# Taylor series associated with a differential-difference operator on the real line 

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#### Abstract

We extend the classical theory of Taylor series to a first-order differential-difference operator $\Lambda$ on the real line which includes as a particular case the Dunkl operator associated with the reflection group $\mathbf{Z}_{2}$ on R. More precisely, we establish first a generalized Taylor formula with integral remainder, and then specify sufficient conditions for a function on $\mathbf{R}$ to be expanded as a generalized Taylor series. Moreover, we provide a criterion of analyticity for functions on $\mathbf{R}$ involving the differential-difference operator $\Lambda$.


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## 0. Introduction

In this paper we consider the first-order differential-difference operator on $\mathbf{R}$

$$
\begin{equation*}
\Lambda f=\frac{\mathrm{d} f}{\mathrm{~d} x}+\frac{A^{\prime}(x)}{A(x)}\left(\frac{f(x)-f(-x)}{2}\right), \tag{1}
\end{equation*}
$$

where

$$
\begin{equation*}
A(x)=|x|^{2 \alpha+1} B(x), \quad \alpha \geqslant-\frac{1}{2}, \tag{2}
\end{equation*}
$$

$B$ being a positive $C^{\infty}$ even function on $\mathbf{R}$. In the case $A(x)=|x|^{2 \alpha+1}, \alpha \geqslant-\frac{1}{2}$, we regain the differential-difference operator

$$
\begin{equation*}
D_{\alpha} f=\frac{\mathrm{d} f}{\mathrm{~d} x}+\left(\alpha+\frac{1}{2}\right) \frac{f(x)-f(-x)}{x}, \tag{3}
\end{equation*}
$$

[^0]which is known as the Dunkl operator of index $\alpha+\frac{1}{2}$ associated with the reflection group $\mathbf{Z}_{2}$ on R. Dunkl operators are parameterized first-order differential-difference operators on some Euclidean space that are related to finite reflection groups. They are introduced in [4,5] in connection with a generalization of the classical theory of spherical harmonics. For the mathematical and physical applications of such operators we refer to the literature cited in [11]. For instance, the one-dimensional Dunkl operator $D_{\alpha}$ plays a major role in the study of quantum harmonic oscillators governed by Wigner's commutation rules [6,9].

A quite new commutative harmonic analysis on the real line related to the differential-difference operator $\Lambda$ was initiated in $[7,8]$ in which several analytic structures on $\mathbf{R}$ were generalized. Through this paper, the classical Taylor series theory on $\mathbf{R}$ is extended to the differential-difference operator ^. More explicitly, we establish in Section 1 the following generalized Taylor formula with integral remainder:

$$
\begin{equation*}
T^{x} f(y)=\sum_{p=0}^{n} b_{p}(y) \Lambda^{p} f(x)+\int_{-|y|}^{|y|} W_{n}(y, z) T^{x} \Lambda^{n+1} f(z) A(z) \mathrm{d} z, \tag{4}
\end{equation*}
$$

where $T^{x}, x \in \mathbf{R}$, stand for the generalized translation operators tied to the differential-difference operator $\Lambda ;\left\{W_{p}\right\}$ and $\left\{b_{p}\right\}, p=0,1,2, \ldots$, being two sequences of functions constructed inductively from the function $A$. In analogy to the classical setting, we determine sufficient conditions under which a $C^{\infty}$ function $f$ on $\mathbf{R}$ may be expanded as a generalized Taylor series in a neighborhood of an arbitrary point $x \in \mathbf{R}$; that is, conditions which ensure that for $|y|$ small enough, the integral remainder in (4) tends to 0 as $n \rightarrow \infty$. Moreover, it turns out that, except for the Dunkl operator case, the generalized Taylor series as discussed here are not power series. In other words, the $b_{p}, p=1,2,3, \ldots$, are in general not polynomials. Nevertheless, we provide in Section 2 a criterion of analyticity on $\mathbf{R}$ involving the differential-difference operator $\Lambda$; that is, a criterion characterizing an analytic function $f$ on $\mathbf{R}$ by means of the sequence $\left\{\Lambda^{p} f\right\}, p=0,1,2, \ldots$. The chief device in the proof of this criterion will be results from the theory of hypo-analytic operators (see [1]).

The notion of Taylor series was first extended in [2] to the Bessel differential operator $L_{\alpha}=\mathrm{d}^{2} / \mathrm{d} x^{2}+$ $((2 \alpha+1) / x) \mathrm{d} / \mathrm{d} x, \alpha \geqslant-\frac{1}{2}$. Such an extension was essentially aimed to allow a formal introduction of a generalized translation operation on the half line tied to the Bessel operator $L_{\alpha}$. Later, Trimèche [12] extended the theory of Delsarte to more general second-order differential operators of the form

$$
\begin{equation*}
L=\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+\frac{A^{\prime}(x)}{A(x)} \frac{\mathrm{d}}{\mathrm{~d} x}, \quad \alpha \geqslant-\frac{1}{2} . \tag{5}
\end{equation*}
$$

It is pointed out that all the results obtained in [12] emerge as easy consequences of those stated in the present article.

## 1. Generalized Taylor series

In the first part of this section we look for a Taylor formula with integral remainder, in which the differential-difference operator $\Lambda$ replaces the usual first derivative operator $\mathrm{d} / \mathrm{d} x$. Naturally, the construction of such a formula will require a number of preliminary results.

Notation. Let $\mathscr{C}$ be the subset of $\mathbf{R}^{2}$ defined by $\mathscr{C}=\left\{(x, y) \in \mathbf{R}^{2}: 0<|y| \leqslant|x|\right\}$.

Define on $\mathscr{C}$ two sequences of functions $\left\{u_{p}(x, y)\right\},\left\{v_{p}(x, y)\right\}, p=0,1,2, \ldots$, via the following recursive integral formulae:

$$
\begin{align*}
& u_{0}(x, y)=\frac{\operatorname{sgn}(x)}{2 A(x)}, v_{0}(x, y)=\frac{\operatorname{sgn}(y)}{2 A(y)}  \tag{6}\\
& u_{p+1}(x, y)=\int_{|y|}^{|x|} v_{p}(x, z) \mathrm{d} z, \quad v_{p+1}(x, y)=\frac{\operatorname{sgn}(y)}{A(y)} \int_{|y|}^{|x|} u_{p}(x, z) A(z) \mathrm{d} z . \tag{7}
\end{align*}
$$

This pair of families of functions enjoys the following properties.

Lemma 1. (i) $u_{p}, v_{p} \in C^{1}(\mathscr{C})$ for all $p=0,1,2, \ldots$, and satisfy for $0<|y|<|x|$ the relations

$$
\begin{array}{ll}
\Lambda_{x} u_{p+1}(x, y)=u_{p}(x, y), & \Lambda_{y} u_{p+1}(x, y)=-v_{p}(x, y), \\
\Lambda_{x} v_{p+1}(x, y)=v_{p}(x, y), & \Lambda_{y} v_{p+1}(x, y)=-u_{p}(x, y) . \tag{9}
\end{array}
$$

(ii) The sequences $\left\{u_{p}(x, y)\right\}$ and $\left\{v_{p}(x, y)\right\}$ may also be computed recursively for $p=2,3,4, \ldots$, by the formulae

$$
\begin{equation*}
u_{p}(x, y)=2 \int_{|y|}^{|x|} u_{p-2}(x, s) u_{1}(s, y) A(s) \mathrm{d} s, \quad v_{p}(x, y)=2 \int_{|y|}^{|x|} v_{p-2}(x, s) v_{1}(s, y) A(s) \mathrm{d} s . \tag{10}
\end{equation*}
$$

Proof. An induction argument gives assertion (i). Assertion (ii) follows easily by combining identities (7).

Remark 1. Appealing to (10), we show inductively that for any $p=0,1,2, \ldots, u_{2 p}(\cdot, y)$ is odd, $u_{2 p+1}(\cdot, y)$ is even, $\left|u_{2 p}(x, y)\right|=u_{2 p}(|x|, y)$, and $u_{2 p+1}(x, y) \geqslant 0$.

Notation. For $x \in \mathbf{R}$ put $\eta(x)=\inf _{|y| \leqslant|x|} B(y), \sigma(x)=\sup _{|y| \leqslant|x|} B(y)$, and $\omega(x)=\sigma(x) / \eta(x)$.
We shall need the following estimates.

Lemma 2. We have

$$
\left.\begin{array}{l}
0 \leqslant u_{1}(x, y) \leqslant \begin{cases}\frac{|y|^{-2 \alpha}-|x|^{-2 \alpha}}{4 \alpha \eta(x)} & \text { if } \alpha \neq 0, \\
\frac{\log (|x| /|y|)}{2 \eta(x)} & \text { if } \alpha=0,\end{cases} \\
\left|u_{2 p}(x, y)\right| \leqslant\left(\frac{x^{2} \omega(x)}{2 \alpha+2}\right)^{p} \frac{u_{1}(x, y)}{|x|} \text { for } p=1,2,3, \ldots,
\end{array}\right\} \begin{aligned}
& 0 \leqslant u_{2 p+1}(x, y) \leqslant\left(\frac{x^{2} \omega(x)}{2 \alpha+2}\right)^{p} u_{1}(x, y) \quad \text { for } p=0,1,2, \ldots
\end{aligned}
$$

Proof. Inequalities (11) follow readily from (6) and (7). Let us check (12). By (10) we have

$$
\begin{align*}
\left|u_{2}(x, y)\right| & =2 \int_{|y|}^{|x|}\left|u_{0}(x, s)\right| u_{1}(s, y) A(s) \mathrm{d} s=\frac{1}{A(x)} \int_{|y|}^{|x|} u_{1}(s, y) A(s) \mathrm{d} s \\
& \leqslant \frac{1}{A(x)} \int_{0}^{|x|} A(s) \mathrm{d} s u_{1}(x, y) \leqslant \frac{|x| \omega(x)}{2 \alpha+2} u_{1}(x, y) \tag{14}
\end{align*}
$$

This implies that (12) holds for $p=1$. Moreover, using (10) and (14) we find

$$
\begin{aligned}
\left|u_{4}(x, y)\right| & =2 \int_{|y|}^{|x|}\left|u_{2}(x, s)\right| u_{1}(s, y) A(s) \mathrm{d} s \leqslant 2 u_{1}(x, y) \int_{|y|}^{|x|}\left|u_{2}(x, s)\right| A(s) \mathrm{d} s \\
& \leqslant u_{1}(x, y) \frac{|x| \omega(x)}{\alpha+1} \int_{0}^{|x|} u_{1}(x, s) A(s) \mathrm{d} s .
\end{aligned}
$$

But by (11),

$$
\begin{equation*}
0 \leqslant \int_{0}^{|x|} u_{1}(x, s) A(s) \mathrm{d} s \leqslant \frac{x^{2} \omega(x)}{8(\alpha+1)} \tag{15}
\end{equation*}
$$

for any $\alpha \geqslant-\frac{1}{2}$. Therefore, (12) is true for $p=2$, and the full result follows by induction. Similarly, the majorization (13) is proved inductively by use of (10) and (15).

Define on $\mathbf{R}$ the family of functions $\left\{b_{p}\right\}$ by setting $b_{0}(x)=1$, and for $p=1,2,3, \ldots$,

$$
b_{p}(x)= \begin{cases}\int_{-|x|}^{|x|} u_{p-1}(x, y) A(y) \mathrm{d} y & \text { if } x \neq 0  \tag{16}\\ 0 & \text { if } x=0\end{cases}
$$

Remark 2. By Remark 1 it follows that for any $p=0,1,2, \ldots, b_{2 p}$ is even, $b_{2 p+1}$ is odd, and $b_{p}$ is positive on $] 0, \infty[$.

For our purpose of a Taylor formula involving the differential-difference operator $\Lambda$, a thorough investigation of the family $\left\{b_{p}\right\}$ seems unavoidable.

Proposition 1. For any $p=0,1,2, \ldots$,

$$
\begin{equation*}
b_{p}(x)=\mathcal{O}\left(x^{p}\right) \quad \text { as } x \rightarrow 0 . \tag{17}
\end{equation*}
$$

Proof. By (16) we have for $x \neq 0,\left|b_{1}(x)\right|=1 / A(x) \int_{0}^{|x|} A(y) \mathrm{d} y \leqslant|x| \omega(x) /(2 \alpha+2)=\mathcal{O}(x)$ as $x \rightarrow 0$. Moreover, from (12), (13) and (16) we deduce the estimates:

$$
\begin{align*}
& \left|b_{2 p+1}(x)\right| \leqslant\left(\frac{x^{2} \omega(x)}{2 \alpha+2}\right)^{p} \frac{b_{2}(x)}{|x|}, \quad p=1,2,3, \ldots  \tag{18}\\
& \left|b_{2 p+2}(x)\right| \leqslant\left(\frac{x^{2} \omega(x)}{2 \alpha+2}\right)^{p} b_{2}(x), \quad p=0,1,2, \ldots \tag{19}
\end{align*}
$$

As inequality (15) says that $0 \leqslant b_{2}(x) \leqslant x^{2} \omega(x) / 4(\alpha+1)$, we see that for any $p=1,2,3, \ldots, b_{p}(x)=$ $\mathcal{O}\left(x^{p}\right)$ when $x \rightarrow 0$.

Proposition 2. The functions $b_{p}, p=0,1,2, \ldots$ are of class $C^{1}$ on $\mathbf{R}$ and satisfy the relation

$$
\begin{equation*}
\Lambda b_{p+1}=b_{p} . \tag{20}
\end{equation*}
$$

Proof. From its expression (16) it is clear that $b_{1}$ is differentiable on $] 0, \infty[$, and for any $x>0$,

$$
\begin{equation*}
b_{1}^{\prime}(x)=1-\frac{A^{\prime}(x)}{A^{2}(x)} \int_{0}^{x} A(y) \mathrm{d} y=1-\left(\frac{2 \alpha+1}{B(x)}+\frac{x B^{\prime}(x)}{B^{2}(x)}\right) \int_{0}^{1} B(t x) t^{2 \alpha+1} \mathrm{~d} t, \tag{21}
\end{equation*}
$$

which tends to $1 /(2 \alpha+2)$ as $x \rightarrow 0^{+}$. This implies that $b_{1}$ is of class $C^{1}$ on $\mathbf{R}$ and $b_{1}^{\prime}(0)=1 /(2 \alpha+2)$. Further, it is immediate from (21) that $\Lambda b_{1}(x)=1$ for all $x \in \mathbf{R}$. Now fix $p=2,3,4, \ldots$. From (8) and (16) it is readily seen that $b_{p} \in C^{1}(\mathbf{R} \backslash\{0\})$ and

$$
\begin{equation*}
\Lambda b_{p}(x)=b_{p-1}(x) \quad \text { for all } x \neq 0 \tag{22}
\end{equation*}
$$

But due to Remark 2, identity (22) becomes $b_{p}^{\prime}(x)=b_{p-1}(x)$ for even $p$, and $b_{p}^{\prime}(x)=b_{p-1}(x)-$ $\left(A^{\prime} / A\right)(x) b_{p}(x)$ for odd $p$. Therefore $b_{p}^{\prime}(x)=\mathcal{O}\left(x^{p-1}\right)$ as $x \rightarrow 0$, by virtue of (2) and (17). This immediately shows that $b_{p} \in C^{1}(\mathbf{R}), b_{p}^{\prime}(0)=0$, and that equality (22) also holds for $x=0$.

Starting from identity (20), we shall prove inductively that the $b_{p}, p=0,1,2, \ldots$, are $C^{\infty}$ functions on $\mathbf{R}$. We begin with the following technical lemma proved by a standard argument.

Lemma 3. Let $f$ be a function of class $C^{n}$ on $[0, \infty[, n=0,1,2, \ldots$. Then the function

$$
H_{\alpha} f(x)= \begin{cases}\frac{1}{x^{\alpha}} \int_{0}^{x} f(t) t^{\alpha} \mathrm{d} t & \text { if } x>0 \\ 0 & \text { if } x=0\end{cases}
$$

is of class $C^{n+1}$ on $\left[0, \infty\left[\right.\right.$. Furthermore, $\left(H_{\alpha} f\right)^{(p)}(0)=p /(\alpha+p) f^{(p-1)}(0)$, for $p=1,2, \ldots, n+1$.
Notation. For a function $f: \mathbf{R} \rightarrow \mathbf{C}$ denote by $f_{\mathrm{e}}(x)=(f(x)+f(-x)) / 2, f_{\mathrm{o}}(x)=(f(x)-f(-x)) / 2$ its even and odd part, respectively. Define $\mathscr{E}$ as the space of $C^{\infty}$ complex-valued functions on $\mathbf{R}$, equipped with the topology of compact convergence of all derivatives. Let $\mathscr{E}_{\mathrm{e}}$ denote the subspace of $\mathscr{E}$ consisting of even functions. $\mathscr{E}_{*}$ stands for the subspace of $\mathscr{E}$ consisting functions $f$ such that $f(0)=0$.

Lemma 4. Let $m, n=0,1,2, \ldots$. Let $f$ be a function of class $C^{m}$ on $\mathbf{R}$ such that $\Lambda^{m} f$ be of class $C^{n}$ on $\mathbf{R}$. Then $f$ is of class $C^{m+n}$ on $\mathbf{R}$.

Proof. It is enough to consider the case where $m=1$. For even $f$, the result is obvious since $\Lambda f=f^{\prime}$ for such functions. For odd $f$, the result follows from Lemma 3, the relation $f(x)=$ $1 / A(x) \int_{0}^{x} \Lambda f(t) A(t) \mathrm{d} t$, and expression (2) of $A$. For arbitrary $f$ the lemma is a consequence of the relations $(\Lambda f)_{\mathrm{e}}=\Lambda\left(f_{\mathrm{o}}\right),(\Lambda f)_{\mathrm{o}}=\Lambda\left(f_{\mathrm{e}}\right)$.

It is now possible to state the following proposition.
Proposition 3. $b_{p} \in \mathscr{E}$ for each $p=0,1,2, \ldots$.
Proof. The result follows inductively by use of Proposition 2 and Lemma 4.
The role of the $b_{p}, p=0,1,2, \ldots$, in our generalized Taylor formula shall be analogous to that of the monomials $x^{p} / p!, p=0,1,2, \ldots$, in the classical Taylor formula. To specify the connection between the families $\left\{b_{p}\right\}$ and $\left\{x^{p} / p!\right\}$, it is useful to recall from [7] the following result.

Theorem 1. There exists a unique isomorphism $V$ of $\mathscr{E}$ such that

$$
\begin{equation*}
V \frac{\mathrm{~d}}{\mathrm{~d} x} f=\Lambda V f \quad \text { and } \quad V f(0)=f(0) \text { for all } f \in \mathscr{E} . \tag{23}
\end{equation*}
$$

The operator $V$ is said to be a transmutation operator between $\Lambda$ and $\mathrm{d} / \mathrm{d} x$ on the space $\mathscr{E}$. For $A(x)=|x|^{2 \alpha+1}, \alpha>-\frac{1}{2}$, this transmutation operator reads

$$
\begin{equation*}
V f(x)=\frac{\Gamma(\alpha+1)}{\sqrt{\pi} \Gamma(\alpha+1 / 2)} \int_{-1}^{1}\left(1-t^{2}\right)^{\alpha-1 / 2}(1+t) f(x t) \mathrm{d} t \tag{24}
\end{equation*}
$$

and is referred to as the Dunkl intertwining operator of index $\alpha+\frac{1}{2}$ associated with the reflection group $\mathbf{Z}_{2}$ on $\mathbf{R}$ (see [5,11]).

We claim the following statement.

Proposition 4. For any $n=0,1,2, \ldots$,

$$
\begin{equation*}
b_{n}(x)=V\left(\frac{y^{n}}{n!}\right)(x), \quad x \in \mathbf{R} \tag{25}
\end{equation*}
$$

In order to prove the proposition, we need the following simple lemma.

Lemma 5. The mapping $f \rightarrow \Lambda f$ is one-to-one from $\mathscr{E}_{*}$ onto $\mathscr{E}$. The inverse mapping is given by

$$
\Lambda^{-1} f(x)=\int_{0}^{x} f_{\mathrm{o}}(y) \mathrm{d} y+\frac{1}{A(x)} \int_{0}^{x} f_{\mathrm{e}}(y) A(y) \mathrm{d} y .
$$

Proof. If $f \in \mathscr{E}$ then (1) leads to $(\Lambda f)_{\mathrm{e}}=f_{\mathrm{o}}^{\prime}+A^{\prime} / A f_{\mathrm{o}}$, and $(\Lambda f)_{\mathrm{o}}=f_{\mathrm{e}}^{\prime}$. That is,

$$
\begin{equation*}
f(x)=\int_{0}^{x}(\Lambda f)_{0}(y) \mathrm{d} y+f(0)+\frac{1}{A(x)} \int_{0}^{x}(\Lambda f)_{\mathrm{e}}(y) A(y) \mathrm{d} y . \tag{26}
\end{equation*}
$$

This makes the result obvious.
Proof of Proposition 4. Set $c_{n}(x)=V\left(y^{n} / n!\right)(x), x \in \mathbf{R}, n=0,1,2, \ldots$. Notice that $c_{0}=b_{0}=1$ by virtue of (23) and (26). Moreover, (23) ensures that $c_{n} \in \mathscr{E}_{*}$ and $\Lambda c_{n}=c_{n-1}$ for all $n=1,2,3, \ldots$. So using Propositions $1-3$ and Lemma 5 we easily check by induction that $c_{n}=b_{n}$ for all $n=0,1,2, \ldots$, which is the required result.

Remark 4. If $A(x)=|x|^{2 \alpha+1}, \alpha \geqslant-\frac{1}{2}$, then identities (24) and (25) entail

$$
b_{2 p}(x)=\frac{1}{(\alpha+1)_{p} p!}\left(\frac{x}{2}\right)^{2 p}, \quad b_{2 p+1}(x)=\frac{1}{(\alpha+1)_{p+1} p!}\left(\frac{x}{2}\right)^{2 p+1}
$$

for all $p=0,1,2, \ldots$. Moreover, it is easily shown that the $b_{n}, n=0,1,2, \ldots$, are polynomials only in the Dunkl operator case.

Before we formulate our generalized Taylor formula, we need to introduce in the space $\mathscr{E}$ certain generalized translation operators $T^{a}, a \in \mathbf{R}$, tied to the differential-difference operator $\Lambda$. Such operators are defined in terms of the transmutation operator $V$ via the formula

$$
T^{a} f(x)=V_{a} V_{x}\left[V^{-1} f(a+x)\right], \quad x \in \mathbf{R} .
$$

Clearly for $A(x)=1$, we regain the ordinary translation operators on $\mathbf{R}: f \rightarrow \tau^{a} f(x)=f(a+x)$. The $T^{a}, a \in \mathbf{R}$, are linear bounded operators from $\mathscr{E}$ into itself, and possess the following fundamental properties:

$$
\begin{equation*}
T^{0}=\text { identity }, \quad T^{a} f(x)=T^{x} f(a) \text { and } \Lambda T^{a}=T^{a} \Lambda \tag{27}
\end{equation*}
$$

For more details about this generalized translation operation we refer to $[7,8]$. We can now state the first central result of this paper.

Theorem 2. Let $f \in \mathscr{E}$ and $a \in \mathbf{R}$. Then for any $n=0,1,2, \ldots$, we have the following generalized Taylor formula with integral remainder:

$$
\begin{equation*}
T^{a} f(x)=\sum_{p=0}^{n} b_{p}(x) \Lambda^{p} f(a)+\int_{-|x|}^{|x|} W_{n}(x, y) T^{a} \Lambda^{n+1} f(y) A(y) \mathrm{d} y, \tag{28}
\end{equation*}
$$

where

$$
\begin{equation*}
W_{n}(x, y)=u_{n}(x, y)+v_{n}(x, y) . \tag{29}
\end{equation*}
$$

In order to simplify the proof of Theorem 2, we first establish the following technical lemma.

Lemma 6. Let $f$ be a function of class $C^{p+1}$ on $\mathbf{R}, p=1,2,3, \ldots$. Then

$$
\int_{-|x|}^{|x|} W_{p}(x, y) \Lambda^{p+1} f(y) A(y) \mathrm{d} y=\int_{-|x|}^{|x|} W_{p-1}(x, y) \Lambda^{p} f(y) A(y) \mathrm{d} y-b_{p}(x) \Lambda^{p} f(0)
$$

Proof. First,

$$
\begin{align*}
\int_{-|x|}^{|x|} u_{p}(x, y) \Lambda^{p+1} f(y) A(y) \mathrm{d} y & =2 \int_{0}^{|x|} u_{p}(x, y)\left(\Lambda^{p+1} f\right)_{\mathrm{e}}(y) A(y) \mathrm{d} y \\
& =2 \int_{0}^{|x|} u_{p}(x, y) \frac{\mathrm{d}}{\mathrm{~d} y}\left[A(y)\left(\Lambda^{p} f\right)_{\mathrm{o}}(y)\right] \mathrm{d} y . \tag{30}
\end{align*}
$$

Observe that by Lemma 2, the function $y \rightarrow A(y) u_{p}(x, y)$ is bounded in a neighborhood of the origin. Further, the classical Taylor formula implies that $\left(\Lambda^{p} f\right)_{\mathrm{o}}(y)=c y+o(y)$ as $y \rightarrow 0$, for some
constant $c$. Therefore, integrating (30) by parts and applying (8) we obtain

$$
\begin{align*}
\int_{-|x|}^{|x|} u_{p}(x, y) \Lambda^{p+1} f(y) A(y) \mathrm{d} y & =-2 \int_{0}^{|x|} \Lambda_{y} u_{p}(x, y)\left(\Lambda^{p} f\right)_{\mathrm{o}}(y) A(y) \mathrm{d} y \\
& =\int_{-|x|}^{|x|} v_{p-1}(x, y) \Lambda^{p} f(y) A(y) \mathrm{d} y \tag{31}
\end{align*}
$$

Moreover,

$$
\begin{align*}
\int_{-|x|}^{|x|} v_{p}(x, y) \Lambda^{p+1} f(y) A(y) \mathrm{d} y & =2 \int_{0}^{|x|} v_{p}(x, y)\left(\Lambda^{p+1} f\right)_{0}(y) A(y) \mathrm{d} y \\
& =2 \int_{0}^{|x|} v_{p}(x, y) \frac{\mathrm{d}}{\mathrm{~d} y}\left(\Lambda^{p} f\right)_{\mathrm{e}}(y) A(y) \mathrm{d} y \tag{32}
\end{align*}
$$

As by (7) and (16), $\lim _{y \rightarrow 0^{+}} 2 A(y) v_{p}(x, y)=b_{p}(x)$; an integration by parts in (32), as well as (9) yield

$$
\begin{align*}
\int_{-|x|}^{|x|} v_{p}(x, y) \Lambda^{p+1} f(y) A(y) \mathrm{d} y & =-2 \int_{0}^{|x|} \Lambda_{y} v_{p}(x, y)\left(\Lambda^{p} f\right) \mathrm{e}(y) A(y) \mathrm{d} y-b_{p}(x) \Lambda^{p} f(0) \\
& =\int_{-|x|}^{|x|} u_{p-1}(x, y) \Lambda^{p} f(y) A(y) \mathrm{d} y-b_{p}(x) \Lambda^{p} f(0) \tag{33}
\end{align*}
$$

The result follows now by combining (31) and (33).
Proof of Theorem 2. Because of (27) it is sufficient to consider the case where $a=0$. By (6) we have

$$
\begin{aligned}
& \int_{-|x|}^{|x|} u_{0}(x, y) \Lambda f(y) A(y)=\frac{\operatorname{sgn}(x)}{A(x)} \int_{0}^{|x|} \frac{\mathrm{d}}{\mathrm{~d} y}\left[A(y) f_{\mathrm{o}}(y)\right] \mathrm{d} y=f_{\mathrm{o}}(x), \\
& \int_{-|x|}^{|x|} v_{0}(x, y) \Lambda f(y) A(y) \mathrm{d} y=\int_{0}^{|x|} f_{\mathrm{e}}^{\prime}(y) \mathrm{d} y=f_{\mathrm{e}}(x)-f(0)
\end{aligned}
$$

This yields the statement for $n=0$. For $n=1,2,3, \ldots$, we get identity (28) inductively by use of Lemma 6.

Remark 5. If $p$ is a polynomial of degree $k(k=1,2,3, \ldots)$, then (3) implies that $D_{\alpha} p$ is a polynomial of degree $k-1$. Therefore in the Dunkl operator case, the integral remainder in (28) will vanish whenever $f$ is a polynomial of degree $n$.

Our next purpose is to determine sufficient conditions under which a function $f$ in $\mathscr{E}$ may be expanded as a generalized Taylor series in the vicinity of an arbitrary point $a \in \mathbf{R}$. This will be achieved with some additional assumptions on the differential-difference operator $\Lambda$. For $\alpha>-\frac{1}{2}$, it was pointed out in [8] that the translation operators $T^{a}, a \in \mathbf{R}$, may be represented as

$$
\begin{equation*}
T^{a} f(x)=\int_{\mathbf{R}} f(y) \mathrm{d} \Omega_{a, x}(y), \quad f \in \mathscr{E} \tag{34}
\end{equation*}
$$

where for each $a, x \in \mathbf{R}, \Omega_{a, x}$ is a distribution on $\mathbf{R}$ with support in $[-|a|-|x|,-\| a|-|x||] \cup$ $\left[\| a|-|x|,|a|+|x|]\right.$. From now on we assume that the distributions $\Omega_{a, x}$ are uniformly norm-bounded measures, i.e, there is a constant $C>0$ such that

$$
\begin{equation*}
\left\|\Omega_{a, x}\right\| \leqslant C \quad \text { for all } a, x \in \mathbf{R} . \tag{35}
\end{equation*}
$$

With this additional assumption we readily check that

$$
\begin{equation*}
\left|T^{a} f(x)\right| \leqslant C \sup _{\|y|-|a \| \leqslant|x|}|f(y)| \quad \text { for all } a, x \in \mathbf{R} \quad \text { and all } f \in \mathscr{E} . \tag{36}
\end{equation*}
$$

Such an estimation will be the key tool in the proof of the next theorem.

Theorem 3. Let $f \in \mathscr{E}$ and $a \in \mathbf{R}$. Suppose that there are $M, \rho>0$ such that

$$
\begin{equation*}
\sup _{\|x|-| a\| \leqslant \rho}\left|\Lambda^{n} f(x)\right| \leqslant M^{n+1} n! \tag{37}
\end{equation*}
$$

for all $n=0,1,2, \ldots$. Then there exists an $r>0$ such that

$$
\begin{equation*}
T^{a} f(x)=\sum_{p=0}^{\infty} b_{p}(x) \Lambda^{p} f(a) \tag{38}
\end{equation*}
$$

for $|x| \leqslant r$. Furthermore, the series in (38) converges uniformly for $|x| \leqslant r$.
The following estimates for the functions $b_{p}, p=0,1,2, \ldots$, sharpen those given by (18) and (19), and may be useful in the proof of Theorem 3.

Proposition 5. For any $p=0,1,2, \ldots$, and $x \geqslant 0$,

$$
\begin{aligned}
& 0 \leqslant b_{2 p}(x) \leqslant \frac{1}{(\alpha+1)_{p} p!}\left(\frac{x^{2} \omega(x)}{4}\right)^{p} \\
& 0 \leqslant b_{2 p+1}(x) \leqslant \frac{1}{(\alpha+1)_{p+1} p!}\left(\frac{x}{2}\right)^{2 p+1}(\omega(x))^{p+1}
\end{aligned}
$$

Proof. Let $\left\{\tilde{u}_{p}\right\}$ denote the family $\left\{u_{p}\right\}$ corresponding to the Dunkl operator. By (6),

$$
\left|u_{0}(x, y)\right|=\frac{1}{2 A(x)} \leqslant \frac{1}{2 \eta(x)|x|^{2 \alpha+1}}=\frac{\left|\tilde{u}_{0}(x, y)\right|}{\eta(x)} .
$$

Further, by (6) and (7),

$$
0 \leqslant u_{1}(x, y)=\int_{|y|}^{|x|} \frac{\mathrm{d} z}{2 A(z)} \leqslant \frac{1}{\eta(x)} \int_{|y|}^{|x|} \frac{\mathrm{d} z}{2 z^{2 \alpha+1}}=\frac{\tilde{u}_{1}(x, y)}{\eta(x)} .
$$

An induction argument shows that for all $p=0,1,2, \ldots$,

$$
\begin{equation*}
\left|u_{p}(x, y)\right| \leqslant \frac{(\omega(x))^{[p / 2]}}{\eta(x)}\left|\tilde{u}_{p}(x, y)\right| \tag{39}
\end{equation*}
$$

where $[p / 2]$ is the integer part of $p / 2$. The proposition follows now by combining (16), (39) and Remark 4.

Proof of Theorem 3. Set $R_{n}(a, x)=\int_{-|x|}^{|x|} W_{n}(x, y) T^{a} \Lambda^{n+1} f(y) A(y) \mathrm{d} y$. By (36) and (37) it follows that $\left|R_{n}(a, x)\right| \leqslant C M^{n+2}(n+1)!\int_{-|x|}^{|x|}\left|W_{n}(x, y)\right| A(y) \mathrm{d} y$ for $0<|x| \leqslant \rho$. As by (29), $\left|W_{n}(x, y)\right| \leqslant$ $\left|u_{n}(x, y)\right|+\left|v_{n}(x, y)\right|=u_{n}(|x|, y)+v_{n}(|x|,|y|)$, we deduce from (7) and (16) that

$$
\begin{equation*}
\int_{-|x|}^{|x|}\left|W_{n}(x, y)\right| A(y) \mathrm{d} y \leqslant b_{n+1}(|x|)+|x| b_{n}(|x|) \tag{40}
\end{equation*}
$$

Therefore, $\left|R_{n}(a, x)\right| \leqslant C M^{n+2}(n+1)!\left(b_{n+1}(|x|)+|x| b_{n}(|x|)\right)$ for $0<|x| \leqslant \rho$. Now using Proposition 5 we get

$$
\left|R_{2 n}(a, x)\right| \leqslant C M^{2}|x|(\omega(x)+2(\alpha+n+1))\left(M^{2} x^{2} \omega(x)\right)^{n} \frac{(2 n+1)!}{2^{2 n+1}(\alpha+1)_{n+1} n!}
$$

and

$$
\left|R_{2 n+1}(a, x)\right| \leqslant C M\left(M^{2} x^{2} \omega(x)\right)^{n+1} \frac{(2 n+3)!}{2^{2 n+2}(\alpha+1)_{n+1}(n+1)!}
$$

for $0<|x| \leqslant \rho$. Applying Stirling's formula, we find

$$
\begin{equation*}
\left|R_{2 n}(a, x)\right|=\left(M^{2} x^{2} \omega(x)\right)^{n} \mathcal{O}\left(n^{-\alpha-1 / 2}\right) \text { and }\left|R_{2 n+1}(a, x)\right|=\left(M^{2} x^{2} \omega(x)\right)^{n+1} \mathcal{O}\left(n^{-\alpha+1 / 2}\right) \tag{41}
\end{equation*}
$$

for $0<|x| \leqslant \rho$ and $n \rightarrow \infty$. Choose an $r \in] 0, \rho\left[\right.$ such that $0<M^{2} r^{2} \omega(r)<1$. As the function $\omega$ is increasing on $\left[0, \infty\left[\right.\right.$, we see by (41) that $\lim _{n \rightarrow \infty} R_{n}(a, x)=0$ uniformly for $0<|x| \leqslant r$. This ends the proof by virtue of Theorem 2.

Remark 6. (i) According to Rösler [10], assumption (35) is satisfied in the Dunkl operator case, with $C=4$.
(ii) Delsarte and Lions [3] have introduced in $\mathscr{E}_{\mathrm{e}}$ translation operators $S^{a}, a \in \mathbf{R}$, tied to the differential operator $L$. It was indicated in [7] that $S^{a} f(x)=\left[T^{a} f(x)+T^{a} f(-x)\right] / 2$ for all $f \in \mathscr{E}_{\mathrm{e}}$. Accordingly, by taking $f$ even in Theorems 2 and 3, we easily regain the corresponding results proved in [12] for the operator $L$.

## 2. Analyticity criterion

As already indicated in Remark 4, the generalized Taylor series (38) are power series only in the Dunkl operator case. Hence it may be useful to provide a criterion of analyticity on $\mathbf{R}$ involving the differential-difference operator $\Lambda$. Throughout this section we assume that the function $B$ given by (2) is analytic on $\mathbf{R}$. Trimèche [12] has obtained the following analyticity criterion for even functions on $\mathbf{R}$ involving the differential operator $L$.

Theorem 4. Let $f$ be a function in $\mathscr{E}_{\mathrm{e}}$. Then $f$ is analytic if and only if for any $\rho>0$ there is an $M_{\rho}>0$ such that

$$
\sup _{0 \leqslant x \leqslant \rho}\left|L^{n} f(x)\right| \leqslant\left(M_{\rho}\right)^{n+1}(2 n)!\quad \text { for all } n=0,1,2, \ldots
$$

The purpose of this section is to establish an analogue of Theorem 4 for the differential-difference operator $\Lambda$. Namely, we claim the following result.

Theorem 5. In order that a function $f$ in $\mathscr{E}$ be analytic it is necessary and sufficient that for any $\rho>0$ there be an $M_{\rho}>0$ such that

$$
\begin{equation*}
\sup _{|x| \leqslant \rho}\left|\Lambda^{n} f(x)\right| \leqslant\left(M_{\rho}\right)^{n+1} n!\quad \text { for all } n=0,1,2, \ldots \tag{42}
\end{equation*}
$$

Proof. Let $f$ be an analytic function. By induction we check that for any $n=0,1,2, \ldots$,

$$
\begin{equation*}
\Lambda^{2 n} f=L^{n}\left(f_{\mathrm{e}}\right)+\frac{\mathrm{d}}{\mathrm{~d} x} L^{n} I\left(f_{\mathrm{o}}\right) \quad \Lambda^{2 n+1} f=L^{n+1} I\left(f_{\mathrm{o}}\right)+\frac{\mathrm{d}}{\mathrm{~d} x} L^{n}\left(f_{\mathrm{e}}\right), \tag{43}
\end{equation*}
$$

where $I$ denotes the map defined on $\mathscr{E}$ by $\operatorname{Ig}(x)=\int_{0}^{x} g(t) \mathrm{d} t, x \in \mathbf{R}$. These relations when combined with Theorem 4, together with the identity $(\mathrm{d} / \mathrm{d} x) g(x)=1 / A(x) \int_{0}^{x} L g(t) A(t) \mathrm{d} t$ which is valid for any $g \in \mathscr{E}_{\mathrm{e}}$, show that condition (42) is necessary. Conversely, suppose that an $f \in \mathscr{E}$ satisfy condition (42). Introduce the function $F(x, y)=\sum_{n=0}^{\infty}(i)^{n}\left(y^{n} / n!\right) \Lambda^{n} f(x)$. By hypothesis (42) and Lemma 4, we can find an $\varepsilon>0$ such that $F(x, y)$ be of class $C^{\infty}$ in the strip $\left.\mathbf{R} \times\right]-\varepsilon, \varepsilon[$. Moreover $F(x, y)$ satisfies for $x \in \mathbf{R}$ and $|y|<\varepsilon$ the differential-difference equation

$$
\begin{equation*}
\Lambda_{x}^{2} F(x, y)+\frac{\partial^{2}}{\partial y^{2}} F(x, y)=0 . \tag{44}
\end{equation*}
$$

Write $F=p+q$ with $p(x, y)=(F(x, y)+F(-x, y)) / 2$ and $q(x, y)=(F(x, y)-F(-x, y)) / 2$. According to (44), the functions $p$ and $q$ satisfy on $\mathbf{R} \times]-\varepsilon, \varepsilon$ [ the homogeneous differential equations

$$
\begin{align*}
& \Delta(x p(x, y))+\left(x \frac{A^{\prime}(x)}{A(x)}-2\right) \frac{\partial p}{\partial x}(x, y)=0  \tag{45}\\
& \Delta\left(x^{2} q(x, y)\right)+\left(x^{2} \frac{A^{\prime}(x)}{A(x)}-4 x\right) \frac{\partial q}{\partial x}(x, y)+\left(x^{2}\left(\frac{A^{\prime}(x)}{A(x)}\right)^{\prime}-2\right) q(x, y)=0 \tag{46}
\end{align*}
$$

where $\Delta=\partial^{2} / \partial x^{2}+\partial^{2} / \partial y^{2}$ is the Laplacian on $\mathbf{R}^{2}$. As by (2) the functions $x A^{\prime}(x) / A(x)$ and $x^{2}\left(A^{\prime}(x) / A(x)\right)^{\prime}$ are analytic in $\mathbf{R}^{2}$, we deduce from [1, Corollary 1.2] that both the left-hand sides of (45) and (46) are hypo-analytic operators on $\mathbf{R}^{2}$. Therefore $p$ and $q$ are analytic in the strip $\mathbf{R} \times]-\varepsilon, \varepsilon[$, and so is $F$. To conclude the proof observe that $f(x)=F(x, 0)$ for all $x \in \mathbf{R}$.

Remark 7. (i) According to identities (43), Theorem 5 immediately implies Theorem 4.
(ii) From (20) and Theorem 5 it follows that the $b_{n}, n=0,1,2, \ldots$, are analytic functions on $\mathbf{R}$.

Corollary 1. $T^{a} f$ is analytic whenever $f$ is analytic and $a \in \mathbf{R}$.
Proof. Let $f$ be analytic and $a \in \mathbf{R}$. By (27) and (36) we have for any $\rho>0$ and $n=0,1,2, \ldots$,

$$
\sup _{|x| \leqslant \rho}\left|\Lambda^{n} T^{a} f(x)\right|=\sup _{|x| \leqslant \rho}\left|T^{a} \Lambda^{n} f(x)\right| \leqslant \sup _{\|x|-| a\| \leqslant \rho}\left|\Lambda^{n} f(x)\right| \leqslant \sup _{|x| \leqslant|a|+\rho}\left|\Lambda^{n} f(x)\right| .
$$

From this and Theorem 5 we deduce that the function $T^{a} f$ is analytic on $\mathbf{R}$.

Remark 8. A combination of Theorems 3, 5 and Corollary 1 shows that any analytic function on $\mathbf{R}$ may be expanded in a generalized Taylor series in a neighborhood of an arbitrary point $a \in \mathbf{R}$.

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