



Neutrino mixing: A_4 variations

Ernest Ma

Physics & Astronomy Department and Graduate Division, University of California, Riverside, CA 92521, USA



ARTICLE INFO

Article history:

Received 12 October 2015
 Received in revised form 12 November 2015
 Accepted 16 November 2015
 Available online 27 November 2015
 Editor: A. Ringwald

ABSTRACT

In the context of the non-Abelian discrete symmetry A_4 , the neutrino mass matrix has been studied extensively. A brief update is presented to focus on the conceptual shift from tribimaximal mixing ($\theta_{13} = 0, \theta_{23} = \pi/4, \tan^2 \theta_{12} = 1/2$) to *cobimaximal* mixing ($\theta_{13} \neq 0, \theta_{23} = \pi/4, \delta_{CP} = \pm\pi/2$) which agrees well with present data. Three specific realistic examples are proposed, two with three and the third with just two parameters.

© 2015 The Author. Published by Elsevier B.V. This is an open access article under the CC BY license (<http://creativecommons.org/licenses/by/4.0/>). Funded by SCOAP³.

The non-Abelian discrete symmetry A_4 is the symmetry of the tetrahedron. It has 12 elements and is the smallest group which admits an irreducible $\underline{3}$ representation. It also has three one-dimensional representations $\underline{1}, \underline{1}', \underline{1}''$. The basic multiplication rule is

$$\underline{3} \times \underline{3} = \underline{1} + \underline{1}' + \underline{1}'' + \underline{3} + \underline{3}. \quad (1)$$

Its application to neutrino mixing began with Ref. [1], where the representation matrices were chosen so that

$$a_1 b_1 + a_2 b_2 + a_3 b_3 \sim \underline{1}, \quad (2)$$

$$a_1 b_1 + \omega a_2 b_2 + \omega^2 a_3 b_3 \sim \underline{1}', \quad (3)$$

$$a_1 b_1 + \omega^2 a_2 b_2 + \omega a_3 b_3 \sim \underline{1}'', \quad (4)$$

$$(a_2 b_3 \pm a_3 b_2, a_3 b_1 \pm a_1 b_3, a_1 b_2 \pm a_2 b_1) \sim \underline{3}, \quad (5)$$

where $a_i, b_i \sim \underline{3}$ and $\omega = \exp(2\pi i/3) = -1/2 + i\sqrt{3}/2$. The three lepton families are assumed to transform as follows:

$$(v_i, l_i)_L \sim \underline{3}, \quad l_{iL}^c \sim \underline{1}, \underline{1}', \underline{1}'', \quad (6)$$

with three Higgs doublets $(\phi_i^+, \phi_i^0) \sim \underline{3}$. Hence the charged-lepton mass matrix is given by

$$\begin{aligned} \mathcal{M}_l &= \begin{pmatrix} f_e v_1^* & f_\mu v_1^* & f_\tau v_1^* \\ f_e v_2^* & f_\mu \omega^2 v_2^* & f_\tau \omega v_2^* \\ f_e v_3^* & f_\mu \omega v_3^* & f_\tau \omega^2 v_3^* \end{pmatrix} \\ &= \begin{pmatrix} v_1^* & 0 & 0 \\ 0 & v_2^* & 0 \\ 0 & 0 & v_3^* \end{pmatrix} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \omega^2 & \omega \\ 1 & \omega & \omega^2 \end{pmatrix} \begin{pmatrix} f_e & 0 & 0 \\ 0 & f_\mu & 0 \\ 0 & 0 & f_\tau \end{pmatrix}. \end{aligned} \quad (7)$$

For $v_1 = v_2 = v_3$, the A_4 symmetry breaks to its residual Z_3 and the unitary transformation linking \mathcal{M}_l to \mathcal{M}_ν is [2,3]

$$U_\omega = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \omega & \omega^2 \\ 1 & \omega^2 & \omega \end{pmatrix}. \quad (8)$$

In the (e, μ, τ) basis, the neutrino mass matrix (assumed Majorana) is

$$\mathcal{M}_\nu^{(e, \mu, \tau)} = U_\omega \mathcal{M}_A U_\omega^T. \quad (9)$$

In general, \mathcal{M}_A is a 3×3 symmetric complex matrix, i.e.

$$\mathcal{M}_A = \begin{pmatrix} a & c & e \\ c & d & b \\ e & b & f \end{pmatrix}. \quad (10)$$

For many years, the neutrino mixing matrix is conjectured to be of the tribimaximal form [4], so that

$$\mathcal{M}_\nu^{(e, \mu, \tau)} = U_B \mathcal{M}_B U_B^T, \quad (11)$$

where

$$U_B = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}. \quad (12)$$

If U_B is indeed the correct neutrino mixing matrix, then \mathcal{M}_B would be diagonal. In general however, it is given by [5,6]

$$\mathcal{M}_B = \begin{pmatrix} m_1 & m_6 & m_4 \\ m_6 & m_2 & m_5 \\ m_4 & m_5 & m_3 \end{pmatrix}, \quad (13)$$

where again $m_{1,2,3,4,5,6}$ are complex. Nonzero $m_{4,5,6}$ indicate thus the deviation from tribimaximal mixing. The A_4 basis is related to the tribimaximal basis through

E-mail address: ma@phyun8.ucr.edu.

$$\mathcal{M}_B = U_A^\dagger \mathcal{M}_A U_A^*, \quad (14)$$

where

$$U_A = U_\omega^\dagger U_B = \begin{pmatrix} 0 & 1 & 0 \\ 1/\sqrt{2} & 0 & i/\sqrt{2} \\ 1/\sqrt{2} & 0 & -i/\sqrt{2} \end{pmatrix}. \quad (15)$$

Their respective parameters are thus related by

$$m_1 = b + (d + f)/2, \quad m_2 = a, \quad m_3 = b - (d + f)/2, \quad (16)$$

$$m_4 = i(f - d)/2, \quad m_5 = i(e - c)/\sqrt{2}, \quad m_6 = (e + c)/\sqrt{2}. \quad (17)$$

To obtain tribimaximal mixing ($\theta_{13} = 0$, $\theta_{23} = \pi/4$, $\tan^2 \theta_{12} = 1/2$), $c = e = 0$ and $f = d$ are required. The remaining three parameters (a, b, d) are in general complex. To obtain *cobimaximal* mixing ($\theta_{13} \neq 0$, $\theta_{23} = \pi/4$, $\delta_{CP} = \pm\pi/2$) which agrees well with present data [7] with $\delta_{CP} = -\pi/2$ [8], what is required [9] is that \mathcal{M}_A be diagonalized by an orthogonal matrix. To see this, let

$$U_{lv} = U_\omega \mathcal{O}, \quad (18)$$

where \mathcal{O} is a real orthogonal matrix, then it is obvious that $U_{\mu i} = U_{\tau i}^*$ for $i = 1, 2, 3$. Comparing this with the Particle Data Group (PDG) convention of the neutrino mixing matrix, i.e.

$$U_{lv}^{PDG} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}, \quad (19)$$

it is obvious that after rotating the phases of the third column and the second and third rows, the two matrices are identical if and only if $s_{23} = c_{23}$ and $\cos \delta = 0$, i.e. $\theta_{23} = \pi/4$ and $\delta_{CP} = \pm\pi/2$. This important insight, i.e. Eq. (18), is a rediscovery of what was actually known already many years ago [10–12]. It is guaranteed if (a, b, c, d, e, f) are all real, so that \mathcal{M}_A is both symmetric and Hermitian.

Another way to arrive at cobimaximal mixing is to use Eqs. (9) and (10), i.e.

$$\mathcal{M}_v^{(e, \mu, \tau)} = U_\omega \begin{pmatrix} a & c & e \\ c & d & b \\ e & b & f \end{pmatrix} U_\omega^T = \begin{pmatrix} A & C & E^* \\ C & D^* & B \\ E^* & B & F \end{pmatrix}, \quad (20)$$

where

$$A = (a + 2b + 2c + d + 2e + f)/3, \quad (21)$$

$$B = (a - b - c + d - e + f)/3, \quad (22)$$

$$C = (a - b - \omega^2 c + \omega d - \omega e + \omega^2 f)/3, \quad (23)$$

$$D^* = (a + 2b + 2\omega c + \omega^2 d + 2\omega^2 e + \omega f)/3, \quad (24)$$

$$E^* = (a - b - \omega c + \omega^2 d - \omega^2 e + \omega f)/3, \quad (25)$$

$$F = (a + 2b + 2\omega^2 c + \omega d + 2\omega e + \omega^2 f)/3. \quad (26)$$

If again (a, b, c, d, e, f) are real, then A, B are real, whereas $E = C$ and $F = D$. This well-known special form was written down already many years ago [13,14], and it was pointed out soon afterward [15] that it is protected by a generalized CP transformation under $\mu - \tau$ exchange, and it guarantees cobimaximal mixing. With the knowledge that $\theta_{13} \neq 0$ [16–18], this extended symmetry is now the subject of many investigations, which began with generalized S_4 [19]. In fact, such remnant residual CP symmetries are under active study [20–22] to reconstruct the neutrino mixing matrix with cobimaximal mixing.

Since tribimaximal mixing is not what the data show, \mathcal{M}_B cannot be diagonal. Many studies are then centered on looking for small off-diagonal terms, i.e. $m_{4,5,6}$ which may be complex. On the other hand, data are perfectly consistent with \mathcal{M}_A as long as it is real. Of course, θ_{13} and θ_{12} are not predicted, but if extra conditions are imposed, they may be correlated. For example, it has been proposed [23] that $c = e = 0$, but $f \neq d$, with a, b, d, f real for \mathcal{M}_A in Eq. (10). This yields cobimaximal mixing together with the prediction that

$$\tan^2 \theta_{12} = \frac{1}{2 - 3 \sin^2 \theta_{13}} > \frac{1}{2}. \quad (27)$$

Using the 2014 Particle Data Group value [7]

$$\sin^2(2\theta_{13}) = (9.3 \pm 0.8) \times 10^{-2}, \quad (28)$$

the value of $\sin^2(2\theta_{12})$ from Eq. (27) is 0.90 with very little deviation, as compared with the PDG value

$$\sin^2(2\theta_{12}) = 0.846 \pm 0.021, \quad (29)$$

which is more than two standard deviations away. This is a generic result corresponding to choosing $m_5 = m_6 = 0$ in Eq. (13).

If $m_4 = m_6 = 0$ is chosen instead, then another generic prediction is

$$\tan^2 \theta_{12} = \frac{1}{2}(1 - 3 \sin^2 \theta_{13}). \quad (30)$$

Again using Eq. (28), $\sin^2(2\theta_{12}) = 0.866 \pm .002$ is obtained, which agrees with Eq. (29) to within one standard deviation. Note that both generic results hold for arbitrary values of δ_{CP} .

In Ref. [6], $e + c = 0$ is assumed so that $m_6 = 0$. In addition, $\delta_{CP} = 0$ and $\theta_{23} = \pi/4$ are assumed, which can be achieved if both m_4 and m_5 are nonzero. In the case $m_4 = m_6 = 0$, but $m_{1,2,3,5}$ complex, an analysis shows [24] that large δ_{CP} correlates with $\theta_{23} \neq \pi/4$ for a fixed nonzero θ_{13} . With the present data, these scenarios are no longer favored. The message now is that cobimaximal mixing should be chosen as the preferred starting point of any improved model of neutrino mass and mixing.

Consider a real \mathcal{M}_A of Eq. (10) with $d = f$ and $c = -e$, i.e.

$$\mathcal{M}_A = \begin{pmatrix} a & -e & e \\ -e & d & b \\ e & b & d \end{pmatrix}. \quad (31)$$

In that case,

$$\mathcal{M}_B = \begin{pmatrix} b+d & 0 & 0 \\ 0 & a & i\sqrt{2}e \\ 0 & i\sqrt{2}e & b-d \end{pmatrix}, \quad (32)$$

i.e. $m_4 = m_6 = 0$, hence the desirable condition of Eq. (30) is obtained. Let \mathcal{M}_B be diagonalized by

$$U_E = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c & is \\ 0 & is & c \end{pmatrix}, \quad (33)$$

so that

$$\mathcal{M}_B = U_E \begin{pmatrix} m'_1 & 0 & 0 \\ 0 & m'_2 & 0 \\ 0 & 0 & m'_3 \end{pmatrix} U_E^T, \quad (34)$$

where $s = \sin \theta_E$, $c = \cos \theta_E$. Then

$$\frac{sc}{c^2 - s^2} = \frac{e\sqrt{2}}{a + b - d}, \quad (35)$$

and the three neutrino mass eigenvalues are

$$m'_1 = b + d, \quad (36)$$

$$m'_2 = \frac{1}{c^2 - s^2} [c^2 a + s^2 (b - d)], \quad (37)$$

$$m'_3 = \frac{1}{c^2 - s^2} [s^2 a + c^2 (b - d)]. \quad (38)$$

The neutrino mixing matrix is now $U_B U_E$, from which

$$s = \sqrt{3} \sin \theta_{13} \quad (39)$$

is obtained. As it is, \mathcal{M}_A has four real parameters (a, b, d, e) to fit three observables ($\theta_{13}, \Delta m_{21}^2, \Delta m_{32}^2$), hence no prediction is possible other than cobimaximal mixing and Eq. (30).

In the case of tribimaximal mixing, i.e. $e = 0$, the simplest A_4 model [25,26] has $d = a$. With this condition, but $e \neq 0$, the three neutrino masses are

$$m'_1 = b + a, \quad m'_2 = a + \frac{s^2 b}{c^2 - s^2}, \quad m'_3 = -a + \frac{c^2 b}{c^2 - s^2}. \quad (40)$$

Using Eq. (28) with the central value $s = 0.2673$, they become

$$m'_1 = b + a, \quad m'_2 = a + 0.08336b, \quad m'_3 = -a + 1.08336b. \quad (41)$$

Using the central values of [7]

$$\begin{aligned} \Delta m_{21}^2 &= 7.53 \pm 0.18 \times 10^{-5} \text{ eV}^2, \\ \Delta m_{32}^2 &= 2.44 \pm 0.06 \times 10^{-3} \text{ eV}^2, \end{aligned} \quad (42)$$

the solution is $b/a = -1.714$ and $a = 0.0183 \text{ eV}$, with $e/a = -0.3642$. Using Eq. (20), the effective neutrino mass in neutrinoless double beta decay is predicted to be

$$m_{ee} = |A| = |a + 2b/3| = 2.6 \times 10^{-3} \text{ eV}, \quad (43)$$

which is very small, as expected from a normal ordering of neutrino masses, and beyond the sensitivity of current and planned experiments.

Another possible three-parameter model is to assume $d = b$, then

$$m'_1 = 2b, \quad m'_2 = 1.08336a, \quad m'_3 = 0.08336a. \quad (44)$$

This implies inverted ordering of neutrino masses with $a = 0.0465 \text{ eV}$, $b = 0.0248 \text{ eV}$, and $e = 0.0099 \text{ eV}$. Hence $m_{ee} = |(a + 4b)/3| = 0.0486 \text{ eV}$ which is presumably verifiable in the future.

As a third example, consider the following *new remarkable model* of just two parameters, with $d = -b = 2a$:

$$\begin{aligned} m'_1 &= 0, \quad m'_2 = \left(\frac{c^2 - 4s^2}{c^2 - s^2} \right) a = 0.75a, \\ m'_3 &= \left(\frac{s^2 - 4c^2}{c^2 - s^2} \right) a = -4.25a. \end{aligned} \quad (45)$$

As a result, $\Delta m_{21}^2 / \Delta m_{32}^2$ is predicted to be 0.032, in excellent agreement with the experimental value of 0.031. (This is a totally new result.) In this case, $a = 0.0116 \text{ eV}$ and $m_{ee} = |a/3| = 3.9 \times 10^{-3} \text{ eV}$, with $e/a = -0.6375$.

In conclusion, it has been pointed out in this paper that A_4 is intimately related to *cobimaximal* mixing ($\theta_{13} \neq 0$, $\theta_{23} = \pi/4$, $\delta_{CP} = \pm\pi/2$) which agrees well with present data, and should replace the previously preferred tribimaximal mixing pattern. In particular, a model is proposed with just *two real parameters*, with the following predictions:

$$\theta_{23} = \pi/4, \quad \delta_{CP} = \pm\pi/2, \quad \tan^2 \theta_{12} = \frac{1}{2} (1 - 3 \sin^2 \theta_{13}), \quad (46)$$

$$\begin{aligned} \frac{\Delta m_{21}^2}{\Delta m_{31}^2} &= \left(\frac{1 - 15 \sin^2 \theta_{13}}{4 - 15 \sin^2 \theta_{13}} \right)^2, \\ m_{ee} &= 3.9 \times 10^{-3} \text{ eV (for } \sin^2 2\theta_{13} = 0.093), \end{aligned} \quad (47)$$

which are all well satisfied by present data (except m_{ee} which is yet to be measured).

This work is supported in part by the U.S. Department of Energy under Grant No. DE-SC0008541.

References

- [1] E. Ma, G. Rajasekaran, Phys. Rev. D 64 (2001) 113012.
- [2] N. Cabibbo, Phys. Lett. B 72 (1978) 333.
- [3] L. Wolfenstein, Phys. Rev. D 18 (1978) 958.
- [4] P.F. Harrison, D.H. Perkins, W.G. Scott, Phys. Lett. B 530 (2002) 167.
- [5] E. Ma, Phys. Rev. D 70 (2004) 031901(R).
- [6] E. Ma, D. Wegman, Phys. Rev. Lett. 107 (2011) 061803.
- [7] Particle Data Group, K.A. Olive, et al., Chin. Phys. C 38 (2014) 090001.
- [8] K. Abe, et al., T2K Collaboration, Phys. Rev. D 91 (2015) 072010.
- [9] E. Ma, Phys. Rev. D 92 (2015) 051301(R).
- [10] K. Fukuura, T. Miura, E. Takasugi, M. Yoshimura, Phys. Rev. D 61 (2000) 073002.
- [11] T. Miura, E. Takasugi, M. Yoshimura, Phys. Rev. D 63 (2001) 013001.
- [12] P.F. Harrison, W.G. Scott, Phys. Lett. B 547 (2002) 219.
- [13] E. Ma, Phys. Rev. D 66 (2002) 117301.
- [14] K.S. Babu, E. Ma, J.W.F. Valle, Phys. Lett. B 552 (2003) 207.
- [15] W. Grimus, L. Lavoura, Phys. Lett. B 579 (2004) 113.
- [16] K. Abe, et al., T2K Collaboration, Phys. Rev. Lett. 107 (2011) 041801.
- [17] F.P. An, et al., Daya Bay Collaboration, Phys. Rev. Lett. 108 (2012) 171803.
- [18] S.-B. Kim, et al., RENO Collaboration, Phys. Rev. Lett. 108 (2012) 191802.
- [19] R.N. Mohapatra, C.C. Nishi, Phys. Rev. D 86 (2012) 073007.
- [20] P. Chen, C.-Y. Yao, G.-J. Ding, Phys. Rev. D 92 (2015) 073002.
- [21] A.S. Joshipura, K.M. Patel, Phys. Lett. B 749 (2015) 159.
- [22] H.-J. He, W. Rodejohann, X.-J. Xu, arXiv:1507.03541 [hep-ph].
- [23] X.-G. He, arXiv:1504.01560 [hep-ph].
- [24] H. Ishimori, E. Ma, Phys. Rev. D 86 (2012) 045030.
- [25] G. Altarelli, F. Feruglio, Nucl. Phys. B 720 (2005) 64.
- [26] E. Ma, Phys. Rev. D 72 (2005) 037301.