# Periodic solutions for $N+2$-body problems with $N+1$ fixed centers 

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#### Abstract

In this paper, we prove the existence of a new periodic solution for $N+2$-body problems with $N+1$ fixed centers and strong-force potentials. In this model, $N$ particles with equal masses are fixed at the vertices of a regular N -gon and the $(N+1)$ th particle is fixed at the center of the $N$-gon, the $(N+2)$ th particle winding around $N$ particles.


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## 1 Introduction and main results

In the eighteenth century, the 2 -fixed center problem was studied by Euler [1-3]. Here, let us consider the $N+1$-fixed center problem: We assume $N$ particles $q_{1}, q_{2}, \ldots, q_{N}$ with equal masses 1 are fixed at the vertices $e^{\sqrt{-1} \frac{2 \pi}{N} j}=\left(\cos \frac{2 \pi j}{N}, \sin \frac{2 \pi j}{N}\right)(j=1, \ldots, N)$ of a regular polygon and the $(N+1)$ th particle $q_{N+1}$ is fixed at the origin $(0,0)$, the $(N+2)$ th particle with mass $m_{N+2}$ is attracted by the other particles, and moves according to Newton's second law and a more general power law than the Newton's universal gravitational square law. In this system, the position $q(t)$ for the $(N+2)$ th particle satisfies the following equation:

$$
\begin{equation*}
m_{N+2} \ddot{q}(t)=\sum_{i=1}^{N+1} \frac{m_{i} m_{N+2}\left(q(t)-q_{i}\right)}{\left|q(t)-q_{i}\right|^{\alpha+2}} . \tag{1.1}
\end{equation*}
$$

Equivalently,

$$
\begin{align*}
& \ddot{q}(t)=\sum_{i=1}^{N} \frac{\left(q(t)-q_{i}\right)}{\left|q(t)-q_{i}\right|^{\alpha+2}}+\frac{m_{N+1}\left(q(t)-q_{N+1}\right)}{\left|q(t)-q_{N+1}\right|^{\alpha+2}},  \tag{1.2}\\
& \ddot{q}(t)=\frac{\partial U(q)}{\partial q} \tag{1.3}
\end{align*}
$$

where

$$
\alpha>0 \quad \text { and } \quad U(q)=\sum_{i=1}^{N} \frac{1}{\left|q(t)-q_{i}\right|^{\alpha}}+\frac{m_{N+1}}{\left|q(t)-q_{N+1}\right|^{\alpha}} .
$$

The type of system (1.2) is called a singular Hamiltonian system which attracts many researchers (see [1-10] and [11-16]).

[^0]Specially, Gordon [10] proved the Keplerian elliptical orbits are the minimizers of Lagrangian action defined on the space for non-zero winding numbers.
In this paper, we use a variational minimizing method to look for a periodic solution for the $(N+2)$ th particle which winds around the $q_{i}(i=1, \ldots, N+1)$.

Definition 1.1 [10] Let $C: x(t):[a, b] \rightarrow R^{2}$ be a given oriented closed curve, and $p \notin C$. Define $\varphi: C \rightarrow S^{1}$ :

$$
\varphi(t)=\frac{x(t)-p}{|x(t)-p|} .
$$

When some point on $C$ goes around the curve once, its image point $\varphi(x(t))$ will go around $S^{1}$ a number of times. This number is defined as the winding number of the curve $C$ relative to the point $p$ and is denoted by $\operatorname{deg}(x(t)-p)$.

Let

$$
\begin{align*}
& f(q)=\int_{0}^{1}\left[\frac{1}{2}|\dot{q}(t)|^{2}+U(q)\right] d t,  \tag{1.4}\\
& q \in \Lambda_{1}=\left\{\begin{array}{l}
q \in W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right), \quad q(t) \neq q_{i}, \quad \text { for } i=1, \ldots, N+1, \\
q\left(t+\frac{k}{N}\right)=\binom{\cos \left(\frac{2 k \pi}{N}\right)-\sin \left(\frac{2 k \pi}{N}\right)}{\sin \left(\frac{2 k \pi}{N}\right) \cos \left(\frac{2 k \pi}{N}\right)} q(t), \\
\operatorname{deg}\left(q(t)-q_{i}\right)=1, \quad \text { for } i=1, \ldots, N, \quad \operatorname{deg}\left(q(t)-q_{N+1}\right)=-1
\end{array}\right\},  \tag{1.5}\\
& \left.q \in \Lambda_{2}=\left\{\begin{array}{l}
q \in W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right), \quad q(t) \neq q_{i}, \quad \text { for } i=1, \ldots, N+1, \\
q\left(t+\frac{k}{N}\right)=\binom{\cos \left(\frac{2 k \pi}{N}\right)-\sin \left(\frac{2 k \pi}{N}\right)}{\sin \left(\frac{2 k \pi}{N}\right)} q(t), \\
\cos \left(\frac{2 k \pi}{N}\right)
\end{array}\right) q\left(q(t)-q_{i}\right)=0, \quad \text { for } i=1, \ldots, N, \quad \operatorname{deg}\left(q(t)-q_{N+1}\right)=1\right\},  \tag{1.6}\\
& q \in \Lambda_{3}=\left\{\begin{array}{l}
q \in W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right), \quad q(t) \neq q_{i}, \quad \text { for } i=1, \ldots, N+1, \\
q\left(t+\frac{k}{N}\right)=\binom{\cos \left(\frac{2 k \pi}{N}\right)-\sin \left(\frac{2 k \pi}{N}\right)}{\sin \left(\frac{2 k \pi}{N}\right) \cos \left(\frac{2 k \pi}{N}\right)} q(t), \\
\operatorname{deg}\left(q(t)-q_{i}\right)=1, \quad \text { for } i=1, \ldots, N, \quad \operatorname{deg}\left(q(t)-q_{N+1}\right)=1
\end{array}\right\},  \tag{1.7}\\
& q \in \Lambda_{4}=\left\{\begin{array}{l}
q \in W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right), \quad q(t) \neq q_{i}, \quad \text { for } i=1, \ldots, N+1, \\
q\left(t+\frac{k}{N}\right)=\left(\begin{array}{c}
\cos \left(\frac{2 k \pi}{N}\right)-\sin \left(\frac{2 k \pi}{N}\right) \\
\sin \left(\frac{2 k \pi}{N}\right) \\
\cos \left(\frac{2 k \pi}{N}\right)
\end{array}\right) q(t), \\
\operatorname{deg}\left(q(t)-q_{i}\right)=1, \quad \text { for } i=1, \ldots, N, \quad \operatorname{deg}\left(q(t)-q_{N+1}\right)=N-1
\end{array}\right\} . \tag{1.8}
\end{align*}
$$

We have the following theorem.
Theorem 1.1 For $\alpha \geq 2$, the minimizer of $f(q)$ on $\bar{\Lambda}_{i}(i=1,2,3,4)$ exists and it is a noncollision periodic solution of (1.1) or (1.2)-(1.3) (please see Figures 1-4 for $N=4$ ).

## 2 The proof of Theorem 1.1

We recall the following famous lemmas, which we need to prove Theorem 1.1.

Lemma 2.1 [9] If $x \in W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right), \alpha \geq 2, a>0$, and there exists $t_{0} \in[0,1]$ such that $x\left(t_{0}\right)=0$, then $\int_{0}^{1}\left[\frac{1}{2}|\dot{x}(t)|^{2}+\frac{a}{|x(t)|^{\alpha}}\right] d t=+\infty$.

If $x_{n} \rightharpoonup x$ in $W^{1,2}\left(R / Z, R^{2}\right)$ and $\exists t_{0}$, s.t. $x\left(t_{0}\right)=0, \alpha \geq 2$, then $\int_{0}^{1} \frac{1}{\left|x_{n}(t)\right|^{\alpha}} d t \rightarrow+\infty$.


Figure $3 q \in \boldsymbol{\Lambda}_{3}$.


Lemma 2.2 (Palais's symmetry principle [17]) Let $\sigma$ be an orthogonal representation of a finite or compact group $G$ on a real Hilbert space $H$, and let $f: H \rightarrow R$ be such that for $\forall \sigma \in G, f(\sigma \cdot x)=f(x)$. Set $H^{G}=\{x \in H: \sigma \cdot x=x, \forall \sigma \in G\}$. Then the critical point off in $H^{G}$ is also a critical point off in $H$.

Lemma 2.3 [5] If $X$ is a reflexive Banach space, $M$ is a weakly closed subset of $X$, and $f: M \rightarrow R \cup\{+\infty\}, f \not \equiv+\infty$ is weakly lower semi-continuous and coercive, then $f$ attains its infimum on $M$.

## Figure $4 q \in \Lambda_{4}$.



Lemma 2.4 (Poincare-Wirtinger inequality) Let $q \in W^{1,2}\left(\mathbb{R} / \mathbb{Z} \mathbb{T}, \mathbb{R}^{d}\right)$ and $\int_{0}^{T} q(t) d t=0$, then $\int_{0}^{T}|\dot{q}(t)|^{2} d t \geq\left(\frac{2 \pi}{T}\right)^{2} \int_{0}^{T} q(t)^{2} d t$. And the inequality takes the equality if and only if $q(t)=\alpha \cos \frac{2 \pi}{T} t+\beta \sin \frac{2 \pi}{T} t, \alpha, \beta \in R^{d}$.

We now prove Theorem 1.1.
Proof By the symmetry of $\Lambda_{i}$, we know for $\forall x \in \Lambda_{i}$,

$$
\begin{equation*}
\int_{0}^{T} q(t) d t=0 \tag{2.1}
\end{equation*}
$$

If $q_{n}(t) \rightharpoonup q(t)$ in $\bar{\Lambda}_{i}$, then by Sobolev's compact embedding theorem, we have $q_{n}(t) \rightarrow q(t)$ in $C[0,1]$.
(i) If $q(t) \in \Lambda_{i}$, then $\lim _{n \rightarrow+\infty} \int_{0}^{1} U\left(q_{n}(t)\right) d t=\int_{0}^{1} U\left(q_{n}(t)\right) d t$. Since $\int_{0}^{1} q_{n} d t=0$, $\frac{1}{2} \int_{0}^{1}\left|\dot{q}_{n}\right|^{2} d t$ can be regarded as the square of an equivalent norm for $W^{1,2}$, so it is weakly lower semi-continuous, so $\underline{\lim f}\left(q_{n}(t)\right) \geq f(q)$.
(ii) If $q(t) \in \partial \Lambda_{i}$, then by Lemma 2.1, $f(q)=+\infty$, we have $\int_{0}^{1} U\left(q_{n}(t)\right) d t \rightarrow+\infty$. So, $\varliminf_{n \rightarrow+\infty} f\left(q_{n}\right)=+\infty \geq f(q)$. Hence $f$ is w.l.s.c.
Using (2.1), we know that $f(q)$ is coercive on $\bar{\Lambda}_{i}$. Lemma 2.3 guarantees that $f(q)$ attains its infimum on $\bar{\Lambda}_{i}$. Let the minimizer be $\tilde{q}$, then

$$
\begin{equation*}
f(\widetilde{q})=\inf _{q \in \bar{\Lambda}_{i}} f(q)<+\infty \tag{2.2}
\end{equation*}
$$

If $\widetilde{q}$ is a collision periodic solution, then there exist $t_{0} \in[0,1]$ and $j \in\{1,2, \ldots, N, N+1\}$ such that $\tilde{q}\left(t_{0}\right)=q_{j}$. Let $x(t)=\widetilde{q}(t)-q_{j}$ and note $x\left(t_{0}\right)=0$. By Lemma 2.1, we have

$$
\begin{align*}
f(\widetilde{q}) & =\int_{0}^{1}\left[\frac{1}{2}|\dot{\tilde{q}}(t)|^{2}+\frac{m_{j}}{\left|\widetilde{q}(t)-q_{j}\right|^{\alpha}}+\sum_{i \neq j}^{N+1} \frac{m_{i}}{\left|\widetilde{q}(t)-q_{i}\right|^{\alpha}}\right] d t \\
& \geq \int_{0}^{1}\left[\frac{1}{2}|\dot{x}(t)|^{2}+\frac{m_{j}}{|x(t)|^{\alpha}}\right] d t=+\infty, \tag{2.3}
\end{align*}
$$

which contradicts the inequality in (2.2). By Lemma $2.2, \widetilde{q}(t)$ is the critical point of $f$ in $W^{1,2}\left(\mathbb{R} / \mathbb{Z}, \mathbb{R}^{2}\right)$; therefore, $\widetilde{q}(t)$ is a non-collision periodic solution.

This completes the proof of Theorem 1.1.

## Competing interests

The authors declare that they have no competing interests.

## Authors' contributions

The authors declare that the study was realized in collaboration with the same responsibility. All authors read, checked and approved the final manuscript.

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