On a Theorem of T. Gneiting on *a*-Symmetric Multivariate Characteristic Functions

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T. Gneiting (1998, J. Multivariate Analysis **64**, 131–147) proved a relation between the primitives of the classes $\Phi_d(2)$ and $\Phi_d(1)$ of 2- and 1-symmetric characteristic functions on \mathbb{R}^d . respectively. We will give a straightforward proof of **CORE**

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1. INTRODUCTION

The characteristic function of a *d*-dimensional distribution is called α -symmetric, $\alpha > 0$, if it can be written in the form

$$\varphi((|v_1|^{\alpha} + \dots + |v_d|^{\alpha})^{1/\alpha}), \qquad v = (v_1, \dots, v_d) \in \mathbb{R}^d, \tag{1}$$

where $\varphi \in C(\overline{\mathbb{R}}_+)$. The set of all these functions will be denoted by $\Phi_d(\alpha)$.

Functions of type (1) are called *radial with respect to the l_{\alpha}-norm*, or α -*radial* for short. Due to Bochner's famous theorem α -symmetric characteristic functions can be interpreted as positive definite α -radial functions.

A basic notion in this context is the so-called *scale mixture* of a function f on \mathbb{R}^d , i.e., a function g on \mathbb{R}^d for which there exists a distribution function G on $\overline{\mathbb{R}}_+$ such that

$$g(x) = \int_0^\infty f(x\tau) \, dG(\tau), \qquad x \in \mathbb{R}^d.$$

The classes we are interested in consist of all functions which can be represented as scale mixtures of a specific function, which we will call the *primitive* generating the class.



In 1938 Schoenberg [9] characterized the class $\Phi_d(2)$ as the set of all scale mixtures of the function Ω_d ; i.e., as the set of all φ satisfying

$$\varphi(t) = \int_0^\infty \Omega_d(tu) \, dF(u), \qquad t \in \mathbb{R}_+, \tag{2}$$

where the primitive Ω_d is defined by

$$\Omega_d(u) = \Gamma\left(\frac{d}{2}\right) \left(\frac{u}{2}\right)^{-(d-2)/2} J_{(d-2)/2}(u), \qquad u \in \mathbb{R}_+.$$
 (3)

 $J_{\nu}(u)$ is the Bessel function of the first kind of order $\nu, \nu \in \mathbb{R}$.

For the class $\Phi_d(1)$, Cambanis *et al.* [4] gave a similar characterization: a function φ is an element of $\Phi_d(1)$ if and only if φ can be expressed in the form

$$\varphi(t) = \int_0^\infty \omega_d(tu) \, dF(u), \qquad t \in \mathbb{R}_+, \tag{4}$$

F, again, being a distribution function on \mathbb{R}_+ . The primitive of the class $\Phi_d(1)$ is given as a Hankel transform by

$$\omega_{d}(u) = \frac{2^{d/2} \Gamma^{2}\left(\frac{d}{2}\right)}{\Gamma\left(\frac{1}{2}\right) \Gamma\left(\frac{d-1}{2}\right)} u^{-(d-2)/2} \int_{1}^{\infty} (\tau^{2} - 1)^{(d-3)/2} \tau^{-(3d-4)/2} \times J_{(d-2)/2}(u\tau) \, d\tau, \qquad u \in \mathbb{R}_{+}.$$
(5)

In this connection, Gneiting [5] recently proved an interesting relation between the two primitives.

THEOREM (Gneiting). For the primitive ω_d of the class $\Phi_d(1)$ and the primitive Ω_{2d-1} of the class $\Phi_{2d-1}(2)$ we have

$$\omega_{d}(t) = \frac{\Gamma^{2}\left(\frac{d}{2}\right)}{\Gamma\left(\frac{1}{2}\right)\Gamma\left(\frac{2d-1}{2}\right)}I^{d-1}\Omega_{2d-1}(t), \qquad t \in \mathbb{R}_{+}, \tag{6}$$

where $If(t) = \int_t^{\infty} f(u) \, du$, $t \in \mathbb{R}_+$, defined for all functions $f: \mathbb{R}_+ \to \mathbb{R}$ for which If(0) exists.

Gneiting proved the theorem by using arguments similar to those of Cambanis. His proof seems to hide the simple structure of the relation (6). Therefore, Gneiting himself asked in [6] on page 24 for a more direct verification just using integration and/or differentiation arguments.

In [2] Berens and the current author gave a representation of the primitives of both classes in terms of *generalized hypergeometric functions*. The advantage of using special functions lies in the fact that there is a powerful calculus available to handle operations like integration and differentiation, as well as integral transforms. We want to use this calculus to give a straightforward proof of the theorem.

2. SOME SPECIAL FUNCTION THEORY

In connection with summability of the inverse Fourier integral on \mathbb{R}^d , Berens and Xu [1] reestablished among others Cambanis' characterization of $\Phi_d(1)$ (cf. [4, Theorem 3.1]); except for the factor $2^d\Gamma^2(d/2)$ they identified the primitive ω_d as the Fourier integral of the (d-1)st B-spline $M_{d-1}(u \mid x_1^2, ..., x_d^2)$, where $x = (x_1, ..., x_d)$ is a point in \mathbb{R}^d and u is a parameter in \mathbb{R}_+ . To be precise, independent of Cambanis they proved

$$M_{d-1}(u \mid (\cdot)_{1}^{2}, ..., (\cdot)_{d}^{2})^{\wedge}(v) = \frac{(\sqrt{u})^{d-2}}{\Gamma\left(\frac{1}{2}\right)\Gamma^{2}\left(\frac{d}{2}\right)}\omega_{d}(\sqrt{u} \mid v \mid_{1}), \qquad v \in \mathbb{R}^{d}.$$
(7)

In [2] the authors picked up the paper [1] and reformulated and reproved the basic results using the calculus of generalized hypergeometric functions. Indeed, the primitives Ω_d and ω_d can be represented as ${}_pF_q$ -symbols. We postpone a detailed discussion of formula (7) for the appendix, since it is not of immediate relevance for the understanding of the following proof.

Using the Pochhammer symbol $(a)_v = a \cdot (a+1) \cdots (a+v-1), v \in \mathbb{N}$, $(a)_0 := 1, a \in \mathbb{C}$, the ${}_pF_q$ -functions are defined for some non-negative integers p and q as

$${}_{p}F_{q}\left[\begin{array}{c}a_{1},...,a_{p}\\b_{1},...,b_{q}\end{array}\right|z\right] = \sum_{k=0}^{\infty} \frac{(a_{1})_{k}\cdots(a_{p})_{k}}{(b_{1})_{k}\cdots(b_{q})_{k}}\cdot\frac{z^{k}}{k!}, \qquad z\in\mathbb{C},$$

where p denotes the number of numerator parameters and q the number of denominator parameters; both groups of parameters can be chosen in the complex plane. In this context the notion *symbol* is used if the hypergeometric series is handled in a formal sense. For a general treatment of the

theory, especially for questions of convergence, asymptotic behavior, etc., the reader is referred to the book by Luke [7].

One example, which will be needed in the following proof, is given by

$$J_{\nu}(z) = \frac{\left(\frac{z}{2}\right)^{\nu}}{\Gamma(\nu+1)} {}_{0}F_{1}\left(\left|\nu+1\right|\right| - \frac{z^{2}}{4}\right)$$

$$\tag{8}$$

[7, Eq. 6.2.7(1)].

To verify Eq. (6) we need two more formulae. The *n*th derivative of the symbol ${}_{p}F_{q}$ is again a hypergeometric ${}_{p}F_{q}$ -series (cf. [7, 3.4.(1)]) and is given by

$$\frac{d^{n}}{dz^{n}} {}_{p}F_{q} \begin{bmatrix} a_{1}, ..., a_{p} \\ b_{1}, ..., b_{q} \end{bmatrix} z \\
= \frac{(a_{1})_{n} \cdots (a_{p})_{n}}{(b_{1})_{n} \cdots (b_{q})_{n}} {}_{p}F_{q} \begin{bmatrix} a_{1}+n, ..., a_{p}+n \\ b_{1}+n, ..., b_{q}+n \end{bmatrix} z].$$
(9)

A second formula gives the derivative of a product of a ${}_{p}F_{q}$ -symbol and a power of the argument (cf. [7, 3.4(2)]):

$$\frac{d^{n}}{dz^{n}} \left\{ z^{\delta}_{\ \ p} F_{q} \begin{bmatrix} a_{1}, ..., a_{p} \\ b_{1}, ..., b_{q} \end{bmatrix} z \right\} \\
= (\delta - n + 1)_{n} z^{\delta - n}_{p+1} F_{q+1} \begin{bmatrix} \delta + 1, a_{1}, ..., a_{p} \\ \delta + 1 - n, b_{1}, ..., b_{q} \end{bmatrix} z \left]. \quad (10)$$

Finally, we need the representation of the primitive ω_d in terms of hypergeometric functions as given in [2]

$$\omega_{d}(\xi) = {}_{1}F_{2} \begin{bmatrix} -\frac{d-2}{2} \\ \frac{1}{2}, \frac{d}{2} \end{bmatrix} -\frac{\xi^{2}}{4} \\ -\frac{\Gamma^{2}\left(\frac{d}{2}\right)}{\Gamma\left(\frac{d+1}{2}\right)\Gamma\left(\frac{d-1}{2}\right)} \xi \cdot {}_{1}F_{2} \begin{bmatrix} -\frac{d-3}{2} \\ \frac{3}{2}, \frac{d+1}{2} \end{bmatrix} -\frac{\xi^{2}}{4} \end{bmatrix}, \quad \xi \in \mathbb{R}_{+}.$$
(11)

3. PROOF OF THE THEOREM

As stated in [5] the operator I can be inverted using differentiation. Instead of Eq. (6) we will therefore verify the equivalent statement

$$\Omega_{2d-1}(t) = (-1)^{d-1} \frac{\Gamma\left(\frac{1}{2}\right) \Gamma\left(\frac{2d-1}{2}\right)}{\Gamma^2\left(\frac{d}{2}\right)} \omega_d^{(d-1)}(t), \qquad t \in \mathbb{R}_+.$$
(12)

One of the two terms in (11) is a polynomial of degree d-2; for d even it's the first term, for d odd it's the second one. The (d-1)th derivative of the polynomial term therefore vanishes.

Let us first consider the case when d is odd. Just for brevity we will use the notation

$$h(\xi) = {}_{1}F_{2} \begin{bmatrix} -\frac{d-2}{2} \\ \frac{1}{2}, \frac{d}{2} \end{bmatrix} -\frac{\xi^{2}}{4}$$

for the nonpolynomial term of the function ω_d . The case d = 1 is quickly verified and omitted. For d > 1 formula (9) gives

$$\frac{d}{d\xi}h(\xi) = \frac{d-2}{d} \cdot 2(-1) i \cdot \left(-\frac{\xi^2}{4}\right)^{1/2} \cdot {}_1F_2 \begin{bmatrix} -\frac{d-2}{2} + 1 \\ \frac{3}{2}, \frac{d}{2} + 1 \end{bmatrix} - \frac{\xi^2}{4} \end{bmatrix},$$

i being the complex unit. Using (10) we continue differentiating to get

$$\frac{d^2}{d\xi^2}h(\xi) = \frac{d-2}{d} \left(\frac{\xi}{2}\right)^{-1} \frac{\xi}{2} \cdot {}_1F_2 \left[\begin{array}{c} -\frac{d-2}{2} + 1 \\ \frac{1}{2}, \frac{d}{2} + 1 \end{array} \right] - \frac{\xi^2}{4}.$$

We remark that the resulting series has a numerator and a denominator parameter in common. The series therefore reduces in order to the above ${}_{1}F_{2}$ -function.

In the next step we use (9) again and continue differentiating. Let us remark that in each odd step the derivative of the argument $\{-(\cdot)^2/4\}$ of the hypergeometric function yields the complex factor which is needed to use formula (10). The negative sign is obtained by inserting the complex unit; this factor vanishes again in the even steps. Because of *d* being odd, differentiation ends in an even step, which yields, after another reduction of order,

$$\left(\frac{d}{d\xi}\right)^{d-1} h(\xi) = C \cdot {}_{1}F_{2} \begin{bmatrix} -\frac{d-2}{2} + \frac{d-1}{2} \\ \frac{1}{2}, \frac{d}{2} + \frac{d-1}{2} \\ \frac{1}{2}, \frac{d}{2} + \frac{d-1}{2} \end{bmatrix} - \frac{\xi^{2}}{4} \end{bmatrix}$$
$$= C \cdot {}_{0}F_{1} \begin{bmatrix} \\ d - \frac{1}{2} \\ - \frac{\xi^{2}}{4} \\ \frac{1}{2} \end{bmatrix}, \quad \xi \in \mathbb{R}_{+},$$

where C denotes a constant which will be determined in the following paragraph. Setting $v = d - \frac{3}{2}$ in the representation (8) of the Bessel function, the $_0F_1$ -series turns out to be the function Ω_{2d-1} . It remains to take a closer look at the constant C. After d-1 differentia-

It remains to take a closer look at the constant C. After d-1 differentiations we have

$$C = \frac{(d-2)(d-4)\cdots(d-(d-1))}{d\cdot(d+2)\cdots(d+(d-3))} = \frac{2^{(d-1)/2}\Gamma\left(\frac{d}{2}\right)\cdot\Gamma\left(\frac{d}{2}\right)}{\Gamma\left(\frac{1}{2}\right)\cdot2^{(d-1)/2}\Gamma\left(\frac{2d-1}{2}\right)}.$$

Observing that $(-1)^{d-1} = 1$, C will be the right constant in (12).

The case of *d* being even is done analogously. Using both differentiation formulae alternatively, the derivatives can again be expressed in terms of hypergeometric functions. Abbreviating again the nonpolynomial term of the function ω_d by

$$h(\xi) = \frac{(-1)\Gamma^2\left(\frac{d}{2}\right)}{\Gamma\left(\frac{d+1}{2}\right)\Gamma\left(\frac{d-1}{2}\right)}\xi \cdot {}_1F_2\left[\frac{-\frac{d-3}{2}}{\frac{3}{2},\frac{d+1}{2}}\right] - \frac{\xi^2}{4}\right],$$

we get

$$\begin{split} \left(\frac{d}{d\xi}\right)^{d-1}h(\xi) \\ &= C \cdot \frac{\left(-1\right)\Gamma^2\left(\frac{d}{2}\right)}{\Gamma\left(\frac{d+1}{2}\right)\Gamma\left(\frac{d-1}{2}\right)} {}_1F_2\left[\begin{array}{c} -\frac{d-3}{2} + \frac{d-2}{2} \\ \frac{1}{2}, \frac{d+1}{2} + \frac{d-2}{2} \\ \frac{1}{2}, \frac{d+1}{2} + \frac{d-2}{2} \\ \end{array}\right], \quad \xi \in \mathbb{R}_+. \end{split}$$

Reduction of order then yields the function Ω_{2d-1} . Here, the constant C has the form

$$C = \frac{(d-3)(d-5)\cdots(d-(d-2)-1)}{(d+1)(d+3)\cdots(d+(d-2)-1)} = \frac{2^{(d-2)/2}\Gamma\left(\frac{d-1}{2}\right)\cdot\Gamma\left(\frac{d+1}{2}\right)}{\Gamma\left(\frac{1}{2}\right)\cdot2^{(d-2)/2}\Gamma\left(\frac{2d-1}{2}\right)}.$$

Since $(-1)^{d-1} = -1$, this too leads to the constant in formula (12).

APPENDIX

The Fourier transform \hat{f} of a function f in $L(\mathbb{R}^d)$ is defined as $\hat{f}(v) = 1/(2\pi)^d \cdot \int_{\mathbb{R}^d} e^{-iv \cdot x} f(x) \, dx, \, v \in \mathbb{R}^d$. Starting with a function φ in $C_0(\overline{\mathbb{R}}_+)$, we want to study nonnegative functions $f \in L(\mathbb{R}^d)$, the Fourier transform of which are positive definite and α -radial; i.e., $\hat{f}(v) = \varphi((|v_1|^{\alpha} + \cdots + |v_d|^{\alpha})^{1/\alpha}) = \varphi(|v|_{\alpha}), \, v \in \mathbb{R}^d$.

Following Richards [8] we get

$$f(x) = \int_0^\infty \varphi(\rho) E_{\rho;d}^{(\alpha)}(x) d\rho, \qquad x \in \mathbb{R}^d,$$
(13)

where the kernel $E_{\rho;d}^{(\alpha)}(x)$ is defined as

$$E_{\rho;d}^{(\alpha)}(x) = \rho^{d-1} \int_{|v|_{\alpha}=1} e^{i\rho x \cdot v} \omega(v),$$

$$\omega(v) = \sum_{j=1}^{d} (-1)^{j-1} v_j dv_1 \cdots dv_{j-1} dv_{j+1} \cdots dv_d.$$

Note that Richards called the kernels $E_{p;d}^{(\alpha)}$ generalized Bessel functions.

At the first glance, the representation (13) is just a formal one, but for $\alpha = 2$ we get the well-known representation

$$E_{\rho;d}^{(2)}(x) = (\sqrt{2\pi})^d |x|_2^{-(d-2)/2} \rho^{d/2} J_{(d-2)/2}(|x|_2 \rho), \tag{14}$$

while for $\alpha = 1$

$$E_{\rho;d}^{(1)}(x) = [x_1^2, ..., x_d^2] H_{\rho;d}(\cdot).$$
(15)

The function $H_{\rho;d}$ is defined as

$$H_{\rho;d}(u) = (-1)^{\lfloor (d-1)/2 \rfloor} 2^d (\sqrt{u})^{d-1} \begin{cases} \sin \rho \sqrt{u}, & \text{for } d \text{ even,} \\ \cos \rho \sqrt{u}, & \text{for } d \text{ odd,} \end{cases} \qquad u \in \mathbb{R}_+.$$

The representation (15) for $\alpha = 1$ was given by Berens and Xu [1], while formula (14) goes back to Schoenberg [9].

Using the formula of Curry and Schoenberg for the *n*th divided difference of a sufficiently smooth function g on \mathbb{R} with knots $x_0, ..., x_n, n \in \mathbb{N}$, not all equal, i.e.,

$$[x_0, ..., x_n] g = \int_{\mathbb{R}} g^{(n)}(u) M_n(u \mid x_0, ..., x_n) du,$$
(16)

we get a second representation of the kernel $E_{\rho;d}^{(1)}$ which enables us to hide the multidimensional structure of $E_{\rho;d}$ into the B-spline kernel,

$$E_{\rho;d}^{(1)}(x) = \int_0^\infty H_{\rho;d}^{(d-1)}(u) M_{d-1}(u \mid x_1^2, ..., x_d^2) du, \qquad x \in \mathbb{R}^d.$$

Compare this with formula (3.6) of [4].

The *B*-spline of order $n, n \in \mathbb{N}$, with knots $x_0, ..., x_n$ on \mathbb{R} hereby is defined as

$$M_n(u \mid x_0, ..., x_n) = [x_0, ..., x_n] \left\{ \frac{(\cdot - u)_+^{n-1}}{(n-1)!} \right\}, \qquad u \in \mathbb{R},$$

where $(u)_{+} = u$, for u > 0, and is 0 otherwise. The spline vanishes outside the largest interval spanned by the knots, is strictly positive inside the interval, and satisfies $\int_{\mathbb{R}} M_n(u | x_0, ..., x_n) du = 1/n!$. Note that formula (16) can also be seen as a definition of the B-spline; cf. deBoor's survey [3] for details. To simplify notation the authors of [1] introduced the functions

$$H_d(u) = (-1)^{\left[(d-1)/2\right]} u^{d-1} \begin{cases} \sin u, & \text{for } d \text{ even,} \\ \cos u, & \text{for } d \text{ odd,} \end{cases} \text{ and}$$
$$h_d(u) = \left(\frac{1}{u} \frac{d}{du}\right)^{d-1} H_d(u).$$

Then

$$H_{\rho;d}(u) = \frac{2^d}{\rho^{d-1}} H_d(\rho \sqrt{u}) \quad \text{and} \quad H_{\rho;d}^{(d-1)}(u) = 2\rho^{d-1} h_d(\rho \sqrt{u}).$$

An obvious transformation of variables then gives

$$f(x) = \int_0^\infty \psi(\sqrt{u}) M_{d-1}(u \mid x_1^2, ..., x_d^2) du, \qquad x \in \mathbb{R}^d,$$
(17)

where

$$\psi(u) = 2 \int_0^\infty \rho^{d-1} \varphi(\rho) h_d(u\rho) d\rho, \qquad u \in \mathbb{R}_+.$$
(18)

In [2] the authors studied the integral transform (18) from the point of view of special functions. They derived a representation of the kernel h_d in terms of generalized hypergeometric functions and determined the kernel, say m_d , of the inverse transform

$$\varphi(\rho) = 2 \int_0^\infty u^{d-1} \psi(u) \, m_d(\rho u) \, du, \qquad \rho \in \mathbb{R}_+.$$

Formally taking the Fourier transform on both sides of Eq. (17) and changing the order of integration yields (7); i.e.,

$$\varphi(|v|_1) = \int_0^\infty \psi(\sqrt{u}) \left(\frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} e^{-iv \cdot x} M_{d-1}(u \mid x_1^2, ..., x_d^2) \, dx \right) du,$$
$$v \in \mathbb{R}^d;$$

see [1] and [2] for details.

For $\alpha = 2$ the situation is more transparent. The reason lies in the fact that the Fourier transform of a 2-radial function is itself 2-radial and vice versa. This is not the case for 1-radial functions, as is shown above. Using $\Omega_d(\rho |x|_2)$ instead of $E_{\rho;d}^{(2)}(x)$, (13) reduces to the classical *Fourier–Bessel*

transform which is self-inverse. To separate the one-dimensional radial part from the multidimensional component we just have to observe that

$$f(x) = \frac{2\pi^{d/2}}{\Gamma\left(\frac{d}{2}\right)} \int_0^\infty \rho^{d-1} \varphi(\rho) \, \Omega_d(|x|_2 \, \rho) \, d\rho, \qquad x \in \mathbb{R}^d.$$

It follows that the transform (17) reduces to the identification $f(x) = \psi(|x|_2), x \in \mathbb{R}^d$.

Theorem 3.1 of Cambanis *et al.* [4] has three equivalent formulations, one in terms of characteristic functions, a second giving a statistical interpretation, and a third using distribution functions. Their central technical lemma (Proposition 3.1) and its proof are purely statistical, but whenever they deal with characteristic functions and/or distribution functions their arguments and the one of Berens and Xu [1] are quite similar.

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