

Bragg-like curve for dark matter searches; binary gases

Akira Hitachi

Molecular Biophysics, Kochi Medical School, Nankoku, Kochi 783-8505, Japan

E-mail: jm-hitachia@kochi-u.ac.jp

Abstract

Bragg-like curves, the electronic energy deposition as a function of the projected range, for recoil ions in N₂, CO₂, CF₄ and CS₂ are presented. The curves are intended for directional search for the dark matter candidate, WIMPs, using the gas TPC. Nuclear quenching factors for very heavy ions produced in α -decay in dry air, N₂, CO₂, CF₄ and CS₂ are also estimated and compared with experimental values.

Keywords: Dark matter; Quenching; Bragg-curve; WIMPs; TPC; LET; head-tail

1. Introduction

Most of the constituent matter of the universe is dark with no emission, no absorption, even no scattering of light. It is, in fact, transparent. In 1933, Zwicky announced that there is a missing mass in the universe by observing the rotational velocity of galaxies in the Coma cluster. Further scientific evidence confirmed his observation: rotational velocity in a galaxy, gravitational lensing, the microwave radiation background, even the shape of thin galaxies require unseen mass to keep a galaxy from collapsing (IDM2006; Sumner 2002). The unseen dark matter accounts for a quarter of the universe. Ordinary matter makes up only 4%. The rest, $\sim 73\%$, is supposed to be dark energy.

In recent years, researchers have gone underground to look for WIMPs (weakly interacting massive particles), the leading candidates for dark matter. They are searching for recoil nuclei of a few tens of keV energy, produced by elastic scattering with WIMPs. However, almost all the signals produced in an underground detector are from the background, mostly γ -rays. It is essential to distinguish the WIMP signal from the background. Here, the radiation physicist and chemist can help to predict the footprints of WIMPs and the like.

One of the methods to identify WIMPs would be the directionality of the recoil ions (Spergel, 1988; CYGNUS2007). The solar system is traveling around the galactic center at about 230 km/sec. The Earth's orbit around the Sun is inclined at an angle of 60° to the galactic plane and the Earth's spin axis is 23.5° away from its orbit. Then, the direction of a "WIMP window" (or the flux of WIMPs) would change annually or even daily. Most of the detectors can observe the annual modulation. However, only detectors with directional capabilities, such as the TPC (time projection chamber), can observe daily fluctuations. Many

new technologies for constructing micro-size-ion tracks have been developed. Typical gases used for these TPCs are electronegative low diffusion gases such as CS₂ (Snowden-Ifft et al., 2000). Another gas often used is CF₄ in which electrons are the charge carriers. CF₄ is chosen because fluorine is considered to be very effective for WIMP search via spin-dependent interaction (Ellis and Flores, 1991).

The nuclear stopping S_n is of the same order of magnitude as that of the electronic stopping S_e for the interaction of ions of a few keV energy with matter (Lindhard et al., 1963). The secondary ions will go into the collisional processes again, and so on. After cascade processes of stopping collisions, considerable amounts of energy go into atomic motion which is wasted as heat in ordinary detectors. Only a part of the energy η goes to the electronic excitation which can be observed as ionization or scintillation signals. The nuclear quenching factor, or the Lindhard factor, is expressed as $q_{nc} = \eta/E$. The quenching factors are defined as the energy ratio that escapes quenching, i.e. $q_{nc} = 1$ for no quenching and $q_{nc} = 0$ for total quenching.

The linear energy transfer (LET) is a major factor in understanding radiation effects. The LET is simply $-dE/dx$ for fast ions, since the total energy loss S_T is due almost exclusively to electronic stopping; $S_T \approx S_e$. However, it is necessary to introduce the electronic LET ($LET_{el} = -d\eta/dx$) for very slow ions since $S_T = S_n + S_e$ (Hitachi, 2005, 2007). In other words, LET_{el} is the specific electronic energy loss along the track of charged particle. A particle of energy E_0 can deposit an electronic energy η_0 in a range R_0 . The particle energy becomes E_1 after an energy loss of ΔE , then it can deposit η_1 in R_1 . LET_{el} can be defined as the electronic energy deposited per unit length travelled by the particle.

$$LET_{el} = -d(qE)/dR = -d\eta/dR \approx -\Delta\eta/\Delta R = -(\eta_1 - \eta_0)/(R_1 - R_0) \quad (1)$$

LET_{el} for slow ions can be compared with LET for fast ions. Then, the knowledge accumulated over many years in radiation physics and chemistry can be used to obtain the scintillation and ionization properties in condensed media such as recoil Xe ions in liquid Xe (Hitachi, 2005).

The evaluation of the range-ionization curve is required for the directional detection of recoil ions in a gaseous TPC. The Bragg curve, the specific energy loss, $-dE/dx$, along the track of a charged particle, gives a range-ionization curve for the fast ions such as α -particles. For fast ions, the track is almost straight and the w -value, the average energy required to produce an ion pair, can be regarded as constant. However, for slow recoil ions, the total energy deposited is different from the electronic energy deposited as discussed above, and the trajectories are tortuous and ramify. Similar to the LET_{el} , we introduce the Bragg-like curve, $-d\eta/dR_{PRJ}$, which is closely related to the range-ionization curve, using the projected range (the distance of penetration), R_{PRJ} in Eq. (1) (Hitachi 2007). The curve can give an averaged,

one dimensional ionization distribution of the recoil ion track from head to tail. The Bragg-like curve was introduced for practical purposes in gas TPC. The values of S_T , S_n , S_e , and LET_{el} , together with the Bragg-like curve, for N ions in nitrogen are compared as a function of energy in Fig. 1. LET_{el} is larger than S_e but smaller than S_T . The values for S_T , S_e , and LET_{el} become the same for fast ions.

2. Recoil ions in binary gases

2.1. The nuclear quenching factor

Lindhard et al. (1963) solved the homogeneous integral equation for $v (= \varepsilon - \eta)$ in a single-element material and gave numerical results for the nuclear quenching factors $q_{nc} = \eta/E$ for $Z_1 = Z_2$ and for $k = 0.1, 0.15$ and 0.2 . Here Z_1 and Z_2 are the atomic numbers for the projectile and the target, and k is a parameter associated with the electronic stopping power, $(d\varepsilon/d\rho)_e = k \cdot \varepsilon^{1/2}$ (ε and ρ are the dimensionless reduced energy and range, respectively). The model has explained the recoil ion to γ , RN/γ , ratios in semiconductors (Gerber 1990) and the ionization measurement for Ar ions in Ar gas (Madsen 1945; Hitachi 2007).

On the other hand, the formulation of the general solutions for q_{nc} becomes quite complicated where the medium contains more than one element and further approximations are necessary in binary gases. The power law approximation, which will be discussed later, was introduced by Lindhard et al. for very heavy ions, but it does not apply to light elements. We will attempt two methods for recoil ions in binary gases: (a) the asymptotic form with $Z_1 \neq Z_2$ and (b) an independent element approach with $Z_1 = Z_2$.

An asymptotic form was introduced by Lindhard et al. (1963) for $Z_1 = Z_2$ with $0.1 < k < 0.2$. It reads

$$v = \frac{\varepsilon}{1 + k \cdot g(\varepsilon)} \quad (2)$$

The function $g(\varepsilon)$ was originally given by a figure and was recently fitted by Lewin and Smith (1996). The asymptotic form reproduces the numerical v within an accuracy of several %. The form was tried here for $Z_1 \neq Z_2$ in binary gases. Also, in the approximation, the target molecule was replaced by an element which has Z and A values close to averaged Z and A of the molecule. Namely, CS_2 molecules were replaced by Al and q_{nc} values for C ions in Al and S ions in Al were obtained by using Eq. (2). Since the average for Z and A for CF_4 are quite close to those of F, the target molecule CF_4 was simply replaced by F and q_{nc} values for C ions and F and F ions in C were calculated. This has also been applied to CO_2 .

The independent element approach was also taken to obtain q_{nc} in binary gases. CX_m may be treated by taking C ions in C and X ions in X, instead of C ions in CX_m and X ions in

CX_m , respectively. Then, the Lindhard factors q_{nc} are calculated using the numerical values for $k = 0.15$, or Eq. (2) when k is not close to 0.10, 0.15 or 0.20.

2.2. The Bragg-like curve

The electron is taken as the charge signal because of its high velocity in most gas detectors. However, its high mobility may be a disadvantage because of its large diffusion in the position determination in TPC for very short ion tracks. Alternatively, electronegative molecules may be used in a TPC for WIMP searches. An ionizing particle enters TPCs filled with electronegative low diffusion gases, such as CS_2 , produces many electrons, ions and fragments. The ejected electron immediately attaches to a surrounding neutral molecule and produces a negative ion. The negative ion will drift towards a two-dimensional-position-sensitive electrode (x-y). The negative ion releases the electron under a high electric field near the anode. The freed electron undergoes the usual electron multiplication process. The time difference between the incidence of the projectile and the arrival times of negative ions give the third dimension (z). The spatial distribution of the charge produced by the projected ion can be reconstructed to give the directionality of WIMPs in a TPC. We calculate here the Bragg-like curve, $d\eta/dR_{PRJ}$, for a one dimensional ionization distribution as a function of projectile direction R_{PRJ} .

The projected range, R_{PRJ} , was obtained from SRIM2006. The stopping power of the compound is obtained using the Bragg rule, in which the stopping power of a compound is simply given by the linear combination of the stopping powers of the individual elements. However, the outer shell electrons have different orbitals in the compound than in the corresponding elements and the core and bond approximation is often taken (Ziegler and Manoyan, 1988). The compound correction may be needed for molecules containing only the light elements, such as H, C, N, F, O, in CO_2 , CF_4 and N_2 . However no correction was applied here because it is not accurate for $Z_1 > 3$ (SRIM2006). The correction for CS_2 is not needed since CS_2 contains a heavy atom, S.

3. Heavy recoil ions in α -decay

In α -decay, the daughter nucleus recoils as a very heavy ion with typically 100-200 keV. The very heavy recoil ions in α -decay produce WIMP-like signals in detector media, and their contribution to the background signal can be very serious. For example, ^{222}Ra decays to produce Po (or Pb) ions and α recoils. Usually, one does not see these as separate particles, because they are produced at the same time. The Po signal is associated with the much larger α signal. However, in some systems Po decays such that an α particle is stopped in the wire or wall and the Pb recoil fully goes into the gas. Then, the recoil Pb ion can produce a

WIMP-like signal. It is important to know what signal will be produced.

We use a power law approximation for $Z_1 \neq Z_2$ at very low energy (Lindhard et al., 1963),

$$\eta = CE^{3/2}, \quad \text{for } E < E_{1c}, E_{2c} \quad (3)$$

where $C = \frac{2}{3}\{E_{1c}^{-1/2} + \frac{1}{2}\gamma^{1/2}E_c^{-1/2}\}$, $\gamma = 4A_1 A_2 / (A_1 + A_2)^2$ and $E_c = \gamma E_{2c}$. Two characteristic energies, E_{1c} and E_{2c} associated with Z and A for the target atom and the projectile atom, set the upper boundary. The energy E_{2c} sets the lower criterion when $Z_1 > Z_2$ as in the cases here.

The values of q_{nc} for Pb ions in α -decays in compounds are obtained using the power law approximation additive property; q_{nc} for a compound is given by the linear combination of q_{nc} of individual elements, e.g., $q_{nc}(\text{CS}_2) = [q_{nc}(\text{C}) + 2 q_{nc}(\text{S})]/3$. Thereby the values for H_2 , CF_4 , CO_2 and dry air ($4\text{N} + \text{O}$) are also calculated.

4. Result and discussion

Stopping powers (Biersack et al., 1975), electronic LET and Bragg-like curves for N ions in nitrogen are shown as a function of the ion energy in Fig. 1. The Bragg peak is at about 6 MeV and far right, outside the figure. The $\text{LET}_{el} = -d\eta/dR_T$, where R_T is the true range, is larger than the electronic stopping power S_e , particularly in the low energy region where the contribution from the nuclear stopping power S_n becomes comparable to or larger than S_e . The Bragg-like curve, $-d\eta/dR_{PRJ}$, is larger than LET_{el} , because the ion track at low energy is tortuous and have some branches, R_{PRJ} becomes shorter than R_T . At high energy, the contribution from the nuclear stopping becomes negligible and the curves for S_T , S_e , LET_{el} , and $d\eta/dR_{PRJ}$ become practically the same for fast ions such as protons, alphas and even for fission fragments. The Bragg-like curve for N ion in N_2 is also plotted as a function of R_{PRJ} in Fig. 2. The N ions enter from the right hand side. The energy η shown here is only the electronic energy that can be used for scintillation or ionization. The energies of ions recoiled by WIMPs are not monochromatic but have a distribution. The projectile energy starts at 200 keV in the figure, however the same curve applies for any energy of 10 keV to 200 keV. The area below the curve (with the unit shown on the axis on the right) expresses the number of ions produced, $N_i = \eta/W$, with a W -value of 36.6 eV for N_2 .

4.1. Recoil nuclei

The q_{nc} values obtained for the recoil ions in CO_2 , CF_4 and CS_2 are shown in Figs. 3-5. The independent element approach ($Z_1 = Z_2$, solid curves) and the asymptotic form ($Z_1 \neq Z_2$, dot-dashed curves) give practically the same values at high energy for all the gases. The difference between the C ion and the other ions (O, F or S) was wider in the asymptotic form than in the independent element approach. The difference is larger at low energies.

The electronic to total stopping power ratio, S_e/S_T are similar for C ions in C and C ions in S. This is also true for S ions in S and S ions in C. Therefore the simple estimate taken here for q_{nc} may be valid in CS₂. The q_{nc} values estimated for recoil C and S ions in CS₂ are shown in Fig. 3. The values calculated for S ions in Al and in S by the asymptotic form are quite close; 0.28 and 0.29 at 10 keV and 0.57 and 0.57 at 200 keV, respectively. Therefore, the results for S ions in C is shown in Fig. 3. Similarly, the results for C ions in Al and in S are very close, therefore the results for C ions in S is shown for the asymptotic form in Fig. 3, as also the results shown for other gases in Figs. 4-5. The q_{nc} values calculated for CS₂ are compared with the experimental work of Snowden-Ifft et al. (2003) which are not pure experimental values but in part simulation. The independent element approach (S ions in S) agrees quite well with the measured values for S ions in CS₂, however, the values calculated for C ions (C in C) give higher values at > 100 keV. The asymptotic form for S ions in C agrees well at high energy but gives a smaller value at low energy. The asymptotic form for C ions in S agrees quite well at low energy but it gives larger values at high energy. The reason why the experimental values for C ions are lower at high energy is not well known.

The values obtained by the power law approximation, C in Al and S in Al, are also shown with broken curves for comparison. The energy criteria set by E_{1c} or E_{2c} are 24 keV and 177 keV, respectively, for C and S ions in Al. The power law approximation gives quite steep curves and the values are large except at very low energy. The energy criterion may be lower than E_{1c} or E_{2c} for light ions.

Fig. 4 shows the results for CF₄. No experimental values for q_{nc} are available for CF₄. The asymptotic form ($Z_1 \neq Z_2$) gives larger values for C ion and smaller values for F ions than the independent approach. However, the differences are small, and both approximations give practically the same values for C and F ions in CF₄ except at a very low energy. This may be due to the Z and A values for C and F atoms being not very different. Similar results were obtained for C and O recoil ions in CO₂ as shown in Fig. 5. The q_{nc} values for N ions in N₂ are also shown in Fig. 5. The values lay between the curves for C and O ions as one might expect.

The Bragg-like curves calculated for recoil ions from 10 to 200 keV in CO₂, CF₄ and CS₂ are shown in Figs. 6-8. C ions are shown with closed circles and the partner ions are closed squares. The value of q_{nc} obtained by the independent element approach was taken. The points are plotted in every 5 keV for 10 to 50 keV, in 10 keV for 50 to 100 keV, and in 20 keV for 100 to 200 keV. W -values of 33 eV, 54 eV and 19 eV were used for CO₂, CF₄ and CS₂, respectively, to obtain ΔN_i . The area below each curve (read with the axis on the right) expresses the number of ions produced, $N_i = \eta/W$.

The Bragg-like curves for C and S ions in CS₂ show strong dependences on energy (Fig. 6) and look suitable for the directional detection of the WIMPs. However, the curves for C and F ions in CF₄ do not have strong energy dependences except below 50 keV. The curve for F ions, which is expected to have a favourable spin dependent interaction, is nearly flat above

50 keV. The number of ions produced in CF₄ is also much smaller than that for CS₂ because of the large W -value (54 eV). The curves for C and O ions in CO₂ are between those for CS₂ and CF₄.

The incident ion can produce high energy recoils of other elements which can ionize the target molecules. For example, the incident C ion may produce S ions in CS₂, and the S ions can contribute to the ionization. This effect may not be included in the independent element approach. It is assumed that the ionization is due to the projectile element in the present calculations for q_{nc} for recoil atoms. It is not clear how much of such effect is included in the asymptotic form. Ling and Knipp (1950) presented a model for Pb ions in Ar similar to the power law approximation by Lindhard. They estimated that about 2/3 of the ionization is due to the secondary ions. If the contribution from the other element ion is large, the curves for C ions and S ions in CS₂ in Fig. 3 may shift nearer to each other. This may explain the behaviour of the experimental q_{nc} curve for C ions. However, it is still strange that the deviation occurs at high energy. The contribution from scattering is smaller at high energy.

4.2. Heavy recoil ions in α -decay

The α -decay from ²¹⁰Po, ThC (²¹²Bi), and ThC' (²¹²Po) produces recoil ions of 103 keV ²⁰⁶Pb, 117 keV ²⁰⁸Tl and 168 keV ²⁰⁸Pb, respectively. Also, ²¹²Rn daughters, ²¹⁸Po gives 112 keV ²¹⁴Pb, and ²¹⁴Po gives 147 keV ²¹⁰Pb. Nuclear quenching factors q_{nc} for heavy recoil ions in α -decay of Po, ThC and ThC' in dry air, N₂, CO₂, CF₄ and CS₂ are obtained by the power law approximation and are shown in Table I. The ionization measurements in a gas basically give the Lindhard factor q_{nc} both for the gaseous and condensed phases: $q_{nc} = W(\alpha)/W(RN)$, where $W(\alpha)$ and $W(RN)$ are the W -values for fast ions and slow recoil ions, respectively. Calculated q_{nc} values are compared with $W(\alpha)/W(RN)$ values measured in gas phase. The values calculated for N₂, CO₂, CF₄ and CS₂ agree well with the experimental values (Stone and Cochran, 1957, Cano, 1968). The general trend is that q_{nc} decreases as the target Z_2 increases.

The power law approximation does not apply for H₂ because E_{2c} becomes too small but it is shown just for comparison. The calculated q_{nc} values are much larger than the measured values. Hydrocarbons have a strange behaviour. The measured values of q_{nc} are 0.457 and 0.543 respectively for 117 keV and 168 keV in H₂ and these are still much higher than those for C, O, F and S. If the general trends with Z_2 apply also for the hydrocarbons, q_{nc} values for CH₄, C₂H₄ and C₃H₆ should be higher than those for CO₂, N₂ and dry air. However, the reverse is observed.

The values obtained by the power law approximation for heavy recoil ions, Pb, in α decay (dashed curve) are also shown for *ca* 100 keV to 170 keV in Figs. 3-5 with dashed

curves. The thin dot-dot dash curves, e.g., Pb ions in C and Pb ions in S, and Pb ions in Al, are shown in Fig. 3.

The Bragg-like curves for Pb ions in N₂, CO₂, CF₄ and CS₂ are obtained using the q_{nc} values given by the same method as recoil ions and are shown in Fig. 9. The ion enters from the right. They give roughly straight lines. This may be because the same energy dependence, $q_{nc} \propto E^{1/2}$, was assumed in the power law approximation. The ionization densities for very heavy recoil ions in α -decay is much higher than those for recoil ions and the range is much shorter. The total amount of the ionization produced by heavy recoil ions, shown as the area below the Bragg-like curve, is in the same order of magnitude as those for recoil ions. Therefore, the very heavy recoil ions in α -decay may give WIMP-like signals in detectors without track shape, or length, discrimination capabilities.

4.3 General remarks

Compounds containing only light elements, hydrocarbons and CF₄, may require careful treatment, since the compound corrections for the stopping power of these compounds are required. The Bragg rule introduces errors of several percent in the stopping power at low energies because the electronic wavefunctions of the target depends on the bonds of each atom with its neighbours. This effect is most important for the low- Z atoms of biological interests (Hobbie, 1987). The W -values may also have considerable energy dependence for these compounds (Combecher, 1980). In addition to this, the triplet states in these molecules are metastable which are not excited by fast particles but can be excited effectively by slow heavy ions. These facts may affect the stopping powers, q_{nc} , and W -values.

The validity of the asymptotic form for $Z_1 \neq Z_2$ with $0.1 < k < 0.2$ is still in question. It looks as if the form is still good for recoil ions in binary gases as shown for CO₂, CF₄ and CS₂. The values for k are within 0.10 to 0.20 for recoil ions in those gases except for CS₂. However, when the projectile becomes very heavy the difference in Z_1 and Z_2 becomes large, then the asymptotic form seems to fail. The values obtained for heavy recoil ions in α -decay are small in most gases. The k values are 0.10 to 0.16 and are generally in the range of 0.1 - 0.2. The q_{nc} values calculated using the asymptotic form for 103 keV Pb ions are 0.20 - 0.21 in He, C, N, O, and S, and those for 168 keV Pb ions are 0.22 - 0.23 and are practically the same, whereas reported experimental values are scattered as shown in Table I. The discrepancy may be due to the fact that the asymptotic form does not include the effects of the secondary ions.

A cylindrical picture of the ion track may be obtained by assuming a Gaussian in the lateral direction and taking lateral straggling (ΔR_{\perp}). It is hard to take the longitudinal straggling (ΔR_{\parallel}) into account but it is possible. Such a two-dimensional picture with cylindrical geometry represents an ensemble average of many ions. However, the shape of each ion track changes because of the scattering. Additionally the amount of ionization

produced is not large. The individual track may stay within the bell-shape contour (with ΔR_{\perp}) or the drop-shape contour (with ΔR_{\perp} and ΔR_{\parallel}). However, the ionization may be distorted and localized. The effect of the straggling on each track is not well known.

Evans et al. (1963) observed ionization density contours for He, N, Ne, and Ar ions in He, N₂, air and Ar gases in the range from *ca* 20 keV to 250 keV. They reported that the attenuation of the beam of ions is found to be approximately exponential in the axial direction and approximately Gaussian in the lateral direction. The Bragg-like curve obtained for N ions in nitrogen compare well with the measured range-ionization curve for 57 keV N ions in air except at the very beginning. However, the range-ionization curve shows less ionization near the end of the track. This may be due to the fact that the present calculation does not take longitudinal straggling into account. The values for 57 keV N ions in N₂ are $R_{PRJ} = 3.25$, $\Delta R_{\parallel} = 0.89$, and $\Delta R_{\perp} = 0.75$, $\times 10^{-3}$ mg/cm² (SRIM). The charge distribution spreads in the longitudinal as well as in the lateral direction. The longitudinal spread reduces the charge density particularly at the end in the range-ionization curve. This is encouraging for the directional detection of WIMP foot prints. The directionality may be clearer in actual range-ionization curve than the Bragg-like curve without the straggling as presented here.

The Bragg-like curve becomes steeper for heavy elements. It may be useful to use gas mixtures such as Ar + M or Xe + M, where M is a electro-negative molecule, in TPC for dark matter searches as far as the directionality is concerned. However, at very low energy, the scattering becomes large for heavy atoms.

5. Summary

The Bragg-like curves, which are closely related to the range-ionization curve, for recoil ions in N₂, CO₂, CF₄ and CS₂ were obtained for the directional detection of recoil ions in gas TPC. The present results show that the information on the change in the nuclear quenching factor, q_{nc} , as a function of the energy is not enough to consider directional capabilities of dark matter detectors. One needs the Bragg-like curve which includes changes in q_{nc} and the stopping powers. The power law approximation for $Z_1 \neq Z_2$ gave satisfactory values of q_{nc} for the heavy recoil ions in α -decay except in the hydrocarbons.

Acknowledgements

The author would like to thank Dr. A. Mozumder and Dr. J. A. LaVerne for helpful discussions. He also thanks Dr. P. Majewski and Dr. K. Miuchi for useful information. He is also grateful to Prof. T.A. King for reading of the manuscript.

References

- Biersack, J.P, Ernst, E., Monge, A., Roth, S., 1975. Tables of Electronic and Nuclear Stopping Powers and Energy Stragglings for Low-Energy Ions, Hahn-Meitner Institut Publication No. HMI-B 175.
- Cano, G.L., 1968. Total ionization and range of low-energy recoil particles in pure and binary gases. *Phys. Rev.* 169, 227-280.
- Combecher, D., 1980. Measurement of W values of low-energy electrons in several gases. *Radiat. Res.* 84, 189-218.
- CYGNUS2007. <http://www.pppa.group.shef.ac.uk/cygnus2007//talks/>
- Ellis, J. , Flores, R.A. 1991. Elastic supersymmetric relic-nucleus scattering revisited. *Phys. Lett. B* 263, 259-266.
- Evans, G.E., Stier, P.M., Barnett C.F., 1953. The stopping of heavy ions in gases. *Phys. Rev.* 90, 825-833.
- Gerbier, G. et al., 1990. Measurement of the ionization of slow silicon nuclei in silicon for the calibration of a silicon dark-matter detector. *Phys. Rev. D.* 42, 3211-14.
- Hitachi, A., 2005. Properties of liquid xenon scintillation for dark matter searches. *Astroparticle Physics*, 24, 247-256.
- Hitachi, A., 2007. Quenching factor and electronic LET in gas at low energy. *Journal of Physics: Conference Series* 65, 012013-1~6.
- Hobbie, R.K., 1987. *Intermediate Physics for Medicine and Biology*, 2nd edition, John Wiley & Sons, New York.
- IDM2006, 2007. The identification of dark matter, ed., Axenides, M., Fanourakis, G., Vergados, J., World Scientific.
- Lewin, J.D., Smith, P.F., 1996. Review of mathematics, numerical factors, and corrections for dark matter experiments based on elastic nuclear recoil. *Astroparticle Phys.* 6, 87-112.
- Lindhard. J., Nielsen, V., Sharff, M., Thomsen P. V., 1963. Integral equations governing radiation effects. *Mat. Fys. Medd. Dan. Vid. Selsk.* 33(10), 1-42.
- Ling R.C., Knipp, J.K., 1950. Ionization by recoil particles from alpha-decay. *Phys. Rev.* 80, 106.
- Madsen, B. S., 1945. Ionization measurements on single recoil particles from Po, ThC, and ThC'. *Mat. Fys. Medd. Dan. Vid. Selsk.* 23(8), 1-16.
- Sharma, A., 1998. Properties of some gas mixtures used in tracking detectors. *SLAC-Journal-JCFA* 16, 3-21.
- Snowden-Ifft. D.P., Ohnuki, T., Rykoff, E.S., Martoff, C.J., 2003. Neutron recoils in the DRIFT detector. *Nucl. Instr. Meth. A* 498, 155-164.
- Snowden-Ifft. D.P., Martoff, C.J., Burwell, J.M., 2000. Low pressure negative ion time projection chamber for dark matter search. *Phys. Rev. D* 61, 101301-1~5.

- Stone, W.G., Cochran, L.W., 1957. Ionization of gases by recoil atoms. *Phys. Rev.* 107, 702-4.
- Spergel, D.N., 1988. Motion of the Earth and the detection of weakly interacting massive particles. *Phys. Rev. D* 37, 1353-5.
- SRIM2006, <http://www.srim.org/>
- Sumner T.J., 2002. Experimental searches for dark matter.
<http://www.livingreviews.org/Articles/Volume5/2002-4sumner/>
- Ziegler, J.F., Manoyan, J.M., 1988. The stopping of ions in compounds. *Nucl. Instr. Methods*, B35, 215-228.
- Zwicky, F., 1933. Die rotverschiebung von extragalaktischen nebeln. *Helv. Phys. Acta.* 6, 110-127.

Table I

Nuclear quenching factors q_{nc} for recoil ions in α -decay of Po, ThC and ThC' in various gases. The experimental values are obtained assuming $q_{nc} = W(\alpha)/W(RN)$. The calculated values were obtained by the power law approximation by Lindhard et al. (1963) which is valid for $E < E_{2c}$.

Source	^{210}Po	ThC (^{212}Bi)		ThC' (^{212}Po)		E_{2c}	
Recoil ion	^{206}Pb	^{208}Tl	^{208}Pb	^{208}Pb			
Energy keV	103	117	168				
gas	expt	calc	expt	calc	expt	calc	keV
H ₂		(0.73)	0.457 ^a	(0.78)	0.543 ^a	(0.93)	26
CH ₄	0.250 ^b		0.265 ^a		0.307 ^a		
C ₂ H ₄	0.236 ^b		0.269 ^a		0.321 ^a		
C ₃ H ₆			0.272 ^a		0.281 ^a		
CO ₂			0.336 ^a		0.347 ^a		
C + 2O		0.297		0.316		0.378	174
N ₂	0.319 ^b	0.302		0.320		0.384	207
Dry air	0.296 ^b						
4N + O		0.298		0.317		0.379	207
CF ₄							
C + 4F		0.280		0.297		0.356	174
CS ₂							
C + 2S		0.246		0.262		0.314	174
Al (for CS ₂)		0.228		0.242		0.290	428
C		0.323		0.344		0.411	174
O		0.284		0.302		0.361	241
F		0.269		0.286		0.342	278
S		0.208		0.221		0.264	550
Source	^{218}Po	^{214}Po		^{214}Po		E_{2c}	
Recoil ion	^{214}Pb	^{214}Pb	^{210}Pb	^{210}Pb			
Energy keV	112	147	147				
gas	expt	calc	expt	calc	expt	calc	keV
CS ₂							
C + 2S		0.253		0.292			174
Al (for CS ₂)		0.237		0.272			430

^a Stone and Cochran (1957).

^b Cano (1968).

Figure Captions

- Fig. 1. The stopping powers, S_T , S_n , S_e , and LET_{el} ($= -d\eta/dx$) for N ions in nitrogen as functions of the energy. The Bragg-like curve ($-d\eta/dR_{PRJ}$) is also shown.
- Fig. 2. The Bragg-like curve for N ions in N_2 . The N ions enter from the right hand side.
- Fig. 3. The nuclear quenching factor, q_{nc} , calculated for C and S ions in CS_2 as a function of energy. The circles are for the measurement and part simulation by Snowden-Ifft (2003). The solid lines are independent element approach and dot-dashed curves are the asymptotic form. The results obtained for C, S and Pb ions using the power law approximation are also shown.
- Fig. 4. The values of q_{nc} for C and F ions in CF_4 as a function of energy. The results for Pb ions are also shown for 103-168 keV.
- Fig. 5. The values of q_{nc} for C and O recoil ions in CO_2 as a function of energy. The results for Pb ions are also shown. The q_{nc} values for N ions in N_2 are also shown (dot-dash curve).
- Fig. 6. The Bragg-like curve, which is closely related to specific ionization, plotted as a function of the projected range $d\eta/dR_{PRJ}$, estimated for 10 to 200 keV recoil C and S ions in CS_2 . The points are plotted in every 5 keV interval for 10 to 50 keV, in 10 keV interval for 50 to 100 keV, and in 20 keV interval for 100 to 200 keV. The area below each curve (with the unit shown on the axis on the right) expresses the number of ions produced, N_i .
- Fig. 7. The Bragg-like curve estimated for recoil C and F ions in CF_4 . The area below each curve expresses the number of ions produced, N_i .
- Fig. 8. The Bragg-like curve estimated for recoil C and O ions in CO_2 .
- Fig. 9. The Bragg-like curve estimated for recoil Pb ions in α -decay in N_2 , CO_2 , CF_4 and CS_2 .

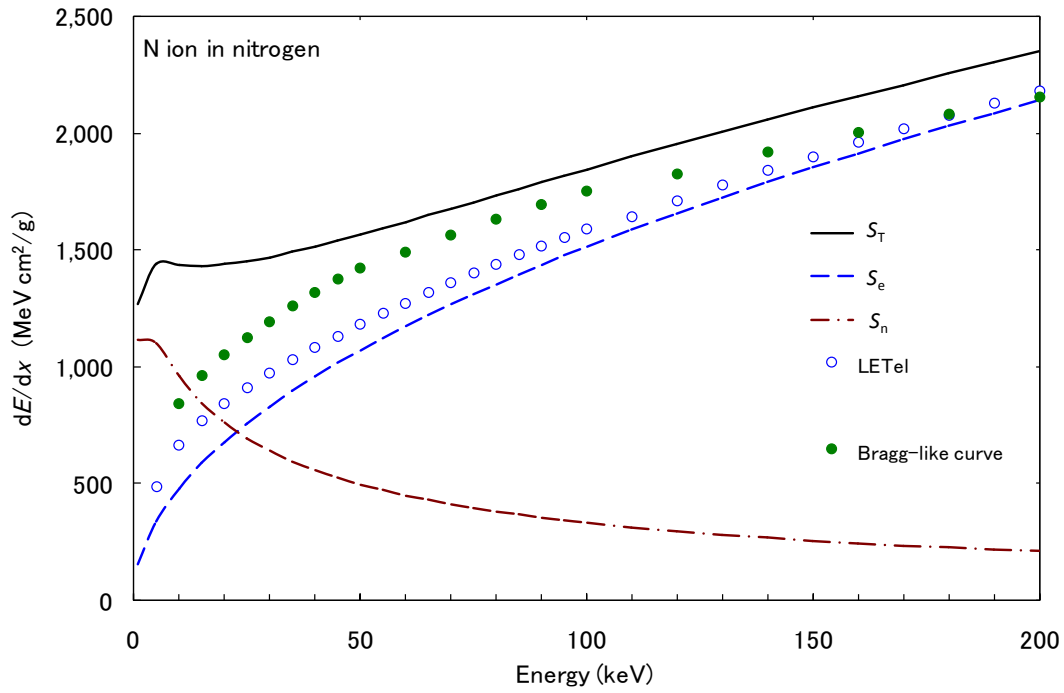


Fig. 1. The stopping powers, S_T , S_n , S_e , and LET_{el} ($= -d\eta/dx$) for N ions in nitrogen as functions of the energy. The Bragg-like curve ($-d\eta/dR_{PRJ}$) is also shown.

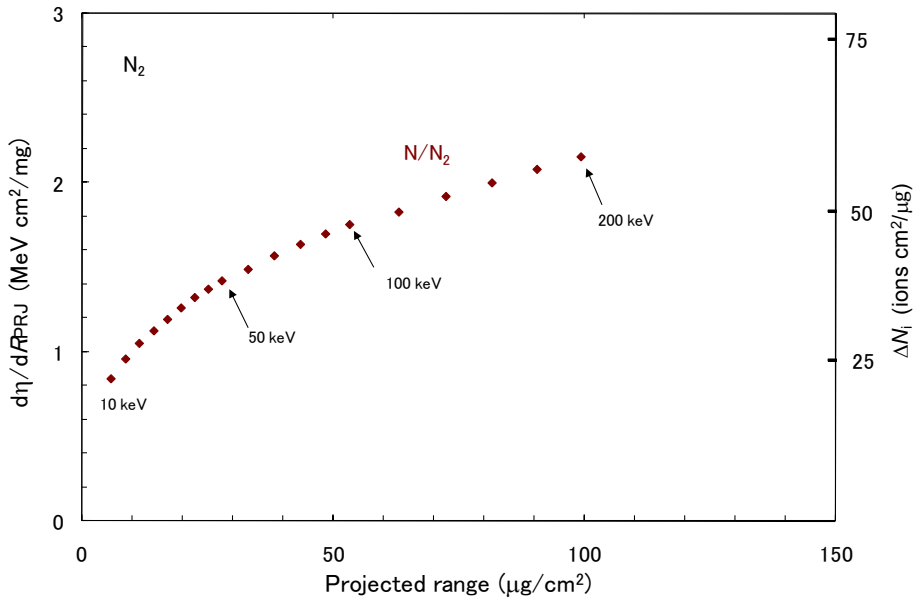


Fig. 2. The Bragg-like curve for N ions in N₂. The N ions enter from the right hand side.

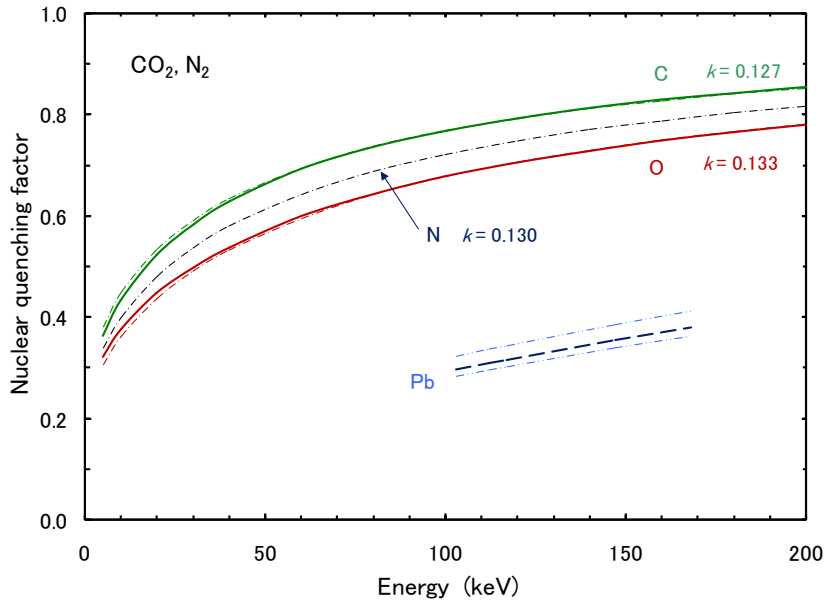


Fig. 5. The values of q_{nc} for C and O recoil ions in CO_2 as a function of energy. The results for Pb ions are also shown. The q_{nc} values for N ions in N_2 are also shown (dot-dash curve).

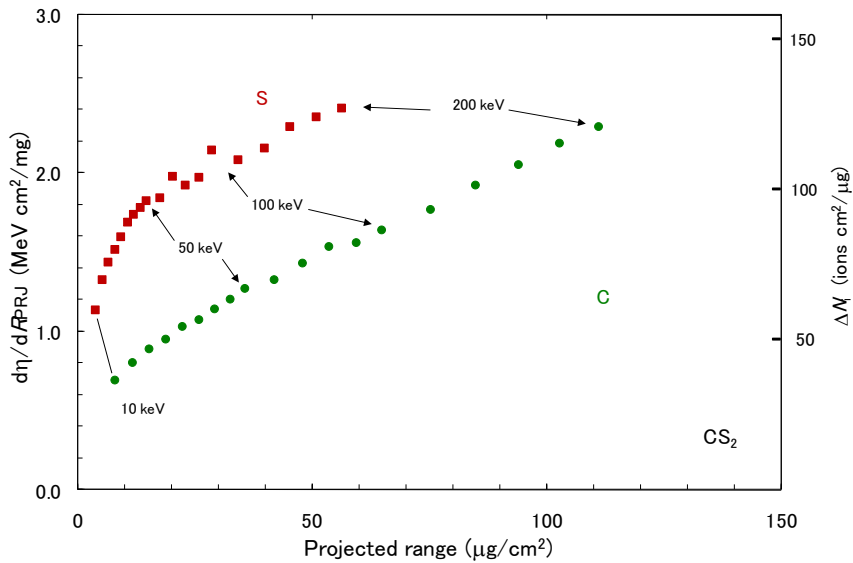


Fig. 6. The Bragg-like curve, which is closely related to specific ionization, plotted as a function of the projected range $d\eta/dR_{PRJ}$, estimated for 10 to 200 keV recoil C and S ions in CS_2 . The points are plotted in every 5 keV interval for 10 to 50 keV, in 10 keV interval for 50 to 100 keV, and in 20 keV interval for 100 to 200 keV. The area below each curve (with the unit shown on the axis on the right) expresses the number of ions produced, N_i .

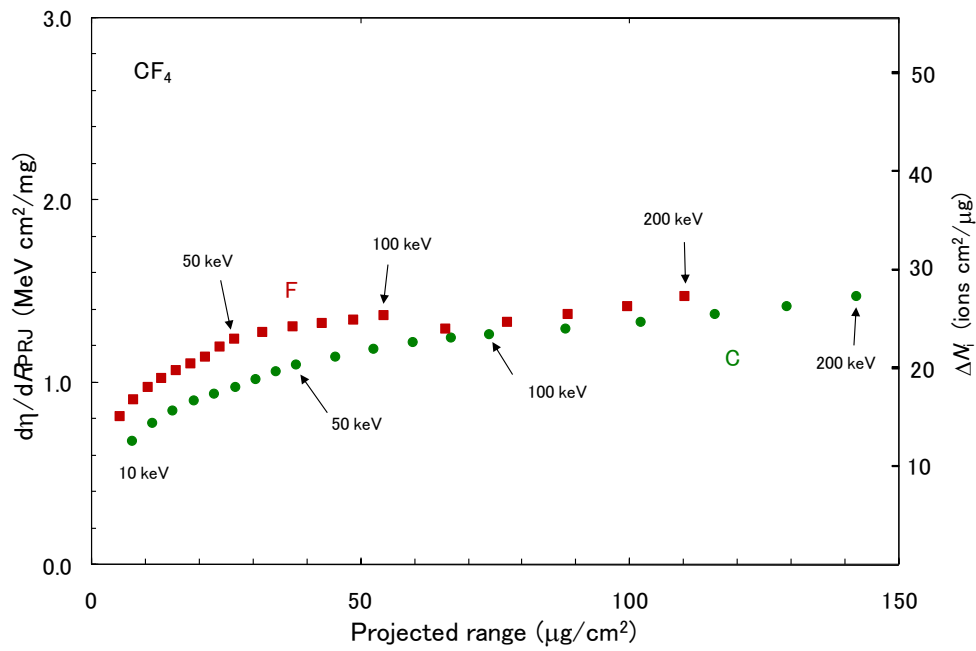


Fig. 7. The Bragg-like curve estimated for recoil C and F ions in CF_4 . The area below each curve expresses the number of ions produced, N_i .

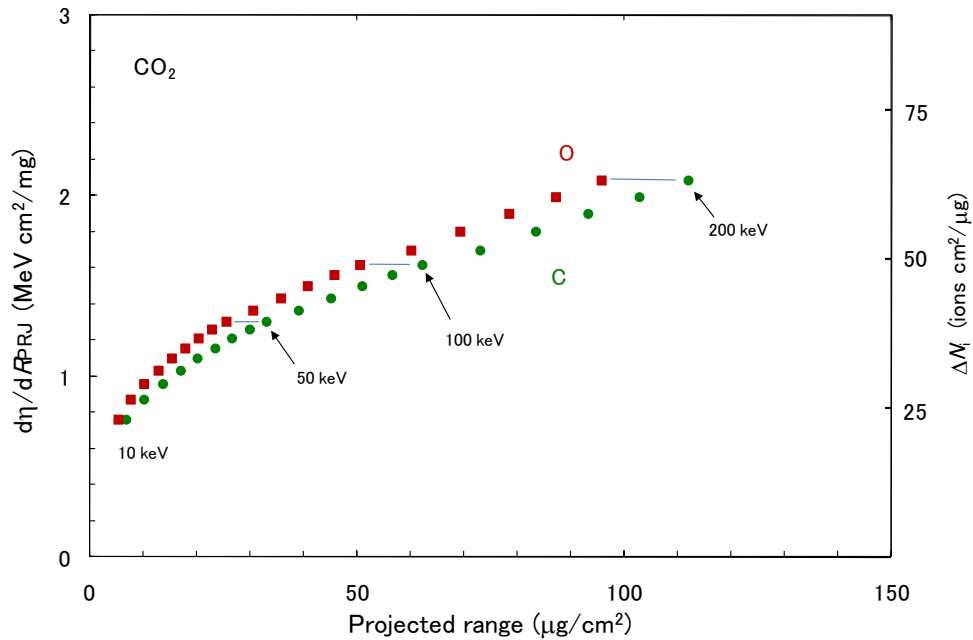


Fig. 8. The Bragg-like curve estimated for recoil C and O ions in CO_2 .

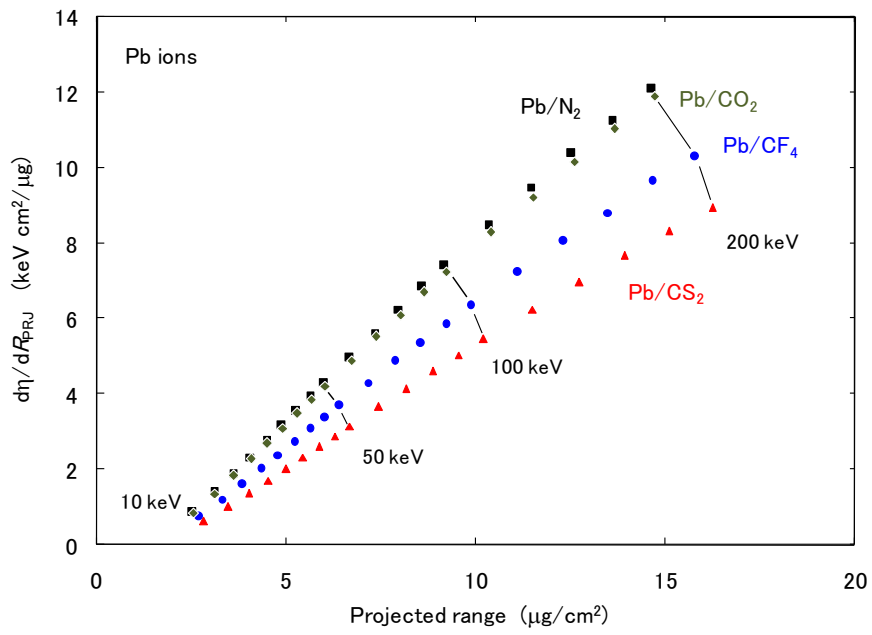


Fig. 9. The Bragg-like curve estimated for recoil Pb ions in α -decay in N₂, CO₂, CF₄ and CS₂.