# A fractional variational iteration method for solving fractional nonlinear differential equations 

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## A R T I CLE INFO

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#### Abstract

Recently, fractional differential equations have been investigated by employing the famous variational iteration method. However, all the previous works avoid the fractional order term and only handle it as a restricted variation. A fractional variational iteration method was first proposed in [G.C. Wu, E.W.M. Lee, Fractional variational iteration method and its application, Phys. Lett. A 374 (2010) 2506-2509] and gave a generalized Lagrange multiplier. In this paper, two fractional differential equations are approximately solved with the fractional variational iteration method.


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## 1. Introduction

Recently the variational iteration method [1] has been widely applied to analytically solve fractional differential equations [2-6], where the term with the fractional derivative was considered as a restricted variation, making the identification of the Lagrange multiplier very inaccurate. To overcome this problem, we investigate the local behavior of fractional differential equations and determine the Lagrange multiplier in a more accurate way with the fractional variation iteration method (FVIM) [7].

## 2. Properties of local fractional calculus

On the basis of Cantor-like sets, Kolwankar and Gangal [8,9] proposed a concept of a local fractional derivative:

$$
\begin{equation*}
D_{x_{0}}^{\alpha} f(x)=\lim _{x \rightarrow x_{0}} \frac{1}{\Gamma(n-\alpha)} \frac{\mathrm{d}^{n}}{\mathrm{~d} x^{n}} \int_{x_{0}}^{x}(x-\xi)^{n-\alpha}\left(f(\xi)-f\left(x_{0}\right)\right) \mathrm{d} \xi \tag{1}
\end{equation*}
$$

where the derivative on the right-hand side is the Riemann-Liouville fractional derivative.
Chen et al. [10] presented the necessary conditions for

$$
\begin{equation*}
D_{x}^{\alpha} f(x)=\lim _{y \rightarrow x} \frac{\Gamma(1+\alpha)(f(y)-f(x))}{(y-x)^{\alpha}}, \quad 0<\alpha \leq 1 \tag{2}
\end{equation*}
$$

We now can derive the following useful properties of the Kolwankar-Gangal derivative.
(a) Integration with respect to $(\mathrm{d} x)^{\alpha}$ :

In the derivation of Eq. (2), for any $\epsilon$, there exists a $d$, where $\left|x_{i+1}-x_{i}\right|<d$, such that

$$
\left|f\left(x_{i+1}\right)-f\left(x_{i}\right)-D_{x_{i}}^{\alpha} f\left(x_{i}\right) \frac{\left(x_{i+1}-x_{i}\right)^{\alpha}}{\Gamma(1+\alpha)}\right|<\epsilon
$$

[^0]Since $f(x)$ is continuous in the closed interval $[a, b]$, if we consider a finite partition, $\left[x_{0}, x_{1}\right], \ldots,\left[x_{i}, x_{i+1}\right], \ldots,\left[x_{n-1}, x_{n}\right]$ where $x_{0}=a, x_{n}=b$, and note that

$$
f(b)-f(a)=\sum_{i=0}^{n-1}\left(f\left(x_{i+1}\right)-f\left(x_{i}\right)\right),
$$

we can choose the maximal $\delta$ such that

$$
\left|f(b)-f(a)-\sum_{i=0}^{n-1} D_{x_{i}}^{\alpha} f\left(x_{i}\right) \frac{\left(x_{i+1}-x_{i}\right)^{\alpha}}{\Gamma(1+\alpha)}\right|<\epsilon .
$$

As a result, we develop a definition of fractional integration.
If $g(x)$ is continuous in the interval $[a, b]$ and the limit of $\sum_{i=0}^{n-1} g\left(x_{i}\right) \frac{\left(x_{i+1}-x_{i}\right)^{\alpha}}{\Gamma(1+\alpha)}$ exists when $n$ tends to infinity, then we say that the function $g(x)$ is $\alpha$-order integrable in the interval $[a, b]$ denoted by

$$
\begin{equation*}
{ }_{a} I_{b}^{\alpha} g(x)=\frac{1}{\Gamma(1+\alpha)} \int_{a}^{b} g(x)(\mathrm{d} x)^{\alpha}=\lim _{n \rightarrow \infty} \sum_{i=0}^{n-1} g\left(x_{i}\right) \frac{\left(x_{i+1}-x_{i}\right)^{\alpha}}{\Gamma(1+\alpha)} . \tag{3}
\end{equation*}
$$

(b) The fractional Leibniz product law:

If $u$ and $v$ are $\alpha$-order differentiable functions, we have the generalized Leibniz product law from Eq. (2)

$$
\begin{equation*}
D_{x}^{(\alpha)}(u v)=u^{(\alpha)} v+u v^{(\alpha)} \tag{4}
\end{equation*}
$$

(c) The fractional Leibniz formulation:

$$
\begin{equation*}
{ }_{o} I_{x}^{\alpha} D_{x}^{\alpha} f(x)=f(x)-f(0), \quad 0<\alpha \leq 1 \tag{5}
\end{equation*}
$$

Therefore, integration by parts can be used in the fractional calculus:

$$
\begin{equation*}
{ }_{a} I_{b}^{\alpha} u^{(\alpha)} v=\left.(u v)\right|_{a} ^{b}-{ }_{a} I_{b}^{\alpha} u v^{(\alpha)} . \tag{6}
\end{equation*}
$$

## 3. The fractional variational iteration method

In this section, two examples are given to illustrate the effect of the proposed method.
Example 1. As the first example, we consider a time-fractional diffusion equation:

$$
\begin{equation*}
\frac{\partial^{\alpha} u(x, t)}{\partial t^{\alpha}}=c \frac{\partial^{2} u(x, t)}{\partial x^{2}}-\frac{\partial(F(x) u(x, t))}{\partial x}, \quad 0<\alpha \leq 1, \tag{7}
\end{equation*}
$$

where $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$ is the Caputo derivative, with initial condition $u(x, 0)=f(x)$.
We replace the fractional Caputo derivative with the local fractional derivative in Eq. (7), and assume $c=1, F(x)=-x$ which leads to

$$
\begin{equation*}
\frac{\partial^{\alpha} u(x, t)}{\partial t^{\alpha}}=\frac{\partial^{2} u(x, t)}{\partial x^{2}}+\frac{\partial(x u(x, t))}{\partial x}, \quad 0<\alpha \leq 1, \tag{8}
\end{equation*}
$$

with the initial condition $u(x, 0)=x^{2}$.
A correction functional for Eq. (8) can be constructed as follows [1,5]:

$$
\begin{equation*}
u_{n+1}(x, t)=u_{n}(x, t)+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{t} \lambda(t, \tau)\left\{\frac{\partial^{\alpha} u_{n}(x, \tau)}{\partial \tau^{\alpha}}-\frac{\partial^{2} \tilde{u}_{n}(x, \tau)}{\partial x^{2}}-\frac{\partial\left(x \tilde{u}_{n}(x, \tau)\right)}{\partial x}\right\}(\mathrm{d} \tau)^{\alpha} \tag{9}
\end{equation*}
$$

with the property, from Eqs. (4)-(6), that $\lambda(t, \tau)$ must satisfy

$$
\begin{equation*}
\frac{\partial^{\alpha} \lambda(t, \tau)}{\partial \tau^{\alpha}}=0, \quad \text { and } \quad 1+\left.\lambda(t, \tau)\right|_{\tau=t}=0 . \tag{10}
\end{equation*}
$$

Therefore, $\lambda(t, \tau)$ can be identified as $\lambda(t, \tau)=-1$.
With the fractional Taylor series [11], we can determine the trial function or the initial value $u_{0}(x, t)=u_{0}(x, 0)=$ $f(x)=x^{2}$ in the iteration formulation as follows:

$$
u_{n+1}(x, t)=u_{n}(x, t)-\frac{1}{\Gamma(1+\alpha)} \int_{0}^{t}\left\{\frac{\partial^{\alpha} u_{n}(x, \tau)}{\partial \tau^{\alpha}}-\frac{\partial^{2} u_{n}(x, t)}{\partial x^{2}}-\frac{\partial\left(x u_{n}(x, t)\right)}{\partial x}\right\}(\mathrm{d} \tau)^{\alpha} .
$$

We can derive

$$
\begin{aligned}
u_{1}(x, t) & =x^{2}-\frac{1}{\Gamma(1+\alpha)} \int_{0}^{t}\left\{\frac{\partial^{\alpha} u_{0}(x, \tau)}{\partial \tau^{\alpha}}-\frac{\partial^{2} u_{0}(x, \tau)}{\partial x^{2}}-\frac{\partial\left(x u_{0}(x, \tau)\right)}{\partial x}\right\}(\mathrm{d} \tau)^{\alpha} \\
& =x^{2}+\frac{\left(2+3 x^{2}\right) t^{\alpha}}{\Gamma(1+\alpha)} \\
u_{2}(x, t) & =x^{2}+\frac{\left(2+3 x^{2}\right) t^{\alpha}}{\Gamma(1+\alpha)}+\frac{\left(8+9 x^{2}\right) t^{2 \alpha}}{\Gamma(1+2 \alpha)} \\
u_{3}(x, t) & =x^{2}+\frac{\left(2+3 x^{2}\right) t^{\alpha}}{\Gamma(1+\alpha)}+\frac{\left(8+9 x^{2}\right) t^{2 \alpha}}{\Gamma(1+2 \alpha)}+\frac{\left(26+27 x^{2}\right) t^{3 \alpha}}{\Gamma(1+3 \alpha)} .
\end{aligned}
$$

As a result, the exact solution can be given in a compact form:

$$
\begin{equation*}
u(x, t)=\lim _{n \rightarrow \infty} u_{n}(x, t)=\lim _{n \rightarrow \infty} \sum_{i=0}^{n} \frac{k^{i} t^{i \alpha}}{\Gamma(1+i \alpha)}=E_{\alpha}\left(k t^{\alpha}\right) \tag{11}
\end{equation*}
$$

where $k^{i}=x^{2}+\left(1+x^{2}\right)\left(3^{i}-1\right)$ and $E_{\alpha}\left(k t^{\alpha}\right)$ is the Mittag-Leffler function. We note that Eq. (12) is also the exact solution of the fractional diffusion equation [5] if we take $\alpha=\frac{1}{2}$.

Example 2. In order to illustrate the FVIM for higher fractional order equations, we only consider an initial value problem given in [12]:

$$
\begin{equation*}
y^{(2 \alpha)}=y^{2}+1, \quad 0<\alpha \leq 1,0 \leq x \leq 1 \tag{12}
\end{equation*}
$$

with $y(0)=0$ and $y^{(\alpha)}(0)=1$, where $y^{(2 \alpha)}=D_{x}^{\alpha} D_{x}^{\alpha} y$.
Construct the following functional:

$$
y_{n+1}(x, t)=y_{n}(x, t)+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \lambda\left\{y_{n}^{(2 \alpha)}-\tilde{y}_{n}^{2}(\xi)-1\right\}(\mathrm{d} \xi)^{\alpha}
$$

We have

$$
\begin{aligned}
\delta y_{n+1} & =\delta y_{n}+\frac{1}{\Gamma(1+\alpha)} \delta \int_{0}^{x} \lambda\left(y_{n}^{(2 \alpha)}-y_{n}^{2}(\xi)-1\right)(\mathrm{d} \xi)^{\alpha} \\
& =\delta y_{n}+\left.\lambda \delta y_{n}^{(\alpha)}\right|_{\xi=x}-\left.\lambda^{(\alpha)}(\tau) \delta y_{n}(\tau)\right|_{\xi=x}+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \lambda^{(2 \alpha)} \delta y_{n}(\mathrm{~d} \xi)^{\alpha} .
\end{aligned}
$$

Similarly, we can set the coefficients of $\delta u_{n}(\tau)$ and $\delta y_{n}^{(\alpha)}$ to zero:

$$
1-\left.\lambda^{(\alpha)}\right|_{\xi=x}=0, \quad \lambda^{(2 \alpha)}=0 \quad \text { and }\left.\quad \lambda\right|_{\xi=x}=0
$$

As a result, the Lagrange multiplier can be identified as

$$
\begin{equation*}
\lambda=\frac{(\xi-x)^{\alpha}}{\Gamma(1+\alpha)} \tag{13}
\end{equation*}
$$

By a similar manipulation, we can derive a more generalized Lagrange multiplier:

$$
\lambda=(-1)^{(m)} \frac{(\xi-x)^{(m-1) \alpha}}{\Gamma(1+(m-1) \alpha)}
$$

for higher fractional nonlinear ordinary differential equations:

$$
y^{(m \alpha)}=N(y), \quad 0<\alpha \leq 1
$$

For Eq. (13), we can check when $\alpha=1$ and $\lambda=\xi-x$ is the multiplier for the differential equation of integer order

$$
\begin{equation*}
y^{(\prime \prime)}=y^{2}+1 \tag{14}
\end{equation*}
$$

The iteration formulation for Eq. (12) can be rewritten as

$$
y_{n+1}(x)=y_{n}(x)+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \frac{(\xi-x)^{\alpha}}{\Gamma(1+\alpha)}\left\{y_{n}^{(2 \alpha)}-y_{n}^{2}-1\right\}(\mathrm{d} \xi)^{\alpha}
$$



Fig. 1. The second-order approximate solution vs. the exact solution when $\alpha=1$. The discontinuous line ( -- ) is the approximate solution when $\alpha=0.9$ and the dotted line (...) is that when $\alpha=0.99$. The continuous line is the exact solution when $\alpha=1$. However, the figure only shows the trend of the approximate solutions on a large scale, which cannot illustrate the local behavior of fractional differential equations.

Taking the initial value

$$
y_{0}=y(0)+\frac{y^{(\alpha)}(0) x^{\alpha}}{\Gamma(1+\alpha)}=\frac{x^{\alpha}}{\Gamma(1+\alpha)},
$$

we can obtain

$$
\begin{aligned}
y_{1}(x)= & y_{0}(x)+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \frac{(\xi-x)^{\alpha}}{\Gamma(1+\alpha)}\left\{y_{0}^{(2 \alpha)}-y_{0}^{2}-1\right\}(\mathrm{d} \xi)^{\alpha} \\
= & \frac{x^{\alpha}}{\Gamma(1+\alpha)}+\frac{x^{2 \alpha}}{\Gamma(1+2 \alpha)}+\frac{\Gamma(1+2 \alpha) x^{4 \alpha}}{\Gamma^{2}(1+\alpha) \Gamma(1+4 \alpha)}, \\
y_{2}(x)= & y_{1}(x)+\frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \frac{(\xi-x)^{\alpha}}{\Gamma(1+\alpha)}\left\{y_{1}^{(2 \alpha)}-y_{1}^{2}-1\right\}(\mathrm{d} \xi)^{\alpha} \\
= & \frac{x^{\alpha}}{\Gamma(1+\alpha)}+\frac{x^{2 \alpha}}{\Gamma(1+2 \alpha)}+\frac{\Gamma(1+2 \alpha) x^{4 \alpha}}{\Gamma^{2}(1+\alpha) \Gamma(1+4 \alpha)}+\frac{2 \Gamma(1+3 \alpha) x^{5 \alpha}}{\Gamma(1+\alpha) \Gamma(1+2 \alpha) \Gamma(1+5 \alpha)} \\
& +\frac{\Gamma(1+4 \alpha) x^{6 \alpha}}{\Gamma^{2}(1+2 \alpha) \Gamma(1+6 \alpha)}+\frac{2 \Gamma(1+5 \alpha) \Gamma(1+2 \alpha) x^{7 \alpha}}{\Gamma^{3}(1+\alpha) \Gamma(1+4 \alpha) \Gamma(1+7 \alpha)}+\frac{2 \Gamma(1+6 \alpha) x^{8 \alpha}}{\Gamma^{2}(1+\alpha) \Gamma(1+4 \alpha) \Gamma(1+8 \alpha)} \\
& +\frac{\Gamma(1+8 \alpha) \Gamma^{2}\left(1+2 \alpha x^{10 \alpha}\right.}{\Gamma^{4}(1+\alpha) \Gamma^{2}(1+4 \alpha) \Gamma(1+10 \alpha)} .
\end{aligned}
$$

If the second-order approximation, $y_{2}(x)$, is sufficient, we can compare the solution with the exact solutions for the case $\alpha=1$ as in Fig. 1.

## 4. Conclusions

In this paper, a fractional variational iteration method is proposed, and proved to be an efficient tool for solving fractional differential equations because the Lagrange multiplier can be identified in a more accurate way using the fractional variational theory. Some other recent work in calculation of variation can be found in Refs. [13-15].

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