

RESEARCH

Open Access

High order Riesz transforms and mean value formula for generalized translate operator

Ismail Ekincioglu^{1*}, Hasan H Sayan² and Cansu Keskin¹

*Correspondence:

ismail.ekincioglu@dpu.edu.tr

¹Department of Mathematics,
Dumlupinar University, Kütahya,
TurkeyFull list of author information is
available at the end of the article**Abstract**

In this paper, the mean value formula depends on the Bessel-generalized shift operator corresponding to the solutions of the boundary value problem related to the multidimensional Bessel operator are studied. In addition, Riesz transforms R_B related to the multidimensional Bessel operators are studied. Since a Bessel-generalized shift operator is a translation operator corresponding to the multidimensional Bessel operator, we construct a family of R_B by using a Bessel-generalized shift operator. Finally, we analyze weighted inequalities involving R_B .

MSC: Primary 31B05; 35G10; secondary 34B30; 47F05

Keywords: Bessel equation; generalized shift operator; mean value theorem

1 Introduction

Singular integral operators are playing an important role in harmonic analysis, the theory of functions and partial differential equations. Singular integrals are associated with the Δ_{B_n} Laplace-Bessel differential operator, which is known as an important operator in analysis and its applications; these have been the research areas of many mathematicians such as Muckenhoupt and Stein [1, 2], Kipriyanov and Klyuchantsev [3, 4], Aliev and Gadjiev [5], Guliev [6], Gadjiev. Also, singular integral operators related to the generalized shift operator were studied by [7] and [8–10] and others.

The Bessel-generalized shift operator is one of the most important generalized shift operator on the half line \mathbb{R}_+ , [11, 12]. The Bessel-generalized translation is used, while studying various problems connected with Bessel operators [13]. Fourier-Bessel harmonic analysis, *i.e.* the part of harmonic analysis addressing various problems on Bessel (Hankel) integral transforms, is closely connected with Bessel-generalized shift operator.

It is well known that the fundamental solutions of the classical stationary of mathematical physics (the harmonic equation, polyharmonic equation, and Helmholtz equation) are radial functions. Therefore, it is natural to seek these solutions as solutions of ordinary differential equations. However, since the spherical coordinate transformation transforms an equation with the Laplace operator in \mathbb{R}^n into an ordinary differential equation with the singular Bessel differential operator, interest arose (probably, a long time ago) in studying methods for constructing the fundamental solutions of singular ordinary differential equations with the Bessel operator in place of the second derivative. In this connection, it might be very useful to prove a theorem on the fundamental solution of an ordinary differential equation involving the Bessel operator with constant coefficients, similar to the well-known theorem on the fundamental solution of an ordinary differential equation.

The result obtained on the fundamental solution of an ordinary differential operator with the Bessel operator has allowed us to analyze equations with the singular differential operator $\Delta_{B_\gamma} = \sum_{i=1}^n B_{\gamma_i} + \sum_{i=n+1}^N \frac{\partial^2}{\partial x_i^2}$ where the different indices γ_i act with respect to part of the variables and these indices may take negative values. The latter fact is essential because the tools used in problems of this kind (the Poisson operator and generalized shift operator of integral nature) are defined only for $\gamma_i > 0$. To find the fundamental solution with a singularity at an arbitrary point, we use the generalized shift operator that acts with respect to the radial variable. Note also that the mixed-type generalized shift operator, which is conventionally used in such problems, coincides with the radial shift operator on radial functions provided that $\gamma_i > 0$. The fundamental solutions (of the B -harmonic and B -polyharmonic equations and of the singular Helmholtz equation) found coincide with the known solutions when $\gamma_i > 0$ and with the classical solutions when $\gamma_i = 0$.

It is well known that harmonic functions satisfy various mean value theorems, which may be considered as generalizations of the Gauss mean value theorem. There have been a number of studies on mean value theorems. Cheng [14] obtained a converse for a different mean value expression. Nicolesco [15] gave an expression in terms of certain iterated means and showed that a converse was also true. The mean value theorems for harmonic functions have also been studied by Pizetti [16], Picone [17], Ekincioglu [8], and Kipriyanova [18] and [19].

The solutions of the boundary value problems for Laplace operator are related to the ordinary shift operator. Also, the solutions of boundary value problems for Laplace-Bessel and Bessel operators are corresponding to the generalized shift operator and Bessel-generalized shift operator, respectively.

In this paper, singular integral operators generated by a Bessel-generalized shift operator are studied. In addition, the mean value formulas related to the Bessel-generalized shift operator are given.

Riesz-Bessel singular integral operators related to generalized shift operator for Laplace-Bessel operator were showed in [7] and [9]. The authors used the mean value theorem related to the generalized shift operator. The mean value formula for the equations $\Delta_{B_{x_n}} u = 0$ and $\Delta_{B_{x_n}} u + \lambda u = 0$ were obtained by [8] and [19], respectively, where $\Delta_{B_{x_n}} = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} + \frac{2\gamma_n}{x_n} \frac{\partial}{\partial x_n}$.

In this study, we introduce the mean value formula for the equation $Bu = 0$ and the high order Riesz-Bessel transform associated with the Bessel-generalized shift operator for the B Bessel differential operator

$$B = \sum_{i=1}^n B_i, \quad B_i = \frac{\partial^2}{\partial x_i^2} + \frac{2\gamma_i}{x_i} \frac{\partial}{\partial x_i}, \quad \gamma_1 > 0, \dots, \gamma_n > 0.$$

Let \mathbb{R}^n be n -dimensional Euclidean space and $x = (x_1, \dots, x_n)$, $\xi = (\xi_1, \dots, \xi_n)$ be vectors in \mathbb{R}^n , then $x \cdot \xi = x_1 \xi_1 + \dots + x_n \xi_n$, $|x| = (x \cdot x)^{1/2}$. Denote $\mathbb{R}_+^n = \{x \in \mathbb{R}^n : x_1 > 0, \dots, x_n > 0\}$, $S_+^n = \{x \in \mathbb{R}_+^n : |x| = 1\}$, $\gamma = (\gamma_1, \dots, \gamma_n)$, $\gamma_1 > 0, \dots, \gamma_n > 0$, $|\gamma| = \gamma_1 + \dots + \gamma_n$ and $d\mu_\gamma(x) = \prod_{i=1}^n x_i^{2\gamma_i} dx$. We shall denote by $L_{p,\gamma}(\mathbb{R}^n, d\mu_\gamma(x))$ -spaces (the Lebesgue space with respect to the measure μ_γ), the set of all measurable functions f on \mathbb{R}_+^n such that the norm

$$\|f\|_{L_{p,\gamma}} \equiv \|f\|_{p,\gamma} = \left(\int_{\mathbb{R}_+^n} |f(x)|^p d\mu_\gamma(x) \right)^{1/p}, \quad 1 \leq p < \infty,$$

is finite. We determine the Bessel-generalized translation $T^\gamma \varphi(x) = u(x)$, $x, y \in \mathbb{R}_+^n$ of a function $\varphi(x) \in C^{(2)}(\mathbb{R}_+^n)$ as the solution to the following initial value problem:

$$\left. \begin{aligned} B_x u(x, y) &= B_y u(x, y), \\ u(x, 0) &= \varphi(x), \quad u_y(x, 0) = 0, \end{aligned} \right\} \tag{1.1}$$

where B is the multidimensional Bessel differential operator [20]. The solution of the initial value problem (1.1) exists, is unique, and can be written explicitly as

$$\begin{aligned} T^\gamma \varphi(x) &:= c_\gamma \int_0^\pi \cdots \int_0^\pi \varphi((x_1, y_1)_{\alpha_1}, \dots, (x_n, y_n)_{\alpha_n}) \\ &\quad \times \left(\prod_{i=1}^n \sin^{2\gamma_i-1} \alpha_i \right) d\alpha_1 \cdots d\alpha_n, \end{aligned} \tag{1.2}$$

where $c_\gamma = \prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})[\Gamma(\frac{1}{2})\Gamma(\gamma_i)]^{-1}$ and $(x_i, y_i)_{\alpha_i} = \sqrt{x_i^2 + y_i^2 - 2x_i y_i \cos \alpha_i}$, $1 \leq i \leq n$. By (1.2), the operator T^γ can be extended to all functions $L_{p,\gamma}(\mathbb{R}_+^n)$. The operator T^γ satisfying (1.1) may be regarded as a Bessel-generalized shift operator Bessel-generalized shift operator (see [4, 11] and [21]). We remark that this shift operator is closely connected with the Bessel differential operator. The convolution operator determined by T^γ is as follows:

$$(f * \varphi)(x) = \int_{\mathbb{R}_+^n} f(y) T^\gamma \varphi(x) d\mu_\gamma(y). \tag{1.3}$$

The convolution (1.3) is known as a B -convolution. We note some properties for the B -convolution and the Bessel-generalized shift operator:

- If $f(x), \varphi(x) \in C(\mathbb{R}_+^n)$, $\varphi(x)$ is a bounded function, $x > 0$, and

$$\int_0^\infty |f(x)| d\mu_\gamma(x) < \infty,$$

then

$$\int_{\mathbb{R}_+^n} T^\gamma f(x) \varphi(y) d\mu_\gamma(y) = \int_{\mathbb{R}_+^n} f(y) T^\gamma \varphi(x) d\mu_\gamma(y).$$

- From the above result, we have the following equality for $\varphi(x) = 1$:

$$\int_{\mathbb{R}_+^n} T^\gamma f(x) d\mu_\gamma(y) = \int_{\mathbb{R}_+^n} f(y) d\mu_\gamma(y).$$

- $(f * \varphi)(x) = (\varphi * f)(x)$.

The Fourier-Bessel transform is defined and invertible on functions $\varphi \in \mathcal{S}(\mathbb{R}_+^n)$,

$$[F_B \varphi](y) = c_\gamma \int_{\mathbb{R}_+^n} \varphi(x) \prod_{i=1}^n j_{\gamma_i-\frac{1}{2}}(x_i y_i) d\mu_\gamma(x),$$

$$[F_B^{-1} \varphi](x) = \int_{\mathbb{R}_+^n} \varphi(y) \prod_{i=1}^n j_{\gamma_i-\frac{1}{2}}(x_i y_i) d\mu_\gamma(y),$$

where $c_\gamma = \prod_{i=1}^n [2^{\gamma_i - \frac{1}{2}} \Gamma(\gamma_i + \frac{1}{2})]^{-1}$ and $j_{\gamma_i - \frac{1}{2}}$ is the normalized Bessel function related to the Bessel function of the first kind by the formula $j_\gamma(r) = 2^\gamma \Gamma(\gamma + 1) J_\gamma(r) r^{-\gamma}$ [11]. However, the following equality for a Fourier-Bessel transformation is true:

$$F_B(f * \varphi)(x) = F_B f(x) F_B \varphi(x)$$

(see [22–28]).

2 The high order Riesz-Bessel transforms R_B associated with Bessel-generalized shift operator

In this section, we consider a Bessel-generalized shift operator related to the multidimensional Bessel differential operator. Then we give the Fourier-Bessel transformation of a homogeneous polynomial which obeys the Bessel equations. Finally, we define the high order Riesz-Bessel transforms related to Bessel-generalized shift operator and so we show that high order Riesz-Bessel transforms obey the condition of classical Riesz transforms, that is, these operators extend to high order Riesz-Bessel transforms [29].

It follows from the general theory of singular integrals that Riesz transforms are bounded on $L_{p,\gamma}(\mathbb{R}^n, d\mu_\gamma(x))$ for all $1 < p < \infty$. In this paper we extend this result to the context of Bessel theory where a similar operator is already defined. It has been noted that the difficulty arises in the application of the classical L_p -theory of Calderon-Zygmund, since Riesz transforms are singular integral operators. In this paper we describe how this theory can be adapted in a Bessel setting and give an $L_{p,\gamma}$ -result for high order Riesz transforms for all $1 < p < \infty$.

We have $\Omega(y) = P_k(y)|y|^{-k}$, $K(y) = \Omega(y)|x|^{-n-2|\gamma|}$, the P_k range over the homogeneous harmonic polynomials the latter arise in special case $k = 1$. Those for $k > 1$, we call the high order Riesz-Bessel transform where we refer to k as the degree of the high order Riesz Bessel transform [29]. They can also be characterized by their invariance properties.

Let P_k be homogeneous polynomial of degree k in \mathbb{R}_+^n . We shall say that P is elliptic if $P(x)$ vanishes only at the origin. For any polynomial P we consider also its corresponding differential polynomial. Thus if $P(x) = \sum a_\alpha x^\alpha$ we write $P(\frac{\partial}{\partial x}) = \sum a_\alpha (\frac{\partial}{\partial x})^\alpha$, where $(\frac{\partial}{\partial x})^\alpha = (\frac{\partial}{\partial x_1})^{\alpha_1} \dots (\frac{\partial}{\partial x_n})^{\alpha_n}$ and with the monomials $x^\alpha = x_1^{\alpha_1} \dots x_n^{\alpha_n}$ (see [30]).

Theorem 2.1 *Suppose that $P_k(x)$ is a homogeneous polynomial of degree k and satisfies for Bessel operator $B[P_k(x)] = 0$ then we have*

$$F_B[P_k(x)e^{-|x|^2}](y) = 2^{-(|\gamma|+k+\frac{n}{2})} i^k P_k(y) e^{-\frac{|y|^2}{4}}.$$

Lemma 2.2 *Let $\theta = (\theta_1, \theta_2, \dots, \theta_n)$. Suppose that*

$$\int_{S_+^n} f(\theta) d\mu_\gamma(\theta) dS = 0$$

and φ is of Schwartz class $\mathcal{S}(\mathbb{R}_+^n)$ then we have the identity

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^n} \frac{f(\frac{x}{|x|})}{|x|^{n+2|\gamma|-\varepsilon}} \varphi(x) d\mu_\gamma(x) = \lim_{\varepsilon \rightarrow 0} \int_{|x|>\varepsilon} \frac{f(\frac{x}{|x|})}{|x|^{n+2|\gamma|}} \varphi(x) d\mu_\gamma(x).$$

Proof The proof follows immediately from the representation

$$\int_{\mathbb{R}_+^n} \frac{f(\frac{x}{|x|})}{|x|^{n+2|\gamma|-\varepsilon}} \varphi(x) d\mu_\gamma(x) = \int_{|x| \leq 1} \frac{f(\frac{x}{|x|})}{|x|^{n+2|\gamma|-\varepsilon}} [\varphi(x) - \varphi(0)] d\mu_\gamma(x) + \int_{|x| > 1} \frac{f(\frac{x}{|x|})}{|x|^{n+2|\gamma|-\varepsilon}} \varphi(x) d\mu_\gamma(x). \quad \square$$

The mean value theorem for multidimensional Bessel differential operators is very convenient for obtaining multidimensional singular integral operators generated by a Bessel-generalized shift operator. Therefore, we studied the mean value formula related to the Bessel-generalized shift operator for the solutions of the boundary value problem for the multidimensional Bessel operator $Bu = 0$.

The Bessel-generalized shift operator is one of the most important generalized shift operator on the half line $\mathbb{R}_+ = [0, +\infty)$ [11–13]. The Bessel-generalized translation is used while studying various problems connected with Bessel operators. Fourier-Bessel harmonic analysis, *i.e.* the part of harmonic analysis addressing various problems on Bessel (Hankel) integral transforms, is closely connected with the Bessel-generalized shift operator.

3 The mean value formula

In this section, we determine the mean value formula for the Bessel-generalized shift operator. Let \mathbb{R}^n be an n -dimensional Euclidian space and $x = (x_1, x_2, \dots, x_n)$, $\xi = (\xi_1, \xi_2, \dots, \xi_n)$ be vectors in \mathbb{R}^n . Then $x \cdot \xi = x_1\xi_1 + \dots + x_n\xi_n$. Denote $\mathbb{R}_+^n = \{x \in \mathbb{R}^n : x_i > 0, 1 \leq i \leq n\}$ and $S_+^{n-1} = \{x \in \mathbb{R}_+^n : |x| = 1\}$, $|\gamma| = \gamma_1 + \dots + \gamma_n$. We assume that $D^+ \subset \mathbb{R}_+^n$ and Ω^+ is its boundary. In this paper, we are mainly concerned with the mean value theorem. Now, we relate this concept in the following theorem.

Theorem 3.1 *Let S_+^{n-1} be a unit sphere centered at the origin, contained in \mathbb{R}_+^n and u be an even regular solution with respect to x_1, \dots, x_n of $Bu = 0$. Then the following formula holds:*

$$\int_{S_+^{n-1}} u(r\theta_1, \dots, r\theta_n) \prod_{i=1}^n \theta_i^{2\gamma_i} dS = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0). \quad (3.1)$$

Let T^γ be the multidimensional Bessel-generalized shift operator and u be an even regular solution of $Bu = 0$ at any interior point of the region $D^+ \subseteq \mathbb{R}_+^n$. Also, the following formula is valid:

$$\int_{S_+^{n-1}} T^\gamma u(r\theta_1, r\theta_2, \dots, r\theta_n) \prod_{i=1}^n \theta_i^{2\gamma_i} dS = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(y). \quad (3.2)$$

Proof We assume that the continuity of u and v in the closed region $\Omega^+ \cup D^+$, continuity of the first and second derivatives of u and v in D^+ , together with continuity of the first derivatives of u in $\Omega^+ \cup D^+$ and the second derivatives of u in D^+ . In addition, the existence of the integrals over D^+ are assumed in Green's formula. The most important tool for the potential theory is provided by this formula in n -dimensional bounded region D^+ with volume element $dg = dx_1 \cdots dx_n$ and its boundary Ω^+ , which we assume to be piecewise

smooth; the two functions u and v are related by Green's formula,

$$\int_{\Omega^+} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega = \int \int_{D^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dg. \tag{3.3}$$

Let $P' = (0, \dots, 0)$ be a boundary point of Ω^+ . We set

$$v = [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} + w(r)$$

such that $r^2 = x_1^2 + x_2^2 + \dots + x_n^2$ and $w \in C^2(D^+)$ denotes the set of twice continuously differentiable functions on D^+ , also an even function with respect to all x_i -variables. We apply formula (3.3) to the region $D^+ \setminus D_\varepsilon^+ \subset D^+$, where D_ε^+ is an upper half sphere centered at P' of radius ε . Letting $\varepsilon \rightarrow 0$ and considering $[(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1}$ for the fundamental solution of B , we deduce that

$$\begin{aligned} & \int \int_{D^+ \setminus D_\varepsilon^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dx \\ &= \int_{\Omega} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega - \int_{\Omega_\varepsilon} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega. \end{aligned} \tag{3.4}$$

Hence, let us write the function v in (3.4), then we have

$$\begin{aligned} & \int \int_{D^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dx - \int \int_{D_\varepsilon^+} [uBw - wBu] \prod_{i=1}^n x_i^{2\gamma_i} dx \\ & \quad - \int \int_{D_\varepsilon^+} [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} Bu \prod_{i=1}^n x_i^{2\gamma_i} dx \\ &= \int_{\Omega} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega + \int_{\Omega_\varepsilon} u \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega \\ & \quad - \int_{\Omega_\varepsilon} \left(u \frac{\partial w}{\partial n} - w \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega \\ & \quad - \int_{\Omega_\varepsilon} [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} \frac{\partial u}{\partial n} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega \end{aligned} \tag{3.5}$$

and

$$\begin{aligned} & \int \int_{D^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dg \\ &= \int_{\Omega} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega \\ & \quad + [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} \int \int_{D_\varepsilon^+} Bu \prod_{i=1}^n x_i^{2\gamma_i} dg \\ & \quad - [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} \int_{\Omega_\varepsilon} \frac{\partial u}{\partial n} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega + \int_{\Omega_\varepsilon} u \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega. \end{aligned}$$

Therefore, we get

$$\int \int_{D^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dg = \int_{\Omega} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega + \int_{\Omega_{\epsilon}} u \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega.$$

Considering $Bv = Bw$, consequently, we have

$$\int \int_{D^+} [uBv - vBu] \prod_{i=1}^n x_i^{2\gamma_i} dg = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0) + \int_{\Omega} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} d\Omega. \tag{3.6}$$

Here,

$$\int_{\Omega_{\epsilon}} \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} = m(\Omega_{\epsilon})$$

and

$$\lim_{\epsilon \rightarrow 0} \left[\int_{\Omega_{\epsilon}} u \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} d\Omega \right] = m(\Omega_{\epsilon}) u(0).$$

Let K_R^+ be a half sphere of radius R centered at P' contained in D^+ and W_R^+ be the surface of K_R^+ . Suppose that

$$v = (n + 2|\gamma| - 2)^{-1} [r^{-(n+2|\gamma|-2)} - R^{-(n+2|\gamma|-2)}]. \tag{3.7}$$

Under our assumptions, we can rewrite the equality (3.6) as

$$\int_{K_R^+} \{uB(n + 2|\gamma| - 2)^{-1} [r^{-(n+2|\gamma|-2)} - R^{-(n+2|\gamma|-2)}] - vBu\} \prod_{i=1}^n x_i^{2\gamma_i} dg = \int_{W_R^+} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} dw + \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0).$$

Since R is constant and $[(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1}$ for the fundamental solution of B , hence we have

$$\begin{aligned} - \int_{K_R^+} vBu \prod_{i=1}^n x_i^{2\gamma_i} dg &= \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0) \\ &+ \int_{W_R^+} \left\{ \left[u \frac{\partial(n + 2|\gamma| - 2)^{-1}}{\partial n} - (n + 2|\gamma| - 2)^{-1} \right] \right. \\ &\times \left. [r^{-(n+2|\gamma|-2)} - R^{-(n+2|\gamma|-2)}] \frac{\partial u}{\partial n} \right\} \prod_{i=1}^n x_i^{2\gamma_i} dw \end{aligned}$$

$$\begin{aligned}
 &= \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0) \\
 &\quad + \int_{W_R^+} u \frac{\partial(n + 2|\gamma| - 2)^{-1} r^{-(n+2|\gamma|-2)}}{\partial n} \prod_{i=1}^n x_i^{2\gamma_i} dw \\
 &= \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})} u(0) - \frac{1}{R^{n+2|\gamma|-1}} \int_{W_R^+} u \prod_{i=1}^n x_i^{2\gamma_i} dw
 \end{aligned}$$

and

$$\begin{aligned}
 &\frac{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2}) R^{n+2|\gamma|-1}} \int_{W_R^+} u(x_1, \dots, x_n) \prod_{i=1}^n x_i^{2\gamma_i} dw \\
 &= u(0) + \frac{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \int_{K_R^+} \nu B u \prod_{i=1}^n x_i^{2\gamma_i} dg. \tag{3.8}
 \end{aligned}$$

Choose the function ν such that

$$\nu = c \frac{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \left[(n + 2|\gamma| - 2) r^{n+2|\gamma|-2} \right]^{-1} + w(r), \tag{3.9}$$

where $w(r) \in C^{(2)}(K_R^+)$, also an even function with respect to all variables and $0 < r \leq R$. Then, under the initial conditions

$$\nu|_{W_R^+} = \frac{\partial \nu}{\partial r} \Big|_{W_R^+} = 0, \tag{3.10}$$

we obtain

$$\begin{aligned}
 \int_{K_R^+} [uB\nu - \nu B u] \prod_{i=1}^n x_i^{2\gamma_i} dg &= \int_{W_R^+} \left(u \frac{\partial \nu}{\partial n} - \nu \frac{\partial u}{\partial n} \right) \prod_{i=1}^n x_i^{2\gamma_i} dw \\
 &= \int_{W_R^+} u \frac{\partial \nu}{\partial n} \prod_{i=1}^n x_i^{2\gamma_i} dw \\
 &= c \frac{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \int_{W_R^+} u \frac{1}{r^{n+2|\gamma|-1}} \prod_{i=1}^n x_i^{2\gamma_i} dw \\
 &= cu(0). \tag{3.11}
 \end{aligned}$$

Since (3.11) is true for all u , we can take an arbitrary function that is even with respect to all variables in $C^{(2)}(K_R^+)$. Therefore, we assume that $u \in C^{(2m+2)}(K_R^+)$. For $\eta \leq m$, we have the following identity:

$$cB^\eta u(0) = \int_{K_R^+} [B^\eta u B \nu - \nu B^{\eta+1} u] \prod_{i=1}^n x_i^{2\gamma_i} dg, \tag{3.12}$$

where B^η denotes the η th order of B , that is, $B^1 u = Bu$, $B^2 u = BBu$, etc. Now, let us define $\{\nu_\eta\}$ as a sequence of functions of the type (3.9). Then the following differential equation

can be written:

$$Bv_{\eta+1} = v''_{\eta+1} + \frac{(n + 2|\gamma| - 1)}{r} v'_{\eta+1} = v_{\eta} \quad (\eta = 0, 1, 2, \dots, m). \tag{3.13}$$

By the initial conditions (3.10) and the function

$$v_0 = \frac{2^{n-1}\Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} (n + 2|\gamma| - 2)^{-1} [r^{-(n+2|\gamma|-2)} - R^{-(n+2|\gamma|-2)}], \tag{3.14}$$

it can be seen that the solution of (3.13) is as follows:

$$v_{\eta+1} = [(n + 2|\gamma| - 2)r^{n+2|\gamma|-2}]^{-1} \int_r^R \rho v_{\eta}(\rho) [\rho^{n+2|\gamma|-2} - r^{n+2|\gamma|-2}] d\rho.$$

Under the initial conditions, for each $\eta = 0, 1, 2, \dots$, the solutions v_0, v_1, \dots corresponding to the constants c_0, c_1, \dots can easily be found and obtained as follows:

$$c_{\eta} = \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{n}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{n}{2})} \quad (\eta = 0, 1, 2, \dots).$$

Replacing v by v_{η} in (3.12), we see that

$$c_{\eta} B^{\eta} u(0) = \int_{K_R^+} [B^{\eta} u B v_{\eta} - v_{\eta} B^{\eta+1} u] \prod_{i=1}^n x_i^{2\gamma_i} dg.$$

Since $Bv_{\eta} = v_{\eta-1}$, we obtain

$$c_{\eta} B^{\eta} u(0) = \int_{K_R^+} [v_{\eta-1} B^{\eta} u - v_{\eta} B^{\eta+1} u] \prod_{i=1}^n x_i^{2\gamma_i} dg.$$

Summing all the equations corresponding to $\eta = 1, 2, \dots, m$, then we conclude

$$\sum_{\eta=1}^m c_{\eta} B^{\eta} u(0) = \int_{K_R^+} [v_0 B u - v_m B^{m+1} u] \prod_{i=1}^n x_i^{2\gamma_i} dg.$$

By (3.8) and the function v_0 , then we get

$$\int_{K_R^+} v_0 B u \prod_{i=1}^n x_i^{2\gamma_i} dg = \sum_{\eta=1}^m c_{\eta} B^{\eta} u(0) + \int_{K_R^+} v_m B^{m+1} u \prod_{i=1}^n x_i^{2\gamma_i} dg$$

and

$$\begin{aligned} & \frac{2^{n-1}\Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \int_{K_R^+} (n + 2|\gamma| - 2)^{-1} [r^{-(n+2|\gamma|-2)} - R^{-(n+2|\gamma|-2)}] B u \prod_{i=1}^n x_i^{2\gamma_i} dg \\ & = \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{n}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{n}{2})} B^{\eta} u(0) + \int_{K_R^+} v_m B^{m+1} u \prod_{i=1}^n x_i^{2\gamma_i} dg. \end{aligned}$$

By the function v in (3.7), we get

$$\begin{aligned} \frac{2^{n-1}\Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \int_{K_R^+} v B u \prod_{i=1}^n x_i^{2\gamma_i} dg &= \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{n}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{n}{2})} B^\eta u(0) \\ &+ \int_{K_R^+} v_m B^{m+1} u \prod_{i=1}^n x_i^{2\gamma_i} dg \end{aligned} \tag{3.15}$$

and from (3.8)

$$\begin{aligned} [c_{n,\gamma} R^{n+2|\gamma|-1}]^{-1} \int_{W_R^+} u(x) \prod_{i=1}^n x_i^{2\gamma_i} dw \\ = u(0) + \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{n}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{n}{2})} B^\eta u(0) + \int_{K_R^+} v_m B^{m+1} u \prod_{i=1}^n x_i^{2\gamma_i} dg. \end{aligned}$$

Thus we have

$$\begin{aligned} [c_{n,\gamma} R^{n+2|\gamma|-1}]^{-1} \int_{W_R^+} u(x) \prod_{i=1}^n x_i^{2\gamma_i} dw \\ = \sum_{\eta=0}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{n}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{n}{2})} B^\eta u(0) + \int_{K_R^+} v_m B^{m+1} u \prod_{i=1}^n x_i^{2\gamma_i} dg, \end{aligned} \tag{3.16}$$

where $c_0 = 1$, $B^0 u_0 = u_0$ and $c_{n,\gamma} = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{n}{2})}$. This formula is valid for all arbitrary function $u \in K_R^+ \subseteq D^+$ which is a continuously differentiable function of order $(2m + 2)$. Let us establish the equality (3.8) with respect to $T^\gamma u(x)$; we get

$$\begin{aligned} \frac{2^{n-1}\Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2}) R^{n+2|\gamma|-1}} \int_{W_R^+} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dw \\ = u(\gamma) + \frac{2^{n-1}\Gamma(|\gamma| + \frac{n}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})} \int_{K_R^+} v B T^\gamma u \prod_{i=1}^n x_i^{2\gamma_i} dg. \end{aligned} \tag{3.17}$$

Setting $v = v_\eta$ and $u(x) = T^\gamma u(x)$ in (3.12), then we obtain

$$\begin{aligned} c B^\eta u(\gamma) &= \int_{K_R^+} [B^\eta T^\gamma u(x) B v - v B^{\eta+1} T^\gamma u(x)] \prod_{i=1}^n x_i^{2\gamma_i} dg, \\ c_\eta B^\eta u(\gamma) &= \int_{K_R^+} [B^\eta T^\gamma u(x) B v_\eta - v_\eta B^{\eta+1} T^\gamma u(x)] \prod_{i=1}^n x_i^{2\gamma_i} dg. \end{aligned}$$

If we sum all the equations corresponding to $\eta = (1, 2, \dots, m)$, we conclude

$$\sum_{\eta=1}^m c_\eta B^\eta u(\gamma) = \int_{K_R^+} [v_0 B T^\gamma u - v_m B^{m+1} T^\gamma u] \prod_{i=1}^n x_i^{2\gamma_i} dg.$$

By (3.8) and the function v_0 , we get

$$\int_{K_R^+} v_0 B T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dg = \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{\eta}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{\eta}{2})} B^\eta u(y) + \int_{W_R^+} v_m B^m T^\gamma u \prod_{i=1}^n x_i^{2\gamma_i} dw \tag{3.18}$$

and

$$\begin{aligned} & \frac{2^{n-1} \Gamma(|\gamma| + \frac{\eta}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2}) R^{n+2|\gamma|-1}} \int_{K_R^+} v B T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dg \\ &= \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{\eta}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{\eta}{2})} B^\eta u(y) + \int_{K_R^+} v_m B^{m+1} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dg, \tag{3.19} \\ & [c_{n,\gamma} R^{n+2|\gamma|-1}]^{-1} \int_{W_R^+} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dw_R \\ &= u(y) + \sum_{\eta=1}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{\eta}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{\eta}{2})} B^\eta u(y) \\ &+ \int_{K_R^+} v_m B^{m+1} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dg, \tag{3.20} \end{aligned}$$

and

$$\begin{aligned} & [c_{n,\gamma} R^{n+2|\gamma|-1}]^{-1} \int_{W_R^+} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dw_R \\ &= \sum_{\eta=0}^m \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{\eta}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{\eta}{2})} B^\eta u(y) + \int_{K_R^+} v_m B^{m+1} T^\gamma u(x) \prod_{i=1}^n x_i^{2\gamma_i} dg, \tag{3.21} \end{aligned}$$

where $B^0 u = u$ and $c_0 = 1$. We note that the second integral on the right-hand side in (3.21) tends to zero for $m \rightarrow \infty$; then we have the result

$$\begin{aligned} & \frac{2^{n-1} \Gamma(|\gamma| + \frac{\eta}{2})}{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2}) R^{n+2|\gamma|-1}} \int_{S_{n-1}^+} T^\gamma u(R\theta_1, \dots, R\theta_n) R^{2|\gamma|} \prod_{i=1}^n \theta_i^{2\gamma_i} R^{n-1} dS \\ &= \sum_{\eta=0}^{\infty} \left(\frac{R}{2}\right)^{2\eta} \frac{\Gamma(|\gamma| + \frac{\eta}{2})}{\eta! \Gamma(\eta + |\gamma| + \frac{\eta}{2})} B^\eta u(y), \end{aligned}$$

where $x = R\theta$ and $dw_R = R^{n-1} dS$. Hence

$$\int_{S_{n-1}^+} T^\gamma u(R\theta_1, R\theta_2, \dots, R\theta_n) \prod_{i=1}^n \theta_i^{2\gamma_i} dS = \frac{\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})}{2^{n-1} \Gamma(|\gamma| + \frac{\eta}{2})} u(y).$$

Thus the proof is completed. □

Proof of Theorem 2.1 We shall need the Fourier-Bessel transforms of the function $e^{-\alpha|x|^2}$. The identity to be proved can be rewritten as

$$F_B(e^{-\alpha|x|^2})(y) = c_\gamma \int_{\mathbb{R}_+^n} e^{-\alpha|x|^2} \prod_{i=1}^n j_{\gamma_i - \frac{1}{2}}(x_i y_i) d\mu_\gamma(x). \tag{3.22}$$

Let $\gamma > -1$, $\alpha > 0$ and $J_\gamma(br)$ be a Bessel function. We recall that

$$\int_0^\infty e^{-\alpha r^2} r^{\gamma+1} J_\gamma(br) dr = b^\gamma (2\alpha)^{-\gamma-1} e^{-\frac{b^2}{4\alpha}},$$

then we get

$$\begin{aligned} J_{\gamma_1 - \frac{1}{2}}(x_1 y_1) &= \left[2^{\gamma_1 - \frac{1}{2}} \Gamma\left(\gamma_1 + \frac{1}{2}\right) \right]^{-1} (x_1 y_1)^{\gamma_1 - \frac{1}{2}} j_{\gamma_1 - \frac{1}{2}}(x_1 y_1), \\ &\vdots \\ J_{\gamma_n - \frac{1}{2}}(x_n y_n) &= \left[2^{\gamma_n - \frac{1}{2}} \Gamma\left(\gamma_n + \frac{1}{2}\right) \right]^{-1} (x_n y_n)^{\gamma_n - \frac{1}{2}} j_{\gamma_n - \frac{1}{2}}(x_n y_n), \end{aligned}$$

where $J_\gamma(r) = [2^\gamma \Gamma(\gamma + 1)]^{-1} r^\gamma j_\gamma(r)$. We may now calculate each of the integrals in (3.22). We have

$$\begin{aligned} I_1 &= \int_0^\infty e^{-\alpha x_1^2} j_{\gamma_1 - \frac{1}{2}}(x_1 y_1) x_1^{2\gamma_1} dx_1 \\ &= \int_0^\infty e^{-\alpha r^2} J_{\gamma_1 - \frac{1}{2}}(rs) r^{2\gamma_1} r^{-\gamma_1 + \frac{1}{2}} s^{-\gamma_1 + \frac{1}{2}} dr \\ &= 2^{\gamma_1 - \frac{1}{2}} \Gamma\left(\gamma_1 + \frac{1}{2}\right) (2\alpha)^{-\gamma_1 - \frac{1}{2}} e^{-\frac{s^2}{4\alpha}}, \end{aligned} \tag{3.23}$$

where $x_1 = r$ and $y_1 = s$. In a similar way, we obtain

$$I_2 = \int_0^\infty e^{-\alpha x_2^2} j_{\gamma_2 - \frac{1}{2}}(x_2 y_2) x_2^{2\gamma_2} dx_2 = \frac{2^{\gamma_2 - \frac{1}{2}} \Gamma\left(\gamma_2 + \frac{1}{2}\right)}{(2\alpha)^{\gamma_2 + \frac{1}{2}}} e^{-\frac{y_2^2}{4\alpha}}, \tag{3.24}$$

$$I_n = \int_0^\infty e^{-\alpha x_n^2} j_{\gamma_n - \frac{1}{2}}(x_n y_n) x_n^{2\gamma_n} dx_n = \frac{2^{\gamma_n - \frac{1}{2}} \Gamma\left(\gamma_n + \frac{1}{2}\right)}{(2\alpha)^{\gamma_n + \frac{1}{2}}} e^{-\frac{y_n^2}{4\alpha}}. \tag{3.25}$$

If we write the equalities (3.23)-(3.25) in (3.22) then we have

$$F_B(e^{-\alpha|x|^2})(y) = e^{-\frac{|y|^2}{4\alpha}} (2\alpha)^{-\frac{2|y|-n}{2}}.$$

Letting $\alpha \rightarrow 1$ and $y \rightarrow -2y$, then we get

$$F_B(e^{-\alpha|x|^2})(y) = e^{-|y|^2} 2^{-n}.$$

So, we can easily obtain the identity

$$F_B(e^{-|x|^2})(y) = c_\gamma \int_{\mathbb{R}_+^n} \prod_{i=1}^n j_{\gamma_i - \frac{1}{2}}(x_i 2y_i) e^{-|x|^2} d\mu_\gamma(x).$$

By similar arguments, we can see that

$$I'_1 = \int_0^\infty e^{-x_1^2} j_{\gamma_1 - \frac{1}{2}}(x_1 2\gamma_1) x_1^{2\gamma_1} dx_1,$$

$$I'_1 = \int_0^\infty e^{-r'^2} j_{\gamma_1 - \frac{1}{2}}(2r' s') r'^{2\gamma_1} dr'.$$

Hence, we get

$$I'_1 = 2^{-1} \Gamma\left(\gamma_1 + \frac{1}{2}\right) e^{-s_1'^2}$$

and

$$I'_n = \int_0^\infty e^{-r'^2} j_{\gamma_n - \frac{1}{2}}(2r' s') r'^{2\gamma_n} dr' = 2^{-1} e^{-s_n'^2} \Gamma\left(\gamma_n + \frac{1}{2}\right).$$

Consequently, we deduce

$$F_B(e^{-\alpha|x|^2})(\gamma) = 2^{-n} \prod_{i=1}^n \Gamma\left(\gamma_i + \frac{1}{2}\right) e^{-|\gamma|^2}. \tag{3.26}$$

When we apply the Bessel differential operator $P_k(B_{t_1}, B_{t_2}, \dots, B_{t_{n-1}}, B_{t_n})$ to both sides of the identity (3.26), then we see

$$F_B[P_k(x)e^{-|x|^2}](t) = \int_0^\infty P_k(x) e^{-x_1^2} j_{\gamma_1 - \frac{1}{2}}(2x_1 t_1) x_1^{2\gamma_1} \\ \dots \int_0^\infty e^{-x_n^2} j_{\gamma_n - \frac{1}{2}}(2x_n t_n) x_n^{2\gamma_n} dx \\ = Q(t) 2^{-n} \prod_{i=1}^n \Gamma\left(\gamma_i + \frac{1}{2}\right) e^{-|t|^2},$$

where $Q(t)$ is a polynomial. The problem is therefore to show that $Q(t) = P_k(iy)$. Now using the identity

$$j_{\gamma - \frac{1}{2}}(r) = \frac{\Gamma(\gamma + \frac{1}{2})}{\Gamma(\gamma)\Gamma(\frac{1}{2})} \int_0^\pi e^{ir \cos \alpha} (\sin \alpha)^{2\gamma - 1} d\alpha$$

we have

$$Q(t) = c_1 \int_{\mathbb{R}_+^n} P_k(x) e^{|\gamma|^2 - |x|^2} \prod_{i=1}^n j_{\gamma_i - \frac{1}{2}}(2x_i t_i) d\mu_\gamma(x) \\ = c_1 \int_{\mathbb{R}_+^n} P_k(x) e^{|\gamma|^2 - (x_1^2 + \dots + x_n^2)} \prod_{i=1}^n j_{\gamma_i - \frac{1}{2}}(2x_i t_i) d\mu_\gamma(x) \\ = c_2 \int_{\mathbb{R}_+^n} P_k(x) e^{|\gamma|^2 - |x|^2} \left(\int_0^\pi e^{2ix_1 t_1 \cos w} \sin^{2\gamma_1 - 1} w dw \right. \\ \left. \dots \int_0^\pi e^{2ix_n t_n \cos w} \sin^{2\gamma_n - 1} w dw \right) d\mu_\gamma(x)$$

$$\begin{aligned}
 &= c_2 \int_{\mathbb{R}_+^n} P_k(x) d\mu_\gamma(x) \int_0^\pi \cdots \int_0^\pi (e^{t_1^2 - x_1^2 - 2ix_1 t_1 \cos w} \sin^{2\gamma_1 - 1} w) \\
 &\quad \cdots (e^{t_n^2 - x_n^2 - 2ix_n t_n \cos w} \sin^{2\gamma_n - 1} w) dw = 2^n \int_{\mathbb{R}_+^n} P_k(x) [T^{-it}(e^{-|x|^2})] d\mu_\gamma(x),
 \end{aligned}$$

where $c_1 = 2^n [\prod_{i=1}^n \Gamma(\gamma_i + \frac{1}{2})]^{-1}$ and $c_2 = \prod_{i=1}^n 2^n \Gamma(\gamma_i + \frac{1}{2}) (\pi^{\frac{n}{2}} \Gamma(\gamma_i))^{-1}$. Replacing t by $-it$ and by the properties of T^y , we obtain

$$Q(-it) = c_1 \int_{\mathbb{R}_+^n} T^{-t} [P_k(x) e^{-|x|^2}] d\mu_\gamma(x). \tag{3.27}$$

Taking the change of variables $x \rightarrow r\theta$ for $0 < r < \infty$ and $\theta \in \mathcal{S}_+^{n-1}$ and applying the polar coordinates in (3.27), this gives the identity

$$Q(-it) = c_1 \int_0^\infty r^{2|\gamma|+n-1} \left(\int_{\mathcal{S}_+^{n-1}} T^{-t} [P_k(r\theta)] \prod_{i=1}^n \theta_i^{2\gamma_i} \right) e^{-r^2} dr d\theta. \tag{3.28}$$

We must calculate the integral $\int_0^\infty r^{2|\gamma|+n-1} e^{-r^2} dr$. Letting $r^2 = u$, then we get

$$\int_0^\infty r^{2|\gamma|+n-1} e^{-r^2} dr = \frac{1}{2} \int_0^\infty u^{|\gamma|+\frac{n}{2}-1} e^{-u} du = \frac{1}{2} \Gamma\left(|\gamma| + \frac{n}{2}\right).$$

Applying Theorem 3.1 to (3.28), we deduce $Q(-it) = P_k(t)$ and

$$\begin{aligned}
 F_B [P_k(x) e^{-\alpha|x|^2}](t) &= c_\gamma \int_{\mathbb{R}_+^n} P_k(x) e^{-|x|^2} \prod_{i=1}^n j_{\gamma_i - \frac{1}{2}}(2x_i t_i) d\mu_\gamma(x) \\
 &= 2^{-(|\gamma| + \frac{n}{2})} P_k(it) e^{-|t|^2}.
 \end{aligned}$$

Since P_k is homogeneous, we have the identity

$$F_B [P_k(x) e^{-|x|^2}](y) = 2^{-(|\gamma|+k+\frac{n}{2})} t^k P_k(y) e^{-\frac{|y|^2}{4}},$$

and so we obtain the desired conclusion. □

We come now to what has been our main goal in this paper.

Theorem 3.2 *Let P_k be homogeneous harmonic polynomial of degree k . Then we get*

$$F_B \left[P_k(x) \frac{P_k(x)}{|x|^{k+n+2|\gamma|}} \right](y) = 2^{-\frac{n-2|\gamma|}{2}} t^k \frac{\Gamma(\frac{k}{2})}{\Gamma(\frac{k+n+2|\gamma|}{2})} \frac{P_k(y)}{|y|^k}.$$

Proof Consider the identity

$$F_B [f(\alpha x)](t) = \alpha^{-n-2|\gamma|} F_B [f(x)] \left(\frac{t}{\alpha} \right). \tag{3.29}$$

By (3.29) and Theorem 2.1, we have

$$F_B [P_k(x) e^{-\alpha|x|^2}](y) = (2\alpha)^{-(k+|\gamma|+\frac{n}{2})} t^k e^{-\frac{|y|^2}{4\alpha}} P_k(y).$$

Assume in addition that $\varphi \in \mathcal{S}(\mathbb{R}_+^n)$. Then

$$\int_{\mathbb{R}_+^n} P_k(x) e^{-\alpha|x|^2} \varphi(x) d\mu_\gamma(x) = (2\alpha)^{-(k+|\gamma|+\frac{n}{2})} i^k \int_{\mathbb{R}_+^n} P_k(x) e^{-\frac{|x|^2}{4\alpha}} \varphi(x) d\mu_\gamma(x).$$

We now integrate both sides of the above with respect to α having multiplied the equation by a suitable power of α ($\alpha^{\beta-1}$, $2\beta = k + n + 2|\gamma| - \varepsilon$). We obtain

$$\begin{aligned} \int_0^\infty \left[\int_{\mathbb{R}_+^n} P_k(x) e^{-\alpha|x|^2} \alpha^{\beta-1} \varphi(x) d\mu_\gamma(x) \right] d\alpha &= \int_0^\infty \left[(2\alpha)^{-(k+|\gamma|+\frac{n}{2})} i^k \alpha^{\beta-1} \right] \\ &\times \left[\int_{\mathbb{R}_+^n} P_k(x) e^{-\frac{|x|^2}{4\alpha}} \varphi(x) d\mu_\gamma(x) \right] d\alpha. \end{aligned}$$

If we use the fact that

$$\int_0^\infty e^{-u} \left(\frac{u}{|x|^2} \right)^{\beta-1} \frac{du}{|x|^2} = \Gamma(\beta) |x|^{-2\beta}$$

we get

$$\int_0^\infty \left[\int_{\mathbb{R}_+^n} P_k(x) e^{-\alpha|x|^2} \alpha^{\beta-1} \varphi(x) d\mu_\gamma(x) \right] d\alpha = \Gamma(\beta) \int_{\mathbb{R}_+^n} \frac{P_k(x)}{|x|^{2\beta}} \varphi(x) d\mu_\gamma(x). \tag{3.30}$$

The corresponding integration for the right side gives

$$\begin{aligned} &\int_0^\infty \left\{ (2\alpha)^{-(k+|\gamma|+\frac{n}{2})} i^k \alpha^{\beta-1} \left[\int_{\mathbb{R}_+^n} P_k(x) e^{-\frac{|x|^2}{4\alpha}} \varphi(x) d\mu_\gamma(x) \right] \right\} d\alpha \\ &= \int_0^\infty 2^{-(k+|\gamma|+\frac{n}{2})} i^k \alpha^{-(\frac{k+\varepsilon}{2}+1)} \int_{\mathbb{R}_+^n} [P_k(x) e^{-\frac{|x|^2}{4\alpha}} \varphi(x) d\mu_\gamma(x)] d\alpha \\ &= 2^{-(k+|\gamma|+\frac{n}{2})} \int_0^\infty \int_{\mathbb{R}_+^n} \alpha^{(-\frac{k+\varepsilon}{2}-1)} P_k(x) e^{-\frac{|x|^2}{4\alpha}} i^k \varphi(x) d\mu_\gamma(x) d\alpha. \end{aligned}$$

Letting $\alpha \rightarrow \frac{1}{4\alpha}$, we obtain

$$\begin{aligned} &\int_0^\infty \left\{ (2\alpha)^{-(k+|\gamma|+\frac{n}{2})} i^k \alpha^{\beta-1} \left[\int_{\mathbb{R}_+^n} P_k(x) e^{-\frac{|x|^2}{4\alpha}} \varphi(x) d\mu_\gamma(x) \right] \right\} d\alpha \\ &= 2^{-(|\gamma|+\frac{n}{2}-\varepsilon)} \Gamma\left(\frac{k+\varepsilon}{2}\right) i^k \int_{\mathbb{R}_+^n} \frac{P_k(x)}{|x|^{k+\varepsilon}} \varphi(x) d\mu_\gamma(x). \end{aligned} \tag{3.31}$$

Thus (3.30) and (3.31) lead to the identity

$$\int_{\mathbb{R}_+^n} \frac{P_k(x)}{|x|^{2\beta}} \varphi(x) d\mu_\gamma(x) = 2^{-\frac{2|\gamma|+n}{2}+\varepsilon} i^k \frac{\Gamma(\frac{k+\varepsilon}{2})}{\Gamma(\beta)} \int_{\mathbb{R}_+^n} \frac{P_k(x)}{|x|^{k+\varepsilon}} \varphi(x) d\mu_\gamma(x).$$

By Lemma 2.2, we have therefore concluded the proof of the theorem. We have

$$F_B \left[p.v. \frac{P_k(x)}{|x|^{k+n+2|\gamma|-\varepsilon}} \right] (y) = 2^{-\frac{2|\gamma|+n}{2}} i^k \frac{\Gamma(\frac{k}{2})}{\Gamma(\frac{k+n+2|\gamma|}{2})} \frac{P_k(y)}{|y|^k}. \quad \square$$

Definition 3.3 Let T^γ be the Bessel-generalized shift operator and let f be a Schwartz function on \mathbb{R}_+^n . We define the high order Riesz-Bessel transforms $R_B^{(k)}$ of order k with respect to the Bessel-generalized shift operator as

$$\begin{aligned} R_B^{(k)}(f)(x) &= c_k(n, \gamma) \left[p.v. \frac{P_k(y)}{|y|^{n+k+2|\gamma|}} * f \right](x) \\ &= c_k(n, \gamma) \lim_{\varepsilon \rightarrow 0} \int_{0 < \varepsilon < |x|} \frac{P_k(y)}{|y|^{k+n+2|\gamma|}} T^\gamma f(x) d\mu_\gamma(y), \end{aligned}$$

where $c_k(n, \gamma) = 2^{-\frac{n+2|\gamma|}{2}} \Gamma(\frac{n+k+2|\gamma|}{2}) [\Gamma(\frac{k}{2})]^{-1}$ ($k = 1, 2, \dots, n$) and $P_k(x)$ is a homogeneous polynomial of degree k in \mathbb{R}_+^n which satisfies $BP_k = 0$.

By Theorem 3.2, we conclude that

$$F_B[R_B^{(k)}(f)](\xi) = i^k P_k(\xi) |\xi|^{-k} F_B[f](\xi). \tag{3.32}$$

One of the important applications of the high order Riesz transforms is that they can be used to mediate between various combinations of partial derivatives of a function. We shall here content ourselves with two very simple illustrations, which examples have an interest on their own and have already the characteristic features of a general type of estimate which can be made in the theory of elliptic differential operators.

Proposition 3.4 *Suppose f is a class of $\mathcal{S}(\mathbb{R}_+^n)$ and has compact support. Let Bf be the Bessel differential operator. Then we have the a priori bound*

$$\|\partial_{x_i} \partial_{x_k} f\|_{p,\gamma} \leq A_p \|Bf\|_{p,\gamma}$$

with A_p independent of f .

In (3.32), we may take $k = 1$. Then this proposition is an immediate consequence of the $L_{p,\gamma}$ boundedness of the Riesz-Bessel transforms generated by a Bessel-generalized shift operator and the identity

$$\partial_{x_i} (\partial_{x_k} f) = -R_B R_B Bf. \tag{3.33}$$

To prove (3.33) we use the Fourier-Bessel transform. Thus if $F_B[f](x)$ is the Fourier-Bessel transform of f , then the Fourier-Bessel transform of $\partial_{x_i} f$ is

$$F_B[\partial_{x_i} f](y) = -x_i F_B[f](y)$$

and

$$F_B[Bf](y) = -|x|^2 F_B[f](y),$$

and so

$$F_B[\partial_{x_i} (\partial_{x_k} f)](y) = -\left(\frac{x_i}{|x|}\right) \left(\frac{x_k}{|x|}\right) |x|^2 F_B[f](y) = -F_B[R_B R_B Bf](y),$$

which gives (3.33). Thus we have

$$\|\partial_{x_i}(\partial_{x_k}f)\|_p = \|R_B R_B Bf\|_p \leq C \|Bf\|_p.$$

By using the Fourier Bessel transformations, we have

$$P(|y|^2)F_B[(Bf)](y) = -|y|^2 F_B[P(B)f](y)$$

(see [9]).

Corollary 3.5 *Suppose P_k is a homogeneous elliptic polynomial of degree k and f is k -times continuously differentiable with compact support. Then we have the priori estimate*

$$\|Bf\|_{p,\gamma} \leq A_p \|P_k(B)f\|_{p,\gamma}, \quad 1 < p < \infty.$$

To prove this inequality, we note that the following relation between Fourier Bessel transform of Bf and $P_k(B)f$ holds:

$$P_k(|y|^2)F_B[Bf](y) = -|y|^2 F_B[P_k(B)f](y).$$

$P_k(|y|^2)$ is non-vanishing except at the origin and let $-\frac{|y|^2}{P_k(|y|^2)}$ be homogeneous of degree zero and indefinitely differentiable on the unit sphere. Then we get

$$Bf = R_B(P_k(B)f).$$

We also have the following $L_{p,\gamma}$ boundedness of the high order Riesz-Bessel transform.

Theorem 3.6 *The high order Riesz-Bessel transforms generated by a Bessel-generalized shift operator are bounded operators from $L_{p,\gamma}(\mathbb{R}_+^n)$ into itself for all $1 < p < \infty$*

$$\|R_B f\|_{p,\gamma} \leq A_p \|f\|_{p,\gamma}.$$

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

All authors contributed equally to the writing of this paper. All authors read and approved the final manuscript.

Author details

¹Department of Mathematics, Dumlupinar University, Kütahya, Turkey. ²Department of Electric-Electronic Engineering, Faculty of Technology, Gazi University, Ankara, Turkey.

Received: 5 December 2013 Accepted: 27 March 2014 Published: 11 Apr 2014

References

- Stein, EM, Weiss, G: Introduction to Fourier Analysis on Euclidean Spaces. Princeton University Press, Princeton (1971)
- Muckenhoupt, B, Stein, EM: Classical expansions and their relation to conjugate harmonic functions. *Trans. Am. Math. Soc.* **118**, 17-92 (1965)
- Kipriyanov, IA: Fourier Bessel transformations and imbedding theorems. *Trudy Math. Inst. Steklov* **89**, 130-213 (1967) (in Russian)
- Kipriyanov, IA, Klyuchantsev, MI: On singular integrals generated by the generalized shift operator II. *Sib. Mat. Zh.* **11**, 1060-1083 (1970) (in Russian)

5. Aliev, IA, Gadjev, AD: Weighted estimates for multidimensional singular integrals generated by the generalized shift operator. *Mat. Sb.* **77**(1), 37-55 (1994)
6. Guliev, VS: Sobolev theorems for B -Riesz potentials. *Dokl. Akad. Nauk, Ross. Akad. Nauk* **358**(4), 450-451 (1998) (in Russian)
7. Aliev, IA: On Riesz transformations generated by generalized shift operator. *Izv. Akad. Nauk Azerb. SSR, Ser. Fiz.-Teh. Mat. Nauk* **8**(1), 7-13 (1987)
8. Ekincioglu, I, Yildirim, H, Akin, O: The mean value theorem for Laplacean-Bessel equation. In: *Invited Lecture Delivered at the Seventh International Colloquium on Differential Equations, vol. II*, pp. 29-37. Academic Publications, Plovdiv (1996)
9. Ekincioglu, I, Ozkin, IK: On higher order Riesz transformations generated by generalized shift operator. *Turk. J. Math.* **21**, 51-60 (1997)
10. Gadjev, AD, Guliev, EV: Two weighted inequality for singular integrals in Lebesgue spaces associated with the Laplace-Bessel differential operator. *Proc. A. Razmadze Math. Inst.* **138**, 1-15 (2005)
11. Levitan, BM: Bessel function expansions in series and Fourier integrals. *Usp. Mat. Nauk*, **6**(2), 102-143 (1951) (in Russian)
12. Kipriyanov, IA: *Singular Elliptic Boundary Value Problems*. Nauka, Moscow (1997) (in Russian)
13. Kipriyanov, IA: Boundary value problems for elliptic partial differential operators. *Sov. Math. Dokl.* **11**, 1416-1419 (1970)
14. Cheng, M: On a theorem of Nicolesco and generalized Laplace operators. *Proc. Am. Math. Soc.* **2**, 77-86 (1951)
15. Nicolesco, M: Les fonctions polyharmoniques. *Actualités Sci. Ind.* **4** (1936)
16. Pizetti, P: Sul significato geometrico del secondo parametro differenziale di una funzione sopra una superficie qualunque. *Rend. Lincei* **18**, 309-316 (1909)
17. Picone, M: Nuovi indirizzi di ricerca teoria e nel calcolo soluzioni di talune equazioni lineari alle derivate parziali della fisica-matematica. *Ann. Sc. Norm. Super. Pisa, Cl. Sci.* **4**, 213-288 (1936)
18. Kipriyanova, NI: Mean value theorems for polyharmonic equations. Non-classic problems. *Math. Phys. Eq. NovoSibirsk*, 81-85 (1984) (in Russian)
19. Kipriyanova, NI: The Mean Value Formula for Singular Differential Operator With Order Second. *Differential Equations* **121**, 11 (1985)
20. Betancor, JJ, Castro, AJ, Curbelo, J: Harmonic analysis operators associated with multidimensional Bessel operators. *Proc. R. Soc. Edinb., Sect. A* **142**(5), 945-974 (2012)
21. Levitan, BM: *The Theory of Generalized Translation Operators*. Nauka, Moscow (1973) (in Russian)
22. Guliev, VS: Sobolev's theorem for the anisotropic Riesz-Bessel potential in Morrey-Bessel spaces. *Dokl. Akad. Nauk, Ross. Akad. Nauk* **367**(2), 155-156 (1999) (in Russian)
23. Guliev, VS: Some properties of the anisotropic Riesz-Bessel potential. *Anal. Math.* **26**(2), 99-118 (2000)
24. Guliev, VS, Garakhanova, NN, Zeren, Y: Pointwise and integral estimates for the Riesz B -potential in terms of B -maximal and B -fractionally maximal functions. *Sib. Mat. Zh.* **49**(6), 1263-1279 (2008) (in Russian). Translation in *Sib. Math. J.* **49**(6), 1008-1022 (2008)
25. Guliev, VS, Garakhanova, NN: The Sobolev-II'in theorem for the Riesz B -potential. *Sib. Mat. Zh.* **50**(1), 63-74 (2009) (in Russian). Translation in *Sib. Math. J.* **50**(1), 49-59 (2009)
26. Kipriyanov, IA, Lyakhov, LN: Multipliers of the mixed Fourier-Bessel transform. *Dokl. Akad. Nauk, Ross. Akad. Nauk* **354**(4), 449-451 (1997) (in Russian)
27. Lyakhov, LN: Multipliers of the mixed conversion of Fourier-Bessel. *Proc. Steklov Inst. Math.* **214**, 234-249 (1997) (in Russian)
28. Lyakhov, LN, Raykhegaur, LB: Singular heat equation with D_{B_γ} -Bessel operator. Fundamental solutions. *Problems in mathematical analysis. No. 67. J. Math. Sci. (N.Y.)* **188**(3), 285-292 (2013)
29. Stempak, K, Torrerá, JL: Higher Riesz transforms and imaginary powers associated to the harmonic oscillator. *Acta Math. Hung.* **111**, 43-64 (2006)
30. Stein, EM: *Singular Integrals and Differentiability Properties of Functions*. Princeton University Press, Princeton (1970)

10.1186/1029-242X-2014-148

Cite this article as: Ekincioglu et al.: High order Riesz transforms and mean value formula for generalized translate operator. *Journal of Inequalities and Applications* 2014, **2014**:148

Submit your manuscript to a SpringerOpen® journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Immediate publication on acceptance
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at ► springeropen.com