# Von Neumann-Jordan Constant of Absolute Normalized Norms on $\mathbb{C}^{2}$ 

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We determine and estimate the von Neumann-Jordan constant of absolute normalized norms on $\mathbb{C}^{2}$ by means of their corresponding continuous convex functions on $[0,1]$. This provides many new interesting examples including those of non- $\ell_{p}$-type as well as some previous ones. It is also shown that all such norms are uniformly non-square except $\ell_{1}$ - and $\ell_{\infty}$-norms. © 2000 Academic Press

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## 1. INTRODUCTION

The notion of the von Neumann-Jordan constant of Banach spaces (hereafter referred to as NJ constant) was introduced by Clarkson in [5]

[^0]and recently it has been studied by several authors (cf. [5, 6, 8-11, 13], etc.). The NJ constant $C_{\mathrm{NJ}}(X)$ of a Banach space $X$ is the smallest constant $C$ for which
$$
\frac{1}{C} \leq \frac{\|x+y\|^{2}+\|x-y\|^{2}}{2\left(\|x\|^{2}+\|y\|^{2}\right)} \leq C
$$
holds for all $x, y \in X$, not both 0 . From Jordan and von Neumann [6], we have $1 \leq C_{\mathrm{NJ}}(X) \leq 2$ for any Banach space $X$, and $X$ is a Hilbert space if and only if $C_{\mathrm{NJ}}(X)=1$. Clarkson [5] calculated the NJ constant of $L_{p}$ by using Clarkson's inequalities. Recently, following his way, it was determined for a sequence of other Banach spaces such as $L_{p}\left(L_{q}\right)\left(L_{q}\right.$-valued $L_{p}$-space), $W_{p}^{k}(\Omega)$ (Sobolev space), $c_{p}$ (Schatten $p$-class operators), and so on ( $[9,11]$, etc.). On the other hand, thanks to the NJ constant we can describe some geometrical and topological structures of Banach spaces. For example, the second and third authors proved that $C_{\mathrm{NJ}}(X)<2$ if and only if $X$ is uniformly non-square, whence $X$ is super-reflexive if and only if $X$ admits an equivalent norm with NJ constant less than 2 ( $[10,13]$ ). We also have that, if $C_{\mathrm{NJ}}(X)<5 / 4$, then the Banach space $X$ has the fixed point property for nonexpansive mappings (cf. [8]). For some other results concerning Rademacher type and cotype we refer the reader to [10].

A norm $\|\cdot\|$ on $\mathbb{C}^{2}$ is said to be absolute if $\|(z, w)\|=\|(|z|,|w|)\|$ for all $z, w \in \mathbb{C}$ and normalized if $\|(1,0)\|=\|(0,1)\|=1$. Let $N_{a}$ denote the family of all absolute normalized norms on $\mathbb{C}^{2}$, and let $\Psi$ denote the family of all continuous convex functions on $[0,1]$ such that $\psi(0)=\psi(1)=1$ and $\max \{1-t, t\} \leq \psi(t) \leq 1(0 \leq t \leq 1)$. Then as in Bonsall and Duncan [3, Section 21, Lemma 3], $N_{a}$ and $\Psi$ are in one-to-one correspondence under the equation $\psi(t)=\|(1-t, t)\|(0 \leq t \leq 1)$. In particular, owing to this we can consider many non- $l_{p}$-type norms easily.

In this paper, we shall determine and estimate the NJ constant of absolute normalized norms on $\mathbb{C}^{2}$ by means of the corresponding convex functions. In particular, this provides a new way to calculate it with no use of Clarkson's inequalities. The main results are stated as follows. Let $\|\cdot\|_{\psi}$ be an absolute normalized norm associated with a convex function $\psi \in \Psi$. Let $M_{1}=\max _{0 \leq t \leq 1} \psi(t) / \psi_{2}(t)$ and $M_{2}=\max _{0 \leq t \leq 1} \psi_{2}(t) / \psi(t)$, respectively, where $\psi_{2}(t):=\left\{(1-t)^{2}+t^{2}\right\}^{1 / 2}$ corresponds to the $\ell_{2}$-norm. First we show that, if $\psi \geq \psi_{2}$ (resp. $\psi \leq \psi_{2}$ ), then $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=M_{1}^{2}$ (resp. $M_{2}^{2}$ ) (Theorem 1). In general, we prove that $\max \left\{M_{1}^{2}, M_{2}^{2}\right\} \leq C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \leq$ $M_{1}^{2} M_{2}^{2}$ (Theorem 2). Theorem 1 gives a class of convex functions for which $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=\max \left\{M_{1}^{2}, M_{2}^{2}\right\}$. We further present a sufficient condition that $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=M_{1}^{2} M_{2}^{2}\left(\max \left\{M_{1}^{2}, M_{2}^{2}\right\}<C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)\right)$ (Theorem 3) and that $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)<M_{1}^{2} M_{2}^{2}$ (Theorem 4), respectively. These results enable us to
present many new interesting examples, especially those of non- $\ell_{p}$-type, for instance, $\ell_{p, 2}$-norm, $p \geq 2$ (Lorentz norm), etc. As a corollary we show that all absolute normalized norms are uniformly non-square except $\ell_{1}$ and $\ell_{\infty}$-norms.

## 2. ABSOLUTE NORMALIZED NORMS ON $\mathbb{C}^{2}$

A norm $\|\cdot\|$ on $\mathbb{C}^{2}$ is said to be absolute if

$$
\|(z, w)\|=\|(|z|,|w|)\| \quad \text { for all } z, w \in \mathbb{C}
$$

and normalized if $\|(1,0)\|=\|(0,1)\|=1$. The $\ell_{p}$-norms $\|\cdot\|_{p}(1 \leq p \leq \infty)$ are basic examples,

$$
\|(z, w)\|_{p}= \begin{cases}\left(|z|^{p}+|w|^{p}\right)^{1 / p} & \text { if } 1 \leq p<\infty \\ \max (|z|,|w|) & \text { if } p=\infty\end{cases}
$$

Let $N_{a}$ denote the family of absolute normalized norms on $\mathbb{C}^{2}$. We recall some basic facts about these norms; for the convenience of the reader we give their proofs following Bonsall and Duncan [3].

Lemma 1 ([3, p. 36]). For any norm $\|\cdot\| \in N_{a}$

$$
\begin{equation*}
\|\cdot\|_{\infty} \leq\|\cdot\| \leq\|\cdot\|_{1} . \tag{1}
\end{equation*}
$$

Indeed, for any $z, w \in \mathbb{C}$

$$
\begin{aligned}
\|(z, w)\|_{\infty} & =\max \{\|(z, 0)\|,\|(0, w)\|\} \\
& =\frac{1}{2} \max \{\|(z, w)+(z,-w)\|,\|(z, w)+(-z, w)\|\} \\
& \leq \frac{1}{2} \max \{\|(z, w)\|+\|(z,-w)\|,\|(z, w)\|+\|(-z, w)\|\} \\
& =\|(z, w)\| \\
& \leq\|(z, 0)\|+\|(0, w)\| \\
& =\|(z, w)\|_{1} .
\end{aligned}
$$

Now let $\Psi$ denote the family of all continuous convex functions $\psi$ on $[0,1]$ with $\psi(0)=\psi(1)=1$ satisfying

$$
\begin{equation*}
\max \{1-t, t\} \leq \psi(t) \leq 1 \quad(0 \leq t \leq 1) \tag{2}
\end{equation*}
$$

Then $N_{a}$ and $\Psi$ are in one-to-one correspondence as follows.
Lemma 2 ([3, p. 37]). (i) Let $\|\cdot\| \in N_{a}$ and let

$$
\begin{equation*}
\psi(t)=\|(1-t, t)\| \quad(0 \leq t \leq 1) . \tag{3}
\end{equation*}
$$

Then $\psi \in \Psi$ : Conversely,
(ii) For a given $\psi \in \Psi$ define

$$
\|(z, w)\|_{\psi}= \begin{cases}(|z|+|w|) \psi\left(\frac{|w|}{|z|+|w|}\right) & \text { if }(z, w) \neq(0,0)  \tag{4}\\ 0 & \text { if }(z, w)=(0,0)\end{cases}
$$

Then $\|\cdot\|_{\psi} \in N_{a}$, and $\|\cdot\|_{\psi}$ satisfies (3).
Proof. (i) This is easy to see ( $\psi$ satisfies (2) by Lemma 1).
(ii) Let $\psi \in \Psi$. We only show the triangle inequality. Let us first see that

$$
\begin{equation*}
\|(p, q)\|_{\psi} \leq\|(r, s)\|_{\psi} \quad \text { if } 0 \leq p \leq r, 0 \leq q \leq s . \tag{5}
\end{equation*}
$$

When $p=0$ or $q=0$, (5) is clear (recall Lemma 1). Thus it is enough to show that

$$
\begin{equation*}
(p+q) \psi\left(\frac{q}{p+q}\right) \leq(r+s) \psi\left(\frac{s}{r+s}\right) \quad \text { if } 0<p \leq r, 0<q \leq s \tag{6}
\end{equation*}
$$

Since $\psi$ is convex and $\psi(t) \geq t$, the function $\psi(t) / t$ is non-increasing. Indeed, let $0<s<t \leq 1$. Then

$$
\psi(t)=\psi\left(\frac{1-t}{1-s} s+\frac{t-s}{1-s} 1\right) \leq \frac{1-t}{1-s} \psi(s)+\frac{t-s}{1-s} \psi(1) .
$$

Hence

$$
\begin{aligned}
\frac{\psi(s)}{s}-\frac{\psi(t)}{t} & \geq \frac{\psi(s)}{s}-\frac{1}{t}\left\{\frac{1-t}{1-s} \psi(s)+\frac{t-s}{1-s}\right\} \\
& \geq \psi(s)\left\{\frac{1}{s}-\frac{1-t}{t(1-s)}\right\}-\frac{t-s}{t(1-s)} \\
& \geq s\left\{\frac{1}{s}-\frac{1-t}{t(1-s)}\right\}-\frac{t-s}{t(1-s)}=0 .
\end{aligned}
$$

Therefore we have

$$
\begin{equation*}
(p+q) \psi\left(\frac{q}{p+q}\right) \leq(r+q) \psi\left(\frac{q}{r+q}\right) . \tag{7}
\end{equation*}
$$

In the same way the function $\psi(t) /(1-t)$ is non-decreasing (use $\psi(t) \geq$ $1-t$ in this case), which implies

$$
\begin{equation*}
(r+q) \psi\left(\frac{q}{r+q}\right) \leq(r+s) \psi\left(\frac{s}{r+s}\right) \tag{8}
\end{equation*}
$$

Combining (7) and (8), we have (6). Now let $(u, v),(z, w) \in \mathbb{C}^{2}$. Then by (5)

$$
\begin{aligned}
\|(u, v)+(z, w)\|_{\psi} & =\|(|u+z|,|v+w|)\|_{\psi} \\
& \leq\|(|u|+|z|,|v|+|w|)\|_{\psi} \\
& =(|u|+|z|+|v|+|w|) \psi\left(\frac{|v|+|w|}{|u|+|z|+|v|+|w|}\right)
\end{aligned}
$$

Noting here that

$$
\begin{aligned}
& \frac{|v|+|w|}{|u|+|v|+|z|+|w|} \\
& \quad=\frac{|u|+|v|}{|u|+|v|+|z|+|w|} \cdot \frac{|v|}{|u|+|v|}+\frac{|z|+|w|}{|u|+|v|+|z|+|w|} \cdot \frac{|w|}{|z|+|w|}
\end{aligned}
$$

we have by the convexity of $\psi$

$$
\begin{aligned}
\|(u, v)+(z, w)\|_{\psi} & \leq(|u|+|v|) \psi\left(\frac{|v|}{|u|+|v|}\right)+(|z|+|w|) \psi\left(\frac{|w|}{|z|+|w|}\right) \\
& =\|(u, v)\|_{\psi}=\|(z, w)\|_{\psi}
\end{aligned}
$$

as desired.
Now let $\psi_{p}(t)=\left\{(1-t)^{p}+t^{p}\right\}^{1 / p} \in \Psi$. Then, as is easily seen, the $\ell_{p}$-norm $\|\cdot\|_{p}$ is associated with $\psi_{p} \in N_{a}$. In what follows we write $\varphi \leq \psi$ if $\varphi(t) \leq \psi(t)$ for all $0 \leq t \leq 1$. We shall need the following simple facts later.

Lemma 3. Let $\varphi, \psi \in \Psi$ and let $\varphi \leq \psi$. Put

$$
M=\max _{0 \leq t \leq 1} \frac{\psi(t)}{\varphi(t)}
$$

Then

$$
\|\cdot\|_{\varphi} \leq\|\cdot\|_{\psi} \leq M\|\cdot\|_{\varphi}
$$

Proof. For any $z, w \in \mathbb{C}$

$$
\begin{aligned}
\|(z, w)\|_{\varphi} & =(|z|+|w|) \varphi\left(\frac{|w|}{|z|+|w|}\right) \\
& \leq(|z|+|w|) \psi\left(\frac{|w|}{|z|+|w|}\right) \\
& =\|(z, w)\|_{\psi} \\
& \leq M(|z|+|w|) \varphi\left(\frac{|w|}{|z|+|w|}\right) \\
& =M\|(z, w)\|_{\varphi} .
\end{aligned}
$$

Lemma 4. Let $\varphi, \psi \in \Psi$ and let $1 / 2 \leq \lambda \leq 1$. Then

$$
\max \left\{\|\cdot\|_{\varphi}, \lambda\|\cdot\|_{\psi}\right\}=\|\cdot\|_{\max \{\varphi, \lambda \psi\}} .
$$

Proof. Note first that $\max \{\varphi, \lambda \psi\} \in \Psi$. Then for any $(z, w) \in \mathbb{C}^{2}$

$$
\begin{aligned}
\|(z, w)\|_{\max \{\varphi, \lambda \psi\}} & =(|z|+|w|) \max \left\{\varphi\left(\frac{|w|}{|z|+|w|}\right), \lambda \psi\left(\frac{|w|}{|z|+|w|}\right)\right\} \\
& =\max \left\{\|(z, w)\|_{\varphi}, \lambda\|(z, w)\|_{\psi}\right\} .
\end{aligned}
$$

## 3. NJ CONSTANT OF ABSOLUTE NORMALIZED NORMSTHE COMPARABLE CASE WITH $\psi_{2}$

The von Neumann-Jordan constant of a Banach (or normed) space $X$ ([5]; cf. [12, p. 550]), we denote by $C_{\mathrm{NJ}}(X)$, is the smallest constant $C$ for which

$$
\frac{1}{C} \leq \frac{\|x+y\|^{2}+\|x-y\|^{2}}{2\left(\|x\|^{2}+\|y\|^{2}\right)} \leq C
$$

holds for all $x, y \in X$, not both 0 .
Let us recall some geometrical notions of a Banach space $X$ (cf. [1]). $X$ or its norm $\|\cdot\|$ is called uniformly convex if for any $\varepsilon>0(0<\varepsilon<2)$ there exists a $\delta>0$ such that $\|x-y\| \geq \varepsilon,\|x\| \leq 1,\|y\| \leq 1$ implies $\|(x+y) / 2\| \leq 1-\delta . X$ is called uniformly non-square provided there exists a $\delta>0$ such that if $\|(x-y) / 2\| \geq 1-\delta,\|x\| \leq 1$ and $\|y\| \leq 1$, then $\|(x+y) / 2\| \leq 1-\delta$. Clearly uniformly convex spaces are uniformly non-square, for the converse uniform non-squareness does not even imply
strict convexity, whereas $X$ admits an equivalent uniformly non-square norm if and only if $X$ is uniformly convexifiable (such a Banach space is precisely super-reflexive).

We summarize some basic facts about the NJ constant.
Proposition A. (i) $1 \leq C_{\mathrm{NJ}}(X) \leq 2$ for any Banach space $X ; C_{\mathrm{NJ}}(X)$ $=1$ if and only if $X$ is a Hilbert space (Jordan and von Neumann [6]).
(ii) $C_{\mathrm{NJ}}(X)<2$ if and only if $X$ is uniformly non-square (Takahashi and Kato [13]; see also [10]).
(iii) $\quad C_{\mathrm{NJ}}\left(L_{p}\right)=C_{\mathrm{NJ}}\left(\ell_{p}\right)=2^{(2 / t)-1}$, where $1 \leq p \leq \infty, 1 / p+1 / p^{\prime}=$ 1, and $t=\min \left\{p, p^{\prime}\right\}$ (Clarkson [5]).

For a norm $\|\cdot\|$ on $\mathbb{C}^{2}$ we write $C_{\mathrm{NJ}}(\|\cdot\|)$ for $C_{\mathrm{NJ}}\left(\left(\mathbb{C}^{2},\|\cdot\|\right)\right)$. We first see that the NJ constant is stable under the symmetric transformation of $\psi$ with respect to the line $t=1 / 2$.

Proposition 1. Let $\psi \in \Psi$ and let $\tilde{\psi}(t)=\psi(1-t)$. Then $C_{N J}\left(\|\cdot\|_{\psi}\right)=$ $C_{\mathrm{NJ}}\left(\|\cdot\|_{\tilde{\psi}}\right)$.

Proof. For $x=(z, w) \in \mathbb{C}^{2}$ put $\tilde{x}=(w, z)$. Then

$$
\|x\|_{\psi}=(|z|+|w|) \psi\left(\frac{|w|}{|z|+|w|}\right)=(|w|+|z|) \tilde{\psi}\left(\frac{|z|}{|w|+|z|}\right)=\|\tilde{x}\|_{\tilde{\psi}}
$$

Therefore we have

$$
\begin{aligned}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) & =\sup _{\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2} \neq 0} \frac{\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2}}{2\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right)} \\
& =\sup _{\|\tilde{x}\|_{\tilde{\psi}}^{2}+\|\tilde{y}\|_{\tilde{\psi}}^{2} \neq 0} \frac{\|\tilde{x}+\tilde{y}\|_{\tilde{\psi}}^{2}+\|\tilde{x}-\tilde{y}\|_{\tilde{\psi}}^{2}}{2\left(\|\tilde{x}\|_{\tilde{\psi}}^{2}+\|\tilde{y}\|_{\tilde{\psi}}^{2}\right)} \\
& =C_{\mathrm{NJ}}\left(\|\cdot\|_{\tilde{\psi}}\right)
\end{aligned}
$$

Theorem 1. Let $\psi \in \Psi$.
(i) Assume that $\psi \geq \psi_{2}$. Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=\max _{0 \leq t \leq 1} \frac{\psi(t)^{2}}{\psi_{2}(t)^{2}} \tag{9}
\end{equation*}
$$

(ii) Assume that $\psi \leq \psi_{2}$. Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=\max _{0 \leq t \leq 1} \frac{\psi_{2}(t)^{2}}{\psi(t)^{2}} \tag{10}
\end{equation*}
$$

Proof. (i) Put $M_{1}=\max _{0 \leq t \leq 1} \psi(t) / \psi_{2}(t)$. Then by Lemma 3

$$
\begin{aligned}
\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2} & \leq M_{1}^{2}\left(\|x+y\|_{2}^{2}+\|x-y\|_{2}^{2}\right) \\
& =2 M_{1}^{2}\left(\|x\|_{2}^{2}+\|y\|_{2}^{2}\right) \\
& \leq 2 M_{1}^{2}\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right) .
\end{aligned}
$$

Now let $\psi / \psi_{2}$ attain the maximum at $t=t_{1}\left(0 \leq t_{1} \leq 1\right)$. Put $x_{1}=(1-$ $\left.t_{1}, 0\right), y_{1}=\left(0, t_{1}\right)$. Then

$$
\begin{align*}
\left\|x_{1}+y_{1}\right\|_{\psi}^{2}+\left\|x_{1}-y_{1}\right\|_{\psi}^{2} & =\left\|\left(1-t_{1}, t_{1}\right)\right\|_{\psi}^{2}+\left\|\left(1-t_{1},-t_{1}\right)\right\|_{\psi}^{2} \\
& =2 \psi\left(t_{1}\right)^{2} \\
& =2 M_{1}^{2}\left\{\left(1-t_{1}\right)^{2}+t_{1}^{2}\right\} \\
& =2 M_{1}^{2}\left(\left\|x_{1}\right\|_{\psi}^{2}+\left\|y_{1}\right\|_{\psi}^{2}\right), \tag{11}
\end{align*}
$$

which implies (9).
(ii) Put $M_{2}=\max _{0 \leq t \leq 1} \psi_{2}(t) / \psi(t)$. Then

$$
\begin{aligned}
\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2} & \leq\|x+y\|_{2}^{2}+\|x-y\|_{2}^{2} \\
& =2\left(\|x\|_{2}^{2}+\|y\|_{2}^{2}\right) \\
& \leq 2 M_{2}^{2}\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right) .
\end{aligned}
$$

Assume $M_{2}=\psi_{2}\left(t_{2}\right) / \psi\left(t_{2}\right)$ with some $t_{2}\left(0 \leq t_{2} \leq 1\right)$. Put $x_{2}=(1-$ $\left.t_{2}, t_{2}\right), y_{2}=\left(1-t_{2},-t_{2}\right)$. Then

$$
\begin{align*}
\left\|x_{2}+y_{2}\right\|_{\psi}^{2}+\left\|x_{2}-y_{2}\right\|_{\psi}^{2} & =4\left\{\left(1-t_{2}\right)^{2}+t_{2}^{2}\right\} \\
& =4 M_{2}^{2} \psi\left(t_{2}\right)^{2} \\
& =2 M_{2}^{2}\left\{\left\|\left(1-t_{2}, t_{2}\right)\right\|_{\psi}^{2}+\left\|\left(1-t_{2},-t_{2}\right)\right\|_{\psi}^{2}\right\} \\
& =2 M_{2}^{2}\left(\left\|x_{2}\right\|_{\psi}^{2}+\left\|y_{2}\right\|_{\psi}^{2}\right), \tag{12}
\end{align*}
$$

whence we have (10). This completes the proof.
Theorem 1 indicates that the NJ constant of $\|\cdot\|_{\psi}$ does not depend on the shape of $\psi$. This is stated in a little more general form:

Corollary 1. Let $\varphi, \psi \in \Psi$ be comparable with $\psi_{2}$.
(i) Let $\varphi \geq \psi_{2}$ and $\psi \geq \psi_{2}$. Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\varphi}\right)=C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \tag{13}
\end{equation*}
$$

if and only if

$$
\max _{0 \leq t \leq 1} \frac{\varphi(t)}{\psi_{2}(t)}=\max _{0 \leq t \leq 1} \frac{\psi(t)}{\psi_{2}(t)}
$$

(ii) Let $\varphi \geq \psi_{2}$ and $\psi \leq \psi_{2}$. Then (13) holds if and only if

$$
\max _{0 \leq t \leq 1} \frac{\varphi(t)}{\psi_{2}(t)}=\max _{0 \leq t \leq 1} \frac{\psi_{2}(t)}{\psi(t)} .
$$

The same is true for the other cases.
Corollary 2 (Clarkson [5]). Let $1 \leq p \leq \infty$ and $1 / p+1 / p^{\prime}=1$. Let $t=\min \left\{p, p^{\prime}\right\}$. Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{p}\right)=2^{(2 / t)-1} . \tag{14}
\end{equation*}
$$

In particular, $C_{\mathrm{NJ}}\left(\|\cdot\|_{1}\right)=C_{\mathrm{NJ}}\left(\|\cdot\|_{\infty}\right)=2$.
Indeed, if $1 \leq p \leq 2$,

$$
\psi_{2}(t) \leq \psi_{p}(t) \leq 2^{(1 / p)-(1 / 2)} \psi_{2}(t) \quad(0 \leq \forall t \leq 1)
$$

where the constant $2^{(1 / p)-(1 / 2)}$ is the best possible. Hence we have (14) by Theorem 1. For the case $2 \leq p \leq \infty$ a parallel argument works.
Remark 1. The only known way to calculate NJ constants needs Clarkson's inequalities (cf. [5, 9, 11]), whereas the above discussion to derive (14) does not require them.
Further, Theorem 1 enables us to obtain many examples easily. Let us present some. The following easy lemma is helpful for applying Theorem 1.

Lemma 5. Let $\varphi(t) \geq \psi(t)>0$ on $[a, b]$. Assume that $\varphi-\psi$ has the maximum, resp. $\psi$ has the minimum, at $t=c$ in $[a, b]$. Then $\varphi / \psi$ attains the maximum at $t=c$.

Indeed, the conclusion is immediate from the identity

$$
\frac{\varphi(t)}{\psi(t)}-1=\frac{\varphi(t)-\psi(t)}{\psi(t)}
$$

Example 1. Let $\|\cdot\|=\max \left\{\|\cdot\|_{2}, \lambda\|\cdot\|_{1}\right\}(1 / \sqrt{2} \leq \lambda \leq 1)$. Then

$$
C_{\mathrm{NJ}}(\|\cdot\|)=2 \lambda^{2} .
$$

In fact by Lemma $4,\|\cdot\|=\|\cdot\|_{\max \left\{\psi_{2}, \lambda \psi_{1}\right\}}$. Then by Theorem 1 we have

$$
\begin{aligned}
C_{\mathrm{NJ}}(\|\cdot\|) & =\max _{0 \leq t \leq 1}\left[\frac{\max \left\{\psi_{2}(t), \lambda \psi_{1}(t)\right\}}{\psi_{2}(t)}\right]^{2} \\
& =\max _{0 \leq t \leq 1}\left\{\frac{\lambda \psi_{1}(t)}{\psi_{2}(t)}\right\}^{2}=\frac{\lambda^{2}}{1 / 2}=2 \lambda^{2} .
\end{aligned}
$$

The following example treats a nonnormalized norm.
Example $2\left(\left[10\right.\right.$, Proposition 1]). Let $\|\cdot\|=\max \left\{\|\cdot\|_{2}, \lambda\|\cdot\|_{\infty}\right\}(1 \leq \lambda$ $\leq \sqrt{2}$ ). Then

$$
C_{\mathrm{NJ}}(\|\cdot\|)=\lambda^{2} .
$$

Indeed, put $\|\cdot\|_{0}=\max \left\{\lambda^{-1}\|\cdot\|_{2},\|\cdot\|_{\infty}\right\}$. Then $\|\cdot\|_{0}=\|\cdot\|_{\max \left\{\lambda^{-1} \psi_{2}, \psi_{\infty}\right\}}$ by Lemma 4 and $\|\cdot\|=\lambda\|\cdot\|_{0}$. Hence we have

$$
\begin{aligned}
C_{\mathrm{NJ}}(\|\cdot\|) & =C_{\mathrm{NJ}}\left(\lambda\|\cdot\|_{0}\right) \\
& =C_{\mathrm{NJ}}\left(\|\cdot\|_{0}\right) \\
& =\max _{0 \leq t \leq 1}\left[\frac{\psi_{2}(t)}{\max \left\{\lambda^{-1} \psi_{2}(t), \psi_{x}(t)\right\}}\right]^{2}
\end{aligned}
$$

by Theorem 1 . Now clearly $\psi_{2}(t) /\left[\max \left\{\lambda^{-1} \psi_{2}(t), \psi_{\infty}(t)\right\}\right]$ is symmetric with respect to $t=1 / 2$. Let $t_{0}$ be such that $\lambda^{-1} \psi_{2}\left(t_{0}\right)=\psi_{o x}\left(t_{0}\right)\left(0 \leq t_{0} \leq 1 / 2\right)$. Then we have

$$
\max _{0 \leq t \leq t_{0}}\left[\frac{\psi_{2}(t)}{\max \left\{\lambda^{-1} \psi_{2}(t), \psi_{\infty}(t)\right\}}\right]^{2}=\lambda^{2}
$$

by Lemma 5, and clearly

$$
\max _{t_{0} \leq t \leq 1 / 2}\left[\frac{\psi_{2}(t)}{\max \left\{\lambda^{-1} \psi_{2}(t), \psi_{\infty}(t)\right\}}\right]^{2}=\max _{t_{0} \leq t \leq 1 / 2}\left[\frac{\psi_{2}(t)}{\lambda^{-1} \psi_{2}(t)}\right]^{2}=\lambda^{2}
$$

Therefore we have $C_{\mathrm{NJ}}(\|\cdot\|)=\lambda^{2}$.
Example 3. Let $2 \leq p<\infty$. Let $\|\cdot\|_{p, 2}$ be the (Lorentz) $\ell_{p, 2}$-norm.

$$
\|(z, w)\|_{p, 2}=\left\{|z|^{* 2}+2^{(2 / p)-1}|w|^{* 2}\right\}^{1 / 2}
$$

where $\left\{|z|^{*},|w|^{*}\right\}$ is the non-increasing rearrangement of $\{|z|,|w|\}$; that is, $|z|^{*} \geq|w|^{*}$. (Note that if $p<2,\|\cdot\|_{p, 2}$ is a quasi-norm; cf. [7, Proposition 1; 14, p. 126; 2, p. 8]). Then

$$
C_{\mathrm{NJ}}\left(\|\cdot\|_{p, 2}\right)=\frac{2}{1+2^{2 / p-1}} .
$$

Indeed, $\|\cdot\|_{p, 2} \in N_{a}$, and the corresponding convex function is given by

$$
\psi_{p, 2}(t)= \begin{cases}\left\{(1-t)^{2}+2^{2 / p-1} t^{2}\right\}^{1 / 2} & \text { if } 0 \leq t \leq 1 / 2 \\ \left\{t^{2}+2^{2 / p-1}(1-t)^{2}\right\}^{1 / 2} & \text { if } 1 / 2 \leq t \leq 1\end{cases}
$$

Since $\psi_{p, 2} \leq \psi_{2}$ and $\psi_{2} / \psi_{p, 2}$ is symmetric with respect to $t=1 / 2$, we find the maximum of $\psi_{2}^{2} / \psi_{p, 2}^{2}$ in the interval [ $0,1 / 2$ ]. The difference $\psi_{2}(t)^{2}-\psi_{p, 2}(t)^{2}=\left(1-2^{2 / p-1}\right) t^{2}$ takes its maximum at $t=1 / 2$, and $\psi_{p, 2}$ has the minimum at $t=1 / 2$. Therefore by Lemma 5 we have

$$
\max _{0 \leq t \leq 1} \frac{\psi_{2}(t)^{2}}{\psi_{p, 2}(t)^{2}}=\frac{\psi_{2}(1 / 2)^{2}}{\psi_{p, 2}(1 / 2)^{2}}=\frac{2}{1+2^{2 / p-1}}
$$

which implies the conclusion by Theorem 1.

## 4. NJ CONSTANT OF ABSOLUTE NORMALIZED NORMSTHE GENERAL CASE

Lemma 6. Let $1 / 2 \leq \alpha \leq 1$ and let

$$
\psi_{\alpha}(t)=\left\{\begin{array}{cl}
\frac{\alpha-1}{\alpha} t+1 & \text { if } 0 \leq t \leq \alpha \\
t & \text { if } \alpha \leq t \leq 1
\end{array}\right.
$$

Then

$$
\begin{align*}
& M_{1}=\max _{0 \leq t \leq 1} \frac{\psi_{\alpha}(t)}{\psi_{2}(t)}=\left\{\left(2-\frac{1}{\alpha}\right)^{2}+1\right\}^{1 / 2},  \tag{15}\\
& M_{2}=\max _{0 \leq t \leq 1} \frac{\psi_{2}(t)}{\psi_{\alpha}(t)}=\left\{\left(\frac{1}{\alpha}-1\right)^{2}+1\right\}^{1 / 2} \tag{16}
\end{align*}
$$

Proof. If $\alpha=1 / 2$ or $\alpha=1$, the conclusion is clear by Lemma 5 . Let $1 / 2<\alpha<1$. Easy calculation shows that the function $\psi_{\alpha} / \psi_{2}$ attains the maximum at $t=(2 \alpha-1) /(3 \alpha-1)$, which gives (15). The function $\psi_{2} / \psi_{\alpha}$ clearly has the maximum at $t=\alpha$, which implies (16).

Although the notation $\psi_{\alpha}$ is not consistent with $\psi_{p}$ corresponding to the $\ell_{p}$-norm, there will be no confusion in the following.

Theorem 2. Let $\psi \in \Psi$ and let

$$
\begin{equation*}
M_{1}=\max _{0 \leq t \leq 1} \frac{\psi(t)}{\psi_{2}(t)} \quad \text { and } \quad M_{2}=\max _{0 \leq t \leq 1} \frac{\psi_{2}(t)}{\psi(t)} . \tag{17}
\end{equation*}
$$

Then

$$
\begin{equation*}
\max \left\{M_{1}^{2}, M_{2}^{2}\right\} \leq C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \leq M_{1}^{2} M_{2}^{2} . \tag{18}
\end{equation*}
$$

Further we have

$$
\begin{equation*}
1 \leq \max \left\{M_{1}^{2}, M_{2}^{2}\right\} \leq M_{1}^{2} M_{2}^{2} \leq 2 \tag{19}
\end{equation*}
$$

Proof. For all $x, y \in \mathbb{C}^{2}$ we have

$$
\begin{align*}
\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2} & \leq M_{1}^{2}\left(\|x+y\|_{2}^{2}+\|x-y\|_{2}^{2}\right) \\
& =2 M_{1}^{2}\left(\|x\|_{2}^{2}+\|y\|_{2}^{2}\right) \\
& \leq 2 M_{1}^{2} M_{2}^{2}\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right) \tag{20}
\end{align*}
$$

which implies that $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \leq M_{1}^{2} M_{2}^{2}$. Next let

$$
\begin{equation*}
M_{1}=\frac{\psi\left(t_{1}\right)}{\psi_{2}\left(t_{1}\right)}, \quad M_{2}=\frac{\psi_{2}\left(t_{2}\right)}{\psi\left(t_{2}\right)} \tag{21}
\end{equation*}
$$

with some $0 \leq t_{1}, t_{2} \leq 1$. Put $x_{1}=\left(1-t_{1}, 0\right), y_{1}=\left(0, t_{1}\right)$. Then by (11) we have $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \geq M_{1}^{2}$. In the same way, by putting $x_{2}=\left(1-t_{2}, t_{2}\right)$, $y_{2}=\left(1-t_{2},-t_{2}\right)$, we have $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \geq M_{2}^{2}$ by (12).

Now we prove (19). The first two inequalities are obvious. If $\psi \geq \psi_{2}$ or $\psi \leq \psi_{2}$, the last inequality in (19) is easy to see (merely note that $\psi_{1} / \psi_{2}$ $\leq \sqrt{2}$ and $\psi_{2} / \psi_{\infty} \leq \sqrt{2}$ ). So assume this is not the case. Then $M_{1}, M_{2}>1$, whence we have (21) with $0<t_{1}, t_{2}<1$. Owing to Proposition 1 we may assume that $t_{1}<t_{2}$. Let $(\alpha, \alpha)$ be the intersection of the line $s=t$ and the line combining the points $(0,1)$ and $\left(t_{2}, \psi\left(t_{2}\right)\right)$. Then evidently $1 / 2<\alpha$ $<1$ and $t_{2}<\alpha$. Hence

$$
M_{2}=\frac{\psi_{2}\left(t_{2}\right)}{\psi\left(t_{2}\right)}=\frac{\psi_{2}\left(t_{2}\right)}{\psi_{\alpha}\left(t_{2}\right)} \leq \frac{\psi_{2}(\alpha)}{\psi_{\alpha}(\alpha)},
$$

where $\psi_{\alpha}$ is as in Lemma 6. On the other hand, since $\psi_{2}\left(t_{1}\right) \leq \psi\left(t_{1}\right) \leq$ $\psi_{\alpha}\left(t_{1}\right)$ by the convexity of $\psi(t)$ and (21), we have

$$
M_{1}=\frac{\psi\left(t_{1}\right)}{\psi_{2}\left(t_{1}\right)} \leq \frac{\psi_{\alpha}\left(t_{1}\right)}{\psi_{2}\left(t_{1}\right)} .
$$

Therefore by Lemma 6 we have

$$
\begin{aligned}
M_{1} M_{2} & =\frac{\psi\left(t_{1}\right)}{\psi_{2}\left(t_{1}\right)} \frac{\psi_{2}\left(t_{2}\right)}{\psi\left(t_{2}\right)} \\
& \leq \max _{0 \leq t \leq 1} \frac{\psi_{\alpha}(t)}{\psi_{2}(t)} \max _{0 \leq t \leq 1} \frac{\psi_{2}(t)}{\psi_{\alpha}(t)} \\
& =\left\{\left(2-\frac{1}{\alpha}\right)^{2}+1\right\}^{1 / 2}\left\{\left(\frac{1}{\alpha}-1\right)^{2}+1\right\}^{1 / 2}
\end{aligned}
$$

Put $u=1 / \alpha-1$. Then, $0<u<1$ and

$$
\begin{align*}
M_{1}^{2} M_{2}^{2} & \leq\left(u^{2}+1\right)\left\{(1-u)^{2}+1\right\} \\
& =u(u-1)\left(u^{2}-u+2\right)+2 \\
& <2 . \tag{22}
\end{align*}
$$

This completes the proof.
Remark 2. (i) In Theorem 2 we have $M_{1}^{2} M_{2}^{2}=2$ if and only if $\alpha=1$ or $\alpha=1 / 2$; in this case $\psi=\psi_{1}$ or $\psi=\psi_{\infty}$. In fact, the "if" part is clear, and the opposite follows directly from (22).
(ii) $\max \left\{M_{1}, M_{2}\right\}=1$ if and only if $\psi=\psi_{2}$.
(iii) $\max \left\{M_{1}, M_{2}\right\}=M_{1} M_{2}$ if and only if $\psi \geq \psi_{2}$ or $\psi \leq \psi_{2}$ : In particular, Theorem 1 is also a result of this fact.

As a consequence of Theorem 2 we have
Corollary 3. Let $\|\cdot\| \in N_{a}$. Then $C_{\mathrm{NJ}}(\|\cdot\|)=2$ if and only if $\|\cdot\|$ is an $\ell_{1}$ - or $\ell_{\infty}$-norm: In other words, all norms in $N_{a}$ except $\ell_{1}$ - and $\ell_{\infty}$-norms are uniformly non-square.
Proof. By Theorem 2, $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=2$ if and only if $M_{1}^{2} M_{2}^{2}=2$, which occurs only when $\psi=\psi_{1}$ or $\psi=\psi_{\infty}$ by Remark 2 (i).
Now, according to Theorem 1, the identity $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=\max \left\{M_{1}^{2}, M_{2}^{2}\right\}$ holds in the estimate (18) of Theorem 2 if $\psi$ is comparable with $\psi_{2}$. The next theorem asserts that for another wide class of convex functions we have $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=M_{1}^{2} M_{2}^{2}$ and $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)>\max \left\{M_{1}^{2}, M_{2}^{2}\right\}$.

Theorem 3. Let $\psi \in \Psi$ and let $\psi(t)=\psi(1-t)$ for all $0 \leq t \leq 1$. Assume that $M_{1}=\max _{0 \leq t \leq 1} \psi(t) / \psi_{2}(t)$ or $M_{2}=\max _{0 \leq t \leq 1} \psi_{2}(t) / \psi(t)$ is taken at $t=1 / 2$. Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=M_{1}^{2} M_{2}^{2} \tag{23}
\end{equation*}
$$

Moreover, if neither $\psi \geq \psi_{2}$ nor $\psi \leq \psi_{2}$,

$$
\begin{equation*}
\max \left\{M_{1}^{2}, M_{2}^{2}\right\}<C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \tag{24}
\end{equation*}
$$

Proof. Suppose first $M_{1}=\psi(1 / 2) / \psi_{2}(1 / 2)$. Take an arbitrary $t$ with $0 \leq t \leq 1$ and put $x=(t, 1-t)$ and $y=(1-t, t)$. Then

$$
\|x\|_{\psi}^{2}=\psi(1-t)^{2}=\psi(t)^{2}, \quad\|y\|_{\psi}^{2}=\psi(t)^{2}
$$

On the other hand

$$
\begin{aligned}
& \|x+y\|_{\psi}^{2}=\|(1,1)\|_{\psi}^{2}=4 \psi(1 / 2)^{2} \\
& \|x-y\|_{\psi}^{2}=\|(2 t-1,1-2 t)\|_{\psi}^{2}=4(2 t-1)^{2} \psi(1 / 2)^{2}
\end{aligned}
$$

Therefore

$$
\begin{aligned}
\frac{\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2}}{2\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right)} & =\frac{4 \psi(1 / 2)^{2}\left\{(2 t-1)^{2}+1\right\}}{4 \psi(t)^{2}} \\
& =\frac{\psi(1 / 2)^{2}\left\{(1-t)^{2}+t^{2}\right\}}{\psi(t)^{2} / 2} \\
& =\frac{\psi(1 / 2)^{2} \psi_{2}(t)^{2}}{\psi_{2}(1 / 2)^{2} \psi(t)^{2}}=M_{1}^{2} \frac{\psi_{2}(t)^{2}}{\psi(t)^{2}}
\end{aligned}
$$

Since $t$ is arbitrary, we have

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right) \geq M_{1}^{2} M_{2}^{2} \tag{25}
\end{equation*}
$$

which, combined with (18), implies (23). In case of $M_{2}=\psi_{2}(1 / 2) / \psi(1 / 2)$, let $x$ and $y$ be as above and put $u=x+y$ and $v=x-y$. Then since

$$
\begin{aligned}
\frac{\|u+v\|_{\psi}^{2}+\|u-v\|_{\psi}^{2}}{2\left(\|u\|_{\psi}^{2}+\|v\|_{\psi}^{2}\right)} & =\frac{2\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right)}{\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2}} \\
& =\frac{\psi_{2}(1 / 2)^{2} \psi(t)^{2}}{\psi(1 / 2)^{2} \psi_{2}(t)^{2}}=M_{2}^{2} \frac{\psi(t)^{2}}{\psi_{2}(t)^{2}}
\end{aligned}
$$

we have (25) and hence (23). The inequality (24) is a direct consequence of (23) and Remark 2 (iii).

Example 4. Let $1 / 2 \leq \beta \leq 1$ Let $\varphi_{\beta}(t)=\max \{1-t, t, \beta\}$ (note that neither $\varphi_{\beta} \geq \psi_{2}$ nor $\varphi_{\beta} \leq \psi_{2}$ if $1 / \sqrt{2}<\beta<1$ ). Then

$$
C_{\mathrm{NJ}}\left(\|\cdot\|_{\varphi_{\beta}}\right)= \begin{cases}\frac{1}{\beta^{2}}\left\{(1-\beta)^{2}+\beta^{2}\right\} & \text { if } \frac{1}{2} \leq \beta \leq \frac{1}{\sqrt{2}} \\ 2\left\{(1-\beta)^{2}+\beta^{2}\right\} & \text { if } \frac{1}{\sqrt{2}} \leq \beta \leq 1\end{cases}
$$

Indeed, by Lemma 5

$$
M_{1}= \begin{cases}1 & \text { if } \frac{1}{2} \leq \beta \leq \frac{1}{\sqrt{2}} \\ \frac{\varphi_{\beta}(1 / 2)}{\psi_{2}(1 / 2)}=\frac{\beta}{1 / \sqrt{2}}=\sqrt{2} \beta & \text { if } \frac{1}{\sqrt{2}} \leq \beta \leq 1\end{cases}
$$

and

$$
M_{2}=\frac{\psi_{2}(\beta)}{\varphi_{\beta}(\beta)}=\frac{1}{\beta}\left\{(1-\beta)^{2}+\beta^{2}\right\}^{1 / 2},
$$

whence we have the conclusion by Theorem 3 .
Finally we see a class of convex functions for which the identity $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)=M_{1}^{2} M_{2}^{2}$ fails to hold.

Theorem 4. Let $\psi \in \Psi$. Let $M_{1}$ and $M_{2}$ be as in Theorem 2. Assume that

$$
\begin{equation*}
\max \left\{t ; M_{1}=\frac{\psi(t)}{\psi_{2}(t)}\right\}<\min \left\{s ; M_{2}=\frac{\psi_{2}(s)}{\psi(s)}\right\} \tag{26}
\end{equation*}
$$

or

$$
\begin{equation*}
\min \left\{t ; M_{1}=\frac{\psi(t)}{\psi_{2}(t)}\right\}>\max \left\{s ; M_{2}=\frac{\psi_{2}(s)}{\psi(s)}\right\} \tag{27}
\end{equation*}
$$

(hence $\psi$ is not symmetric with respect to $t=1 / 2$ ). Then

$$
\begin{equation*}
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)<M_{1}^{2} M_{2}^{2} . \tag{28}
\end{equation*}
$$

Proof. It is enough to show (28) in the case (26) by Proposition 1. Put

$$
t_{0}=\max \left\{t ; M_{1}=\frac{\psi(t)}{\psi_{2}(t)}\right\}, \quad s_{0}=\min \left\{s ; M_{2}=\frac{\psi_{2}(s)}{\psi(s)}\right\} .
$$

Then clearly $0<t_{0}<s_{0}<1$. Assume that (28) is not valid. Then there are $x, y \in \mathbb{C}^{2}$, not both 0 , such that

$$
\|x+y\|_{\psi}^{2}+\|x-y\|_{\psi}^{2}=2 M_{1}^{2} M_{2}^{2}\left(\|x\|_{\psi}^{2}+\|y\|_{\psi}^{2}\right) .
$$

It should be noted that all of $x, y, x+y$, and $x-y$ are not 0 because $\psi \neq \psi_{2}$. Then by (20) we have

$$
\begin{equation*}
\|x+y\|_{\psi}=M_{1}\|x+y\|_{2}, \quad\|x-y\|_{\psi}=M_{1}\|x-y\|_{2} \tag{29}
\end{equation*}
$$

and

$$
\begin{equation*}
\|x\|_{2}=M_{2}\|x\|_{\psi}, \quad\|y\|_{2}=M_{2}\|y\|_{\psi} \tag{30}
\end{equation*}
$$

Put $x=(a, b)$ and $y=(c, d)$. Then by (29)

$$
\begin{aligned}
& \psi\left(\frac{|b+d|}{|a+c|+|b+d|}\right)=M_{1} \psi_{2}\left(\frac{|b+d|}{|a+c|+|b+d|}\right) \\
& \psi\left(\frac{|b-d|}{|a-c|+|b-d|}\right)=M_{1} \psi_{2}\left(\frac{|b-d|}{|a-c|+|b-d|}\right)
\end{aligned}
$$

from which it follows that

$$
\frac{|b+d|}{|a+c|+|b+d|} \leq t_{0}, \quad \frac{|b-d|}{|a-c|+|b-d|} \leq t_{0}
$$

In the same way, by (30) we have

$$
\frac{|b|}{|a|+|b|} \geq s_{0}, \quad \frac{|d|}{|c|+|d|} \geq s_{0} .
$$

Therefore

$$
\begin{aligned}
& |b+d| \leq \frac{t_{0}}{1-t_{0}}|a+c| \\
& |b-d| \leq \frac{t_{0}}{1-t_{0}}|a-c|
\end{aligned}
$$

and

$$
|b| \geq \frac{s_{0}}{1-s_{0}}|a|, \quad|d| \geq \frac{s_{0}}{1-s_{0}}|c|
$$

Hence we have

$$
\begin{aligned}
|b+d|^{2}+|b-d|^{2} & \leq\left(\frac{t_{0}}{1-t_{0}}\right)^{2}\left(|a+c|^{2}+|a-c|^{2}\right) \\
& =2\left(\frac{t_{0}}{1-t_{0}}\right)^{2}\left(|a|^{2}+|c|^{2}\right)
\end{aligned}
$$

and

$$
\begin{aligned}
|b+d|^{2}+|b-d|^{2} & =2\left(|b|^{2}+|d|^{2}\right) \\
& \geq 2\left(\frac{s_{0}}{1-s_{0}}\right)^{2}\left(|a|^{2}+|c|^{2}\right) .
\end{aligned}
$$

Consequently we have $t_{0} /\left(1-t_{0}\right) \geq s_{0} /\left(1-s_{0}\right)$ because $|a|^{2}+|c|^{2} \neq 0$, and hence $t_{0} \geq s_{0}$, which contradicts our assumption. This completes the proof.

Corollary 4. Let $\psi \in \Psi$. Let $M_{1}$ and $M_{2}$ be as in Theorem 2. If there exists uniquely one point such that $\psi / \psi_{2}\left(\right.$ resp. $\left.\psi_{2} / \psi\right)$ attains $M_{1}\left(\right.$ resp. $\left.M_{2}\right)$, then

$$
C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi}\right)<M_{1}^{2} M_{2}^{2} .
$$

In fact, these points at which $M_{1}$ and $M_{2}$ are attained are different by the assumption.
Remark 3. (i) In Theorem 3 the condition $\psi(t)=\psi(1-t)$ is essential. Indeed, modify the function $\varphi_{\beta}$ in Example 4 as follows: Let $1 / \sqrt{2}<\gamma<1$. Let $t_{0}$ be the smaller solution of the equation, $\psi_{2}(t)=\gamma$. Define

$$
\omega_{\gamma}(t)= \begin{cases}\psi_{2}(t) & \text { if } 0 \leq t \leq t_{0} \\ \gamma & \text { if } t_{0} \leq t \leq \gamma \\ t & \text { if } \gamma \leq t \leq 1\end{cases}
$$

Then, $\omega_{\gamma} / \psi_{2}$ has the maximum at $t=1 / 2$, but $\omega_{\gamma}$ is not symmetric with respect to $t=1 / 2$. On the other hand, $\omega_{\gamma}$ satisfies the condition in Corollary 4, and hence we have $C_{\mathrm{NJ}}\left(\|\cdot\|_{\omega_{\gamma}}\right)<M_{1}^{2} M_{2}^{2}$.
(ii) For $\psi_{\alpha}$ in Lemma 6 we have $C_{\mathrm{NJ}}\left(\|\cdot\|_{\psi_{\alpha}}\right)<M_{1}^{2} M_{2}^{2}$ by Corollary 4, where $M_{1}$ and $M_{2}$ are as in Lemma 6 .

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