Theoretical study of BTF/TNA cocrystal: Effects of hydrostatic pressure and temperature

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Abstract

Cocrystallization is a promising technique for the design and preparation of new explosives, and the stability of cocrystal is highly concerned by the researchers. In order to make a better understanding of the behavior of cocrystal under the extreme conditions, DFT (density functional theory) calculation is performed to investigate the effect of hydrostatic pressure on geometrical and electronic structures of the cocrystal BTF (benzotrifuroxan)/TNA (2,4,6-trinitroaniline). When the hydrostatic pressure is applied, the lattice constants, volume, density and total energy change gradually except at the pressures of 40 GPa and 79–83 GPa. It is noteworthy that new chemical bonds form when the pressure is up to 83 GPa. The band gap of the cocrystal becomes smaller when the pressure is applied, and finally the cocrystal shows a characteristic of metal. The mechanical property of cocrystal is calculated by MD (molecular dynamics) simulation. The results show that the cocrystal has a better ductibility at low temperature, and has the best tenacity at 295 K.

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Keywords: Cocrystal; BTF/TNA; DFT; High pressure; Crystal structure

1. Introduction

Cocrystal is defined as a new crystal with distinct solid-state properties created by combining two or more neutral species in a definite ratio [1,2]. Cocrystallization is an effective approach to alter the dissolution rate, thermal stability and bioavailability of pharmaceuticals in pharmaceutical industry [3–5]. Its great success in pharmaceutical industry has attracted much attention in other fields, such as optical, semiconducting materials and energetic materials [6–8]. Cocrystallization can modify the performance of energetic material at molecular level compared with coating, crystal morphology modification and other traditional methods [9–11]. It is feasible to tailor the different molecules to get the definite performance according to some rules with cocrystallization method, which has been verified in pharmaceutical fields. Seventeen cocrystals of TNT (2,4,6-trinitrotoluene) were reported and the safety performance has been greatly enhanced [12]. However, the energy is diluted because of the introduction of non-energetic materials. Bolton et al. [13] used the solvent evaporation method to prepare the cocrystal of TNT/CL-20 (2,4,6,8,10,12-hexanitrohexaazaisowurtzitane). The cocrystal has a poorer detonation performance but a reduced sensitivity compared with pure CL-20. The cocrystal of HMX (1,3,5,7-tetranitro-1,3,5,7-tetrazocane)/CL-20 is predicted to be more powerful than HMX [14]. Nevertheless, the sensitivity of the cocrystal is similar to that of HMX. Millar et al. [15] changed the physicochemical properties of CL-20 by cocrystallizing CL-20 with four solvents. It is notable that the CL-20 molecule adopts different conformations before and after desolvation, and this may provide a highly-selective means of polymorph screening.

When a detonation takes place, the high velocity shock waves can result in very high temperatures and pressures [16],
which may lead to a phase transition of energetic material to form new polymorphs or decomposition reactions [17,18]. As a result, the crystal structure would deform. It is well known that the sensitivity and detonation performance of the explosive are intimately bound up with the crystal packing. Therefore it is necessary to investigate the behavior of explosives under the extreme conditions. In recent years, a lot of efforts have been laid both in experimental and theoretical calculation in this respect. Diamond anvil cell, which can generate a pressure of hundreds of GPa, was used to study the changes of materials at high pressure, such as cell parameters, phase transition, and so on [19,20]. Compared with the changes of materials at high pressure, such as cell parameters, calculation in this respect. Diamond anvil cell, which can Therefore it is necessary to investigate the behavior of explosives at high pressure, such as cell parameters, phase transition, and so on [19,20]. Compared with the experiment, theoretical calculations show much more details to help us to understand the changing mechanism. Microscopic changes of geometrical structure and electronic distribution can be studied by theoretical calculation. To date, a lot of widely used explosives, such as HMX, CL-20 [21], RDX (hexahydro-1,3,5-trinitro-1,3,5-triazine) [22], and so on [23–27], at different hydrostatic pressures were studied using the computational methods. As for the cocrystal explosives, it is of great importance to consider their stability at high pressure. However, the formation and decomposition mechanism of the cocrystal explosive is not clear, and its behavior under the extreme conditions is not mentioned. Furthermore, whether the molecular interactions are destroyed in the process of detonation is concerned by the researchers, especially considering the advantages of cocrystal explosives that have been demonstrated. However, to our knowledge, no work has been reported on the study of the behavior of cocrystal explosive at high pressure.

By the previous study, BTF can be cocrystallized with a series of explosives, including TNA, TNT, CL-20, and so on. Generally, the sensitivity of the cocrystal is between those of both BTF and TNA. However, the sensitivity of BTF/TNA [28] is higher than those of BTF and TNA. Hence, BTF/ TNA is more worthy of theoretical investigation. For comparison, a blank test was also performed in the absence of the other partner of the couple. As we all know, the mechanical property is an important indicator of the explosive property, and the shear modulus is related to the sensitivity of explosive. Therefore, the mechanical properties of the cocrystal at different temperatures were calculated by MD simulation.

2. Computational methods

The calculations in this study were performed within the framework of DFT based on CASTEP code [29]. The LDA functional proposed by Ceperley and Alder [30] and parameterized by Perdew and Zunder [31], named CA-PZ, was employed. The core electrons in each atomic species were replaced with the norm-conserved pseudopotentials. The primitive cell was relaxed by using the BFGS method. The cutoff energy of plane wave was set to 600 eV. Brillouin zone sampling was performed by using the Monkhost-Pack scheme with a k-point grid of \(2 \times 2 \times 1\). The values of the kinetic energy cutoff and the k-point grid were determined to ensure the convergence of total energy. In the geometry relaxation, the total energy of the system was less than \(2.0 \times 10^{-5}\) eV, the residual force was less than 0.05 eV/Å, the displacement of atoms was less than 0.002 Å, and the residual bulk stress was less than 0.1 GPa.

All MD simulations were carried out in an isothermal-isobaric ensemble (NPT) with the COMPASS [32] force field in Discover code. The unit cell was taken from single X-ray diffraction. A supercell consisting of \(3 \times 3 \times 2\) unit cells was prepared and put into the periodic simulation cell. Each cell contains 72 BTF molecules, 72 TNA molecules and 2376 atoms in total. Temperature was controlled by using the stochastic collision method proposed by Andersen [33]. The temperatures were set to 195 K, 245 K, 295 K, 345 K and 395 K. The initial velocities were set using the Maxwell–Boltzmann profiles at given temperatures. The velocity Verlet time integration method was used with a time step of 1 fs. The Verlet leapfrog scheme was used to integrate the equations of motion. The total simulation time was set to 2 ns, the equilibration was achieved in the first 1 ns, and the data was collected in the next 1 ns.

3. Results and discussion

3.1. Crystal structures and properties

LDA and GGA often give contradictory results when the crystal structure is optimized [16]. Therefore, in order to benchmark the performances of LDA and GGA, both those functionals were applied to the cocrystal of BTF/TNA as a test. As shown in Table 1. After accomplishing total energy minimization for each unit cell, the optimized lattice parameters and the experimental results shows that the LDA results are in well agreement with the experimental ones compared to the GGA results. As for LDA functional, the relative errors of the lattice constants \(a, b, c, \) and \(\beta\) are \(-3.06\%, \ -0.66\%, +1.05\%, \) and \(+0.07\%, \) respectively. The little difference between the calculated and experimental lattice parameters demonstrates that the method and the parameters used in this calculation are reliable. However, the difference may be due to the intermolecular interactions occurring in the crystal lattice, which are not well described by DFT [34].

The experimental crystallization and structure of BTF/TNA cocrystal are presented in Figs. 1 and 2. Each unit cell has four BTF molecules and four TNA molecules.

The relaxed lattice constants of the cocrystal at different hydrostatic pressures are presented in Fig. 3. On the whole, the

<table>
<thead>
<tr>
<th>Method</th>
<th>(a/\text{Å})</th>
<th>(b/\text{Å})</th>
<th>(c/\text{Å})</th>
<th>(\beta/\text{°})</th>
</tr>
</thead>
<tbody>
<tr>
<td>LDA</td>
<td>9.114 (−3.06)</td>
<td>12.591 (−0.66)</td>
<td>14.478 (+1.05)</td>
<td>97.478 (+0.07)</td>
</tr>
<tr>
<td>PW91</td>
<td>12.054 (+28.2)</td>
<td>14.361 (+13.3)</td>
<td>16.440 (+14.7)</td>
<td>96.667 (−0.76)</td>
</tr>
</tbody>
</table>

\(a\) The values given in the parentheses mean the relative error of the calculated values to the experimental ones in percentage.
lattice constants reduce with the increase in pressure. It is worth noting that all the lattice constants increase abnormally at 40 GPa, while in the region of 79–83 GPa, b and c increase again, but a decreases sharply. It can be deduced that the structural transformation may happen in those regions. Obviously, the largest compression of the unit cell takes place at low pressure below 10 GPa. The higher the pressure is, the smaller the compression ratio is. For example, in the region of 0–10 GPa, the compressions of a, b and c are 9.50%, 6.30% and 9.03%, respectively, which are much larger than 1.80%, 1.03% and 0.92% in the region of 50–60 GPa. In the low-pressure zone, the distance between the molecules is far, so the intermolecular repulsion is not large. Therefore, the crystal is easily compressed compared to chemical bond. On the contrary, in the high-pressure zone, the intermolecular repulsion is much larger, and the crystal is difficult to be compressed. As is evident, the compressibilities along three directions are different, which indicates that the compressibility of the crystal is anisotropic. As the pressure goes up, the compressibility in each direction converges. Additionally, the crystal symmetry maintains the same monocline space group P21/c under all hydrostatic compressions.

The volume, density, and total energy of optimized unit cell as a function of pressure are depicted in Fig. 4. With the increase in pressure, the volume of the unit cell decreases constantly. Obviously, like the lattice parameters, the unit cell volume has the largest compression ratio at low pressure. For instance, the volume is reduced by 3.69% in the region of 50–60 GPa, which is much smaller than that (22.86%) in the region of 10–20 GPa. The change in energy is very similar to that in density, and they all increase with the increase in pressure. The energy and density decrease significantly at the pressure of 40 GPa. As shown in Fig. 5, the distance of N1O1 $$H7A decreases gradually with the increase in pressure. However, at 40 GPa, the length of hydrogen bond is 1.668 Å, which is smaller than 1.790 Å at 30 GPa or 1.764 Å at 50 GPa. Hence, the strength of hydrogen bond increases greatly due to the decrease in bond lengths. The crystal becomes more stable and the total energy decreases. The density and total energy raise linearly with the increase in pressure below 40 GPa. The higher total energy means that the more the explosive releases blasting energy during the explosion, the more powerful the explosive is [16]. Meanwhile, the pressure-induced density enlargement would enhance the detonation performance. As a result, they show the same changing tendency with the increase in pressure. For pure BTF crystal, the lattice constants, bond lengths and some other
parameters of BTF change gradually with the increase in pressure. No isomerization or other particular changes were observed even at the applied pressure of 100 GPa.

3.2. Molecule structure

Both the unit cell and the molecular geometry are affected by the external pressure. Some main geometrical parameters, including bond length and torsion angle, are presented to illuminate the effect of hydrostatic pressure.

As shown in Fig. 6, the imposed compression mostly squeezes out the intermolecular space and causes only a little change in intramolecular geometry in the low-pressure zone. Some bonds, such as N1–O2, C1–C2, C7–C8, C8–N8, C8–C9 and C7–N7, shorten gradually and obviously, while some other bonds, such as N2–O2 and N1–O1, change a little bit. This indicates that an endocyclic ring is more likely to be compressed. N1–O2, N9–O9, N1–O2 and N9–O9 shorten abruptly at 40 GPa. Conversely, C1–C2 and N2–O2 lengthen all of a sudden. At 79 GPa, N1–O2 shorten while N2–O2 lengthen. This indicates that N2–O2 may be broken when the explosive decomposes. Simultaneously, in the TNA molecule, the C–NO2 bonds adjacent to the amino lengthen suddenly, indicating that the C–NO2 bond is the trigger bond of TNA, which is consistent with the previous study [35].

The torsion angles of C1–C6–C5–N5 and C9–C8–C7–N7 are 179.208° and 177.444°, respectively, at 0 GPa, which are close to 180° (Fig. 7). Therefore, the five-membered furoxan ring and the six-membered ring in BTF are almost on one plane, amino and benzene rings in TNA are also on the same plan. In the region of 0–30 GPa, the torsion angle of C1–C6–C5–N5 is sharply decreased to 144.528° because of the effect of pressure, that is to say, the five-membered furoxan ring forms an angle with the benzene ring. On the other hand, the torsion angle of C9–C8–C7–N7 decreases significantly in the region of 0–20 GPa, then it increases at the following 10 GPa. Following this, the torsion angle decreases significantly at 40 GPa, and then the angle maintains around 170°. When the pressure gets higher, the distance between the molecules becomes smaller. So the distance between O1 and H7A decreases from 2.279 Å to 1.790 Å. Consequently, they form a strong hydrogen bonding. With the attraction of O1, H7A moves from the original position. As a result, C7–N7 forms a torsion angle with the benzene ring. After that, the hydrogen bonding becomes stable so that the angle does not show obvious change.

Fig. 8 shows the perspective views of one layer of BTF/TNA cocrystal unit cell at the selected pressures. When the
When the pressure reaches to 64 GPa, O in the nitro group forms a chemical bond with C in the benzene ring.

3.3. Electronic structure

Density of state (DOS) analysis is useful to understand better the change of electronic structure caused by external pressure [36]. The calculated total density of states (TDOS) of BTF/TNA cocrystal at different pressures is shown in Fig. 9. Fermi level is plotted by the dash line. The main characteristics of DOS are summarized as follows: (1) The top of valence band and the bottom of conduction band are mainly due to p orbit, indicating that p orbit plays a key role in chemical reaction; (2) When the pressure is low, the curves of DOS are characterized by obvious peaks, but the peaks widen gradually with the increase in pressure, which means that the band splitting and band dispersion are accompanied by a broadening of DOS; and (3) Moreover, the conduction bands have a tendency of shifting to the lower energy. This leads to the band gap be smaller, and the crystal shows the characteristic of metal [37].

The energy gap between the highest occupied crystal orbit (HOCO) and lowest unoccupied crystal orbit (LUCO) is a main parameter to characterize the electronic structure of a solid [17]. It can be seen from Fig. 10 that the band gap decreases from 2.019 eV to 1.134 eV in the region of 0–30 GPa. This is because of the decrease in intermolecular space under compression which leads to an increase in the overlap of different groups of bands and hence increases the charge overlap and delocalization in the system. The tendency of band gap is similar to that of volume of the bulk. The band gap decreases gradually in the region of 30–39 GPa. Because of the structural transformation, the band gap is increased to 1.432 eV at 40 GPa. The band gap decreases gradually in the region of 41–78 GPa. However, the reduction rate is much smaller than that of it in the region of 0–30 GPa. At 79 GPa, the band gap decreases suddenly with the change of bond length. In the region of 80–82 GPa, the band gap rises perpendicularly. Then the band gap decreases again at 83 GPa, which is a turning point to form the unexpected chemical bonds. The formation of the new chemical bonds results in the enhancement of the delocalization of the crystal. In another word, the electron is more likely to transfer and the band gap reduces. Compared with the lattice constants and geometrical parameters, the band gap is more sensitive to pressure. According to the principle of easiest transition (PET) [33], the smaller the band gap is, the more easily the electron transits from the occupied valence band to the empty conduction band. That means the explosive is more likely to detonate. Therefore, the BTF/TNA cocrystal has a higher sensitivity when the external pressure is applied.

3.4. MD simulation results

BTF/TNA cocrystal crystallizes in a monoclinic P2₁/c space group and contains 8 molecules per unit cell. Since the pressure is higher than 82 GPa, something unreasonable changes happen in the geometry. At 83 GPa, O in the nitro group of TNA forms a chemical bond with C in the benzene ring. More chemical bonds can be found in the unit cell among TNA molecules at 100 GPa. The same situation was also found in TNA crystal.
data of single X-ray diffraction were collected at 135 K, we carried out the calculation at 135 K to confirm the applicability and accuracy of the method and parameters applied in the latter calculations. It can be known from Table 2 that the errors among the optimized lattice constants \(a\), \(b\), \(c\) and the experimental values are \(+0.61\%\), \(+0.60\%\), \(+0.60\%\), respectively, after the system reached to the equilibrium state. The result shows that the method and the parameters are reliable.

The lattice parameters at different temperatures were calculated and presented in Table 2. It is obvious that the lattice parameters rise with the increase in temperature.
The experimental and calculated lattice parameters at different temperatures.

<table>
<thead>
<tr>
<th>$T/K$</th>
<th>135</th>
<th>195</th>
<th>245</th>
<th>295</th>
<th>345</th>
<th>395</th>
</tr>
</thead>
<tbody>
<tr>
<td>$b/\AA$</td>
<td>12.675</td>
<td>12.751</td>
<td>(+0.60)</td>
<td>12.795</td>
<td>12.835</td>
<td>12.868</td>
</tr>
</tbody>
</table>

Note that for the most general crystal structure where all 21 constants are independent, the Reuss modulus depends on only nine of the single-crystal compliances, others were omitted.

<table>
<thead>
<tr>
<th>$T/K$</th>
<th>195</th>
<th>245</th>
<th>295</th>
<th>345</th>
<th>395</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E$</td>
<td>3.88</td>
<td>3.17</td>
<td>2.45</td>
<td>2.23</td>
<td>1.65</td>
</tr>
<tr>
<td>$\nu$</td>
<td>0.43</td>
<td>0.43</td>
<td>0.44</td>
<td>0.44</td>
<td>0.43</td>
</tr>
<tr>
<td>$K$</td>
<td>9.07</td>
<td>8.16</td>
<td>7.26</td>
<td>6.05</td>
<td>4.20</td>
</tr>
<tr>
<td>$G$</td>
<td>1.36</td>
<td>1.11</td>
<td>0.85</td>
<td>0.78</td>
<td>0.57</td>
</tr>
<tr>
<td>$K/G$</td>
<td>6.67</td>
<td>7.35</td>
<td>8.54</td>
<td>7.76</td>
<td>7.37</td>
</tr>
</tbody>
</table>

4. Conclusions

In this study, the theoretical simulations were performed to investigate the behaviors of cocrystal under the extreme conditions. The results show that, when the hydrostatic pressure is applied, the lattice constants, volume, density and total energy change gradually except at 40 GPa and 79–83 GPa. A new chemical bond is found when the pressure is up to 83 GPa. Meanwhile, with the pressure augments, the electron transfers and the band gap decreases significantly. That is to say, the sensitivity of cocrystal becomes higher when the hydrostatic pressure rises. The MD simulation suggests that the cocrystal has a better ductibility at low temperature, and has the best tenacity at 295 K.

Acknowledgments

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References


