Spanning Trees on Hypercubic Lattices and Nonorientable Surfaces

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Abstract—We consider the problem of enumerating spanning trees on lattices. Closed-form expressions are obtained for the spanning tree generating function for a hypercubic lattice in \( d \) dimensions under free, periodic, and a combination of free and periodic boundary conditions. Results are also obtained for a simple quartic net embedded on two nonorientable surfaces, a M"obius strip and the Klein bottle. Our results are based on the use of a formula expressing the spanning tree generating function in terms of the eigenvalues of an associated tree matrix. An elementary derivation of this formula is given. © 2000 Elsevier Science Ltd. All rights reserved.

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1. INTRODUCTION

The problem of enumerating spanning trees on a graph was first considered by Kirchhoff [1] in his analysis of electrical networks. Consider a graph \( G = \{V, E\} \) consisting of a vertex set \( V \) and an edge set \( E \). We shall assume that \( G \) is connected. A subset of edges \( T \subset E \) is a spanning tree if it has \(|V| - 1\) edges with at least one edge incident at each vertex. Therefore, \( T \) has no cycles.

In ensuing discussions, we shall use \( T \) to also denote the spanning tree.

Number the vertices from 1 to \(|V|\) and associate to the edge \( e_{ij} \) connecting vertices \( i \) and \( j \) a weight \( x_{ij} \), with the convention of \( x_{ii} = 0 \). The enumeration of spanning trees concerns with the evaluation of the tree generating function

\[
T(G; \{x_{ij}\}) = \sum_{T \subseteq E} \prod_{e_{ij} \in T} x_{ij},
\]

where \( T \subseteq E \) is a spanning tree of \( G \).
where the summation is taken over all spanning trees $T$. Particularly, the number of spanning
trees on $G$ is obtained by setting $x_{ij} = 1$ as

$$N_{SP T}(G) = T(G; 1). \quad (2)$$

Considerations of spanning tree also arise in statistical physics [4] in the enumeration of close-packed dimers (perfect matchings) [5]. Using a similar consideration, for example, one of us [6]
has evaluated the number of spanning trees for the simple quartic, triangular, and honeycomb
lattices in the limit of $|V| \to \infty$. In this letter, we report new results on the evaluation of the
generating function equation (1) for finite hypercubic lattices in arbitrary dimensions. Results
are also obtained for a simple quartic net embedded on two nonorientable surfaces, the Möbius
strip and the Klein bottle. As the main formula used in this letter is a relation expressing the
tree generating function in terms of the eigenvalues of an associated tree matrix, for completeness
we give an elementary derivation of this formula.

2. THE TREE MATRIX

For a given graph $G = \{V, E\}$ consider a $|V| \times |V|$ matrix $M(G)$ with elements

$$M_{ij}(G) = \begin{cases} \sum_{k=1}^{\lfloor V \rfloor} x_{ik}, & i = j = 1, 2, \ldots, |V|, \\ -x_{ij}, & \text{if vertices } i, j, i \neq j, \text{are connected by an edge}, \\ 0, & \text{otherwise}. \end{cases} \quad (3)$$

We shall refer to $M(G)$ simply as the tree matrix. It is well known [7,8] that the tree generating
function, equation (1), is given by the cofactor of any element of the tree matrix, and that the
cofactor is the same for all elements. Namely, we have the identity

$$T(G; \{x_{ij}\}) = \text{the cofactor of any element of the matrix } M(G). \quad (4)$$

The tree generating function can also be expressed in terms of the eigenvalues of the tree
matrix $M(G)$ [2, p. 39]. We give here an elementary derivation of this result which we use in
subsequent sections.

Let $M(G)$ be the tree matrix of a graph $G = \{V, E\}$. Since the sum of all elements in a row
of $M(G)$ equals to zero, $M(G)$ has 0 as an eigenvalue and, by definition, we have

$$\det |M_{ij}(G) - \lambda \delta_{ij}| = -\lambda F(\lambda), \quad (5)$$

where

$$F(\lambda) = \prod_{i=2}^{\lfloor V \rfloor} (\lambda_i - \lambda), \quad (6)$$

$\lambda_2, \lambda_3, \ldots, \lambda_{|V|}$ being the remaining eigenvalues.

Now the sum of all elements in a row of the determinant $|M_{ij}(G) - \lambda \delta_{ij}|$ is $-\lambda$. This permits
us to replace the first column of $\det |M_{ij}(G) - \lambda \delta_{ij}|$ by a column of elements $-\lambda$ without affecting
its value. Next we carry out a Laplace expansion of the resulting determinant along the modified
column, obtaining

$$\det |M_{ij}(G) - \lambda \delta_{ij}| = -\lambda \sum_{i=1}^{\lfloor V \rfloor} C_{i1}(\lambda), \quad (7)$$

where $C_{i1}(\lambda)$ is the cofactor of the $(i1)^{th}$ element of the determinant. Combining equations
(5)–(7), we are led to the identity

$$F(\lambda) = \sum_{i=1}^{\lfloor V \rfloor} C_{i1}(\lambda). \quad (8)$$
Now, \( C_{i1}(0) \) is precisely the cofactor of the \((i1)^{th}\) element of \( M(G) \) which, by equation (4), is equal to the tree generating function \( T(G; \{x_{ij}\}) \). It follows that, after setting \( \lambda = 0 \) in equation (8), we obtain the expression

\[
T(G; \{x_{ij}\}) = \frac{1}{|V|} \prod_{i=2}^{V} \lambda_i. \tag{9}
\]

This result can also be deduced by considering the tree matrix of a graph obtained from \( G \) by adding an auxiliary vertex connected to all vertices with edges of weight \( x \), followed by taking the limit of \( x \to 0 \) [9].

3. HYPERCUBIC LATTICES

We now deduce the closed-form expression for the tree generating function for a hypercubic lattice in \( d \) dimensions under various boundary conditions.

3.1. Free Boundary Conditions

**Theorem 1.** Let \( Z_d \) be a \( d \)-dimensional hypercubic lattice of size \( N_1 \times N_2 \times \cdots \times N_d \) with edge weights \( x_i \) along the \( i^{th} \) direction, \( i = 1, 2, \ldots, d \). The tree generating function for \( Z_d \) is

\[
T(Z_d; \{x_i\}) = \frac{2^{N-1}}{N} \prod_{n_1=0}^{N_1-1} \cdots \prod_{n_d=0}^{N_d-1} \left[ \sum_{i=1}^{d} x_i \left( 1 - \cos \frac{n_i \pi}{N_i} \right) \right], \tag{10}
\]

where \( N = N_1N_2\ldots N_d \).

**Proof.** The tree matrix of \( Z_d \) assumes the form of a linear combination of direct products of smaller matrices,

\[
M(Z_d) = \sum_{i=1}^{d} x_i [2I_{N_1} \otimes I_{N_2} \otimes \cdots \otimes I_{N_d} - I_{N_1} \otimes \cdots \otimes I_{N_{i-1}} \otimes H_{N_i} \otimes I_{N_{i+1}} \otimes \cdots \otimes I_{N_d}],
\]

where \( I_N \) is an \( N \times N \) identity matrix and \( H_N \) is the \( N \times N \) tri-diagonal matrix

\[
H_N = \begin{pmatrix}
1 & 0 & 0 & \cdots & 0 & 0 & 0 \\
1 & 1 & 0 & \cdots & 0 & 0 & 0 \\
0 & 1 & 1 & \cdots & 0 & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\
0 & 0 & 0 & \cdots & 1 & 0 & 1 \\
0 & 0 & 0 & \cdots & 0 & 1 & 1
\end{pmatrix}. \tag{12}
\]

It is readily verified that \( H_N \) is diagonalized by the similarity transformation

\[
S_NS_N^{-1} = \Lambda_N, \tag{13}
\]

where \( S_N \) and \( S_N^{-1} \) are \( N \times N \) matrices with elements

\[
(S_N)_{mn} = (S_N^{-1})_{nm} = \sqrt{\frac{2}{N}} \cos \left[ \left( 2n + 1 \right) \frac{m\pi}{2N} \right] + \left( \sqrt{\frac{1}{N}} - \sqrt{\frac{2}{N}} \right) \delta_{mn}, \tag{14}
\]

\( m, n = 0, 1, \ldots, N - 1. \)
and \( \Lambda_N \) is an \( N \times N \) diagonal matrix with diagonal elements
\[
\lambda_n = 2 \cos \frac{n\pi}{N}, \quad n = 0, 1, \ldots, N - 1.
\] (15)

Here \( \delta_{m,n} \) is the Kronecker delta. It follows that \( M(Z_d) \) is diagonalized by the similarity transformation
\[
S_N M(Z_d) S_N^{-1} = \Lambda_N,
\] (16)

where
\[
S_N = S_{N_1} \otimes S_{N_2} \otimes \cdots \otimes S_{N_d},
\] (17)

and \( \Lambda_N \) is an \( N \times N \) diagonal matrix with diagonal elements
\[
\lambda_{n_1, \ldots, n_d} = 2 \sum_{i=1}^{d} x_i \left[ 1 - \cos \frac{n_i \pi}{N_i} \right], \quad n_i = 0, 1, \ldots, N_i - 1.
\] (18)

Now, we have \( \lambda_{n_1, \ldots, n_d} = 0 \) for \( n_1 = n_2 = \cdots = n_d = 0 \). This establishes Theorem 1 after using equation (9).

**Remark.** The result equation (18) generalizes the \( d = 2 \) eigenvalues of \( M(Z_2) \) for \( x_i = 1 \) reported in [2, p. 74].

### 3.2. Periodic Boundary Conditions

In applications in physics, one often requires periodic boundary conditions depicted by the condition that two "boundary" vertices at coordinates \( (\ldots, n_i = 1, \ldots) \) and \( (\ldots, n_i = N_i, \ldots) \), \( i = 1, 2, \ldots, d \), are connected by an extra edge. This leads to a lattice \( Z_{d}^{\text{Per}} \) which is a regular graph with degree \( 2d \) at all vertices. For \( d = 2 \), for example, \( Z_2^{\text{Per}} \) can be regarded as being embedded on the surface of a torus.

**Theorem 2.** Let \( Z_{d}^{\text{Per}} \) be a hypercubic lattice in \( d \) dimensions of size \( N_1 \times N_2 \times \cdots \times N_d \) with edge weights \( x_i \) along the \( i \)th direction, \( i = 1, 2, \ldots, d \) with periodic boundary conditions. The tree generating function for \( Z_{d}^{\text{Per}} \) is
\[
T(Z_{d}^{\text{Per}}, \{x_i\}) = \frac{2^{N-1}}{N} \prod_{n_1=0}^{N_1-1} \cdots \prod_{n_d=0}^{N_d-1} \left[ \sum_{i=1}^{d} x_i \left( 1 - \cos \frac{2n_i \pi}{N_i} \right) \right],
\] (19)

\((n_1, \ldots, n_d) \neq (0, \ldots, 0)\).

**Proof.** The tree matrix assumes the form
\[
M(Z_{d}^{\text{Per}}) = \sum_{i=1}^{d} x_i \left[ 2I_{N_1} \otimes I_{N_2} \otimes \cdots \otimes I_{N_d} - I_{N_1} \otimes \cdots \otimes I_{N_d} \right]
\] (20)

\( \otimes I_{N_{i-1}} \otimes G_N \otimes I_{N_{i+1}} \otimes \cdots \otimes I_{N_d} \),

where \( G_N \) is the \( N \times N \) cyclic matrix
\[
G_N = \begin{pmatrix}
0 & 1 & 0 & \cdots & 0 & 0 & 1 \\
1 & 0 & 1 & \cdots & 0 & 0 & 0 \\
0 & 1 & 0 & \cdots & 0 & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\
0 & 0 & 0 & \cdots & 1 & 0 & 1 \\
1 & 0 & 0 & \cdots & 0 & 1 & 0
\end{pmatrix}
\] (21)
As in equation (16), the matrix $M(Z_{d}^{\text{Per}})$ can be diagonalized by a similarity transformation generated by

$$R_{N} = R_{N_{1}} \otimes R_{N_{2}} \otimes \cdots \otimes R_{N_{d}},$$  \hspace{1cm} (22)

where $R_{N}$ is an $N \times N$ matrix with elements

$$(R_{N})_{nm} = (R_{N}^{-1})^{*}_{nm} = N^{-1/2}e^{i2\pi mn/N},$$  \hspace{1cm} (23)

where $^*$ denotes the complex conjugate, yielding eigenvalues of $G_{N}$ as

$$\lambda_{n} = 2\cos \frac{2n\pi}{N}, \quad n = 0, 1, \ldots, N - 1.$$  \hspace{1cm} (24)

This establishes Theorem 2 after using equation (9).

### 3.3. Periodic Boundary Conditions Along $m \leq d$ Directions

**Corollary.** Let $Z_{d}^{\text{Per}(m)}$ be a hypercubic lattice in $d$ dimensions of size $N_{1} \times N_{2} \times \cdots \times N_{d}$ with periodic boundary conditions in directions $1, 2, \ldots, m, < d$ and free boundaries in the remaining $d - m$ directions. The tree generating function is

$$T(Z_{d}^{\text{Per}(m)}; \{x_{i}\}) = \frac{2^{N-1}}{N} \prod_{n_{1}=0}^{N_{1}-1} \cdots \prod_{n_{m}=0}^{N_{m}-1} \left[ \sum_{i=1}^{m} x_{i} \left( 1 - \cos \frac{2n_{i}\pi}{N_{i}} \right) \right] + \sum_{i=m+1}^{d} x_{i} \left( 1 - \cos \frac{n_{i}\pi}{N_{i}} \right), \quad (n_{1}, \ldots, n_{d}) \neq (0, \ldots, 0).$$  \hspace{1cm} (25)

### 4. THE MÖBIUS STRIP AND THE KLEIN BOTTLE

Due to the interplay with the conformal field theory [10], it is of current interest in statistical physics to study lattice systems on nonorientable surfaces [11,12]. Here, we consider two such surfaces, the Möbius strip and the Klein bottle, and obtain the respective tree generating functions.

#### 4.1. The Möbius Strip

**Theorem 3.** Let $Z_{2}^{\text{Mob}}$ be an $M \times N$ simple quartic net embedded on a Möbius strip forming a Möbius net of width $M$ and twisted in the direction $N$, with edge weights $x_{1}$ and $x_{2}$ along directions $M$ and $N$, respectively. The tree generating function for $Z_{2}^{\text{Mob}}$ is

$$T(Z_{2}^{\text{Mob}}; \{x_{1}, x_{2}\}) = \frac{2^{MN-1}}{MN} \prod_{m=0}^{M-1} \prod_{n=0}^{N-1} \left[ x_{1} \left( 1 - \cos \frac{mn\pi}{M} \right) \right] + x_{2} \left( 1 - \cos \frac{4n + 1 - (-1)^{m}}{2N} \pi \right), \quad (m, n) \neq (0, 0).$$  \hspace{1cm} (26)

**Proof.** Specifically, let the two vertices at coordinates $\{m, 1\}$ and $\{M - m, N\}$, $m = 1, 2, \ldots, M$ be connected with a lattice edge of weight $x_{2}$. Then the tree matrix assumes the form

$$M_{2}^{\text{Mob}} = 2(x_{1} + x_{2}) I_{M} \otimes I_{N} - x_{1} H_{M} \otimes I_{N} - x_{2} [I_{M} \otimes F_{N} + J_{M} \otimes K_{N}],$$  \hspace{1cm} (27)
where

\[
F_N = \begin{pmatrix}
0 & 1 & 0 & 0 & \cdots & 0 & 0 & 0 \\
1 & 0 & 1 & 0 & \cdots & 0 & 0 & 0 \\
0 & 1 & 0 & 1 & \cdots & 0 & 0 & 0 \\
\vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\
0 & 0 & 0 & 0 & \cdots & 1 & 0 & 1 \\
0 & 0 & 0 & 0 & \cdots & 0 & 0 & 0
\end{pmatrix}, \quad J_M = \begin{pmatrix}
0 & 0 & \cdots & 0 & 0 \\
0 & 0 & \cdots & 0 & 0 \\
0 & 0 & \cdots & 0 & 0 \\
1 & 0 & \cdots & 0 & 0 \\
0 & 1 & \cdots & 0 & 0
\end{pmatrix}, \quad K_N = \begin{pmatrix}
0 & 0 & \cdots & 0 & 1 \\
0 & 0 & \cdots & 0 & 0 \\
0 & 0 & \cdots & 0 & 0 \\
\vdots & \vdots & \cdots & \vdots & \vdots \\
0 & 0 & \cdots & 0 & 0 \\
1 & 0 & \cdots & 0 & 0
\end{pmatrix}.
\]

Since \(H_M\) and \(J_M\) commute, they can be simultaneously diagonalized by applying the similarity transformation equation (13). The transformed matrix \(S_N \text{M}(Z_2^{M_0})S_N^{-1}\) is block diagonal with \(N \times N\) blocks

\[
2 \left( x_1 - x_1 \cos \frac{m\pi}{M} + x_2 \right) I_N - x_2 (F_N + (-1)^m K_N), \quad m = 0, 1, \ldots, M - 1.
\]

Now, the eigenvalues of \(G_N = F_N + K_N\) and \(F_N - K_N\) are, respectively, \(2 \cos \left[\frac{2m + 1}{N}\pi\right]\) and \(2 \cos \left[\frac{2n + 1}{N}\pi\right]\), \(m = 0, 1, \ldots, N - 1\). Theorem 3 is established by combining these results with equation (9).

REMARK. For \(M = 2\) and \(x_1 = x_2 = 1\), equation (26) gives the number of spanning trees on a \(2 \times N\) Möbius ladder as

\[
N_{SPT} = \frac{1}{2N} \prod_{j=1}^{N-1} \left[ 3 - (-1)^j - 2 \cos \frac{j\pi}{N} \right] = \frac{N}{2} \left[ 2 + (2 + \sqrt{3})^N + (2 - \sqrt{3})^N \right].
\]

These two equivalent expressions have previously been given by [2, p. 218] and by Guy and Harary [3], respectively.

4.2. The Klein Bottle

The embedding of an \(M \times N\) simple quartic net on a Klein bottle is accomplished by further imposing a periodic boundary condition to \(Z_2^{M_0}\) in the \(M\) direction, namely, by connecting vertices of \(Z_2^{M_0}\) at coordinates \(\{1, n\}\) and \(\{M, n\}\), \(n = 1, 2, \ldots, N\) with an edge of weight \(x_1\). This leads to a lattice \(Z_2^{Klein}\) of the topology of a Klein bottle.

THEOREM 4. The tree generating function for \(Z_2^{Klein}\) (described in the above) is

\[
T \left( Z_2^{Klein}; \{x_1, x_2\} \right) = \frac{2^{MN-1}}{MN-N} \left[ \prod_{n=1}^{N-1} x_2 \left( 1 - \cos \frac{2n\pi}{N} \right) \right] \times \prod_{m=1}^{\lfloor M/2 \rfloor} \left[ \prod_{n=0}^{2N-1} x_1 \left( 1 - \cos \frac{2m\pi}{M} \right) + x_2 \left( 1 - \cos \frac{n\pi}{N} \right) \right]
\]

\[
\times \begin{cases} 
\prod_{n=0}^{N-1} \left[ 2x_1 - x_2 \left( 1 - \cos \frac{2n + 1}{N}\pi \right) \right], & \text{for } M \text{ even}, \\
1, & \text{for } M \text{ odd},
\end{cases}
\]

where \(\lfloor n \rfloor\) is the integral part of \(n\).
Proof. The tree matrix of $Z_2^{\text{Klein}}$ assumes the form
\[
M(Z_2^{\text{Klein}}) = 2(x_1 + x_2) I_M \otimes I_N - x_1 G_M \otimes I_N - x_2 [I_M \otimes F_N + J_M \otimes K_N].
\]
(31)
To obtain its eigenvalues, we first apply the similarity transformation generated by $R_M$ in the $M$ subspace. While this diagonalizes $G_M$ with eigenvalues $2\cos(2m\pi/M)$, $m = 0, 1, \ldots, M - 1$, it transforms the tree matrix $M(Z_2^{\text{Klein}})$ into
\[
\begin{pmatrix}
A_0 + B_0 & 0 & 0 & \cdots & 0 & 0 & 0 \\
0 & A_1 & 0 & \cdots & 0 & B_1 & 0 \\
0 & 0 & A_2 & \cdots & 0 & B_2 & 0 \\
\vdots & \vdots & \vdots & \cdots & \vdots & \vdots & \vdots \\
0 & 0 & B_{-2} & \cdots & 0 & A_{-2} & 0 \\
0 & B_{-1} & 0 & \cdots & 0 & 0 & A_{-1}
\end{pmatrix},
\]
(32)
where $A_m$ and $B_m$ are $N \times N$ matrices given by
\[
A_m = 2 \left[ x_1 + x_2 - x_1 \cos \frac{2m\pi}{M} \right] I_N - x_2 F_N,
\]
\[
B_m = -e^{2\pi im/M} x_2 K_N, \quad m = 0, 1, \ldots, M - 1.
\]
The matrix equation (32) is block diagonal with blocks $A_0 + B_0, (A_m B_m)$, $m = 1, 2, \ldots, [(M - 1)/2]$ and, for $m = \text{even}$, $A_{M/2} + B_{M/2}$. The eigenvalues of individual blocks can be deduced from those of $F_N \pm K_N$. We are led to the theorem after using equation (9).

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