

BIFURCATION OF LIMIT CYCLES FROM A n -DIMENSIONAL LINEAR CENTER INSIDE A CLASS OF PIECEWISE LINEAR DIFFERENTIAL SYSTEMS

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ABSTRACT. Let n be an even integer. We study the bifurcation of limit cycles from the periodic orbits of the n -dimensional linear center given by the differential system

$$\dot{x}_1 = -x_2, \quad \dot{x}_2 = x_1, \quad \dots, \quad \dot{x}_{n-1} = -x_n, \quad \dot{x}_n = x_{n-1},$$

perturbed inside a class of piecewise linear differential systems. Our main result shows that at most $(4n - 6)^{n/2-1}$ limit cycles can bifurcate up to first-order expansion of the displacement function with respect to a small parameter. For proving this result we use the averaging theory in a form where the differentiability of the system is not needed.

1. INTRODUCTION AND STATEMENT OF THE MAIN RESULT

Piecewise linear differential systems appear in a natural way in control theory, and in the study of electrical circuits. These systems can present complicated dynamical phenomena such as those exhibited by general non-linear differential systems. One of the main ingredients in the qualitative description of the dynamical behavior of a differential system is the number and the distribution of its limit cycles.

The goal of this paper is to study, in \mathbb{R}^n for all n even, the existence of limit cycles of the control systems of the form

$$(1) \quad \dot{x} = A_0 x + \varepsilon F(x),$$

with $|\varepsilon| \neq 0$ a sufficiently small real parameter, where A_0 is equal to

$$A_0 = \begin{pmatrix} 0 & -1 & 0 & 0 & \dots & 0 & 0 \\ 1 & 0 & 0 & 0 & \dots & 0 & 0 \\ 0 & 0 & 0 & -1 & \dots & 0 & 0 \\ 0 & 0 & 1 & 0 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 & -1 \\ 0 & 0 & 0 & 0 & \dots & 1 & 0 \end{pmatrix}$$

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and $F : \mathbb{R}^n \longrightarrow \mathbb{R}^n$ is given by

$$F(x) = Ax + \varphi(k^T x)b,$$

with $A \in \mathcal{M}_n(\mathbb{R})$, $k, b \in \mathbb{R}^n \setminus \{0\}$ and $\varphi : \mathbb{R} \longrightarrow \mathbb{R}$ is the piecewise linear function

$$(2) \quad \varphi(x_1) = \begin{cases} -1 & \text{if } x_1 \in (-\infty, -1), \\ x_1 & \text{if } x_1 \in [-1, 1], \\ 1 & \text{if } x_1 \in (1, \infty), \end{cases}$$

where $x = (x_1, \dots, x_n)^T$. The independent variable is denoted by t , vectors of \mathbb{R}^n are column vectors, and k^T denotes a transposed vector.

For $\varepsilon = 0$ system (1) becomes

$$(3) \quad \dot{x}_1 = -x_2, \quad \dot{x}_2 = x_1, \quad \dots, \quad \dot{x}_{n-1} = -x_n, \quad \dot{x}_n = x_{n-1}.$$

We note that the origin of (3) is a *global isochronous center*, i.e. all its orbits different from the origin are periodic with the same period 2π . A *limit cycle* of a differential system is an isolated periodic orbit in the set of all periodic orbits of the system. The *Poincaré map* (or equivalently, the *displacement map*) is a suitable tool for studying limit cycles of autonomous systems (detailed explanations can be found in [6] or [8]; see also section 3). We recall that a limit cycle of a system corresponds to an isolated zero of the displacement function.

We must mention that there are many papers studying the bifurcation of limit cycles from the periodic orbits of a center, but almost all of them perturb a 2-dimensional center, see for instance the book [7] and the references quoted there. Of course, there are papers dedicated to perturb centers of dimension > 2 , but not too much. Later on we shall comment some of those papers.

The limit cycles of system (1) for $\varepsilon \neq 0$ sufficiently small that we shall study will be the ones bifurcating from the periodic orbits of the n -dimensional center (3) (i.e. of system (1) with $\varepsilon = 0$). As we shall see later on this study has been made for $n = 2, 4$. Here we shall do it for $n \geq 6$ even. Our main result can be stated as follows.

Theorem 1. *For all $n \geq 6$ even at most $(4n-6)^{n/2-1}$ limit cycles of system (1) with $\varepsilon \neq 0$ sufficiently small can bifurcate from the periodic orbits of the n -dimensional center (3), up to first-order expansion of the displacement function of system (1) with respect to the small parameter ε .*

Theorem 1 will be proved in section 5.

In section 2 we shall present the scheme of the proof of Theorem 1, mainly based in four lemmas.

We emphasize that the bifurcation from $\varepsilon = 0$ to $\varepsilon \neq 0$ in Theorem 1 takes place for $\varepsilon > 0$ and for $\varepsilon < 0$ sufficiently small, i.e. on both sides of the value $\varepsilon = 0$. We remark that in a Hopf bifurcation the limit cycle only appears on one side of the bifurcation value of the parameter, but in our

case in which the limit cycles bifurcate from periodic orbits of the period annulus of a center they appear in both sides of $\varepsilon = 0$.

The proof of Theorem 1 is based on the first-order averaging method. We will present this method in section 3, in the form obtained in [1]. The advantage of this result is that the smoothness assumptions for the vector field of the differential system are minimal. In particular, it can be applied to piecewise linear differential systems, which are not C^2 (not even C^1), as it is required in its classical version, see for instance, Theorem 11.5 of [14]. This non-differential application of the averaging method to control systems was used for the first time in [2]. This method has been used frequently for computing periodic orbits; see for instance [4, 11]. From the paper [3] we can study the stability of the limit cycles of Theorem 1; for more details see remark 16.

Reference [9] can be seen for a theoretical discussion about suitable transformations of high dimensional differential systems which are small perturbations of a center, into the standard form for averaging. The general idea is to relate this change of variables with the first integrals of the unperturbed center.

We would like to add some comments related to our approach to the problem of counting the limit cycles of piecewise linear differential systems. We have chosen here to study bifurcation with respect to a small parameter from the periodic orbits of a center, up to first-order expansion of the displacement map. For some values of the coefficients, this order is sufficient for finding the exact number of limit cycles. But in some cases the first-order expansion of the displacement map can be identically zero, then a higher order averaging theory is needed. The study can be done by using second-, third-, ... order averaging theory. A key point in these studies is the relation between the averaging theory and the displacement map due to the fact that the displacement map of a piecewise linear differential system is analytic in a neighborhood of a limit cycle.

2. SCHEME OF THE PROOF OF THEOREM 1

First in Lemma 2 we shall reduce the n parameters of the vector b in the definition of the function $F(x)$ to one.

Lemma 2. *By a linear change of variables system (1) can be transformed into the system*

$$(4) \quad \dot{x} = A_0x + \varepsilon \bar{A}x + \varepsilon \varphi(x_1) \bar{b},$$

where $\bar{A} \in M_n(\mathbb{R})$ is an arbitrary matrix and $\bar{b} = (b_1, \dots, b_n) = e_l$ for some $l = 1, 3, \dots, n-1$, where e_l is the l -th vector of the canonical basis of \mathbb{R}^n .

The proof of Theorem 1 uses the averaging theory. One of the main difficulties for applying this theory is to write the differential system that we want to study in the normal form of the averaging theory. For our differential system (1) Lemma 3 provides the change of variables which writes system

(1) in the normal form of the averaging theory. The change of variables of Lemma 3 is defined using the first integrals of system (1) with $\varepsilon = 0$.

Lemma 3. *Changing the variables (x_1, x_2, \dots, x_n) to $(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)})$ by using*

$$\begin{aligned} x_1 &= r \cos \theta, & x_2 &= r \sin \theta, \\ x_3 &= r_1 \cos(\theta + \theta_1), & x_4 &= r_1 \sin(\theta + \theta_1), \\ & & & \vdots \\ x_{n-1} &= r_{(n/2-1)} \cos(\theta + \theta_{(n/2-1)}), & x_n &= r_{(n/2-1)} \sin(\theta + \theta_{(n/2-1)}), \end{aligned}$$

system (4) is transformed into a system of the form

$$(5) \quad \begin{aligned} \frac{dr}{d\theta} &= \varepsilon H_1(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2), \\ \frac{dr_1}{d\theta} &= \varepsilon H_2(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2), \\ \frac{d\theta_1}{d\theta} &= \varepsilon H_3(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2), \\ & \vdots \\ \frac{dr_{(n/2-1)}}{d\theta} &= \varepsilon H_{n-2}(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2), \\ \frac{d\theta_{(n/2-1)}}{d\theta} &= \varepsilon H_{n-1}(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2). \end{aligned}$$

where

$$\begin{aligned} H_1 &= \cos \theta F_1 + \sin \theta F_2, \\ H_j &= \cos(\theta + \theta_{\frac{j}{2}}) F_{j+1} + \sin(\theta + \theta_{\frac{j}{2}}) F_{j+2}, \\ H_{j+1} &= -\frac{1}{r} \cos \theta F_2 + \frac{1}{r} \sin \theta F_1 + \frac{1}{r_{\frac{j}{2}}} \cos(\theta + \theta_{\frac{j}{2}}) F_{j+2} - \frac{1}{r_{\frac{j}{2}}} \sin(\theta + \theta_{\frac{j}{2}}) F_{j+1}, \end{aligned}$$

for $j = 2, 4, \dots, n-2$, and for $i = 1, 2, \dots, n$ we have that $F_i = (a_{i1} \cos \theta + a_{i2} \sin \theta)r + \sum_{l=3}^{\frac{n}{2}+1} (a_{i(2l-3)} \cos(\theta + \theta_{l-2}) + a_{i(2l-2)} \sin(\theta + \theta_{l-2})) r_{l-2} + \varphi(r \cos \theta) b_i$.

We take ε_0 sufficiently small, m arbitrarily large and

$$D_m = \left\{ (r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) \in \left(\frac{1}{m}, m \right) \times \left[\mathbb{S}^1 \times \left(\frac{1}{m}, m \right) \right]^{\frac{n}{2}-1} \right\}.$$

Then the vector field of system (5) is well defined and continuous on $\mathbb{S}^1 \times D_m \times (-\varepsilon_0, \varepsilon_0)$. Moreover the system is 2π -periodic with respect to variable θ and locally Lipschitz with respect to variables $(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)})$.

We will apply the averaging theory stated in Theorem 6 (see Section 3) to system (5). Our next step is to find the corresponding averaged function (8). We will denote it by $h : D_m \rightarrow \mathbb{R}^{n-1}$, $h = (h_1, \dots, h_{n-1})^T$.

Lemma 4. *The components of the averaged function h are given by*

$$\begin{aligned}
 h_1 &= c_1 r + \sum_{l=3}^{\frac{n}{2}+1} \left(c_{1,(2l-4)} \cos \theta_{l-2} + c_{1,(2l-3)} \sin \theta_{l-2} \right) r_{l-2} + b_1 I_1(r), \\
 (6) \quad h_j &= (c_j \cos \theta_{\frac{j}{2}} + c_{(j+1)} \sin \theta_{\frac{j}{2}}) r + \\
 &\quad \sum_{l=3}^{\frac{n}{2}+1} \left(c_{j,(2l-4)} \cos(\theta_{\frac{j}{2}} - \theta_{l-2}) + c_{j,(2l-3)} \sin(\theta_{\frac{j}{2}} - \theta_{l-2}) \right) r_{l-2} + \\
 &\quad b_{j+1} \cos \theta_{\frac{j}{2}} I_1(r), \\
 h_{j+1} &= c + \sum_{l=3}^{\frac{n}{2}+1} \left(c_{1,(2l-3)} \cos \theta_{l-2} - c_{1,(2l-4)} \sin \theta_{l-2} \right) \frac{r_{l-2}}{r} + \\
 &\quad \left(c_{(j+1)} \cos \theta_{\frac{j}{2}} - c_j \sin \theta_{\frac{j}{2}} \right) \frac{r}{r_{\frac{j}{2}}} + \\
 &\quad \sum_{l=3}^{\frac{n}{2}+1} \left(c_{j,(2l-3)} \cos(\theta_{\frac{j}{2}} - \theta_{l-2}) - c_{j,(2l-4)} \sin(\theta_{\frac{j}{2}} - \theta_{l-2}) \right) \frac{r_{l-2}}{r_{\frac{j}{2}}} - \\
 &\quad \frac{b_{j+1}}{r_{\frac{j}{2}}} \sin \theta_{\frac{j}{2}} I_1(r),
 \end{aligned}$$

for $j = 2, 4, \dots, n-2$, where the coefficients c 's are given by

$$\begin{aligned}
 c &= (a_{12} - a_{21})\pi, \quad c_1 = (a_{11} + a_{22})\pi, \\
 c_j &= (a_{(j+1)1} + a_{(j+2)2})\pi, \quad c_{(j+1)} = (a_{(j+2)1} - a_{(j+1)2})\pi, \\
 c_{1,(2l-4)} &= (a_{1(2l-3)} + a_{2(2l-2)})\pi, \quad c_{1,(2l-3)} = (a_{1(2l-2)} - a_{2(2l-3)})\pi, \\
 c_{j,(2l-4)} &= (a_{(j+1)(2l-3)} + a_{(j+2)(2l-2)})\pi, \\
 c_{j,(2l-3)} &= (a_{(j+2)(2l-3)} - a_{(j+1)(2l-2)})\pi,
 \end{aligned}$$

for $l = 3, \dots, n/2 + 1$, and the coefficients a_{rs} are the elements of the matrix \overline{A} of Lemma 2.

Lemmas 2, 3 and 4 are proved in Section 4.

Let $f : D \rightarrow \mathbb{R}^n$ be a C^1 function, with $f(a) = 0$, where D is an open subset of \mathbb{R}^n and $a \in D$. The point a is a *simple zero* of f if the Jacobian $J_f(a)$ of f at a is not zero.

The following result is a key point for proving Theorem 1.

Lemma 5. *From each periodic orbit of the differential system (3), with $n \geq 2$ even, which corresponds to a simple zero in D_m of the function*

$h = (h_1, \dots, h_{n-1})$ given in (6), a branch of limit cycles bifurcates for the differential system (4).

In [10] the authors studied the bifurcation of limit cycles from a two-dimensional center inside \mathbb{R}^n . In this case they proved that the maximum number of limit cycles that can bifurcate from the periodic orbits of the linear differential center is 1, up to first-order expansion of the displacement function of the perturbed system with respect to the small parameter ε .

The main result of [2] states that for $n = 4$ at most three limit cycles of the piecewise linear differential system (1) bifurcate from the periodic orbits of system (3), up to first-order expansion of the displacement function of (1) with respect to the small parameter ε . Moreover, this upper bound is reached.

In [5] we study the bifurcation of limit cycles from the periodic orbits of a 2-dimensional and 4-dimensional linear center in \mathbb{R}^n perturbed inside a class of discontinuous piecewise linear differential systems.

Now for $n \geq 6$ it is very difficult to obtain the exact number of limit cycles of system (1) that bifurcate from the periodic orbits of system (3). Thus the study of the zeros of the averaged function $h = (h_1, \dots, h_{n-1})$ for $n \geq 6$ becomes very difficult with respect to the cases $n = 2, 4$. So for $n \geq 6$ we can only provide the upper bound given in Theorem 1 for the number of limit cycles that can bifurcate from the periodic orbits of system (3).

The proof of Lemma 5 and Theorem 1 will be the subject of sections 4 and 5 respectively.

3. FIRST-ORDER AVERAGING METHOD

The aim of this section is to present the first-order averaging method as obtained in [1]. Differentiability of the vector field is not needed. The specific conditions for the existence of a simple isolated zero of the averaged function are given in terms of the Brouwer degree. In fact, the Brouwer degree theory is the key point in the proof of this theorem. We remind here that continuity of some finite dimensional function is a sufficient condition for the existence of its Brouwer degree (see [12] for precise definitions).

Theorem 6. *We consider the following differential system*

$$(7) \quad \dot{x}(t) = \varepsilon H(t, x) + \varepsilon^2 R(t, x, \varepsilon),$$

where $H : \mathbb{R} \times D \rightarrow \mathbb{R}^n$, $R : \mathbb{R} \times D \times (-\varepsilon_f, \varepsilon_f) \rightarrow \mathbb{R}^n$ are continuous functions, T -periodic in the first variable, and D is an open subset of \mathbb{R}^n . We define $h : D \rightarrow \mathbb{R}^n$ as

$$(8) \quad h(z) = \int_0^T H(s, z) ds,$$

and assume that:

- (i) H and R are locally Lipschitz with respect to x ;

- (ii) for $a \in D$ with $h(a) = 0$, there exists a neighborhood V of a such that $h(z) \neq 0$ for all $z \in \overline{V} \setminus \{a\}$ and $d_B(h, V, 0) \neq 0$ (here $d_B(h, V, 0)$ denote the Brouwer degree of h at 0).

Then, for $|\varepsilon| > 0$ sufficiently small, there exists an isolated T -periodic solution $\psi(\cdot, \varepsilon)$ of system (7) such that $\psi(0, \varepsilon) \rightarrow a$ as $\varepsilon \rightarrow 0$.

Here we will need some facts from the proof of Theorem 6. Hypothesis (i) assures the existence and uniqueness of the solution of each initial value problem on the interval $[0, T]$. Hence, for each $z \in D$, it is possible to denote by $x(\cdot, z, \varepsilon)$ the solution of (7) with the initial value $x(0, z, \varepsilon) = z$. We consider also the displacement function $\zeta : D \times (-\varepsilon_f, \varepsilon_f) \rightarrow \mathbb{R}^n$ defined by

$$\zeta(z, \varepsilon) = \int_0^T [\varepsilon H(t, x(t, z, \varepsilon)) + \varepsilon^2 R(t, x(t, z, \varepsilon), \varepsilon)] dt.$$

From the proof of Theorem 6 we extract the following facts.

Remark 7. For every $z \in D$ the following relation holds

$$x(T, z, \varepsilon) - x(0, z, \varepsilon) = \zeta(z, \varepsilon).$$

The function ζ can be written in the form

$$\zeta(z, \varepsilon) = \varepsilon h(z) + \varepsilon^2 O(1),$$

where h is given by (8) and the symbol $O(1)$ denotes a bounded function on every compact subset of $D \times (-\varepsilon_f, \varepsilon_f)$. Moreover, for $|\varepsilon|$ sufficiently small, $z = \psi(0, \varepsilon)$ is an isolated zero of $\zeta(\cdot, \varepsilon)$.

Note that from Remark 7 it follows that a zero z of the displacement function $\zeta(z, \varepsilon)$ at time T provides initial conditions for a periodic orbit of the system of period T . We also remark that $h(z)$ is the displacement function up to terms of order ε . Consequently the zeros of $h(z)$, when $h(z)$ is not identically zero, also provides periodic orbits of period T .

For a given systems there is the possibility that the function ζ is not globally differentiable, but the function h is C^1 as we shall see in Section 5. In fact, only differentiability in some neighborhood of a fixed isolated zero of h could be enough. When this is the case, one can use the following remark in order to verify the hypothesis (ii) of Theorem 6.

Remark 8. Let $h : D \rightarrow \mathbb{R}^n$ be a C^1 function, with $h(a) = 0$, where D is an open subset of \mathbb{R}^n and $a \in D$. Whenever a is a simple zero of h then there exists a neighborhood V of a such that $h(z) \neq 0$ for all $z \in \overline{V} \setminus \{a\}$ and $d_B(h, V, a) \in \{-1, 1\}$.

4. PROOF OF LEMMA 5

In this section we prove Lemmas 2, 3 and 4 and we give some preliminary results necessary for the proof of Lemma 5 and Theorem 1.

Proof of Lemma 2. Since the linear change of variables $x = Jy$, with J an invertible matrix, transforms system (1) into

$$\dot{y} = J^{-1}A_0Jy + \varepsilon J^{-1}AJy + \varepsilon\varphi(k^T Jy)J^{-1}b,$$

we have to find J such that

$$(9) \quad J^{-1}A_0J = A_1,$$

$$(10) \quad k^T = e_1^T J^{-1},$$

$$(11) \quad J^{-1}b = \bar{b}.$$

We denote by z_{ij} , for $i, j = 1, \dots, n$ the elements of the matrix J^{-1} . Using equations (9) and (10), easy computations show that J^{-1} is given by

$$\begin{pmatrix} k_1 & k_2 & k_3 & k_4 & \dots & k_{n-1} & k_n \\ -k_2 & k_1 & -k_4 & k_3 & \dots & -k_n & k_{n-1} \\ z_{31} & z_{32} & z_{33} & z_{34} & \dots & z_{3(n-1)} & z_{3n} \\ -z_{32} & z_{31} & -z_{34} & z_{33} & \dots & -z_{3n} & z_{3(n-1)} \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ z_{(n-1)1} & z_{(n-1)2} & z_{(n-1)3} & z_{(n-1)4} & \dots & z_{(n-1)(n-1)} & z_{(n-1)n} \\ -z_{(n-1)2} & z_{(n-1)1} & -z_{(n-1)4} & z_{(n-1)3} & \dots & -z_{(n-1)n} & z_{(n-1)(n-1)} \end{pmatrix}.$$

If we take $z_{rr} = 1$, for $r = 3, 5, \dots, n-1$, $z_{ij} = 0$ (of course, except for the elements mentioned previously), $k_1 = b_1/(b_1^2 + b_2^2)$ and $k_2 = b_2/(b_1^2 + b_2^2)$ in the expression of J^{-1} , then equation (11) is satisfied with $\bar{b} = e_1$. In this case we obtain a matrix J^{-1} whose determinant is $k_1^2 + k_2^2$.

On the other hand, if we take $z_{31} = z_{53} = \dots = z_{l(l-2)} = 1$, $z_{rr} = 1$, for $r = l+2, l+4, \dots, n-1$, $z_{1(l-2)} = k_{(l-2)} = 1$, $z_{1(l-1)} = k_{(l-1)} = 0$, $k_l = -b_l/(b_l^2 + b_{(l+1)}^2)$, $k_{(l+1)} = -b_{(l+1)}/(b_l^2 + b_{(l+1)}^2)$ and $z_{ij} = 0$ for the remaining elements, in the expression of J^{-1} , then equation (11) is satisfied with $\bar{b} = e_l$, where $l = 3, 5, \dots, n-1$. In this case we obtain for each $l = 3, 5, \dots, n-1$ a matrix J^{-1} whose determinant is $-(k_l^2 + k_{(l+1)}^2)$.

We recall that we work in the hypotheses that neither k , nor b is the null vector. So at least one of the above statements (i.e. when $\bar{b} = e_l$ for some $l = 1, 3, \dots, n-1$) is well defined. Moreover at least one of the above expressions for the determinant of J^{-1} is nonzero. Hence there exists the change of variables $x = Jy$. This completes the proof of Lemma 2. \square

Now we prove Lemma 3.

Proof of Lemma 3. System (4) in the variables $(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)})$ becomes

$$\begin{aligned}\dot{\theta} &= 1 + \frac{\varepsilon}{r}(\cos \theta F_2 - \sin \theta F_1), \\ \dot{r} &= \varepsilon H_1(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}), \\ \dot{r}_1 &= \varepsilon H_2(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}), \\ \dot{\theta}_1 &= \varepsilon H_3(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}), \\ &\vdots \\ \dot{r}_{(n/2-1)} &= \varepsilon H_{n-2}(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}), \\ \dot{\theta}_{(n/2-1)} &= \varepsilon H_{n-1}(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}).\end{aligned}$$

We note that for $|\varepsilon|$ sufficiently small $\dot{\theta}(t) > 0$ for each t when $(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) \in \mathbb{S}^1 \times D_m$. Now we eliminate the variable t in the above system by considering θ as the new independent variable. It is easy to see that the right-hand side of the new system is well defined and continuous on $\mathbb{S}^1 \times D_m \times (-\varepsilon_0, \varepsilon_0)$, it is 2π -periodic with respect to the independent variable θ and locally Lipschitz with respect to $(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)})$. Form (5) is obtained after an expansion with respect to the small parameter ε . \square

Our next step is to prove Lemma 4. In order to calculate the expression of h given by (6), we will use the following formulas

$$\begin{aligned}\int_0^{2\pi} \cos^2 \theta d\theta &= \pi, & \int_0^{2\pi} \sin^2 \theta d\theta &= \pi, & \int_0^{2\pi} \cos(\theta + s) \sin(\theta + s) d\theta &= 0, \\ \int_0^{2\pi} \cos \theta \cos(\theta + s) d\theta &= \pi \cos s, & \int_0^{2\pi} \cos \theta \sin(\theta + s) d\theta &= \pi \sin s, \\ \int_0^{2\pi} \sin \theta \cos(\theta + s) d\theta &= -\pi \sin s, & \int_0^{2\pi} \sin \theta \sin(\theta + s) d\theta &= \pi \cos s, \\ \int_0^{2\pi} \cos^2(\theta + s) d\theta &= \pi, & \int_0^{2\pi} \sin^2(\theta + s) d\theta &= \pi, & \int_0^{2\pi} \cos \theta \sin \theta d\theta &= 0, \\ \int_0^{2\pi} \cos(\theta + s_1) \sin(\theta + s_2) d\theta &= \pi \sin(s_2 - s_1), \\ \int_0^{2\pi} \sin(\theta + s_1) \cos(\theta + s_2) d\theta &= -\pi \sin(s_2 - s_1), \\ \int_0^{2\pi} \sin(\theta + s_1) \sin(\theta + s_2) d\theta &= \int_0^{2\pi} \cos(\theta + s_1) \cos(\theta + s_2) d\theta = \\ &= \pi \cos(s_2 - s_1),\end{aligned}$$

for all $s, s_1, s_2 \in \mathbb{R}$.

For each $r > 0$ we define

$$I_1(r) = \int_0^{2\pi} \varphi(r \cos \theta) \cos \theta d\theta,$$

$$I_2(r) = \int_0^{2\pi} \varphi(r \cos \theta) \sin \theta d\theta,$$

where φ is the piecewise linear function given by (2).

The next lemma is proved in [2].

Lemma 9. *The integrals I_1 and I_2 satisfy*

$$I_2(r) = 0 \text{ for all } r > 0,$$

and

$$(12) \quad I_1(r) = \begin{cases} \pi r & \text{if } 0 < r \leq 1, \\ \pi r + \frac{2}{r} \sqrt{r^2 - 1} - 2r \arctan(\sqrt{r^2 - 1}) & \text{if } r > 1. \end{cases}$$

Proof of Lemma 4. We must compute

$$(13) \quad h_j(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) = \int_0^{2\pi} H_j(\theta, r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) d\theta,$$

for $j = 1, \dots, n-1$. Using Lemma 9, from (13) we obtain the expressions given in (6) for the components of the function h . This ends the proof of Lemma 4. \square

For each point $(r^*, \theta_1^*, r_1^*, \dots, \theta_{(n/2-1)}^*, r_{(n/2-1)}^*) \in D_m$ of system (6) we have a periodic orbit of system (3), because

$$r = \sqrt{x_1^2 + x_2^2}, \quad r_i = \sqrt{x_{(2i+1)}^2 + x_{(2i+2)}^2}, \quad \theta_i = \arctan \frac{x_{(2i+2)}}{x_{(2i+1)}} - \arctan \frac{x_2}{x_1},$$

for $i = 1, \dots, n/2 - 1$ are first integrals of system (3).

Of course any isolated 2π -periodic solution of (5) with $|\varepsilon| \neq 0$ sufficiently small, corresponds to a limit cycle of (4).

Proof of Lemma 5. Lemma 3 states that the hypotheses of Theorem 6 are fulfilled for system (5), where the function h is given by (6). Using also Remark 8 we conclude that for $|\varepsilon|$ sufficiently small and for each simple zero $(r^*, \theta_1^*, r_1^*, \dots, \theta_{(n/2-1)}^*, r_{(n/2-1)}^*) \in D_m$ of h , there exists an isolated 2π -periodic solution $\varphi(\cdot, \varepsilon)$ of system (5) such that $\varphi(0, \varepsilon) \rightarrow (r^*, \theta_1^*, r_1^*, \dots, \theta_{(n/2-1)}^*, r_{(n/2-1)}^*)$ as $\varepsilon \rightarrow 0$. In short we have proved the Theorem 5. \square

For a proof of the next result see Lemma 3.4 of [2].

Lemma 10. *The displacement function of system (4) for the transversal section $x_2 = 0$, written in the coordinates of Lemma 3, has the form*

$$\varepsilon h(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) + O(\varepsilon^2).$$

5. PROOF OF THEOREM 1

The proof of Theorem 1 will be the subject of this section. We note that in order to find the zeros of h in D_m it is sufficient to look for them in $(0, \infty) \times \left[\mathbb{S}^1 \times (0, \infty) \right]^{\frac{n}{2}-1}$. This is due to the fact that m can be chosen arbitrarily large, and h , as well as the transformation of Lemma 3 are 2π -periodic with respect to the variables $\theta_1, \dots, \theta_{(n/2-1)}$.

Before proving Theorem 1 for arbitrary $n \geq 6$, we shall provide a more detailed proof for the particular case $n = 6$.

Proposition 11. *Consider system (6) with $n = 6$. Let $h : (0, \infty) \times [\mathbb{S}^1 \times (0, \infty)]^2 \rightarrow \mathbb{R}^5$ be the function $h = (h_1, h_2, h_3, h_4, h_5)$ whose components are given by (6). Then*

- (i) h is of class C^1 ;
- (ii) the maximum number of isolated zeros of h in $(0, \infty) \times [\mathbb{S}^1 \times (0, \infty)]^2$ is 324.

5.1. Proof of Proposition 11. Function h is a composition of some elementary functions and the function I_1 . A direct study of I_1 shows that it is of class C^1 on $(0, \infty)$. Thus statement (i) holds. We will divide the proof of the statement (ii) into Lemmas 13, 14 and 15.

The following result will be needed later on. For a proof see Lemma 4.2 of [2].

Lemma 12. *We consider the equation*

$$(14) \quad I_1(r) = \lambda r, \quad r > 0,$$

where I_1 is given by (12) and λ is a real parameter. Then we have the following situations.

- (a) If $0 < \lambda < \pi$, then (14) has a unique solution $r^* > 1$.
- (b) If $\lambda = \pi$, then (14) has the interval $(0, 1]$ as the set of solutions.
- (c) If $\lambda > \pi$ or $\lambda \leq 0$, then (14) has no solution.

Our next goal is to find zeros of $h = 0$ when $n = 6$. The components of h are given by

$$\begin{aligned}
(15) \quad h_1 &= c_1 r + (c_{1,2} \cos \theta_1 + c_{1,3} \sin \theta_1) r_1 + \\
&\quad + (c_{1,4} \cos \theta_2 + c_{1,5} \sin \theta_2) r_2 + b_1 I_1(r) \\
h_2 &= (c_2 \cos \theta_1 + c_3 \sin \theta_1) r + c_{2,2} r_1 + \\
&\quad (c_{2,4} \cos(\theta_1 - \theta_2) + c_{2,5} \sin(\theta_1 - \theta_2)) r_2 + b_3 \cos \theta_1 I_1(r), \\
h_3 &= (c + c_{2,3}) + (c_{1,3} \cos \theta_1 - c_{1,2} \sin \theta_1) \frac{r_1}{r} + \\
&\quad (c_{1,5} \cos \theta_2 - c_{1,4} \sin \theta_2) \frac{r_2}{r} + (c_3 \cos \theta_1 - c_2 \sin \theta_1) \frac{r}{r_1} + \\
&\quad (c_{2,5} \cos(\theta_1 - \theta_2) - c_{2,4} \sin(\theta_1 - \theta_2)) \frac{r_2}{r_1} - \frac{b_3}{r_1} \sin \theta_1 I_1(r), \\
h_4 &= (c_4 \cos \theta_2 + c_5 \sin \theta_2) r + (c_{4,2} \cos(\theta_2 - \theta_1) + c_{4,3} \sin(\theta_2 - \theta_1)) r_1 + \\
&\quad c_{4,4} r_2 + b_5 \cos \theta_2 I_1(r), \\
h_5 &= (c + c_{4,5}) + (c_{1,3} \cos \theta_1 - c_{1,2} \sin \theta_1) \frac{r_1}{r} + \\
&\quad (c_{1,5} \cos \theta_2 - c_{1,4} \sin \theta_2) \frac{r_2}{r} + (c_5 \cos \theta_2 - c_4 \sin \theta_2) \frac{r}{r_2} + \\
&\quad (c_{4,3} \cos(\theta_2 - \theta_1) - c_{4,2} \sin(\theta_2 - \theta_1)) \frac{r_1}{r_2} - \frac{b_5}{r_2} \sin \theta_2 I_1(r).
\end{aligned}$$

By Lemma 2 we get (b_1, b_3, b_5) equal to $(1, 0, 0)$ or $(0, 1, 0)$ or $(0, 0, 1)$. Each one of the three next lemmas make reference to one of these cases.

Lemma 13. *Consider $(b_1, b_3, b_5) = (1, 0, 0)$. So system (15) can have at most 324 zeros.*

Proof. Using $h_2 = h_4 = 0$ we obtain

$$r_1 = \frac{f_1(\theta_1, \theta_2)}{g_1(\theta_1, \theta_2)} r \quad \text{and} \quad r_2 = \frac{f_2(\theta_1, \theta_2)}{g_1(\theta_1, \theta_2)} r,$$

and we get $r \neq 0$. Consider $\tilde{r}_1 = g_1(\theta_1, \theta_2)^2 r_1 / r$ and $\tilde{r}_2 = g_1(\theta_1, \theta_2)^2 r_2 / r$.

With the notations $\cos \theta_1 = x$, $\sin \theta_1 = z$, $\cos \theta_2 = y$, $\sin \theta_2 = w$, multiplying $h_3 = 0$ by \tilde{r}_1 and $h_5 = 0$ by \tilde{r}_2 , we obtain a system with four equations and four unknowns (x, y, z, w) given by

$$\begin{aligned}
(16) \quad \tilde{h}_3 &= 0 \\
\tilde{h}_5 &= 0 \\
x^2 + z^2 - 1 &= 0 \\
y^2 + w^2 - 1 &= 0.
\end{aligned}$$

Using the software *Mathematica* we get that the two first equations of (16) are polynomial functions of degree 9. By Bezout's Theorem (see [13]), the

polynomial system (16) has at most 324 solutions. Using the first equation of (15) and Lemma 12, to each solution $(x^*, y^*, z^*, w^*) \in (-1, 1) \times (-1, 1) \times (-1, 1) \times (-1, 1)$ corresponds, at most, one solution r^* of (14) and consequently, one solution r_1^* and one solution r_2^* . \square

Lemma 14. *Consider $(b_1, b_3, b_5) = (0, 1, 0)$. So system (15) can have at most 224 zeros.*

Proof. Using $h_1 = h_4 = 0$ we obtain

$$r_1 = \frac{f_3(\theta_1, \theta_2)}{g_2(\theta_1, \theta_2)}r \quad \text{and} \quad r_2 = \frac{f_4(\theta_1, \theta_2)}{g_2(\theta_1, \theta_2)}r,$$

and we get $r \neq 0$. Consider $\tilde{r}_1 = g_2(\theta_1, \theta_2)^2 r_1 / r$ and $\tilde{r}_2 = g_2(\theta_1, \theta_2)^2 r_2 / r$.

With the notations $\cos \theta_1 = x$, $\sin \theta_1 = z$, $\cos \theta_2 = y$, $\sin \theta_2 = w$, multiplying $h_5 = 0$ by \tilde{r}_2 and considering $\tilde{h}_2 = g_2(\theta_1, \theta_2)^2 z h_2 / r + \tilde{r}_1 x h_3$, we obtain a system with four equations and four unknowns (x, y, z, w) given by

$$(17) \quad \begin{aligned} \tilde{h}_2 &= 0 \\ \tilde{h}_5 &= 0 \\ x^2 + z^2 - 1 &= 0 \\ y^2 + w^2 - 1 &= 0. \end{aligned}$$

Using the software *Mathematica* we get that the first and the second equations of (17) are polynomial functions of degree 8 and 7 respectively. By Bezout's Theorem, the polynomial system (17) has at most 224 solutions. Using the second equation of (15) and Lemma 12, to each solution $(x^*, y^*, z^*, w^*) \in (-1, 1) \times (-1, 1) \times (-1, 1) \times (-1, 1)$ corresponds, at most, one solution r^* of (14) and consequently, one solution r_1^* and one solution r_2^* . \square

Lemma 15. *Consider $(b_1, b_3, b_5) = (0, 0, 1)$. So system (15) can have at most 224 zeros.*

Proof. Using $h_1 = h_2 = 0$ we obtain

$$r_1 = \frac{f_5(\theta_1, \theta_2)}{g_3(\theta_1, \theta_2)}r \quad \text{and} \quad r_2 = \frac{f_6(\theta_1, \theta_2)}{g_3(\theta_1, \theta_2)}r,$$

and we get $r \neq 0$. Consider $\tilde{r}_1 = g_3(\theta_1, \theta_2)^2 r_1 / r$ and $\tilde{r}_2 = g_3(\theta_1, \theta_2)^2 r_2 / r$.

With the notations $\cos \theta_1 = x$, $\sin \theta_1 = z$, $\cos \theta_2 = y$, $\sin \theta_2 = w$, multiplying $h_3 = 0$ by \tilde{r}_1 and considering $\tilde{h}_4 = g_3(\theta_1, \theta_2)^2 w h_4 / r + \tilde{r}_2 y h_5$, we obtain a system with four equations and four unknowns (x, y, z, w) given by

$$(18) \quad \begin{aligned} \tilde{h}_3 &= 0 \\ \tilde{h}_4 &= 0 \\ x^2 + z^2 - 1 &= 0 \\ y^2 + w^2 - 1 &= 0. \end{aligned}$$

Using the software *Mathematica* we get that the first and the second equations of (18) are polynomial functions of degree 7 and 8 respectively. By Bezout's Theorem, the polynomial system (18) has at most 224 solutions.

Using the fourth equation of (15) and Lemma 12, to each solution $(x^*, y^*, z^*, w^*) \in (-1, 1) \times (-1, 1) \times (-1, 1) \times (-1, 1)$ corresponds, at most, one solution r^* of (14) and consequently, one solution r_1^* and one solution r_2^* . \square

5.2. Proof of Theorem 1 for the general case. As for the case $n = 6$ we note that the function h is a composition of some elementary functions and the C^1 function I_1 .

Proof of Theorem 1. Consider $n \geq 6$ and $\bar{b} = e_1$. The subsystem, obtained from (6), formed by all the equations $h_j = 0$, where j is even, is a $(n/2 - 1) \times (n/2 - 1)$ linear system in the variables r_i , $i = 1, \dots, n/2 - 1$. By Cramer's Rule we get $r_i = \Delta r_i / \Delta_n$, where Δ_n is the determinant of the matrix (P_{ji}) with $P_{ji} = P_{ji}(y_{j/2}, w_{j/2}, y_i, w_i)$, a polynomial of degree 2, being the coefficient of r_i in h_j and Δr_i is equal to Δ_n replacing the column i by the vector $(B_j r)$, where $-B_j$ is the coefficient of r in h_j . An easy computation shows that $r_i = r f_i / \Delta_n$ where f_i is an appropriated function.

It is easy to see that, for each n , $\deg \Delta_n = (\deg \Delta_{n-2}) + 2 = n - 2$, where $\deg p$ denotes the degree of the polynomial p (note that Δ_n has $n/2 - 1$ rows).

Multiplying h_{j+1} by $(r_{j/2}(\Delta_n)^2)/r$, we obtain a polynomial of degree $2n - 3$ in the variables $(y_1, w_1, \dots, y_{n/2-1}, w_{n/2-1})$ where $y_k = \cos \theta_k$ and $w_k = \sin \theta_k$ for $k = 1, \dots, n/2 - 1$. By Bezout's Theorem, the subsystem, obtained of (6), formed by all the equations $h_{j+1} = 0$, where $j+1 \geq 3$ is odd, together with the equations $y_k^2 + w_k^2 - 1 = 0$ for $k = 1, \dots, n/2 - 1$ has $(4n - 6)^{n/2-1}$ solutions $(y_1^*, w_1^*, \dots, y_{n/2-1}^*, w_{n/2-1}^*)$. Using the first equation of (6) and Lemma 12, for each one of the previous solutions, we obtain at most one r^* , and consequently at most one r_i^* , $i = 1, \dots, n/2 - 1$.

If $\bar{b} = e_j$ we replace h_j by h_1 in the evaluation to obtain r_i , $i = 1, \dots, n/2 - 1$. As consequence we apply the Bezout's Theorem for polynomials of smaller degree. So the number of solutions when $\bar{b} \neq e_1$ is smaller. \square

Remark 16. Using the main result of [3] the stability of the limit cycles associated with the solution $(r^*, \theta_1^*, r_1^*, \dots, \theta_{(n/2-1)}^*, r_{(n/2-1)}^*)$ is given by the eigenvalues of the matrix

$$\frac{\partial(h_1, \dots, h_{n-1})}{\partial(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)})} \Big|_{(r, \theta_1, r_1, \dots, \theta_{(n/2-1)}, r_{(n/2-1)}) = (r^*, \theta_1^*, r_1^*, \dots, \theta_{(n/2-1)}^*, r_{(n/2-1)}^*)}$$

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