brought to you by **CORE**

RESEARCH

 Boundary Value Problems a SpringerOpen Journal

Open Access

Periodic solutions for N + 2-body problems with N + 1 fixed centers

Furong Zhao^{1,2*}, Fengying Li² and Jian Chen^{2,3}

*Correspondence: zhaofurong2006@163.com ¹Department of Mathematics and Computer Science, Mianyang Normal University, Mianyang, Sichuan 621000, P.R. China ²Department of Mathematics, Sichuan University, Chengdu, 610064, P.R. China Full list of author information is available at the end of the article

Abstract

In this paper, we prove the existence of a new periodic solution for N + 2-body problems with N + 1 fixed centers and strong-force potentials. In this model, N particles with equal masses are fixed at the vertices of a regular N-gon and the (N + 1)th particle is fixed at the center of the N-gon, the (N + 2)th particle winding around N particles.

MSC: 34C15; 34C25; 70F10

Keywords: N + 2-body problems with N + 1-fixed centers; minimizing variational methods; strong-force potentials

1 Introduction and main results

In the eighteenth century, the 2-fixed center problem was studied by Euler [1–3]. Here, let us consider the N + 1-fixed center problem: We assume N particles q_1, q_2, \ldots, q_N with equal masses 1 are fixed at the vertices $e^{\sqrt{-1}\frac{2\pi}{N}j} = (\cos\frac{2\pi j}{N}, \sin\frac{2\pi j}{N})$ $(j = 1, \ldots, N)$ of a regular polygon and the (N+1)th particle q_{N+1} is fixed at the origin (0, 0), the (N+2)th particle with mass m_{N+2} is attracted by the other particles, and moves according to Newton's second law and a more general power law than the Newton's universal gravitational square law. In this system, the position q(t) for the (N + 2)th particle satisfies the following equation:

$$m_{N+2}\ddot{q}(t) = \sum_{i=1}^{N+1} \frac{m_i m_{N+2}(q(t) - q_i)}{|q(t) - q_i|^{\alpha+2}}.$$
(1.1)

Equivalently,

$$\ddot{q}(t) = \sum_{i=1}^{N} \frac{(q(t) - q_i)}{|q(t) - q_i|^{\alpha + 2}} + \frac{m_{N+1}(q(t) - q_{N+1})}{|q(t) - q_{N+1}|^{\alpha + 2}},$$
(1.2)

$$\ddot{q}(t) = \frac{\partial U(q)}{\partial q},\tag{1.3}$$

where

$$\alpha > 0$$
 and $U(q) = \sum_{i=1}^{N} \frac{1}{|q(t) - q_i|^{\alpha}} + \frac{m_{N+1}}{|q(t) - q_{N+1}|^{\alpha}}.$

The type of system (1.2) is called a singular Hamiltonian system which attracts many researchers (see [1-10] and [11-16]).



© 2013 Zhao et al.; licensee Springer. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/2.0), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited. Specially, Gordon [10] proved the Keplerian elliptical orbits are the minimizers of Lagrangian action defined on the space for non-zero winding numbers.

In this paper, we use a variational minimizing method to look for a periodic solution for the (N + 2)th particle which winds around the q_i (i = 1, ..., N + 1).

Definition 1.1 [10] Let $C : x(t) : [a, b] \to R^2$ be a given oriented closed curve, and $p \notin C$. Define $\varphi : C \to S^1$:

$$\varphi(t)=\frac{x(t)-p}{|x(t)-p|}.$$

When some point on *C* goes around the curve once, its image point $\varphi(x(t))$ will go around S^1 a number of times. This number is defined as the winding number of the curve *C* relative to the point *p* and is denoted by deg(x(t) - p).

Let

$$f(q) = \int_0^1 \left[\frac{1}{2} \left| \dot{q}(t) \right|^2 + U(q) \right] dt, \tag{1.4}$$

$$deg(q(t) - q_i) = 1, \quad \text{for } i = 1, \dots, N, \qquad deg(q(t) - q_{N+1}) = -1$$

$$q \in \Lambda_2 = \begin{cases} q \in W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^2), \quad q(t) \neq q_i, \quad \text{for } i = 1, \dots, N+1, \\ q(t + \frac{k}{N}) = \begin{pmatrix} \cos(\frac{2k\pi}{N}) - \sin(\frac{2k\pi}{N}) \\ \sin(\frac{2k\pi}{N}) - \cos(\frac{2k\pi}{N}) \end{pmatrix} q(t), \qquad (1.6)$$

$$q \in \Lambda_{3} = \begin{cases} q \in W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^{2}), & q(t) \neq q_{i}, & \text{for } i = 1, \dots, N + 1, \\ q(t + \frac{k}{N}) = \begin{pmatrix} \cos(\frac{2k\pi}{N}) - \sin(\frac{2k\pi}{N}) \\ \sin(\frac{2k\pi}{N}) \cos(\frac{2k\pi}{N}) \end{pmatrix} q(t), \\ \deg(q(t) - q_{i}) = 1, & \text{for } i = 1, \dots, N, & \deg(q(t) - q_{N+1}) = 1 \end{cases},$$
(1.7)

$$q \in \Lambda_{4} = \begin{cases} q \in W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^{2}), & q(t) \neq q_{i}, & \text{for } i = 1, \dots, N+1, \\ q(t + \frac{k}{N}) = \begin{pmatrix} \cos(\frac{2k\pi}{N}) - \sin(\frac{2k\pi}{N}) \\ \sin(\frac{2k\pi}{N}) & \cos(\frac{2k\pi}{N}) \end{pmatrix} q(t), \\ \deg(q(t) - q_{i}) = 1, & \text{for } i = 1, \dots, N, & \deg(q(t) - q_{N+1}) = N-1 \end{cases}$$
(1.8)

We have the following theorem.

Theorem 1.1 For $\alpha \ge 2$, the minimizer of f(q) on $\overline{\Lambda}_i$ (i = 1, 2, 3, 4) exists and it is a noncollision periodic solution of (1.1) or (1.2)-(1.3) (please see Figures 1-4 for N = 4).

2 The proof of Theorem 1.1

We recall the following famous lemmas, which we need to prove Theorem 1.1.

Lemma 2.1 [9] If $x \in W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^2)$, $\alpha \ge 2$, a > 0, and there exists $t_0 \in [0,1]$ such that $x(t_0) = 0$, then $\int_0^1 [\frac{1}{2} |\dot{x}(t)|^2 + \frac{a}{|x(t)|^{\alpha}}] dt = +\infty$. If $x_n \rightharpoonup x$ in $W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^2)$ and $\exists t_0$, s.t. $x(t_0) = 0$, $\alpha \ge 2$, then $\int_0^1 \frac{1}{|x_n(t)|^{\alpha}} dt \rightarrow +\infty$.







Lemma 2.2 (Palais's symmetry principle [17]) Let σ be an orthogonal representation of a finite or compact group G on a real Hilbert space H, and let $f : H \to R$ be such that for $\forall \sigma \in G, f(\sigma \cdot x) = f(x)$. Set $H^G = \{x \in H : \sigma \cdot x = x, \forall \sigma \in G\}$. Then the critical point of f in H^G is also a critical point of f in H.

Lemma 2.3 [5] If X is a reflexive Banach space, M is a weakly closed subset of X, and $f: M \to R \cup \{+\infty\}, f \neq +\infty$ is weakly lower semi-continuous and coercive, then f attains its infimum on M.



Lemma 2.4 (Poincare-Wirtinger inequality) Let $q \in W^{1,2}(\mathbb{R}/\mathbb{Z}\mathbb{T},\mathbb{R}^d)$ and $\int_0^T q(t) dt = 0$, then $\int_0^T |\dot{q}(t)|^2 dt \ge (\frac{2\pi}{T})^2 \int_0^T q(t)^2 dt$. And the inequality takes the equality if and only if $q(t) = \alpha \cos \frac{2\pi}{T} t + \beta \sin \frac{2\pi}{T} t$, $\alpha, \beta \in \mathbb{R}^d$.

We now prove Theorem 1.1.

Proof By the symmetry of Λ_i , we know for $\forall x \in \Lambda_i$,

$$\int_{0}^{T} q(t) dt = 0.$$
 (2.1)

If $q_n(t) \rightharpoonup q(t)$ in $\overline{\Lambda}_i$, then by Sobolev's compact embedding theorem, we have $q_n(t) \rightarrow q(t)$ in C[0,1].

- (i) If $q(t) \in \Lambda_i$, then $\lim_{n \to +\infty} \int_0^1 U(q_n(t)) dt = \int_0^1 U(q_n(t)) dt$. Since $\int_0^1 q_n dt = 0$, $\frac{1}{2} \int_0^1 |\dot{q}_n|^2 dt$ can be regarded as the square of an equivalent norm for $W^{1,2}$, so it is weakly lower semi-continuous, so $\lim_{n \to +\infty} f(q_n(t)) \ge f(q)$.
- (ii) If $q(t) \in \partial \Lambda_i$, then by Lemma 2.1, $f(q) = +\infty$, we have $\int_0^1 U(q_n(t)) dt \to +\infty$. So, $\underline{\lim}_{n \to +\infty} f(q_n) = +\infty \ge f(q)$. Hence *f* is w.l.s.c.

Using (2.1), we know that f(q) is coercive on $\overline{\Lambda}_i$. Lemma 2.3 guarantees that f(q) attains its infimum on $\overline{\Lambda}_i$. Let the minimizer be \tilde{q} , then

$$f(\widetilde{q}) = \inf_{q \in \overline{\Lambda}_i} f(q) < +\infty.$$
(2.2)

If \tilde{q} is a collision periodic solution, then there exist $t_0 \in [0,1]$ and $j \in \{1, 2, ..., N, N + 1\}$ such that $\tilde{q}(t_0) = q_j$. Let $x(t) = \tilde{q}(t) - q_j$ and note $x(t_0) = 0$. By Lemma 2.1, we have

$$f(\widetilde{q}) = \int_0^1 \left[\frac{1}{2} |\widetilde{q}(t)|^2 + \frac{m_j}{|\widetilde{q}(t) - q_j|^{\alpha}} + \sum_{i \neq j}^{N+1} \frac{m_i}{|\widetilde{q}(t) - q_i|^{\alpha}} \right] dt$$

$$\geq \int_0^1 \left[\frac{1}{2} |\dot{x}(t)|^2 + \frac{m_j}{|x(t)|^{\alpha}} \right] dt = +\infty, \qquad (2.3)$$

which contradicts the inequality in (2.2). By Lemma 2.2, $\tilde{q}(t)$ is the critical point of f in $W^{1,2}(\mathbb{R}/\mathbb{Z}, \mathbb{R}^2)$; therefore, $\tilde{q}(t)$ is a non-collision periodic solution.

This completes the proof of Theorem 1.1.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

The authors declare that the study was realized in collaboration with the same responsibility. All authors read, checked and approved the final manuscript.

Author details

¹Department of Mathematics and Computer Science, Mianyang Normal University, Mianyang, Sichuan 621000, P.R. China. ²Department of Mathematics, Sichuan University, Chengdu, 610064, P.R. China. ³School of Science, Southwest University of Science and Technology, Mianyang, Sichuan 621000, P.R. China.

Acknowledgements

The authors sincerely thank the referees for their many helpful comments and suggestions and also express their sincere gratitude to Professor Zhang Shiqing for his discussions and corrections. This work is supported by NSF of China and Youth Fund of Mianyang Normal University.

Received: 26 January 2013 Accepted: 2 May 2013 Published: 17 May 2013

References

- 1. Euler, M: De motu coproris ad duo centra virium fixa attracti. Nov. Commun. Acad. Sci. Imp. Petrop. 10, 207-242 (1766)
- 2. Euler, M: De motu coproris ad duo centra virium fixa attracti. Nov. Commun. Acad. Sci. Imp. Petrop. 11, 152-184 (1767)
- 3. Euler, M: Probleme un corps etant attire en raison reciproque quarree des distances vers deux points fixes donnes
- trouver les cas ou la courbe decrite par ce corps sera algebrique. Hist. Acad. R. Sci. Bell. Lett. Berlin 2, 228-249 (1767)
 Ambrosetti, A, Coti Zelati, V: Critical points with lack of compactness and applications to singular dynamical systems. Ann. Mat. Pura Appl. 149, 237-259 (1987)
- 5. Ambrosetti, A, Coti Zelati, V: Periodic Solutions for Singular Lagrangian Systems. Springer, Boston (1993)
- Bahri, A, Rabinowitz, PH: A minimax method for a class of Hamiltonian systems with singular potentials. J. Funct. Anal. 82, 412-428 (1989)
- Benci, V, Giannoni, G: Periodic solutions of prescribed energy for a class of Hamiltonian system with singular potentials. J. Differ. Equ. 82, 60-70 (1989)
- Degiovanni, M, Giannoni, F: Dynamical systems with Newtonian type potentials. Ann. Sc. Norm. Super. Pisa, Cl. Sci. 15, 467-494 (1988)
- 9. Gordon, W: Conservative dynamical systems involving strong forces. Trans. Am. Math. Soc. 204, 113-135 (1975)
- 10. Gordon, W: A minimizing property of Keplerian orbits. Am. J. Math. 99, 961-971 (1977). doi:10.2307/2373993
- Rabinowitz, PH: A note on periodic solutions of prescribed energy for singular Hamiltonian systems. J. Comput. Appl. Math. 52, 147-154 (1994)
- 12. Siegel, C, Moser, J: Lectures on Celestial Mechanics. Springer, Berlin (1971)
- 13. Wang, XR, He, S: Lagrangian actions on 3-body problems with two fixed centers. Bound. Value Probl. 2012, 28 (2012)
- Tanaka, K: A prescribed energy problem for a singular Hamiltonian system with weak force. J. Funct. Anal. 113, 351-390 (1993)
- 15. Tanaka, K: A prescribed energy problem for conservative singular Hamiltonian system. Arch. Ration. Mech. Anal. 128, 127-164 (1994)
- 16. Zhang, SQ: Multiple geometrically distinct closed noncollision orbits of fixed energy for *N*-body type problems with strong force potentials. Proc. Am. Math. Soc. **124**, 3039-3046 (1996)
- 17. Palais, R: The principle of symmetric criticality. Commun. Math. Phys. 69, 19-30 (1979)

doi:10.1186/1687-2770-2013-129

Cite this article as: Zhao et al.: **Periodic solutions for** *N* + 2**-body problems with** *N* + 1 **fixed centers**. *Boundary Value Problems* 2013 **2013**:129.

Submit your manuscript to a SpringerOpen[®] journal and benefit from:

- ► Convenient online submission
- ► Rigorous peer review
- Immediate publication on acceptance
- Open access: articles freely available online
- ► High visibility within the field
- ▶ Retaining the copyright to your article

Submit your next manuscript at > springeropen.com