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Periodic Orbits on Discrete Dynamical Systems

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Abstract—In this paper, we discuss the discrete dynamical system

$$x_{n+1} = \beta x_n - g(x_n), \qquad n = 0, 1, \dots,$$
 (*)

arising as a discrete-time network of single neuron, where β is the internal decay rate, g is a signal function. First, we consider the case where g is of McCulloch-Pitts nonlinearity. Periodic orbits are discussed according to different range of β . Moreover, we can construct periodic orbits. Then, we consider the case where g is a sigmoid function. Sufficient conditions are obtained for (*) has periodic orbits of arbitrary periods and an example is also given to illustrate the theorem. © 2003 Elsevier Science Ltd. All rights reserved.

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1. INTRODUCTION

There are a lot of problems which can be described, at least to a crude approximation, by a simple first-order difference equation

$$x_{n+1} = f(x_n), \qquad n = 0, 1, \dots$$
 (1.1)

Studies of the dynamical properties of such models can provide us with many useful messages. There are many works in this area [1-5]. Since Li and York [6] and May [7] noticed that even simple mathematical models can demonstrate very complicated dynamics, people have paid more attention to the periodic solutions and "chaotic" behavior of discrete dynamical systems; we refer to [8–10]. However, most works of (1.1) are based on the assumption that f is continuous. But in reality, for example, in some models of neural networks, f may be not continuous [11–13].

In this paper, we consider the following discrete-time equation:

$$x_{n+1} = \beta x_n - g(x_n), \qquad n = 0, 1, \dots,$$
 (1.2)

where $\beta \in (0, \infty)$, g is a nonlinear function.

Equation (1.2) arises as a discrete-time network of single neuron where β is the internal decay rate, g is a signal function [11]. We consider two cases of the signal function:

- (I) g is of McCulloch-Pitts nonlinearity;
- (II) g is a sigmoid function.

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Now we recall some definitions [8].

DEFINITION 1.1. A point x_s is called a stationary state of (1.1) if

 $f(x_s) = x_s.$

DEFINITION 1.2. A stationary state x_s of (1.1) is stable if for every $\epsilon > 0$, there exists $\delta > 0$ such that

$$|x_0 - x_s| \leq \delta$$
 implies that $|x_n - x_s| \leq \epsilon$ for all $n \geq 1$.

A stationary state x_s that is not stable is said to be unstable.

DEFINITION 1.3. An orbit $O(x_0) = \{x_0, x_1, x_2, ...\}$ of (1.1) is said to be periodic of period $p \ge 2$ if

$$x_p = x_0$$
 and $x_i \neq x_0$, for $1 \le i \le p-1$.

DEFINITION 1.4. A periodic orbit $\{x_0, x_1, \ldots, x_{p-1}, \ldots\}$ of period p is stable if each point x_i , $i = 0, 1, \ldots, p-1$, is a stable stationary state of dynamical system $x_{n+1} = f^p(x_n)$. A periodic orbit $\{x_0, x_1, \ldots, x_{p-1}, \ldots\}$ of period p which is not stable is said to be unstable.

DEFINITION 1.5. A point z is said to be a limit point of $O(x_0)$ if there exists a subsequence $\{x_{n_k} : k = 0, 1, ...\}$ of $O(x_0)$ such that $|x_{n_k} - z| \to 0$ as $k \to \infty$. The limit set $L(x_0)$ of the orbit $O(x_0)$ is the set of all limit points of the orbit.

DEFINITION 1.6. An orbit $O(x_0)$ is said to be asymptotically periodic if its limit set is a periodic orbit. An orbit $O(x_0)$ such that $x_{n+p} = x_n$ for some $n \ge 1$ and some $p \ge 2$ is said to be eventually periodic.

In Section 2, we shall discuss the periodic orbits of (1.2) when g is of McCulloch-Pitts nonlinearity. We shall show that (1.2) has a unique stable periodic orbit of period 2 when $\beta \in (0, 1)$; equation (1.2) has infinitely many stable periodic orbits of period 2 when $\beta = 1$; equation (1.2) has an unstable stational state and unstable periodic orbits of period 2,4 when $\beta \in (1, \sqrt{2})$; equation (1.2) has unstable periodic orbits of all even periods when $\beta \in [\sqrt{2}, (1 + \sqrt{5})/2)$, and (1.2) has periodic orbits of arbitrary periods when $\beta \in [(1 + \sqrt{5})/2, \infty)$. Moreover, we can construct the periodic orbits of (1.2). In Section 3, we consider (1.2) when g is a sigmoid function. By using Li and York's theorem [6], we get a sufficient condition for (1.2) to have periodic orbits of arbitrary periods. We also give an example to illustrate the result.

2. g IS OF MCCULLOCH-PITTS NONLINEARITY

Throughout this section, we will assume that g is of McCulloch-Pitts nonlinearity and

$$g(x) = \begin{cases} 1, & \text{if } x \ge 0, \\ -1, & \text{if } x < 0. \end{cases}$$
(2.1)

First, we consider (1.2) when $\beta \in (0, 1)$, and we have the following theorem.

THEOREM 2.1. Assume that $\beta \in (0,1)$. Then the periodic orbit $O(1/(\beta+1))$ is a stable periodic orbit with period 2. And for every $x_0 \in R$, the orbit $O(x_0)$ is asymptotically periodic with $L(x_0) = \{1/(\beta+1), -1/(\beta+1)\}.$

PROOF. Let $h(x) = \beta x - g(x)$. Then (1.2) can be rewritten

$$x_{n+1} = h(x_n), \qquad n = 0, 1, \dots$$
 (2.2)

Clearly,

$$h\left(\frac{1}{\beta+1}\right) = -\frac{1}{\beta+1}, \qquad h\left(-\frac{1}{\beta+1}\right) = \frac{1}{\beta+1};$$

this implies $O(1/(\beta+1)) = \{1/(\beta+1), -1/(\beta+1), \dots\}$ is a periodic orbit of (1.2) with period 2.

Now we show that $O(1/(\beta + 1))$ is stable. In fact, for each $\epsilon > 0$, let $\delta = \min\{1/(\beta + 1), \epsilon\}$. Then

$$\left|x_0 - \frac{1}{\beta+1}\right| < \delta \text{ implies } \left|h^{2n}(x_0) - \frac{1}{\beta+1}\right| = \beta^{2n} \left|x_0 - \frac{1}{\beta+1}\right| < \epsilon$$

This shows that $1/(\beta+1)$ is a stable stationary state of the dynamical $x_{n+1} = h^2(x_n)$. Similarly, $-1/(\beta+1)$ is a stable stationary state of the dynamical $x_{n+1} = h^2(x_n)$.

For every $x_0 \in R$, there is no harm in assuming that $x_0 \ge 0$. Since $x_{n+1} = \beta x_n - 1 \le x_n - 1$ for $x_n \ge 0$, there exists a nonnegative integer n_0 such that $x_i \ge 0$ for $0 \le i \le n_0$ and $x_{n_0+1} < 0$. Thus,

$$x_{n_0+2} = \beta x_{n_0+1} + 1 = \beta^2 x_{n_0} + 1 - \beta$$

This leads to

$$x_{n_0+2} - \frac{1}{\beta+1} = \beta^2 \left(x_{n_0} - \frac{1}{\beta+1} \right).$$

In general, we have

$$x_{n_0+2n} - \frac{1}{\beta+1} = \beta^{2n} \left(x_{n_0} - \frac{1}{\beta+1} \right).$$

Therefore, $\lim_{n\to\infty} x_{n_0+2n} = 1/(\beta+1)$.

Similarly, we get $\lim_{n\to\infty} x_{n_0+1+2n} = -1/(\beta+1)$. Thus, $L(x_0) = \{1/(\beta+1), -1/(\beta+1)\}$, and the proof is complete.

Similar to proof of Theorem 2.1, we can show the following theorem.

THEOREM 2.2. Assume $\beta = 1$. Then, for every $x_0 \in R$, the orbit $O(x_0)$ of (1.2) is eventually periodic with period 2.

Now we consider the case when $\beta > 1$.

THEOREM 2.3. Let $\beta \in (1, \infty)$. Then (1.2) has a stationary state, periodic orbits of period 2,4. Suppose there is a positive integer k such that

$$\beta^{2k} - 2\beta^{2k-2} + 1 > 0. \tag{2.3}$$

Then, (1.2) has a periodic orbit of period 2k. Suppose there is a positive integer k such that

$$\beta^{2k+1} - 2\beta^{2k-1} - 1 \ge 0. \tag{2.4}$$

Then, (1.2) has a periodic orbit of period 2k + 1.

PROOF. $1/(\beta - 1)$ is a stationary state of (1.2); $O(1/(\beta + 1))$ is a 2-periodic orbit of (1.2) and $O((\beta + 1)/(\beta^2 + 1))$ is a 4-periodic orbit of (1.2).

Now suppose (2.3) holds. We will construct a periodic orbit $O(x_0)$ of period 2k such that

$$\begin{aligned} x_0 > 0, \quad x_1 \ge 0, \quad x_2 < 0, \quad x_3 < 0, \quad (-1)^i x_i > 0, \\ \text{for } i = 4, \dots, 2k - 1, \quad x_{2k} = x_0. \end{aligned}$$
 (2.5)

Then

$$\begin{aligned} x_1 &= \beta x_0 - 1 \ge 0, \\ x_2 &= \beta^2 x_0 - \beta - 1 < 0, \\ x_3 &= \beta^3 x_0 - \beta^2 - \beta + 1 < 0, \\ x_4 &= \beta^4 x_0 - \beta^3 - \beta^2 + \beta + 1 > 0, \\ x_5 &= \beta^5 x_0 - \beta^4 - \beta^3 + \beta^2 + \beta - 1 < 0, \\ \vdots \\ x_{2k-1} &= \beta^{2k-1} x_0 - \beta^{2k-2} - \beta^{2k-3} + \beta^{2k-4} + \beta^{2k-5} - \beta^{2k-6} + \dots - 1, \\ x_{2k} &= \beta^{2k} x_0 - \beta^{2k-1} - \beta^{2k-2} + \beta^{2k-3} + \beta^{2k-4} - \beta^{2k-5} + \beta^{2k-6} - \dots + 1. \end{aligned}$$

Since $x_{2k} = x_0$, we get

$$x_0 = \frac{\beta^{2k-1} + \beta^{2k-2} - \beta^{2k-3} - \beta^{2k-4} + \beta^{2k-5} - \beta^{2k-6} + \dots - 1}{\beta^{2k} - 1}.$$
 (2.6)

By (1.2), we see that

$$\min\{x_0, x_1, x_4, x_6, \dots, x_{2k-2}\} = x_1 = \frac{\beta^{2k-1} - \beta^{2k-2} - \beta^{2k-3} + \beta^{2k-4} - \beta^{2k-5} + \beta^{2k-6} - \dots + 1}{\beta^{2k} - 1}$$
$$= \frac{(\beta - 1) \left(\beta^{2k-2} - \beta^{2k-4} - \beta^{2k-6} - \dots - 1\right)}{\beta^{2k} - 1}$$
$$= \frac{\beta^{2k} - 2\beta^{2k-2} + 1}{(\beta^{2k} - 1) (\beta + 1)} > 0$$

and

$$\max\{x_2, x_3, x_5, x_7, \dots, x_{2k-1}\} = x_{2k-1}$$

= $\frac{-\beta^{2k-1} + \beta^{2k-2} + \beta^{2k-3} - \beta^{2k-4} - \beta^{2k-5} + \beta^{2k-6} - \dots + 1}{\beta^{2k} - 1}$
< 0.

Thus, if (2.3) holds, the orbit $O(x_0)$ where x_0 is defined by (2.6) satisfies (2.5) and is a periodic orbit of period 2k for (1.2).

Similarly, if (2.4) holds, we can show that the orbit $O(x_0)$ where

$$x_0 = \frac{\beta^{2k} - \beta^{2k-1} - \beta^{2k-2} + \beta^{2k-3} - \beta^{2k-4} + \beta^{2k-5} - \beta^{2k-6} + \dots - 1}{\beta^{2k+1} - 1}$$
(2.7)

is a periodic orbit of period 2k + 1 for (1.2) and satisfies $x_0 \ge 0$, $x_1 < 0$, $x_2 < 0$, $(-1)^i x_i < 0$ for i = 3, ..., 2k.

REMARK 2.1. If $\beta > 1$, all periodic orbits of (1.2) are unstable and (2.3) always holds for k = 1, 2.

By a simple computation, we see that the equation

$$x^{2k} - 2x^{2k-2} + 1 = 0 (2.8)$$

has a unique positive real root α_k for $k \ge 3$. Clearly, (2.3) holds for $\beta > \alpha_k$. Let $\alpha_1 = \alpha_2 = 1$ since (2.3) always holds for k = 1 and k = 2. Then, $\{\alpha_k\}_{k\ge 2}$ is an increasing sequence in $[1, \sqrt{2})$ and $\lim_{k\to\infty} \alpha_k = \sqrt{2}$.

The equation

$$x^{2k+1} - 2x^{2k-1} - 1 = 0 (2.9)$$

has a unique positive real root β_k for $k \ge 1$ and $\beta_1 = (1 + \sqrt{5})/2$. Clearly, (2.4) holds for $\beta \ge \beta_k$, $\{\beta_k\}$ is a decreasing sequence in $(\sqrt{2}, 2)$, and $\lim_{k\to\infty} \beta_k = \sqrt{2}$.

Based on the above discussion and Theorem 2.3, we have the following corollary.

COROLLARY 2.4. If (2.3) holds, (1.2) has periodic orbits of period 2k, 2k - 2, ..., 4, 2 and a stationary state. If (2.4) holds, (1.2) has periodic orbits of period 2k + 1, 2k + 3, ... Especially, for $\beta \in [\sqrt{2}, \infty)$, (1.2) has periodic orbits of arbitrary even periods; for $\beta \in [(1 + \sqrt{5})/2, \infty)$, (1.2) has periodic orbits of arbitrary periods.

REMARK 2.2. $\beta \in [(1+\sqrt{5})/2, \infty)$ is a necessary and sufficient condition for (1.2) has a periodic orbit of period 3.

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3. g IS A SIGMOID FUNCTION

In this section, motivated by [14], we assume that there exist positive constants r, ϵ such that

$$\begin{aligned} |g(x) - 1| &\leq \epsilon, & \text{if } x \geq r, \\ |g(x) + 1| &\leq \epsilon, & \text{if } x \leq -r. \end{aligned}$$
(3.1)

THEOREM 3.1. Assume that $\beta \in ((1+\sqrt{5})/2, \infty)$, $r \in (0, (\beta^2-\beta-1)/(\beta^3-1))$, g(x) is continuous in R and satisfies (3.1). If

$$\frac{1}{\beta^{2}}\left(1+\frac{1}{\beta^{2}}+\frac{1}{\beta^{4}}\right)\left(\frac{\beta^{2}-\beta-1}{\beta^{3}-1}-r\right)^{2}+\frac{\epsilon^{2}}{\beta^{2}}\left(3+\frac{4}{\beta}+\frac{4}{\beta^{2}}+\frac{2}{\beta^{3}}+\frac{1}{\beta^{4}}\right) +\frac{2\epsilon}{\beta^{2}}\left(1+\frac{1}{\beta}+\frac{2}{\beta^{2}}+\frac{1}{\beta^{3}}+\frac{1}{\beta^{4}}\right)\left(\frac{\beta^{2}-\beta-1}{\beta^{3}-1}-r\right)\leq\left(\frac{\beta^{2}-\beta-1}{\beta^{3}-1}-r\right)^{2},$$
(3.2)

for every m = 1, 2, ..., (1.2) has a periodic orbit of period m.

PROOF. Since $\beta x - g(x)$ is continuous in R, by Theorem 1 of [6], it suffices if we can prove that (1.2) has a periodic orbit of period 3. We rewrite (1.2) as

$$x_n = \frac{1}{\beta} x_{n+1} + \frac{1}{\beta} g(x_n).$$
(3.3)

Now, we show that

$$y_{n+1} = \frac{1}{\beta} y_n + \frac{1}{\beta} g(y_{n-2}) \tag{3.4}$$

has a 3-periodic solution $\{y_n^*\}$. Then $O(y_0^*)$ is a 3-periodic orbit of (1.2).

If y_{-2} , y_{-1} , y_0 are given, by (3.4), we have

$$y_{1} = \frac{1}{\beta}y_{0} + \frac{1}{\beta}g(y_{-2}),$$

$$y_{2} = \frac{1}{\beta^{2}}y_{0} + \frac{1}{\beta^{2}}g(y_{-2}) + \frac{1}{\beta}g(y_{-1}),$$

$$y_{3} = \frac{1}{\beta^{3}}y_{0} + \frac{1}{\beta^{3}}g(y_{-2}) + \frac{1}{\beta^{2}}g(y_{-1}) + \frac{1}{\beta}g(y_{0}).$$
(3.5)

We define a map H from R^3 to R^3 by $H(y_1, y_2, y_3) = (z_1, z_2, z_3)$ where

$$z_{1} = \frac{1}{\beta} y_{3} + \frac{1}{\beta} g(y_{1}),$$

$$z_{2} = \frac{1}{\beta^{2}} y_{3} + \frac{1}{\beta^{2}} g(y_{1}) + \frac{1}{\beta} g(y_{2}),$$

$$z_{3} = \frac{1}{\beta^{3}} y_{3} + \frac{1}{\beta^{3}} g(y_{1}) + \frac{1}{\beta^{2}} g(y_{2}) + \frac{1}{\beta} g(y_{3}).$$
(3.6)

Let $\bar{y} = (\bar{y}_1, \bar{y}_2, \bar{y}_3)$ where

$$\bar{y}_1 = \frac{-\beta^2 + \beta + 1}{\beta^3 - 1}, \qquad \bar{y}_2 = \frac{\beta^2 - \beta + 1}{\beta^3 - 1}, \qquad \bar{y}_3 = \frac{\beta^2 + \beta - 1}{\beta^3 - 1}.$$

Here \bar{y}_3 , \bar{y}_2 , \bar{y}_1 are the first three terms of the 3-periodic orbit $O((\beta^2 + \beta - 1)/(\beta^3 - 1))$ in Theorem 2.3. Clearly, $\bar{y}_1 < 0$, $\bar{y}_2 > 0$, $\bar{y}_3 > 0$ and we have

$$\begin{split} \bar{y}_1 &= \frac{1}{\beta} \bar{y}_3 - \frac{1}{\beta}, \\ \bar{y}_2 &= \frac{1}{\beta^2} \bar{y}_3 - \frac{1}{\beta^2} + \frac{1}{\beta}, \\ \bar{y}_3 &= \frac{1}{\beta^3} \bar{y}_3 - \frac{1}{\beta^3} + \left(\frac{1}{\beta^2} + \frac{1}{\beta}\right). \end{split}$$
(3.7)

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$$\begin{split} \text{If } \sqrt{(y_1 - \bar{y}_1)^2 + (y_2 - \bar{y}_2)^2 + (y_3 - \bar{y}_3)^2} &\leq \delta_0 = (\beta^2 - \beta - 1)/(\beta^3 - 1) - r, \\ y_1 &\leq \bar{y}_1 + \delta_0 = -r, \\ y_2 &\geq \bar{y}_2 - \delta_0 = \frac{\beta^2 - \beta + 1}{\beta^3 - 1} - \frac{\beta^2 - \beta - 1}{\beta^3 - 1} + r \geq r, \\ y_3 &\geq \bar{y}_3 - \delta_0 = \frac{\beta^2 + \beta - 1}{\beta^3 - 1} - \frac{\beta^2 - \beta - 1}{\beta^3 - 1} + r \geq r. \end{split}$$

In view of (3.6), (3.7), we get

$$z_1 - \bar{y}_1 = \frac{1}{\beta} (y_3 - \bar{y}_3) + \frac{1}{\beta} (g(y_1) + 1).$$

Noticing (3.1), we have

$$|z_1 - ar{y}_1| \leq rac{1}{eta} \, |y_3 - ar{y}_3| + rac{1}{eta} |g(y_1) + 1| \leq rac{1}{eta} \delta_0 + rac{1}{eta} \epsilon.$$

Similarly, we get

$$|z_2 - \bar{y}_2| \le \frac{1}{\beta^2} \delta_0 + \left(\frac{1}{\beta^2} + \frac{1}{\beta}\right) \epsilon, \qquad |z_3 - \bar{y}_3| \le \frac{1}{\beta^3} \delta_0 + \left(\frac{1}{\beta^3} + \frac{1}{\beta^2} + \frac{1}{\beta}\right) \epsilon$$

Therefore, by (3.2),

$$\begin{aligned} (z_1 - \bar{y}_1)^2 + (z_2 - \bar{y}_2)^2 + (z_3 - \bar{y}_3)^2 &\leq \frac{1}{\beta^2} \left(1 + \frac{1}{\beta^2} + \frac{1}{\beta^4} \right) \delta_0^2 + \frac{\epsilon^2}{\beta^2} \left(3 + \frac{4}{\beta} + \frac{4}{\beta^2} + \frac{2}{\beta^3} + \frac{1}{\beta^4} \right) \\ &+ \frac{2\delta_0 \epsilon}{\beta^2} \left(1 + \frac{1}{\beta} + \frac{2}{\beta^2} + \frac{1}{\beta^3} + \frac{1}{\beta^4} \right) \\ &\leq \left(\frac{\beta^2 - \beta - 1}{\beta^3 - 1} - r \right)^2 = \delta_0^2. \end{aligned}$$

Thus, H maps $N(\bar{y}, \delta_0) = \{(y_1, y_2, y_3) \in \mathbb{R}^3 : (y_1 - \bar{y}_1)^2 + (y_2 - \bar{y}_2)^2 + (y_3 - \bar{y}_3)^2 \leq \delta_0^2\}$ into $N(\bar{y}, \delta_0)$. By the Brouwer fixed-point theorem, H has a fixed point $(y_1^*, y_2^*, y_3^*) \in N(\bar{y}, \delta_0)$. The solution of (3.4) with initial condition $y_{-2} = y_1^*$, $y_{-1} = y_2^*$, $y_0 = y_3^*$ is a periodic solution of period 3. The theorem is now complete.

REMARK 3.1. It is obvious that $(1/\beta^2)(1 + 1/\beta^2 + 1/\beta^4) < 1$ for $\beta \in ((1 + \sqrt{5})/2, \infty)$, and thus, (3.2) can be satisfied if we choose small ϵ .

EXAMPLE. Let c > 0 and $g_0(x) = (1 - e^{-cx})/(1 + e^{-cx})$. Then

$$\begin{aligned} |g_0(x) - 1| &= \frac{2e^{-cx}}{1 + e^{-cx}} \le 2e^{-cr}, & \text{if } x \ge r, \\ |g_0(x) + 1| &= \frac{2}{1 + e^{-cx}} \le 2e^{-cr}, & \text{if } x \le -r. \end{aligned}$$

Assuming $\beta \in ((1 + \sqrt{5})/2, \infty)$, $r \in (0, (\beta^2 - \beta - 1)/(\beta^3 - 1))$, (3.2) holds, and $c \ge (1/r) \ln(2/\epsilon)$ $(\epsilon < 2)$, then (3.1) holds for $g = g_0$. According to Theorem 3.1,

$$x_{n+1} = \beta x_n - g_0(x_n), \qquad n = 0, 1, \ldots,$$

has periodic orbits of arbitrary periods.

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