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Yellow Protein Is Controlled by the External pH

# The Transient Accumulation of the Signaling State of Photoactive

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ABSTRACT The signaling state of the photoreceptor photoactive yellow protein is the long-lived intermediate I<sub>2</sub>'. The pH dependence of the equilibrium between the transient photocycle intermediates  $I_2$  and  $I_2'$  was investigated. The formation of  $I_2'$  from I<sub>2</sub> is accompanied by a major conformational change. The kinetics and intermediates of the photocycle and of the photoreversal were measured by transient absorption spectroscopy from pH 4.6 to 8.4. Singular value decomposition (SVD) analysis of the data at pH 7 showed the presence of three spectrally distinguishable species: I1, I2, and I2'. Their spectra were determined using the extrapolated difference method. I<sub>2</sub> and I<sub>2</sub>' have electronic absorption spectra, with maxima at 370  $\pm$  5 and 350  $\pm$  5 nm, respectively. Formation of the signaling state is thus associated with a change in the environment of the protonated chromophore. The time courses of the I<sub>1</sub>, I<sub>2</sub>, and I<sub>2</sub>' intermediates were determined from the wavelength-dependent transient absorbance changes at each pH, assuming that their spectra are pH-independent. After the formation of  $I_{2}$  (~2 ms), these three intermediates are in equilibrium and decay together to the initial dark state. The equilibrium between  $I_2$  and  $I_2'$  is pH dependent with a pK<sub>a</sub> of 6.4 and with  $l_2'$  the main species above this pK<sub>a</sub>. Measurements of the pH dependence of the photoreversal kinetics with a second flash of 355 nm at a delay of 20 ms confirm this pK<sub>a</sub> value. I<sub>2</sub> and I<sub>2</sub>' are photoreversed with reversal times of  $\sim$ 55  $\mu$ s and several hundred microseconds, respectively. The corresponding signal amplitudes are pH dependent with a pK<sub>a</sub> of ~6.1. Photoreversal from  $l_2$ dominates above the pKa. The transient accumulation of I2', the active state of photoactive yellow protein, is thus controlled by the proton concentration. The rate constant  $k_3$  for the recovery to the initial dark state also has a pK<sub>a</sub> of ~6.3. This equality of the equilibrium and kinetic p $K_a$  values is not accidental and suggests that  $k_3$  is proportional to  $[I_2']$ .

### INTRODUCTION

Photoactive yellow protein (PYP) is significant as the structural prototype for a large and diverse superfamily of signaling proteins that share a common structural motif termed the PAS domain (1,2). PAS domains are found in all kingdoms of life, generally as the sensory component of multi-domain proteins. The structure of PYP has a central six-stranded antiparallel  $\beta$ -sheet, with four structural features: an N-terminal cap, a PAS core with the first three  $\beta$ -strands of the central  $\beta$ -sheet, a helical connector, and a so-called  $\beta$ -scaffold consisting of the last three  $\beta$ -strands (3–7). PYP from the halophilic purple phototrophic bacterium Halorhodospira halophila is a small 14-kDa soluble cytoplasmic protein. The physiological function of H. halophila PYP was reported to be in negative phototaxis, resulting in movement away from blue/UV light (8). Its chromophore is p-hydroxycinnamic acid covalently bound via a thioester linkage to cysteine-69. In the dark, the phenol group of the chromophore is deprotonated, and the  $C_7=C_8$  bond is *trans*. The ionization of the chromophore and its hydrogen bonding to E-46 and Y-42 are mainly responsible for the observed spectral tuning in the dark state ( $\lambda_{max} \approx 446$  nm). Lightinduced *trans-cis* isomerization around the  $C_7=C_8$  bond is

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rapid (<3 ps) and is followed by a sequence of slower relaxations in the dark that ultimately lead to recovery of the dark state in less than 1 s. This photocycle has already been studied in considerable detail (9-11). The first long-lived intermediate, after the two very short-lived intermediates  $I_0$ and  $I_0^{\ddagger}$ , is  $I_1$ .  $I_1$  forms in ~3 ns and has a red-shifted absorption spectrum ( $\lambda_{max} \approx 460$  nm). In several hundred microseconds it decays to I<sub>2</sub>. I<sub>2</sub> has a protonated chromophore and a blue-shifted absorption spectrum that is commonly believed to have its  $\lambda_{\text{max}}$  value at 355 nm. In I<sub>2</sub> the chromophore phenol is partially exposed to the aqueous medium and hydrogen-bonded to the side chain of R-52 (6). The protonation of the chromophore occurs either intramolecularly from E-46 (12) or from the external medium (13). In several milliseconds  $I_2$  is transformed into  $I_2'$ .  $I_2'$  is believed to be the signaling state. It also has a protonated chromophore with a  $\lambda_{max}$  value similar to that of I<sub>2</sub>. The I<sub>2</sub> to I<sub>2</sub>' transition is associated with a major global structural change, which has been documented by NMR (14,15), CD (16), small-angle x-ray scattering (17), and FTIR (12,18). Formation of  $I_2'$  is associated with exposure of a hydrophobic surface patch (19), presumably the recognition and binding site for a response regulator. Experiments with hydrophobic dyes showed that these bind transiently to  $I_2'$  but not to  $I_2$  (13). PYP shares a number of features, such as spectral tuning, photoisomerization, transient chromophore (de)protonation, and photoreversal, with other photoreceptors such as rhodopsin and phytochrome. This together with the availability

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of high-resolution structures of intermediates (3–7) makes PYP an attractive model system for signal transduction (1).

For an understanding of the mechanism of PYP, a detailed characterization of its photocycle is essential. Such studies have been carried out by both electronic (9-11,20) and vibrational (21,22) spectroscopy. The photocycle kinetics are pH (23–25) and salt (26,27) dependent. Whereas early models assumed a unidirectional sequential mechanism (9,20), it is becoming increasingly clear that back reactions and equilibria between intermediates play essential roles (11,24,25,28). These equilibria are also pH (24-26,28-30) and salt dependent (26). An example of the pH- and salt-dependent equilibrium between  $I_1$  and  $I_2/I_2'$  was recently described for the mutant Y98Q (26). At alkaline pH, the  $I_1'$  and  $I_2'$  intermediates are in equilibrium with a pK<sub>a</sub> of  $\sim$ 9.9 (25,28). Most photocycle intermediates can be photoreversed when excited with light of the appropriate wavelength at the right time (31–33). Using this double-flash method, we showed that the  $I_2$  and  $I_2'$  intermediates are in equilibrium (33).  $I_2$  partly decays to  $I_2'$  in  $\sim 2$  ms and then remains in equilibrium with  $I_2'$  until the end of the cycle (33). Our laboratory showed from time-resolved fluorescence and photostationary absorption measurements that this  $I_2/I_2'$  equilibrium is pH dependent with a  $pK_a$  of 6.3 (34). Interestingly the rate constant for the kinetics of the ground state recovery has a bell-shaped pH dependence with a lower  $pK_a$  of 6.4 (23), i.e., identical to the value for the  $I_2/I_2'$  equilibrium (34). The nature of the group(s) responsible for this pK<sub>a</sub> is still not entirely certain, but it is commonly ascribed to E-46 (35,36).

The I<sub>2</sub> and I<sub>2</sub>' intermediates were originally characterized by time-resolved FTIR spectroscopy (12,18), but they can also be distinguished by transient absorption spectroscopy in the UV/visible. Indeed, our laboratory recently obtained  $\lambda_{max}$  values of 372 and 352 nm for I<sub>2</sub> and I<sub>2</sub>', respectively, from measurements of the pH dependence of the photostationary absorbance in the presence of background illumination (34). Here we characterize the equilibrium and spectra of the I<sub>2</sub> and I<sub>2</sub>' intermediates by kinetic methods using transient absorption spectroscopy in the UV/visible and the extrapolated difference method (37). These methods allow a direct determination of the time-dependent intermediate populations and their pH dependencies. We find that the transient accumulation of the signaling state (I<sub>2</sub>') increases with pH with a pK<sub>a</sub> of ~6.4.

#### MATERIALS AND METHODS

#### Protein production and purification

*H. halophila* holo-PYP was produced by coexpression with the biosynthetic enzymes TAL and pCL and subsequently purified from *Escherichia coli* BL 21 (DE3) as described (38).

#### Transient absorption spectroscopy

Time-resolved absorption spectroscopy with single- and double-flash excitation was performed as described (13,24,33,39). To resolve the photoreversal kinetics, the data acquisition was triggered on the second flash (33,39). SVD methods were used as described (13,24,40–42).

#### Data analysis

Spectra and time courses of intermediates were obtained using the extrapolated difference method as described (25,37). The transient absorbance change  $\Delta A (\lambda, t)$  is given by a sum of contributions from the spectral intermediates *i*, whose relative concentration is given by  $n_i(t)$ . In matrix notation:

$$\Delta \mathbf{A} = (\mathbf{A} - \mathbf{A}_{p})\mathbf{n} \tag{1}$$

where the columns of **A** are the spectra of intermediates and each of the identical columns of  $A_p$  is the spectrum of the dark state P. Assuming first-order kinetics,  $\Delta A$  can be represented by a sum of exponentials. The wavelength-dependent coefficients of the exponentials are the amplitude spectra  $B_i(\lambda)$ . The  $B_i(\lambda)$  are given in matrix notation by:

$$\mathbf{B} = (\mathbf{A} - \mathbf{A}_{p})\mathbf{C} \tag{2}$$

where  $C_{ii}$  is the weight of the *j*th exponential in the time dependence of  $n_i(t)$ .

The amplitude spectra are ordered from low to high apparent time constants in the matrix **B**. New matrices  $\tilde{\mathbf{B}}$  and  $\tilde{\mathbf{C}}$  are formed from the columns of **B** and **C** by adding up columns (25,37). The columns of  $\tilde{\mathbf{B}}$  represent the extrapolated absorption difference spectra; the columns of  $\tilde{\mathbf{C}}$  contain the relative contributions of the intermediates in these difference spectra. In the case of PYP at acid and neutral pH, only three intermediates contribute, as we will see, in the time range investigated: I<sub>1</sub>, I<sub>2</sub>, and I<sub>2</sub>'. Their relative contributions to the *i*th extrapolated absorption difference spectrum will be called (*x<sub>i</sub>*, *y<sub>i</sub>*, *z<sub>i</sub>*) in the order I<sub>1</sub>, I<sub>2</sub>, I<sub>2</sub>'.

The intermediate spectrum  $(\mathbf{A})_i$  can then be calculated from column *i* of matrix  $\tilde{\mathbf{C}}^{-1}$ :

$$\left(\mathbf{A}\right)_{i} = \mathbf{\tilde{B}}(\mathbf{\tilde{C}}^{-1})_{i} + \left(\mathbf{A}_{p}\right)_{i}$$
(3)

To calculate the elements of  $\tilde{C}$  the following two constraints are introduced:

- 1. The sum of the relative intermediate concentrations of I<sub>1</sub>, I<sub>2</sub>, and I<sub>2</sub>' in the extrapolated difference spectra is constant at all times before ground state recovery and equals the fraction of molecules cycling,  $\eta$ . This means that the sum of *x*, *y*, and *z* for each extrapolated difference spectrum equals  $\eta$  or that the sum of the matrix elements of each column of  $\tilde{\mathbf{C}}$  equals  $\eta$ , and that of  $\tilde{\mathbf{C}}^{-1}$  is  $1/\eta$ . This is simply the conservation law for the number of cycling molecules in a sequential cycle.
- 2. The absorption of  $I_2{'}$  is identical to zero for wavelengths larger than or equal to 410 nm.  $I_2{'}$  has its absorption maximum for the longest wavelength  $S_0\text{-}S_1$  transition around 355 nm. Because the spectra of all intermediates have similar bandwidths, we can estimate from the spectrum of P, which does not absorb beyond 500 nm, that  $I_2{'}$  does not absorb beyond 410 nm.

Matrix calculations were performed with Matlab version R12.1. Fits with sum of exponentials were carried out with Microcal Origin 7.5.

#### RESULTS

# Spectra and time courses of intermediates at pH 7: extrapolated difference method

Transient absorbance changes were measured at 19 wavelengths, ranging from 330 to 510 nm, in the time domain from 50 ns to 5 s. Fig. 1 *A* shows data at pH 7. For clarity, time traces at only eight wavelengths are shown. The complete data set was subjected to SVD analysis. The first



FIGURE 1 (*A*) Transient absorption changes after excitation at 430 nm at 19 wavelengths varying from 330 to 510 nm. For clarity, only the traces at the indicated wavelengths are shown. The vertical dashed lines indicate the time constants for a global fit to the weighted SVD time traces with a sum of three exponentials.  $\tau_1 = 270 \ \mu s$  is the rise time of  $I_2$ ,  $\tau_2 = 2.0 \ ms$  is the rise time of  $I_2$ ,  $\tau_3 = 260 \ ms$  is the return to P. The dotted lines, only distinguishable from the data in the microsecond time range, are the fits. Conditions: pH 7, 20°C, 50 mM KCl, and 50 mM Tris. PYP concentration 35  $\mu$ M. (*B*) Amplitude spectra  $B_i(\lambda)$  calculated from the amplitudes of the exponential fits to the SVD time traces and the corresponding basis spectra of the data in *A*. The three amplitude spectra correspond to the following time constants:  $\tau_1 = 270 \ \mu s$  ( $\odot$ ),  $\tau_2 = 2.0 \ ms$  ( $\Box$ ),  $\tau_3 = 260 \ ms$  ( $\odot$ ). The solid curve is a scaled and inverted ground-state spectrum. (*C*) Extrapolated difference spectra obtained from the amplitude spectra of *B* as described in the text:  $\mathbf{\tilde{B}}_1$  ( $\blacktriangle$ ),  $\mathbf{\tilde{B}}_2$  ( $\diamondsuit$ ), and  $\mathbf{\tilde{B}}_3$  ( $\bigcirc$ ).

six singular values were 11.2, 1.8, 0.14, 0.05, 0.03, and 0.02. We consider the first three to be significant, suggesting the presence of only three spectrally distinguishable intermediates. The additional components ( $s_4$ ,  $s_5$ , and  $s_6$ ) show very

noisy time traces and were therefore neglected. The three weighted time traces from SVD were fitted simultaneously starting at 10  $\mu$ s with a sum of three exponentials with time constants  $\tau_1 = 270 \ \mu s$ ,  $\tau_2 = 2.0 \ ms$ , and  $\tau_3 = 260 \ ms$ . These times are marked by vertical dashed lines in Fig. 1 A. The dotted lines in Fig. 1 A, which can barely be distinguished from the data, represent these fit curves for the individual time traces. From the fit to the SVD time traces and the corresponding basis spectra, the amplitude spectra  $B_i(\lambda)$ were calculated as described (42). These are presented in Fig. 1 B. The amplitude spectra provide considerable insight into the spectra of the intermediates.  $B_1(\lambda)$  clearly describes the transition from I<sub>1</sub> ( $\lambda_{max} \approx 460 \text{ nm}$ ) to I<sub>2</sub> with a  $\lambda_{max}$  value above 360 nm.  $B_2(\lambda)$  is apparently a transition from an equilibrium of  $I_1$  and  $I_2$  to  $I_2'$  with  $I_2'$  blue-shifted with respect to I<sub>2</sub>.  $B_3(\lambda)$  represents the ground state recovery and suggests that  $I_2'$  has its  $\lambda_{max}$  value near 350 nm. Comparison of the negative minimum of  $\mathbf{B}_1$  with the positive maximum of  $\mathbf{B}_3$  suggests that the transition from  $I_2$  to  $I_2'$  is associated with a blue shift of the order of  $\sim 20$  nm. So qualitatively the conclusion that there is a blue shift between the early UV intermediate  $I_2$  and the later UV intermediate  $I_2'$  is already apparent from the data alone (amplitude spectra) without any model-dependent assumptions. The following quantitative analysis, which does use two plausible constraints, strengthens this conclusion.

We recently described in detail how to construct the intermediate spectra and time courses from the **B** spectra using the extrapolated difference method (25,37). The main equations were briefly summarized in Materials and Methods. The three amplitude spectra **B**<sub>1</sub>, **B**<sub>2</sub>, and **B**<sub>3</sub> were used to construct the  $\tilde{\mathbf{B}}$  matrix as described (25,37). The three columns,  $\tilde{\mathbf{B}}_1$ ,  $\tilde{\mathbf{B}}_2$ , and  $\tilde{\mathbf{B}}_3$ , representing the extrapolated difference spectra are presented in Fig. 1 *C*.  $\tilde{\mathbf{B}}_1$  equals the initial absorbance change right after the flash and suggests that the initial bleach led to the formation of the I<sub>1</sub> intermediate (positive absorbance change near 480 nm).

We now use the second constraint as described in Materials and Methods, that  $I_2'$  does not absorb beyond 410 nm, and consider Eq. 3 only in the range  $\lambda > 410$  nm; i.e., we drop the rows for the shorter wavelengths. Then the third column of the reduced matrix **A**, corresponding to the spectrum of  $I_2'$ , is the null vector:  $(\mathbf{A})_3 = \vec{0}$ . Because  $(\mathbf{A})_3$  is zero, we can solve Eq. 3 for  $(\tilde{\mathbf{C}}^{-1})_3$ . In this way, we determine the third column of  $\tilde{\mathbf{C}}^{-1}$ . Under the first constraint, the sum of these matrix elements equals  $\eta^{-1}$ . In this way, we find, for the fraction of molecules cycling,  $\eta = 0.371$ . Finally, using  $\tilde{\mathbf{B}}$  and  $\mathbf{A}_p$  for the whole spectral range allows us to calculate the spectrum of  $I_2'$  from  $(\tilde{\mathbf{C}}^{-1})_3$  using Eq. 3. The result is shown in Fig. 2 *A* (*solid squares*). The  $\lambda_{max}$  value of the spectrum is at ~350 ± 5 nm.

Because only  $I_1$  contributes to  $\tilde{\mathbf{B}}_1$  (see Fig. 1 *C*), the elements  $\tilde{\mathbf{C}}_{21} = y_1$  and  $\tilde{\mathbf{C}}_{31} = z_1$  of  $\tilde{\mathbf{C}}$  are given by  $y_1 = z_1 = 0$ . This allows us to calculate the elements of the first column of  $\tilde{\mathbf{C}}^{-1}$ . The result is  $\tilde{\mathbf{C}}_{11}^{-1} = 1/x_1, \tilde{\mathbf{C}}_{21}^{-1} = 0, \tilde{\mathbf{C}}_{31}^{-1} = 0$ .

Because the sum of these elements equals  $1/\eta$  (conservation constraint), we have  $x_1 = \eta = 0.371$ . With  $(\tilde{\mathbf{C}}^{-1})_1$  now completely known, we can calculate the spectrum of  $I_1$  from  $(\tilde{\mathbf{C}}^{-1})_1, \tilde{\mathbf{B}}$ , and  $\mathbf{A}_p$  using Eq. 3. The result is shown in Fig. 2 *A* (*open squares*). This spectrum of  $I_1$  is in good agreement with that obtained in previous work (25,26).

To calculate the spectrum of the third spectral species,  $I_2$ , we proceeded as follows. From Fig. 1 *B*, we note that **B**<sub>1</sub>



FIGURE 2 (*A*) Intermediate spectra  $I_1$  ( $\Box$ ),  $I_2$  ( $\bigcirc$ ), and  $I_2'$  ( $\blacksquare$ ) calculated from the extrapolated difference spectra of Fig. 1 C. The solid curve represents the spectrum of dark state *P* for comparison. Vertical dashed lines indicate the wavelengths of the blue (430 nm) and violet (355 nm) excitation flashes used. (*B*)  $I_2$  spectrum for various allowed values of  $y_2$  as described in the text.  $y_2 = 0.12$  ( $\bigcirc$ ),  $y_2 = 0.15$  ( $\blacksquare$ ),  $y_2 = 0.18$  ( $\blacktriangle$ ),  $y_2 = 0.22$  ( $\bigcirc$ ),  $y_2 = 0.27$  ( $\Box$ ), and  $y_2 = 0.37$  ( $\triangle$ ). (*C*) Time courses of the relative concentrations of  $I_1$  (*solid line*),  $I_2$  (*dashed line*), and  $I_2'$  (*dotted line*) calculated according to Eq. 1. The time course of the sum of the relative concentrations of  $I_1$ ,  $I_2$ , and  $I_2'$  is indicated by dashed-dot-dashed line. The vertical dashed lines indicate the time constants from the global SVD fit of Fig. 1 *A*.

reflects a transition between two intermediates with  $\lambda_{max}$ values of ~460 nm (decay of  $I_1$ ) and 370 nm (rise of  $I_2$ ). Because there is apparently no contribution from the more blue-shifted species  $I_2'$  ( $\lambda_{max} \approx 350 \text{ nm}$ ) in **B**<sub>1</sub>, which is well known to be formed from  $I_2$  in the next transition (12),  $I_2'$  does not contribute to  $\mathbf{\tilde{B}}_2$  either. Dye-binding experiments also showed that the formation of the signaling state  $I_2'$  is delayed with respect to the formation of  $I_2$  (13). Therefore, we conclude that  $I_2'$  is not involved in the first transition, and thus,  $z_2 = 0$ . The elements  $x_2$ ,  $y_2$ ,  $z_2$  of the second column of  $\tilde{C}$ can now be expressed in terms of  $x_2$ ,  $y_2$ , and  $\eta$  with the help of  $\mathbf{\hat{C}}^{-1}\mathbf{\hat{C}} = \mathbf{I}$ . Using the conservation constraint,  $x_2 + y_2 = \eta$ , we finally obtain for the elements of the second column of  $\tilde{\mathbf{C}}^{-1}$ ,  $\tilde{\mathbf{C}}_{12}^{-1} = -(\eta - \mathbf{y}_2)/\eta \mathbf{y}_2$ ,  $\tilde{\mathbf{C}}_{22}^{-1} = 1/y_2$ ,  $\tilde{\mathbf{C}}_{32}^{-1} = 0$ . So we have now determined all elements of the second column of  $\tilde{\mathbf{C}}^{-1}$ , the only free parameter remaining is  $y_2$ . Because  $x_i$ ,  $y_i$ , and  $z_i$  can assume only positive values and  $x_2 + y_2 = \eta$ ,  $y_2$  is restricted to values between 0 and  $\eta$ . The spectrum of I<sub>2</sub> can now be calculated from  $(\tilde{\mathbf{C}}^{-1})_2$ ,  $\tilde{\mathbf{B}}$ , and  $\mathbf{A}_p$  using Eq. 3. The results are shown in Fig. 2 *B* for six values of  $y_2$  from 0.12 to 0.37 ( $\sim \eta$ ). Because the extinction coefficient has to be positive, physically meaningful absorption spectra are obtained only for  $y_2 \ge 0.22$ . Of the spectra remaining in Fig. 2 *B*, we pick the one associated with  $y_2 = 0.22$  because it has the smallest spectral bandwidth. For  $y_2$  considerably larger than 0.22, the spectral bandwidth becomes much larger than for P and I<sub>1</sub>, and a secondary absorption maximum develops near 460 nm. This contradicts the original assumption that the UV transitions of I<sub>2</sub> and I<sub>2</sub>' are the longest-wavelength transitions of these intermediates, which precludes transitions at higher wavelengths. The spectrum of  $I_2$  for  $y_2 = 0.22$  is redrawn in Fig. 2 A (open circles). Its  $\lambda_{\text{max}}$  value is at  $\sim$  370  $\pm$ 5 nm. We note that the value of  $\lambda_{max}$  is independent of the choice of  $y_2$ .

With the spectra of I<sub>1</sub>, I<sub>2</sub>, and I<sub>2</sub>', the time courses of the intermediates were calculated from the experimental  $\Delta A(\lambda, t)$  data by matrix inversion of Eq. 1. The time dependence of the relative concentrations of the I<sub>1</sub>, I<sub>2</sub>, and I<sub>2</sub>' intermediates at pH 7 are shown in Fig. 2 *C*. I<sub>1</sub> partially decays to I<sub>2</sub> in 270  $\mu$ s. I<sub>1</sub> and I<sub>2</sub> then further decay around 2 ms to an I<sub>1</sub>/I<sub>2</sub>/I<sub>2</sub>' equilibrium. This equilibrium finally decays to *P* in 260 ms. Also shown is the sum of the relative concentrations of these intermediates (*dash-dot line*). To a good approximation, this sum is constant over the entire time range before the decay to *P*, validating the data analysis. Its value is very close to  $\eta = 0.371$ , the fraction cycling, showing the internal consistency of the analysis.

#### pH dependence of photocycle kinetics

To learn more about the nature of the transition between the acid and the neutral pH regimes, the photocycle kinetics were measured at the following 15 pH values: 4.6, 4.8, 5.1, 5.4, 5.7, 6.0, 6.3, 6.6, 6.75, 6.9, 7.35, 7.7, 7.9, 8.1, 8.4. With excitation at 430 nm, time traces were collected at the seven wavelengths

340, 370, 390, 410, 450, 490, and 500 nm over the time range from 50 ns to 50 s. Results for selected wavelengths are shown in Fig. 3. Note that the panels of Fig. 3 have very different vertical scales and correspondingly different signal/noise ratios. The smallest pH-induced absorbance changes are at 500 nm. The initial absorbance change is almost pH independent at every wavelength, suggesting that the amount of  $I_1$ formed is independent of pH in this range. At each pH, the absorbance changes at all wavelengths could be fitted simultaneously with a sum of three exponentials. The first time constant was virtually constant in this pH range, varying between 200 and 350  $\mu$ s. The second time constant varied between 1.3 (pH 8.4) and 10.6 (pH 5.1) ms. As we saw above, the first transition results the decay of  $I_1$  to  $I_2$ , and the second transition is the decay of  $I_1/I_2$  to the  $I_1/I_2/I_2'$  equilibrium. The data of Fig. 3 show that the third time constant, the return of the  $I_1/I_2/$  $I_2'$  equilibrium to P, is strongly pH dependent, slowing down with decreasing pH.

Some preliminary conclusions on the pH dependence of the I<sub>1</sub>, I<sub>2</sub>, I<sub>2</sub>' intermediate populations may be drawn by inspection of these data. At 340 nm, the extinction coefficient of I<sub>2</sub>' is larger than that of I<sub>2</sub> (Fig. 2 *A*). Although the sequence of time traces is not entirely regular, the absorbance at 340 nm around 10 ms (Fig. 3 *A*) seems to increase with pH, suggesting an increase in the relative amount of I<sub>2</sub>'. At 370 nm, the extinction coefficient of I<sub>2</sub> is larger than that of I<sub>2</sub>' (Fig. 2 *A*). The decrease in absorbance at 370 nm with pH in the millisecond time range (panel B of Fig. 3) may thus be interpreted as a decrease in the relative amount of I2. At 390 nm, the difference in extinction coefficient between  $I_2$  and  $I_2'$ is even larger. This wavelength is therefore diagnostic for the  $I_2$  and  $I_2'$  transition and for pH effects on this transition. Panel C of Fig. 3 shows the increase in absorbance below 1 ms caused by the  $I_1$ -to- $I_2$  transition (the extinction coefficient of  $I_2$  is larger than that of  $I_1$  at this wavelength, see Fig. 2 A). The amount of  $I_2$  formed apparently decreases with increasing pH, judging from the amplitude of the absorption change below 1 ms in Fig. 3 C. Around 2-3 ms, there is a large decrease in absorbance caused by the I2-to-I2' transition. The amplitude of this transition increases with pH, suggesting that more  $I_2'$  is formed at alkaline pH. Panel D shows time traces at 410 nm. This wavelength is appropriate for monitoring the I<sub>1</sub>' intermediate, which occurs at alkaline pH (25). These traces indicate that this intermediate is absent in this pH range. The traces at 500 nm (panel F) are characteristic of I<sub>1</sub>. They suggest that, with increasing pH, the  $I_1$ -to- $I_2$  transition slows down somewhat and that more  $I_1$ remains after the I<sub>1</sub>/I<sub>2</sub>-to-I<sub>2</sub>' transition. This is also supported by the traces at 450 nm (panel E) indicating that the ground state depletion decreases with pH. In the next section these qualitative conclusions about the pH dependence of the intermediate populations, which were drawn from the data themselves in a model-independent way without any assumptions, are confirmed by a quantitative analysis.



FIGURE 3 pH dependence of the transient absorbance changes after excitation at 430 nm at various wavelengths: (A) 340 nm (characteristic for  $I_2'$ ), (B) 370 nm (characteristic for  $I_2$ ), (C) 390 nm (characteristic for I<sub>2</sub>), (D) 410 nm (characteristic for  $I_1'$ ), (E) 450 nm (characteristic for P), and (F) 500 nm (characteristic for  $I_1$ ). The color codes for the pH values in each panel are the following: black, pH 8.4; red, pH 6.9; green, pH 6.6; blue, pH 6.0; light blue, pH 5.7; pink, pH 5.4; dark blue, pH 4.8. Conditions: 50 mM MES, 50 mM KCl, 20°C. PYP concentration 53 µM.

To obtain the time courses of the intermediate populations, it was assumed that the spectra of  $I_1$ ,  $I_2$ , and  $I_2'$  of Fig. 2 A are pH independent and that no other intermediates contribute in the pH range from 4.6 to 8.4. Equation 1 was then used to calculate the time traces  $n_i(t)$  for each intermediate at each pH value from the absorbance changes  $\Delta A(\lambda, t)$  and the spectra  $A_i(\lambda)$  by matrix inversion. The time dependencies of the populations of  $I_1$ ,  $I_2$ , and  $I_2'$  at 7 of the 15 pH values are shown in panels A, B, and C of Fig. 4. They confirm what was suggested by the data of Fig. 3:  $I_1$  decays partially to  $I_2$ ;  $I_2$  then partially decays to  $I_2'$ ; beyond 10 ms  $I_1$ ,  $I_2$ , and  $I_2'$  are in equilibrium and return together to P. Fig. 4 D shows that the sum of the populations is approximately constant in time and equal to the fraction cycling. Whereas the population of  $I_1$  in equilibrium with  $I_2$  and  $I_2'$  is only slightly pH dependent (see traces in Fig. 4 A around 10 ms), the concentrations of I<sub>2</sub> and I<sub>2</sub>' show a strong and opposite pH dependence. With increasing pH, the amount of I2' increases at the expense of a corresponding decrease in the I<sub>2</sub> population. To quantify this pH dependence of the equilibrium populations, their concentrations at 10 ms were taken from Fig. 4, B and C (vertical dashed lines). The corresponding concentrations of  $I_2$  and  $I_2'$  are plotted in Fig. 5 A as a function of pH. The solid curves are the results of a simultaneous fit with the Henderson-Hasselbalch equation. The fit parameters were  $pK_a \approx 6.4$  and  $n \approx 0.98$ .  $I_2$  and  $I_2'$  are thus in a pH-dependent equilibrium.

The ground state recovery slows down at acid pH. The rate constant  $k_3$  for this recovery was determined from a global fit of the data of Fig. 3 at all seven wavelengths. Its pH dependence is plotted in Fig. 5 *B*. A fit with the Henderson-Hasselbalch equation results in a pK<sub>a</sub> of ~6.3 and *n* of ~0.8. The apparent decay rate for the dark-state recovery thus seems to be proportional to the I<sub>2</sub>' population in the I<sub>1</sub>/I<sub>2</sub> /I<sub>2</sub>' equilibrium. This relationship is expected in the framework of a simple model presented in the discussion.

#### pH dependence of photoreversal kinetics

Recently, we measured and analyzed the kinetics of photo reversal from  $I_2$ ,  $I_2'$ , and  $I_1$  at pH 6 in detail (33). The time delay between the two flashes was varied from 1  $\mu$ s to 3 s, and the photoreversal kinetics were measured at 26 wavelengths from 330 to 510 nm (33). The photoreversal time traces at pH 6 required two exponentials for an adequate fit with well-separated time constants of  $\tau_1 = 60$  and  $\tau_2 = 400$  $\mu$ s. These time constants were assigned to photoreversal from  $I_2$  (60  $\mu$ s) and  $I_2'$  (400  $\mu$ s), respectively, on the basis of the delay dependence of the amplitudes (33). Moreover, we found that  $I_2$  and  $I_2'$  are in equilibrium (33). Here, our focus is on the pH dependence. The photoreversal kinetics were therefore measured at only two wavelengths (340 and 450 nm) and at the fixed delay of 20 ms. At this time delay all three intermediates  $I_2$ ,  $I_2'$ , and  $I_1$  are present and in equilibrium (Figs. 2 C and 4).



FIGURE 4 Time courses of the relative concentrations of the  $I_1(A)$ ,  $I_2(B)$ , and  $I_{2'}(C)$  intermediates at various pH values calculated using Eq. 1 as explained in the text. (*D*) Time course of the sum of the populations of  $I_1$ ,  $I_2$ , and  $I_{2'}$ . Color code as in Fig. 3.

The photoreversal signals at pH 5.1 and 8.1 are shown in Fig. 6 *A*. For clarity, only the traces at 2 of the 15 pH values are shown together with their simultaneous fits (*solid lines*). As at pH 6 (33), the kinetics required two exponentials over the pH range from 4.6 to 6.9. These two phases can be clearly discerned in the data at pH 5.1. From pH 7.3 on, only the



FIGURE 5 (A) pH dependence of the equilibrium concentrations of the I<sub>2</sub> ( $\bigcirc$ ) and I<sub>2</sub>' ( $\bullet$ ) intermediates at 10 ms derived from Fig. 4, *B* and *C*, respectively. The solid curves are simultaneous fits of these titration curves with the Henderson-Hasselbalch equation with pK<sub>a</sub> = 6.4 and  $n \approx 0.98$ . (*B*) pH dependence of the decay rate  $k_3$  of the ground-state recovery. For every pH, the decay rate  $k_3$  was derived from the simultaneous fit of the measured transient absorbance changes at 340 nm, 370 nm, 390 nm, 410 nm, 450 nm, 490 nm, and 500 nm with a sum of three exponentials. The solid curve in *B* is the fit of the decay rate  $k_3$  with the Henderson-Hasselbalch equation with pK<sub>a</sub> = 6.3,  $n \approx 0.84$ .

slow component was required. The fast time constant was virtually pH independent, varying between 48 and 63  $\mu$ s. The second time constant varied between 330 and 770  $\mu$ s in the pH range investigated. Comparison of the time traces in Fig. 6 A makes it clear that there are significant differences between pH 5.1 and 8.1. The total initial photoreversal signal is somewhat larger at pH 5.1 than at 8.1, suggesting that more  $I_2/I_2'$  can be photoreversed. The amplitude data provide further insight. The pH dependence of the amplitudes  $A_1$ and  $A_2$  (for the corresponding time constants  $\tau_1$  and  $\tau_2$ ) are plotted in Fig. 6, B (340 nm) and C (450 nm). These figures confirm that, with increasing pH,  $A_1$  becomes smaller, approaching zero around pH 7 at both wavelengths, whereas  $A_2$  shows a corresponding increase. The solid curves are simultaneous fits of the amplitudes at 340 and 450 nm with the Henderson-Hasselbalch equation, with  $pK_a = 6.1$  and n =1.9. We assigned  $A_1$  and  $A_2$  to photoreversal from  $I_2$  and  $I_2'$ , respectively (33). The results of Fig. 6, B and C thus suggest that, with increasing pH, the  $I_2/I_2'$  equilibrium shifts from  $I_2$  at pH 4.6 to  $I_2'$  at pH 8.4. These results, from the pH dependence of the photoreversal amplitudes, thus support our observations from the photocycle, where a pK<sub>a</sub> of 6.4 was obtained. We note that in the photocycle experiments, n = 0.98 (I<sub>2</sub>/I<sub>2</sub>' equilibrium) or 0.8 (recovery rate), whereas we obtained n = 1.9 from the photoreversal kinetics. The photoreversal absorbance changes required various corrections and are, moreover, quite small. The errors are correspondingly large. The photocycle data points, on the other hand, display much less scattering, in particular for  $k_3$ , and did



FIGURE 6 (*A*) Photoreversal signals, at 340 nm (*positive*) and 450 nm (*negative*) at pH 5.1 (*black*) and 8.1 (*red*), calculated as described in the text. For clarity, the data obtained at 13 other pH values ranging from 4.6 to 8.4 are not shown. The solid curves represent a simultaneous exponential fit to the 340-nm and 450-nm traces. Conditions: 20°C, 50 mM KCl and 50 mM MES, PYP concentration 53  $\mu$ M. The delay between first (*blue*, 430 nm) and second (*violet*, 355 nm) flashes is 20 ms. (*B* and *C*) pH dependence of the photoreversal amplitudes of the fast (48–60  $\mu$ s) and slow (300–770  $\mu$ s) components, respectively, obtained from the simultaneous fit of the 340-nm and 450-nm traces of panel *A*. The solid curves of panels *B* and *C* represent a common fit to the pH dependence of all four amplitudes with the Henderson-Hasselbalch equation, with a pK<sub>a</sub> of 6.1 (*dashed vertical lines*) and a Hill coefficient of 1.9.

not require complex corrections. We believe, therefore, that the number of protons involved is one in both experiments.

## DISCUSSION

From measurements of the pH dependence of the photocycle and photoreversal kinetics of PYP between 4.6 and 8.4, we obtained the following results: 1) the spectra of the signaling state  $I_2$  (350 nm) and its precursor  $I_2$  (370 nm) differ by ~20 nm; 2) from several milliseconds (formation of  $I_2$ ) to the end of the cycle, the three intermediates  $I_1$ ,  $I_2$ , and  $I_2'$  are in equilibrium; and 3) the pK<sub>a</sub> of the pH-dependent equilibrium between  $I_2'$  and its precursor  $I_2$  is ~6.4 from photocycle kinetics and  $\sim 6.1$  from photoreversal kinetics. The pH is thus an important parameter that controls the amount of receptor in the active state. This is analogous to the case of the photoreceptor rhodopsin, where the equilibrium between the signaling state M<sub>II</sub> and its precursor M<sub>I</sub> is also strongly pH dependent (e.g., (43)). We recently showed (26) that the  $I_2/I_2$ equilibrium, like the M<sub>I</sub>/M<sub>II</sub> equilibrium (44), also depends on the salt concentration. In both photoreceptors, high salt favors the signaling state. The proposed reaction scheme for the kinetics and equilibria of the photocycle in this pH range is presented in Fig. 7.

The existence of two distinguishable  $I_2$  intermediates with protonated chromophores was first demonstrated by timeresolved FTIR (12,18).  $I_2$  decays to  $I_2'$  in ~2–3 ms (12,18). This transition to the signaling state  $I_2'$  is characterized by a global conformational change (12,18). Here, we showed that this transition may also be monitored by transient electronic absorption spectroscopy and determined the absorption



FIGURE 7 Proposed model for the kinetics of the photocycle and photoreversal of PYP in the pH range from 4 to 8. The  $I_2^{cis}$  and  $I_2^{rcis}$  intermediates are in a pH-dependent equilibrium and photoreverse to  $P^{trans}$  with exponential time constants of 57 and 380  $\mu$ s. The pK<sub>a</sub> of the  $I_2^{cis}/I_2^{rcis}$  equilibrium is 6.4. Note the equilibria among  $I_1$ ,  $I_2$ , and  $I_2'$ . These three intermediates decay together to P. For clarity the short-lived intermediates  $I_0$  and  $I_0^{\ddagger}$  between P\* and  $I_1$  are not shown. The proton arrows indicate the transient uptake and release of protons as detected by a pH-sensitive dye in solution (13).

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spectra of I<sub>2</sub> and I<sub>2</sub>' by the extrapolated difference method (25,37). The  $\lambda_{max}$  values of I<sub>2</sub> and I<sub>2</sub>' are 370 ± 5 and 350 ± 5 nm, respectively. Absorption spectra for I<sub>2</sub> and I<sub>2</sub>' at pH 8.1 were presented by Hendriks et al. (11). No  $\lambda_{max}$  values were provided, but the spectrum of I<sub>2</sub> was said to be "slightly red-shifted" with respect to I<sub>2</sub>' in agreement with our results. The spectrum of I<sub>2</sub> presented by Hendriks et al. (11) is so noisy that it is difficult to estimate  $\lambda_{max}$ . The poor quality of this spectrum is probably caused by the fact that at pH 8.1 the contribution of I<sub>2</sub> in the equilibrium is very low, as we showed here (Fig. 5 *A*).

Using the spectra of  $I_2$  and  $I_2'$  from Fig. 2 *A* and assuming that they are pH independent in the pH range from 4.6 to 8.4, we obtained the time dependence of the concentrations of  $I_1$ ,  $I_2$ , and  $I_2'$  (Fig. 4). From ~5 ms onward,  $I_1$ ,  $I_2$ , and  $I_2'$  are in equilibrium and decay together to the initial dark state P. The equilibrium intermediate populations of  $I_2$  and  $I_2'$  are pH dependent with a pK<sub>a</sub> of ~6.4. Below the pK<sub>a</sub>,  $I_2$  is the major species; above the pK<sub>a</sub> the opposite holds. We note that these transitions are not complete (Fig. 5 *A*).

The sum of the populations of the  $I_1$ ,  $I_2$ , and  $I_2'$  intermediates in Fig. 4 D is not as constant in time for every pH as it should be and as it is at pH 7 (Fig. 2 C). One cause could be that the intermediate spectra are not exactly pH independent over the whole pH range, so that the spectra derived at pH 7 are not quite correct for every pH. So we derived a second set of intermediate spectra from a joint SVD analysis over the whole pH range as in Joshi et al. (25). These are the best pH-averaged intermediate spectra. The differences with the set at pH 7 (Fig. 2 A) are small. A third set was obtained from the scaled subtraction method (25). All three sets of intermediate spectra lead to the same pK<sub>a</sub> value. The results differed only somewhat in terms of the end values of the titrations at high pH. The sum of the populations remained slightly time-dependent, however, and other inadequacies developed, such as slightly negative populations. Because the spectra of I2 and I2' are quite similar, the analysis is sensitive to the small differences between them. We believe this is the likely source of the error in Fig. 4 D.

Measurements of the pH dependence of the photoreversal kinetics of  $I_2$  and  $I_2'$  at a delay of 20 ms (Fig. 6) provide further support for this  $pK_a$ . The kinetics are characterized by a fast (50–60  $\mu$ s) and a slow (300–800  $\mu$ s) component, which result from photoreversal from I2 and I2', respectively (33). The corresponding amplitudes are pH dependent (Fig. 6, B and C) with a pK<sub>a</sub> of 6.1, in good agreement with the value of 6.4 obtained from the photocycle kinetics. The amplitude for I<sub>2</sub> goes to zero beyond pH 7.3. The photoreversal data thus suggest that no I<sub>2</sub> remains at alkaline pH, whereas the photocycle data indicate that some  $I_2$  remains because  $[I_2]/([I_2] + [I_2'])$  is ~0.25 at pH 8.6 (Fig. 5 A). We note that the end value of the I<sub>2</sub> population at high pH in Fig. 5 A is quite sensitive to the exact choice of the  $I_1$  spectrum. In a recent investigation on the intermediates and spectra at alkaline pH (25), we found that only three intermediates,  $I_1$ ,  $I_1'$ , and  $I_2'$ , are present above pH 8, i.e., no  $I_2$ .

The pH dependence of the absorption spectrum of a photostationary mixture of P, I2, and I2' produced by background illumination was recently analyzed by SVD (34). It was shown that  $I_2$  and  $I_2'$  are in a pH-dependent equilibrium with a pK<sub>a</sub> of 6.3 and that the spectrum of the high-pH species  $I_2'$  is blue-shifted with respect to that of the low-pH species  $I_2$  (34). The  $I_2$  and  ${I_2}^\prime$  intermediates also differ with regard to the fluorescence lifetime of the single tryptophan of PYP, W-119 (34). In  $I_2$ , the lifetime is long (0.82 ns). In  $I_2'$ , the lifetime is much shorter (0.04 ns). Using background illumination, these authors showed that the fluorescence decay in the photostationary state is pH dependent with I2 dominating at low pH and  $I_2'$  at high pH (34). The pK<sub>a</sub> was ~6.3. Combining the pH dependence of the fluorescence amplitudes with that of the photostationary absorption, absorption spectra were calculated for the  $I_2$  and  $I_2'$  species. These had  $\lambda_{max}$  values of 372 and 352 nm for  $I_2$  and  ${I_2}'$  , respectively (34), in good agreement with the results reported in this work.

The results on the pH dependence of the  $I_2/I_2'$  equilibrium (34) were recently confirmed (29). Analysis of the photostationary absorption spectra by a scaled subtraction procedure yielded a pK<sub>a</sub> of 6.4 and  $\lambda_{max}$  values of 367 and 356 nm for  $I_2$  and  $I_2'$ , respectively (29). These authors showed, moreover, from CD and small-angle x-ray scattering experiments that the global structural transition occurs between these two intermediates with a pK<sub>a</sub> of 6.4. Together with the kinetics results from time-resolved absorption spectroscopy presented here, these complementary methods lead to a comprehensive picture of the  $I_2$ -to- $I_2'$  equilibrium.

As is well known (23) and confirmed here (Fig. 5 C), the rate  $k_3$ , for the ground state recovery, is also pH dependent with a pK<sub>a</sub> of 6.3, i.e., within experimental error equal to that for the  $I_2/I_2'$  equilibrium (6.4 and 6.1, from single flash photolysis and photoreversal measurements, respectively). Thus,  $k_3$  seems to be proportional to the  $I_2'$  population (compare Fig. 5, A and B). Such a proportionality is expected under the following conditions: 1) the equilibration rates among  $I_1$ ,  $I_2$ , and  $I_2'$  are rapid compared to the microscopic rates of return from each intermediate to the ground state; 2) the latter are pH independent; and 3) the rate from  $I_2'$  to P is much larger than from the other intermediates. A similar model was recently proposed to explain the pH dependence of the rate of ground state recovery at alkaline pH (25). At alkaline pH, the  $I_1$ ,  $I_1'$ , and  $I_2'$  intermediates are in equilibrium. The pK<sub>a</sub> of the  $I_1'$ -to- $I_2'$  equilibrium is ~9.9, and the ground state recovery rate constant  $k_3$  has a pK<sub>a</sub> of 9.7 and is proportional to the  $I_2'$  population. There is thus a striking similarity between the behavior at high and low pH. For both branches of the bell-shaped pH dependence of  $k_3$ , it seems that  $k_3$  is proportional to  $[I_2']$ . This proportionality is not exact however, because  $k_3$  approaches zero at low pH, whereas  $[I_2']$ approaches a constant value unequal to zero. A more detailed model thus seems to be required. Nevertheless, this symmetry between low and high pH behavior is worth pointing out, and the underlying model provides a lowest-order explanation.

An important question concerns the group responsible for the pK<sub>a</sub> of ~6.4. The similar pK<sub>a</sub> for the recovery rate  $k_3$  is commonly attributed to the carboxyl group of E-46 (35,36). We now need to discuss this pK<sub>a</sub> in the context of the underlying  $I_2/I_2'$  equilibrium. What is the mechanism whereby a change in protonation of E-46 shifts the equilibrium from  $I_2$ to  $I_2'$ . In  $I_2$  the chromophore is already protonated and has moved away from E-46 toward the surface (6). If E-46 is the internal proton donor for the chromophore, its carboxyl group is presumably already deprotonated in I2, in accordance with some observations from time-resolved FTIR (12). In  $I_2'$  the chromophore remains protonated, but the protein structure is changed in a major way. It is unclear how the deprotonated E-46 could affect the  $I_2/I_2'$  conformational equilibrium. If, however, E-46 is not the internal proton donor, and the chromophore is protonated from the external medium, as suggested (13), E-46 could remain protonated in  $I_2$  and be deprotonated in  $I_2'$ . In other time-resolved FTIR measurements (18), a positive band was observed at 1759  $cm^{-1}$  with a risetime of 113  $\mu$ s (formation of I<sub>2</sub>) and assigned to an environmental shift of the protonated E-46. Brudler et al. (18) were not aware of the  $I_2/I_2'$  equilibrium and concluded from the fact that the amplitude of the positive band was significantly smaller than that of the negative band due to the initial dark state, that only a fraction of the molecules cycling had a protonated E-46 in I2. Their experiments were, however, performed in buffer at pH 7. At this pH, the  $I_2/I_2'$  equilibrium is far on the side of  $I_2'$  (see Fig. 5 A), so that only a minority of molecules would have been in the I<sub>2</sub> state. It is thus consistent with the time-resolved data of (18) to conclude that in  $I_2$  the carboxyl group of E-46 is protonated. Recent photostationary FTIR measurements also indicate that E-46 is at least partially protonated in I<sub>2</sub> (29). A role of E-46 in controlling the  $I_2/I_2$ equilibrium is thus plausible. In the absence of the carboxyl group, in the mutant E46Q, the conformational change in  $I_2'$  is much smaller or absent (12), and the absorption maximum at pH 7 is at 368 nm (45), i.e., I<sub>2</sub>-like. These results suggest that in the absence of E-46 the  $I_2/I_2'$  equilibrium is predominantly or entirely on the side of I2 and further support the idea that E-46 is responsible for the wild-type  $pK_a$  of 6.4. We note that this reinterpretation of the FTIR results is consistent with a mechanism of chromophore protonation from the external medium (13).

Another residue that might be involved is H-108. The pH dependence of the steady-state proton uptake was investigated (46), and from the observed pH dependence, a pK<sub>a</sub> of 6.6 was obtained, which was attributed to histidine-108. This residue is located on the central  $\beta$ -scaffold and may be involved in the interaction between the  $\beta$ -scaffold and the N-terminal domain. It was recently postulated, on the basis of the observation that the I<sub>2</sub>-to-I<sub>2</sub>' transition is blocked at low salt concentrations, that the loss of this interaction is a prerequisite for the formation of I<sub>2</sub>' (26).

Recently two forms of  $I_1$  could be distinguished on the basis of their resonance Raman spectra, which are in a

pH-dependent equilibrium with a pK<sub>a</sub> of ~6.2 (30). The lowpH form  $(I_1^{\ell})$  lacks the hydrogen bond with E-46, whereas the high-pH form  $(I_1^h)$  has both hydrogen bonds. It is possible that the pH dependence observed here for the  $I_2/I_2'$  equilibrium results from the pH dependence of the preceding  $I_1^{\ell}/I_1^h$  equilibrium.

In conclusion, we have shown from measurements of the time-dependent intermediate populations that the equilibrium between  $I_2$  and  $I_2'$  and the formation of the signaling state  $I_2'$  are pH dependent with a pK<sub>a</sub> of ~6.4. The intracellular pH may thus regulate the amount of active receptor. The value of this pK<sub>a</sub> provides an explanation for the similar well-known pK<sub>a</sub> of the rate constant for the ground-state recovery. We find, moreover, that  $I_2'$  is blue-shifted with respect to  $I_2$  by ~20 nm, suggesting a different chromophore environment for the exposed chromophore in the signaling state.

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