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Models

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**DISCUSSION  
PAPER**

# Asymptotic Results for GMM Estimators of Stochastic Volatility Models

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## Abstract

We derive closed-form expressions for the optimal weighting matrix for GMM estimation of the stochastic volatility model with AR(1) log-volatility, and for the asymptotic covariance matrix of the resulting estimator. The moment conditions considered are generated by the absolute observations (which is the standard approach in this literature) or by the log-squared observations. We use the expressions to compare the performances of GMM and other estimators that have been proposed, and to optimally select small sets of moment conditions from very large sets.

JEL: C13, C22

Keywords: Stochastic volatility, GMM

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# 1 Introduction

Over the last two decades there has been an increasing interest in stochastic volatility (SV), which was introduced by Clark (1973) and extended by Tauchen and Pitts (1983), as a framework for the analysis of time-varying volatility in financial markets. This interest is partly due to an important contribution by Hull and White (1987), where SV models arise as discrete time approximations to continuous time volatility diffusions used in option pricing. More generally, it is recognized that SV models constitute a valuable alternative to GARCH-type models for analysing financial time series (Ghysels, Harvey, and Renault (1996), Shephard (1996)).

Due to the fact that in SV models the mean and the volatility are driven by separate stochastic processes (implying that volatility is unobservable), SV models are much harder to estimate than GARCH models. This paper presents analytical results that may be used to improve and assess the quality of GMM-based estimation of SV models. GMM, while not asymptotically efficient, is still the simplest estimation method for SV models currently available. It has been proposed by Taylor (1986) and Melino and Turnbull (1990), and its properties have been studied using Monte Carlo methods by Jacquier, Polson, and Rossi (1994), Andersen and Sørensen (1996, 1997), and Andersen, Chung, and Sørensen (1999). Other available estimation methods for SV models include quasi-maximum likelihood (Nelson (1988), Harvey, Ruiz, and Shephard (1994), Ruiz (1994)), simulated maximum likelihood (Danielsson and Richard (1993), Danielsson (1994)), simulation-based GMM (Duffie and Singleton (1993)), indirect inference (Gouriéroux, Monfort, and Renault (1993), Monfardini (1998)), Markov chain Monte Carlo methods (Jacquier, Polson, and Rossi (1994), Kim, Shephard, and Chib (1998), Chib, Nardari, and Shephard (2002)), efficient method of moments (Gallant, Hsieh, and Tauchen (1997), Andersen, Chung, and

Sørensen (1999)), Monte Carlo maximum likelihood (Sandmann and Koopman (1998)), and (approximate) maximum likelihood (Fridman and Harris (1998)). Apart from quasi-maximum likelihood, all of these methods are computationally more demanding, as they rely – often quite heavily – on numerical simulation and/or integration techniques both for obtaining point estimates and for assessing the accuracy of the latter. In view of its simplicity, we consider GMM estimation as a useful alternative to the more elaborate methods.

In this paper we derive closed-form expressions for the optimal weighting matrix for GMM estimation of the basic SV model, and for the asymptotic covariance matrix of the optimal GMM estimator, for a large class of moment conditions. To date, applications of GMM in this context have typically relied on a nonparametrically estimated weighting matrix, because an expression for the optimal weighting matrix (as a function of the parameters) was not available.

The moment conditions that we consider fall into two categories. The first set of conditions is obtained by considering the first two moments and the auto-covariances of any order of the log-squared observations. These conditions have recently been considered by Wright (1999), in connection with the fractionally integrated SV model. The second set of moment conditions are derived from the absolute observations and are more standard in this literature. We study moment conditions that involve the product of any number of absolute observations, each one raised to any positive integer power and lagged any number of periods. This set considerably extends the set of moment conditions that have been employed so far. The results that we present pertain to any selection of moment conditions from these two sets.

In Section 2, we present the basic SV model and the moment conditions. Expressions for the optimal weighting matrix and the asymptotic covariance matrix of the GMM estimator are derived in Section 3. Section 4 presents some comparative evidence on the relative efficiencies of the GMM and other

estimators (partly compiled from the literature). We also show how the analytical results of this paper provide a fast and accurate tool to select a small set of highly informative moment conditions from very large sets of moment conditions. Section 5 concludes. Proofs are given in the Appendix.

## 2 Moment conditions for the SV model

The basic SV model is given by

$$y_t = \exp(h_t/2) u_t, \tag{1}$$

$$h_{t+1} = \mu + \phi(h_t - \mu) + \sigma\sqrt{1 - \phi^2}v_t, \tag{2}$$

where  $y_t$  is observable,  $h_t$  is latent log-volatility,  $(u_t, v_t)$  is i.i.d.  $N(0, I)$ , and  $\theta = (\mu, \phi, \sigma)'$  is a vector of parameters. The restriction  $|\phi| < 1$  is imposed, ensuring that  $y_t$  is stationary and ergodic. While it is more common to parameterise the model in terms of  $\lambda = (\alpha, \phi, \omega)'$ , with  $\alpha = \mu(1 - \phi)$  and  $\omega = \sigma\sqrt{1 - \phi^2}$ , we prefer the parameterisation in terms of  $\theta$  for algebraic reasons and because of an invariance with respect to  $\mu$  given below. For comparison with earlier studies, however, numerical standard errors will be presented in terms of  $\lambda$ .

From the point of view of inference, the fundamental problem with the SV model is the latent character of  $h_t$ , which makes it difficult to compute the values of the likelihood function and hence to estimate  $\theta$  by maximum likelihood. It is easy, however, to derive moment conditions implied by the SV model and then to apply the Generalized Method of Moments (Hansen, 1982). The moment conditions considered in this paper relate either to the log-squared observations,  $\log y_t^2$ , or to the absolute observations,  $|y_t|$ . The latter class of moment conditions constitutes the standard approach to GMM estimation of SV models (Taylor (1986), Melino and Turnbull (1990), Jacquier, Polson, and Rossi (1994), Andersen and Sørensen (1996, 1997), Andersen, Chung, and Sørensen (1999)). The former class of moment conditions is suggested in passing by Jacquier,

Polson, and Rossi (1994), and is effectively employed by Wright (1999) in the context of the fractionally integrated SV model.

Moment conditions related to  $\log y_t^2$  are easily obtained. It follows from (1) that  $\log y_t^2 = h_t + \log u_t^2$ . The mean and variance of  $\log u_t^2$  are known to be  $c_1 = -\log 2 - \gamma = -1.2704$  and  $c_2 = \frac{1}{2}\pi^2 = 4.9348$ , respectively, where  $\gamma = 0.5772$  is Euler's constant. Let

$$\begin{aligned} z_t &= \log y_t^2 - \mu - c_1 \\ &= h_t - \mu + \log u_t^2 - c_1. \end{aligned}$$

Since  $h_t \sim N(\mu, \sigma^2)$ ,  $\text{Cov}(h_t, h_{t-i}) = \phi^{|i|}\sigma^2$ , and  $u_t$  is i.i.d. and independent of  $h_t$ , it follows that

$$E[z_t] = 0, \tag{3}$$

$$E[z_t z_{t-i}] = \phi^i \sigma^2 + I_{(i=0)} c_2, \quad i \geq 0, \tag{4}$$

where  $I_{(\cdot)}$  is the indicator function. It can be shown that none of these conditions is redundant in the sense of Breusch et al. (1999).

We now derive the class of moment conditions generated by the expectation of  $|y_{t_1}^{i_1} \dots y_{t_p}^{i_p}|$ , where  $i_1, \dots, i_p$  are positive integers and  $t_1 > \dots > t_p$ . Let  $\nu_i$  be the  $i$ -th absolute moment of a standard normal random variate, i.e.

$$\nu_i = E|u_t|^i = \frac{2^{i/2}}{\sqrt{\pi}} \Gamma\left(\frac{i+1}{2}\right) = \begin{cases} \sqrt{\frac{2^i}{\pi}} \left(\frac{i-1}{2}\right)! & \text{for } i \text{ odd,} \\ \frac{i!}{(i/2)! 2^{i/2}} & \text{for } i \text{ even,} \end{cases}$$

where  $\Gamma(z)$  is the gamma function. Then, because  $t_1 > \dots > t_p$ ,

$$E \prod_{j=1}^p \left| \nu_{i_j}^{-1} u_{t_j}^{i_j} \right| = 1.$$

Furthermore,  $\sum_{j=1}^p i_j h_{t_j}$  is normally distributed with mean  $\mu \sum_{j=1}^p i_j$  and variance  $\sigma^2 \sum_{j,j'=1}^p i_j i_{j'} \phi^{|t_j - t_{j'}|}$ . So, by property that  $E \exp(X) = \exp(a + \frac{1}{2}b^2)$

when  $X \sim N(a, b^2)$ , we have

$$E \exp \left( \frac{1}{2} \sum_{j=1}^p i_j h_{t_j} \right) = \exp \left( \delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right),$$

where

$$\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} = \frac{\mu}{2} \sum_{j=1}^p i_j + \frac{\sigma^2}{8} \sum_{j, j'=1}^p i_j i_{j'} \phi^{|t_j - t_{j'}|}.$$

Hence, defining

$$\begin{aligned} Y_{t_1, \dots, t_p}^{i_1, \dots, i_p} &= \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right) \prod_{j=1}^p \left| \nu_{i_j}^{-1} g_{t_j}^{i_j} \right| \\ &= \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right) \exp \left( \frac{1}{2} \sum_{j=1}^p i_j h_{t_j} \right) \prod_{j=1}^p \left| \nu_{i_j}^{-1} u_{t_j}^{i_j} \right|, \end{aligned}$$

it follows that

$$E \left[ Y_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right] = 1, \quad i_1, \dots, i_p \geq 1; \quad t_1 > \dots > t_p. \quad (5)$$

It is obvious that adding the same integer to  $t_1, \dots, t_p$  yields the same moment condition. As far as we know, within the class of moment conditions defined by (5), only moment conditions where  $p = 1$  or where  $p = 2$  and  $i_1 = i_2 \in \{1, 2\}$  have so far been considered in the literature.

### 3 Optimal GMM

Let  $E(f_t) = f$  be a finite selection of the set of moment conditions given by (3)–(5) that identifies  $\theta$ . Let  $g_t = f_t - f$ . By assumption, the observations on  $y_t$  permit us to calculate  $g_1, \dots, g_T$  as functions of  $\theta$ . The optimal GMM estimator (Hansen (1982)) of  $\theta$  based on this selection is  $\hat{\theta} = \arg \min_{\theta} \bar{g}' \hat{V}^{-1} \bar{g}$ , where  $\bar{g} = T^{-1} \sum_{t=1}^T g_t$  and  $\hat{V}$  consistently estimates  $V$ , where

$$V = \sum_{l=-\infty}^{\infty} E(g_t g_{t-l}') = \sum_{l=-\infty}^{\infty} \text{Cov}(f_t, f_{t-l}').$$

The asymptotic covariance matrix of  $\sqrt{T}(\hat{\theta} - \theta)$  is  $(D'V^{-1}D)^{-1}$ , where  $D = E(\frac{\partial g_t}{\partial \theta'})$ . Expressions for  $D$  and  $V$ , for an arbitrary selection of moment conditions, are presented below. These expressions make it possible to compute the optimal weighting matrix  $V^{-1}$  and the asymptotic covariance matrix  $(D'V^{-1}D)^{-1}$  of the GMM estimator as functions of the parameter values. Substituting estimates for these parameter values yields estimates of  $V^{-1}$  and  $(D'V^{-1}D)^{-1}$ , which will generally be more precise than the nonparametric estimator based on Bartlett weights that is routinely used in a GMM context. The Monte Carlo results of Andersen and Sørensen (1996) show that the latter estimator of  $V$  may be imprecise even in samples of size 50,000. Using the expressions presented here avoids such problems. Furthermore, the expression for  $V$  makes it also possible to estimate  $\theta$  by the continuous-updating GMM estimator of Hansen, Heaton, and Yaron (1996), that is, by solving  $\min_{\theta} \bar{g}'V^{-1}\bar{g}$ .

Some straightforward calculus shows that the rows of  $D$  are to be selected (according to the selection of moments) from

$$\begin{pmatrix} -1 & 0 & 0 \\ 0 & -i\phi^{i-1}\sigma^2 & -2\phi^i\sigma \\ -\nabla_{\mu}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} & -\nabla_{\phi}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} & -\nabla_{\sigma}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} \end{pmatrix} \begin{matrix} i \geq 0; \\ i_1, \dots, i_p \geq 1; \quad t_1 > \dots > t_p; \end{matrix}$$

where

$$\begin{aligned} \nabla_{\mu}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} &= \frac{1}{2} \sum_{j=1}^p i_j, \\ \nabla_{\phi}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} &= \frac{\sigma^2}{8} \sum_{j, j'=1}^p i_j i_{j'} |t_j - t_{j'}| \phi^{|t_j - t_{j'}| - 1}, \\ \nabla_{\sigma}\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} &= \frac{\sigma}{4} \sum_{j, j'=1}^p i_j i_{j'} \phi^{|t_j - t_{j'}|}. \end{aligned}$$

The main result of this paper is an expression for the elements of  $V$ , given in Theorem 1 below. Let  $c_i = E(\log u_t^2 - c_1)^i$ ,  $i = 3, 4$ . It is shown in the Appendix



that

$$c_3 = -14\zeta(3) = -16.829,$$

$$c_4 = \frac{7}{4}\pi^4 = 170.47,$$

where  $\zeta(z)$  is the Riemann zeta function. For  $i \geq 1$ , let

$$\kappa_i = \log 2 + \psi\left(\frac{i+1}{2}\right) - c_1,$$

$$\xi_i = \kappa_i^2 + \psi'\left(\frac{i+1}{2}\right) - c_2,$$

where  $\psi(z) = \frac{d}{dz} \log \Gamma(z)$ , the digamma function, and  $\psi'(z) = \frac{d}{dz} \psi(z)$ , the trigamma function.

**Theorem 1** For any  $a_t$  and  $b_t$ , let  $V(a_t, b_t) = \sum_{l=-\infty}^{\infty} \text{Cov}(a_t, b_{t-l})$ . Let  $i, j \geq 0$ . Then

$$V(z_t, z_t) = \frac{1+\phi}{1-\phi} \sigma^2 + c_2, \quad (6)$$

$$V(z_t, z_t z_{t-j}) = I_{(j=0)} c_3, \quad (7)$$

$$V(z_t z_{t-i}, z_t z_{t-j}) = A_1 \sigma^4 + A_2 c_2 \sigma^2 + I_{(i=j \neq 0)} c_2^2 + I_{(i=j=0)} (c_4 - c_2^2), \quad (8)$$

where

$$A_1 = |i-j| \phi^{|i-j|} + |i+j| \phi^{|i+j|} + \left( \phi^{|i-j|} + \phi^{|i+j|} \right) \frac{1+\phi^2}{1-\phi^2},$$

$$A_2 = 2 \left( \phi^{|i-j|} + \phi^{|i+j|} \right).$$

Let  $i_1, \dots, i_{p+q} \geq 1$ ;  $t_1 > \dots > t_p$ ;  $t_{p+1} > \dots > t_{p+q}$ ; and let

$$L = \{l | \{t_1, \dots, t_p\} \cap \{t_{p+1} - l, \dots, t_{p+q} - l\} \neq \emptyset\}.$$

Then

$$V\left(Y_{t_1, \dots, t_p}^{i_1, \dots, i_p}, Y_{t_{p+1}, \dots, t_{p+q}}^{i_{p+1}, \dots, i_{p+q}}\right) = B + \sum_{l \in L} (B_l + 1) C_l, \quad (9)$$

where

$$B = \sum_{l=-\infty}^{\infty} B_l, \quad B_l = \exp \left( \frac{\sigma^2}{4} \sum_{j=1}^p \sum_{j'=p+1}^{p+q} i_j i_{j'} \phi^{|t_j - t_{j'} + l|} \right) - 1,$$

$$C_l = \left( \prod_{j=1}^{p+q} \nu_{i_j}^{-1} \right) E \left( \prod_{j=1}^p |u_{t_j}^{i_j}| \prod_{j=p+1}^{p+q} |u_{t_j - l}^{i_j}| \right) - 1,$$

and

$$V \left( z_t, Y_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right) = D_1 \sigma^2 + D_2, \quad (10)$$

$$V \left( z_t z_{t-i}, Y_{t_1, \dots, t_p}^{i_1, \dots, i_p} \right) = D'_1 \sigma^4 + D'_2 \sigma^2 + D'_3, \quad (11)$$

where

$$D_1 = \frac{1}{2} \left( \frac{1+\phi}{1-\phi} \right) \sum_{j=1}^p i_j, \quad D_2 = \sum_{j=1}^p \kappa_{i_j},$$

$$D'_1 = \frac{1}{4} \sum_{j, j'=1}^p i_j i_{j'} \phi^{|t_j - t_{j'} + i|} \left( |t_j - t_{j'} + i| + \frac{1+\phi^2}{1-\phi^2} \right),$$

$$D'_2 = \frac{1}{2} \sum_{j, j'=1}^p i_j \kappa_{i_{j'}} \left( \phi^{|t_j - t_{j'} + i|} + \phi^{|t_j - t_{j'} - i|} \right),$$

$$D'_3 = I_{(i=0)} \sum_{j=1}^p \xi_{i_j} + \sum_{j, j'=1}^p I_{(i=t_j - t_{j'} \neq 0)} \kappa_{i_j} \kappa_{i_{j'}}.$$

We see that, not unexpectedly, the optimal weighting matrix,  $V^{-1}$ , and the GMM asymptotic covariance matrix,  $(D'V^{-1}D)^{-1}$ , do not depend on  $\mu$ . From a computational point of view, notice that  $L$  has at most  $pq$  elements, so computing  $\sum_{l \in L} (B_l + 1)C_l$  requires a finite number of steps. Furthermore,  $B$  can be approximated by  $B(I) = \sum_{l=-I}^I B_l$ , where  $I$  is a positive integer. As the following lemma shows, the error of approximation  $|B - B(I)|$  is bounded by an exponentially decaying function in  $I$ , and this bound can be inverted to determine  $I$  as a function of the desired accuracy of the approximation.

**Lemma 1** *Let  $i_1, \dots, i_{p+q} \geq 1$ ;  $t_1 > \dots > t_p$ ;  $t_{p+1} > \dots > t_{p+q}$ ; and let  $I$  be a positive integer. Then*

$$|B - B(I)| \leq 2 \frac{\exp(a|\phi|^I) - 1}{1 - |\phi|},$$

where

$$a = \frac{\sigma^2}{4} \sum_{j=1}^p \sum_{j'=p+1}^{p+q} i_j i_{j'} |\phi|^{-|t_j - t_{j'}|}.$$

If one is interested in  $\lambda$  rather than  $\theta$ , one may apply the transformation  $\theta \mapsto \lambda(\theta)$  to yield  $\hat{\lambda} = \lambda(\hat{\theta})$ , the optimal GMM estimator of  $\lambda$ , which has asymptotic covariance matrix  $(\frac{\partial \lambda}{\partial \theta})(D'V^{-1}D)^{-1}(\frac{\partial \lambda}{\partial \theta})'$ , with

$$\frac{\partial \lambda}{\partial \theta'} = \begin{pmatrix} 1 - \phi & -\mu & 0 \\ 0 & 1 & 0 \\ 0 & -\sigma\phi(1 - \phi^2)^{-1/2} & (1 - \phi^2)^{1/2} \end{pmatrix}.$$

## 4 Comparison of GMM and other estimators

In this section we first compare the relative efficiencies of GMM and other estimators, for two sets of values of  $\lambda$ , namely  $(\alpha, \phi, \omega) = (-0.736, 0.90, 0.363)$  and  $(\alpha, \phi, \omega) = (-0.1472, 0.98, 0.1657)$ . These parameter values have been used in earlier Monte Carlo studies (Jacquier, Polson, and Rossi (1994), Andersen and Sørensen (1996), Fridman and Harris (1998), Sandmann and Koopman (1998), Andersen, Chung, and Sørensen (1999)). Tables 1 and 2 present the results. The asymptotic standard errors of the GMM estimators were computed using the expressions derived above. The moment conditions were selected from the set related to the log-squared observations, or from the set related to the absolute observations, or from both. For comparability with other studies, from (5) we only selected moment conditions for which  $p = 1$  or for which  $p = 2$  and  $i_1 = i_2 \in \{1, 2\}$ . The (finite sample) standard errors of the other estimators were taken from the aforementioned Monte Carlo studies and multiplied by  $\sqrt{T}$ . The relative asymptotic efficiency of the GMM estimators is seen to increase rapidly

with the number of moments, at least when this number is small. Using a large number of moment conditions yields asymptotic standard errors slightly above those of the MCMC method, which is known to be asymptotically efficient. In this respect, it appears that some of the published standard errors regarding ML and Monte Carlo ML are not in line with those of the MCMC method.

In Monte Carlo studies it is often found that the small sample bias of the GMM estimator grows with the number of moment conditions. Newey and Smith (2000, 2001) show that the number of terms of the second-order bias increases linearly with the number of moment conditions. Thus, rather than using a large number of moment conditions (relative to the sample size), it is in terms of bias often safer to select only a small number of them. It is important, then, to choose the moments judiciously, in the sense that they contain as much information as possible for the estimand. Several authors have addressed the question of how to select the moment conditions to estimate the SV model, essentially by resorting to Monte Carlo simulation of the accuracy of the GMM estimator for any given choice of moments. The results of the previous section provide a more precise and much faster tool to guide the choice of moments. To illustrate this point, consider the sets  $M_L$  and  $M_A$  of log-moment and absolute moment conditions, respectively, defined as

$$M_L : (3)-(4) \text{ with } i \leq 50,$$

$$M_A : (5) \text{ with } \max_{j,j'} |t_j - t_{j'}| \leq 15 \text{ and } \sum_{j=1}^p i_j \leq \begin{cases} 20 & \text{for } p = 1, \\ 4 & \text{for } p = 2, 3, 4. \end{cases}$$

The sets  $M_L$  and  $M_A$  comprise 52 and 985 moment conditions, respectively. We performed a search for the set of  $k$  moment conditions, selected from either  $M_L$ ,  $M_A$ , or  $M_L \cup M_A$ , that yield the smallest asymptotic standard error of  $\hat{\phi}$ . Global optimisation, by enumeration, was performed over  $M_L$  for  $k = 3, 4, 5$ , and over  $M_A$  and  $M_L \cup M_A$  for  $k = 3$ . Global optimisation over  $M_A$  and  $M_L \cup M_A$  for  $k = 4$  and  $k = 5$  turned out to be infeasible in terms of computation time, and

**Table 1: Standard errors of  $\sqrt{T}\hat{\lambda}$**

$T$	# Moments	Method of estimation	Standard Error of		
			$\sqrt{T}\hat{\alpha}$	$\sqrt{T}\hat{\phi}$	$\sqrt{T}\hat{\omega}$
$\infty$	3	GMM (log-moments) <sup>a</sup>	127.52	17.31	32.66
$\infty$	12	GMM (log-moments) <sup>a</sup>	12.04	1.63	3.80
$\infty$	27	GMM (log-moments) <sup>a</sup>	10.06	1.36	3.22
$\infty$	102	GMM (log-moments) <sup>a</sup>	10.04	1.36	3.22
$\infty$	3	GMM (absolute moments) <sup>b</sup>	178.46	24.18	46.78
$\infty$	15	GMM (absolute moments) <sup>b</sup>	11.34	1.53	2.96
$\infty$	30	GMM (absolute moments) <sup>b</sup>	8.14	1.10	2.18
$\infty$	75	GMM (absolute moments) <sup>b</sup>	7.55	1.02	2.03
$\infty$	14	GMM (joint moments) <sup>c</sup>	16.92	2.29	4.27
$\infty$	22	GMM (joint moments) <sup>c</sup>	11.30	1.53	2.92
$\infty$	42	GMM (joint moments) <sup>c</sup>	8.12	1.10	2.14
$\infty$	102	GMM (joint moments) <sup>c</sup>	7.53	1.02	1.99
10000	14	Infeasible GMM <sup>d</sup> (true weight)	11.4	1.6	3.1
4000	14	Infeasible GMM <sup>d</sup> (true weight)	10.6	1.5	3.1
4000	4	EMM: GARCH(1,1) <sup>e</sup>	9.51	1.2	3.1
4000	6	EMM: GARCH(1,1) - Kz(2) <sup>e</sup>	9.68	1.3	3.2
4000	8	EMM: GARCH(1,1) - Kz(4) <sup>e</sup>	8.28	1.1	2.1
2000	24	GMM <sup>f</sup>	18	3	3.8
2000	-	Quasi-ML <sup>f</sup>	20	3	4.8
2000	-	MCMC <sup>f</sup>	6.6	1	1.5
500	-	ML <sup>g</sup>	9.1	1	2
500	-	Monte Carlo ML <sup>h</sup>	0.5	2.2	2

Parameter values:  $(\alpha, \phi, \omega) = (-0.736, 0.90, 0.363)$ .

GMM conditions are selected from Eqs. (3)–(5), as indicated below. Most footnotes refer to multiple lines in the Table.

a. Eqs. (3)–(4) with  $i$  running from 0 to 1, 10, 25, and 100, respectively.

b. Eqs. (5) with  $p = 1$ ,  $i_1$  running from 1 to 1, 5, 10, and 25 respectively; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 1, 5, 10, and 25 respectively.

c. Eqs. (3)–(4) with  $i$  running from 0 to 3, 5, 10, and 25, respectively; Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 3, 5, 10, and 25, respectively; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 3, 5, 10, and 25, respectively.

d. Andersen and Sørensen (1996), Table 3: Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 4; Eq. (5) with  $p = 2$ ,  $i_1 = i_2 = 1$ ,  $t_1 - t_2 \in \{6, 8, 10, 12, 14\}$ ; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 = 2$ ,  $t_1 - t_2 \in \{15, 17, 19, 21, 23\}$ . ‘Infeasible GMM’ uses a nonparametric estimate of the weighting matrix based on a large sample of simulated data using true parameter values.

e. Andersen, Chung, and Sørensen (1999), Table 3.

f. Jacquier, Polson, and Rossi (1994), Tables 5–7. For GMM: Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 4; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 10.

g. Fridman and Harris (1998), Table 1.

h. Sandmann and Koopman (1998), Table 2.

**Table 2: Standard errors of  $\sqrt{T}\hat{\lambda}$**

$T$	# Moments	Method of estimation	Standard Error of		
			$\sqrt{T}\hat{\alpha}$	$\sqrt{T}\hat{\phi}$	$\sqrt{T}\hat{\omega}$
$\infty$	3	GMM (log-moments) <sup>a</sup>	136.37	18.53	77.30
$\infty$	12	GMM (log-moments) <sup>a</sup>	6.67	0.90	4.00
$\infty$	27	GMM (log-moments) <sup>a</sup>	2.96	0.40	1.71
$\infty$	52	GMM (log-moments) <sup>a</sup>	2.51	0.34	1.39
$\infty$	102	GMM (log-moments) <sup>a</sup>	2.49	0.34	1.37
$\infty$	3	GMM (absolute moments) <sup>b</sup>	264.71	35.95	150.79
$\infty$	15	GMM (absolute moments) <sup>b</sup>	8.49	1.15	4.79
$\infty$	30	GMM (absolute moments) <sup>b</sup>	4.15	0.56	2.28
$\infty$	75	GMM (absolute moments) <sup>b</sup>	2.48	0.34	1.23
$\infty$	14	GMM (joint moments) <sup>c</sup>	14.95	2.03	8.43
$\infty$	22	GMM (joint moments) <sup>c</sup>	8.45	1.15	4.76
$\infty$	42	GMM (joint moments) <sup>c</sup>	4.12	0.56	2.26
$\infty$	102	GMM (joint moments) <sup>c</sup>	2.44	0.33	1.20
4000	4	EMM: GARCH(1,1) <sup>d</sup>	2.8	0.37	1.3
4000	4	EMM: GARCH(1,1) - Kz(2) <sup>d</sup>	2.9	0.39	1.8
4000	6	EMM: GARCH(1,1) - Kz(4) <sup>d</sup>	2.7	0.36	1.0
500	24	GMM <sup>e</sup>	5.8	0.80	2
500	-	Quasi-ML <sup>e</sup>	12	2	3.1
500	-	MCMC <sup>e</sup>	2.7	0.4	1
500	-	ML <sup>f</sup>	0.4	0.30	0.8
500	-	Monte Carlo ML <sup>g</sup>	0.2	2	1

Parameter values:  $(\alpha, \phi, \omega) = (-0.1472, 0.98, 0.1657)$ .

GMM conditions are selected from Eqs. (3)–(5), as indicated below. Most footnotes refer to multiple lines in the Table.

a. Eqs. (3)–(4) with  $i$  running from 0 to 1, 10, 25, 50, and 100, respectively.

b. Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 1, 5, 10, and 25 respectively; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 1, 5, 10, and 25 respectively.

c. Eqs. (3)–(4) with  $i$  running from 0 to 3, 5, 10, and 25, respectively; Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 3, 5, 10, and 25, respectively; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 3, 5, 10, and 25, respectively.

d. Andersen, Chung, and Sørensen (1999), Table 3.

e. Jacquier, Polson, and Rossi (1994), Tables 5–7. For GMM: Eq. (5) with  $p = 1$ ,  $i_1$  running from 1 to 4; and Eq. (5) with  $p = 2$ ,  $i_1 = i_2 \in \{1, 2\}$ ,  $t_1 - t_2$  running from 1 to 10.

f. Fridman and Harris (1998), Table 1.

g. Sandmann and Koopman (1998), Table 2.

**Table 3: Asymptotic standard errors of  $\sqrt{T}\hat{\phi}$  for parsimoniously selected moments**

Moment set	$k$	Selected moments	Standard Error of		
			$\sqrt{T}\hat{\alpha}$	$\sqrt{T}\hat{\phi}$	$\sqrt{T}\hat{\omega}$
$M_L$	3	$z_t; z_t z_{t-1}; z_t z_{t-11}$	18.31	2.49	5.41
$M_A$	3	$Y_t^2; Y_{t,t-7}^{1,2}; Y_{t,t-5,t-14}^{1,1,1}$	10.59	1.44	4.72
$M_L \cup M_A$	3	$z_t z_{t-10}; Y_t^2; Y_{t,t-7,t-15}^{1,1,1}$	10.08	1.37	4.07
$M_L$	4	$z_t; z_t z_{t-1}; z_t z_{t-10}; z_t z_{t-12}$	14.78	2.01	4.62
$M_A$	4	$Y_t^1; Y_t^2; Y_{t,t-10}^{1,1}; Y_{t,t-8,t-15}^{1,1,1}$	9.65	1.31	2.55
$M_L \cup M_A$	4	$z_t z_{t-10}; Y_t^2; Y_{t,t-5,t-14}^{1,1,1}; Y_{t,t-7,t-13}^{1,1,1}$	9.46	1.28	4.16
$M_L$	5	$z_t; z_t z_{t-1}; z_t z_{t-9}; z_t z_{t-11}; z_t z_{t-14}$	13.37	1.82	4.31
$M_A$	5	$Y_t^1; Y_t^2; Y_{t,t-7}^{1,1}; Y_{t,t-9,t-13}^{1,1}$	9.07	1.23	2.47
$M_L \cup M_A$	5	$Y_t^1; Y_t^2; Y_{t,t-7}^{1,1}; Y_{t,t-9,t-13}^{1,1}$	9.07	1.23	2.47

Parameter values:  $(\alpha, \phi, \omega) = (-0.736, 0.90, 0.363)$ .

in these cases we experimented with the Point Exchange algorithm (Fedorov (1972)). This algorithm does not necessarily yield the global optimum, and its output depends on the starting selection of moment conditions as input. By picking the starting selection at random and repeating this a couple of times, the algorithm was able to reproduce the global optimum in all cases where enumeration was possible. We therefore applied it in those cases where global optimisation was not feasible, without the guarantee of having found the globally optimal selection of moments from the specified sets. The parameter values were fixed at  $(\alpha, \phi, \omega) = (-0.736, 0.90, 0.363)$ , as in Table 1. Table 3 reports the selected moments and the asymptotic standard errors of the corresponding GMM estimators. Comparing Table 3 with Table 1 yields the following conclusions: (i) there is a dramatic increase in efficiency by selecting the moments in an

optimal way; (ii) given that the efficiency bound for the asymptotic standard error of  $\sqrt{T}\hat{\phi}$  (which is asymptotically attained by the MCMC estimator) appears to be around 1, the efficiency loss of the GMM estimator with optimal moment selection from  $M_L \cup M_A$  is not excessively large, even in the just-identified case ( $k = 3$ ); (iii) while  $M_A$  contains a richer (also a much larger) set of moment conditions than  $M_L$  – as is reflected by the smaller asymptotic standard errors – the combination of  $M_A$  and  $M_L$  may yield an improvement upon  $M_A$ , as is the case here for  $k = 3$  and  $k = 4$ . Finally, we remark that the optimal selection of moment conditions from any given set generally depends on the parameter values.

## 5 Conclusion

The standard approach in the literature on GMM estimation of SV models has been to derive closed-form moment conditions from the expectations of  $|y_t^i|$ ,  $|y_{t_1}y_{t_2}|$  and  $|y_{t_1}^2y_{t_2}^2|$  for any  $i$ ,  $t_1$ , and  $t_2$ . We have extended this class of conditions to include the expectation of  $|y_{t_1}^{i_1}\dots y_{t_p}^{i_p}|$  for arbitrary  $i_1, \dots, i_p$  and  $t_1, \dots, t_p$ , and, following Wright (1999), the first two moments and the autocovariances of  $\log y_t^2$ . A closed-form expression for the optimal weighting matrix for any subset of those conditions has been derived and, as a by-product, an expression for the GMM asymptotic covariance matrix. These expressions can be used for improved GMM estimation of the SV model with AR(1) log-volatility and to compute GMM standard errors more accurately. It is also of interest to note that, upon redefining  $c_i$ ,  $\nu_i$ ,  $\kappa_i$ , and  $\xi_i$  appropriately, all expressions are generalised to SV models where the multiplicative shocks in the mean equation (1) are non-normal.



## Appendix

**Calculation of  $c_3$  and  $c_4$ .** For any positive integer  $n$ , upon substituting  $t = x^2/2$ ,

$$\begin{aligned} \int_{-\infty}^{\infty} \left(\log \frac{x^2}{2}\right)^n \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx &= 2 \int_0^{\infty} \left(\log \frac{x^2}{2}\right)^n \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \frac{1}{\sqrt{\pi}} \int_0^{\infty} (\log t)^n t^{-1/2} e^{-t} dt \\ &= \frac{\Gamma^{(n)}\left(\frac{1}{2}\right)}{\Gamma\left(\frac{1}{2}\right)}, \end{aligned}$$

where  $\Gamma^{(n)}(z)$  is the  $n$ -th derivative of  $\Gamma(z)$ . See Abramowitz and Stegun (1970) for properties and values of the gamma and related functions that are used below. Now,  $c_1 = -\log 2 - \gamma = \psi\left(\frac{1}{2}\right) + \log 2$ , where  $\psi(z) = \frac{d}{dz} \log \Gamma(z) = \frac{\Gamma'(z)}{\Gamma(z)}$ . Hence, for  $n = 3, 4$ ,

$$\begin{aligned} c_n &= \int_{-\infty}^{\infty} (\log x^2 - c_1)^n \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \int_{-\infty}^{\infty} \left(\log \frac{x^2}{2} - \psi\left(\frac{1}{2}\right)\right)^n \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= g_n\left(\frac{1}{2}\right), \end{aligned}$$

where

$$g_n(z) = \sum_{i=0}^n \binom{n}{i} \frac{\Gamma^{(i)}(z)}{\Gamma(z)} (-\psi(z))^{n-i}. \quad (12)$$

Taking successive derivatives of  $\Gamma^{(1)}(z) = \Gamma(z)\psi(z)$  gives, upon rewriting,

$$\begin{aligned} \Gamma^{(2)}(z) &= \Gamma(z) \left\{ \psi'(z) + [\psi(z)]^2 \right\}, \\ \Gamma^{(3)}(z) &= \Gamma(z) \left\{ \psi''(z) + 3\psi'(z)\psi(z) + [\psi(z)]^3 \right\}, \\ \Gamma^{(4)}(z) &= \Gamma(z) \left\{ \psi'''(z) + 4\psi''(z) + 3\psi'(z) \left[ \psi'(z) + 2[\psi(z)]^2 \right] + [\psi(z)]^3 \right\}, \end{aligned}$$

where primes denote derivatives. Substituting these expressions into (12) yields

$$g_3(z) = \psi''(z)$$

and

$$g_4(z) = \psi'''(z) + 3[\psi'(z)]^2.$$

Now,  $\psi'(\frac{1}{2}) = \frac{\pi^2}{2}$ ,  $\psi''(\frac{1}{2}) = -14\zeta(3)$ , and  $\psi'''(\frac{1}{2}) = \pi^4$ , where  $\zeta(3) = 1.202$ . Hence  $c_3 = -14\zeta(3) = -16.83$  and  $c_4 = \frac{7}{4}\pi^4 = 170.5$ .

The proof of Theorem 1 makes use of the following lemmas.

**Lemma 2** *Let  $X \sim N(0, 1)$  and let  $a$  be a positive integer. Then*

$$\text{Cov}(\log X^2, |X^a|) = \nu_a \kappa_a \quad (13)$$

and

$$\text{Cov}((\log X^2 - c_1)^2, |X^a|) = \nu_a \xi_a. \quad (14)$$

**Proof.** Upon substituting  $z = x^2/2$ ,

$$\begin{aligned} \text{Cov}(\log X^2, |X^a|) &= 2 \int_0^\infty (\log x^2 - c_1) x^a \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \frac{2^{a/2}}{\sqrt{\pi}} \int_0^\infty (\log z + \log 2 - c_1) z^{(a-1)/2} e^{-z} dz \\ &= \frac{2^{a/2}}{\sqrt{\pi}} \Gamma\left(\frac{a+1}{2}\right) \left(\psi\left(\frac{a+1}{2}\right) + \log 2 - c_1\right) \\ &= \nu_a \kappa_a, \end{aligned}$$

and, using  $\frac{\Gamma''(z)}{\Gamma(z)} = \psi'(z) + (\psi(z))^2$  (with primes denoting derivatives),

$$\begin{aligned} \text{Cov}((\log X^2 - c_1)^2, |X^a|) &= 2 \int_0^\infty [(\log x^2 - c_1)^2 - c_2] x^a \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \frac{2^{a/2}}{\sqrt{\pi}} \int_0^\infty [(\log z + \log 2 - c_1)^2 - c_2] z^{(a-1)/2} e^{-z} dz \\ &= \frac{2^{a/2}}{\sqrt{\pi}} \Gamma\left(\frac{a+1}{2}\right) \left(\frac{\Gamma''\left(\frac{a+1}{2}\right)}{\Gamma\left(\frac{a+1}{2}\right)} + (\log 2 - c_1)^2 - c_2\right. \\ &\quad \left.+ 2(\log 2 - c_1)\psi\left(\frac{a+1}{2}\right)\right) \\ &= \nu_a \xi_a. \end{aligned}$$

**Lemma 3** *Let  $X_1, X_2,$  and  $X_3$  be jointly normal with  $\mu_i = EX_i$  and  $\sigma_{ij} = \text{Cov}(X_i, X_j)$ . Then*

$$\text{Cov}(X_1, \exp X_3) = \sigma_{13} \exp\left(\mu_3 + \frac{1}{2}\sigma_{33}\right) \quad (15)$$

and

$$\text{Cov}(X_1 X_2, \exp X_3) = (\sigma_{13} \sigma_{23} + \mu_1 \sigma_{23} + \mu_2 \sigma_{13}) \exp\left(\mu_3 + \frac{1}{2} \sigma_{33}\right). \quad (16)$$

**Proof.** Assume first that  $\mu_i = 0$  and  $\sigma_{ii} > 0$  for all  $i$ . Let  $\mu_{i|j} = \sigma_{ij} \sigma_{jj}^{-1} X_j$  be the conditional mean of  $X_i$ , given  $X_j$ , and  $\sigma_{ij|k} = \sigma_{ij} - \sigma_{ik} \sigma_{jk} \sigma_{kk}^{-1}$  the conditional covariance between  $X_i$  and  $X_j$ , given  $X_k$ . Then,

$$\begin{aligned} \text{Cov}(X_1, \exp X_3) &= \text{Cov}(E(X_1|X_3), \exp X_3) \\ &= \sigma_{13} \sigma_{33}^{-1} \text{Cov}(X_3, \exp X_3) \\ &= \sigma_{13} \exp\left(\frac{1}{2} \sigma_{33}\right), \end{aligned}$$

$$\begin{aligned} \text{Cov}(X_1 X_2, \exp X_3) &= \text{Cov}(E(X_1 X_2|X_3), \exp X_3) \\ &= \text{Cov}(\mu_{1|3} \mu_{2|3} + \sigma_{12|3}, \exp X_3) \\ &= \sigma_{13} \sigma_{23} \sigma_{33}^{-2} \text{Cov}(X_3^2, \exp X_3) \\ &= \sigma_{13} \sigma_{23} \exp\left(\frac{1}{2} \sigma_{33}\right), \end{aligned}$$

using the fact that, for a standard normal variate  $X$ ,

$$\begin{aligned} \text{Cov}(X, \exp(bX)) &= \int_{-\infty}^{\infty} x e^{bx} \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \int_{-\infty}^{\infty} x \left(\frac{1}{\sqrt{2\pi}}\right) e^{-(x-b)^2/2+b^2/2} dx \\ &= b \exp\left(\frac{1}{2} b^2\right) \end{aligned}$$

and

$$\begin{aligned} \text{Cov}(X^2, \exp(bX)) &= \int_{-\infty}^{\infty} (x^2 - 1) e^{bx} \left(\frac{1}{\sqrt{2\pi}}\right) e^{-x^2/2} dx \\ &= \int_{-\infty}^{\infty} (x^2 - 1) \left(\frac{1}{\sqrt{2\pi}}\right) e^{-(x-b)^2/2+b^2/2} dx \\ &= b^2 \exp\left(\frac{1}{2} b^2\right). \end{aligned}$$

The extension to the case where  $\mu_i \neq 0$  for some  $i$  is straightforward, and any degenerate case follows upon taking the appropriate limit in the non-degenerate case.

**Proof of Theorem 1.** Write  $z_t = k_t + w_t$ , where  $k_t = h_t - \mu$  and  $w_t = \log u_t^2 - c_1$ . Then,  $w_t$  and  $k_t$  have zero mean and are independent, and, for any integers  $i, j, l$ , we have  $\text{Cov}(k_t, k_{t-i}) = \phi^{|i|}\sigma^2$ ,  $\text{Cov}(k_t k_{t-i}, k_{t-j}) = 0$ , and  $\text{Cov}(k_t k_{t-i}, k_{t-j} k_{t-l}) = (\phi^{|j|+|i-l|} + \phi^{|l|+|i-j|})\sigma^4$ . Using these properties and the equalities

$$\sum_{l=-\infty}^{\infty} \phi^{|l|} = \frac{1+\phi}{1-\phi},$$

$$\sum_{l=-\infty}^{\infty} \phi^{|i+l|+|j+l|} = \phi^{|i-j|} \left( |i-j| + \frac{1+\phi^2}{1-\phi^2} \right),$$

we obtain, for  $i, j \geq 0$ ,

$$\begin{aligned} V(z_t, z_t) &= \sum_{l=-\infty}^{\infty} [\text{Cov}(k_t, k_{t-l}) + \text{Cov}(w_t, w_{t-l})] \\ &= \sum_{l=-\infty}^{\infty} \left( \phi^{|l|}\sigma^2 + I_{(l=0)}c_2 \right) \\ &= \frac{1+\phi}{1-\phi}\sigma^2 + c_2, \\ V(z_t, z_t z_{t-j}) &= \sum_{l=-\infty}^{\infty} \text{Cov}(w_t, w_{t-l} w_{t-j-l}) \\ &= \sum_{l=-\infty}^{\infty} I_{(l=j=0)}c_3 \\ &= I_{(j=0)}c_3, \end{aligned}$$

and

$$\begin{aligned}
V(z_t z_{t-i}, z_t z_{t-j}) &= \sum_{l=-\infty}^{\infty} [\text{Cov}(k_t k_{t-i}, k_{t-l} k_{t-j-l}) + \text{Cov}(w_t k_{t-i}, k_{t-l} w_{t-j-l}) \\
&\quad + \text{Cov}(w_t k_{t-i}, w_{t-l} k_{t-j-l}) + \text{Cov}(k_t w_{t-i}, k_{t-l} w_{t-j-l}) \\
&\quad + \text{Cov}(k_t w_{t-i}, w_{t-l} k_{t-j-l}) + \text{Cov}(w_t w_{t-i}, w_{t-l} w_{t-j-l})] \\
&= \sum_{l=-\infty}^{\infty} \left[ \left( \phi^{|l|+|j-i+l|} + \phi^{|j+l|+|l-i|} \right) \sigma^4 + I_{(l=-j)} \phi^{|i+j|} c_2 \sigma^2 \right. \\
&\quad + I_{(l=0)} \phi^{|i-j|} c_2 \sigma^2 + I_{(l=i-j)} \phi^{|i-j|} c_2 \sigma^2 + I_{(l=i)} \phi^{|i+j|} c_2 \sigma^2 \\
&\quad \left. + I_{(i=j \neq 0)} I_{(l=0)} c_2^2 + I_{(i=j=0)} I_{(l=0)} (c_4 - c_2^2) \right] \\
&= \left( |i-j| \phi^{|i-j|} + |i+j| \phi^{|i+j|} + \left( \phi^{|i-j|} + \phi^{|i+j|} \right) \frac{1+\phi^2}{1-\phi^2} \right) \sigma^4 \\
&\quad + 2 \left( \phi^{|i-j|} + \phi^{|i+j|} \right) c_2 \sigma^2 + I_{(i=j \neq 0)} c_2^2 + I_{(i=j=0)} (c_4 - c_2^2),
\end{aligned}$$

giving (6)–(8). To establish (9), recall the definition of  $Y_{t_1, \dots, t_p}^{i_1, \dots, i_p}$ , from which

$$\begin{aligned}
&\text{Cov} \left( Y_{t_1, \dots, t_p}^{i_1, \dots, i_p}, Y_{t_{p+1}-l, \dots, t_{p+q}-l}^{i_{p+1}, \dots, i_{p+q}} \right) \\
&= E \left( Y_{t_1, \dots, t_p}^{i_1, \dots, i_p} Y_{t_{p+1}-l, \dots, t_{p+q}-l}^{i_{p+1}, \dots, i_{p+q}} \right) - 1 \\
&= E \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} - \delta_{t_{p+1}, \dots, t_{p+q}}^{i_{p+1}, \dots, i_{p+q}} + \frac{1}{2} \sum_{j=1}^p i_j h_{t_j} + \frac{1}{2} \sum_{j=p+1}^{p+q} i_j h_{t_j-l} \right) \\
&\quad \times E \left( \prod_{j=1}^p \left| \nu_{i_j}^{-1} u_{t_j}^{i_j} \right| \prod_{j=p+1}^{p+q} \left| \nu_{i_j}^{-1} u_{t_j-l}^{i_j} \right| \right) - 1.
\end{aligned}$$

Now,

$$\begin{aligned}
&E \exp \left( \frac{1}{2} \sum_{j=1}^p i_j h_{t_j} + \frac{1}{2} \sum_{j=p+1}^{p+q} i_j h_{t_j-l} \right) \\
&= \exp \left( \delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} + \delta_{t_{p+1}, \dots, t_{p+q}}^{i_{p+1}, \dots, i_{p+q}} + \frac{\sigma^2}{4} \sum_{j=1}^p \sum_{j'=p+1}^{p+q} i_j i_{j'} \phi^{|t_j - t_{j'} + l|} \right),
\end{aligned}$$

so

$$\begin{aligned}
\text{Cov} \left( Y_{t_1, \dots, t_p}^{i_1, \dots, i_p}, Y_{t_{p+1}-l, \dots, t_{p+q}-l}^{i_{p+1}, \dots, i_{p+q}} \right) &= (B_l + 1)(C_l + 1) - 1 \\
&= B_l + (B_l + 1)C_l.
\end{aligned}$$

Summing over  $l$  gives (9), because  $C_l = 0$  whenever  $l \notin L$ . Furthermore, by (13) and (15),

$$\begin{aligned}
& \text{Cov} \left( z_t, Y_{t_1-l, \dots, t_p-l}^{i_1, \dots, i_p} \right) \\
&= E \left( k_t \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} + \frac{1}{2} \sum_{j=1}^p i_j h_{t_j-l} \right) \right) \\
&\quad + \left( E \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} + \frac{1}{2} \sum_{j=1}^p i_j h_{t_j-l} \right) \right) E \left( w_t \prod_{j=1}^p \left| \nu_{i_j}^{-1} u_{t_j-l}^{i_j} \right| \right) \\
&= \frac{\sigma^2}{2} \sum_{j=1}^p i_j \phi^{|t-t_j+l|} + \sum_{j=1}^p I_{(t=t_j-l)} \kappa_{i_j},
\end{aligned}$$

which, upon summing over  $l$ , gives (10). Finally,

$$\begin{aligned}
\text{Cov} \left( z_t z_{t-i}, Y_{t_1-l, \dots, t_p-l}^{i_1, \dots, i_p} \right) &= \text{Cov} \left( k_t k_{t-i} + k_t w_{t-i} + w_t k_{t-i} + w_t w_{t-i}, \right. \\
&\quad \left. \exp \left( -\delta_{t_1, \dots, t_p}^{i_1, \dots, i_p} + \frac{1}{2} \sum_{j=1}^p i_j h_{t_j-l} \right) \prod_{j=1}^p \left| \nu_{i_j}^{-1} u_{t_j-l}^{i_j} \right| \right) \\
&= T_1 + T_2 + T_3 + T_4,
\end{aligned}$$

say, with, using Lemma 2 and Lemma 3,

$$\begin{aligned}
T_1 &= \left( \frac{\sigma^2}{2} \sum_{j=1}^p i_j \phi^{|t-t_j+l|} \right) \left( \frac{\sigma^2}{2} \sum_{j=1}^p i_j \phi^{|t-i-t_j+l|} \right), \\
T_2 &= \left( \frac{\sigma^2}{2} \sum_{j=1}^p i_j \phi^{|t-t_j+l|} \right) \sum_{j=1}^p I_{(t-i=t_j-l)} \kappa_{i_j}, \\
T_3 &= \left( \frac{\sigma^2}{2} \sum_{j=1}^p i_j \phi^{|t-i-t_j+l|} \right) \sum_{j=1}^p I_{(t=t_j-l)} \kappa_{i_j}, \\
T_4 &= I_{(i=0)} \sum_{j=1}^p I_{(t=t_j-l)} \xi_{i_j} + I_{(i \neq 0)} \sum_{j, j'=1}^p I_{(t=t_j-l)} I_{(t-i=t_{j'}-l)} \kappa_{i_j} \kappa_{i_{j'}}.
\end{aligned}$$

Summing over  $l$  gives (11), which concludes the proof.

**Proof of Lemma 1.** Since  $|\phi|^{t_j - t_{j'} + l} \leq |\phi|^{-|t_j - t_{j'}| + |l|}$ ,

$$B_l \leq \exp\left(\frac{\sigma^2}{4} \sum_{j=1}^p \sum_{j'=p+1}^{p+q} i_j i_{j'} |\phi|^{t_j - t_{j'} + l}\right) - 1 \leq \exp(a|\phi|^{|l|}) - 1.$$

By an argument of symmetry,

$$\exp(-a|\phi|^{|l|}) - 1 \leq B_l \leq \exp(a|\phi|^{|l|}) - 1$$

and so, because  $\exp(z) - 1 \geq 1 - \exp(-z)$  for any  $z$ ,

$$|B_l| \leq \exp(a|\phi|^{|l|}) - 1.$$

Therefore,

$$\begin{aligned} |B - B(I)| &\leq \sum_{l=I+1}^{\infty} (|B_{-l}| + |B_l|) \leq 2 \sum_{l=I}^{\infty} (\exp(a|\phi|^l) - 1) \\ &= 2 \sum_{l=I}^{\infty} \sum_{k=1}^{\infty} \frac{a^k |\phi|^{kl}}{k!} = 2 \sum_{k=1}^{\infty} \left(\frac{a^k}{k!}\right) \left(\frac{|\phi|^{kI}}{1 - |\phi|^k}\right) \\ &< \frac{2}{1 - |\phi|} \sum_{k=1}^{\infty} \frac{(a|\phi|^I)^k}{k!} = 2 \frac{\exp(a|\phi|^I) - 1}{1 - |\phi|}. \end{aligned}$$

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