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Nonlinear Schrödinger equation with unbounded or vanishing potentials: Solutions concentrating on lower dimensional spheres

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ABSTRACT

We study positive bound states for the equation

$$-\varepsilon^2 \Delta u + V(x)u = K(x)f(u), \quad x \in \mathbb{R}^N,$$

where $\varepsilon > 0$ is a real parameter and *V* and *K* are radial positive potentials. We are especially interested in solutions which concentrate on a *k*-dimensional sphere, $1 \le k \le N - 1$, as $\varepsilon \to 0$. We adopt a purely variational approach which allows us to consider broader classes of potentials than those treated in previous works. For example, *V* and *K* might be singular at the origin or vanish superguadratically at infinity.

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1. Introduction

We consider the nonlinear Schrödinger equation

$$i\hbar\frac{\partial\psi}{\partial t} = -\frac{\hbar^2}{2m}\Delta\psi + W(x)\psi - |\psi|^{p-1}\psi, \quad (t,x) \in \mathbb{R} \times \mathbb{R}^N,$$
(1)

which appears for instance in nonlinear optics or condensed matter physics. A *standing wave* solution of (1) is a solution of the form

$$\psi(t, x) = e^{-iEt/\hbar}u(x),$$

where *E* is the energy of the wave. The function ψ is a standing wave solution of (1) if and only if *u* is a solution of the semilinear elliptic equation

$$-\varepsilon^2 \Delta u + V(x)u = |u|^{p-1}u, \quad x \in \mathbb{R}^N,$$
⁽²⁾

where $\varepsilon^2 = \hbar^2/2m$ and V(x) = W(x) - E. It is a *bound state* if $u \in H^1(\mathbb{R}^N)$. From a physical point of view, one expects to recover the laws of classical mechanics when $\hbar \to 0$. It is thus interesting to study the behavior of the solutions of (2) as ε tends to 0. The bound states of (2) with ε small are referred to as *semiclassical states*.

It is well known that problem (2) possesses solutions which exhibit concentration phenomena as $\varepsilon \rightarrow 0$. More precisely, these solutions converge uniformly to 0 outside some concentration set, while remaining uniformly positive in the concentration set. This concentration set can be either a point, a finite set of points or a manifold.

The solutions concentrating around one or several isolated points have been intensively studied (see for example [1,8] and their bibliographies).

On the other hand, one can ask if there exist solutions of (2) concentrating on a higher dimensional set. This problem has been solved for some specific higher dimensional sets. Solutions concentrating on curves have been found recently in [13], see also [12] for the case N = 2 and [14,15] for a Neumann singularly perturbed problem. Here we shall restrict ourselves to the problem of solutions concentrating around spheres. In several recent papers [3–7,9], solutions concentrating on (N - 1)-dimensional spheres have been found. In [16], solutions concentrating on (N - 2)-dimensional spheres are investigated.

We focus on solutions concentrating around a *k*-dimensional sphere in \mathbb{R}^N , $1 \le k \le N - 1$. The existence of such solutions has been discussed in remarks in [1,2,9]. Particular problems arise in the *critical frequency* case, namely when $\inf_{\mathbb{R}^N} V = 0$. These problems have been tackled in [4] and [9].

Theorem 1. (See Ambrosetti and Ruiz [4].) Assume that p > 1, that $V \in C^1(\mathbb{R}^N)$ is a positive bounded radially symmetric potential, that ∇V is bounded and that

$$\liminf_{|x|\to\infty} V(x)|x|^2 > 0.$$

If there exists r^* such that the function $\mathcal{M}: (0, \infty) \to \mathbb{R}$ defined for r > 0 by

$$\mathcal{M}(r) := r^{N-1} \left[V(r) \right]^{\frac{p+1}{p-1} - \frac{1}{2}}$$
(3)

has an isolated local maximum or minimum at $r = r^*$, then, for $\varepsilon > 0$ small enough, Eq. (2) has a positive radially symmetric solution $u_{\varepsilon} \in H^1(\mathbb{R}^N)$ that concentrates at the sphere $|x| = r^*$.

The problem in [9] is rather different. The potential V vanishes and the solutions concentrate around zeroes of V. The asymptotic behavior depends on the shape of V around 0.

Theorem 1 relies on a Lyapunov–Schmidt reduction method. The aim of this note is to examine possible improvements in the previous results that can be obtained by using the penalization method, a variational method originally due to Del Pino and Felmer [11] and adapted to our framework in the papers [8,17]. This method permits us to treat superquadratically decaying potentials, or even compactly supported potentials.

Our results include the following simple particular case.

Theorem 2. Let $N \ge 3$, $p > \frac{N}{N-2}$ and $V \in C(\mathbb{R}^N \setminus \{0\}, \mathbb{R}^+)$ be a radial potential. If there exists $r^* > 0$ such that the function $\mathcal{M}(r)$ defined by (3) has an isolated local minimum at $r = r^*$ such that $\mathcal{M}(r^*) > 0$, then for ε small enough, Eq. (2) has a positive radially symmetric solution u_{ε} that concentrates on the sphere of radius r^* .

If $N \ge 5$, one has also that $u_{\varepsilon} \in L^2(\mathbb{R}^N)$ (see Corollary 5.7).

In contrast with Theorem 1, we do not require any boundedness assumption on V or its derivatives, and we treat potentials V which are singular at the origin or vanish superquadratically at infinity.

Theorem 2 is a particular case of Theorem 3 below, which deals with a nonlinearity which is neither necessarily homogeneous nor autonomous, see Eq. (4) below. Furthermore, we will find solutions concentrating on a *k*-dimensional sphere, $1 \le k \le N - 1$. In this case, the critical exponent to be taken into consideration is $p_k = \frac{N-k+2}{N-k-2}$ if $N-k \ge 3$, $p_k = \infty$ if N-k = 1, 2. We also obtain results for N = 2 with a little more care, see Section 6.

Let us point out that if V is compactly supported and $p \leq \frac{N}{N-2}$, then Eq. (2) has no positive solution in the neighborhood of infinity, see the discussion in [17].

Assuming that the potential V is cylindrically symmetric, we can reduce (2) to a problem in \mathbb{R}^{N-k} . The single-peaked solutions of this problem can then be extended to \mathbb{R}^N by symmetry. In this way, we obtain a solution of (2) concentrating around a k-dimensional sphere. Observe that since the reduced problem is in \mathbb{R}^{N-k} , the critical exponent to be considered is the one in dimension N - k. This allows for example to treat critical or supercritical problems by looking for cylindrically symmetric (non-necessarily radial) solutions. Another kind of solutions to (2) with p supercritical and V decaying superquadratically have been found in [10].

2. Assumptions and main result

We shall study the equation with a more general nonlinearity

$$-\varepsilon^2 \Delta u + V(x)u = K(x)f(u), \quad x \in \mathbb{R}^N.$$
(4)

Let *k* be a fixed integer such that $1 \le k \le N - 1$. This number *k* is the dimension of the sphere on which we want to construct concentrating solutions. Let us choose any (N - k - 1)-dimensional linear subspace $\mathcal{H} \subset \mathbb{R}^N$. We denote by \mathcal{H}^{\perp} the orthogonal complement of \mathcal{H} .

2.1. The potentials

We consider a nonnegative potential $V \in C(\mathbb{R}^N \setminus \{0\})$ and a nonnegative competing function $K \in C(\mathbb{R}^N \setminus \{0\})$, $K \neq 0$. We assume that for every $R \in \mathbf{O}(N)$ such that $R(\mathcal{H}) = \mathcal{H}$, we have $V \circ R = V$ and $K \circ R = K$. This will be the case if for example V and K are radial functions. Notice that if we write x = (x', x'') with $x' \in \mathcal{H}$ and $x'' \in \mathcal{H}^{\perp}$, this assumption is equivalent to the fact that V and K depend only on |x'| and |x''|.

2.2. The nonlinearity

We make classical assumptions on f that lead to a good minimax characterization of the infimum on the Nehari manifold. Namely, we assume that $f: \mathbb{R}^+ \to \mathbb{R}^+$ is continuous and that

(*f*₁) there exists q > 1 such that $f(s) = O(s^q)$ as $s \to 0^+$, (*f*₂) there exists p > 1 such that $\frac{1}{p+1} > \frac{1}{2} - \frac{1}{N-k}$ and $f(s) = O(s^p)$ as $s \to \infty$,

(*f*₃) there exists $2 < \theta \le p + 1$ such that

$$0 < \theta F(s) \leq f(s)s$$
 for $s > 0$,

where $F(s) := \int_0^s f(\sigma) d\sigma$, (f_4) the function

$$s \mapsto \frac{f(s)}{s}$$

is nondecreasing.

Notice that (f_2) is nothing but the subcriticality condition in dimension N - k.

2.3. The growth conditions

Following [8,17] we impose one of the three sets of growth conditions at infinity:

 (\mathcal{G}_{∞}^1) there exists $\sigma < (N-2)q - N$ such that

$$\limsup_{|x|\to\infty}\frac{K(x)}{|x|^{\sigma}}<\infty;$$

 (\mathcal{G}^2_{∞}) there exists $\sigma \in \mathbb{R}$ such that

$$\liminf_{|x|\to\infty} V(x)|x|^2 > 0 \quad \text{and} \quad \limsup_{|x|\to\infty} \frac{K(x)}{|x|^{\sigma}} < \infty;$$

 (\mathcal{G}^3_{∞}) there exist $\alpha < 2$ and $\sigma \in \mathbb{R}$ such that

$$\liminf_{|x|\to\infty} V(x)|x|^{\alpha} > 0 \quad \text{and} \quad \limsup_{|x|\to\infty} \frac{K(x)}{\exp(\sigma|x|^{\frac{2-\alpha}{2}})} < \infty.$$

Note that in comparison with [8], in (\mathcal{G}_{∞}^2) and (\mathcal{G}_{∞}^3) , V might vanish somewhere. We also impose one of the three sets of growth conditions at the origin, which mirror those at infinity:

 (\mathcal{G}_0^1) there exists $\tau > -2$, such that

$$\limsup_{|x|\to 0}\frac{K(x)}{|x|^{\tau}}<\infty;$$

 (\mathcal{G}_0^2) there exists $\tau \in \mathbb{R}$ such that

$$\liminf_{|x|\to 0} V(x)|x|^2 > 0 \quad \text{and} \quad \limsup_{|x|\to 0} \frac{K(x)}{|x|^{\tau}} < \infty;$$

 (\mathcal{G}_0^3) there exist $\gamma > 2$ and $\tau \in \mathbb{R}$ such that

$$\liminf_{|x| \to 0} V(x)|x|^{\gamma} > 0 \quad \text{and} \quad \limsup_{|x| \to 0} \frac{K(x)}{\exp(\tau |x|^{-\frac{\gamma-2}{2}})} < \infty.$$

By Kelvin transform, there is a duality between the conditions at the origin and the conditions at infinity, at least in the case where $f(t) = t^p$. If one defines \hat{u} to be the Kelvin transform of u, i.e.,

$$\hat{u}(x) = \frac{1}{|x|^{N-2}} u\left(\frac{x}{|x|^2}\right)$$

and the transformed potentials

$$\hat{V}(x) = \frac{1}{|x|^4} V\left(\frac{x}{|x|^2}\right)$$

and

$$\hat{K}(x) = \frac{1}{|x|^{N+2-p(N-2)}} K\left(\frac{x}{|x|^2}\right),$$

the function u_{ε} solves (4) if and only if \hat{u}_{ε} solves the same problem with \hat{V} and \hat{K} in place of V and K. One sees that V, K satisfy (\mathcal{G}_0^i) if and only if \hat{V} and \hat{K} satisfy (\mathcal{G}_{∞}^i) .

The problem at the origin is in a sense in duality with the one at infinity. Whereas a slow decay of V at infinity does allow a lot of freedom for K, a strong singularity at the origin allows for very singular K's too. The critical threshold growth is $1/|x|^2$ both at the origin and at infinity. This can be made clearer if we observe that the optimal barrier functions at the origin are the optimal one at infinity mapped by Kelvin transform.

2.4. The auxiliary potential

Before we can state our last assumption, we need a few preliminaries. Let a, b > 0. The equation

$$-\Delta u + au = bf(u) \quad \text{in } \mathbb{R}^{N-k}$$
(5)

is called the *limit equation* associated with (4). The weak solutions of (5) are critical points of the functional $\mathcal{I}_{a,b}: H^1(\mathbb{R}^{N-k}) \to \mathbb{R}$ defined by

$$\mathcal{I}_{a,b}(u) := \frac{1}{2} \int_{\mathbb{R}^{N-k}} \left(|\nabla u|^2 + au^2 \right) dx - b \int_{\mathbb{R}^{N-k}} F(u) \, dx.$$
(6)

Any nontrivial critical point $u \in H^1(\mathbb{R}^{N-k})$ of $\mathcal{I}_{a,b}$, belongs to the Nehari manifold

$$\mathcal{N}_{a,b} := \left\{ u \in H^1(\mathbb{R}^{N-k}) \mid u \neq 0 \text{ and } \left\langle \mathcal{I}'_{a,b}(u), u \right\rangle = 0 \right\}.$$

A solution $u \in H^1(\mathbb{R}^{N-k})$ is a least-energy solution of (5) if

$$\mathcal{I}_{a,b}(u) = \inf_{\nu \in \mathcal{N}_{a,b}} \mathcal{I}_{a,b}(\nu).$$

The ground-energy function is defined by

$$\mathcal{E}: \mathbb{R}^+ \times \mathbb{R}^+ \to \mathbb{R}^+: (a, b) \mapsto \mathcal{E}(a, b) := \inf_{u \in \mathcal{N}_{a, b}} \mathcal{I}_{a, b}(u),$$

and the *auxiliary potential* $\mathcal{M}: \mathbb{R}^N \to (0, +\infty]$ by

$$(x', x'') \mapsto \mathcal{M}(x', x'') := \begin{cases} |x''|^k \mathcal{E}(V(x), K(x)) & \text{if } K(x) > 0, \\ +\infty & \text{if } K(x) = 0. \end{cases}$$

The following lemma states some properties of the ground-energy function, see [8, Lemma 3].

Lemma 2.1. Assume $f : \mathbb{R}^+ \to \mathbb{R}^+$ is a continuous function that fulfills assumptions $(f_1)-(f_4)$. Then, for every $(a, b) \in \mathbb{R}^+_0 \times \mathbb{R}^+_0$, $\mathcal{E}(a, b)$ is a critical value of $\mathcal{I}_{a,b}$ and we have

$$\mathcal{E}(a,b) = \inf_{\substack{u \in H^1(\mathbb{R}^N) \\ u \neq 0}} \max_{t \ge 0} \mathcal{I}_{a,b}(tu).$$

If $u \in \mathcal{N}_{a,b}$ and $\mathcal{E}(a,b) = \mathcal{I}_{a,b}(u)$, then $u \in C^1(\mathbb{R}^N)$ and up to a translation, u is a radial function such that $\nabla u(x) \cdot x < 0$ for every $x \in \mathbb{R}^N \setminus \{0\}$. Moreover, the following properties hold:

(i) \mathcal{E} is continuous in $\mathbb{R}_0^+ \times \mathbb{R}_0^+$; (ii) for every $b^* \in \mathbb{R}_0^+$, $a \to \mathcal{E}(a, b^*)$ is strictly increasing; (iii) for every $a^* \in \mathbb{R}_0^+$, $b \to \mathcal{E}(a^*, b)$ is strictly decreasing; (iv) for every $\lambda > 0$, $\mathcal{E}(\lambda a, \lambda b) = \lambda^{1-N/2} \mathcal{E}(a, b)$; (v) if $f(u) = u^p$ with $\frac{1}{2} - \frac{1}{N-k} < \frac{1}{p+1} < \frac{1}{2}$, then

$$\mathcal{E}(a,b) = \mathcal{E}(1,1)a^{\frac{p+1}{p-1} - \frac{N}{2}}b^{-\frac{2}{p-1}}.$$

If $f(u) = u^p$ with $\frac{1}{2} - \frac{1}{N-k} < \frac{1}{p+1} < \frac{1}{2}$, the last property of the preceding lemma implies the following explicit form of the auxiliary potential:

$$\mathcal{M}(x', x'') = \mathcal{E}(1, 1) |x''|^k [V(x)]^{\frac{p+1}{p-1} - \frac{N-k}{2}} [K(x)]^{\frac{-2}{p-1}}$$

Due to the symmetry that we shall impose on the solution (see (14)), the concentration can only occur in the space \mathcal{H}^{\perp} . We assume that there exists a smooth bounded open set $\Lambda \subset \mathbb{R}^N$ such that

$$\bar{\Lambda} \cap \mathcal{H} = \emptyset, \qquad \Lambda \cap \mathcal{H}^{\perp} \neq \emptyset, \tag{7}$$

for every $R \in \mathbf{O}(N)$ such that $R(\mathcal{H}) = \mathcal{H}$,

$$R(\Lambda) = \Lambda \tag{8}$$

and

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$$0 < \inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M} < \inf_{\partial \Lambda \cap \mathcal{H}^{\perp}} \mathcal{M}.$$
 (9)

In the case where k = N - 2, we shall need the condition

$$\inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M} < 2 \inf_{\Lambda} \mathcal{M}.$$
(10)

By continuity of \mathcal{M} in Λ , this condition is not restrictive. Similarly, we can also assume that V > 0 on $\overline{\Lambda}$ and that \mathcal{M} is continuous on $\overline{\Lambda}$.

Our main result is the following theorem.

Theorem 3. Let $N \ge 2$, $V, K \in C(\mathbb{R}^N \setminus \{0\}, \mathbb{R}^+)$ satisfy one set (\mathcal{G}_0^i) of growth conditions at the origin and one set (\mathcal{G}_∞^j) of growth conditions at infinity, and f satisfy assumptions $(f_1)-(f_4)$. Assume there exists an open bounded set $\Lambda \subset \mathbb{R}^N$ such that (7), (8), (9) and, if k = N - 2, (10) hold. Then there exists $\varepsilon_0 > 0$ such that for every $0 < \varepsilon < \varepsilon_0$, problem (2) has at least one positive solution u_{ε} . Moreover, for every $0 < \varepsilon < \varepsilon_0$, there exists $x_{\varepsilon} \in \Lambda \cap \mathcal{H}^{\perp}$ such that u_{ε} attains its maximum at x_{ε} ,

$$\begin{split} \liminf_{\varepsilon \to 0} u_{\varepsilon}(x_{\varepsilon}) &> 0, \\ \lim_{\varepsilon \to 0} \mathcal{M}(x_{\varepsilon}) &= \inf_{A \cap \mathcal{H}^{\perp}} \mathcal{M} \end{split}$$

and there exist C > 0 and $\lambda > 0$ such that

$$u_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d(x, S_{\varepsilon}^{k})}{1 + d(x, S_{\varepsilon}^{k})}\right) \left(1 + |x|^{2}\right)^{\frac{-(N-2)}{2}}, \quad \forall x \in \mathbb{R}^{N},$$

where S_{ε}^{k} is the k-sphere centered at the origin and of radius $|x_{\varepsilon}''|$.

In the special case where $x_0 \in \Lambda \cap \mathcal{H}^{\perp}$ is the unique minimizer of \mathcal{M} on $\Lambda \cap \mathcal{H}^{\perp}$, then $x_{\varepsilon} \to x_0$, and the solution concentrates around a *k*-dimensional sphere of radius $|x_0|$ centered at the origin.

One should note that the theorem is valid in dimension 2, but the solutions that are obtained do not decay at infinity in general.

Theorem 2 follows from Theorem 3 by taking $K \equiv 1$, $f(u) = u^p$ and k = N - 1. Indeed, we notice that the growth condition (\mathcal{G}_0^1) is always satisfied whereas the condition (\mathcal{G}_∞^1) holds if and only if (N-2)p - N > 0, i.e. $p > \frac{N}{N-2}$.

The sequel of the paper is devoted to the proof of Theorem 3. In Section 3, we introduce a penalized problem and prove that it has a least-energy solution. In Section 4, we study the asymptotics of this solution and in Section 5, we obtain decay estimates of the solution and show that it also solves the original problem. In all these sections, we assume that $N \ge 3$. The modifications for the case N = 2 will be addressed in Section 6.

3. The penalization scheme

We assume that $N \ge 3$. The homogeneous Sobolev space $\mathcal{D}^{1,2}(\mathbb{R}^N)$ is the closure of the set of compactly supported smooth functions $\mathcal{D}(\mathbb{R}^N)$ with respect to the norm

$$\left(\int\limits_{\mathbb{R}^N}|\nabla u|^2\,dx\right)^{\frac{1}{2}}.$$

Thanks to the Sobolev inequality, we have $\mathcal{D}^{1,2}(\mathbb{R}^N) \subset L^{2^*}(\mathbb{R}^N)$. Let us also recall Hardy's inequality in \mathbb{R}^N . One has

$$\left(\frac{N-2}{2}\right)^2 \int_{\mathbb{R}^N} \frac{|u(x)|^2}{|x|^2} dx \leqslant \int_{\mathbb{R}^N} |\nabla u|^2,$$

for all $u \in \mathcal{D}^{1,2}(\mathbb{R}^N)$.

Following [17], we define the *penalization potential* $H : \mathbb{R}^N \to \mathbb{R}$ by

$$H(x) := \frac{\kappa}{|x|^2 ((\log |x|)^2 + 1)^{\frac{1+\beta}{2}}}$$

where $\beta > 0$ and $0 < \kappa < (\frac{N-2}{2})^2$. Notice that for all $x \in \mathbb{R}^N$, we have

$$H(x) \leqslant \frac{\kappa}{|x|^2}.$$

By Hardy's inequality, we deduce that the quadratic form associated to $-\Delta - H$ is positive, i.e.

$$\int_{\mathbb{R}^N} \left(|\nabla u|^2 - Hu^2 \right) \ge \left(\left(\frac{N-2}{2} \right)^2 - \kappa \right) \int_{\mathbb{R}^N} \frac{|u(x)|^2}{|x|^2} \, dx \ge 0, \tag{11}$$

for all $u \in \mathcal{D}^{1,2}(\mathbb{R}^N)$.

This inequality implies the following comparison principle.

Proposition 3.1. Let $\Omega \subset \mathbb{R}^N \setminus \{0\}$ be a smooth domain. Let $v, w \in H^1_{loc}(\Omega) \cap C(\overline{\Omega})$ be such that $\nabla(w-v)_- \in L^2(\Omega), (w-v)_-/|x| \in L^2(\Omega)$ and

$$-\Delta w - Hw \ge -\Delta v - Hv, \quad \forall x \in \Omega.$$
⁽¹²⁾

If $\partial \Omega \neq \emptyset$, assume also that $w \ge v$ on $\partial \Omega$. Then $w \ge v$ in Ω .

Proof. It suffices to multiply the inequality (12) by $(w - v)_{-}$, integrate by parts and use (11).

Fix $\mu \in (0, 1)$. We define the penalized nonlinearity $g_{\varepsilon} : \mathbb{R}^N \times \mathbb{R}^+ \to \mathbb{R}$ by

$$g_{\varepsilon}(x,s) := \chi_{\Lambda}(x)K(x)f(s) + (1 - \chi_{\Lambda}(x))\min\{(\varepsilon^2 H(x) + \mu V(x))s, K(x)f(s)\}$$

Let $G_{\varepsilon}(x,s) := \int_0^s g_{\varepsilon}(x,\sigma) d\sigma$. One can check that g_{ε} is a Carathéodory function with the following properties:

 $(g_1) \ g_{\varepsilon}(x,s) = o(s), \ s \to 0^+$, uniformly in compact subsets of \mathbb{R}^N , (g_2) there exists p > 1 such that $\frac{1}{p+1} > \frac{1}{2} - \frac{1}{N-k}$ and

$$\lim_{s\to\infty}\frac{g_{\varepsilon}(x,s)}{s^p}=0,$$

(g₃) there exists $2 < \theta \leq p + 1$ such that

$$0 < \theta G_{\varepsilon}(x, s) \leq g_{\varepsilon}(x, s)s \quad \forall x \in \Lambda, \ \forall s > 0,$$
$$0 < 2G_{\varepsilon}(x, s) \leq g_{\varepsilon}(x, s)s \leq (\varepsilon^{2}H(x) + \mu V(x))s^{2} \quad \forall x \notin \Lambda, \ \forall s > 0,$$

 (g_4) the function

$$s\mapsto \frac{g_{\varepsilon}(x,s)}{s}$$

is nondecreasing for all $x \in \mathbb{R}^N$.

We look for a positive solution of the penalized equation

$$-\varepsilon^2 \Delta u + V(x)u = g_{\varepsilon}(x, u) \quad \text{in } \mathbb{R}^N$$
 ($\mathcal{P}_{\varepsilon}$)

in the Hilbert space

$$H^1_V(\mathbb{R}^N) := \left\{ u \in \mathcal{D}^{1,2}(\mathbb{R}^N) \mid \int_{\mathbb{R}^N} V u^2 < \infty \right\}$$

endowed with the norm

$$\|u\|_{\varepsilon}^{2} := \int_{\mathbb{R}^{N}} \left(\varepsilon^{2} |\nabla u|^{2} + V u^{2}\right).$$
(13)

We will search for a solution of $(\mathcal{P}_{\varepsilon})$ in the closed subspace

$$H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N}) := \left\{ u \in H^{1}_{V}(\mathbb{R}^{N}) \mid \forall R \in \mathbf{O}(N) \text{ s.t. } R(\mathcal{H}) = \mathcal{H}, \ u \circ R = u \right\}.$$
(14)

In other words, if x = (x', x'') with $x' \in \mathcal{H}$ and $x'' \in \mathcal{H}^{\perp}$, the functions in $H^1_{V, \mathcal{H}}(\mathbb{R}^N)$ depend only on |x'| and |x''|.

Define $J_{\varepsilon}: H^1_V \xrightarrow{\mathcal{H}} (\mathbb{R}^N) \to \mathbb{R}$ by

$$J_{\varepsilon}(u) := \frac{1}{2} \int_{\mathbb{R}^{N}} \left(\varepsilon^{2} \left| \nabla u(x) \right|^{2} + V(x) \left| u(x) \right|^{2} \right) dx - \int_{\mathbb{R}^{N}} G_{\varepsilon} \left(x, u(x) \right) dx$$

The functional J_{ε} is well defined and of class $C^{1}(H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N}),\mathbb{R})$. By the principle of symmetric criticality [18], critical points are weak solutions of $(\mathcal{P}_{\varepsilon})$. Furthermore, J_{ε} has the mountain pass geometry. It remains to show that J_{ε} satisfies the Palais–Smale condition. The proof below is inspired from [8]. Recall that a sequence $(u_{n})_{n} \subset H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N})$ is a Palais–Smale sequence for J_{ε} if

$$J_{\varepsilon}(u_n) \leqslant C$$
 and $J'_{\varepsilon}(u_n) \to 0$, $n \to \infty$.

Proposition 3.2. For ε sufficiently small, every Palais–Smale sequence for J_{ε} contains a convergent subsequence.

Proof. Let $(u_n)_n \subset H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be a Palais–Smale sequence for J_{ε} . It is standard to check, using (g_3) , that for ε sufficiently small, the sequence $(u_n)_n$ is bounded in $H^1_{V,\mathcal{H}}(\mathbb{R}^N)$. We infer that, up to a subsequence, $u_n \rightharpoonup u$ in $H^1_{V,\mathcal{H}}(\mathbb{R}^N)$.

For $\lambda \in \mathbb{R}^+$, set $A_{\lambda} := \overline{B(0, e^{\lambda})} \setminus B(0, e^{-\lambda})$. Note that

$$H(x) \leqslant \frac{\kappa}{|x|^2 |\log |x||^{1+\beta}}$$

By Hardy's inequality, we have for $\lambda \ge 0$,

$$\int_{\mathbb{R}^N\setminus A_{\lambda}} Hu_n^2 \leqslant \frac{\kappa}{\lambda^{1+\beta}} \int_{\mathbb{R}^N} \frac{|u_n(x)|^2}{|x|^2} dx \leqslant \frac{\kappa}{\lambda^{1+\beta}} \left(\frac{2}{N-2}\right)^2 \int_{\mathbb{R}^N} |\nabla u_n|^2.$$

Since $(u_n)_n$ is bounded in $H^1_{V \mathcal{H}}(\mathbb{R}^N)$, for every $\delta > 0$, there exists $\bar{\lambda} \ge 0$ such that

$$\limsup_{n \to \infty} \int_{\mathbb{R}^N \setminus A_{\bar{\lambda}}} Hu_n^2 < \delta.$$
(15)

Now we claim that for all $\delta > 0$, there exists $\tilde{\lambda} > 0$ such that

$$\limsup_{n \to \infty} \int_{\mathbb{R}^N \setminus A_{\bar{\lambda}}} V u_n^2 < \delta.$$
(16)

We only sketch the proof, since the arguments are similar to those in [8, Lemma 6]. Since $\overline{\Lambda} \subset \mathbb{R}^N \setminus \{0\}$ is compact, there exists $\lambda_0 \ge 0$ such that

$$\bar{\Lambda} \subset A_{\lambda_0}$$

Let $\zeta \in C^{\infty}(\mathbb{R})$ be such that $0 \leq \zeta \leq 1$ and

$$\zeta(s) = \begin{cases} 0 & \text{if } |s| \leq \frac{1}{2}, \\ 1 & \text{if } |s| \geq 1. \end{cases}$$

Define a cut-off function $\eta_{\lambda} \in C^{\infty}(\mathbb{R}^N, \mathbb{R})$ by

$$\eta_{\lambda}(x) := \zeta\left(\frac{\log|x|}{\lambda}\right).$$

Since $\langle J'_{\varepsilon}(u_n), \eta_{\lambda}u_n \rangle = o(1)$ as $n \to \infty$, we deduce that

$$\int_{\mathbb{R}^{N}} \left(\varepsilon^{2} |\nabla u_{n}|^{2} + V u_{n}^{2} \right) \eta_{\lambda} = \int_{\mathbb{R}^{N}} g_{\varepsilon} \left(x, u_{n}(x) \right) u_{n}(x) \eta_{\lambda}(x) \, dx - \varepsilon^{2} \int_{\mathbb{R}^{N}} u_{n} \nabla u_{n} \cdot \nabla \eta_{\lambda} + o(1), \quad (17)$$

as $n \to \infty$. If $\lambda \ge 2\lambda_0$, $\eta_{\lambda} = 0$ on Λ and it follows from (g_3) that

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$$\int_{\mathbb{R}^{N}} g_{\varepsilon} (x, u_{n}(x)) u_{n}(x) \eta_{\lambda}(x) dx \leq \int_{\mathbb{R}^{N}} (\varepsilon^{2} H + \mu V) u_{n}^{2} \eta_{\lambda}.$$
(18)

On the other hand, using Hardy's inequality, we can show as in [8] that

$$\left|\int_{\mathbb{R}^N} u_n \nabla u_n \cdot \nabla \eta_\lambda\right| \leqslant \frac{C}{\lambda} \|u_n\|_{\varepsilon}^2.$$
(19)

Combining (17), (18) and (19), we get, for every $\lambda \ge 2\lambda_0$,

$$\int_{\mathbb{R}^N \setminus A_{\lambda}} \left(\varepsilon^2 |\nabla u_n|^2 + (1-\mu) V u_n^2 \right) \eta_{\lambda} \leqslant \int_{\mathbb{R}^N} \left(\varepsilon^2 |\nabla u_n|^2 + (1-\mu) V u_n^2 \right) \eta_{\lambda}$$
$$\leqslant \int_{\mathbb{R}^N} \varepsilon^2 H u_n^2 \eta_{\lambda} + \frac{C}{\lambda} \|u_n\|_{\varepsilon}^2 + o(1).$$

By (15), for λ large enough,

$$\limsup_{n\to\infty}\int_{\mathbb{R}^N}Hu_n^2\eta_\lambda\leqslant\limsup_{n\to\infty}\int_{\mathbb{R}^N\setminus A_{\bar{\lambda}}}Hu_n^2<\frac{\delta}{2}$$

the claim follows.

Conclusion. We can write

$$\|u_n - u\|_{\varepsilon}^2 = J_{\varepsilon}'(u_n)(u_n - u) - J_{\varepsilon}'(u)(u_n - u) + \int_{\mathbb{R}^N} \left(g_{\varepsilon}(x, u_n(x)) - g_{\varepsilon}(x, u(x))\right) \left(u_n(x) - u(x)\right) dx.$$
(20)

We notice that the first two terms in the right-hand side tend to 0 as $n \to \infty$. Fix $\delta > 0$ and let $\lambda > 0$ be such that (15) and (16) hold. We evaluate the integral in the third term of (20) separately on Λ , $A_{\lambda} \setminus \Lambda$ and $\mathbb{R}^N \setminus A_{\lambda}$, where $\lambda = \max\{\tilde{\lambda}, \bar{\lambda}\}$.

By (g_2) , one has $|g_{\varepsilon}(x, u_n(x))| \leq C |u_n(x)|^p$. By Rellich Theorem, the embedding $H^1_{V,\mathcal{H}}(\Lambda) \hookrightarrow L^q(\Lambda)$ is compact for all q > 1 such that $\frac{1}{q} > \frac{1}{2} - \frac{1}{N-k}$. We can thus assume that $u_n \to u$ in $L^{p+1}(\Lambda)$. We deduce that $g_{\varepsilon}(x, u_n) \to g_{\varepsilon}(x, u)$ in $L^q(\Lambda)$ as $n \to \infty$, where $q := \frac{p+1}{p}$. We conclude from Hölder's inequality that

$$\int_{\Lambda} \left(g_{\varepsilon} \left(x, u_n(x) \right) - g_{\varepsilon} \left(x, u(x) \right) \right) \left(u_n(x) - u(x) \right) dx \to 0, \quad \text{as } n \to \infty.$$

By (g_3) , one has $|g_{\varepsilon}(x, u_n(x))| \leq (\varepsilon^2 H(x) + \mu V(x))|u_n(x)|$ for $x \in A_{\lambda} \setminus \Lambda$. By Rellich Theorem, we can assume that $u_n \to u$ in $L^2(A_{\lambda} \setminus \Lambda)$. We deduce that $g_{\varepsilon}(x, u_n) \to g_{\varepsilon}(x, u)$ in $L^2(A_{\lambda} \setminus \Lambda)$ as $n \to \infty$. We conclude as above that

$$\int_{A_{\lambda}\setminus A} \left(g_{\varepsilon} \left(x, u_{n}(x) \right) - g_{\varepsilon} \left(x, u(x) \right) \right) \left(u_{n}(x) - u(x) \right) dx \to 0, \quad \text{as } n \to \infty.$$

Finally, using (g_3) , (15) and (16), we obtain

$$\begin{split} &\limsup_{n \to \infty} \int_{\mathbb{R}^N \setminus A_{\lambda}} \left| g_{\varepsilon} \left(x, u_n(x) \right) - g_{\varepsilon} \left(x, u(x) \right) \right| \left| u_n(x) - u(x) \right| dx \\ &\leqslant 2 \limsup_{n \to \infty} \int_{\mathbb{R}^N \setminus A_{\lambda}} \left(\left| g_{\varepsilon} \left(x, u_n(x) \right) u_n(x) \right| + \left| g_{\varepsilon} \left(x, u(x) \right) u(x) \right| \right) dx \\ &\leqslant 2 \limsup_{n \to \infty} \int_{\mathbb{R}^N \setminus A_{\lambda}} \left(\varepsilon^2 H + \mu V \right) \left(u_n^2 + u^2 \right) \\ &\leqslant 4(1 + \mu) \delta, \end{split}$$

since $\lambda \ge \overline{\lambda}$ and $\lambda \ge \overline{\lambda}$.

Since $\delta > 0$ is arbitrary, we conclude

$$\lim_{n\to\infty}\|u_n-u\|_{\varepsilon}=0$$

which ends the proof. \Box

We can now state an existence theorem for the penalized problem ($\mathcal{P}_{\varepsilon}$). The proof follows from standard arguments.

Theorem 4. Let $g: \mathbb{R} \times \mathbb{R}^+ \to \mathbb{R}$ be a Carathéodory function satisfying $(g_1)-(g_4)$ and $V \in C(\mathbb{R}^N \setminus \{0\})$ be a nonnegative function. Then, for all $\varepsilon > 0$, the functional J_{ε} possesses a nontrivial critical point $u_{\varepsilon} \in H^1_{V,\mathcal{H}}(\mathbb{R}^N)$, which is characterized by

$$c_{\varepsilon} := J_{\varepsilon}(u_{\varepsilon}) = \inf_{\substack{u \in H^1_{V,\mathcal{H}}(\mathbb{R}^N) \setminus \{0\}}} \max_{t>0} J_{\varepsilon}(tu).$$
(21)

The function u_{ε} found in Theorem 4 is called a *least-energy solution* of $(\mathcal{P}_{\varepsilon})$. By standard regularity theory, if $u \in H^1_{loc}(\mathbb{R}^N)$ is a solution of $(\mathcal{P}_{\varepsilon})$, then $u \in W^{2,q}_{loc}(\mathbb{R}^N)$ for every $q \in (1, \infty)$. In particular, $u \in C^{1,\alpha}_{loc}(\mathbb{R}^N)$ for every $\alpha \in (0, 1)$. Since g_{ε} is not continuous, we cannot achieve a better regularity. Notice also that, by the strong maximum principle, any nontrivial nonnegative solution $u \in C^{1,\alpha}_{loc}(\mathbb{R}^N)$ of $(\mathcal{P}_{\varepsilon})$ is positive in \mathbb{R}^N .

4. Asymptotics of solutions

In this section we study the asymptotic behavior as $\varepsilon \to 0$ of the solution found in Theorem 4. We follow closely the arguments in [8, §6]. We first prove an energy estimate which is the counterpart of [8, Lemma 12]. Let $\mathbb{R}^{N-k}_+ := \mathbb{R}^{N-k-1} \times \mathbb{R}^+$.

Proposition 4.1 (Upper estimate of the critical value). Suppose that the assumptions of Theorem 4 are satisfied. For ε small enough, the critical value c_{ε} defined in (21) satisfies

$$c_{\varepsilon} \leqslant \varepsilon^{N-k} \Big(\omega_k \inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M} + o(1) \Big) \quad \text{as } \varepsilon \to 0,$$

where ω_k is the volume of the unit sphere in \mathbb{R}^{k+1} . Moreover, the solution u_{ε} of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4 satisfies, for some C > 0,

$$\|u_{\varepsilon}\|_{\varepsilon}^2 \leq C\varepsilon^{N-k}.$$

Proof. Let $x_0 = (0, x''_0) \in \Lambda \cap \mathcal{H}^{\perp}$ be such that $\mathcal{M}(x_0) = \inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M}$. Denote by I_0 the functional defined by (6) with $a = V(x_0)$ and $b = K(x_0)$ and w a ground state of (5). Take $\eta \in \mathcal{D}(\mathbb{R}^{N-k}_+)$ to be a cut-off function such that $0 \leq \eta \leq 1$, $\eta = 1$ in a neighborhood of $(0, |x''_0|)$ and $\|\nabla \eta\|_{\infty} \leq C$. Consider the test function

$$u(x) := \eta(x', |x''|) w\left(\frac{x'}{\varepsilon}, \frac{|x''| - |x_0''|}{\varepsilon}\right).$$

Setting

$$u(x',x'') =: v\left(\frac{x'}{\varepsilon},\frac{|x''|-|x_0''|}{\varepsilon}\right),$$

we compute by a change of variable

$$\begin{split} J_{\varepsilon}(tu) &= \omega_{k} \frac{t^{2}}{2} \int\limits_{\mathbb{R}^{N-k-1}} \int\limits_{-\frac{|x_{0}'|}{\varepsilon}}^{\infty} \left(|\nabla v|^{2} + V\left(\varepsilon y', \varepsilon \rho + |x_{0}''|\right) v^{2} \right) \left(\varepsilon \rho + |x_{0}''|\right)^{k} \varepsilon \, d\rho \, \varepsilon^{N-k-1} \, dy' \\ &- \omega_{k} \int\limits_{\mathbb{R}^{N-k-1}} \int\limits_{-\frac{|x_{0}''|}{\varepsilon}}^{\infty} G\left(\varepsilon y', \varepsilon \rho + |x_{0}''|, tv\right) \left(\varepsilon \rho + |x_{0}''|\right)^{k} \varepsilon \, d\rho \, \varepsilon^{N-k-1} \, dy'. \end{split}$$

For ε small enough, we obtain

$$\varepsilon^{-(N-k)} J_{\varepsilon}(tu) \leqslant \omega_k |x_0''|^k I_0(tw) + o(1).$$
⁽²²⁾

We deduce from (21) that

$$\varepsilon^{-(N-k)}c_{\varepsilon} \leq \max_{t>0} \varepsilon^{-(N-k)} J_{\varepsilon}(tu)$$
$$\leq \omega_k |\mathbf{x}_0''|^k \max_{t>0} I_0(tw) + o(1)$$
$$= \omega_k \mathcal{M}(\mathbf{x}_0) + o(1),$$

which is the desired conclusion. \Box

Proposition 4.2 (No uniform convergence to 0 in Λ). Suppose that the assumptions of Theorem 4 are satisfied and let $(u_{\varepsilon})_{\varepsilon} \subset H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ obtained in Theorem 4. Then there exists $\delta > 0$ such that

$$||u_{\varepsilon}||_{L^{\infty}(\Lambda)} \geq \delta.$$

Proof. Suppose by contradiction that $||u_{\varepsilon}||_{L^{\infty}(\Lambda)} \to 0$ as $\varepsilon \to 0$. Then, (f_1) implies that, for all ε sufficiently small, $Kf(u_{\varepsilon}) \leq \mu V u_{\varepsilon}$ in Λ . By (g₃), we deduce that

$$-\varepsilon^2(\Delta u_{\varepsilon} + Hu_{\varepsilon}) + (1-\mu)Vu_{\varepsilon} \leq 0 \quad \text{in } \mathbb{R}^N.$$

Proposition 3.1 then implies that $u_{\varepsilon} \equiv 0$ for all ε sufficiently small, which is impossible. \Box

By the symmetry imposed on u_{ε} , one can write $u_{\varepsilon}(x', x'') = \tilde{u}_{\varepsilon}(x', |x''|)$ with $\tilde{u}_{\varepsilon} : \mathbb{R}^{N-k}_{+} \to \mathbb{R}$. Since the H^{1}_{V} -norm of u_{ε} is of the order $\varepsilon^{(N-k)/2}$, it is natural to rescale $\tilde{u}_{\varepsilon}(x', |x''|)$ as $\tilde{u}_{\varepsilon}(x'_{\varepsilon} + \varepsilon y', |x''_{\varepsilon}| + \varepsilon |y''|)$ around a well-chosen family of points $x_{\varepsilon} = (x'_{\varepsilon}, x''_{\varepsilon}) \in \mathbb{R}^{N}$.

The next lemma shows that the sequences of rescaled solutions converge, up to a subsequence, in $C^1_{loc}(\mathbb{R}^{N-k})$ to a function $v \in H^1(\mathbb{R}^{N-k})$.

Lemma 4.3. Suppose that the assumptions of Theorem 4 are satisfied. Let $u_{\varepsilon} \in H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N})$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4, $(\varepsilon_{n})_{n} \subset \mathbb{R}^{+}$ and $(x_{n})_{n} \subset \mathbb{R}^{N}$ be sequences such that $\varepsilon_{n} \to 0$ and $x_{n} = (x'_{n}, x''_{n}) \to \overline{x} = (\overline{x}', \overline{x}'') \in \overline{A}$ as $n \to \infty$. Set

$$\Omega_n := \mathbb{R}^{N-k-1} \times \left] - \frac{|x_n''|}{\varepsilon_n}, + \infty \right[$$

and let $v_n : \Omega_n \to \mathbb{R}$ be defined by

$$v_n(y,z) := \tilde{u}_{\varepsilon_n} \left(x'_n + \varepsilon_n y, \left| x''_n \right| + \varepsilon_n z \right), \tag{23}$$

where $\tilde{u}_{\varepsilon_n}: \mathbb{R}^{N-k}_+ \to \mathbb{R}$ is such that $u_{\varepsilon_n}(x', x'') = \tilde{u}_{\varepsilon_n}(x', |x''|)$. Then, there exists $v \in H^1(\mathbb{R}^{N-k})$ such that, along a subsequence that we still denote by $(v_n)_n$,

$$v_n \xrightarrow{C^1_{\text{loc}}(\mathbb{R}^{N-k})} v$$

Proof. First observe that each v_n solves the equation

$$-\Delta v_n - \frac{\varepsilon_n k}{z} \frac{\partial v_n}{\partial z} + V(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z) v_n = g_{\varepsilon_n}(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z, v_n),$$
(24)

in Ω_n . We infer from Proposition 4.1 that for all $n \in \mathbb{N}$,

$$\int_{\Omega_n} \left(\left| \nabla v_n(y,z) \right|^2 + V\left(x'_n + \varepsilon_n y, \left| x''_n \right| + \varepsilon_n z \right) \left| v_n(y,z) \right|^2 \right) dy \, dz \leqslant C,$$

with C > 0 independent of *n*.

Define a cut-off function $\eta_R \in \mathcal{D}(\mathbb{R}^{N-k})$ such that $0 \leq \eta_R \leq 1$, $\eta_R(x) = 1$ if $|x| \leq R/2$, $\eta_R(x) = 0$ if $|x| \geq R$ and $\|\nabla \eta_R\|_{\infty} \leq C/R$ for some C > 0. Choose $(R_n)_n$ such that $R_n \to \infty$ and $\varepsilon_n R_n \to 0$. Since $\bar{x} \in \Lambda$ and $\bar{\Lambda} \cap \mathcal{H} = \emptyset$, one has $\varepsilon_n R_n \leq |x_n''|$ if n is large enough. Define $w_n \in H^1_{\text{loc}}(\mathbb{R}^{N-k})$ by

$$w_n(y) := \eta_{R_n}(y) v_n(y).$$

On the one hand, we notice that

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$$\int_{\mathbb{R}^{N-k}} w_n^2 \leqslant \int_{B(0,R_n)} v_n^2 \leqslant \frac{1}{\inf_{B(x_n,\varepsilon_n R_n)} V} \int_{\Omega_n} V(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z) |v_n(y,z)|^2 dy dz$$

Since V is positive on $\overline{\Lambda}$ and continuous on \mathbb{R}^N , the convergence of x_n to a point in $\overline{\Lambda}$ implies that

$$\int_{\mathbb{R}^{N-k}} w_n^2 \leqslant C.$$
(25)

On the other hand, we compute in the same way as in [8, Lemma 13]

$$\int_{\mathbb{R}^{N-k}} |\nabla w_n|^2 \leqslant C \|v_n\|_{H^1(B(0,R_n))}^2.$$
 (26)

Since

$$\|v_n\|_{H^1(B(0,R_n))} \leq C \|u_{\varepsilon_n}\|_{\varepsilon},$$

we deduce from (25) and (26) that $(w_n)_n$ is bounded in $H^1(\mathbb{R}^{N-k})$. Since w_n solves Eq. (24) on $B(0, R_n)$ for all n, classical regularity estimates yield that for every R > 0 and every q > 1,

$$\sup_{n \in \mathbb{N}} \|v_n\|_{W^{2,q}(B(0,R))} < \infty.$$
(27)

Up to a subsequence, we can now assume that $(w_n)_n$ converges weakly in $H^1(\mathbb{R}^{N-k})$ to some function $v \in H^1(\mathbb{R}^{N-k})$. By (27), for every compact $K \subset \mathbb{R}^{N-k}$, w_n converges to v in $C^1(K)$. Moreover, for n large enough, $w_n = v_n$ in K so that $v_n \to v$ in $C^1(K)$. \Box

In the next two lemmas, we will estimate from below the action of u_{ε} inside and outside neighborhoods of points. Since we expect the concentration set to be a *k*-sphere in \mathbb{R}^N , the following distance will be useful. For $x, y \in \mathbb{R}^N$, let

$$d_{\mathcal{H}}(x, y) := \sqrt{|x' - y'|^2 + (|x''| - |y''|)^2}.$$

Thus $d_{\mathcal{H}}(x, y)$ represents the distance between the *k*-spheres centered at x' and y', and of radius |x''| and |y''| respectively. We denote by $B_{\mathcal{H}}$ the balls for the distance $d_{\mathcal{H}}$, i.e.,

$$B_{\mathcal{H}}(x,r) = \left\{ y \in \mathbb{R}^N \colon d_{\mathcal{H}}(x,y) < r \right\}.$$

Lemma 4.4. Suppose that the assumptions of Theorem 4 are satisfied. Let $u_{\varepsilon} \in H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N})$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4, $(\varepsilon_{n})_{n} \subset \mathbb{R}^{+}$ and $(x_{n})_{n} \subset \mathbb{R}^{n}$ be sequences such that $\varepsilon_{n} \to 0$ and $x_{n} = (x'_{n}, x''_{n}) \to \overline{x} = (\overline{x}', \overline{x}'') \in \overline{A}$ as $n \to \infty$. If

$$\liminf_{n \to \infty} u_{\varepsilon_n}(x_n) > 0, \tag{28}$$

then we have, up to a subsequence,

$$\liminf_{R\to\infty}\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}\left(\int\limits_{T_n(R)}\frac{1}{2}(\varepsilon_n^2|\nabla u_{\varepsilon_n}|^2+Vu_{\varepsilon_n}^2)-G_{\varepsilon_n}(x,u_{\varepsilon_n})\right)\geq\omega_k\mathcal{M}(\bar{x}),$$

where $T_n(R) := B_{\mathcal{H}}(x_n, \varepsilon_n R)$.

Proof. Let v_n be defined by (23). Passing to a subsequence if necessary, we may assume that there exists $v \in H^1(\mathbb{R}^{N-k})$ such that $v_n \to v$ in $C^1_{loc}(\mathbb{R}^{N-k})$. Since Λ is smooth, we can also assume that the sequence of characteristic functions $\chi_n(y, z) = \chi_\Lambda(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z)$ converges almost everywhere to a measurable function χ satisfying $0 \leq \chi \leq 1$. We then deduce that v solves the limiting equation

$$-\Delta v + V(\bar{x})v = \tilde{g}(y, v)$$
 in \mathbb{R}^{N-k} ,

where

$$\tilde{g}(y,s) := \chi(y)K(\bar{x})f(s) + (1 - \chi(y))\min\{\mu V(\bar{x})s, K(\bar{x})f(s)\}.$$

By (28), we know that $v(0) = \lim_{n \to \infty} v_n(0) > 0$, so that v is not identically zero.

It was shown in [8, Lemma 14] that

$$\begin{split} \liminf_{R \to \infty} \liminf_{n \to \infty} & \int_{B(0,R)} \left(\frac{1}{2} \left(\left| \nabla v_n(y,z) \right|^2 + V \left(x'_n + \varepsilon_n y, \left| x''_n \right| + \varepsilon_n z \right) \left| v_n(y,z) \right|^2 \right) \right. \\ & - G_{\varepsilon_n} \left(x'_n + \varepsilon_n y, \left| x''_n \right| + \varepsilon_n z, v_n(y,z) \right) \right) dz \, dy \\ & \ge \frac{1}{2} \int_{\mathbb{R}^{N-k}} \left(\left| \nabla v \right|^2 + V \left(\bar{x} \right) v^2 \right) - \int_{\mathbb{R}^{N-k}} \tilde{G} \left(y, v(y) \right) dy, \end{split}$$

where $\tilde{G}(x, s) := \int_0^s \tilde{g}(x, \sigma) d\sigma$.

Set $B_n(R) := B((x', |x''|), \varepsilon_n R) \subset \mathbb{R}^{N-k}$. By a computation similar to the one leading to (22), we have

$$\begin{split} &\int\limits_{T_n(R)} \left(\frac{1}{2} \left(\varepsilon_n^2 |\nabla u_{\varepsilon_n}(x)|^2 + V(x) |u_{\varepsilon_n}(x)|^2 \right) - G_{\varepsilon_n}(x, u_{\varepsilon_n}(x)) \right) dx \\ &= \omega_k \int\limits_{B_n(R)} \left(\frac{1}{2} \left(\varepsilon_n^2 |\nabla \tilde{u}_{\varepsilon_n}(x', r)|^2 + V(x', r) |\tilde{u}_{\varepsilon_n}(x', r)|^2 \right) - G_{\varepsilon_n}(x', r, \tilde{u}_{\varepsilon_n}(x', r)) \right) r^k dr dx' \\ &= \omega_k |\bar{x}''|^k \varepsilon_n^{N-k} \int\limits_{B(0,R)} \left(\frac{1}{2} \left(|\nabla v_n(y, z)|^2 + V(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z) |v_n(y, z)|^2 \right) \right. \\ &- G_{\varepsilon_n}(x'_n + \varepsilon_n y, |x''_n| + \varepsilon_n z, v_n(y, z)) \right) dz \, dy + o(1). \end{split}$$

The conclusion follows. \Box

Lemma 4.5. Suppose that the assumptions of Theorem 4 are satisfied. Let $u_{\varepsilon} \in H^{1}_{V,\mathcal{H}}(\mathbb{R}^{N})$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4, $(\varepsilon_{n})_{n} \subset \mathbb{R}^{+}$ and $(x_{n}^{i})_{n} \subset \mathbb{R}^{N}$ be sequences such that $\varepsilon_{n} \to 0$ and for $1 \leq i \leq M$, $x_{n}^{i} \to \overline{x}^{i} \in \overline{A}$ as $n \to \infty$. Then, up to a subsequence, we have

$$\liminf_{R\to\infty}\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}\left(\int\limits_{\mathbb{R}^N\setminus\mathcal{T}_n(R)}\frac{1}{2}(\varepsilon_n^2|\nabla u_{\varepsilon_n}|^2+Vu_{\varepsilon_n}^2)-G_{\varepsilon_n}(x,u_{\varepsilon_n})\right)\geq 0,$$

where $\mathcal{T}_n(R) := \bigcup_{i=1}^M B_{\mathcal{H}}(x_n^i, \varepsilon_n R).$

Proof. See [8, Lemma 15]. □

Proposition 4.6 (Lower estimate of the critical value). Suppose that the assumptions of Theorem 4 are satisfied. Let $u_{\varepsilon} \in H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4, $(\varepsilon_n)_n \subset \mathbb{R}^+$ and $(x_n^i)_n \subset \mathbb{R}^N$ be sequences such that $\varepsilon_n \to 0$ and for $1 \leq i \leq M, x_n^i \to \overline{x}^i \in \overline{\Lambda}$ as $n \to \infty$. If for every $1 \leq i < j \leq M$, we have

$$\limsup_{n \to \infty} \frac{d_{\mathcal{H}}(x_n^i, x_n^j)}{\varepsilon_n} = \infty$$

and if for every $1 \leq i \leq M$,

$$\liminf_{n\to\infty}u_{\varepsilon_n}(x_n^i)>0,$$

then the critical value c_{ε} defined in (21) satisfies

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n} \ge \omega_k\sum_{i=1}^M \mathcal{M}(\bar{x}^i).$$

Proof. This is a consequence of the two previous lemmas, see [8, Proposition 16] for the details. \Box

The following proposition is the key result for the next section.

Proposition 4.7 (Uniform convergence to 0 outside small balls). Suppose that the assumptions of Theorem 4 are satisfied and that Λ satisfies the assumptions of Section 2.4. Let $(u_{\varepsilon})_{\varepsilon} \subset H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ obtained in Theorem 4. If $(x_{\varepsilon})_{\varepsilon>0} \subset \Lambda$ is such that

$$\liminf_{\varepsilon\to 0} u_{\varepsilon}(x_{\varepsilon}) > 0,$$

then

- (i) $\lim_{\varepsilon \to 0} \mathcal{M}(x_{\varepsilon}) = \inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M},$
- (ii) $\limsup_{\varepsilon \to 0} \frac{\operatorname{dist}(x_{\varepsilon}, \mathcal{H}^{\perp})}{\varepsilon} < \infty$,
- (iii) $\liminf_{\varepsilon \to 0} d_{\mathcal{H}}(x_{\varepsilon}, \partial \Lambda) > 0$,

(iv) for every $\delta > 0$, there exist $\varepsilon_0 > 0$ and R > 0 such that, for every $\varepsilon \in (0, \varepsilon_0)$,

$$\|u_{\varepsilon}\|_{L^{\infty}(\Lambda\setminus B_{\mathcal{H}}(x_{\varepsilon},\varepsilon R))} \leq \delta.$$

Proof. The first assertion is a direct consequence of Propositions 4.1 and 4.6, see [8, Proposition 33] for the details.

For the second assertion, since $\overline{\Lambda}$ is compact, we can assume by contradiction that there exist sequences $(\varepsilon_n)_n \subset \mathbb{R}^+$ and $(x_n)_n \subset \Lambda$ such that $\varepsilon_n \to 0$, $x_n \to \overline{x} \in \overline{\Lambda}$,

$$\liminf_{n\to\infty}u_{\varepsilon_n}(x_n)>0,$$

and

$$\limsup_{n\to\infty}\frac{\operatorname{dist}(x_n,\mathcal{H}^{\perp})}{\varepsilon_n}=\infty$$

If k = N - 2, let $R \in \mathbf{O}(N)$ denote the reflexion with respect to \mathcal{H}^{\perp} . By definition of $H^1_{V,\mathcal{H}}(\mathbb{R}^N)$, $u \circ R = u$, and thus

$$\liminf_{n\to\infty}u_{\varepsilon_n}\big(R(x_n)\big)>0.$$

The preceding assumption implies that $\lim_{n\to\infty} \frac{d_{\mathcal{H}}(x_n, R(x_n))}{\varepsilon_n} = \infty$. By Proposition 4.6, we obtain

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n} \ge \omega_k \big(\mathcal{M}(\bar{x}) + \mathcal{M}\big(R(\bar{x})\big)\big) \ge 2\omega_k \inf_{\Lambda} \mathcal{M},$$

which, together with Proposition 4.1

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n}\leqslant\omega_k\inf_{\Lambda\cap\mathcal{H}^{\perp}}\mathcal{M}$$

is in contradiction with (10).

In the case where k < N - 2, since $\inf_{\Lambda} \mathcal{M} > 0$, choose $\ell \in \mathbb{N}$ such that

$$\inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M} < \ell \inf_{\Lambda} \mathcal{M}.$$
⁽²⁹⁾

There exist isometries R_1, \ldots, R_l of \mathbb{R}^N such that $R_i(\mathcal{H}) = \mathcal{H}$ and $R_i(\bar{x}) \neq R_j(\bar{x})$, for every $i, j \in \{1, \ldots, \ell\}$ with $i \neq j$. One has hence

$$\limsup_{n\to\infty}\frac{d_{\mathcal{H}}(R_i(x_n),R_j(x_n))}{\varepsilon_n}=\infty.$$

By Proposition 4.6, we get

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n} \ge \omega_k \sum_{i=1}^{\ell} \mathcal{M}(R_i(\bar{x})) \ge l\omega_k \inf_{\Lambda} \mathcal{M},$$

so that, in view of the upper estimate of Proposition 4.1, we have a contradiction with (29).

For the third assertion, suppose by contradiction that there exist sequences $(\varepsilon_n)_n \subset \mathbb{R}^+$ and $(x_n)_n \subset \mathbb{R}^N$ such that $\varepsilon_n \to 0$,

$$\liminf_{n\to\infty}u_{\varepsilon_n}(x_n)>0,$$

and $x_n \to \bar{x} \in \partial \Lambda$. We have just proven that $\bar{x} \in \mathcal{H}^{\perp}$. By Proposition 4.6, we have

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n} \geqslant \omega_k \mathcal{M}(\bar{x}) \geqslant \omega_k \inf_{\partial \Lambda \cap \mathcal{H}^{\perp}} \mathcal{M}.$$

This inequality, along with Proposition 4.1, contradicts (9).

In order to obtain the last assertion, suppose by contradiction that there exist sequences $(\varepsilon_n)_n \subset \mathbb{R}^+$, $(x_n)_n$ and $(y_n)_n \subset \Lambda$ such that $\varepsilon_n \to 0$,

$$u_{\varepsilon_n}(y_n) \geq \delta$$
,

and

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$$\lim_{n\to\infty}\frac{d_{\mathcal{H}}(x_n,\,y_n)}{\varepsilon_n}=\infty.$$

Up to a subsequence, we can assume that $x_n \to \bar{x} \in \Lambda$ and $y_n \to \bar{y} \in \Lambda$. In view of the second assertion, one has $\bar{x} \in \mathcal{H}^{\perp}$ and $\bar{y} \in \mathcal{H}^{\perp}$. Therefore, by Proposition 4.6,

$$\liminf_{n\to\infty}\varepsilon_n^{-(N-k)}c_{\varepsilon_n} \ge \omega_k \big(\mathcal{M}(\bar{x}) + \mathcal{M}(\bar{y})\big) \ge 2\omega_k \inf_{\Lambda \cap \mathcal{H}^{\perp}} \mathcal{M}.$$

In view of the assumption of (9), this would contradict Proposition 4.1. \Box

5. Barrier functions

5.1. Linear inequation outside small balls

In this section we prove that for ε small enough, the solutions of the penalized problem ($\mathcal{P}_{\varepsilon}$) are also solutions of the initial problem (2). We follow the arguments of [17]. First we notice that the solutions of ($\mathcal{P}_{\varepsilon}$) satisfy a linear inequation outside small balls.

Lemma 5.1. Suppose that the assumptions of Proposition 4.7 are satisfied and let $(u_{\varepsilon})_{\varepsilon>0} \subset H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4 and $(x_{\varepsilon})_{\varepsilon>0} \subset \Lambda$ be such that

$$\liminf_{\varepsilon\to 0} u_{\varepsilon}(x_{\varepsilon}) > 0.$$

Then there exist $\rho > 0$ and $\varepsilon_0 > 0$ such that for all $\varepsilon \in (0, \varepsilon_0)$,

$$-\varepsilon^{2}(\Delta u_{\varepsilon} + Hu_{\varepsilon}) + (1-\mu)Vu_{\varepsilon} \leq 0 \quad in \mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon\rho).$$
(30)

Proof. Set

$$\eta := \inf_{x \in \Lambda} \frac{\mu V(x)}{K(x)}.$$

Since *V* and *K* are bounded positive continuous functions on $\bar{\Lambda}$, $\eta > 0$. By (f_1) , there exists $\delta > 0$ such that

$$\frac{f(s)}{s} \leqslant \eta \quad \text{for all } s \leqslant \delta.$$

By Proposition 4.7, we can find $\varepsilon_0 > 0$ and $\rho > 0$ such that for all $\varepsilon \in (0, \varepsilon_0]$, one has

$$u_{\varepsilon}(x) \leq \delta$$
 for all $x \in \Lambda \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho)$.

Hence

$$K(x)f(u_{\varepsilon}(x)) \leq \mu V(x)u_{\varepsilon}(x) \text{ in } \Lambda \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho).$$

We conclude that

$$-\varepsilon^2 \Delta u_{\varepsilon} + (1-\mu) V u_{\varepsilon} \leqslant -\varepsilon^2 \Delta u_{\varepsilon} + V u_{\varepsilon} - K f(u_{\varepsilon}) = 0 \quad \text{in } \Lambda \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho).$$

The fact that u_{ε} satisfies (30) in $\mathbb{R}^N \setminus \Lambda$ follows directly from the definition of the penalized nonlinearity. \Box

This lemma suggests that we can compare the solution u_{ε} with supersolutions of the operator $-\varepsilon^2(\Delta + H) + (1 - \mu)V$ in order to obtain decay estimates of u_{ε} .

5.2. Comparison functions

The next lemma provides a minimal positive solution of the operator $-\Delta - H$ in $\mathbb{R}^N \setminus \overline{\Lambda}$.

Lemma 5.2. For every $\varepsilon > 0$, there exists $\Psi_{\varepsilon} \in C^2((\mathbb{R}^N \setminus \{0\}) \setminus \Lambda)$ such that

$$\begin{cases} -\varepsilon^2 (\Delta \Psi_{\varepsilon} + H\Psi_{\varepsilon}) + (1-\mu)V\Psi_{\varepsilon} = 0 & \text{in } \mathbb{R}^N \setminus \bar{\Lambda}, \\ \Psi_{\varepsilon} = 1 & \text{on } \partial \Lambda, \end{cases}$$

and

$$\int_{\mathbb{R}^N \setminus \Lambda} \left(\left| \nabla \Psi_{\varepsilon}(x) \right|^2 + \frac{|\Psi_{\varepsilon}(x)|^2}{|x|^2} \right) dx < \infty.$$
(31)

Moreover, there exists C > 0 *such that, for every* $x \in \mathbb{R}^N \setminus \Lambda$ *and every* $\varepsilon > 0$ *,*

$$0 < \Psi_{\varepsilon}(x) \leqslant \frac{C}{(1+|x|)^{N-2}}.$$
(32)

Proof. The function Ψ_{ε} is obtained by minimizing

$$\int_{\mathbb{R}^N\setminus\Lambda} \left(\varepsilon^2 \left(|\nabla u|^2 - Hu^2\right) + (1-\mu)Vu^2\right) dx$$

on the set

$$\{u \in H^1_V(\mathbb{R}^N): u = 1 \text{ on } \partial \Lambda\}.$$

By classical elliptic regularity theory, $\Psi_{\varepsilon} \in C^2((\mathbb{R}^N \setminus \{0\}) \setminus \Lambda)$. The estimate (31) follows from (11). In order to obtain the estimate (32) consider the problem

$$\begin{cases} -\Delta \Psi - H\Psi = 0 & \text{in } \mathbb{R}^N \setminus \bar{\Lambda}, \\ \Psi = 1 & \text{on } \partial \Lambda. \end{cases}$$

We have just proved that this problem has a solution $\Psi \in C^2((\mathbb{R}^N \setminus \{0\}) \setminus \Lambda)$ such that

$$\int_{\mathbb{R}^N \setminus A} \left(\left| \nabla \Psi(x) \right|^2 + \frac{|\Psi(x)|^2}{|x|^2} \right) dx < \infty.$$
(33)

Now set for $\rho \in (0, 1)$ and $x \in B(0, \rho)$,

$$W(x) := (N-2)\beta - \kappa \left(\log \frac{1}{|x|}\right)^{-\beta}$$

We compute that

$$-\Delta W(x) = \frac{\kappa \beta}{|x|^2} \left[(N-2) \left(\log \frac{1}{|x|} \right)^{-(1+\beta)} + (1+\beta) \left(\log \frac{1}{|x|} \right)^{-(2+\beta)} \right].$$

Since for $|x| \leq 1$,

$$H(x) \leqslant \frac{\kappa}{(|x|^2 \log \frac{1}{|x|})^{1+\beta}}$$

the function W is a supersolution of $-\Delta - H$ in B(0, 1). Moreover, if one takes $\rho < 1$ such that

$$(N-2)\beta\left(\log\frac{1}{\rho}\right)^{\beta} > \kappa,$$

W is positive on $\partial B(0, \rho)$. In view of (33) Proposition 3.1 implies that Ψ is bounded from above by a positive multiple of *W* in $B(0, \rho)$. Since Ψ is continuous and *W* is bounded in B(0, 1), we obtain that Ψ is bounded in B(0, 1). By similarly considering

$$W(x) := \frac{1}{|x|^{N-2}} \left((N-2)\beta - \kappa \left(\log |x| \right)^{-\beta} \right)$$

(see [17, Lemma 3.4]), we obtain that $\Psi(x) \sim |x|^{N-2}$. We have thus proven that

$$\Psi(x) \leqslant \frac{C}{(1+|x|)^{N-2}}.$$

Now, note that since V is nonnegative,

$$-\Delta \Psi_{\varepsilon} - H\Psi_{\varepsilon} \leq 0.$$

In view of (33) and (31), Proposition 3.1 is applicable, and for every $x \in \mathbb{R}^N \setminus \Lambda$,

$$\Psi_{\varepsilon}(x) \leqslant \Psi(x) \leqslant \frac{C}{(1+|x|)^{N-2}}.$$

As explained in [17], the estimate (32) is the best one can hope for if V decays rapidly at infinity or is compactly supported. However, if V decays quadratically or subquadratically at infinity, we can improve (32).

Lemma 5.3. Let Ψ_{ε} be given by Lemma 5.2.

(1) If $\liminf_{|x|\to\infty} V(x)|x|^2 > 0$, then there exist $\lambda > 0$, R > 0 and C > 0 such that for every $\varepsilon > 0$ and $x \in \mathbb{R}^N \setminus B(0, R)$,

$$\Psi_{\varepsilon}(x) \leqslant C\left(\frac{R}{|x|}\right)^{\frac{N-2}{2}+\sqrt{\left(\frac{N-2}{2}\right)^2-\kappa+\frac{\lambda^2}{\varepsilon^2}}}$$

(2) If $\liminf_{|x|\to\infty} V(x)|x|^{\alpha} > 0$ with $\alpha < 2$, then there exist $\lambda > 0$, R > 0, C > 0 and $\varepsilon_0 > 0$ such that for every $\varepsilon \in (0, \varepsilon_0)$ and $x \in \mathbb{R}^N \setminus B(0, R)$,

$$\Psi_{\varepsilon}(\mathbf{x}) \leqslant C \exp\left(-\frac{\lambda}{\varepsilon}\left(|\mathbf{x}|^{\frac{2-\alpha}{2}} - R^{\frac{2-\alpha}{2}}\right)\right).$$

(3) If $\liminf_{|x|\to 0} V(x)|x|^2 > 0$, then there exist $\lambda > 0$, r > 0 and C > 0 such that for every $\varepsilon > 0$ and $x \in B(0, r)$,

$$\Psi_{\varepsilon}(x) \leqslant C\left(\frac{|x|}{r}\right)^{\sqrt{(\frac{N-2}{2})^2 - \kappa + \frac{\lambda^2}{\varepsilon^2} - \frac{N-2}{2}}}$$

(4) If $\liminf_{|x|\to 0} V(x)|x|^{\alpha} > 0$ with $\alpha < 2$, then there exist $\lambda > 0$, R > 0, C > 0 and $\varepsilon_0 > 0$ such that for every $\varepsilon \in (0, \varepsilon_0)$ and $x \in B(0, r)$,

$$\Psi_{\varepsilon}(\mathbf{x}) \leq C \exp\left(-\frac{\lambda}{\varepsilon}\left(|\mathbf{x}|^{-\frac{\alpha-2}{2}}-r^{-\frac{\alpha-2}{2}}\right)\right).$$

Proof. For (1), there exist R > 0 and $\lambda > 0$ such that for $x \in \mathbb{R}^N \setminus B(0, R)$

$$(1-\mu)V(x) \ge \frac{\lambda^2}{|x|^2}.$$

One then checks that

$$W(x) = \left(\frac{R}{|x|}\right)^{\frac{N-2}{2} + \sqrt{\left(\frac{N-2}{2}\right)^2 - \kappa + \frac{\lambda^2}{\varepsilon^2}}}$$

is a supersolution in $\mathbb{R}^N \setminus B(0, R)$.

For (2), there exist R > 0 and $\eta > 0$ such that for $x \in \mathbb{R}^N \setminus B(0, R)$

$$(1-\mu)V(x) \ge \frac{\eta}{|x|^{\alpha}}$$

One then checks that

$$W(x) = \exp\left(-\frac{\lambda}{\varepsilon}\left(|x|^{\frac{2-\alpha}{2}} - R^{\frac{2-\alpha}{2}}\right)\right)$$

is a supersolution in $\mathbb{R}^N \setminus B(0, R)$ with $\lambda^2 < (\frac{2}{2-\alpha})^2 \nu$ and ε small enough. The proofs of the other assertions are similar. \Box

Another important tool is a function that describes the exponential decay of u_{ε} inside Λ .

Lemma 5.4. Let $\bar{x} \in \Lambda$ and R > 0 be such that

$$B_{\mathcal{H}}(\bar{x},R) \subset \Lambda. \tag{34}$$

Define

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$$\Phi_{\varepsilon}^{\bar{x}}(x) := \cosh\left(\lambda \frac{R - d_{\mathcal{H}}(x, \bar{x})}{\varepsilon}\right).$$
(35)

There exist $\lambda > 0$ *and* $\varepsilon_0 > 0$ *such that for every* $\varepsilon \in (0, \varepsilon_0)$ *, one has*

$$-\varepsilon^2 \Delta \Phi_{\varepsilon}^{\bar{x}} + (1-\mu) V \Phi_{\varepsilon}^{\bar{x}} \ge 0 \quad in \ B_{\mathcal{H}}(\bar{x}, R).$$

Proof. First one computes

$$-\varepsilon^{2}\Delta\Phi_{\varepsilon}^{\bar{x}}(x) = -\lambda^{2}\cosh\left(\frac{\lambda}{\varepsilon}\left(R - d_{\mathcal{H}}(x,\bar{x})\right)\right) + \frac{\varepsilon\lambda}{d_{\mathcal{H}}(x,\bar{x})}\left(N - 1 - k\frac{|\bar{x}''|}{|x''|}\right)\sinh\left(\frac{\lambda}{\varepsilon}\left(R - d_{\mathcal{H}}(x,\bar{x})\right)\right).$$

Let us choose $\lambda > 0$ such that $\lambda^2 < (1 - \mu) \inf_{\Lambda} V$. In view of (34), one has for $x \in B_{\mathcal{H}}(\bar{x}, R)$,

$$-\varepsilon^{2} \Delta \Phi_{\varepsilon}^{x}(x) + (1-\mu) V \Phi_{\varepsilon}^{x}(x)$$

$$\geq \frac{\varepsilon \lambda}{d_{\mathcal{H}}(x,\bar{x})} \left(N - 1 - k \frac{|\bar{x}''|}{|x''|} \right) \sinh\left(\frac{\lambda}{\varepsilon} \left(R - d_{\mathcal{H}}(x,\bar{x}) \right) \right)$$

$$+ \left((1-\mu) \inf_{\Lambda} V - \lambda^{2} \right) \cosh\left(\frac{\lambda}{\varepsilon} \left(R - d_{\mathcal{H}}(x,\bar{x}) \right) \right).$$

This last expression is positive if ε is sufficiently small. \Box

Lemma 5.5. Let $(x_{\varepsilon})_{\varepsilon} \subset \Lambda$ be such that

$$\liminf_{\varepsilon\to 0} d_{\mathcal{H}}(x_{\varepsilon}, \partial \Lambda) > 0$$

and $\rho > 0$. Then, there exist $\varepsilon_0 > 0$ and a family of functions $(W_{\varepsilon})_{0 < \varepsilon < \varepsilon_0} \subset C^{1,1}_{loc}((\mathbb{R}^N \setminus \{0\}) \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho))$ such that for all $\varepsilon \in (0, \varepsilon_0)$, one has

(i) W_{ε} satisfies the inequation

$$-\varepsilon^{2}(\Delta+H)W_{\varepsilon}+(1-\mu)VW_{\varepsilon} \geq 0 \quad in \mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon},\varepsilon\rho),$$

(ii) $\nabla W_{\varepsilon} \in L^{2}(\mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho))$ and $\frac{W_{\varepsilon}}{|x|} \in L^{2}(\mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho))$,

- (iii) $W_{\varepsilon} \ge 1$ on $\partial B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho)$,
- (iv) for every $x \in B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon \rho)$,

$$W_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d_{\mathcal{H}}(x, x_{\varepsilon})}{1 + d_{\mathcal{H}}(x, x_{\varepsilon})}\right) (1 + |x|)^{-(N-2)}, \quad x \in \mathbb{R}^{N}.$$

Moreover,

(1) if $\liminf_{|x|\to\infty} V(x)|x|^2 > 0$, then there exist $\lambda > 0$, $\nu > 0$ and C > 0 such that for $\varepsilon > 0$ small enough,

$$W_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d_{\mathcal{H}}(x, x_{\varepsilon})}{1 + d_{\mathcal{H}}(x, x_{\varepsilon})}\right) \left(1 + |x|\right)^{-\frac{\nu}{\varepsilon}};$$

(2) if $\liminf_{|x|\to\infty} V(x)|x|^{\alpha} > 0$ with $\alpha > 2$, then there exist $\lambda > 0$ and C > 0 such that for $\varepsilon > 0$ small enough,

$$W_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d_{\mathcal{H}}(x, x_{\varepsilon})}{1 + d_{\mathcal{H}}(x, x_{\varepsilon})} \left(1 + |x|\right)^{\frac{2-\alpha}{2}}\right);$$

(3) if $\liminf_{|x|\to 0} V(x)|x|^2 > 0$, then there exist $\lambda > 0$, $\nu > 0$ and C > 0 such that for $\varepsilon > 0$ small enough,

$$W_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d_{\mathcal{H}}(x, x_{\varepsilon})}{1 + d_{\mathcal{H}}(x, x_{\varepsilon})}\right) \left(\frac{|x|}{1 + |x|}\right)^{\frac{\nu}{\varepsilon}};$$

(4) if $\liminf_{|x|\to 0} V(x)|x|^{\alpha} > 0$ with $\alpha > 2$, then there exist $\lambda > 0$ and C > 0 such that for $\varepsilon > 0$ small enough,

$$W_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d_{\mathcal{H}}(x, x_{\varepsilon})}{1 + d_{\mathcal{H}}(x, x_{\varepsilon})} \left(\frac{|x|}{1 + |x|}\right)^{\frac{\alpha-2}{2}}\right).$$

Proof. Let Ψ_{ε} be given by Lemma 5.2. Choose a set $U \subset \mathbb{R}^N$ such that $\overline{\Lambda} \subset U$, $0 \notin \overline{U}$ and \overline{U} is compact. Choose $\tilde{\Psi}_{\varepsilon} \in C^2(\mathbb{R}^N \setminus \{0\}) \cap H^1_{\text{loc}}(\mathbb{R}^N)$ such that $\tilde{\Psi}_{\varepsilon} = \Psi_{\varepsilon}$ in $\mathbb{R}^N \setminus U$ and $\tilde{\Psi}_{\varepsilon} = 1$ in Λ . In view of the estimate of Lemma 5.2, one can also ensure that $\sup_{\varepsilon > 0} \|\tilde{\Psi}_{\varepsilon}\|_{L^{\infty}(U)} < \infty$. Choose R > 0 such that

$$R < \liminf_{\varepsilon \to 0} \operatorname{dist}(x_{\varepsilon}, \partial \Lambda).$$
(36)

Let $\Phi_{\varepsilon}^{x_{\varepsilon}}$ be given by (35) and set

$$w_{\varepsilon}(x) := \begin{cases} \Phi_{\varepsilon}^{x_{\varepsilon}}(x) & \text{if } x \in B_{\mathcal{H}}(x_{\varepsilon}, R), \\ \tilde{\Psi}_{\varepsilon}(x) & \text{if } x \in \mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon}, R). \end{cases}$$

By (36), for ε small enough, $B_{\mathcal{H}}(x_{\varepsilon}, R) \subset \Lambda$ so that $w_{\varepsilon} \in C^{1,1}(\mathbb{R}^N)$. Moreover, if ε is small enough, Lemma 5.4 is applicable and in $B_{\mathcal{H}}(x_{\varepsilon}, R) \setminus B_{\mathcal{H}}(x_{\varepsilon}, \varepsilon\rho)$, we have

$$-\varepsilon^{2}(\Delta+H)w_{\varepsilon}+(1-\mu)Vw_{\varepsilon} \geq -\varepsilon^{2}\Delta\Phi_{\varepsilon}^{x_{\varepsilon}}+(1-\mu)V\Phi_{\varepsilon}^{x_{\varepsilon}} \geq 0.$$

In $\Lambda \setminus B_{\mathcal{H}}(x_{\varepsilon}, R)$, one has

$$-\varepsilon^{2}(\Delta+H)w_{\varepsilon}+(1-\mu)Vw_{\varepsilon}=-\varepsilon^{2}H+(1-\mu)\left(\inf_{\Lambda}V\right)\geq0,$$

for ε small enough. In $U \setminus \Lambda$, one has

$$-\varepsilon^{2}(\Delta+H)w_{\varepsilon}+(1-\mu)Vw_{\varepsilon}=-\varepsilon^{2}(\Delta+H)\tilde{\Psi}_{\varepsilon}+(1-\mu)V\tilde{\Psi}_{\varepsilon}\geq 0,$$

for ε small enough since $V \tilde{\Psi}_{\varepsilon}$ is positive on \overline{U} . Finally, in $\mathbb{R}^N \setminus U$, one has

$$-\varepsilon^{2}(\Delta + H)w_{\varepsilon} + (1-\mu)Vw_{\varepsilon} = -\varepsilon^{2}(\Delta + H)\Psi_{\varepsilon} + (1-\mu)V\Psi_{\varepsilon} = 0.$$

We set

$$W_{\varepsilon}(x) := \frac{W_{\varepsilon}(x)}{\cosh \lambda(\frac{R}{\varepsilon} - \rho)},$$

where λ is chosen as in the previous lemma. It is standard to see that W_{ε} satisfies properties (ii) and (iii). Statement (iv) follows from Lemma 5.2. The other conclusions follow from Lemma 5.3.

Thanks to the previous lemma, we obtain an upper bound on the solutions $(u_{\varepsilon})_{\varepsilon>0}$ of $(\mathcal{P}_{\varepsilon})$.

Proposition 5.6. Suppose that the assumptions of Proposition 4.7 are satisfied. Let $(u_{\varepsilon})_{\varepsilon>0} \subset H^1_{V,\mathcal{H}}(\mathbb{R}^N)$ be the positive solutions of $(\mathcal{P}_{\varepsilon})$ found in Theorem 4 and $(x_{\varepsilon})_{\varepsilon>0} \subset \Lambda$ be such that

$$\liminf_{\varepsilon\to 0} u_{\varepsilon}(x_{\varepsilon}) > 0.$$

Then there exist C > 0, $\lambda > 0$ and $\varepsilon_0 > 0$ such that for all $\varepsilon \in (0, \varepsilon_0)$,

$$u_{\varepsilon}(x) \leq C \exp\left(-\frac{\lambda}{\varepsilon} \frac{d(x, S_{\varepsilon}^{k})}{1 + d(x, S_{\varepsilon}^{k})}\right) \left(1 + |x|\right)^{-(N-2)}, \quad x \in \mathbb{R}^{N}.$$
(37)

...

Moreover, (1), (2), (3) and (4) in Lemma 5.5 hold with u_{ε} in place of W_{ε} .

Proof. By Lemma 5.1, there exist $\rho > 0$ and $\varepsilon_0 > 0$ such that for all $\varepsilon \in (0, \varepsilon_0)$, the solution u_{ε} satisfies inequation (30). Further, $\|u_{\varepsilon}\|_{L^{\infty}(B_{\mathcal{H}}(x_{\varepsilon},\varepsilon\rho))}$ is bounded as $\varepsilon \to 0$ in view of Lemma 4.3. Let $(W_{\varepsilon})_{\varepsilon}$ be the family of barrier functions given by Lemma 5.5. By Proposition 3.1, we have

$$u_{\varepsilon}(x) \leq ||u_{\varepsilon}||_{L^{\infty}(B_{\mathcal{H}}(x_{\varepsilon},\varepsilon\rho))}W_{\varepsilon}(x) \text{ in } \mathbb{R}^{N} \setminus B_{\mathcal{H}}(x_{\varepsilon},\varepsilon\rho),$$

and the conclusion comes from Lemma 5.5. \Box

We are now in a position to prove Theorem 3.

Proof of Theorem 3. We know from Theorem 4 that the modified equation $(\mathcal{P}_{\varepsilon})$ possesses a positive solution $u_{\varepsilon} \in H^1_{V,\mathcal{H}}(\mathbb{R}^N)$. In order to prove that for ε small enough, this solution actually solves (2), it suffices to show that, for every $x \in (\mathbb{R}^N \setminus \{0\}) \setminus \Lambda$, one has

$$K(x)\frac{f(u_{\varepsilon}(x))}{u_{\varepsilon}(x)} \leq \varepsilon^{2}H(x) + \mu V(x).$$

Assume that *V* and *K* satisfy (\mathcal{G}_{∞}^1) and (\mathcal{G}_0^1) , by Proposition 5.6 and assumptions (f_4) and (f_1) , if $\varepsilon > 0$ is small enough, we have for all $x \in \mathbb{R}^N \setminus \Lambda$,

$$K(x)\frac{f(u_{\varepsilon}(x))}{u_{\varepsilon}(x)} \leq K(x)\frac{f(Ce^{-\frac{\lambda}{\varepsilon}}(1+|x|)^{-(N-2)})}{Ce^{-\frac{\lambda}{\varepsilon}}(1+|x|)^{-(N-2)}}$$
$$\leq Ce^{-\frac{\lambda}{\varepsilon}(q-1)}(1+|x|)^{\sigma-(N-2)(q-1)}$$
$$\leq \frac{\varepsilon^{2}\kappa}{|x|^{2}((\log|x|)^{2}+1)^{\frac{1+\beta}{2}}} = \varepsilon^{2}H(x).$$

The other cases can be treated in a similar way. \Box

In some settings, it is interesting to determine whether the solutions are in L^2 . We obtain as a byproduct the following

Corollary 5.7. Let u_{ε} be the solution of (2) found in Theorem 3. If $N \ge 5$ or $\liminf_{|x|\to\infty} |x|^2 V(x) > 0$, then, for ε small enough, $u_{\varepsilon} \in L^2(\mathbb{R}^N)$.

Proof. This follows immediately from Proposition 5.6. □

6. The two-dimensional case

In dimension N = 2, the method has to be modified because the classical Hardy inequality fails on unbounded domains of \mathbb{R}^2 . Let us recall the Hardy-type inequality that was proved in [17, Lemma 6.1]:

Lemma 6.1. Let R > r. Then there exists C > 0 such that for every $u \in \mathcal{D}(\mathbb{R}^2)$,

$$\int_{\mathbb{R}^2} |\nabla u|^2 + C \int_{B(0,R)\setminus B(0,r)} u^2 \ge \frac{1}{4} \int_{\mathbb{R}^2\setminus B(0,R)} \frac{u^2(x)}{|x|^2(\log\frac{|x|}{r})^2} dx.$$

We deduce therefrom

Lemma 6.2. If $V \in C(\mathbb{R}^2 \setminus \{0\})$ is nonnegative and non-identically 0, then there exists $\kappa_0 > 0$ such that for $\varepsilon > 0$ sufficiently small, for every $u \in \mathcal{D}(\mathbb{R}^2)$,

$$\kappa_0 \int_{\mathbb{R}^2} \frac{u^2(x)}{|x|^2 (1 + (\log |x|)^2)} \, dx \leq \int_{\mathbb{R}^2} \varepsilon^2 |\nabla u|^2 + V u^2.$$

Proof. One sees that by the conformal transformation $x \mapsto \frac{x}{|x|^2}$, Lemma 6.1 becomes

$$\int_{\mathbb{R}^2} |\nabla u|^2 + C \int_{B(0,R)\setminus B(0,r)} u^2 \ge \frac{1}{4} \int_{B(0,r)} \frac{u^2(x)}{|x|^2 (\log |x|)^2} \, dx.$$

Therefore, one has

$$\int_{\mathbb{R}^2} \frac{u^2(x)}{|x|^2(1+(\log|x|)^2)} \, dx \leqslant C \bigg(\int_{\mathbb{R}^2} |\nabla u|^2 + \int_{B(0,2)\setminus B(0,1/2)} u^2 \bigg).$$

Since V is continuous and does not vanish identically, there exist $\bar{x} \in \mathbb{R}^2$ and $\bar{r} > 0$, such that $\inf_{B(\bar{x},\bar{r})} V > 0$. Hence, there exists C > 0 such that

$$\int_{B(0,2)\setminus B(0,1/2)} u^2 \leqslant C\bigg(\int_{\mathbb{R}^2} |\nabla u|^2 + V|u|^2\bigg).$$

Bringing the inequalities together, there exists C > 0 such that

$$\int_{\mathbb{R}^2} \frac{u^2(x)}{|x|^2(1+(\log|x|)^2)} \, dx \leq C \int_{\mathbb{R}^2} (|\nabla u|^2 + V u^2).$$

This brings the conclusion when $\varepsilon > 0$ is small enough. \Box

The space $H^1_V(\mathbb{R}^2)$ can thus be defined as in the case N > 2 as the closure of $\mathcal{D}(\mathbb{R}^2)$ with respect to the norm defined by (13).

The penalization potential $H : \mathbb{R}^2 \to \mathbb{R}$ is defined by

$$H(x) := \frac{\kappa}{|x|^2 (1 + (\log |x|)^2)^{\frac{2+\beta}{2}}},$$

where $\beta > 0$ and $\kappa \in (0, \kappa_0)$. We see that

$$H(x) \leqslant \frac{\kappa}{|x|^2(1+(\log|x|)^2)}$$

Together with Lemma 6.2, this ensures the positivity of the quadratic form associated to $-\varepsilon^2(\Delta + H) + V$.

As in the case N > 2, this inequality implies the following comparison principle.

Proposition 6.3. Let $\Omega \subset \mathbb{R}^2$ be a smooth domain. Let $v, w \in H^1_{loc}(\Omega) \cap C(\overline{\Omega})$ be such that $\nabla(w - v)_- \in L^2(\Omega)$, $(w - v)_-/(|x|(1 + |\log |x||)) \in L^2(\Omega)$ and

$$-\varepsilon^{2}(\Delta+H)w + Vw \ge -\varepsilon^{2}(\Delta+H)v + Vv, \quad in \ \Omega.$$

If $\partial \Omega \neq \emptyset$, assume also that $w \ge v$ on $\partial \Omega$. Then $w \ge v$ in Ω .

One continues the proof of Theorem 3 as in the case N > 2. In Proposition 3.2, one takes $A_{\lambda} := \overline{B(0, e^{e^{\lambda}})} \setminus B(0, e^{-e^{\lambda}})$ and

$$\eta_{\lambda}(x) := \zeta\left(\frac{\log|\log|x||}{\lambda}\right).$$

One then obtains estimate (19) by using Lemma 6.2 instead of Hardy's inequality. The only other notable difference lies in the choice of the function W in the proof of Lemma 5.2, where one follows the construction of [17, Lemma 6.3], i.e.

$$W(x) = \beta(\beta + 1) - \kappa \left| \log |x| \right|^{-\beta}.$$

References

- Antonio Ambrosetti, Andrea Malchiodi, Perturbation Methods and Semilinear Elliptic Problems on ℝⁿ, Progr. Math., vol. 240, Birkhäuser Verlag, 2006.
- [2] Antonio Ambrosetti, Andrea Malchiodi, Concentration phenomena for NLS: recent results and new perspectives, in: Perspectives in Nonlinear Partial Differential Equations, in: Contemp. Math., vol. 446, Amer. Math. Soc., Providence, RI, 2007, pp. 19–30.
- [3] Antonio Ambrosetti, Andrea Malchiodi, Wei-Ming Ni, Singularly perturbed elliptic equations with symmetry: Existence of solutions concentrating on spheres, part I, Comm. Math. Phys. 235 (2003) 427–466.
- [4] Antonio Ambrosetti, David Ruiz, Radial solutions concentrating on spheres of nonlinear Schrödinger equations with vanishing potentials, Proc. Roy. Soc. Edinburgh Sect. A 136 (5) (2006) 889–907.
- [5] Marino Badiale, Teresa D'Aprile, Concentration around a sphere for a singularly perturbed Schrödinger equations, Nonlinear Anal. 49 (2002) 947–985.
- [6] Thomas Bartsch, Shuangjie Peng, Semiclassical symmetric Schrödinger equations: existence of solutions concentrating simultaneously on several spheres, Z. Angew. Math. Phys. 58 (5) (2007) 778–804.
- [7] Vieri Benci, Teresa D'Aprile, The semiclassical limit of the nonlinear Schrödinger equation in a radial potential, J. Differential Equations 184 (2002) 109–138.
- [8] Denis Bonheure, Jean Van Schaftingen, Bound state solutions for a class of nonlinear Schrödinger equations, Rev. Mat. Iberoam. 24 (2008) 297-351.

- [9] Jaeyoung Byeon, Zhi-Qiang Wang, Spherical semiclassical states of a critical frequency for Schrödinger equations with decaying potentials, J. Eur. Math. Soc. (JEMS) 8 (2) (2006) 217–228.
- [10] Juan Dávila, Manuel del Pino, Monica Musso, Juncheng Wei, Standing waves for supercritical nonlinear Schrödinger equations, J. Differential Equations 236 (1) (2007) 164–198.
- [11] Manuel del Pino, Patricio Felmer, Local mountain passes for semilinear elliptic problems in unbounded domains, Calc. Var. Partial Differential Equations 4 (2) (1996) 121-137.
- [12] Manuel del Pino, Michal Kowalczyk, Juncheng Wei, Concentration on curves for nonlinear Schrödinger equations, Comm. Pure Appl. Math. 60 (1) (2007) 113–146.
- [13] Fethi Mahmoudi, Andrea Malchiodi, Marcelo Montenegro, Solutions to the nonlinear Schrödinger equation carrying momentum along a curve, Comm. Pure Appl. Math. 62 (9) (2009) 1155-1264.
- [14] Fethi Mahmoudi, Andrea Malchiodi, Concentration on minimal submanifolds for a singularly perturbed Neumann problem, Adv. Math. 209 (2) (2007) 460–525.
- [15] Andrea Malchiodi, Construction of multidimensional spike-layers, Discrete Contin. Dyn. Syst. 14 (1) (2006) 187-202.
- [16] Riccardo Molle, Donato Passaseo, Concentration phenomena for solutions of superlinear elliptic problems, Ann. Inst. H. Poincaré Anal. Non Linéaire 23 (2006) 63–84.
- [17] Vitaly Moroz, Jean Van Schaftingen, Semiclassical stationary states for nonlinear Schrödinger equations with fast decaying potentials, Calc. Var. Partial Differential Equations 37 (1–2) (2010) 1–27.
- [18] Richard S. Palais, The principle of symmetric criticality, Comm. Math. Phys. 69 (1) (1979) 19-30.