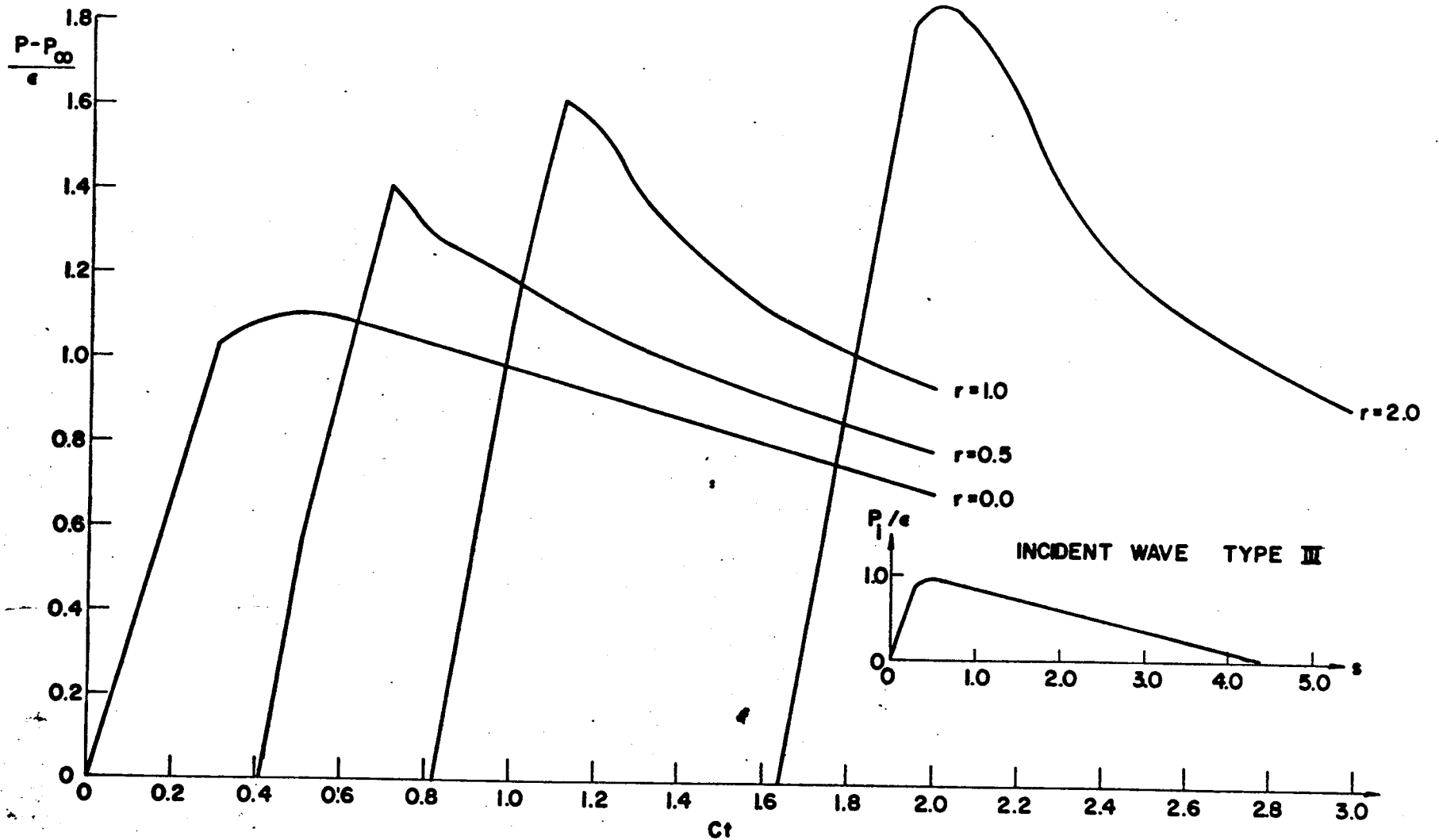


Fig. 8 Pressure signature received at points along the line $\theta = \pi/2$, $\phi = \pi$ for plane wave incident at equal angles with the edges and with wave form of type III

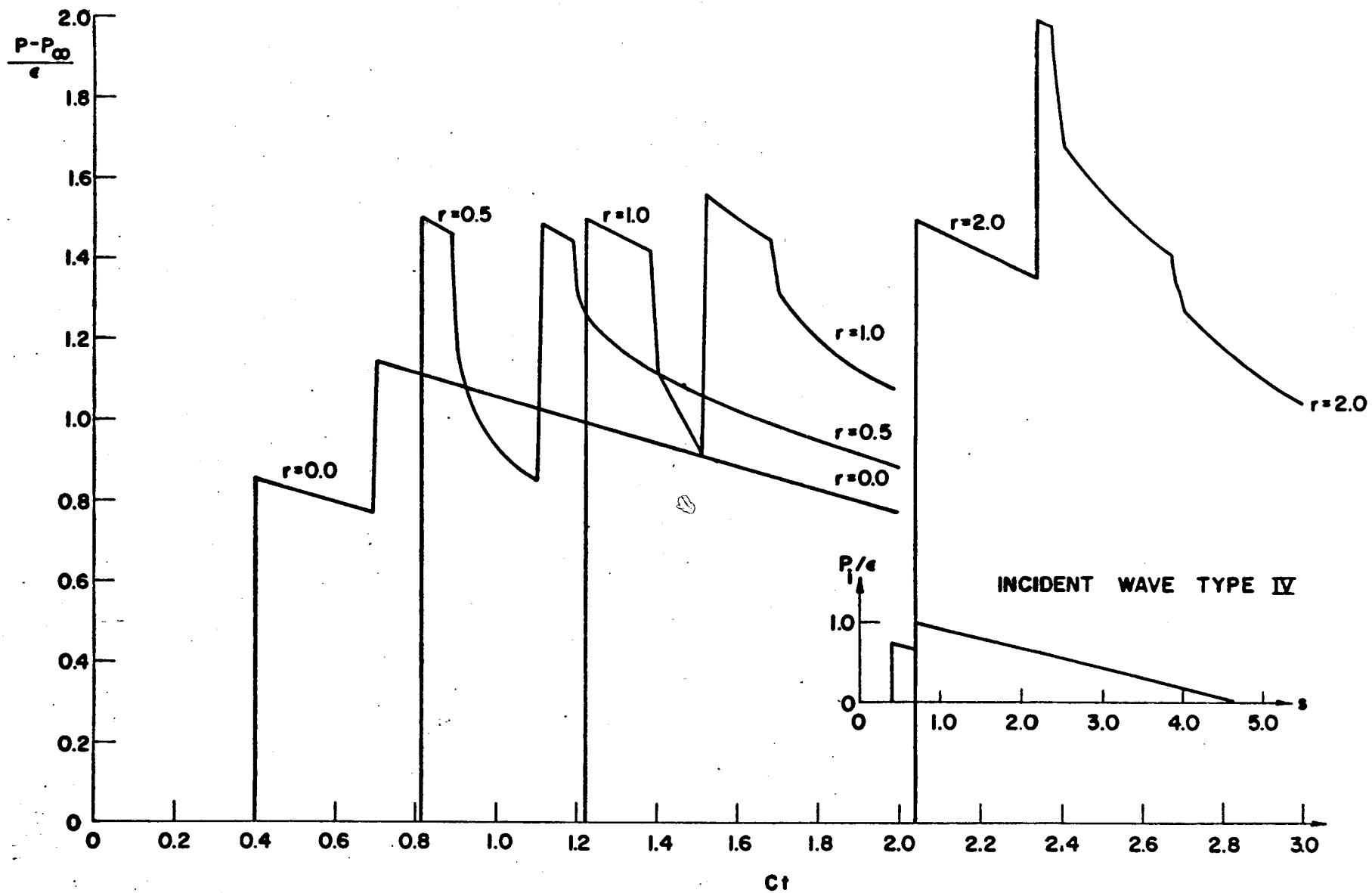


Fig. 9 Pressure signature received at points along the line $\theta = \pi/2$, $\phi = \pi$ for plane wave incident at equal angles with the edges and with wave form of type IV

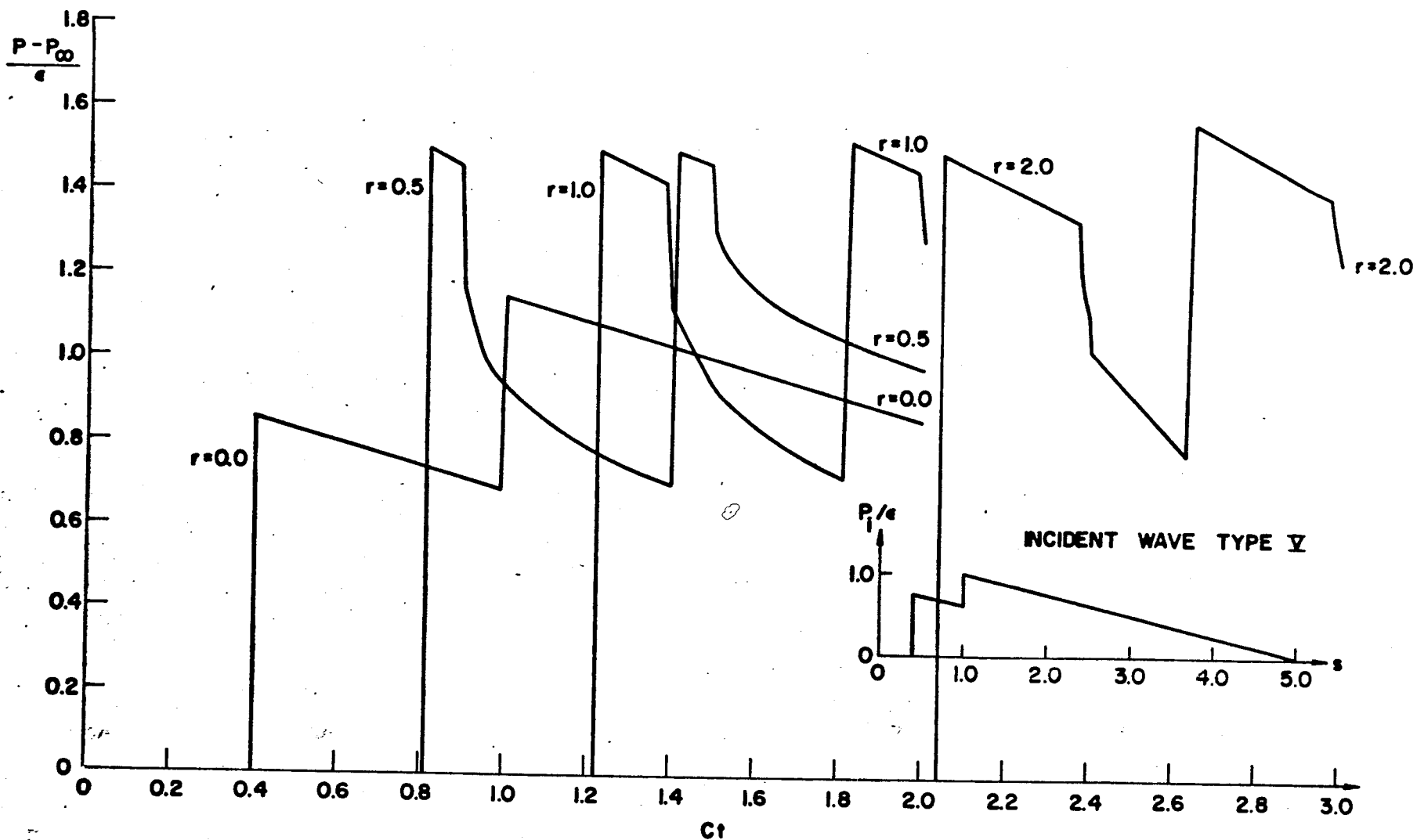


Fig. 10 Pressure signature received at points along line $\theta = \pi/2$, $\varphi = \pi$ for plane wave incident at equal angles with the edges and with wave form at type V