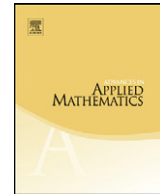




ELSEVIER

Contents lists available at SciVerse ScienceDirect

Advances in Applied Mathematics

www.elsevier.com/locate/yaama

A discrete Fourier transform associated with the affine Hecke algebra [☆]

J.F. van Diejen ^{*}, E. Emsiz

Facultad de Matemáticas, Pontificia Universidad Católica de Chile, Casilla 306, Correo 22, Santiago, Chile

ARTICLE INFO

Article history:

Received 10 October 2011

Accepted 22 December 2011

Available online 6 March 2012

MSC:

primary 42A38

secondary 20C08, 33D80

Keywords:

Discrete Fourier transform

Affine Hecke algebra

Spherical function

ABSTRACT

We introduce an explicit representation of the double affine Hecke algebra (of type A_1) at $q = 1$ that gives rise to a periodic counterpart of a well-known Fourier transform associated with the affine Hecke algebra.

© 2012 Elsevier Inc. All rights reserved.

1. Introduction

It is a classroom observation that any orthogonal basis of a separable Hilbert space gives rise to a corresponding (generalized) Fourier transform (establishing an isomorphism between the Hilbert space and an l^2 space over the basis indices). Of particular interest are in this context the Fourier series corresponding to bases of classical (basic) hypergeometric orthogonal polynomials and their generalizations given by Fourier-type transforms with kernels governed by non-polynomial (basic) hypergeometric functions [10,12]. Remarkably, such (basic) hypergeometric Fourier transforms can be fruitfully understood as generalized spherical transforms associated with double affine Hecke algebras and their degenerations [4,9,11,15,17].

[☆] Work was supported in part by the *Fondo Nacional de Desarrollo Científico y Tecnológico (FONDECYT)* Grants # 1090118 and # 11100315, and by the *Anillo ACT56 'Reticulados y Simetrías'* financed by the *Comisión Nacional de Investigación Científica y Tecnológica (CONICYT)*.

^{*} Corresponding author.

E-mail addresses: diejen@mat.puc.cl (J.F. van Diejen), eemsiz@mat.puc.cl (E. Emsiz).

In this paper we focus on a limiting situation with the Fourier basis given by one-dimensional Hall–Littlewood polynomials. The corresponding Fourier transform is a well-known parameter-deformation associated with the affine Hecke algebra (of type A_1) interpolating between spherical transforms for rank-one p -adic symmetric spaces [2,16,21]. By extending the affine Hecke algebra to a double affine Hecke algebra (at $q = 1$), we arrive at a finite-dimensional discrete Fourier transform that may be seen as a natural periodic analog of this deformed p -adic spherical transform.

The material is organized as follows. Section 2 recalls the definition of the A_1 -type double affine Hecke algebra (at $q = 1$) and outlines some of its key properties. Section 3 introduces two representations of this double affine Hecke algebra on the space of complex functions over \mathbb{Z} in terms of difference-reflection operators and integral-reflection operators, respectively. In Section 4 an explicit operator intertwining the difference-reflection representation and the integral-reflection representation is presented. The unilarity of the difference-reflection representation with respect to an appropriate Hilbert space structure is studied in Section 5 and a discrete Laplacian corresponding to a central element in the double affine Hecke algebra is introduced. The construction of a (generalized) spherical function associated with the double affine Hecke algebra in Section 6 gives rise to a basis of eigenfunctions diagonalizing the discrete Laplacian. The eigenfunctions in question constitute the kernel of a discrete Fourier transform associated with the affine Hecke algebra whose Plancherel formula is determined in Section 7.

Our discrete Laplacian and its eigenfunctions turn out to be special instances of a quantum Hamiltonian and its Bethe–Ansatz eigenfunctions modeling an integrable n -particle system discretizing the delta Bose gas on the circle [5] (cf. Remark 7.5 below for further details). From this perspective, our results provide the first step towards the construction of an appropriate Hecke-algebraic framework in the spirit of [7] for the discrete quantum model introduced in Ref. [5].

2. Double affine Hecke algebra

The *double affine Hecke algebra* \mathbb{H} with parameter $\tau \in (0, 1)$ (and $q = 1$) of type A_1 is the complex unital associative algebra generated by invertible elements T, U and X subject to the relations [15, Ch. 6.1]

$$(T - \tau)(T + \tau^{-1}) = 0, \quad U^2 = 1, \quad \text{and} \quad UXU = X^{-1} = T^{-1}XT^{-1}. \tag{2.1}$$

The subalgebra $H := \langle T, U \rangle \subset \mathbb{H}$ is referred to as the (extended) *affine Hecke algebra*.

A convenient linear basis for \mathbb{H} can be constructed by means of the underlying (extended) *affine Weyl group* $W := \langle s, u \mid s^2 = 1, u^2 = 1 \rangle$. Indeed, any group element $w \in W$ can be written uniquely in the form $w = u^{i_0} s_{i_1} \cdots s_{i_r}$, where $s_1 := s, s_0 := usu, i_0, \dots, i_r \in \{0, 1\}$, and $r \in \{0, 1, 2, \dots\}$ is minimal. Such a decomposition of w is called *reduced expression* and its length r defines the *length* $\ell(w)$ of the group element. Upon setting $T_w := U^{i_0} T_{i_1} \cdots T_{i_r}$, with $T_1 := T$ and $T_0 := UTU$, one has that the algebra elements $T_w X^k, w \in W, k \in \mathbb{Z}$ form a linear basis of \mathbb{H} and that the algebra elements $T_w, w \in W$ form a linear basis of the subalgebra H . The quadratic relation for T in Eq. (2.1) implies the following elementary multiplication rules in H

$$T_w T = T_{ws} + \frac{1}{2}(1 - \eta(w))(\tau - \tau^{-1})T_w, \tag{2.2a}$$

$$T T_w = T_{sw} + \frac{1}{2}(1 - \eta(w^{-1}))(\tau - \tau^{-1})T_w, \tag{2.2b}$$

where $\eta : W \rightarrow \{1, -1\}$ is defined by

$$\eta(w) := \ell(ws) - \ell(w). \tag{2.2c}$$

For future reference we also exhibit the multiplicative action of the generators of H on the abelian subalgebra (X) and vice versa. In the former situation only the action of T is nontrivial and given by the following classical formula

$$\begin{aligned}
 TX^k &= X^{-k}T + (\tau - \tau^{-1})\frac{X^k - X^{-k}}{1 - X^{-2}} \quad (k \in \mathbb{Z}) \\
 &= X^{-k}T + \text{sign}(k)(\tau - \tau^{-1}) \sum_{j=0}^{|k|-1} X^{|k|-2j}.
 \end{aligned}
 \tag{2.3}$$

For $k = 1$ and $k = -1$ this formula becomes $TX = X^{-1}T + (\tau - \tau^{-1})X$ and $TX^{-1} = XT - (\tau - \tau^{-1})X$, respectively, which are both manifest upon multiplying the relation $TX^{-1}T = X$ by T (from the left or from the right, respectively) and successive elimination of the T^2 factor with the aid of the quadratic relation for T . The general case then readily follows by induction in k .

The multiplicative action of X on the basis of H is more intricate. To give a precise description, let us note that the affine Weyl group decomposes as $W = \Omega \rtimes W_S$, where $\Omega := \langle u \rangle = \{1, u\}$ and W_S denotes the normal subgroup $\langle s_0, s_1 \rangle \subset W$. This decomposition gives rise to the following partial Bruhat order on W :

$$w < w' \iff w^{-1}w' \in W_S \quad \text{and} \quad \ell(w) < \ell(w').
 \tag{2.4}$$

Proposition 2.1. *Let $w = u^r w'$ with $r \in \{0, 1\}$ and $w' \in W_S$, and let $\varepsilon \in \{-1, 1\}$. Then one has that*

$$\begin{aligned}
 T_w X^\varepsilon &= X^{\varepsilon(-1)^{\ell(w)+r}} T_w \\
 &\quad + \varepsilon \eta(w) X^{\eta(w)(-1)^{\ell(w)+r}} \sum_{\substack{v \in W \\ v < w}} a(\ell(w) - \ell(v)) T_v,
 \end{aligned}
 \tag{2.5a}$$

with $\eta(w)$ given by Eq. (2.2c) and

$$a(k) := \left(\frac{1 - \tau^2}{1 + \tau^2} \right) (\tau^{-k} + (-1)^{k+1} \tau^k).
 \tag{2.5b}$$

Proof. It is sufficient to consider the case that $r = 0$, $\varepsilon = 1$ and $\eta(w) = 1$:

$$T_w X = X^{(-1)^{\ell(w)}} \left(T_w + \sum_{v \in W, v < w} a(\ell(w) - \ell(v)) T_v \right)
 \tag{2.6}$$

(for $w \in W_S$ with $ws > w$). Indeed, the cases $r = 0$ and $r = 1$ are related via the multiplication of both sides from the left by U , and the formulas corresponding to the other three combinations of signs $(\varepsilon, \eta(w)) = (-1, 1), (1, -1), (-1, -1)$ follow from the case $(\varepsilon, \eta(w)) = (1, 1)$ upon invoking the relations $T_w X^{-1} = (X + X^{-1})T_w - T_w X$ (cf. Remark 2.2 below) and $T_{uwu}X = UT_w X^{-1}U$. (Notice in this respect that $\eta(uwu) = -\eta(w)$ for $w \in W \setminus \Omega$.) We will now proceed to prove Eq. (2.6) by induction on $\ell(w)$ starting from the trivial case $\ell(w) = 0$. Let $w \in W_S$ with $ws > w$ and let $i = \ell(w) \bmod 2$ (i.e. $i \in \{0, 1\}$) is chosen such that $s_i w > w$ (and $\eta(s_i w) = 1$). Assuming that the relation in Eq. (2.6) holds for w , one readily deduces that

$$\begin{aligned}
 T_{s_i w} X &= T_{s_i} X^{(-1)^{\ell(w)}} \left(T_w + \sum_{v \in W, v < w} a(\ell(w) - \ell(v)) T_v \right) \\
 &\stackrel{(i)}{=} X^{(-1)^{\ell(s_i w)}} (T_{s_i} - (\tau - \tau^{-1})) \left(T_w + \sum_{v \in W, v < w} a(\ell(w) - \ell(v)) T_v \right) \\
 &\stackrel{(ii)}{=} X^{(-1)^{\ell(s_i w)}} \left(T_{s_i w} + \sum_{\substack{v \in W, v < s_i w \\ s_i v < v}} a(\ell(s_i w) - \ell(v)) T_v \right. \\
 &\quad \left. + \sum_{\substack{v \in W, v < s_i w \\ s_i v > v}} (a(\ell(s_i w) - \ell(v) - 2) - (\tau - \tau^{-1})a(\ell(s_i w) - \ell(v) - 1)) T_v \right) \\
 &\stackrel{(iii)}{=} X^{(-1)^{\ell(s_i w)}} \left(T_{s_i w} + \sum_{v \in W, v < s_i w} a(\ell(s_i w) - \ell(v)) T_v \right).
 \end{aligned}$$

Here we have used (i) the relation in Eq. (2.3) with $k = -1$ (upon conjugation by U if $i = 0$), (ii) the property that $T_{s_i} T_v = T_{s_i v}$ if $s_i v > v$ and $T_{s_i} T_v = T_{s_i v} + (\tau - \tau^{-1}) T_v$ if $s_i v < v$ (cf. Eq. (2.2b)), and (iii) the identity $a(k - 2) - (\tau - \tau^{-1})a(k - 1) = a(k)$. \square

Remark 2.2. It is immediate from the relation in Eq. (2.3) (for $k = 1$ and $k = -1$) in combination with the commutation relation $UX^{\pm 1} = X^{\mp 1}U$ that the element $X + X^{-1}$ lies in the center $\mathcal{Z}(\mathbb{H})$ of the double affine Hecke algebra.

3. Representations of \mathbb{H} on $\mathcal{C}(\mathbb{Z})$

For any (fixed) positive integer M , the action $sn := -n$, $un := M - n$ on $n \in \mathbb{Z}$ determines a faithful representation of W on the space $\mathcal{C}(\mathbb{Z})$ of functions $f : \mathbb{Z} \rightarrow \mathbb{C}$ given by $(wf)(n) := f(w^{-1}n)$. A fundamental domain with respect to the action of $W_S \subset W$ on \mathbb{Z} is given by $\Lambda_M := \{0, 1, 2, \dots, M\}$. For any $n \in \mathbb{Z}$, let us denote by w_n the unique shortest element in W_S mapping n into Λ_M , i.e.

$$w_n = \begin{cases} \cdots s_0 s_1 s_0 & \text{with } \ell(w_n) = p & \text{if } n = pM + r, \quad p \in \mathbb{Z}_{>0}, \quad 0 < r \leq M, \\ \cdots s_1 s_0 s_1 & \text{with } \ell(w_n) = p & \text{if } n = -pM + r, \quad p \in \mathbb{Z}_{>0}, \quad 0 \leq r < M, \\ 1 & & \text{if } 0 \leq n \leq M. \end{cases} \tag{3.1}$$

With the aid of this representation of W , we now introduce two explicit representations of the double affine Hecke algebra \mathbb{H} on $\mathcal{C}(\mathbb{Z})$. For convenience, we often employ the shorthand conventions f_n for $f(n)$ and n_+ for $w_n n$, and it will moreover be assumed that $M > 1$ from now on.

3.1. Difference-reflection representation

Let $\hat{T} : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ be the operator

$$\hat{T} := \tau + \tau^{\text{sign}}(s - 1), \tag{3.2a}$$

where τ and τ^{sign} act by multiplication, i.e. $(\tau f)_n := \tau f_n$ and $(\tau^{\text{sign}} f)_n := \tau^{\text{sign}(n)} f_n$ (with the convention that $\text{sign}(0) := 0$). More specifically, one has that

$$(\hat{T}f)_n = \begin{cases} \tau f_{-n} & \text{if } n \geq 0, \\ \tau^{-1} f_{-n} + (\tau - \tau^{-1}) f_n & \text{if } n < 0. \end{cases} \tag{3.2b}$$

Similarly, let $\hat{X} : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ denote the operator characterized by

$$\begin{aligned}
 (\hat{X}f)_n &:= \tau^{(\ell(w_{n-1})-\ell(w_n))(1+\eta(w_{un}))} f_{n-1} \\
 &+ \text{sign}(n) \sum_{v \in W, v < w_n} \tau^{-\ell(w_n)-\ell(v)} a(\ell(w_n) - \ell(v)) f_{v(n-\text{sign}(n))_+}, \tag{3.3}
 \end{aligned}$$

with $a(\cdot)$ taken from Eq. (2.5b).

Theorem 3.1. *The assignment $T \rightarrow \hat{T}, U \rightarrow u, X \rightarrow \hat{X}$ extends (uniquely) to a representation of the double affine Hecke algebra \mathbb{H} on $\mathcal{C}(\mathbb{Z})$, i.e.*

$$(\hat{T} - \tau)(\hat{T} + \tau^{-1}) = 0 \quad \text{and} \quad u\hat{X}u = \hat{X}^{-1} = \hat{T}^{-1}\hat{X}\hat{T}^{-1}. \tag{3.4}$$

It is not hard to verify the quadratic relation for \hat{T} in Eq. (3.4) directly, however, to check this way the validity of the commutation relations involving u, \hat{X} and \hat{T}, \hat{X} turns out to be very cumbersome. Instead, Theorem 3.1 will follow below in Section 4 as a consequence of intertwining relations connecting to a second representation of \mathbb{H} in terms of integral-reflection operators.

3.2. Integral-reflection representation

Let I be a discrete integral-reflection operator on $\mathcal{C}(\mathbb{Z})$ of the form

$$I := \tau s + (\tau - \tau^{-1})J, \tag{3.5a}$$

where

$$(Jf)_n := \begin{cases} -f_{n-2} - f_{n-4} - \dots - f_{-n} & \text{if } n > 0, \\ 0 & \text{if } n = 0, \\ f_n + f_{n+2} + \dots + f_{-n-2} & \text{if } n < 0 \end{cases} \tag{3.5b}$$

(thus in essence integrating f from n to $-n$ with step 2), and let D denote the translation operator on lattice functions given by $(Df)_n := f_{n-1}$.

Theorem 3.2. *The assignment $T \rightarrow I, U \rightarrow u, X \rightarrow D$ extends (uniquely) to a representation of the double affine Hecke algebra \mathbb{H} on $\mathcal{C}(\mathbb{Z})$, i.e.*

$$(I - \tau)(I + \tau^{-1}) = 0 \quad \text{and} \quad uDu = D^{-1} = I^{-1}DI^{-1}. \tag{3.6}$$

Proof. From the definition of I it is clear that

$$I^2 - (\tau - \tau^{-1})I - 1 = (\tau - \tau^{-1})^2(J^2 - J) + (\tau^2 - 1)(sJ + Js - s + 1).$$

To verify the quadratic relation for I it is therefore sufficient to show that (a) $J^2 = J$ and (b) $sJ + Js = s - 1$. The latter two relations follow from the following elementary observations: (i) if f is odd then $Jf = f$, (ii) Jf is odd for any f , (iii) $sf - f$ is odd for any f . Indeed, one has that $J^2 f \stackrel{(i),(ii)}{=} Jf$ and $(sJ + Js)f \stackrel{(ii)}{=} J(s-1)f \stackrel{(i),(iii)}{=} (s-1)f$ for any $f \in \mathcal{C}(\mathbb{Z})$.

It is immediate from the definitions that $uDu = D^{-1}$, and to infer the last relation involving I and D it is convenient to first recast it in the form $D^{-1}I = I^{-1}D$. Upon remarking that $I^{-1} = I - (\tau - \tau^{-1})$ (by the quadratic relation for I), substitution of Eq. (3.5a), and invoking of the elementary property

$sDs = D^{-1}$, one ends up with the relation $D = JD - D^{-1}J$, which is readily verified by acting with both sides on an arbitrary function $f \in \mathcal{C}(\mathbb{Z})$. \square

Remark 3.3. The relation $\hat{X}^{-1} = u\hat{X}u$ gives rise to the following explicit formula for the action of the inverse operator \hat{X}^{-1} on $f \in \mathcal{C}(\mathbb{Z})$:

$$\begin{aligned}
 (\hat{X}^{-1}f)_n &= \tau^{(\ell(w_{n+1})-\ell(w_n))(1+\eta(w_n))} f_{n+1} \\
 &\quad - \text{sign}(n) \sum_{v \in W, v < w_n} \tau^{-\ell(w_n)-\ell(v)} a(\ell(w_n) - \ell(v)) f_{v(n-\text{sign}(n))_+}. \tag{3.7}
 \end{aligned}$$

4. Intertwining operator

For $h \in \mathbb{H}$, we will denote by $\hat{T}(h)$ and $I(h)$ the images of h under the representations of Theorems 3.1 and 3.2, respectively, and we will furthermore employ the shorthand notations $\hat{T}_w := \hat{T}(T_w)$ and $I_w := I(T_w)$ ($w \in W$).

Let $\mathcal{J} : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ be the operator of the form

$$(\mathcal{J}f)_n := \tau^{-\ell(w_n)} (I_{w_n}f)_{n_+} \quad (f \in \mathcal{C}(\mathbb{Z}), n \in \mathbb{Z}), \tag{4.1}$$

where (recall) $n_+ := w_n n$.

Proposition 4.1. *The operator \mathcal{J} (4.1) constitutes a linear automorphism of the space $\mathcal{C}(\mathbb{Z})$.*

Proposition 4.2. *The operator \mathcal{J} (4.1) satisfies the intertwining relations*

$$\hat{T}\mathcal{J} = \mathcal{J}I, \quad u\mathcal{J} = \mathcal{J}u, \quad \hat{X}\mathcal{J} = \mathcal{J}D. \tag{4.2}$$

It is immediate from Propositions 4.1 and 4.2 that the double affine Hecke algebra relations in Theorem 3.1 follow from those in Theorem 3.2. In other words, the operator $\mathcal{J} : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ in Eq. (4.1) provides an explicit intertwining operator between both representations of the double affine Hecke algebra:

$$\hat{T}(h)\mathcal{J} = \mathcal{J}I(h) \quad (\forall h \in \mathbb{H}). \tag{4.3}$$

4.1. Proof of Proposition 4.1

We need to show that $\mathcal{J} : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ is a bijection, which is immediate from the following triangularity property:

$$(I_{w_n}f)_{n_+} = \tau^{-\ell(w_n)} f_n + \sum_{k \in \mathbb{Z}, k < n} *f_k \quad (f \in \mathcal{C}(\mathbb{Z}), n \in \mathbb{Z}), \tag{4.4}$$

with $k < n$ iff either (i) $w_k < w_n$ or (ii) $w_k = w_n$ and $|k| < |n|$. Here and below the star symbols $*$ refer to the expansion coefficients of lower terms (with respect to the partial order $<$) whose precise values are not relevant for the argument of the proof. Indeed, it is clear from the triangularity in Eq. (4.4) that for any $g \in \mathcal{C}(\mathbb{Z})$ the linear equation $(\mathcal{J}f)_n = g_n$ ($n \in \mathbb{Z}$) can be uniquely solved by induction in n with respect to $<$. We will now prove the triangularity in question by induction on $\ell(w_n)$ starting from the straightforward case that $\ell(w_n) \leq 1$. (The case $\ell(w_n) = 0$ is in fact trivial since then $n \in \Delta_M$ and $(I_{w_n}f)_{n_+} = f_n$.) To this end we first observe that for $\ell(w_n) = 1$ the triangularity in Eq. (4.4) is a direct consequence of the following lemma.

Lemma 4.3. For $i \in \{0, 1\}$, the action of I_{s_i} on $f \in \mathcal{C}(\mathbb{Z})$ is of the form

$$(I_{s_i} f)_n = \begin{cases} \tau^{-1} f_{s_i n} + \sum_{\substack{k \in \mathbb{Z} \\ w_k < w_{s_i n}} * f_k + \sum_{\substack{k \in \mathbb{Z}, |k| < |s_i n| \\ w_k = w_{s_i n}} * f_k & \text{if } \ell(w_{s_i n}) > \ell(w_n), \\ \sum_{\substack{k \in \mathbb{Z} \\ w_k < w_n}} * f_k + \sum_{\substack{k \in \mathbb{Z}, |k| \leq |n| \\ w_k = w_n}} * f_k & \text{if } \ell(w_{s_i n}) \leq \ell(w_n). \end{cases}$$

Proof. From the formula for the action in Eqs. (3.5a), (3.5b) and its conjugation with the action of u one deduces that

$$(I_{s_1} f)_n = (If)_n = \begin{cases} \tau^{-1} f_{-n} - (\tau - \tau^{-1}) \sum_{k=1}^{n-1} f_{n-2k} & \text{if } n > 0, \\ \tau f_{-n} + (\tau - \tau^{-1}) \sum_{k=0}^{-n-1} f_{n+2k} & \text{if } n \leq 0, \end{cases}$$

and

$$(I_{s_0} f)_n = (uIf)_n = \begin{cases} \tau^{-1} f_{2M-n} - (\tau - \tau^{-1}) \sum_{k=1}^{M-n-1} f_{n+2k} & \text{if } n < M, \\ \tau f_{2M-n} + (\tau - \tau^{-1}) \sum_{k=0}^{n-M-1} f_{n-2k} & \text{if } n \geq M. \end{cases}$$

These explicit formulas are readily seen to imply the lemma (upon employing the formula in Eq. (3.1) to determine the lengths of the relevant group elements w_{n-2k} and w_{n+2k} , together with the equivalences $n > 0 \Leftrightarrow \ell(w_{s_1 n}) > \ell(w_n)$ and $n < M \Leftrightarrow \ell(w_{s_0 n}) > \ell(w_n)$). \square

Now let us assume that $n \in \mathbb{Z}$ with $\ell(w_n) \geq 1$ is such that the triangularity in Eq. (4.4) holds for all $k \in \mathbb{Z}$ with $\ell(w_k) \leq \ell(w_n)$, and let $m = s_i n$ with $\ell(w_m) = \ell(w_n) + 1$ (so $w_m = w_n s_i$). Our induction hypothesis guarantees that

$$(I_{w_m} f)_{m_+} = (I_{w_n} I_{s_i} f)_{n_+} = \tau^{-\ell(w_n)} (I_{s_i} f)_n + \sum_{\substack{l \in \mathbb{Z} \\ w_l < w_n}} * (I_{s_i} f)_l + \sum_{\substack{l \in \mathbb{Z}, |l| < |n| \\ w_l = w_n}} * (I_{s_i} f)_l. \tag{4.5}$$

Application of Lemma 4.3 to all terms on the RHS of Eq. (4.5) shows that $(I_{w_m} f)_{m_+}$ is equal to $\tau^{-\ell(w_m)} f_m$ plus a linear combination of terms involving evaluations f_k with $k < m$, which completes the induction step (and thus the proof of the proposition). Indeed, one has that (i) $\tau^{-\ell(w_n)} (I_{s_i} f)_n = \tau^{-\ell(w_n)-1} f_{s_i n} + \text{lower terms} = \tau^{-\ell(w_m)} f_m + \text{lower terms}$, (ii) if $w_l < w_n$ then $w_{s_i l} \leq w_l s_i < w_n s_i = w_{s_i n} = w_m$, and (iii) if $w_l = w_n$ with $|l| < |n|$ then $\ell(w_l s_i) = \ell(w_l) + 1$ and $|s_i l| < |s_i n|$, whence $w_{s_i l} = w_l s_i = w_n s_i = w_m$ and $|s_i l| < |m|$.

4.2. Proof of Proposition 4.2

Let $f \in \mathcal{C}(\mathbb{Z})$ and $n \in \mathbb{Z}$.

4.2.1. Proof of the relation $\hat{T}\mathcal{J} = \mathcal{J}I$

Straightforward manipulations reveal that

$$\begin{aligned} (\hat{T}\mathcal{J}f)_n &= \tau(\mathcal{J}f)_n + \tau^{\text{sign}(n)}((\mathcal{J}f)_{-n} - (\mathcal{J}f)_n) \\ &\stackrel{\text{(i)}}{=} \tau^{1-\ell(w_n)}(I_{w_n} f)_{n_+} + \tau^{\text{sign}(n)}(\tau^{-\ell(w_n s)}(I_{w_n s} f)_{n_+} - \tau^{-\ell(w_n)}(I_{w_n} f)_{n_+}) \\ &\stackrel{\text{(ii)}}{=} \tau^{-\ell(w_n)} \left((I_{w_n s} f)_{n_+} + \frac{1}{2}(1 - \eta(w_n))(\tau - \tau^{-1})(I_{w_n} f)_{n_+} \right) \\ &\stackrel{\text{(iii)}}{=} \tau^{-\ell(w_n)}(I_{w_n} If)_{n_+} = (\mathcal{J}If)_n \end{aligned}$$

(with $\eta : W \rightarrow \{-1, 1\}$ as defined in Eq. (2.2c)). In steps (i) and (ii) we exploited that for $n \neq 0$: $w_{-n} = w_n s$ and $\ell(w_n s) = \ell(w_n) + \text{sign}(n)$, whereas for $n = 0$: $\tau^{-\ell(w_n)}(I_{w_n} f)_{n_+} = f_0 = \tau^{-1}(I f)_0 = \tau^{-\ell(w_n s)}(I_{w_n s} f)_{n_+}$. Step (iii) hinges in turn on the Hecke algebra relation in Eq. (2.2a).

4.2.2. Proof of the relation $u\mathcal{J} = \mathcal{J}u$

It is immediate from the definitions that

$$\begin{aligned} (u\mathcal{J}f)_n &= (\mathcal{J}f)_{un} = \tau^{-\ell(w_{un})}(I_{w_{un}} f)_{w_{un}un} \\ &= \tau^{-\ell(w_n)}(uI_{w_n} u f)_{un_+} = \tau^{-\ell(w_n)}(I_{w_n} u f)_{n_+} = (\mathcal{J}u f)_n \end{aligned}$$

(where it was used that $w_{un} = u w_n u$ and $I_{u w_n u} = u I_{w_n} u$).

4.2.3. Proof of the relation $\hat{X}\mathcal{J} = \mathcal{J}D$

The proof hinges on the Hecke algebra relation of Proposition 2.1 (with $r = 0$ and $\varepsilon = 1$):

$$\begin{aligned} (\mathcal{J}Df)_n &= \tau^{-\ell(w_n)}(I_{w_n} Df)_{n_+} \\ &= \tau^{-\ell(w_n)}(D^{(-1)^{\ell(w_n)}} I_{w_n} f)_{n_+} \\ &\quad + \eta(w_n) \tau^{-\ell(w_n)} \sum_{v \in W, v < w_n} a(\ell(w_n) - \ell(v)) (D^{\eta(w_n)(-1)^{\ell(w_n)}} I_v f)_{n_+} \\ &\stackrel{(i)}{=} \tau^{-\ell(w_n)}(I_{w_n} f)_{n_+ - (-1)^{\ell(w_n)}} \\ &\quad + \text{sign}(n) \tau^{-\ell(w_n)} \sum_{v \in W, v < w_n} a(\ell(w_n) - \ell(v)) (I_v f)_{n_+ - \text{sign}(n)(-1)^{\ell(w_n)}} \\ &\stackrel{(ii)}{=} \tau^{(\ell(w_{n-1}) - \ell(w_n))(1 + \eta(w_{un}))} (\mathcal{J}f)_{n-1} \\ &\quad + \text{sign}(n) \sum_{v \in W, v < w_n} \tau^{-\ell(w_n) - \ell(v)} a(\ell(w_n) - \ell(v)) (\mathcal{J}f)_{v^{-1}(n - \text{sign}(n))_+} \\ &= (\hat{X}\mathcal{J}f)_n. \end{aligned}$$

Here we used in step (i) that $\eta(w_n) = \text{sign}(n)$ for $n \neq 0$, and in step (ii) that

$$\begin{aligned} n_+ - (-1)^{\ell(w_n)} &= w_n(n - 1) \\ &= \begin{cases} s(n - 1)_+ (= -1) & \text{if } n \equiv 0 \pmod{2M} \text{ with } n \leq 0, \\ s_0(n - 1)_+ (= M + 1) & \text{if } n \equiv M \pmod{2M} \text{ with } n < 0, \\ (n - 1)_+ & \text{otherwise} \end{cases} \end{aligned}$$

and $n_+ - \text{sign}(n)(-1)^{\ell(w_n)} = w_n(n - \text{sign}(n)) = (n - \text{sign}(n))_+$. Notice in this connection that $\ell(w_{n-1}) = \ell(w_n) + 1$ if $n \equiv 0 \pmod{M}$ with $n \leq 0$, and that (for any $f \in \mathcal{C}(\mathbb{Z})$): $f_{-1} = \tau(I f)_1$, $f_{M+1} = \tau(I_{s_0} f)_{M-1}$, and moreover $\tau^{-\ell(v)}(I_v f)_{vn} = (\mathcal{J}f)_n$ for (any) $v \in W$ with $vn = n_+$ (since $\tau^{-1}(I f)_0 = f_0$ and $\tau^{-1}(I_{s_0} f)_M = f_M$).

5. Unitarity and the Laplacian

The assignment $T_W \rightarrow T_W^* := T_{W^{-1}}$ ($w \in W$) extends to an antilinear anti-involution $*$ of the affine Hecke algebra $H \subset \mathbb{H}$ turning it into a $*$ -algebra. Let $l^2(\mathbb{Z}, \delta)$ be the Hilbert space inside $\mathcal{C}(\mathbb{Z})$ deter-

mined by the inner product

$$\langle f, g \rangle_\delta := \sum_{n \in \mathbb{Z}} f_n \bar{g}_n \delta_n, \quad \delta_n := \tau^{2\ell(w_n)} \tag{5.1}$$

(where the bar refers to the complex conjugate).

Proposition 5.1. *The difference-reflection representation $h \rightarrow \hat{T}(h)$ ($h \in \mathbb{H}$) restricts to a unitary representation of the affine Hecke algebra $H \subset \mathbb{H}$ into the space of bounded operators on $l^2(\mathbb{Z}, \delta)$, i.e.*

$$\langle \hat{T}(h)f, g \rangle_\delta = \langle f, \hat{T}(h^*)g \rangle_\delta \quad (h \in H, f, g \in l^2(\mathbb{Z}, \delta)). \tag{5.2}$$

Proof. Clearly it is sufficient to prove the proposition for the generators u and $\hat{T} = \tau + \tau^{\text{sign}(s-1)}$. The corresponding operators are bounded in $l^2(\mathbb{Z}, \delta)$ because $\langle uf, uf \rangle_\delta = \langle f, f \rangle_\delta$ as $u\delta = \delta$ and $\langle sf, sf \rangle_\delta = \langle \tau^{2\text{sign} f}, f \rangle_\delta$ as $s\delta = \tau^{2\text{sign} \delta}$. (Notice in this connection also that $\langle \tau^{\text{sign} f}, f \rangle_\delta \leq \tau^{-1} \langle f, f \rangle_\delta$.) In an analogous manner it is seen that $\langle uf, g \rangle_\delta = \langle f, ug \rangle_\delta$ and that $\langle \hat{T}f, g \rangle_\delta = \langle (\tau - \tau^{\text{sign}})f + \tau^{\text{sign}}sf, g \rangle_\delta = \langle f, (\tau - \tau^{\text{sign}})g + \delta^{-1} s \tau^{\text{sign} \delta} g \rangle_\delta = \langle f, \hat{T}g \rangle_\delta$ (using also that $s\tau^{\text{sign}} = \tau^{-\text{sign} s}$). \square

Let us define the *Laplacian* on $\mathcal{C}(\mathbb{Z})$ associated with the (center of the) double affine Hecke algebra \mathbb{H} as the operator $L : \mathcal{C}(\mathbb{Z}) \rightarrow \mathcal{C}(\mathbb{Z})$ of the form

$$L := \hat{T}(X + X^{-1}) = \hat{X} + \hat{X}^{-1}. \tag{5.3}$$

It is immediate from the action of \hat{X} and \hat{X}^{-1} in Eqs. (3.3) and (3.7), that the Laplacian L (5.3) is a discrete difference operator whose explicit action on $f \in \mathcal{C}(\mathbb{Z})$ reads

$$(Lf)_n = \tau^{(\ell(w_{n+1})-\ell(w_n))(1+\eta(w_n))} f_{n+1} + \tau^{(\ell(w_{n-1})-\ell(w_n))(1+\eta(w_{un}))} f_{n-1} \tag{5.4}$$

($n \in \mathbb{Z}$), or equivalently:

$$(Lf)_n = a_n f_{n+1} + b_n f_{n-1} \quad (n \in \mathbb{Z}), \tag{5.5a}$$

with

$$a_n = \begin{cases} \tau^2 & \text{if } n \in M\mathbb{Z}_{>0}, \\ 1 & \text{otherwise,} \end{cases} \quad b_n = \begin{cases} \tau^2 & \text{if } n \in M\mathbb{Z}_{\leq 0}, \\ 1 & \text{otherwise.} \end{cases} \tag{5.5b}$$

For $\tau \rightarrow 1$, the action of L (5.3) degenerates to the action of the standard (undeformed) discrete Laplacian on $\mathcal{C}(\mathbb{Z})$ (shifted by an additive constant such that the diagonal terms vanish).

Proposition 5.2. *The Laplacian L (5.3) constitutes a (bounded) self-adjoint operator on $l^2(\mathbb{Z}, \delta)$, i.e.*

$$\langle Lf, g \rangle_\delta = \langle f, Lg \rangle_\delta \quad (f, g \in l^2(\mathbb{Z}, \delta)). \tag{5.6}$$

Proof. It is manifest from Eqs. (5.5a), (5.5b) that the operator L is bounded on $l^2(\mathbb{Z}, \delta)$. Moreover, one has that

$$\begin{aligned} \langle Lf, g \rangle_\delta &= \sum_{n \in \mathbb{Z}} (a_n f_{n+1} + b_n f_{n-1}) \bar{g}_n \delta_n \\ &= \sum_{n \in \mathbb{Z}} f_n (a_{n-1} \bar{g}_{n-1} \delta_{n-1} + b_{n+1} \bar{g}_{n+1} \delta_{n+1}) \\ &= \sum_{n \in \mathbb{Z}} f_n (a_n \bar{g}_{n+1} + b_n \bar{g}_{n-1}) \delta_n = \langle f, Lg \rangle_\delta, \end{aligned}$$

where it was used that $a_n \delta_n = b_{n+1} \delta_{n+1}$ ($\forall n \in \mathbb{Z}$). \square

6. Spherical function

It is well known that the eigenfunctions of the (standard) discrete Laplacian in $\mathcal{C}(\mathbb{Z})$ restricted to the finite-dimensional invariant subspace of W_S -symmetric functions are given by the elements of the (appropriate) discrete Fourier-cosine basis [22] (cf. also Remark 7.4 below). For our deformed Laplacian L (5.3), the role of the cosine function will be taken over by a spherical function associated with the double affine Hecke algebra.

Parameterized by an auxiliary spectral variable $\xi \in \mathbb{R}$, we define the spherical function in $\mathcal{C}(\mathbb{Z})$ associated with \mathbb{H} as

$$\Phi_\xi := \mathcal{J} \phi_\xi \quad \text{with } \phi_\xi := (1 + \tau I) \mathbf{e}^{i\xi}, \tag{6.1}$$

where $\mathbf{e}^{i\xi} \in \mathcal{C}(\mathbb{Z})$ represents the plane wave function $\mathbf{e}^{i\xi}(n) := e^{i\xi n}$, $n \in \mathbb{Z}$. From the explicit action of I (3.5a), (3.5b) on $\mathbf{e}^{i\xi}$:

$$I \mathbf{e}^{i\xi} = \tau \mathbf{e}^{-i\xi} + (\tau - \tau^{-1}) \frac{\mathbf{e}^{i\xi} - \mathbf{e}^{-i\xi}}{1 - e^{2i\xi}},$$

it is readily seen that ϕ_ξ can be expressed as a linear combination of plane waves

$$\phi_\xi(n) = c(\xi) e^{i\xi n} + c(-\xi) e^{-i\xi n}, \tag{6.2a}$$

with

$$c(\xi) := \frac{1 - \tau^2 e^{-2i\xi}}{1 - e^{-2i\xi}}. \tag{6.2b}$$

This representation of ϕ_ξ can be conveniently rewritten as a linear combination of two Chebyshev polynomials of the second kind:

$$\phi_\xi(n) = U_n(\cos(\xi)) + \tau^2 U_{-n}(\cos(\xi)) \quad \text{with } U_n(\cos(\xi)) = \sin((n+1)\xi) / \sin(\xi),$$

whence it is a polynomial of degree $|n|$ in $\cos(\xi)$. These trigonometric polynomials in the spectral parameter are often referred to as the one-dimensional Hall–Littlewood polynomials. It is manifest from the above explicit expressions that ϕ_ξ and thus also Φ_ξ are real-valued: $\overline{\phi_\xi} = \phi_\xi$, $\overline{\Phi_\xi} = \Phi_\xi$.

Let $\mathcal{C}(\mathbb{Z})^W := \{f \in \mathcal{C}(\mathbb{Z}) \mid wf = f, \forall w \in W\}$ and $\mathcal{C}(\mathbb{Z})^{W_S} := \{f \in \mathcal{C}(\mathbb{Z}) \mid wf = f, \forall w \in W_S\}$.

Proposition 6.1. *The spherical function Φ_ξ (6.1) lies in the subspace $\mathcal{C}(\mathbb{Z})^{W_S}$ of W_S -invariant functions if the spectral parameter $\xi \in \mathbb{R}$ satisfies the equation*

$$e^{iM\xi} = \varepsilon \left(\frac{1 - \tau^2 e^{2i\xi}}{\tau^2 - e^{2i\xi}} \right), \quad \varepsilon \in \{1, -1\}, \tag{6.3}$$

with the case $\varepsilon = 1$ corresponding to the situation that $\Phi_\xi \in \mathcal{C}(\mathbb{Z})^W \subset \mathcal{C}(\mathbb{Z})^{W_S}$.

Proof. From the quadratic relation $T(1 + \tau T) = \tau(1 + \tau T)$ together with the intertwining relations in Proposition 4.2, it is seen that $\hat{T}\Phi_\xi = \hat{T}\mathcal{J}\phi_\xi = \mathcal{J}I\phi_\xi = \mathcal{J}I(1 + \tau I)\mathbf{e}^{i\xi} = \mathcal{J}\tau(1 + \tau I)\mathbf{e}^{i\xi} = \tau\Phi_\xi$, whence $s\Phi_\xi = \Phi_\xi$ (for any $\xi \in \mathbb{R}$). Moreover, from Propositions 4.1, 4.2 it is clear that $u\Phi_\xi = \varepsilon\Phi_\xi$ if and only if $u\phi_\xi = \varepsilon\phi_\xi$ ($\varepsilon \in \{1, -1\}$). The explicit expansion in Eqs. (6.2a), (6.2b) reveals that this relation for ϕ_ξ holds if $e^{iM\xi} = \varepsilon c(-\xi)/c(\xi) = \varepsilon(1 - \tau^2 e^{2i\xi})/(\tau^2 - e^{2i\xi})$. \square

Eq. (6.3) for the spectral parameter turns out to be the simplest case of a Bethe–Ansatz equation for the discrete Bose gas on the circle studied in Ref. [5] (cf. Remark 7.5 below). We will now describe what the solutions of these Bethe–Ansatz equations in [5] amount to in the present situation (with proofs included to keep the presentation self-contained). Let

$$V(\xi) := M\xi + \theta(\xi), \tag{6.4a}$$

with

$$\theta(\xi) := (1 - \tau^4) \int_0^{2\xi} \frac{dx}{1 + \tau^4 - 2\tau^2 \cos x} \tag{6.4b}$$

$$= 2 \arctan\left(\frac{1 + \tau^2}{1 - \tau^2} \tan \frac{\xi}{2}\right) = i \log\left(\frac{1 - \tau^2 e^{2i\xi}}{e^{2i\xi} - \tau^2}\right). \tag{6.4c}$$

Here the branches of $\arctan(\cdot)$ and $\log(\cdot)$ are assumed to be chosen such that θ is quasi-periodic: $\theta(\xi + \pi) = \theta(\xi) + 2\pi$ and $\theta(\xi)$ varies from $-\pi$ to π as ξ varies from $-\pi/2$ to $\pi/2$ (which corresponds to the principal branch). Our parameter regime $\tau \in (0, 1)$ moreover guarantees that the odd function θ is smooth and strictly monotonously increasing on \mathbb{R} . The following proposition characterizes those solutions of Eq. (6.3) that are relevant for our purposes.

Proposition 6.2.

- (i) For any $m \in \Lambda_M$, the equation $V(\xi) = (m + 1)\pi$ has a unique real-valued solution ξ_m .
- (ii) We have that $0 < \xi_0 < \xi_1 < \dots < \xi_M < \pi$ and that $\xi_{um} = \xi_{M-m} = \pi - \xi_m$.
- (iii) The spectral value ξ_m solves Eq. (6.3) with $\varepsilon = 1$ if $m \in \Lambda_M^o := 2\mathbb{Z} \cap \Lambda_M$ and with $\varepsilon = -1$ if $m \in \Lambda_M^e := (2\mathbb{Z} + 1) \cap \Lambda_M$.

Proof. Part (i) follows from the fact that $V : \mathbb{R} \rightarrow \mathbb{R}$ is bijective. Indeed, the function $V(\xi)$ is strictly monotonously increasing on \mathbb{R} with $V(\xi) \rightarrow -\infty$ as $\xi \rightarrow -\infty$ and $V(\xi) \rightarrow +\infty$ as $\xi \rightarrow +\infty$.

The first statement of part (ii) is clear from the strict monotonicity of $V(\xi)$ and the fact that $V(0) = 0$ and $V(\pi) = (M + 2)\pi$; the second statement is manifest from the equality $V(\xi_{M-m}) = (M - m + 1)\pi = V(\pi - \xi_m)$.

By definition, one has that $M\xi_m + \theta(\xi_m) = (m + 1)\pi$. Upon multiplication by i and exponentiating both sides of the equation one deduces that $e^{iM\xi_m} = (-1)^m(1 - \tau^2 e^{2i\xi_m})/(\tau^2 - e^{2i\xi_m})$, which implies part (iii). \square

For the spectral values ξ_m , $m \in \Lambda_M$ the spherical function Φ_ξ (6.1) can now be written explicitly in terms of the one-dimensional Hall–Littlewood polynomials:

$$\Phi_{\xi_m}(n) = c(\xi_m)e^{i\xi_m n_+} + c(-\xi_m)e^{-i\xi_m n_+} \quad (m \in \Lambda_M, n \in \mathbb{Z}). \tag{6.5}$$

Indeed, for any $f \in \mathcal{C}(\mathbb{Z})$ and $n \in \Lambda_M$ one has that $(\mathcal{J}f)_n = f_n$ (by definition). Hence, Eq. (6.5) is immediate from Eqs. (6.2a), (6.2b), first for $n \in \Lambda_M$ and then (trivially) extended to arbitrary $n \in \mathbb{Z}$ using the W_S -invariance of Φ_{ξ_m} , $m \in \Lambda_M$ (following from Propositions 6.1 and 6.2).

7. Discrete Fourier transform

Since Λ_M is a fundamental domain for the action of W_S on \mathbb{Z} , the W_S -invariant subspace $l^2(\mathbb{Z}, \delta)^{W_S} := l^2(\mathbb{Z}, \delta) \cap C(\mathbb{Z})^{W_S}$ can be identified with the finite-dimensional space $l^2(\Lambda_M, \Delta)$ of functions $f : \Lambda_M \rightarrow \mathbb{C}$ endowed with the inner product

$$\langle f, g \rangle_\Delta := \sum_{n \in \Lambda_M} f_n \bar{g}_n \Delta_n \quad (f, g \in l^2(\Lambda_M, \Delta)), \tag{7.1a}$$

with respect to a (suitably normalized) symmetrized weight function $\Delta : \Lambda_M \rightarrow (0, \infty)$ of the form

$$\Delta_n := \frac{1}{W_S(\tau^2)} \sum_{m \in W_S n} \delta_m = \begin{cases} (1 + \tau^2)^{-1} & \text{if } n = 0, M, \\ 1 & \text{if } 0 < n < M. \end{cases} \tag{7.1b}$$

Here we have employed the shorthand notation

$$W_S(\tau^2) := \sum_{w \in W_S} \tau^{2\ell(w)} = 1 + 2 \sum_{k=1}^\infty \tau^{2k} = (1 + \tau^2)/(1 - \tau^2).$$

The computation of Δ_n (7.1b) hinges on the following elementary identity:

$$\begin{aligned} \sum_{m \in W_S n} \delta_m &= \sum_{m \in W_S n} \tau^{2\ell(w_m)} \\ &= \begin{cases} \sum_{k=0}^\infty \tau^{2k} = 1/(1 - \tau^2) & \text{if } n = 0, M, \\ (1 + \tau^2) \sum_{k=0}^\infty \tau^{2k} = (1 + \tau^2)/(1 - \tau^2) & \text{if } 0 < n < M. \end{cases} \end{aligned}$$

Since $X + X^{-1}$ is a central element in \mathbb{H} (cf. Remark 2.2 above), it is clear that $L\hat{T} = \hat{T}L$ and $Lu = uL$, which implies that L maps the invariant subspace $C(\mathbb{Z})^{W_S} = \{f \in C(\mathbb{Z}) \mid sf = f, uf = \pm f\} = \{f \in C(\mathbb{Z}) \mid \hat{T}f = \tau f, uf = \pm f\}$ into itself. Indeed, with the aid of Proposition 5.2 and Eqs. (5.5a), (5.5b) one infers that L restricts to a self-adjoint in operator in $l^2(\Lambda_M, \Delta)$ whose explicit action on $f \in l^2(\Lambda_M, \Delta)$ is given by

$$(Lf)_n = \begin{cases} (1 + \tau^2)f_1 & \text{if } n = 0, \\ f_{n+1} + f_{n-1} & \text{if } 0 < n < M, \\ (1 + \tau^2)f_{M-1} & \text{if } n = M. \end{cases} \tag{7.2}$$

Theorem 7.1. *The spherical functions $\Phi_{\xi_m} : \Lambda_M \rightarrow \mathbb{R}$, $m \in \Lambda_M$ are eigenfunctions of the Laplacian L in $l^2(\Lambda_M, \Delta)$:*

$$L\Phi_{\xi_m} = 2 \cos(\xi_m)\Phi_{\xi_m} \quad (m \in \Lambda_M).$$

Proof. Elementary manipulations involving the explicit expansion in Eq. (6.5) reveal that $L\Phi_{\xi_m} = (\hat{X} + \hat{X}^{-1})\mathcal{J}\phi_{\xi_m} = \mathcal{J}(D + D^{-1})(c(\xi_m)\mathbf{e}^{i\xi_m} + c(-\xi_m)\mathbf{e}^{-i\xi_m}) = 2 \cos(i\xi_m)\Phi_{\xi_m}$, since $(D + D^{-1})\mathbf{e}^{i\xi} = 2 \cos(\xi)\mathbf{e}^{i\xi}$. To verify that $\Phi_{\xi_m} \neq 0$, it suffices to compute its value at the origin: $\Phi_{\xi_m}(0) = c(\xi_m) + c(-\xi_m) = 1 + \tau^2 > 0$. \square

It follows from Theorem 7.1 and part (ii) of Proposition 6.2 that the eigenvalues of L in $l^2(\Lambda_M, \Delta)$ are non-degenerate. The spherical functions Φ_{ξ_m} , $m \in \Lambda_M$ thus form an orthogonal basis of $l^2(\Lambda_M, \Delta)$.

To describe the corresponding Plancherel formula it is convenient to introduce the dual space $l^2(\Lambda_M, \hat{\Delta})$ of functions $f : \Lambda_M \rightarrow \mathbb{C}$ endowed with the inner product

$$\langle f, g \rangle_{\hat{\Delta}} := \sum_{m \in \Lambda_M} f_m \bar{g}_m \hat{\Delta}_m \quad (f, g \in l^2(\Lambda_M, \hat{\Delta})), \tag{7.3a}$$

with respect to a weight function $\hat{\Delta} : \Lambda_M \rightarrow (0, \infty)$ of the form

$$\begin{aligned} \hat{\Delta}_m &:= \frac{1}{2c(\xi_m)c(-\xi_m)V'(\xi_m)} \\ &= \frac{(1 - \cos(2\xi_m))}{(1 + \tau^4 - 2\tau^2 \cos(2\xi_m))} \left(M + \frac{2(1 - \tau^4)}{1 + \tau^4 - 2\tau^2 \cos(2\xi_m)} \right)^{-1}. \end{aligned} \tag{7.3b}$$

Let $\Phi : \Lambda_M \times \Lambda_M \rightarrow \mathbb{R}$ denote the kernel function given by $\Phi_{m;n} := \Phi_{\xi_m}(n)$ ($m, n \in \Lambda_M$).

Theorem 7.2.

(i) The functions $\Phi_{m;\cdot} : \Lambda_M \rightarrow \mathbb{R}$, $m \in \Lambda_M$ form an orthogonal basis of $l^2(\Lambda_M, \Delta)$, viz. for any $m, m' \in \Lambda_M$

$$\langle \Phi_{m;\cdot}, \Phi_{m';\cdot} \rangle_{\Delta} = \sum_{n=0}^M \Phi_{m;n} \Phi_{m';n} \Delta_n = \begin{cases} 0 & \text{if } m \neq m', \\ \hat{\Delta}_m^{-1} & \text{if } m = m'. \end{cases}$$

(ii) The functions $\Phi_{\cdot;n} : \Lambda_M \rightarrow \mathbb{R}$, $n \in \Lambda_M$ form an orthogonal basis of $l^2(\Lambda_M, \hat{\Delta})$, viz. for any $n, n' \in \Lambda_M$

$$\langle \Phi_{\cdot;n}, \Phi_{\cdot;n'} \rangle_{\hat{\Delta}} = \sum_{m=0}^M \Phi_{m;n} \Phi_{m;n'} \hat{\Delta}_m = \begin{cases} 0 & \text{if } n \neq n', \\ \Delta_n^{-1} & \text{if } n = n'. \end{cases}$$

Proof. Parts (i) and (ii) are equivalent, since both can be reformulated as the affirmation that the matrix $[\hat{\Delta}_m^{1/2} \Phi_{m;n} \Delta_n^{1/2}]_{0 \leq m, n \leq M}$ is orthogonal. In view of the above mentioned orthogonality of the basis $\Phi_{m;\cdot} = \Phi_{\xi_m}$, $m \in \Lambda_M$ in $l^2(\Lambda_M, \Delta)$, it is thus sufficient to verify the norm formula in part (i) for $m = m'$. An explicit computation reveals that

$$\begin{aligned} \sum_{n=0}^M \Phi_{\xi_m}^2(n) \Delta_n &= \sum_{n=0}^M (c(\xi_m)e^{i\xi_m n} + c(-\xi_m)e^{-i\xi_m n})^2 \Delta_n \\ &= c^2(\xi_m) \left(\sum_{n=1}^{M-1} e^{2i\xi_m n} + \frac{1 + e^{2i\xi_m M}}{1 + \tau^2} \right) + c^2(-\xi_m) \left(\sum_{n=1}^{M-1} e^{-2i\xi_m n} + \frac{1 + e^{-2i\xi_m M}}{1 + \tau^2} \right) \\ &\quad + 2c(\xi_m)c(-\xi_m) \left(M + \frac{1 - \tau^2}{1 + \tau^2} \right), \end{aligned}$$

which is readily seen to simplify to $2c(\xi_m)c(-\xi_m)V'(\xi_m)$ upon summation of the two terminating geometric series and subsequent elimination of all occurrences of exponentials of the form $e^{\pm 2iM\xi_m}$ via the relation $e^{2iM\xi_m} = c^2(-\xi_m)/c^2(\xi_m)$ (cf. Propositions 6.1 and 6.2). \square

Finally, we are now in the position to define the following *discrete Fourier transform* $\mathcal{F} : l^2(\Lambda_M, \Delta) \rightarrow l^2(\Lambda_M, \hat{\Delta})$ associated with \mathbb{H} :

$$\hat{f}_m := (\mathcal{F}f)_m := \langle f, \Phi_{m;\cdot} \rangle_\Delta = \sum_{n=0}^M f_n \Phi_{m;n} \Delta_n \quad (f \in l^2(\Lambda_M, \Delta), m \in \Lambda_M). \tag{7.4a}$$

According to Theorem 7.2 the transform in question constitutes a unitary isomorphism of $l^2(\Lambda_M, \Delta)$ onto $l^2(\Lambda_M, \hat{\Delta})$ with an inversion formula given by

$$f_n = (\mathcal{F}^{-1}\hat{f})_n = \langle \hat{f}, \Phi_{\cdot;n} \rangle_{\hat{\Delta}} = \sum_{m=0}^M \hat{f}_m \Phi_{m;n} \hat{\Delta}_m \quad (\hat{f} \in l^2(\Lambda_M, \hat{\Delta}), n \in \Lambda_M). \tag{7.4b}$$

Furthermore, according to Theorem 7.1 the discrete Fourier transform diagonalizes the Laplacian L in $l^2(\Lambda_M, \Delta)$ in the sense that $L = \mathcal{F}^{-1}\hat{E}\mathcal{F}$, where \hat{E} denotes the real multiplication operator in $l^2(\Lambda_M, \hat{\Delta})$ given by $(\hat{E}f)_m := 2 \cos(\xi_m) f_m$ ($f \in l^2(\Lambda_M, \hat{\Delta}), m \in \Lambda_M$).

Remark 7.3. The Hilbert space $l^2(\Lambda_M, \Delta)$ admits the following natural orthogonal decomposition: $l^2(\Lambda_M, \Delta) = l^2(\Lambda_M, \Delta)^u \oplus l^2(\Lambda_M, \Delta)^{-u}$, where $l^2(\Lambda_M, \Delta)^u := \{f \in l^2(\Lambda_M, \Delta) \mid uf = f\}$ and $l^2(\Lambda_M, \Delta)^{-u} := \{f \in l^2(\Lambda_M, \Delta) \mid uf = -f\}$ (so $l^2(\Lambda_M, \Delta)^u \cong l^2(\mathbb{Z}, \delta) \cap \mathcal{C}(\mathbb{Z})^W$). Similarly, the Hilbert space $l^2(\Lambda_M, \hat{\Delta})$ decomposes orthogonally as $l^2(\Lambda_M, \hat{\Delta}) = l^2(\Lambda_M^e, \hat{\Delta}) \oplus l^2(\Lambda_M^o, \hat{\Delta})$. The discrete Fourier transform \mathcal{F} maps $l^2(\Lambda_M, \Delta)^u$ onto $l^2(\Lambda_M^e, \hat{\Delta})$ and $l^2(\Lambda_M, \Delta)^{-u}$ onto $l^2(\Lambda_M^o, \hat{\Delta})$. Indeed, it follows from Theorem 7.2 combined with (the proof of) Proposition 6.1 and part (iii) of Proposition 6.2, that the spherical functions $\Phi_{\xi_m}, m \in \Lambda_M^e$ and $\Phi_{\xi_m}, m \in \Lambda_M^o$ constitute orthogonal bases of the subspaces $l^2(\Lambda_M, \Delta)^u$ and $l^2(\Lambda_M, \Delta)^{-u}$, respectively.

Remark 7.4. It is not difficult to infer that for $\tau \rightarrow 1$ one has that: $\xi_m \rightarrow \frac{m\pi}{M}, \Phi_{m;n} \rightarrow 2 \cos(\frac{mn\pi}{M})$, and

$$\Delta_n, 2M\hat{\Delta}_n \rightarrow \begin{cases} 1/2 & \text{if } n = 0, M, \\ 1 & \text{if } 0 < n < M. \end{cases}$$

In other words, in this limit our discrete Fourier transform (diagonalizing the Laplacian L (5.3)) degenerates as expected to the discrete cosine transform (diagonalizing the undeformed discrete Laplacian [22,1]).

Remark 7.5. The orthogonality relations in Theorem 7.2(ii) are a finite counterpart of well-known orthogonality relations for the one-dimensional Hall–Littlewood polynomials $\phi_\xi(n), n \in \mathbb{Z}_{\geq 0}$:

$$\frac{1}{2\pi} \int_0^\pi \phi_\xi(n) \phi_\xi(n') \frac{d\xi}{c(\xi)c(-\xi)} = \begin{cases} 0 & \text{if } n \neq n', \\ \mathcal{N}_n & \text{if } n = n', \end{cases}$$

with $\mathcal{N}_0 = 1 + \tau^2$ and $\mathcal{N}_n = 1$ for $n > 0$. They constitute the simplest instances of (conjectural) discrete orthogonality relations for the Hall–Littlewood polynomials [6] originating from the Bethe–Ansatz method for the periodic quantum integrable particle model with pairwise delta-potential interactions known as the delta Bose gas on the circle [5,14,8,13]. From this perspective, the spectral problem for our Laplacian L (5.3) may be viewed as a discrete analog of the textbook spectral problem for a quantum particle on the circle in the presence of a delta-potential. Based on the results presented in this paper, it is natural to expect that the generalization of the representations in Section 3 to the case of double affine Hecke algebras associated with arbitrary Weyl groups will provide an appropriate algebraic framework encoding the symmetries of the integrable discretization of the delta Bose gas on the circle introduced in [5]. This would generalize (i.e. discretize) corresponding recent constructions in terms of degenerations of the double affine Hecke algebra for the (continuous) delta Bose gas on the circle [7]. It is furthermore plausible that the spectral problem for L (5.3) is in turn a singular limiting case of a transcendental spectral problem for Ruijsenaars’ analytic difference version of the

one-dimensional Schrödinger operator with Weierstrass \wp -potential [18–20], which can be captured through a very general double affine Hecke-algebraic construction due to Cherednik [3]. From this viewpoint, the representations of the double affine Hecke algebra at $q = 1$ exhibited here in rank one are expected to interpolate between the corresponding representations of the degenerate double affine Hecke algebra in [7] and the more general Hecke-algebraic settings of Ref. [3].

Acknowledgments

We thank the referees for some helpful suggestions improving the presentation.

References

- [1] L. Auslander, F.A. Grünbaum, The Fourier transform and the discrete Fourier transform, *Inverse Problems* 5 (1989) 149–164.
- [2] P. Cartier, Harmonic analysis on trees, in: C.C. Moore (Ed.), *Harmonic Analysis on Homogeneous Spaces*, American Mathematical Society, Providence, RI, 1973, pp. 419–424.
- [3] I. Cherednik, Difference-elliptic operators and root systems, *Int. Math. Res. Not. IMRN* 1995 (1995) 43–58.
- [4] I. Cherednik, *Double Affine Hecke Algebras*, Cambridge University Press, Cambridge, 2005.
- [5] J.F. van Diejen, Diagonalization of an integrable discretization of the repulsive delta Bose gas on the circle, *Comm. Math. Phys.* 267 (2006) 451–476.
- [6] J.F. van Diejen, Finite-dimensional orthogonality structures for Hall–Littlewood polynomials, *Acta Appl. Math.* 99 (2007) 301–308.
- [7] E. Emsiz, E.M. Opdam, J.V. Stokman, Periodic integrable systems with delta-potentials, *Comm. Math. Phys.* 264 (2006) 191–225.
- [8] M. Gaudin, *La fonction d'onde de Bethe*, Masson, Paris, 1983.
- [9] W. Groenevelt, Fourier transforms related to the root system of rank 1, *Transform. Groups* 12 (2007) 77–116.
- [10] R. Koekoek, P.A. Lesky, R.F. Swarttouw, *Hypergeometric Orthogonal Polynomials and Their q -Analogues*, Springer-Verlag, Berlin, 2010.
- [11] H.T. Koelink, q -Krawtchouk polynomials as spherical functions on the Hecke algebra of type B , *Trans. Amer. Math. Soc.* 352 (2000) 4789–4813.
- [12] E. Koelink, J.V. Stokman, The Askey–Wilson function transform scheme, in: J. Bustoz, M.E.H. Ismail, S.K. Suslov (Eds.), *Special Functions 2000: Current Perspective and Future Directions*, Kluwer Academic Publishers, Dordrecht, 2001, pp. 221–241.
- [13] V.E. Korepin, N.M. Bogoliubov, A.G. Izergin, *Quantum Inverse Scattering Method and Correlation Functions*, Cambridge University Press, Cambridge, 1993.
- [14] E.H. Lieb, W. Liniger, Exact analysis of an interacting Bose gas. I. The general solution and the ground state, *Phys. Rev.* (2) 130 (1963) 1605–1616.
- [15] I.G. Macdonald, *Affine Hecke Algebras and Orthogonal Polynomials*, Cambridge University Press, Cambridge, 2003.
- [16] F.I. Mautner, Spherical functions over p -adic fields, I, *Amer. J. Math.* 80 (1958) 441–457;
F.I. Mautner, Spherical functions over p -adic fields, II, *Amer. J. Math.* 86 (1964) 171–200.
- [17] M. Noumi, J.V. Stokman, Askey–Wilson polynomials: an affine Hecke algebra approach, in: R. Alvarez-Nodarse, F. Marcellán, W. Van Assche (Eds.), *Laredo Lectures on Orthogonal Polynomials and Special Functions*, Nova Science Publishers, Hauppauge, NY, 2004, pp. 111–144.
- [18] S.N.M. Ruijsenaars, Complete integrability of relativistic Calogero–Moser systems and elliptic function identities, *Comm. Math. Phys.* 110 (1987) 191–213.
- [19] S.N.M. Ruijsenaars, Generalized Lamé functions. I. The elliptic case, *J. Math. Phys.* 40 (1999) 1595–1626.
- [20] S.N.M. Ruijsenaars, Relativistic Lamé functions: completeness vs. polynomial asymptotics, *Indag. Math. (N.S.)* 14 (2003) 515–544.
- [21] A.J. Silberger, *PGL_2 over the p -Adics: Its Representations, Spherical Functions, and Fourier Analysis*, Springer-Verlag, Berlin, New York, 1970.
- [22] G. Strang, The discrete cosine transform, *SIAM Rev.* 41 (1999) 135–147.