

Finite vertex algebras and nilpotence

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Abstract

I show that simple finite vertex algebras are commutative, and that the Lie conformal algebra structure underlying a reduced (= without nilpotent elements) finite vertex algebra is nilpotent.

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1. Introduction

In this paper I investigate the effect of a finiteness assumption on the singular part of the Operator Product Expansion of quantum fields belonging to a vertex algebra. The vertex algebra structure encodes algebraic properties of chiral fields in a 2-dimensional Conformal Field Theory. Its axiomatic definition was given by Borcherds in [4], and amounts to associating with each element v of a vector space V a *vertex operator*, or quantum field,

$$Y(v, z) \in (\text{End}V)[[z, z^{-1}]],$$

satisfying singular generalizations of the commutativity and unit axioms for the left multiplication operators in an associative algebra. This structure captures all algebraic properties of families of mutually local fields, containing the identity field, acting on some vector space of *physical states* V .

The products encoding the algebraic properties of a vertex algebra structure are typically written in terms of a formal expansion

$$Y(a, z)Y(b, w) = \sum_{j=0}^{\infty} \frac{Y(a_{(j)}b, w)}{(z-w)^{j+1}} + :Y(a, z)Y(b, w): \quad (1.1)$$

of the composition of quantum fields, called the Operator Product Expansion (OPE for short). The singular part

$$\sum_{j=0}^{\infty} \frac{Y(a_{(j)}b, w)}{(z-w)^{j+1}}$$

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of the OPE (1.1) only depends on the commutation properties of fields $Y(\cdot, z)$, and is in many ways remindful of a Lie algebra; the regular part $: Y(a, z)Y(b, w) :$ is called *normally ordered product* or Wick product, and essentially depends on the action of quantum fields on V . It is to some extent similar to an associative commutative product.

One can disregard the latter part of the structure, and axiomatize the singular part of the OPE only. The thus obtained algebraic structure is called “(Lie) conformal algebra”, but it is also known in the literature as vertex Lie algebra [8,13], or Lie pseudoalgebra [1] over $\mathbb{C}[\partial]$. Lie conformal algebras were introduced by Kac [11] to characterize algebraic properties of pairwise local formal distributions (in z and z^{-1}) with values in a Lie algebra. Hence every vertex algebra is thus also a Lie conformal algebra, and the latter structure measures the failure of the normally ordered product from being associative and commutative.

Indeed, if the Lie conformal algebra underlying a vertex algebra is trivial, then the product defined as $a \circ b = Y(a, z)b|_{z=0}$ gives [4] a commutative associative algebra structure on V from which it is possible to recover the vertex algebra product. In fact, in this case $Y(a, z)b = (e^{zT}a) \circ b$ – where T is the (infinitesimal) translation operator $Y(Ta, z) = dY(a, z)/dz$ – and is therefore completely determined by \circ .

The main interest in the role of Lie conformal algebras in vertex algebra theory is due to the existence [11–14] of a *universal enveloping vertex algebra* functor, which is adjoint to the forgetful functor from vertex to Lie conformal algebras. Many interesting vertex algebras are in fact obtained as simple quotients of the universal enveloping vertex algebra associated with a suitably chosen (and typically much smaller) Lie conformal algebra. However, a different strategy is possible: one could, in principle, analyse possible vertex algebra structures by first studying the singular OPE (i.e. the underlying Lie conformal algebra), and then inserting the normally ordered product on top of this. My aim in this note is to show that this strategy gives interesting results when V is a finitely generated $\mathbb{C}[T]$ -module.

The physically interesting vertex algebras which are usually considered are called Vertex Operator Algebras (=VOAs). They are graded vector fields that are often endowed with some additional structure (e.g. a Virasoro field inducing the grading); however, all known examples are very large objects — typically of superpolynomial growth. A first explanation for this is Borchers’ observation [5] that in a finite-dimensional vertex algebra all products $Y(a, z)b$ are necessarily regular in z , as the underlying Lie conformal algebra must be trivial (because all elements are torsion, see [7,11]). The vertex algebra structure reduces, as mentioned above, to that of a unital finite-dimensional commutative associative algebra with a derivation T . Also, in the presence of a grading induced by a Virasoro field the dimension of the homogeneous component of degree n is at least the number of partitions of n , which grows super-polynomially. However, no interesting (e.g. simple) examples are known of a vertex algebra structure on a graded vector space of polynomial growth, even in the absence of the additional requirements for a VOA.

After finite-dimensional vector spaces, the next to easiest case is that of finitely generated modules over $\mathbb{C}[T]$: they are vector spaces of linear growth when a grading is given, and this is the case I handle in this note. There seems to be no previously known description of algebraic properties of such *finite vertex algebras* that are not finite-dimensional.

Section 2 contains a list of definitions and basic results in the theory of Lie conformal algebras and vertex algebras. I also exhibit the motivating [Example 2.2](#), showing that there exist finite vertex algebras that do not reduce to associative commutative algebras. However, the vertex algebra provided in this example is constructed by means of nilpotent elements: as in the case of commutative algebras, such elements form an ideal of the vertex algebra which I call *nilradical*. The quotient of V by its nilradical has no nonzero (strongly) nilpotent elements.

Sections 3.1 and 3.2 apply results from [6] to the case of finite vertex algebras in order to show that finite simple vertex algebras are commutative ([Theorem 3.2](#)), hence that the Lie conformal algebra structure underlying a finite vertex algebra is always solvable ([Theorem 3.3](#)).

In Section 4, I study how a finite vertex algebra decomposes under the adjoint action of a Lie conformal subalgebra. After describing the representation theory of finite solvable Lie conformal algebras I show, in [Theorem 4.4](#), that the generalized weight submodule with respect to any nonzero weight is an abelian ideal of the vertex algebra structure. The presence of abelian ideals witnesses the existence of nilpotent elements, therefore there can be no nonzero weights in the absence of nilpotent elements. This strong algebraic fact is the basis for the results presented in Section 5, and is proved by means of the identity (3.2) introduced in Section 3.1.

The main result from last section is [Theorem 5.2](#) stating that any element s lying in a finite vertex algebra V with trivial nilradical has a nilpotent adjoint conformal action on V . By a Lie conformal algebra analogue of Engel’s theorem, developed in [7], the Lie conformal algebra underlying V must indeed ([Theorem 5.1](#)) be nilpotent. This statement essentially depends on both the finiteness assumption and the presence of a vertex algebra structure: it

basically means that every finite vertex algebra may be described as an extension of a nilpotent (as a Lie conformal algebra) vertex algebra by an ideal only containing nilpotent elements (i.e. contained in the nilradical).

The representation theory of the Virasoro Lie algebra or of affine Kac–Moody algebras is never used. The spirit of this paper is that the interplay between the Operator Product Expansion and the λ -bracket, in the case of a vertex algebra linearly generated by a finite number of quantum fields together with their derivatives, is strong enough to allow one to prove a number of results in a totally elementary way, even in the absence of a grading.

Some of the ideas contained in this work originate from an old manuscript written between 1998 and 1999 while I was visiting Université de Paris VI “Pierre et Marie Curie” and Université de Strasbourg “Louis Pasteur” as a European Union TMR post-doc. I would like to thank both institutions for hospitality.

2. Generalities on vertex and Lie conformal algebras

2.1. Vertex algebras

In what follows I quote some well-known facts about vertex algebras: precise statements and proofs can be found in [11]. Let V be a complex vector space. A *field* on V is an element $\phi(z) \in (\text{End}V)[[z, z^{-1}]]$ with the property that $\phi(v) \in V((z)) = V[[z]][z^{-1}]$ for every $v \in V$. In other words, if

$$\phi(z) = \sum_{i \in \mathbb{Z}} \phi_i z^{-i-1}$$

then $\phi_n(v) = 0$ for sufficiently large n .

Definition 2.1. A vertex algebra is a (complex) vector space V endowed with a linear *state-field correspondence* $Y : V \rightarrow (\text{End}V)[[z, z^{-1}]]$, a *vacuum element* $\mathbf{1}$ and a linear endomorphism $T \in \text{End}V$ satisfying the following properties:

- **Field axiom:** $Y(v, z)$ is a field for all $v \in V$.
- **Locality axiom:** For every $a, b \in V$

$$(z - w)^N [Y(a, z), Y(b, w)] = 0$$

for sufficiently large N .

- **Vacuum axiom:** The vacuum element $\mathbf{1}$ is such that

$$Y(\mathbf{1}, z) = \text{id}_V, \quad Y(a, z)\mathbf{1} \equiv a \pmod{zV[[z]]},$$

for all $a \in V$.

- **Translation invariance:** T satisfies

$$[T, Y(a, z)] = Y(Ta, z) = \frac{d}{dz} Y(a, z),$$

for all $a \in V$.

Note that the vector space V carries a natural $\mathbb{C}[T]$ -module structure. Fields $Y(a, z)$ are called *vertex operators*, or *quantum fields*.

A vertex algebra is a family of pairwise local fields acting on V containing the identity (constant) field. Indeed, every family of pairwise local fields containing the identity field can be realized as a vertex algebra up to changing the vector space V of physical states (see [11]). The vertex algebra structure therefore captures all algebraic aspects of families of pairwise local fields. A vertex algebra V is *finite* if V is a finitely generated $\mathbb{C}[T]$ -module.

There are two basic constructions of new vertex operators from two given ones. The first one is given by rephrasing what we earlier called “singular OPE”: since $(z - w)^N$ kills the commutator $[Y(a, z), Y(b, w)]$, the latter may be expanded into a linear combination:

$$\sum_{j=0}^{N-1} c_j(w) \frac{\delta^{(j)}(z-w)}{j!}, \tag{2.1}$$

where

$$\delta(z - w) = \sum_{j \in \mathbb{Z}} w^j z^{-j-1}$$

is the Dirac delta formal distribution and $\delta^{(j)}$ its j .th derivative with respect to w . The uniquely determined fields $c_j(w)$ are then vertex operators $Y(c_j, w)$ corresponding to elements $c_j = a_{(j)}b = a_{(j)}(b)$ where the $a_{(j)} \in \text{End}V$ are the coefficients of

$$Y(a, z) = \sum_{j \in \mathbb{Z}} a_{(j)}z^{-j-1}.$$

It is customary to view the \mathbb{C} -bilinear maps $a \otimes b \mapsto a_{(j)}b, j \in \mathbb{Z}$, as products describing the vertex algebra structure. Locality can be rephrased by stating that commutators between coefficients of quantum fields satisfy:

$$[a_{(m)}, b_{(n)}] = \sum_{j \geq 0} \binom{m}{j} (a_{(j)}b)_{(m+n-j)}, \tag{2.2}$$

for all $a, b \in V$.

Another way to put together quantum fields to produce new ones is given by the normally ordered product (or Wick product) defined as

$$: Y(a, z)Y(b, z) := Y(a, z)_+Y(b, z) + Y(b, z)Y(a, z)_-, \tag{2.3}$$

where

$$Y(a, z)_- = \sum_{j \in \mathbb{N}} a_{(j)}z^{-j-1}, \quad Y(a, z)_+ = Y(a, z) - Y(a, z)_-. \tag{2.4}$$

Then $: Y(a, z)Y(b, z) :$ is also a vertex operator, and it equals $Y(a_{(-1)}b, z)$.

Example 2.1. Let V be a unital associative commutative algebra, T a derivation of V . Then setting $Y(a, z)b = (e^{zT}a)b$ and choosing the unit $\mathbf{1} \in V$ to be the vacuum element makes V into a vertex algebra.

Such a vertex algebra is called *holomorphic*¹ in [11], and is the “uninteresting” case of a vertex algebra structure. It occurs whenever all vertex operators are regular in z : in this case one can always construct an associative commutative algebra, together with a derivation, inducing the vertex algebra structure as in Example 2.1. This is always the case when V is finite-dimensional [4,5]. One of the consequences of the vertex algebra axioms is the following:

- **Skew-commutativity:** $Y(a, z)b = e^{zT}Y(b, -z)a$

for all choices of a, b .

If A and B are subsets of V , then we may define $A \cdot B$ as the \mathbb{C} -linear span of all products $a_{(j)}b$, where $a \in A, b \in B, j \in \mathbb{Z}$. If B is a $\mathbb{C}[T]$ -submodule of V , then $A \cdot B$ is also a $\mathbb{C}[T]$ -submodule of V , as by translation invariance T is a derivation of all j -products, and $(Ta)_{(j)} = -ja_{(j-1)}$.

Notice that by skew-commutativity, $A \cdot B$ is contained in the $\mathbb{C}[T]$ -submodule generated by $B \cdot A$; equality $A \cdot B = B \cdot A$ then holds whenever $A \cdot B$ and $B \cdot A$ are both $\mathbb{C}[T]$ -submodules of V . Also, observe that $A \subset A \cdot V$ by the vacuum axiom and that $A \cdot V$ is always a $\mathbb{C}[T]$ -submodule of V , as $a_{(-2)}\mathbf{1} = Ta$. In particular, $a \cdot V = \mathbb{C}a \cdot V$ is a $\mathbb{C}[T]$ -module of V containing a .

Before proceeding, recall that a *subalgebra* of a vertex algebra V is a $\mathbb{C}[T]$ -submodule U containing $\mathbf{1}$ such that $U \cdot U = U$. In other words, all coefficients of $Y(a, z)b$ belong to U whenever a and b do. Similarly, a $\mathbb{C}[T]$ -submodule I of V is an *ideal* if $I \cdot V \subset I$; skew-commutativity then shows that $V \cdot I = I \cdot V$. A proper ideal can never contain the vacuum $\mathbf{1}$, and if M is an ideal of V , then $M + \mathbb{C}\mathbf{1}$ is a subalgebra, whose rank as a $\mathbb{C}[T]$ -module equals that of M .

¹ Notation and terminology are often contrasting in the vertex algebra world, and this is no exception. Notice that in most of the literature vertex operator algebras are known to be holomorphic if they have a semisimple representation theory, and the adjoint representation is the unique irreducible module, see [9].

The quotient V/I of a vertex algebra V by an ideal I has a unique vertex algebra structure making the canonical projection $\pi : V \rightarrow V/I$ a vertex algebra homomorphism, i.e. a $\mathbb{C}[T]$ -homomorphism such that $\pi(a_{(n)}b) = \pi(a)_{(n)}\pi(b)$ for every $a, b \in V, n \in \mathbb{Z}$.

A vertex algebra V is *commutative* if all quantum fields $Y(a, z), a \in V$ commute with one another; equivalently, if $a_{(n)}b = 0$ for all $a, b \in V, n \geq 0$. Commutative vertex algebras are all as in Example 2.1. The *centre* of V is the subspace of all elements $c \in V$ such that $a_{(n)}c = 0 = c_{(n)}a$ for all $a \in V, n \geq 0$. Then, by (2.2), coefficients of $Y(c, z)$ commute with coefficients of all quantum fields.

Lemma 2.1. *Assume that c lies in the centre of a vertex algebra V . Then the subspace $\ker c_{(-1)}$ is stable under the action of coefficients of all quantum fields $Y(a, z), a \in V$. In particular, $\ker c_{(-1)}$ is an ideal of V as soon as it is a $\mathbb{C}[T]$ -submodule, e.g. when $Tc = 0$.*

Proof. By (2.2) we have

$$a_{(m)}(c_{(-1)}x) - c_{(-1)}(a_{(m)}x) = \sum_{j \geq 0} \binom{m}{j} (a_{(j)}c)_{(m-j-1)}x,$$

for all $a \in V, m \in \mathbb{Z}$. Since c lies in the centre of V , the right-hand side vanishes; hence, if $x \in \ker c_{(-1)}$, it follows that $a_{(m)}x \in \ker c_{(-1)}$ as well. The last claim follows from the fact that $c_{(-1)}(Tx) = T(c_{(-1)}x) - (Tc)_{(-1)}x$. \square

2.2. A noncommutative finite vertex algebra

The following is an example of a vertex algebra structure on a finitely generated $\mathbb{C}[T]$ -module for which some positive products $u_{(j)}v, j \geq 0$ are nonzero:

Example 2.2. Let $V = \mathbb{C}[T]a \oplus \mathbb{C}[T]b \oplus \mathbb{C}\mathbf{1}$. Define $\mathbf{1}$ to be the vacuum element of V and set

$$\begin{aligned} Y(a, z)b &= Y(b, z)a = Y(b, z)b = 0, \\ Y(\mathbf{1}, z) &= \text{id}_V \\ Y(a, z)\mathbf{1} &= e^{zT}a \quad Y(b, z)\mathbf{1} = e^{zT}b \\ Y(a, z)a &= e^{zT/2}\psi(z)b, \end{aligned}$$

where $\psi(z) = \psi(-z)$ is any even Laurent series in z . Extend by \mathbb{C} -linearity the state-field correspondence Y to all of V after setting

$$Y(Tu, z)v = \frac{d}{dz}Y(u, z)v$$

and

$$Y(u, z)(Tv) = \left(T - \frac{d}{dz}\right)(Y(u, z)v),$$

so that translation invariance is satisfied. The only vertex algebra axiom still to check is locality, and the only nontrivial statement to prove is

$$(z - w)^n[Y(a, z), Y(a, w)]\mathbf{1} = 0,$$

for some n . However, we have

$$\begin{aligned} [Y(a, z), Y(a, w)]\mathbf{1} &= Y(a, z)e^{wT}a - Y(a, w)e^{zT}a \\ &= e^{(z+w)T/2} (\iota_{z,w}\psi(z-w) - \iota_{w,z}\psi(w-z))b, \end{aligned}$$

where $\iota_{z,w}$ (resp. $\iota_{w,z}$) is a prescription to consider the expansion in the domain $|z| > |w|$ (resp. in the domain $|w| > |z|$), see [11]. If we choose n so that $z^n\psi(z)$ is regular in z , multiplication by $(z - w)^n$ makes the above expression zero, due to the fact that $\psi(z) = \psi(-z)$: it is an *expansion of zero* in the sense of [10]. Notice that if we choose a nonregular $\psi(z)$, then at least one of the products $a_{(j)}a, j \geq 0$, is nonzero, so that V is noncommutative.

In the example above, elements a and b are *nilpotent*, in a sense that we are about to clarify.

2.3. The nilradical

An ideal I of a vertex algebra V is *abelian* if $I^2 = I \cdot I = 0$. An element $a \in V$ is *strongly nilpotent* of degree n if every product of elements in V containing a at least n times, under all j -products and any parenthesization, gives 0.

Let $x \in V$ be a strongly nilpotent element of degree $n > 2$, and a be a nonzero product of $[(n + 1)/2]$ copies of x . Then a is strongly nilpotent of degree two.

Lemma 2.2. *Let $a \in V$ be a strongly nilpotent element of degree two. Then a generates an abelian ideal of V .*

Proof. Clear. \square

Remark 2.1. An element $a \in V$ is (nonstrongly) nilpotent of degree n if every product of at least n copies of a , under any product and parenthesization gives 0. If V is either commutative or graded, then it is easy to show that every nilpotent element is strongly nilpotent; this holds in general, as can be proved using [2, Remark 7.8]. In particular, $Y(a, z)a = 0$ guarantees that $a \cdot V$ is an abelian ideal of V .

Corollary 2.1. *A vertex algebra V possesses nonzero strongly nilpotent elements if and only if it contains a nonzero abelian ideal.*

Proof. Every nonzero element in an abelian ideal is strongly nilpotent of degree 2. The converse is Lemma 2.2. \square

Let us now denote

$$I^1 = I, \quad I^{n+1} = I^n \cdot I^n, \quad n > 0.$$

Then I is a *nil-ideal* if $I^n = 0$ for sufficiently large values of n .

Lemma 2.3. *Let V be a vertex algebra, $N \subset V$ a nil-ideal, $\pi : V \rightarrow V/N$ the natural projection. Then $I \subset V$ is a nil-ideal if and only if $\pi(I)$ is a nil-ideal of V/N .*

Proof. We have $\pi(I^n) = \pi(I)^n$, hence if I is a nil-ideal, $\pi(I)$ is too. On the other hand, if $\pi(I)^n = \pi(I^n) = 0$, then $I^n \subset N$. Thus, $I^{n+k} \subset N^k$ which is zero for sufficiently large k . \square

Corollary 2.2. *The sum of nil-ideals is a nil-ideal.*

Proof. Let I, J be nil-ideals of V , and let $\pi : V \rightarrow V/I$ be the natural projection. Then $\pi(I + J)$ equals $\pi(J)$ which is a nil-ideal. \square

Corollary 2.3. *Let V be a finite vertex algebra. Then V has a unique maximal nil-ideal.*

Proof. Existence of some maximal nil-ideal follows from finiteness of the $\mathbb{C}[T]$ -module V , which is therefore Noetherian; uniqueness from Corollary 2.2. \square

The unique maximal nil-ideal of a finite vertex algebra V is called the *nilradical* $N(V)$ of V . It is clear that every strongly nilpotent element of V lies in $N(V)$. Furthermore, the quotient $V/N(V)$ has no strongly nilpotent elements, hence it has a trivial nilradical. We will call a vertex algebra with a trivial nilradical, or equivalently with no nontrivial abelian ideal, a *reduced* vertex algebra.

Remark 2.2. It is true in general that an element lies in the nilradical $N(V)$ if and only if it is nilpotent, hence strongly nilpotent; however, we will not need this fact.

2.4. Lie conformal algebras

Algebraic properties of commutators of quantum fields are encoded in the notion of Lie conformal algebra. See [7] for generalities on (Lie) conformal algebras and λ -brackets.

Definition 2.2. A *Lie conformal algebra* is a $\mathbb{C}[\partial]$ -module R with a \mathbb{C} -bilinear product $(a, b) \mapsto [a_\lambda b] \in V[\lambda]$ satisfying the following axioms:

- (C1) $[a_\lambda b] \in R[\lambda]$,
- (C2) $[\partial a_\lambda b] = -\lambda[a_\lambda b]$, $[a_\lambda \partial b] = (\partial + \lambda)[a_\lambda b]$,
- (C3) $[a_\lambda b] = -[b_{-\partial-\lambda} a]$,
- (C4) $[a_\lambda [b_\mu c]] - [b_\mu [a_\lambda c]] = [[a_\lambda b]_{\lambda+\mu} c]$,

for every $a, b, c \in V$.

Any vertex algebra V can be given a $\mathbb{C}[\partial]$ -module structure by setting $\partial = T$. Then defining

$$[a_\lambda b] = \sum_{n \in \mathbb{N}} \frac{\lambda^n}{n!} a_{(n)} b$$

endows V with a Lie conformal algebra structure. Indeed, (C1) follows from the field axiom, (C2) from translation invariance, (C3) from skew-commutativity, and (C4) from (2.2). In all that follows we will denote the infinitesimal translation operator T in a vertex algebra by ∂ .

If A and B are subspaces of a Lie conformal algebra R , then we may define $[A, B]$ as the \mathbb{C} -linear span of all λ -coefficients in the products $[a_\lambda b]$, where $a \in A, b \in B$. It follows from axiom (C2) that if B is a $\mathbb{C}[\partial]$ -submodule of R , then $[A, B]$ is also a $\mathbb{C}[\partial]$ -submodule of R . Notice that if A and B are both $\mathbb{C}[\partial]$ -submodules, then $[A, B] = [B, A]$ by axiom (C3). A *subalgebra* of a Lie conformal algebra R is a $\mathbb{C}[\partial]$ -submodule $S \subset R$ such that $[S, S] \subset S$.

A Lie conformal algebra R is *solvable* if, after defining

$$R^{(0)} = R, \quad R^{(n+1)} = [R^{(n)}, R^{(n)}], \quad n \geq 0,$$

we find that $R^{(N)} = 0$ for sufficiently large N . R is solvable iff it contains a solvable ideal S such that R/S is again solvable. Solvability of a nonzero Lie conformal algebra R trivially fails if R equals its *derived subalgebra* $R' = [R, R]$. Similarly, R is *nilpotent* if, after defining

$$R^{[0]} = R, \quad R^{[n+1]} = [R, R^{[n]}], \quad n \geq 0, \tag{2.5}$$

we find that $R^{[N]} = 0$ for sufficiently large N .

An *ideal* of a Lie conformal algebra R is a $\mathbb{C}[\partial]$ -submodule $I \subset R$ such that $[R, I] \subset I$. If I, J are ideals of R , then $[I, J]$ is an ideal as well. An ideal I is said to be *central* if $[R, I] = 0$, i.e. if it is contained in the *centre* $Z = \{r \in R \mid [r_\lambda s] = 0 \text{ for all } s \in R\}$ of R . R is *abelian* if it coincides with its centre, i.e. if $[R, R] = 0$.

A Lie conformal algebra R is *simple* if its only ideals are trivial, and R is not abelian. An interesting such example occurs when, for each choice of $0 \neq r \in R$, it occurs that $[r, R] = [Cr, R] = R$. In this case R is a *strongly simple* Lie conformal algebra.

Notice that, when V is a vertex algebra, we should distinguish between ideals of the vertex algebra structure and ideals of the underlying Lie conformal algebra. Indeed, ideals of the vertex algebra are also ideals of the Lie conformal algebra, but the converse is generally false, as it can be seen by noticing that $\mathbb{C}\mathbf{1}$ is always a central ideal of the Lie conformal algebra structure, but it is never an ideal of the vertex algebra.

In order to avoid confusion, we will denote by V^{Lie} the Lie conformal algebra structure underlying a vertex algebra V ; similarly, if $S \subset V$ is a $\mathbb{C}[\partial]$ -submodule closed under all nonnegative products (n) , $n \in \mathbb{N}$, we will denote by S^{Lie} the corresponding Lie conformal algebra structure. The reader should pay special attention to the fact that a vertex algebra V is commutative if and only if the Lie conformal algebra V^{Lie} is abelian, and that claiming that I is an abelian ideal of V is a stronger statement than saying that I is an abelian ideal in V^{Lie} . We will say that V is solvable (resp. nilpotent), whenever V^{Lie} is.

2.5. Finite simple Lie conformal algebras

Every Lie conformal algebra R has a maximal solvable ideal, called *radical* of R and denoted by $Rad R$. A Lie conformal algebra is called *semisimple* if it has no solvable ideal; the quotient $R/Rad R$ is always semisimple.

An investigation of Lie conformal algebra structures on finitely generated $\mathbb{C}[\partial]$ -modules was undertaken in [7], where a classification of simple and semisimple ones, together with generalizations of standard theorems in Lie representation theory, are presented.

It turns out that the only (up to isomorphism) simple Lie conformal algebra structures over finitely generated $\mathbb{C}[\partial]$ -modules are the Virasoro conformal algebra and current conformal algebras over a finite-dimensional simple Lie

algebra, which are described below. Semisimple instances are direct sums of Lie conformal algebras that are either simple or nontrivial semidirect sums of a Virasoro conformal algebra with a simple current one.

Example 2.3. Let R be a free $\mathbb{C}[\partial]$ -module of rank one, generated by an element L . Then

$$[L_\lambda L] = (\partial + 2\lambda)L \quad (2.6)$$

uniquely extends to a Lie conformal algebra structure on R , which is easily seen to be strongly simple. $R = \text{Vir}$ is called *Virasoro conformal algebra*.

Example 2.4. Let \mathfrak{g} be a finite-dimensional complex Lie algebra, and let $R = \mathbb{C}[\partial] \otimes \mathfrak{g}$. There exists a unique λ -bracket on R extending

$$[g_\lambda h] = [g, h], \quad (2.7)$$

for $g, h \in \mathfrak{g} \simeq 1 \otimes \mathfrak{g} \subset R$, and satisfying all axioms for a Lie conformal algebra. R is called *current conformal algebra* and is denoted by $\text{Cur } \mathfrak{g}$. It is a simple Lie conformal algebra whenever \mathfrak{g} is a simple Lie algebra. However, $\text{Cur } \mathfrak{g}$ is never strongly simple, as for no choice of $g \in \mathfrak{g}$ does $\text{ad } g \in \text{End } \mathfrak{g}$ satisfy surjectivity.

2.6. Centre and torsion

We now need a statement on subalgebras of a finite Lie conformal algebra R which have the same rank as R .

Definition 2.3. Let U, V be $\mathbb{C}[\partial]$ -modules. A *conformal linear map* from U to V is a \mathbb{C} -linear map $f_\lambda : U \rightarrow V[\lambda]$ such that $f_\lambda(\partial u) = (\partial + \lambda)f_\lambda u$ for all $u \in U$.

The space of all conformal linear maps from U to V is denoted by $\text{Chom}(U, V)$. It can be turned into a $\mathbb{C}[\partial]$ -module via

$$(\partial f)_\lambda u = -\lambda f_\lambda u.$$

Remark 2.3. Let U, V, W be $\mathbb{C}[\partial]$ -modules. A $\mathbb{C}[\partial]$ -linear homomorphism $\phi : V \rightarrow W$ induces a corresponding $\mathbb{C}[\partial]$ -linear homomorphism $\phi_* : V[\lambda] \rightarrow W[\lambda]$ mapping $v\lambda^n$ to $\phi(v)\lambda^n$. Then if $f \in \text{Chom}(U, V)$, the composition $\phi_* \circ f$ lies in $\text{Chom}(U, W)$.

Lemma 2.4. Let U, V be $\mathbb{C}[\partial]$ -modules, $f \in \text{Chom}(U, V)$. If $U_0 \subset U, V_0 \subset V$ are $\mathbb{C}[\partial]$ -submodules such that $f_\lambda(u_0) \in V_0[\lambda]$ for all $u_0 \in U_0$, then f induces a unique $\bar{f} \in \text{Chom}(U/U_0, V/V_0)$ satisfying $\bar{f}_\lambda[u] = [f_\lambda u]$.

Proof. Let $\pi : V \rightarrow V/V_0$ be the natural projection. Then $\pi_* \circ f$ is a conformal linear map from U to V/V_0 which kills all elements from U_0 . \square

The most typical example of a conformal linear map comes from the adjoint action in Lie conformal algebras. Indeed, if R is a Lie conformal algebra, and $r \in R$, then

$$(\text{ad } r)_\lambda x = [r_\lambda x]$$

defines a conformal linear map from R into itself, due to axiom (C2).

Lemma 2.5 ([7]). If $f \in \text{Chom}(U, V)$ and $u \in \text{Tor } U$, then $f_\lambda u = 0$.

Corollary 2.4. The torsion of a Lie conformal algebra is contained in its centre.

Proof. Let R be a Lie conformal algebra, $r \in R, t \in \text{Tor } R$. The adjoint action of r is a conformal linear map from R into itself, hence it maps the torsion element t to $[r_\lambda t] = 0$ by Lemma 2.5. \square

Lemma 2.6. Let $S \subset R$ be finite Lie conformal algebras, such that R/S is a torsion $\mathbb{C}[\partial]$ -module. Then S is an ideal of R containing R' , i.e. R/S is abelian.

Proof. Since S is a subalgebra of R , the adjoint action of S on R stabilizes the $\mathbb{C}[\partial]$ -submodule S . By Lemma 2.4, S acts on the quotient R/S , which is torsion. By Lemma 2.5, the action of S on R/S is trivial, or in other words $[S, R] \subset S$, which amounts to saying that S is an ideal of R .

Thus, the adjoint action of R on itself stabilizes S , and we may repeat the above argument to conclude that $R' = [R, R] \subset S$. It immediately follows that R/S is abelian. \square

2.7. Irreducible central extensions of the Virasoro conformal algebra

In this section we compute all finite irreducible central extensions of the Virasoro conformal algebra. Finite central extensions of Vir are described, up to equivalence, by cohomology classes [3] of $H^2(Vir, Z)$, where Z is the finitely generated $\mathbb{C}[\partial]$ -module describing the centre.²

The centre Z being a finitely generated $\mathbb{C}[\partial]$ -module, we can decompose it (non-canonically) into a direct sum of its torsion with a free $\mathbb{C}[\partial]$ -module. This leads to a corresponding direct sum decomposition of the related cohomology. In order to understand finite central extensions of Vir , it is thus sufficient to compute $H^2(Vir, Z)$ when Z is either a free $\mathbb{C}[\partial]$ -module of rank one, or an indecomposable torsion $\mathbb{C}[\partial]$ -module. The following facts were proved in [7] and [3] respectively:

Proposition 2.1. All central extensions of Vir by a free $\mathbb{C}[\partial]$ -module of rank one are trivial.

Proposition 2.2. Let $\mathbb{C}_\alpha, \alpha \in \mathbb{C}$, be the 1-dimensional $\mathbb{C}[\partial]$ -module on which the action of ∂ is given via scalar multiplication by α . Then:

- if $\alpha \neq 0$ all central extensions of Vir by \mathbb{C}_α are trivial;
- if $\alpha = 0$ then there is a unique (up to isomorphism and scalar multiplication) nontrivial central extension of Vir by $\mathbb{C} = \mathbb{C}_0$ given by

$$[L_\lambda L] = (\partial + 2\lambda)L + \lambda^3. \tag{2.8}$$

Remark 2.4. A computation of 2-cocycles of Vir with values in the trivial Vir -module \mathbb{C}_0 shows that they are of the form $p(\lambda) = c_1\lambda + c_3\lambda^3$, whereas trivial 2-cocycles (i.e. 2-coboundaries) are of the form $p(\lambda) = c_1\lambda$.

Recall that a central extension is called *irreducible* if it equals its derived algebra. Clearly, no nonzero trivial central extension is irreducible. My aim is to show that the non-trivial central extension (2.8) is the unique (nonzero) irreducible finite central extension of Vir .

Proposition 2.3. Let C be a finitely generated torsion $\mathbb{C}[\partial]$ -module on which ∂ acts invertibly. Then every central extension of Vir by C is trivial.

Proof. Let the central extension be given by

$$[L_\lambda L] = (\partial + 2\lambda)L + p(\lambda),$$

for some $p(\lambda) \in C[\lambda]$. By a computation similar to that in [7, Lemma 8.11], one obtains $\partial p(\lambda) = (\partial + 2\lambda)p(0)$, whence $p(\lambda) = (\partial + 2\lambda)\partial^{-1}p(0)$. Then $L + \partial^{-1}p(0)$ is a standard generator of a Virasoro conformal algebra, hence it splits the central extension. \square

Lemma 2.7. Solutions $p(\partial, x) \in \mathbb{C}[\partial, x]/(\partial^{N+1})$ of

$$(\lambda - \mu)p(\partial, \lambda + \mu) = (\partial + \lambda + 2\mu)p(\partial, \lambda) - (\partial + 2\lambda + \mu)p(\partial, \mu) \text{ mod } \partial^{N+1} \tag{2.9}$$

are all of the form $p(\partial, \lambda) = (\partial + 2\lambda)q(\partial) + c\lambda^3\partial^N \text{ mod } \partial^{N+1}, c \in \mathbb{C}$.

² It is clear that every finite central extension of Vir splits as an extension of $\mathbb{C}[\partial]$ -modules, as Vir is a free $\mathbb{C}[\partial]$ -module.

Proof. By induction on $N \geq 0$. The basis of induction follows from Remark 2.4. Next assume $N > 0$. Then (2.9) also holds modulo ∂^N , and inductive assumption gives $p(\partial, \lambda) = (\partial + 2\lambda)q(\partial) + c_0\lambda^3\partial^{N-1} \pmod{\partial^N}$. As a consequence:

$$p(\partial, \lambda) = (\partial + 2\lambda)q(\partial) + c_0\lambda^3\partial^{N-1} + \alpha(\lambda)\partial^N \pmod{\partial^{N+1}}.$$

We can substitute this into (2.9) and get

$$(\lambda - \mu)\alpha(\lambda + \mu) - (\lambda + 2\mu)\alpha(\lambda) + (2\lambda + \mu)\alpha(\mu) = c_0(\lambda^3 - \mu^3).$$

The left-hand side is linear in α and homogeneous with respect to the joint degree in λ and μ . Hence we can solve it degree by degree, looking for solutions of the form $\alpha(x) = ax^n$. It is then easy to check that solutions only exist when $c_0 = 0$, and are of the form $\alpha(x) = q_N\lambda + c\lambda^3$. We conclude that

$$\begin{aligned} p(\partial, \lambda) &= (\partial + 2\lambda)q(\partial) + (q_N\lambda + c\lambda^3)\partial^N \\ &= (\partial + 2\lambda) \left(q(\partial) + \frac{q_N}{2}\partial^N \right) + c\lambda^3\partial^N \pmod{\partial^{N+1}}. \quad \square \end{aligned}$$

Proposition 2.4. Let C_N be a finitely generated torsion $\mathbb{C}[\partial]$ -module isomorphic to $\mathbb{C}[\partial]/(\partial^{N+1})$, $N \geq 1$. Then there is a unique (up to isomorphism and scalar multiplication) nontrivial central extension of Vir by C_N given by $[L_\lambda L] = (\partial + 2\lambda)L + \lambda^3\partial^N$.

Proof. The 2-cocycle property for $p(\partial, \lambda)$ as in

$$[L_\lambda L] = (\partial + 2\lambda)L + p(\partial, \lambda),$$

leads to solving (2.9); hence Lemma 2.7 gives $p(\partial, \lambda) = (\partial + 2\lambda)q(\partial) + c\lambda^3\partial^N$. However, a 2-cocycle is trivial if and only if it is of the form $(\partial + 2\lambda)q(\partial)$, whence the claim. \square

Remark 2.5. All of the above nontrivial central extensions of Vir are equivalent to one of the form $[L_\lambda L] = (\partial + 2\lambda)L + \lambda^3c$, where $\partial c = 0$.

Theorem 2.1. A finite non-zero irreducible central extension of Vir is isomorphic to that given in (2.8).

Proof. We already know that a non-trivial central extension of Vir by the $\mathbb{C}[\partial]$ -module C is only possible if C contains nonzero torsion elements. A torsion finitely generated $\mathbb{C}[\partial]$ -module is a finite-dimensional vector space, on which ∂ acts as a \mathbb{C} -linear endomorphism. Then C decomposes into a direct sum of a submodule on which ∂ acts invertibly, and of summands as in Proposition 2.4.

Irreducibility and Proposition 2.3 prove that the summand on which ∂ acts invertibly is trivial. On the other hand, Proposition 2.4 shows that we may choose a lifting L of the standard Virasoro generator so that

$$[L_\lambda L] = (\partial + 2\lambda)L + \lambda^3c,$$

for some $c \in C$ such that $\partial c = 0$. Using irreducibility and Proposition 2.1 gives $C = \mathbb{C}c$. \square

3. Finite vertex algebras

3.1. Commutativity of finite simple vertex algebras

We have seen that coefficients of vertex operators in a vertex algebra

$$Y(a, z) = \sum_{j \in \mathbb{Z}} a_{(j)} z^{-j-1} \tag{3.1}$$

satisfy the Lie bracket (2.2):

$$[a_{(m)}, b_{(n)}] = \sum_{j \in \mathbb{N}} \binom{m}{j} (a_{(j)}b)_{(m+n-j)}$$

for every $a, b \in V, m, n \in \mathbb{Z}$. Multiplying both sides of (2.2) by $\lambda^m z^{-n-1}/m!$ and adding up over all $m \in \mathbb{N}, n \in \mathbb{Z}$, after applying both sides to $c \in V$, gives

$$[a_\lambda Y(b, z)c] = e^{\lambda z} Y([a_\lambda b], z)c + Y(b, z)[a_\lambda c], \tag{3.2}$$

for all $a, b, c \in V$.

This allows one to explicitly write down the λ -bracket of a vertex operator with the normally ordered product of two others – it suffices to take the constant term in z in both sides – but the formula is definitely more useful in the above form. Eq. (3.2) can be used in order to prove the following statements (see [6]):

Lemma 3.1. *Let V be a vertex algebra, $U \subset V$ a subspace. Then $[U, V]$ is a (vertex) ideal of V .*

Remark 3.1. It is important to realize that, by Lemma 3.1, elements of the descending sequence (2.5) are indeed ideals of the vertex algebra V and not only of the Lie conformal algebra V^{Lie} .

The lemma above has the following immediate and striking consequence:

Theorem 3.1. *Let V be a noncommutative simple vertex algebra. Then V^{Lie} is an irreducible central extension of a strongly simple Lie conformal algebra.*

Remark 3.2. The strong simplicity property is clearly expressed in the proof but not explicitly stated in [6].

We will use Theorem 3.1, and our knowledge of finite simple Lie conformal algebras, in order to show that all simple vertex algebra structures over finitely generated $\mathbb{C}[\partial]$ -modules are commutative.

Proposition 3.1. *There is no vertex algebra V such that $V = \mathbb{C}[\partial]L + \mathbb{C}\mathbf{1}$, where $L \notin \text{Tor } V$ and $[L_\lambda L] = (\partial + 2\lambda)L + c\lambda^3\mathbf{1}$, for some $c \in \mathbb{C}$.*

Proof. We proceed by contradiction. We know that

$$[L_\lambda L] = (\partial + 2\lambda)L + c\lambda^3\mathbf{1}, \quad [L_\lambda \mathbf{1}] = 0, \quad [\mathbf{1}_\lambda \mathbf{1}] = 0,$$

and that

$$Y(\mathbf{1}, z)\mathbf{1} = \mathbf{1}, \quad Y(\mathbf{1}, z)L = L, \quad Y(L, z)\mathbf{1} = e^{z\partial}L.$$

All that we need to determine is $Y(L, z)L$. Let us write

$$Y(L, z)L = a(\partial, z)L + b(z)\mathbf{1}. \tag{3.3}$$

Then (3.2) gives

$$[L_\lambda Y(L, z)L] = e^{\lambda z} Y([L_\lambda L], z)L + Y(L, z)[L_\lambda L], \tag{3.4}$$

which, after expanding and comparing coefficients of L , yields

$$(e^{\lambda z} - 1) \frac{da(\partial, z)}{dz} + (2\lambda(e^{\lambda z} + 1) + \partial)a(\partial, z) = -c\lambda^3(e^{\lambda z} + e^{z\partial}) + (\partial + 2\lambda)a(\partial + \lambda, z). \tag{3.5}$$

Using (3.3) and substituting $\lambda = -\partial/2$, this becomes

$$(e^{-z\partial/2} - 1) \frac{da(\partial, z)}{dz} - \partial e^{-z\partial/2} a(\partial, z) = \frac{c\partial^3}{8}(e^{z\partial} + e^{-z\partial/2}), \tag{3.6}$$

i.e., a linear differential equation in $a(\partial, z)$ whose solutions are of the form

$$a(\partial, z) = \frac{K(\partial)e^{z\partial}}{(e^{z\partial/2} - 1)^2} - \frac{c\partial^2}{8}(1 + e^{z\partial}).$$

In order for $a(\partial, z)$ to be compatible with $[L_\lambda L] = (\partial + 2\lambda)L$, one needs

$$a(\partial, z) = 2/z^2 + \partial/z + (\text{regular in } z).$$

This forces $K(\partial) = \partial^2/2$, hence the only solution of (3.6) satisfying this additional condition is

$$a(\partial, z) = \frac{\partial^2 e^{z\partial}}{2(e^{z\partial/2} - 1)^2} - \frac{c\partial^2}{8}(1 + e^{z\partial}). \quad (3.7)$$

Checking that this value of $a(\partial, z)$ is not a solution of (3.5) is a rather lengthy but straightforward computation. One may also observe that substituting (3.7) into the right-hand side of (3.5) gives a denominator of the form $(e^{z(\partial+\lambda)/2} - 1)^2$ which cannot be obtained from the left-hand side. \square

Theorem 3.2. *Every simple finite vertex algebra is commutative*

Proof. Let V be a finite simple vertex algebra. By Theorem 3.1, either V is commutative or V^{Lie} is an irreducible central extension of a strongly simple Lie conformal algebra. It is then enough to address the latter case, showing it leads to a contradiction.

We have seen that every finite strongly simple Lie conformal algebra is isomorphic to Vir . Moreover, Theorem 2.1 gives a description of all finite nonzero irreducible central extensions of Vir . Thus we know that $V = \mathbb{C}[\partial]L + \mathbb{C}\mathbf{1}$, with $[L_\lambda L] = (\partial + 2\lambda)L + c\lambda^3\mathbf{1}$, for some $0 \neq c \in \mathbb{C}$. Then Proposition 3.1 leads to a contradiction. \square

The following claim is a technical statement that we will use later on:

Lemma 3.2. *Let V be a vertex algebra, and $M \subset V$ a minimal ideal such that $M = \mathbb{C}[\partial]L + \mathbb{C}k$, where L is a nontorsion element, $\partial k = 0$ and*

$$[L_\lambda L] = (\partial + 2\lambda)L + \lambda^3 k.$$

Then M can be endowed with a vertex algebra structure by choosing the vacuum element $\mathbf{1}_M$ to be a suitable scalar multiple of k .

Proof. The vacuum element $\mathbf{1}$ of V lies outside of M , so the only thing we need to prove is that we may choose an element inside M whose quantum field acts as the identity on M . As $\partial k = 0$, then $Y(k, z)$ does not depend on z . Moreover $Y(k, z)k \in M$ is a torsion element as

$$\partial(Y(k, z)k) = Y(\partial k, z)k + Y(k, z)(\partial k) = 0.$$

Then $Y(k, z)k = \alpha k$ for some $\alpha \in \mathbb{C}$, hence $Y(k - \alpha\mathbf{1}, z)k = 0$. The element $c = k - \alpha\mathbf{1}$ satisfies $\partial c = 0$, hence is a torsion element, contained in the centre of V . By Lemma 2.1, $\ker c_{(-1)}$ is an ideal of V containing k , therefore it must contain all of M by the minimality assumption. Thus, $Y(k - \alpha\mathbf{1}, z)$ has zero restriction on all of M , and $Y(k, z)|_M = \alpha \text{id}_M$. If $\alpha \neq 0$, then we are done by setting $\mathbf{1}_M = \alpha^{-1}k$.

The case $\alpha = 0$ can be ruled out as follows: we know that $k_{(-1)}V \subset M$ as M is an ideal containing k . Moreover, we just showed that $k_{(-1)}M = 0$. Now, by (2.2), for every choice of $a, b \in V$, one has:

$$(k_{(-1)}a)_{(m)}(k_{(-1)}b) = k_{(-1)}((k_{(-1)}a)_{(m)}b) + \sum_{j \geq 0} \binom{m}{j} ((k_{(-1)}a)_{(j)}k)_{(m+j-1)}b, \quad (3.8)$$

where both summands on the right hand side vanish. This shows that $k_{(-1)}V$ is a subalgebra contained in M in which all products vanish, therefore $k_{(-1)}V \subset \mathbb{C}k$. In other words, $\mathbb{C}k$ is an ideal of V , which contradicts the minimality of M . \square

3.2. Solvability of finite vertex algebras

Lemma 3.3. *Let V be a finite vertex algebra, and S be the intersection of all vertex subalgebras $U \subset V$ such that $\text{rk } U = \text{rk } V$. Then V is solvable if and only if S is.*

Proof. Let U be a vertex subalgebra of V such that $\text{rk } U = \text{rk } V$. U is clearly a subalgebra of V^{Lie} ; hence, by Lemma 2.6, an ideal containing the derived subalgebra of V^{Lie} .

The intersection S of all such vertex subalgebras is then itself an ideal of V^{Lie} containing its derived subalgebra, hence V^{Lie}/S is abelian. Therefore, V is solvable if and only if S is. \square

Remark 3.3. Observe that if the vertex subalgebra S in the above lemma is such that $\text{rk } S = \text{rk } V$, then it is the minimal vertex subalgebra of V of rank equal to $\text{rk } V$. In particular, S possesses no proper vertex subalgebras of equal rank.

Lemma 3.4. *Let V be a finite vertex algebra, and N be the sum of all vertex ideals of V contained in $\text{Tor } V$. Then V/N has no nonzero torsion ideal and V is solvable if and only if V/N is. Moreover, V contains proper vertex subalgebras of rank $\text{rk } V$ if and only if V/N does.*

Proof. The sum of ideals in a vertex algebra is again an ideal. Also, the sum of torsion elements lies in $\text{Tor } V$. Hence, the sum N of all vertex ideals of V contained in $\text{Tor } V$ is the maximal such ideal of V . As $\text{Tor } V$ lies in the centre of V^{Lie} , N^{Lie} is abelian, so V is solvable if and only if V/N is.

The other claims follow from the correspondence between ideals (resp. subalgebras) of V/N and ideals (resp. subalgebras) of V containing N , and the fact that torsion modules are of zero rank. \square

Theorem 3.3. *Every finite vertex algebra is solvable.*

Proof. Assume by way of contradiction that V is a counter-example of minimal rank. By Lemmas 3.3 and 3.4 and Remark 3.3, we may assume that V has no proper vertex subalgebra of equal rank, and no nonzero torsion ideal.

Now, observe that if $I \subset V$ is a nonzero ideal with $\text{rk } I < \text{rk } V$, then $\text{rk } V/I < \text{rk } V$ as I cannot lie in $\text{Tor } V$. Then the vertex algebras V/I and $I + \mathbb{C}\mathbf{1}$ are both solvable by the minimality assumption, hence $I^{\text{Lie}} \subset (I + \mathbb{C}\mathbf{1})^{\text{Lie}}$ is solvable and V^{Lie} is an extension of solvable Lie conformal algebras, a contradiction. Therefore, all nonzero vertex ideals of V are of the same rank as V .

As a consequence, either V is simple, or has a unique nonzero proper ideal M which is a complement to $\mathbb{C}\mathbf{1}$. Indeed, if M is a nonzero ideal, then $\text{rk } M = \text{rk } V$, and $M + \mathbb{C}\mathbf{1}$ is a vertex subalgebra of V , hence $M + \mathbb{C}\mathbf{1} = V$.

We already know that finite simple vertex algebras are commutative, hence solvable, so it is enough to address the non-simple case. Let V be a nonsolvable finite vertex algebra whose only nonzero vertex ideal $M \neq V$ is such that $V = M + \mathbb{C}\mathbf{1}$. Then $[V, V] = [M, M]$ is a vertex ideal of V , hence it equals M . If $U \subset M$ is a subspace, then Lemma 3.1 shows that $[U, V]$ is a (vertex) ideal of V , hence it equals either 0 or M . As a consequence, if $u \in M$ then either u is central in V or $[u, M] = [u, V] = [\mathbb{C}u, V] = M$. This shows that M^{Lie} is either strongly simple or a central extension of a strongly simple conformal algebra. As $[M, M] = M$, the central extension must be irreducible.

If M^{Lie} is strongly simple, then we conclude that V is as in Proposition 3.1, with $c = 0$, whence a contradiction. If, on the other hand, M^{Lie} is a central extension of a strongly simple Lie conformal algebra, then Lemma 3.2 shows that M can be given a vertex algebra structure which contradicts Proposition 3.1. \square

4. Conformal adjoint decomposition

4.1. Finite modules over finite solvable Lie conformal algebras

In this paper, we will need some basic results from representation theory of solvable and nilpotent Lie conformal algebras. A representation of a Lie conformal algebra R is a $\mathbb{C}[\partial]$ -module V along with a λ -action $R \otimes V \ni r \otimes v \mapsto r_\lambda v \in V[\lambda]$ such that

$$(\partial r)_\lambda v = -\lambda r_\lambda v, \quad r_\lambda(\partial v) = (\partial + \lambda)r_\lambda v, \tag{4.1}$$

$$r_\lambda(s_\mu v) - s_\mu(r_\lambda v) = [r_\lambda s]_{\lambda+\mu} v, \tag{4.2}$$

for all $r, s \in R, v \in V$. The action of $r \in R$ on V is nilpotent if

$$r_{\lambda_1}(r_{\lambda_2}(\dots(r_{\lambda_n} v) \dots)) = 0$$

for all $v \in V$ and sufficiently large n . The following conformal versions of Engel's and Lie's Theorems were proved in [7].

Theorem 4.1. *Let R be a finite Lie conformal algebra for which every element $r \in R$ has a nilpotent adjoint action. Then R is a nilpotent Lie conformal algebra.*

Theorem 4.2. Let R be a finite Lie solvable conformal algebra, V its finite module. Then there exists $0 \neq v \in V$ and $\phi : R \ni r \rightarrow \phi_r(\lambda) \in \mathbb{C}[\lambda]$ such that

$$r_\lambda v = \phi_r(\lambda)v, \quad (4.3)$$

for all $r \in R$.

An element v such as that in Theorem 4.2 is a *weight vector*. Then ϕ is the *weight* of v , and it necessarily satisfies $\phi_{\partial r}(\lambda) = -\lambda\phi_r(\lambda)$. The set of all weight vectors of a given weight ϕ , along with zero, is the *weight subspace* V_ϕ .

Remark 4.1. The statement in [7] only deals with R -modules that are free as $\mathbb{C}[\partial]$ -modules, but clearly extends to nonfree modules, since $\text{Tor } V$ is a submodule of V which is killed by R .

Lemma 4.1. Let V be a representation of the Lie conformal algebra R . Then V_ϕ is always a vector subspace of V . Also, it is a $\mathbb{C}[\partial]$ -submodule whenever $\phi \equiv 0$.

Now set:

$$V_0^\phi = 0, \quad V_{i+1}^\phi = \{v \in V \mid r_\lambda v - \phi(r)v \in V_i^\phi \text{ for all } r \in R\}, \quad i \geq 0.$$

Then $V_1^\phi = V_\phi$, and $V_1^\phi \subset V_2^\phi \subset \dots$ is an ascending chain of subspaces of V . The subspace $\bigcup V_i^\phi = V^\phi$ is the *generalized weight subspace* of weight ϕ . Clearly, r acts nilpotently on V exactly when V coincides with the generalized 0-weight space for the action of (the Lie conformal algebra generated by) r .

Proposition 4.1 ([1,7]). Let V be a representation of the Lie conformal algebra R . Then:

- V^ϕ is a $\mathbb{C}[\partial]$ -submodule of V ;
- V/V^0 has no 0-weight vectors: in particular, it is torsion-free;
- if V is torsion-free, then V/V^ϕ is too;
- if $\phi \neq \psi$, then $V^\phi \cap V^\psi = 0$;
- the sum of all generalized weight spaces for the action of R on V is direct.

The sum of all generalized weight spaces may fail to coincide with the R -module V . However, the following generalized-weight-space decomposition, proved in the context of Lie pseudoalgebras [1], holds for nilpotent Lie conformal algebras.

Theorem 4.3. Let R be a finite nilpotent Lie conformal algebra, V its finite module. Then V decomposes into a direct sum of generalized weight subspaces for the action of R .

In practice, we will often consider weight spaces and generalized weight spaces with respect to the action of a single element $s \in R$. If S is the subalgebra generated by s , we will say a weight for the action of S on some module V is a *weight* of s . This abuse of notation is justified by the fact that in the case $S = \langle s \rangle = \mathbb{C}[\partial]s + S'$, any weight ϕ for the action of S on some module V satisfies $\phi(S') = 0$.

4.2. Matrix form

Let R be a Lie conformal algebra, and V be an R -module. Then the map $V \ni v \rightarrow r_\lambda v \in V[\lambda]$ is conformal linear for all $r \in R$. The $\mathbb{C}[\partial]$ -module structure built on $\text{Chom}(V, V)$ is such that the map $r \mapsto \{v \mapsto r_\lambda v\}$ is $\mathbb{C}[\partial]$ -linear.

One may indeed build up a Lie conformal algebra structure on $\text{Chom}(V, V)$ in such a way that the above map is always a homomorphism of Lie conformal algebras. It suffices to define:

$$[f_\lambda g]_\mu v = f_\lambda(g_{\mu-\lambda}v) - g_{\mu-\lambda}(f_\lambda v), \quad (4.4)$$

whenever $f, g \in \text{Chom}(V, V)$. This Lie conformal algebra structure is usually denoted by $gc(V)$, or simply gc_n when V is a free $\mathbb{C}[\partial]$ -module of rank n . The standard way to represent elements of gc_1 is by identifying it with $\mathbb{C}[\partial, x]$, with the $\mathbb{C}[\partial]$ -module structure given via multiplication by ∂ , and the conformal linear action on $\mathbb{C}[\partial]$ given on its free generator 1 by:

$$x_\lambda^n 1 = (\partial + \lambda)^n.$$

Then the λ -bracket $[p(\partial, x)_\lambda q(\partial, x)]$ equals

$$p(-\lambda, x + \partial + \lambda)q(\partial + \lambda, x) - q(\partial + \lambda, x - \lambda)p(-\lambda, x). \tag{4.5}$$

However, in this paper I will employ a different choice, and denote elements of gc_1 by the effect they have on the free generator. This identifies gc_1 with $\mathbb{C}[\partial, \lambda]$, and has two major inefficiencies: first of all, the $\mathbb{C}[\partial]$ -module structure is obtained via multiplication by $-\lambda$; moreover, as we are already employing λ to denote elements, we will have to compute α - rather than λ -bracket of elements. However, this choice is by far more readable than the standard one. The bracket expressed in (4.5) then becomes:

$$[a(\partial, \lambda)_\alpha b(\partial, \lambda)] = a(\partial, \alpha)b(\partial + \alpha, \lambda - \alpha) - b(\partial, \lambda - \alpha)a(\partial + \lambda - \alpha, \alpha). \tag{4.6}$$

Now let V, W be $\mathbb{C}[\partial]$ -modules. A conformal linear map $f \in \text{Chom}(V, W)$ is determined by its values on a set of $\mathbb{C}[\partial]$ -generators of V . If V and W are free, then, for any given choice of $\mathbb{C}[\partial]$ -bases $(v^1, \dots, v^m), (w^1, \dots, w^n)$ of V and W respectively, we can establish a correspondence between $\text{Chom}(V, W)$ and $n \times m$ matrices with coefficients in $\mathbb{C}[\partial, \lambda]$, similarly to what done above in the case of gc_1 .

In general, $\mathbb{C}[\partial]$ -modules fail to be free. If M is a finitely generated $\mathbb{C}[\partial]$ -module, M can be (noncanonically) decomposed as the direct sum of a free module and of its torsion submodule $\text{Tor } M$. By Lemma 2.5, $f \in \text{Chom}(M, N)$ always maps $\text{Tor } M$ to zero.

Since we need to employ a matrix representation for any conformal linear map $f \in \text{Chom}(M, N)$ between finitely generated modules that may (and typically will) fail to be free, we can proceed as follows. Decompose M and N as a direct sum of a free module and their torsion submodule. If we pick a free $\mathbb{C}[\partial]$ -basis of the free part, and a \mathbb{C} -basis of the torsion part, we can use this set of generators to represent conformal linear maps through matrices: we will call such a set of generators a *base*. As a conformal linear map in $\text{Chom}(M, N)$ always factors via $M/\text{Tor } M$, which is free, special care is only needed for the treatment of torsion in the range module.

Note that if we agree that the $\mathbb{C}[\partial]$ -linear combination expressing elements of N in terms of a given base is such that coefficients multiplying torsion elements lie in \mathbb{C} (rather than in $\mathbb{C}[\partial]$) then all coefficients are uniquely determined. This unique expression enables us to write down a well-behaved matrix representing the conformal linear map. Matrix coefficients corresponding to torsion elements then lie in $\mathbb{C}[\lambda]$ rather than in $\mathbb{C}[\partial, \lambda]$.

If $f, g \in gc(M, M)$ and the matrices representing them are given by

$$F = (f_{ij}(\partial, \lambda)), G = (g_{ij}(\partial, \lambda)),$$

respectively, then by (4.4), the matrix representing $[f_\alpha g]$ is given by

$$F(\partial, \alpha)G(\partial + \alpha, \lambda - \alpha) - G(\partial, \lambda - \alpha)F(\partial + \lambda - \alpha, \alpha), \tag{4.7}$$

where multiplication of matrices is the usual row-by-column product. Notice that, according to such a matrix representation of conformal linear maps, Theorem 4.2 guarantees the existence of a base in which matrices representing the action of the solvable Lie conformal algebra R are simultaneously upper triangular. Similarly, Theorem 4.3 means that matrices can be put in block diagonal form, where each block represents the action on a single generalized weight submodule.

Later, we will call the diagonal entries of a triangular matrix representing the action of some $s \in S$, S solvable, *eigenvalues* of the element s .

4.3. Adjoint action on a vertex algebra of a Lie conformal subalgebra

In what follows V will be a finite vertex algebra, unless otherwise stated. If S is a Lie conformal subalgebra of V , then S is solvable by Theorem 3.3. Using formula (3.2) I want to show that

Theorem 4.4. *If ψ is a nonzero weight for the adjoint action of a Lie conformal subalgebra S on the finite vertex algebra V , then the generalized weight space V^ψ is a vertex ideal of V , and it satisfies $V^\psi \cdot V^\psi = 0$.*

I will divide the proof of Theorem 4.4 in a few easy steps. Let α be a weight for the action of S on V , β for its action on V/V^α . Denote by $V^{\alpha, \beta}$ the preimage of $(V/V^\alpha)^\beta$ via the canonical projection $\pi : V \rightarrow V/V^\alpha$. Then we have:

Lemma 4.2. Let $U \subset V$ be a proper S -submodule of V with the property that $U \cdot V^\psi \subset V^\psi$, and choose an element $w \in V$ such that $\bar{w} = [w] \in V/U$ is a weight vector of weight ϕ . Then $w \cdot V^\psi \subset V^{\psi, \psi+\phi}$.

Proof. As $[s_\lambda w] = \phi_s(\lambda)w \pmod U$, then we have $s_{(h)}w = \phi_s^h w + u_s^h$, for some $u_s^h \in U$, where the ϕ_s^h are such that

$$\phi_s(\lambda) = \sum_h \phi_s^h \frac{\lambda^h}{h!}.$$

I will prove that

$$w \cdot V_n^\psi \subset V^{\psi, \psi+\phi}$$

by induction on n — the basis of induction $n = 0$ being trivial, as $V_0^\psi = 0$.

Let $b \in V_{n+1}^\psi$, and set $s_{(h)}b = \psi_s^h b + v_s^h$ with $v_s^h \in V_n^\psi$. We know that $w_{(N)}b = 0$ for sufficiently large N . So if $Y(w, z)b \notin V^{\psi, \psi+\phi}[[z, z^{-1}]]$ we choose k maximal with respect to the property that $w_{(k)}b \notin V^{\psi, \psi+\phi}$. Let us compute by means of (2.2):

$$\begin{aligned} s_{(m)}(w_{(k)}b) - w_{(k)}(s_{(m)}b) &= \sum_{j=0}^m \binom{m}{j} (s_{(j)}w)_{(m+k-j)}b \\ &= (s_{(m)}w)_{(k)}b + \sum_{j=0}^{m-1} \binom{m}{j} (s_{(j)}w)_{(m+k-j)}b, \end{aligned} \tag{4.8}$$

hence

$$s_{(m)}(w_{(k)}b) - (\psi_s^m + \phi_s^m)w_{(k)}b = w_{(k)}v_s^m + (u_s^m)_{(k)}b + \sum_{j=0}^{m-1} \binom{m}{j} (\phi_s^j (w_{(m+k-j)}b) + (u_s^j)_{(m+k-j)}b). \tag{4.9}$$

Now, $v_s^h \in V_n^\psi$, so $Y(w, z)v_s^h \in V^{\psi, \psi+\phi}[[z, z^{-1}]]$. Also, $u_s^h \in U$, hence $Y(u_s^h, z)b \in V^\psi[[z, z^{-1}]]$. Moreover, each $w_{(m+k-j)}b$ in the summation lies in $V^{\psi, \psi+\phi}$ by the maximality of k . Therefore, $(s_{(m)} - (\psi + \phi)_s^m)(w_{(k)}b) \in V^{\psi, \psi+\phi}$, showing $w_{(k)}b \in V^{\psi, \psi+\phi}$, a contradiction. \square

Lemma 4.3. Under the same hypotheses as in Lemma 4.2, $w \cdot V^\psi \subset V^\psi$.

Proof. The statement is clear if $\phi = 0$, as $V^{\psi, \psi} = V^\psi$. Otherwise, choose a base $\{r^i\}$ of $V^{\psi, \psi+\phi}/V^\psi$ on which the action of S is triangular, and lift it to $V^{\psi, \psi+\phi}$. Then, if $b \in V^\psi$, we can express $Y(w, z)b$ as some (depending on z) element from V^ψ plus a $\mathbb{C}[\partial]((z))$ -linear combination of elements from this base:

$$Y(w, z)b = v(z) + \sum_i A^i(\partial, z)r^i. \tag{4.10}$$

My aim is to show that all A^i are zero. I will prove that it is so for $b \in V_k^\psi$ by induction on k . If not all of the A^i are zero, choose N maximal such that A^N is nonzero. Then (3.2) gives

$$[s_\lambda Y(w, z)b] = e^{\lambda z} Y([s_\lambda w], z)b + Y(w, z)[s_\lambda b],$$

and using triangularity of the action of s on the chosen base, along with the induction assumption, shows that

$$(\phi_s(\lambda) + \psi_s(\lambda))A^N(\partial + \lambda, z) = (e^{\lambda z}\phi_s(\lambda) + \psi_s(\lambda))A^N(\partial, z), \tag{4.11}$$

as $[s_\lambda b] - \psi_s(\lambda)b$ lies inside V_{k-1}^ψ .

Now, since neither ϕ nor ψ is identically zero, there must be some s such that ϕ_s and ψ_s are both non-zero. If for such an s we get $\phi_s + \psi_s = 0$, then A^N must be zero, giving a contradiction. If instead $\phi_s + \psi_s \neq 0$, then

$$\Gamma(\lambda, z) = \frac{e^{\lambda z}\phi_s(\lambda) + \psi_s(\lambda)}{\phi_s(\lambda) + \psi_s(\lambda)} \tag{4.12}$$

is a non-zero element of $\mathbb{C}(\lambda)[[z]]$ satisfying

$$\Gamma(\lambda + \mu, z) = \Gamma(\lambda, z)\Gamma(\mu, z). \tag{4.13}$$

It is then easy to show that Γ 's constant term as a power series in z must be one. $\Gamma(\lambda, z)$ is indeed of the form $e^{\lambda\gamma(z)}$ for some power series $\gamma(z) = \gamma_1 z + \gamma_2 z^2 + \dots$

By comparing coefficients of z and z^2 in (4.13) one concludes that $\phi_s(\lambda)/(\phi_s(\lambda) + \psi_s(\lambda)) = 0$ or 1. But this is only possible if either ϕ_s or ψ_s is zero, contrary to the assumption that they are both nonzero. We obtain a contradiction, which proves that all A^i vanish. \square

Lemma 4.4. V^ψ is an ideal of the vertex algebra V .

Proof. Let U be maximal among all S -submodules of V such that $U \cdot V^\psi \subset V^\psi$. If $U \neq V$, choose a weight vector w in V/U . Then $(U + \mathbb{C}[\partial]w) \cdot V^\psi \subset V^\psi$ by Lemma 4.3, against the maximality of U . Hence, U must equal V , and V^ψ is an ideal. \square

Lemma 4.5. Let V be a (not necessarily finite) vertex algebra, V^ϕ and V^ψ generalized weight subspaces for the adjoint action of the conformal subalgebra S of V . Then $V^\phi \cdot V^\psi \subset V^{\phi+\psi}$.

Proof. I will show that $V_i^\phi \cdot V_j^\psi \subset V_j^{\phi+\psi}$ by induction on $n = i + j$.

Say $v \in V_i^\phi, w \in V_j^\psi, i + j = n + 1$. Set $s_{(h)}v = \phi_s^h v + v_s^h, s_{(h)}w = \psi_s^h w + w_s^h$. Then $v_s^h \in V_{i-1}^\phi, w_s^h \in V_{j-1}^\psi$. If $Y(v, z)w \notin V^{\phi+\psi}[[z, z^{-1}]]$, then choose a maximal k with the property that $v_{(k)}w \notin V^{\phi+\psi}$. Then

$$s_{(m)}(v_{(k)}w) = v_{(k)}(s_{(m)}w) + (s_{(m)}v)_{(k)}w + \sum_{j=0}^{m-1} \binom{m}{j} (s_{(j)}v)_{(m+k-j)}w,$$

whence

$$s_{(m)}(v_{(k)}w) - (\phi_s^m + \psi_s^m)(v_{(k)}w) = v_{(k)}w_s^m + (v_s^m)_{(k)}w + \sum_{j=0}^{m-1} \binom{m}{j} (\phi_s^j(v_{(m+n-j)}w) + (\psi_s^j)_{(m+n-j)}w). \tag{4.14}$$

The right-hand side of (4.14) lies in $V^{\phi+\psi}$, hence $v_{(k)}w$ does too, giving a contradiction. \square

Proof of Theorem 4.4. V^ψ is an ideal by Lemma 4.4. On the other hand, Lemma 4.5 shows $V^\psi \cdot V^\psi \subset V^{2\psi}$. As ψ is a nonzero weight, Proposition 4.1 gives $V^\psi \cap V^{2\psi} = 0$, hence $V^\psi \cdot V^\psi = 0$. \square

Recall that an ideal I of a vertex algebra V is *abelian* if $I \cdot I = 0$, and that V is *reduced* if it has no abelian ideals or equivalently if its nilradical is trivial.

Corollary 4.1. Let V be a finite reduced vertex algebra, N be a nilpotent Lie conformal subalgebra of V . Then the adjoint action of N on V is achieved via nilpotent conformal linear maps.

Proof. By Theorem 4.4, any nonzero weight ϕ for the adjoint action of N on V would give an abelian vertex ideal V^ϕ . Since this is not possible, the only weight is 0. But N is a nilpotent Lie conformal algebra so, by Theorem 4.3, every N -module decomposes as a direct sum of generalized weight spaces, showing $V = V^0$. This means that the action of N on V is nilpotent. \square

Corollary 4.1 has the following immediate consequence:

Corollary 4.2. If V is a finite reduced vertex algebra, then $\text{Cur } \mathfrak{g}$ can arise as a subalgebra of V^{Lie} only when \mathfrak{g} is a nilpotent Lie algebra.

Proof. Every element $g \in 1 \otimes \mathfrak{g} \subset \text{Cur } \mathfrak{g}$ spans an abelian (hence nilpotent) Lie conformal subalgebra of V . By Corollary 4.1, g must act nilpotently on all of V , and in particular on $\text{Cur } \mathfrak{g}$ itself. Then \mathfrak{g} is a finite-dimensional Lie algebra on which every element is ad-nilpotent, and \mathfrak{g} is nilpotent by the usual Engel theorem for Lie algebras. \square

Remark 4.2. We observed in Theorem 3.3 that every finite vertex algebra is solvable, hence we knew already that $\text{Cur } \mathfrak{g}$ arises as a subalgebra of V^{Lie} only when \mathfrak{g} is solvable.

5. Nilpotence of finite reduced vertex algebras

The main result of this section is the following:

Theorem 5.1. *Any finite reduced vertex algebra is nilpotent.*

Before we prove this, we show a stronger result characterizing the conformal adjoint action of elements from V .

Lemma 5.1. *Let V be a finite vertex algebra, $s \in V$. If the (conformal) adjoint action of s is not nilpotent on V , then there exists a nonzero \bar{s} whose adjoint action has a weight vector w of nonzero weight.*

Proof. In what follows, by “action of s ”, I will always mean the conformal adjoint action of $s \in V^{\text{Lie}}$ on V . Notice that the finite vertex algebra V is solvable, hence all subalgebras of V^{Lie} are solvable Lie conformal algebras, for which the adjoint action on V satisfies the conditions of Theorem 4.2. In particular, the subalgebra $\langle s \rangle \subset V^{\text{Lie}}$ generated by s is solvable, and s acts triangularly in a suitably chosen base of V .

If s has a weight vector of nonzero weight then the statement holds with $\bar{s} = s$. We can thus assume, without loss of generality that the only weight of the adjoint action of s on V is zero. Since s does not act nilpotently, V^0 cannot equal the whole V and Proposition 4.1 shows that the action of s on V/V^0 only has nonzero weights.

Therefore, let us consider a weight vector \bar{w} in V/V^0 of non-zero weight $\phi(\lambda)$. Without loss of generality, we can assume that the degree of ϕ in λ be odd. In fact, we can always replace s by ∂s , which must be nonzero, otherwise s would be a torsion element, with a trivial adjoint action. Eigenvalues of ∂s are then obtained by multiplying those of s by $-\lambda$, and either $\phi(\lambda)$ or $-\lambda\phi(\lambda)$ is of odd degree. Moreover, the adjoint action of s is nilpotent if and only if that of ∂s is.

So, let s be an element of V for which there exists an element $\bar{w} \in V/V^0$ of nonzero weight $\phi(\lambda)$ having odd degree n in λ . My plan is to find an element $\bar{s} = s + s'$, $s' \in \langle s \rangle'$, and a lifting $w \in V$ of \bar{w} in such a way that w will be a (non-zero) weight vector in V of weight ϕ for \bar{s} .

Choose any lifting w of \bar{w} . The $\mathbb{C}[\partial]$ -submodule W spanned by V^0 together with w is preserved by the action of S . Let us fix a base of W consisting of some S -triangular base (v_1, v_2, \dots, v_h) for V^0 along with the lifting w . According to the matrix representation introduced in Section 4.2, the action of s on W will be represented by the following matrix:

$$\begin{pmatrix} 0 & * & \dots & * & X^1(\lambda) \\ 0 & 0 & * & * & X^2(\lambda) \\ \vdots & & \ddots & \vdots & \\ 0 & \dots & \dots & 0 & X^n(\lambda) \\ 0 & \dots & \dots & 0 & \phi(\lambda) \end{pmatrix}, \tag{5.1}$$

where $\phi(\lambda)$ is the weight of \bar{w} . Notice that coefficients X^i may depend on ∂ if they refer to an element of the $\mathbb{C}[\partial]$ -basis of the free part of V^0 , but only depend on λ when they refer to basis elements of $\text{Tor } V^0$.

If all of the X^i are zero, then w is a weight vector, and we are done. If instead some of the X^i are nonzero, let us choose i to be maximal with the property that X^i is nonzero. I will show that I can find an element $\bar{s} = s + s'$, $s' \in S'$ and a lifting w' of \bar{w} such that, in the matrix representation of $s + s'$ with respect to the base (v_1, \dots, v_h, w') of W , all X^j , $j \geq i$ vanish. An easy induction will then prove the statement.

If all X^j , $j > i$ are zero and $X^i \neq 0$, then we can compute the corresponding matrix coefficient in the commutator $[s_\alpha s]$. By (4.7), it is given by

$$X^i(\partial, \alpha)\phi(\lambda - \alpha) - X^i(\partial, \lambda - \alpha)\phi(\alpha).$$

Let us write

$$\phi(\lambda) = \sum_{i=0}^n \phi_i \lambda^i, \quad X^i(\partial, \lambda) = \sum_{j=0}^m X_j(\partial) \lambda^j.$$

Say m is even, and recall that we chose n to be odd. Then the i th entry in the last column of the matrix representing the α^{m+n} coefficient (call it t_1) in $[s_\alpha s]$ equals $-2X_m(\partial)\phi_n$. Hence the matrix representing the element

$s_1 = (-\partial)^m t_1 / 2\phi_n$ is upper triangular with zero eigenvalues and the i .th entry in the last column is precisely $-X_m(\partial)\lambda^m$, opposite to the highest degree in λ of $X^i(\partial, \lambda)$. Thus, the matrix representing $s + s_1$ has the same eigenvalues as s , all X^j , $j > i$ vanish, and the degree in λ of X^i is lower. Notice that $s_1 \in S'$.

On the other hand, if m is odd, assume $m \neq n$. Then the i .th entry in the last column of the matrix representing the coefficient (call it t_2) multiplying α^{m+n-1} in $[s_\alpha s]$ is given by

$$2X_m(\partial)\phi_{n-1} - 2X_{m-1}(\partial)\phi_n + (n - m)\lambda X_m(\partial)\phi_n.$$

Then the matrix representing the element $s_2 = (-\partial)^{m-1} t_2 / (m - n)\phi_n$ is upper triangular with zero eigenvalues. Moreover, the i .th entry in the last column has the same top degree term in λ as X^i , with opposite sign. As before, the element $s + s_2$ has the same eigenvalues as s , all X^j , $j > i$ vanish, and the degree of X^i is lower. Notice that $s_2 \in S'$ as well. Finally, when $m = n$, it is enough to replace w with $w + X_m(\partial)v_i/\phi_m$ in order to kill the term of top degree of X^i in the matrix representation of s .

By induction, we can then find an $\bar{s} = s + s'$, $s' \in S'$ and a lifting w of \bar{w} in such a way that the corresponding matrix is upper triangular with the same eigenvalues as s , and all X^j , $j \geq i$ vanish. \square

Theorem 5.2. *Let V be a finite reduced vertex algebra, $s \in V$. Then the adjoint action of s on V is nilpotent.*

Proof. If there is some $s \in S$ for which the adjoint conformal action on V is not nilpotent, then Lemma 5.1 finds an element $\bar{s} \in \langle s \rangle$ possessing a weight vector of non-zero weight ϕ . Then Theorem 4.4 shows that the generalized weight space V^ϕ with respect to the conformal subalgebra generated by s is a nonzero abelian ideal of the vertex algebra V , which is a contradiction, as V is reduced. \square

Proof of Theorem 5.1. Every element $v \in V^{\text{Lie}}$ has a nilpotent adjoint conformal action. By Theorem 4.1, V^{Lie} is then a nilpotent Lie conformal algebra. \square

Corollary 5.1. *Let V be a finite vertex algebra, and define*

$$V^{[0]} = V, \quad V^{[n+1]} = [V, V^{[n]}], \quad n \geq 0.$$

Then all $V^{[n]}$ are ideals of V , and the descending sequence

$$V = V^{[0]} \supset \dots \supset V^{[n]} \supset \dots$$

stabilizes on an ideal contained in the nilradical of V .

Proof. The quotient of V by its nilradical N is nilpotent, hence the above sequence for V/N stabilizes to zero. By lifting it back to V , this only happens if the claim holds. \square

In other words, every finite vertex algebra can be expressed as an extension of a nilpotent vertex algebra by a nil-ideal.

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