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A new homotopy perturbation method for solving systems of partial differential equations

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ABSTRACT

In this paper, a new homotopy perturbation method (NHPM) is introduced for obtaining solutions of systems of non-linear partial differential equations. Theoretical considerations are discussed. To illustrate the capability and reliability of the method three examples are provided. Comparison of the results of applying NHPM with those of applying HPM reveal the effectiveness and convenience of the new technique.

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1. Introduction

In this work, we study systems of partial differential equations (PDEs). Such systems arise in many areas of mathematics, engineering and physical sciences. These equations are often too complicated to be solved exactly and even if an exact solution is obtained, the required calculations may be too complicated. Very recently, many powerful methods have been presented, such as the Adomian decomposition method [1,2], the variational iteration method [3,4], the homotopy perturbation method (HPM) [5–21], and the differential transform method [22–24].

The general form of a system of PDEs can be considered as the following:

$$\frac{\partial u_j}{\partial t} + N_j(x_1, \dots, x_{n-1}, t, u_1, \dots, u_n) = g_j(x_1, x_2, \dots, x_{n-1}, t), \quad j = 1, \dots, n,$$
(1)

with the following initial conditions:

$$u_i(x_1, x_2, \dots, x_{n-1}, t_0) = f_i(x_1, x_2, \dots, x_{n-1}), \quad j = 1, \dots, n$$

where N_1, \ldots, N_n are non-linear operators, which usually depend on the functions u_i and their derivatives, and g_1, g_2, \ldots, g_n are inhomogeneous terms.

This paper is arranged as follows. In Section 2, the new modification of HPM, called NHPM, for solving systems of partial differential equations is presented. The efficiency of the new method is verified by the numerical results for three sample examples in Section 3. Comparisons between this method and HPM are illustrated in this section. The conclusions appear in Section 4.

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2. The basic idea of NHPM

For solving system (1), by using NHPM, we construct the following homotopies:

$$(1-p)\left(\frac{\partial U_j}{\partial t}-u_{j,0}\right)+p\left(\frac{\partial U_j}{\partial t}+N_j(x_1,\ldots,x_{n-1},t,U_1,\ldots,U_n)-g_j\right)=0, \quad j=1,\ldots,n$$

or

$$\frac{\partial U_j}{\partial t} = u_{j_0} - p \left(u_{j_0} + N_j(x_1, x_2, \dots, x_{n-1}, t, U_1, \dots, U_n) - g_j \right), \quad j = 1, \dots, n.$$
(3)

Applying the inverse operator, $L^{-1} = \int_{t_0}^{t} (.) dt$, to both sides of Eq. (3), we obtain

$$U_{j}(x_{1}, x_{2}, ..., x_{n-1}, t) = U_{j}(x_{1}, x_{1}, ..., x_{n-1}, t_{0}) + \int_{t_{0}}^{t} u_{j_{0}} dt$$

$$-p \int_{t_{0}}^{t} (u_{j_{0}} + N_{j}(x_{1}, x_{2}, ..., x_{n-1}, t, U_{1}, ..., U_{n}) - g_{j}) dt, \quad j = 1, ..., n$$

$$(4)$$

where

$$U_i(x_1, x_2, \dots, x_{n-1}, t_0) = u_i(x_1, x_2, \dots, x_{n-1}, t_0), \quad j = 1, \dots, n.$$

Let us present the solution of the system (4) as the following:

$$U_i = U_{i0} + pU_{i1} + p^2U_{i2} + \cdots, \quad j = 1, \dots, n$$
 (5)

where U_{ij} , i = 1, ..., n, j = 0, ..., n, are functions which should be determined. Suppose that the initial approximations of the solutions of Eqs. (1) are in the following form:

$$u_{i,0}(x_1, x_2, \dots, x_{n-1}, t) = \sum_{i=0}^{\infty} a_{i,i}(x_1, x_2, \dots, x_{n-1}) P_j(t), \quad i = 1, \dots, n$$
(6)

where $a_{i,j}(x_1, x_2, \dots, x_{n-1}), i = 1, \dots, n, j = 0, \dots, n$, are unknown coefficients and $P_0(t), P_1(t), P_2(t), \dots$ are specific functions.

Substituting (5) and (6) into (4) and equating the coefficients of p with the same powers leads to

$$p^{0}: U_{i,0}(x_{1}, x_{2}, \dots, x_{n-1}, t) = f_{i}(x_{1}, x_{2}, \dots, x_{n-1}) + \sum_{j=0}^{\infty} a_{ij} \int_{t_{0}}^{t} P_{j}(t) dt,$$

$$p^{1}: U_{i,1}(x_{1}, x_{2}, \dots, x_{n-1}, t) = -\sum_{j=0}^{\infty} a_{ij} \int_{t_{0}}^{t} P_{j}(t) dt - \int_{t_{0}}^{t} \left(N_{i}(x_{1}, x_{2}, \dots, x_{n-1}, t, U_{1,0}, \dots, U_{n,0}) - g_{i} \right) dt,$$

$$p^{2}: U_{i,2}(x_{1}, x_{2}, \dots, x_{n-1}, t) = -\int_{t_{0}}^{t} \left(N_{i}(x_{1}, x_{2}, \dots, x_{n-1}, t, U_{1,0}, \dots, U_{n,0}, U_{1,1}, \dots, U_{n,1}) \right) dt,$$

$$\vdots$$

$$p^{j}: U_{i,j}(x_{1}, x_{2}, \dots, x_{n-1}, t) = -\int_{t_{0}}^{t} \left(N_{i}(x_{1}, x_{2}, \dots, x_{n-1}, t, U_{1,0}, \dots, U_{n,0}, \dots, U_{1,j-1}, \dots, U_{n,j-1}) \right) dt,$$

$$\vdots$$

Now if we solve these equations in such a way that $U_{i,1}(x_1x_2,\ldots,x_{n-1},t)=0$, then Eqs. (7) yield

$$U_{i,2}(x_1x_2,\ldots,x_{n-1},t)=U_{i,3}(x_1x_2,\ldots,x_{n-1},t)=\cdots=0.$$

Therefore the exact solution may be obtained as the following:

$$u_i(x_1, x_2, \dots, x_{n-1}, t) = U_{i,0}(x_1, x_2, \dots, x_{n-1}, t) = f_i(x_1, x_2, \dots, x_{n-1}) + \sum_{i=0}^{\infty} a_{i,j} \int_{t_0}^{t} P_j(t) dt, \quad i = 1, \dots, n.$$
 (8)

It is worth mentioning that if $g_i(x_1, x_2, \dots, x_{n-1}, t)$, and $u_{i,0}(x_1, x_2, \dots, x_{n-1}, t)$, are analytic around $t = t_0$, then their Taylor series can be defined as

$$u_{i,0}(x_1, x_2, \dots, x_{n-1}, t) = \sum_{j=0}^{\infty} a_{i,j}(x_1, x_2, \dots, x_{n-1})(t - t_0)^n,$$

$$g_i(x_1, x_2, \dots, x_{n-1}, t) = \sum_{j=0}^{\infty} a_{i,j}^*(x_1, x_2, \dots, x_{n-1})(t - t_0)^n,$$
(9)

which can be used in Eqs. (7), where $a_{i,j}(x_1, x_2, \dots, x_{n-1}), i = 1, \dots, n, j = 0, \dots, n$, are unknown coefficients which must be computed, and $a_{i,j}^*(x_1, x_2, \dots, x_{n-1}), i = 1, \dots, n, j = 0, \dots, n$, are known ones.

To show the capability of the method, NHPM has been applied to some examples in the next section.

3. Numerical results

To demonstrate the effectiveness of the method three examples of systems of non-linear partial differential equations are presented.

Example 1. Consider the following system of three-dimensional partial differential equations:

$$\begin{cases}
\frac{\partial u}{\partial t} - v \frac{\partial u}{\partial x} - \frac{\partial v}{\partial t} \frac{\partial u}{\partial y} = 1 - x + y + t, \\
\frac{\partial v}{\partial t} - u \frac{\partial v}{\partial x} - \frac{\partial u}{\partial t} \frac{\partial v}{\partial y} = 1 - x - y - t,
\end{cases} (10)$$

with initial conditions

$$u(x, y, 0) = x + y - 1,$$

 $v(x, y, 0) = x - y + 1.$

The exact solutions are

$$u(x, y, t) = x + y + t - 1,$$

 $v(x, y, t) = x - y - t + 1.$

The HPM and NHPM methods are used to approximate the solutions.

First we apply the HPM method.

According to the homotopy perturbation method, we have

$$\begin{cases} U_0 = x + y - 1, \\ V_0 = x - y + 1, \end{cases}$$

$$\begin{cases} U_1 = 2t + \frac{1}{2}t^2, \\ V_1 = -\frac{1}{2}t^2, \end{cases}$$

$$\begin{cases} U_2 = -\frac{1}{6}t^3 - \frac{1}{2}t^2, \\ V_2 = \frac{1}{2}t^2 + \frac{1}{6}t^3 - 2t, \end{cases}$$

$$\begin{cases} U_3 = \frac{1}{3}t^3 + \frac{1}{24}t^4 - \frac{1}{2}t^2 - 2t, \\ V_3 = -\frac{1}{24}t^4 + \frac{1}{2}t^2, \end{cases}$$

Therefore, the solution will be as follows:

$$\begin{cases} u = -1 + x + y + 2t - \frac{1}{5040}t^7 + \frac{1}{40}t^5 - \frac{1}{2}t^3 + \cdots, \\ v = 1 + x - y + 2t + \frac{1}{5040}t^7 - \frac{1}{40}t^5 + \frac{1}{2}t^3 + \cdots. \end{cases}$$

To solve Eq. (10) by using NHPM, we construct the following homotopies:

$$\frac{\partial U}{\partial t}(x, y, t) = u_0(x, y, t) - p\left(u_0(x, y, t) - V\frac{\partial U}{\partial x} - \frac{\partial V}{\partial t}\frac{\partial U}{\partial y} - 1 + x - y - t\right),$$

$$\frac{\partial V}{\partial t}(x, y, t) = v_0(x, y, t) - p\left(v_0(x, y, t) - U\frac{\partial V}{\partial x} - \frac{\partial U}{\partial t}\frac{\partial V}{\partial y} - 1 + x + y + t\right).$$
(11)

Applying the inverse operator, $L^{-1} = \int_0^t (.) dt$, to both sides of these equations, we obtain

$$U(x,y,t) = U(x,y,0) + \int_0^t u_0(x,y,t)dt - p \int_0^t \left(u_0(x,y,t) - V \frac{\partial U}{\partial x} - \frac{\partial V}{\partial t} \frac{\partial U}{\partial y} - 1 + x - y - t \right) dt,$$

$$V(x,y,t) = V(x,y,0) + \int_0^t v_0(x,y,t)dt - p \int_0^t \left(v_0(x,y,t) - U \frac{\partial V}{\partial x} - \frac{\partial U}{\partial t} \frac{\partial V}{\partial y} - 1 + x + y + t \right) dt.$$
(12)

Suppose that the solutions of system (12) are as assumed in (5); substituting Eqs. (5) into Eqs. (12), collecting the same powers of p, and equating each coefficient of p to zero yields

$$\begin{split} p^0: & \begin{cases} U_0(x,y,t) = U(x,y,0) + \int_0^t u_0(x,y,t) \mathrm{d}t, \\ V_0(x,y,t) = V(x,y,0) + \int_0^t v_0(x,y,t) \mathrm{d}t, \end{cases} \\ p^1: & \begin{cases} U_1(x,y,t) = \int_0^t \left(-u_0(x,y,t) + V_0 \frac{\partial U_0}{\partial x} + \frac{\partial V_0}{\partial t} \frac{\partial U_0}{\partial y} + 1 - x + y + t \right) \mathrm{d}t, \\ V_1(x,y,t) = \int_0^t \left(-v_0(x,y,t) + U_0 \frac{\partial V_0}{\partial x} + \frac{\partial U_0}{\partial t} \frac{\partial V_0}{\partial y} + 1 - x - y - t \right) \mathrm{d}t, \end{cases} \\ p^2: & \begin{cases} U_2(x,y,t) = \int_0^t \left(V_0 \frac{\partial U_1}{\partial x} + V_1 \frac{\partial U_0}{\partial x} + \frac{\partial V_0}{\partial t} \frac{\partial U_1}{\partial y} + \frac{\partial V_1}{\partial t} \frac{\partial U_0}{\partial y} \right) \mathrm{d}t, \\ V_2(x,y,t) = \int_0^t \left(U_0 \frac{\partial V_1}{\partial x} + U_1 \frac{\partial V_0}{\partial x} + \frac{\partial U_0}{\partial t} \frac{\partial V_1}{\partial y} + \frac{\partial U_1}{\partial t} \frac{\partial V_0}{\partial y} \right) \mathrm{d}t, \end{cases} \\ \vdots \\ p^j: & \begin{cases} U_{j+1}(x,y,t) = \int_0^t \sum_{k=0}^j \left(V_k \frac{\partial U_{j-k}}{\partial x} + \frac{\partial V_k}{\partial t} \frac{\partial U_{j-k}}{\partial y} \right) \mathrm{d}t, \end{cases} \end{aligned}$$

Assume $\begin{cases} u_0(x,y,t) = \sum_{n=0}^{\infty} a_n(x,y) P_n(t), \ P_k(t) = t^k, \ U(x,y,0) = u(x,y,0), \\ v_0(x,y,t) = \sum_{n=0}^{\infty} b_n(x,y) P_n(t), \ P_k(t) = t^k, \ V(x,y,0) = v(x,y,0). \end{cases}$ Solving the above equations for $U_1(x,y,t), V_1(x,y,t)$ leads to the result

$$U_{1}(x, y, t) = (-a_{0}(x, y) + b_{0}(x, y) + 2)t + \left(-\frac{1}{2}a_{1}(x, y) + \frac{1}{2}b_{0}(x, y) + \frac{1}{2}(x - y + 1)a_{0x}(x, y) + \frac{1}{2}b_{1}(x, y) + \frac{1}{2}a_{0,y}(x, y)b_{0}(x, y) + \frac{1}{2}\right)t^{2} + \left(-\frac{1}{3}a_{2}(x, y) + \frac{1}{6}b_{1}(x, y) + \frac{1}{6}(x - y + 1)a_{1x}(x, y) + \frac{1}{3}a_{0x}(x, y)b_{0}(x, y) + \frac{1}{3}b_{2}(x, y) + \frac{1}{3}a_{0y}(x, y)b_{1}(x, y) + \frac{1}{6}a_{1y}(x, y)b_{0}(x, y) + \frac{1}{8}a_{0x}(x, y)b_{1}(x, y) + \frac{1}{8}a_{1x}(x, y)b_{0}(x, y) + \frac{1}{4}a_{0y}(x, y)b_{2}(x, y) + \frac{1}{8}a_{1y}(x, y)b_{1}(x, y) + \frac{1}{12}a_{2y}(x, y)b_{0}(x, y) + \frac{1}{12}a_{2y$$

$$\begin{split} V_{1}(x,y,t) &= (-b_{0}(x,y) - a_{0}(x,y))t \\ &+ \left(-\frac{1}{2}b_{1}(x,y) + \frac{1}{2}a_{0}(x,y) + \frac{1}{2}(x+y-1)b_{0x}(x,y) - \frac{1}{2}a_{1}(x,y) + \frac{1}{2}a_{0}(x,y)b_{0y}(x,y) - \frac{1}{2}\right)t^{2} \\ &+ \left(-\frac{1}{3}b_{2}(x,y) + \frac{1}{6}a_{1}(x,y) + \frac{1}{6}(x+y-1)b_{1x}(x,y) + \frac{1}{3}a_{0}(x,y)b_{0x}(x,y) - \frac{1}{2}\right)t^{2} \\ &+ \left(-\frac{1}{3}a_{2}(x,y) + \frac{1}{6}a_{0}(x,y)b_{1y}(x,y) + \frac{1}{3}a_{1}(x,y)b_{0y}(x,y) - \frac{1}{3}a_{0}(x,y)b_{0x}(x,y) + \frac{1}{4}a_{0}(x,y)b_{1y}(x,y) + \frac{1}{4}a_{1}(x,y)b_{0y}(x,y) + \frac{1}{8}a_{1}(x,y)b_{0x}(x,y) + \frac{1}{8}a_{0}(x,y)b_{1x}(x,y) - \frac{1}{4}a_{0}(x,y)b_{1y}(x,y) + \frac{1}{4}a_{0}(x,y)b_{0y}(x,y) + \frac$$

By the vanishing of $U_1(x, y, t)$, $V_1(x, y, t)$, the coefficients $a_n(x, y)$, $b_n(x, y)$ (n = 1, 2, 3, ...) are determined as

$$a_0(x, y) = 1,$$
 $a_1(x, y) = a_2(x, y) = a_3(x, y) = a_4(x, y) = \cdots 0,$
 $b_0(x, y) = -1,$ $b_1(x, y) = b_2(x, y) = b_3(x, y) = b_4(x, y) = \cdots 0$

Therefore we obtain the solution of Eq. (14) as

$$u(x, y, t) = U_0(x, y, t) = x + y - 1 + a_0(x, y)t + \frac{1}{2}a_1(x, y)t^2 + \frac{1}{3}a_2(x, y)t^3 + \frac{1}{4}a_3(x, y)t^4 + \dots = x + y + t - 1,$$

$$v(x, y, t) = V_0(x, y, t) = x - y + 1 + b_0(x, y)t + \frac{1}{2}b_1(x, y)t^2 + \frac{1}{3}b_2(x, y)t^3 + \frac{1}{4}b_3(x, y)t^4 + \dots = x - y - t + 1,$$

which is an exact solution.

Example 2. Consider the following system of two non-linear equations:

$$\begin{cases} \frac{\partial u}{\partial x} - v \frac{\partial u}{\partial t} + u \frac{\partial v}{\partial t} = -1 + e^{x} \sin t, \\ \frac{\partial v}{\partial x} + \frac{\partial u}{\partial t} \frac{\partial v}{\partial x} + \frac{\partial v}{\partial t} \frac{\partial u}{\partial x} = -1 - e^{-x} \cos t, \end{cases}$$
(13)

with boundary conditions

$$u(0, t) = \sin t,$$

$$v(0, t) = \cos t.$$

The exact solutions of this system are

$$u(x, t) = e^{x} \sin t,$$

$$v(x, t) = e^{-x} \cos t.$$

This system will be solved by using HPM and NHPM.

First we apply the HPM method.

According to the homotopy perturbation method, we have

$$\begin{cases} U_0 = \sin t, \\ V_0 = \cos t, \end{cases}$$

$$\begin{cases} U_1 = e^x \sin t - \sin t, \\ V_1 = -x + e^{-x} \cos t - \cos t, \end{cases}$$

$$\begin{cases} U_2 = e^x - e^{-x} - 2x - \frac{1}{2}x^2 \cos(t), \\ V_2 = -e^{-x} \cos^2 t + e^x \sin^2 t + 2\cos^2 t + x\cos t - 1, \end{cases}$$

$$\vdots$$

Therefore, the solution will be as follows:

$$\begin{cases} u = e^{x} \sin t + e^{x} - e^{-x} - 2x - \frac{1}{2}x^{2} \cos(t) + \cdots, \\ v = -e^{-x} \cos^{2} t + e^{x} \sin^{2} t + e^{-x} \cos t + 2 \cos^{2} t + x \cos t - x - 1 + \dots \end{cases}$$

To solve system (12) by using NHPM, we construct the following homotopy:

$$\frac{\partial U}{\partial x}(x,t) = u_0(x,t) - p\left(u_0(x,t) - V\frac{\partial U}{\partial t} + U\frac{\partial V}{\partial t} + 1 - e^x \sin t\right),$$

$$\frac{\partial V}{\partial x}(x,t) = v_0(x,t) - p\left(v_0(x,t) + \frac{\partial U}{\partial t}\frac{\partial V}{\partial x} + \frac{\partial V}{\partial t}\frac{\partial U}{\partial x} + 1 + e^{-x}\cos t\right).$$
(14)

Applying the inverse operator, $L^{-1} = \int_0^x (.) dx$, to both sides of the above equations, we obtain

$$U(x,t) = U(0,t) + \int_0^x u_0(x,t) dx - p \int_0^x \left(u_0(x,t) - V \frac{\partial U}{\partial t} + U \frac{\partial V}{\partial t} + 1 - e^x \sin t \right) dx,$$

$$V(x,t) = V(0,t) + \int_0^x v_0(x,t) dx - p \int_0^x \left(v_0(x,t) + \frac{\partial U}{\partial t} \frac{\partial V}{\partial x} + \frac{\partial V}{\partial t} \frac{\partial U}{\partial x} + 1 + e^{-x} \cos t \right) dx.$$
(15)

Suppose that the solutions of system (15) have the form (5); substituting Eqs. (5) into Eqs. (15), collecting terms with the same powers of p, and equating each coefficient of p to zero results in

$$p^{0}: \begin{cases} U_{0}(x,t) = U(0,t) + \int_{0}^{x} u_{0}(x,t) dx, \\ V_{0}(x,t) = V(0,t) + \int_{0}^{x} v_{0}(x,t) dx, \end{cases}$$

$$p^{1}: \begin{cases} U_{1}(x,t) = \int_{0}^{x} \left(-u_{0}(x,t) + V_{0} \frac{\partial U_{0}}{\partial t} - U_{0} \frac{\partial V_{0}}{\partial t} - 1 + e^{x} \sin t\right) dx, \\ V_{1}(x,t) = \int_{0}^{x} \left(-v_{0}(x,t) - \frac{\partial U_{0}}{\partial t} \frac{\partial V_{0}}{\partial x} - \frac{\partial V_{0}}{\partial t} \frac{\partial U_{0}}{\partial x} - 1 - e^{-x} \cos t\right) dx, \end{cases}$$

$$p^{2}: \begin{cases} U_{2}(x,t) = \int_{0}^{x} \left(V_{0} \frac{\partial U_{1}}{\partial t} + V_{1} \frac{\partial U_{0}}{\partial t} - U_{0} \frac{\partial V_{1}}{\partial t} - U_{1} \frac{\partial V_{1}}{\partial t}\right) dx, \\ V_{2}(x,t) = \int_{0}^{x} \left(-\frac{\partial U_{0}}{\partial t} \frac{\partial V_{1}}{\partial x} - \frac{\partial U_{1}}{\partial t} \frac{\partial V_{0}}{\partial x} - \frac{\partial V_{0}}{\partial t} \frac{\partial U_{1}}{\partial x} - \frac{\partial V_{1}}{\partial t} \frac{\partial U_{0}}{\partial x}\right) dx, \end{cases}$$

$$\vdots$$

$$p^{j}: \begin{cases} U_{j+1}(x,t) = \int_{0}^{x} \sum_{k=0}^{j} V_{k} \left(\frac{\partial U_{j-k}}{\partial t} - U_{k} \frac{\partial V_{j-k}}{\partial t} \right) dx, \\ V_{j+1}(x,t) = \int_{0}^{x} \sum_{k=0}^{j} \left(-\frac{\partial U_{k}}{\partial t} \frac{\partial V_{j-k}}{\partial x} - \frac{\partial V_{k}}{\partial t} \frac{\partial U_{j-k}}{\partial y} \right) dx, \end{cases}$$

Assuming that $\begin{cases} u_0(x,t) = \sum_{n=0}^{\infty} a_n(t) x^k, & U(0,t) = u(0,t), \\ v_0(x,t) = \sum_{n=0}^{\infty} b_n(t) x^k, & V(0,t) = v(0,t), \end{cases}$ and solving equations $U_1(x,t), V_1(x,t)$ leads to the following results:

$$\begin{split} U_1(x,t) &= (-a_0(t) + \sin t) \, x + \left(-\frac{1}{2} a_1(t) + \frac{a_0'(t) \cos t}{2} + \frac{b_0(t) \cos t}{2} - \frac{b_0'(t) \sin t}{2} + \frac{a_0(t) \sin t}{2} + \frac{\sin t}{2} \right) x^2 \\ &+ \left(-\frac{1}{3} a_2(t) + \frac{a_1'(t) \cos t}{6} + \frac{b_1(t) \cos t}{6} + \frac{b_0(t) a_0'(t)}{3} - \frac{b_1'(t) \sin t}{6} + \frac{a_1(t) \sin t}{6} - \frac{b_0'(t) a_0(t)}{3} + \frac{\sin t}{6} \right) x^3 \\ &+ \left(-\frac{1}{4} a_3(t) + \frac{a_2'(t) \cos t}{12} + \frac{b_2(t) \cos t}{12} + \frac{b_1(t) a_0'(t)}{8} + \frac{b_0(t) a_1'(t)}{8} + \frac{b_0(t) a_1'(t)}{8} \right) x^4 + \cdots, \\ &- \frac{b_2'(t) \sin t}{12} + \frac{a_2(t) \sin t}{12} - \frac{b_0'(t) a_1(t)}{8} - \frac{b_1'(t) a_0(t)}{8} + \frac{\sin t}{24} \end{split}$$

$$V_1(x,t) &= (-b_0(t) - b_0(t) \cos t + a_0(t) \sin t - \cos t - 1) x \\ &+ \left(-\frac{1}{2} b_1(t) - \frac{b_1(t) \cos t}{2} - \frac{b_0(t) a_0'(t)}{2} + \frac{a_1(t) \sin t}{2} - \frac{b_0'(t) a_0(t)}{2} + \frac{\cos t}{2} \right) x^2 \end{split}$$

$$+\left(-\frac{1}{3}b_{2}(t)-\frac{b_{2}(t)\cos t}{3}-\frac{b_{0}(t)a_{1}'(t)}{6}-\frac{b_{1}(t)a_{0}'(t)}{3}+\frac{a_{2}(t)\sin t}{3}-\frac{b_{1}'(t)a_{0}(t)}{6}-\frac{b_{0}'(t)a_{1}(t)}{3}-\frac{\cos t}{6}\right)x^{3}\\+\left(-\frac{1}{4}b_{3}(t)-\frac{b_{3}(t)\cos t}{4}-\frac{b_{0}(t)a_{2}'(t)}{12}-\frac{b_{1}(t)a_{1}'(t)}{8}-\frac{b_{2}(t)a_{0}'(t)}{4}\\+\frac{a_{3}(t)\sin t}{4}-\frac{b_{0}'(t)a_{2}(t)}{4}-\frac{b_{1}'(t)a_{1}(t)}{8}-\frac{b_{2}'(t)a_{0}(t)}{12}+\frac{\cos t}{24}\right)x^{4}+\cdots$$

By the vanishing of $U_1(x, t)$, $V_1(x, t)$, the coefficients $a_n(t)$, $b_n(t)$ (n = 1, 2, 3, ...) are obtained as follows:

$$a_0(t) = \sin t$$
, $a_1(t) = \sin t$, $a_2(t) = \frac{1}{2}\sin t$, $a_3(t) = \frac{1}{6}\sin t$, $a_4(t) = \frac{1}{24}\sin t$, ... $b_0(t) = -\cos t$, $b_1(t) = \cos t$, $b_2(t) = -\frac{1}{2}\cos t$, $b_3(t) = \frac{1}{6}\cos t$, $b_4(t) = -\frac{1}{24}\cos t$, ...

This implies that

$$u(x,t) = U_0(x,t) = \sin t + a_0(t)x + \frac{1}{2}a_1(t)x^2 + \frac{1}{3}a_2(t)x^3 + \frac{1}{4}a_3(t)x^4 + \dots = e^x \sin t,$$

$$v(x,t) = V_0(x,t) = \cos t + b_0(t)x + \frac{1}{2}b_1(t)x^2 + \frac{1}{3}b_2(t)x^3 + \frac{1}{4}b_3(t)x^4 + \dots = e^{-x}\cos t.$$

In this example the exact solutions are gained.

Example 3. Consider the following non-linear system of inhomogeneous partial differential equations:

$$\begin{cases}
\frac{\partial u}{\partial t} - \frac{\partial w}{\partial x} \frac{\partial v}{\partial t} - \frac{1}{2} \frac{\partial w}{\partial t} \frac{\partial^{2} u}{\partial x^{2}} = -4xt, \\
\frac{\partial v}{\partial t} - \frac{\partial w}{\partial t} \frac{\partial^{2} u}{\partial x^{2}} = 6t, \\
\frac{\partial w}{\partial t} - \frac{\partial^{2} u}{\partial x^{2}} - \frac{\partial v}{\partial x} \frac{\partial w}{\partial t} = 4xt - 2t - 2,
\end{cases} \tag{16}$$

subject to the initial condition

$$u(x, 0) = x^{2} + 1,$$

 $v(x, 0) = x^{2} - 1,$
 $w(x, 0) = x^{2} - 1.$

The exact solutions are

$$u(x, t) = x^{2} - t^{2} + 1,$$

 $v(x, t) = x^{2} + t^{2} - 1,$
 $w(x, t) = x^{2} - t^{2} - 1.$

Using the homotopy perturbation method leads to

$$\begin{cases} U_0 = x^2 + 1, \\ V_0 = x^2 - 1, \\ W_0 = x^2 - 1, \end{cases}$$

$$\begin{cases} U_1 = -2xt^2, \\ V_1 = 3t^2, \\ W_1 = 2xt^2 - t^2, \end{cases}$$

$$\begin{cases} U_2 = \frac{1}{2}(16x - 2)t^2, \\ V_2 = \frac{1}{2}(8x - 4)t^2, \\ W_2 = x(4x - 2)t^2, \end{cases}$$

$$\begin{cases} U_3 = 3t^4 + (x(8x - 4) + x(4x - 2))t^2, \\ V_3 = 2x(4x - 2)t^2, \\ W_3 = 2x^2(4x - 2)t^2, \end{cases}$$

:

Then, the approximate solutions, in series form, are

$$\begin{cases} u \approx 1 + x^2 - t^2, \\ v \approx x^2 - 1 + t^2, \\ w \approx x^2 - 1 - t^2. \end{cases}$$

For solving this system by using NHPM, we consider the following homotopy:

$$\frac{\partial U}{\partial t}(x,t) = u_0(x,t) - p\left(u_0(x,t) - \frac{\partial W}{\partial x}\frac{\partial V}{\partial t} - \frac{1}{2}\frac{\partial W}{\partial t}\frac{\partial^2 U}{\partial x^2} + 4xt\right),$$

$$\frac{\partial V}{\partial t}(x,t) = v_0(x,t) - p\left(v_0(x,t) - \frac{\partial W}{\partial t}\frac{\partial^2 U}{\partial x^2} - 6t\right),$$

$$\frac{\partial W}{\partial t}(x,t) = w_0(x,t) - p\left(w_0(x,t) - \frac{\partial^2 U}{\partial x^2} - \frac{\partial V}{\partial x}\frac{\partial W}{\partial t} - 4xt + 2t + 2\right).$$
(17)

Applying the inverse operator, $L^{-1} = \int_0^t (.) dt$, to both sides of the system (17) leads to

$$U(x,t) = U(x,0) + \int_0^t u_0(x,t) dt - p \int_0^t \left(u_0(x,t) - \frac{\partial W}{\partial x} \frac{\partial V}{\partial t} - \frac{1}{2} \frac{\partial W}{\partial t} \frac{\partial^2 U}{\partial x^2} + 4xt \right) dt,$$

$$V(x,t) = V(x,0) + \int_0^t v_0(x,t) dt - p \int_0^t \left(v_0(x,t) - \frac{\partial W}{\partial t} \frac{\partial^2 U}{\partial x^2} - 6t \right) dt,$$

$$W(x,t) = W(x,0) + \int_0^t w_0(x,t) dt - p \int_0^t \left(w_0(x,t) - \frac{\partial^2 U}{\partial x^2} - \frac{\partial V}{\partial x} \frac{\partial W}{\partial t} - 4xt + 2t + 2 \right) dt.$$

$$(18)$$

Suppose that the solutions of system (18) are in the form (5); substituting Eqs. (5) into Eqs. (18), collecting the terms with the same powers of p, and equating each coefficient of p to zero results in

$$p^{0}: \begin{cases} U_{0}(x,t) = U(x,0) + \int_{0}^{t} u_{0}(x,t) dt, \\ V_{0}(x,t) = V(x,0) + \int_{0}^{t} v_{0}(x,t) dt, \\ W_{0}(x,t) = W(x,0) + \int_{0}^{t} w_{0}(x,t) dt, \end{cases}$$

$$p^{1}: \begin{cases} U_{1}(x,t) = \int_{0}^{t} \left(-u_{0}(x,t) + \frac{\partial W_{0}}{\partial x} \frac{\partial V_{0}}{\partial t} + \frac{1}{2} \frac{\partial W_{0}}{\partial t} \frac{\partial^{2} U_{0}}{\partial x^{2}} - 4xt\right) dt, \\ V_{1}(x,t) = \int_{0}^{t} \left(-v_{0}(x,t) + \frac{\partial W_{0}}{\partial t} \frac{\partial^{2} U_{0}}{\partial x^{2}} + 6t\right) dt, \\ W_{1}(x,t) = \int_{0}^{t} \left(-w_{0}(x,t) + \frac{\partial^{2} U_{0}}{\partial x^{2}} + \frac{\partial V_{0}}{\partial x} \frac{\partial W_{0}}{\partial t} + 4xt - 2t - 2\right) dt, \end{cases}$$

$$p^{2}: \begin{cases} U_{2}(x,t) = \int_{0}^{t} \left(\frac{\partial W_{1}}{\partial x} \frac{\partial V_{0}}{\partial t} + \frac{\partial W_{0}}{\partial x} \frac{\partial V_{1}}{\partial t} + \frac{1}{2} \frac{\partial W_{0}}{\partial t} \frac{\partial^{2} U_{1}}{\partial x^{2}} + \frac{1}{2} \frac{\partial W_{1}}{\partial t} \frac{\partial^{2} U_{0}}{\partial x^{2}}\right) dt, \\ V_{2}(x,t) = \int_{0}^{t} \left(\frac{\partial W_{0}}{\partial x} \frac{\partial^{2} U_{1}}{\partial x^{2}} + \frac{\partial W_{0}}{\partial x} \frac{\partial^{2} U_{0}}{\partial x^{2}}\right) dt, \\ W_{2}(x,t) = \int_{0}^{t} \left(\frac{\partial^{2} U_{1}}{\partial x^{2}} + \frac{\partial V_{0}}{\partial x} \frac{\partial W_{1}}{\partial t} + \frac{\partial V_{1}}{\partial x} \frac{\partial W_{0}}{\partial t}\right) dt \end{cases}$$

:

$$p^{j}: \begin{cases} U_{j+1}(x,t) = \int_{0}^{t} \left(\sum_{k=0}^{j} \left(\frac{\partial W_{k}}{\partial x} \frac{\partial V_{j-k}}{\partial x} + \frac{1}{2} \frac{\partial W_{k}}{\partial t} \frac{\partial^{2} U_{j-k}}{\partial x^{2}} \right) \right) dt, \\ V_{j+1}(x,t) = \int_{0}^{t} \left(\sum_{k=0}^{j} \frac{\partial W_{k}}{\partial x} \frac{\partial^{2} U_{j-k}}{\partial x^{2}} \right) dt, \\ W_{j+1}(x,t) = \int_{0}^{t} \left(\frac{\partial^{2} U_{j}}{\partial x^{2}} + \sum_{k=0}^{j} \frac{\partial V_{k}}{\partial x} \frac{\partial W_{j-k}}{\partial t} \right) dt \end{cases}$$

:

Assume that

$$\begin{cases} u_0(x,t) = \sum_{n=0}^{\infty} a_n(x)t^k, & U(x,0) = u(x,0), \\ v_0(x,t) = \sum_{n=0}^{\infty} b_n(x)t^k, & V(x,0) = v(x,0), \\ w_0(x,t) = \sum_{n=0}^{\infty} c_n(x)t^k, & W(x,0) = w(x,0) \end{cases}$$

If we set the Taylor series of $U_1(x, t)$, $V_1(x, t)$, and $W_1(x, t)$ at t = 0 equal to zero, then

$$\begin{split} U_1(x,y) &= \left(-a_0(x) + 2xb_0(x) + c_0(x) \right)t + \left(-\frac{1}{2}a_1(x) + xb_1(x) + \frac{1}{2}b_0(x)c_0'(x) + \frac{1}{2}c_1(x) + \frac{1}{4}a_0''(x)c_0(x) - 2x \right)t^2 \\ &+ \left(-\frac{1}{3}a_2(x) + \frac{2}{3}xb_2(x) + \frac{1}{3}b_1(x)c_0'(x) + \frac{1}{6}b_0(x)c_1'(x) + \frac{1}{3}c_2(x) + \frac{1}{12}a_1''(x)c_0(x) + \frac{1}{6}a_0''(x)c_1(x) \right)t^3 \\ &+ \left(-\frac{1}{4}a_3(x) + \frac{1}{2}xb_3(x) + \frac{1}{4}b_2(x)c_0'(x) + \frac{1}{8}b_1(x)c_1'(x) + \frac{1}{12}b_0(x)c_2'(x) \right)t^4 + \dots = 0, \\ V_1(x,y) &= \left(-b_0(x) + 2c_0(x) \right)t + \left(-\frac{1}{2}b_1(x) + c_1(x) + \frac{1}{2}c_0(x)a_0''(x) + 3 \right)t^2 \\ &+ \left(-\frac{1}{3}b_2(x) + \frac{2}{3}c_2(x) + \frac{1}{6}c_0(x)a_1''(x) + \frac{1}{3}c_1(x)a_0''(x) \right)t^3 \\ &+ \left(-\frac{1}{4}b_3(x) + \frac{1}{2}c_3(x) + \frac{1}{12}c_0(x)a_2''(x) + \frac{1}{8}c_1(x)a_1''(x) + \frac{1}{4}c_2(x)a_0''(x) \right)t^4 + \dots = 0, \\ W_1(x,y) &= \left(-c_0(x) + 2xc_0(x) \right)t + \left(-\frac{1}{2}c_1(x) + \frac{1}{2}a_0''(x) + xc_1(x) + b_0'(x)c_0(x) + 2x - 1 \right)t^2 \\ &+ \left(-\frac{1}{3}c_3(x) + \frac{1}{6}a_1''(x) + \frac{2}{3}xc_2(x) + \frac{1}{3}b_0'(x)c_1(x) + \frac{1}{6}b_1'(x)c_0(x) \right)t^3 \\ &+ \left(-\frac{1}{4}c_3(x) + \frac{1}{12}a_2''(x) + \frac{1}{2}xc_3(x) + \frac{1}{4}b_0'(x)c_2(x) + \frac{1}{8}b_1'(x)c_1(x) + \frac{1}{12}b_2'(x)c_0(x) \right)t^4 + \dots = 0. \end{split}$$

It follows easily that

$$a_0(x) = 0,$$
 $a_1(x) = -2,$ $a_2(x) = 0,$ $a_3(x) = 0,$ $a_4(x) = 0, \dots$
 $b_0(x) = 0,$ $b_1(x) = 2,$ $b_2(x) = 0,$ $b_3(x) = 0,$ $b_4(x) = 0, \dots$
 $c_0(x) = 0,$ $c_1(x) = -2,$ $c_2(x) = 0,$ $c_3(x) = 0,$ $c_4(x) = 0, \dots$

Therefore, the exact solutions of the system of partial differential equations can be expressed as follows:

$$u(x,t) = U_0(x,t) = x^2 + 1 + a_0(x)t + \frac{1}{2}a_1(x)t^2 + \frac{1}{3}a_2(x)t^3 + \frac{1}{4}a_3(x)t^4 + \dots = x^2 - t^2 + 1,$$

$$v(x,t) = V_0(x,t) = x^2 - 1 + b_0(x)t + \frac{1}{2}b_1(x)t^2 + \frac{1}{3}b_2(x)t^3 + \frac{1}{4}b_3(x)t^4 + \dots = x^2 + t^2 - 1,$$

$$w(x,t) = W_0(x,t) = x^2 - 1 + c_0(x)t + \frac{1}{2}c_1(x)t^2 + \frac{1}{3}c_2(x)t^3 + \frac{1}{4}c_3(x)t^4 + \dots = x^2 - t^2 - 1.$$

4. Conclusions

In this article, a new modification of HPM, called NHPM, has been introduced for solving systems of non-linear partial differential equations. This method has been applied to three examples successfully, and exact solutions of the equations are achieved, where traditional HPM leads to an approximate solution. Numerical results reveal that NHPM is a powerful tool for solving linear and non-linear initial and boundary value problems. The basic idea described in this paper is strong enough to be employed to solve other functional equations. The convergence of the method is under study by our research group. The computations associated with the examples were performed using Maple 13.

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