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# Energy decay and nonexistence of solution for a reaction-diffusion equation with exponential nonlinearity

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In this work we consider the energy decay result and nonexistence of global solution for a reaction-diffusion equation with generalized Lewis function and nonlinear exponential growth. There are very few works on the reaction-diffusion equation with exponential growth  $f$  as a reaction term by potential well theory. The ingredients used are essentially the Trudinger-Moser inequality.

**Keywords:** reaction-diffusion equation; stable and unstable set; exponential reaction term; decay rate; global nonexistence

**1 Introduction**

In this paper, we study the following initial boundary value problem with generalized Lewis function  $a(x, t)$  which depends on both spacial variable and time:

$$a(x, t)u_t - \Delta u = f(u), \quad x \in \Omega, t > 0, \quad (1)$$

$$u(x, t) = 0, \quad x \in \partial\Omega, t > 0, \quad (2)$$

$$u(x, 0) = u_0(x), \quad x \in \Omega, \quad (3)$$

here  $f(s)$  is a reaction term with exponential growth at infinity to be specified later,  $\Omega$  is a bounded domain with smooth boundary  $\partial\Omega$  in  $R^2$ .

For the reaction-diffusion equation with polynomial growth reaction terms (that is, equation (1) with  $a(x, t) = 1$  and  $f(u) = |u|^{p-1}u$ ), there have been many works in the literature; one can find a review of previous results in [1, 2] and references therein, which are not listed in this paper just for concision. Problem (1)-(3) with  $a(x, t) > 0$  describes the chemical reaction processes accompanied by diffusion [2]. The author of work [1] proved the existence and asymptotic estimates of global solutions and finite time blow-up of problem (1)-(3) with  $a(x, t) > 0$  and the critical Sobolev exponent  $p = \frac{n+2}{n-2}$  for  $f(u) = u^p$ .

In this paper we assume that  $f(s)$  is a reaction term with exponential growth like  $e^{s^2}$  at infinity. When  $a(x, t) = 1$ ,  $f(u) = e^u$ , model (1)-(3) was proposed by [3] and [4]. In this case, Fujita [5] studied the asymptotic stability of the solution. Peral and Vazquez [6] and Pulkkinen [7] considered the stability and blow-up of the solution. Tello [8] and Ioku [9] considered the Cauchy problem of heat equation with  $f(u) \approx e^{u^2}$  for  $|u| \geq 1$ .

Recently, Alves and Cavalcanti [10] were concerned with the nonlinear damped wave equation with exponential source. They proved global existence as well as blow-up of so-

lutions in finite time by taking the initial data inside the potential well [11]. Moreover, they also got the optimal and uniform decay rates of the energy for global solutions.

Motivated by the ideas of [1, 10], we concentrate on studying the uniform decay estimate of the energy and finite time blow-up property of problem (1)-(3) with generalized Lewis function  $a(x, t)$  and exponential growth  $f$  as a reaction term. To the authors' best knowledge, there are very few works in the literature that take into account the reaction-diffusion equation with exponential growth  $f$  as a reaction term by potential well theory. The majority of works in the literature make use of the potential well theory when  $f$  possesses polynomial growth. See, for instance, the works [12–16] and a long list of references therein. The ingredients used in our proof are essentially the Trudinger-Moser inequality (see [17, 18]). We establish decay rates of the energy by considering ideas from the work of Messaoudi [15]. The case of nonexistence results is also treated, where a finite time blow-up phenomenon is exhibited for finite energy solutions by the standard concavity method adapted for our context.

The remainder of our paper is organized as follows. In Section 2 we present the main assumptions and results, Section 3 and Section 4 are devoted to the proof of the main results.

Throughout this study, we denote by  $\|\cdot\|$ ,  $\|\cdot\|_p$ ,  $\|\cdot\|_{H_0^1}$  the usual norms in spaces  $L^2(\Omega)$ ,  $L^p(\Omega)$  and  $H_0^1(\Omega)$ , respectively.

## 2 Assumptions and preliminaries

In this section, we present the main assumptions and results. We always assume that:

- (A1)  $a(x, t)$  is a positive differentiable function and is bounded for  $t \in [0, +\infty)$ ,  $x \in \Omega$ .
- (A2)  $f : R \rightarrow R$  is a  $C^1$  function. The function  $f(t)/t$  is increasing in  $(0, \infty)$ , and for each  $\beta > 0$ , there exists a positive constant  $C_\beta$  such that

$$|f(t)| \leq C_\beta e^{\beta t^2}, \quad |f'(t)| \leq C_\beta e^{\beta t^2}. \tag{4}$$

- (A3) For each  $\varepsilon > 0$ ,  $\beta > 0$  and  $p > 1$  fixed, there exists a positive constant  $C(\varepsilon, \beta)$  such that

$$|f(t)| \leq \varepsilon |t| + C(\varepsilon, \beta) |t|^{p-1} e^{\beta t^2}, \tag{5}$$

$$|F(t)| \leq \varepsilon |t|^2 + C(\varepsilon, \beta) |t|^p e^{\beta t^2}, \tag{6}$$

where  $F(t) = \int_0^t f(s) ds$ .

- (A4) There exists a positive constant  $\theta > 2$  such that

$$0 < \theta F(t) < f(t)t, \quad t \in R \setminus \{0\}. \tag{7}$$

A typical example of functions satisfies (A2)-(A4) is  $f(t) = C|t|^{p-1}te^{Mt^\alpha}$ , with given  $p > 1$ ,  $M > 0$ ,  $C > 0$ , and  $\alpha \in (1, 2)$ .

Now we define some functional as follows:

$$E(t) = E(u) = \frac{1}{2} \|\nabla u\|^2 - \int_{\Omega} F(u) dx, \tag{8}$$

$$I(t) = I(u) = \|\nabla u\|^2 - \int_{\Omega} uf(u) dx, \tag{9}$$

then the ‘potential depth’ given by

$$d = \inf \left\{ \sup_{\lambda \in \mathbb{R}} E(\lambda u), u \in H_0^1 \setminus \{0\} \right\}$$

is a positive constant [10]. Hence, we are able to define stable and unstable sets respectively as follows:

$$W_1 = \{u \in H_0^1, E(u) < d, I(u) > 0\},$$

$$W_2 = \{u \in H_0^1, E(u) < d, I(u) < 0\}.$$

We also need the following lemmas.

**Lemma 2.1** [17, 18] *Let  $\Omega$  be a bounded domain in  $\mathbb{R}^2$ . For all  $u \in H_0^1(\Omega)$ ,*

$$e^{\alpha|u|^2} \in L^1(\Omega) \quad \text{for all } \alpha > 0, \tag{10}$$

*and there exist positive constants  $m_2$  such that*

$$\sup_{u \in H_0^1(\Omega), \|\nabla u\| \leq 1} \int_{\Omega} e^{\alpha|u|^2} dx = m_2 < \infty \quad \text{for all } \alpha \leq 4\pi. \tag{11}$$

**Lemma 2.2** [19] *Let  $\phi(t)$  be a nonincreasing and nonnegative function on  $[0, \infty)$ , such that*

$$\sup_{s \in [t, t+1]} \phi(s) \leq C(\phi(t) - \phi(t+1)), \quad t > 0, \tag{12}$$

*then*

$$\phi(t) \leq Ce^{-\omega t},$$

*where  $C, \omega$  are positive constants depending on  $\phi(0)$  and other known qualities.*

**Lemma 2.3** [20] *Suppose that a positive, twice-differentiable function  $H(t)$  satisfies on  $t \geq 0$  the inequality*

$$H''(t)H(t) - (\delta + 1)(H'(t))^2 \geq 0, \tag{13}$$

*where  $\delta > 0$ , then there is  $t_1 < t_2 = \frac{H(0)}{\delta H'(0)}$  such that  $H(t) \rightarrow \infty$  as  $t \rightarrow t_1$ .*

In order to state and prove our main results, we remind that by the embedding theorem there exists a constant  $C_0$  depending on  $p$  and  $\Omega$  only such that

$$\|u\|_p \leq C_0 \|\nabla u\|. \tag{14}$$

By multiplying equation (1) by  $u_t$ , integrating over  $\Omega$ , using integration by parts and  $a(x, t) > 0$ , we get

$$E'(t) = - \int_{\Omega} a(x, t) u_t^2(x, t) dx \leq 0. \tag{15}$$

Our main results read as follows.

**Theorem 2.1** *Let (A1)-(A4) hold. Assume further that  $u_0 \in W_1$  satisfies*

$$\varepsilon_0 C_0^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \left( \frac{2\theta}{\theta - 2} E(0) \right)^{p-2} < 1 \tag{16}$$

*for some sufficiently small  $\varepsilon_0 > 0$  and  $C_{\varepsilon_0} > 0$ . Then there exist positive constants  $K$  and  $k$  such that the energy  $E(t)$  satisfies the decay estimates for large  $t$*

$$E(t) \leq K e^{-kt}. \tag{17}$$

**Theorem 2.2** *Let (A1)-(A4) hold. Assume further that  $a_t(x, t) \leq 0$ ,  $u_0 \in W_2$  and  $E(0) < \frac{(\theta-2)d}{\theta} < d$ , then the solutions of (1)-(3) blow up in finite time.*

### 3 Proof of decay of the energy

In this section we prove Theorem 2.1. We divide the proof into two lemmas.

**Lemma 3.1** *Under the assumptions of Theorem 2.1, we have, for all  $t \geq 0$ ,  $u(t) \in W_1$ .*

*Proof* Since  $I(u_0) \geq 0$ , then there exists (by continuity)  $T_m < T$  such that

$$I(u(t)) \geq 0, \quad \forall t \in [0, T_m].$$

This and (A4) give

$$\begin{aligned} E(t) &= \left( \frac{1}{2} - \frac{1}{\theta} \right) \|\nabla u\|^2 + \frac{1}{\theta} I(u) + \int_{\Omega} \left( \frac{1}{\theta} u f(u) - F(u) \right) dx \\ &\geq \left( \frac{1}{2} - \frac{1}{\theta} \right) \|\nabla u\|^2, \quad \forall t \in [0, T_m]. \end{aligned} \tag{18}$$

So, by (15) we have

$$\|\nabla u\|^2 \leq \frac{2\theta}{\theta - 2} E(t) \leq \frac{2\theta}{\theta - 2} E(0) \leq \frac{2\theta d}{\theta - 2}, \quad \forall t \in [0, T_m]. \tag{19}$$

We then use (5), the Holder inequality and the embedding theorem to obtain, for each  $t \in [0, T_m]$ ,

$$\begin{aligned} \int_{\Omega} u f(u) dx &\leq \int_{\Omega} [\varepsilon |u|^2 + C(\varepsilon, \beta) |u|^p e^{\beta u^2}] dx \\ &\leq \varepsilon C_0^2 \|\nabla u\|^2 + C(\varepsilon, \beta) \left( \int_{\Omega} |u|^{2p} dx \right)^{\frac{1}{2}} \left( \int_{\Omega} e^{2\beta u^2} dx \right)^{\frac{1}{2}}. \end{aligned} \tag{20}$$

Once  $\|\nabla u\|^2 \leq \frac{2\theta d}{\theta - 2}$ , we choose  $\beta$  such that  $\frac{\theta \beta d}{\theta - 2} < \pi$ , then, from Trudinger-Moser inequality (11),

$$\int_{\Omega} e^{2\beta |u|^2} dx \leq \int_{\Omega} e^{2\beta \|\nabla u\|^2 \left( \frac{|u|}{\|\nabla u\|} \right)^2} dx \leq m_2,$$

and therefore, by (16) for  $\varepsilon_0 > 0$  and  $C_{\varepsilon_0} > 0$ , we have

$$\begin{aligned} \int_{\Omega} uf(u) \, dx &\leq \varepsilon_0 C_0^2 \|\nabla u\|^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \|\nabla u\|^p \\ &= \varepsilon_0 C_0^2 \|\nabla u\|^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \left(\frac{2\theta}{\theta-2} E(t)\right)^{p-2} \|\nabla u\|^2 \\ &\leq \left[\varepsilon_0 C_0^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \left(\frac{2\theta}{\theta-2} E(0)\right)^{p-2}\right] \|\nabla u\|^2 < \|\nabla u\|^2. \end{aligned} \tag{21}$$

By virtue of (21) and the definition of  $I(t)$ , we have

$$I(t) = \|\nabla u\|^2 - \int_{\Omega} uf(u) \, dx > 0.$$

This shows that  $u(t) \in W_1$  for all  $t \in [0, T_m]$ . By repeating this procedure and the fact that  $E(t) \leq E(0)$ , we obtain

$$\lim_{t \rightarrow T_m} \left( \varepsilon_0 C_0^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \left(\frac{2\theta}{\theta-2} E(t)\right)^{p-2} \right) \leq \rho < 1.$$

This is extended to  $T$ . □

**Lemma 3.2** *Under the assumptions of Theorem 2.1, we have, for  $\eta = 1 - [\varepsilon_0 C_0^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \times (\frac{2\theta}{\theta-2} E(0))^{p-2}]$ ,*

$$\eta \|\nabla u\|^2 < I(t). \tag{22}$$

*Proof* It suffices to rewrite (21) as

$$\begin{aligned} \int_{\Omega} uf(u) \, dx &\leq \left[\varepsilon_0 C_0^2 + C_{\varepsilon_0} C_0^p m_2^{\frac{1}{2}} \left(\frac{2\theta}{\theta-2} E(0)\right)^{p-2}\right] \|\nabla u\|^2 \\ &= (1 - \eta) \|\nabla u\|^2 = \|\nabla u\|^2 - \eta \|\nabla u\|^2. \end{aligned} \tag{23}$$

Thus (22) follows from (23). □

*Proof of Theorem 2.1* We integrate (15) over  $[t, t + 1]$  to obtain

$$E(t) - E(t + 1) = \int_t^{t+1} \int_{\Omega} a(x, s) u_t^2(x, s) \, dx \, ds = D^2(t). \tag{24}$$

Now we multiply (1) by  $u$  and integrate over  $\Omega \times [t, t + 1]$  to arrive at

$$\begin{aligned} \int_t^{t+1} I(s) \, ds &= \int_t^{t+1} \left[ \|\nabla u\|^2 - \int_{\Omega} uf(u) \, dx \right] ds \\ &= \int_t^{t+1} \int_{\Omega} a(x, t) u_t(x, t) u(x, t) \, dx \, ds \\ &\leq A \int_t^{t+1} \|a^{\frac{1}{2}} u_t(s)\| \|u(s)\| \, ds, \end{aligned} \tag{25}$$

where  $A^2 = \sup_{(x,t) \in \Omega \times [0,+\infty)} |a(x,t)|$ . Exploiting (14) and (19), we obtain

$$\begin{aligned} \int_t^{t+1} I(s) ds &\leq AC_0 \left( \frac{2\theta}{\theta-2} \right)^{\frac{1}{2}} \left( \sup_{s \in [t,t+1]} E^{\frac{1}{2}}(t) \right) \int_t^{t+1} \|a^{\frac{1}{2}} u_t\| ds \\ &\leq AC_0 \left( \frac{2\theta}{\theta-2} \right)^{\frac{1}{2}} \left( \sup_{s \in [t,t+1]} E^{\frac{1}{2}}(t) \right) D(t). \end{aligned} \tag{26}$$

Using (7), (23) and (22), we have

$$\begin{aligned} E(t) &= \frac{\theta-2}{2\theta} \|\nabla u\|^2 + \frac{1}{\theta} I(t) + \int_{\Omega} \left( \frac{1}{\theta} u f(u) - F(u) \right) dx \\ &\leq \frac{\theta-2}{2\theta} \|\nabla u\|^2 + \frac{1}{\theta} I(t) + \int_{\Omega} \frac{2}{\theta} u f(u) dx \\ &\leq \frac{\theta-2}{2\theta} \|\nabla u\|^2 + \frac{1}{\theta} I(t) + \frac{2}{\theta} (1-\eta) \|\nabla u\|^2 \\ &\leq \left[ \frac{\theta-2}{2\theta\eta} + \frac{1}{\theta} + \frac{2}{\theta} (1-\eta) \right] I(t). \end{aligned} \tag{27}$$

Integrating both sides of (27) over  $[t, t+1]$  and using (26), one can write

$$\int_t^{t+1} E(s) ds = \left[ \frac{\theta-2}{2\theta\eta} + \frac{1}{\theta} + \frac{2}{\theta} (1-\eta) \right] AC_0 \left( \frac{2\theta}{\theta-2} \right)^{\frac{1}{2}} \left( \sup_{s \in [t,t+1]} E^{\frac{1}{2}}(t) \right) D(t). \tag{28}$$

By using (15) again, we have  $E(s) \geq E(t+1), \forall s \leq t+1$ , hence

$$\int_t^{t+1} E(s) ds \geq E(t+1). \tag{29}$$

Inserting (29) in (24) and using (27), we easily have

$$\begin{aligned} E(t) &\leq \int_t^{t+1} E(s) ds + \int_t^{t+1} \int_{\Omega} a(x,s) u_t^2(x,s) dx ds \\ &\leq \left[ \frac{\theta-2}{2\theta\eta} + \frac{1}{\theta} + \frac{2}{\theta} (1-\eta) \right] AC_0 \left( \frac{2\theta}{\theta-2} \right)^{\frac{1}{2}} \left( \sup_{s \in [t,t+1]} E^{\frac{1}{2}}(t) \right) D(t) + D^2(t) \\ &\leq C_1 [E^{\frac{1}{2}}(t) D(t) + D^2(t)] \end{aligned} \tag{30}$$

for  $C_1$  a constant depending on  $C_0, A, \theta, \eta$  only. We then use Young's inequality to get from (30) and (24)

$$\sup_{s \in [t,t+1]} E(t) \leq C_2 D^2(t) \leq C_2 (E(t) - E(t+1)). \tag{31}$$

By (12) in Lemma 2.2 we then get the results. □

#### 4 Proof of the blow-up result

In this section, we shall prove Theorem 2.2 by adapting the concavity method (see Levine [20]). We recall the following lemma in [10].

**Lemma 4.1** [10] *Assume that  $u_0 \in W_2$  and  $E(0) < d$ , then it holds that*

$$u(t) \in W_2 \quad \text{for } t \in [0, T_{\max}), \tag{32}$$

$$\|\nabla u\|^2 \geq 2d \quad \text{for } t \in [0, T_{\max}). \tag{33}$$

*Proof of Theorem 2.2* Assume by contradiction that the solution is global. Then, for any  $T > 0$ , we consider the function  $H(\cdot) : [0, T] \rightarrow R^+$  defined by

$$\begin{aligned} H(t) &= \int_0^t \int_{\Omega} a(x,s)u^2(x,s) \, dx \, ds + \int_0^t \int_{\Omega} (s-t)a_t(x,s)u^2(x,s) \, dx \, ds \\ &\quad + (T-t) \int_{\Omega} a(x,0)u_0^2(x) \, dx + \rho(t+t_0)^2, \end{aligned} \tag{34}$$

where  $t_0, T, \rho$  are positive constants which will be fixed later. Direct computations show that

$$\begin{aligned} H'(t) &= \int_{\Omega} a(x,t)u^2(x,t) \, dx - \int_0^t \int_{\Omega} a_t(x,s)u^2(x,s) \, dx \, ds - \int_{\Omega} a(x,0)u_0^2 \, dx + 2\rho(t+t_0) \\ &= 2 \int_0^t \int_{\Omega} a(x,s)u_t(x,s)u(x,s) \, dx \, ds + 2\rho(t+t_0), \end{aligned} \tag{35}$$

$$H''(t) = 2 \int_{\Omega} a(x,t)u(x,t)u_t(x,s) \, dx + 2\rho. \tag{36}$$

Then, due to equations (1), (7) and (33), we have

$$\begin{aligned} H''(t) &= -2\|\nabla u\|^2 + 2 \int_{\Omega} uf(u) \, dx + 2\rho \\ &\geq -2\|\nabla u\|^2 + 2\theta \int_{\Omega} F(u) \, dx + 2\rho = (\theta-2)\|\nabla u\|^2 - 2\theta E(t) + 2\rho \\ &= (\theta-2)\|\nabla u\|^2 - 2\theta E(0) + 2\theta \int_0^t \int_{\Omega} a(x,s)u_t^2(x,s) \, dx \, ds + 2\rho \\ &\geq 2(\theta-2)d - 2\theta E(0) + 2\theta \int_0^t \int_{\Omega} a(x,s)u_t^2(x,s) \, dx \, ds + 2\rho. \end{aligned} \tag{37}$$

Now we take  $0 < \rho < \frac{(\theta-2)d - \theta E(0)}{\theta-1}$  such that  $2(\theta-2)d - 2\theta E(0) + 2\rho > 2\theta\rho$  (this  $\rho$  can be chosen since  $E(0) < \frac{(\theta-2)d}{\theta}$ ), and then

$$H''(t) \geq 2\theta\rho + 2\theta \int_0^t \int_{\Omega} a(x,s)u_t^2(x,s) \, dx \, ds. \tag{38}$$

We also note that

$$H(0) = T \int_{\Omega} a(x,0)u_0^2(x) \, dx + \rho t_0^2 > 0,$$

$$H'(0) = 2\rho t_0 > 0,$$

$$H''(t) \geq 2\theta\rho > 0, \quad t \geq 0.$$

Therefore  $H(t)$  and  $H'(t)$  are both positive. Since  $a_t(x, t) \leq 0$  for all  $x \in \Omega$  and  $t \geq 0$ , by the construction of  $H(t)$ , it is clear that

$$H(t) \geq \int_0^t \int_{\Omega} a(x, s) u^2(x, s) dx ds + \rho(t + t_0)^2. \quad (39)$$

Thus, for all  $(\xi, \eta) \in R^2$ , from (35), (38) and (39) it follows that

$$\begin{aligned} & H(t)\xi^2 + H'(t)\xi\eta + \frac{1}{2\theta}H''(t)\eta^2 \\ & \geq \left( \int_0^t \int_{\Omega} a(x, s) u(x, s)^2 dx ds + \rho(t + t_0)^2 \right) \xi^2 \\ & \quad + 2\xi\eta \int_0^t \int_{\Omega} a(x, s) u(x, s) u_t(x, s) dx ds + 2\rho(t + t_0)\xi\eta \\ & \quad + \rho\eta^2 + \eta^2 \int_0^t \int_{\Omega} a(x, s) u_t^2(x, s) dx ds \geq 0, \end{aligned}$$

which implies

$$(H'(t))^2 - \frac{2}{\theta}H(t)H''(t) \leq 0.$$

That is,

$$H(t)H''(t) - \frac{\theta}{2}(H'(t))^2 \geq 0.$$

Then we complete the proof by the standard concavity method (Lemma 2.3) since  $\theta > 2$ . □

#### Competing interests

The authors declare that they have no competing interests.

#### Authors' contributions

The authors declare that the study was realized in collaboration with the same responsibility. All authors read and approved the final manuscript.

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