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Status of the CUORE and results from the CUORE-0 neutrinoless double beta decay experiments

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Abstract

CUORE is a 741 kg array of TeO₂ bolometers for the search of neutrinoless double beta decay of 130 Te. The detector is being constructed at the Laboratori Nazionali del Gran Sasso, Italy, where it will start taking data in 2015. If the target background of 0.01 counts/(keV·kg·y) will be reached, in five years of data taking CUORE will have a 1σ half life sensitivity of 10^{26} y. CUORE-0 is a smaller experiment constructed to test and demonstrate the performances expected for CUORE. The detector is a single tower of 52 CUORE-like bolometers that started taking data in spring 2013. The status and perspectives of CUORE will be discussed, and the first CUORE-0 data will be presented.

Keywords: Double beta decay, Neutrino mass, Bolometers

1. Introduction

Neutrinos are massive particles. A beautiful proof of this important property was obtained by neutrino oscillation experiments more than a decade ago. Since then, the key role of neutrinoless double beta decay searches has been established, as attested by the growing number of experimental proposals in the last years.

Neutrinoless double beta decay $(\beta\beta0\nu)$ is a proposed very rare nuclear process in which a nucleus transforms into its (A, Z+2) isobar with the emission of two electrons. While the two neutrino channel $(\beta\beta 2\nu)$ – where two neutrinos are contemporary emitted in the decay – is allowed by the Standard Model of Particle Physics and has been observed experimentally in a dozen of isotopes with half-lives of the order $10^{18} - 10^{21}$ y, the neutrinoless mode is a lepton number violating process which can occur only if the Majorana character of the neutrino is allowed. Therefore, $\beta\beta0\nu$ offers a unique experimental chance to investigate still unresolved fundamental questions, since its observation would undoubtedly unveil the neutrino character, confirm lepton number violation and allow to assess the absolute neutrino mass scale with high sensitivity, thus helping point us towards the proper extension of the Standard Model [1–3].

Neutrinoless double beta decay can proceed via different mechanisms: in the case of a virtual exchange of a light Majorana neutrino between the two nucleons, the decay rate is proportional to the square of the so-called effective Majorana mass $|\langle m_{\beta\beta} \rangle|$ (a coherent sum of neutrino mass eigenstates) [4]:

$$[T_{1/2}^{0\nu}]^{-1} = \frac{|\langle m_{\beta\beta}\rangle|^2}{m_*^2} G^{0\nu} |M^{0\nu}|^2 \tag{1}$$

where $T_{1/2}^{0\nu}$ is the decay half-life, $G^{0\nu}$ is the two-body phase-space integral, $M^{0\nu}$ is the $\beta\beta0\nu$ Nuclear Matrix Element (NME), and m_e is the electron mass. The product $F_N^{0\nu} = G^{0\nu}|M^{0\nu}|^2$ includes all the nuclear details of the decay and it is usually referred to as *nuclear factor of merit*. While $G^{0\nu}$ can be calculated with reasonable accuracy, the NME value is strongly dependent on the nuclear model used for its evaluation so that discrepancies of about a factor 2-3 among the various theoretical calculations may be found [5–14]. Such uncertainties are of course reflected on the $|\langle m_{B\beta} \rangle|$ inferred values.

Experimentally one measures the energy deposited by the two electrons, which result in a peak centered at the $\beta\beta0\nu$ candidate isotope transition energy $(Q_{\beta\beta})$. For neutrino masses in the Inverted Hierarchy region, halflives in the range $10^{26} - 10^{27}$ y are expected for several $\beta\beta0\nu$ candidates [15]. This implies a few decays per 100 kg of candidate isotope per year. In a realistic experiment this faint signal must be singled out among the background events in the energy region around $Q_{\beta\beta}$. Hence, the sensitivity of a given experiment critically depends on the number of spurious counts in the region of interest: for a 68% confidence level, it is defined as the decay half-life corresponding to the maximum signal that could be hidden by a 1σ background fluctuation $n_B = \sqrt{BT\Delta M}$, where Δ is the FWHM energy resolution, M is the detector mass, T is the measuring time, and B is the background per unit mass, energy, and time. The experimental figure of merit is thus given by:

$$F^{0\nu} = T_{1/2}^{Back,Fluct.} = \frac{\ln 2N_{\beta\beta}T}{n_B} \propto \eta \epsilon \sqrt{\frac{MT}{B\Delta}}$$
 (2)

where $N_{\beta\beta}$ is the number of $\beta\beta0\nu$ decaying nuclei under observation, η is the $\beta\beta0\nu$ candidate isotopic abun-

dance and ϵ is the detection efficiency. This clearly shows that the sensitivity to the $\beta\beta0\nu$ signal goes linearly with the isotopic fraction and the detection efficiency, as the square root of mass and measuring time and, much worse – given the relation between half-life and $|\langle m_{\beta\beta}\rangle|$ – the improvements on the assessment of the neutrino mass have only a fourth root dependence on the experimental parameters.

2. The CUORE experiment

CUORE (Cryogenic Underground Observatory for Rare Events) is an array of 988 natTeO2 low temperature detectors (or bolometers) with the primary goal of searching for the $\beta\beta0\nu$ of ¹³⁰Te at Laboratori Nazionali del Gran Sasso (LNGS), Italy [16]. The array is a compact structure of 19 towers, each one containing 52 detectors arranged on 13 planes (see Figures 1 and 2). Every detector is a low temperature calorimeter, i.e. a very sensitive device operated at mK temperatures in which any energy release in the absorber is measured via its temperature rise through a suitable thermal sensor (the thermometer) [17]. In principle any dielectric and diamagnetic crystal can act as a particle absorber, therefore a wide choice of materials is made available by such technology. This characteristics together with the excellent energy resolution make low temperature calorimetry an ideal technique for experiments in the rare event physics field. In particular, TeO₂ crystals are very convenient low temperature absorbers because of their low specific heat, their intrinsic radiopurity and their crystallographic properties, which allow to grow excellent crystals of relatively large mass. Moreover, TeO₂ contains the isotope 130 Te, a quite appealing $\beta\beta0\nu$ candidate due to its large natural isotopic abundance (34%), its high transition energy ($Q_{\beta\beta}$ = 2528 keV, above most γ background lines from natural radioactivity), and its encouraging nuclear matrix element evaluations. In the CUORE array each detector is composed by a $5 \times 5 \times 5 \text{ cm}^3$ TeO₂ crystal (750 g of mass) on whose surface a Neutron Transmutation Doped (NTD) Ge thermistor is glued to measure any temperature change of the absorber and convert it into an electric signal. An electro-thermal link restores the operating temperature after every interaction. A Joule heater is also glued on the surface of the TeO₂ crystals to inject a known amount of energy in each detector for off-line thermal gain corrections (all detectors will be operated independently). The TeO₂ crystals are secured in groups of four inside low background copper frames by means of Teflon supports; each of the 19 CUORE towers is composed by 13 such elements. The total mass of the

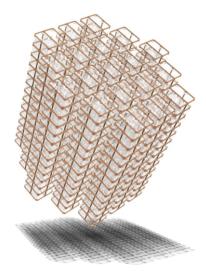


Figure 1: The CUORE detector array scheme.

CUORE array amounts to $741\,\mathrm{kg}$ of TeO_2 (206 kg of $^{130}\mathrm{Te}$).

The CUORE detectors will be installed in a complex cryogenic apparatus where they will be operated at a temperature of $\sim 10 \, \text{mK}$. The low temperature setup includes a custom cryostat (with six nested copper shields), a cryogen free cooling system with five pulse tubes, a very powerful dilution refrigerator, and a fast cooling system for pre-cooling. All materials used to build the experimental apparatus have been selected following stringent requirements concerning their radiopurity. The detector towers will be hanging on a thick copper disk thermally anchored to the 10 mK stage and mechanically decoupled from the rest of the setup to minimize vibrational noise. A set of cold roman lead [18] radioactive shields will screen the detectors from the radioactive contaminations of the cryogenic apparatus. A detector calibration system consisting of twelve thoriated wires, to be guided in proximity of the towers inside small copper tubes, has been designed to uniformily irradiate the detectors during the monthly energy calibrations with 232 Th γ rays. Finally, heavy room temperature shields made by layers of borated polyethylene, boric-acid powder, and lead bricks surround the apparatus to reduce the background contribution from neutrons and γ rays. The CUORE experimental site is located in Hall A of LNGS, under a rock overburden of ~ 3600 m of water equivalent, which reduces the muon flux to $\sim 3 \times 10^{-8} \,\mu/\text{s/cm}^2$ [19]. Presently CUORE is in the final phase of its construction (see Section 5) and it is expected to start taking data in 2015.

TeO₂ low temperature detectors (or bolometers) have

a long history in $\beta\beta0\nu$ searches of ¹³⁰Te [20, 21], the best sensitivity being achieved by Cuoricino, a ~ 40 kg tower of 62 bolometers (~ 11 kg of ¹³⁰Te) operated at LNGS site in the years 2003–2008: after an exposure of 19.75 kg y of ¹³⁰Te, Cuoricino set a lower limit on the decay half-life of this isotope of 2.8×10^{24} y (90% C.L.) [21]. To improve that result, as already discussed (see 2), one needs to increase the experimental mass, improve the energy resolution and, what's even more tricky, reduce the radioactive background contribution in the energy region of interest (ROI – centered on $Q_{\beta\beta}$).

CUORE design goals, besides the increase of the detector mass of a factor of ~ 20 with respect to Cuoricino, are a FWHM resolution of 5 keV at an energy around 2.5 MeV, and a background rate in the ROI of 10^{-2} counts/(keV·kg·y). This translates in a 1σ (90% C.L.) sensitivity of 1.6×10^{26} y (9.5 $\times 10^{25}$ y) in 5 y of measuring time, corresponding to an effective Majorana mass sensitivity in the range 50-130 meV.

3. The CUORE background reduction strategy

Cuoricino final sensitivity was limited by the background rate in the ROI, which was dominated by the surface radioactive contaminations of the detector materials – TeO₂ crystals and copper structures. Therefore, in preparing the CUORE array, several improvements have been implemented. First of all the detector design has been conceived to minimize the amount of passive materials close to the TeO₂ crystals: only selected radiopure copper and Teflon have been used for the tower skeleton and the total mass of the copper frame was reduced by a factor of 2.3 with respect to Cuoricino. Besides decreasing the copper surface background contribution to the ROI, this fact helps also optimizing the time coincidence analysis between the detectors (unlike background interactions, a $\beta\beta$ 0 ν event is fully contained inside a single $5 \times 5 \times 5 \text{ cm}^3$ TeO₂ crystal, with a calculated efficiency of 87%). To further mitigate the α background component a specially developed surface treatment protocol was applied to all copper parts directly facing the detectors. This procedure, called TECM (for Tumbling, Electropolishing, Chemical etching, and Magnetron plasma etching) has demonstrated to efficiently reduce the concentration of both ²³⁸U and ²³²Th on the copper surface to better than 1.3×10^{-7} Bq/cm² [22]. Moreover, to limit cosmogenic activation, all the copper material of the CUORE experiment has been stored underground at LNGS in between production steps.

Finally, a strict production protocol to limit bulk and

surface contamination of the TeO₂ crystals has been developed in collaboration with the crystal manufacturer at the Shangai Institute of Ceramics, Chinese Academy of Science [23]. After production, the TeO₂ crystals were transported to LNGS at sea level to minimize cosmogenic activation. Cryogenic tests on few crystals from different production batches demonstrated bulk and surface contamination levels of 6.7×10^{-7} Bq/kg and 8.9×10^{-9} Bq/cm² at 90% C.L., respectively, for 238 U, and of 8.4×10^{-7} Bq/kg and 2.0×10^{-9} Bq/cm² at 90% C.L., respectively, for 232 Th [24].

Special care was also devoted to the CUORE assembly line. A dedicated class 1000 clean room area was constructed at the first floor of the underground CUORE experimental building and is equipped with a set of specially designed glove boxes, inside which all operations needed to assemble a detector tower are performed under radon-free atmosphere. All materials used to construct the glove boxes were carefully selected, and all tools used inside the glove boxes, especially those that come into physical contact with the detector components, were cleaned and certified for their radiopurity. The clean area is divided into four rooms, each one devoted to specific tasks. In the gluing area thermistors and heaters are attached to the crystals by means of a semi-automated robotic system, to guarantee the uniformity and repeatability of the glue joints. In the assembly area, the instrumented crystals are secured inside the copper frames with the teflon spacers and then mechanically assembled into a tower. After that, a set of flexible printed circuit board cables is attached to the tower and in situ wire bonding is performed for the electrical connections of the detectors. The completed towers are then put up in the storage area inside specially designed nitrogen-flushed boxes to await the installation in the cryostat. The fourth room of the clean area provides direct access to the experimental area of the cryogenic setup, so as to always preserve the detectors in a protected environment. More details on the CUORE detector assembly can be found in [25].

In order to check the effectiveness of all the actions undertaken to improve Cuoricino results, we decided to run the CUORE-0 experiment. CUORE-0 is one CUORE-like tower built using the same detector components, cleaning protocols and assembly procedures defined for CUORE. CUORE-0 is therefore a proof of concept of the CUORE detector at all stages, including data acquisition and data analysis frameworks. Besides that, CUORE-0 is an independent $\beta\beta0\nu$ experiment on its own, which should extend the physics reach beyond Cuoricino while CUORE is being assembled.

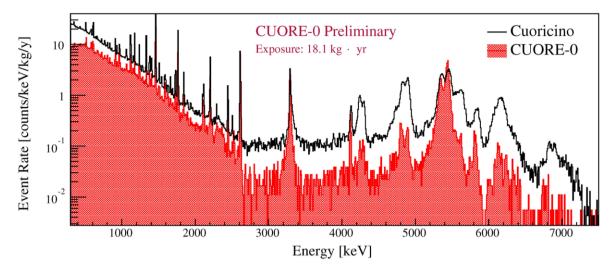


Figure 3: Cuoricino (line) and CUORE-0 (shaded) background spectra comparison. Only events with a single crystal hit are considered (anti-coincidence cut). The strong background reduction achieved with the CUORE-0 detector is evident. It amounts to a factor of \sim 6 in the α region (i.e. energies above 2.7 MeV) and to a factor of \sim 2 in the γ region (i.e. energies below 2.7 MeV).



Figure 2: The CUORE-0 tower before its installation in the cryostat.

4. CUORE-0

CUORE-0 is an array of $52^{\text{nat}}\text{TeO}_2$ detectors of $750\,\text{g}$ each, for a total mass of $39\,\text{kg}$ (~11 kg of ^{130}Te). It is operated in the same experimental setup that hosted Cuoricino, i.e. a 25 year-old cryostat located in Hall A of LNGS, with an external shielding made of lead and borated-polyethylene, and enclosed inside a Faraday cage [26]. Also the front-end electronics and data acquisition hardware are the same used for Cuoricino. The waveforms coming from the detectors are amplified, filtered with a six-pole active Bessel filter, and then fed

into a National Instrument PXI analog-to-digital converter (ADC) with 18-bit resolution and 125 S/s sampling frequency. All samples are continuously saved to disk and then processed with Apollo, the data acquisition software developed for CUORE, which identifies pulses with a constant fraction software trigger. In addition, a heater pulse is injected every $300 \, \mathrm{s}$ to each bolometer for stabilization monitoring, and baseline samples are recorded at $200 \, \mathrm{s}$ intervals to keep under control the working temperature and noise conditions of each detector. The signals have a typical time evolution of $\sim 5 \, \mathrm{s}$, with rise times in the range $40-80 \, \mathrm{ms}$ and decay times in the range $100-700 \, \mathrm{ms}$.

The energy calibration of the detectors is obtained by inserting two thoriated tungsten wires in between the external shielding and the cryostat. Usually, 2-3 days of acquisition time are enough to independently calibrate each detector in the energy region between 511 and 2615 keV with 232 Th γ rays. Physics (or low background) runs are usually 3-4 weeks long. A CUORE-0 data set is composed by an initial calibration, a low background run, and a final calibration. The raw data are then processed offline with Diana, the analysis software developed for CUORE. The analysis procedure is basically the same applied to the Cuoricino data [21], the main steps being amplitude evaluation, gain correction, energy calibration, and time coincidence analysis among the detectors. Especially this last point serves as active background rejection technique, which will be even more effective with the CUORE geometry: if any two or more detectors register signal pulses within 100 ms from each other, the events are tagged as coincident and discarded. In fact the coincidence events are mostly attributed to background interactions, such as Compton-scattered γ rays or α decays on the surface of nearby crystals, since the majority (87.4±1.1% [21]) of the $\beta\beta0\nu$ decay electrons are expected to be fully contained in a single crystal. Events are selected also on the basis of their shape: noisy pulses are discarded as a result of a comparison with a set of template signals. The last step is the production of the total energy spectra (calibration and background) by summing all the single detectors spectra for the entire acquired statistics.

CUORE-0 data taking started in March 2013, though in not ideal conditions: the first results (Phase I) were published in summer 2013 [27, 28]. Then a long stop was needed for cryostat maintainance. The data taking resumed in November 2013, with strongly improved performance, and is still going on (Phase II). The total spectrum obtained by summing all calibration spectra of all active channels (49 out of 52) acquired during Phase II shows a FWHM resolution at 2.615 MeV (²⁰⁸Tl line) equal to 4.8 keV, better than aimed for CUORE (see Section 2). The total background exposure sums up to 18.1 kg·y (6.2 kg·y of ¹³⁰Te) and covers the period from March 2013 until May 2014 (Phase I + Phase II). The overall detection efficiency is $77.6 \pm 1.3\%$, and includes the anti-coincidence cut, the pulse selection, and the containment efficiency. The resulting spectrum is shown in Figure 3. It is possible to recognize the γ lines of ²⁰⁸Tl, ⁴⁰K and ⁶⁰Co, attributed to a ²³²Th contamination of the cryostat, and the ones of ²¹⁴Bi, attributed to the presence of ²²²Rn in the air around the cryostat during the initial runs. For what concerns the peaks in the α region of the background spectrum, the main sources are bulk contaminations of the TeO₂ crystals, while other structures are attributed to surface contaminations of the TeO₂ crystals or of the passive materials close to them. The peak at ~3.2 MeV is attributed to ¹⁹⁰Pt, a bulk contamination of the TeO₂ crystals presumably originating by the platinum crucible during crystal growth. The continuum between 2.7 MeV and 3.9 MeV, excluding the ¹⁹⁰Pt peak, is attributed to α particles of degraded energy, i.e. particles that deposit a fraction of their energy in a crystal and the rest in an inactive material, like the copper structure (therefore no coincidence cut is possible). This continuum extends down into the gamma region and in Cuoricino was responsible for a significant fraction ($\sim 50 \pm 20\%$) of the background in the ROI. Figure 3 shows the tremendous improvement achieved with CUORE-0: the background in the α contin-

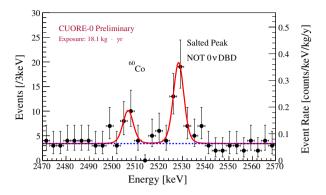


Figure 4: The CUORE-0 $\beta\beta$ 0 ν energy region of interest (ROI). The line at 2.505 MeV is the 60 Co sum peak while the one at 2.528 MeV is the salted peak (see text).

uum now amounts to 0.020 ± 0.001 counts/(keV·kg·y) - was 0.110 ± 0.001 counts/(keV·kg·y) in Cuoricino a factor ~6 improvement thanks to the background reduction strategy undertaken for CUORE. This translates in a background reduction in the ROI of a factor of ~2.5 with respect to Cuoricino: in this case the count rate still accounts for a significant contribution from the environmental γ s coming from the old cryostat bulk contaminations, consistent with what was observed in Cuoricino. A closer look to the energy spectrum in the ROI is shown in Figure 4: the backgound rate is $0.063 \pm 0.006 \text{ counts/(keV} \cdot \text{kg} \cdot \text{y})$ – was 0.153 ± 0.006 counts/(keV·kg·y) in Cuoricino. The peak at the $\beta\beta0\nu$ energy is a false one intentionally produced by our blinding procedure: a small blinded fraction of the events within $\pm 10 \,\text{keV}$ of the $^{208}\text{Tl} \, 2.615 \,\text{MeV}$ background line is exchanged with the events within $\pm 10 \,\mathrm{keV}$ of the $^{130}\mathrm{Te}~Q_{\beta\beta}$ value. The fit function (solid line of Fig. 4) is a flat background with two gaussians, one for the ⁶⁰Co 2.505 MeV sum peak and one for the salted peak. The data unblinding is expected in Spring 2015, when the CUORE-0 sensitivity should surpass the Cuoricino one. CUORE-0 data taking will instead continue until CUORE starts.

5. CUORE status and prospects

CUORE is currently at an advanced state of construction. The detector is fully assembled: the 19 towers are ready and safely stored inside nitrogen flushed canisters, while waiting to be installed inside the cryostat.

Major efforts are now devoted to the last phases of the cryostat preparation. Given the complexity of the system a phased installation was chosen, with a limited number of items being added at each step.

The commissioning plan started in 2012. Currently, the cryogenic system is fully assembled and the first successfull cooldown was recently completed, showing a base temperature (~7 mK) safely below the design value of 10 mK. Next steps foresee the installation of the calibration system, of the wiring system, of the cold shields, and of the detector suspension plate. The final step will be the tower installation.

The CUORE experiment data taking is expected to start at the end of 2015. The encouraging results obtained with CUORE-0 seem to indicate that the CUORE design goal of 0.01 counts/(keV·kg·y) is within reach.

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References

- [1] S. Bilenky and C. Giunti, Mod. Phys. Lett. A 27 (2012) 1230015
- [2] S.R. Elliott and P. Vogel, Ann. Rev. Nucl. Part. Sci. 52 (2002) 115.
- [3] F.T. Avignone III et al., Rev. Mod. Phys. 80 (2008) 481.
- [4] J. Beringer et al. (Particle Data Group), Phys. Rev. D86 (2012) 010001
- [5] J. Suhonen and O. Civitarese, Phys. Rept. 300 (1998) 123.
- [6] F. Simkovic et al., Phys. Rev. C60 (1999) 055502.
- [7] J.Kotila and F. Iachello, Phys. Rev. C85 (2012) 034316.
- [8] J. Menendez et al., Nucl. Phys. A 818 (2009) 139.
- [9] A. Faessler et al., JoP G: Nucl. Part. Phys. **39** (2012) 124006.
- [10] D. Fang et al., Phys. Rev. C83 (2011) 034320.
- [11] J. Suhonen and O. Civitarese, JoP G: Nucl. Part. Phys. 39 (2012) 124005.
- [12] J. Barea et al., Phys. Rev. C87 (2013) 014315.
- [13] P.K. Rath et al., Phys. Rev. C82 (2010) 064310.
- [14] T.R. Rodriguez et al., Phys. Rev. Lett. 105 (2010) 252503.
- [15] D.R. Artusa et al., Eur. Phys. J. C 74 (2014) 3096.

- [16] C. Arnaboldi et al., Nucl. Inst. Meth. A 518 (2004) 775.
- [17] E. Fiorini and T. Niinikowski, Nucl. Inst. Meth. A 224 (1984) 83
- [18] A. Alessandrello et al., Nucl. Inst. Meth. B 142 (1998) 163.
- [19] D. Mei and A. Hime, Phys. Rev. D 73 (2006) 053004.
- [20] C. Arnaboldi et al., Phys. Rev. C 78 (2008) 035502.
- [21] E. Andreotti et al., Astropart. Phys. 34 (2011) 822.
- [22] F. Alessandria et al., Astropart. Phys. 45 (2013) 13.
- [23] C. Arnaboldi et al., J. Cryst. Growth 312 (2010) 2999.
- [24] F. Alessandria et al., Astropart. Phys. 35 (2012) 839.[25] D.R. Artusa et al., accepted for publication by AHEP.
- arXiv:1402.6072. [26] C. Arnaboldi et al., Phys. Lett. B 557 (2003) 167.
- [27] M. Vignati et al., TAUP 2013 proceedings, to be published by Elsevier B.V.
- [28] D.R. Artusa et al., Eur. Phys. J. C 74 (2014) 2956.