# Concavity of Eigenvalue Sums and the Spectral Shift Function 

Vadim Kostrykin<br>Fraunhofer-Institut für Lasertechnik, Steinbachstraße 15, D-52074 Aachen, Germany E-mail: kostrykin@t-online.de, kostrykin@ilt.fhg.de

Communicated by R. B. Melrose
Received January 18, 2000; revised May 3, 2000; accepted May 4, 2000

DEDICATED TO ROBERT SCHRADER ON THE OCCASION OF HIS 60TH BIRTHDAY
iew metadata, citation and similar papers at core.ac.uk

Hermitian matrix $V$ is concave (convex) with respect to $V$. Using the theory of the spectral shift function we generalize this property to self-adjoint operators on a separable Hilbert space with an arbitrary spectrum. More precisely, we prove that the spectral shift function integrated with respect to the spectral parameter from $-\infty$ to $\lambda$ (from $\lambda$ to $+\infty$ ) is concave (convex) with respect to trace class perturbations. The case of relative trace class perturbations is also considered. © 2000 Academic Press
Key Words: eigenvalue problems; spectral shift function; perturbation theory.

## 1. INTRODUCTION AND MAIN RESULTS

Consider an arbitrary Hermitian matrix $V$. Let $\lambda_{j}(V), j \in \mathbb{N}$ denote its eigenvalues enumerated in the increasing order and repeated according to their multiplicity. Consider the eigenvalue sums

$$
S_{\lambda}^{(-)}(V)=\sum_{j: \lambda_{j}(V) \leqslant \lambda}\left(\lambda_{j}(V)-\lambda\right), \quad S_{\lambda}^{(+)}(V)=\sum_{j: \lambda_{j}(V) \geqslant \lambda}\left(\lambda_{j}(V)-\lambda\right),
$$

which equivalently can be written in the form

$$
S_{\lambda}^{(-)}(V)=-\int_{-\infty}^{\lambda} N^{(-)}\left(\lambda^{\prime} ; V\right) d \lambda^{\prime}, \quad S_{\lambda}^{(+)}(V)=\int_{\lambda}^{\infty} N^{(+)}\left(\lambda^{\prime} ; V\right) d \lambda^{\prime},
$$

with $N^{( \pm)}$being the counting functions, i.e. $N^{( \pm)}(\lambda ; V)=\#\left\{j: \pm \lambda_{j}(V) \geqslant\right.$ $\pm \lambda\}$. By means of the min-max principle it can be easily proved (see, e.g.,
[17, 26]) that $S^{(-)}(V)$ is concave and $S^{(+)}(V)$ is convex with respect to $V$, i.e., for any Hermitian matrices $V_{1}$ and $V_{2}$ and any $\alpha \in[0,1]$

$$
\begin{equation*}
\pm S_{\lambda}^{( \pm)}\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \leqslant \pm\left(\alpha S_{\lambda}^{( \pm)}\left(V_{1}\right)+(1-\alpha) S_{\lambda}^{( \pm)}\left(V_{2}\right)\right) . \tag{1}
\end{equation*}
$$

These inequalities play an important role in several problems of quantum and statistical physics (see, e.g., references cited in [17]).

In the present note we show that for a wide class of self-adjoint operators on a separable Hilbert space $\mathfrak{H}$, which need not have purely discrete spectrum, the properties (1) are valid for properly regularized $S_{\lambda}^{( \pm)}$. More precisely, instead of $V$ compared to the zero operator we consider pairs $\left(A_{0}+V, A_{0}\right)$. For an arbitrary self-adjoint operator $A_{0}$ and any self-adjoint trace class operator $V$ we define

$$
\begin{equation*}
\zeta^{(-)}\left(\lambda ; A_{0}+V, A_{0}\right):=\int_{-\infty}^{\lambda} \xi\left(\lambda^{\prime} ; A_{0}+V, A_{0}\right) d \lambda^{\prime} \tag{2}
\end{equation*}
$$

and

$$
\begin{equation*}
\zeta^{(+)}\left(\lambda ; A_{0}+V, A_{0}\right):=\int_{\lambda}^{\infty} \xi\left(\lambda^{\prime} ; A_{0}+V, A_{0}\right) d \lambda^{\prime}, \tag{3}
\end{equation*}
$$

where $\xi\left(\lambda ; A_{0}+V, A_{0}\right)$ is the spectral shift function for the pair of operators $\left(A_{0}+V, A_{0}\right)$. Recall that for an arbitrary self-adjoint operator $A_{0}$ and any self-adjoint trace class operator $V$ the spectral shift function $\xi\left(\lambda ; A_{0}+V, A_{0}\right)$ exists such that $\xi\left(\cdot ; A_{0}+V, A_{0}\right) \in L^{1}(\mathbb{R})$. Let $F \in C_{\text {loc }}^{1}(\mathbb{R})$ be such that its derivative $F^{\prime}$ belongs to the Wiener class $W(\mathbb{R})$, i.e. $F^{\prime}(\lambda)$ is representable in the form

$$
F^{\prime}(\lambda)=\int_{\mathbb{R}} e^{-i \lambda t} d \sigma(t),
$$

where $\sigma(\cdot)$ is a finite complex-valued Borel measure on $\mathbb{R},|\sigma(\mathbb{R})|<\infty$. Then

$$
\begin{equation*}
\operatorname{tr}\left(F\left(A_{0}+V\right)-F\left(A_{0}\right)\right)=\int_{\mathbb{R}} F^{\prime}(\lambda) \xi\left(\lambda ; A_{0}+V, A_{0}\right) d \lambda . \tag{4}
\end{equation*}
$$

This last equation may be used as a definition of the spectral shift function. A wider class of functions for which the trace formula (4) remains valid is discussed in [2]. A review on the spectral shift function is the paper by Birman and Yafaev [3] (see also the book [27] and [5-7, 19] for recent results).

In the sequel we use the notation $\mathscr{J}_{p}, p \geqslant 1$ for the von NeumannSchatten ideals of compact operators such that in particular $\mathscr{I}_{1}$ denotes the set of the trace class operators (see, e.g., [9]). $\operatorname{spec}(A)$ denotes the spectrum of the operator $A . \mathscr{2}(A)$ is the domain of the quadratic form associated with the self-adjoint operator $A$.

If in some open interval $(a, b)$ the spectrum of $A_{0}$ is purely discrete then $\xi\left(b-0 ; A_{0}+V, A_{0}\right)-\xi\left(a+0 ; A_{0}+V, A_{0}\right)$ equals the difference of the total multiplicities of the spectra of $A_{0}$ and $A_{0}+V$ lying in $(a, b)$. Thus if we take $A_{0}=\lambda_{+} I$ with some $\lambda_{+}>\sup \operatorname{spec}(V)$, then $\zeta^{(-)}\left(\lambda ; A_{0}+V, A_{0}\right)=S_{\lambda}^{(-)}(V)$ for all $\lambda<\lambda_{+}$. Similarly $A_{0}=\lambda_{-} I$ with some $\lambda_{-}<\inf \operatorname{spec}(V)$ leads to $\zeta^{(+)}\left(\lambda ; A_{0}+V, A_{0}\right)=S_{\lambda}^{(+)}(V)$ for all $\lambda>\lambda_{-}$.

Theorem 1. Let $A_{0}$ and $V$ be self-adjoint operators on a separable Hilbert space $\mathfrak{H}, V \in \mathscr{g}_{1}$. For an arbitrary real-valued nonincreasing $f$ of bounded total variation the functional

$$
\begin{equation*}
g(V)=\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{0}+V, A_{0}\right) d \lambda \tag{5}
\end{equation*}
$$

is concave with respect to the perturbation $V$, i.e., for arbitrary $V_{1}, V_{2} \in \mathscr{J}_{1}$ the inequality

$$
\begin{equation*}
g\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \geqslant \alpha g\left(V_{1}\right)+(1-\alpha) g\left(V_{2}\right) \tag{6}
\end{equation*}
$$

holds for all $\alpha \in[0,1]$.
In particular we can take $f(\lambda)=\chi_{\left(-\infty, \lambda_{0}\right]}(\lambda)$, the characteristic function of $\left(-\infty, \lambda_{0}\right]$ with arbitrary $\lambda_{0} \in \mathbb{R}$, such that $g(V)=\zeta^{(-)}\left(\lambda_{0} ; A_{0}+V, A_{0}\right)$, the integrated spectral shift function (2). From Theorem 1 it follows that $\zeta^{(-)}\left(\lambda ; A_{0}+V, A_{0}\right)$ is concave with respect to $V$.

It is known that

$$
\int_{\mathbb{R}} \xi\left(\lambda ; A_{0}+V, A_{0}\right) d \lambda=\operatorname{tr} V,
$$

which is obviously linear in $V$. Since an arbitrary nondecreasing function $\tilde{f}$ of bounded total variation can be represented as a difference of a constant and a nonincreasing $f$ of bounded total variation we obtain

Corollary 1. Let $A_{0}$ and $V$ be as in Theorem 1. For an arbitrary real-valued nondecreasing $\tilde{f}$ of bounded total variation the functional

$$
\tilde{g}(V)=\int_{\mathbb{R}} \tilde{f}(\lambda) \xi\left(\lambda ; A_{0}+V, A_{0}\right) d \lambda
$$

is convex with respect to the perturbation $V$, i.e. for arbitrary $V_{1}, V_{2} \in \mathscr{g}_{1}$ the inequality

$$
\tilde{g}\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \leqslant \alpha \tilde{g}\left(V_{1}\right)+(1-\alpha) \tilde{g}\left(V_{2}\right)
$$

holds for all $\alpha \in[0,1]$.
In particular $\tilde{f}(\lambda)=\chi_{\left[\lambda_{0},+\infty\right)}(\lambda)$ satisfies the conditions of the corollary and thus $\zeta^{(+)}\left(\lambda ; A_{0}+V, A_{0}\right)$ defined by (3) is convex with respect to $V$.

Corollary 2. Let the function $f$ satisfy the conditions of Theorem 1. Let $V(\alpha)$ be a $\mathscr{I}_{1}$-valued operator family concave (in the operator sense) with respect to $\alpha$. Then the real-valued function $\alpha \mapsto g(V(\alpha))$ is concave. Similarly, if $f$ satisfies the conditions of Corollary 1 and $V(\alpha)$ is convex, then $\alpha \mapsto \tilde{g}(V(\alpha))$ is also convex.

Theorem 1 and Corollary 2 will be proved in Section 2 below.
We note that a special case of this result was proved recently by Gesztesy et al. [8, Corollary 1.9] by different methods.

In the article [4] written by the present author in collaboration with R. Geisler and R. Schrader we have proven that the integrated spectral shift function for the pair of Schrödinger operators is concave with respect to the perturbation potential. Here we will prove that this property holds for an arbitrary pair ( $A, A_{0}$ ) of self-adjoint semibounded operators on a separable Hilbert space $\mathfrak{G}$ provided that their difference is a relative trace class perturbation of $A_{0}$.

More precisely, we suppose that $A_{0}$ is a self-adjoint operator, semibounded from below, and $V$ is also self-adjoint and $A_{0}$-compact in the form sense, i.e. for all $a>-\inf \operatorname{spec}\left(A_{0}\right)$ the operator $\left(A_{0}+a\right)^{-1 / 2}$ $V\left(A_{0}+a\right)^{-1 / 2}$ is compact. Then the operator $A_{V}=A_{0}+V$, defined in the form sense, is self-adjoint with $\mathscr{2}\left(A_{V}\right)=\mathscr{2}\left(A_{0}\right)$ and also semibounded from below. Suppose that for some $p \geqslant 1$ and for all sufficiently large $a$

$$
\begin{equation*}
\left(A_{V}+a\right)^{-p}-\left(A_{0}+a\right)^{-p} \in \mathscr{I}_{1} . \tag{7}
\end{equation*}
$$

If $\mathscr{I}$ is an interval of the real axis such that $\mathscr{I} \supset \operatorname{spec}\left(A_{V}\right) \cup \operatorname{spec}\left(A_{0}\right)$ and for some real-valued strictly monotone $\varphi \in C_{\text {loc }}^{2}(\mathscr{F})$ the difference $\varphi\left(A_{V}\right)-\varphi\left(A_{0}\right)$ is trace class then the spectral shift function $\xi\left(\lambda ; A_{V}, A_{0}\right)$ for the pair of operators ( $A_{V}, A_{0}$ ) can be defined by means of the relation

$$
\begin{equation*}
\xi\left(\lambda ; A_{V}, A_{0}\right):=\varepsilon \xi\left(\varphi(\lambda) ; \varphi\left(A_{V}\right), \varphi\left(A_{0}\right)\right), \quad \varepsilon=\operatorname{sign} \varphi^{\prime}(\lambda), \tag{8}
\end{equation*}
$$

which turns out to be independent of $\varphi$. Obviously, $\xi\left(\lambda ; A_{V}, A_{0}\right)$ satisfies the trace formula (4) for some class of admissible functions $F$. This construction is known in the literature as the "invariance principle" for the
spectral shift function (see e.g. [3, 27]). Setting in (8) $\varphi(\lambda)=(\lambda+a)^{-p}$ we obtain

$$
\begin{equation*}
\xi\left(\lambda ; A_{V}, A_{0}\right)=-\xi\left((\lambda+a)^{-p} ;\left(A_{V}+a\right)^{-p},\left(A_{0}+a\right)^{-p}\right) . \tag{9}
\end{equation*}
$$

It vanishes for all $\lambda<\inf \left\{\operatorname{spec}\left(A_{V}\right), \operatorname{spec}\left(A_{0}\right)\right\}$.
In the special case $A_{0}=-\Delta$ in $L^{2}\left(\mathbb{R}^{v}\right)$ and $V$ being the multiplication operator by a real-valued measurable function $V(x)$ the conditions above are satisfied with any $p>(v-1) / 2$ for $v \geqslant 4$ and with $p=1$ for $v \leqslant 3$ provided that

$$
\begin{array}{llll}
V \in L^{v / 2}\left(\mathbb{R}^{v}\right) \cap l^{1}\left(L^{2}\left(\mathbb{R}^{v}\right)\right) & \text { for } & v \geqslant 5, & \\
V \in L^{r}\left(\mathbb{R}^{v}\right) \cap l^{1}\left(L^{2}\left(\mathbb{R}^{v}\right)\right) & \text { for } & v=4 \quad \text { and some } & r>2, \\
V \in L^{2}\left(\mathbb{R}^{v}\right) \cap L^{1}\left(\mathbb{R}^{v}\right) & \text { for } & v=2,3, & \\
V \in L^{1}(\mathbb{R}) & \text { for } & v=1 .
\end{array}
$$

For the definition of the Birman-Solomyak classes $l^{p}\left(L^{q}\right)$ see, e.g., [24].
Let $\mathscr{C}\left(A_{0}, a_{0}, p\right), a_{0} \in \mathbb{R}, p \geqslant 1$ denote a set of self-adjoint operators on the separable Hilbert space $\mathfrak{G}$ satisfying the following properties:
(i) every $V \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ is $A_{0}$-compact in the form sense;
(ii) $a_{0}>-\inf \operatorname{spec}\left(A_{V}\right)$ for all $V \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ and the condition [7] is satisfied for all $V \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ and all $a \geqslant a_{0}$;
(iii) the set $\mathscr{C}\left(A_{0}, a_{0}, p\right)$ is convex, i.e., $V_{1}, V_{2} \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ implies that $\alpha V_{1}+(1-\alpha) V_{2} \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ for all $\alpha \in[0,1]$.

We will say that a set possessing these properties for some $a_{0} \in \mathbb{R}$ and $p \geqslant 1$ is $A_{0}$-convex. Obviously $\mathscr{C}\left(A_{0}, a_{0}, p\right)$ is also $A$-convex for any operator $A$ such that $A-A_{0} \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$.

As an example consider two self-adjoint operators $V_{j}$ which are $A_{0}$-compact in the form sense and satisfy

$$
\left(A_{0}+a\right)^{-1 / 2} V_{j}\left(A_{0}+a\right)^{-p-1 / 2} \in \mathscr{J}_{1}, \quad j=1,2
$$

for some $a>-\inf \operatorname{spec}\left(A_{0}\right)$ and $p \geqslant 1$. Any operator lying in the convex hull $\left\{\alpha V_{1}+(1-\alpha) V_{2}, \alpha \in[0,1]\right\}$ of $\left\{V_{1}, V_{2}\right\}$ is obviously also $A_{0}$-compact in the form sense. Take $a_{0}>a$ such that

$$
\left\|\left(A_{0}+a_{0}\right)^{-1 / 2} V_{j}\left(A_{0}+a_{0}\right)^{-1 / 2}\right\|<1
$$

for both $j=1,2$. By [22, Theorem XI.12] we obtain that the condition (7) is satisfied for all $V=\alpha V_{1}+(1-\alpha) V_{2}$ with $\alpha \in[0,1]$ and arbitrary $a \geqslant a_{0}$. Thus, the convex hull of $\left\{V_{1}, V_{2}\right\}$ is $A_{0}$-convex.

Theorem 2. Let $A_{0}$ be a self-adjoint operator semibounded from below and $\mathscr{C}\left(A_{0}, a_{0}, p\right)$ be some $A_{0}$-convex set. Let $q$ equal $p$ if $p=1$ and the smallest odd integer larger than $p$ if $p>1$. Let $A_{V}$ with $V \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ denote the operator $A_{0}+V$ defined in the form sense. For an arbitrary real-valued nonnegative nonincreasing $f$ of bounded total variation on $\left[-a_{0},+\infty\right)$ such that

$$
\sup _{\lambda \in\left[-a_{0},+\infty\right)}(1+|\lambda|)^{q+1}|f(\lambda)|<\infty
$$

the functional

$$
g(V)=\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{V}, A_{0}\right) d \lambda
$$

is concave on $\mathscr{C}\left(A_{0}, a_{0}, p\right)$, i.e., for arbitrary $V_{1}, V_{2} \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$ the inequality

$$
\begin{equation*}
g\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \geqslant \alpha g\left(V_{1}\right)+(1-\alpha) g\left(V_{2}\right) \tag{10}
\end{equation*}
$$

holds for all $\alpha \in[0,1]$.
The proof of this theorem will be given in Section 3 below.
As discussed in [4] (see also Proposition 3.1 below) the concavity (convexity) of $g(V)(\tilde{g}(V)$, respectively) implies that $g(\alpha V)$ is subadditive and $\tilde{g}(\alpha V)$ is superadditive with respect to $\alpha \in \mathbb{R}_{+}$. Subadditivity and superadditivity properties with respect to the perturbation (rather than with respect to the coupling constant) do not hold generally. In the special case of the Schrödinger operators this was observed in [4, 14, 15]. Subadditivity and superadditivity properties of the spectral shift function play an important role in some problems related to random Schrödinger operators [14, 15]. Also they allow one to study the strong coupling limit. In particular, in Section 3 we will prove

Corollary 3. Let $A_{0}$ be an arbitrary self-adjoint operator and $V \geqslant 0$. Assume that either
(i) $V$ is trace class or
(ii) $A_{0}$ is semibounded from below, $V$ is $A_{0}$-compact in the form sense and $\left(A_{0}+a\right)^{-1 / 2} V \cdot\left(A_{0}+a\right)^{-p-1 / 2} \in \mathscr{J}_{1}$ for some $p \geqslant 1$ and some $a>$ $-\inf \operatorname{spec}\left(A_{0}\right)$.

Then for any nonincreasing function $f$ of bounded total variation, which in the case (ii) satisfies additionally the conditions of Theorem 2, the limit

$$
\lim _{\alpha \rightarrow \infty} \frac{1}{\alpha} \int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{0}+\alpha V, A_{0}\right) d \lambda
$$

exists and is finite.
For other results related to the strong coupling limit we refer to [20, 21, 23].
Most of the results of the present note have appeared previously in [13] in a slightly less general form.

## 2. TRACE CLASS PERTURBATIONS

The proof of Theorem 1 relies on the following result of Birman and Solomyak [2]:

Lemma 2.1. Let $f \geqslant 0$ be a nonincreasing function with bounded total variation. Then for any self-adjoint operators $A_{0}$ and $V$ on $\mathfrak{G}, V \in \mathscr{\mathscr { A }}_{1}$
(i) the real-valued function $\alpha \mapsto \operatorname{tr}\left[f\left(A_{0}+\alpha V\right) V\right]$ is nonincreasing, i.e., for $\alpha_{1} \leqslant \alpha_{2}$ the inequality

$$
\operatorname{tr}\left[f\left(A_{0}+\alpha_{1} V\right) V\right] \geqslant \operatorname{tr}\left[f\left(A_{0}+\alpha_{2} V\right) V\right]
$$

holds,
(ii)

$$
\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{0}+\alpha V, A_{0}\right) d \lambda=\int_{0}^{\alpha} \operatorname{tr}\left[f\left(A_{0}+s V\right) V\right] d s .
$$

Remark 2.1. The proof in [2] of part (i) relies on the theory of the double Stieltjes operator integral. An alternative proof not using this formalism is given by Gesztesy et al. [8]. Part (ii) of the lemma is proven in [2] for the case $f(\lambda)=\chi_{\left(-\infty, \lambda_{0}\right]}(\lambda)$. The present extension is immediate. Alternative proofs of (ii) have appeared in [5, 25]. An operator-valued version of this formula for sign-definite perturbations is given in [5].

From Lemma 2.1 (i) it follows that

$$
G: \alpha \mapsto \int_{0}^{\alpha} \operatorname{tr}\left[f\left(A_{0}+s V\right) V\right] d s
$$

is concave. Indeed a necessary and sufficient condition for $G(\cdot)$ to be concave is

$$
\begin{equation*}
2 G(\alpha)-G(\alpha+h)-G(\alpha-h) \geqslant 0 \tag{2.1}
\end{equation*}
$$

for all $\alpha \in \mathbb{R}, h \geqslant 0$. Since $\alpha \mapsto \operatorname{tr}\left[f\left(A_{0}+\alpha V\right) V\right]$ is nonincreasing we have

$$
\int_{\alpha}^{\alpha+h} \operatorname{tr}\left[f\left(A_{0}+s V\right) V\right] d s-\int_{\alpha-h}^{\alpha} \operatorname{tr}\left[f\left(A_{0}+s V\right) V\right] d s \leqslant 0,
$$

which is equivalent to (2.1). Now by the claim (ii) of Lemma 2.1 it follows that the functional $g(V)(5)$ is concave with respect to the coupling constant.

By the chain rule for the spectral shift function (see, e.g., [3])

$$
\xi\left(\lambda ; A_{1}+\alpha V, A_{1}\right)=\xi\left(\lambda ; A_{1}+\alpha V, A_{0}\right)+\xi\left(\lambda ; A_{0}, A_{1}\right),
$$

we have that

$$
\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{1}+\alpha V, A_{0}\right) d \lambda
$$

is also concave with respect to $\alpha$ for arbitrary $A_{0}$ and $A_{1}$ such that $A_{1}-A_{0} \in \mathscr{I}_{1}$. Thus for arbitrary $t_{1}, t_{2} \in \mathbb{R}$ and arbitrary $V \in \mathscr{I}_{1}$ we have

$$
\begin{aligned}
& \int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{1}+\alpha t_{1} V+(1-\alpha) t_{2} V, A_{0}\right) d \lambda \\
& \quad \geqslant \alpha \int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{1}+t_{1} V, A_{0}\right) d \lambda \\
& \quad+(1-\alpha) \int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{1}+t_{2} V, A_{0}\right) d \lambda
\end{aligned}
$$

for all $\alpha \in[0,1]$. Taking $t_{1}=0, t_{2}=1, A_{1}=A_{0}+V_{1}$, and $V=V_{2}-V_{1}$ we obtain

$$
g\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \geqslant \alpha g\left(V_{1}\right)+(1-\alpha) g\left(V_{2}\right),
$$

thus proving the claim of Theorem 1, however, under the additional requirement that $f \geqslant 0$. To eliminate this requirement let us consider the
function $f_{1}$, which differs from $f \geqslant 0$ by a negative constant $c$. Since $\operatorname{tr} V$ is linear in $V$, the induced functional

$$
g_{1}(V)=\int_{\mathbb{R}} f_{1}(\lambda) \xi\left(\lambda ; A_{0}+V, A_{0}\right) d \lambda=g(V)+c \operatorname{tr} V
$$

is also concave in $V$. This completes the proof of Theorem 1.

## Proof of Corollary 2. By Theorem 1

$$
\begin{equation*}
g\left(\alpha V_{1}+(1-\alpha) V_{2}\right) \geqslant \alpha g\left(V_{1}\right)+(1-\alpha) g\left(V_{2}\right) \tag{2.2}
\end{equation*}
$$

for all $\alpha \in[0,1]$. By the monotonicity of the spectral shift function with respect to the perturbation $g(V)$ is nondecreasing with respect to $V$, i.e., $g\left(V_{1}\right) \geqslant g\left(V_{2}\right)$ for $V_{1} \geqslant V_{2}$.

Let now $V_{1}=V\left(\beta_{1}\right)$ and $V_{2}=V\left(\beta_{2}\right)$. By the concavity of $V(\alpha)$, i.e., by

$$
V\left(\alpha \beta_{1}+(1-\alpha) \beta_{2}\right) \geqslant \alpha V\left(\beta_{1}\right)+(1-\alpha) V\left(\beta_{2}\right),
$$

and by the monotonicity of $g(V)$, from (2.2) it follows that

$$
g\left(V\left(\alpha \beta_{1}+(1-\alpha) \beta_{2}\right)\right) \geqslant \alpha g\left(V\left(\beta_{1}\right)\right)+(1-\alpha) g\left(V\left(\beta_{2}\right)\right) .
$$

The second part of the claim can be proved similarly.

## 3. RELATIVE TRACE CLASS PERTURBATIONS

We turn to the case of relative trace class perturbations of $A_{0}$ and prove Theorem 2. The conditions of this theorem imply that

$$
\begin{equation*}
\sup _{\lambda \in\left[-a_{0},+\infty\right)}\left(\lambda+a_{0}\right)^{q+1}|f(\lambda)|<\infty \tag{3.1}
\end{equation*}
$$

with $q$ being equal to $p$ if $p=1$ and to the smallest odd integer larger than $p$ if $p>1$. Choose arbitrary $V_{1}, V_{2} \in \mathscr{C}\left(A_{0}, a_{0}, p\right)$. Obviously, $\{(1-\alpha)$ $\left.\left(V_{2}-V_{1}\right), \alpha \in[0,1]\right\} \subseteq \mathscr{C}\left(A_{0}, a_{0}, p\right)-V_{1}$ and the set $\mathscr{C}_{1}\left(A_{0}, a_{0}, p\right):=$ $\mathscr{C}\left(A_{0}, a_{0}, p\right)-V_{1}$ is $A_{0}$-convex. Note that $0 \in \mathscr{C}_{1}\left(A_{0}, a_{0}, p\right)$. Thus, as in the case of trace class perturbations it suffices to prove that for any $V \in \mathscr{C}_{1}\left(A_{0}, a_{0}, p\right)$ the function $g(\alpha V)$ is concave with respect to $\alpha \in[0,1]$.

We start with the simplest case $p=1$ in the condition (7). For all $a \geqslant a_{0}$ the resolvents $\left(A_{\alpha V}+a\right)^{-1}$ and $\left(A_{0}+a\right)^{-1}$ are bounded nonnegative operators. By assumption their difference is trace class and therefore the
spectral shift function $\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)$ can be defined by means of the invariance principle as given by (9). For all $a \geqslant a_{0}$ it satisfies the inequality

$$
\begin{equation*}
\int_{\mathbb{R}} \frac{\left|\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)\right|}{(\lambda+a)^{2}} d \lambda \leqslant\left\|\left(A_{\alpha V}+a\right)^{-1}-\left(A_{0}+a\right)^{-1}\right\|_{\mathscr{I}_{1}} . \tag{3.2}
\end{equation*}
$$

Lemma 3.1. Let $f(\lambda) \geqslant 0$ be a nonincreasing function of bounded total variation. Then for all $a \geqslant a_{0}$

$$
g_{a}(\alpha)=\int_{\mathbb{R}} \frac{f(\lambda)}{(\lambda+a)^{2}} \xi\left(\lambda ; A_{\alpha V}, A_{0}\right) d \lambda
$$

is concave with respect to $\alpha$.
Proof. We change the integration variable $\lambda \rightarrow t=(\lambda+a)^{-1}$ and use the invariance principle (9) to obtain

$$
\begin{equation*}
g_{a}(\alpha)=-\int_{0}^{\infty} f\left(\frac{1-a t}{t}\right) \xi\left(t ;\left(A_{\alpha V}+a\right)^{-1},\left(A_{0}+a\right)^{-1}\right) d t \tag{3.3}
\end{equation*}
$$

It is easy to see that $f((1-a t) / t)$ is nondecreasing with respect to $t$. It is well known (see, e.g., [1, 16; 18 Proposition 1.3.11]) that the function $x \mapsto x^{-1}$ is concave on the set of invertible positive operators; i.e., for arbitrary invertible positive operators $X$ and $Y$ the inequality ( $\beta X+$ $(1-\beta) Y)^{-1} \leqslant \beta X^{-1}+(1-\beta) Y^{-1}$ holds in the operator sense for all $\beta \in[0,1]$. Taking $X=A_{\alpha_{1} V}+a$ and $Y=A_{\alpha_{2} V}+a$ with an arbitrary $a \geqslant a_{0}$ and using the fact that $\beta X+(1-\beta) Y=A_{\left(\beta \alpha_{1}+(1-\beta) \alpha_{2}\right) V}+a$ we obtain that

$$
\left(A_{\left(\beta \alpha_{1}+(1-\beta) \alpha_{2}\right) V}+a\right)^{-1} \leqslant \beta\left(A_{\alpha_{1} V}+a\right)^{-1}+(1-\beta)\left(A_{\alpha_{2} V}+a\right)^{-1}
$$

for all $\beta \in[0,1]$; i.e., the operator $\left(A_{\alpha V}+a\right)^{-1}$ is convex with respect to $\alpha \in \mathbb{R}$. Therefore by Corollary 2 the integral in (3.3) is convex with respect to $\alpha$ and thus $g_{a}(\alpha)$ is concave.

From (3.1) it follows that the function $f$ satisfies the condition

$$
\sup _{\lambda \in\left[-a_{0},+\infty\right)}\left(\lambda+a_{0}\right)^{2}|f(\lambda)|<\infty .
$$

Since $\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)=0$ for all $\lambda \leqslant-a_{0}$ we may suppose that $\lambda \geqslant-a_{0}$. Thus for all $a \geqslant 2 a_{0}$ we have $a^{2}(a+\lambda)^{-2} \leqslant 4$. Obviously,

$$
\left|\frac{a^{2}}{(\lambda+a)^{2}} f(\lambda) \xi\left(\lambda ; A_{\alpha V}, A_{0}\right)\right| \leqslant 4 \frac{\left|\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)\right|}{\left(\lambda+a_{0}\right)^{2}} \sup _{\lambda \in\left[-a_{0},+\infty\right)}\left(\lambda+a_{0}\right)^{2}|f(\lambda)| .
$$

Therefore, by (3.2) and by the Lebesgue dominated convergence theorem we have

$$
\lim _{a \rightarrow+\infty} a^{2} g_{a}(\alpha)=\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{\alpha V}, A_{0}\right) d \lambda .
$$

From Lemma 3.1 it follows now that the integral on the r.h.s. is concave with respect to $\alpha$. As noted above the concavity with respect to the coupling constant implies the concavity with respect to the perturbation. This remark completes the proof of Theorem 2 in the case $p=1$.

We turn to the case $p>1$ in the condition (7) and note that the operator $\left(A_{\alpha V}+a_{0}\right)^{-p}$ is neither convex nor concave with respect to $\alpha[1,16]$. To treat this case we need the following

Lemma 3.2 ([12 Theorem 1]; [27 Theorem 8.10.4]). Assume that $A^{p}-A_{0}^{p} \in \mathscr{I}_{1}$ for some $p>1$. Then $A-A_{0} \in \mathscr{I}_{q}$ for any $q>p$. Let $\left\{P_{n}\right\}_{n \in \mathbb{N}}$ be a strictly monotone family of finite dimensional orthogonal projections converging strongly to the identity operator I. Then

$$
\xi\left(\lambda ; A, A_{0}\right)=\lim _{n \rightarrow \infty} \xi\left(\lambda ; A_{0}+P_{n}\left(A-A_{0}\right) P_{n}, A_{0}\right)
$$

in $L^{1}\left(\mathbb{R} ; \lambda^{q-1} d \lambda\right)$, where $q$ is the smallest odd integer greater than $p$.
The following lemma generalizes Lemma 3.1 to the case $p>1$ :

Lemma 3.3. Let $f(\lambda) \geqslant 0$ be a nonincreasing function of bounded total variation. Then for all $a \geqslant a_{0}$

$$
g_{a}(\alpha)=\int_{\mathbb{R}} \frac{f(\lambda)}{(\lambda+a)^{q+1}} \xi\left(\lambda ; A_{\alpha V}, A_{0}\right) d \lambda
$$

is concave with respect to $\alpha$.
Proof. By the invariance principle the spectral shift function $\xi\left(\lambda ; A_{V}, A_{0}\right)$ can be represented in the form

$$
-\xi\left((\lambda+a)^{-1} ;\left(A_{V}+a\right)^{-1},\left(A_{0}+a\right)^{-1}\right) .
$$

We introduce the operator $W(\alpha)=\left(A_{\alpha V}+a\right)^{-1}-\left(A_{0}+a\right)^{-1}$. Let $\left\{P_{n}\right\}_{n \in \mathbb{N}}$ be a family of finite dimensional orthogonal projections as in Lemma 3.2. Consider $W_{n}(\alpha)=P_{n} W(\alpha) P_{n} \in \mathscr{J}_{1}$ and define

$$
\begin{aligned}
g_{a}^{(n)}(\alpha) & =-\int_{\mathbb{R}} \frac{f(\lambda)}{(\lambda+a)^{q+1}} \xi\left((\lambda+a)^{-1} ;\left(A_{0}+a\right)^{-1}+W_{n}(\alpha),\left(A_{0}+a\right)^{-1}\right) d \lambda \\
& =-\int_{0}^{\infty} f\left(\frac{1-a t}{t}\right) t^{q-1} \xi\left(t ;\left(A_{0}+a\right)^{-1}+W_{n}(\alpha),\left(A_{0}+a\right)^{-1}\right) d t
\end{aligned}
$$

with $q$ being defined as in Theorem 2. Recall that $W(\alpha)$ is convex with respect to $\alpha$ and therefore $P_{n} W(\alpha) P_{n}$ is also convex. Thus by Corollary 2 the function $g_{a}^{(n)}(\alpha)$ is concave for every $n \in \mathbb{N}$. To prove that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} g_{a}^{(n)}(\alpha)=g_{a}(\alpha) \tag{3.4}
\end{equation*}
$$

we estimate as

$$
\begin{aligned}
\left|g_{a}(\alpha)-g_{a}^{(n)}(\alpha)\right| \leqslant & \sup _{\lambda \in\left[-a_{0},+\infty\right)}|f(\lambda)| \\
& \times \int_{0}^{\infty} t^{q-1} \mid \xi\left(t ;\left(A_{0}+a\right)^{-1}+W(\alpha),\left(A_{0}+a\right)^{-1}\right) \\
& -\xi\left(t ;\left(A_{0}+a\right)^{-1}+W_{n}(\alpha),\left(A_{0}+a\right)^{-1}\right) \mid d t .
\end{aligned}
$$

By Lemma 3.2 the r.h.s. tends to zero thus proving (3.4) and completing the proof of the lemma.

To complete the proof of Theorem 2 as in the case $p=1$ we consider the limit $a \rightarrow+\infty$ of $a^{q+1} g_{a}(\alpha)$. By the inequality

$$
\begin{aligned}
\int_{\mathbb{R}} \frac{\left|\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)\right|}{(\lambda+a)^{q+1}} d \lambda & \leqslant \int_{\mathbb{R}} \frac{\left|\xi\left(\lambda ; A_{\alpha V}, A_{0}\right)\right|}{(\lambda+a)^{p+1}} d \lambda \\
& \leqslant\left\|\left(A_{\alpha V}+a\right)^{-p}-\left(A_{0}+a\right)^{-p}\right\|_{\mathscr{J}_{1}}
\end{aligned}
$$

valid for all $a \geqslant a_{0}+1$ and again by the Lebesgue dominated convergence theorem we obtain

$$
\lim _{a \rightarrow+\infty} a^{q+1} g_{a}(\alpha)=\int_{\mathbb{R}} f(\lambda) \xi\left(\lambda ; A_{\alpha V}, A_{0}\right) d \lambda .
$$

Now from Lemma 3.3 it follows that the integral on the r.h.s. is concave with respect to $\alpha$. This completes the proof of Theorem 2 .

We turn to the proof of Corollary 3.

Proposition 3.1. Under the assumptions of Corollary 3 (but without the restriction $V \geqslant 0$ ) the functional $g(V)$ is subadditive in the coupling constant in the sense that for arbitrary $\alpha_{1}, \alpha_{2} \geqslant 0$

$$
\begin{equation*}
g\left(\left(\alpha_{1}+\alpha_{2}\right) V\right) \leqslant g\left(\alpha_{1} V\right)+g\left(\alpha_{2} V\right) . \tag{3.5}
\end{equation*}
$$

Moreover, for arbitrary $\alpha_{1}, \alpha_{2} \geqslant 0$ the inequality

$$
\begin{equation*}
g\left(\left(\alpha_{1}-\alpha_{2}\right) V\right) \geqslant g\left(\alpha_{1} V\right)+g\left(-\alpha_{2} V\right) \tag{3.6}
\end{equation*}
$$

## holds.

Proof. The assumption (ii) of the Corollary 3 and the proof of Theorem XI. 12 in [22] imply that for an arbitrary finite interval $[a, b] \subset \mathbb{R}$ there is finite $a_{0} \in \mathbb{R}$ such that $\left(A_{\alpha V}+a\right)^{-p}-\left(A_{0}+a\right)^{-p} \in \mathscr{F}_{1}$ for all $a \geqslant a_{0}$. Thus we may set $\mathscr{C}\left(A_{0}, a_{0}, p\right)=\{\alpha V, \alpha \in[a, b]\}$. By Theorem 2 we obtain that $g(\alpha V)$ is concave with respect to $\alpha \in[a, b]$. Since $a$ and $b$ are arbitrary the function $g(\alpha V)$ is concave on $\mathbb{R}$. In the case of assumption (i) the concavity of $g(\alpha V)$ for all $\alpha \in \mathbb{R}$ is guaranteed directly by Theorem 1 .

Recall that the necessary and sufficient condition [11, Theorem 6.2.4] for a measurable concave function $\phi(\alpha)$ to be subadditive on $\mathbb{R}_{+}$is that $\phi(+0) \geqslant 0$. This proves (3.5). To prove (3.6) we use the fact (see, e.g., [ 10 , Theorem 110]) that any continuous concave function $\phi(x)$ satisfies the inequality

$$
\begin{equation*}
\phi\left(x-h^{\prime}\right)+\phi\left(x+h^{\prime}\right) \geqslant \phi(x-h)+\phi(x+h) \tag{3.7}
\end{equation*}
$$

provided that $|h| \geqslant\left|h^{\prime}\right|$. We set $x=\left(\alpha_{1}-\alpha_{2}\right) / 2, h^{\prime}=\left(\alpha_{2}-\alpha_{1}\right) / 2, h=\left(\alpha_{2}+\alpha_{1}\right) / 2$ and apply the inequality (3.7) to the function $g(\alpha V)$. Since $g(0)=0$ we arrive at the claim (3.6).

Proof of Corollary 3. Let $\gamma=\inf _{\alpha>0} \alpha^{-1} \phi(\alpha)$. Recall (see, e.g,. [11, Theorem 6.6.1]) that if $\phi(\alpha)$ is a measurable subadditive function, which is finite for all finite $\alpha$, then $-\infty \leqslant \gamma<\infty$ and

$$
\lim _{\alpha \rightarrow+\infty} \frac{\phi(\alpha)}{\alpha}=\gamma .
$$

We take $\phi(\alpha)=g(\alpha V)$. By Proposition 3.1 it is subadditive on $\mathbb{R}_{+}$. By the monotonicity propertiy of the spectral shift function the condition $V \geqslant 0$ implies that $\phi(\alpha) \geqslant 0$ for all $\alpha \in \mathbb{R}_{+}$. Therefore $\gamma \geqslant 0$, thus proving the corollary.

## ACKNOWLEDGMENTS

The author is indebted to V. Enss and R. Schrader for useful remarks. Stimulating communications with F. Gesztesy and K. Makarov are also gratefully acknowledged.

## REFERENCES

1. J. Bendat and S. Sherman, Monotone and convex operator functions, Trans. Amer. Math. Soc. 79 (1955), 58-71.
2. M. Sh. Birman and M. Z. Solomyak, Remarks on the spectral shift function, J. Soviet. Math. 3 (1975), 408-419.
3. M. Sh. Birman and D. R. Yafaev, The spectral shift function: Work by M. G. Krein and its further development, St. Petersburg Math. J. 4 (1993), 833-870.
4. R. Geisler, V. Kostrykin, and R. Schrader, Concavity properties of Krein's spectral shift function, Rev. Math. Phys. 7, (1995), 161-181.
5. F. Gesztesy, K. A. Makarov, and S. N. Naboko, The spectral shift operator, in "Mathematical Results in Quantum Mechanics, (J. Dittrich, P. Exner, and M. Tater, Eds.), Operator Theory: Advances and Applications, Vol. 108, pp. 59-90, Birkhäuser, Basel, 1999.
6. F. Gesztesy and K. A. Makarov, Some applications of the spectral shift operator, in "Fields Institute Communication Series" (A. G. Ramm, P. N. Shivakumar, and A. V. Strauss, Eds.), Amer. Math. Soc., Providence, in press.
7. F. Gesztesy and K. A. Makarov, The $\Xi$ operator and its relation to Krein's spectral shift operator, J. Anal. Math. in press; available from http://www.ma.utexas.edu/mp_arc/ as 99-112.
8. F. Gesztesy, K. A. Makarov, and A. K. Motovilov, Monotonicity and concavity properties of the spectral shift function, in "Canadian Mathematical Society's Conference Proceedings Series" (F. Gesztesy, H. Holden, J. Jost, S. Paycha, M. Röckner, and S. Scarlatti, Eds.), Canad. Math. Soc. and Amer. Math. Soc., Providence, in press; available from http://www.ma.utexas.edu/mp_arc/ as 99-339.
9. I. C. Gohberg and M. G. Krein, "Introduction to the Theory of Linear Nonselfadjoint Operators in Hilbert Space," Amer. Math. Soc. Transl. of Math. Monogr., Vol 18, Amer. Math. Soc., Providence, RI, 1969.
10. G. H. Hardy, J. E. Littlewood, and G. Pólya, "Inequalities," Cambridge Univ. Press, Cambridge, UK, 1934.
11. E. Hille, "Functional Analysis and Semigroups," Amer. Math. Soc. Colloq. Publ., Vol. 31, Amer. Math. Soc., New York, 1948.
12. L. S. Koplienko, On the theory of the spectral shift function, Problemy Mat. Fiz. 5 (1971), 62-79. [In Russian].
13. V. Kostrykin, "The Spectral Shift Function and Its Applications to Random Schrödinger Operators," Habilitationsschrift, RWTH Aachen, 1999.
14. V. Kostrykin and R. Schrader, Scattering theory approach to random Schrödinger operators in one dimension, Rev. Math. Phys. 11 (1999), 187-242.
15. V. Kostrykin and R. Schrader, The density of states and the spectral shift density of random Schrödinger operators, Rev. Math. Phys. 12 (2000), in press; available from http://www.ma.utexas.edu/mp_arc/ as 99-239.
16. F. Kraus, Über konvexe Matrixfunktionen, Math. Z. 41 (1936), 18-420.
17. E. H. Lieb and H. Siedentop, Convexity and concavity of eigenvalue sums, J. Statist. Phys. 63 (1991), 811-816.
18. G. K. Pedersen, " $C^{*}$-Algebras and Their Automorphism Groups," Academic Press, London, 1979.
19. A. B. Pushnitski, Representation of the spectral shift function for perturbations of fixed signs, St. Petersburg Math. J. 9 (1998), 1181-1194.
20. A. B. Pushnitski, Integral estimates for the spectral shift function, St. Petersburg Math. J. 10 (1999), 1047-1070.
21. A. B. Pushnitski, Spectral shift function of the Schrödinger operator in the large coupling constant limit, Comm. Partial Differential Equations 25 (2000), 703-736; available from http://www.ma.utexas.edu/mp_arc/ as 98-417.
22. M. Reed and B. Simon, "Methods of Modern Mathematical Physics, III. Scattering Theory," Academic Press, New York, 1979.
23. O. Safronov, Spectral shift function in the large coupling constant limit, preprint, 1999; available from http://www.ma.utexas.edu/mp_arc/ as 99-1332.
24. B. Simon, "Trace Ideals and Their Applications," Cambridge Univ. Press, New York, 1979.
25. B. Simon, Spectral averaging and the Krein spectral shift, Proc. Amer. Math. Soc. 126 (1998), 1409-1413.
26. W. Thirring, "Lehrbuch der mathematischen Physik Vol. III: Quantenmechanik von Atomen und Molekulen," Springer, Vienna, 1979.
27. D. R. Yafaev, "Mathematical Scattering Theory: General Theory," Amer. Math. Soc. Transl. Math. Monographs, Vol. 105, Amer. Math. Soc., Providence, 1992.
