



## Erratum

Erratum to “Rapidity and transverse momentum dependence of inclusive  $J/\psi$  production in pp collisions at  $\sqrt{s} = 7$  TeV” [Phys. Lett. B 704 (5) (2011) 442] <sup>☆</sup>

ALICE Collaboration

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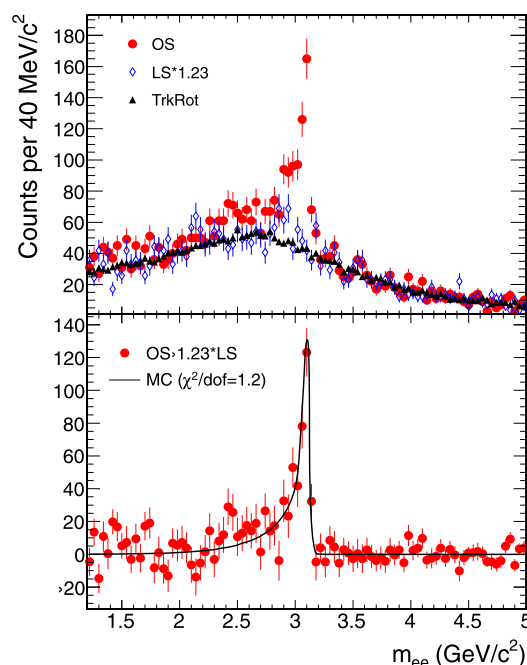
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We have identified a bias in the calculation of the cross section of  $J/\psi$  production measured in the  $e^+e^-$  channel at central rapidity ( $|y| < 0.9$ ) [1]. The acceptance and efficiency corrections were evaluated using a Monte Carlo simulation, based on PYTHIA [2,3], which did not include the radiative decays of  $J/\psi$  ( $J/\psi \rightarrow e^+e^- \gamma$ ). In these decays the reconstructed dielectron invariant mass,  $m_{e^+e^-}$ , is biased towards smaller values than the nominal  $J/\psi$  mass, since the photon contribution is neglected. As a consequence, the fraction of signal events in the  $m_{e^+e^-}$  range 2.92–3.16  $\text{GeV}/c^2$  was overestimated by about 10%. Moreover, the requirement of the transverse momentum of the daughter electrons being larger than 1  $\text{GeV}/c$  is more selective for radiative than non-radiative  $J/\psi$  decays. Therefore, the  $J/\psi$  acceptance was also overestimated by about 5%.

We have now evaluated the acceptance and efficiency corrections with a simulation where the decay of the  $J/\psi$  particles is handled by the EvtGen package [4], and where the final state radiation is described using PHOTOS [5,6]. The new acceptance times efficiency value ( $A \times \epsilon$ ) after all analysis cuts is a factor 1.155 smaller than that previously evaluated, independently of  $p_T$ . Neglecting the effect of radiative decays therefore results in underestimating both the  $p_T$ -integrated and the differential cross sections by 15.5%.

For the dimuon channel, where no invariant mass cut is applied and the occurrence of final state radiation is reduced (by about a factor of three [7]), the differences in the  $A \times \epsilon$  values obtained with the new and previous simulations are about 1–2%, well within the systematic uncertainty associated with the signal extraction.

We have further verified that in the dielectron channel the procedure used to derive the  $p_T$ -differential cross section, which is based on the computation of the  $A \times \epsilon$  values, produces a result fully compatible with that obtained by applying an unfolding cor-



**Fig. 1.** Top panel: invariant mass distributions for opposite-sign (OS) and like-sign (LS) electron pairs ( $|y| < 0.9$ , all  $p_T$ ), as well as for pairs obtained with one track randomly rotated (TrkRot). Bottom panel: the difference of the OS and LS distributions with a fit to the Monte Carlo (MC) signal superimposed.

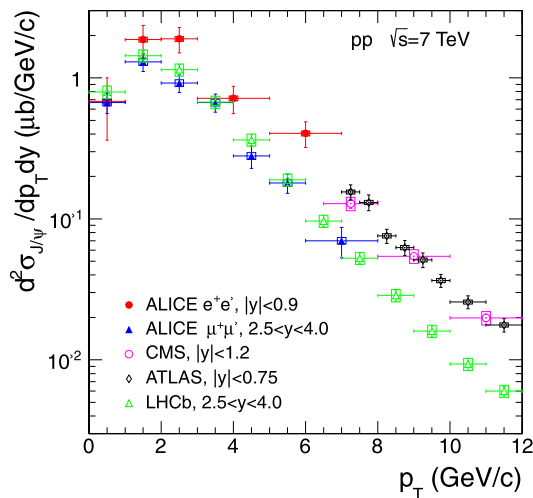
rection procedure (see, e.g., review [8]), even when considering the radiative decays.

In Fig. 1 the invariant mass distributions of electron pairs are shown. In particular, in the bottom panel the new Monte Carlo line shape is superimposed on the difference of the opposite and like sign distributions. The fraction of the signal within the invariant mass range 2.92–3.16  $\text{GeV}/c^2$  estimated using this Monte Carlo

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**Table 1**Summary of the results on the  $J/\psi$  differential cross sections.

$p_T$ (GeV/c)	$N_{J/\psi}$	$A \times \epsilon$	$d^2\sigma_{J/\psi}/dp_T dy$ ( $\mu\text{b}/(\text{GeV}/c)$ )	Systematic uncertainties			
				Correl. ( $\mu\text{b}/(\text{GeV}/c)$ )	Non-correl. ( $\mu\text{b}/(\text{GeV}/c)$ )	Polariz., CS ( $\mu\text{b}/(\text{GeV}/c)$ )	Polariz., HE ( $\mu\text{b}/(\text{GeV}/c)$ )
$ y  < 0.9$							
[0; 1]	$50 \pm 17$	0.122	$0.68 \pm 0.24$	0.02	0.21	+0.16, -0.18	+0.08, -0.12
[1; 2]	$86 \pm 17$	0.076	$1.87 \pm 0.37$	0.07	0.31	+0.42, -0.50	+0.28, -0.39
[2; 3]	$79 \pm 13$	0.069	$1.89 \pm 0.31$	0.08	0.23	+0.33, -0.43	+0.35, -0.44
[3; 5]	$75 \pm 13$	0.086	$0.72 \pm 0.13$	0.02	0.09	+0.06, -0.08	+0.16, -0.13
[5; 7]	$50 \pm 9$	0.104	$0.40 \pm 0.07$	0.01	0.05	+0.001, -0.005	+0.06, -0.08
$y$							
			$d\sigma_{J/\psi}/dy$ ( $\mu\text{b}$ )	( $\mu\text{b}$ )	( $\mu\text{b}$ )	( $\mu\text{b}$ )	( $\mu\text{b}$ )
[-0.9; 0.9]	$352 \pm 32$	0.085	$6.90 \pm 0.62$	0.28	0.96	+0.9, -1.3	+1.0, -1.5



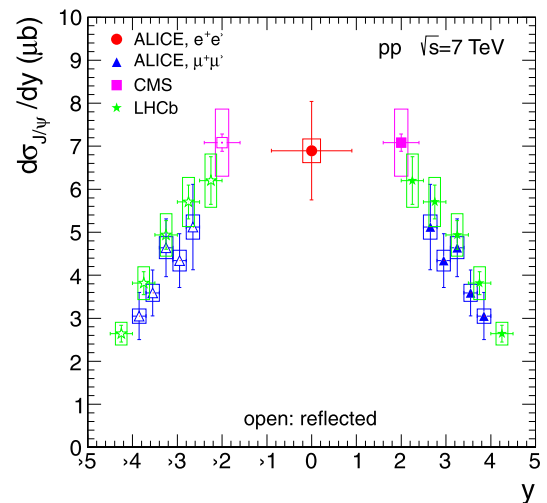
**Fig. 2.** Double differential  $J/\psi$  production cross section as a function of  $p_T$  for the midrapidity range and for the forward rapidity data, compared with results from the other LHC experiments [10–12], obtained in similar rapidity ranges. The error bars represent the quadratic sum of the statistical and systematic errors, while the systematic uncertainties on luminosity are shown as boxes. The symbols are plotted at the center of each bin.

is  $66.8 \pm 1.9\%$ . The main contribution to the uncertainty on this quantity comes from the accuracy of the description of the detector material, as discussed in [1]. A smaller contribution (1%, in terms of the relative error) is attributed to the small discrepancies between the invariant mass distribution as provided by QED at next to leading order [9] and by the event generator (EvtGen + PHOTOS); the latter contribution remains even after taking into account the detector resolution.

The corrected value of the production cross section is  $\sigma_{J/\psi}(|y| < 0.9) = 12.4 \pm 1.1$  (stat.)  $\pm 1.8$  (syst.)  $+ 1.8$  ( $\lambda_{\text{HE}} = 1$ )  $- 2.7$  ( $\lambda_{\text{HE}} = -1$ )  $\mu\text{b}$ . In Table 1 the resulting differential cross sections are summarized. Finally, in Figs. 2 and 3 we have updated accordingly the ALICE data points at central rapidity.

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**Fig. 3.**  $J/\psi$  cross section as a function of rapidity, compared with results from the other LHC experiments [10–12]. The error bars represent the quadratic sum of the statistical and systematic errors, while the systematic uncertainties on luminosity are shown as boxes. The symbols are plotted at the center of each bin.

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### ALICE Collaboration

K. Aamodt<sup>1</sup>, A. Abrahantes Quintana<sup>2</sup>, D. Adamová<sup>3</sup>, A.M. Adare<sup>4</sup>, M.M. Aggarwal<sup>5</sup>, G. Aglieri Rinella<sup>6</sup>, A.G. Agocs<sup>7</sup>, A. Agostinelli<sup>8</sup>, S. Aguilar Salazar<sup>9</sup>, Z. Ahammed<sup>10</sup>, N. Ahmad<sup>11</sup>, A. Ahmad Masoodi<sup>11</sup>, S.U. Ahn<sup>12,i</sup>, A. Akindinov<sup>13</sup>, D. Aleksandrov<sup>14</sup>, B. Alessandro<sup>15</sup>, R. Alfaro Molina<sup>9</sup>, A. Alici<sup>16</sup>,

A. Alkin<sup>17</sup>, E. Almaráz Aviña<sup>9</sup>, J. Alme<sup>18</sup>, T. Alt<sup>19</sup>, V. Altini<sup>20,ii</sup>, I. Altsybeev<sup>21</sup>, C. Andrei<sup>22</sup>,  
 A. Andronic<sup>23</sup>, V. Anguelov<sup>19,iii</sup>, C. Anson<sup>24</sup>, T. Antičić<sup>25</sup>, F. Antinori<sup>26</sup>, P. Antonioli<sup>27</sup>, L. Aphecetche<sup>28</sup>,  
 H. Appelshäuser<sup>29</sup>, N. Arbor<sup>30</sup>, S. Arcelli<sup>8</sup>, A. Arend<sup>29</sup>, N. Armesto<sup>31</sup>, R. Arnaldi<sup>15</sup>, T. Aronsson<sup>4</sup>,  
 I.C. Arsene<sup>23</sup>, A. Asryan<sup>21</sup>, A. Augustinus<sup>6</sup>, R. Averbeck<sup>23</sup>, T.C. Awes<sup>32</sup>, J. Åystö<sup>33</sup>, M.D. Azmi<sup>11</sup>,  
 M. Bach<sup>19</sup>, A. Badalà<sup>34</sup>, Y.W. Baek<sup>12,i</sup>, R. Bailhache<sup>29</sup>, R. Bala<sup>15</sup>, R. Baldini Ferroli<sup>16</sup>, A. Baldisseri<sup>35</sup>,  
 A. Baldit<sup>36</sup>, J. Bán<sup>37</sup>, R. Barbera<sup>38</sup>, F. Barile<sup>20</sup>, G.G. Barnaföldi<sup>7</sup>, L.S. Barnby<sup>39</sup>, V. Barret<sup>36</sup>, J. Bartke<sup>40</sup>,  
 M. Basile<sup>8</sup>, N. Bastid<sup>36</sup>, B. Bathen<sup>41</sup>, G. Batigne<sup>28</sup>, B. Batyunya<sup>42</sup>, C. Baumann<sup>29</sup>, I.G. Bearden<sup>43</sup>,  
 H. Beck<sup>29</sup>, I. Belikov<sup>44</sup>, F. Bellini<sup>8</sup>, R. Bellwied<sup>45</sup>, E. Belmont-Moreno<sup>9</sup>, S. Beole<sup>46</sup>, I. Berceau<sup>22</sup>,  
 A. Bercuci<sup>22</sup>, E. Berdermann<sup>23</sup>, Y. Berdnikov<sup>47</sup>, C. Bergmann<sup>41</sup>, L. Betev<sup>6</sup>, A. Bhasin<sup>48</sup>, A.K. Bhati<sup>5</sup>,  
 L. Bianchi<sup>46</sup>, N. Bianchi<sup>49</sup>, C. Bianchin<sup>50</sup>, J. Bielčik<sup>51</sup>, J. Bielčíková<sup>3</sup>, A. Bilandzic<sup>52</sup>, E. Biolcati<sup>46</sup>,  
 A. Blanc<sup>36</sup>, F. Blanco<sup>53</sup>, F. Blanco<sup>45</sup>, D. Blau<sup>14</sup>, C. Blume<sup>29</sup>, N. Bock<sup>24</sup>, A. Bogdanov<sup>54</sup>, H. Bøggild<sup>43</sup>,  
 M. Bogolyubsky<sup>55</sup>, L. Boldizsár<sup>7</sup>, M. Bombara<sup>56</sup>, C. Bombonati<sup>50</sup>, J. Book<sup>29</sup>, H. Borel<sup>35</sup>, A. Borissov<sup>57</sup>,  
 C. Bortolin<sup>50,iv</sup>, S. Bose<sup>58</sup>, F. Bossú<sup>6,v</sup>, M. Botje<sup>52</sup>, S. Böttger<sup>59</sup>, B. Boyer<sup>60</sup>, P. Braun-Munzinger<sup>23</sup>,  
 L. Bravina<sup>61</sup>, M. Bregant<sup>28</sup>, T. Breitner<sup>59</sup>, M. Broz<sup>62</sup>, R. Brun<sup>6</sup>, E. Bruna<sup>4</sup>, G.E. Bruno<sup>20,\*</sup>, D. Budnikov<sup>63</sup>,  
 H. Buesching<sup>29</sup>, S. Bufalino<sup>46</sup>, O. Busch<sup>64</sup>, Z. Buthelezi<sup>65</sup>, D. Caffarri<sup>50</sup>, X. Cai<sup>66</sup>, H. Caines<sup>4</sup>,  
 E. Calvo Villar<sup>67</sup>, P. Camerini<sup>68</sup>, V. Canoa Roman<sup>69,vi</sup>, G. Cara Romeo<sup>27</sup>, F. Carena<sup>6</sup>, W. Carena<sup>6</sup>,  
 F. Carminati<sup>6</sup>, A. Casanova Díaz<sup>49</sup>, M. Caselle<sup>6</sup>, J. Castillo Castellanos<sup>35</sup>, V. Catanescu<sup>22</sup>, C. Cavicchioli<sup>6</sup>,  
 J. Cepila<sup>51</sup>, P. Cerello<sup>15</sup>, B. Chang<sup>33</sup>, S. Chapeland<sup>6</sup>, J.L. Charvet<sup>35</sup>, S. Chattopadhyay<sup>58</sup>,  
 S. Chattopadhyay<sup>10</sup>, M. Cherney<sup>70</sup>, C. Cheshkov<sup>71</sup>, B. Cheynis<sup>71</sup>, E. Chiavassa<sup>46</sup>, V. Chibante Barroso<sup>6</sup>,  
 D.D. Chinellato<sup>72</sup>, P. Chochula<sup>6</sup>, M. Chojnacki<sup>73</sup>, P. Christakoglou<sup>73</sup>, C.H. Christensen<sup>43</sup>,  
 P. Christiansen<sup>74</sup>, T. Chujo<sup>75</sup>, C. Cicalo<sup>76</sup>, L. Cifarelli<sup>8,ii</sup>, F. Cindolo<sup>27</sup>, J. Cleymans<sup>65</sup>, F. Coccetti<sup>16</sup>,  
 J.-P. Coffin<sup>44</sup>, G. Conesa Balbastre<sup>30</sup>, Z. Conesa del Valle<sup>44,ii</sup>, P. Constantin<sup>64</sup>, G. Contin<sup>68</sup>,  
 J.G. Contreras<sup>69</sup>, T.M. Cormier<sup>57</sup>, Y. Corrales Morales<sup>46</sup>, I. Cortés Maldonado<sup>77</sup>, P. Cortese<sup>78</sup>,  
 M.R. Cosentino<sup>72</sup>, F. Costa<sup>6</sup>, M.E. Cotallo<sup>53</sup>, E. Crescio<sup>69</sup>, P. Crochet<sup>36</sup>, E. Cuautle<sup>79</sup>, L. Cunqueiro<sup>49</sup>,  
 G. D'Erasmus<sup>20</sup>, A. Dainese<sup>26</sup>, H.H. Dalsgaard<sup>43</sup>, A. Danu<sup>80</sup>, D. Das<sup>58</sup>, I. Das<sup>58</sup>, A. Dash<sup>81</sup>, S. Dash<sup>15</sup>,  
 S. De<sup>10</sup>, A. De Azevedo Moregula<sup>49</sup>, G.O.V. de Barros<sup>82</sup>, A. De Caro<sup>83</sup>, G. de Cataldo<sup>84</sup>, J. de Cuveland<sup>19</sup>,  
 A. De Falco<sup>85</sup>, D. De Gruttola<sup>83</sup>, N. De Marco<sup>15</sup>, S. De Pasquale<sup>83</sup>, R. de Rooij<sup>73</sup>, E. Del Castillo  
 Sanchez<sup>6</sup>, H. Delagrange<sup>28</sup>, Y. Delgado Mercado<sup>67</sup>, G. Dellacasa<sup>78,vii</sup>, A. Deloff<sup>86</sup>, V. Demanov<sup>63</sup>,  
 E. Dénes<sup>7</sup>, A. Deppman<sup>82</sup>, D. Di Bari<sup>20</sup>, C. Di Giglio<sup>20</sup>, S. Di Liberto<sup>87</sup>, A. Di Mauro<sup>6</sup>, P. Di Nezza<sup>49</sup>,  
 T. Dietel<sup>41</sup>, R. Divià<sup>6</sup>, Ø. Djuvsland<sup>1</sup>, A. Dobrin<sup>57</sup>, T. Dobrowolski<sup>86</sup>, I. Domínguez<sup>79</sup>, B. Dönigus<sup>23</sup>,  
 O. Dordic<sup>61</sup>, O. Driga<sup>28</sup>, A.K. Dubey<sup>10</sup>, L. Ducroux<sup>71</sup>, P. Dupieux<sup>36</sup>, A.K. Dutta Majumdar<sup>58</sup>, M.R. Dutta  
 Majumdar<sup>10</sup>, D. Elia<sup>84</sup>, D. Emschermann<sup>41</sup>, H. Engel<sup>59</sup>, H.A. Erdal<sup>18</sup>, B. Espagnon<sup>60</sup>, M. Estienne<sup>28</sup>,  
 S. Esumi<sup>75</sup>, D. Evans<sup>39</sup>, S. Evrard<sup>6</sup>, G. Eyyubova<sup>61</sup>, D. Fabris<sup>26</sup>, J. Faivre<sup>30</sup>, D. Falchieri<sup>8</sup>, A. Fantoni<sup>49</sup>,  
 M. Fasel<sup>23</sup>, R. Fearick<sup>65</sup>, A. Fedunov<sup>42</sup>, D. Fehlker<sup>1</sup>, V. Fekete<sup>62</sup>, D. Felea<sup>80</sup>, G. Feofilov<sup>21</sup>, A. Fernández  
 Téllez<sup>77</sup>, E.G. Ferreira<sup>31</sup>, A. Ferretti<sup>46</sup>, R. Ferretti<sup>78</sup>, M.A.S. Figueredo<sup>82</sup>, S. Filchagin<sup>63</sup>, R. Fini<sup>84</sup>,  
 D. Finogeev<sup>88</sup>, F.M. Fionda<sup>20</sup>, E.M. Fiore<sup>20</sup>, M. Floris<sup>6</sup>, S. Foertsch<sup>65</sup>, P. Foka<sup>23</sup>, S. Fokin<sup>14</sup>,  
 E. Fragiaco<sup>89</sup>, M. Fragkiadakis<sup>90</sup>, U. Frankenfeld<sup>23</sup>, U. Fuchs<sup>6</sup>, F. Furano<sup>6</sup>, C. Furget<sup>30</sup>, M. Fusco  
 Girard<sup>83</sup>, J.J. Gaardhøje<sup>43</sup>, S. Gadrat<sup>30</sup>, M. Gagliardi<sup>46</sup>, A. Gago<sup>67</sup>, M. Gallio<sup>46</sup>, P. Ganoti<sup>32</sup>,  
 C. Garabatos<sup>23</sup>, E. Garcia-Solis<sup>91</sup>, R. Gemme<sup>78</sup>, J. Gerhard<sup>19</sup>, M. Germain<sup>28</sup>, C. Geuna<sup>35</sup>, A. Gheata<sup>6</sup>,  
 M. Gheata<sup>6</sup>, B. Ghidini<sup>20</sup>, P. Ghosh<sup>10</sup>, P. Gianotti<sup>49</sup>, M.R. Girard<sup>92</sup>, P. Giubellino<sup>46,viii</sup>,  
 E. Gladysz-Dziadus<sup>40</sup>, P. Glässel<sup>64</sup>, R. Gomez<sup>93</sup>, L.H. González-Trueba<sup>9</sup>, P. González-Zamora<sup>53</sup>,  
 S. Gorbunov<sup>19</sup>, S. Gotovac<sup>94</sup>, V. Grabski<sup>9</sup>, L.K. Graczykowski<sup>92</sup>, R. Grajcarek<sup>64</sup>, A. Grelli<sup>73</sup>, A. Grigoras<sup>6</sup>,  
 C. Grigoras<sup>6</sup>, V. Grigoriev<sup>54</sup>, A. Grigoryan<sup>95</sup>, S. Grigoryan<sup>42</sup>, B. Grinyov<sup>17</sup>, N. Grion<sup>89</sup>, P. Gros<sup>74</sup>,  
 J.F. Grosse-Oetringhaus<sup>6</sup>, J.-Y. Grossiord<sup>71</sup>, F. Guber<sup>88</sup>, R. Guernane<sup>30</sup>, C. Guerra Gutierrez<sup>67</sup>,  
 B. Guerzoni<sup>8</sup>, K. Gulbrandsen<sup>43</sup>, H. Gulkanyan<sup>95</sup>, T. Gunji<sup>96</sup>, A. Gupta<sup>48</sup>, R. Gupta<sup>48</sup>, H. Gutbrod<sup>23</sup>,  
 Ø. Haaland<sup>1</sup>, C. Hadjidakis<sup>60</sup>, M. Haiduc<sup>80</sup>, H. Hamagaki<sup>96</sup>, G. Hamar<sup>7</sup>, L.D. Hanratty<sup>39</sup>,  
 Z. Harmanova<sup>56</sup>, J.W. Harris<sup>4</sup>, M. Hartig<sup>29</sup>, D. Hasegan<sup>80</sup>, D. Hatzifotiadou<sup>27</sup>, A. Hayrapetyan<sup>95,ii</sup>,  
 M. Heide<sup>41</sup>, M. Heinz<sup>4</sup>, H. Helstrup<sup>18</sup>, A. Herghelegiu<sup>22</sup>, G. Herrera Corral<sup>69</sup>, N. Herrmann<sup>64</sup>,  
 K.F. Hetland<sup>18</sup>, B. Hicks<sup>4</sup>, P.T. Hille<sup>4</sup>, B. Hippolyte<sup>44</sup>, T. Horaguchi<sup>75</sup>, Y. Hori<sup>96</sup>, P. Hristov<sup>6</sup>,  
 I. Hřivnáčová<sup>60</sup>, M. Huang<sup>1</sup>, S. Huber<sup>23</sup>, T.J. Humanic<sup>24</sup>, D.S. Hwang<sup>97</sup>, R. Ilkaev<sup>63</sup>, I. Ilkiv<sup>86</sup>,  
 M. Inaba<sup>75</sup>, E. Incani<sup>85</sup>, G.M. Innocenti<sup>46</sup>, M. Ippolitov<sup>14</sup>, M. Irfan<sup>11</sup>, C. Ivan<sup>23</sup>, A. Ivanov<sup>21</sup>,  
 M. Ivanov<sup>23</sup>, V. Ivanov<sup>47</sup>, A. Jachołkowski<sup>6</sup>, P.M. Jacobs<sup>98</sup>, L. Jancurová<sup>42</sup>, S. Jangal<sup>44</sup>, M.A. Janik<sup>92</sup>,

R. Janik<sup>62</sup>, P.H.S.Y. Jayarathna<sup>45,ix</sup>, S. Jena<sup>99</sup>, L. Jirde<sup>6</sup>, G.T. Jones<sup>39</sup>, P.G. Jones<sup>39</sup>, P. Jovanović<sup>39</sup>, H. Jung<sup>12</sup>, W. Jung<sup>12</sup>, A. Jusko<sup>39</sup>, S. Kalcher<sup>19</sup>, P. Kaliňák<sup>37</sup>, M. Kalisky<sup>41</sup>, T. Kalliokoski<sup>33</sup>, A. Kalweit<sup>100</sup>, R. Kamermans<sup>73,vii</sup>, K. Kanaki<sup>1</sup>, E. Kang<sup>12</sup>, J.H. Kang<sup>101</sup>, V. Kaplin<sup>54</sup>, A. Karasu Uysal<sup>6</sup>, O. Karavichev<sup>88</sup>, T. Karavicheva<sup>88</sup>, E. Karpechev<sup>88</sup>, A. Kazantsev<sup>14</sup>, U. Kebschull<sup>59</sup>, R. Keidel<sup>102</sup>, M.M. Khan<sup>11</sup>, P. Khan<sup>58</sup>, A. Khanzadeev<sup>47</sup>, Y. Kharlov<sup>55</sup>, B. Kileng<sup>18</sup>, D.J. Kim<sup>33</sup>, D.S. Kim<sup>12</sup>, D.W. Kim<sup>12</sup>, J.H. Kim<sup>97</sup>, J.S. Kim<sup>12</sup>, M. Kim<sup>101</sup>, S. Kim<sup>97</sup>, S.H. Kim<sup>12</sup>, S. Kirsch<sup>6,x</sup>, I. Kisel<sup>19</sup>, S. Kiselev<sup>13</sup>, A. Kisiel<sup>6</sup>, J.L. Klay<sup>103</sup>, J. Klein<sup>64</sup>, C. Klein-Bösing<sup>41</sup>, M. Kliemant<sup>29</sup>, A. Kluge<sup>6</sup>, M.L. Knichel<sup>23</sup>, K. Koch<sup>64</sup>, M.K. Köhler<sup>23</sup>, A. Kolojvari<sup>21</sup>, V. Kondratiev<sup>21</sup>, N. Kondratyeva<sup>54</sup>, A. Konevskih<sup>88</sup>, E. Kornaš<sup>40</sup>, C. Kottachchi Kankanamge Don<sup>57</sup>, R. Kour<sup>39</sup>, M. Kowalski<sup>40</sup>, S. Kox<sup>30</sup>, G. Koyithatta Meethalevedu<sup>99</sup>, K. Kozlov<sup>14</sup>, J. Kral<sup>33</sup>, I. Králik<sup>37</sup>, F. Kramer<sup>29</sup>, I. Kraus<sup>23</sup>, T. Krawutschke<sup>64,xi</sup>, M. Kretz<sup>19</sup>, M. Krivda<sup>39,xii</sup>, F. Krizek<sup>33</sup>, M. Krus<sup>51</sup>, E. Kryshen<sup>47</sup>, M. Krzewicki<sup>52</sup>, Y. Kucheriaev<sup>14</sup>, C. Kuhn<sup>44</sup>, P.G. Kuijter<sup>52</sup>, P. Kurashvili<sup>86</sup>, A. Kurepin<sup>88</sup>, A.B. Kurepin<sup>88</sup>, A. Kuryakin<sup>63</sup>, S. Kushpil<sup>3</sup>, V. Kushpil<sup>3</sup>, M.J. Kweon<sup>64</sup>, Y. Kwon<sup>101</sup>, P. La Rocca<sup>38</sup>, P. Ladrón de Guevara<sup>79</sup>, V. Lafage<sup>60</sup>, I. Lakomov<sup>21</sup>, C. Lara<sup>59</sup>, D.T. Larsen<sup>1</sup>, C. Lazzeroni<sup>39</sup>, Y. Le Bornec<sup>60</sup>, R. Lea<sup>68</sup>, M. Lechman<sup>6</sup>, K.S. Lee<sup>12</sup>, S.C. Lee<sup>12</sup>, F. Lefèvre<sup>28</sup>, J. Lehnert<sup>29</sup>, L. Leistam<sup>6</sup>, M. Lenhardt<sup>28</sup>, V. Lenti<sup>84</sup>, I. León Monzón<sup>93</sup>, H. León Vargas<sup>29</sup>, P. Lévai<sup>7</sup>, X. Li<sup>104</sup>, R. Lietava<sup>39</sup>, S. Lindal<sup>61</sup>, V. Lindenstruth<sup>19</sup>, C. Lippmann<sup>23</sup>, M.A. Lisa<sup>24</sup>, L. Liu<sup>1</sup>, V.R. Loggins<sup>57</sup>, V. Loginov<sup>54</sup>, S. Lohn<sup>6</sup>, D. Lohner<sup>64</sup>, C. Loizides<sup>98</sup>, K.K. Loo<sup>33</sup>, X. Lopez<sup>36</sup>, M. López Noriega<sup>60</sup>, E. López Torres<sup>2</sup>, G. Løvhøiden<sup>61</sup>, X.-G. Lu<sup>64</sup>, P. Luettig<sup>29</sup>, M. Lunardon<sup>50</sup>, G. Luparello<sup>46</sup>, L. Luquin<sup>28</sup>, C. Luzzi<sup>6</sup>, K. Ma<sup>66</sup>, R. Ma<sup>4</sup>, D.M. Madagodahettige-Don<sup>45</sup>, A. Maevskaya<sup>88</sup>, M. Mager<sup>6</sup>, D.P. Mahapatra<sup>81</sup>, A. Maire<sup>44</sup>, M. Malaev<sup>47</sup>, I. Maldonado Cervantes<sup>79</sup>, D. Mal'Kevich<sup>13</sup>, P. Malzacher<sup>23</sup>, A. Mamonov<sup>63</sup>, L. Manceau<sup>36</sup>, V. Manko<sup>14</sup>, F. Manso<sup>36</sup>, V. Manzari<sup>84</sup>, Y. Mao<sup>66,xiii</sup>, M. Marchisone<sup>46</sup>, J. Mareš<sup>105</sup>, G.V. Margagliotti<sup>68</sup>, A. Margotti<sup>27</sup>, A. Marín<sup>23</sup>, C. Markert<sup>106</sup>, I. Martashvili<sup>107</sup>, P. Martinengo<sup>6</sup>, M.I. Martínez<sup>77</sup>, A. Martínez Davalos<sup>9</sup>, G. Martínez García<sup>28</sup>, Y. Martynov<sup>17</sup>, A. Mas<sup>28</sup>, S. Masciocchi<sup>23</sup>, M. Maserà<sup>46</sup>, A. Masoni<sup>76</sup>, L. Massacrier<sup>71</sup>, M. Mastromarco<sup>84</sup>, A. Mastroserio<sup>6</sup>, Z.L. Matthews<sup>39</sup>, A. Matyja<sup>40</sup>, D. Mayani<sup>79</sup>, M.A. Mazzone<sup>87</sup>, F. Meddi<sup>108</sup>, A. Menchaca-Rocha<sup>9</sup>, P. Mendez Lorenzo<sup>6</sup>, J. Mercado Pérez<sup>64</sup>, M. Meres<sup>62</sup>, Y. Miake<sup>75</sup>, J. Midori<sup>109</sup>, L. Milano<sup>46</sup>, J. Milosevic<sup>61,xiv</sup>, A. Mischke<sup>73</sup>, D. Miśkowiec<sup>6,xv</sup>, C. Mitu<sup>80</sup>, J. Mlynarz<sup>57</sup>, B. Mohanty<sup>10</sup>, L. Molnar<sup>6</sup>, L. Montaño Zetina<sup>69</sup>, M. Monteno<sup>15</sup>, E. Montes<sup>53</sup>, M. Morando<sup>50</sup>, D.A. Moreira De Godoy<sup>82</sup>, S. Moretto<sup>50</sup>, A. Morsch<sup>6</sup>, V. Muccifora<sup>49</sup>, E. Mudnic<sup>94</sup>, H. Müller<sup>6</sup>, S. Muhuri<sup>10</sup>, M.G. Munhoz<sup>82</sup>, L. Musa<sup>6</sup>, A. Musso<sup>15</sup>, B.K. Nandi<sup>99</sup>, R. Nania<sup>27</sup>, E. Nappi<sup>84</sup>, C. Nattrass<sup>107</sup>, F. Navach<sup>20</sup>, S. Navin<sup>39</sup>, T.K. Nayak<sup>10</sup>, S. Nazarenko<sup>63</sup>, G. Nazarov<sup>63</sup>, A. Nedosekin<sup>13</sup>, F. Nendaz<sup>71</sup>, M. Nicassio<sup>20</sup>, B.S. Nielsen<sup>43</sup>, S. Nikolaev<sup>14</sup>, V. Nikolic<sup>25</sup>, S. Nikulin<sup>14</sup>, V. Nikulin<sup>47</sup>, B.S. Nilsen<sup>70</sup>, M.S. Nilsson<sup>61</sup>, F. Noferini<sup>27</sup>, G. Nooren<sup>73</sup>, N. Novitzky<sup>33</sup>, A. Nyanin<sup>14</sup>, A. Nyatha<sup>99</sup>, C. Nygaard<sup>43</sup>, J. Nystrand<sup>1</sup>, H. Obayashi<sup>109</sup>, A. Ochirov<sup>21</sup>, H. Oeschler<sup>100</sup>, S.K. Oh<sup>12</sup>, J. Oleniacz<sup>92</sup>, C. Oppedisano<sup>15</sup>, A. Ortiz Velasquez<sup>79</sup>, G. Ortona<sup>6,v</sup>, A. Oskarsson<sup>74</sup>, P. Ostrowski<sup>92</sup>, I. Otterlund<sup>74</sup>, J. Otwinowski<sup>23</sup>, G. Øvrebek<sup>1</sup>, K. Oyama<sup>64</sup>, K. Ozawa<sup>96</sup>, Y. Pachmayer<sup>64</sup>, M. Pachr<sup>51</sup>, F. Padilla<sup>46</sup>, P. Pagano<sup>83</sup>, G. Paić<sup>79</sup>, F. Painke<sup>19</sup>, C. Pajares<sup>31</sup>, S. Pal<sup>35</sup>, S.K. Pal<sup>10</sup>, A. Palaha<sup>39</sup>, A. Palmeri<sup>34</sup>, G.S. Pappalardo<sup>34</sup>, W.J. Park<sup>23</sup>, V. Patricchio<sup>84</sup>, A. Pavlinov<sup>57</sup>, T. Pawlak<sup>92</sup>, T. Peitzmann<sup>73</sup>, D. Peresunko<sup>14</sup>, C.E. Pérez Lara<sup>52</sup>, D. Perini<sup>6</sup>, D. Perrino<sup>20</sup>, W. Peryt<sup>92</sup>, A. Pesci<sup>27</sup>, V. Peskov<sup>6,xvi</sup>, Y. Pestov<sup>110</sup>, A.J. Peters<sup>6</sup>, V. Petráček<sup>51</sup>, M. Petran<sup>51</sup>, M. Petris<sup>22</sup>, P. Petrov<sup>39</sup>, M. Petrovici<sup>22</sup>, C. Petta<sup>38</sup>, S. Piano<sup>89</sup>, A. Piccotti<sup>15</sup>, M. Pikna<sup>62</sup>, P. Pillot<sup>28</sup>, O. Pinazza<sup>6</sup>, L. Pinsky<sup>45</sup>, N. Pitz<sup>29</sup>, F. Piuz<sup>6</sup>, D.B. Piyarathna<sup>57,xvii</sup>, R. Platt<sup>39</sup>, M. Płoskoń<sup>98</sup>, J. Pluta<sup>92</sup>, T. Pocheptsov<sup>42,xviii</sup>, S. Pochybova<sup>7</sup>, P.L.M. Podesta-Lerma<sup>93</sup>, M.G. Poghosyan<sup>46</sup>, B. Polichtchouk<sup>55</sup>, A. Pop<sup>22</sup>, V. Pospíšil<sup>51</sup>, B. Potukuchi<sup>48</sup>, S.K. Prasad<sup>57</sup>, R. Preghenella<sup>16</sup>, F. Prino<sup>15</sup>, C.A. Pruneau<sup>57</sup>, I. Pshenichnov<sup>88</sup>, G. Puddu<sup>85</sup>, A. Pulvirenti<sup>38,ii</sup>, V. Punin<sup>63</sup>, M. Putiš<sup>56</sup>, J. Putschke<sup>4</sup>, E. Quercigh<sup>6</sup>, H. Qvigstad<sup>61</sup>, A. Rachevski<sup>89</sup>, A. Rademakers<sup>6</sup>, S. Radomski<sup>64</sup>, T.S. Rähä<sup>33</sup>, J. Rak<sup>33</sup>, A. Rakotozafindrabe<sup>35</sup>, L. Ramello<sup>78</sup>, A. Ramírez Reyes<sup>69</sup>, M. Rammler<sup>41</sup>, R. Raniwala<sup>111</sup>, S. Raniwala<sup>111</sup>, S.S. Räsänen<sup>33</sup>, D. Rathee<sup>5</sup>, K.F. Read<sup>107</sup>, J.S. Real<sup>30</sup>, K. Redlich<sup>86,xix</sup>, R. Renfordt<sup>29</sup>, A.R. Reolon<sup>49</sup>, A. Reshetin<sup>88</sup>, F. Rettig<sup>19</sup>, J.-P. Revol<sup>6</sup>, K. Reygers<sup>64</sup>, H. Ricard<sup>100</sup>, L. Riccati<sup>15</sup>, R.A. Ricci<sup>112</sup>, M. Richter<sup>1,xx</sup>, P. Riedler<sup>6</sup>, W. Riegler<sup>6</sup>, F. Riggi<sup>38</sup>, M. Rodríguez Cahuantzi<sup>77</sup>, D. Rohr<sup>19</sup>, D. Röhrich<sup>1</sup>, R. Romita<sup>23</sup>, F. Ronchetti<sup>49</sup>, P. Rosinský<sup>6</sup>, P. Rosnet<sup>36</sup>, S. Rossegger<sup>6</sup>, A. Rossi<sup>50</sup>, F. Roukoutakis<sup>90</sup>, S. Rousseau<sup>60</sup>, C. Roy<sup>44</sup>, P. Roy<sup>58</sup>, A.J. Rubio Montero<sup>53</sup>, R. 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A. Sandoval<sup>9</sup>, M. Sano<sup>75</sup>, S. Sano<sup>96</sup>, R. Santo<sup>41</sup>, R. Santoro<sup>84</sup>, J. Sarkamo<sup>33</sup>, P. Saturnini<sup>36</sup>, E. Scapparone<sup>27</sup>, F. Scarlassara<sup>50</sup>, R.P. Scharenberg<sup>113</sup>, C. Schiaua<sup>22</sup>, R. Schicker<sup>64</sup>, C. Schmidt<sup>23</sup>, H.R. Schmidt<sup>23,xxi</sup>, S. Schreiner<sup>6</sup>, S. Schuchmann<sup>29</sup>, J. Schukraft<sup>6</sup>, Y. Schutz<sup>28,ii</sup>, K. Schwarz<sup>23</sup>, K. Schweda<sup>64</sup>, G. Scioli<sup>8</sup>, E. Scomparin<sup>15</sup>, P.A. Scott<sup>39</sup>, R. Scott<sup>107</sup>, G. Segato<sup>50</sup>, S. Senyukov<sup>78</sup>, J. Seo<sup>12</sup>, S. Serici<sup>85</sup>, E. Serradilla<sup>53</sup>, A. Sevcenco<sup>80</sup>, I. Sgura<sup>84</sup>, G. Shabratova<sup>42</sup>, R. Shahoyan<sup>6</sup>, N. Sharma<sup>5</sup>, S. Sharma<sup>48</sup>, K. Shigaki<sup>109</sup>, M. Shimomura<sup>75</sup>, K. Shtejer<sup>2</sup>, Y. Sibiriyak<sup>14</sup>, M. 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Urbán<sup>56</sup>, G.M. Urciuoli<sup>87</sup>, G.L. Usai<sup>85</sup>, M. Vajzer<sup>51</sup>, M. Vala<sup>42,xii</sup>, L. Valencia Palomo<sup>60</sup>, S. Vallero<sup>64</sup>, N. van der Kolk<sup>52</sup>, M. van Leeuwen<sup>73</sup>, P. Vande Vyvre<sup>6</sup>, L. Vannucci<sup>112</sup>, A. Vargas<sup>77</sup>, R. Varma<sup>99</sup>, M. Vasileiou<sup>90</sup>, A. Vasiliev<sup>14</sup>, V. Vechernin<sup>21</sup>, M. Veldhoen<sup>73</sup>, M. Venaruzzo<sup>68</sup>, E. Vercellin<sup>46</sup>, S. Vergara<sup>77</sup>, D.C. Vernekohl<sup>41</sup>, R. Vernet<sup>116</sup>, M. Verweij<sup>73</sup>, L. Vickovic<sup>94</sup>, G. Viesti<sup>50</sup>, O. Vikhlyantsev<sup>63</sup>, Z. Vilakazi<sup>65</sup>, O. Villalobos Baillie<sup>39</sup>, A. Vinogradov<sup>14</sup>, L. Vinogradov<sup>21</sup>, Y. Vinogradov<sup>63</sup>, T. Virgili<sup>83</sup>, Y.P. Viyogi<sup>10</sup>, A. Vodopyanov<sup>42</sup>, K. Voloshin<sup>13</sup>, S. Voloshin<sup>57</sup>, G. Volpe<sup>20</sup>, B. von Haller<sup>6</sup>, D. Vranic<sup>23</sup>, J. Vrláková<sup>56</sup>, B. Vulpescu<sup>36</sup>, A. Vyushin<sup>63</sup>, B. Wagner<sup>1</sup>, V. Wagner<sup>51</sup>, R. Wan<sup>44,xxii</sup>, D. Wang<sup>66</sup>, M. Wang<sup>66</sup>, Y. Wang<sup>64</sup>, Y. Wang<sup>66</sup>, K. Watanabe<sup>75</sup>, J.P. Wessels<sup>41,viii</sup>, U. Westerhoff<sup>41</sup>, J. Wiechula<sup>64,xxiii</sup>, J. Wikne<sup>61</sup>, M. Wilde<sup>41</sup>, A. Wilk<sup>41</sup>, G. Wilk<sup>86</sup>, M.C.S. Williams<sup>27</sup>, B. Windelband<sup>64</sup>, H. Yang<sup>35</sup>, S. Yasnopolskiy<sup>14</sup>, J. Yi<sup>115</sup>, Z. Yin<sup>66</sup>, H. Yokoyama<sup>75</sup>, I.-K. Yoo<sup>115</sup>, X. Yuan<sup>66</sup>, I. Yushmanov<sup>14</sup>, E. Zabrodin<sup>61</sup>, C. Zach<sup>51</sup>, C. Zampolli<sup>6</sup>, S. Zaporozhets<sup>42</sup>, A. Zarochentsev<sup>21</sup>, P. Závada<sup>105</sup>, N. Zaviyalov<sup>63</sup>, H. Zbroszczyk<sup>92</sup>, P. Zelnick<sup>59,ii</sup>, A. Zenin<sup>55</sup>, I. Zgura<sup>80</sup>, M. Zhalov<sup>47</sup>, X. Zhang<sup>66,i</sup>, D. Zhou<sup>66</sup>, F. Zhou<sup>66</sup>, Y. Zhou<sup>73</sup>, X. Zhu<sup>66</sup>, A. Zichichi<sup>8,xxiv</sup>, G. Zinovjev<sup>17</sup>, Y. Zoccarato<sup>71</sup>, M. Zynovyev<sup>17</sup>

<sup>1</sup> Department of Physics and Technology, University of Bergen, Bergen, Norway

<sup>2</sup> Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba

<sup>3</sup> Nuclear Physics Institute, Academy of Sciences of the Czech Republic, Řež u Prahy, Czech Republic

<sup>4</sup> Yale University, New Haven, CT, United States

<sup>5</sup> Physics Department, Panjab University, Chandigarh, India

<sup>6</sup> European Organization for Nuclear Research (CERN), Geneva, Switzerland

<sup>7</sup> KFKI Research Institute for Particle and Nuclear Physics, Hungarian Academy of Sciences, Budapest, Hungary

<sup>8</sup> Dipartimento di Fisica dell'Università and Sezione INFN, Bologna, Italy

<sup>9</sup> Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico

<sup>10</sup> Variable Energy Cyclotron Centre, Kolkata, India

<sup>11</sup> Department of Physics, Aligarh Muslim University, Aligarh, India

<sup>12</sup> Gangneung-Wonju National University, Gangneung, South Korea

<sup>13</sup> Institute for Theoretical and Experimental Physics, Moscow, Russia

<sup>14</sup> Russian Research Centre Kurchatov Institute, Moscow, Russia

<sup>15</sup> Sezione INFN, Turin, Italy

<sup>16</sup> Centro Fermi – Centro Studi e Ricerche e Museo Storico della Fisica “Enrico Fermi”, Rome, Italy

<sup>17</sup> Bogolyubov Institute for Theoretical Physics, Kiev, Ukraine

<sup>18</sup> Faculty of Engineering, Bergen University College, Bergen, Norway

<sup>19</sup> Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany

<sup>20</sup> Dipartimento Interateneo di Fisica ‘M. Merlin’ and Sezione INFN, Bari, Italy

<sup>21</sup> V. Fock Institute for Physics, St. Petersburg State University, St. Petersburg, Russia

<sup>22</sup> National Institute for Physics and Nuclear Engineering, Bucharest, Romania

<sup>23</sup> Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt, Germany

<sup>24</sup> Department of Physics, Ohio State University, Columbus, OH, United States

<sup>25</sup> Rudjer Bošković Institute, Zagreb, Croatia

<sup>26</sup> Sezione INFN, Padova, Italy

<sup>27</sup> Sezione INFN, Bologna, Italy

<sup>28</sup> SUBATECH, Ecole des Mines de Nantes, Université de Nantes, CNRS-IN2P3, Nantes, France

<sup>29</sup> Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany

<sup>30</sup> Laboratoire de Physique Subatomique et de Cosmologie (LPSC), Université Joseph Fourier, CNRS-IN2P3, Institut Polytechnique de Grenoble, Grenoble, France

<sup>31</sup> Departamento de Física de Partículas and IGFAE, Universidad de Santiago de Compostela, Santiago de Compostela, Spain

<sup>32</sup> Oak Ridge National Laboratory, Oak Ridge, TN, United States

<sup>33</sup> Helsinki Institute of Physics (HIP) and University of Jyväskylä, Jyväskylä, Finland

- 34 Sezione INFN, Catania, Italy
- 35 Commissariat à l'Energie Atomique, IRFU, Saclay, France
- 36 Laboratoire de Physique Corpusculaire (LPC), Clermont Université, Université Blaise Pascal, CNRS-IN2P3, Clermont-Ferrand, France
- 37 Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia
- 38 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Catania, Italy
- 39 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- 40 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Cracow, Poland
- 41 Institut für Kernphysik, Westfälische Wilhelms-Universität Münster, Münster, Germany
- 42 Joint Institute for Nuclear Research (JINR), Dubna, Russia
- 43 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- 44 Institut Pluridisciplinaire Hubert Curien (IPHC), Université de Strasbourg, CNRS-IN2P3, Strasbourg, France
- 45 University of Houston, Houston, TX, United States
- 46 Dipartimento di Fisica Sperimentale dell'Università and Sezione INFN, Turin, Italy
- 47 Petersburg Nuclear Physics Institute, Gatchina, Russia
- 48 Physics Department, University of Jammu, Jammu, India
- 49 Laboratori Nazionali di Frascati, INFN, Frascati, Italy
- 50 Dipartimento di Fisica dell'Università and Sezione INFN, Padova, Italy
- 51 Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
- 52 Nikhef, National Institute for Subatomic Physics, Amsterdam, Netherlands
- 53 Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
- 54 Moscow Engineering Physics Institute, Moscow, Russia
- 55 Institute for High Energy Physics, Protvino, Russia
- 56 Faculty of Science, P.J. Šafárik University, Košice, Slovakia
- 57 Wayne State University, Detroit, MI, United States
- 58 Saha Institute of Nuclear Physics, Kolkata, India
- 59 Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- 60 Institut de Physique Nucléaire d'Orsay (IPNO), Université Paris-Sud, CNRS-IN2P3, Orsay, France
- 61 Department of Physics, University of Oslo, Oslo, Norway
- 62 Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia
- 63 Russian Federal Nuclear Center (VNIIEF), Sarov, Russia
- 64 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- 65 Physics Department, University of Cape Town, iThemba LABS, Cape Town, South Africa
- 66 Hua-Zhong Normal University, Wuhan, China
- 67 Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
- 68 Dipartimento di Fisica dell'Università and Sezione INFN, Trieste, Italy
- 69 Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
- 70 Physics Department, Creighton University, Omaha, NE, United States
- 71 Université de Lyon, Université Lyon 1, CNRS/IN2P3, IPN-Lyon, Villeurbanne, France
- 72 Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
- 73 Nikhef, National Institute for Subatomic Physics and Institute for Subatomic Physics of Utrecht University, Utrecht, Netherlands
- 74 Division of Experimental High Energy Physics, University of Lund, Lund, Sweden
- 75 University of Tsukuba, Tsukuba, Japan
- 76 Sezione INFN, Cagliari, Italy
- 77 Benemérita Universidad Autónoma de Puebla, Puebla, Mexico
- 78 Dipartimento di Scienze e Technologie Avanzate dell'Università del Piemonte Orientale and Gruppo Collegato INFN, Alessandria, Italy
- 79 Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 80 Institute of Space Sciences (ISS), Bucharest, Romania
- 81 Institute of Physics, Bhubaneswar, India
- 82 Universidade de São Paulo (USP), São Paulo, Brazil
- 83 Dipartimento di Fisica 'E.R. Caianiello' dell'Università and Gruppo Collegato INFN, Salerno, Italy
- 84 Sezione INFN, Bari, Italy
- 85 Dipartimento di Fisica dell'Università and Sezione INFN, Cagliari, Italy
- 86 Soltan Institute for Nuclear Studies, Warsaw, Poland
- 87 Sezione INFN, Rome, Italy
- 88 Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
- 89 Sezione INFN, Trieste, Italy
- 90 Physics Department, University of Athens, Athens, Greece
- 91 Chicago State University, Chicago, IL, United States
- 92 Warsaw University of Technology, Warsaw, Poland
- 93 Universidad Autónoma de Sinaloa, Culiacán, Mexico
- 94 Technical University of Split FESB, Split, Croatia
- 95 Yerevan Physics Institute, Yerevan, Armenia
- 96 University of Tokyo, Tokyo, Japan
- 97 Department of Physics, Sejong University, Seoul, South Korea
- 98 Lawrence Berkeley National Laboratory, Berkeley, CA, United States
- 99 Indian Institute of Technology, Mumbai, India
- 100 Institut für Kernphysik, Technische Universität Darmstadt, Darmstadt, Germany
- 101 Yonsei University, Seoul, South Korea
- 102 Zentrum für Technologietransfer und Telekommunikation (ZIT), Fachhochschule Worms, Worms, Germany
- 103 California Polytechnic State University, San Luis Obispo, CA, United States
- 104 China Institute of Atomic Energy, Beijing, China
- 105 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- 106 The University of Texas at Austin, Physics Department, Austin, TX, United States
- 107 University of Tennessee, Knoxville, TN, United States
- 108 Dipartimento di Fisica dell'Università 'La Sapienza' and Sezione INFN, Rome, Italy
- 109 Hiroshima University, Hiroshima, Japan
- 110 Budker Institute for Nuclear Physics, Novosibirsk, Russia
- 111 Physics Department, University of Rajasthan, Jaipur, India
- 112 Laboratori Nazionali di Legnaro, INFN, Legnaro, Italy

<sup>113</sup> *Purdue University, West Lafayette, IN, United States*

<sup>114</sup> *Lawrence Livermore National Laboratory, Livermore, CA, United States*

<sup>115</sup> *Pusan National University, Pusan, South Korea*

<sup>116</sup> *Centre de Calcul de l'IN2P3, Villeurbanne, France*

\* Corresponding author.

*E-mail address: giuseppe.bruno@ba.infn.it (G.E. Bruno).*

<sup>i</sup> Also at Laboratoire de Physique Corpusculaire (LPC), Clermont Université, Université Blaise Pascal, CNRS-IN2P3, Clermont-Ferrand, France.

<sup>ii</sup> Also at European Organization for Nuclear Research (CERN), Geneva, Switzerland.

<sup>iii</sup> Now at Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany.

<sup>iv</sup> Also at Dipartimento di Fisica dell'Università, Udine, Italy.

<sup>v</sup> Also at Dipartimento di Fisica Sperimentale dell'Università and Sezione INFN, Turin, Italy.

<sup>vi</sup> Also at Benemérita Universidad Autónoma de Puebla, Puebla, Mexico.

<sup>vii</sup> Deceased.

<sup>viii</sup> Now at European Organization for Nuclear Research (CERN), Geneva, Switzerland.

<sup>ix</sup> Also at Wayne State University, Detroit, Michigan, United States.

<sup>x</sup> Also at Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany.

<sup>xi</sup> Also at Fachhochschule Köln, Köln, Germany.

<sup>xii</sup> Also at Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia.

<sup>xiii</sup> Also at Laboratoire de Physique Subatomique et de Cosmologie (LPSC), Université Joseph Fourier, CNRS-IN2P3, Institut Polytechnique de Grenoble, Grenoble, France.

<sup>xiv</sup> Also at "Vinča" Institute of Nuclear Sciences, Belgrade, Serbia.

<sup>xv</sup> Also at Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt, Germany.

<sup>xvi</sup> Also at Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico.

<sup>xvii</sup> Also at University of Houston, Houston, Texas, United States.

<sup>xviii</sup> Also at Department of Physics, University of Oslo, Oslo, Norway.

<sup>xix</sup> Also at Institute of Theoretical Physics, University of Wrocław, Wrocław, Poland.

<sup>xx</sup> Now at Department of Physics, University of Oslo, Oslo, Norway.

<sup>xxi</sup> Also at Eberhard Karls Universität Tübingen, Tübingen, Germany.

<sup>xxii</sup> Also at Hua-Zhong Normal University, Wuhan, China.

<sup>xxiii</sup> Now at Eberhard Karls Universität Tübingen, Tübingen, Germany.

<sup>xxiv</sup> Also at Centro Fermi – Centro Studi e Ricerche e Museo Storico della Fisica "Enrico Fermi", Rome, Italy.