# On the conservativeness and the recurrence of symmetric jump-diffusions 

Jun Masamune ${ }^{\text {a }}$, Toshihiro Uemura ${ }^{\text {b }}$, Jian Wang ${ }^{\text {c,* }}$<br>${ }^{\text {a }}$ Department of Mathematics and Statistics, Penn State, Altoona, 3000 Ivyside Park, Altoona, PA 16601, USA<br>b Department of Mathematics, Faculty of Engineering Science, Kansai University, Suita-shi, Osaka 564-8680, Japan<br>${ }^{\text {c }}$ School of Mathematics and Computer Science, Fujian Normal University, 350007, Fuzhou, PR China

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#### Abstract

Sufficient conditions for a symmetric jump-diffusion process to be conservative and recurrent are given in terms of the volume of the state space and the jump kernel of the process. A number of examples are presented to illustrate the optimality of these conditions; in particular, the situation is allowed to be that the state space is topologically disconnected but the particles can jump from a connected component to the other components.


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## 1. Introduction and main results

Let $(X, d, m)$ be a metric measure space. We assume that every metric ball $B(x, r)=$ $\{z \in X: d(x, z)<r\}$ centered at $x \in X$ with radius $r>0$ is pre-compact, and the measure $m$ is a Radon measure with full support. In particular, $X$ is locally compact and separable. Let $(\mathcal{E}, \mathcal{F})$ be a regular symmetric Dirichlet form in $L^{2}(X ; m)$. We denote the extended Dirichlet space of $(\mathcal{E}, \mathcal{F})$ by $\mathcal{F}_{e}$, and a quasi-continuous version of $u \in \mathcal{F}_{e}$ by $\tilde{u}$. According to the Beurling-Deny theorem, see, e.g., [8, Theorem 3.2.1 and Lemma 4.5.4], we can express $(\mathcal{E}, \mathcal{F})$ as follows

$$
\begin{aligned}
\mathcal{E}(u, v)= & \mathcal{E}^{(c)}(u, v)+\iint_{x \neq y}(\tilde{u}(x)-\tilde{u}(y))(\tilde{v}(x)-\tilde{v}(y)) J(d x, d y) \\
& +\int_{X} \tilde{u}(x) \tilde{v}(x) k(d x) \quad \text { for any } u, v \in \mathcal{F}_{e},
\end{aligned}
$$

where $\left(\mathcal{E}^{(c)}, C_{0}(X) \cap \mathcal{F}\right)$ is a strongly-local symmetric form and $C_{0}(X)$ is the space of all real-valued continuous functions on $X$ with compact support; $J$ is a symmetric positive Radon measure on the product space $X \times X$ off the diagonal $\{(x, x): x \in X\}$; and $k$ is a positive Radon measure on $X$.

Let $\mu_{(\cdot, \cdot\rangle}$ be a bounded signed measure, see [8, Lemma 3.2.3], such that

$$
\mathcal{E}^{(c)}(u, v)=\frac{1}{2} \mu_{\langle u, v\rangle}(X)=\frac{1}{2} \int_{X} \mu_{\langle u, v\rangle}(d x) \quad \text { for } u, v \in \mathcal{F}_{e} .
$$

Throughout the paper, we assume the following set (A) of conditions:
(A-1) The killing measure $k$ does not appear; that is, the corresponding process is no killing inside.
(A-2) For each $u, v \in \mathcal{F}_{e}$, the measure $\mu_{\langle u, v\rangle}$ is absolutely continuous with respect to $m$. We denote the corresponding Radon-Nikodym density by $\Gamma^{(c)}(u, v)$; namely,

$$
\mu_{\langle u, v\rangle}(d x)=\Gamma^{(c)}(u, v)(x) m(d x) .
$$

(A-3) The jump measure $J$ has a symmetric kernel $j(x, d y)$ over $X \times \mathcal{B}(X)$ such that

$$
J(d x, d y)=j(x, d y) m(d x)(=j(y, d x) m(d y)=J(d y, d x))
$$

For $u, v \in \mathcal{F}_{e}$, define

$$
\Gamma^{(j)}(u, v)(x)=\int_{x \neq y}(\tilde{u}(x)-\tilde{u}(y))(\tilde{v}(x)-\tilde{v}(y)) j(x, d y),
$$

and

$$
\mathcal{E}^{(j)}(u, v)=\int \Gamma^{(j)}(u, v)(x) m(d x)
$$

Therefore, the form $\mathcal{E}$ has the following expression for any $u, v \in \mathcal{F}_{e}$ :

$$
\begin{aligned}
\mathcal{E}(u, v) & =\mathcal{E}^{(c)}(u, v)+\mathcal{E}^{(j)}(u, v) \\
& =\frac{1}{2} \int_{X} \Gamma^{(c)}(u, v)(x) m(d x)+\int_{X} \Gamma^{(j)}(u, v)(x) m(d x) \\
& =\frac{1}{2} \int_{X} \Gamma^{(c)}(u, v)(x) m(d x)+\iint_{x \neq y}(\tilde{u}(x)-\tilde{u}(y))(\tilde{v}(x)-\tilde{v}(y)) j(x, d y) m(d x) .
\end{aligned}
$$

Let $\psi_{K}$ be the distance function from a compact set $K$ of $X$, i.e., $\psi_{K}(\cdot)=\inf _{y \in K} d(\cdot, y)$. For every $r>0$, we denote $B(K, r)=\left\{x \in X: \psi_{K}<r\right\}$ and its closure $\left\{x \in X: \psi_{K} \leqslant r\right\}$ by $\bar{B}(K, r)$. Clearly, $B(K, r)$ is pre-compact. Let $\mathcal{F}_{\text {loc }}$ be the set of measurable functions $u$ such that for each relatively compact open set $G$ of $X$ there exists $w \in \mathcal{F}$ which satisfies that $\left.u\right|_{G}=\left.w\right|_{G}$ $m$-a.e. Additionally, we assume the following set (M) of conditions so that both $\mathcal{E}^{(c)}$ and $\mathcal{E}^{(j)}$ are compatible with the distance $d$ :
(M-1) $\psi_{K} \in \mathcal{F}_{\text {loc }}$ for every compact set $K \subset X$,
(M-2) $M_{c}:=\operatorname{ess}_{\sup _{x \in X^{(c)}}} \Gamma^{(c)}(d, d)(x)<\infty$,
(M-3) $M_{j}:=\operatorname{ess} \sup _{x \in X^{(j)}} \int_{x \neq y}\left(1 \wedge d^{2}(x, y)\right) j(x, d y)<\infty$,
where $X^{(c)}=\left\{x \in X: \Gamma^{(c)} \neq 0\right\}$ and $X^{(j)}=\left\{x \in X: \Gamma^{(j)} \neq 0\right\}$.
There are many classical examples of symmetric diffusions or symmetric pure jump processes whose Dirichlet form satisfies conditions (A) and (M): for instance, strongly-local Dirichlet forms on a metric measure space, whose distance is the Carnot-Carathéodori distance associated with the Dirichlet form. This includes canonical Dirichlet forms on Riemannian manifolds, CR manifolds, sub-Riemannian manifolds, and weighted manifolds; divergence type operators with bounded coefficients on Euclidean spaces; the sum of squares of vector fields satisfying Hörmader's condition, the quantum graphs, and pre-fractals. Other examples are symmetric $\alpha$ stable Lévy processes with $\alpha \in(0,2)$ on Euclidean spaces, and symmetric random walks on graphs.

Let $A$ be the generator of $(\mathcal{E}, \mathcal{F})$ in $L^{2}(X ; m)$. We denote the associated semigroup and the resolvent by $\left(T_{t}\right)_{t \geqslant 0}=\left(e^{t A}\right)_{t \geqslant 0}$ and $G=\int_{0}^{\infty} T_{t} d t$, respectively. The Dirichlet form $(\mathcal{E}, \mathcal{F})$ is called conservative if

$$
T_{t} 1 \equiv 1, \quad m \text {-a.e. for any } t>0
$$

and recurrent if

$$
G f(x) \equiv 0 \text { or } \infty \quad \text { for any } f \in L_{+}^{1}(X ; m) \text { and } m \text {-a.e. } x \in X
$$

It is a classical result that Brownian motion on $\mathbb{R}^{n}$ is conservative for any $n \geqslant 1$ and is recurrent if and only if $n=1,2$. This result has been generalized to the Wiener process of complete Riemannian manifolds, and one of the most important discoveries is that a certain bound of the volume at infinity - rather than the dimension - implies these properties. This fact was first found by M.P. Gaffney [10] for the conservativeness, and it has been refined by various methods in [1,23,36,17,5,14]. Especially, R. Azencott [1] and A. Grigor'yan [14] demonstrated that the conservativeness may fail without a condition on the curvature or volume. On the other hand, the recurrence of the Wiener process of Riemannian manifolds or jump processes has been investigated by several authors in [4,22,38,11,12,28]. Furthermore, K.-T. Sturm [35] extended the theory to a general strongly-local regular Dirichlet form on a metric measure space equipped with the Carnot-Carathéodori distance.

Recently, there has been a tremendous amount of work devoted to the conservation property of a non-local Dirichlet form; for instance, the physical Laplacian on an infinite graph [7,6,39-41,24,18-20] and non-local Dirichlet forms [26,15,33]; however, as far as the authors know, there is only one result by Z.-Q. Chen and T. Kumagai [3] for the Dirichlet form which has both the strongly-local and non-local terms. Due to its nature, the associated process is called a jump-diffusion process.

Our first main purpose is to investigate the conservative property of a jump-diffusion process. For any $x \in X$ and $r>0$, the volume of $\bar{B}(x, r)$ is denoted by $V(x, r)$.

## Theorem 1.1. If

$$
\begin{equation*}
\liminf _{r \rightarrow \infty} \frac{\ln V\left(x_{0}, r\right)}{r \ln r}<\infty, \tag{1.1}
\end{equation*}
$$

for some $x_{0} \in X$, then $(\mathcal{E}, \mathcal{F})$ is conservative.
This result was obtained for a non-local Dirichlet form in [15, Theorem 1.1], where the lefthand side of (1.1) is required to be less than $1 / 2$. Let us explain the significance of removing the constant $1 / 2$ by comparing the uniqueness class with the conservation property. Let $\mathcal{U}$ be the set of the solutions to the Cauchy problem of the heat equation with zero initial data. If any $u \in \mathcal{U}$ is identically 0 , then $\mathcal{U}$ is called a uniqueness class. Under an integrability assumption, determining the uniqueness class implies the conservativeness of Riemannian manifolds [13], Dirichlet forms [35], and graphs [20]. In fact, A. Grigor'yan [13] and K.-T. Sturm [35] established the sharp conservation test for complete Riemannian manifolds and strongly-local Dirichlet forms, respectively, in this way. However, X. Huang [20, Section 3.3] constructed an example of a graph, which verifies that the constant $1 / 2$ is indeed needed for the uniqueness class. Therefore, Theorem 1.1 together with Huang's example demonstrates that the uniqueness class condition is really stronger than the conservation property for a graph.

Next, we turn to the recurrence. For any $x \in X$ and $r>0$, the volumes of the closed ball $\bar{B}(x, r)$ intersected with $X^{(c)}$ and $X^{(j)}$ are denoted by $V^{(c)}(x, r)$ and $V^{(j)}(x, r)$, respectively. For $r>0$, define

$$
\omega(r)=\sup _{x \in X^{(j)}} \int_{x \neq y}(d(x, y) \wedge r)^{2} j(x, d y)
$$

Our second main result is
Theorem 1.2. If

$$
\begin{equation*}
\liminf _{r \rightarrow \infty} \frac{1}{r^{2}}\left[V^{(c)}\left(x_{0}, r\right)+V^{(j)}\left(x_{0}, r\right) \omega(r)\right]<\infty \tag{1.2}
\end{equation*}
$$

for some $x_{0} \in X$, then $(\mathcal{E}, \mathcal{F})$ is recurrent.
Theorem 1.2 was proven in the case of the Wiener process (namely, the process does not jump) on a complete Riemannian manifold by S.Y. Cheng and S.T. Yau [4]. Theorem 1.2 is sharp for an isotropic symmetric $\alpha$-stable Lévy process on $\mathbb{R}^{n}$, see, e.g., [30, Corollary 37.17 and Theorem 37.18] or Example 5.2 in Section 5. Here, let us mention that [30, Corollary 37.17 and Theorem 37.18] are derived from the characteristic functions of the associated processes, see [32] for the recent development on this topic; while Theorem 1.2 is based on the theory of Dirichlet forms.

This paper is organized as follows. Section 2 is devoted to the preliminaries. Here we establish an integral-derivation type property for a Dirichlet form of jump-process type, which is a technical key to prove the conservation property. The main results, Theorems 1.1 and 1.2, are proved in Sections 3 and 4, respectively. Finally, in Section 5 we present some examples of symmetric jump-diffusions to illustrate the power of our main theorems.

## 2. Preliminaries: the integral-derivation property

In this section, we first prepare the preliminaries and then proceed to establish an integralderivation type property for a Dirichlet form with jump-diffusion type. This will be used to prove the conservation property in the next section.

We begin with the following quite elementary fact.
Lemma 2.1. If $u \in \mathcal{F}_{\text {loc }} \cap L^{\infty}$ has compact support, where $L^{\infty}=L^{\infty}(X)$ is the space of realvalued bounded measurable functions on $X$, then $u \in \mathcal{F} \cap L^{\infty}$.

Proof. Suppose that supp $u \subset K$ with a compact set $K$. Let $\eta \in \mathcal{F} \cap L^{\infty}$ agree with $u$ on $B(K, 1)$. Because of the regularity and the fact that the constant function belongs to $\mathcal{F}_{\text {loc }}$, see the remark in [8, p. 117], there is a function $\chi \in \mathcal{F} \cap L^{\infty}$ such that $\left.\chi\right|_{K}=1$ and supp $\chi \subset B(K, 1)$. Since $\eta \chi \in \mathcal{F}$ and $u=\eta \chi$, the statement follows.

For the sake of simplicity, hereafter we denote $\Gamma[\cdot]=\Gamma(\cdot, \cdot), \mathcal{E}[\cdot]=\mathcal{E}(\cdot, \cdot)$, etc. We say that the jump range of $\mathcal{E}$ or $\mathcal{E}^{(j)}$ is uniformly bounded, if there exists a constant $a>0$ such that $\operatorname{supp}(j(x, \cdot)) \subset B(x, a)$ for every $x \in X$.

Lemma 2.2. Suppose that the jump range of $\mathcal{E}$ is uniformly bounded. If $u \in \mathcal{F}_{\text {loc }} \cap L^{\infty}$ is constant outside a compact set, then for any $v \in \mathcal{F} \cap L^{\infty}, u v \in \mathcal{F} \cap L^{\infty}$.

Proof. Let $K \subset X$ be a compact set such that $u$ is constant outside it. Consider the sequence of cut-off functions $\left(\chi_{l}\right)_{l \in \mathbb{N}}$, where for $l \geqslant 1$,

$$
\chi_{l}=\left(\left(2-l^{-1} \psi\right) \wedge 1\right)_{+}
$$

By Lemma 2.1, the function $\chi_{l}$ belongs to $\mathcal{F}$ for any $l \geqslant 1$. Obviously, $\chi_{l}=1$ on $B(K, l)$ and $\operatorname{supp}\left(\chi_{l}\right) \subset \bar{B}(K, 2 l)$.

We set for any $l \geqslant 1, v_{l}=u v \chi_{l}$. Since $u \in \mathcal{F}_{\text {loc }} \cap L^{\infty}$ and $v \in \mathcal{F} \cap L^{\infty}, v_{l}$ belongs to $\mathcal{F}_{\text {loc }} \cap L^{\infty}$ and has compact support. Hence, Lemma 2.1 shows that $v_{l} \in \mathcal{F}$ for any $l \geqslant 1$.

Next, we claim that the sequence $\left(v_{l}\right)_{l \geqslant 1}$ is $\mathcal{E}$-Cauchy. Set $\chi_{l, l^{\prime}}=\chi_{l}-\chi_{l^{\prime}}$ for $l, l^{\prime} \geqslant 1$. Since the jump range of $\mathcal{E}$ is uniformly bounded, for large enough $l$ and $l^{\prime}$,

$$
\mathcal{E}\left[v_{l}-v_{l^{\prime}}\right]=\mathcal{E}\left[\left(\chi_{l}-\chi_{l^{\prime}}\right) u v\right]=\kappa \cdot \mathcal{E}\left[\chi_{l, l^{\prime}} v\right],
$$

where $\kappa=\left.u\right|_{K^{c}}$. By [8, Lemma 3.2.5],

$$
\mathcal{E}^{(c)}\left[\chi_{l, l^{\prime}} v\right] \leqslant 2 \int v^{2} \Gamma^{(c)}\left[\chi_{l, l^{\prime}}\right] d m+2 \int \chi_{l, l^{\prime}}^{2} \Gamma^{(c)}[v] d m
$$

Because of (M) and the chain rule of the strongly-local Dirichlet form, see, e.g., [35, p. 190], $\Gamma^{(c)}\left[\chi_{l, l^{\prime}}\right] \rightarrow 0$ as $l, l^{\prime} \rightarrow \infty$. This together with the fact $\chi_{l, l^{\prime}} \rightarrow 0$ as $l, l^{\prime} \rightarrow \infty$ yields that $\mathcal{E}^{(c)}\left[\chi_{l, l^{\prime}} v\right]$ tends to zero as $l, l^{\prime} \rightarrow \infty$.

On the other hand,

$$
\begin{aligned}
\mathcal{E}^{(j)}\left[\chi_{l, l^{\prime}} v\right] \leqslant & 2 \int v^{2}(x) \int\left(\chi_{l, l^{\prime}}(x)-\chi_{l, l^{\prime}}(y)\right)^{2} j(x, d y) m(d x) \\
& +2 \iint \chi_{l, l^{\prime}}^{2}(y)(v(x)-v(y))^{2} j(x, d y) m(d x) \\
= & (I)+(I I) .
\end{aligned}
$$

For any $x \in X$,

$$
\begin{aligned}
\int & \left(\chi_{l, l^{\prime}}(x)-\chi_{l, l^{\prime}}(y)\right)^{2} j(x, d y) \\
& =\int\left(\left(\chi_{l}(x)-\chi_{l}(y)\right)-\left(\chi_{l^{\prime}}(x)-\chi_{l^{\prime}}(y)\right)\right)^{2} j(x, d y) \\
& \leqslant 2 \int\left(\chi_{l}(x)-\chi_{l}(y)\right)^{2} j(x, d y)+2 \int\left(\chi_{l^{\prime}}(x)-\chi_{l^{\prime}}(y)\right)^{2} j(x, d y) \\
& \leqslant 2\left(l^{-2}+l^{\prime-2}\right) \int d(x, y)^{2} j(x, d y)
\end{aligned}
$$

Combining the fact that $\operatorname{supp}(j(x, d y)) \subset B(x, a)$ for all $x \in X$ and some $a>0$ with the assumption $(\mathrm{M})$, the last term in the right-hand side of the equation above is dominated by

$$
2\left(1+a^{2}\right) M_{j}\left(l^{-2}+l^{\prime-2}\right)
$$

which tends to 0 as $l, l^{\prime} \rightarrow \infty$. Hence $(I) \rightarrow 0$ as $l, l^{\prime} \rightarrow \infty$. Since $\chi_{l, l^{\prime}} \rightarrow 0, m$-a.e. as $l, l^{\prime} \rightarrow \infty$, $(I I) \rightarrow 0$ as $l, l^{\prime} \rightarrow \infty$. Thus, $\mathcal{E}^{(j)}\left[\chi_{l, l^{\prime}} v\right] \rightarrow 0$ as $l, l^{\prime} \rightarrow \infty$, and so the desired claim follows.

Finally, since $v_{l} \rightarrow u v, m$-a.e. as $l \rightarrow \infty, u v \in \mathcal{F}_{e}$. This together with the fact $u v \in L^{2}$ and [8, Theorem 1.5.2(iii)] yields that $u v \in \mathcal{F}$.

The following is the integral-derivation property for our Dirichlet form.
Lemma 2.3. Suppose that the jump range of $\mathcal{E}$ is uniformly bounded. If $u \in \mathcal{F} \cap L^{\infty}$ and $\phi \in$ $\mathcal{F}_{\text {loc }} \cap L^{\infty}$ is constant outside a compact set, then

$$
\begin{equation*}
\mathcal{E}(u, u \phi)=\int u \Gamma(u, \phi) d m+\int \phi \Gamma[u] d m \tag{2.3}
\end{equation*}
$$

where $\Gamma=\frac{1}{2}\left(\Gamma^{(c)}+\Gamma^{(j)}\right)$.
Proof. According to Lemma 2.2, $u \phi \in \mathcal{F}$. By the derivation property of $\mathcal{E}^{(c)}$, see, e.g., [8, Lemma 3.2.5 and the note on p. 117],

$$
\int \Gamma^{(c)}(u, u \phi) d m=\int u \Gamma^{(c)}(u, \phi) d m+\int \phi \Gamma^{(c)}[u] d m
$$

Next, by the integral property of a non-local Dirichlet form, see [27, Proposition 2.2], we have

$$
\int \Gamma^{(j)}(u, u \phi) d m=\int u \Gamma^{(j)}(u, \phi) d m+\int \phi \Gamma^{(j)}[u] d m .
$$

Combining the two identities, we obtain (2.3).

## 3. Proof of Theorem 1.1: the conservation property

The aim of this section is to prove Theorem 1.1. For any $a>0$, consider a symmetric form $\left(\mathcal{E}^{(j, a)}, \mathcal{F}\right)$ defined by

$$
\mathcal{E}^{(j, a)}[u]=\iint(u(x)-u(y))^{2} \mathbb{1}_{\{d(x, y) \leqslant a\}} j(x, d y) m(d x) \quad \text { for } u \in \mathcal{F}
$$

Under the condition $(\mathrm{M}),\left(\mathcal{E}^{(j, a)}+\mathcal{E}^{(c)}, \mathcal{F}\right)$ is a regular Dirichlet form, and it is conservative if and only if so is $(\mathcal{E}, \mathcal{F})$, see [31, Section 4] and [26, Section 3]. Clearly, $\left(\mathcal{E}^{(j, a)}, \mathcal{F}\right)$ has uniformly bounded range. Therefore, in order to prove the conservation property, we may and do assume that $\mathcal{E}$ has uniformly bounded jump range. More precisely, we suppose that there exists a constant $a>0$ such that

$$
j(x, d y)=\mathbb{1}_{B(x, a)}(y) j(x, d y) \quad \text { for all } x \in X .
$$

Our proof is basically the Davies method [5], which was used also in [15]; however, we are able to get a better result because of the choice of $a$. In this section, the constant $a$ will be

$$
\begin{equation*}
a=a\left(x_{0}, m\right):=\left[8 \liminf _{r \rightarrow \infty} \frac{\log V\left(x_{0}, r\right)}{r \log r}+9\right]^{-1} \tag{3.4}
\end{equation*}
$$

where $x_{0} \in X$ is the reference point in Theorem 1.1. For $f \in C_{0}(X)$ with $f \geqslant 0$, set

$$
\psi(x)=d(x, \operatorname{supp}(f))
$$

and

$$
\phi(x)=e^{\alpha \psi(x)}
$$

where $\alpha>0$ is a constant determined later. Note that if $n \geqslant 1$ and $x \in X$ satisfy

$$
n \geqslant a^{-1}\left[4 a+2 d\left(x_{0}, \operatorname{supp}(f)\right)\right] \quad \text { and } \quad(n-2) a \leqslant d\left(x, x_{0}\right) \leqslant(n+1) a
$$

then

$$
\psi(x) \geqslant d\left(x, x_{0}\right)-d\left(x_{0}, \operatorname{supp}(f)\right) \geqslant(n-2) a-d\left(x_{0}, \operatorname{supp}(f)\right) \geqslant a n / 2
$$

and so

$$
\begin{equation*}
\phi(x)=e^{\alpha \psi(x)} \geqslant e^{a \alpha n / 2} \tag{3.5}
\end{equation*}
$$

For the function $f$ above and any $t \geqslant 0$, we denote $u_{t}=T_{t} f$. Since $\left(T_{t}\right)_{t \geqslant 0}$ is analytic, $u_{t}$ belongs to the domain of the $L^{2}$-generator $A$ of $(\mathcal{E}, \mathcal{F})$; in particular, $u_{t} \in \mathcal{F} \cap L^{\infty}$ for any $t>0$.

The following lemma provides the key estimate.
Lemma 3.1. Using the notations above, for any $t \geqslant 0$,

$$
\begin{equation*}
\int_{0}^{t} \int \phi \Gamma\left[u_{s}\right] d m d s \leqslant 2 e^{\gamma t}\left\|\phi^{1 / 2} f\right\|_{2}^{2} \tag{3.6}
\end{equation*}
$$

where $\gamma=\alpha^{2}\left(e^{2 \alpha a}+1\right) M / 2$ and $M=M_{c} \vee M_{j}$.
Proof. In the following, we denote the norm and the inner product of $L^{2}(X ; m)$ by $\|\cdot\|_{2}$ and $\langle\cdot, \cdot\rangle$, respectively. For any $n \geqslant 1$, set

$$
\phi_{n}(x)=e^{\alpha(\psi(x) \wedge n)}
$$

Since $\psi \in \mathcal{F}_{\text {loc }}$, we may apply an argument in [8, pp. 116-117] to deduce that $\phi_{n} \in \mathcal{F}_{\text {loc }}$ for every $n \geqslant 1$. Taking into account that $\psi \in L^{\infty}$ is constant outside a compact set, Lemma 2.2 shows that for every $t>0$ and $n \geqslant 1, u_{t} \phi_{n} \in \mathcal{F}$. Therefore, by Lemma 2.3, for all $t>0$,

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t}\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2} & =\left\langle\dot{u}_{t}, \phi_{n} u_{t}\right\rangle \\
& =-\mathcal{E}\left(u_{t}, \phi_{n} u_{t}\right) \\
& =-\int \phi_{n} \Gamma\left[u_{t}\right] d m-\int u_{t} \Gamma\left(u_{t}, \phi_{n}\right) d m \\
& \leqslant-\int \phi_{n} \Gamma\left[u_{t}\right] d m+\left|\int u_{t} \Gamma\left(u_{t}, \phi_{n}\right) d m\right|
\end{aligned}
$$

where $\dot{u}_{t}=\frac{d}{d t} u_{t}$. This is,

$$
\begin{equation*}
\int \phi_{n} \Gamma\left[u_{t}\right] d m \leqslant\left|\int u_{t} \Gamma\left(u_{t}, \phi_{n}\right) d m\right|-\frac{1}{2} \frac{d}{d t}\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2} \tag{3.7}
\end{equation*}
$$

Next, we estimate the first term on the right side of this equation. For every $x \in X$, according to the Cauchy-Schwarz inequality,

$$
\begin{aligned}
\left|\Gamma^{(j)}\left(u_{t}, \phi_{n}\right)(x)\right| & =\left|\int\left(u_{t}(x)-u_{t}(y)\right)\left(\phi_{n}(x)-\phi_{n}(y)\right) j(x, d y)\right| \\
& \leqslant \sqrt{\int\left(u_{t}(x)-u_{t}(y)\right)^{2} j(x, d y)} \sqrt{\int\left(\phi_{n}(x)-\phi_{n}(y)\right)^{2} j(x, d y)} \\
& =\sqrt{\Gamma^{(j)}\left[u_{t}\right](x)} \sqrt{\Gamma^{(j)}\left[\phi_{n}\right](x)} .
\end{aligned}
$$

By the Cauchy-Schwarz inequality again,

$$
\begin{aligned}
\left|\int u_{t} \Gamma^{(j)}\left(u_{t}, \phi_{n}\right) d m\right| & \leqslant \int \phi_{n}^{1 / 2} \sqrt{\Gamma^{(j)}\left[u_{t}\right]} \phi_{n}^{-1 / 2} \sqrt{u_{t}^{2} \Gamma^{(j)}\left[\phi_{n}\right]} d m \\
& \leqslant \sqrt{\int \phi_{n} \Gamma^{(j)}\left[u_{t}\right] d m} \sqrt{\int \phi_{n}^{-1} u_{t}^{2} \Gamma^{(j)}\left[\phi_{n}\right] d m}
\end{aligned}
$$

Since

$$
\left|e^{\alpha r}-1\right| \leqslant \alpha e^{\alpha a}|r| \quad \text { for any } r \in(0, a]
$$

it follows that

$$
\left|\phi_{n}(x)-\phi_{n}(y)\right| \leqslant \alpha e^{\alpha a} \phi_{n}(x) d(x, y) \quad \text { for any } x, y \in X \text { with } d(x, y) \leqslant a,
$$

and so

$$
\Gamma^{(j)}\left[\phi_{n}\right](x) \leqslant\left(\alpha e^{\alpha a} \phi(x)\right)^{2} \int d^{2}(x, y) j(x, d y) \quad \text { for every } x \in X
$$

Since $\operatorname{supp}(j(x, d y)) \subset B(x, a)$ for any $x \in X$ and some constant $a \in(0,1)$, we get

$$
\begin{aligned}
\int \phi_{n}^{-1} u_{t}^{2} \Gamma^{(j)}\left[\phi_{n}\right] d m & \leqslant \alpha^{2} e^{2 \alpha a} \int \phi_{n}(x) u_{t}^{2}(x) \int d(x, y)^{2} j(x, d y) m(d x) \\
& \leqslant \alpha^{2} e^{2 \alpha a} \int \phi_{n}(x) u_{t}^{2}(x) \int(d(x, y) \wedge a)^{2} j(x, d y) m(d x) \\
& \leqslant M_{j} \alpha^{2} e^{2 \alpha a} \int \phi_{n} u_{t}^{2} d m
\end{aligned}
$$

Therefore, for any $\lambda>0$,

$$
\begin{aligned}
\left|\int u_{t} \Gamma^{(j)}\left(u_{t}, \phi_{n}\right) d m\right| & \leqslant \sqrt{M_{j} \int \phi_{n} \Gamma^{(j)}\left[u_{t}\right] d m} \sqrt{\alpha^{2} e^{2 \alpha a} \int \phi_{n} u_{t}^{2} d m} \\
& \leqslant \frac{M_{j}}{2 \lambda} \int \phi_{n} \Gamma^{(j)}\left[u_{t}\right] d m+\frac{\lambda \alpha^{2} e^{2 \alpha a}}{2} \int \phi_{n} u_{t}^{2} d m \\
& =\frac{M_{j}}{2 \lambda} \int \phi_{n} \Gamma^{(j)}\left[u_{t}\right] d m+\frac{\lambda \alpha^{2} e^{2 \alpha a}}{2}\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2}
\end{aligned}
$$

where in the last inequality we have used the fact that $2 \xi \eta \leqslant \lambda^{-1} \xi^{2}+\lambda \eta^{2}$ for any $\xi, \eta \geqslant 0$ and $\lambda>0$.

On the other hand, we apply the argument above for the local term to get that

$$
\left|\int u_{t} \Gamma^{(c)}\left(u_{t}, \phi_{n}\right) d m\right| \leqslant \sqrt{\int \phi_{n} \Gamma^{(c)}\left[u_{t}\right] d m} \sqrt{\int \phi_{n}^{-1} u_{t}^{2} \Gamma^{(c)}\left[\phi_{n}\right] d m}
$$

According to the chain rule for a strongly-local Dirichlet form, see, e.g., [35, p. 190],

$$
\int \phi_{n}^{-1} u_{t}^{2} \Gamma^{(c)}\left[\phi_{n}\right] d m \leqslant \alpha^{2} \int u_{t}^{2} \phi_{n} \Gamma^{(c)}[d] d m
$$

which along with the assumption (M) gives us

$$
\int \phi_{n}^{-1} u_{t}^{2} \Gamma^{(c)}\left[\phi_{n}\right] d m \leqslant M_{c} \alpha^{2} \int u_{t}^{2} \phi_{n} d m .
$$

We again follow the argument above to obtain the estimate:

$$
\left|\int u_{t} \Gamma^{(c)}\left(u_{t}, \phi_{n}\right) d m\right| \leqslant \frac{M_{c}}{2 \lambda} \int \phi_{n} \Gamma^{(c)}\left[u_{t}\right] d m+\frac{\lambda \alpha^{2}}{2}\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2} \quad \text { for any } \lambda>0
$$

Combining the estimates for the non-local and strongly-local terms, we get that

$$
\left|\int u_{t} \Gamma\left(u_{t}, \phi_{n}\right) d m\right| \leqslant \frac{M}{2 \lambda} \int \phi_{n} \Gamma\left[u_{t}\right] d m+\frac{\lambda \alpha^{2}\left(e^{2 \alpha a}+1\right)}{2}\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2}
$$

By applying this inequality for (3.7), we have

$$
\begin{equation*}
\left(2-\frac{M}{\lambda}\right) \int \phi_{n} \Gamma\left[u_{s}\right] d m \leqslant \lambda \alpha^{2}\left(e^{2 \alpha a}+1\right)\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2}-\frac{d}{d s}\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2} \tag{3.8}
\end{equation*}
$$

If we integrate this with respect to $s$ over $[0, t]$, then

$$
\begin{align*}
& \left(2-\frac{M}{\lambda}\right) \iint_{0}^{t} \int \phi_{n} \Gamma\left[u_{s}\right] d m \\
& \quad \leqslant \lambda \alpha^{2}\left(e^{2 \alpha a}+1\right) \int_{0}^{t}\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2} d s-\left(\left\|\phi_{n}^{1 / 2} u_{t}\right\|_{2}^{2}-\left\|\phi_{n}^{1 / 2} f\right\|_{2}^{2}\right) \tag{3.9}
\end{align*}
$$

We estimate $\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2}$ for any $s \leqslant t$ by first letting $\lambda=M / 2$ in (3.8),

$$
\frac{d}{d s}\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2} \leqslant \frac{M \alpha^{2}\left(e^{2 \alpha a}+1\right)}{2}\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2}
$$

and then, by applying the Gronwall inequality:

$$
\left\|\phi_{n}^{1 / 2} u_{s}\right\|_{2}^{2} \leqslant \exp \left(\frac{M \alpha^{2}\left(e^{2 \alpha a}+1\right) s}{2}\right)\left\|\phi_{n}^{1 / 2} f\right\|_{2}^{2}
$$

Substituting this into (3.9), we have

$$
\begin{aligned}
& \left(2-\frac{M}{\lambda}\right) \int_{0}^{t} \int \phi_{n} \Gamma\left[u_{s}\right] d m d s \\
& \quad \leqslant\left\|\phi_{n}^{1 / 2} f\right\|_{2}^{2}+\frac{2 \lambda}{M}\left[\exp \left(M \alpha^{2}\left(e^{2 \alpha a}+1\right) t / 2\right)-1\right]\left\|\phi_{n}^{1 / 2} f\right\|_{2}^{2}
\end{aligned}
$$

Setting $\lambda=M$, this becomes

$$
\int_{0}^{t} \int \phi_{n} \Gamma\left[u_{s}\right] d m d s \leqslant 2 \exp \left(M \alpha^{2}\left(e^{2 \alpha a}+1\right) t / 2\right)\left\|\phi_{n}^{1 / 2} f\right\|_{2}^{2}
$$

The required assertion (3.6) follows by letting $n \rightarrow \infty$.
We are in a position to prove Theorem 1.1.
Proof of Theorem 1.1. We adopt the notations in the proof of Lemma 3.1. Define a cut-off function $g_{n}$ for any $n \geqslant 1$ as follows

$$
g_{n}(x):=\left(\left(n-a^{-1} d\left(x, x_{0}\right)\right) \wedge 1\right)_{+}
$$

By Lemma 2.1, $g_{n}$ belongs to $\mathcal{F}$. To the end of the proof, we show that there exists a sequence $\left(n_{k}\right)_{k \geqslant 0}$ such that $n_{k} \rightarrow \infty$ as $k \rightarrow \infty$, and for every $t>0$,

$$
\int_{0}^{t}\left\langle\dot{u}_{s}, g_{n_{k}}\right\rangle d s \rightarrow 0 \quad \text { as } k \rightarrow \infty
$$

Indeed, we can deduce from this and the dominated convergence theorem that

$$
\left\langle T_{t} f, 1\right\rangle=\lim _{k \rightarrow \infty}\left\langle u_{t}, g_{n_{k}}\right\rangle=\lim _{k \rightarrow \infty}\left\langle f, g_{n_{k}}\right\rangle+\lim _{k \rightarrow \infty} \int_{0}^{t}\left\langle\dot{u}_{s}, g_{n_{k}}\right\rangle d s=\langle f, 1\rangle
$$

which immediately implies the conservation property.

Since $\left(u_{s}\right)_{s>0}$ solves the heat equation and $g_{n} \in \mathcal{F}$,

$$
\begin{equation*}
\int_{0}^{t}\left\langle\dot{u}_{s}, g_{n}\right\rangle d s=-\int_{0}^{t} \mathcal{E}\left(u_{s}, g_{n}\right) d s=-\int_{0}^{t}\left(\mathcal{E}^{(c)}\left(u_{s}, g_{n}\right)+\mathcal{E}^{(j)}\left(u_{s}, g_{n}\right)\right) d s \tag{3.10}
\end{equation*}
$$

First, we estimate the second term, the harder one, on the right side. For any $t>0$,

$$
\begin{align*}
\left|\int_{0}^{t} \mathcal{E}^{(j)}\left(u_{s}, g_{n}\right) d s\right| & \leqslant \int_{0}^{t}\left|\int \Gamma^{(j)}\left(u_{s}, g_{n}\right) d m\right| d s \\
& \leqslant \int_{0}^{t}\left[\int \sqrt{\Gamma^{(j)}\left[u_{s}\right]} \sqrt{\Gamma^{(j)}\left[g_{n}\right]} d m\right] d s \\
& =\int_{0}^{t}\left[\int \sqrt{\phi \Gamma^{(j)}\left[u_{s}\right]} \sqrt{\phi^{-1} \Gamma^{(j)}\left[g_{n}\right]} d m\right] d s \\
& \leqslant \int_{0}^{t} \sqrt{\int \phi \Gamma^{(j)}\left[u_{s}\right] d m} \sqrt{\int \phi^{-1} \Gamma^{(j)}\left[g_{n}\right] d m d s} \\
& \leqslant \sqrt{\iint \phi \Gamma^{(j)}\left[u_{s}\right] d m d s} \sqrt{\iint \phi_{0}^{-1} \Gamma^{(j)}\left[g_{n}\right] d m d s} \\
& =\sqrt{\int_{0}^{t} \int \phi \Gamma^{(j)}\left[u_{s}\right] d m d s} \sqrt{t \int \phi^{-1} \Gamma^{(j)}\left[g_{n}\right] d m} \tag{3.11}
\end{align*}
$$

where all the inequalities above follow from the Cauchy-Schwarz inequality. For any $n>0$, let $A_{n}$ denote the following annulus associated with the constant $a$

$$
A_{n}=A_{n}(a)=\bar{B}\left(x_{0},(n+1) a\right) \backslash B\left(x_{0},(n-2) a\right)
$$

Since $\operatorname{supp}\left(g_{n}\right) \subset B\left(x_{0}, n a\right)$ and $\operatorname{supp}(j(x, d y)) \subset B(x, a)$ for all $x \in X$, it holds that if $x \notin A_{n}$,

$$
\Gamma^{(j)}\left[g_{n}\right](x)=\int\left(g_{n}(x)-g_{n}(y)\right)^{2} j(x, d y)=0
$$

if $x \in A_{n}$,

$$
\begin{aligned}
\Gamma^{(j)}\left[g_{n}\right](x) & \leqslant a^{-2} \int d(x, y)^{2} j(x, d y) \\
& \leqslant a^{-2} \int(d(x, y) \wedge a)^{2} j(x, d y) \\
& \leqslant a^{-2} M_{j}
\end{aligned}
$$

where in the last inequality we have used the fact that $0<a<1$. Choosing $n$ large enough so that $n \geqslant a^{-1}\left[4 a+2 d\left(x_{0}, \operatorname{supp}(f)\right)\right]$, we get from (3.5) that

$$
\begin{aligned}
\int \phi^{-1} \Gamma^{(j)}\left[g_{n}\right] d m & =\int_{A_{n}} \phi^{-1} \Gamma^{(j)}\left[g_{n}\right] d m \\
& \leqslant a^{-2} M_{j} e^{-a \alpha n / 2} m\left(A_{n}\right)
\end{aligned}
$$

Therefore, by (3.11),

$$
\left|\int_{0}^{t} \mathcal{E}^{(j)}\left(u_{s}, g_{n}\right) d s\right|^{2} \leqslant a^{-2} t M_{j} e^{-a \alpha n / 2} m\left(A_{n}\right) \int_{0}^{t} \int \phi \Gamma^{(j)}\left[u_{s}\right] d m d s
$$

In a similar way, we can prove that

$$
\left|\int_{0}^{t} \mathcal{E}^{(c)}\left(u_{s}, g_{n}\right) d s\right|^{2} \leqslant a^{-2} t M_{c} e^{-a \alpha n / 2} m\left(A_{n}\right) \iint_{0}^{t} \int \phi \Gamma^{(c)}\left[u_{s}\right] d m d s
$$

Therefore,

$$
\begin{equation*}
\left|\int_{0}^{t} \mathcal{E}\left(u_{s}, g_{n}\right) d s\right|^{2} \leqslant 2 a^{-2} t M e^{-a \alpha n / 2} m\left(A_{n}\right) \int_{0}^{t} \int \phi \Gamma\left[u_{s}\right] d m d s . \tag{3.12}
\end{equation*}
$$

We now apply (3.12) and Lemma 3.1 for (3.10) to get that

$$
\begin{align*}
& \left|\int_{0}^{t}\left\langle\dot{u}_{s}, g_{n}\right\rangle d s\right|^{2} \\
& \quad \leqslant 2 a^{-2} t M e^{-a \alpha n / 2} m\left(A_{n}\right) \int_{0}^{t} \int \phi \Gamma\left[u_{s}\right] d m d s \\
& \quad \leqslant 4 a^{-2} t M\left\|\phi^{1 / 2} f\right\|_{2}^{2} \exp \left(\frac{M \alpha^{2}\left(e^{2 \alpha a}+1\right) t}{2}-\frac{\alpha a n}{2}+\log m\left(A_{n}\right)\right) . \tag{3.13}
\end{align*}
$$

Finally, we estimate (3.13) by applying the volume assumption (1.1). Indeed, according to (1.1), there exists a sequence $\left(n_{k}\right)_{k \geqslant 1}$ such that $n_{k} \rightarrow \infty$ as $k \rightarrow \infty$, and for a large enough $k \geqslant 1$,

$$
\begin{aligned}
\log m\left(A_{n_{k}}\right) & \leqslant \log V\left(x_{0},\left(n_{k}+1\right) a\right) \\
& \leqslant\left(c_{3}-1 / 2\right)\left(\left(n_{k}+1\right) a\right) \log \left(\left(n_{k}+1\right) a\right) \\
& \leqslant a c_{3} n_{k} \log n_{k},
\end{aligned}
$$

where

$$
c_{3}=\liminf _{r \rightarrow \infty} \frac{\log V\left(x_{0}, r\right)}{r \log r}+1 .
$$

Taking $\alpha=4 c_{3} \log n_{k}$ and $k$ large enough such that $n_{k} \geqslant a^{-1}\left[4 a+2 d\left(x_{0}, \operatorname{supp}(f)\right)\right]$, we estimate the right side of (3.13) to get

$$
\begin{aligned}
\left|\int_{0}^{t}\left\langle\dot{u}_{s}, g_{n_{k}}\right\rangle d s\right|^{2} \leqslant & 4 a^{-2} t M\left\|\phi^{1 / 2} f\right\|_{2}^{2} \\
& \times \exp \left(\frac{M \alpha^{2}\left(e^{2 \alpha a}+1\right) t}{2}-2 a c_{3} n_{k} \log n_{k}+a c_{3} n_{k} \log n_{k}\right) \\
= & 4 a^{-2} t M\left\|\phi^{1 / 2} f\right\|_{2}^{2} \exp \left(\frac{M \alpha^{2}\left(e^{2 \alpha a}+1\right) t}{2}-a c_{3} n_{k} \log n_{k}\right)
\end{aligned}
$$

Since $e^{2 \alpha a}=n_{k}^{8 a c_{3}}$ and $8 a c_{3}<1$, the inequality above implies that for any $t>0$

$$
\lim _{k \rightarrow \infty} \int_{0}^{t}\left\langle\dot{u}_{s}, g_{n_{k}}\right\rangle d s=0
$$

This completes the proof.

## 4. Proof of Theorem 1.2: the recurrence

This section is devoted to the proof of the recurrence test, Theorem 1.2.
Proof of Theorem 1.2. Let $x_{0} \in X$ be the reference point in Theorem 1.2. For $R>2$, set

$$
\theta_{R}(x)=\left(\left(\frac{R-d\left(x, x_{0}\right)}{R-1}\right) \wedge 1\right)_{+}
$$

Since $\theta_{R}$ belongs to $\mathcal{F}_{\text {loc }} \cap L^{\infty}$ and has compact support, by Lemma 2.1, $\theta_{R}$ belongs to $\mathcal{F}$. According to the condition $(\mathrm{M})$ and the chain-rule for a strongly-local Dirichlet form,

$$
\begin{aligned}
\mathcal{E}^{(c)}\left[\theta_{R}\right] & =\int_{X} \Gamma^{(c)}\left[\theta_{R}\right] d m \\
& =\left(\frac{1}{R-1}\right)^{2} \int_{\bar{B}\left(x_{0}, R\right)} \Gamma^{(c)}[d] d m \\
& \leqslant M_{c}\left(\frac{1}{R-1}\right)^{2} V^{(c)}\left(x_{0}, R\right) \\
& \leqslant \frac{4 M_{c} V^{(c)}\left(x_{0}, R\right)}{R^{2}}
\end{aligned}
$$

On the other hand, we find that for any $c_{1}>2$

$$
\begin{aligned}
\mathcal{E}^{(j)}\left[\theta_{R}\right]= & \iint\left(\theta_{R}(x)-\theta_{R}(y)\right)^{2} j(x, d y) m(d x) \\
\leqslant & \frac{2}{(R-1)^{2}} \int_{B\left(x_{0}, R\right)} \int_{B\left(x_{0}, c_{1} R\right)} d(x, y)^{2} j(x, d y) m(d x) \\
& +2 \int_{B\left(x_{0}, R\right)} \int_{B\left(x_{0}, c_{1} R\right)^{c}} j(x, d y) m(d x) \\
\leqslant & \frac{2}{(R-1)^{2}} \int_{B\left(x_{0}, R\right)} \int_{d(x, y) \leqslant 2 c_{1} R} d(x, y)^{2} j(x, d y) m(d x) \\
& +2 \int_{B\left(x_{0}, R\right)} \int_{d(x, y) \geqslant\left(c_{1}-1\right) R} j(x, d y) m(d x)
\end{aligned}
$$

where we used the facts that $d(x, y) \leqslant R+c_{1} R \leqslant 2 c_{1} R$ if $x \in B\left(x_{0}, R\right)$ and $y \in B\left(x_{0}, c_{1} R\right)$; $d(x, y) \geqslant c_{1} R-R \geqslant R_{1}$ if $x \in B\left(x_{0}, R\right)$ and $y \notin B\left(x_{0}, c_{1} R\right)$. The last expression is bounded from above by

$$
\begin{aligned}
\leqslant & \frac{8 c_{1}^{2}}{(R-1)^{2}} \iint_{B\left(x_{0}, R\right)} \int(d(x, y) \wedge R)^{2} j(x, d y) m(d x) \\
& +\frac{2}{R^{2}} \int_{B\left(x_{0}, R\right)} \int(d(x, y) \wedge R)^{2} j(x, d y) m(d x) \\
\leqslant & \frac{33 c_{1}^{2}}{R^{2}} \int_{B\left(x_{0}, R\right)} \int(d(x, y) \wedge R)^{2} j(x, d y) m(d x)
\end{aligned}
$$

Therefore, under the assumption (M), we have that for $c_{2}=4 M_{c}+33 c_{1}^{2}$

$$
\begin{aligned}
\mathcal{E}\left[\theta_{R}\right] & \leqslant \frac{1}{R^{2}}\left[4 M_{c} V^{(c)}\left(x_{0}, R\right)+33 c_{1}^{2} V^{(j)}\left(x_{0}, R\right) \sup _{x \in X^{(j)}} \int(d(x, y) \wedge R)^{2} j(x, d y)\right] \\
& \leqslant \frac{c_{2}}{R^{2}}\left[V^{(c)}\left(x_{0}, R\right)+V^{(j)}\left(x_{0}, R\right) \sup _{x \in X^{(j)}} \int(d(x, y) \wedge R)^{2} j(x, d y)\right]
\end{aligned}
$$

According to the volume condition (1.2), there exists a sequence $\left(n_{k}\right)_{k \geqslant 0}$ such that $n_{k} \rightarrow \infty$ as $k \rightarrow \infty$, and

$$
\liminf _{k \rightarrow \infty} \mathcal{E}\left[\theta_{R_{n_{k}}}\right]<\infty
$$

Applying [8, Theorem 1.6.3] and [34, (1.6.1) and (1.6.1')], this completes the proof.

## 5. Examples

In this section we present some examples to illustrate the power of Theorems 1.1 and 1.2. Throughout the section, we denote the space of real-valued Lipschitz continuous functions with
compact support on a metric space $X$ by $C_{0}^{\text {Lip }}(X)$. For a measure space $(X, m)$ and a quadratic form $\mathcal{E}$ defined in $L^{2}(X ; m)$, we denote

$$
\mathcal{E}_{1}[u]=\|u\|_{L^{2}}^{2}+\mathcal{E}[u]
$$

whenever the right side makes sense. We start with the following remark for the volume test in Theorem 1.1.

Remark 5.1. Let ( $X, d, m$ ) be a complete metric measure space such that $m$ is a Radon measure with full support. Assume that there is a point $x_{0} \in X$ such that

$$
\sup _{r>0} \frac{V\left(x_{0}, 2 r\right)}{V\left(x_{0}, r\right)}<\infty
$$

where $V\left(x_{0}, r\right)$ denotes the volume of the closed ball centered at $x_{0}$ with radius $r>0$. This assumption is called the volume doubling condition at point $x_{0}$, and it implies that there is a constant $\kappa>0$ such that

$$
\sup _{r>0} \frac{V\left(x_{0}, r\right)}{r^{\kappa}}<\infty
$$

In particular, condition (1.1) in Theorem 1.1 is satisfied. A typical example which fulfills the volume doubling condition is a Riemannian manifold with non-negative Ricci curvature.

### 5.1. Sharpness examples

In the following example, we consider two classes of symmetric jump processes on the so called $\kappa$-set.

Example 5.2. Let $(X,|\cdot|, m)$ be a closed $\kappa$-set in $\mathbb{R}^{n}$ with $0<\kappa \leqslant n$, i.e., $|\cdot|$ is the Euclidean distance, and for all $x \in X$ and $r>0$,

$$
m(B(x, r)) \asymp r^{\kappa} .
$$

Here, the symbol $\asymp$ means that the ratio of the left and the right-hand sides is pinched by two positive constants. Assume that the jump kernel $j(x, d y)$ has a density $j(x, y)$ with respect to the measure $m(d y)$ such that one of the following two conditions is satisfied with a constant $\alpha \in(0,2)$ :

$$
\begin{gather*}
j(x, y) \asymp \frac{1}{|x-y|^{\kappa+\alpha}} \mathbb{1}_{\{|x-y| \leqslant 1\}}+\frac{1}{|x-y|^{\kappa+\beta}} \mathbb{1}_{\{|x-y|>1\}}, \quad \text { where } 0<\beta<\infty ;  \tag{i}\\
j(x, y) \asymp \frac{1}{|x-y|^{\kappa+\alpha}} \mathbb{1}_{\{|x-y| \leqslant 1\}}+\frac{e^{-c|x-y|}}{|x-y|^{\kappa+\alpha}} \mathbb{1}_{\{|x-y|>1\}}, \quad \text { where } c>0 . \tag{ii}
\end{gather*}
$$

For $u, v \in C_{0}^{\text {Lip }}(X)$, define

$$
\mathcal{E}(u, v)=\iint_{x \neq y}(u(x)-u(y))(v(x)-v(y)) j(x, y) m(d x) m(d y) .
$$

Let $\mathcal{F}$ be the closure of $C_{0}^{\text {Lip }}(X)$ with respect to the $\sqrt{\mathcal{E}_{1}}$-norm. The symmetric form $(\mathcal{E}, \mathcal{F})$ is a regular Dirichlet form in $L^{2}(X, m)$, see, e.g., [37]. According to Theorems 1.1 and 1.2, the Dirichlet form $(\mathcal{E}, \mathcal{F})$ is conservative, and it is recurrent if additionally $0<\kappa \leqslant \beta \wedge 2$ and $0<\kappa \leqslant 2$ for the cases (i) and (ii), respectively.

Remark 5.3. Example 5.2 is motivated by recent developments for layered stable processes [16] and tempering stable processes [29]. In particular, in case (i) if $\beta=\alpha$, then the associated Hunt process is called a stable-like process [2].

### 5.2. Disconnected space

The following example shows that the state space may be topologically disconnected, and the particles jump between different connected components and it behaves as a jump-diffusion inside a connected component.

Example 5.4. Let $X=\bigcup_{i \in \mathbb{Z}} X_{i}$, where for each $i \in \mathbb{Z}, X_{i}=\left\{\left(x_{i}, i\right) \in \mathbb{R}^{n+1}: x_{i} \in \mathbb{R}^{n}\right\}$. Any point $x$ in $X$ can be expressed uniquely as $x=\left(x_{i}, i\right)$ with $x_{i} \in \mathbb{R}^{n}$ and $i \in \mathbb{Z}$, and we denote the associated projections by $p: X \rightarrow \mathbb{R}^{n}$ and $q: X \rightarrow \mathbb{Z}$. For any $x, y \in X$, the distance $d$ is given by

$$
d(x, y)=|p(x)-p(y)|+|q(x)-q(y)|,
$$

where $|\cdot|$ is the Euclidean distance. Let $m(d x)=\sum_{i \in \mathbb{Z}} m_{i}\left(d x_{i}\right)$ be a measure on $X$ such that for each $i \geqslant 1, m_{i}\left(d x_{i}\right)=\Psi\left(x_{i}\right) d x_{i}$ is a measure on $X_{i}$, where $\Psi \in C\left(\mathbb{R}^{n}\right)$ is a positive function, and $d x_{i}$ is the $n$-dimensional Lebesgue measure. Clearly, $m$ is a Radon measure on $X$. The state space is the triple $(X, d, m)$.

For any $u \in C_{0}^{\text {Lip }}(X)$, define

$$
\mathcal{E}[u]=\mathcal{E}^{(c)}[u]+\mathcal{E}^{(j)}[u],
$$

where

$$
\begin{gathered}
\mathcal{E}^{(c)}[u]=\int_{X}|\nabla u|^{2} d m \\
\mathcal{E}^{(j)}[u]=\int_{X} \int_{x \neq y}(u(x)-u(y))^{2} j(x, y) m(d x) m(d y),
\end{gathered}
$$

and

$$
j(x, y) \asymp \frac{d(x, y)^{-(n+\alpha)} \mathbb{1}_{\{d(x, y)<1\}}+d(x, y)^{-(n+\beta+1)} \mathbb{1}_{\{d(x, y) \geqslant 1\}}}{\Psi(p(x))+\Psi(p(y))}, \quad x, y \in X
$$

with some constants $0<\alpha<2$ and $\beta>0$. Let $\mathcal{F}$ be the closure of $C_{0}^{\mathrm{Lip}}(X)$ with respect to the $\sqrt{\mathcal{E}_{1}}$-norm. Since for any $x \in X$

$$
\begin{aligned}
& \int_{x \neq y}\left(1 \wedge d(x, y)^{2}\right) j(x, y) m(d y) \\
& \quad \leqslant \int_{0<d(x, y)<1} \frac{d(x, y)^{-(n+\alpha-2)} \Psi(p(y)) d p(y)}{\Psi(p(x))+\Psi(p(y))}+\int_{d(x, y) \geqslant 1} \frac{d(x, y)^{-(n+\beta+1)} \Psi(p(y)) d p(y)}{\Psi(p(x))+\Psi(p(y))} \\
& \quad \leqslant \int_{0<d(x, y)<1} d(x, y)^{-(n+\alpha-2)} d p(y)+2 \sum_{k \geqslant 0} \int_{|p(x)-p(y)| \geqslant k+1}|p(x)-p(y)|^{-(n+\beta+1)} d p(y),
\end{aligned}
$$

which is bounded from above by some absolute constant $c>0$, it follows form the proof of [37] that $(\mathcal{E}, \mathcal{F})$ is a regular Dirichlet form in $L^{2}(X, m)$.

According to the arguments above, we can easily claim that the condition (M) is satisfied. Therefore, by Theorem 1.1, if there is a constant $c>0$ such that for $r>0$ large enough

$$
\begin{equation*}
\sum_{0 \leqslant k \leqslant[r]_{B(0,[r]-k)}} \int \Psi(z) d z \leqslant r^{c r}, \tag{5.14}
\end{equation*}
$$

where $d z$ is the $n$-dimensional Euclidean measure and $[r]$ is the least integer such that $[r] \geqslant r$, then the Dirichlet form $(\mathcal{E}, \mathcal{F})$ is conservative. For instance, (5.14) is satisfied, if $\Psi(x) \leqslant|x|^{|x|} \ln |x|$ for $|x|$ large enough.

For the recurrence, we additionally assume that there are two constants $c_{0}, c_{1}>0$ such that

$$
\begin{equation*}
j(x, y) \leqslant \frac{\mathbb{1}_{\left\{d(x, y) \leqslant c_{0}\right\}}}{d(x, y)^{1+\alpha}} \tag{5.15}
\end{equation*}
$$

and

$$
\begin{equation*}
\Psi(x) \leqslant c_{1}|x|^{1-n} \quad \text { for }|x| \text { large enough. } \tag{5.16}
\end{equation*}
$$

Condition (5.16) will imply that for any point $x_{0} \in X$,

$$
\liminf _{r \rightarrow \infty} \frac{V\left(x_{0}, r\right)}{r^{2}} \leqslant 2 \liminf _{r \rightarrow \infty} \frac{1}{r^{2}} \sum_{0 \leqslant k \leqslant[r]} \int_{B\left(x_{0},[r]-k\right)} \Psi(x) d x<\infty
$$

Next, by (5.15), there is a constant $c_{2}>0$ depending only on the dimension such that

$$
\begin{aligned}
\omega(r) & \leqslant \sup _{x \in X} \int_{X} d(x, y)^{2} j(x, y) \Psi(p(y)) d p(y) \\
& \leqslant c_{1} \sup _{x \in X} \int_{d(x, y) \leqslant c_{0}} d(x, y)^{1-\alpha}|p(y)|^{1-n} d p(y) \\
& \leqslant 2 c_{1} c_{2} \sum_{0 \leqslant k \leqslant\left[c_{0}\right]} \int_{0}^{\left[c_{0}\right]-k} r^{1-\alpha} d r<\infty .
\end{aligned}
$$

Therefore, $(\mathcal{E}, \mathcal{F})$ is recurrent by Theorem 1.2.

### 5.3. Volume tests

The first volume test for non-local Dirichlet forms to be conservative was obtained in [26, Main Result], and then refined in [15, Theorem 1.1]. It is easy to construct an example, which is not covered by these tests but by Theorem 1.1. Here, we illustrate this by using a weighted Euclidean space as well as a model manifold.

Example 5.5. Let $(\mathbb{R},|\cdot|, m)$ be a weighted Euclidean space, where $|\cdot|$ is the Euclidean distance and the measure is $m(d x)=e^{2 \lambda|x|} d x$ for some $\lambda>0$. For $u \in C_{0}^{\text {Lip }}(\mathbb{R})$, define

$$
\mathcal{E}[u]=\iint_{x \neq y}(u(x)-u(y))^{2} j(x, y) m(d x) m(d y)
$$

where

$$
j(x, y)=\left(e^{-\lambda(|x|+|y|)}\right) \mathbb{1}_{\{|x-y| \leqslant 1\}} .
$$

Let $\mathcal{F}$ be the closure of $C_{0}^{\text {Lip }}(\mathbb{R})$ with respect to the $\sqrt{\mathcal{E}_{1}}$-norm. The symmetric form $(\mathcal{E}, \mathcal{F})$ becomes a regular Dirichlet form in $L^{2}(\mathbb{R}, m)$, see, e.g., [37]. Let $j(x, d y)=j(x, y) m(d y)$. It holds that

$$
\begin{aligned}
\sup _{x \in \mathbb{R}} \int\left(1 \wedge|x-y|^{2}\right) j(x, d y) & =\sup _{x \in \mathbb{R}} \int_{\{|y-x| \leqslant 1\}}|x-y|^{2} j(x, y) m(d y) \\
& =\sup _{x \in \mathbb{R}} e^{-\lambda|x|} \int_{\{|z| \leqslant 1\}} z^{2} e^{\lambda|x-z|} d z \\
& \leqslant \int_{\{|z| \leqslant 1\}} z^{2} e^{\lambda|z|} d z<\infty .
\end{aligned}
$$

On the other hand, it is easy to see that in this example (1.1) is also satisfied. Therefore, according to Theorem 1.1, the Dirichlet form $(\mathcal{E}, \mathcal{F})$ is conservative.

However, since $x \mapsto e^{-r|x|} \notin L^{1}(\mathbb{R}, m)$ for any $r \leqslant 2 \lambda$, this example is not covered by [26, Main Result].

Example 5.6 (Model manifolds). (See, e.g., [14].) Let ( $\mathbb{S}^{n}, g$ ) be the $n$-dimensional unit sphere with $n \geqslant 1$. A model manifold $M=(0,+\infty) \times \mathbb{S}^{n}$ is a Riemannian manifold with Riemannian tensor

$$
d r^{2}+\sigma^{2}(r) g
$$

where $\sigma$ is a locally-Lipschitz continuous positive function on $[0,+\infty)$ such that $\sigma(0)=0$ and $\sigma^{\prime}(+0)=0$. Thanks to these two conditions, the manifold $M$ is geodesically complete, and so it satisfies the assumption for the state space as explained in Introduction. Let $d m=\omega_{n} \sigma^{n}(r) d r$ be a measure on $M$, where $\omega_{n}$ is the volume of $\mathbb{S}^{n}$.

For any $u \in C_{0}^{\text {Lip }}(M)$, define

$$
\mathcal{E}[u]=\mathcal{E}^{(c)}[u]+\mathcal{E}^{(j)}[u],
$$

where

$$
\begin{gathered}
\mathcal{E}^{(c)}[u]=\int_{M}|\nabla u|^{2} d m, \\
\mathcal{E}^{(j)}[u]:=\iint_{M \times M \backslash \text { diag }}(u(x)-u(y))^{2} j(x, y) m(d y) m(d x)
\end{gathered}
$$

and

$$
j(x, y)=\left[\frac{\mathbb{1}_{\{d(x, y)<1\}}}{\sigma(r(x)) \sigma(r(y))}\right]^{n} .
$$

Let $\mathcal{F}$ be the closure of $C_{0}^{\mathrm{Lip}}(M)$ with respect to the $\sqrt{\mathcal{E}_{1}}$-norm. It is easy to check that the symmetric form $(\mathcal{E}, \mathcal{F})$ is a regular Dirichlet form in $L^{2}(M, m)$.

By [9], it is known that (M-2) is satisfied. On the other hand, since

$$
\sup _{x, y \in M} j(x, y) \sigma^{n}(r(y)) \leqslant 1,
$$

we obtain that

$$
\begin{aligned}
M_{j} & =\sup _{x \in M} \int_{M}\left(1 \wedge d(x, y)^{2}\right) j(x, d y) \\
& \leqslant \sup _{x \in M} \int_{M} d(x, y)^{2} j(x, y) m(d y) \\
& \leqslant \sup _{x \in M} \int_{d(x, y)<1} d(x, y)^{2} \omega_{n} d y \\
& \leqslant \omega_{n} .
\end{aligned}
$$

Therefore, (M-3) is also satisfied. Since (M-1) clearly follows, we can apply our main theorem. For example, if $\sigma$ satisfies

$$
\sigma(r) \asymp\left[r^{r}(1+\ln r) \vee 1\right]^{1 / n},
$$

then for any fixed $x_{0} \in M$,

$$
r^{r / 2}<V\left(x_{0}, r\right)<2 r^{r} \quad \text { for large } r>0 .
$$

Therefore, $(\mathcal{E}, \mathcal{F})$ is conservative by Theorem 1.1. We note that this model manifold $M$ does not satisfy the volume tests in [26,15].

### 5.4. A mixed-type Laplacian on graphs

A graph admits natural different "Laplacians"; namely, a physical Laplacian, a combinatorial Laplacian, and a quantum Laplacian. The former two are non-local operators, and the last one is a local operator. The combinatorial Laplacian is bounded, and so the corresponding process always is conservative. The conservativeness of the process associated with the physical Laplacian was studied in $[6,7,39,40,15]$. The conservativeness and recurrence of the process generated by the quantum Laplacian was studied in [35]. In the following example, we consider the sum of a physical Laplacian and a quantum Laplacian, and study its conservativeness.

Let $X=(V, E)$ be a locally finite graph, where $V$ and $E$ are the sets of vertices and edges, respectively. Let $\mu$ be a positive function on $X$, and $\omega: X \times X \rightarrow[0, \infty)$ be a symmetric nonnegative function, such that $\omega(x, y)=0$ whenever $x=y$ for $x, y \in X$ or at least one of $x$ and $y$ does not belong to $V$. Now, we recall the standard adapted distance $d$ in [15]. For any $x, y \in X$, $x \sim y$ means that $x, y$ are neighbors; that is, $(x, y) \in E$. For all $x, y \in V$ with $x \sim y$, define

$$
\sigma(x, y)=\min \left\{\frac{1}{\sqrt{\operatorname{deg}(x)}}, \frac{1}{\sqrt{\operatorname{deg}(y)}}, 1\right\},
$$

where

$$
\operatorname{deg}(x)=\frac{1}{\mu(x)} \sum_{y: y \sim x} \omega(x, y)
$$

It naturally induces a metric $d$ on $V$ as

$$
d(x, y)=\inf \left\{\sum_{i=0}^{n-1} \sigma\left(x_{i}, x_{i}+1\right): x_{0}, \ldots, x_{n} \text { is a chain connecting } x \text { and } y\right\} .
$$

The metric $d$ can be extended to $X$ by linear interpolation. We assume that the lengths of all edges $e \in E$ are uniformly bounded from below by a positive constant. This implies that ( $X, d$ ) is a metrically complete space; in particular, our assumption on the space is satisfied.

We further assume that each edge $e \in E$ is isometric to an interval of $\mathbb{R}$, which yields the measure $d x$ on $e$. The space $(X, d)$ is a metric graph. Consider the following measure $m$ on $X$ :

$$
m:=\delta_{E} \phi d x+\delta_{V} \mu
$$

where $\phi$ is a continuous positive function on $E$.
For $u \in C_{0}^{\text {Lip }}(X)$, define

$$
\mathcal{E}[u]:=\mathcal{E}^{(c)}[u]+\mathcal{E}^{(j)}[u],
$$

where

$$
\mathcal{E}^{(c)}[u]=\int_{E}\left(\frac{\partial u}{\partial x}\right)^{2} d m,
$$

and

$$
\mathcal{E}^{(j)}[u]=\sum_{x, y \in V}(u(x)-u(y))^{2} \omega(x, y) .
$$

The generators associated with $\mathcal{E}^{(c)}$ and $\mathcal{E}^{(j)}$ are called the quantum graph, see, e.g. [25] and the physical Laplacian, respectively. Let $\mathcal{F}$ be the closure of $C_{0}^{\text {Lip }}(X)$ with respect to the $\sqrt{\mathcal{E}_{1}}$-norm. We have

Lemma 5.7. The form $(\mathcal{E}, \mathcal{F})$ is a regular Dirichlet form.
Proof. First, we claim that $C_{0}^{\mathrm{Lip}}(X)$ is dense in $L^{2}(X ; m)$. Let $x_{0}$ be a fixed point in $V$. For any $u \in L^{2}(X ; m)$ and any $\epsilon>0$, choose $R>0$ so large that there is a function $v_{\epsilon} \in C_{0}^{\infty}(B(R) \cap E)$ which satisfies

$$
\left\|v_{\epsilon}-\left.u\right|_{E}\right\|_{L^{2}(E ; d x)}<\epsilon,
$$

and that the function $w_{\epsilon}=\mathbb{1}_{B(R)} u$ satisfies that

$$
\left\|w_{\epsilon}-\left.u\right|_{V}\right\|_{L^{2}(V ; \mu)}<\epsilon,
$$

where $B(R):=B\left(x_{0}, R\right)$. Set $\tilde{u}_{\epsilon}=\delta_{E} v_{\epsilon}+\delta_{V} w_{\epsilon}$. For any $x \in B(R)$ and $e \in E$ with $x \sim e$ (i.e., $x \in e)$, let $\delta=\delta(x, e)$ be a positive number such that $\delta<|e| / 2$, and modify $\tilde{u}_{\epsilon}$ on $e \cap B(x, \delta)$ so that $\tilde{u}_{\epsilon}$ is linear and continuous on $e \cap B(x, \delta)$. Furthermore, since $B(R) \cap V$ is finite, by the Hopf-Rinow type property of locally finite graphs [21], we are able to do this modification for any $x \in B(R) \cap V$ and any $e \in E$ with $x \sim e$. Consequently, we obtain a sequence of functions $u_{\epsilon}^{\delta} \in C_{0}^{\text {Lip }}(B(R))$ which converges to $u$ in $L^{2}(X ; m)$ as $\delta, \epsilon \rightarrow 0$. The required claim is proved.

Next, we verify that $\left(\mathcal{E}, C_{0}^{\mathrm{Lip}}(X)\right)$ is closable. Let $\left(u_{n}\right)_{n \geqslant 1} \subset C_{0}^{\text {Lip }}(X)$ be an $\mathcal{E}_{1}$-Cauchy sequence such that $u_{n} \rightarrow 0$ in $L^{2}(X ; m)$ as $n \rightarrow \infty$. One can easily prove that $\mathcal{E}^{(c)}\left[\left.u_{n}\right|_{E}\right] \rightarrow 0$ as $n \rightarrow \infty$, since $\mathcal{E}^{(c)}$ is equivalent to the Dirichlet integral of an open interval. Moreover, if $v \in C_{0}^{\mathrm{Lip}}(X)$, then

$$
\mathcal{E}^{(j)}\left(\left.u_{n}\right|_{V},\left.v\right|_{V}\right)=\sum_{x, y \in V}\left(u_{n}(x)-u_{n}(y)\right)(v(x)-v(y)) \omega(x, y) \rightarrow 0 \quad \text { as } n \rightarrow \infty .
$$

Therefore, the desired claim follows and we denote the closure of $\left(\mathcal{E}, C_{0}^{\mathrm{Lip}}(X)\right)$ by $(\mathcal{E}, \mathcal{F})$.
The Markov property of $(\mathcal{E}, \mathcal{F})$ follows immediately from the definition of $\mathcal{E}$. Finally, since $C_{0} \cap \mathcal{F}$ is both dense in $C_{0}$ and $\mathcal{F}$ with respect to the sup-norm and the $\mathcal{E}_{1}$-norm, respectively, $(\mathcal{E}, \mathcal{F})$ is regular.

It is easy to see that the conditions (M-1) and (M-2) are satisfied since $X^{(c)}=E$. Moreover, since $\mathcal{E}^{(j)}$ can be expressed as

$$
\mathcal{E}^{(j)}[u]=\iint_{X \times X}(u(x)-u(y))^{2} \frac{\omega(x, y)}{\mu(x) \mu(y)} m(d y) m(d x),
$$

the associated jump kernel $j$ and $\Gamma_{j}$ have the forms

$$
j(x, d y)=\frac{\omega(x, y)}{\mu(x) \mu(y)} m(d y)
$$

and

$$
\Gamma_{j}[u](x)=\int_{X}(u(x)-u(y))^{2} \frac{\omega(x, y)}{\mu(x) \mu(y)} m(d y) \quad \text { for any } x \in X .
$$

Clearly, (M-3) is satisfied. Therefore the Dirichlet form $(\mathcal{E}, \mathcal{F})$ satisfies the condition (M).
To state our main result in this subsection, we need some notations. Denote by $\rho$ the graph distance extended to $X$, and by $B_{\rho}\left(x_{0}, R\right)$ the associated ball at $x_{0} \in V$ with radius $R>0$. For any $n \in \mathbb{N}$, let $S_{\rho}\left(x_{0}, n\right)$ be the "boundary" $B_{\rho}\left(x_{0}, n\right) \backslash B_{\rho}\left(x_{0}, n-1\right)$.

Proposition 5.8. If $\mu$ is the counting measure and there are a point $x_{0} \in V$ and a constant $C>0$ such that

$$
\begin{equation*}
m\left(S_{\rho}\left(x_{0}, n\right)\right) \leqslant C n^{2} \quad \text { for all large enough } n \in \mathbb{N}, \tag{5.17}
\end{equation*}
$$

then $(\mathcal{E}, \mathcal{F})$ is conservative.
Proof. The condition (5.17) implies that for any $x \in V$,

$$
\begin{equation*}
d\left(x_{0}, x\right) \geqslant \delta \log \rho\left(x_{0}, x\right) \tag{5.18}
\end{equation*}
$$

where $\delta>0$ is a constant depending only on $C$ in (5.17) (see [15]). Let $\overline{x x^{\prime}}$ be the edge with boundary $\left\{x, x^{\prime}\right\}$. Let $y \in X$ and $x, x^{\prime} \in V$ such that $y \in \overline{x x^{\prime}}$. Without loss of generality, we assume that $\rho\left(x_{0}, y\right) \leqslant \rho\left(x_{0}, x^{\prime}\right)$. By using (5.18), the triangle inequality and the fact that $d\left(x, x^{\prime}\right) \leqslant \rho\left(x, x^{\prime}\right)=1$, we find that

$$
\rho\left(x_{0}, y\right) \leqslant e^{d\left(x_{0}, x^{\prime}\right) / \delta} \leqslant e^{1 / \delta} e^{d\left(x_{0}, x\right) / \delta} .
$$

Since $d\left(x_{0}, y\right) \geqslant d\left(x_{0}, x\right) \wedge d\left(x_{0}, x^{\prime}\right)$, we obtain that there is a constant $c>0$ such that

$$
\rho\left(x_{0}, y\right) \leqslant c e^{d\left(x_{0}, y\right) / \delta \quad \text { for any } y \in X . ~}
$$

It follows that there exists a constant $b>0$ such that

$$
m\left(B_{d}\left(x_{0}, r\right)\right) \leqslant m\left(B_{\rho}\left(x_{0}, c e^{r / \delta}\right)\right) \leqslant \exp (b r) \quad \text { for all large enough } r>0
$$

Therefore, $(\mathcal{E}, \mathcal{F})$ is conservative by Theorem 1 .
Remark 5.9. By an example of R. Wojciechowski [41], the boundary volume growth of quadratic rate (5.17) is sharp. The second part of Proposition 5.8 was obtained in [15] for a physical Laplacian on a graph.

On the other hand, it is easy to check that the condition (5.17) is satisfied, if there is a constant $C>0$ such that
(1) $\mu\left(S_{\rho}\left(x_{0}, n\right)\right) \leqslant C n^{2}$ for all large enough $n \in \mathbb{N}$,
(2) $\phi(x) \leqslant C \rho\left(x_{0}, x\right)^{-2}$ for every $x \in X$.

Indeed, the first condition implies that there are at most $\left(\mathrm{Cn}^{2}\right)^{2}$-many edges in $S_{\rho}\left(x_{0}, n\right)$ connecting vertices in $S_{\rho}(n)$ and $S_{\rho}(n-1)$. The second condition then implies that there is a constant $c>0$ such that

$$
m\left(S_{\rho}\left(x_{0}, n\right) \cap E\right) \leqslant \frac{C^{3} n^{4}}{(n-1)^{2}} \leqslant c n^{2} \quad \text { for all large enough } n
$$

This together with the first condition yields (5.17).

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[^0]:    * Corresponding author.

    E-mail addresses: jum35@psu.edu (J. Masamune), t-uemura@kansai-u.ac.jp (T. Uemura), jianwang@fjnu.edu.cn (J. Wang).

