



## Enhanced collectivity in $^{74}\text{Ni}$

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### ABSTRACT

The neutron-rich nucleus  $^{74}\text{Ni}$  was studied with inverse-kinematics inelastic proton scattering using a  $^{74}\text{Ni}$  radioactive beam incident on a liquid hydrogen target at a center-of-mass energy of 80 MeV. From the measured de-excitation  $\gamma$  rays, the population of the first  $2^+$  state was quantified. The angle-integrated excitation cross section was determined to be 14(4) mb. A deformation length of  $\delta = 1.04(16)$  fm was extracted in comparison with distorted wave theory, which suggests that the enhancement of collectivity established for  $^{70}\text{Ni}$  continues up to  $^{74}\text{Ni}$ . A comparison with results of shell model and quasi-particle random phase approximation calculations indicates that the magic character of  $Z = 28$  or  $N = 50$  is weakened in  $^{74}\text{Ni}$ .

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### 1. Introduction

The evolution of the nuclear shell structure is one of the most fundamental topics in the investigation of unstable nuclei. Recent studies have shown that the shell structure established for stable nuclei can change in nuclei far from the valley of  $\beta$ -stability [1]. Some of the traditional magic numbers lose their magicity in neutron-rich regions, while some other proton or neutron num-

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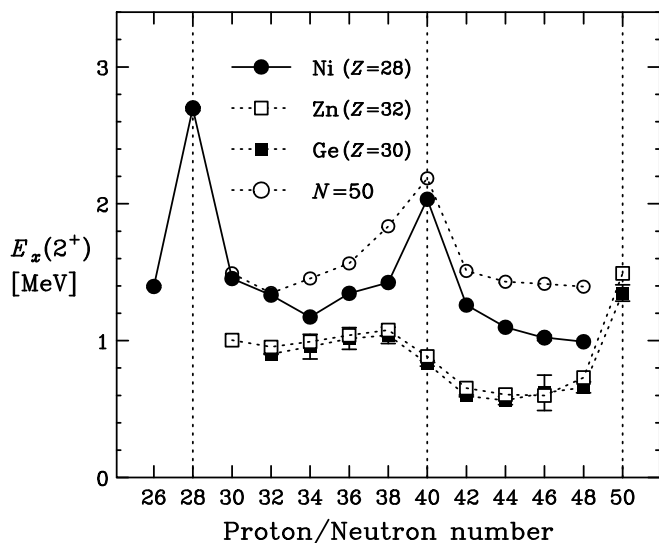
neutron numbers  $N = 8$  and  $N = 20$ , respectively, which are well-established magic numbers in stable nuclei [2–8]. It is interesting to extend the study to heavier regions of the nuclear chart to investigate whether such structural changes occur there as well. One particular interest is the behavior of the magicity at proton number  $Z = 28$  and neutron number  $N = 50$  in the neutron-rich nucleus  $^{78}\text{Ni}$  and neighbouring isotopes. In order to address this question, we have studied the collectivity of  $^{74}\text{Ni}$  with  $Z = 28$  and  $N = 46$  by measuring the proton inelastic scattering to the first excited  $2^+$  ( $2_1^+$ ) state.

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The Ni isotopes have reduced collectivity reflecting the magic character of the closed proton shell at  $Z = 28$ . The excitation energies of the  $2_1^+$  states ( $E_x(2_1^+)$ ) along the Ni isotopic chain are higher than those of the neighbouring even- $Z$  isotopes such as Zn ( $Z = 30$ ) or Ge ( $Z = 32$ ) (Fig. 1). However, for the Ni isotopes between neutron numbers  $N = 40$  and  $N = 50$ , where the valence neutrons occupy the  $g_{9/2}$  orbital, the  $E_x(2^+)$  value decreases continuously. This feature does not match with the expectations that single-closed-shell can be described by a constant  $E_x(2^+)$  values along the isotopic chain. Indeed, for the  $N = 50$  isotones with valence protons in the  $g_{9/2}$  orbital, the  $E_x(2^+)$  values are almost constant between  $^{92}\text{Mo}$  ( $N = 50$ ,  $Z = 42$ ) and  $^{98}\text{Cd}$  ( $N = 50$ ,  $Z = 48$ ) as shown by open circles in Fig. 1. The monotonic decrease of  $E_x(2^+)$  in the Ni isotopes up to  $^{76}\text{Ni}$  may be an indication of the magic character of  $Z = 28$  or  $N = 50$  being weakened.

The degree of collectivity can be characterized more directly by the quadrupole deformation parameter  $\beta_2$  or the deformation length  $\delta$  which is the product of  $\beta_2$  and the nuclear radius  $R$ . These quantities can be derived from the transition strength between the ground  $0^+$  state and  $2_1^+$  state using various probes, such as electromagnetic transition or proton inelastic scattering. When extracting  $\beta_2$  and  $\delta$  from the cross sections of the proton inelastic



**Fig. 1.** Experimental  $E_x(2^+)$  values for Ni isotopes ( $Z=28$ ) and the neighbouring even- $Z$  isotopes, Zn ( $Z=30$ ) and Ge ( $Z=32$ ), as a function of the neutron number. The  $E_x(2^+)$  values for the  $N=50$  isotones are also shown as a function of proton number. The values are taken from Refs. [16–21].

scattering, the deformation length  $\delta$  is less model dependent [4] and subsequently used as a measure of collectivity in this work. The  $\delta$  value derived from the reduced electric quadrupole transition probability  $B(E2)$ , which is known up to  $^{70}\text{Ni}$ , is very small in  $^{68}\text{Ni}$  [11] and increases abruptly at  $^{70}\text{Ni}$  [12]. In Ref. [12] this is interpreted as an indication that the shell gap at  $Z=28$  in  $^{70}\text{Ni}$  is reduced by the occupation of the neutron  $g_{9/2}$  orbital. Beyond  $^{70}\text{Ni}$ , where further enhancement of the collectivity is suggested by the decreasing  $E_x(2_1^+)$  value,  $\delta$  is not known due to the difficulty of producing such neutron-rich nuclei at sufficient intensity for spectroscopy. In this work, we have succeeded to obtain the deformation length  $\delta$  for  $^{74}\text{Ni}$  from the angle-integrated cross section for the excitation of the first  $2^+$  state in the proton inelastic scattering.

The difficulty caused by the low  $^{74}\text{Ni}$  beam intensity has been overcome by incorporating an efficient combination of proton inelastic scattering and  $\gamma$ -ray spectroscopy. The excited states of  $^{74}\text{Ni}$  were identified by measuring the de-excitation  $\gamma$  rays [4] instead of the recoil protons [13,14]. Consequently, a very thick proton target could be used without loss of energy resolution. Furthermore, by using a liquid hydrogen target, very highly sensitive measurements can be realized as demonstrated with the case of  $^{30}\text{Ne}$  [15], because hydrogen has the largest number density of atoms within the target material with a given energy-loss-equivalent thickness. High efficiency is also obtained thanks to the small scattering angle in the laboratory reference frame due to the use of inverse kinematics with the lightest-mass target, where 100% acceptance for the scattered projectiles can easily be achieved in an experimental setup. Another advantage of using a liquid hydrogen target is the significant reduction of  $\gamma$ -ray background due to the absence of  $\gamma$  rays originating from the excitation of the target nuclei. This high efficiency and low background conditions enabled us to observe the  $2_1^+$  state in the very neutron-rich nucleus  $^{74}\text{Ni}$  in spite of the very low beam intensity of less than 1 count per second (cps).

## 2. Experiment

The experiment was performed at the National Superconducting Cyclotron Laboratory (NSCL) at the Michigan State University using the large-acceptance A1900 fragment separator [22] and the

high-resolution S800 spectrograph [23]. A  $^{86}\text{Kr}$  primary beam accelerated to 140 MeV/nucleon by the Coupled Cyclotron Facility was directed onto a 399-mg/cm<sup>2</sup>-thick  $^9\text{Be}$  fragmentation target at a typical intensity of 20 particle nA. The produced fragments were selected and separated by the A1900. The momentum acceptance was set to 2%. At the second intermediate image of the A1900, an achromatic wedge degrader of 45-mg/cm<sup>2</sup>-thick aluminum was placed to purify the secondary beam. Within a secondary cocktail beam,  $^{74}\text{Ni}$  was obtained with an average intensity of 0.5 cps while the total intensity of all transmitted fragments amounted to 80 cps.

The secondary beam was delivered to the analysis line of the S800 spectrograph operated in focus mode; the beam was achromatically focused on the target. The particle identification was performed on an event-by-event basis by measuring the energy loss ( $\Delta E$ ), magnetic rigidity ( $B\rho$ ), and time-of-flight (TOF) of each nucleus. A Si-PIN detector with a thickness of 74  $\mu\text{m}$  was used to measure  $\Delta E$  of the incoming projectile beam at the entrance of the S800 analysis line. The  $B\rho$  value was obtained from the angle at the intermediate focal plane of the S800 analysis line measured by a pair of parallel-plate avalanche counters. The TOF was obtained from the time difference between the signals of plastic scintillators at the exit of the A1900 and the final focal plane of the S800 spectrograph.

The secondary beam bombarded the liquid hydrogen target of thickness 210 mg/cm<sup>2</sup> to induce proton inelastic scattering. The liquid hydrogen was contained in a cylindrical cell, a part of “CRYogenic ProTON and Alpha target system” (CRYPTA) [24]. The target cell had aluminum entrance and exit windows with thicknesses of 0.22 mm and diameters of 30 mm. The contribution of the target windows to the inelastic scattering was estimated to be 3% of that by hydrogen. The average mid-target energy of  $^{74}\text{Ni}$  was 81 MeV/nucleon, which corresponds to center-of-mass energy of 80 MeV.

The reaction residues were analyzed with the S800 spectrograph, whose magnetic field was set to transmit the (in)elastically scattered  $^{74}\text{Ni}$  ions. The particle identification was performed by combining  $B\rho$ ,  $\Delta E$ , and TOF measured by the focal-plane detectors of the S800 [25]. The  $B\rho$  value was obtained from the position of the ion at the final focal plane measured by two cathode read-out drift chambers (CRDCs). An ionization chamber located downstream of the CRDCs provided the  $\Delta E$  information. For the TOF, the flight time differences between the target position and the final focal plane were used, the former being extracted from the time measured by the plastic scintillator at the exit of the A1900 and  $B\rho$  of the incident beam. The mass distribution for the Ni isotopes was obtained as shown in Fig. 2 with a resolution of 0.6% in full width at half maximum (FWHM). With this resolution, contaminations in the  $^{74}\text{Ni}$  mass gate set in the present analysis (indicated by the dotted lines in Fig. 2) are negligible.

The  $\gamma$  rays emitted from the excited nuclei were detected by a barrel array of NaI(Tl) scintillators with an inner diameter of 43 cm [26,27] surrounding the liquid hydrogen target. The array consisted of 21 NaI(Tl) scintillator bars, each of which was 55.0 cm long and had a trapezoidal cross section with a height of 6.0 cm and upper and lower bases of 5.5 cm and 7.0 cm, respectively. The barrel covered a polar-angle range from 38 to 122 degrees with respect to the beam axis. This asymmetric geometry was adopted to enhance the detection efficiency of  $\gamma$  rays emitted in flight with forward-shifted angular distributions due to the Lorentz boost. The signals were read out by two photomultiplier tubes optically coupled to both ends of each scintillator bar. The energy of the  $\gamma$  ray in the laboratory frame was obtained from the geometrical average of the two signals from each pair of photomultiplier tubes. The detection position was obtained from the logarithm of the ratio of the two signals. Using the position information, a Doppler-shift correction

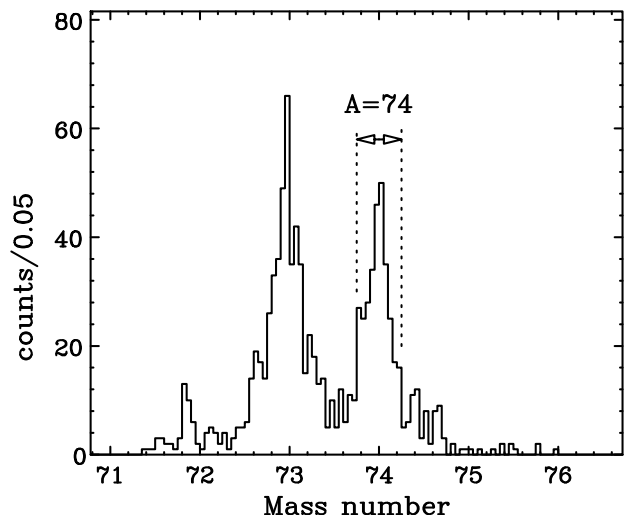


Fig. 2. Mass distribution for Ni isotopes behind the secondary target as measured with the S800 spectrograph. The dotted lines indicate the region of the mass gate used in the present analysis.

was applied to the energy of the  $\gamma$  rays emitted from the moving nucleus at a velocity of  $v/c \sim 0.4$ . Energy calibrations were performed with  $^{22}\text{Na}$  and  $^{137}\text{Cs}$  standard sources. The intrinsic energy resolution was 13% (FWHM) for the 1275-keV  $\gamma$  ray. The position was calibrated by using a collimated  $^{60}\text{Co}$  source and a position resolution of 3 cm (FWHM) was obtained.

In order to extract the energy and intensity of the  $\gamma$ -ray transitions, the spectral shape expected for mono-energetic  $\gamma$  rays from a moving source was computed with a Monte Carlo simulation using the GEANT code [28]. The simulation took into account the geometries of the NaI(Tl) array, CRYPTA, and the beam pipe, as well as the energy and position resolution and the energy loss of the beam in the target. The overall energy resolution for the 1024 keV  $\gamma$  rays emitted from  $^{74}\text{Ni}$  moving with  $v/c \sim 0.4$  was 18% (FWHM). The photo-peak efficiency was calculated to be 16(2)%, where the error is a systematic one estimated by comparing the measured  $\gamma$ -ray spectra of the standard sources with the simulated one.

### 3. Result

Fig. 3 shows the Doppler-shift corrected energy spectrum of the prompt  $\gamma$  rays detected in coincidence with the inelastically-scattered  $^{74}\text{Ni}$  ions. A prominent peak is seen at 1020(11) keV, which is ascribed to the  $2_1^+ \rightarrow 0_{\text{g.s.}}^+$  transition from the selectivity of proton inelastic scattering. This energy is consistent with the value (1024(1) keV) reported in the  $\beta$ -decay study of  $^{74}\text{Co}$  [20]. The peak intensity was extracted by fitting the spectrum. The best fit result is shown by the solid curve, which is composed of the simulated detector response for a 1020-keV  $\gamma$  ray (dashed curve) and an exponential background (dotted curve). In addition to the 1020-keV peak, a peak-like structure is visible at 786(30) keV. The statistical significance of this peak was evaluated to be one standard deviation based on a fit that incorporated the 786-keV  $\gamma$ -ray peak.

The angle-integrated cross section for the  $2_1^+$  excitation was determined from the yield of the  $\gamma$  rays after considering possible feeding from higher excited states. In the observed spectrum, the 786-keV peak is the only candidate for a feeding transition. Its contribution amounts to 20(20)% of the yield of the 1020-keV peak. Since the statistical significance is low and the coincidence relation with the  $2_1^+ \rightarrow 0_{\text{g.s.}}^+$  transition is uncertain, the cross sections

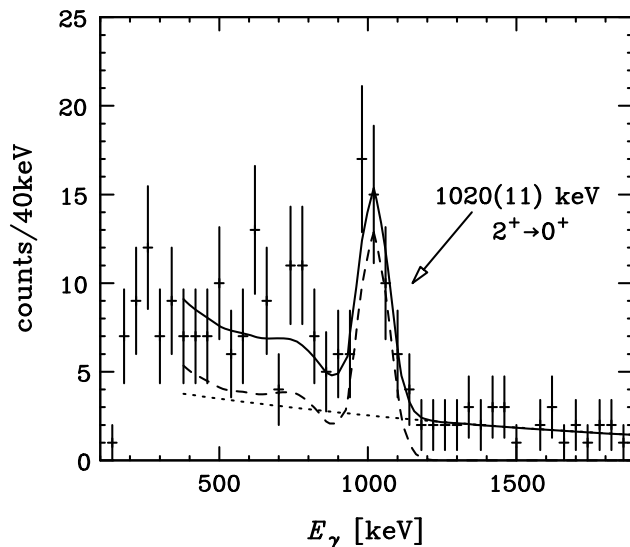


Fig. 3. Doppler-shift corrected  $\gamma$ -ray spectrum for the proton inelastic scattering of  $^{74}\text{Ni}$ . The solid curve shows the best fit result, which consists of the simulated detector response for the 1020-keV  $\gamma$  ray (dashed curve) and an exponential background (dotted curve).

were derived with and without assuming the 786-keV peak as a feeding transition. Their average was adopted while the difference was taken into account as a systematic error.

In addition, we considered the possible feeding from a  $3^-$  state, because proton inelastic scattering generally has a large cross section for the excitation to low-lying, collective  $3^-$  states. In the case of the stable Ni isotopes, the feeding transitions from the  $3^-$  state to the  $2_1^+$  state have intensities of about 20 ~ 30% of the direct excitation to the  $2_1^+$  state [29,30], and have energies around 2 ~ 3 MeV. Such a transition cannot be observed in the present measurement because of limited sensitivity. We therefore assumed the feeding contribution from the  $3^-$  state to be 25(5)%, although the corresponding photo-peak was not observed. The  $0_{\text{g.s.}}^+ \rightarrow 2_1^+$  excitation cross section was then determined to be 14(4) mb, where the quoted error includes both the statistical (20%) and systematic uncertainties. The systematic error is attributed to the uncertainties in the efficiency of the  $\gamma$ -ray detection (10%), the target thickness (3%), and uncertainty of the feeding corrections (20%).

The deformation length  $\delta^{p,p'}$  for the  $0_{\text{g.s.}}^+ \rightarrow 2_1^+$  transition was extracted from calculations based on the distorted wave theory using the ECIS97 code [31]. We adopted two different approaches. The first one employed the global optical potential set KD02 [32]. By adopting the collective vibrational model, the transition potential was obtained from the derivative of the optical potential with an amplitude of  $\delta^{p,p'}$ . The  $\delta^{p,p'}$  value was determined so that the calculated total inelastic scattering cross section reproduces the measured cross section. The  $\delta^{p,p'}$  value obtained in this way was 1.07(16) fm, where the error includes the uncertainties of the experimental cross section. The corresponding deformation parameter  $\beta_2 (= \delta^{p,p'}/R)$  is 0.21(3) using a radius of  $R = 5.04$  fm which is taken from  $R = r_0 A^{1/3}$  with  $r_0 = 1.2$  fm.

In the second approach the optical potential and transition potential were derived by folding the effective nucleon–nucleon interaction with the nucleon densities and transition densities, respectively, using the MINC code [33]. For the effective nucleon–nucleon interaction, the density-dependent effective interaction proposed by Jeukenne, Lejeune, and Mahaux (JLM) [34] was used with normalization factors of 0.95 for the real and imaginary parts [35,36]. The density distributions of protons and neutrons

are of Woods–Saxon form with a common radius parameter of 4.51 fm, which was chosen to reproduce the root mean square (rms) radius of the matter distribution (4.03 fm) obtained from Hartree–Fock–Bogoliubov (HFB) calculations using the SkM\* parameterization [37]. The transition densities were obtained from the derivatives of the proton and neutron densities. The  $\delta^{p,p'}$  value obtained with this approach is 1.02(15) fm. The  $\delta^{p,p'}$  values extracted using the global optical potential and folding model potential agree with each other. We take the average of 1.04(16) fm as the adopted value for the subsequent discussions.

Due to the large neutron excess,  $^{74}\text{Ni}$  may have a neutron-skin structure where the neutron distribution has a larger radius than the proton distribution. Indeed, the above HFB calculation gives different rms radii for neutrons (4.13 fm) and protons (3.87 fm). In order to evaluate the effect of the possible neutron-skin structure, a microscopic calculation using Woods–Saxon distributions with different radius parameters for neutrons (4.66 fm) and protons (4.27 fm) was performed. These values were chosen to reproduce the rms radii obtained from the HFB calculation. The result is 1.03(16) fm for  $\delta^{p,p'}$ , which agrees with the one obtained without incorporating the neutron-skin effect. This shows that the possible neutron-skin structure does not have a significant effect on the present analysis.

#### 4. Discussion

The extracted  $\delta^{p,p'}$  value is shown in Fig. 4(a) together with those of stable Ni isotopes. The open circles indicate deformation lengths  $\delta^C$  obtained from  $B(E2)$  values using the following equation,

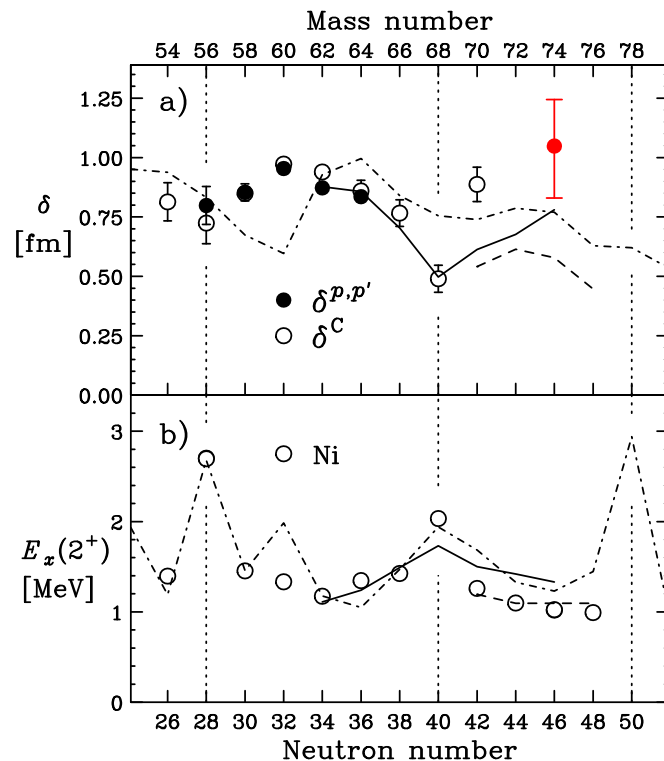
$$\delta^C = \beta_2^C R = \frac{4\pi}{3ZeR} \sqrt{B(E2; 0^+ \rightarrow 2^+)}$$

where  $R = r_0 A^{1/3}$  with  $r_0 = 1.2$  fm. The  $\delta^{p,p'}$  value is sensitive to the quadrupole collectivity of both protons and neutrons, while  $\delta^C$  is sensitive only to that of the protons. The two quantities  $\delta^{p,p'}$  and  $\delta^C$  are identical if the proton and neutron motions are the same. Indeed, the experimental values of  $\delta^{p,p'}$  and  $\delta^C$  are very close in the region where both values are available. In the following, we assume that the  $\delta^{p,p'}$  and  $\delta^C$  are close to each other also in the neutron-rich region, and both  $\delta^{p,p'}$  and  $\delta^C$  are referred to as  $\delta$ .

The  $\delta$  value for  $^{74}\text{Ni}$  obtained in the present study is about twice as large as that for  $^{68}\text{Ni}$ , and is comparable to or may even exceed that for  $^{70}\text{Ni}$ . This indicates that the enhancement of collectivity also persists in  $^{74}\text{Ni}$ , which is consistent with the trend of the  $E_x(2^+)$  values discussed in Section 1.

The enhancement of quadrupole collectivity is inferred also for  $^{72}\text{Ni}$ , which could be estimated from the  $B(E2)$  value of  $^{73}\text{Cu}$  [38]. It is suggested that the  $3/2^-$  ground states and the low-lying  $7/2^-$  states in the Cu isotopes in this region can be described in terms of a particle-core model. The evolution of the  $B(E2)$  values between those states is very similar to the behavior of the  $B(E2)$  values in the adjacent Ni isotopes [38]. The  $B(E2)$  value for  $^{72}\text{Ni}$  can then be estimated from that of  $^{73}\text{Cu}$ , yielding a  $\delta$  value of 1.09 fm, which is almost the same as that of  $^{74}\text{Ni}$ . This is consistent with the above interpretation of the enhanced collectivity in the Ni isotopes in this region.

Fig. 4 compares the experimental  $\delta$  and  $E_x(2^+)$  values with theoretical ones. The dashed curve shows a shell-model calculation with the neutron configuration space of the  $p_{3/2}$ ,  $p_{1/2}$ ,  $f_{5/2}$ , and  $g_{9/2}$  orbitals outside of a  $^{56}\text{Ni}$  inert core [39]. In order to take into account the character of the  $Z = 28$  proton core, which is rather soft [40], a large neutron effective charge of  $e_n = 1.0e$  was used. However, the calculation cannot reproduce the large  $\delta$  values of  $^{70}\text{Ni}$  and  $^{74}\text{Ni}$ . This means that the model space in this calculation



**Fig. 4.** (a) The  $\delta$  values for Ni isotopes. The open circles show  $\delta^C$  obtained from  $B(E2)$  values [11,12,16,18]. The filled circles show  $\delta^{p,p'}$  obtained from the proton inelastic scattering [13,47]. (b) The  $E_x(2^+)$  systematics for Ni isotopes. In both figures the dashed and solid curves show the  $\delta$  and  $E_x(2^+)$  values from the shell-model calculations of Refs. [39] and [11], respectively. The dot-dashed curve shows the results of HFB + QRPA calculations [41].

is not large enough to account for the collectivity observed in the Ni isotopes.

Results of a shell-model calculation with a larger model space reported in Ref. [11] are shown by the solid curve. The model space comprises the full  $fp$ -shell for protons and the  $p_{3/2}$ ,  $p_{1/2}$ ,  $f_{5/2}$ , and  $g_{9/2}$  orbitals for neutrons. For effective charges, standard values of  $1.5e$  and  $0.5e$  were used for protons and for neutrons, respectively. The calculated  $\delta$  values are larger than those of Ref. [39] and the agreement with the experimental value has improved. However, the calculation still shows smaller  $\delta$  values than the experimental values for  $^{70}\text{Ni}$  and  $^{74}\text{Ni}$ . This underestimation of collectivity is consistent with the overestimation of the  $E_x(2^+)$  values. This is in contrast to the results in the lighter region below neutron number  $N = 40$ , where a good agreement is obtained for both  $\delta$  and  $E_x(2^+)$ .

Results of a calculation based on the HFB model and quasi-particle random phase approximation (QRPA) using the SkM\* effective interaction [41] are shown by the dot-dashed curve. This calculation also underestimates the  $\delta$  values and overestimates the  $E_x(2^+)$  values in  $^{70}\text{Ni}$  and  $^{74}\text{Ni}$ . The fact that these calculations fail to reproduce the large  $\delta$  values for the Ni isotopes in the  $N > 40$  region infers a modification of structure due to a mechanism characteristic for neutron-rich nuclei.

One explanation for this phenomenon might lie in the softening of the  $Z = 28$  core. The shell gap at  $Z = 28$  is caused by the spin-orbit interaction and can easily be influenced by subtle changes of the potential due to unbalanced neutron and proton numbers. Indeed, the shell gap at  $Z = 28$  is predicted to be narrowed by the tensor interaction between protons and neutrons [42–45,12]. Since the tensor interaction is strongly attractive between  $\pi f_{5/2}$  and  $\nu g_{9/2}$  and strongly repulsive between  $\pi f_{7/2}$  and  $\nu g_{9/2}$ , the

energy gap between  $\pi f_{5/2}$  and  $\pi f_{7/2}$  is effectively reduced with increasing numbers of the valence neutrons in the  $g_{9/2}$  orbital. As a result, the probability of proton excitations from the  $f_{7/2}$  orbital into the upper orbitals increases, leading to the weakening of the shell closure.

The enhancement of collectivity could be due to neutrons, namely, to the enhanced probability of neutron excitations from the  $g_{9/2}$  orbital into the  $d_{5/2}$  orbital across the  $N = 50$  shell gap, which could occur if the gap is narrowed. In the neighbouring isotopes,  $N = 50$  is known to be a good magic number down to Zn isotopes [21,46]. Therefore, this interpretation points to a weakening of the  $N = 50$  neutron magicity characteristic in the neutron-rich Ni isotopes. In both scenarios, the collectivity could be more enhanced in the more neutron-rich region and thus the doubly magic nature of  $^{78}\text{Ni}$  may be weakened. To further clarify the origin and mechanism of the enhanced collectivity observed in the present study, an investigation of nuclei closer to  $^{78}\text{Ni}$  is necessary.

## 5. Summary

In summary, the neutron-rich nucleus  $^{74}\text{Ni}$  was studied by proton inelastic scattering in inverse kinematics at a center-of-mass energy of 80 MeV. The hitherto impossible measurement of this extremely exotic nucleus was realized by combining the high beam intensity available at NSCL with the efficient use of a thick liquid hydrogen target (CRYPTA) in conjunction with  $\gamma$ -ray spectroscopy to tag the inelastic excitation. The observation of the  $\gamma$ -ray peak at 1020(11) keV confirms the energy of the  $2^+ \rightarrow 0^+$  transition reported in an earlier  $\beta$ -decay study. The angle-integrated cross section for the excitation of the  $2^+_1$  state was determined to be  $\sigma = 14(4)$  mb. From the comparison with the distorted wave theory, the  $\delta$  value of 1.04(16) fm was extracted. This large  $\delta$  value suggests that the enhancement of collectivity reported for  $^{70}\text{Ni}$  continues up to  $^{74}\text{Ni}$ . The obtained  $\delta$  value is larger than all values obtained by shell-model and HFB + QRPA calculations, although both approaches reproduce  $\delta$  in the lighter Ni isotopes with  $N < 40$ . This infers a modification of the nuclear structure due to a mechanism characteristic for neutron-rich nuclei. The enhanced collectivity of  $^{74}\text{Ni}$  located near the doubly-magic nucleus  $^{78}\text{Ni}$  may be an indication of the weakening of the magicity at  $Z = 28$  or  $N = 50$ . In order to clarify the mechanism driving the enhancement in collectivity, a study of the structure of nuclei closer to  $^{78}\text{Ni}$  is needed.

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