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## On the Solution of Equations Containing Radicals by the Decomposition Method

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Equations containing radicals are solved using Adomian's decomposition method. © 1985 Academic Press, Inc.

The decomposition method [1, 2] has been used to solve wide classes of equations [2]. Adomian [2] and Adomian and Rach [3] have investigated use of the method when composite nonlinearities of the form  $N_0(N_1x)$  or  $N_0(N_1(N_2x))$ , etc., are involved. We now deal specifically with algebraic or differential equations containing radicals using the above referenced works.

We will consider an example in the standard form used in [1, 2]  $Ly + Ny = x(t)$ , where  $L$  is a linear operator and  $N$  is a nonlinear operator. First suppose we are considering an algebraic equation with  $L$  simply being multiplication by a constant  $m$ , with  $x(t)$  a constant  $k$ , and  $Ny = (\alpha y^2 + \beta)^{1/2}$ .  $L^{-1}$  will now mean division by  $m$  and  $y = y_0 - L^{-1}Ny$  or

$$y = \sum_{n=0}^{\infty} y_n = y_0 - (1/m)(\alpha y^2 + \beta)^{1/2} \quad (1)$$

where  $y_0 = k/m$ . We treat the nonlinear or radical term  $(\alpha y^2 + \beta)^{1/2}$  as a composite nonlinearity [2, 3] with  $N_0(u^0) = \sqrt{u^0} = \sum_{n=0}^{\infty} A_n^0$ , where  $u^0 = \alpha(N_1 u^1) + \beta$  and  $N_1 u^1 = \sum_{n=0}^{\infty} A_n^1$  with  $u^1 = y$ . Now

$$\sum_{n=0}^{\infty} u_n^0 = \alpha \sum_{n=0}^{\infty} A_n^1 + \beta \quad (2)$$

so that

$$u_0^0 = \beta + \alpha A_0^1 = A_0^0$$

For  $n \geq 1$  we have

$$u_n^0 = \alpha A_n^1 = A_n^0.$$

Calculating the  $A_n^0$  polynomials

$$\begin{aligned}
 A_0^0 &= (u_0^0)^{1/2} = (\beta + \alpha A_0^1)^{1/2} \\
 A_1^0 &= (\frac{1}{2})(u_0^0)^{-1/2}(u_1^0) = (\frac{1}{2})(\beta + \alpha A_0^1)^{-1/2}(\alpha A_1^1) \\
 A_2^0 &= (\frac{1}{2})(u_0^0)^{-1/2}(u_2^0) - (\frac{1}{8})(u_0^0)^{-3/2}(u_1^0)^2 \\
 A_3^0 &= (\frac{1}{2})(u_0^0)^{-1/2}(u_3^0) - (\frac{1}{4})(u_0^0)^{-3/2}(u_1^0)(u_2^0) \\
 &\quad + (\frac{1}{16})(u_0^0)^{-5/2}(u_1^0)^3 \\
 A_4^0 &= (\frac{1}{2})(u_0^0)^{-1/2}(u_4^0) - (\frac{1}{4})(u_0^0)^{-3/2}\{(u_2^0)^2/2 + (u_1^0)(u_3^0)\} \\
 &\quad + (\frac{3}{16})(u_0^0)^{-5/2}(u_1^0)^2(u_2^0) - (\frac{5}{128})(u_0^0)^{-7/2}(u_1^0)^4 \\
 A_5^0 &= (\frac{1}{2})(u_0^0)^{-1/2}(u_5^0) - (\frac{1}{4})(u_0^0)^{-3/2}\{u_2^0 u_3^0 + u_1^0 u_4^0\} \\
 &\quad + (\frac{3}{16})(u_0^0)^{-5/2}\{(u_1^0)(u_2^0)^2 + (u_1^0)^2(u_3^0)\} \\
 &\quad - (\frac{5}{32})(u_0^0)^{-7/2}(u_1^0)^3(u_2^0) + (\frac{7}{256})(u_0^0)^{-9/2}(u_1^0)^5 \\
 &\quad \vdots
 \end{aligned}$$

Since  $y^2 = \sum_{n=0}^{\infty} A_n^1$ , we calculate

$$\begin{aligned}
 A_0^1 &= y_0^2 \\
 A_1^1 &= 2y_0 y_1 \\
 A_2^1 &= y_1^2 + 2y_0 y_2 \\
 A_3^1 &= 2y_1 y_2 + 2y_0 y_3 \\
 A_4^1 &= y_2^2 + 2y_1 y_3 + 2y_0 y_4 \\
 A_5^1 &= 2y_2 y_3 + 2y_1 y_4 + 2y_0 y_5 \\
 &\quad \vdots
 \end{aligned}$$

Now

$$\begin{aligned}
 y_0 &= k/m \\
 y_1 &= -(1/m)(\beta + \alpha(k^2/m^2))^{1/2} \\
 y_2 &= -(\alpha/m)(k/m)\{(-1/m)(\beta + (\alpha k^2/m^2))^{1/2}(\beta + \alpha y_0^2)\}^{-1/2} = \alpha k/m^3.
 \end{aligned}$$

Indicating by  $\phi_n$  the  $n$  term approximation  $\sum_{v=0}^{n-1} y_v$ , we have the following results if we choose  $\alpha = \frac{1}{4}$ ,  $\beta = 1$ ,  $m = 10$ ,  $k = 5$ :

$$\begin{aligned}
 y_0 &= 0.5000000000 & \phi_1 &= 0.5000000000 \\
 y_1 &= -0.1030776406 & \phi_2 &= 0.3969223594 \\
 y_2 &= 0.0012500000 & \phi_3 &= 0.3981723594
 \end{aligned}$$

which is already correct to three decimal places ( $y_3$  is  $1.364 \times 10^{-4}$ ). (The correct result is 0.4000000 so that the error by  $\phi_3$  is less than half a percent.) Since this result is obtained with only three terms, and we are solving the *nonlinear* problem, not a linearized version, it is rather remarkable.

Now let us consider a *differential equation in the same standard form*. Let  $L = d/dt$ , and  $x = 0$  (for convenience only [2]).  $L^{-1}$  is now the definite integral from 0 to  $t$ . Let  $y(0) = k$  and  $Ny = (\alpha y^2 + \beta)^{1/2}$  as before. Again,

$$y = y_0 - L^{-1}Ny.$$

Letting  $y = \sum_{n=0}^{\infty} y_n$ , we have  $y_0 = y(0) = k$  and since we have the same  $A_n$  polynomials,

$$\begin{aligned} y_0 &= k \\ y_1 &= -L^{-1}(\beta + \alpha k^2)^{1/2} = -(\beta + \alpha k^2)^{1/2} t \\ y_2 &= -L^{-1}(\beta + \alpha k^2)^{-1/2} \alpha y_0 y_1 = \alpha k t^2 / 2 \\ y_3 &= -L^{-1}(1/2)(\beta + \alpha k^2)^{-1/2} \alpha (y_1^2 + 2y_0 y_2) \\ &\quad - (1/8)(\beta + \alpha k^2)^{-3/2} (2\alpha y_0 y_1)^2 \\ &= -\alpha(\beta + \alpha k^2)^{1/2} t^3 / 6 \\ &\vdots \\ y &= k - (\beta + \alpha k^2)^{1/2} t + \alpha k t^2 / 2 - \alpha(\beta + \alpha k^2)^{1/2} t^3 / 6 \dots, \end{aligned}$$

as the solution for

$$\frac{dy}{dt} + \sqrt{\beta + \alpha y^2} = 0 \quad y(0) = k.$$

As shown in [2], cases where  $L$  is a higher order operator or  $x(t) \neq 0$  are straightforward generalizations. It is interesting to choose some values for the constants. Let  $k = 5$ ,  $\alpha = \frac{1}{4}$ ,  $\beta = 1$ . Then

$$\begin{aligned} y_0 &= 5.0 \\ y_1 &= -2.692t \\ y_2 &= 0.625t^2 \\ y_3 &= -0.112t^3. \end{aligned}$$

The solution as a function of  $t$  is a rapidly damped oscillation for small  $t$ , e.g., at  $t = 0.1$ , the 1-term approximation of  $\phi$  is 5.0, the 2-term approximation  $\phi_2$  is 4.73, and further changes by going to  $\phi_3$  or  $\phi_4$  are in

the third decimal place. For  $t = 1.0$ ,  $\phi_1 = 5$ ,  $\phi_2 = 2.31$ ,  $\phi_3 = 2.93$ ,  $\phi_4 = 2.82, \dots$ , so we expect  $2.82 < \phi_n < 2.93$  for  $n > 4$  and can go to  $\phi_6$  with the already computed polynomials.

#### REFERENCES

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