# Radial Symmetry of Self-Similar Solutions for Semilinear Heat Equations

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The symmetry properties of positive solutions of the equation

$$\varDelta u + \frac{1}{2} x \cdot \nabla u + \frac{1}{p-1} u + u^p = 0 \quad \text{in } \mathbf{R}^n,$$

where  $n \ge 2$ , p > (n+2)/n, was studied. It was proved that u must be radially symmetric about the origin provided  $u(x) = o(|x|^{-2/(p-1)})$  as  $|x| \to \infty$ , and that there exist non-radial solutions u satisfying  $\limsup_{|x|\to\infty} |x|^{2/(p-1)}u(x) > 0$ . © 2000 Academic Press

### 1. INTRODUCTION

In this paper we consider the symmetry properties of positive solutions of the equation

$$\Delta u + \frac{1}{2} x \cdot \nabla u + \frac{1}{p-1} u + u^p = 0 \quad \text{in } \mathbf{R}^n,$$
(1.1)

where  $n \ge 2$  and p > 1. This equation arises in the study of (forward) selfsimilar solutions of the semilinear heat equation

$$w_t = \Delta w + w^p \qquad \text{in } \mathbf{R}^n \times (0, \infty). \tag{1.2}$$

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It is well known that if w(x, t) satisfies (1.2), then, for  $\mu > 0$  the rescaled functions

$$w_{\mu}(x, t) = \mu^{2/(p-1)} w(\mu x, \mu^2 t)$$

define a one parameter family of solutions to (1.2). A solution w is said to be self-similar, when  $w_{\mu}(x, t) = w(x, t)$  for all  $\mu > 0$ . It can be easily checked that w is a self-similar solution to (1.2) if and only if w has the form

$$w(x, t) = t^{-1/(p-1)} u(x/\sqrt{t}), \qquad (1.3)$$

where *u* satisfies the elliptic Eq. (1.1). Moreover, if *u* has spherical symmetry, that is if u = u(r), r = |x|, then *u* satisfies the ordinary differential equation

$$u'' + \left(\frac{n-1}{r} + \frac{r}{2}\right)u' + \frac{1}{p-1}u + u^p = 0, \qquad r > 0.$$
(1.4)

Such self-similar solutions are often used to describe the large time behavior of global solutions to the Cauchy problem, see, e.g., [11, 13, 3, 4, 14, 5, and 15], and to show nonuniqueness of solution to (1.2) with zero initial data in a certain functional space, see [12].

First we state the result concerning the symmetry properties of the solution of (1.1).

THEOREM 1.1. Let  $u \in C^2(\mathbb{R}^n)$  be a positive solution of (1.1) such that

$$u(x) = o(|x|^{-2/(p-1)})$$
 as  $|x| \to \infty$ . (1.5)

Then u must be radially symmetric about the origin.

The proof of Theorem 1.1 is based on the moving planes argument. This technique was developed by Serrin [18] in PDE theory, and extended and generalized by Gidas, Ni, and Nirenberg [9, 10]. We remark that with a change of variables we are still able to prove a radial symmetry result for Eq. (1.1).

Let us consider the problem

$$\begin{cases} u'' + \left(\frac{n-1}{r} + \frac{r}{2}\right)u' + \frac{1}{p-1}u + |u|^{p-1}u = 0, \quad r > 0, \\ u'(0) = 0 \quad \text{and} \quad u(0) = \alpha \in \mathbf{R}. \end{cases}$$
(1.6)

The problem (1.6) has been investigated extensively in [12, 16, 20, and 2]. We denote by  $u(r; \alpha)$  the unique solution of (1.6). We recall that  $u(r; \alpha)$  has the following properties:

(i)  $\lim_{r \to \infty} r^{2/(p-1)} u(r; \alpha) = L(\alpha)$  exists and is finite for every  $\alpha \in \mathbf{R}$  (see [12, Theorem 5]);

(ii) if  $L(\alpha) = 0$ , then there exists a constant  $A \neq 0$  such that

$$u(r; \alpha) = Ae^{-r^2/4}r^{2/(p-1)-n}\{1 + O(r^{-2})\} \quad \text{as} \quad r \to \infty$$

(see [16, Theorem 1]);

(iii) if  $p \ge (n+2)/(n-2)$ , then  $u(r; \alpha)$  is positive on  $[0, \infty)$  and  $L(\alpha) > 0$  for every  $\alpha > 0$  (see [12, Theorem 5]);

(iv) if  $(n+2)/n , then there exists a unique <math>\alpha > 0$  such that  $u(r; \alpha)$  is positive on  $[0, \infty)$  and  $L(\alpha) = 0$  (see [20, Theorem 1] and [2, Theorem 1.2 and Corollary 1.3]).

By virtue of Theorem 1.1 we obtain the following:

COROLLARY 1.1. (i) Assume that  $p \ge (n+2)/(n-2)$ . Then there exists no positive solution u of (1.1) satisfying (1.5).

(ii) Assume that (n+2)/n . Then there exists a unique positive solution <math>u(x) satisfying (1.5). Moreover, the solution u is radially symmetric about the origin.

*Remark.* The result (i) is differently proven by [3, Proposition 4.3] based on the Pohozaev identity.

Following the notations in [3] and [14], we define

$$L^{2}(K) = \left\{ u: \mathbf{R}^{n} \to \mathbf{R}; \int_{\mathbf{R}^{n}} |u|^{2} K(x) \, dx < \infty \right\} \quad \text{and}$$
$$H^{1}(K) = \left\{ u: \mathbf{R}^{n} \to \mathbf{R}; \int_{\mathbf{R}^{n}} \left( |u|^{2} + |\nabla u|^{2} \right) K(x) \, dx < \infty \right\},$$

where  $K(x) = \exp(|x|^2/4)$ . Escobedo and Kavian have shown in [3, Proposition 3.5] that if  $1 and if <math>u \in H^1(K)$  is a solution of (1.1), then  $u \in C^2(\mathbb{R}^n)$  and satisfies  $u(x) = O(\exp(-|x|^2/8))$  as  $|x| \to \infty$ . As a consequence of Corollary 1.1, we obtain the following:

COROLLARY 1.2. Assume that (n+2)/n . Then the problem

$$\begin{cases} \Delta u + \frac{1}{2} x \cdot \nabla u + \frac{1}{p-1} u + u^p = 0 & \text{in } \mathbf{R}^n, \\ u \in H^1(K) & \text{and} \quad u > 0 & \text{in } \mathbf{R}^n, \end{cases}$$
(1.7)

has a unique solution.

Let us consider the Cauchy problem

$$\begin{cases} w_t = \Delta w + w^p & \text{in } \mathbf{R}^n \times (0, \infty), \\ w(x, 0) = \tau w_0 & \text{in } \mathbf{R}^n, \end{cases}$$
(1.8)

where  $w_0 \in L^2(K) \cap L^{\infty}(\mathbb{R}^n)$ ,  $w_0 \ge 0$ , and  $\tau > 0$  is a parameter. We denote by  $w(x, t; \tau)$  the unique solutions of (1.8) (see [15]). Combining the result by Kawanago [15, Theorem 1] and Corollary 1.2, we obtain the following, where the asymptotic behavior of  $w(\cdot, t; \tau)$  as  $t \to \infty$  becomes clearer.

COROLLARY 1.3. Assume that  $(n+2)/n . Then there exists a unique <math>\tau_0 > 0$  such that the solution  $w(x, t; \tau)$  is a global solution if  $\tau \in (0, \tau_0]$ , and  $w(x, t; \tau)$  blows up in finite time if  $\tau \in (\tau_0, \infty)$ . Moreover,  $w(x, t; \tau_0)$  satisfies

$$\lim_{t \to \infty} \|t^{1/(p-1)} w(\cdot, t; \tau_0) - u_0(\cdot/\sqrt{t})\|_{L^{\infty}(\mathbf{R}^n)} = 0,$$

where  $u_0$  is a unique solution of the problem (1.7).

Next we consider the existence of nonradial solutions of (1.1). Let p > (n+2)/n and let U(r) be a positive solution of (1.4) satisfying

$$U'(0) = 0$$
 and  $\lim_{r \to \infty} r^{2/(p-1)} U(r) > 0.$  (1.9)

The existence of such U is obtained by [12, Theorem 5]. Define  $\ell = \ell(U) > 0$  as

$$\ell = \lim_{r \to \infty} r^{2/(p-1)} U(r).$$
(1.10)

We investigate the Cauchy problem for Eq. (1.2) with

$$w(x, 0) = w_0 \in L^1_{\text{loc}}(\mathbf{R}^n), \tag{1.11}$$

where

$$0 \le w_0(x) \le \ell |x|^{-2/(p-1)}, \qquad w_0 \ne 0, \qquad x \in \mathbf{R}^n \setminus \{0\}.$$
(1.12)

Relation (1.11) is taken in the sense of  $L^1_{loc}(\mathbf{R}^n)$ , that is,

$$\int_{K} |w(x, t) - w_0(x)| \, dx \to 0 \qquad \text{as} \quad t \to 0$$

for any compact subset K of  $\mathbb{R}^n$ . We note that  $w_0 \in L^1_{loc}(\mathbb{R}^n)$  if (1.12) holds with p > (n+2)/n.

THEOREM 1.2. Let p > (n+2)/n. Assume that (1.12) holds, where  $\ell$  is the constant in (1.10). Then there exists a positive solution  $w \in C^{2, 1}(\mathbb{R}^n \times (0, \infty))$  of (1.2) and (1.11). Assume, furthermore, that  $w_0 \in C(\mathbb{R}^n \setminus \{0\})$ , then w satisfies

$$w(x, t) \rightarrow w_0(x)$$
 as  $t \rightarrow 0$  uniformly in  $|x| \ge r$  for every  $r > 0$ . (1.13)

Moreover, w is self-similar if  $\mu^{2/(p-1)}w_0(\mu x) = w_0(x)$  for every  $\mu > 0$ .

COROLLARY 1.4. Let p > (n+2)/n. Assume that  $A: S^{n-1} \to \mathbf{R}$  is continuous and satisfies

$$0 \leqslant A(\sigma) \leqslant \ell, \qquad A \not\equiv 0, \qquad \sigma \in S^{n-1}. \tag{1.14}$$

Then there exists a positive self-similar solution  $w \in C^{2, 1}(\mathbb{R}^n \times (0, \infty))$  of (1.2) satisfying (1.11) and (1.13) with  $w_0(x) = A(x/|x|) |x|^{-2/(p-1)}$ .

Recall that self-similar solutions w to (1.2) have the form (1.3) with u satisfying (1.1). Therefore,  $w(\sigma, t) = r^{2/(p-1)}u(r\sigma)$  for  $\sigma \in S^{n-1}$ , where  $r = 1/\sqrt{t}$ . Then we obtain the following corollary, which shows that condition (1.5) in Theorem 1.1 is crucial.

COROLLARY 1.5. Let p > (n+2)/n. Assume that  $A: S^{n-1} \to \mathbf{R}$  is continuous and satisfies (1.14). Then there exists a positive non-radial solution u of (1.1) satisfying

$$r^{2/(p-1)}u(r\sigma) \to A(\sigma)$$
 as  $r \to \infty$  uniformly in  $\sigma \in S^{n-1}$ .

*Remark.* (i) If 1 , no time global, non-negative, and nontrivial solution exists in (1.2) (see, e.g., [7, 19, and 14]). Therefore, (1.1) admits a positive solution only if <math>p > (n+2)/n.

(ii) We find that the solution w of (1.2) and (1.11) obtained in Theorem 1.2 is a minimal solution of the integral equation

$$w(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y; t) w_0(y) \, dy + \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y; t - s) [w(y, s)]^p \, dy \, ds,$$

where  $\Gamma(x; t) = (4\pi t)^{-n/2} e^{-|x|^2/4t}$ . See the proof of Theorem 1.2 below.

(iii) Galaktionov and Vazquez [8] studied the Cauchy problem (1.2) and (1.11) with singular initial values for the case p > n/(n-2).

This paper is organized as follows: in Section 2, we treat the symmetry properties of the solutions and prove Theorem 1.1. In Section 3, we state some propositions concerning the properties of solutions to the Cauchy

problem with singular initial data, and prove Theorem 1.2. We prove these propositions in Section 4.

### 2. RADIAL SYMMETRY

In this section we investigate more general equation

$$\Delta u + \frac{1}{2}x \cdot \nabla u + ku + f(u) = 0 \quad \text{in } \mathbf{R}^n, \tag{2.1}$$

where  $n \ge 2$ , k is a nonnegative constant, and  $f \in C^1[0, \infty)$ . Theorem 1.1 is a consequence of the following result.

THEOREM 2.1. Assume that

$$f(s) = O(s^{\sigma})$$
 as  $s \to 0$  for some  $\sigma > 1$ . (2.2)

Let u be a positive solution of (2.1) such that

$$u(x) = o(|x|^{-2k}) \qquad as \quad |x| \to \infty.$$
(2.3)

Then u must be radially symmetric about the origin.

We obtain Theorem 2.1 by the following two propositions.

PROPOSITION 2.1. Assume that (2.2) holds. Let u be a positive solution satisfying (2.3). Then, for every m > 0,  $u(x) = o(|x|^{-m})$  as  $|x| \to \infty$ .

**PROPOSITION 2.2.** Let u be a positive solution such that

$$u(x) = o(|x|^{-2\alpha}) \quad as \ |x| \to \infty \quad for \ some \ \alpha > 0. \tag{2.4}$$

Assume that

$$\alpha > k + \max\{ |f'(s)| : 0 \le s \le ||u||_{L^{\infty}(\mathbf{R}^{n})} \}.$$
(2.5)

Then u must be radially symmetric.

To prove Proposition 2.1 we prepare the following lemma.

LEMMA 2.1. Assume that (2.2) holds. Let u be a positive solution of (2.1) such that

$$u(x) = o(|x|^{-\beta})$$
 as  $|x| \to \infty$  for some  $\beta > 2k$ .

Then, for every m > 0,  $u(x) = o(|x|^{-m})$  as  $|x| \to \infty$ .

*Proof.* Set  $v(x) = |x|^{\beta} u(x)$ . Then  $v(x) \to 0$  as  $|x| \to \infty$  and v satisfies

$$\Delta v + \frac{1}{2} x \cdot \nabla v - \frac{2\beta}{|x|^2} x \cdot \nabla v - \left(\frac{\beta}{2} - k\right) v + \frac{\beta(\beta + 2 - n)}{|x|^2} v + |x|^{\beta} f(|x|^{-\beta} v) = 0.$$

Defining

$$Lv \equiv \varDelta v + \frac{1}{2} x \cdot \nabla v - \frac{2\beta}{|x|^2} x \cdot \nabla v,$$

we have

$$Lv = \left[ \left( \frac{\beta}{2} - k \right) - \frac{\beta(\beta + 2 - n)}{|x|^2} - \frac{f(u)}{u} \right] v.$$

Note that  $\beta > k/2$  and  $f(s)/s \to 0$  as  $s \to 0$  by (2.2). Then there exists a  $R_1 > 0$  such that  $Lv \ge 0$  for  $|x| \ge R_1$ .

Fix m > 0 and define  $w(r) = r^{-m}$ , r = |x|. Then it follows that

$$Lw = w_{rr} + \frac{n-1-2\beta}{r} w_r + \frac{1}{2} rw_r = r^{-1} \left[ -\frac{m}{2} + \frac{m(m+1) - m(n-1-2\beta)}{r^2} \right].$$

Then there exists a  $R_2 > 0$  such that  $Lw \le 0$  for  $|x| \ge R_2$ . Let  $R_0 = \max\{R_1, R_2\}$ . Take C > 0 so large that  $Cw - v \ge 0$  on  $|x| = R_0$ . Then Cw - v satisfies

$$L(Cw-v) \leq 0$$
 for  $|x| > R_0$  and  $Cw-v \to 0$  as  $|x| \to \infty$ .

By the maximum principle, we obtain  $Cw \ge v$  for  $|x| \ge R_0$ , i.e.,  $u(x) \le C |x|^{-m-\beta}$  for  $|x| \ge R_0$ . Since m > 0 is arbitrary, we obtain the conclusion.

**Proof** of Proposition 2.1. Set  $v(x) = |x|^{2k} u(x)$ . Then  $v(x) \to 0$  as  $|x| \to \infty$  and v satisfies

$$\Delta v + \frac{1}{2}x \cdot \nabla v - \frac{4k}{|x|^2}x \cdot \nabla v + \frac{2k(2k+2-n)}{|x|^2}v + |x|^{2k}f(|x|^{-2k}v) = 0.$$

Defining

$$Lv \equiv \varDelta v + \frac{1}{2} x \cdot \nabla v - \frac{4k}{|x|^2} x \cdot \nabla v,$$

we have

$$Lv = -\frac{2k(2k+2-n)}{|x|^2}v - |x|^{2k}f(u).$$

We note here that

$$|x|^{2k} f(u) = \frac{f(u)}{u^{\sigma}} \frac{v^{\sigma}}{|x|^{2k(\sigma-1)}}.$$

Then, by (2.2), we obtain

$$Lv = o(|x|^{-\delta_0})$$
 as  $|x| \to \infty$ ,

where  $\delta_0 = \min\{2, 2k(\sigma - 1)\}$ . Let  $w(x) = |x|^{-\delta}$ , where  $\delta = \delta_0/2$ . Then we have

$$Lw = |x|^{-\delta} \left[ -\frac{\delta}{2} + \frac{\delta(\delta+1) - \delta(n-1-2k)}{|x|^2} \right] \leqslant -\frac{\delta}{4} |x|^{-\delta}$$

for |x| sufficiently large. Then there exists a  $R_0 > 0$  such that  $Lw - Lv \le 0$ for  $|x| \ge R_0$ . Take  $C \ge 1$  so large that  $Cw \ge v$  on  $|x| = R_0$ . Then Cw - vsatisfies

$$L(Cw-v) \leq 0$$
 for  $|x| > R_0$  and  $Cw-v \to 0$  as  $|x| \to \infty$ .

By the maximum principle, we obtain  $Cw \ge v$  for  $|x| \ge R_0$ , i.e.,  $v(x) \le C |x|^{-\delta}$  for  $|x| \ge R_0$ . This implies that  $u(x) = o(|x|^{-\beta})$  as  $|x| \to \infty$  for  $\beta = 2k + \delta > 2k$ . By Lemma 2.1, we conclude that  $u(x) = o(|x|^{-m})$  as  $|x| \to \infty$  for every m > 0.

Next we prove Proposition 2.2. Let u be a positive solution of (2.1) satisfying (2.4). Define

$$w(x, t) = t^{-\alpha} u(x/\sqrt{t}), \qquad (x, t) \in \mathbf{R}^n \times (0, \infty).$$
(2.6)

LEMMA 2.2. (i) For every T > 0,  $w(x, t) \rightarrow 0$  as  $|x| \rightarrow \infty$  uniformly in  $t \in (0, T]$ ;

(ii) For every  $\lambda > 0$ ,  $w(x, t) \to 0$  as  $t \to 0$  uniformly in  $|x| \ge \lambda$ .

*Proof.* From (2.4), for any  $\varepsilon > 0$ , there exists a R > 0 such that

$$|x|^{2\alpha} u(x) < \varepsilon, \qquad |x| \ge R. \tag{2.7}$$

By virtue of (2.6) we have

$$|x|^{2\alpha} w(x,t) = |x/\sqrt{t}|^{2\alpha} u(x/\sqrt{t}).$$
(2.8)

(i) Fix T > 0. From (2.7) and (2.8) it follows that

$$|x|^{2\alpha} w(x, t) < \varepsilon$$
 for  $|x| \ge RT^{1/2}$ ,  $t \in (0, T]$ .

Since  $\varepsilon > 0$  is arbitrary, we have  $w(x, t) \to 0$  as  $|x| \to \infty$  uniformly in  $t \in (0, T]$ .

(ii) From (2.7) and (2.8) we obtain

$$\lambda^{2\alpha} w(x, t) \leq |x|^{2\alpha} w(x, t) < \varepsilon, \qquad \text{for} \quad |x| \geq \lambda, \quad 0 < t < (\lambda/R)^2.$$

Then  $w(x, t) \to 0$  as  $t \to 0$  uniformly in  $|x| \ge \lambda$ .

For  $\lambda \in \mathbf{R}$ , we define  $T_{\lambda}$  and  $\Omega_{\lambda}$  as

$$T_{\lambda} = \left\{ x = (x_1, ..., x_n) \in \mathbf{R}^n : x_1 = \lambda \right\} \qquad \text{and} \qquad \mathcal{Q}_{\lambda} = \left\{ x \in \mathbf{R}^n : x_1 < \lambda \right\},$$

respectively. For  $x = (x_1, ..., x_n) \in \mathbf{R}^n$  and  $\lambda \in \mathbf{R}$ , let  $x^{\lambda}$  be the reflection of x with respect to  $T_{\lambda}$ , i.e.,  $x^{\lambda} = (2\lambda - x_1, x_2, ..., x_n)$ . It is easy to see that, if  $\lambda > 0$ ,

 $|x^{\lambda}| > |x| \quad \text{for } x \in \Omega_{\lambda} \qquad \text{and} \qquad \left\{ x^{\lambda} : x \in \Omega_{\lambda} \right\} = \left\{ x : x_1 > \lambda \right\} \subset \left\{ x : |x| \ge \lambda \right\}.$ 

By Lemma 2.2 we have the following:

LEMMA 2.3. Let  $\lambda > 0$ .

(i) For every T > 0,  $w(x^{\lambda}, t) \to 0$  as  $|x| \to \infty$ ,  $x \in \Omega_{\lambda}$ , uniformly in  $t \in (0, T]$ ;

(ii)  $w(x^{\lambda}, t) \to 0 \text{ as } t \to 0 \text{ uniformly in } x \in \Omega_{\lambda}.$ 

LEMMA 2.4. Define  $\phi$  as  $\phi(x, t) = w(x, t) - w(x^{\lambda}, t)$ . Then we have

 $\phi_t = \varDelta \phi + c(x, t) \phi \quad in \ \Omega_\lambda \times (0, \infty) \qquad and \qquad \phi = 0 \quad on \ T_\lambda \times (0, \infty),$ 

where

$$c(x, t) = \left[ -(\alpha - k) + \int_0^1 f'(u^{\lambda} - s(u^{\lambda} - u)) \, ds \right] t^{-1}$$
(2.9)

with  $u = u(x/\sqrt{t})$  and  $u^{\lambda} = u(x^{\lambda}/\sqrt{t})$ .

*Proof.* From (2.6) we have

$$w_t = \Delta w - (\alpha - k) t^{-1} w + t^{-\alpha - 1} f(t^{\alpha} w), \qquad (x, t) \in \mathbf{R}^n \times (0, \infty).$$

Let  $w^{\lambda}(x, t) = w(x^{\lambda}, t)$ . Then  $w^{\lambda}$  satisfies

$$w_t^{\lambda} = \varDelta w^{\lambda} - (\alpha - k) t^{-1} w^{\lambda} + t^{-\alpha - 1} f(t^{\alpha} w^{\lambda}), \qquad (x, t) \in \Omega_{\lambda} \times (0, \infty).$$

Then we have  $\phi_t = \Delta \phi + c \phi$ , where c is the function in (2.9).

LEMMA 2.5. For 
$$\lambda > 0$$
, we have  $w(x, t) \ge w(x^{\lambda}, t)$  for  $(x, t) \in \Omega_{\lambda} \times (0, \infty)$ .

*Proof.* Let  $\phi(x, t) = w(x, t) - w(x^{\lambda}, t)$ . We show that  $\phi(x, t) \ge 0$  for  $(x, t) \in \Omega_{\lambda} \times (0, \infty)$ . Assume to the contrary that there exists a  $(x_0, t_0) \in \Omega_{\lambda} \times (0, \infty)$  such that  $\phi(x_0, t_0) < 0$ . Take  $\varepsilon > 0$  so small that  $\phi(x_0, t_0) < -\varepsilon < 0$ . By Lemma 2.3(ii), we can take  $T_0 \in (0, t_0)$  so that  $w(x^{\lambda}, T_0) < \varepsilon$  for  $x \in \Omega_{\lambda}$ . Then it follows that

$$\phi(x, T_0) > -\varepsilon, \qquad x \in \Omega_{\lambda}. \tag{2.10}$$

Fix  $T > t_0$ . By Lemma 2.3(i), we can take  $R > |x_0|$  so large that  $w(x^{\lambda}, t) < \varepsilon$  for  $|x| \ge R$ ,  $x \in \Omega_{\lambda}$ ,  $t \in [T_0, T]$ . Then we obtain

$$\phi(x,t) > -\varepsilon, \qquad |x| \ge R, \qquad x \in \Omega_{\lambda}, \qquad t \in [T_0,T]. \tag{2.11}$$

Define  $Q = \{x \in \Omega_{\lambda} : |x| < R\}$ . Let  $\Sigma$  be a parabolic boundary of  $Q \times (T_0, T)$ , i.e.,

$$\Sigma = (Q \times \{T_0\}) \cup (\partial Q \times (T_0, T)).$$

From Lemma 2.4, (2.10), and (2.11) we have

$$\phi_t = \Delta \phi + c \phi$$
 in  $Q \times (T_0, T)$  and  $\phi \ge -\varepsilon$  on  $\Sigma$ .

Note that  $c(x, t) \leq 0$  from (2.5) and (2.9). Define  $\psi$  as  $\psi(x, t) = \phi(x, t) + \varepsilon$ . Then  $\psi$  satisfies

$$\psi_t \ge \Delta \psi + c \psi$$
 in  $Q \times (T_0, T)$  and  $\psi \ge 0$  on  $\Sigma$ .

By the maximum principle ([17, Chapter 3, Theorem 7]), we have  $\psi \ge 0$  on  $\overline{Q} \times [T_0, T]$ , which implies that

$$\phi(x,t) \ge -\varepsilon, \qquad (x,t) \in \overline{Q} \times [T_0,T]. \tag{2.12}$$

On the other hand,  $(x_0, t_0) \in Q \times (T_0, T)$  and  $\phi(x_0, t_0) < -\varepsilon$ . This contradicts to (2.12). Hence  $\phi(x, t) \ge 0$  for  $(x, t) \in \Omega_\lambda \times (0, \infty)$ .

*Proof of Proposition* 2.2. By Lemma 2.5, we have  $w(x, t) \ge w(x^{\lambda}, t)$  for  $\lambda > 0$ ,  $(x, t) \in \Omega_{\lambda} \times (0, \infty)$ . From the continuity of w, we have  $w(x, t) \ge w(x^0, t)$  for  $(x, t) \in \Omega_0 \times (0, \infty)$ .

We can repeat the previous arguments for the negative  $x_1$ -direction to conclude that  $w(x, t) \leq w(x^0, t)$  for  $(x, t) \in \Omega_0 \times (0, \infty)$ . Hence, w(x, t) must be symmetric about the plane  $x_1 = 0$ . Since the Eq.(2.1) is invariant under the rotation and conclude that w(x, t) is symmetric in every direction. Therefore, u must be radially symmetric about the origin.

### 3. PROOF OF THEOREM 1.2

We assume, henceforth, that p > (n+2)/n. Recall that U(r) is a positive solution of (1.4) satisfying (1.9), and that  $\ell$  is a positive constant defined by (1.10). Let W = W(x, t) be a self-similar solution of (1.2) of the form

$$W(x, t) = t^{-1/(p-1)} U(|x|/\sqrt{t}).$$
(3.1)

To prove Theorem 1.2, we consider the Cauchy problem

$$\begin{cases} w_t = \Delta w + g & \text{in } \mathbf{R}^n \times (0, \infty), \\ w(x, 0) = w_0 \in L^1_{\text{loc}}(\mathbf{R}^n). \end{cases}$$
(3.2)

The initial condition in (3.2) is taken in the sense of  $L^1_{loc}(\mathbf{R}^n)$ . We always assume that  $w_0$  satisfies (1.12), and that g satisfies

$$0 \leq g(x, t) \leq [W(x, t)]^{p} \quad \text{in } \mathbf{R}^{n} \times (0, \infty).$$
(3.3)

We consider the function w satisfying the growth condition:

$$\begin{cases} \text{there exists positive constants } C, \alpha, \text{ and } r \text{ such that} \\ |w(x, t)| \leq Ce^{\alpha |x|^2} \text{ for all } |x| \geq r \text{ and all } t \in [0, \infty). \end{cases}$$
(3.4)

We obtain the following propositions. Proofs will be given in the next section.

**PROPOSITION 3.1.** The function

$$\phi(x, t) = \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y; t - s) g(y, s) \, dy \, ds$$

is well defined on  $\mathbb{R}^n \times (0, \infty)$ . Assume, furthermore, that g is continuous on  $\mathbb{R}^n \times (0, \infty)$ , then  $\phi \in C^{1,0}(\mathbb{R}^n \times (0, \infty))$ .

**PROPOSITION 3.2.** Assume that

 $\begin{cases} g \in C(\mathbf{R}^n \times (0, \infty)) \text{ and } x \mapsto g(x, t) \text{ is locally Hölder continuous} \\ uniformly \text{ in } t \in [t_0, T] \text{ for every } t_0 \text{ and } T \text{ with } 0 < t_0 < T < \infty. \end{cases}$ (3.5)

Then the function  $w \in C^{2,1}(\mathbb{R}^n \times (0, \infty))$  is a solution of the problem (3.2) satisfying the growth condition (3.4) if and only if

$$w(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y; t) w_0(y) \, dy + \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y; t - s) g(y, s) \, dy \, ds.$$
(3.6)

Assume, furthermore, that  $w_0 \in C(\mathbb{R}^n \setminus \{0\})$ , then w defined by (3.6) satisfies

$$w(x, t) \rightarrow w_0(x)$$
 as  $t \rightarrow 0$  uniformly in  $|x| \ge r$  for every  $r > 0$ . (3.7)

We show that W defined by (3.1) satisfies

$$W(x, 0) = \ell |x|^{-2/(p-1)}$$
 in the sense of  $L^{1}_{loc}(\mathbf{R}^{n})$ . (3.8)

In fact, we have  $|x|^{2/(p-1)} W(x, t) = |x/\sqrt{t}|^{2/(p-1)} U(|x|/\sqrt{t})$ . By virtue of (1.10), we obtain  $0 \le |x|^{2/(p-1)} W(x, t) \le C$  with a constant *C* and

$$|x|^{2/(p-1)} W(x, t) \to \ell$$
 as  $t \to 0$ ,  $|x| \neq 0$ .

Fix a compact set  $K \subset \mathbf{R}^n$ . By the Lebesgue dominated convergence theorem, we have

$$\int_{K} |W(x, t) - \ell |x|^{-2/(p-1)} |dx$$
  
=  $\int_{K} |x|^{-2/(p-1)} ||x|^{2/(p-1)} W(x, t) - \ell |dx \to 0$  as  $t \to 0$ .

Hence, (3.8) holds. A quick check implies that W is a solution to the problem (3.2) with  $g = W^p$  and  $w_0(x) = \ell |x|^{-2/(p-1)}$ . By Propositions 3.1 and 3.2, we have

$$\int_0^t \int_{\mathbf{R}^n} \Gamma(x-y:t-s) [W(y,s)]^p \, dy \, ds < \infty$$

and

$$W(x, t) = \int_{\mathbf{R}^{n}} \Gamma(x - y : t) [\ell |y|^{-2/(p-1)}] dy$$
  
+  $\int_{0}^{t} \int_{\mathbf{R}^{n}} \Gamma(x - y : t - s) [W(y, s)]^{p} dy ds$  (3.9)

for  $(x, t) \in \mathbf{R}^n \times (0, \infty)$ .

*Proof of Theorem* 1.2. Define  $w_i$ , i = 1, 2, ..., inductively by

$$\begin{cases} w_1(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y : t) w_0(y) \, dy, \\ w_{i+1}(x, t) = w_1(x, t) + \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y : t - s) [w_i(y, s)]^p \, dy \, ds, \quad i = 1, 2, \dots. \end{cases}$$
(3.10)

By virtue of (1.12) and (3.9) we have

$$0 \leq w_1(x, t) \leq \int_{\mathbf{R}^n} \Gamma(x - y : t) [\ell | y|^{-2/(p-1)}] dy$$
$$\leq W(x, t) \quad \text{for} \quad (x, t) \in \mathbf{R}^n \times (0, \infty).$$

Then, by induction,  $w_i$  is well defined and satisfies

$$0 \leqslant w_1(x, t) \leqslant \cdots \leqslant w_i(x, t) \leqslant w_{i+1}(x, t) \leqslant \cdots \leqslant W(x, t) \quad \text{in } \mathbf{R}^n \times (0, \infty).$$

Define  $w(x, t) = \lim_{i \to \infty} w_i(x, t)$ . Letting  $i \to \infty$  in (3.10), by the Lebesgue dominated convergence theorem, we obtain

$$w(x,t) = \int_{\mathbf{R}^n} \Gamma(x-y:t) w_0(y) \, dy + \int_0^t \int_{\mathbf{R}^n} \Gamma(x-y:t-s) [w(y,s)]^p \, dy \, ds.$$
(3.11)

Observe that w is continuous and satisfies  $w \le W$ . Then  $g = w^p$  is continuous and satisfies (3.3). By Proposition 3.1, we have  $w \in C^{1,0}(\mathbb{R}^n \times (0, \infty))$ , and then  $g = w^p$  satisfies (3.5). Therefore, by Proposition 3.2,  $w \in C^{2,1}(\mathbb{R}^n \times (0, \infty))$ and w is a solution of (3.2) with  $g = w^p$ , that is, a solution of (1.1) and (1.11). Moreover, if  $w_0 \in C(\mathbb{R}^n \setminus \{0\})$ , then w satisfies (1.13). We easily see that w is a minimal solution of the integral equation (3.11).

Assume that  $\mu^{2/(p-1)}w_0(\mu x) = w_0(x)$  for all  $\mu > 0$ . Then we have  $\mu^{2/(p-1)}w(\mu x, \mu^2 t) = w(x, t)$  for all  $\mu > 0$  by the uniqueness of the minimal solutions. This implies that w is a self-similar solution.

### 4. PROOF OF PROPOSITIONS 3.1 AND 3.2

#### **4.1.** Define $\psi$ by

$$\psi(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y : t) w_0(y) \, dy,$$

where  $w_0 \in L^1_{loc}(\mathbf{R}^n)$  satisfies (1.12). Then  $\psi \in C^{2,1}(\mathbf{R}^n \times (0, \infty))$  and satisfies

 $\psi_t = \Delta \psi$  in  $\mathbf{R}^n \times (0, \infty)$  and  $\psi(\cdot, t) \to w_0(\cdot)$  in  $L^1_{\text{loc}}(\mathbf{R}^n)$  as  $t \to 0$ . (4.1)

(See, e.g., [1, Chapter 5, Theorem 6.1].) We obtain the following:

**PROPOSITION 4.1.** For every r > 0, there exists a constant C = C(r) > 0 such that

$$0 \leq \psi(x, t) \leq C$$
 for  $|x| \geq r$  and  $0 < t < \infty$ . (4.2)

Assume, furthermore, that  $w_0 \in C(\mathbb{R}^n \setminus \{0\})$ , then  $\psi$  satisfies

$$\psi(x, t) \to w_0(x)$$
 as  $t \to 0$  uniformly in  $|x| \ge r$  for every  $r > 0$ . (4.3)

*Proof.* Let r > 0 be arbitrary but fixed. We note that, by (1.12), there exists a constant  $C_1 > 0$  such that

$$0 \leqslant w_0(x) \leqslant C_1 \qquad \text{for} \quad |x| \ge r. \tag{4.4}$$

First, we show that there exists a constant C > 0 satisfying

$$0 \leq \psi(x, t) \leq C$$
 for  $|x| \geq 2r$  and  $0 < t < \infty$ . (4.5)

In fact, we write

$$\psi(x, t) = \int_{|y| \leq r} + \int_{|y| \geq r} \Gamma(x - y : t) w_0(y) dy \equiv I_1 + I_2.$$

Observe that  $|x - y| \ge r$  for  $|x| \ge 2r$  and  $|y| \le r$ . Then, since  $w_0 \in L^1_{loc}(\mathbb{R}^n)$ ,

$$I_1 \leq \Gamma_r \int_{|y| \leq r} w_0(y) \, dy < \infty,$$

where  $\Gamma_r = \sup \{ \Gamma(x - y : t) : |x - y| \ge r \} = \Gamma(r : r^2/2n)$ . From (4.4) we have

$$I_2 \leqslant C_1 \int_{|y| \geqslant r} \Gamma(x - y : t) \, dy \leqslant C_1.$$

Therefore we obtain (4.5). Since r > 0 is arbitrary, (4.2) holds.

Next we show (4.3). Since  $\int_{\mathbf{R}^n} \Gamma(x - y : t) dy = 1$ , we have

$$\psi(x, t) - w_0(x) = \int_{\mathbf{R}^n} \Gamma(x - y; t) [w_0(y) - w_0(x)] \, dy$$

For  $|x| \ge r$ , we write

$$\begin{split} |\psi(x, t) - w_0(x)| \\ \leqslant & \int_{\substack{|x-y| \ge \rho \\ |y| \ge r}} + \int_{\substack{|x-y| \ge \rho \\ |y| \le r}} + \int_{\substack{|x-y| < \rho \\ |y| < r}} \Gamma(x-y:t) |w_0(y) - w_0(x)| \, dy \\ \equiv & I_1 + I_2 + I_3, \end{split}$$

where  $\rho \in (0, r)$  is arbitrary but be fixed. From (4.4) we have

$$I_1 \leq 2C_1 \int_{|x-y| \ge \rho} \Gamma(x-y:t) \, dy \to 0$$
 as  $t \to 0$ .

Since  $w_0 \in L^1_{loc}(\mathbf{R}^n)$ , we have

$$I_2 \leq \int_{|y| \leq r} (w_0(y) + C_1) dy \sup_{|x-y| \geq \rho} \Gamma(x-y:t) \to 0$$
 as  $t \to 0$ .

We estimate  $I_3$  as

$$I_{3} \leq \sup_{|x-y| < \rho} |w_{0}(y) - w_{0}(x)| \int_{\mathbf{R}^{n}} \Gamma(x-y:t) \, dy = \sup_{|x-y| < \rho} |w_{0}(y) - w_{0}(x)|.$$

Therefore, for arbitrary  $\rho \in (0, r)$ ,

$$\lim_{t \to 0} (\sup_{|x| \ge r} |\psi(x, t) - w_0(x)|) \le \sup_{\substack{|x| \ge r \\ |x - y| < \rho}} |w_0(y) - w_0(x)|.$$

We see that  $w_0(x)$  is uniformly continuous in  $|x| \ge r$  since  $w_0(x) \to 0$  as  $|x| \to \infty$ . Hence, we obtain (4.3).

## 4.2. In this subsection we show the following:

**PROPOSITION 4.2.** Assume that g satisfies (3.3). Then the function  $\phi$  defined by

$$\phi(x, t) = \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y : t - s) g(y, s) \, dy \, ds$$

is well defined on  $\mathbf{R}^n \times (0, \infty)$  and satisfies the following properties:

- (i)  $\phi(x, t) \rightarrow 0$  as  $t \rightarrow 0$  uniformly in  $|x| \ge r$  for every r > 0;
- (ii) For every r > 0, there exists a constant C > 0 such that

$$0 \leq \phi(x, t) \leq C$$
 for  $|x| \geq r$  and  $0 < t < \infty$ ;

- (iii)  $\phi(\cdot, t) \to 0$  in  $L^1_{\text{loc}}(\mathbf{R}^n)$  as  $t \to 0$ ;
- (iv) Assume that  $g \in C(\mathbb{R}^n \times (0, \infty))$ . Then  $\phi \in C^{1,0}(\mathbb{R}^n \times (0, \infty))$ ;
- (v) Assume that (3.5) holds. Then  $\phi \in C^{2,1}(\mathbb{R}^n \times (0,\infty))$  and satisfies

$$\phi_t = \Delta \phi + g$$
 in  $\mathbf{R}^n \times (0, \infty)$ .

To prove Proposition 4.2, define J and  $\Phi$  as

$$\begin{cases} J(x, t:s) = \int_{\mathbf{R}^n} \Gamma(x - y: t - s) [W(y, s)]^p \, dy \\ \Phi(x, t) = \int_0^t J(x, t:s) \, ds. \end{cases}$$

We note that  $\phi(x, t) \leq \Phi(x, t)$  by (3.3). We recall that  $[W(x, t)]^p = t^{-p/(p-1)} [U(x/\sqrt{t})]^p$ . Since  $\mu^{2/(p-1)} W(\mu x, \mu^2 t) = W(x, t)$  for all  $\mu > 0$ , we have

$$\mu^{2/(p-1)} \Phi(\mu x, \mu^2 t) = \Phi(x, t) \quad \text{for all} \quad \mu > 0$$
(4.6)

by direct calculation. Since U(r) is bounded on  $[0, \infty)$ , there exists a constant  $C_2 > 0$  such that, for  $(x, t) \in \mathbf{R}^n \times (0, \infty)$ ,

$$J(x, t:s) \leqslant C_2 s^{-p/(p-1)}, \qquad 0 < s < t.$$
(4.7)

We obtain the following estimates of J(x, t:s).

LEMMA 4.1. For every  $t_0 > 0$ , there exist constants  $\sigma \in [0, 1)$  and  $C_3 = C_3(t_0) > 0$  such that

$$J(x,t:s) \leq C_3 s^{-\sigma} \qquad for \ all \quad (x,t,s) \in \mathbf{R}^n \times [t_0,\infty) \times (0,t_0/2]. \tag{4.8}$$

LEMMA 4.2. There exist constants  $\sigma \in [0, 1)$  and  $C_4 > 0$  such that

$$J(x, t:s) \leq C_4 s^{-\sigma} \qquad for \quad 0 < s < t \leq \frac{1}{4} \quad and \quad |x| \ge 1.$$

$$(4.9)$$

To prove Lemmas 4.1 and 4.2, we prepare the following lemma.

LEMMA 4.3. Let p > (n+2)/n. Then there exists  $q \ge p$  satisfying

$$q > \frac{n}{n-2}$$
 and  $\frac{p}{p-1} - \frac{q}{q-1} < 1.$  (4.10)

*Proof.* If p > n/(n-2), then q = p satisfies (4.10). Assume that (n+2)/n . We see that <math>(n+2)/n < p implies (n-2)/n > 2-p. Choose q > 0 as (n-2)/n > 1/q > 2-p. Then we obtain q > n/(n-2) and pq - 2q + 1 > 0. The latter implies the second of (4.10).

*Proof of Lemma* 4.1. By Lemma 4.3 we can choose  $q \ge p$  satisfying (4.10). By virtue of (1.9) there exists a constant  $C_5 > 0$  such that

$$\begin{bmatrix} U(r) \end{bmatrix}^p \leq C_5 r^{-2p/(p-1)} \quad \text{for} \quad r \geq 1.$$

Since 2p/(p-1) is decreasing, we obtain

$$[U(r)]^{p} \leq C_{5} r^{-2q/(q-1)} \quad \text{for} \quad r \geq 1.$$
 (4.11)

With no loss of generality, we may assume  $t_0 < 1$ . Fix  $(x, t, s) \in \mathbf{R}^n \times [t_0, \infty) \times (0, t_0/2]$ . Note that  $0 < \sqrt{s} < 1$  for  $s \in (0, t_0/2]$ . We write

$$\begin{aligned} J(x, t:s) &= s^{-p/(p-1)} \left[ \int_{|y| \ge 1} + \int_{\sqrt{s} \le |y| \le 1} \\ &+ \int_{|y| \le \sqrt{s}} \Gamma(x-y:t-s) [U(|y|/\sqrt{s})]^p \, dy \right] \\ &\equiv s^{-p/(p-1)} [I_1 + I_2 + I_3]. \end{aligned}$$

From (4.11) we have

$$I_1 \leqslant C_5 s^{q/(q-1)} \int_{|y| \ge 1} \Gamma(x-y:t-s) \, dy \leqslant C_5 s^{q/(q-1)}$$

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$$I_2 \leq C_5 s^{q/(q-1)} \int_{\sqrt{s} \leq |y| \leq 1} \Gamma(x-y:t-s) |y|^{-2q/(q-1)} dy$$

Since  $\Gamma(x, t) \leq (4\pi t)^{-n/2}$  and  $t - s \geq t_0/2$ , we obtain

$$I_2 \leqslant C_5 (2\pi t_0)^{-n/2} s^{q/(q-1)} \int_{|y| \leqslant 1} |y|^{-2q/(q-1)} dy.$$

By virtue of q > n/(n-2), the integral on the right-hand side is convergent. To estimate  $I_3$ , we perform the change of variable  $z = y/\sqrt{s}$  to obtain

$$I_{3} \leq (2\pi t_{0})^{-n/2} \int_{|y| \leq \sqrt{s}} \left[ U(|y|/\sqrt{s}) \right]^{p} dy$$
$$\leq (2\pi t_{0})^{-n/2} s^{n/2} \int_{|z| \leq 1} \left[ U(|z|) \right]^{p} dz.$$

Therefore, we obtain

$$J(x, t:s) \leq C[s^{-p/(p-1)+q/(q-1)} + s^{-p/(p-1)+n/2}] \quad \text{for some} \quad C > 0.$$

Note that p > (n+2)/n implies p/(p-1) - n/2 < 1. Hence, we obtain (4.8) with

$$\sigma = \min\left\{\frac{p}{p-1} - \frac{q}{q-1}, \frac{p}{p-1} - \frac{n}{2}\right\} < 1.$$

This completes the proof.

*Proof of Lemma* 4.2. By the same argument in the proof of Lemma 4.1, we obtain (4.11). Let  $|x| \ge 1$  and let  $0 < s < t \le 1/4$ . Note that  $\sqrt{s} < \sqrt{t} \le 1/2$ . We write

$$\begin{aligned} J(x, t:s) &= s^{-p/(p-1)} \left[ \int_{|y| \ge 1/2} + \int_{\sqrt{s} \le |y| \le 1/2} \\ &+ \int_{|y| \le \sqrt{s}} \Gamma(x-y:t-s) \left[ U(|y|/\sqrt{s}) \right]^p dy \right] \\ &\equiv s^{-p/(p-1)} [I_1 + I_2 + I_3]. \end{aligned}$$

From (4.11) we have

$$I_1 \leq C_5 s^{q/(q-1)} \int_{|y| \ge 1/2} \Gamma(x-y:t-s) |y|^{-2q/(q-1)} \, dy \leq 2^{2q/(q-1)} C_5 s^{-q/(q-1)}$$

and

$$I_2 \leq C_5 s^{q/(q-1)} \int_{\sqrt{s} \leq |y| \leq 1/2} \Gamma(x-y:t-s) |y|^{-2q/(q-1)} dy.$$

Observe that  $|x - y| \ge 1/2$  for  $|x| \ge 1$  and  $|y| \le 1/2$ . Then we obtain

$$I_2 \leq C_5 \Gamma_{1/2} s^{-q/(q-1)} \int_{|y| \leq 1/2} |y|^{-2q/(q-1)} dy,$$

where  $\Gamma_{1/2} = \sup_{t>0} \{\Gamma(x-y:t): |x-y| \ge 1/2\} = \Gamma(1/2:n/8)$ . Since q > n/(n-2), the integral on the right-hand side is convergent. To estimate  $I_3$ , we perform the change of variable  $z = y/\sqrt{s}$  to obtain

$$I_{3} \leq \Gamma_{1/2} \int_{|y| \leq \sqrt{s}} \left[ U(|y|/\sqrt{s}) \right]^{p} dy \leq \Gamma_{1/2} s^{n/2} \int_{|z| \leq 1} \left[ U(|z|) \right]^{p} dz.$$

Therefore, by the same argument of the proof of Lemma 4.1, we obtain (4.9).

LEMMA 4.4. The function  $\Phi$  is well defined and continuous on  $\mathbb{R}^n \times (0, \infty)$ . Moreover, for every  $t_0 > 0$ , there exists a constant  $C = C(t_0) > 0$  such that

$$0 \leq \Phi(x, t) \leq C$$
 for  $x \in \mathbf{R}^n$  and  $t \geq t_0$ . (4.12)

*Proof.* For  $t \ge t_0 > 0$ , we write

$$\Phi(x, t) = \int_0^{t_0/2} J(x, t:s) \, ds + \int_{t_0/2}^t J(x, t:s) \, ds \equiv I_1 + I_2.$$

From (4.8), we have  $I_1 \leq (1-\sigma)^{-1} 2^{\sigma-1} C_3 t_0^{1-\sigma}$ . By (4.7) it follows that

$$I_2 \leq \int_{t_0/2}^{\infty} J(x, t:s) \, ds \leq (p-1) \, 2^{1/(p-1)} C_2 t_0^{-1/(p-1)}.$$

Thus we obtain (4.12). Since  $t_0 > 0$  is arbitrarily,  $\Phi$  is well defined on  $\mathbb{R}^n \times (0, \infty)$ .

LEMMA 4.5. (i)  $\Phi(x, t) \to 0$  as  $t \to 0$  uniformly in  $|x| \ge r$  for every r > 0; (ii) There exists a constant C > 0 such that

$$|x|^{2/(p-1)} \Phi(x,t) \leq C \quad for \quad x \in \mathbf{R}^n \setminus \{0\}, \quad 0 < t < \infty;$$

(iii)  $\Phi(\cdot, t) \to 0$  in  $L^1_{\text{loc}}(\mathbf{R}^n)$  as  $t \to 0$ .

*Proof.* First, we show that

$$\Phi(x, t) \to 0$$
 as  $t \to 0$  uniformly on  $|x| \ge 1$ , (4.13)

and that there exists a constant  $C_6 > 0$  satisfying

$$\Phi(x, t) \leqslant C_6 \qquad \text{for} \quad |x| \ge 1, \quad 0 < t < \infty.$$
(4.14)

In fact, by Lemma 4.2 we have

$$\Phi(x, t) \leq \frac{C_4}{1-\sigma} t^{1-\sigma} \quad \text{for} \quad 0 < t \leq \frac{1}{4} \quad \text{and} \quad |x| \ge 1.$$

This implies (4.13). From (4.12), we obtain (4.14).

By virtue of (4.6), we find that (4.13) implies (i). Moreover,

$$|\mu x|^{2/(p-1)} \Phi(\mu x, \mu^2 t) = |x|^{2/(p-1)} \Phi(x, t)$$
 for all  $\mu > 0$ .

Let |x| = 1. Then, by (4.14), we obtain

$$|\mu x|^{2/(p-1)} \Phi(\mu x, \mu^2 t) \leq C_6$$
 for all  $\mu > 0$ ,  $|x| = 1$ ,  $0 < t < \infty$ .

This implies (ii). Fix a compact set  $K \subset \mathbb{R}^n$ . By the Lebesgue dominated convergence theorem, we observe that

$$\int_{K} \Phi(x, t) \, dx = \int_{K} |x|^{-2/(p-1)} \left| |x|^{2/(p-1)} \Phi(x, t) \right| \, dx \to 0 \qquad \text{as} \quad t \to 0,$$

which implies (iii).

*Proof of Proposition* 4.2. Since  $\phi \leq \Phi$ , the function  $\phi$  is well defined on  $\mathbb{R}^n \times (0, \infty)$  and satisfies (i)–(iii) by Lemmas 4.4 and 4.5(i)–(iii). To show (iv) and (v), fix  $t_0$  and T with  $0 < t_0 < T$ . By the Fubini theorem and the property of the heat kernel  $\Gamma$ , we obtain

$$\phi(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y : t - t_0) \,\phi(y, t_0) \,dy + \int_{t_0}^t \int_{\mathbf{R}^n} \Gamma(x - y; t - s) \,g(y, s) \,dy \,ds$$

for  $(x, t) \in \mathbf{R}^n \times [t_0, T]$ . If g is continuous, then  $\phi \in C^{1,0}(\mathbf{R}^n \times [t_0, T])$ , and if (3.5) holds, then  $\phi \in C^{2,1}(\mathbf{R}^n \times [t_0, T])$  and satisfies  $\phi_t = \Delta \phi + g$  in

 $\mathbb{R}^n \times [t_0, T]$ . (See, e.g., [6, Chapter 1, Theorem 9].) Since  $t_0$  and T are arbitrarily, (iv) and (v) hold.

**4.3.** Observe that Proposition 3.1 is included by Proposition 4.2. We prove Proposition 3.2 by employing Propositions 4.1 and 4.2.

*Proof of Proposition* 3.2. Assume that *w* is defined by (3.6), that is,  $w = \psi + \phi$ . Then,  $w \in C^{2, 1}(\mathbb{R}^n \times (0, \infty))$  and is a solution to the problem of (3.2) by (4.1) and Proposition 4.2(iii) and (v). By (4.2) and Proposition 4.2(ii), *w* satisfies the growth condition (3.4).

Conversely, assume that *w* is a solution of (3.2) satisfying (3.4). Define  $\tilde{w}$  as

$$\tilde{w}(x, t) = \int_{\mathbf{R}^n} \Gamma(x - y : t) w_0(y) \, dy + \int_0^t \int_{\mathbf{R}^n} \Gamma(x - y : t - s) g(y, s) \, dy \, ds.$$

Then, by the argument above,  $\tilde{w}$  is a solution to the problem (3.2) satisfying (3.4). Let  $v = w - \tilde{w}$ . Then v satisfies the growth condition (3.4) and is a solution of the problem

 $v_t = \Delta v$  in  $\mathbf{R}^n \times (0, \infty)$  and v(x, 0) = 0 in the sense of  $L^1_{loc}(\mathbf{R}^n)$ .

By the uniqueness theorem ([1, Chapter 5, Theorem 6.1]), we have  $v \equiv 0$ , that is,  $w \equiv \tilde{w}$ . This implies (3.6).

Moreover, if  $w_0 \in C(\mathbb{R}^n \setminus \{0\})$ , then, by (4.3) and Proposition 4.2(i), we obtain (3.7).

### REFERENCES

- 1. E. DiBenedetto, "An Introduction to Partial Differential Equations," Birkhäuser, Boston, 1995.
- C. Dohmen and M. Hirose, Structure of positive radial solutions to the Haraux–Weissler equation, *Nonlinear Anal. TMA* 33 (1998), 51–69.
- M. Escobedo and O. Kavian, Variational problems related to self-similar solutions for the heat equation, *Nonlinear Anal. TMA* 11 (1987), 1103–1133.
- M. Escobedo and O. Kavian, Asymptotic behavior of positive solutions of a nonlinear heat equation, *Houston J. Math.* 13 (1987), 39–50.
- M. Escobedo, O. Kavian, and H. Matano, Large time behavior of solutions of a dissipative semilinear heat equation, *Comm. Partial Differential Equations* 27 (1995), 1427–1452.
- 6. A. Friedman, "Partial Differential Equations of Parabolic Type," Prentice-Hall, New Jersey, 1964.
- 7. H. Fujita, On the blowing up of solutions of the Cauchy problem for  $u_i = \Delta u + u^{1+\alpha}$ , J. Fac. Sci. Univ. Tokyo, Sect. I 13 (1966), 109–124.
- V. A. Galaktionov and J. L. Vazquez, Continuation of blowup solutions of nonlinear heat equations in several space dimensions, *Comm. Pure Appl. Math.* 50 (1997), 1–67.

- B. Gidas, W.-M. Ni, and L. Nirenberg, Symmetry and related properties via the maximum principle, *Comm. Math. Phys.* 68 (1979), 209–243.
- B. Gidas, W.-M. Ni, and L. Nirenberg, Symmetry of positive solutions of nonlinear elliptic equations in R", in "Mathematical Analysis and Applications, Part A" (L. Nachbin, Ed.), Adv. Math. Suppl. Stud., Vol. 7, pp. 369–402, Academic Press, New York, 1981.
- A. Gmira and L. Veron, Large time behaviour of the solutions of a semilinear parabolic equation in R<sup>N</sup>, J. Differential Equations 53 (1984), 258–276.
- A. Haraux and F. B. Weissler, Non-uniqueness for a semilinear initial value problem, Indiana Univ. Math. J. 31 (1982), 167–189.
- S. Kamin and L. A. Peletier, Large time behaviour of solutions of the heat equation with absorption, Ann. Scuola Norm. Sup. Pisa 12 (1985), 393–408.
- O. Kavian, Remarks on the large time behavior of a nonlinear diffusion equation, Annal. Institut Henri Poincaré-Analyse Nonlinéaire 4 (1987), 423–452.
- T. Kawanago, Asymptotic behavior of solutions of a semilinear heat equation with subcritical nonlinearity, Annal. Institut Henri Poincaré-Analyse Nonlinéaire 13 (1996), 1–15.
- 16. L. A. Peletier, D. Terman, and F. B. Weissler, On the equation  $\Delta u + \frac{1}{2}x \cdot \nabla u + f(u) = 0$ , Arch. Rational Mech. Anal. 94 (1986), 83–99.
- M. Protter and H. Weinberger, "Maximum Principles in Differential Equations," Prentice–Hall, Englewood Cliffs, New Jersey, 1967.
- J. Serrin, A symmetry problem in potential theory, Arch. Rational Mech. Anal. 43 (1971), 304–318.
- 19. F. B. Weissler, Existence and non-existence of global solutions for a semilinear heat equation, *Israel J. Math.* **38** (1981), 29–40.
- E. Yanagida, Uniqueness of rapidly decaying solutions to the Haraux–Weissler equation, J. Differential Equations 127 (1996), 561–570.