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Models for Noncommuting Operators

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This paper develops a model theory for a pair of noncommuting operators. Using backward shift operators on a Fock space Rota's Theorem is generalized, i.e., it is shown that any two bounded operators on a Hilbert space are simultaneously similar to part of a pair of backward shift operators on a Fock space. These shift operators and the Fock space framework are also used to develop a dilation theory for two noncommuting operators.

1. INTRODUCTION

Rota [4] proved that any bounded operator A , on a Hilbert space \mathcal{H} , with spectral radius less than one, is similar to part of a backward shift operator. Another result along this line is given in [1, 3, 5]. It states that any contraction is unitarily equivalent to part of a co-isometry, i.e., if A on \mathcal{H} is a contraction then $HA = (V|_{\mathcal{W}})H$, where H is a unitary operator, from \mathcal{H} onto \mathcal{W} , V is a co-isometry, and \mathcal{W} is an invariant subspace for V . (An operator V on \mathcal{H} is an isometry if $V^*V = I$, the identity on \mathcal{H} . A co-isometry is the adjoint of an isometry.) In this paper we generalize the above results to a pair of bounded operators, A, N on \mathcal{H} . First it is shown that A, N are simultaneously similar to part of two shift operators. Then this result is refined; if $A^*A + N^*N \leq I$, then it is shown that A, N are simultaneously unitarily equivalent to part of two co-isometries. We say that A, N are *simultaneously similar* to [unitarily equivalent to] *part of* R, T , if (1) R, T are operators on a Hilbert space \mathcal{V} , (2) there exists an invariant subspace \mathcal{W} , for both R and T , (3) there exists a similarity [unitary] transformation H mapping \mathcal{H} onto \mathcal{W} such that $HA = (R|_{\mathcal{W}})H$ and $HN = (T|_{\mathcal{W}})H$, respectively. It is emphasized that the same operator H is used to intertwine both A with $R|_{\mathcal{W}}$ and N with $T|_{\mathcal{W}}$. (Note: A, N is *simultaneously similar* to [unitarily equivalent to] R, T if (1) holds, and there exists a similarity [unitary] transformation H mapping \mathcal{H} onto \mathcal{V} such that $HA = RH$ and $HN = TH$, respectively.)

Our model theory for noncommuting operators is motivated by problems arising in nonlinear systems [2]. It can also be viewed as a representation theory for an operator A perturbed by N . The models obtained are shift operators defined on a Fock space. The result is a generalization of the existing dilation theory for one operator [1, 3–5], a deeper understanding of how noncommuting operators interact, and a solution to certain problems in mathematical systems theory [2].

2. THE SHIFT OPERATORS S AND E

In this section we introduce several different shift operators on a Fock space. These operators will be used to develop a model theory.

First some notation is established. Throughout, all spaces are Hilbert spaces, and A , N are bounded linear operators on \mathcal{X} . The adjoint of an operator A , is denoted by A^* , the open unit disc by D , and the (n -fold) unit polydisc by $D^n = D \times D \times \cdots \times D$. The Hardy space, $\mathcal{H}_n(\mathcal{X})$ is the space of all analytic functions, f in D^n with values in the Hilbert space \mathcal{X} , such that the Taylor coefficients are square summable. Each f in $\mathcal{H}_n(\mathcal{X})$ has a power series expansion given by

$$f(\lambda_1, \lambda_2, \dots, \lambda_n) = \sum_{i_1 > 0, \dots, i_n > 0} f_{i_1, i_2, \dots, i_n} \lambda_1^{i_1} \lambda_2^{i_2} \cdots \lambda_n^{i_n}, \quad (2.1)$$

where the series converges uniformly in D^n , all f_{i_1, i_2, \dots, i_n} are elements in \mathcal{X} , and the norm is

$$\|f\|_{\mathcal{H}_n}^2 \doteq \sum_{i_n > 0, \dots, i_1 > 0} \|f_{i_1, \dots, i_n}\|_{\mathcal{X}}^2. \quad (2.2)$$

Clearly $\mathcal{H}_n(\mathcal{X})$ is a Hilbert space. The Fock space $\mathcal{F}_1(\mathcal{X})$ is the Hilbert space defined as the orthogonal direct sum of the \mathcal{H}_n 's:

$$\mathcal{F}_1(\mathcal{X}) \doteq \bigoplus_{n=1}^{\infty} \mathcal{H}_n(\mathcal{X}). \quad (2.3)$$

For convenience elements in $\mathcal{F}_1(\mathcal{X})$ are represented by two different notations: both $\bigoplus_1^{\infty} f_n$ and $\{f_1, f_2, \dots\}$ represent the same element in $\mathcal{F}_1(\mathcal{X})$.

The *backward shift operator*, S_n mapping $\mathcal{H}_n(\mathcal{X})$ into $\mathcal{H}_n(\mathcal{X})$ is the linear operator defined by

$$S_n f(\lambda_1, \dots, \lambda_n) \doteq \frac{1}{\lambda_1} [f(\lambda_1, \dots, \lambda_n) - f(0, \lambda_2, \lambda_3, \dots, \lambda_n)]. \quad (2.4)$$

Note the operator S_n only acts on the Taylor coefficients of λ_1 in the power

series expansion of f . The *generalized backward shift operator* $S_{\mathcal{X}}$ mapping $\mathcal{F}_1(\mathcal{X})$ into $\mathcal{F}_1(\mathcal{X})$ is defined by:

$$S_{\mathcal{X}} \bigoplus_1^{\infty} f_n \doteq \bigoplus_1^{\infty} S_n f_n \quad \left(\bigoplus_1^{\infty} f_n \in \mathcal{F}_1(\mathcal{X}) \right). \quad (2.5)$$

The adjoint $S_{\mathcal{X}}^*$ is

$$S_{\mathcal{X}}^* \bigoplus_1^{\infty} f_n = \bigoplus_1^{\infty} \lambda_1 f_n \quad \left(\bigoplus_1^{\infty} f_n \in \mathcal{F}_1(\mathcal{X}) \right). \quad (2.6)$$

Clearly $S_{\mathcal{X}}^*$ is an isometry. Thus, $S_{\mathcal{X}}$ is a co-isometry.

The *evaluation operator*, E_n mapping $\mathcal{H}_n(\mathcal{X})$ into $\mathcal{H}_{n-1}(\mathcal{X})$ for $n \geq 1$ is given by

$$E_n f(\lambda_1, \lambda_2, \dots, \lambda_{n-1}) \doteq f(0, \lambda_1, \lambda_2, \dots, \lambda_{n-1}) \quad (f \in \mathcal{H}_n(\mathcal{X})). \quad (2.7)$$

The E_n operator evaluates λ_1 at zero and relabels the complex variables $\lambda_i \rightarrow \lambda_{i-1}$. By convention $\mathcal{H}_0(\mathcal{X}) \doteq \mathcal{X}$. Thus $E_1 f(\lambda_1) = f(0) \in \mathcal{X}$ if $f \in \mathcal{H}_1(\mathcal{X})$. The *generalized evaluation operator* $E_{\mathcal{X}}$ mapping $\mathcal{F}_1(\mathcal{X})$ into $\mathcal{F}_1(\mathcal{X})$ is

$$E_{\mathcal{X}} \bigoplus_1^{\infty} f_n \doteq \bigoplus_{n=1}^{\infty} E_{n+1} f_{n+1} \quad (2.8)$$

$$\doteq \{f_2(0, \lambda_1), f_3(0, \lambda_1, \lambda_2), f_4(0, \lambda_1, \lambda_2, \lambda_3), f_5(0, \lambda_1, \lambda_2, \lambda_3, \lambda_4), \dots\},$$

where $\bigoplus_1^{\infty} f_n = \{f_1(\lambda_1), f_2(\lambda_1, \lambda_2), \dots\} \in \mathcal{F}_1(\mathcal{X})$. The adjoint of $E_{\mathcal{X}}$ is given by

$$E_{\mathcal{X}}^* \bigoplus_1^{\infty} f_n \doteq \{0, f_1(\lambda_2), f_2(\lambda_2, \lambda_3), f_4(\lambda_2, \lambda_3, \lambda_4), \dots\}. \quad (2.9)$$

Clearly, $E_{\mathcal{X}}^*$ is an isometry and $E_{\mathcal{X}}$ is a co-isometry. The subscript \mathcal{X} is dropped from S and E when the underlying space is understood.

The operators S and E are the models we use. It turns out that “any” pair of operators A, N are simultaneously similar to part of S and E ; see Proposition 1. Furthermore, these operators have several interesting properties. The spectrum of S and E is the closed unit disc; the point spectrum of S and E is the open unit disc (Problem 67 of [3]). It is easy to verify that $\text{ran } S^*$ is orthogonal to $\text{ran } E^*$ (ran denotes the range). Furthermore, $\mathcal{F}_1(\mathcal{X})$ is the orthogonal direct sum of $\text{ran } S^*$, $\text{ran } E^*$ and \mathcal{X} (identifying \mathcal{X} with the obvious subspace of $\mathcal{H}_1(\mathcal{X})$). Thus $S^*S + E^*E \leq I$. The dimension of \mathcal{X} is called the *multiplicity* of S and E , in accordance

with the usual definition for shift operators, since \mathcal{L} is cyclic for the algebra generated by S^* and E^* .

Finally, we introduce an operator Φ . Let A, N be two bounded linear operators on \mathcal{L} and let

$$F_i \doteq (I - \lambda_i A)^{-1} = \sum_{n=0}^{\infty} A^n \lambda_i^n. \quad (2.10)$$

Define a sequence of mappings $\Phi_i: \mathcal{L} \rightarrow \mathcal{H}_i(\mathcal{L})$ by $\Phi_1(\lambda_1) \doteq F_1$; $\Phi_2(\lambda_1, \lambda_2) \doteq F_2 N F_1$ and generally

$$\Phi_n(\lambda_1, \lambda_2, \dots, \lambda_n) \doteq F_n N F_{n-1} N \cdots N F_1 = F_n N \Phi_{n-1} \quad (n \geq 2). \quad (2.11)$$

In the case that

$$\sum_{i=1}^{\infty} \|\Phi_i x\|^2 \leq (\text{const.}) \|x\|^2 \quad (\text{for all } x \in \mathcal{L}) \quad (2.12)$$

define the bounded linear operator Φ from \mathcal{L} into $\mathcal{F}_1(\mathcal{L})$ by

$$\Phi x = \bigoplus_{n=1}^{\infty} \Phi_n x = \{\Phi_1 x, \Phi_2 x, \Phi_3 x, \dots\}. \quad (2.13)$$

Throughout, A and N are fixed, and Φ always refers to the above transformation. Clearly $\|\Phi x\| \geq \|x\|$ for all $x \in \mathcal{L}$. Hence Φ (if defined) always has closed range and is a similarity transformation from \mathcal{L} onto its range.

3. SIMULTANEOUS SIMILARITY

PROPOSITION 1. *Let A, N be operators on \mathcal{L} such that (2.12) holds. Then A and N are simultaneously similar to part of S and E on $\mathcal{F}_1(\mathcal{L})$.*

Proof. First we show that $S\Phi = \Phi A$ and $E\Phi = \Phi N$. The former equality follows from $S_1 F_1 = F_1 A$ (see (2.10)) and

$$\begin{aligned} S\Phi &= \bigoplus_1^{\infty} S_n \Phi_n = \bigoplus_1^{\infty} F_n N F_{n-1} N \cdots N S_1 F_1 \\ &= \bigoplus_1^{\infty} F_n N F_{n-1} N \cdots N F_1 A \\ &= \bigoplus_1^{\infty} \Phi_n A = \Phi A. \end{aligned} \quad (3.1)$$

The other equality follows from $F_1(0) = I$ and

$$\begin{aligned}
 E\Phi &= \bigoplus_1^{\infty} E_{n+1} \Phi_{n+1} \\
 &= \bigoplus_{n=1}^{\infty} E_{n+1} F_{n+1} N F_n N \cdots N F_1 \\
 &= \bigoplus_1^{\infty} F_n N F_{n-1} N \cdots N F_1 N F_1(0) = \Phi N.
 \end{aligned} \tag{3.2}$$

Since the ran Φ is invariant for both S and E , the proof is complete.

Clearly (2.12) does not hold for all A and N . For instance choose $A = I$ and $N = 0$. However, there always exists a $\varepsilon > 0$ such that the corresponding condition for εA and εN holds. Therefore our models S, E for A, N are perfectly general.

COROLLARY 1. *If A, N are bounded operators on \mathcal{E} , and $A^*A + N^*N \leq rI$, where $r < 1$ then A, N are simultaneously similar to part of S, E on $\mathcal{F}_1(\mathcal{E})$.*

Proof. We must verify that Φ is a bounded operator. Let P be any positive self-adjoint operator, and L be the transformation mapping positive operators into positive operators defined by $LP \doteq A^*PA + N^*PN$. Clearly $LI \leq rI$ and $LP \leq LQ$ if $P \leq Q$. Thus

$$L^n I = LL^{n-1} I \leq rL^{n-1} I \leq r^n I. \tag{3.3}$$

This implies that $\sum_{i=0}^n L^i I$ is an increasing sequence of positive operators bounded by $(1-r)^{-1}$. By Problem 94 of [3], this sequence has a limit

$$\begin{aligned}
 R &= \sum_{i=0}^{\infty} L^i I = I + A^*A + N^*N + A^{*2}A^2 + A^*N^*NA \\
 &\quad + N^*A^*AN + N^{*2}N^2 + \cdots
 \end{aligned} \tag{3.4}$$

in the strong operator topology.

The expansion for $\|\Phi x\|^2$ is

$$\begin{aligned}
 \|\Phi x\|^2 &= \|x\|^2 + \sum_{i=1}^{\infty} (A^{*i} A^i x, x) \\
 &\quad + \sum_{i>0, j>0} (A^{*j} N^* A^{*i} A^i N A^j x, x) + \cdots.
 \end{aligned} \tag{3.5}$$

It is easy to show that (3.4) and (3.5) contain exactly the same terms, i.e., $(Rx, x) = \|\Phi x\|^2$ for all $x \in \mathcal{X}$. Since R is bounded

$$\|\Phi x\|^2 = \sum_{i=1}^{\infty} \|\Phi_i x\|^2 = (Rx, x) \leq M \|x\|^2 \quad (3.6)$$

and the proof is complete.

Equation (3.6) also proves

COROLLARY 2. *If $A^*A + N^*N \leq rI$ for some $r < 1$, then (2.12) holds.*

A converse to Corollary 2 is

COROLLARY 3. *Let A, N be bounded operators on \mathcal{X} . If (2.12) holds, then there exists a Hilbert norm $\|\cdot\|_0$ on \mathcal{X} equivalent to $\|\cdot\|$ such that*

$$\|Ax\|_0^2 + \|Nx\|_0^2 \leq r \|x\|_0^2 \quad (x \in \mathcal{X}) \quad (3.7)$$

for some $r < 1$.

Proof. The proof is omitted. It is almost identical to problem 122 in [3]; the other norm on \mathcal{X} is defined by $\|x\|_0^2 \doteq \|\Phi x\|_{\mathcal{F}}^2$. ■

Corollaries 2 and 3 show that (2.12) holds if and only if A, N are simultaneously similar to a pair A_0, N_0 such that for some $r < 1$ and all $x \in \mathcal{X}$,

$$\|A_0 x\|^2 + \|N_0 x\|^2 \leq r \|x\|^2.$$

If $N = 0$ the above reduces to the following standard result (Problem 122 in [3]): A on \mathcal{X} is similar to a strict contraction if and only if

$$\sum_{i=0}^{\infty} \|A^i x\|^2 \leq M \|x\|^2 \quad (x \in \mathcal{X}). \quad (3.8)$$

(T is a strict contraction if $\|T\| < 1$.) In other words, the spectral radius of A is strictly less than one if and only if (3.8) holds.

Remark. In Proposition 1 and Corollaries 2, 3 the condition (2.12) plays an important role. One can express this condition through a Lyapunov equation. We claim that (2.12) holds if and only if there exists a positive operator P such that $0 < P < \infty$ and

$$P - A^*PA - N^*PN = I \quad (0 < P < \infty). \quad (3.9)$$

Assume (2.12). Then $P \doteq \Phi^* \Phi$ satisfies (3.9). This follows from the expansion of $\Phi^* \Phi$:

$$\begin{aligned} \Phi^* \Phi &\doteq \sum_{i>0} A^*{}^i A + \sum_{j>0, i>0} A^*{}^j N^* A^*{}^i A^i N A^j \\ &+ \sum_{k>0, j>0, i>0} A^*{}^k N^* A^*{}^j N^* A^*{}^i A^i N A^j N A^k + \dots \end{aligned} \quad (3.10)$$

Assume P satisfies (3.9). Let \mathcal{X}_0 be the Hilbert space \mathcal{X} equipped with the following inner product $(x, x)_0 \doteq (Px, x)$. Clearly $\|\cdot\|$ and $\|\cdot\|_0$ are equivalent norms. Set $Q = A^*PA + N^*PN$. Using (3.9) and $P = I + Q$ a simple calculation gives

$$\begin{aligned} \frac{\|Ax\|_0^2 + \|Nx\|_0^2}{\|x\|_0^2} &= \frac{(A^*PA + N^*PNx, x)}{(Px, x)} \\ &= \frac{(Qx, x)}{(x, x) + (Qx, x)} \\ &= \frac{(Qx, x)}{(x, x)} \leq \frac{\|Q\|}{1 + \|Q\|} < 1. \end{aligned} \quad (3.11)$$

Hence (3.7) holds. Corollary 2 gives (2.12). In many applications, obtaining a solution P to (3.9) is easier than proving that (2.12) holds. Finally, it is noted that the solution to (3.9) (if it exists) is unique.

4. UNITARY EQUIVALENCE

If $N = 0$ then Proposition 1 reduces to Rota's Theorem [4]. Problem 121 of [3] is a refinement of Rota's Theorem. In our more general setting, this refinement becomes

PROPOSITION 2. *Let A, N be bounded operators on \mathcal{X} , such that $A^*A + N^*N \leq I$, and let Φ_n for $n \geq 1$ be defined by (2.11). If $A^n \rightarrow 0$ strongly and $\|N\Phi_n x\|_{\mathcal{F}_n} \rightarrow 0$ for all $x \in \mathcal{X}$, as $n \rightarrow \infty$, then A, N are simultaneously unitarily equivalent to part of the shifts $S_{\mathcal{D}}, E_{\mathcal{D}}$ on $\mathcal{F}_1(\mathcal{D})$, for some closed linear subspace \mathcal{D} of \mathcal{X} .*

The proof depends on the following

LEMMA 1. *Let A and N be bounded operators on \mathcal{X} such that*

$A^*A + N^*N \leq I$, and D be the positive square root of $I - A^*A - N^*N$. Let \mathcal{D} be the closure of the range of D . Then

- (i) $A^{*n}A^n$ strongly converges to the positive operator A_∞^2 , as $n \rightarrow \infty$.
- (ii) For each $x \in \mathcal{X}$ the sequence $\|N\Phi_n x\|_{\mathcal{F}_n}$ is decreasing.
- (iii)

$$\sum_{n=1}^{\infty} \|D\Phi_n x\|^2 \leq \|x\|^2 \quad (\text{for all } x \in \mathcal{X})$$

so that the operator $D\Phi$ mapping \mathcal{X} into $\mathcal{F}_1(\mathcal{D})$ defined by $D\Phi x = \bigoplus_1^\infty D\Phi_n x$, is well defined. In fact

$$\|D\Phi x\|_{\mathcal{F}_1}^2 + \|A_\infty x\|_{\mathcal{X}}^2 + \sum_1^\infty \|A_\infty N\Phi_n x\|_{\mathcal{F}_n}^2 + \lim_{n \rightarrow \infty} \|N\Phi_n x\|_{\mathcal{F}_n}^2 = \|x\|^2 \quad (4.1)$$

for all $x \in \mathcal{X}$.

Proof. Part (i) follows because A is a contraction, i.e., $A^{*n}A^n$ is a sequence of decreasing positive operators.

Consulting (2.11) gives

$$\begin{aligned} \|D\Phi_1 x\|_{\mathcal{F}_1}^2 &= \lim_{k \rightarrow \infty} \sum_{i=0}^k \|DA^i x\|_{\mathcal{X}}^2 \\ &= \lim_{k \rightarrow \infty} \sum_{i=0}^k (A^{*i}(I - A^*A - N^*N)A^i x, x) \quad (4.2) \\ &= \|x\|^2 - \lim_{k \rightarrow \infty} \|A^k x\|^2 - \|N\Phi_1 x\|_{\mathcal{F}_1}^2. \end{aligned}$$

Therefore,

$$\|D\Phi_1 x\|^2 = \|x\|^2 - \|A_\infty x\|^2 - \|N\Phi_1 x\|^2. \quad (4.3)$$

Following the same procedure on the general term $n > 1$ gives

$$\|D\Phi_n x\|^2 = \|N\Phi_{n-1} x\|^2 - \|A_\infty N\Phi_{n-1} x\|^2 - \|N\Phi_n x\|^2 \quad (4.4)$$

summing to n on (4.3), (4.4) and rearranging terms:

$$\begin{aligned} \sum_{i=1}^n \|D\Phi_i x\|^2 + \|A_\infty x\|^2 + \sum_{i=1}^{n-1} \|A_\infty N\Phi_i x\|^2 \\ = \|x\| - \|N\Phi_n x\|^2. \end{aligned} \quad (4.5)$$

Since the left-hand side is positive and increasing in n , the $\|N\Phi_n x\|^2$ are decreasing. Part (iii) follows by taking limits in (4.5).

Proof of Proposition 2. Let $D\Phi$ be the operator given in the lemma. Following (3.1), (3.2), it is easy to verify that

$$S_{\mathcal{D}}D\Phi = D\Phi A \quad \text{and} \quad E_{\mathcal{D}}D\Phi = D\Phi N. \quad (4.6)$$

The hypothesis of the Proposition and (4.1) guarantees that $D\Phi$ is an isometry. Since the $\text{ran } D\Phi$ is invariant under $S_{\mathcal{D}}$ and $E_{\mathcal{D}}$, the proof is complete.

COROLLARY 4. *If A, N are bounded operators on \mathcal{X} and $A^*A + N^*N \leq rI$ where $r < 1$ then A and N are simultaneously unitarily equivalent to part of S and E on $\mathcal{F}_1(\mathcal{X})$.*

Proof. We verify that the hypothesis of the proposition are satisfied. Clearly $A^n \rightarrow 0$. Equation (3.6) and Corollary 2 guarantees that $\|N\Phi_n x\|_{\mathcal{X}_n}^2 \rightarrow 0$ as $n \rightarrow \infty$ for all $x \in \mathcal{X}$. Since $r < 1$ we have $\mathcal{D} = \mathcal{X}$ and the proof is complete.

By employing a trick found in [1, 5] the hypothesis $A^n \rightarrow 0$ and $N\Phi_n \rightarrow 0$ strongly as $n \rightarrow \infty$ in Proposition 2 are removed. This begins with

PROPOSITION 3. *Let A, N be operators on \mathcal{X} , and $A^*A + N^*N \leq I$. If $\|N\Phi_n x\| \rightarrow 0$ as $n \rightarrow \infty$ for all $x \in \mathcal{X}$, then A, N are simultaneously unitarily equivalent to part of a pair of co-isometries.*

Proof. Throughout the notation of Lemma 1 is used. Let \mathcal{A} be the closure of the $\text{ran } A_\infty$ and W the operator mapping \mathcal{A} into \mathcal{A} defined by $WA_\infty x \doteq A_\infty Ax$. It is easy to show that W is an isometry (see p. 51 of [1] or p. 39 of [5]). By Proposition (2.3), p. 6 of [5], W can be extended to a unitary operator S_0 on some larger Hilbert space \mathcal{Y} , i.e., \mathcal{A} is a subspace of \mathcal{Y} and $W = S_0|_{\mathcal{A}}$. Further, $S_0 A_\infty = A_\infty A$.

Let $\mathcal{F}_0(\mathcal{Y})$ be the following Fock space

$$\mathcal{F}_0(\mathcal{Y}) \doteq \bigoplus_{n=0}^{\infty} \mathcal{H}_n(\mathcal{Y}). \quad (4.7)$$

(Recall $\mathcal{H}_0(\mathcal{Y}) \doteq \mathcal{Y}$.) Define the co-isometry $S_{\mathcal{Y}}$ on $\mathcal{F}_0(\mathcal{Y})$ by

$$S_{\mathcal{Y}} \bigoplus_0^{\infty} f_n \doteq \bigoplus_0^{\infty} S_n f_n = \{S_0 f_0, S_1 f_1, S_2 f_2, \dots\}, \quad (4.8)$$

where S_0 is the above unitary operator and S_n is the usual backward shift operator on $\mathcal{H}_n(\mathcal{Y})$ for $n \geq 1$, (see (2.4)). The co-isometry $E_{\mathcal{Y}}$ on $\mathcal{F}_0(\mathcal{Y})$ is defined by

$$E_{\mathcal{Y}} \bigoplus_0^{\infty} f_n \doteq \bigoplus_{n=0}^{\infty} E_{n+1} f_{n+1} = \{E_1 f_1, E_2 f_2, E_3 f_3, \dots\}, \quad (4.9)$$

where E_n for $n \geq 1$ is the evaluation operator mapping $\mathcal{H}_n(\mathcal{Y})$ into $\mathcal{H}_{n-1}(\mathcal{Y})$, (see (2.7)).

Consider the operator Φ_0 mapping \mathcal{X} into $\mathcal{F}_0(\mathcal{Y})$ defined by

$$\begin{aligned}\Phi_\infty x &\doteq A_\infty x \oplus A_\infty N \Phi x \\ &= \{A_\infty x, A_\infty N \Phi_1 x, A_\infty N \Phi_2 x, \dots\} \quad (x \in \mathcal{X}).\end{aligned}\tag{4.10}$$

By following the calculations in (3.1), (3.2) with the definition of S_0 it is easy to verify that

$$S_{\mathcal{Y}} \Phi_\infty = \Phi_\infty A \quad \text{and} \quad E_{\mathcal{Y}} \Phi_\infty = \Phi_\infty N.\tag{4.11}$$

To complete the proof we combine the above with the proof of Proposition 2. Consider the operator $D\Phi \oplus \Phi_\infty$ mapping \mathcal{X} into $\mathcal{F}_1(\mathcal{D}) \oplus \mathcal{F}_0(\mathcal{Y})$ defined by $D\Phi x \oplus \Phi_\infty x$ when $x \in \mathcal{X}$. This operator is an isometry, by (4.1) and (4.10). Clearly the operators $S_{\mathcal{D}} \oplus S_{\mathcal{Y}}$ and $E_{\mathcal{D}} \oplus E_{\mathcal{Y}}$ on $\mathcal{F}_1(\mathcal{D}) \oplus \mathcal{F}_0(\mathcal{Y})$ are co-isometries. Further (4.6), (4.11) give $(S_{\mathcal{D}} \oplus S_{\mathcal{Y}})(D\Phi x \oplus \Phi_\infty x) = D\Phi Ax \oplus \Phi_\infty Ax$ and $(E_{\mathcal{D}} \oplus E_{\mathcal{Y}})(D\Phi x \oplus \Phi_\infty x) = D\Phi Nx \oplus \Phi_\infty Nx$, where $x \in \mathcal{X}$. Since the range of $D\Phi \oplus \Phi_\infty$ is an invariant subspace for both $S_{\mathcal{D}} \oplus S_{\mathcal{Y}}$ and $E_{\mathcal{D}} \oplus E_{\mathcal{Y}}$ the proof is complete.

Finally we are ready to prove

PROPOSITION 4. *If A, N are operators on \mathcal{X} such that $A^*A + N^*N \leq I$, then A, N are simultaneously unitarily equivalent to part of a pair of co-isometries.*

Proof. Since $\|N\Phi_n x\|^2$ is a decreasing sequence (see Lemma 1), there exists a positive operator P on \mathcal{X} such that

$$(P^2 x, x) = \|Px\|^2 = \lim_{n \rightarrow \infty} (\Phi_n^* N^* N \Phi_n x, x) = \lim_{n \rightarrow \infty} \|N\Phi_n x\|_{\mathcal{F}_n}^2$$

By the definitions (2.2) and (2.11) we have

$$\|Px\|^2 = \lim_{n \rightarrow \infty} \sum_{i_1 > 0, \dots, i_n > 0} \|NA^{i_n} NA^{i_{n-1}} N \dots NA^{i_1} x\|_{\mathcal{F}_n}^2\tag{4.12}$$

and

$$\|Px\|_{\mathcal{Z}}^2 = \|PN\Phi_1 x\|_{\mathcal{F}_1}^2\tag{4.13}$$

for all $x \in \mathcal{X}$. Let \mathcal{Z} be the closure of the range of P , and let Z be any unitary operator on the Fock space $\mathcal{F}_0(\mathcal{Z})$ such that

$$Z\{Px, 0, 0, 0, \dots\} = \{0, PN\Phi_1 x, 0, 0, \dots\}.\tag{4.14}$$

Such an operator exists by (4.13).

Let S_x be the co-isometry on $\mathcal{F}_0(\mathcal{L})$ defined by

$$S_x \bigoplus_0^\infty f_n \doteq \bigoplus_0^\infty S_n f_n \quad \left(\bigoplus_0^\infty f_n \in \mathcal{F}_0(\mathcal{L}) \right), \quad (4.15)$$

where S_n is the backward shift operator in $\mathcal{H}_n(\mathcal{L})$, for $n > 0$, and $S_0 \doteq I$. Define \tilde{E}_x to be the co-isometry on $\mathcal{F}_0(\mathcal{Z})$ given by $\tilde{E}_x \doteq ZE_x$, where E_x is the generalized evaluation operator, (replace \mathcal{Y} by \mathcal{L} in (4.9)). Consider the operator Φ_x mapping \mathcal{X} into $\mathcal{F}_0(\mathcal{L})$ defined by

$$\Phi_x x \doteq \{0, PN\Phi_1 x, 0, 0, \dots\} \quad (x \in \mathcal{X}). \quad (4.16)$$

From (4.13)

$$\|\Phi_x x\|^2 = \lim_{n \rightarrow \infty} \|N\Phi_n x\|^2 \quad (x \in \mathcal{X}). \quad (4.17)$$

Using (4.14), (2.4), (2.7) a simple calculation verifies that

$$S_x \Phi_x = \Phi_x A \quad \text{and} \quad \tilde{E}_x \Phi_x = \Phi_x N. \quad (4.18)$$

At this point the proof is exactly the same as Proposition 3 except one uses the co-isometries $S_{\mathcal{D}} \oplus S_{\mathcal{Y}} \oplus S_x$ and $E_{\mathcal{D}} \oplus E_{\mathcal{Y}} \oplus \tilde{E}_x$ on $\mathcal{F}_1(\mathcal{D}) \oplus \mathcal{F}_0(\mathcal{Y}) \oplus \mathcal{F}_0(\mathcal{L})$, along with the isometry mapping \mathcal{X} into $\mathcal{F}_1(\mathcal{D}) \oplus \mathcal{F}_0(\mathcal{Y}) \oplus \mathcal{F}_0(\mathcal{L})$ defined by $D\Phi x \oplus \Phi_\infty x \oplus \Phi_x x$ where $x \in \mathcal{X}$. Note (4.1), (4.10) and (4.17) guarantee that the last operator preserves the norm. The intertwining property follows from (4.6), (4.11), and (4.18).

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