# Subnormality for arbitrary powers of 2-variable weighted shifts whose restrictions to a large invariant subspace are tensor products ${ }^{\star}$ 

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#### Abstract

The Lifting Problem for Commuting Subnormals (LPCS) asks for necessary and sufficient conditions for a pair of subnormal operators on Hilbert space to admit commuting normal extensions. We study LPCS within the class of commuting 2 -variable weighted shifts $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ with subnormal components $T_{1}$ and $T_{2}$, acting on the Hilbert space $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$ with canonical orthonormal basis $\left\{e_{\left(k_{1}, k_{2}\right)}\right\}_{k_{1}, k_{2} \geqslant 0}$. The core of a commuting 2-variable weighted shift $\mathbf{T}, c(\mathbf{T})$, is the restriction of $\mathbf{T}$ to the invariant subspace generated by all vectors $e_{\left(k_{1}, k_{2}\right)}$ with $k_{1}, k_{2} \geqslant 1$; we say that $c(\mathbf{T})$ is of tensor form if it is unitarily equivalent to a shift of the form $\left(I \otimes W_{\alpha}, W_{\beta} \otimes I\right)$, where $W_{\alpha}$ and $W_{\beta}$ are subnormal unilateral weighted shifts. Given a 2-variable weighted shift $\mathbf{T}$ whose core is of tensor form, we prove that LPCS is solvable for $\mathbf{T}$ if and only if LPCS is solvable for any power $\mathbf{T}^{(m, n)}:=\left(T_{1}^{m}, T_{2}^{n}\right)(m, n \geqslant 1)$.


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## 1. Introduction

The Lifting Problem for Commuting Subnormals (LPCS) asks for necessary and sufficient conditions for a pair of subnormal operators on Hilbert space to admit commuting normal extensions. In previous work $[6-9,14-16,27]$ we have studied LPCS from a number of different approaches. One such approach is to consider commuting pairs $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ of subnormal operators and to ask to what extent the existence of liftings for the powers $\mathbf{T}^{(m, n)}:=\left(T_{1}^{m}, T_{2}^{n}\right)$ ( $m, n \geqslant 1$ ) can guarantee a lifting for $\mathbf{T}$. For the class of 2 -variable weighted shifts $\mathbf{T}$, it is often the case that the powers of $\mathbf{T}$ are less complex than the initial pair; thus it becomes especially significant to unravel how subnormality behaves under the action $(m, n) \mapsto \mathbf{T}^{(m, n)}$ ( $h, \ell \geqslant 1$ ).

Within the class of 2 -variable weighted shifts, we consider the subclass $\mathcal{T C}$ consisting of pairs whose cores are of tensor form; that is

$$
\mathcal{T C}:=\left\{\mathbf{T} \in \mathfrak{H}_{0}: c(\mathbf{T}) \text { is of tensor form }\right\} .
$$

This subclass has proved to be particularly attractive, since it is possible to separate, within it, subnormality from $k$-hyponormality; thus, results about LPCS for pairs in $\mathcal{T C}$ are especially useful. The class $\mathcal{T C}$ is small enough to allow for a simple description of its pairs, yet large enough to be used as test ground for many significant hypotheses.

Before we proceed, we briefly pause to establish our terminology. For $\alpha \equiv\left\{\alpha_{n}\right\}_{n=0}^{\infty}$ a bounded sequence of positive real numbers (called weights), let $W_{\alpha}: \ell^{2}\left(\mathbb{Z}_{+}\right) \rightarrow \ell^{2}\left(\mathbb{Z}_{+}\right)$be the associated unilateral weighted shift, defined by $W_{\alpha} e_{n}:=\alpha_{n} e_{n+1}$ (all $n \geqslant 0$ ), where $\left\{e_{n}\right\}_{n=0}^{\infty}$ is the canonical orthonormal basis in $\ell^{2}\left(\mathbb{Z}_{+}\right)$. Similarly, consider double-indexed positive bounded sequences $\alpha_{\mathbf{k}}, \beta_{\mathbf{k}} \in \ell^{\infty}\left(\mathbb{Z}_{+}^{2}\right), \mathbf{k} \equiv\left(k_{1}, k_{2}\right) \in \mathbb{Z}_{+}^{2}$, and let $\left\{e_{\mathbf{k}}\right\}_{\mathbf{k} \in \mathbb{Z}_{+}^{2}}$ be the canonical orthonormal basis in $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$. We define the 2-variable weighted shift $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ acting on $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$ by

$$
T_{1} e_{\mathbf{k}}:=\alpha_{\mathbf{k}} e_{\mathbf{k}+\varepsilon_{1}} \quad \text { and } \quad T_{2} e_{\mathbf{k}}:=\beta_{\mathbf{k}} e_{\mathbf{k}+\varepsilon_{2}}
$$

where $\varepsilon_{1}:=(1,0)$ and $\varepsilon_{2}:=(0,1)$. The core of a commuting 2 -variable weighted shift $\mathbf{T}$ (in symbols, $c(\mathbf{T})$ ) is the restriction of $\mathbf{T}$ to the invariant subspace generated by all vectors $e_{\left(k_{1}, k_{2}\right)}$ with $k_{1}, k_{2} \geqslant 1$; we say that $c(\mathbf{T})$ is of tensor form if it is unitarily equivalent to a shift of the form $\left(I \otimes W_{\alpha}, W_{\beta} \otimes I\right)$, where $W_{\alpha}$ and $W_{\beta}$ are subnormal unilateral weighted shifts. Fig. 1 shows both the weight and Berger measure diagrams of a typical pair in $\mathcal{T C}$. As shown in [8], each $\mathbf{T} \in \mathcal{T C}$ is completely determined by five parameters, i.e., the 1 -variable measures $\sigma, \tau, \xi$ and $\eta$, and the positive number $a \equiv \alpha_{(0,1)}$. As we mentioned before, $\mathcal{T C}$ is of substantial interest to us, since it provides a fertile ground to test results on subnormality and $k$-hyponormality, and in particular about the solubility of LPCS.

Let us now denote the class of commuting pairs of subnormal operators on Hilbert space by $\mathfrak{H}_{0}$, the class of subnormal pairs by $\mathfrak{H}_{\infty}$, and for an integer $k \geqslant 1$, the class of $k$-hyponormal pairs in $\mathfrak{H}_{0}$ by $\mathfrak{H}_{k}$. Clearly, $\mathfrak{H}_{\infty} \subseteq \cdots \subseteq \mathfrak{H}_{k} \subseteq \cdots \subseteq \mathfrak{H}_{1} \subseteq \mathfrak{H}_{0}$; the main results in [14] and [6] show that these inclusions are all proper; moreover, examples illustrating these proper inclusions can be found in $\mathcal{T C}$.

In this paper we show that for $\mathbf{T} \in \mathcal{T C}$, the subnormality of any power $\mathbf{T}^{(m, n)}$ implies the subnormality of $\mathbf{T}$. To accomplish this, we first show that every power of $\mathbf{T} \in \mathcal{T C}$ is the orthogonal direct sum of 2 -variable weighted shifts in $\mathcal{T C}$. Since each 2 -variable weighted shift in $\mathcal{T C}$


Fig. 1. Weight and Berger measure diagrams of a typical 2-variable weighted shift in $\mathcal{T C}$.
is completely determined by five parameters, we then study how the five parameters of each direct summand in a power are related to the five parameters in the initial 2 -variable weighted shift. Next, we recall from [8] that each $\mathbf{T} \in \mathcal{T C}$ is associated with a pair of linear functionals $\varphi \equiv \varphi(\mathbf{T})$ and $\psi \equiv \psi(\mathbf{T})$ (each depending on the five parameters), and that $\mathbf{T}$ is subnormal if and only if $\varphi \geqslant 0$ and $\psi \geqslant 0$. With all of this at our disposal, we proceed to establish a connection between the pair $(\varphi, \psi)$ associated with $\mathbf{T}$ and those associated with the summands in the orthogonal direct sum decomposition of $\mathbf{T}^{(m, n)}$. This eventually leads to the proof of our main result (Theorem 3.1).

This result provides a complete generalization of Theorem 3.9 in [7]. At the time we wrote [7], the techniques available to us allowed us to deal only with the quadratic powers $\mathbf{T}^{(1,2)}$ and $\mathbf{T}^{(2,1)}$; with the aid of a number of additional examples, together with the main result in [8], we have now been able to handle the case of arbitrary powers.

As an application of Theorem 3.1, we can exhibit a hyponormal 2-variable weighted shift such that none of its powers is subnormal. We describe the shift in Example 8.1. This provides a striking and concrete example of the big gap that exists between hyponormality and subnormality for 2-variable weighted shifts, even within a relatively simple class like $\mathcal{T C}$.

## 2. Notation and preliminaries

To describe our results in detail we need some notation; we further expand on our terminology later in this section. For $\mathbf{k}=\left(k_{1}, k_{2}\right) \in \mathbb{Z}_{+}^{2}$, we shall let $\mathcal{M}_{i}\left(\operatorname{resp} . \mathcal{N}_{j}\right)$ be the subspace of $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$ which is spanned by the canonical orthonormal basis vectors associated to indices $\mathbf{k}$ with $k_{1} \geqslant 0$ and $k_{2} \geqslant i$ (resp. $k_{1} \geqslant j$ and $k_{2} \geqslant 0$ ). The core of a 2 -variable weighted shift $\mathbf{T}$ is the restriction of $\mathbf{T}$ to $\mathcal{M}_{1} \cap \mathcal{N}_{1}$; in symbols, $c(\mathbf{T}):=\left.\mathbf{T}\right|_{\mathcal{M}_{1} \cap \mathcal{N}_{1}}$. A 2-variable weighted shift $\mathbf{T}$ is said to be of tensor form if $\mathbf{T} \cong\left(I \otimes W_{\xi}, W_{\eta} \otimes I\right)$, where $W_{\xi}$ and $W_{\eta}$ are unilateral weighted shifts. The class of all 2-variable weighted shifts $\mathbf{T} \in \mathfrak{H}_{0}$ whose cores are of tensor form is denoted by $\mathcal{T C}$; that is,

$$
\mathcal{T C}:=\left\{\mathbf{T} \in \mathfrak{H}_{0}: c(\mathbf{T}) \text { is of tensor form }\right\}
$$

(see Fig. 1(i)).
It is well known that the commutativity of a pair of subnormals is necessary but not sufficient for the existence of a lifting [1,20-22], and it has recently been shown that the joint hyponormality of the pair is necessary but not sufficient [14]. Our previous work [14-17,6-9,26,27] has revealed that the nontrivial aspects of the LPCS are best detected within the class $\mathfrak{H}_{0}$, especially within $\mathcal{T C}$; we thus focus our attention on this class.

For a single operator $T$, the subnormality of all powers $T^{n}(n \geqslant 2)$ does not imply the subnormality of $T$, even if $T$ is a unilateral weighted shift [24, pp. 378-379]. Thus one might guess that if we were to impose a further condition such as the subnormality of a restriction of $T$ to an invariant subspace, e.g., the subnormality of $\left.T\right|_{\bigvee\left\{e_{k} \in \ell^{2}\left(\mathbb{Z}_{+}\right): k \geqslant i\right\}}$ (for some $i \geqslant 1$ ), that $T$ would then be subnormal. However, even if we assume this for $i=1$, the subnormality of $T$ is not guaranteed. For example, let $T:=\operatorname{shift}\left(\frac{1}{3}, \frac{1}{2}, 1,1, \ldots\right)$. Then $\left.T\right|_{\bigvee\left\{e_{k} \in \ell^{2}\left(\mathbb{Z}_{+}\right): k \geqslant 1\right\}} \equiv \operatorname{shift}\left(\frac{1}{2}, 1,1, \ldots\right)$ is subnormal, and also all powers $T^{n}(n \geqslant 2)$ are subnormal, but $T$ is not subnormal. As a mater of fact, no backward extension $\operatorname{shift}\left(\alpha_{0}, \frac{1}{2}, 1,1, \ldots\right)$ can be subnormal [5, Corollary 6]. More generally, the necessary and sufficient conditions for a unilateral weighted shift $W_{\alpha}$ to be subnormal when we assume that $\left.W_{\alpha}\right|_{\bigvee\left\{e_{k} \in \ell^{2}\left(\mathbb{Z}_{+}\right): k \geqslant 1\right\}}$ is subnormal (with Berger measure $\mu$ ) were obtained in [5, Proposition 8]: $W_{\alpha}$ is subnormal if and only if $\frac{1}{t} \in L^{1}(\mu)$ and $\alpha_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\mu)} \leqslant 1$.

In the multivariable case, the analogous results are highly nontrivial, if one further assumes that each component is subnormal. In 1-variable, the subspace $\bigvee\left\{e_{k} \in \ell^{2}\left(\mathbb{Z}_{+}\right): k \geqslant 1\right\}$ can be regarded as "the core of $T$ "; as we move into two variables it is therefore natural to consider the condition $\mathbf{T} \in \mathcal{T} \mathcal{C}$.

To prove our results, we require a number of tools and techniques introduced in previous work, e.g., the Six-point Test (Lemma 4.1), the Backward Extension Theorem for 2-variable weighted shifts (Lemma 4.2) and the formula to reconstruct the Berger measure of a unilateral weighted shift (4.1), together with a new direct sum decomposition for powers of 2-variable weighted shifts which parallels the decomposition used in [11] to analyze $k$-hyponormality for powers of (1-variable) weighted shifts. Concretely, to analyze the power $\mathbf{T}^{(m, n)}$ of a 2-variable weighted shift $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$, we split the ambient space $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$ as an orthogonal direct sum $\bigoplus_{p=0}^{m-1} \bigoplus_{q=0}^{n-1} \mathcal{H}_{(p, q)}^{(m, n)}$, where for $p=0,1, \ldots, m-1$, and $q=0,1, \ldots, n-1$,

$$
\begin{equation*}
\mathcal{H}_{(p, q)}^{(m, n)}:=\bigvee\left\{e_{(m \ell+p, n k+q)}: k, \ell \geqslant 0\right\} . \tag{2.1}
\end{equation*}
$$

Each of the subspaces $\mathcal{H}_{(p, q)}^{(m, n)}$ reduces $T_{1}^{m}$ and $T_{2}^{n}$, and $\mathbf{T}^{(m, n)}$ is subnormal if and only if each summand $\left.\mathbf{T}^{(m, n)}\right|_{\mathcal{H}_{(p, q)}^{(m, n)}}$ is subnormal. For a set of pairs $\mathcal{X}$, let $\bigoplus \mathcal{X}$ denote the set of pairs that can be written as orthogonal sums of pairs in $\mathcal{X}$. We will show in Section 5 that $\bigoplus \mathcal{T C}$ is invariant under the action $(m, n) \mapsto \mathbf{T}^{(m, n)}(m, n \geqslant 1)$.

Briefly stated, our strategy to prove our main result (Theorem 3.1) is as follows: (i) if $\mathbf{T} \in$ $\mathcal{T C}$ then each power $\mathbf{T}^{(m, n)} \in \bigoplus \mathcal{T C}$; (ii) without loss of generality, we can always assume $m=1$; (iii) the pair $(\varphi, \psi)$ associated with $\mathbf{T}$ is directly related to the pairs $\left(\varphi^{(p, q)}, \psi^{(p, q)}\right)$ associated to the direct summands in the orthogonal decomposition of $\mathbf{T}^{(1, n)}$; (iv) if a power $\mathbf{T}^{(1, n)}$ is subnormal, the functionals $\varphi^{(0,0)}$ and $\psi^{(0,1)}$ are both positive; and (v) it then follows that $\varphi$ and $\psi$ are both positive, and therefore $\mathbf{T}$ is subnormal.

We devote the rest of this section to establishing some additional terminology and notation. Let $\mathcal{H}$ be a complex Hilbert space and let $\mathcal{B}(\mathcal{H})$ denote the algebra of bounded linear operators on $\mathcal{H}$. Recall that a bounded linear operator $T \in \mathcal{B}(\mathcal{H})$ is normal if $T^{*} T=T T^{*}$, and subnormal if $T=\left.N\right|_{\mathcal{H}}$, where $N$ is normal and $N(\mathcal{H}) \subseteq \mathcal{H}$. An operator $T$ is said to be hyponormal if $T^{*} T \geqslant T T^{*}$. For $S, T \in \mathcal{B}(\mathcal{H})$, let $[S, T]:=S T-T S$. An $n$-tuple $\mathbf{T}:=\left(T_{1}, \ldots, T_{n}\right)$ of operators on $\mathcal{H}$ is said to be (jointly) hyponormal if the operator matrix

$$
\left[\mathbf{T}^{*}, \mathbf{T}\right]:=\left(\begin{array}{cccc}
{\left[T_{1}^{*}, T_{1}\right]} & {\left[T_{2}^{*}, T_{1}\right]} & \cdots & {\left[T_{n}^{*}, T_{1}\right]} \\
{\left[T_{1}^{*}, T_{2}\right]} & {\left[T_{2}^{*}, T_{2}\right]} & \cdots & {\left[T_{n}^{*}, T_{2}\right]} \\
\vdots & \vdots & \ddots & \vdots \\
{\left[T_{1}^{*}, T_{n}\right]} & {\left[T_{2}^{*}, T_{n}\right]} & \cdots & {\left[T_{n}^{*}, T_{n}\right]}
\end{array}\right)
$$

is positive semidefinite on the direct sum of $n$ copies of $\mathcal{H}$ (cf. [2,10,12,13,23]). For instance, if $n=2$,

$$
\left[\mathbf{T}^{*}, \mathbf{T}\right]:=\left(\begin{array}{cc}
{\left[T_{1}^{*}, T_{1}\right]} & {\left[T_{2}^{*}, T_{1}\right]} \\
{\left[T_{1}^{*}, T_{2}\right]} & {\left[T_{2}^{*}, T_{2}\right]}
\end{array}\right) .
$$

The $n$-tuple $\mathbf{T} \equiv\left(T_{1}, T_{2}, \ldots, T_{n}\right)$ is said to be normal if $\mathbf{T}$ is commuting and each $T_{i}$ is normal, and $\mathbf{T}$ is subnormal if $\mathbf{T}$ is the restriction of a normal $n$-tuple to a common invariant subspace. In particular, a commuting pair $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ is said to be $k$-hyponormal $(k \geqslant 1)$ [6] if

$$
\mathbf{T}(k):=\left(T_{1}, T_{2}, T_{1}^{2}, T_{2} T_{1}, T_{2}^{2}, \ldots, T_{1}^{k}, T_{2} T_{1}^{k-1}, \ldots, T_{2}^{k}\right)
$$

is hyponormal, or equivalently

$$
\left.\left[\mathbf{T}(k)^{*}, \mathbf{T}(k)\right]=\left(\left[\left(T_{2}^{q} T_{1}^{p}\right)^{*}, T_{2}^{m} T_{1}^{n}\right]\right)\right)_{\substack{1 \leqslant p+q \leqslant k \\ 1 \leqslant n+m \leqslant k}} \geqslant 0
$$

Clearly, normal $\Rightarrow$ subnormal $\Rightarrow k$-hyponormal. For $\alpha \equiv\left\{\alpha_{n}\right\}_{n=0}^{\infty}$ a bounded sequence of positive real numbers (called weights), let $W_{\alpha}: \ell^{2}\left(\mathbb{Z}_{+}\right) \rightarrow \ell^{2}\left(\mathbb{Z}_{+}\right)$be the associated unilateral weighted shift, defined by $W_{\alpha} e_{n}:=\alpha_{n} e_{n+1}$ (all $n \geqslant 0$ ), where $\left\{e_{n}\right\}_{n=0}^{\infty}$ is the canonical orthonormal basis in $\ell^{2}\left(\mathbb{Z}_{+}\right)$. For a weighted shift $W_{\alpha}$, the moments of $\alpha$ are given as

$$
\gamma_{k} \equiv \gamma_{k}(\alpha):= \begin{cases}1 & \text { if } k=0  \tag{2.2}\\ \alpha_{0}^{2} \cdots \alpha_{k-1}^{2} & \text { if } k>0\end{cases}
$$

It is easy to see that $W_{\alpha}$ is never normal, and that it is hyponormal if and only if $\alpha_{0} \leqslant$ $\alpha_{1} \leqslant \cdots$. Similarly, consider double-indexed positive bounded sequences $\alpha_{\mathbf{k}}, \beta_{\mathbf{k}} \in \ell^{\infty}\left(\mathbb{Z}_{+}^{2}\right)$, $\mathbf{k} \equiv\left(k_{1}, k_{2}\right) \in \mathbb{Z}_{+}^{2}$. We define the 2-variable weighted shift $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ by

$$
T_{1} e_{\mathbf{k}}:=\alpha_{\mathbf{k}} e_{\mathbf{k}+\varepsilon_{1}} \quad \text { and } \quad T_{2} e_{\mathbf{k}}:=\beta_{\mathbf{k}} e_{\mathbf{k}+\varepsilon_{2}}
$$

where $\varepsilon_{1}:=(1,0)$ and $\varepsilon_{2}:=(0,1)$. Clearly,

$$
\begin{equation*}
T_{1} T_{2}=T_{2} T_{1} \quad \Leftrightarrow \quad \beta_{\mathbf{k}+\varepsilon_{1}} \alpha_{\mathbf{k}}=\alpha_{\mathbf{k}+\varepsilon_{2}} \beta_{\mathbf{k}} \quad\left(\text { all } \mathbf{k} \in \mathbb{Z}_{+}^{2}\right) \tag{2.3}
\end{equation*}
$$

In an entirely similar way one can define multivariable weighted shifts. Given $\mathbf{k} \equiv\left(k_{1}, k_{2}\right) \in \mathbb{Z}_{+}^{2}$, the moment of $(\alpha, \beta)$ of order $\mathbf{k}$ is

$$
\gamma_{\mathbf{k}} \equiv \gamma_{\mathbf{k}}(\alpha, \beta):= \begin{cases}1 & \text { if } k_{1}=0 \text { and } k_{2}=0 \\ \alpha_{(0,0)}^{2} \cdots \alpha_{\left(k_{1}-1,0\right)}^{2} & \text { if } k_{1} \geqslant 1 \text { and } k_{2}=0 \\ \beta_{(0,0)}^{2} \cdots \beta_{\left(0, k_{2}-1\right)}^{2} & \text { if } k_{1}=0 \text { and } k_{2} \geqslant 1 \\ \alpha_{(0,0)}^{2} \cdots \alpha_{\left(k_{1}-1,0\right)}^{2} \cdot \beta_{\left(k_{1}, 0\right)}^{2} \cdots \beta_{\left(k_{1}, k_{2}-1\right)}^{2} & \text { if } k_{1} \geqslant 1 \text { and } k_{2} \geqslant 1\end{cases}
$$

We remark that, due to the commutativity condition (2.3), $\gamma_{\mathbf{k}}$ can be computed using any nondecreasing path from $(0,0)$ to $\left(k_{1}, k_{2}\right)$.

We now recall a well-known characterization of subnormality for multivariable weighted shifts [19] (due to C. Berger [3, II.6.10] and independently established by R. Gellar and L.J. Wallen [18] in the single variable case): $\mathbf{T}$ admits a commuting normal extension if and only if there is a probability measure $\mu$ (which we call the Berger measure of $\mathbf{T}$ ) defined on the 2-dimensional rectangle $R=\left[0, a_{1}\right] \times\left[0, a_{2}\right]$ (where $a_{i}:=\left\|T_{i}\right\|^{2}$ ) such that $\gamma_{\mathbf{k}}=$ $\int_{R} s^{k_{1}} t^{k_{2}} d \mu(s, t)$, for all $\mathbf{k} \equiv\left(k_{1}, k_{2}\right) \in \mathbb{Z}_{+}^{2}$.

The following well-known result, which links the Berger measure of a subnormal unilateral weighted shift with the Berger measure of its restriction to a suitable invariant subspace, will be needed in Section 5.

Lemma 2.1. (See [14, p. 5140].) Let $W_{\alpha}$ be a subnormal unilateral weighted shift and let $\xi$ denote its Berger measure. For $n \geqslant 1$ let $\mathcal{L}_{n}:=\bigvee\left\{e_{h}: h \geqslant n\right\}$ denote the invariant subspace obtained by removing the first $n$ vectors in the canonical orthonormal basis of $\ell^{2}\left(\mathbb{Z}_{+}\right)$. Then the Berger measure of $\left.W_{\alpha}\right|_{\mathcal{L}_{n}}$ is

$$
\begin{equation*}
d \xi_{n}(s):=\frac{s^{n}}{\gamma_{n}} d \xi(s) \tag{2.4}
\end{equation*}
$$

where $\gamma_{n}$ is the $n$-th moment of $\alpha$, given by (2.2).
We will occasionally write $\operatorname{shift}\left(\alpha_{0}, \alpha_{1}, \ldots\right)$ to denote the weighted shift with weight sequence $\left\{\alpha_{k}\right\}_{k=0}^{\infty}$. We also denote by $U_{+}:=\operatorname{shift}(1,1, \ldots)$ the (unweighted) unilateral shift, and for $0<a<1$ we let $S_{a}:=\operatorname{shift}(a, 1,1, \ldots)$. Observe that $U_{+}$and $S_{a}$ are subnormal, with Berger measures $\delta_{1}$ and $\left(1-a^{2}\right) \delta_{0}+a^{2} \delta_{1}$, respectively, where $\delta_{p}$ denotes the point-mass probability measure with support the singleton $\{p\}$.

Given integers $i$ and $m(m \geqslant 1,0 \leqslant i \leqslant m-1)$, consider $\mathcal{H} \equiv \ell^{2}\left(\mathbb{Z}_{+}\right)=\bigvee\left\{e_{n}: n \geqslant 0\right\}$ and define $\mathcal{H}_{i}:=\bigvee\left\{e_{m j+i}: j \geqslant 0\right\}$, so $\mathcal{H}=\bigoplus_{i=0}^{m-1} \mathcal{H}_{i}$. For a sequence $\alpha \equiv\left\{\alpha_{n}\right\}_{n=0}^{\infty}$, let $\alpha(m: i):=\left\{\Pi_{k=0}^{m-1} \alpha_{m j+i+k}\right\}_{j=0}^{\infty}$, that is, $\alpha(m: i)$ denotes the sequence of products of numbers in adjacent packets of size $m$, beginning with the product $\alpha_{i} \cdots \alpha_{i+m-1}$. For example, $\alpha(2: 0): \alpha_{0} \alpha_{1}, \alpha_{2} \alpha_{3}, \alpha_{4} \alpha_{5}, \ldots$, and $\alpha(3: 2): \alpha_{2} \alpha_{3} \alpha_{4}, \alpha_{5} \alpha_{6} \alpha_{7}, \ldots$ Then for $m \geqslant 1$ and $0 \leqslant$ $i \leqslant m-1, W_{\alpha(m: i)}$ is unitarily equivalent to $\left.W_{\alpha}^{m}\right|_{\mathcal{H}_{i}}$. Therefore, $W_{\alpha}^{m}$ is unitarily equivalent to $\bigoplus_{i=0}^{m-1} W_{\alpha(m: i)}$. This analysis naturally leads to the following result, which will be needed in Section 6.

Lemma 2.2. (See [11, Theorem 2.9].) Let $W_{\alpha}$ be a subnormal unilateral weighted shift with Berger measure $\mu$. Then $W_{\alpha(m, i)}$ is subnormal with Berger measure $\mu_{(m, i)}$, where

$$
d \mu_{(m, 0)}(s)=d \mu\left(s^{\frac{1}{m}}\right) \quad \text { and } \quad d \mu_{(m, i)}(s)=\frac{s^{\frac{i}{m}}}{\gamma_{i}} d \mu\left(s^{\frac{1}{m}}\right) \quad \text { for } 1 \leqslant i \leqslant m-1
$$

## 3. Statement of the main result

In [7] we showed that if $\mathbf{T} \in \mathcal{T C}$ (see Fig. 1(i)), then

$$
\begin{equation*}
\mathbf{T}^{(1,2)} \in \mathfrak{H}_{\infty} \quad \Leftrightarrow \quad \mathbf{T}^{(2,1)} \in \mathfrak{H}_{\infty} \quad \Leftrightarrow \quad \mathbf{T} \in \mathfrak{H}_{\infty} . \tag{3.1}
\end{equation*}
$$

The main result in this paper, which follows, is based on (3.1) and a myriad of examples that have arisen in our previous research.

Theorem 3.1. Let $\mathbf{T} \in \mathcal{T C}$. The following statements are equivalent.
(i) $\mathbf{T} \in \mathfrak{H}_{\infty}$;
(ii) $\mathbf{T}^{(m, n)} \in \bigoplus \mathfrak{H}_{\infty}$ for all $m, n \geqslant 1$;
(iii) $\mathbf{T}^{(m, n)} \in \bigoplus \mathfrak{H}_{\infty}$ for some $m, n \geqslant 1$.

## 4. Some basic facts

For the reader's convenience, in this section we list several well-known auxiliary results and definitions which are needed for the proof of the main result. First, to detect hyponormality for 2 -variable weighted shifts we use a simple criterion involving a base point $\mathbf{k}$ in $\mathbb{Z}_{+}^{2}$ and its five neighboring points in $\mathbf{k}+\mathbb{Z}_{+}^{2}$ at path distance at most 2 .

Lemma 4.1 (Six-point Test). (See [4].) Let $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ be a 2 -variable weighted shift, with weight sequences $\alpha$ and $\beta$. Then

$$
\begin{aligned}
& {\left[\mathbf{T}^{*}, \mathbf{T}\right] \geqslant 0} \\
& \quad \Leftrightarrow \quad H_{\left(k_{1}, k_{2}\right)}(1):=\left(\begin{array}{cc}
\alpha_{\mathbf{k}+\varepsilon_{1}}^{2}-\alpha_{\mathbf{k}}^{2} & \alpha_{\mathbf{k}+\varepsilon_{2}} \beta_{\mathbf{k}+\varepsilon_{1}}-\alpha_{\mathbf{k}} \beta_{\mathbf{k}} \\
\alpha_{\mathbf{k}+\varepsilon_{2}} \beta_{\mathbf{k}+\varepsilon_{1}}-\alpha_{\mathbf{k}} \beta_{\mathbf{k}} & \beta_{\mathbf{k}+\varepsilon_{2}}^{2}-\beta_{\mathbf{k}}^{2}
\end{array}\right) \geqslant 0 \quad\left(\text { all } \mathbf{k} \in \mathbb{Z}_{+}^{2}\right) .
\end{aligned}
$$

Next, we present a criterion to detect the subnormality of 2-variable weighted shifts. First, we need some definitions.
(i) Let $\mu$ and $v$ be two positive measures on $\mathbb{R}_{+}$. We say that $\mu \leqslant v$ on $X:=\mathbb{R}_{+}$, if $\mu(E) \leqslant$ $\nu(E)$ for all Borel subset $E \subseteq \mathbb{R}_{+}$; equivalently, $\mu \leqslant \nu$ if and only if $\int f d \mu \leqslant \int f d \nu$ for all $f \in C(X)$ such that $f \geqslant 0$ on $\mathbb{R}_{+}$.
(ii) Let $\mu$ be a probability measure on $X \times Y$, and assume that $\frac{1}{t} \in L^{1}(\mu)$. The extremal measure $\mu_{\text {ext }}$ (which is also a probability measure) on $X \times Y$ is given by $d \mu_{\text {ext }}(s, t):=$ $\left(1-\delta_{0}(t)\right) \frac{1}{t \|^{\frac{1}{t} \|_{L^{1}}(\mu)}} d \mu(s, t)$.
(iii) Given a measure $\mu$ on $X \times Y$, the marginal measure $\mu^{X}$ is given by $\mu^{X}:=\mu \circ \pi_{X}^{-1}$, where $\pi_{X}: X \times Y \rightarrow X$ is the canonical projection onto $X$. Thus $\mu^{X}(E)=\mu(E \times Y)$, for every $E \subseteq X$.

To state the following result, recall the notation in (2.1), and let $\mathcal{M}:=\mathcal{M}_{1} \equiv \mathcal{H}_{(0,1)}^{(1,1)}$.
Lemma 4.2 (Subnormal backward extension). (See [14, Proposition 3.10].) Let $\mathbf{T} \equiv\left(T_{1}, T_{2}\right)$ be a 2-variable weighted shift, and assume that $\left.\mathbf{T}\right|_{\mathcal{M}}$ is subnormal with associated measure $\mu_{\mathcal{M}}$
and that $W_{0}:=\operatorname{shift}\left(\alpha_{00}, \alpha_{10}, \ldots\right)$ is subnormal with associated measure $\sigma$. Then $\mathbf{T}$ is subnormal if and only if
(i) $\frac{1}{t} \in L^{1}\left(\mu_{\mathcal{M}}\right)$;
(ii) $\beta_{00}^{2} \leqslant\left(\left\|\frac{1}{t}\right\|_{L^{1}\left(\mu_{\mathcal{M}}\right)}\right)^{-1}$;
(iii) $\beta_{00}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\mu_{\mathcal{M}}\right)}\left(\mu_{\mathcal{M}}\right)_{e x t}^{X} \leqslant \sigma$.

Moreover, if $\beta_{00}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\mu_{\mathcal{M}}\right)}=1$, then $\left(\mu_{\mathcal{M}}\right)_{\text {ext }}^{X}=\sigma$. In the case when $\mathbf{T}$ is subnormal, its Berger measure $\mu$ is given by

$$
\begin{align*}
d \mu(s, t)= & \beta_{00}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\mu_{\mathcal{M}}\right)} d\left(\mu_{\mathcal{M}}\right)_{\text {ext }}(s, t) \\
& +\left(d \sigma(s)-\beta_{00}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\mu_{\mathcal{M}}\right)} d\left(\mu_{\mathcal{M}}\right)_{\text {ext }}^{X}(s)\right) d \delta_{0}(t) \tag{4.1}
\end{align*}
$$

## 5. The structure of powers of $\mathbf{2}$-variable weighted shifts in $\mathcal{T} \mathcal{C}$

Consider a 2-variable weighted shift $\mathbf{T} \equiv\left(T_{1}, T_{2}\right) \in \mathcal{T C}$ (see Fig. 1(i)). Since $T_{1}$ (resp. $T_{2}$ ) is subnormal, we know that $\operatorname{shift}\left(\alpha_{1}, \alpha_{2}, \alpha_{3} \cdots\right)$ (resp. $\operatorname{shift}\left(\beta_{1}, \beta_{2}, \beta_{3} \cdots\right)$ ) is subnormal; let $\xi$ (resp. $\eta$ ) be its Berger measure. Similarly, let $\sigma$ (resp. $\tau$ ) denote the Berger measure of $\operatorname{shift}\left(x_{0}, x_{1}, x_{2}, \ldots\right)\left(\operatorname{resp} . \operatorname{shift}\left(y_{0}, y_{1}, y_{2}, \ldots\right)\right)$. Finally, let $\tau_{1}$ be the Berger measure of $\left.\operatorname{shift}\left(y_{1}, y_{2}, y_{3}, \ldots\right) \equiv \operatorname{shift}\left(y_{0}, y_{1}, y_{2}, \ldots\right)\right|_{\backslash\left\{e_{k}: k \geqslant 1\right\}}$. Fig. 1 shows the general form of a pair in $\mathcal{T C}$, and that it is uniquely determined by the five parameters $\sigma, \tau, a, \xi$ and $\eta$. Thus, in what follows we will identify a pair $\mathbf{T} \in \mathcal{T C}$ with the 5-tuple $\langle\sigma, \tau, a, \xi, \eta\rangle$. We shall also let $[a, \xi]$ denote the Berger measure of the subnormal unilateral weighted shift $W$ whose 0 -th weight is $a$ and with $\xi$ as the Berger measure of $\left.W\right|_{\mathcal{L}_{1}}$, where $\mathcal{L}_{1}:=\bigvee\left\{e_{k}: k \geqslant 1\right\}$. For instance, in Fig. 1(i) the Berger measure for the first row is $[a, \xi]$, and for the second row is $\left[\frac{a \beta_{1}}{y_{1}}, \xi\right]$. Finally, we shall let $z_{j} \equiv z_{j}(\eta)$ denote the $j$-th weight of the unilateral weighted shift whose Berger measure is $\eta$; that is, $\operatorname{shift}\left(z_{0}, z_{1}, \ldots\right)$ has Berger measure $\eta$.

Lemma 5.1. Let $\mathbf{T} \in \mathcal{T C}$, and let $m, n \geqslant 1$. Then $\mathbf{T}^{(m, n)} \in \bigoplus \mathcal{T C}$.
Proof. First recall from (2.1) that $\ell^{2}\left(\mathbb{Z}_{+}^{2}\right)$ can be written as an orthogonal direct sum $\bigoplus_{p=0}^{m-1} \bigoplus_{q=0}^{n-1} \mathcal{H}_{(p, q)}^{(m, n)}$, where for $p=0,1, \ldots, m-1$ and $q=0,1, \ldots, n-1$ we have $\mathcal{H}_{(p, q)}^{(m, n)}:=$ $\bigvee\left\{e_{(m \ell+p, n k+q)}: k, \ell \geqslant 0\right\}$. Now write $\mathbf{T} \equiv\langle\sigma, \tau, a, \xi, \eta\rangle$. We shall establish that

$$
\begin{equation*}
\mathbf{T}^{(m, n)}=\bigoplus_{p=0}^{m-1} \bigoplus_{q=0}^{n-1}\left\langle\sigma^{(p, q)}, \tau^{(p, q)}, a^{(p, q)}, \xi^{(p, q)}, \eta^{(p, q)}\right\rangle \tag{5.1}
\end{equation*}
$$

where $\left\langle\sigma^{(p, q)}, \tau^{(p, q)}, a^{(p, q)}, \xi^{(p, q)}, \eta^{(p, q)}\right\rangle$ is the 5-tuple associated to the restriction of $\mathbf{T}^{(m, n)}$ to the reducing subspace $\mathcal{H}_{(p, q)}^{(m, n)}$. Since $\mathbf{T}^{(m, n)}=\left(\mathbf{T}^{(m, 1)}\right)^{(1, n)}$, and since $\mathbf{T}^{(m, 1)} \in \bigoplus \mathcal{T} \mathcal{C}$ if and only if $\mathbf{T}^{(1, m)} \in \bigoplus \mathcal{T C}$, it suffices to prove (5.1) in the case $m=1$. The proof is simple but a bit tedious, and it entails careful diagram chasing in Fig. 3. Visual inspection of that figure reveals that when $m=1$ we have $\sigma^{(0,0)}=\sigma, \sigma^{(0,1)}=[a, \xi], \ldots, \sigma^{(0, n-1)}=\left[\frac{a z_{0} \cdots z_{n-3}}{y_{1} \cdots y_{n-2}}, \xi\right](n \geqslant 3)$; moreover,


Fig. 2. Berger measure diagram and weight diagram of the 2-variable weighted shifts in Lemma 5.3 and Example 5.4, respectively.
$\tau^{(0, q)}=\tau_{(n, q)}$, where the latter notation was introduced in Lemma 2.2. Still looking at Fig. 1, we observe that $a^{(0, q)}=\frac{a z_{0} \cdots z_{n-2+q}}{y_{1} \cdots y_{n-1+q}}$ and that $\xi^{(0, q)}=\xi, \eta^{(0,0)}=\eta_{(n, n-1)}$ and $\eta^{(0, q)}=\left(\eta_{(n, q-1)}\right)_{1}$ $(q \geqslant 1)$. This completes the proof.

Corollary 5.2. For $m, n \geqslant 1$ we have $(\bigoplus \mathcal{T C})^{(m, n)} \subseteq \bigoplus \mathcal{T C}$.
We now restate the main result in [8]. First, recall that if $\tau$ is the Berger measure of $\operatorname{shift}\left(y_{0}, y_{1}, \ldots\right)$, we denote by $\tau_{1}$ the Berger measure of $\operatorname{shift}\left(y_{1}, y_{2}, \ldots\right)$. As described in Lemma 2.1, we know that $d \tau_{1}(t) \equiv \frac{t}{y_{0}^{2}} d \tau(t)$.

Lemma 5.3. (See [8, Theorem 2.3].) Let $\mathbf{T} \equiv\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$ be as in Fig. 2(i) and let

$$
\begin{align*}
\psi & :=\tau_{1}-a^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} \eta, \\
\varphi & :=\sigma-y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\psi)} \delta_{0}-a^{2} y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\eta)} \frac{\xi}{s}, \tag{5.2}
\end{align*}
$$

where $y_{0} \equiv \beta_{00}:=\sqrt{\int t d \tau(t)}$. Then $\mathbf{T}$ is subnormal if and only if $\psi \geqslant 0$ and $\varphi \geqslant 0$.
Example 5.4. Assume the very simple case of $\mathbf{T} \equiv\langle\sigma, \tau, a, \xi, \eta\rangle$, where $\sigma:=\left[x, \delta_{1}\right], \tau:=$ $\left[y, \delta_{1}\right], 0<a<1, \xi:=\delta_{1}$ and $\eta:=\delta_{1}$ (cf. Fig. 2(ii)). Then $\psi=\delta_{1}-a^{2} \delta_{1}=\left(1-a^{2}\right) \delta_{1}$ and $\varphi=\left(1-x^{2}\right) \delta_{0}+x^{2} \delta_{1}-y^{2}\left(1-a^{2}\right) \delta_{0}-a^{2} y^{2} \delta_{1}=\left\{\left(1-x^{2}\right)-y^{2}\left(1-a^{2}\right)\right\} \delta_{0}+\left(x^{2}-a^{2} y^{2}\right) \delta_{1}$. Thus, $\mathbf{T}$ is subnormal if and only if $\left(1-x^{2}\right)-y^{2}\left(1-a^{2}\right) \geqslant 0$, a condition identical to that in [14, Proposition 2.11].


Fig. 3. Weight diagrams of $\left.\mathbf{T}^{(1, n)}\right|_{\mathcal{H}_{(0,0)}^{(1, n)}}$ and $\left.\mathbf{T}^{(1, n)}\right|_{\mathcal{H}_{(0,1)}^{(1, n)}}$.
6. The pairs $(\psi, \varphi)$ for $\langle\sigma, \tau, a, \xi, \eta\rangle$ and $\left\langle\sigma^{(p, q)}, \tau^{(p, q)}, a^{(p, q)}, \xi^{(p, q)}, \eta^{(p, q)}\right\rangle$

In this section we establish a direct relationship between the pair $(\psi, \varphi)$ associated to $\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$ and some of the pairs $\left(\psi^{(p, q)}, \varphi^{(p, q)}\right)$ associated to the direct summands $\left\langle\sigma^{(p, q)}, \tau^{(p, q)}, a^{(p, q)}, \xi^{(p, q)}, \eta^{(p, q)}\right\rangle$ in $\langle\sigma, \tau, a, \xi, \eta\rangle^{(m, n)}$.

Proposition 6.1. Let $\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$, and let $n \geqslant 2$. Consider the decomposition $\langle\sigma, \tau, a$, $\xi, \eta\rangle^{(1, n)}=\bigoplus_{q=0}^{n-1}\left\langle\sigma^{(0, q)}, \tau^{(0, q)}, a^{(0, q)}, \xi^{(0, q)}, \eta^{(0, q)}\right\rangle$, and let $(\psi, \varphi)\left(r e s p .\left(\psi^{(0, q)}, \varphi^{(0, q)}\right)\right)$ be the associated pair in Lemma 5.3. Then
(i) $\psi^{(0,1)} \geqslant 0 \Leftrightarrow \psi \geqslant 0$; and
(ii) $\varphi^{(0,0)} \geqslant 0 \Leftrightarrow \varphi \geqslant 0$.

Proof. We refer the reader to Fig. 3. By Lemma 5.1, we have $\sigma^{(0,0)}=\sigma, \sigma^{(0,1)}=[a, \xi], \tau^{(0,0)}=$ $\tau_{(n, 0)}, \tau^{(0,1)}=\tau_{(n, 1)}, a^{(0,0)}=\frac{a z_{0} \cdots z_{n-2}}{y_{1} \cdots y_{n-1}}, a^{(0,1)}=\frac{a z_{0} \cdots z_{n-1}}{y_{1} \cdots y_{n}}, \xi^{(0,0)}=\xi^{(0,1)}=\xi, \eta^{(0,0)}=\eta_{(n, n-1)}$ and $\eta^{(0,1)}=\left(\eta_{(n, 0)}\right)_{1}$. Then

$$
\begin{align*}
\psi^{(0,1)} & =\left(\tau^{(0,1)}\right)_{1}-\left(a^{(0,1)}\right)^{2}\left\|\frac{1}{s}\right\|_{L^{1}\left(\xi^{(0,1)}\right)} \eta^{(0,1)} \\
& =\left(\tau_{(n, 1)}\right)_{1}-\frac{a^{2} z_{0}^{2} \cdots z_{n-1}^{2}}{y_{1}^{2} \cdots y_{n}^{2}}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)}\left(\eta_{(n, 0)}\right)_{1} . \tag{6.1}
\end{align*}
$$

We now calculate $\left(\tau_{(n, 1)}\right)_{1}$ and $\left(\eta_{(n, 0)}\right)_{1}$. From Lemma 2.2 we know that

$$
d \tau_{(n, 1)}(t)=\frac{t^{\frac{1}{n}}}{y_{0}^{2}} d \tau\left(t^{\frac{1}{n}}\right)
$$

so that

$$
d\left(\tau_{(n, 1)}\right)_{1}(t)=\frac{t}{y_{1}^{2} \cdots y_{n}^{2}} d \tau_{(n, 1)}(t)=\frac{t^{1+\frac{1}{n}}}{y_{0}^{2} y_{1}^{2} \cdots y_{n}^{2}} d \tau\left(t^{\frac{1}{n}}\right) \quad(\text { by Lemma 2.1 })
$$

On the other hand, and again using Lemma 2.2, we have

$$
d \eta_{(n, 0)}(t)=d \eta\left(t^{\frac{1}{n}}\right)
$$

so that

$$
d\left(\eta_{(n, 0)}\right)_{1}(t)=\frac{t}{z_{0}^{2} \cdots z_{n-1}^{2}} d \eta_{(n, 0)}(t)=\frac{t}{z_{0}^{2} \cdots z_{n-1}^{2}} d \eta\left(t^{\frac{1}{n}}\right) \quad \text { (again by Lemma 2.1). }
$$

It follows from (6.1) that

$$
\begin{aligned}
d \psi^{(0,1)}(t) & =\frac{t^{1+\frac{1}{n}}}{y_{0}^{2} y_{1}^{2} \cdots y_{n}^{2}} d \tau\left(t^{\frac{1}{n}}\right)-\frac{a^{2} z_{0}^{2} \cdots z_{n-1}^{2}}{y_{1}^{2} \cdots y_{n}^{2}}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} \cdot \frac{t}{z_{0}^{2} \cdots z_{n-1}^{2}} d \eta\left(t^{\frac{1}{n}}\right) \\
& =\frac{t}{y_{1}^{2} \cdots y_{n}^{2}}\left\{\frac{t^{\frac{1}{n}}}{y_{0}^{2}} d \tau\left(t^{\frac{1}{n}}\right)-a^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} d \eta\left(t^{\frac{1}{n}}\right)\right\} \\
& =\frac{t}{y_{1}^{2} \cdots y_{n}^{2}}\left\{d \tau_{1}\left(t^{\frac{1}{n}}\right)-a^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} d \eta\left(t^{\frac{1}{n}}\right)\right\} \\
& =\frac{t}{y_{1}^{2} \cdots y_{n}^{2}} d \psi\left(t^{\frac{1}{n}}\right) .
\end{aligned}
$$

It now readily follows that $\psi^{(0,1)} \geqslant 0$ if and only if $\psi \geqslant 0$, which establishes (i).
To prove (ii), we begin by calculating $\psi^{(0,0)}$. As in (6.1), we have

$$
\begin{aligned}
d \psi^{(0,0)}(t) & =d\left(\tau^{(0,0)}\right)_{1}(t)-\left(a^{(0,0)}\right)^{2}\left\|\frac{1}{s}\right\|_{L^{1}\left(\xi^{(0,0)}\right)} d \eta^{(0,0)}(t) \\
& =d\left(\tau_{(n, 0)}\right)_{1}(t)-\frac{a^{2} z_{0}^{2} \cdots z_{n-2}^{2}}{y_{1}^{2} \cdots y_{n-1}^{2}}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} d \eta_{(n, n-1)}(t) \\
& =\frac{t}{y_{0}^{2} \cdots y_{n-1}^{2}} d \tau\left(t^{\frac{1}{n}}\right)-\frac{a^{2} z_{0}^{2} \cdots z_{n-2}^{2}}{y_{1}^{2} \cdots y_{n-1}^{2}}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} \frac{t^{\frac{n-1}{n}}}{z_{0}^{2} \cdots z_{n-2}^{2}} d \eta\left(t^{\frac{1}{n}}\right)
\end{aligned}
$$

(by Lemmas 2.1 and 2.2)

$$
=\frac{1}{y_{0}^{2} y_{1}^{2} \cdots y_{n-1}^{2}}\left\{t d \tau\left(t^{\frac{1}{n}}\right)-a^{2} y_{0}^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} t^{\frac{n-1}{n}} d \eta\left(t^{\frac{1}{n}}\right)\right\} .
$$

It follows that

$$
\begin{aligned}
y_{0}^{2} y_{1}^{2} \cdots y_{n-1}^{2} \int \frac{1}{t} d \psi^{(0,0)}(t) & =\int d \tau\left(t^{\frac{1}{n}}\right)-a^{2} y_{0}^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} \int t^{-\frac{1}{n}} d \eta\left(t^{\frac{1}{n}}\right) \\
& =1-a^{2} y_{0}^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)}\left\|\frac{1}{t}\right\|_{L^{1}(\eta)} .
\end{aligned}
$$

On the other hand,

$$
\begin{aligned}
y_{0}^{2} \int \frac{1}{t} d \psi(t) & =y_{0}^{2}\left\{\int \frac{1}{t} d \tau_{1}(t)-a^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)} \int \frac{1}{t} d \eta(t)\right\} \\
& =y_{0}^{2} \int \frac{1}{t} \cdot \frac{t}{y_{0}^{2}} d \tau(t)-a^{2} y_{0}^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)}\left\|\frac{1}{t}\right\|_{L^{1}(\eta)} \\
& =1-a^{2} y_{0}^{2}\left\|\frac{1}{s}\right\|_{L^{1}(\xi)}\left\|\frac{1}{t}\right\|_{L^{1}(\eta)} .
\end{aligned}
$$

Thus,

$$
\begin{equation*}
y_{0}^{2} y_{1}^{2} \cdots y_{n-1}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\psi^{(0,0)}\right)}=y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\psi)} . \tag{6.2}
\end{equation*}
$$

Consider now

$$
\varphi^{(0,0)}=\sigma^{(0,0)}-y_{0}^{2} \cdots y_{n-1}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\psi^{(0,0)}\right)} \delta_{0}-\left(a^{(0,0)}\right)^{2} y_{0}^{2} \cdots y_{n-1}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\eta^{(0,0)}\right)} \frac{\xi^{(0,0)}}{s} .
$$

We know that $\sigma^{(0,0)}=\sigma$, that $a^{(0,0)}=\frac{a z_{0} \cdots z_{n-2}}{y_{1} \cdots y_{n-1}}$ and that $\xi^{(0,0)}=\xi$, so using (6.2) we obtain

$$
\varphi^{(0,0)}=\sigma-y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\psi)} \delta_{0}-a^{2} y_{0}^{2} z_{0}^{2} \cdots z_{n-2}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\eta^{(0,0)}\right)} \frac{\xi}{s} .
$$

Since $\varphi=\sigma-y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\psi)} \delta_{0}-a^{2} y_{0}^{2}\left\|\frac{1}{t}\right\|_{L^{1}(\eta) \frac{\xi}{s}}$, it is easy to see that it suffices to prove that $z_{0}^{2} \cdots z_{n-2}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\eta^{(0,0)}\right)}=\left\|\frac{1}{t}\right\|_{L^{1}(\eta)}$. We know that $\eta^{(0,0)}=\eta_{(n, n-1)}$, so

$$
\begin{aligned}
z_{0}^{2} \cdots z_{n-2}^{2}\left\|\frac{1}{t}\right\|_{L^{1}\left(\eta^{(0,0)}\right)} & =z_{0}^{2} \cdots z_{n-2}^{2} \int \frac{1}{t} d \eta_{(n, n-1)}(t) \\
& =z_{0}^{2} \cdots z_{n-2}^{2} \int \frac{1}{t} \frac{t^{\frac{n-1}{n}}}{z_{0}^{2} \cdots z_{n-2}^{2}} d \eta\left(t^{\frac{1}{n}}\right) \\
& =\int t^{-\frac{1}{n}} d \eta\left(t^{\frac{1}{n}}\right)=\left\|\frac{1}{t}\right\|_{L^{1}(\eta)},
\end{aligned}
$$

as desired.
Corollary 6.2. Let $\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$, and let $n \geqslant 2$. Assume that $\langle\sigma, \tau, a, \xi, \eta\rangle^{(1, n)}$ is subnormal. Then $\langle\sigma, \tau, a, \xi, \eta\rangle$ is subnormal.

Proof. Assume that $\langle\sigma, \tau, a, \xi, \eta\rangle^{(1, n)}$ is subnormal, and recall that the power of a 2-variable weighted shift splits as an orthogonal direct sum of 2-variable weighted shifts. Moreover, each summand is in $\mathcal{T} \mathcal{C}$ (because $\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$ ). The fact that $\langle\sigma, \tau, a, \xi, \eta\rangle^{(1, n)}$ is subnormal readily implies that each direct summand is subnormal, and then Lemma 5.3 says that $\psi^{(0, q)} \geqslant 0$ and $\varphi^{(0, q)} \geqslant 0$ (all $q \geqslant 0$ ). In particular, $\psi^{(0,1)} \geqslant 0$ and $\varphi^{(0,0)} \geqslant 0$. It follows from Proposition 6.1 that $\psi \geqslant 0$ and $\varphi \geqslant 0$. Applying Lemma 5.3 once again, we see that $\langle\sigma, \tau, a, \xi, \eta\rangle$ is subnormal.

Corollary 6.3. Let $\langle\sigma, \tau, a, \xi, \eta\rangle \in \mathcal{T C}$, and let $m \geqslant 2$. Assume that $\langle\sigma, \tau, a, \xi, \eta\rangle^{(m, 1)}$ is subnormal. Then $\langle\sigma, \tau, a, \xi, \eta\rangle$ is subnormal.

## 7. Proof of the main theorem

We are now ready to prove our main result, which we restate for the reader's convenience.
Theorem 7.1. Let $\mathbf{T} \in \mathcal{T C}$. The following statements are equivalent.
(i) $\mathbf{T} \in \mathfrak{H}_{\infty}$;
(ii) $\mathbf{T}^{(m, n)} \in \bigoplus \mathfrak{H}_{\infty}$ for all $m, n \geqslant 1$;
(iii) $\mathbf{T}^{(m, n)} \in \bigoplus \mathfrak{H}_{\infty}$ for some $m, n \geqslant 1$.

Proof. It is clear that (i) $\Rightarrow$ (ii) and that (ii) $\Rightarrow$ (iii). Assume that (iii) holds, with $n \geqslant 2$. Since $\mathbf{T}^{(m, n)}=\left(\mathbf{T}^{(m, 1)}\right)^{(1, n)}$, we can use Corollary 6.2 to conclude that $\mathbf{T}^{(m, 1)}$ is subnormal. If we now apply Corollary 6.3 , we obtain that $\mathbf{T}$ is subnormal, as desired.

## 8. An application

In our previous work [14-17,6-9,26,27], we have shown that there are many different families of commuting pairs of subnormal operators, jointly hyponormal but not admitting commuting normal extensions, that is, $\mathbf{T} \in \mathfrak{H}_{1}$ but $\mathbf{T} \notin \mathfrak{H}_{\infty}$ (all $m, n \geqslant 1$ ). As a simple application of Theorem 7.1, we now show that $\mathfrak{H}_{1} \cap \mathcal{T C} \neq \mathfrak{H}_{\infty} \cap \mathcal{T C}$; moreover, there exists $\mathbf{T} \in \mathcal{T C}$, such that $\mathbf{T} \in \mathfrak{H}_{1}$ but $\mathbf{T}^{(m, n)} \notin \bigoplus \mathfrak{H}_{\infty}($ all $m, n \geqslant 1)$. We recall that $\operatorname{shift}\left(x_{0}, x_{1}, \ldots\right)$ and $\operatorname{shift}\left(y_{0}, y_{1}, \ldots\right)$ are subnormal unilateral weighted shifts with Berger measures $\sigma$ and $\tau$, respectively. Consider a contractive 2 -variable weighted shift $\mathbf{T} \in \mathfrak{H}_{0}$ whose weight diagram is given by Fig. 4(i); that is, in the 5-tuple $\langle\sigma, \tau, a, \xi, \eta\rangle$, we have

$$
\begin{aligned}
& * d \sigma(t):=\left(1-\kappa^{2}\right) d \delta_{0}(t)+\frac{\kappa^{2}}{2} d t+\frac{\kappa^{2}}{2} d \delta_{1}(t), \\
& * \tau \text { is the Berger measure of shift }\left(y_{0}, y_{1}, \ldots\right) \text {, with } \tau_{1} \text { the } 2 \text {-atomic Berger measure } \\
& \quad \rho_{0} \delta_{t_{0}}+\rho_{1} \delta_{t_{1}} \text { of the Stampfli subnormal completion of } \sqrt{\omega_{0}}<\sqrt{\omega_{1}}<\sqrt{\omega_{2}}, \\
& * a \text { is a positive number, } \\
& * \xi:=\delta_{1}, \text { and } \\
& * \eta:=\delta_{1} .
\end{aligned}
$$

Example 8.1. Let $\mathbf{T} \equiv\langle\sigma, \tau, a, \xi, \eta\rangle$ be the 2-variable weighted shift given by Fig. 4(i), with $\sigma$, $\tau_{1}, a, \xi$ and $\eta$ as above. Then $\mathbf{T} \in \mathfrak{H}_{1}$ and $\mathbf{T}^{(m, n)} \notin \mathfrak{H}_{\infty}($ all $m, n \geqslant 1)$ if and only if $s(\kappa)<y_{0}<$ $h(\kappa)$, where


Fig. 4. Weight diagram of the 2-variable weighted shift in Example 8.1 and graphs of $s(\kappa)$ and $h(\kappa)$ on the interval $[0,1]$, respectively; here $\omega_{0}:=\frac{5}{6}, \omega_{1}:=\frac{6}{7}, \omega_{2}:=\frac{7}{8}$ and $y_{0}:=\beta_{0}$.

$$
s(\kappa):=\min \left\{\frac{\sqrt{t_{1}}}{a} \sqrt{\rho_{1}}, \sqrt{\frac{\left(1-\kappa^{2}\right)}{\left\|\frac{1}{t}\right\|_{L^{1}\left(\tau_{1}\right)}-\frac{a^{2}}{t_{1}}}}, \frac{\sqrt{t_{1}}}{a} \sqrt{\frac{\kappa^{2}}{2}}, \sqrt{\frac{1}{\left\|\frac{1}{t}\right\|_{L^{1}\left(\tau_{1}\right)}}}\right\}
$$

and

$$
h(\kappa):=\sqrt{\frac{x_{0}^{2} y_{1}^{2}\left(x_{1}^{2}-x_{0}^{2}\right)}{x_{0}^{2}\left(x_{1}^{2}-x_{0}^{2}\right)+\left(a^{2}-x_{0}^{2}\right)^{2}}} .
$$

Fig. 4(ii) specifies a region in the ( $\kappa, \beta_{0}$ ) plane where $\mathbf{T}$ is hyponormal but none of its powers is subnormal. A detailed analysis of Example 8.1 and of other applications of Theorem 3.1 will be discussed elsewhere.

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