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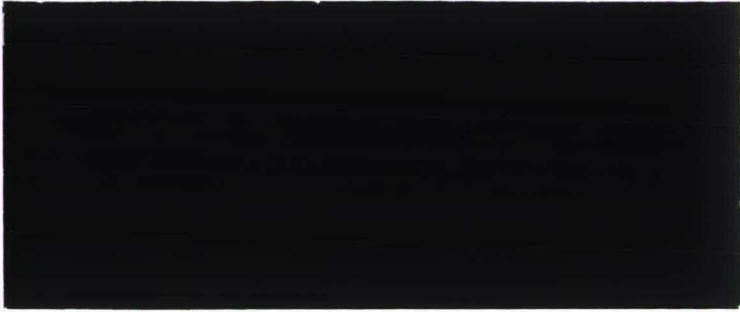


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DEPARTMENT OF ECONOMICS
RESEARCH MEMORANDUM

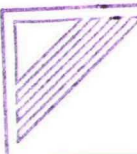
DENSITY OF THE LEAST SQUARES ESTIMATOR
IN THE MULTIVARIATE LINEAR MODEL WITH
ARBITRARILY NORMAL VARIABLES

B.B. van der Genugten

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DENSITY OF THE LEAST SQUARES ESTIMATOR IN THE MULTIVARIATE LINEAR MODEL WITH ARBITRARILY NORMAL VARIABLES

B.B. VAN DER GENUGTEN*

Abstract

An expression is obtained for the density of the least squares estimator $\hat{B} = (X_2'X_2)^{-1}X_2'X_1$ in the multivariate linear model $X_1 = X_2B + E$ for the case that $X = [X_1 \ X_2]$ has a normal distribution.

This expression is a multi-dimensional integral with a dimension independent of the number of rows of X .

Two particular cases are worked out in detail.

Keywords

multivariate linear model, least squares estimator, numerical evaluation of probability densities.

AMS classification: 60E05, 62J05.

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1 Introduction and statement of the result

Consider the multivariate linear model

$$X_1 = X_2 B + E \quad (1)$$

with dependent variables $X_1 \in \mathbf{R}^{n \times p}$ and explanatory variables $X_2 \in \mathbf{R}^{n \times k}$. The least squares estimator for B is given by

$$\hat{B} = (X_2' X_2)^{-1} X_2' X_1. \quad (2)$$

Assume that $X = [X_1 \ X_2] \in \mathbf{R}^{n \times m}$ (with $m = p + k$) has a non-singular normal distribution:

$$\text{vec}(X) \sim N_{nm}(\mu, \Omega), \quad \Omega > 0. \quad (3)$$

The theorem below gives the density g of \hat{B} as a multi-dimensional integral of the form

$$g(B) = \int_{\mathbf{R}^{k \times p}} \varphi(B, U_{21}) dU_{21}. \quad (4)$$

The dimension does not depend on n .

In section 2 the particular case $\mu = 0$, $\Omega = I_{nm}$ is considered. It is shown that the expression (4) leads to the density of the (central) matrix t -distribution.

In section 3 a detailed expression is given for the two-dimensional case $p = 1$, $k = 2$. This result can be used immediately for numerical calculations.

In the integrand φ of (4) the concept of the determinant- differential operator is used. Such an operator D works on functions $\varphi : \mathbf{R}^{k \times k} \rightarrow \mathbf{R}$ and is defined by $D\psi(u) = (\det \partial/\partial U = \{\partial/\partial u_{rs}\}, r, s = 1, \dots, k$. So $D\psi(U)$ is a linear combination of $k!$ partial derivatives of ψ of order k . As usual higher orders are defined by $D^p = D^{p-1}$ ($p = 2, 3, \dots$). The particular value at $U = U_0$ is denoted by $D^p\psi(U_0)$.

The operator D appears in a natural way in the differentiation of characteristic functions. For $X, U \in \mathbf{R}^{k \times k}$ we have

$$\begin{aligned}
\int |X| f(X) \operatorname{etr}(iU'X) dx &= \int |X| f(X) \exp\{i \sum_r \sum_s u_{rs} x_{rs}\} dX = \\
&= i^{-k} \sum_{(i_1, \dots, i_k)} \int f(X) (\partial^k / \partial u_{i_k} \dots \partial u_{i_1}) \operatorname{etr}(iU'X) dX \\
&= i^{-k} D \int f(X) \operatorname{etr}(iU'X) dx
\end{aligned}$$

and more general for $p = 1, 2, \dots$:

$$\int_{\mathbf{R}^{k \times k}} |X|^p f(X) \operatorname{etr}(iU'x) dX = i^{-kp} D^p \int_{\mathbf{R}^{k \times k}} f(X) \operatorname{etr}(iU'X) dX \quad (5)$$

Theorem

Under the condition that the integral in (4) is absolutely convergent we have:

$$\varphi(B, U_{21}) = \left(\frac{1}{2\pi i}\right)^{kp} D_{22}^p \psi(U_{21}, -(BU'_{21} + U_{21}B')/2) \quad (6)$$

with determinant-differential operator

$$D_{22} = \det(\partial/\partial U_{22}), \quad U_{22} \in \mathbf{R}^{k \times k} \quad (7)$$

and with

$$\psi(U_{21}, U_{22}) = |\Omega|^{1/2} \exp\left(-\frac{1}{2}\mu'\Omega^{-1}\mu\right) \cdot |\Phi|^{-1/2} \exp\left(\frac{1}{2}\mu'\Omega^{-1}\Phi^{-1}\Omega^{-1}\mu\right) \quad (8)$$

$$\Phi = \Phi(U_{21}, U_{22}) = \Omega^{-1} - i(V + V') \quad (9)$$

$$V = V(U_{21}, U_{22}) = U \otimes I_n \quad (10)$$

$$U = \begin{bmatrix} 0 & 0 \\ U_{21} & U_{22} \end{bmatrix} \in \mathbf{R}^{m \times m} \quad (11)$$

The theorem shows that for the numerical calculation of the integral (4) an expression for $D_{22}^p \psi$ in (6) is needed. Therefore such a calculation is only straightforward possible with computer packages having the possibility of alpha-numeric evaluation (e.g. Mathematica)

2 The particular case $\mu = 0$, $\Omega = I_{nm}$

Substitution of $\mu = 0$ and $\Omega = I_{nm} = I_m \otimes I_n$ into (9) gives

$$\Phi = \left(\left[\begin{array}{cc} I_p & 0 \\ 0 & I_k \end{array} \right] - i \left[\begin{array}{cc} 0 & U'_{21} \\ U_{21} & U_{22} + U'_{22} \end{array} \right] \right) \otimes I_n$$

So, with (8)

$$\psi(U_{21}, U_{22}) = |\Phi|^{-1/2} = |I_k + U_{21}U'_{21} - i(U_{22} + U'_{22})|^{-n/2} \quad (12)$$

It is not easy to calculate $D_{22}^p \psi(U_{21}, U_{22})$ from (12). Using some known results about central Wishart-distributions it can be shown that

$$D_{22}^p \psi(U_{21}, U_{22}) = (2i)^{kp} \frac{\Gamma_k(n/2 + p)}{\Gamma_k(n/2)} \cdot |I_k + U_{21}U'_{21} - i(U_{22} + U'_{22})|^{-n/2-p}, \quad (13)$$

where Γ_k stands for the multivariate gammafunction

$$\Gamma_k(n/2) = \pi^{k(k-1)/4} \prod_{j=1}^k \Gamma((n-j+1)/2). \quad (14)$$

The derivation of (13) from (12) can be found at the end of this section.

Substitution of (13) into (6) leads with (4) to

$$g(B) = \left(\frac{1}{\pi}\right)^{kp} \frac{\Gamma_k(n/2 + p)}{\Gamma_k(n/2)} \int_{\mathbf{R}^{k \times p}} |I_k + U_{21}U'_{21} + i(BU'_{21} + U_{21}B')|^{-n/2-p} dU_{21}$$

The integral can be reduced in the same way as in Phillips (1985), p. 160. Transform with $W = (I_k + BB')^{-1/2}(U_{21} + iB)$. Then

$$I_k + U_{21}U'_{21} + i(BU'_{21} + U_{21}B') = (I_k + BB')^{1/2}(I_k + WW')(I_k + BB')^{1/2}.$$

The transformation has Jacobian $|I_k + BB'|^{p/2}$. The domain of integration can be taken to be \mathbf{R}^{ki} as before using some standard arguments in contour-integration. This leads to

$$g(B) = \left(\frac{1}{\pi}\right)^{kp} \frac{\Gamma_k(n/2 + p)}{\Gamma_k(n/2)} |I_k + BB'|^{-(n+p)/2} \int_{\mathbf{R}^{k \times p}} |I_k + WW'|^{-n/2-p} dW$$

Since (e.g. see Dickey (1967), p. 512)

$$\int_{\mathbf{R}^{k \times p}} |I_k + WW'|^{-n/2-p} dW = \pi^{kp/2} \frac{\Gamma_k((n+p)/2)}{\Gamma_k(n/2+p)}$$

we get finally

$$g(B) = \pi^{-kp/2} \frac{\Gamma_k((n+p)/2)}{\Gamma_k(n/2)} |I_k + BB'|^{-(n+p)/2}. \quad (15)$$

This is the density of the central standardized matrix t distribution $t_{n,k,p}$. By definition $T = J^{-1}C \sim t_{n,k,p}$ if $C \in \mathbf{R}^{k \times p}$ and $J \in \mathbf{R}^{k \times k}$ are independent with C standard normal and $JJ' = \sum_1^n u_j u_j' \sim \chi_{n,m}^2 = \chi_{n,m}^2(I_m)$, the central standard Wishart distribution with n degrees of freedom and dimension m (i.e. $u_j \sim N_m(0, \Sigma)$ with $\Sigma = I_m$ and independent). For details see Johnson and Kotz (1972), p. 151.

Of course in this particular case there is a more direct way in showing that $\hat{B} \sim t_{n,k,p}$. Since $\text{vec}(\hat{B}) = \text{vec}((X_2'X_2)^{-1}X_2'X_1) = (I_p \otimes (X_2'X_2)^{-1}X_2') \text{vec}(X_1)$ we have $\mathcal{L}(\text{vec}(\hat{B})|X_2) = N_{kp}(0, I_p \otimes (X_2'X_2)^{-1})$. So $\mathcal{L}(\text{vec}((X_2'X_2)^{1/2}\hat{B}|X_2) = \mathcal{L}(I_p \otimes (X_2'X_2)^{1/2} \text{vec}(\hat{B})|X_2) = N_{kp}(0, I_{kp})$. Hence, $J = (X_2'X_2)^{1/2}$ and $C = (X_2'X_2)^{1/2}\hat{B}$ are independent with C standard normal and $JJ' = X_2'X_2 \sim \chi_{n,k}^2$. Then, by definition, $\hat{B} = J^{-1}C \sim t_{n,k,p}$.

We conclude this section with the derivation of (13). Let $\hat{W} \sim \chi_{n,k}^2(\Sigma)$. Then the characteristic function $\varphi(U) = E\{\text{etr}(iU'\hat{W})\}$, $U \in \mathbf{R}^{k \times k}$ is given by

$$\varphi(U) = \{|\Sigma^{-1}|/|\Sigma^{-1} - i(U + U')|\}^{n/2}$$

and the density $f(W)$ for $W > 0$ by

$$f(W) = K(n, \Sigma^{-1}) |W|^{(n-k-1)/2} \text{etr}(-\frac{1}{2}\Sigma^{-1}W),$$

where

$$K(n, \Sigma^{-1}) = \{2^{nk/2} \Gamma_k(n/2)\}^{-1} |\Sigma|^{n/2}$$

(see Johnson and Kotz (1972), p. 162-163).

It follows from (12) that $|\Sigma|^{n/2}\psi(U_{21}, U_{22})$ is the characteristic function of $\chi_{n,k}^2(\Sigma)$ at U_{22} if $\Sigma = (I_k + U_{21}U'_{21})^{-1}$. Hence,

$$|\Sigma|^{n/2}\psi(U_{21}, U_{22}) = \int K(n, \Sigma^{-1})|W|^{(n-k-1)/2} \text{etr} \left(-\frac{1}{2}\Sigma^{-1}W + iU'_{22}W \right) dW$$

and so, using (6),

$$\begin{aligned} |\Sigma|^{n/2}D_{22}^p\psi(U_{21}, U_{22}) &= i^{kp} \int K(n, \Sigma^{-1})|W|^{(n-k-1)/2+p} \text{etr} \left(-\frac{1}{2}\Sigma^{-1}W + iU'_{22}W \right) dW = \\ &= i^{kp} K(n, \Sigma^{-1})/K(n+2p, \Sigma^{-1})|\Sigma|^{n/2+p}|\Sigma^{-1} - i(U_{22} + U'_{22})|^{-n/2-p} \end{aligned}$$

or

$$D_{22}^p\psi(U_{21}, U_{22}) = (2i)^{kp} \frac{\Gamma_k(n/2+p)}{\Gamma_k(n/2)} |\Sigma^{-1} - i(U_{22} + U'_{22})|^{-n/2-p}$$

This proves (12).

3 The particular case $p = 1, k = 2$

For this most simple non-trivial case explicit expressions for $D_{22}^p\psi = D_{22}\psi$ are manageable. Write $U_{22} = \{u_{rs}\}$, $r, s = 1, 2$. Then

$$D_{22}\psi = \left\{ \det \begin{bmatrix} \partial/\partial u_{11} & \partial/\partial u_{12} \\ \partial/\partial u_{21} & \partial/\partial u_{22} \end{bmatrix} \right\} \psi = (\partial^2/\partial u_{11}\partial u_{22} - \partial^2/\partial u_{21}\partial u_{12})\psi.$$

The lemma below gives an explicit expression for $\partial\psi/\partial u_{rs}$ and $\partial^2\psi/\partial u_{rs}\partial u_{pq}$.

Lemma

Let

$$J_{rs} = \frac{1}{2}\partial/\partial u_{rs}(V + V'), \tag{16}$$

then

$$\partial\psi/\partial u_{rs} = ia_{rs}\psi \quad (17)$$

with

$$a_{rs} = \text{tr}(\Phi^{-1}I_{rs}) + 2\mu'\Omega^{-1}\Phi^{-1}J_{rs}\Phi^{-1}\Omega^{-1}\mu. \quad (18)$$

Furthermore, let

$$\Phi_{rs pq} = J_{rs}\Phi^{-1}J_{pq} + J_{pq}\Phi^{-1}J_{rs}, \quad (19)$$

then

$$\partial^2\psi/\partial u_{rs}\partial u_{pq} = -(a_{rs pq} + a_{rs}a_{pq})\psi \quad (20)$$

with

$$a_{rs pq} = \text{tr}(\Phi^{-1}\Phi_{rs pq}) + 4\mu'\Omega^{-1}\Phi^{-1}\Phi_{rs pq}\Phi^{-1}\Omega^{-1}\mu. \quad (21)$$

Proof

We have

$$\begin{aligned} \partial|\Phi|^{-1/2}/\partial u_{rs} &= -\frac{1}{2}|\Phi|^{-3/2}\partial|\Phi|/\partial u_{rs} = \\ &= -\frac{1}{2}|\Phi|^{-3/2} \cdot |\Phi| \text{tr}\{\Phi^{-1}\partial\Phi/\partial u_{rs}\} = i|\Phi|^{-1/2} \text{tr}(\Phi^{-1}J_{rs}) \end{aligned}$$

$$\partial\Phi^{-1}/\partial u_{rs} = -\Phi^{-1}(\partial\Phi/\partial u_{rs})\Phi^{-1} = 2i\Phi^{-1}J_{rs}\Phi^{-1}.$$

Writing $c = |\Omega|^{1/2} \exp(-\frac{1}{2}\mu'\Omega^{-1}\mu)$ this gives

$$\begin{aligned} \partial\psi/\partial u_{rs} &= ci\{|\Phi|^{-1/2} \text{tr}(\Phi^{-1}J_{rs}) \cdot \exp(\mu'\Omega^{-1}\Phi^{-1}\Omega^{-1}\mu) + \\ &\quad + |\Phi|^{-1/2}(2i\mu'\Omega^{-1}J_{rs}\Phi^{-1}\Omega^{-1}\mu) \cdot \exp(\mu'\Omega^{-1}\Phi^{-1}\Omega^{-1}\mu)\} = ia_{rs}\psi, \end{aligned}$$

proving (18). Furthermore,

$$\begin{aligned}\partial \operatorname{tr}(\Phi^{-1} J_{rs}) / \partial u_{pq} &= \operatorname{tr}\{(\partial \Phi^{-1} / \partial u_{pq}) J_{rs}\} = \\ &= 2i \operatorname{tr}(\Phi^{-1} J_{pq} \Phi^{-1} J_{rs}) = i \operatorname{tr}(\Phi^{-1} \Phi_{rspq}) \\ \\ \partial \Phi^{-1} J_{rs} \Phi^{-1} / \partial u_{pq} &= 2i \Phi^{-1} J_{pq} \Phi^{-1} J_{rs} \Phi^{-1} + 2i \Phi^{-1} J_{rs} \Phi^{-1} J_{pq} \Phi^{-1} = \\ &= 2i \Phi^{-1} \Phi_{rspq} \Phi^{-1}\end{aligned}$$

Therefore, using (18),

$$\begin{aligned}\partial^2 \psi / \partial u_{rs} \partial u_{pq} &= i\{(\partial a_{rs} / \partial u_{pq}) \psi + a_{rs} (\partial \psi / \partial u_{pq})\} = \\ &= i\{i \operatorname{tr}(\Phi^{-1} \Phi_{rspq}) + 2\mu' \Omega^{-1} (2i \Phi^{-1} \Phi_{rspq} \Phi^{-1}) \Omega^{-1} \mu + i a_{rs} a_{pq}\} \psi = \\ &= -(a_{rspq} + a_{rs} a_{pq}) \psi,\end{aligned}$$

proving (21).

4 Proof of the theorem

The proof is given with some preparatory lemma's.

Lemma 1

Let $X \sim N_n(\mu, \Omega)$, $\Omega > 0$. Set $Y_j = X' A_j X$ with symmetric $A_j \in \mathbb{R}^{n \times n}$, $j = 1, \dots, s$ and $Y = (Y_1, \dots, Y_s)'$. Then for $u = (u_1, \dots, u_s)'$:

$$\begin{aligned}\varphi(u) &= E\{\exp(iu'Y)\} = \\ &= |\Omega|^{-1/2} \exp(-\frac{1}{2}\mu' \Omega^{-1} \mu) \cdot |\Phi|^{-1/2} \exp(\frac{1}{2}\mu' \Omega^{-1} \Phi^{-1} \Omega^{-1} \mu)\end{aligned}\tag{22}$$

with

$$\Phi = \Phi(u) = \Omega^{-1} - 2i \sum_{j=1}^s u_j A_j.\tag{23}$$

Proof

Follows after a reformulation of Magnus (1986), Lemma 5, p. 102.

Lemma 2

Let $X \in \mathbf{R}^{n \times m}$ with $\text{vec}(X) \sim N_{nm}(\mu, \Omega)$, $\Omega > 0$. Set $Y = X'AX$ with symmetric $A \in \mathbf{R}^{n \times n}$. Then for $U \in \mathbf{R}^{m \times m}$:

$$\begin{aligned} \varphi(U) &= E\{\text{etr}(iU'Y)\} = \\ &= |\Omega|^{-1/2} \exp(-\frac{1}{2}\mu'\Omega^{-1}\mu) \cdot |\Phi|^{-1/2} \exp(\frac{1}{2}\mu'\Omega^{-1}\Phi^{-1}\Omega^{-1}\mu) \end{aligned} \quad (24)$$

with

$$\Phi = \Phi(U) = \Omega^{-1} - i(U + U') \otimes A \quad (25)$$

Proof

Write $X = [x_1 \dots x_m]$, $x = \text{vec}(X) = [x'_1 \dots x'_m]'$, $Y = (Y_{ij})$. Then $Y_{ij} = (X'AX)_{ij} = x'_i A_{ij} x_j$ with $A_{ij} \in \mathbf{R}^{n \times n}$ a matrix with A for block (i, j) and zeros elsewhere. Since for $U = (u_{ij})$:

$$\text{tr}(U'Y) = \sum \sum u_{ij} Y_{ij} = \sum \sum u_{ij} x'_i A_{ij} x_j = \frac{1}{2} \sum \sum u_{ij} x'_i (A_{ij} + A_{ji}) x_j$$

it follows from (22), (23) with $s = m^2$ that $\varphi(u)$ is given by (24) with

$$\begin{aligned} \Phi &= \Phi(U) = \Omega^{-1} - i\{\sum \sum u_{ij} (A_{ij} + A_{ji})\} = \\ &= \Omega^{-1} - i\{\sum \sum (u_{ij} + u_{ji}) A_{ij}\} = \Omega^{-1} - i(U + U') \otimes A, \end{aligned}$$

proving (25).

Lemma 3

Let $C_{21} \in \mathbf{R}^{k \times p}$, $C_{22} \in \mathbf{R}^{k \times k}$ with $C_{22} > 0$ a.s. If $E|C_{22}|^p < \infty$ and if $[C_{21} \ C_{22}]$ is absolutely continuous (with respect to C_{22} in the obvious sense) then $\hat{B} = C_{22}^{-1}C_{21}$ has a continuous density g , for $B \in \mathbf{R}^{k \times p}$ given by

$$g(B) = \left(\frac{1}{2\pi i}\right)^{kp} \int_{\mathbf{R}^{k \times p}} D_{22}^p \psi(U_{21}, -(BU'_{21} + U_{21}B')/2) dU_{21} \quad (26)$$

provided that the integral is absolutely convergent.

Proof

The Lemma is a slight correction of Phillips (1985), theorem, p. 185.

Proof of the theorem

Follows from Lemma 2 and Lemma 3. Take $A = I_n$ and

$$U = \begin{bmatrix} 0 & 0 \\ U_{21} & U_{22} \end{bmatrix}.$$

Then combination of (24)-(26) leads to (8)-(11).

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