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On the BKS pairing for Kähler quantizations of the cotangent bundle of a Lie group

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Abstract

A natural one-parameter family of Kähler quantizations of the cotangent bundle T^*K of a compact Lie group K , taking into account the half-form correction, was studied in [C. Florentino, P. Matias, J. Mourão, J.P. Nunes, Geometric quantization, complex structures and the coherent state transform, *J. Funct. Anal.* 221 (2005) 303–322]. In the present paper, it is shown that the associated Blattner–Kostant–Sternberg (BKS) pairing map is unitary and coincides with the parallel transport of the quantum connection introduced in our previous work, from the point of view of [S. Axelrod, S. Della Pietra, E. Witten, Geometric quantization of Chern–Simons gauge theory, *J. Differential Geom.* 33 (1991) 787–902]. The BKS pairing map is a composition of (unitary) coherent state transforms of K , introduced in [B.C. Hall, The Segal–Bargmann coherent state transform for compact Lie groups, *J. Funct. Anal.* 122 (1994) 103–151]. Continuity of the Hermitian structure on the quantum bundle, in the limit when one of the Kähler polarizations degenerates to the vertical real polarization, leads to the unitarity of the corresponding BKS pairing map. This is in agreement with the unitarity up to scaling (with respect to a rescaled inner product) of this pairing map, established by Hall.

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Contents

1. Introduction 181

2. The prequantum BKS pairing 183

 2.1. The prequantum bundle \mathcal{H}^{prQ} 184

 2.2. The prequantum BKS pairing and its associated connection on \mathcal{H}^{prQ} 186

3. Unitarity of the quantum BKS pairing map and the CST 188

 3.1. The quantum BKS pairing between Kähler polarizations and the vertical polarization 189

 3.2. Unitarity of the quantum BKS pairing map 195

 3.3. Relation between the quantum BKS pairing map and the CST 196

1. Introduction

Let K be a compact, connected Lie group of dimension n and let T^*K be its cotangent bundle. We start by recalling some aspects of [2] where, in connection with work of Hall in [6], the geometric quantization of T^*K was studied using a natural one-parameter family of Kähler structures. These Kähler structures are induced on T^*K via the following natural identifications of T^*K with the complexified group $K_{\mathbb{C}}$. Consider, for any real parameter $s > 0$, the diffeomorphisms

$$\begin{aligned} \psi_s : T^*K &\rightarrow K_{\mathbb{C}} \\ (x, Y) &\mapsto \psi_s(x, Y) = x e^{isY}. \end{aligned} \tag{1.1}$$

Here, $x \in K$, $Y \in \mathfrak{K} \equiv \text{Lie}(K)$, and we identify T^*K with $K \times \mathfrak{K}^*$ using left invariant forms and then with $K \times \mathfrak{K}$ by means of a fixed Ad -invariant inner product on \mathfrak{K} . The diffeomorphisms ψ_s endow T^*K with a family of complex structures J_s and one can check that, together with the canonical symplectic structure ω on T^*K , the pair (ω, J_s) defines a Kähler structure on T^*K for every $s \in \mathbb{R}_+$ [2]. This family includes the Kähler structure on T^*K considered by Hall in [6].

In this paper, we consider the Blattner–Kostant–Sternberg (BKS) pairing between two different Kähler quantizations of T^*K . To describe the results, let us consider the framework used in [2]. Let L denote the trivial complex line bundle on T^*K , with trivial Hermitian structure (its sections are therefore identified with C^∞ functions on T^*K). Following the geometric quantization program with half-form correction, let us introduce the half-form bundle δ_s , which is a square root of the (trivial) J_s -canonical bundle κ_s over T^*K . Choosing canonical trivializing J_s -holomorphic sections Ω_s of $\kappa_s = \delta_s^2$ and $\sqrt{\Omega_s}$ of δ_s (we refer to the next section for precise formulas) one introduces a natural Hermitian structure on $L \otimes \delta_s$ so that, for a smooth section σ_s of the form

$$\sigma_s = f \sqrt{\Omega_s}, \tag{1.2}$$

with $f \in C^\infty(T^*K)$, one has

$$|\sigma_s|^2 := |f|^2 |\Omega_s|, \tag{1.3}$$

where $|\Omega_s|$ is defined by $\overline{\Omega_s} \wedge \Omega_s = |\Omega_s|^2 b \epsilon$, $b = (2i)^n (-1)^{n(n-1)/2}$ and $\epsilon = \frac{1}{n!} \omega^n$ is the Liouville measure on T^*K .

The prequantum Hilbert space $\mathcal{H}_s^{\text{prQ}}$, depending on s , is then the (norm completion of the) space of C^∞ sections σ_s of $L \otimes \delta_s$ which are square-integrable with respect to the Hermitian structure (1.3), that is

$$\langle \sigma_s, \sigma_s \rangle^{\text{prQ}} := \int_{T^*K} |\sigma_s|^2 \epsilon < +\infty. \tag{1.4}$$

Proceeding with the quantization program, the quantum Hilbert space \mathcal{H}_s^{Q} is defined to be the subspace of $\mathcal{H}_s^{\text{prQ}}$ consisting of polarized (J_s -holomorphic) sections of $L \otimes \delta_s$. This is naturally a sub-Hilbert space of $\mathcal{H}_s^{\text{prQ}}$.

Both families of Hilbert spaces can be collected to form the Hilbert prequantum bundle $\mathcal{H}^{\text{prQ}} \rightarrow \mathbb{R}_+$ and the quantum bundle $\mathcal{H}^{\text{Q}} \rightarrow \mathbb{R}_+$, which is naturally a sub-bundle of \mathcal{H}^{prQ} .

In the spirit of [1], we consider a natural connection δ^{prQ} on \mathcal{H}^{prQ} and use it to induce, by orthogonal projection, an Hermitian connection δ^{Q} on \mathcal{H}^{Q} . Note that, from (1.2)–(1.4), \mathcal{H}^{prQ} has a natural global (inner product preserving) trivializing morphism defined by

$$\begin{aligned} L^2(T^*K, \epsilon) \times \mathbb{R}_+ &\rightarrow \mathcal{H}^{\text{prQ}} \\ (f, s) &\mapsto \frac{f}{\sqrt{|\Omega_s|}} \sqrt{\Omega_s}. \end{aligned}$$

The prequantum connection δ^{prQ} on \mathcal{H}^{prQ} is defined to be the connection induced via this map from the trivial bundle $L^2(T^*K, \epsilon) \times \mathbb{R}_+ \rightarrow \mathbb{R}_+$ equipped with the trivial connection. Then, the quantum connection δ^{Q} is defined to be simply the orthogonal projection of δ^{prQ} to the quantum bundle. This construction is a natural generalization, to the case with the half-form correction and with a natural trivializing section of the square root of the canonical bundle, of the framework considered in [1]. One of the advantages of this approach consists in the fact that it ensures automatically that the quantum connection is Hermitian.

In the present paper, the framework of [1] is extended in another natural way. Namely, we consider the Blattner–Kostant–Sternberg (BKS) pairing between two different fibers of \mathcal{H}^{Q} , corresponding to two different quantizations of T^*K (see, for instance, [12]). This is the restriction of a Hermitian pairing between the corresponding prequantum Hilbert spaces, which we call the prequantum BKS pairing, and which is defined by

$$\langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}} = \int_{T^*K} \bar{f} f' \sqrt{\frac{\Omega_s \wedge \Omega_{s'}}{b\epsilon}} \epsilon,$$

for $\sigma_s = f \sqrt{\Omega_s} \in \mathcal{H}_s^{\text{prQ}}$ and $\sigma_{s'} = f' \sqrt{\Omega_{s'}} \in \mathcal{H}_{s'}^{\text{prQ}}$. This prequantum BKS pairing defines a map

$$B_{s s'}^{\text{prQ}} : \mathcal{H}_{s'}^{\text{prQ}} \rightarrow \mathcal{H}_s^{\text{prQ}}, \tag{1.5}$$

whose infinitesimal version induces a (also natural) connection δ on \mathcal{H}^{prQ} . We prove in Theorem 1 that δ coincides with the connection δ^{prQ} . Therefore, the quantum connection δ^{Q} will also be induced from the BKS pairing on \mathcal{H}^{Q} . This result shows that the approach of the infinitesimal

BKS pairing coincides with that of [1] considered in [2]. We note, however, that the prequantum BKS pairing map does not coincide with the parallel transport associated to δ^{PrQ} (see Theorem 2.)

In [6], Hall considered two polarizations for the quantization of T^*K . Fixing the same inner product as Hall on \mathfrak{K} , one of these polarizations, the Kähler one, corresponds to the case of $s = 1$, and the other, the vertical real polarization, corresponds, in our framework, to letting s go to 0. He proved that the BKS pairing map between the Hilbert spaces associated to these two quantizations coincides (up to scale) with the Segal–Bargmann coherent state transform (CST) for K introduced in [4], provided that one takes into account the half-form correction. In Theorem 5, see Section 3, we elaborate on this relation, by showing that the (quantum) BKS pairing map

$$B_{ss'}^{\text{Q}} : \mathcal{H}_{s'}^{\text{Q}} \rightarrow \mathcal{H}_s^{\text{Q}}$$

obtained by restriction of (1.5) to the quantum bundle is, in fact, unitary. Moreover, and unlike the prequantum case, we show in Theorem 6 that $B_{ss'}^{\text{Q}}$ coincides with the parallel transport of δ^{Q} and corresponds also to the CST. In particular this implies that the BKS pairing maps satisfy

$$B_{ss'}^{\text{Q}} \circ B_{s's''}^{\text{Q}} = B_{ss''}^{\text{Q}}.$$

The quantum bundle \mathcal{H}^{Q} can be naturally extended to a bundle over $[0, \infty)$ by including a fiber over $s = 0$ corresponding to the space of vertically polarized sections. Continuity at $s = 0$ of the Hermitian structure on the extended quantum bundle motivates a particular choice of scale of the inner product on the space of vertically polarized sections. We show that with this choice the BKS pairing map B_{s0}^{Q} becomes unitary. We note that this is in agreement with [6], where this map was shown to be unitary up to scale but for a different choice of scale in the inner product. We study this fact in Section 3.1.

Therefore, we have obtained a class of examples of quantization where the BKS pairing map, between half-form corrected Hilbert spaces corresponding to two different Kähler quantizations, is an isometric isomorphism. We also confirm in a different way the result of [2] that the parallel transport of the quantum connection coincides with the one given by the CST. We stress that we prove unitarity of the quantum BKS pairing map and not just unitarity up to scale. Therefore, for the family of polarizations that we consider, cotangent bundles of compact Lie groups provide new examples in which the half-form correction leads to the unitarity of the BKS pairing map (which is the case for Fock spaces, see [12, Chapter 10]).

For other related works on the dependence of the quantization on the choice of polarization, the unitarity of the BKS pairing map and the relation between vertical and Kähler polarizations, see also [3, 8–11].

2. The prequantum BKS pairing

Let K be a compact, connected Lie group of real dimension n , and let us fix once and for all the Ad -invariant inner product (\cdot, \cdot) on its Lie algebra \mathfrak{K} , for which the corresponding Riemannian measure is the normalized Haar measure on K . Let $\{X_i\}_{i=1}^n$ be left-invariant vector fields on

K forming an orthonormal basis for \mathfrak{K} . The cotangent bundle $T^*K \cong K \times \mathfrak{K}$ is naturally a symplectic manifold with a canonical symplectic 2-form defined by $\omega = -d\theta$, where

$$\theta = \sum_{i=1}^n y^i w^i.$$

Here, (y^1, \dots, y^n) are the global coordinates on \mathfrak{K} corresponding to the basis $\{X_i\}_{i=1}^n$, and $\{w^i\}_{i=1}^n$ is the basis of left-invariant 1-forms on K dual to $\{X_i\}_{i=1}^n$, pulled-back to T^*K by the canonical projection. We let ϵ denote the Liouville volume form on T^*K , given by

$$\epsilon = \frac{1}{n!} \omega^n.$$

For each $s > 0$, let $\psi_s : T^*K \rightarrow K_{\mathbb{C}}$ denote the diffeomorphism defined in (1.1) by $\psi_s(x, Y) = x e^{i s Y}$, which induces on T^*K a Kähler structure (ω, J_s) .

2.1. The prequantum bundle \mathcal{H}^{prQ}

We start by recalling from [2], some of the formulas that will be needed later on. Let \tilde{X}_j , $j = 1, \dots, n$, be the vector fields on T^*K generating the right action of K lifted to T^*K and given by

$$\psi_{s*} \tilde{X}_j = X_{j, \mathbb{C}},$$

where $X_{j, \mathbb{C}}$ denotes the extension of X_j from a left-invariant vector field on $K \subset K_{\mathbb{C}}$ to the corresponding left-invariant vector field on $K_{\mathbb{C}}$. Let $\{\tilde{w}^j\}$ be the basis of left invariant 1-forms defined by $\tilde{w}^j(\tilde{X}_k) = \delta_k^j$ and $\tilde{w}^j(J_s \tilde{X}_k) = 0$, for $j, k = 1, \dots, n$. For every $s \in \mathbb{R}_+$, a frame of left invariant J_s -holomorphic 1-forms is then

$$\{\tilde{\eta}_s^j = \tilde{w}^j - i J_s \tilde{w}^j\}_{j=1}^n,$$

where $(J_s w)(X) = w(J_s X)$, for a vector field X and a 1-form w on T^*K .

Consider now the J_s -canonical bundle $\kappa_s \rightarrow T^*K$ whose sections are J_s -holomorphic n -forms with natural Hermitian structure defined as follows. For a J_s -holomorphic n -form α_s , let $|\alpha_s|$ be the unique non-negative C^∞ function on T^*K such that $\bar{\alpha}_s \wedge \alpha_s = |\alpha_s|^2 b \epsilon$, where $b = (2i)^n (-1)^{n(n-1)/2}$. Following [6] we write

$$|\alpha_s|^2 = \frac{\bar{\alpha}_s \wedge \alpha_s}{b \epsilon}.$$

Given the inner product on \mathfrak{K} , a canonical trivializing J_s -holomorphic section of κ_s is given by

$$\Omega_s := \tilde{\eta}_s^1 \wedge \dots \wedge \tilde{\eta}_s^n$$

and has norm [2]

$$|\Omega_s|^2 = s^n \eta^2(sY), \tag{2.1}$$

where $\eta(Y)$ is the Ad_K -invariant function, defined for Y in a Cartan subalgebra by the following product over a set R^+ of positive roots of \mathfrak{K} ,

$$\eta(Y) = \prod_{\alpha \in R^+} \frac{\sinh \alpha(Y)}{\alpha(Y)}. \tag{2.2}$$

The following proposition generalizes (2.1).

Proposition 1. *Let $s, s' > 0$. We have*

$$\bar{\Omega}_s \wedge \Omega_{s'} = \left(\frac{s+s'}{2}\right)^n \eta^2\left(\frac{s+s'}{2}Y\right) b\epsilon = |\Omega_{(s+s')/2}|^2 b\epsilon.$$

Proof. The result follows by direct computation. From [5] and the definition of ψ_s , we can write $D\psi_s : TT^*K \rightarrow TK_{\mathbb{C}}$ as

$$D\psi_s(x, Y) = \begin{bmatrix} \cos \operatorname{ad} s Y & \frac{1 - \cos \operatorname{ad} s Y}{\operatorname{ad} Y} \\ -\sin \operatorname{ad} s Y & \frac{\sin \operatorname{ad} s Y}{\operatorname{ad} Y} \end{bmatrix},$$

using the (x, Y) coordinate basis on T^*K and the basis $\{X_{j,\mathbb{C}}, JX_{j,\mathbb{C}}\}_{j=1,\dots,n}$ on $K_{\mathbb{C}}$, where J is the complex structure on $K_{\mathbb{C}}$. From the explicit expressions for the forms $\tilde{\eta}_s^j$ we then find, for $s, s' > 0$,

$$\bar{\Omega}_s \wedge \Omega_{s'} = \tilde{\eta}_s^1 \wedge \dots \wedge \tilde{\eta}_s^n \wedge \tilde{\eta}_{s'}^1 \wedge \dots \wedge \tilde{\eta}_{s'}^n = \det \begin{bmatrix} \bar{M}_s & \bar{N}_s \\ M_{s'} & N_{s'} \end{bmatrix} (-1)^{n(n-1)/2} \epsilon,$$

where the endomorphisms M_s and N_s are defined by

$$M_s = e^{-i \operatorname{ad} s Y}, \quad N_s = \frac{1 - e^{-i \operatorname{ad} s Y}}{\operatorname{ad} Y}.$$

From the left invariance of the forms, this determinant can be evaluated for Y in the Cartan subalgebra which, after taking care of the contribution from the null space of $\operatorname{ad} Y$, yields

$$\bar{\Omega}_s \wedge \Omega_{s'} = \left(\frac{s+s'}{2}\right)^n \prod_{\alpha \in R} \frac{(e^{s'\langle \alpha, Y \rangle} - e^{-s\langle \alpha, Y \rangle})}{(s+s')\langle \alpha, Y \rangle} b\epsilon,$$

where the product runs over the set R of all roots of \mathfrak{K} . The result then follows from definition (2.2). \square

As in the introduction, let δ_s be the J_s -holomorphic bundle of half-forms on T^*K , with trivializing section whose square is Ω_s . Following [6], we will denote this section by $\sqrt{\Omega_s}$. Recall that the prequantum Hilbert space \mathcal{H}^{prQ} is the completion of the space of C^∞ -sections of $L \otimes \delta_s$ of finite norm (1.4), where L denotes the trivial complex line bundle on T^*K , and its Hermitian structure $\langle \cdot, \cdot \rangle^{\text{prQ}}$ can be written as

$$\langle \sigma_s, \sigma'_s \rangle^{\text{prQ}} = \int_{T^*K} \bar{f} f' |\Omega_s| \epsilon = \int_{T^*K} \bar{f} f' \sqrt{\frac{\overline{\Omega_s} \wedge \Omega_s}{b\epsilon}} \epsilon,$$

for two smooth sections of $L \otimes \delta_s$ written as $\sigma_s = f \sqrt{\Omega_s}$, $\sigma'_s = f' \sqrt{\Omega_s}$, with $f, f' \in C^\infty(T^*K)$.

The smooth Hilbert bundle structure on \mathcal{H}^{prQ} and the prequantum connection δ^{prQ} are chosen to be the ones compatible with the global trivializing map

$$\begin{aligned} L^2(T^*K, \epsilon) \times \mathbb{R}_+ &\rightarrow \mathcal{H}^{\text{prQ}} \\ (f, s) &\mapsto \frac{f}{\sqrt{|\Omega_s|}} \sqrt{\Omega_s}, \end{aligned} \tag{2.3}$$

and the trivial connection on $L^2(T^*K, \epsilon) \times \mathbb{R}_+$. Note that the section of δ_s given by $\sqrt{\Omega_s}/\sqrt{|\Omega_s|}$ has unit norm, and its use would simplify some of the formulas. However, since this section is not holomorphic, we use $\sqrt{\Omega_s}$ as trivializing section of δ_s , which will be more suited for describing the polarized sections in Section 3.

2.2. The prequantum BKS pairing and its associated connection on \mathcal{H}^{prQ}

The general procedure, as described for instance in [12], for defining a BKS pairing, suggests, in our setting, the definition of a Hermitian pairing which we call the prequantum BKS pairing, as follows.

Definition 1. Let $\sigma_s = f \sqrt{\Omega_s} \in \mathcal{H}_s^{\text{prQ}}$ and $\sigma'_{s'} = f' \sqrt{\Omega_{s'}} \in \mathcal{H}_{s'}^{\text{prQ}}$. Their BKS pairing is defined by

$$\langle \sigma_s, \sigma'_{s'} \rangle^{\text{BKS}} = \int_{T^*K} \bar{f} f' \sqrt{\frac{\overline{\Omega_s} \wedge \Omega_{s'}}{b\epsilon}} \epsilon. \tag{2.4}$$

Note that the integral above exists for $\sigma_s, \sigma'_{s'}$ satisfying the conditions for sections of the prequantum bundle in (1.4), which in this case read $f \sqrt{|\Omega_s|}, f' \sqrt{|\Omega_{s'}|} \in L^2(T^*K, \epsilon)$. This can be readily checked by using Proposition 1 to write (2.4) as

$$\langle \sigma_s, \sigma'_{s'} \rangle^{\text{BKS}} = \int_{T^*K} \bar{f} f' |\Omega_{(s+s')/2}| \epsilon, \tag{2.5}$$

and using the definition of $\eta(Y)$ in (2.2) to prove that the smooth function

$$\phi(s, s', Y) := \frac{|\Omega_{(s+s')/2}|^2}{|\Omega_s| |\Omega_{s'}|} = \left(\frac{s + s'}{2\sqrt{ss'}} \right)^n \frac{\eta^2\left(\frac{s+s'}{2} Y\right)}{\eta(sY)\eta(s'Y)} \tag{2.6}$$

is real, positive and bounded, for fixed s, s' .

We remark that, in the case when s' and s coincide, the prequantum BKS pairing is equal to the Hermitian structure $\langle \cdot, \cdot \rangle^{\text{prQ}}$ on $\mathcal{H}_s^{\text{prQ}}$. Also, we note that the pairing (2.4) is nondegenerate.

As mentioned above, δ^{prQ} was defined as the Hermitian connection for which sections \mathcal{H}^{prQ} of the form $\sigma_s = \frac{f}{\sqrt{|\Omega_s|}} \sqrt{\Omega_s}$, with $f \in L^2(T^*K, \epsilon)$, are parallel. It is also natural to consider a connection, δ , on \mathcal{H}^{prQ} , induced from the infinitesimal prequantum BKS pairing by the formula

$$\langle \sigma_s, \delta_{\partial/\partial s} \sigma'_s \rangle^{\text{prQ}} = \left. \frac{\partial}{\partial s'} \right|_{s'=s} \langle \sigma_s, \sigma'_s \rangle^{\text{BKS}}. \tag{2.7}$$

Let us prove that these two connections are the same.

Theorem 1. *The connections δ^{prQ} and δ on \mathcal{H}^{prQ} are equal.*

Proof. Let σ be a smooth section of \mathcal{H}^{prQ} , such that $\sigma_s = \frac{f}{\sqrt{|\Omega_s|}} \sqrt{\Omega_s} \in \mathcal{H}_s^{\text{prQ}}$, for $f \in L^2(T^*K, \epsilon)$, $s > 0$. Since sections of this type can be used to form a global moving frame for \mathcal{H}^{prQ} , to prove the theorem, it suffices to show that σ is parallel with respect to the connection δ . Let $\tau_s = g \sqrt{\Omega_s} \in \mathcal{H}_s^{\text{prQ}}$. We have from (2.5) and (2.6),

$$\langle \tau_s, \delta_{\partial/\partial s} \sigma_s \rangle^{\text{prQ}} = \left. \frac{\partial}{\partial s'} \right|_{s'=s} \langle \tau_s, \sigma'_s \rangle^{\text{BKS}} = \int_{T^*K} \bar{g} f |\Omega_s|^{1/2} \left(\left. \frac{\partial \sqrt{\phi}}{\partial s'} \right|_{s'=s} \right) \epsilon.$$

From a straightforward computation using (2.6) we obtain

$$\left. \frac{\partial \phi(s, s', Y)}{\partial s'} \right|_{s'=s} = 0,$$

which implies that $\delta_{\partial/\partial s} \sigma_s = 0$, as required. \square

One immediate consequence of this theorem is that δ^{prQ} induced from the infinitesimal prequantum BKS pairing is unitary. We note that the calculation in the proof of Theorem 1 relies just on the expression for $\Omega_s \wedge \Omega_{s'}$ in Proposition 1. While in this case this is related to the fact that the continuous family of complex polarizations is generated by the flux of a particular vector field [2], it is not clear to us whether this holds more generally.

In the next section, we will show that the restriction of the prequantum BKS pairing to \mathcal{H}^{Q} yields the (unitary) parallel transport associated to δ^{Q} . This, however, does not hold for the prequantum bundle. In fact, let us define the prequantum BKS pairing map $B_{ss'}^{\text{prQ}} : \mathcal{H}_{s'}^{\text{prQ}} \rightarrow \mathcal{H}_s^{\text{prQ}}$ by

$$\langle \sigma_s, \sigma'_s \rangle^{\text{BKS}} = \langle \sigma_s, B_{ss'}^{\text{prQ}} \sigma'_s \rangle^{\text{prQ}}. \tag{2.8}$$

Theorem 2. *The prequantum BKS pairing map (2.8), $B_{ss'}^{\text{prQ}} : \mathcal{H}_{s'}^{\text{prQ}} \rightarrow \mathcal{H}_s^{\text{prQ}}$ is given by*

$$B_{ss'}^{\text{prQ}} \left(f' \frac{\sqrt{\Omega_{s'}}}{\sqrt{|\Omega_{s'}|}} \right) = \sqrt{\phi} f' \frac{\sqrt{\Omega_s}}{\sqrt{|\Omega_s|}},$$

for $f' \sqrt{\Omega_{s'}} / \sqrt{|\Omega_{s'}|} \in \mathcal{H}_{s'}^{\text{prQ}}$, where ϕ is given by (2.6). Consequently, $B_{ss'}^{\text{prQ}}$ does not coincide with the parallel transport of the connection δ^{prQ} .

Proof. Let us write

$$B_{ss'}^{\text{prQ}} \left(f' \frac{\sqrt{\Omega_{s'}}}{\sqrt{|\Omega_{s'}|}} \right) = g \frac{\sqrt{\Omega_s}}{\sqrt{|\Omega_s|}},$$

for some vector $g \frac{\sqrt{\Omega_s}}{\sqrt{|\Omega_s|}} \in \mathcal{H}_s^{\text{prQ}}$. Then, using (2.5) we have for any $f \sqrt{\Omega_s} \in \mathcal{H}_s^{\text{prQ}}$, that

$$\int_{T^*K} \bar{f} f' \frac{|\Omega_{(s+s')/2}|}{\sqrt{|\Omega_{s'}|}} \epsilon = \left\langle f \sqrt{\Omega_s}, B_{ss'}^{\text{prQ}} \left(f' \frac{\sqrt{\Omega_{s'}}}{\sqrt{|\Omega_{s'}|}} \right) \right\rangle^{\text{prQ}} = \int_{T^*K} \bar{f} g \sqrt{|\Omega_s|} \epsilon.$$

Since $f \in L^2(T^*K, |\Omega_s|^{1/2} \epsilon)$ is arbitrary, we obtain $g = \sqrt{\phi} f'$, as wanted.

On the other hand, the parallel transport $P_{ss'} : \mathcal{H}_{s'}^{\text{prQ}} \rightarrow \mathcal{H}_s^{\text{prQ}}$ of δ^{prQ} is determined by (2.3) giving

$$P_{ss'} \left(f' \frac{\sqrt{\Omega_{s'}}}{\sqrt{|\Omega_{s'}|}} \right) = f' \frac{\sqrt{\Omega_s}}{\sqrt{|\Omega_s|}}.$$

Clearly, this is not equal to $B_{ss'}^{\text{prQ}}$ since ϕ is not identically equal to 1. \square

It follows from Theorem 2 that $B_{ss'}^{\text{prQ}}$ is not unitary. We will see, however, that its restriction to the quantum Hilbert bundle is unitary.

3. Unitarity of the quantum BKS pairing map and the CST

Consider the quantum Hilbert bundle $\mathcal{H}^Q \rightarrow \mathbb{R}_+$, whose fiber over each $s > 0$ consists of J_s -polarized sections of $L \otimes \delta_s$. Recall from [2] that \mathcal{H}_s^Q is given by

$$\mathcal{H}_s^Q := \left\{ \sigma_s = F e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s}, F \text{ is } J_s\text{-holomorphic and } \|\sigma_s\|_s^Q < \infty \right\},$$

where the norm refers to the Hermitian structure on \mathcal{H}^Q as a sub-bundle of \mathcal{H}^{prQ} . The pairing (2.4) between fibers of \mathcal{H}^{prQ} restricted to polarized sections defines the BKS pairing between fibers of \mathcal{H}^Q , for which we use the same notation.

Definition 2. For $s, s' > 0$, write

$$\sigma_s = F e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s} \in \mathcal{H}_s^Q \quad \text{and} \quad \sigma'_{s'} = F' e^{-\frac{s'|Y|^2}{2\hbar_0}} \sqrt{\Omega_{s'}} \in \mathcal{H}_{s'}^Q.$$

Their (quantum) BKS pairing is given by

$$\langle \sigma_s, \sigma'_{s'} \rangle^{\text{BKS}} = \int_{T^*K} \bar{F} F' e^{-\frac{(s+s')|Y|^2}{2\hbar_0}} \sqrt{\frac{\bar{\Omega}_s \wedge \Omega_{s'}}{b\epsilon}} \epsilon. \tag{3.1}$$

As in (2.7), the BKS pairing induces a connection δ^{BKS} on $\mathcal{H}^{\mathbb{Q}}$ defined by the same formula

$$\langle \sigma_s, \delta_{\partial/\partial s}^{\text{BKS}} \sigma_{s'} \rangle^{\mathbb{Q}} = \frac{\partial}{\partial s'} \Big|_{s'=s} \langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}}$$

for $\sigma_s \in \mathcal{H}_s^{\mathbb{Q}}$ and $\sigma_{s'} \in \mathcal{H}_{s'}^{\mathbb{Q}}$. Recall that, along the lines of [1], the quantum bundle is equipped with a quantum connection $\delta^{\mathbb{Q}}$, obtained from the orthogonal projection $P : \mathcal{H}^{\text{pr}\mathbb{Q}} \rightarrow \mathcal{H}^{\mathbb{Q}}$. More precisely, we have $\delta^{\mathbb{Q}} = P \circ \delta^{\text{pr}\mathbb{Q}}$, which implies immediately

Theorem 3. *The connections δ^{BKS} and $\delta^{\mathbb{Q}}$ on $\mathcal{H}^{\mathbb{Q}}$ coincide.*

Proof. This is an obvious corollary of Theorem 1, since the pairing on $\mathcal{H}^{\mathbb{Q}}$ is obtained by restriction from $\mathcal{H}^{\text{pr}\mathbb{Q}}$. \square

The BKS pairing formalism in this case provides results consistent with the approach of [1]. We will establish below the fact that the BKS pairing map for $\mathcal{H}^{\mathbb{Q}}$ is unitary and coincides with the parallel transport of $\delta^{\mathbb{Q}}$.

3.1. The quantum BKS pairing between Kähler polarizations and the vertical polarization

In [2], we showed that the parallel transport associated to $\delta^{\mathbb{Q}}$ corresponds to Hall’s CST, so that, in fact, $\delta^{\mathbb{Q}}$ -parallel sections of $\mathcal{H}^{\mathbb{Q}}$ satisfy a heat equation on $K_{\mathbb{C}}$. This provides a better understanding of some of the results of [6].

In [6], Hall has also shown that, up to a constant, the CST corresponds to the quantum BKS pairing map between the vertical real polarized and the Kähler polarized quantum Hilbert spaces. In this section, we establish in Theorem 4 an important identity for the quantum BKS pairing between Kähler polarizations. This identity will help us understand better the $s \rightarrow 0$ limit and, in Section 3.3, the relation between the quantum BKS pairing and the CST.

Let now Δ be the invariant Laplacian on K and \mathcal{C} denote the analytic continuation from K to $K_{\mathbb{C}}$. We recall from [4], that the CST is a unitary isomorphism of Hilbert spaces defined by

$$\begin{aligned} C_{\hbar} : L^2(K, dx) &\rightarrow \mathcal{H}L^2(K_{\mathbb{C}}, dv_{\hbar}) \\ f &\mapsto C_{\hbar}(f) = \mathcal{C} \circ e^{\frac{\hbar}{2}\Delta} f, \end{aligned}$$

where dx is the normalized Haar measure on K , $dv_{\hbar}(g) = v_{\hbar}(g) dg$ is the K -averaged heat kernel measure of [4], and dg is the Haar measure on $K_{\mathbb{C}}$. Recall from [5] the explicit form of the K -averaged heat kernel measure for $\hbar = s\hbar_0$

$$v_{\hbar}(g) = (a_s s^{\frac{n}{2}} \eta(Y))^{-1} e^{-\frac{|Y|^2}{\hbar}}, \tag{3.2}$$

where

$$a_s = (\pi \hbar_0)^{\frac{n}{2}} e^{|\rho|^2 \hbar_0 s}, \tag{3.3}$$

and ρ is the Weyl vector given by half the sum of the positive roots of \mathfrak{K} . Using the relation of parallel sections of $\mathcal{H}^{\mathbb{Q}}$ with holomorphic functions on $K_{\mathbb{C}}$ we will now obtain an explicit

formula for the BKS pairing as an integral over $K_{\mathbb{C}}$. As in Definition 2 of the quantum BKS pairing, let

$$\sigma_s = F e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s} \in \mathcal{H}_s^Q \quad \text{and} \quad \sigma_{s'} = F' e^{-\frac{s'|Y|^2}{2\hbar_0}} \sqrt{\Omega_{s'}} \in \mathcal{H}_{s'}^Q, \quad s, s' > 0, \quad (3.4)$$

where $F = \hat{F} \circ \psi_s$ and $F' = \hat{F}' \circ \psi_{s'}$ are, respectively, J_s -holomorphic and $J_{s'}$ -holomorphic functions on T^*K , with \hat{F}, \hat{F}' holomorphic functions on $K_{\mathbb{C}}$.

Proposition 2. *The BKS pairing on \mathcal{H}^Q is given by*

$$\langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}} = a_{(s+s')/2} \int_{K_{\mathbb{C}}} \overline{\hat{F}(gZ)} \hat{F}'(gZ^{-1}) dv_{\hbar''}(g),$$

where $g = xe^{iY}$, $Z = e^{i\frac{s-s'}{s+s'}Y}$ and $\hbar'' = \frac{s+s'}{2}\hbar_0$.

Proof. From Definition 1 and Proposition 1 we have

$$\langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}} = \int_{T^*K} \overline{\hat{F}(xe^{isY})} \hat{F}'(xe^{is'Y}) e^{-\frac{(s+s')|Y|^2}{2\hbar_0}} \left(\frac{s+s'}{2}\right)^{\frac{n}{2}} \eta\left(\frac{s+s'}{2}Y\right) \epsilon. \quad (3.5)$$

We recall from [2,5] that $\psi_s^* dg = |\Omega_s|^2 \epsilon = s^n \eta^2(sY) \epsilon$. Using (3.2) and (3.3), the formula then follows from a change of variables of integration, from (x, Y) to $g = xe^{iY'}$ where $Y' = \frac{s+s'}{2}Y$, and where at the end we renamed Y' as Y again. \square

The computation of the BKS pairing, and also its relation to the CST, is made very explicit with the following result.

Theorem 4. *Let $\sigma_s \in \mathcal{H}_s^Q$ and $\sigma_{s'} \in \mathcal{H}_{s'}^Q$ be given by*

$$\begin{aligned} \sigma_s &= (C_{s\hbar_0} f) \circ \psi_s e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s}, \\ \sigma_{s'} &= (C_{s'\hbar_0} f') \circ \psi_{s'} e^{-\frac{s'|Y|^2}{2\hbar_0}} \sqrt{\Omega_{s'}}, \end{aligned}$$

where $s, s' > 0$ and $f, f' \in L^2(K, dx)$. We have

$$\langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}} = a_{(s+s')/2} \langle f, f' \rangle_{L^2(K, dx)}.$$

We remark that, from Theorem 3, one knows that the connection δ^Q is also the connection induced from the BKS pairing between infinitesimally close fibers of \mathcal{H}^Q . Note, however, that of course one can give many different explicit Hermitian pairings between fibers of \mathcal{H}^Q , such that they all induce the same connection δ^Q on \mathcal{H}^Q . Theorem 4 implies that, among those, the pairing map that corresponds to the CST is the BKS pairing map naturally defined from geometric

quantization. This is the map $B_{s s'}^Q : \mathcal{H}_{s'}^Q \rightarrow \mathcal{H}_s^Q$ obtained by restriction of the prequantum BKS map (2.8),

$$\langle \sigma_s, \sigma'_{s'} \rangle^{\text{BKS}} = \langle \sigma_s, B_{s s'}^Q \sigma'_{s'} \rangle^Q, \tag{3.6}$$

for $\sigma_s \in \mathcal{H}_s^Q$ and $\sigma'_{s'} \in \mathcal{H}_{s'}^Q$.

Postponing the proof of Theorem 4 to the end of this section, let us now address the problem of taking the $s' \rightarrow 0$ limit in the BKS pairing, which corresponds to degenerating one of the Kähler polarizations into the vertical polarization.

From [2] we obtain, for σ_s as in (3.4),

$$\sigma_s \in \mathcal{H}_s^Q \iff \hat{F} \in \mathcal{H}L^2(K_C, d\nu_{\hbar})$$

with $\hbar = s\hbar_0$. Let us use the CST to define a family of continuous sections of \mathcal{H}^Q labelled by vectors $f \in L^2(K, dx)$,

$$\sigma_s = (C_{s\hbar_0} f) \circ \psi_s e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s}. \tag{3.7}$$

Due to the fact that the coherent state transform is an isomorphism, the values of the sections σ in (3.7) span, for every $s > 0$, the whole fiber \mathcal{H}_s^Q . This defines a global trivialization of the vector bundle \mathcal{H}^Q , which is particularly suitable for studying the $s \rightarrow 0$ limit. Let $\Omega_0 = \lim_{s \rightarrow 0} \Omega_s = dx$, so that $\sqrt{\Omega_0} = \sqrt{dx}$, and let \mathcal{H}_0^Q be the space

$$\mathcal{H}_0^Q = \{ f \sqrt{\Omega_0}, f \in L^2(K, dx) \}.$$

We consider the extension of the bundle $\mathcal{H}^Q \rightarrow \mathbb{R}_+$ to a bundle

$$\widehat{\mathcal{H}^Q} \rightarrow [0, \infty),$$

with a given trivialization

$$\begin{aligned} L^2(K, dx) \times [0, \infty) &\rightarrow \widehat{\mathcal{H}^Q} \\ (f, s) &\mapsto (C_{s\hbar_0} f) \circ \psi_s e^{-\frac{s|Y|^2}{2\hbar_0}} \sqrt{\Omega_s}, \quad s > 0, \\ (f, 0) &\mapsto f \sqrt{\Omega_0}. \end{aligned} \tag{3.8}$$

Sections (3.7) are then extended by continuity to $s = 0$ by setting

$$\sigma_0 = f \sqrt{\Omega_0}. \tag{3.9}$$

The BKS pairing between a Kähler polarized section $\sigma_s \in \mathcal{H}_s^Q$ as in (3.7) and a vertically polarized section σ'_0 of the form (3.9), with fixed $f' \in L^2(K, dx)$, is then defined by

$$\langle \sigma_s, \sigma'_0 \rangle^{\text{BKS}} = \int_{T^*K} \overline{(C_{s\hbar_0} f)(x e^{isY})} f'(x) e^{-\frac{s}{2\hbar_0} |Y|^2} \left(\frac{s}{2}\right)^{\frac{n}{2}} \eta\left(\frac{s}{2} Y\right) \epsilon, \tag{3.10}$$

if the integral is absolutely convergent, which corresponds formally to taking the $s' \rightarrow 0$ limit inside the integral in (3.5).

Remark 1. In [6], it is shown that (after a trivial rescaling of the variable Y to sY) this integral defines on dense subspaces a Hermitian pairing between the Hilbert space of Kähler polarized sections and the Hilbert space of vertically polarized sections. Setting $s = 1$ in (3.10) and taking care of the different normalizations, we obtain the same expression as in [6, Theorem 2.6] for this pairing. (In our definition of the pairing there is an extra factor of $2^{-n/2}$ as compared to [6].)

Using our notation and normalization conventions, Hall’s result implies that for any $s > 0$,

$$\langle \sigma_s, \sigma'_0 \rangle^{\text{BKS}} = a_{s/2} \langle f, f' \rangle_{L^2(K, dx)}. \tag{3.11}$$

Remark 2. Notice that, even though the integral expression (3.10) can only be defined on dense subspaces, from (3.11) it follows that the pairing can be extended by continuity to the respective completions [6]. This is in contrast to the pairing between two Kähler polarized spaces, where the integral defining the BKS pairing (2.4) is absolutely convergent.

We recall from [6] that the pairing (3.10) does coincide with the general prescription for the geometric quantization BKS pairing between a Kähler polarized Hilbert space and a vertically polarized one. (See [12, Section 10.4].) This can also be easily deduced from (3.1) and Proposition 1.

Corollary 1. Upon identifying $\mathcal{H}_s^{\mathbb{Q}}$ with $L^2(K, dx)$ as in (3.8), the BKS pairing (3.1) between polarized sections has a well-defined (weak) $s' \rightarrow 0$ limit. This limit coincides with (3.10).

Proof. This follows immediately from Theorem 4 by taking the limit $s' \rightarrow 0$, with fixed $f, f' \in L^2(K, dx)$, which gives $\lim_{s' \rightarrow 0} \langle \sigma_s, \sigma'_{s'} \rangle^{\text{BKS}} = a_{s/2} \langle f, f' \rangle_{L^2(K, dx)}$ as in (3.11). \square

The associated pairing map between the Hilbert space of vertically polarized sections and the Hilbert space of Kähler polarized sections (say for $s = 1$) is unitary only up to scale as shown by Hall.

To further clarify this issue, let us address the choice of inner product for vertically polarized sections which will be important for the unitarity of the BKS pairing maps $B_{s,0}^{\mathbb{Q}}$. The inner product, $\langle \sigma_0, \sigma'_0 \rangle^{\mathbb{Q}}$, that is usually considered for vertically polarized sections of the form (3.9) reads $\int_K \bar{f} f' dx$ (see, e.g., [12]). However, in order to have continuity of the Hermitian structure on $\widehat{\mathcal{H}}^{\mathbb{Q}}$ at $s = 0$, we will define instead the inner product as

$$\langle \sigma_0, \sigma'_0 \rangle^{\mathbb{Q}} := \int_K \bar{f} f' (\pi \hbar_0)^{n/2} dx. \tag{3.12}$$

We then have,

Corollary 2. With the definition of the inner product (3.12), the Hermitian structure on the bundle $\widehat{\mathcal{H}}^{\mathbb{Q}}$ is continuous at $s = 0$.

Proof. The result follows from Theorem 4 with $s = s'$ and from (3.3). \square

The reason for the appearance of the coefficient $(\pi \hbar_0)^{n/2}$ in (3.12) is related to the fact that in the inner product of the holomorphic sections there is a factor associated to the contribution of the integration along \mathfrak{R} , which does not tend to 1 when $s \rightarrow 0$. Corollary 2 is an indication that, for real polarizations which correspond to boundary points of spaces of holomorphic polarizations, and perhaps more generally, the inner product should be rescaled as in (3.12). As we will see, this choice, motivated by the continuity of the Hermitian structure, will also lead to the unitarity of the BKS pairing map, $B_{s_0}^Q$ (rather than just unitarity up to scale).

For the proof of Theorem 4, we will need two auxiliary lemmas which prove analogues of [5, Eq. (4)]. Let ρ_\hbar denote the analytic continuation of the heat kernel on K to $K_{\mathbb{C}}$, as in [4]. Moreover, let $*$: $K_{\mathbb{C}} \rightarrow K_{\mathbb{C}}$ denote the unique anti-holomorphic anti-automorphism of $K_{\mathbb{C}}$ which extends the map $x \rightarrow x^{-1}$ on K (see also [7]). We now show that the Dirac delta distribution on K , with respect to the Haar measure dx , can be informally written as

$$\delta(x) = \int_{K_{\mathbb{C}}} \rho_{2\hbar}(x^{-1}g^*g) dv_{\hbar}(g).$$

More precisely, let $C(K)$ be the space of continuous functions on K with the supremum norm and let \mathcal{F} denote the dense subspace of finite linear combinations of matrix elements of irreducible representations of K . Then, we have

Lemma 1. *Let $\{f_n\}_{n \in \mathbb{N}}$ be a sequence in \mathcal{F} converging to $f \in C(K)$. Then*

$$\lim_{n \rightarrow \infty} \int_{K_{\mathbb{C}}} \left(\int_K \rho_{2\hbar}(x^{-1}g^*g) f_n(x) dx \right) dv_{\hbar}(g) = f(e), \tag{3.13}$$

where $e \in K$ is the identity.

Proof. Since evaluation at the identity is a continuous linear functional on $C(K)$, it suffices to show that for any matrix element R_{ij} of any irreducible representation R of K , we have

$$\int_{K_{\mathbb{C}}} \left(\int_K \rho_{2\hbar}(x^{-1}g^*g) R_{ij}(x) dx \right) dv_{\hbar}(g) = R_{ij}(e) = \delta_{ij}. \tag{3.14}$$

To show this, recall that $\rho_\hbar = \sum_R d_R e^{-\frac{\hbar}{2}c_R} \chi_R$, where d_R is the dimension of R , χ_R denotes its character and c_R is the negative of the eigenvalue of Δ corresponding to the eigenvector χ_R . Now, recall Weyl’s classical orthogonality relations

$$\int_K \overline{R_{ij}(x)} R'_{lk}(x) dx = \frac{\delta_{RR'}}{d_R} \delta_{jk} \delta_{il}. \tag{3.15}$$

From the unitarity of the CST of [4], we then obtain the following identities,

$$e^{-\frac{\hbar}{2}(c_R+c_{R'})} \int_{K_{\mathbb{C}}} \overline{R_{ij}(g)} R'_{lk}(g) dv_{\hbar}(g) = \frac{\delta_{RR'}}{d_R} \delta_{jk} \delta_{il}, \tag{3.16}$$

where we have used the fact that $C_{\hbar} R_{ij} = e^{-\frac{\hbar}{2}c_R} R_{ij}$. (We are denoting by the same symbol a function on K and its analytic continuation to $K_{\mathbb{C}}$.)

Upon substituting the expression for $\rho_{2\hbar}$ in (3.14), using $\chi_R(g) = \sum_{i=1}^{d_R} R_{ii}(g)$ and the orthogonality relations (3.15), (3.14) becomes

$$e^{-\hbar c_R} \sum_{k=1}^{d_R} \int_{K_{\mathbb{C}}} \overline{R_{ki}(g)} R_{kj}(g) dv_{\hbar}(g) = R_{ij}(e).$$

Using (3.16), we obtain the result. \square

The Dirac delta distribution can also be written informally in a different way as

$$\delta(x_1^{-1}x_2) = \int_{K_{\mathbb{C}}} \overline{\rho_{\hbar}(ge^{itY}x_1^{-1})} \rho_{\hbar'}(ge^{-itY}x_2^{-1}) dv_{\hbar''}(g),$$

which should be interpreted as above. More precisely,

Lemma 2. *Let $\{f_n\}_{n \in \mathbb{N}}$ be a sequence in \mathcal{F} converging to $f \in C(K)$. Then,*

$$\lim_{n \rightarrow \infty} \int_{K_{\mathbb{C}}} \left(\int_K \overline{\rho_{\hbar}(ge^{itY}x_1^{-1})} \rho_{\hbar'}(ge^{-itY}x_2^{-1}) f_n(x_1) dx_1 \right) dv_{\hbar''}(g) = f(x_2), \tag{3.17}$$

for any $x_2 \in K$, $\hbar + \hbar' = 2\hbar''$, and any real number t .

Proof. To prove the lemma, we substitute the explicit expressions for the heat kernels in (3.17) and use the fact that, with $g = xe^{iY} \in K_{\mathbb{C}}$, we have

$$\overline{\chi_R(xe^{i(1+t)Y}x_1^{-1})} = \chi_R(x_1e^{i(1+t)Y}x^{-1}).$$

Using again $\rho_{\hbar} = \sum_R d_R e^{-\frac{\hbar}{2}c_R} \chi_R$ and rewriting the characters as a sum of products of matrix elements, we express the integral as an integral on T^*K . One then finds, upon integration along K , and use of the orthogonality relations (3.15), that the t dependence in (3.17) is apparent and cancels out. Finally, from the fact that $\hbar + \hbar' = 2\hbar''$ and using $\int_K dx = 1$, one rewrites the expression as an integral on $K_{\mathbb{C}}$ so that Eq. (3.17) becomes

$$\lim_{n \rightarrow \infty} \int_{K_{\mathbb{C}}} \left(\int_K \rho_{2\hbar}(x_2^{-1}x_1g^*g) f_n(x_1) dx_1 \right) dv_{\hbar}(g) = f(x_2),$$

which, by (3.13), proves the lemma. \square

We now prove Theorem 4.

Proof of Theorem 4. Let $\sigma_s, \sigma_{s'}$ be as described in the theorem. From Proposition 2, we have

$$\langle \sigma_s, \sigma_{s'} \rangle^{\text{BKS}} = a_{(s+s')/2} \int_{K_{\mathbb{C}}} \overline{\hat{F}(gZ)} \hat{F}'(gZ^{-1}) d\nu_{\hbar''}(g), \tag{3.18}$$

where $Z = e^{i\frac{s-s'}{s+s'}Y}$, $\hbar'' = \frac{s+s'}{2}\hbar_0$ and where $\hat{F} = C_{\hbar}f$ and $\hat{F}' = C_{\hbar'}f'$, with $\hbar = s\hbar_0$ and $\hbar' = s'\hbar_0$. By definition,

$$\hat{F}(gZ) = \int_K \rho_{\hbar}(gZx_1^{-1})f(x_1) dx_1 \tag{3.19}$$

and

$$\hat{F}'(gZ^{-1}) = \int_K \rho_{\hbar'}(gZ^{-1}x_2^{-1})f'(x_2) dx_2. \tag{3.20}$$

Substituting (3.19) and (3.20) in (3.18), we see that to prove the result for $f, f' \in \mathcal{F}$, it is enough to show that

$$\int_{K_{\mathbb{C}}} \left(\int_K \overline{\rho_{\hbar}(gZx_1^{-1})} \rho_{\hbar'}(gZ^{-1}x_2^{-1})f(x_1) dx_1 \right) d\nu_{\hbar}(g) = f(x_2),$$

which follows easily from Lemma 2. To prove the result for all $f, f' \in L^2(K, dx)$, we note that from (2.5) it follows that the prequantum BKS pairing is continuous on $\mathcal{H}_s^{\text{prQ}} \times \mathcal{H}_{s'}^{\text{prQ}}$ and, therefore, the quantum BKS pairing on $\mathcal{H}_s^{\text{Q}} \times \mathcal{H}_{s'}^{\text{Q}}$ is also continuous. The theorem then follows from the fact that \mathcal{F} is dense in $L^2(K, dx)$, which implies, from the isomorphisms given by (3.8), that the space of sections σ_s of the form above with $f \in \mathcal{F}$, is also dense in \mathcal{H}_s^{Q} for all $s > 0$. \square

3.2. Unitarity of the quantum BKS pairing map

In this section, we prove the unitarity of the quantum BKS pairing map (3.6) between quantum Hilbert spaces \mathcal{H}_s^{Q} and $\mathcal{H}_{s'}^{\text{Q}}$ using a direct calculation (the unitarity of these maps also follows from Theorem 4).

Theorem 5. *The map $B_{ss'}^{\text{Q}} : \mathcal{H}_{s'}^{\text{Q}} \rightarrow \mathcal{H}_s^{\text{Q}}$ is unitary for $s, s' \geq 0$.*

In order to prove Theorem 5 we establish an auxiliary lemma. Let

$$\sigma_{sij}^R = R_{ij}(xe^{isY})e^{-\frac{s|Y|^2}{2\hbar_0}}\sqrt{\Omega_s},$$

where R_{ij} denotes a matrix element of an irreducible representation R of K , of dimension d_R , and $i, j = 1, \dots, d_R$.

Lemma 3. For any $s, s' \geq 0$, we have

$$B_{ss'}^Q(\sigma_{s'ij}^R) = e^{-\frac{s-s'}{2}\hbar_0(c_R+|\rho|^2)}\sigma_{sij}^R.$$

Proof. The lemma is an immediate consequence of the equality

$$\begin{aligned} \langle \sigma_{sij}^R, B_{ss'}^Q\sigma_{s'ij'}^{R'} \rangle^Q &= \langle \sigma_{sij}^R, \sigma_{s'ij'}^{R'} \rangle^{\text{BKS}} \\ &= \frac{\delta_{RR'}\delta_{ii'}\delta_{jj'}}{d_R}(\pi\hbar_0)^{\frac{n}{2}}e^{\frac{s+s'}{2}\hbar_0(c_R+|\rho|^2)}. \end{aligned} \tag{3.21}$$

To establish it, recall that $\epsilon = dx dY$. For $s = s' = 0$ the equality is trivial. Let either s or s' be positive. We then obtain

$$\begin{aligned} \langle \sigma_{sij}^R, \sigma_{s'ij'}^{R'} \rangle^{\text{BKS}} &= \int_{K \times \mathfrak{R}} \sum_{k=1}^{d_R} \overline{R_{ik}(x)R_{kj}(e^{isY})} \sum_{k'=1}^{d_{R'}} R'_{i'k'}(x)R'_{k'j'}(e^{is'Y}) \\ &\quad \times e^{-\frac{(s+s')}{2\hbar_0}|Y|^2} \left(\frac{s+s'}{2}\right)^{\frac{n}{2}} \eta\left(\frac{s+s'}{2}Y\right) dx dY \\ &= \delta_{RR'}\delta_{ii'}\frac{1}{d_R} \int_{\mathfrak{R}} R_{jj'}(e^{i(s+s')Y})e^{-\frac{s+s'}{2\hbar_0}|Y|^2} \left(\frac{s+s'}{2}\right)^{\frac{n}{2}} \eta\left(\frac{s+s'}{2}Y\right) dY. \end{aligned}$$

From (3.16) we obtain (3.21) and the lemma. \square

Proof of Theorem 5. The sections $\{\sigma_{sij}^R\}$, with R running over all irreducible representations of K and $i, j = 1, \dots, d_R$, form an orthogonal basis of $\mathcal{H}_{s'}^Q$. Therefore, using Lemma 3, the unitarity of the map $B_{ss'}^Q$ is equivalent to the condition

$$\frac{\langle B_{ss'}^Q\sigma_{s'ij}^R, B_{ss'}^Q\sigma_{s'ij}^R \rangle^Q}{\langle \sigma_{s'ij}^R, \sigma_{s'ij}^R \rangle^Q} = \frac{(\langle \sigma_{sij}^R, B_{ss'}^Q\sigma_{s'ij}^R \rangle^Q)^2}{\langle \sigma_{sij}^R, \sigma_{sij}^R \rangle^Q \langle \sigma_{s'ij}^R, \sigma_{s'ij}^R \rangle^Q} = 1.$$

This follows directly from (3.21). \square

Note that the unitarity of B_{s0}^Q follows immediately from [6] if one chooses the inner product for the real polarization as in (3.12).

3.3. Relation between the quantum BKS pairing map and the CST

In [2], it was shown that the parallel transport between two fibers of \mathcal{H}^Q , \mathcal{H}_s^Q and $\mathcal{H}_{s'}^Q$, for the quantum connection δ^Q corresponds to $C_{\hbar} \circ C_{\hbar'}^{-1}$, where $\hbar = s\hbar_0$ and $\hbar' = s'\hbar_0$. More explicitly, we defined the CST bundle $\mathcal{H}^H \rightarrow \mathbb{R}_+$ with $\mathcal{H}_s^H = \mathcal{H}L^2(K_{\mathbb{C}}, d\nu_{\hbar})$. The CST bundle comes equipped with a connection δ^H for which the parallel transport $U_{\hbar\hbar'}: \mathcal{H}_{s'}^H \rightarrow \mathcal{H}_s^H$ is given

by the CST as $U_{hh'} = C_h \circ C_{h'}^{-1}$. Moreover, there exists a natural unitary bundle isomorphism $S: \mathcal{H}^H \rightarrow \mathcal{H}^Q$ such that $\delta^Q = S \circ \delta^H \circ S^{-1}$. At the fiber over $s > 0$ this is given by

$$S_s(\hat{F}) = \hat{F} \circ \psi_s \frac{e^{-\frac{s|Y|^2}{2h_0}}}{\sqrt{a_s}} \sqrt{\Omega_s}, \quad (3.22)$$

for $\hat{F} \in \mathcal{HL}^2(K_{\mathbb{C}}, dv_h)$.

As a corollary of the theorems above, we obtain the following result, which adds to [2, Theorem 4] the explanation of the role of the BKS pairing, from the point of view of the CST.

Theorem 6. For $s, s' > 0$, the BKS pairing map $B_{s s'}^Q: \mathcal{H}_{s'}^Q \rightarrow \mathcal{H}_s^Q$ is given by

$$B_{s s'}^Q = S_s \circ C_h \circ C_{h'}^{-1} \circ S_{s'}^{-1} = S_s \circ U_{hh'} \circ S_{s'}^{-1},$$

where $h = sh_0$ and $h' = s'h_0$. This pairing coincides with the parallel transport of δ^Q .

Proof. The result follows from direct computation and from Theorem 4 and formula (3.22), where we use the unitarity of C_h and also $a_{(s+s')/2} = \sqrt{a_s a_{s'}}$. Since $\delta^Q \circ S = S \circ \delta^H$, the theorem states that the parallel transport of the quantum connection is given by the BKS pairing map which is clearly unitary. \square

This is in agreement with the results of [2] where, however, the quantum connection was defined in a different way by orthogonal projection. Recall also that sections of \mathcal{H}^Q which are horizontal with respect to δ^Q satisfy a heat equation on $K_{\mathbb{C}}$. The corresponding explicit form for the connection δ^Q can also be found directly from Theorem 6, using [2, Lemma 1].

As a final remark, note that in the case when the group K is abelian, all the formulas in the paper remain valid upon setting $\rho = 0$ and $\eta \equiv 1$.

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