# Estimate of the solution of the Dirichlet problem for parabolic equations and applications 

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#### Abstract

In the present paper we consider the Dirichlet problem for quasilinear nonuniformly parabolic equations. A new sufficient condition which guarantees the a priori estimate of the maximum of the modulus of the solution is formulated. A several applications of this estimate are given. © 2002 Elsevier Science (USA). All rights reserved.


## 1. Introduction and main result

Consider the following problem:

$$
\begin{equation*}
u_{t}-a_{i j}(t, \mathbf{x}, u, \nabla u) u_{x_{i} x_{j}}=F(t, \mathbf{x}, u, \nabla u) \quad \text { in } Q_{T}=\Omega \times(0, T) \tag{1.1}
\end{equation*}
$$

(we assume the usual summation convention),

$$
\begin{equation*}
u(0, \mathbf{x})=\phi(\mathbf{x}) \quad \text { in } \Omega \quad \text { and } \quad u=\chi(t, \mathbf{x}) \quad \text { on } S_{T}=\partial \Omega \times[0, T] \tag{1.2}
\end{equation*}
$$

where $\Omega \subset \mathbf{R}^{n}$ is a bounded domain, $\mathbf{x}=\left(x_{1}, x_{2}, \ldots, x_{n}\right), \nabla u=\left(u_{x_{1}}, \ldots, u_{x_{n}}\right)$, $a_{i j}=a_{j i}, i, j=1, \ldots, n$. Without loss of generality suppose that $\Omega$ lies in the strip $-l_{1}<x_{1}<l_{1}$. Assume that the functions $a_{i j}(t, \mathbf{x}, u, \mathbf{p}), F(t, \mathbf{x}, u, \mathbf{p})$ are

[^0]defined on the set $\bar{Q}_{T} \times \mathbf{R} \times \mathbf{R}^{n}$ take finite values for $(t, \mathbf{x}) \in Q_{T}$ and finite $u, \mathbf{p}$ and
\[

$$
\begin{equation*}
a_{i j}(t, \mathbf{x}, u, \mathbf{p}) \xi_{i} \xi_{j} \geqslant 0 \quad \text { for all } \xi \in \mathbf{R}^{n},(t, \mathbf{x}, u, \mathbf{p}) \in \bar{Q}_{T} \times \mathbf{R} \times \mathbf{R}^{n} . \tag{1.3}
\end{equation*}
$$

\]

There are several sufficient conditions which guarantee the boundedness of a classical solution of problem (1.1), (1.2) (see [1-5]). Remind that a classical solution is a solution belonging to $C^{0}\left(\bar{Q}_{T}\right) \cap C_{t, \mathbf{\mathbf { x }}}^{1,2}\left(Q_{T}\right)$. Here $C_{t, \mathbf{\mathbf { x }}}^{1,2}\left(Q_{T}\right)$ is the set of functions having the first derivative with respect to $t$ and the second derivatives with respect to $\mathbf{x}$ continuous in $Q_{T} ; C^{0}\left(\bar{Q}_{T}\right)$ is the set of continuous in $\bar{Q}_{T}$ functions. In the present paper we give a new sufficient condition guaranteeing the a priori estimate of $|u|$.

Suppose that the right side of the equation can be represented in the form

$$
\begin{equation*}
F(t, \mathbf{x}, u, \mathbf{p})=f_{1}(t, \mathbf{x}, u, \mathbf{p})+f_{2}(t, \mathbf{x}, u, \mathbf{p}), \tag{1.4}
\end{equation*}
$$

where $f_{1}$ and $f_{2}$ have different properties. Suppose that $f_{1}(t, \mathbf{x}, u, \mathbf{p})$ for $(t, \mathbf{x}) \in$ $Q_{T}$ and any $u, p_{1}$ satisfies the following restriction:

$$
\begin{equation*}
\left|f_{1}\left(t, \mathbf{x}, u, p_{1}, 0, \ldots, 0\right)\right| \leqslant a_{11}\left(t, \mathbf{x}, u, p_{1}, 0, \ldots, 0\right) \psi\left(\left|p_{1}\right|\right) \tag{1.5}
\end{equation*}
$$

Here $\psi(\rho)$ is a continuously differentiable function, $\psi(\rho)>0$ for $\rho>0, \psi(0) \geqslant$ 0 and we assume that

$$
\begin{equation*}
\int_{0}^{+\infty} \frac{d \rho}{\psi(\rho)}>2 l_{1} \tag{1.6}
\end{equation*}
$$

In order to formulate the conditions on $f_{2}(t, \mathbf{x}, u, \mathbf{p})$ let us introduce the function $h\left(x_{1}\right)$ as a solution of the following problem:

$$
\begin{equation*}
h^{\prime \prime}+\psi\left(\left|h^{\prime}\right|\right)=0, \quad h\left(-l_{1}\right)=M, \quad h\left(l_{1}\right)=H . \tag{1.7}
\end{equation*}
$$

Here $M \geqslant m \equiv \max \left\{\sup _{\Omega}|\phi|, \sup _{S_{T}}|\chi|\right\}$, the constant $H$ will be defined below. Represent the solution of problem (1.7) in parametric form using the substitution $q(h)=h^{\prime}\left(x_{1}\right), q_{x_{1}}^{\prime}\left(h\left(x_{1}\right)\right)=q(h) q^{\prime}(h):$

$$
h(q)=\int_{q}^{q_{1}} \frac{\rho d \rho}{\psi(\rho)}+M, \quad x_{1}(q)=\int_{q}^{q_{1}} \frac{d \rho}{\psi(\rho)}-l_{1},
$$

where $q \in\left[q_{0}, q_{1}\right]$ and $q_{0}, q_{1}$ are chosen such that $0<q_{0}<q_{1}<+\infty$ and

$$
x_{1}(q)=\int_{q_{0}}^{q_{1}} \frac{d \rho}{\psi(\rho)}=2 l_{1}
$$

This is possible due to (1.6). Put

$$
H=\int_{q_{0}}^{q_{1}} \frac{\rho d \rho}{\psi(\rho)}+M
$$

Assume that $f_{2}$ satisfies the following conditions:

$$
\begin{align*}
& f_{2}\left(t, \mathbf{x}, u, p_{1}, 0, \ldots, 0\right) \leqslant 0 \quad \text { for } u \geqslant M, p_{1} \in\left[q_{0}, q_{1}\right]  \tag{1}\\
& f_{2}\left(t, \mathbf{x}, u,-p_{1}, 0, \ldots, 0\right) \geqslant 0 \quad \text { for } u \leqslant-M, p_{1} \in\left[q_{0}, q_{1}\right] . \tag{2}
\end{align*}
$$

Let us formulate now the main result.

Theorem. Let $u(t, \mathbf{x})$ be a classical solution of problem (1.1), (1.2). Suppose that conditions (1.3)-(1.6), (1.8) are fulfilled; then

$$
\sup _{Q_{T}}|u| \leqslant h\left(x_{1}\right) \leqslant H .
$$

The proof will be given in the second section.
In the third section we give examples of applications of the theorem. In particular, from the theorem one can obtain the following fact. Consider the linear heat equation for the anisotropic media, i.e. we suppose that in different directions the heat conductivity is different. In $x_{i}$ direction the heat conductivity coefficient is $a_{i i}(t, \mathbf{x}) \geqslant 0$ :

$$
\begin{equation*}
u_{t}=\operatorname{div}\left(a_{11} u_{x_{1}}, \ldots, a_{n n} u_{x_{n}}\right)+f(t, \mathbf{x}), \tag{1.9}
\end{equation*}
$$

here $u(t, \mathbf{x}) \geqslant 0$ is an absolute temperature, $f(t, \mathbf{x}) \geqslant 0$ is a source. Consider problem (1.9), (1.2) and assume that $\phi(\mathbf{x}) \equiv \chi(t, \mathbf{x}) \equiv 0$. Suppose for simplicity that $a_{11}$ is constant. From the theorem it follows that

$$
\sup _{Q_{T}}|u| \leqslant \frac{3}{2} l_{1}^{2} \frac{\sup f}{a_{11}} .
$$

We observe here the phenomenon of cooling of a body for fixed $l_{1}$ and $f(t, \mathbf{x})$ when the heat conductivity increases in one direction. In fact, $u(t, \mathbf{x}) \rightarrow 0$ as $a_{11} \rightarrow+\infty$. Similar effect we have in the nonlinear case. This phenomenon has a simple physical interpretation (see Section 3).

Moreover, from the theorem we can easily obtain the standard a priori estimate of sup $|u|$ for linear equation (see Section 3, estimate (3.4)) as well as for the nonlinear one (see Section 2, Remark 4).

## 2. Proof of the theorem

Let $u(t, \mathbf{x})$ be a classical solution of problem (1.1), (1.2). Define the operator $L$ by the following:

$$
L(u) \equiv u_{t}-A_{11}\left(u_{x_{1} x_{1}}+\psi\left(\left|u_{x_{1}}\right|\right)\right)-\sum_{i+j>2} A_{i j} u_{x_{i} x_{j}},
$$

where $A_{i j}(t, \mathbf{x}) \equiv a_{i j}(t, \mathbf{x}, u, \nabla u)$. Obviously $h^{\prime}\left(x_{1}\right)=q \geqslant 0$ (see (1.7)) and hence $u-h$ is nonpositive on $\Omega \cup S_{T}$. It is clear that $L(h)=0$ and thus for $w=u-h$ we have

$$
\begin{aligned}
L(u)-L(h) & \equiv L_{0}(w) \equiv w_{t}-A_{11}\left(w_{x_{1} x_{1}}+\beta w_{x_{1}}\right)-\sum_{i+j>2} A_{i j} w_{x_{i} x_{j}} \\
& =F(t, \mathbf{x}, u, \nabla u)-a_{11}(t, \mathbf{x}, u, \nabla u) \psi\left(\left|u_{x_{1}}\right|\right)
\end{aligned}
$$

where (from the mean value theorem)

$$
\beta=\psi^{\prime}\left(\rho^{*}\right) \frac{\left|u_{x_{1}}\right|-\left|h_{x_{1}}\right|}{u_{x_{1}}-h_{x_{1}}} .
$$

Due to the fact that $\psi(\rho)$ is a continuously differentiable function and that $u(t, \mathbf{x})$ is a classical solution we conclude that $|\beta| \leqslant\left|\psi^{\prime}\left(\rho^{*}\right)\right|<+\infty$ in $\bar{Q}_{T} \backslash\left(S_{T} \cup \Omega\right)$.

Consider the function $\widetilde{w}=w e^{-t}$. One can easily see that

$$
\begin{aligned}
\tilde{w}_{t} & +\widetilde{w}-A_{11}\left(\widetilde{w}_{x_{1} x_{1}}+\beta \tilde{w}_{x_{1}}\right)-\sum_{i+j>2} A_{i j} \tilde{w}_{x_{i} x_{j}} \\
& =e^{-t}\left(F(t, \mathbf{x}, u, \nabla u)-a_{11}(t, \mathbf{x}, u, \nabla u) \psi\left(\left|u_{x_{1}}\right|\right)\right) .
\end{aligned}
$$

Suppose that $\widetilde{w}$ achieves positive maximum at the point $N \in \bar{Q}_{T} \backslash\left(S_{T} \cup \Omega\right)$. Then at this point $\widetilde{w}>0$ and $\nabla \widetilde{w}=0$, i.e. $u>h \geqslant M$ and $u_{x_{1}}=h^{\prime}>0, u_{x_{i}}=0$ for $i=2,3, \ldots, n$. Due to (1.4), (1.5), (1.81) we have

$$
\begin{aligned}
\widetilde{w}_{t}+\widetilde{w} & -A_{11}\left(\widetilde{w}_{x_{1} x_{1}}+\beta \tilde{w}_{x_{1}}\right)-\left.\sum_{i+j>2} A_{i j} \tilde{w}_{x_{i} x_{j}}\right|_{N} \\
=e^{-t} & \left(f_{1}\left(t, \mathbf{x}, u, u_{x_{1}}, 0, \ldots, 0\right)-a_{11}\left(t, \mathbf{x}, u, u_{x_{1}}, 0, \ldots, 0\right) \psi\left(\left|u_{x_{1}}\right|\right)\right. \\
& \left.\quad+f_{2}\left(t, \mathbf{x}, u, h^{\prime}, 0, \ldots, 0\right)\right)\left.\right|_{N} \leqslant 0
\end{aligned}
$$

This contradicts the assumption that $\widetilde{w}$ attains positive maximum in $\bar{Q}_{T}$ \} $\left(S_{T} \cup \Omega\right)$. From the nonpositivity of $\widetilde{w}$ on $S_{T} \cup \Omega$ we conclude that $\widetilde{w} \leqslant 0$ in $\bar{Q}_{T}$ and hence

$$
w=u(t, \mathbf{x})-h\left(x_{1}\right) \leqslant 0 \quad \text { in } \bar{Q}_{T} .
$$

Now consider the function $v \equiv u+h$. Obviously $v$ is nonnegative on $S_{T} \cup \Omega$, because $h^{\prime} \geqslant 0$. Define operator $L_{1}$ :

$$
L_{1}(u) \equiv u_{t}-A_{i j} u_{x_{i} x_{j}} .
$$

It is clear that

$$
L_{1}(u)=F(t, \mathbf{x}, u, \nabla u) \quad \text { and } \quad L_{1}(h)=a_{11}(t, \mathbf{x}, u, \nabla u) \psi\left(\left|h^{\prime}\right|\right) .
$$

For $\tilde{v}=v e^{-t}$ we have

$$
\tilde{v}_{t}+\tilde{v}-A_{i j} \tilde{v}_{x_{i} x_{j}}=e^{-t}\left(f(t, \mathbf{x}, u, \nabla u)+a_{11}(t, \mathbf{x}, u, \nabla u) \psi\left(\left|h^{\prime}\right|\right)\right) .
$$

Suppose that function $\tilde{v}$ attains negative minimum at the point $N \in \bar{Q}_{T} \backslash$ $\left(S_{T} \cup \Omega\right)$. At this point $\tilde{v}<0$ and $\nabla \tilde{v}=0$, i.e. $u<-h \leqslant-M$ and $u_{x_{1}}=-h^{\prime}<0$, $u_{x_{i}}=0$ for $i=2,3, \ldots, n$. Due to (1.4), (1.5), (1.82) we obtain

$$
\begin{aligned}
& \tilde{v}_{t}+\tilde{v}-\left.A_{i j} \tilde{v}_{x_{i} x_{j}}\right|_{N} \\
& =e^{-t}\left(f_{1}\left(t, \mathbf{x}, u,-h^{\prime}, 0, \ldots, 0\right)+a_{11}\left(t, \mathbf{x}, u,-h^{\prime}, 0, \ldots, 0\right) \psi\left(\left|h^{\prime}\right|\right)\right. \\
& \left.\quad \quad+f_{2}\left(t, \mathbf{x}, u,-h^{\prime}, 0, \ldots, 0\right)\right)\left.\right|_{N} \geqslant 0
\end{aligned}
$$

This contradicts the assumption that $\tilde{v}$ attains negative minimum at $N$. Taking into account that $\tilde{v}=(u+h) e^{-t} \geqslant 0$ on $S_{T} \cup \Omega$, we conclude that $\tilde{v} \geqslant 0$ in $\bar{Q}_{T}$ and hence

$$
v=u(t, \mathbf{x})+h\left(x_{1}\right) \geqslant 0 \quad \text { in } \bar{Q}_{T} .
$$

Thus we obtain that $|u(t, \mathbf{x})| \leqslant h\left(x_{1}\right) \leqslant h\left(l_{1}\right) \equiv H$. The theorem is proved.
Let us formulate several remarks.
Remark 1. Instead of conditions (1.8) we can take the following ones:

$$
\begin{align*}
& f_{2}\left(t, \mathbf{x}, u,-p_{1}, 0, \ldots, 0\right) \leqslant 0 \quad \text { for } u \geqslant M, p_{1} \in\left[q_{0}, q_{1}\right],  \tag{1}\\
& f_{2}\left(t, \mathbf{x}, u, p_{1}, 0, \ldots, 0\right) \geqslant 0 \quad \text { for } u \leqslant-M, p_{1} \in\left[q_{0}, q_{1}\right] . \tag{2}
\end{align*}
$$

In this case the barrier is a solution of the same equation as in (1.7) but with the other boundary conditions, namely $h\left(-l_{1}\right)=H, h\left(l_{1}\right)=M$. The estimate here is $\sup |u(t, \mathbf{x})| \leqslant h\left(x_{1}\right) \leqslant h\left(-l_{1}\right) \equiv H$.

Remark 2. If conditions (1.8) and (2.1) are fulfilled then $\sup |u| \leqslant h(0)$.
Remark 3. The choice of the quantity $H$ in (1.7) actually results from the necessity of the fulfillment of condition $h^{\prime}\left(x_{1}\right)>0$ for $\left|x_{1}\right|<l_{1}$.

Consider the following problem. Let $f_{1}=f_{1}(t, \mathbf{x})$ and $a_{11} \equiv 1$. Denote by $f_{0}$ the $\sup \left|f_{1}(t, \mathbf{x})\right|$. As a barrier $h\left(x_{1}\right)$ we take the solution of the equation $h^{\prime \prime}=-f_{0}$. The first boundary condition is $h\left(-l_{1}\right)=M$, instead of the second one we take the condition $h^{\prime}\left(x_{1}\right)>0$ for $\left|x_{1}\right| \leqslant l_{1}$. We obtain

$$
h\left(x_{1}\right)=-\frac{x_{1}^{2}}{2} f_{0}+l_{1} f_{0} x_{1}+M+\frac{3 l_{1}^{2}}{2} f_{0}
$$

The estimate in this case takes the following form

$$
\sup |u| \leqslant h(0)=M+2 l_{1}^{2} f_{0} .
$$

Remark 4. Assume that $f_{1}(t, \mathbf{x}, u, \mathbf{0}) \equiv 0$ and conditions (1.3), (1.4) are fulfilled. Suppose that for some constant $M \geqslant m$

$$
u f_{2}(t, \mathbf{x}, u, 0,0, \ldots, 0) \leqslant 0 \quad \text { for }|u|>M .
$$

Then

$$
\sup _{Q_{T}}|u| \leqslant M .
$$

In fact, here as a barrier we can take $h \equiv M$. Obviously $w^{0} \equiv(u-M) e^{-t} \leqslant 0$ on $S_{T} \cup \Omega$. Moreover

$$
L_{2}\left(w^{0}\right) \equiv w_{t}^{0}+w^{0}-A_{i j} w_{x_{i} x_{j}}^{0}=e^{-t} f(t, \mathbf{x}, u, \nabla u)
$$

Suppose that $w^{0}$ achieves positive maximum at the point $N \in \bar{Q}_{T} \backslash(S \cup \Omega)$. Then at this point $w^{0}>0$ and $\nabla w^{0}=0$, i.e. $u>M \geqslant 0$ and $\nabla u=0$. Due to (1.4), (1.5), (1.81) we have

$$
w_{t}^{0}+w^{0}-\left.A_{i j} w_{x_{i} x_{j}}^{0}\right|_{N}=\left.e^{-t} f_{2}(t, \mathbf{x}, u, 0,0, \ldots, 0)\right|_{N} \leqslant 0
$$

This contradicts the assumption that $w^{0}$ attains positive maximum in $\bar{Q}_{T}$ \} $\left(S_{T} \cup \Omega\right)$. From the nonpositivity of $w^{0}$ on $S_{T} \cup \Omega$ we conclude that $w^{0} \leqslant 0$ in $\bar{Q}_{T}$ and hence $u-h \leqslant 0$ in $\bar{Q}_{T}$.

Function $v^{0} \equiv(u+M) e^{-t}$ is nonnegative on $S_{T} \cup \Omega$. It is clear that

$$
v_{t}^{0}+v^{0}-A_{i j} v_{x_{i} x_{j}}^{0}=e^{-t} F(t, \mathbf{x}, u, \nabla u) .
$$

Suppose that function $v^{0}$ attains negative minimum at the point $N \in \bar{Q}_{T} \backslash$ $\left(S_{T} \cup \Omega\right)$. At this point $v^{0}<0$ and $\nabla v^{0}=0$, i.e. $u<-M \leqslant 0$ and $\nabla u=0$. Due to (1.4), (1.5), (1.82) we obtain

$$
v_{t}^{0}+v^{0}-\left.A_{i j} v_{x_{i} x_{j}}^{0}\right|_{N}=\left.e^{-t} f_{2}(t, \mathbf{x}, u, \mathbf{0})\right|_{N} \geqslant 0
$$

This contradicts the assumption that $v^{0}$ attains negative minimum at $N$. Taking into account that $v^{0} \geqslant 0$ on $S_{T} \cup \Omega$, we conclude that $v^{0} \geqslant 0$ in $\bar{Q}_{T}$ and hence $u \geqslant-M$ in $\bar{Q}_{T}$.

So we obtain that $|u(t, \mathbf{x})| \leqslant M$.

## 3. Examples

Let us first consider the linear equation

$$
\begin{align*}
& u_{t}-a_{i j}(t, \mathbf{x}) u_{x_{i} x_{j}}=f(t, \mathbf{x})+b_{i}(t, \mathbf{x}) u_{x_{i}}+c(t, \mathbf{x}) u \\
& \quad \text { in } Q_{T}=\Omega \times(0, T) . \tag{3.1}
\end{align*}
$$

Denote

$$
f_{1} \equiv f(t, \mathbf{x})+b_{i}(t, \mathbf{x}) p_{i}, \quad f_{2} \equiv c(t, \mathbf{x}) u
$$

assume that

$$
a_{11}(t, \mathbf{x}) \geqslant a_{0}>0 \quad \text { and } \quad c(t, \mathbf{x}) \leqslant 0
$$

where $a_{0}$ is some positive constant. One can easily see that conditions (1.8) as well as (2.1) are fulfilled. Condition (1.5) is satisfied with $\psi\left(\left|p_{1}\right|\right) \equiv K\left(1+\left|p_{1}\right|\right)$, where $K=\max \left\{\sup |f|, \sup \left|b_{1}\right|\right\} a_{0}^{-1}$. In this case we can easily construct the barrier in the explicit form

$$
h\left(x_{1}\right)=M-l_{1}+e^{2 K l_{1}} K^{-1}-e^{K\left(l_{1}-x_{1}\right)} K^{-1}-x_{1}
$$

as a solution of the problem

$$
\begin{aligned}
& h^{\prime \prime}\left(x_{1}\right)+K\left|h^{\prime}\left(x_{1}\right)\right|=-K, \quad h\left(-l_{1}\right)=M \\
& h\left(l_{1}\right)=M-2 l_{1}+K^{-1}\left(e^{2 K l_{1}}-1\right)
\end{aligned}
$$

Instead of the boundary condition at $l_{1}$ we can take condition $h^{\prime}\left(x_{1}\right)>0$ for $\left|x_{1}\right|<l_{1}$ (see Remark 3). Thus for the solution of problem (3.1), (1.2) from Remark 2 we obtain that

$$
\begin{equation*}
\sup _{Q_{T}}|u| \leqslant h(0)=M-l_{1}+e^{K l_{1}} K^{-1}\left(e^{K l_{1}}-1\right) \tag{3.2}
\end{equation*}
$$

If in Eq. (3.1) the coefficient $b_{1} \equiv 0$ then as a function $\psi$ we can take the constant $K_{1}=\sup |f| a_{0}^{-1}$, i.e. $h^{\prime \prime}\left(x_{1}\right)=-K_{1}$ and hence $h\left(x_{1}\right)=M-x_{1}^{2} K_{1} / 2+K_{1} l_{1} x_{1}+$ $3 K_{1} l_{1}^{2} / 2$. The estimate has the following form:

$$
\begin{equation*}
\sup _{Q_{T}}|u| \leqslant h(0)=\frac{3}{2} K_{1} l_{1}^{2}+M \tag{1}
\end{equation*}
$$

If $c(t, \mathbf{x}) \leqslant \lambda$, where $\lambda$ is some positive constant then $|u| \leqslant h(0) e^{\lambda T}$.
Consider the heat equation under the assumption that the coefficient of heat conductivity is different in different directions

$$
\begin{equation*}
u_{t}-\operatorname{div}\left(a_{11} u_{x_{1}}, \ldots, a_{n n} u_{x_{n}}\right)=f(t, \mathbf{x}) \tag{3.3}
\end{equation*}
$$

Here $u(t, \mathbf{x})$ is temperature. For the solution of problem (3.3), (1.2) the estimate (3.2) holds, here $b_{1}=a_{11 x_{1}}$. If $a_{11 x_{1}}=0$ then estimate (3.21) is valid.

Let us mention here that estimate (3.2) ((3.21)) depends not on the intensity of the source $f(t, \mathbf{x})$ but on the ratio

$$
\max \left\{\sup |f|, \sup \left|b_{1}\right|\right\} / \inf a_{11} \quad\left(\sup |f| / \inf a_{11}\right)
$$

where $a_{11}$ is a coefficient of the heat conductivity in the $x_{1}$ direction.
Let $m=0$. For the arbitrary fixed source $f$ by choosing the constant $a_{0}$ to be sufficiently big we can make $|u|$ arbitrary small. As it has been already
mentioned in the first section this fact has a simple physical explanation. The heat flow through the boundary in the $x_{1}$ direction increases as $a_{0}=\inf a_{11}$ increases assisting to the cooling of the body. Remind that according to the Fourier law the heat flow through the boundary $S_{T} \equiv \partial \Omega \times(0, T)$ is given by the integral

$$
\int_{S_{T}}\left(a_{11} u_{x_{1}}, \ldots, a_{n n} u_{x_{n}}\right) \cdot \mathbf{n} d \mathbf{x} d t
$$

where $\mathbf{n}$ is an external normal vector to the boundary.
If on the boundary we require the absence of the heat flow $\partial u / \partial \mathbf{n}=0$ (homogeneous Neumann problem), then the value of $u$ is determined only by the intensity of the sources and by the initial data and does not depend on the conductivity of the media.

Using the theorem we can obtain an estimate independent of the coefficients of the principal part of the equation. Consider Eq. (3.1), let

$$
a_{i j} \xi_{i} \xi_{j} \geqslant 0 \quad \text { and } \quad c(t, \mathbf{x}) \leqslant-c_{0}<0
$$

Put

$$
f_{1} \equiv 0, \quad f_{2} \equiv f(t, \mathbf{x})+b_{i}(t, \mathbf{x}) p_{i}+c(t, \mathbf{x}) u
$$

As a barrier $h$ here we take the constant $M=\max \left\{m, \sup |f| / c_{0}\right\}$. In that case obviously

$$
\begin{aligned}
& f_{2}\left(t, \mathbf{x}, u, h^{\prime}, 0, \ldots, 0\right) \equiv f_{2}(t, \mathbf{x}, u, \mathbf{0}) \leqslant 0 \quad \text { for } u>M \\
& f_{2}\left(t, \mathbf{x}, u,-h^{\prime}, 0, \ldots, 0\right) \equiv f_{2}(t, \mathbf{x}, u, \mathbf{0}) \geqslant 0 \quad \text { for } u<-M .
\end{aligned}
$$

Hence we conclude that for the solution of problem (3.1), (1.2) the following estimate takes place:

$$
\sup _{Q_{T}}|u| \leqslant \max \left\{m, \frac{\sup |f|}{c_{0}}\right\} .
$$

If $c(t, \mathbf{x})<\lambda$, where $\lambda$ is a positive constant, then

$$
\begin{equation*}
\sup _{Q_{T}}|u| \leqslant \inf _{\lambda>c_{0}}\left(e^{\lambda T} \max \left\{m, \frac{\sup |f|}{\lambda-c_{0}}\right\}\right), \tag{3.4}
\end{equation*}
$$

where $c_{0}=\sup c(t, \mathbf{x})$. This is a standard a priori estimate for the solution of problem (3.1), (3.2).

Let us pass to the nonlinear case. Consider the following semilinear equation

$$
\begin{equation*}
u_{t}-k(u) \Delta u=Q(u), \quad k(u)>0 . \tag{3.5}
\end{equation*}
$$

We suppose that $Q(u)$ does not satisfy condition $u Q(u) \leqslant 0$ for $|u|>M$. It is well known that generally speaking the solution of the Dirichlet problem for
that equation blows-up, i.e. there exists $t^{*}<+\infty$ such that $\sup _{\mathbf{x}}|u(t, \mathbf{x})| \rightarrow+\infty$ when $t \rightarrow t^{*}$.

From the theorem it follows that if in (3.5) function $Q(u)$ satisfies the inequality $|Q(u)| \leqslant C_{0} k(u)$, where $C_{0}$ is some positive constant, then the solution is bounded for all $t>0$. If instead of (3.5) we consider equation

$$
u_{t}-k_{i}(u) u_{x_{i} x_{i}}=Q(u), \quad k_{i}(u)>0, \quad i=1, \ldots, n,
$$

then for the boundedness of the solution it is sufficient to require the fulfillment of the inequality $|Q(u)| \leqslant C_{0} k_{i}(u)$ only for one value of index $i$.

Consider the nonlinear heat equation

$$
\begin{equation*}
u_{t}-\operatorname{div}(k(u) \nabla u)=Q(u), \tag{3.6}
\end{equation*}
$$

where $k(u)>0$ is continuously differentiable function.
Write this equation in the following form:

$$
u_{t}-k(u) \Delta u=f_{1},
$$

where $f_{1} \equiv Q(u)+k^{\prime}(u)|\nabla u|^{2}$.
Suppose that $|Q(u)| \leqslant C_{0} k(u)$ and $\left|k^{\prime}(u)\right| \leqslant C_{0} k(u)$. In that case condition (1.5) is fulfilled with $\psi\left(\left|p_{1}\right|\right)=C_{0}\left(1+p_{1}^{2}\right)$. The integral in (1.6) is equal to $\pi / 2 C_{0}$. Thus in order to obtain the estimate of $|u|$ we require the constant $C_{0}$ to be less than $\pi / 4 l_{1}$.

If $u k^{\prime}(u) \leqslant 0$ then we write Eq. (3.6) in the form

$$
u_{t}-k(u) \Delta u=Q(u)+f_{2}
$$

where $f_{2} \equiv k^{\prime}(u)|\nabla u|^{2}$. In that case we obtain the estimate of $|u|$ without supplementary assumptions on $C_{0}$.

Suppose that $Q(z) \geqslant 0$ when $z \leqslant 1$ and $u \geqslant 1$ on the parabolic boundary of the domain. One can easily see that in this case $u(t, \mathbf{x}) \geqslant 1$. In fact, let

$$
L u \equiv u_{t}-\tilde{k} \Delta u-\tilde{k}_{1}|\nabla u|^{2},
$$

where $\tilde{k}=\tilde{k}(t, \mathbf{x})=k(u), \tilde{k}_{1}=\tilde{k}_{1}(t, \mathbf{x})=k^{\prime}(u)$. Obviously $L u=Q(u)$. For $v \equiv u-1$ we obtain

$$
L u \equiv v_{t}-\tilde{k} \Delta v-\tilde{k}_{1}|\nabla v|^{2}=Q(u)
$$

We conclude that $v$ cannot attain negative values at the internal points of the domain. Due to the fact that $v \geqslant 0$ on the parabolic boundary of the domain we conclude that $v \geqslant 0$ in the whole domain.

Consider the case when the coefficient of the heat conductivity has the form $k(u)=k_{0} u^{\alpha}$, where $k_{0}>0$ and $\alpha>0$ are some constants. Suppose that $Q(z) \geqslant 0$ for $z \leqslant 1$ and $u \geqslant 1$ on the parabolic boundary of the domain. Obviously $k^{\prime}(u) \leqslant$ $\alpha k(u)$. Let $|Q(u)| \leqslant \alpha k(u)$. We put $\psi\left(\left|p_{1}\right|\right)=\alpha\left(1+p_{1}^{2}\right)$, to fulfill the conditions of the theorem it is necessary for $\alpha$ to be less than $\pi / 4 l_{1}$.

Similarly we can investigate the following equation (anisotropic case):

$$
u_{t}-\operatorname{div}\left(k_{1}(u) u_{x_{1}}, \ldots, k_{n}(u) u_{x_{n}}\right)=Q(u) .
$$

Let us demonstrate the application of the theorem on one more example. Consider the following problem (for simplicity we restrict ourselves by onedimensional case):

$$
\begin{align*}
& u_{t}-u_{x x}=l^{2} u_{x}^{2}-u^{2} \quad \text { in } Q_{T}=(-l, l) \times(0, T)  \tag{3.7}\\
& u(0, x)=u_{0}(x), \quad u(t, \pm l)=0, \quad u_{0}( \pm l)=0 . \tag{3.8}
\end{align*}
$$

Suppose that $\left|u_{0 x}(x)\right| \leqslant K$ and consequently $\left|u_{0}(x)\right| \leqslant K(l-|x|)$. Consider the auxiliary equation

$$
\begin{equation*}
u_{t}-u_{x x}=l^{2} u_{x}^{2}-f(u) \quad \text { in } Q_{T} \tag{3.9}
\end{equation*}
$$

with conditions (3.8), where

$$
f(u)= \begin{cases}u^{2}, & \text { for }|u| \leqslant K l, \\ K^{2} l^{2}, & \text { for }|u|>K l .\end{cases}
$$

Let us obtain the estimate $|u(t, x)| \leqslant K l$ for the solution of auxiliary problem (3.9), (3.8). Consider the function $v(t, x) \equiv u(t, x)+h(x)$, where $h(x)=K$. $(l+x)$. It is clear that

$$
v_{t}-v_{x x}=l^{2} u_{x}^{2}-f(u)
$$

For $\tilde{v}(t, x)=v(t, x) e^{-t}$ we obtain

$$
\tilde{v}_{t}+\tilde{v}-\tilde{v}_{x x}=\left(l^{2} u_{x}^{2}-f(u)\right) e^{-t}
$$

If the function $\tilde{v}$ attains negative minimum at the point $N \in \bar{Q}_{T} \backslash \Gamma$ ( $\Gamma$ is parabolic boundary of the domain $Q_{T}$ ), then at this point $\tilde{v}_{x}=0$, i.e. $u_{x}=-h^{\prime}=-K$ and hence

$$
\tilde{v}_{t}+\tilde{v}-\left.\tilde{v}_{x x}\right|_{N}=\left.\left(l^{2} K^{2}-f(u)\right) e^{-t}\right|_{N} \geqslant 0 .
$$

This contradicts the assumption that at the point $N$ we have negative minimum. It is clear that on $\Gamma$ the function $\tilde{v}$ is nonnegative, hence $\tilde{v} \geqslant 0$ in $\bar{Q}_{T}$. This gives us the estimate

$$
\begin{equation*}
u(t, x) \geqslant-K(l+x) \tag{3.10}
\end{equation*}
$$

Now consider the function $w(t, x) \equiv u(t, x)+h_{1}(x)$ where $h_{1}(x)=K(l-x)$. For $\widetilde{w}(t, x)=w(t, x) e^{-t}$ we obtain

$$
\widetilde{w}_{t}+\widetilde{w}-\widetilde{w}_{x x}=\left(l^{2} u_{x}^{2}-f(u)\right) e^{-t} .
$$

At the point $N_{1} \in \bar{Q}_{T} \backslash \Gamma$ of the negative minimum of function $\widetilde{w}$ we have $u_{x}=$ $h^{\prime}=K$ and

$$
\tilde{w}_{t}+\widetilde{w}-\left.\widetilde{w}_{x x}\right|_{N_{1}}=\left.\left(l^{2} K^{2}-f(u)\right) e^{-t}\right|_{N_{1}} \geqslant 0
$$

From this contradiction and from the fact that $\widetilde{w} \geqslant 0$ on the parabolic boundary of the domain $Q_{T}$ we conclude that $\widetilde{w} \geqslant 0$ in $\bar{Q}_{T}$. Hence

$$
\begin{equation*}
u(t, x) \geqslant-K(l-x) \tag{3.11}
\end{equation*}
$$

From (3.10) and (3.11) we obtain

$$
u(t, x) \geqslant-K(l-|x|) \geqslant-K l .
$$

Similarly we can obtain the estimate $u \leqslant K l$. As a consequence we conclude that Eqs. (3.9) and (3.7) coincide and the estimate $|u(t, x)| \leqslant M=K l$ holds as well for the solution of problem (3.7), (3.8).

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