# Compactly supported solutions to stationary degenerate diffusion equations 

Stan Alama ${ }^{*, 1}$, Qiuping Lu ${ }^{2}$<br>Dept. of Mathematics and Statistics, McMaster Univ., Hamilton, ON, Canada L8S 4K1

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#### Abstract

We consider non-negative solutions of the semilinear elliptic equation in $\mathbb{R}^{n}$ with $n \geqslant 3$ : $$
-\Delta u=a(x) u^{q}+b(x) u^{p}
$$ where $0<q<1, p>q, a(x)$ sign-changing, $a=a^{+}-a^{-}$and $b(x) \leqslant 0$ is non-positive. Under appropriate growth assumption on $a^{-}$at infinity, we prove that all solutions in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ are compactly supported and their support is contained in a large ball with radius determined by $a$. When $\Omega^{0+}=\left\{x \in \mathbb{R}^{n} \mid a(x) \geqslant 0\right\}$ has several compact connected components, we give conditions under which there may or may not exist solutions which vanish identically on one or more of the components of $\Omega^{0+}$. For instance, we introduce a positive parameter $\lambda$ and replace $a$ by $\lambda a^{+}-a^{-}$. We then show that for $\lambda$ small, all solutions have compact support and there exist solutions with supports in any combination of these connected components of $\Omega^{0+}$. For $\lambda$ large and $p \leqslant 1$ the solution is unique and supported in all of $\Omega^{0+}$. We also prove the existence of the limit $\lambda \rightarrow \infty$ of this solution, which solves $-\Delta w=a^{+} w^{q}$ and $\lim _{|x| \rightarrow \infty} w(x)=0$. The analysis is based on comparison arguments and a priori bounds.


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## 1. Introduction

In this paper, we study the semilinear elliptic problem in $\mathbb{R}^{n}$ with $n \geqslant 3$ :

$$
\begin{cases}-\Delta u=a(x) u^{q}+b(x) u^{p} & \text { in } \mathbb{R}^{n}, 0<q<1, p>q,  \tag{1.1}\\ u \geqslant 0 & \text { in } \mathbb{R}^{n}, u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)\end{cases}
$$

where $a(x)$ and $b(x)$ are locally Hölder continuous, and $b(x) \leqslant 0$. By $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ we mean the completion of $C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ under the Dirichlet seminorm, $\left(\int_{\mathbb{R}^{n}}|\nabla u|^{2} d x\right)^{1 / 2}$. The important feature of this equation is that it combines a non-Lipschitz nonlinearity with a sign-changing coefficient $a(x)$, and it was originally observed by Schatzman [22] that solutions could vanish on large sets and in fact that, under appropriate hypotheses on $a(x)$, there exist solutions with compact support. Our main goal in this paper is to give general conditions on the coefficients $a(x)$ and $b(x)$ which guarantee that solutions must have compact support in $\mathbb{R}^{n}$, and more generally to study a phenomenon related to "dead cores", regions in $\mathbb{R}^{n}$ in which solutions vanish identically (see $[2,4,5]$ ).

Equations of this type (1.1) arise as stationary solutions to the degenerate reaction-diffusion equations of the form,

$$
w_{t}-\Delta\left(w^{m}\right)=w\left(a(x)+b(x) w^{s}\right), \quad m>1, s>0
$$

with nonlinearity of "logistic" type. Assuming time-independence and making the change of variable $u=w^{m}$, we arrive at (1.1) with $q=1 / m$ and $p=(s+1) / m>q$. If $w(x, t)$ represents a population density, then $a(x)$ represents a sort of growth rate, and the region where $a(x)>0$ is favorable to population growth, whereas the region where $a(x)<0$ is hostile to the species. It is therefore natural to assume that

$$
\begin{equation*}
\Omega^{+}=\left\{x \in \mathbb{R}^{n} \mid a(x)>0\right\} \text { is not empty and bounded. } \tag{1.2}
\end{equation*}
$$

Under this hypothesis alone we show that a solution exists:
Theorem 1.1. Assume (1.2), there exists a classical solution $U \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{\infty}\left(\mathbb{R}^{n}\right)$ of (1.1). Moreover $U \leqslant \omega$, where $\omega$ is the unique positive solution to

$$
\begin{equation*}
-\Delta \omega=a^{+} \omega^{q} \quad \text { in } \mathbb{R}^{n}, \quad \lim _{|x| \rightarrow \infty} \omega=0 \tag{1.3}
\end{equation*}
$$

Here we denote as usual, $a^{+}(x)=\max (0, a(x))$ and $a^{-}(x)=\max (0,-a(x))$. The existence and uniqueness of $\omega$ follows from Theorem $2^{\prime}$ in [8].

In order for solutions to have compact support, we require more information about $a(x)$. Let $\Omega^{0+}=\left\{x \in \mathbb{R}^{n} \mid a(x) \geqslant 0\right\}$. In case $\Omega^{0+}$ is unbounded, the strong maximum principle leads us to expect that in general solutions will not have compact support. So to obtain compactly supported solutions we will assume

$$
\begin{equation*}
\Omega^{0+} \text { is bounded and } \Omega^{+} \text {is non-empty. } \tag{1.4}
\end{equation*}
$$

However, this hypothesis is not sufficient, and in addition we must impose some conditions on the decay of the negative part $a^{-}(x)$. We prove:

Theorem 1.2. Assume (1.4). There exist $\eta>0$ and $\rho>0$ so that if

$$
\begin{equation*}
\liminf _{|x| \rightarrow \infty} a^{-}|x|^{(n-2)(1-q)}>\eta, \tag{1.5}
\end{equation*}
$$

every weak solution $u$ of (1.1) is classical and compactly supported, with $\operatorname{supp}(u)$ contained in $B(0, \rho)$.

To understand the interdependence of the constants in Theorem 1.2, think of $a^{+}(x)$ as being fixed and consider how the asymptotic behavior of $a^{-}(x)$ affects the support of the solution. By the hypothesis (1.4) we can choose $\rho_{1}>0$ so that $\Omega^{0+} \Subset B\left(0, \rho_{1}\right)$. Then, the constant $\eta$ will depend on $q, p$, $n, \rho_{1}$ and $\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$, and the constant $\rho$ will depend on $q, p, n, \rho_{1},\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$ and the asymptotic behavior of $a^{-}$near infinity. If $a^{-}$is decaying too rapidly to zero, then we may not have compact supported solutions (see Remark 2.12 at the end of Section 2). Nevertheless, we cannot claim that condition (1.5) is sharp, as we will see in Theorem 1.9 below.

In case $a(x)$ is bounded away from zero at infinity the result is somehow simpler:
Corollary 1.3. Assume (1.4), and suppose that there exist $\alpha>0$ and $\rho_{1}>0$ such that

$$
\begin{equation*}
a^{-}(x) \geqslant \alpha \quad \text { for all }|x| \geqslant \rho_{1} . \tag{1.6}
\end{equation*}
$$

Then, there exists $R>0$ so that every weak solution $u$ has support $\operatorname{supp}(u) \subset B(0, R)$. Moreover, $R$ depends only on $q, p, n, \rho_{1}$, and $\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$.

Note that we cannot expect the result of Theorem 1.2 to hold in $\mathbb{R}^{1}$ or $\mathbb{R}^{2}$. Solutions to (1.1) in $\mathbb{R}^{1}$ or $\mathbb{R}^{2}$ will be compactly supported under the hypothesis (1.6), but we cannot prove the uniform control on the support in terms of the coefficients as in Corollary 1.3. Note also that if $q \geqslant 1$, by the strong maximum principle there could not be any compactly supported solution at all, so the sublinearity of $q$ is crucial for the existence of solution with compact support.

We also study the structure of the solution set of (1.1) in case the favorable domain $\Omega^{+}$has several components. We make the following assumption on $\Omega^{+}$:

$$
\left\{\begin{array}{l}
(1.4) \text { holds, } \Omega^{+} \text {has } k<\infty \text { connected components with } \Omega^{+}=\bigcup_{i=1}^{k} \Omega_{i}^{+},  \tag{1.7}\\
\text {and each connected component } \Omega_{i}^{+} \text {satisfies an interior ball condition. }
\end{array}\right.
$$

Set $M=\{1,2,3, \ldots, k\}$. Under (1.7), for any solution $u(x)$ of (1.1), by the strong maximum principle it is easy to see that solution $u(x)$ is either positive in $\Omega_{i}^{+}$or completely vanishes in $\Omega_{i}^{+}$for any $i \in M$ (see Lemma 3.11). We make the following definition:

Definition 1.4. (1) For any non-empty subset $I \subset M=\{1,2, \ldots, k\}$, denote by $S_{I}$ the class of solutions of (1.1) which are positive in $\Omega_{I}^{+}=\bigcup_{i \in I} \Omega_{i}^{+}$.
(2) $N_{I}$ denotes the set $\left\{u \in S_{I} \mid u \equiv 0\right.$ in $\left.\Omega^{+}-\Omega_{I}^{+}\right\}$.

For any non-empty $I \subset M$, a solution in $N_{I}$ vanishes identically in $\Omega^{+}-\Omega_{I}^{+}$, part of the favorable set $\Omega^{+}$. We will call such solutions dead core solutions, in analogy with the classical use of the term [ $2,5,19$ ]. This is a slight abuse of this term: in previous usage, it refers to solutions of boundary-value problems which vanish on compact subdomains, while our solutions will vanish on much larger, non-compact regions. Nevertheless, the phenomena are clearly related via the failure of the strong maximum principle for degenerate equations or non-Lipschitz nonlinearities.

If dead core solutions (as defined above) do exist, the class $S_{I}$ can contain many elements: see the following proposition and Theorem 1.9. However, in many cases the class $N_{I}$ can have at most one solution. Following the idea in [3], we present a generalization of the uniqueness result of Spruck [23]:

Theorem 1.5. Assume (1.7), if $p \geqslant 1$, then the number of elements in $N_{I}$ is at most 1 for any non-empty I. In particular if $k=1$, then the solution to (1.1) is unique and its support is connected.

We may also prove uniqueness of solutions in class $N_{I}$ with $q<p<1$ under some additional hypotheses on $b(x)$; see Theorem 3.14.

When Theorem 1.5 applies, the solution space of (1.1) is completely characterized by the support properties of the solutions. Consider the following example: let $\Omega_{i}^{+}, i=1, \ldots, k$, be any smooth,


Fig. 1. $a(x)$ as in the example (1.8).
compact and connected open sets in $\mathbb{R}^{n}$, with $\min _{i \neq j} \operatorname{dist}\left(\Omega_{i}^{+}, \Omega_{j}^{+}\right)>0$. Let $b(x) \equiv-1$ and define $a(x)=a_{\lambda}(x)$ by

$$
a(x)= \begin{cases}\lambda, & \text { if } x \in \Omega_{i}^{+}, i=1, \ldots, k  \tag{1.8}\\ -1, & \text { if } x \notin \bigcup_{i=1}^{k} \Omega_{i}^{+}\end{cases}
$$

where $\lambda>0$ is a fixed constant. (See Fig. 1.) Combining the results of Corollary 1.3 and Theorem 1.5, we have:

Proposition 1.6. Assume $a(x)$ is defined as in (1.8) with fixed $\lambda>0, b(x) \equiv-1$, and $p \geqslant 1$.

1. There exists $\delta^{*}=\delta^{*}\left(\Omega^{+}, \lambda\right)>0$ so that if $\min _{i \neq j} \operatorname{dist}\left(\Omega_{i}^{+}, \Omega_{j}^{+}\right) \geqslant \delta^{*}$, then $N_{I}$ contains exactly one solution for each $I \neq \emptyset$.
2. There exists $\delta_{*}=\delta_{*}\left(\Omega^{+}, \lambda\right)>0$ so that if $\max _{i \neq j} \operatorname{dist}\left(\Omega_{i}^{+}, \Omega_{j}^{+}\right) \leqslant \delta_{*}$, then (1.1) admits exactly one solution in $\mathbb{R}^{n}$. This solution is positive on the set $\Omega^{0+}$.

Note that in case (1), Eq. (1.1) admits exactly $2^{k}-1$ nontrivial solutions in all.
In the general case, we can assert the existence of a minimal and maximal element in $S_{I}$ for each $I$ :

Theorem 1.7. (a) Under the hypothesis (1.7), $S_{I}$ has a minimal element $u_{I} \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{\infty}\left(\mathbb{R}^{n}\right)$ for any non-empty $I \subset M$.
(b) Assuming (1.2), (1.1) always has a maximal solution $U$ in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{\infty}\left(\mathbb{R}^{n}\right)$, which is positive in $\Omega^{+}$ and bounded above by $\omega$ from (1.3). In particular under (1.7), $S_{I}$ has the same maximal element $U$ for any non-empty $I \subset M$.

There is a connection between the maximal solution of (1.1) and the minimal energy solution of the functional $E: \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{q+1}\left(\mathbb{R}^{n}\right) \rightarrow \mathbb{R}$, defined by

$$
E(v)=\int_{\mathbb{R}^{n}}|\nabla v|^{2} d x-\frac{1}{q+1} \int_{\mathbb{R}^{n}} a\left(v^{+}\right)^{q+1} d x-\frac{1}{p+1} \int_{\mathbb{R}^{n}} b\left(v^{+}\right)^{p+1} d x
$$

If we assume $a(x)$ and $b(x)$ are uniformly bounded in $\mathbb{R}^{n}$, then $E$ is smooth (see [12]). The interesting fact is that $\inf _{v \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{q+1}\left(\mathbb{R}^{n}\right)} E(v)$ is achieved at a non-negative function $U$, which is a solution of (1.1). By minimization it is easy to see that $U>0$ in all connected components of $\Omega^{+}$, that is $U \in S_{M}$.

By Theorem 1.5, when $p \geqslant 1$ this minimal energy solution is the (unique) maximal solution in $S_{M}$. On the other hand, as pointed out in [1], if there do exist dead core solutions (in some class $N_{I}, I \neq \emptyset$ ) then these solutions cannot be energy minimizers even modulo any finite dimensional subspace. As a result, our study of dead core solutions to (1.1) will rely on comparison arguments, monotone iteration, and a priori estimates.

Another way to influence the support properties of the solutions is by varying the relative strengths of the positive and negative parts of $a(x)$. We introduce a parameter $\lambda>0$, and set

$$
a_{\lambda}(x)=\lambda a^{+}(x)-a^{-}(x) .
$$

Clearly this does not affect the geometry of the favorable and unfavorable regions, and $\Omega^{+}, \Omega^{0+}$ and $\Omega^{-}$remain the same for the whole family of $a_{\lambda}$. Intuitively, we expect that the size of the support of the solutions of

$$
\begin{equation*}
-\Delta u=a_{\lambda}(x) u^{q}+b u^{p} \quad \text { in } \mathbb{R}^{n}, u \geqslant 0 \text { and } u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right), \tag{1.9}
\end{equation*}
$$

should grow with increasing $\lambda$. We will study the asymptotic behavior of solutions for large and small $\lambda$, but we first require a further condition on the geometry of the set $\Omega^{0+}$ :

Definition 1.8. We say $a(x)$ is admissible if (1.7) holds, and:

1. $\Omega^{0+}$ also has exactly $k$ connected components with $\Omega^{0+}=\bigcup_{i=1}^{k} \Omega_{i}^{0+}$,
2. $\Omega_{i}^{+} \subset \Omega_{i}^{0+}$ for $i \in M$,
3. $\operatorname{dist}\left(\Omega_{i}^{0+}, \Omega_{j}^{0+}\right)>0$ for $i \neq j$.

For admissible $a$ we have the following theorem:
Theorem 1.9. Assume $a(x)$ is admissible, there exists $\lambda_{*}>0$ so that for all $\lambda<\lambda_{*}$, all solutions of (1.9) are compactly supported and $N_{I}^{q} \neq \emptyset$ for any non-empty collection $I \subset M$.

Note that for $\lambda<\lambda_{*}$ the compact support of the solutions follows even in the absence of asymptotic conditions on $a(x)$, showing that condition (1.5) cannot be sharp. To emphasize the dependence on $\lambda$, Eq. (1.9) is often referred as (1.9) (the subscript $\lambda$ is omitted if no confusion arises). For $\lambda$ big, we have the following theorem in case $q<p \leqslant 1$.

Theorem 1.10. Assume (1.7), if $q<p \leqslant 1$, there exists $\lambda_{1}^{*}>0$ so that Eq. (1.9) has a unique solution $u_{\lambda}$, which is positive in $\overline{\Omega^{+}}$for all $\lambda>\lambda_{1}^{*}$. Moreover,

$$
\begin{cases}\text { if } p<1, & \lim _{\lambda \rightarrow \infty} \lambda^{\frac{1}{q-1}} u_{\lambda}=\omega, \\ \text { if } p=1, & \text { where } \omega \text { is from (1.3), } \\ \lim _{\lambda \rightarrow \infty} \lambda^{\frac{1}{q-1}} u_{\lambda}=\bar{\omega}, & \text { where } \bar{\omega} \text { is in (4.3). }\end{cases}
$$

In particular, combining Theorem 1.2 and the above theorem we conclude that as $\lambda \rightarrow \infty$ the support must grow.

Corollary 1.11. Assume (1.7), if $q<p \leqslant 1$ and $\liminf _{|x| \rightarrow \infty} a^{-}|x|^{(n-2)(1-q)}=\infty$, then there exists $\lambda_{2}^{*}>0$ so that problem (1.9) has a unique compactly supported solution $u_{\lambda}$ with $u_{\lambda}>0$ in $\Omega^{+}$for all $\lambda \geqslant \lambda_{2}^{*}$. Moreover, $\operatorname{supp}\left(u_{\lambda}\right)$ expands to infinity as $\lambda$ goes to infinity, that is

$$
\bigcup_{\lambda>\lambda_{2}^{*}} \operatorname{supp}\left(u_{\lambda}\right)=\mathbb{R}^{n} .
$$

For the case $p>1$, the asymptotic behavior is more complicated, and depends strongly on the form of $b(x)$. Some specific results are proven in Theorem 4.7 in Section 4.

To illustrate our results on the parametrized problem (1.9), we return to the previous piecewise constant $a(x)$ from our example (1.8). For simplicity, assume $b(x) \equiv 0$. For $\lambda$ small enough, Theorem 1.9 applies and we conclude that (1.9) admits a unique solution in $N_{I}$ for all $I \neq \emptyset$, so (1.9) admits exactly $2^{k}-1$ solutions in all, and all but one has dead cores. For $\lambda$ sufficiently large, by Theorem 1.10 the equation has exactly one solution which is positive in all of $\Omega^{+}$.

The existence of compactly supported solutions for equations on $\mathbb{R}^{n}$ with non-Lipschitz nonlinearities was originally proven by Schatzman [22], using Puel's existence theorem [20]. Spruck [23] imposed a monotonicity condition, $x \cdot \nabla a(x)<0$, and considered (1.1) with $b(x) \equiv 0$. Spruck proved uniqueness of compactly supported solutions by means of a special version of Hopf's boundary lemma. He also proved that, under the same hypotheses, the support of the solution is star-shaped with Lipschitz boundary. In [3,4] C. Bandle, M.A. Pozio and A. Tesei studied dead core solutions for this problem in a bounded domain with both Dirichlet and Neumann boundary condition, and S. Alama [1] used a bifurcation analysis for dead cores in the Neumann problem for similar equations. Most other previous work on equations of the form (1.1) has been for bounded domains or for constant coefficient equations in the whole space $\mathbb{R}^{n}$ : see $[6,9,11,13,15,17,18]$ and the references therein.

Our method for proving compact support is inspired by the paper of Cortázar, Elgueta and Felmer [11] on the constant coefficient equation $-\Delta u=u^{p}-u^{q}$ in $\mathbb{R}^{n}$. In this paper, they prove that all $H^{1}$ solutions of the equation must have compact support. (See Gui [15] for an alternative approach to this problem.) In this paper we follow [11] to prove Theorem 1.2, by deriving a priori estimates and applying comparison theorems. This is the content of Section 2. In Section 3 we deal with existence and uniqueness questions and prove Theorems 1.5 and 1.7 and Proposition 1.6. Section 4 concerns the parametrized problems (1.9). Finally, in Section 5 we discuss solutions which are not in the space $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ and present some examples related to our results.

Finally, we thank the reviewer for the care and patience taken in reading this paper and for the many helpful suggestions.

## 2. Compact support

In this section we prove Theorem 1.2. The method we use is derived from the approach of Cortázar, Elgueta and Felmer [11] on the constant coefficient equation $-\Delta u=u^{p}-u^{q}$.

First we develop a few useful lemmas. Assume (1.4), and pick $\rho_{1}>0$ such that $\Omega^{0+} \Subset B\left(0, \rho_{1}\right)$.
Lemma 2.1. Assume (1.4). Then, any weak solution $u$ of (1.1) is a classical solution and $\lim _{|x| \rightarrow \infty} u(x)=0$.
Proof. The regularity of $u$ follows from standard bootstrap arguments; see Appendix B in Struwe [24] or Theorem 0 in Brezis [7]. Since $u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$, then $u \in L^{2^{*}}\left(\mathbb{R}^{n}\right)$. Hence for any $\epsilon>0$, there exists $R(\epsilon)>\rho_{1}$ such that

$$
\|u\|_{L^{2^{*}}\left(\mathbb{R}^{n}-B(0, R)\right)}<\epsilon \text { for all } R>R(\epsilon) .
$$

So for any $x \in \mathbb{R}^{n}-B(0, R(\epsilon)+2)$, we have $B(x, 1) \Subset \mathbb{R}^{n}-B(0, R(\epsilon))$ and $-\Delta u(y) \leqslant 0$ in $B(x, 1)$. Therefore, by the mean value property of subharmonic function, we conclude:

$$
0 \leqslant u(x) \leqslant \frac{1}{|B(x, 1)|} \int_{B(x, 1)} u(y) d y \leqslant C\|u\|_{L^{2^{*}}(B(x, 1))} \leqslant C \epsilon,
$$

that is $\lim _{|x| \rightarrow \infty} u(x)=0$.
Remark 2.2. Note that we only need to assume the minimal hypothesis (1.2) and $u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$. On the other hand, obtaining $\lim _{|x| \rightarrow \infty} u(x)=0$ is very crucial to the succeeding arguments. In Section 5, we will discuss solutions whose Dirichlet energy is not assumed to be finite.

The next lemma shows that $u$ not only uniformly tends to zero, but also goes to zero with certain speed, as $|x|$ goes to infinity.

Lemma 2.3. Assume (1.4), then $u(x) \leqslant \frac{C}{|x|^{n-2}}$, where $C=\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \rho_{1}^{n-2}$ and $x \in \mathbb{R}^{n}-B\left(0, \rho_{1}\right)$.
Proof. Let $v=\frac{C}{|x|^{n-2}}$, then by the special choice of $C$ we have

$$
-\Delta v(x)=0 \quad \text { in } \mathbb{R}^{n}-B\left(0, \rho_{1}\right) \quad \text { and } \quad v(x) \geqslant u(x) \quad \text { on } \partial B\left(0, \rho_{1}\right)
$$

Now consider $w=u-v$, then $w$ satisfies:

$$
-\Delta w(x) \leqslant 0 \quad \text { in } \mathbb{R}^{n}-B\left(0, \rho_{1}\right) \quad \text { and } \quad w(x) \leqslant 0 \quad \text { on } \partial B\left(0, \rho_{1}\right)
$$

Moreover we see that $\lim _{|x| \rightarrow \infty} w(x)=0$. We now claim that $w(x) \leqslant 0$ in $\mathbb{R}^{n}-B\left(0, \rho_{1}\right)$. Indeed, otherwise there would exist $x_{0} \in \mathbb{R}^{n}-\overline{B\left(0, \rho_{1}\right)}$ such that $u\left(x_{0}\right)>v\left(x_{0}\right)>0$. We also notice that $w(x) \leqslant 0$ on $\partial B\left(0, \rho_{1}\right)$ and $\lim _{|x| \rightarrow \infty} w(x)=0$, we may assume $w$ attains its maximum at $x_{0}$. Therefore we would have

$$
0 \leqslant-\Delta w\left(x_{0}\right)=-\Delta u\left(x_{0}\right)<0
$$

a contradiction. Hence, we have $u \leqslant v=\frac{C}{|x|^{n-2}}$ in $\mathbb{R}^{n}-B\left(0, \rho_{1}\right)$.
We must estimate the $L^{\infty}$-norm of the solution. We prove:

Proposition 2.4. Assume (1.4). There exists a constant $C$ so that for any solution $u$ of (1.1) we have

$$
\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \leqslant C\left(q, p, n,\left\|a^{+}\right\|_{L^{\infty}\left(\Omega^{+}\right)}, \Omega^{+}\right),
$$

where this $C$ tends to zero as $\left\|a^{+}\right\|_{L^{\infty}\left(\Omega^{+}\right)}$tends to zero.
The maximum principle yields:
Lemma 2.5. Assume (1.2), then $\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$ is attained in $\overline{\Omega^{+}}$, i.e. there exists $x_{0} \in \overline{\Omega^{+}}$such that $\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}=$ $u\left(x_{0}\right)$.

Proof. Since $\lim _{|x| \rightarrow \infty} u(x)=0$, we may assume $\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$ is attained at $x_{1}$, which is not in $\overline{\Omega^{+}}$. Let $\Omega$ be the connected component of $\mathbb{R}^{n}-\overline{\Omega^{+}}$, which contains $x_{1}$. By the strong maximum principle, $u(x)=u\left(x_{1}\right)$ in $\bar{\Omega}$. Since $\overline{\Omega^{+}} \cap \bar{\Omega}$ is not empty, we are done.

We next estimate the Dirichlet energy of the solutions. The essential step in the proof (see Claim below) will also be useful in our existence proofs in the next section:

Lemma 2.6. Assume (1.4), then we have

$$
\int_{\mathbb{R}^{n}}|\nabla u|^{2} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} u^{q+1} d x
$$

Proof. Since $u$ satisfies the equation $-\Delta u=a u^{q}+b u^{p}$, multiply both sides of the equation by $u$ and integrate by parts. We have for $R \geqslant \rho_{1}$

$$
\begin{align*}
\int_{B(0, R)}|\nabla u|^{2} d x-\int_{\partial B(0, R)} \frac{\partial u}{\partial n} u d S & =\int_{B(0, R)} a^{+} u^{q+1} d x+\int_{B(0, R)} b u^{p+1} d x \\
& \leqslant \int_{\mathbb{R}^{n}} a^{+} u^{q+1} d x . \tag{2.1}
\end{align*}
$$

Claim. There exists a sequence $\left\{R_{n}\right\}$ and $\lim _{n \rightarrow \infty} R_{n}=\infty$, such that $\int_{\partial B\left(0, R_{n}\right)} \frac{\partial u}{\partial n} u d S \rightarrow 0$ as $n \rightarrow \infty$.
Indeed, we have the following estimate:

$$
\begin{aligned}
\left|\int_{\partial B(0, R)} \frac{\partial u}{\partial n} u d S\right| & \leqslant \frac{\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \rho_{1}^{n-2}}{R^{n-2}} \int_{\partial B(0, R)}|\nabla u| d S \\
& \leqslant \frac{\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \rho_{1}^{n-2}}{R^{n-2}}\|1\|_{L^{2}(\partial B(0, R))}\|\nabla u\|_{L^{2}(\partial B(0, R))} \\
& \leqslant \frac{C R^{\frac{n-1}{2}}}{R^{n-2}}\|\nabla u\|_{L^{2}(\partial B(0, R))} \\
& \leqslant \frac{C}{R^{\frac{n-3}{2}}}\|\nabla u\|_{L^{2}(\partial B(0, R))} .
\end{aligned}
$$

Notice $\infty>\|\nabla u\|_{L^{2}\left(\mathbb{R}^{n}\right)}^{2}=\int_{0}^{\infty} \int_{\partial B(0, R)}|\nabla u|^{2} d S d r$, so there exists a sequence $\left\{R_{n}\right\}$ with $\lim _{n \rightarrow \infty} R_{n}=$ $\infty$, such that $\|\nabla u\|_{L^{2}\left(\partial B\left(0, R_{n}\right)\right)} \rightarrow 0$ as $n \rightarrow \infty$. Therefore

$$
\int_{\partial B\left(0, R_{n}\right)} \frac{\partial u}{\partial n} u d S \rightarrow 0 \quad \text { as } n \rightarrow \infty
$$

and the claim is proven.
Applying this to (2.1), we have

$$
\int_{\mathbb{R}^{n}}|\nabla u|^{2} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} u^{q+1} d x
$$

The next step is a bootstrap argument. Recall $\rho_{1}$ is chosen such that $\Omega^{0+} \Subset B\left(0, \rho_{1}\right)$.
Lemma 2.7. Assume (1.2). For any positive integer $s \geqslant 2$, there exists a constant

$$
C^{\prime}=C^{\prime}\left(q, p, n, s,\left\|a^{+}\right\|_{L^{\infty}\left(\Omega^{+}\right)},\|u\|_{L^{2^{*}}\left(\mathbb{R}^{n}\right)}, \Omega^{+}\right)
$$

so that

$$
\|u\|_{L^{(s+1) \frac{n}{n-2}\left(\mathbb{R}^{n}\right)}} \leqslant C^{\prime}
$$

Moreover, $C^{\prime}$ tends to zero as $\left\|a^{+}\right\|_{L^{\infty}\left(\Omega^{+}\right)}$and $\|u\|_{L^{2^{*}}\left(\mathbb{R}^{n}\right)}$ tend to zero.

Proof. We rewrite Eq. (1.1) as

$$
-\Delta u+a^{-} u^{q}-b u^{p}=a^{+} u^{q}
$$

Then we just follow the steps in Appendix B in Struwe [24] (or Theorem 0 in Brezis [7]). The term $a^{-} u^{q}-b u^{p}$ are non-negative and so they may be neglected. Note that the existence of the integrals is assured a priori without truncation, since $a^{+}$is continuous and has compact support.

The uniform estimate on $u$ will be derived by comparison with the solution $\omega \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ of

$$
\begin{equation*}
-\Delta \omega=a^{+}(x) \omega^{q}, \quad \omega(x) \rightarrow 0 \quad \text { as }|x| \rightarrow \infty \tag{2.2}
\end{equation*}
$$

The existence and uniqueness of $\omega \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ follows from Theorem $2^{\prime}$ in Brezis and Kamin [8], using only the hypothesis that $\Omega^{+}$is bounded. The decay at infinity is also clear from Lemma 2.1. Note that the conclusions of Lemmas $2.5,2.6$, and 2.7 also hold with $\omega$ replacing $u$.

We are now ready to complete the proof of Proposition 2.4.
Proof of Proposition 2.4. As remarked above, $\omega$ satisfies the conclusions of the above Lemmas 2.5 , 2.6 and 2.7. From Lemma 2.6, we know that

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|\nabla \omega|^{2} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x \tag{2.3}
\end{equation*}
$$

By the Sobolev embedding we also have

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|\nabla \omega|^{2} d x \geqslant C\|\omega\|_{L^{2^{*}}\left(\mathbb{R}^{n}\right)}^{2} \tag{2.4}
\end{equation*}
$$

for some constant $C$ independent of $\omega$. Also by Hölder we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x \leqslant\left\|a^{+}\right\|_{L^{t}\left(\mathbb{R}^{n}\right)}\left\|\omega^{q+1}\right\|_{L^{2^{*}}\left(\mathbb{R}^{n+1}\right.}=\left\|a^{+}\right\|_{L^{t}\left(\mathbb{R}^{n}\right)}\|\omega\|_{L^{*}\left(\mathbb{R}^{n}\right)}^{q+1} \tag{2.5}
\end{equation*}
$$

where $t$ is the conjugate of $\frac{2^{*}}{q+1}$. Therefore combine (2.3), (2.4) and (2.5), we get

$$
C\|\omega\|_{L^{2^{*}}\left(\mathbb{R}^{n}\right)}^{2} \leqslant\left\|a^{+}\right\|_{L^{t}\left(\mathbb{R}^{n}\right)}\|\omega\|_{L^{2}\left(\mathbb{R}^{n}\right)}^{q+1},
$$

that is

$$
\|\omega\|_{L^{2}}\left(\mathbb{R}^{n}\right)<C\left(\left\|a^{+}\right\|_{L^{t}\left(\mathbb{R}^{n}\right)}\right)^{\frac{1}{1-q}}
$$

Choosing $s$ so that $(s+1) \frac{n}{n-2} \geqslant \frac{(n+1)}{2}$, from Lemma 2.7, we know that $\|\omega\|_{L^{(s+1)} \frac{n}{n-2}\left(\mathbb{R}^{n}\right)}$ is uniformly bounded. Apply the standard elliptic estimate (see [14]) on the domain $\Omega^{+} \Subset \hat{O}=\left\{x \in \mathbb{R}^{n} \mid\right.$ $\left.\operatorname{dist}\left(x, \Omega^{0+}\right)<\epsilon\right\}$, where $\epsilon$ is chosen so small that $\left\|a^{+}\right\|_{L^{\infty}(\hat{O})}=\left\|a^{+}\right\|_{L^{\infty}\left(\Omega^{+}\right)}$and volume $(\hat{O}) \leqslant$ 2volume ( $\Omega^{+}$), we have

$$
\|\omega\|_{W^{2, \frac{n+1}{2}}\left(\Omega^{+}\right)} \leqslant C\left(\|\Delta \omega\|_{L^{\frac{n+1}{2}}(\hat{o})}+\|\omega\|_{L^{\frac{n+1}{2}}(\hat{o})}\right) .
$$

Then by Sobolev embedding theorem, in view of Lemma 2.5 we have shown that there exists a constant $C^{\prime \prime}=C^{\prime \prime}\left(q, p, n,\left\|a^{+}\right\|_{\infty}, \Omega^{+}\right)$so that $\|\omega\|_{\infty} \leqslant C^{\prime \prime}$.

To conclude, we use the following comparison result, which will be useful throughout the paper:

Lemma 2.8. Assume (1.7), and let $u_{1}, u_{2} \in C_{\text {loc }}^{1}\left(\mathbb{R}^{n}\right) \cap W_{\text {loc }}^{2, s}\left(\mathbb{R}^{n}\right)$, $s>n$, be two functions such that for some $I \subset M$,

1. $u_{1}, u_{2}$ are positive in $\overline{\Omega_{I}^{+}}$;
2. $u_{1}=0$ in $\overline{\Omega^{+}-\Omega_{I}^{+}}$;
3. $\lim _{|x| \rightarrow \infty} u_{1}(x)=\lim _{|x| \rightarrow \infty} u_{2}(x)=0$;
4. For a.e. $x \in \mathbb{R}^{n}$,

$$
\begin{aligned}
& -\Delta u_{1} \leqslant a u_{1}^{q}+b u_{1}^{p}, \\
& -\Delta u_{2} \geqslant a u_{2}^{q}+b u_{2}^{p} .
\end{aligned}
$$

Then we must have $u_{1} \leqslant u_{2}$ in $\mathbb{R}^{n}$.
This lemma is proven for $C^{2}$ solutions in a bounded domain in [3]. Essentially the same proof shows that it is true for solutions in $C_{l o c}^{1}\left(\mathbb{R}^{n}\right) \cap W_{l o c}^{2, s}\left(\mathbb{R}^{n}\right)$ for $s>n$ by using Serrin's maximum principle [21], but for completeness we include the proof in Appendix A.

Applying Lemma 2.8 we thus obtain $0 \leqslant u(x) \leqslant \omega(x)$ holds for all $x \in \mathbb{R}^{n}$, for any solution $u$ of (1.1), and the proposition is proven.

Corollary 2.9. Assume (1.4), then

$$
u(x) \leqslant \frac{\eta_{1} \rho_{1}^{n-2}}{|x|^{n-2}}
$$

for any $x \in \mathbb{R}^{n}-B\left(0, \rho_{1}\right)$ and $\eta_{1}$ is a number depending on $q, p, n, \Omega^{+}$and $\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$.
Now we are ready to prove Theorem 1.2.
Proof of Theorem 1.2. This is a comparison argument using the method from [11]. For positive numbers $M, c$, let $w(s)$ be the function defined implicitly by

$$
\int_{w(s)}^{M} \frac{d t}{\sqrt{\frac{c}{q+1} t^{q+1}}}=\sqrt{2} s,
$$

with constants $M$ and $c$ to be chosen later. Indeed we can write $w(s)$ explicitly in terms of $s$

$$
w(s)=\left[M^{1-\frac{q+1}{2}}-\sqrt{2} s\left(1-\frac{q+1}{2}\right) \sqrt{\frac{c}{q+1}}\right]^{\frac{2}{1-q}} .
$$

Notice that since $0<q<1$, then $\frac{2}{1-q}>2$. So $w(s)$ is at least twice continuously differentiable in $[0, B]$, where $B$ is defined by

$$
M^{1-\frac{q+1}{2}}=\sqrt{2} B\left(1-\frac{q+1}{2}\right) \sqrt{\frac{c}{q+1}}
$$

It is easy to see that $w(s)$ satisfies

$$
w^{\prime \prime}(s)-c w^{q}(s)=0 \quad \text { in }(0, B)
$$

Moreover $w(s)$ is a decreasing function in $s, w(B)=w^{\prime}(B)=w^{\prime \prime}(B)=0$. Therefore, by defining $w(s) \equiv 0$ for $s \in[B, \infty)$, we obtain a nonincreasing solution of

$$
w^{\prime \prime}(s)-c w^{q}(s)=0 \quad \text { in }(0, \infty)
$$

with $w(0)=M$ and $\operatorname{supp}(w)=[0, B]$.
We know from the previous corollary that $u(x) \leqslant \frac{\eta_{1} \rho_{1}^{n-2}}{|x|^{n-2}}$ for $|x|>\rho_{1}$. Consider the function $f(s):[0,1] \rightarrow \mathbb{R}$ defined by

$$
f(s)=w^{\frac{1-q}{2}}(s)=M^{1-\frac{q+1}{2}}-\sqrt{2} s\left(1-\frac{q+1}{2}\right) \sqrt{\frac{c}{q+1}}
$$

where for some constant $\rho>\rho_{1}+1$, we choose our values of $M, c$ as

$$
M=\frac{\eta_{1} \rho_{1}^{n-2}}{(\rho-1)^{n-2}} \quad \text { and } \quad c=\frac{\eta}{\rho^{(n-2)(1-q)}},
$$

and $\eta$ is such that

$$
\sqrt{2}\left(1-\frac{q+1}{2}\right) \sqrt{\frac{\eta}{q+1}}=\left(\eta_{1} \rho_{1}^{n-2}\right)^{1-\frac{q+1}{2}}+1 .
$$

We now claim that there exists $\rho>\rho_{1}+1$ such that

$$
a^{-} \geqslant \frac{\eta}{|x|^{(n-2)(1-q)}} \quad \text { for }|x| \geqslant \rho-1 \quad \text { and } \quad f(1)<0
$$

Actually we can rewrite $f(1)$ as the following form

$$
f(1)=\frac{1}{\left(\rho-1 \frac{(n-2)(1-q)}{2}\right.}\left[\left(\eta_{1} \rho_{1}^{n-2}\right)^{1-\frac{q+1}{2}}-\sqrt{2}\left(1-\frac{q+1}{2}\right) \sqrt{\frac{\eta}{q+1}}\left(\frac{\rho-1}{\rho}\right)^{\frac{(n-2)(1-q)}{2}}\right] .
$$

By the assumption that $\liminf _{|x| \rightarrow \infty} a^{-}|x|^{(n-2)(1-q)}>\eta$ and the choice of $\eta$, the claim is true for some large $\rho$.

Since $f(0)>0$ and $f(1)<0$, then according to the Mean Value Theorem

$$
0<\sup _{0 \leqslant t \leqslant 1}\{t \mid f(s) \geqslant 0 \text { for } s \in[0, t]\}<1 .
$$

Therefore for the choice of $M$ and $c, B$ is well defined and $0<B<1$.
Let $v(x)=w(|x|-(\rho-1))$, then we see that $v$ satisfies

$$
\begin{gathered}
\Delta v-c v^{q} \leqslant 0 \quad \text { in } \mathbb{R}^{n}-B(0, \rho-1), \\
v=M \quad \text { on } \partial\left(\mathbb{R}^{n}-B(0, \rho-1)\right) .
\end{gathered}
$$

Also notice that for $|x| \in[\rho-1, \rho), a^{-}(x) \geqslant \frac{\eta}{|x|^{(n-2)(1-q)}}>\frac{\eta}{\rho^{(n-2)(1-q)}}=c$.
For $u$ we have

$$
\begin{gathered}
\Delta u+a u^{q}+b u^{p}=0 \text { in } \mathbb{R}^{n}-B(0, \rho-1), \\
u \leqslant M \text { on } \partial\left(\mathbb{R}^{n}-B(0, \rho-1)\right) .
\end{gathered}
$$

By subtracting them, we have

$$
-\Delta(v-u) \geqslant-a(x) u^{q}-c v^{q} \quad \text { for } x \in \mathbb{R}^{n}-B(0, \rho-1) .
$$

We now claim that $v \geqslant u \geqslant 0$ for $x \in \mathbb{R}^{n}-B(0, \rho-1)$. Otherwise there would exist $x_{0} \in\left(\mathbb{R}^{n}-\right.$ $\overline{B(0, \rho-1)})$ such that $u\left(x_{0}\right)>v\left(x_{0}\right) \geqslant 0$, which implies that $v-u$ attains its global minimal value at some point $x_{0} \in \mathbb{R}^{n}-\overline{B(0, \rho-1)}$. At $x_{0}$,

$$
\begin{aligned}
0 & \geqslant-\Delta(v-u)\left(x_{0}\right) \\
& \geqslant\left(-a\left(x_{0}\right) u^{q}\left(x_{0}\right)-c v^{q}\left(x_{0}\right)\right) \\
& \geqslant \begin{cases}-a\left(x_{0}\right) u^{q}\left(x_{0}\right)>0 & \text { if } v\left(x_{0}\right)=0, \\
\left(-a\left(x_{0}\right)-c\right) v^{q}\left(x_{0}\right)>0 & \text { if } v\left(x_{0}\right)>0,\end{cases}
\end{aligned}
$$

a contradiction, and so the claim is proven.
So we must have $v \geqslant u \geqslant 0$ for $x \in\left(\mathbb{R}^{n}-B(0, \rho-1)\right)$, which implies $u$ has compact support. Therefore supp $u \Subset B(0, \rho)$.

In the end we note that the main ingredient in the above proof is the decay estimate on the solution in the exterior of $B\left(0, \rho_{1}\right)$. Any improvement on the required decay (1.5) of $a^{-}(x)$ would require a sharper estimate in Corollary 2.9.

Remark 2.10. As noted in the Introduction, the one- and two-dimensional cases must be handled differently. In particular, the decay estimates of Lemma 2.3 are no longer valid in $\mathbb{R}^{n}, n=1,2$, principally because the fundamental solution does not decay to zero at infinity. Nevertheless, Lemma 2.1 still holds for any classical solution in $L^{t}\left(\mathbb{R}^{n}\right)$ for $t \geqslant 1$, so we can prove in very similar way that all classical solutions of (1.1) in $L^{t}\left(\mathbb{R}^{n}\right)$ have compact supports under strong assumption $\liminf |x| \rightarrow \infty a^{-}(x)>0$. However, we cannot uniformly control the size of the support because the Sobolev inequalities are domain-dependent in dimensions $n=1,2$. The statement we can make in any dimension is the following:

Theorem 2.11. Under (1.4), if $\liminf _{|x| \rightarrow \infty} a^{-}>0$, all classical solutions in $L^{t}\left(\mathbb{R}^{n}\right)$ for $t \geqslant 1$ must have compact support.

Proof. The proof is much simpler since we do not need to choose the place where we make the comparison. We may just pick $M=\|u\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$ and compare $w$ with $u$ outside $B\left(0, \rho_{1}\right)$.

Remark 2.12. If $a^{-}(x)$ decays too fast at infinity, solutions may not have compact support. To see this, recall that in a bounded domain with Neumann condition there is a necessary condition for the existence of solutions (see [1,4]),

$$
\begin{equation*}
\int_{\operatorname{supp}(U)} a(x)+b U^{p-q} d x<0 . \tag{2.6}
\end{equation*}
$$

When Eq. (1.1) has solutions with compact support, they also solve the Neumann problem inside a big ball, so this condition must hold a posteriori. However, if we choose $b \equiv 0$ and $a(x) \in L^{1}\left(\mathbb{R}^{n}\right)$ satisfying (1.4) with $\int_{\mathbb{R}^{n}} a(x)>0$ then a compactly supported solution could never satisfy (2.6), and thus no solution can have compact support.

## 3. Existence and uniqueness

In this section we present the proof of the basic existence theorems, Theorems 1.1 and 1.7 , and a more general form of the uniqueness result Theorem 1.5. Throughout we assume the dimension $n \geqslant 3$.

### 3.1. Existence of maximal and minimal solutions

We use the method of monotone iteration (see Hess [16]). Existence results of this type are well known in the setting of bounded domains (see [3], for example) and we adapt the technique here for entire solutions in $\mathbb{R}^{n}$.

Recall from (2.2) that $\omega \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ is the unique solution of the equation

$$
-\Delta \omega=a^{+} \omega^{q}, \quad \omega(x) \rightarrow 0 \quad \text { as }|x| \rightarrow \infty .
$$

To begin the iteration, we start with the following equation:

$$
\begin{equation*}
-\Delta z+a^{-} z^{q}-b z^{p}=a^{+} f^{q} \quad \text { in } \mathbb{R}^{n}, z \geqslant 0 \text { and } z \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right), \tag{3.1}
\end{equation*}
$$

where $f$ is Hölder continuous and $0 \leqslant f \leqslant \omega$.
Lemma 3.1. Assume (1.2), then (3.1) has a solution $Z \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) \cap L^{\infty}\left(\mathbb{R}^{n}\right)$ and $Z \leqslant \omega$.
Proof. We use monotone iteration to prove this result. Let us consider the following equation

$$
\begin{equation*}
-\Delta z_{n}+a^{-} z_{n}^{q}-b z_{n}^{p}=a^{+} f^{q} \quad \text { in } B(0, n), z \geqslant 0 \text { and } z \in H_{0}^{1}(B(0, n)) . \tag{3.2}
\end{equation*}
$$

It is easy to see that $\underline{z}=0$ is a subsolution, $\bar{z}=\int_{\mathbb{R}^{n}} \Phi(x-y) a^{+}(y) f^{q}(y) d y$ is a supersolution, where $\Phi$ is the fundamental solution of Laplace's equation. Furthermore we see that $\bar{z} \leqslant \omega$ since $0 \leqslant f \leqslant \omega$. Therefore by sub-supersolution method the above equation has a solution $Z_{n} \leqslant \omega$.

Claim 1. The solution $Z_{n}$ is unique.
Indeed if there is another solution $\overline{Z_{n}}$, then we have

$$
-\Delta\left(Z_{n}-\overline{Z_{n}}\right) d x+a^{-}\left(Z_{n}^{q}-{\overline{Z_{n}}}^{q}\right) d x-b\left(Z_{n}^{p}-{\overline{Z_{n}}}^{p}\right) d x=0 .
$$

Then multiply both sides by $Z_{n}-\overline{Z_{n}}$ and integrate by parts, we have

$$
\int_{B(0, n)}\left|\nabla\left(Z_{n}-\overline{Z_{n}}\right)\right|^{2} d x+\int_{B(0, n)} a^{-}\left(Z_{n}^{q}-{\overline{Z_{n}}}^{q}\right)\left(Z_{n}-\overline{Z_{n}}\right) d x-\int_{B(0, n)} b\left(Z_{n}^{p}-{\overline{Z_{n}}}^{p}\right)\left(Z_{n}-\overline{Z_{n}}\right) d x=0 .
$$

Since they are all non-negative, we conclude that $Z_{n}=\overline{Z_{n}}$.
It is easy to see that $Z_{n+1}$ is a supersolution for (3.2), so by the uniqueness $Z_{n+1} \geqslant Z_{n}$. Moreover we have the following estimate

$$
\int_{B(0, n)}\left|\nabla Z_{n}\right|^{2} d x+\int_{B(0, n)} a^{-} Z_{n}^{q+1} d x-\int_{B(0, n)} b Z_{n}^{p+1} d x=\int_{B(0, n)} a^{+} f^{q} Z_{n} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x .
$$

Therefore let $Z=\lim _{n \rightarrow \infty} Z_{n}$, then $Z$ is a solution to (3.1) and it satisfies

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|\nabla Z|^{2} d x+\int_{\mathbb{R}^{n}} a^{-} Z^{q+1} d x-\int_{\mathbb{R}^{n}} b Z^{p+1} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x \tag{3.3}
\end{equation*}
$$

and $Z \leqslant \omega$.

Next, in order to prove the uniqueness of the solution $Z$, we need to improve Lemma 2.3. Let

$$
\begin{equation*}
V=\frac{[n(n-2)]^{\frac{n-2}{4}}}{\left(1+|x|^{2}\right)^{\frac{n-2}{2}}} . \tag{3.4}
\end{equation*}
$$

Of course $V(x)$ is (up to scaling and translation) the unique solution of the familiar critical Sobolev exponent equation in $\mathbb{R}^{n}, \Delta V+V^{\frac{n+2}{n-2}}=0$ (see [10]). Under assumption (1.2), $\Omega^{+}$is bounded, we can pick $\rho_{1}>0$ so that $\Omega^{+} \Subset B\left(0, \rho_{1}\right)$. We have the following lemma.

Lemma 3.2. Assume $Z$ in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ is a non-negative smooth solution of $-\Delta Z \leqslant 0$ in $\mathbb{R}^{n}-B\left(0, \rho_{1}\right)$. Then, there exists a constant $C>0$ such that $Z \leqslant C V$ in $\mathbb{R}^{n}-B\left(0, \rho_{1}\right)$. Moreover there exists an increasing sequence $\left\{R_{n}\right\}$ with $\lim _{n \rightarrow \infty} R_{n}=\infty$ so that $\int_{\partial B\left(0, R_{n}\right)} \frac{\partial Z}{\partial n} Z \rightarrow 0$ as $n \rightarrow \infty$.

Proof. From Lemma 2.1, we have $\lim _{|x| \rightarrow \infty} Z(x)=0$. Following the proof of Lemma 2.3, replace $v$ in Lemma 2.3 by $C V$ for some big constant $C$. Notice that $-\Delta(C V)>0$, we can show that $Z \leqslant C V$ in $\mathbb{R}^{n}-B\left(0, \rho_{1}\right)$. Therefore we have $Z \leqslant C_{1}|x|^{-(n-2)}$ for some $C_{1}$, then the last part follows from the claim in Lemma 2.6.

Lemma 3.3. Under hypothesis (1.2), the solution obtained from previous lemma is unique.
Proof. Suppose there are two solutions $Z, \bar{Z}$, which satisfy the estimates (3.3), then we have

$$
-\Delta(Z-\bar{Z}) d x+a^{-}\left(Z^{q}-\bar{Z}^{q}\right) d x-b\left(Z^{p}-\bar{Z}^{p}\right) d x=0
$$

Then multiply both sides by $Z-\bar{Z}$ and integrate by parts over $B(0, R)$, we have

$$
\begin{aligned}
& \int_{B(0, R)}|\nabla(Z-\bar{Z})|^{2} d x+\int_{B(0, R)} a^{-}\left(Z^{q}-\bar{Z}^{q}\right)(Z-\bar{Z}) d x-\int_{B(0, R)} b\left(Z^{p}-\bar{Z}^{p}\right)(Z-\bar{Z}) d x \\
& \quad+\int_{\partial B(0, R)} \frac{\partial(Z-\bar{Z})}{\partial n}(Z-\bar{Z}) d S=0 .
\end{aligned}
$$

From Lemma 3.2, there exists a sequence $\left\{R_{n}\right\}$ and $\lim _{n \rightarrow \infty} R_{n}=\infty$, such that $\int_{\partial B\left(0, R_{n}\right)} \frac{\partial(Z-\bar{Z})}{\partial n}(Z-$ $\bar{Z}) \rightarrow 0$ as $n \rightarrow \infty$. Let $n \rightarrow \infty$, we have

$$
\int_{\mathbb{R}^{n}}|\nabla(Z-\bar{Z})|^{2} d x+\int_{\mathbb{R}^{n}} a^{-}\left(Z^{q}-\bar{Z}^{q}\right)(Z-\bar{Z}) d x-\int_{\mathbb{R}^{n}} b\left(Z^{p}-\bar{Z}^{p}\right)(Z-\bar{Z}) d x=0 .
$$

So we can conclude $Z=\bar{Z}$.
Next we go to the main iteration process. Let us consider the following iteration equation

$$
\begin{equation*}
-\Delta u_{n+1}+a^{-} u_{n+1}^{q}-b u_{n+1}^{p}=a^{+} u_{n}^{q} \quad \text { in } \mathbb{R}^{n}, u_{n} \geqslant 0 \text { and } u_{n} \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right), \tag{3.5}
\end{equation*}
$$

where $u_{1}=\underline{u}_{\rho}$ for small $\rho$ such that $0 \leqslant \underline{u}_{\rho} \leqslant \omega$, and $\underline{u}_{\rho}$ is constructed in the following manner: take small ball $B \Subset \Omega^{+}$, let $\underline{a}=\inf _{x \in B} a(x)$, define

$$
\underline{u}_{\rho}= \begin{cases}\rho \xi_{i} & \text { in } B, \\ 0, & \text { else },\end{cases}
$$

where $\xi>0$ is the eigenfunction corresponding to the first eigenvalue of the following problem

$$
-\Delta \xi=\lambda a \xi \quad \text { in } B \quad \text { and } \quad \xi=0 \quad \text { on } \partial B
$$

for details, see [3]. Notice that $\underline{u}_{\rho}$ is Lipschitz in $\mathbb{R}^{n}$ and satisfies

$$
\int_{\mathbb{R}^{n}} \nabla \underline{u}_{\rho} \nabla \phi d x \leqslant \int_{\mathbb{R}^{n}} a \underline{u}_{\rho}^{q} \phi d x+\int_{\mathbb{R}^{n}} b \underline{u}_{\rho}^{p} \phi d x \quad \text { for } \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right),
$$

since $\underline{u}_{\rho}$ is a subsolution.
Lemma 3.4. Assume (1.2), then $\omega \geqslant u_{2} \geqslant u_{1}=\underline{u}_{\rho}$ and

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left|\nabla u_{2}\right|^{2} d x+\int_{\mathbb{R}^{n}} a^{-} u_{2}^{q+1} d x-\int_{\mathbb{R}^{n}} b u_{2}^{p+1} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x . \tag{3.6}
\end{equation*}
$$

Proof. The proof is simple, we just go back to proof of Lemma 3.1. It is easy to see that $\underline{u}_{\rho}$ is a subsolution to Eq. (3.2) and $\omega$ is a supersolution, therefore we have desired results.

Base on the above lemma, we can start our induction process.
Lemma 3.5. Assume (1.2), then $\omega \geqslant u_{n+1} \geqslant u_{n} \geqslant \underline{u}_{\rho}$ and

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left|\nabla u_{n}\right|^{2} d x+\int_{\mathbb{R}^{n}} a^{-} u_{n}^{q+1} d x-\int_{\mathbb{R}^{n}} b u_{n}^{p+1} d x \leqslant \int_{\mathbb{R}^{n}} a^{+} \omega^{q+1} d x . \tag{3.7}
\end{equation*}
$$

Proof. From the above lemma we know the initial step is true, now assume $u_{n} \geqslant u_{n-1}$, we have

$$
-\Delta u_{n+1}+a^{-} u_{n+1}^{q}-b u_{n+1}^{p} \geqslant-\Delta u_{n}+a^{-} u_{n}^{q}-b u_{n}^{p},
$$

that is

$$
-\Delta\left(u_{n}-u_{n+1}\right)+a^{-}\left(u_{n}^{q}-u_{n+1}^{q}\right)-b\left(u_{n}^{p}-u_{n+1}^{p}\right) \leqslant 0 .
$$

Then multiply both sides by $\left(u_{n}-u_{n+1}\right)^{+}$, integrate over $B(0, R)$, and follow the steps in claim of Lemma 3.3, we will have

$$
\int_{\mathbb{R}^{n}}\left|\left(u_{n}-u_{n+1}\right)^{+}\right|^{2} d x+\int_{\mathbb{R}^{n}} a^{-}\left(u_{n}^{q}-u_{n+1}^{q}\right)\left(u_{n}-u_{n+1}\right)^{+} d x-\int_{\mathbb{R}^{n}} b\left(u_{n}^{p}-u_{n+1}^{p}\right)\left(u_{n}-u_{n+1}\right)^{+} d x \leqslant 0 .
$$

From the above we conclude $\left(u_{n}-u_{n+1}\right)^{+}=0$, which means $u_{n+1} \geqslant u_{n}$.
Now we are in position to prove Theorem 1.1:
Proof of Theorem 1.1. We simply take $U=\lim _{n \rightarrow \infty} u_{n}$, then in view of the estimates from the previous we know $U$ is the required solution of (1.1).

Remark 3.6. Assumption (1.2) is not essential for the existence of $U$; for a more general existence result see [8].

Corollary 3.7. Under hypothesis (1.4), (1.1) has a classical compactly supported solution $U$ with its support contained in $B(0, \rho)$ if $\lim _{|x| \rightarrow \infty} a^{-}(x)|x|^{(n-2)(1-q)}>\eta$, where $\eta$ and $\rho$ are from Theorem 1.2.

In terms of the solution class $S_{I}$ (defined in (1.4)), we obtain the following existence result:
Corollary 3.8. Assume (1.7), then $S_{I} \neq \emptyset$ for any non-empty $I \in M$.
Proof. Construct a subsolution $\underline{u}_{\rho}$ as a superposition of disjointly supported subsolutions, one for each component of $\Omega_{I}^{+}$, as in the proof of Theorem 1.1. Monotone iteration then produces a solution which is positive in each component.

With some minor modifications of these arguments we may now prove the existence of minimal and maximal elements in $S_{I}$ as announced in Theorem 1.7.

Proof of Theorem 1.7. We first assume that hypothesis (1.7) holds. To prove assertion (a) we may argue as in Theorem 4 of [18], and define $u_{I}=\inf _{u \in S_{I}} u(x)$, which (by [18]) is the desired minimal solution.

To prove part (b) of Theorem 1.7, we refine our choice of supersolution. To achieve this, we bring in the function $V(x)$ from (3.4).

Claim. Assuming only (1.2), there exists a constant $C>0$ such that

$$
\begin{equation*}
-\Delta(C V)+a^{-}(C V)^{q}-b(C V)^{p} \geqslant a^{+}(C V)^{q} \tag{3.8}
\end{equation*}
$$

and $C V \geqslant u$ for any solution $u$ of (1.1).
Indeed, for (3.8) to hold, we only need $C$ big enough to achieve $C^{1-q} V^{2^{*}-1-q} \geqslant a^{+}$since $-\Delta V=$ $V^{2^{*}-1}$. Notice that by assuming only (1.2), we still have $L^{\infty}$-estimate as in Lemma 2.4 because of Lemma 3.2, so for $C$ big enough, we have $C V \geqslant u$.

We then restart our iteration process (3.5) as in the proof of Theorem 1.1, starting with the supersolution $u_{1}=C V \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$. By similar proof to Lemma 3.5 , we deduce that a decreasing sequence $\left\{u_{n}\right\}$ such that $u_{n+1} \leqslant u_{n} \leqslant u_{1}=C V$ and (3.7) holds with CV in place of $\omega$. Moreover for any solution $u$ of (1.1), we must have $u \leqslant u_{n}$ for all $n$, by an induction process. Hence we conclude that $u_{n} \rightarrow U$, a solution of (1.1), and with $U \geqslant u$ holding for any solution $u$ of (1.1). Therefore, $U$ must be the maximal solution. Since $\omega>0$ in $\mathbb{R}^{n}$, it is an easy consequence of case 1 in Lemma 2.8 that $U<\omega$ in $\mathbb{R}^{n}$.

Finally, we also conclude that, for the parametrized family of problems (1.9), the maximal solution $U_{\lambda} \in S_{M}$ is monotone.

Corollary 3.9. Assuming (1.7), the maximal solution $U_{\lambda} \in S_{M}$ of (1.9) is increasing as $\lambda$ increases.
Proof. For $0<\lambda_{1}<\lambda_{2}$, we pick $C$ big enough so that $C V \geqslant \max \left(U_{\lambda_{1}}, U_{\lambda_{2}}\right)$. Then we iterate this supersolution $C V$ for (1.9) $)_{\lambda_{2}}$. As above, it results in a monotone decreasing sequence $\left\{u_{n}\right\}$ and $u_{n} \geqslant$ $\max \left(U_{\lambda_{1}}, U_{\lambda_{2}}\right)$. But then $U_{\lambda_{1}} \leqslant u_{n} \rightarrow U_{\lambda_{2}}$.

### 3.2. Uniqueness and support

We now turn to questions of uniqueness and the characterization of the solution space of (1.1) in terms of the supports of the solutions. First, we note that every solution of (1.1) must be positive in at least one connected component of $\Omega^{+}$.

Lemma 3.10. Assume (1.7), if $I=\emptyset$, then $N_{I}=\emptyset$.
Proof. If $I=\emptyset$, then for any $u \in N_{I}$ it is a subharmonic function. Thus, for any $x \in \mathbb{R}^{n}$, we have

$$
0 \leqslant u(x) \leqslant \frac{1}{|\partial B(x, R)|} \int_{\partial B(x, R)} u(y) d y \rightarrow 0 \quad \text { as } R \rightarrow \infty
$$

since $\lim _{|x| \rightarrow \infty} u(x)=0$, which implies $u(x)=0$.
Lemma 3.11. The classical solution $u$ of (1.1) is either positive in $\Omega_{i}^{+}$or entirely zero in $\Omega_{i}^{+}$for any $i \in M$.
Proof. Let us consider the set $S=\left\{x \in \Omega_{i}^{+} \mid u(x)=0\right\}$. First, we claim that $S$ is open in $\Omega_{i}^{+}$. Indeed, if $S \neq \emptyset$, then pick any $x_{0} \in S$, we have $u\left(x_{0}\right)=0$ and $a\left(x_{0}\right)>0$. Therefore by continuity

$$
a(x)+b(x) u^{p-q}(x)>0 \quad \text { in } B\left(x_{0}, \epsilon\right)
$$

for small $\epsilon>0$. Hence we have

$$
-\Delta u=a(x) u^{q}+b(x) u^{p}=\left(a(x)+b(x) u^{p-q}\right) u^{q} \geqslant 0 \quad \text { in } B\left(x_{0}, \epsilon\right), \quad u(x) \geqslant 0 \quad \text { in } B\left(x_{0}, \epsilon\right) .
$$

So by maximum principle $u(x) \equiv 0$ in $B\left(x_{0}, \epsilon\right)$, which means $S$ is open in $\Omega_{i}^{+}$.
It is also clear by continuity that $S$ is closed in $\Omega_{i}^{+}$, therefore $u>0$ in $\Omega_{i}^{+}$or $u \equiv 0$ in $\Omega_{i}^{+}$due to the connectivity of $\Omega_{i}^{+}$.

For uniqueness questions it is not enough to know that any solution $u$ of (1.1) is positive in $\Omega_{i}^{+}$, we require $u$ to be positive up to and including the boundary of $\Omega_{i}^{+}$. In case $p \geqslant 1$, this follows from the classical Hopf Lemma (see below) but the question is more delicate for $p<1$. Therefore we define a function class $P$ which incorporates this property,

$$
\begin{aligned}
P= & \left\{v \in C_{l o c}^{1}\left(\mathbb{R}^{n}\right) \cap W_{l o c}^{2, s}\left(\mathbb{R}^{n}\right) \mid s>n,\right. \\
& \left.v \ngtr 0 \text { in } \mathbb{R}^{n}, v>0 \text { in } \overline{\Omega_{i}^{+}} \text {if } v>0 \text { in } \Omega_{i}^{+} \text {for } i \in M\right\} .
\end{aligned}
$$

Lemma 3.12. Assume (1.7) and solution $u$ of (1.1) is positive in $\Omega_{i}^{+}$for some $i \in M$. Then if $p \geqslant 1, u>0$ in $\overline{\Omega_{i}^{+}}$; if $p<1$ and $\frac{b(x)}{a(x)}$ is uniformly bounded in $\overline{\Omega^{+}}, u>0$ in $\overline{\Omega_{i}^{+}}$.

Proof. For $p \geqslant 1$, hypothesis (1.7) ensures that an interior ball condition is satisfied by $\Omega_{i}^{+}$, so we may directly apply Hopf's Lemma to the equation

$$
-\Delta u-b(x) u^{p}=a(x) u^{q} \geqslant 0 \quad \text { in } \Omega_{i}^{+}, \quad u \geqslant 0 \quad \text { in } \overline{\Omega_{i}^{+}}
$$

We conclude that $u>0$ in $\overline{\Omega^{+}}$.
For $p<1$, if $x_{0} \in \partial \Omega_{i}^{+}$and $u\left(x_{0}\right)=0$. Since $\Omega_{i}^{+}$satisfies an interior ball condition, we take a small ball $B_{\epsilon} \subset \Omega_{i}^{+}$with radius $\epsilon$ and $x_{0} \in \partial B_{\epsilon}$. For $\epsilon$ small we have

$$
-\Delta u=a(x)\left(1+\frac{b(x)}{a(x)} u^{p-q}\right) u^{q} \geqslant 0 \quad \text { in } B_{\epsilon}
$$

Hopf's Lemma implies that $\nabla u\left(x_{0}\right) \neq 0$. Since $u$ attains minimal value at $x_{0}, \nabla u\left(x_{0}\right)=0$, a contradiction.

Definition 3.13. We say that $b(x)$ is compatible with $a(x)$ if any nonzero solution $u$ of (1.1) lies in the function set $P$.

From the above lemma, we see that when $p \geqslant 1$, any non-positive $b(x)$ is compatible with $a(x)$, but for $p<1$ an extra hypothesis must be imposed on $b$ near $\partial \Omega_{i}^{+}$.

In Spruck [23], it was shown that if $\lim \sup _{|x| \rightarrow \infty} a(x)<0$ and $\nabla a(x) \cdot x<0$, Eq. (1.1) with $b(x) \equiv 0$ has a unique compactly supported solution. With a very mild assumption on $a(x)$, we show a uniqueness result, which generalizes the uniqueness result of Spruck and includes Theorem 1.5.

Theorem 3.14. Assume (1.7), if $b(x)$ is compatible with $a(x)$, then the number of elements in $N_{I}$ is at most 1 for any non-empty I. In particular if $k=1$, then the solution to (1.1) is unique and its support is connected.

Proof. Let us take two element $u_{1}$ and $u_{2}$ from $N_{I}$ and apply Lemma 2.8. We check all the conditions required by Lemma 2.8. Since $b(x)$ is compatible with $a(x)$, then $u_{1}, u_{2} \in P$, which means that the first condition is satisfied. The second condition is automatic because $u_{1}, u_{2} \in N_{I}$. The third condition is assured by Lemma 2.1 and the last condition is also automatic. We apply Lemma 2.8 twice and get $u_{1} \leqslant u_{2}, u_{1} \geqslant u_{2}$, hence $u_{1} \equiv u_{2}$.

Immediately we combine the existence of a maximal solution from Theorem 1.7(b) with the uniqueness result obtained above to get the following corollary.

Corollary 3.15. Assume (1.7), if b is compatible with a(x), the maximal solution $U$ is the unique element in $S_{M}$.
Here we present the proof for Proposition 1.6.
Proof of Proposition 1.6. Recall that in Proposition 1.6 we assume $p \geqslant 1$. For $b(x) \equiv-1$, define $a_{i}(x)$ by

$$
a_{i}(x)= \begin{cases}\lambda, & \text { if } x \in \Omega_{i}^{+}, \\ -1, & \text { if } x \notin \Omega_{i}^{+} .\end{cases}
$$

Then we consider the following equation:

$$
\begin{equation*}
-\Delta u=a_{i}(x) u^{q}-u^{p} \quad \text { in } \mathbb{R}^{n}, u \geqslant 0 \text { and } u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) . \tag{3.9}
\end{equation*}
$$

By elliptic regularity and the strong maximum principle, we see that solution to (3.9) lies in the set $P$. By Lemma 2.8 the solution to (3.9) is unique, denoted by $u_{i}$, and compactly supported by Theorem 1.2. So if we choose $\delta^{*}$ big enough so that $\operatorname{supp}\left(u_{i}\right) \cap \operatorname{supp}\left(u_{j}\right)=\emptyset$ for any $i \neq j$, then by Theorem 3.14 $N_{I}$ contains exactly one solution for each $I \neq \emptyset$. This proves the first assertion of Proposition 1.6.

For the second statement in the proposition, notice that $\Omega_{i}^{+} \Subset \operatorname{supp}\left(u_{i}\right)$. For non-empty $I$, choose any $i \in I$, then the minimal solution $u_{I}$ in $S_{I}$ is supersolution for problem (3.9). By uniqueness we find that $u_{I} \geqslant u_{i}$ for any $i \in I$. So if we choose $\delta_{*}$ so small that $\operatorname{supp}\left(u_{i}\right) \cap \Omega_{j}^{+} \neq \emptyset$ for any $i \neq j$, which implies that $\operatorname{supp}\left(u_{I}\right) \cap \Omega_{j}^{+} \neq \emptyset$. Therefore by uniqueness result above, (1.1) has only one solution, which is positive in $\Omega^{+}$.

In the case where $a(x)$ and $b(x)$ are radially symmetric we obtain more precise information, via uniqueness:

Corollary 3.16. Assume (1.7), if $a$ and $b$ are radially symmetric and $b$ is compatible with $a$, then the unique element in $N_{I}$ is radially symmetric. If in addition $\Omega^{+}$is a ball centered at the origin and $b(x)=0$ in $\Omega^{+}$, then the unique solution of (1.1) is radially decreasing.

Proof. If $a$ and $b$ are radially symmetric, since Laplace's operator is invariant under rotation, Lemma 2.8 ensures that the element in $N_{I}$ must be radially symmetric. If $\Omega^{+}$is a ball centered at the origin and $b(x)=0$ in $\Omega^{+}$, then $b$ is compatible with $a$. The result follows from maximum principle and the fact that the unique solution uniformly converges to zero at the infinity.

Remark 3.17. Note that we do not need to apply the moving planes method in this setting since we have the uniqueness result. Furthermore, the moving plane process would require more stringent hypotheses on $a$ and $b$.

## 4. The parametrized equation

In this section we consider the effect of the parameter $\lambda$ on the shape and multiplicity of solutions to the parametrized family,

$$
\begin{equation*}
-\Delta u=a_{\lambda}(x) u^{q}+b u^{p} \quad \text { in } \mathbb{R}^{n}, u \geqslant 0 \text { and } u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right), \tag{4.1}
\end{equation*}
$$

where $a_{\lambda}(x)=\lambda a^{+}(x)-a^{-}(x)$.
We consider both asymptotic limits, $\lambda \rightarrow 0$ and $\lambda \rightarrow \infty$, for the maximal solution $U_{\lambda} \in L^{\infty}\left(\mathbb{R}^{n}\right) \cap$ $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$. Assuming that $b$ is compatible with $a$, we show that as $\lambda$ decreases to zero, $U_{\lambda}$ decreases (see Corollary 3.9), and ultimately the support of $U_{\lambda}$ breaks into several disjoint components. This is the content of Theorem 1.9, stated in the Introduction. Before proving the theorem we need two lemmas.

Lemma 4.1. For any positive constant $c$, the equation $\Delta \bar{u}=c \bar{u}^{q}$ in $\mathbb{R}^{n}$ has a radial solution $\bar{U}=\theta r^{\frac{2}{1-q}}$, where $\theta^{1-q}=\frac{(1-q)^{2} c}{2[n-q(n-2)]}$.

Lemma 4.2. For any ball $B \Subset \mathbb{R}^{n}-\Omega^{0+}$, any $g(x) \geqslant 0$, the following problem has at most one non-negative classical solution,

$$
-\Delta v=a(x) v^{q}+b v^{p} \quad \text { in } B, \quad v=g \quad \text { on } \partial B .
$$

Proof. Lemma 4.1 follows from direct calculation. For Lemma 4.2, suppose there were two classical solutions $v_{1}, v_{2}$ and $v_{1} \neq v_{2}$. Assuming that ( $v_{1}-v_{2}$ ) achieves a positive maximum value at some $x_{0} \in B$, we have

$$
0 \leqslant-\Delta\left(v_{1}-v_{2}\right)\left(x_{0}\right)=a\left(x_{0}\right)\left(v_{1}^{q}\left(x_{0}\right)-v_{2}^{q}\left(x_{0}\right)\right)+b\left(x_{0}\right)\left(v_{1}^{p}\left(x_{0}\right)-v_{2}^{p}\left(x_{0}\right)\right)<0,
$$

which is a contradiction. Hence we must have $v_{1} \equiv v_{2}$, and Lemma 4.2 is verified.
With the help of the above two lemmas, we give the proof of Theorem 1.9.
Proof of Theorem 1.9. By assumption $a(x)$ is admissible, then $\operatorname{dist}\left(\Omega_{i}^{0+}, \Omega_{j}^{0+}\right)>0$ for any $i \neq j$. Let $\delta=\frac{1}{16} \inf _{i \neq j} \operatorname{dist}\left(\Omega_{i}^{0+}, \Omega_{j}^{0+}\right)$, then $\delta>0$. As before, take $\rho_{1}>0$ so that $\Omega^{0+} \Subset B\left(0, \rho_{1}\right)$. Let

$$
C_{i}=\left\{x \in B\left(0, \rho_{1}+48 \delta\right) \mid \operatorname{dist}\left(x, \Omega_{i}^{0+}\right) \leqslant \delta\right\} .
$$

It is easy to see that $C_{i} \cap C_{j}=\emptyset$ for any $i \neq j$. Let $C=\bigcup_{i \in M} C_{i}$. We define

$$
N=\left\{x \in B\left(0, \rho_{1}+32 \delta\right) \mid \operatorname{dist}\left(x, \Omega^{0+}\right) \geqslant 4 \delta\right\} .
$$

For any $x \in N, \overline{B(x, \delta)} \cap C_{i}=\emptyset$ for any $i \in M$. Finally set

$$
\underline{a}=\inf _{x \in B(0, R+48 \delta)-C} a^{-}(x) .
$$

By Lemma 4.1 we find that the following equation

$$
\Delta \bar{u}=\underline{a} \bar{u}^{q} \quad \text { in } B(x, \delta), \quad \bar{u}=\theta(\delta)^{\frac{2}{1-q}} \quad \text { on } \partial B(x, \delta)
$$

has a solution $\bar{U}=\theta|y-x|^{\frac{2}{1-q}}$, where $x \in N$.
We now claim that the equation

$$
-\Delta v=a^{-}(y) v^{q}+b(y) v^{p} \quad \text { in } B(x, \delta), \quad v=U_{\lambda} \quad \text { on } \partial B(x, \delta)
$$

has a unique solution $v=U_{\lambda}$ if $\left\|U_{\lambda}\right\|_{L^{\infty}\left(R^{n}\right)} \leqslant \theta(\delta)^{\frac{2}{1-q}}$, where $U_{\lambda}$ is the maximal solution of (4.1). Indeed, from Lemma 2.4, we have $\lim _{\lambda \rightarrow 0} U_{\lambda}=0$. Therefore there exists $\lambda_{*}>0$ such that $\left\|U_{\lambda}\right\|_{L^{\infty}\left(R^{n}\right)} \leqslant$ $\theta(\delta)^{\frac{2}{1-q}}$ for $\lambda \leqslant \lambda_{*}$. Hence $\bar{U}$ is a supersolution for the above equation and 0 is a subsolution, so the above equation has a solution $v$. But it is clear the maximal solution $U_{\lambda}$ is also a solution, by Lemma 4.2 we know that the above equation has a unique solution $v=U_{\lambda}$. Therefore we have $0 \leqslant U_{\lambda}(x) \leqslant \bar{U}(x)=0$, which means $U_{\lambda}(x)=0$ for all $x \in N$. Hence we can write

$$
U_{\lambda}=\sum_{i \in M} u_{i} \quad \text { and } \quad \operatorname{supp} u_{i} \cap \operatorname{supp} u_{j}=\emptyset \quad \text { for } i \neq j,
$$

where $\Omega^{+} \subset \operatorname{supp}\left(u_{i}\right)$, which means that $\sum_{i \in I} u_{i} \in N_{I}$.
Remark 4.3. If we also assume that $b$ is compatible with $a$, then solutions in $N_{I}$ will be unique for all $I \neq \emptyset$. Thus, for $\lambda$ small, the support of the maximal solution $U_{\lambda}$ will break into $k$ disjoint components, and these disjoint compactly supported solutions then generate the unique element in each $N_{I}$ by superposition.

Next we consider the limit $\lambda \rightarrow \infty$. We will prove that for $p<1$ and $\lambda$ sufficiently large, (4.1) has a unique solution. In particular, it will be positive in each of the connected components of $\Omega^{0+}$.

Under the assumption (1.7), fix $I \subset M$, by Theorem 1.7 the minimal element of $S_{I}$ exists, denoted by $u_{\lambda}$. Let $v_{\lambda}=\lambda^{\frac{1}{q-1}} u_{\lambda}$, then $v_{\lambda}$ satisfies the following equation in $\mathbb{R}^{n}$,

$$
\begin{equation*}
-\Delta v=a^{+} v^{q}-\frac{a^{-}}{\lambda} v^{q}+b \lambda^{\frac{1-p}{q-1}} v^{p}, \quad v \geqslant 0 \text { and } v \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right) . \tag{4.2}
\end{equation*}
$$

For each fixed $\lambda$, let $\bar{S}_{I}$ be the corresponding set of $S_{I}$, associated with the above equation (4.2). Let us begin the proof of Theorem 1.10 with a few lemmas. Recall that Theorem 1.10 concerns the case $p<1$.

Lemma 4.4. Assume (1.7), then $v_{\lambda}$ is the minimal element in $\bar{S}_{I}$. Moreover $v_{\lambda}$ is increasing as $\lambda$ increases and $v_{\lambda} \leqslant \omega$, where $\omega$ is as in (2.2).

Proof. Suppose $0<\lambda_{1}<\lambda_{2}$, since $-\frac{1}{\lambda}$ and $b \lambda^{\frac{1-p}{q-1}}$ are increasing in $\lambda$, then from Theorem 1.1 we see that $v_{\lambda_{2}}$ is a supersolution for Eq. (4.2) at $\lambda=\lambda_{1}$, by choosing suitable $\underline{u}_{\rho}$ as subsolution, we get a solution $v \in \bar{S}_{I}$ of (4.2) at $\lambda=\lambda_{1}$ such that $v_{\lambda_{1}} \leqslant v \leqslant v_{\lambda_{2}}$. The last part is from Theorem 1.7.

Lemma 4.5. There exists a unique solution $\bar{\omega}$ to the following problem:

$$
\begin{equation*}
-\Delta \bar{\omega}=a^{+} \bar{\omega}^{q}+b \bar{\omega}, \quad \bar{\omega}(x) \rightarrow 0 \quad \text { as }|x| \rightarrow \infty . \tag{4.3}
\end{equation*}
$$

Proof. Again the existence and uniqueness of $\bar{\omega}$ follows from [8].
Lemma 4.6. Assume (1.7), we have the following:

> for $p<1, \quad v_{\lambda}$ converges uniformly to $\omega, \quad$ where $\omega$ is from (2.2),
> for $p=1, \quad v_{\lambda}$ converges uniformly to $\bar{\omega}, \quad$ where $\bar{\omega}$ is from (4.3).

Proof. From the above lemma we see that $v_{\lambda}$ is increasing in $\lambda$ and uniformly bounded. Let $V=\lim _{\lambda \rightarrow \infty} v_{\lambda}$, then it is easy to see that $\lim _{|x| \rightarrow \infty} V(x)=0$. Moreover from Eq. (4.2) we have $\left\|v_{\lambda}\right\|_{C_{1, \alpha(B(0, R))}}$ is uniformly bounded for fixed $R$, therefore we have $v_{\lambda}$ uniformly converges to $V$. Hence $V$ solves the following equation

$$
-\Delta w=a^{+} w^{q} \text { if } p>1 \text { or }-\Delta w=a^{+} w^{q}+b w \quad \text { if } p=1 .
$$

Since $\lim _{|x| \rightarrow \infty} V(x)=0$, by uniqueness of the solution (in either case) we obtain our conclusion.
Now we are ready to prove Theorem 1.10.
Proof of Theorem 1.10. Since $u_{\lambda}=\lambda^{\frac{1}{1-q}} v_{\lambda}$ and $v_{\lambda}$ uniformly converges to $V>0$ in $\mathbb{R}^{n}$. Since $u_{\lambda}$ is minimal element in $S_{I}$ and the choices for $I$ are finite, all solutions of (4.1) $\lambda_{\lambda}$ are positive in $\overline{\Omega^{+}}$when $\lambda$ is large. Therefore Lemma 2.8 implies that $S_{M}$ has only one element.

In the case $p>1$, it is easy to see that the above process could not work, partially because now $\lim _{\lambda \rightarrow \infty} \lambda^{\frac{1-p}{q-1}}=\infty$. To obtain some asymptotic results in this regime, we must impose some conditions on $b(x)$ and around the set $\Omega^{+}$.

Let $u_{\lambda}$ be any solution of (1.9). We prove:
Theorem 4.7. Assume (1.7) and $p>1$.

1. For any $\sigma>0, \lim _{\lambda \rightarrow \infty}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q+\sigma} \lambda^{-1}=\infty$.
2. If inf $x_{x \in \Omega^{+}}|b(x)|>0$, then $\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q} \leqslant C \lambda$ for some constant $C>0$ independent of $\lambda$.
3. If $b(x)=0$ in a ball $B \subset \Omega_{i}^{+}$for some $i \in M$, then $\liminf _{\lambda \rightarrow \infty}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \lambda^{\frac{1}{q-1}}>0$, for $I=\{i\}$ and $u_{\lambda} \in S_{I}$.
4. If $\left\{x \in \mathbb{R}^{n} \mid b(x)=0\right\} \cap \Omega^{+}$has an open connected component $O$ such that $\Omega_{i}^{+} \cap O \neq \emptyset$ for any $i \in M$, then the problem (1.9) has a unique solution, which is positive in $\Omega^{+}$, for $\lambda$ large enough.

We first require some lemmas. Let $v_{\lambda}=\lambda^{\frac{1}{q-1}} u_{\lambda}$, so that $v_{\lambda}$ satisfies the following modified equation:

$$
\begin{equation*}
-\Delta v=a^{+} v^{q}-\frac{a^{-}}{\lambda} v^{q}+\frac{b}{\lambda^{1-\frac{\epsilon}{1-q}}} u_{\lambda}^{p-q-\epsilon} v^{q+\epsilon} \quad \text { in } \mathbb{R}^{n}, \tag{4.4}
\end{equation*}
$$

where we pick $\epsilon>0$ small so that $\epsilon<1-q$. Under the assumption (1.7) Eq. (4.4) admits a minimal solution $\overline{v_{\lambda}}$ in the class $\overline{S_{I}}$, where $\overline{S_{I}}$ is the corresponding set of $S_{I}$, associated with Eq. (4.4).

Lemma 4.8. Assume (1.7), iffor some $\sigma>0, \liminf _{\lambda \rightarrow \infty}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q+\sigma} \lambda^{-1}<\infty$. Then there exists an increasing sequence $\left\{\lambda_{n}\right\}$ with $\lim _{n \rightarrow \infty} \lambda_{n}=\infty$ so that $\left\|\overline{\bar{V}_{n}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}$ has a positive lower bound $C\left(\Omega_{I}^{+}\right)$independent of $\lambda$ for large $\lambda$.

Proof. By assumption we can pick an increasing sequence $\left\{\lambda_{n}\right\}$ with $\lim _{n \rightarrow \infty} \lambda_{n}=\infty$ so that $\left\|u_{\lambda_{n}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q+\sigma} \lambda_{n}^{-1} \leqslant C$ for some $C>0$. Hence we have $\left\|u_{\lambda_{n}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \leqslant C \lambda_{n}^{\frac{1}{p-q+\sigma}}$, so

$$
\frac{u_{\lambda_{n}}^{p-q-\epsilon}}{\lambda_{n}^{1-\frac{\epsilon}{1-q}}} \leqslant C \lambda_{n}^{\frac{p-q-\epsilon}{p-q+\sigma}-\left(1-\frac{\epsilon}{1-q}\right)} .
$$

For this fixed $\sigma>0$, we can choose $\epsilon>0$ small so that $\frac{p-q-\epsilon}{p-q+\sigma}<1-\frac{\epsilon}{1-q}$. Therefore $\frac{b}{\lambda_{n}^{1-\epsilon} 1-q} u_{\lambda_{n}}^{p-q-\epsilon} \rightarrow 0$ uniformly in $\overline{\Omega^{+}}$. By the same proof for Lemma 4 in [18], the result follows.

We are ready to prove the first two assertions in Theorem 4.7.
Proof of Theorem 4.7. For the first, since under hypothesis (1.7) the choices for $I$ are finite, we only need to show the theorem is true for $u_{\lambda}$, here $u_{\lambda}$ is the minimal element in $S_{I}$ for each fixed $I \subset M$.

We proceed by contradiction: suppose that for some $\sigma>0$,

$$
\begin{equation*}
\liminf _{\lambda \rightarrow \infty}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q+\sigma} \lambda^{-1} \leqslant c_{0}<\infty . \tag{4.5}
\end{equation*}
$$

Take $\epsilon$ and $\left\{\lambda_{n}\right\}$ as in the proof of the above lemma. There exists $C>0$, for example $C=C\left(\Omega_{I}^{+}\right) / 2$ from the above lemma, so that for $n$ large

$$
\left\|v_{\lambda_{n}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \geqslant \| \overline{{\overline{\lambda_{\lambda}}} \|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \geqslant C . . . . . . . . .}
$$

This implies $\left\|u_{\lambda_{n}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \geqslant C \lambda_{n}^{\frac{1}{1-q}}$, which contradicts (4.5). Therefore, we must have $\liminf \lambda_{\lambda \rightarrow \infty}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q+\sigma} \lambda^{-1}=\infty$ for any $\sigma>0$.

For the second part, let us assume $\inf _{x \in \Omega^{+}}|b(x)|=\bar{b}>0$. Lemma 2.5 implies that

$$
b\left(x_{0}\right) u_{\lambda}^{p}\left(x_{0}\right)+a\left(x_{0}\right) u_{\lambda}^{q}\left(x_{0}\right)=-\Delta u_{\lambda}\left(x_{0}\right) \geqslant 0,
$$

where $u_{\lambda}$ attains its global maximum at $x_{0} \in \overline{\Omega^{+}}$. Hence we have

$$
\lambda\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{q} \geqslant\left(-b\left(x_{0}\right)\right)\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p} \geqslant \bar{b}\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p},
$$

that is $\left\|u_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)}^{p-q} \leqslant \lambda\left\|a^{+}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \bar{b}^{-1}$. This proves the second assertion of the theorem.
To prove the third assertion, we require the following lemma from [18]:
Lemma 4.9. (See [18, Lemma 4].) Assume (1.7), if $b(x)=0$ in a ball $B \subset \Omega_{i}^{+}$for some $i \in M$, then $\left\|\overline{\bar{v}_{\lambda}}\right\|_{L^{\infty}(B)}$ has a positive lower bound $C(B)$ independent of $\lambda$ for large $\lambda$.

The proof of the third statement is then very simple. Take $c=\frac{1}{2} C(B)$, where $C(B)$ is from the above lemma. For large $\lambda$,

$$
\left\|v_{\lambda}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \geqslant\left\|\overline{v_{\lambda}}\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \geqslant c,
$$

which leads to our conclusion.

Finally, we prove the fourth statement of Theorem 4.7 by contradiction. Since $\Omega^{+}$has finitely many connected components, we can assume that there exists an increasing sequence $\left\{\lambda_{n}\right\}$ with $\lim _{n \rightarrow \infty} \lambda_{n}=\infty$ so that $u_{\lambda_{n}}>0$ in $\Omega_{i}^{+}$for some $i \in M$ and $u_{\lambda_{n}}=0$ in $\Omega_{j}^{+}$for some $j \neq i$. Take $I=\{i\}$, assume $u_{\lambda_{n}}$ is the minimal element in $S_{I}$, then like the proof for Theorem 1.10, restricting to a subsequence if necessary, we have $v_{\lambda_{n}} \geqslant \overline{V_{\lambda_{n}}} \rightarrow \bar{V}$ in $O$, which solves $-\Delta w=a^{+} w^{q}$ in $O$, by the above lemma and maximum principle $\bar{V}>0$ in $O$, which contradicts our assumption. This concludes the proof of Theorem 4.7.

## 5. Solutions not in $\mathcal{D}^{\mathbf{1 , 2}}\left(\mathbb{R}^{\boldsymbol{n}}\right)$

In this section we consider the possibility that $u$ is an entire solution to the PDE $-\Delta u=a(x) u^{q}+$ $b(x) u^{p}$ in $\mathbb{R}^{n}$, without requiring $u$ to lie in the finite energy space $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$. When $u \in \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ we prove the uniform decay estimates in Lemmas 2.1 and 2.3 . Since these estimates are crucial for the proof of compact support (Theorem 1.2), it is useful to have some discussion in this direction.

In general we do expect that there are solutions, which are not in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$. For example, consider the following equation

$$
\begin{equation*}
\Delta z=c_{1} z^{q}+b_{1}(x) z^{p} \quad \text { in } \mathbb{R}^{n}-B\left(0, r_{1}\right), \tag{5.1}
\end{equation*}
$$

where $c_{1}>0,0<q<p$ and $r_{1}>0$, moreover $b_{1}(x)=c_{2}\left(r-r_{1}\right)^{\frac{1+q-2 p}{1-q}} r^{-1}$ for $r \geqslant r_{1}$ and $c_{2}>0$. We seek the form of solution $z=\theta\left(r-r_{1}\right)^{\frac{2}{1-q}}$ for some $\theta>0$. Plug it into Eq. (5.1) and calculate. It is easy to see that if the following is satisfied

$$
c_{1}=\frac{2}{1-q}\left(\frac{2}{1-q}-1\right) \theta^{1-q} \quad \text { and } \quad c_{2}=\frac{2(n-1)}{1-q} \theta^{1-p} .
$$

Hence $z=\theta\left(r-r_{1}\right)^{\frac{2}{1-q}}$ is a solution of Eq. (5.1). Let us look at the equation

$$
-\Delta u=a(x) u^{q}+b(x) u^{p} \quad \text { in } \mathbb{R}^{n},
$$

where $a(x)=-c_{1}$ in $\mathbb{R}^{n}-B\left(0, R_{1}\right)$ for some $R_{1}>0$ and $b(x) \leqslant 0$.
From Theorem 1.2 we find that there exists $\rho>R_{1}$ such that $\operatorname{supp}(u) \subset B(0, \rho)$ for any compactly supported solution $u$ of the above equation and $\rho$ is independent of $b$. Now if we assume $r_{1}=\rho$ and $b(x)=-b_{1}(x)$ for $|x| \geqslant \rho$, then the above equation has a solution $U$ of the form

$$
U= \begin{cases}\theta(r-\rho)^{\frac{2}{1-q}} & \text { for } r \geqslant \rho, \\ u & \text { for } r \leqslant \rho\end{cases}
$$

and it is clear that $U \notin \mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$.
The first step to prove the compactness of the solution is to show that solution uniformly converges to zero at infinity. By Lemma 2.1, solutions in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$ have this kind of property. But without assuming in $\mathcal{D}^{1,2}\left(\mathbb{R}^{n}\right)$, we can still have this vanishing property. Here is a result in radial case.

Theorem 5.1. Assume $a(x)=a(|x|), b(x)=b(|x|)$ and $\liminf _{|x| \rightarrow \infty} a^{-}(x)=c>0$. Then any smooth radial solution $u(x)=u(|x|)$ has the property that

$$
\lim _{|x| \rightarrow \infty} u(x)=0
$$

if $u \in L^{\infty}\left(\mathbb{R}^{n}\right)$ and $\lim _{r \rightarrow \infty} u_{r} r^{-1}=0$.

Proof. First pick $R>0$ big enough so that $a^{-}(x) \geqslant \frac{c}{2}$ for any $|x| \geqslant R$, then we discuss the behavior of $u(r)$ in the domain $[R, \infty)$, we divide into three cases.

Case 1. There exists $r_{1} \geqslant R$ such that $u_{r}\left(r_{1}\right)=0$.
First we see that $u(r)$ must attain local minimal value at $r=r_{1}$. Indeed if $u\left(r_{1}\right)=0$, then we are done! So assume $u\left(r_{1}\right)>0$, then we have

$$
u_{r r}\left(r_{1}\right)=u_{r r}\left(r_{1}\right)+(n-1) u_{r}\left(r_{1}\right) r^{-1}=\Delta u\left(r_{1}\right)=-a\left(r_{1}\right) u^{q}\left(r_{1}\right)-b\left(r_{1}\right) u^{p}\left(r_{1}\right)>0 .
$$

Therefore $u(r)$ attains local minimal value at $r=r_{1}$.
We now claim that $u(r)=0$ for $r \geqslant r_{1}$. Otherwise, there would exist $r_{2}>r_{1}$ such that $u\left(r_{2}\right)>0$. So we should have that $u_{r}(r)>0$ for $r>r_{2}$. If not, there would exist $r_{3}>r_{2}$ such that $u_{r}\left(r_{3}\right)=0$, it is easy to see that $u(r)$ attains local minimal value at $r=r_{3}$, therefore $u(r)$ achieves local maximal value at some $r_{4} \in\left(r_{1}, r_{3}\right)$, but this is impossible because $u_{r r}\left(r_{4}\right)=-a\left(r_{4}\right) u^{q}\left(r_{4}\right)-b\left(r_{4}\right) u^{p}\left(r_{4}\right)>0$. So we must have $u_{r}(r)>0$ for $r>r_{2}$. Since $u \in L^{\infty}\left(\mathbb{R}^{n}\right)$ we have $u(r) \uparrow M$ as $r \rightarrow \infty$ for some positive constant $M$, this leads to a sequence $\left\{r_{n}\right\}$ and $\lim _{n \rightarrow \infty} r_{n}=\infty$ such that $\lim _{n \rightarrow \infty} u_{r r}\left(r_{n}\right)=0$. But since $\lim _{r \rightarrow \infty} u_{r} r^{-1}=0$, then we should have

$$
\liminf _{n \rightarrow \infty} u_{r r}\left(r_{n}\right)=\liminf _{n \rightarrow \infty}\left(-a\left(r_{n}\right) u^{q}\left(r_{n}\right)-b\left(r_{n}\right) u^{p}\left(r_{n}\right)\right)>0 .
$$

This is a contradiction! So we have $u(r)=0$ for $r \geqslant r_{1}$.
Case 2. $u_{r}(r)>0$ for any $r \in[R, \infty)$.
From the above proof we see that this case is impossible.
Case 3. $u_{r}(r)<0$ for any $r \in[R, \infty)$.
Let $m=\lim _{r \rightarrow \infty} u(r)$, then we must have $m=0$. Otherwise $m>0$, then we can find a sequence $\left\{r_{n}\right\}$ and $\lim _{n \rightarrow \infty} r_{n}=\infty$ such that

$$
\lim _{n \rightarrow \infty} u_{r r}\left(r_{n}\right)=0
$$

But

$$
\liminf _{n \rightarrow \infty} u_{r r}\left(r_{n}\right)=\liminf _{n \rightarrow \infty}\left(-a\left(r_{n}\right) u^{q}\left(r_{n}\right)-b\left(r_{n}\right) u^{p}\left(r_{n}\right)-\frac{n-1}{r_{n}} u_{r}\left(r_{n}\right)\right) \geqslant \liminf _{n \rightarrow \infty}\left(-a\left(r_{n}\right) u^{q}\left(r_{n}\right)\right)>0 .
$$

This is a contradiction, and hence we must have $\lim _{r \rightarrow \infty} u(r)=0$ !
It should be interesting and possible to prove some results in non-radial settings. In particular, Brezis and Kamin prove (Lemma 6 of [8]) that under hypothesis (1.2), if we could show that

$$
\frac{1}{|\partial B(0, R)|} \int_{\partial B(0, R)} u \rightarrow 0 \quad \text { as } R \rightarrow 0,
$$

then any smooth solution $u \in L^{\infty}\left(\mathbb{R}^{n}\right)$ of (1.1) has the property that $\lim _{|x| \rightarrow \infty} u(x)=0$. Many open questions remain.

## Appendix A

In this section we give a detailed proof of Lemma 2.8. The analogous result for bounded domains appeared in [3] and originated from [23] (see also [8]). Although the proof is almost the same as in [3], we provide a proof here since the lemma is so crucial to our methods.

For our version we will use Serrin's maximum principle for strong solutions in $W_{l o c}^{2, s}$ [21]. Hence, we note that for solutions $u \in W_{l o c}^{2, s}\left(\mathbb{R}^{n}\right), s>n, \Delta u$ exists a.e. pointwise and coincides with its distributional derivative. In particular, $-\Delta u \geqslant f$ a.e. pointwise is equivalent to $-\Delta u \geqslant f$ in the sense of distributions.

Following [3], we transform our solution $u$ by $U=\frac{1}{1-q} u^{1-q}$, and require also the regularized form $U_{\epsilon}=\frac{1}{1-q}(u+\epsilon)^{1-q}$. For later use, we note that these transformed functions satisfy:

$$
\begin{gathered}
\nabla U_{\epsilon}=(u+\epsilon)^{-q} \nabla u \\
\Delta U_{\epsilon}=(-q)(u+\epsilon)^{-q-1}|\nabla u|^{2}+(u+\epsilon)^{-q} \Delta u \\
\Delta U_{\epsilon}=(-q)(u+\epsilon)^{q-1}\left|\nabla U_{\epsilon}\right|^{2}-a \frac{u^{q}}{(u+\epsilon)^{q}}-b \frac{u^{p}}{(u+\epsilon)^{q}},
\end{gathered}
$$

and similarly for $U$ by setting $\epsilon=0$. We are now ready to prove the comparison result, Lemma 2.8.
Proof of Lemma 2.8. Suppose the contrary, $D:=\left\{x \in \mathbb{R}^{n} \mid u_{1}>u_{2}\right\} \neq \emptyset$. Let $U_{1}=\frac{1}{1-q} u_{1}^{1-q}, U_{2}=$ $\frac{1}{1-q} u_{2}^{1-q}$, so that

$$
U_{1}>U_{2} \text { in } D .
$$

Since $\lim _{|x| \rightarrow \infty} u_{1}(x)=\lim _{|x| \rightarrow \infty} u_{2}(x)=0$, there exists a point $x_{0} \in D$ where the difference $\delta:=U_{1}-$ $U_{2}$ attains its maximal value. Let us now distinguish two cases.

Case 1. Suppose that $U_{2}\left(x_{0}\right)>0$ for some $x_{0} \in D$ where $\delta$ takes its maximal value. Denote by $V$ the maximal connected component of the set $D_{1}:=\left\{x \in D: U_{2}(x)>0\right\}$ containing $x_{0}$, then $\delta$ belongs to $C^{2}(V)$ and the above calculations show that

$$
\Delta \delta=\Delta U_{1}-\Delta U_{2} \geqslant(-q) u_{1}^{q-1}\left|\nabla U_{1}\right|^{2}+q u_{2}^{q-1}\left|\nabla U_{2}\right|^{2}-b\left(u_{1}^{p-q}-u_{2}^{p-q}\right) .
$$

It is easy to see that from $u_{1}>u_{2}$ in $D$ we have

$$
u_{2}^{q-1}>u_{1}^{q-1} \quad \text { in } D, \quad u_{1}^{p-q}>u_{2}^{p-q} \quad \text { in } D .
$$

So

$$
\begin{aligned}
& \Delta \delta+q u_{2}^{q-1}\left(\left(\nabla U_{1}+\nabla U_{1}\right), \nabla \delta\right) \\
& \quad \geqslant(-q) u_{1}^{q-1}\left|\nabla U_{1}\right|^{2}+q u_{2}^{q-1}\left|\nabla U_{2}\right|^{2}+q u_{2}^{q-1}\left|\nabla U_{1}\right|^{2}-q u_{2}^{q-1}\left|\nabla U_{2}\right|^{2}-b\left(u_{1}^{p-q}-u_{2}^{p-q}\right) \\
& \quad \geqslant\left(q u_{2}^{q-1}-q u_{1}^{q-1}\right)\left|\nabla U_{1}\right|^{2}-b\left(u_{1}^{p-q}-u_{2}^{p-q}\right) \\
& \\
& \quad \geqslant 0 .
\end{aligned}
$$

Since $\delta$ assumes its maximal value at an interior point of $V$, the weak maximum principle of Serrin [21] ensures that $\delta \equiv$ constant in $V$. It then follows that

$$
0=\nabla \delta=\nabla U_{1}-\nabla U_{2} \quad \text { and } \quad \Delta \delta \equiv 0 \quad \text { in } V .
$$

But by assumption in $V$ we have

$$
0 \equiv \Delta \delta=q\left|\nabla U_{1}\right|^{2}\left(u_{2}^{q-1}-u_{1}^{q-1}\right)-b\left(u_{1}^{p-q}-u_{2}^{p-q}\right) \geqslant 0,
$$

which implies $\nabla U_{1}=0$ in $V$, in turn we have $\nabla U_{2}=0$ in $V$, so $u_{1}$ and $u_{2}$ must be constant in $V$. But since $\lim _{|x| \rightarrow \infty} u_{1}(x)=\lim _{|x| \rightarrow \infty} u_{2}(x)=0$, we must have $u_{1} \equiv u_{2}$ in $V$. This is a contradiction, so Case 1 is impossible.

Case 2. Suppose $U_{2}\left(x_{0}\right)=0$ for all $x_{0}$ where $\delta$ achieves its maximal value. Let

$$
C:=\left\{x \in D: \delta(x)=\delta\left(x_{0}\right)\right\} .
$$

Note by assumption, $U_{2} \equiv 0$ in $C$. Since $\delta=\delta\left(x_{0}\right)>0$ in $C$, we have $U_{1}>0$ in $C$. On the other hand also by assumption $u_{1} \equiv 0$ in $\overline{\Omega^{+}-\Omega_{I}^{+}}$. Hence we have

$$
c \cap \overline{\left(\Omega^{+}-\Omega_{I}^{+}\right)}=\emptyset .
$$

From definition of the set $P, U_{2}>0$ in $\overline{\Omega_{I}^{+}}$,

$$
\subset \cap \overline{\Omega_{I}^{+}}=\emptyset .
$$

So we have

$$
c \cap \overline{\Omega^{+}}=\emptyset,
$$

which means that $C$ and $\overline{\Omega^{+}}$are at a positive distance to each other. Therefore there exists a neighborhood $U$ of $C$ such that $\bar{U} \cap \overline{\Omega^{+}}=\emptyset$ and $\delta(x)>0$ in $\bar{U}$. Then by monotonicity we obtain

$$
\min _{\bar{W}}\left(u_{1}-u_{2}\right)>0,
$$

where $W$ is a connected component of $U$.
Thus there exists $b>0$ such that $\delta(x) \leqslant b<\delta\left(x_{0}\right)$ for $\forall x \in \partial W$. For $\forall \epsilon>0$, we define

$$
U_{2 \epsilon}:=\frac{1}{1-q}\left(u_{2}+\epsilon\right)^{1-q}, \quad \delta_{\epsilon}:=U_{1}-U_{2 \epsilon} .
$$

Clearly $\delta_{\epsilon} \leqslant \delta$ in $D$. We can pick positive $\epsilon$ small enough such that

$$
u_{1}>u_{2}+\epsilon \text { and } \delta_{\epsilon}\left(x_{0}\right)>0 .
$$

It follows that

$$
\delta_{\epsilon}(x) \leqslant \delta(x) \leqslant b<\delta_{\epsilon}\left(x_{0}\right) \quad \forall x \in \partial W .
$$

Hence $\delta_{\epsilon}$ attains its maximum value at some interior point in $W$ and is not constant in $\bar{W}$. On the other hand, from assumption (1.7) we have $a \leqslant 0$ in $\bar{W}$, therefore

$$
\begin{aligned}
\Delta \delta_{\epsilon} & \geqslant(-q) u_{1}^{q-1}\left|\nabla U_{1}\right|^{2}+q\left(u_{2}+\epsilon\right)^{q-1}\left|\nabla U_{2 \epsilon}\right|^{2}+a \frac{u_{2}^{q}}{\left(u_{2}+\epsilon\right)^{q}}-a-b\left(u_{1}^{p-q}-\frac{u_{2}^{p}}{\left(u_{2}+\epsilon\right)^{q}}\right) \\
& \geqslant(-q)\left(u_{2}+\epsilon\right)^{q-1}\left(\left|\nabla U_{1}\right|^{2}-\left|\nabla U_{2 \epsilon}\right|^{2}\right)+a\left(\frac{u_{2}^{q}}{\left(u_{2}+\epsilon\right)^{q}}-1\right)-b\left(u_{1}^{p-q}-\frac{\left(u_{2}+\epsilon\right)^{p}}{\left(u_{2}+\epsilon\right)^{q}}\right) \\
& \geqslant(-q)\left(u_{2}+\epsilon\right)^{q-1}\left(\nabla\left(U_{1}+U_{2, \epsilon}\right) \cdot \nabla \delta_{\epsilon}\right)
\end{aligned}
$$

in $W$.

$$
-\Delta \delta_{\epsilon}+(-q)\left(u_{2}+\epsilon\right)^{q-1}\left(\nabla\left(U_{1}+U_{2, \epsilon}\right) \cdot \nabla \delta_{\epsilon}\right) \leqslant 0
$$

But by the weak maximum principle of Serrin [21], $\delta_{\epsilon}$ cannot achieve its maximum in $W$ unless it is constant. This is a contradiction, whence the result follows.

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[^0]:    * Corresponding author.

    E-mail addresses: alama@mcmaster.ca (S. Alama), qlu@dim.uchile.cl (Q. Lu).
    ${ }^{1}$ Supported by an NSERC Research Grant.
    2 Current address: Centro de Modelamiento Matematico, Av. Blanco Encalada 2120, Santiago, Chile.

