J Mater Sci (2013) 48:7392–7398 DOI 10.1007/s10853-013-7554-8

Template-based synthesis and discontinuous hysteresis loops of cobalt nanotube arrays

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Received: 6 April 2013/Accepted: 25 June 2013/Published online: 4 July 2013 © The Author(s) 2013. This article is published with open access at Springerlink.com

Abstract Co nanotubes were successfully prepared in the pores of anodic aluminium oxide templates using a DC electrodeposition method. The dependence of the product morphology on the applied potential was studied over the range -1.0 to -3.0 V. The results showed that the wall thickness of the nanotubes became thinner as the applied potential was reduced, and that there existed a critical potential (v_c) related to the electrodeposition parameters, below or above which the electrodeposition process is dominated by kinetics and thermodynamics, respectively. The formation of nanotubes is the result of kinetics dominating the electrodeposition process. Magnetic measurements showed that the hysteresis loops of the nanotubes were discontinuous. Theoretical analysis suggested the existence of some peculiar magnetization configuration such as a vertical Bloch line, which was responsible for the discontinuity of the hysteresis loops.

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Introduction

An exciting research subject in material science is the fabrication and investigation of nanostructures such as nanowires and nanotubes because of their unique properties and potential applications in various areas, such as catalysis [1], sensing [2], field emission [3], hybrid solar cells and electronic devices [4, 5]. Ferromagnetic nanotubes have drawn special attention in recent years because their magnetic anisotropy and magnetic reversal mechanism are much different from those of nanowires [6-10]. Ferromagnetic Co nanotubes have single axis magnetic anisotropy, which makes them a promising material for future applications in magnetic sensors and ultrahigh density magnetic recording. The synthesis of Co nanotubes has focused on anodic aluminium oxide (AAO) template-based electrodeposition because it is not only a convenient inexpensive technique, but also modification of the AAO template is not necessary for the fabrication of nanotubes, which eliminates the introduction of impurities in the nanotubes [11]. To obtain nanotubes, the common method used in template-based electrodeposition is to control the applied current densities to a level higher than 20 mA/cm². However, this often results in a lower filling rate and inhomogeneous tube length [12, 13]. In this paper, we report the fabrication of Co nanotube arrays embedded in AAO templates with a high filling rate and uniform growth. The dependence of Co morphology on the applied potential has also been studied. A critical potential (v_c) for obtaining Co nanotube by template-based electrodeposition method is represented. The experimental results show that the formation of Co nanotubes is the result of kinetics dominating the electrodeposition process. Magnetic measurements were carried out using a physical property measurement system (PPMS) at room temperature. The

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results showed that hysteresis loops of the Co nanotubes are discontinuous. A possible reason is discussed. This work will benefit understanding of the electrochemical growth mechanism of metal nanotubes as well as the magnetization reversal of ferromagnetic nanotubes.

Experimental

We fabricated anodic oxide templates as follows: in order to eliminate internal stress, a high-purity (99.999 %) aluminium (Al) foil was thermally treated at 400 °C for 2 h and then brought to room temperature. The whole process is in an argon atmosphere to prevent the Al from oxidation. Next, the heat-treated Al foil was electropolished in a mixed solution of alcohol and perchloric acid (ratio by volume 4:1) for 5 min to smooth the surface. To obtain highly ordered pores, a two-step anodization was used. In the first step, the Al foil was anodized at 5 °C and 50 V DC in 0.3 M oxalic acid for about 8 h. The resulting aluminium oxide layer was then removed by immersing the anodized Al in a mixed solution of 0.4 M chromic acid and 0.6 M phosphoric acid (ratio by volume 1:1) at room temperature. Subsequently, the samples were reanodized under the same anodization conditions as in the first step. This resulted in the formation of AAO nanopores with diameters of 50 nm. After the second anodization, the remaining Al was removed using saturated calcium chloride. The AAO film was then etched in a 0.6 M phosphoric acid solution to remove the continuous barrier layer present at the bottom of each nanopore and enlarge the nanopore diameters. Finally, an AAO film with a thickness of about 70 µm and pore diameters of about 70 nm was obtained. Such AAO templates were used as described below, to fabricate Co nanotubes by electrochemical deposition.

A Cu conductive layer was sputtered onto one side of the AAO template to provide an electrical contact. Electrodeposition was performed in a three-electrode cell under constant voltage at room temperature, where the sputtered conducting layer served as the working electrode. A saturated calomel electrode (SCE) was used as the reference, and a graphite pole was used as the counter electrode. The electrolyte used to prepare the Co nanotubes was prepared from reagent grade chemicals and deionized water. It consisted of 0.1 M CoSO₄·7H₂O and 0.1 M H₃BO₃. The pH of the electrolyte was 2.5. The electrodeposition was carried out for 20 min, at room temperature (25 °C) with different DC voltages ranging from -1.0 to -3.0 V. The total charge current density at any applied potential is less than 1 mA/cm². As a result, nanostructures with a length about 6 µm were obtained.

The morphology of the Co nanotubes was characterized using a field-emission scanning electron microscope (SEM, Hitachi S-4800) and a transmission electron microscope (TEM, Hitachi H-7650). The crystallographic structure and composition of the nanostructures were examined using an X-ray diffractometer (XRD) with Cu K_{α} radiation and using X-ray energy dispersive spectroscopy (EDS). The magnetic properties were measured at room temperature using a physical properties measurement system (PPMS-6700) with the magnetic field parallel and perpendicular to the nanostructures. Except for the TEM sample whose AAO template was completely removed by alkaline treatment, all samples used in the measurements reported here were partially embedded in the AAO.

Results and discussion

Figure 1a, b shows the SEM images of an AAO template before and after sputtering the Cu layer. We can see from Fig. 1a that the AAO pores had an average diameter of about 70 nm. To serve efficiently as a cathode, a Cu layer was first sputtered onto one side of the AAO. The sputtering power and time were 2 kW and 20 min, respectively. It can be seen from Fig. 1b that almost all the AAO pores were covered with Cu. Our Co nanotubes and nanowires were fabricated using such an AAO template coated with Cu. Figure 1c, d are the TEM image and EDS spectrum of a typical as-prepared nanotubes, respectively. It may be seen that the diameter of the Co nanotubes is 70 nm which is consistent with the AAO template structure, as shown in Fig. 1a. The EDS spectrum of the Co nanotubes confirms that the as-grown nanotubes consist of pure cobalt. The observed O, Al, and C peaks are generated from the alumina film and the conductive tape.

However, several groups reported that the electrodeposition product consisted of only nanowires when using an AAO template with its pores blocked [14–17]. This is inconsistent with our experiment results. We attributed it to differences in the electrodeposition parameters. The different results also imply that it is possible for a more universal growth mechanism to explain template-based electrochemical deposition of metal nanotubes. In view of this point, we have studied in detail the dependence of the Co morphology on the applied potential.

SEM images of the Co nanostructures are shown in Figs. 1c and 2a–e. The images demonstrate that our Co nanotubes have a large aspect ratio and they take on a uniform growth. The wall width of the Co nanotubes is in the range 15–18 nm. The filling rate of the nanostructures is nearly 100 %. The SEM images also demonstrate that either nanowires or nanotubes can be successfully deposited by adjusting the deposition potential between -1.0 and -3.0 V. The figures indicate that there exists a critical potential, $v_c \approx -1.5$ V. The deposit is in the form of a

Fig. 1 a Front side SEM images of AAO, b backside SEM images of AAO with Cu sputtering for 20 min, c TEM images of single Co nanotubes formed with applied potential -1.3 V, d EDS spectrum of a typical as-prepared nanostructure



nanotube when $v \le v_c$. Otherwise, when $v > v_c$, nanowires are deposited. We also found in our experiment that v_c may vary with the electrodeposition parameters (such as temperature, concentration and constituent) in a given system. It therefore appears that v_c can be defined as a critical state, below or above which the electrodeposition process is dominated by two distinct factors. According to the Reaction Principles of Physical Chemistry [18], a reaction process is determined by the two factors of kinetics and thermodynamics. Once the reaction takes place, the crystal develops in a manner determined by ambient condition. Our XRD results indicated that the Co nanotubes or nanowires all have a similar growth habit, respectively. The nanotubes are oriented growth along kinetics crystallographic plane, but the nanowires are oriented growth along thermodynamics crystallographic plane. Therefore we assumed that v_c is a critical state separating the deposition process into two sections dominated by kinetics or thermodynamics, respectively.

The XRD patterns of typical Co nanostructures (shown in Fig. 3) demonstrate that both Co nantubes and nanowires have an hcp (PDF00-001-1278) structure, but they have different crystal orientations. The sharp diffraction peaks at 75.764° and 44.600° correspond, respectively, to the kinetics crystallographic plane (110) and the thermodynamics crystallographic plane (002). This indicates that the kinetics crystallographic plane (110) can fully develop

when $v \leq v_c$, and that the thermodynamics crystallographic plane (002) can fully develop when $v > v_c$. These can be attributed to the different arrangement of atom under different conditions. A foreign atom has a preferential rank among thermodynamics crystallographic planes (002) when $v > v_c$, because the energy of the atom is so low that it has enough time to perform close packing among atoms. But when $v \leq v_c$, the foreign atom has high energy, and the relaxation time is short. Therefore it preferentially ranks among kinetic crystallographic plane (110) which is not a closely packed crystallographic plane. Moreover, Fig. 2c-f clearly shows that the wall thickness of the Co nanotubes becomes increasingly thin as the applied negative potential decreases. This is due to the fact that the growth rate of the kinetics crystallographic plane (110) becomes much faster than that of the thermodynamics crystallographic plane (002) with an increase in the kinetics factor in electrodeposition process. These results suggest that Co nanotubes and nanowires are formed under conditions where kinetics or thermodynamics dominate the electrodeposition process, respectively. This is in good agreement with Ref. [19]. Thus, we can conclude that the electrodeposition process is dominated by kinetics or thermodynamics when $v \leq v_c$ and $v > v_{\rm c}$, respectively.

The magnetic properties of Co nanostructures were studied at room temperature using a PPMS with the magnetic fields applied perpendicular (\perp) and parallel (||) to the



Fig. 2 SEM images of Co nanostructures: **a**, **b** nanowires formed with applied potentials of -1.0 and -1.2 V, respectively; **c**-**f** nanotubes formed with applied potentials of -1.25, -1.5, -2.0 and -3.0 V, respectively



Fig. 3 XRD patterns of typical Co nanostructures, **a** nanotubes formed with an applied potential of -2.0 V, **b** nanowires formed with an applied potential of -1.0 V

long axis of the nanostructures. The magnetic parameters of the Co nanotubes with different applied potentials are shown in Table 1. These data show that the remanent squareness (SQ) and coercivities (H_c) parallel to the direction of the presented nanotubes are large compared to values for the perpendicular direction, which shows that the easy axis of the Co nanotubes is parallel to the tube axis. This conclusion also follows from the magnetic hysteresis curves (shown in Fig. 4a, b), and can be attributed to the strong shape anisotropy of the nanostructures. In addition, Table 1 also shows that the value for H_c first increases up to -1.25 V and then decreases. We attributed the former to the structure difference between nanowires and nanotubes, and the latter to the variation of Co particles stacking density.

Table 1 Magnetic parameters, coercivity (H_c) and remanent squareness (SQ) of Co nanotubes and nanowires embedded in AAO templates with the field applied parallel (||) and perpendicular (\perp) to the tube long axis

Potential (V)	SQ^{\perp} (emu)	SQ^{\parallel} (emu)	$\mathrm{H}_{\mathrm{c}}^{\perp}$ (Oe)	H_c^{\parallel} (Oe)
-1.0	0.00025	0.001267	175	691
-1.2	0.000276	0.001424	192	710
-1.25	0.00038	0.00285	210	760
-1.3	0.000394	0.00297	203	726
-1.5	0.000362	0.00215	180	690
-2.0	0.0006	0.00283	170	749
-3.0	0.0011	0.0023	135	563

The hysteresis loops investigated for Co nanotubes grown with different applied potentials indicated that they have similar characteristics: (1) discontinuous intervals appear in the curves, (2) the values of ΔH (defined below) in each of the hysteresis loops for the nanotubes are similar. These characteristics are interesting because the hysteresis loops of Co nanowires do not have such features. The hysteresis loops of typical samples are shown in Fig. 4, where the above differences for Co nanotubes and nanowires may be clearly seen. The results imply that a reversal processes takes place in Co nanotubes that is distinct from what occurs in Co nanowire. The inset of Fig. 4d shows an enlargement of the magnetic moment dependence on the field in the discontinuous interval from Hs₁ (1493 Oe) to Hs_2 (1676 Oe). It can be seen a field range exists, $\Delta H = H_{s2} - H_{s1} = 183$ Oe, over which the magnetic moment remains constant during the reversal processes. This suggests that there exists some peculiar magnetic Fig. 4 Hysteresis loops of Co nanostructures formed with an applied potential of $\mathbf{a} - 1.0 \text{ V}$ and $\mathbf{b} - 2.0 \text{ V}$, measured at room temperature, \mathbf{c} and \mathbf{d} the partial amplification of \mathbf{a} and \mathbf{b} with the applied field perpendicular to the nanostructures. *Inset* enlargement of \mathbf{d} over the interval from 1493 to 1676 Oe

Fig. 5 Schematic diagram of

both CW and CCW

vortex wall configurations in Co nanotubes: **a** CW alone existence, **b** CCW alone existence and **c** coexistence of





one vortex wall. In the latter situation, a peculiar configu-

ration arises in the regions where CW and CCW vortices

configuration that remains stable over a certain field range in the Co nanotubes. Once this configuration is formed, it is hard to magnetize, which leads to the hysteresis loops being discontinuous in the corresponding field range.

The existence of vortex configurations in magnetic nanotubes with a high aspect ratio has been reported in recent years [20–22]. When the applied field is perpendicular to the long axis of the Co nanotubes, the distribution of magnetic moments takes on the vortex state and corresponds to the curling mode [23]. The magnetization configuration of the vortex wall in Co nanotubes is presented schematically in Fig. 5. As shown in Fig. 5a–c, the vortex state in Co nanotubes has two possible chiralities: clockwise (CW) and counterclockwise (CCW). The CW and CCW vortex states may exist separately or coexist in meet or detach. Since the configuration is similar to that of the Bloch line that has been investigated previously [24– 26], we refer below to this configuration as a Bloch line. With reference to magnetic bubble material, Malozemoff et al reported that the plane field needed for a single vertical Bloch line magnetization reversal is eight times of saturation moment (M_s) if the external in-plane field exerted on magnetic bubble material is parallel, but oppositely directed, to the vertical Bloch line (shown in Fig. 6) [27]. The value of M_s is given by the expression:

$$M_{\rm s}=\frac{M}{\pi(r_1^2-r_2^2)hn}$$



Fig. 6 Schematic diagram of the location relationship between the in-plane field and the vertical Bloch line

where *h* and *n* are the height and number of the nanotubes per square millimetre; r_1 and r_2 are the inner and outer radius of the nanotubes; M_s is saturation moment. We obtain the value $M_s \approx 214 \text{ emu/cm}^3$ for our Co nanotubes formed with an applied potential of -2.0 V. In addition, $H_{s2} \approx 1676$ Oe. Therefore Fig. 4d, from $8M_s =$ 1712 \approx H_{s2}. This is in agreement with the theory of the Bloch line mentioned above. This suggests that the peculiar configuration identified above in Co nanotubes may be a vertical Bloch line (shown in Fig. 5c). Of course, the relationship between 8Ms and Hs2 may vary with the location of the vertical Bloch line in the Co nanotubes when the applied field is perpendicular (\bot) to the long axis of the nanostructures. Nevertheless, the reversal process of these vertical Bloch lines is similar, which accounts for the discontinuous interval appearing at low applied field. During the reversal process, when $H < H_{s1}$ or $H > H_{s2}$ (Fig. 4d, inset), the CCW or CW vortices are immediately broken up as the external field increases, so that the hysteresis loops are continuous over these ranges. Breaking up of the vertical Bloch line is, however, very hard, and the Bloch line can remain stable in the reversal field range, $H_{s1} < H < H_{s2}$. Once H increases to H_{s2} , the vertical Bloch line instantly disappears. This give rise to the discontinuity in the hysteresis loops for the Co nanotubes.

Conclusion

Co nanotubes were fabricated using AAO templates by a DC electrodeposition method. According to the dependence of the morphology of the Co nanostructures on the applied potential, we determined a critical potential v_c for fabrication of Co nanostructures through DC electrodeposition. We defined v_c as a critical state separating the deposition process into two sections dominated by kinetics or thermodynamics, respectively. When the applied negative potential was equal to or below v_c , the deposition process is dominated by kinetics. To the contrary, when the applied negative potential was above v_c , the deposition process is dominated by thermodynamics. The experimental results demonstrated that Co nanotubes are the

products of kinetics dominating the electrodeposition process, and that the wall thickness of the Co nanotubes became thinner as the applied potential was reduced. Magnetic measurements showed that the hysteresis loops of the Co nanotubes were discontinuous. Calculations indicate that the in-plane reversal field needed to break up the peculiar configuration responsible for the discontinuity is about eight times the saturation moment, which suggests that the existence of a vertical Bloch line is possible in Co nanotubes.

Acknowledgements This work is supported by the Natural Science Foundation of Hebei Province (Grant No. A2012205038).

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