

Low-Loss and High-Voltage III-Nitride Transistors for Power Switching Applications

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(*Invited Paper*)

Abstract—This paper describes recent technological advances on III-nitride-based transistors for power switching applications. Focuses are placed on the progress toward enhancing the breakdown voltage, lowering the ON-resistance, suppressing current collapse, and reducing the leakage current in AlGaIn/GaN high-electron mobility transistors (HEMTs). Recent publications revealed that the tradeoff relation between ON-resistance and breakdown voltage in AlGaIn/GaN HEMTs exceeded the SiC limit and was getting close to the GaN limit; however, the breakdown voltage achieved was still lower than the theoretical impact ionization limit. A novel process featuring strain-controlled annealing with a metal stack, including Al gave rise to significant reduction in the sheet resistance in AlGaIn/GaN heterostructures, suggesting the possibility of dramatic reduction in ON-resistance of GaN-based power devices. Some of the interesting approaches to suppress current collapse indicated that surface trapping effects must be controlled by the optimization of surface processing as well as by the reduction of bulk traps in the epitaxial layers. Close correlation between the local gate leakage current and point defects exposed on the free-standing GaN substrate demonstrated that further reduction of defects on bulk GaN substrates is truly required as future challenges.

Index Terms—Breakdown voltage, current collapse, field plate (FP), GaN, gate leakage, high-electron mobility transistor (HEMT), ON-resistance.

I. INTRODUCTION

III-NITRIDE-BASED transistors represented by AlGaIn/GaN high-electron mobility transistors (HEMTs) are promising as low-loss and high-voltage switching devices to be utilized for a variety of power conversion circuits. Since GaN is a material having a wide bandgap of 3.4 eV with a direct transition band structure, it inherently shows a high critical electric field, low intrinsic carrier concentration, high-electron mobility, and high-drift velocity. In addition, capability of growing high-quality heterojunctions, such as AlGaIn/GaN and InAlN/GaN, has enabled us to achieve high-density 2-D electron gas (2-DEG) of $>10^{13}$ cm $^{-2}$ with a high-electron mobility of exceeding 2000 cm 2 /V s.

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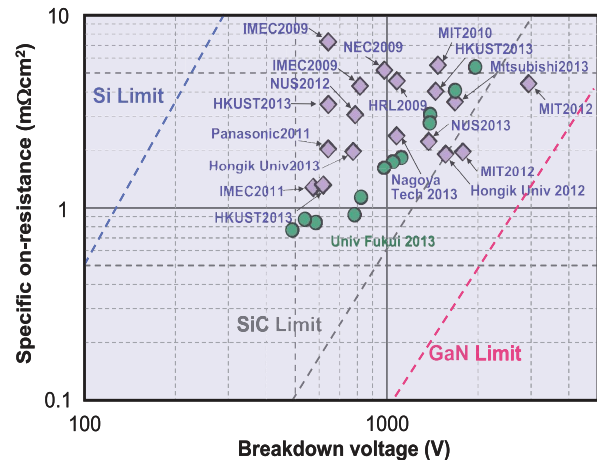


Fig. 1. Present status of specific ON-resistance versus breakdown voltage. Referred data can be found in [1]–[18].

Great efforts have been made over the last decade to improve the tradeoff relationship between breakdown voltage and specific ON-resistance of nitride-based transistors. In AlGaIn/GaN HEMTs, the progress has been achieved by optimizing the device geometry, reducing the defect density in the epitaxial layers, improving the quality of buffer layers, and introducing newly developed process technologies. Fig. 1 shows reported values of specific ON-resistance as a function of breakdown voltage for GaN-based HEMTs, including our developed AlGaIn/GaN HEMTs with varied gate-to-drain spacing [1]–[18]. It is evident that the power switching capability of AlGaIn/GaN HEMTs is already far beyond the Si limit, partly exceeds the SiC limit, and is getting close to the GaN limit.

In this paper, technological advances on AlGaIn/GaN HEMTs have been described, including our recent results for improving dc and pulsed performance of AlGaIn/GaN HEMTs. Focuses are placed on the progress toward enhancing the breakdown voltage, lowering the ON-resistance, suppressing current collapse, and reducing the leakage current by introducing free-standing GaN substrates. Section II discusses experimental and simulated results performed to increase breakdown voltages of AlGaIn/GaN HEMTs. Section III presents several technologies to reduce the ON-resistance of AlGaIn/GaN HEMTs by lowering contact resistances and access resistances. In Section IV, experimental attempts are presented for the suppression of current collapse. Section V reports the advantage of using free-standing GaN substrate to reduce gate leakage current in AlGaIn/GaN HEMTs. Finally, the conclusion is drawn in Section VI.

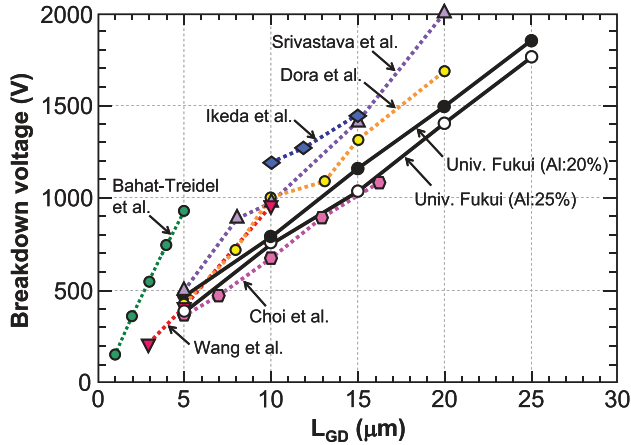


Fig. 2. Breakdown voltage plotted as a function of gate-to-drain distance. Referred data can be found in [19]–[24].

II. BREAKDOWN VOLTAGE

A large number of experimental results have been reported on the measured three-terminal breakdown voltage for AlGaIn/GaN HEMTs. Fig. 2 shows examples of measured results plotted as a function of the distance between gate and drain (L_{gd}) for AlGaIn/GaN HEMTs. The results indicate that the breakdown voltage is linearly proportional to L_{gd} in the breakdown voltage range up to at least 2000 V [19]–[24]. The averaged breakdown electric field ranges from 0.6 to 1.6 MV/cm [10]–[13], [25], which is still much lower than the expected impact ionization limit of 3 MV/cm. In general, the electric field concentrates at the gate edge in the drain side, as calculated by the 2-D computer simulation [26]. Such nonuniform electric field distribution does not guarantee the linear increase in the breakdown voltage, but is more likely to exhibit gradually saturated characteristics. Thus, the results shown in Fig. 2 imply that the lateral electric field along gate-to-drain direction is not concentrated at the gate edge, but more widely distributed or almost constant [27], [28].

Provided that magnitudes of positive and negative polarization charges (including spontaneous and piezoelectric charges), located at the interface and the surface of the AlGaIn barrier layer, respectively, are mutually balanced, more uniform lateral electric field distributions could be accomplished along gate-to-drain. By applying strong drain voltages, 2-DEG charges in the channel are almost totally depleted, accompanied by the reduction in the ionized surface donor charge [29]. Hence, remaining charges between gate and drain would be a pair of balanced positive and negative polarization charges, resulting in constant electric field distribution, which is similar to the superjunction concept used in Si MOSFETs [30] and GaN-based diodes [31].

Fig. 3(a) and (b) shows the electric field distribution between gate and drain, based on 2-D computer simulation by assuming balanced polarization charges and depletion of 2-DEG charges. When positive and negative polarization charges are balanced, the electric field distribution in the channel becomes almost constant throughout the gate-to-drain region regardless of drain bias voltages assumed. However, since passivation films and/or GaN cap layer are formed in the

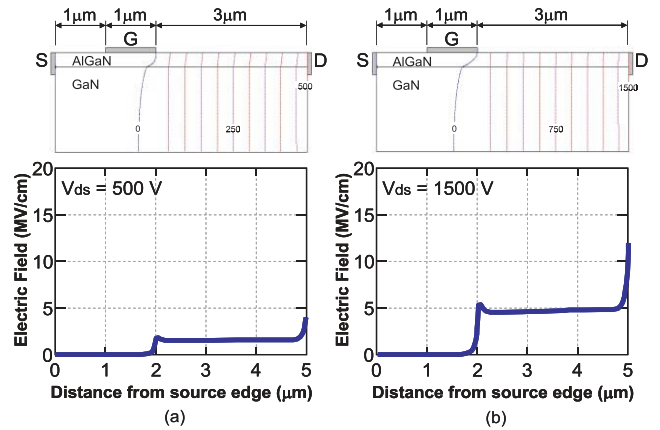


Fig. 3. Simulated electric field distribution between gate and drain in AlGaIn/GaN HEMT at drain voltage of (a) 500 and (b) 1500 V.

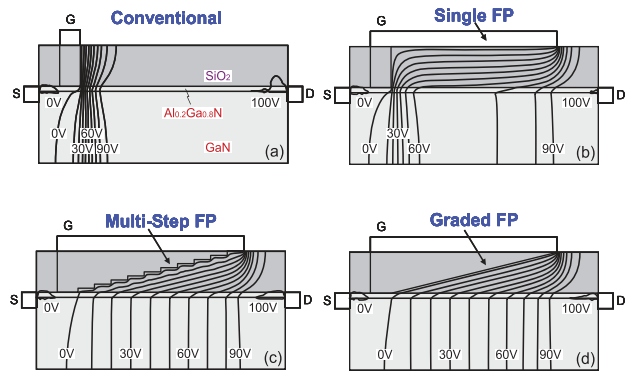


Fig. 4. Simulated equipotential contour lines in AlGaIn/GaN HEMTs for (a) without FP, (b) single FP, (c) multistep FP, and (d) graded FP. The applied drain voltage is 100 V.

actual AlGaIn/GaN HEMTs, further study is needed whether the net surface negative charges balance with the positive polarization ones at the AlGaIn/GaN interface. It is also an open question why the experimentally estimated lateral critical electric field is much lower than the theoretical limit. To solve such discrepancy, it would be of importance to investigate undesirable leakage current mechanisms related to the high density of crystal defects.

In addition, to make a gate-to-drain spacing longer, introducing field plate (FP) is also an effective way for increasing the breakdown voltage. The validity of FP has been recognized since early 2000s [32]–[34], and HEMTs with multi-FP were also reported [35], [36]. The simulation results on how the electric field distribution are modified by the number of FP steps are shown in Fig. 4(a)–(d), where equipotential contour lines are depicted. By increasing the number of step [Fig. 4(b) and (c)], the peak electric field along the gate-to-drain direction is decreased and the field distribution spreads, resulting in improved breakdown characteristics. The most suitable structure is the graded FP [Fig. 4(d)], although a special process is required for fabricating such a structure [19], [37].

It is known that AlGaIn/GaN HEMTs tend to suffer from a high-gate leakage current originating from the leaky

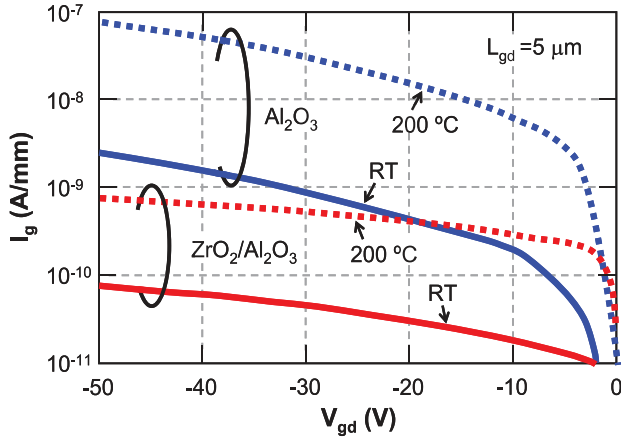


Fig. 5. Two-terminal reverse gate I - V characteristics of AlGaIn/GaN MIS-HEMTs with gate insulator of composite ZrO_2/Al_2O_3 and single Al_2O_3 measured at room temperature and 200 °C.

Schottky barrier at the metal–semiconductor interface. One of the effective ways to suppress such a high-gate leakage is the use of metal–insulator–semiconductor (MIS) structure for the gate. As a gate insulator, material properties of high resistivity, high-breakdown field, low-interface state density, and high permittivity are required. A large number of insulators, such as SiO_2 , SiN , Al_2O_3 , ZrO_2 , HfO_2 , and AlN , have been developed for fabricating MIS-HEMTs [38]–[43], and have been deposited with a variety of methods, including chemical vapor deposition (CVD), atomic layer deposition (ALD), sputtering, electron beam evaporation, and *in situ* SiN growth by metal–organic CVD [44]–[48]. Using ALD, we have developed a ZrO_2/Al_2O_3 double layered gate insulator with Al_2O_3 being the bottom and have shown that the composite insulator is advantageous in terms of low leakage current even at high temperatures. As shown in Fig. 5, a composite ZrO_2/Al_2O_3 insulator resulted in a reduced gate leakage current by two orders of magnitude at 200 °C, as compared with those with a single-layered insulator of ZrO_2 or Al_2O_3 [49].

III. ON-RESISTANCE

Lowering the ON-resistance is an important challenge for achieving high-efficiency operation in power switching circuits with AlGaIn/GaN HEMTs. The ON-resistance between source and drain is expressed as the sum of contact resistance of source/drain electrodes, access resistance between source and gate, channel resistance beneath the channel, and access resistance between gate and drain.

Various kinds of metal stacks and corresponding annealing conditions have been attempted to obtain lower contact resistances in AlGaIn/GaN HEMTs [50]–[54]. The commonly used metal stacks are $Ti/Al/Ni/Au$ and $Ti/Al/Mo/Au$, and a contact resistance of 0.2–0.5 Ωmm has been typically reported. Gold-free ohmic metals have also been developed to ensure compatibility with the Si fabrication process [15], [55], [56]. Although good ohmic contact behaviors are obtained for an AlGaIn barrier with relatively a low-Al composition of 0.15–0.3, it becomes rather difficult to obtain a reasonable contact resistance for AlGaIn barriers with a high-Al composition

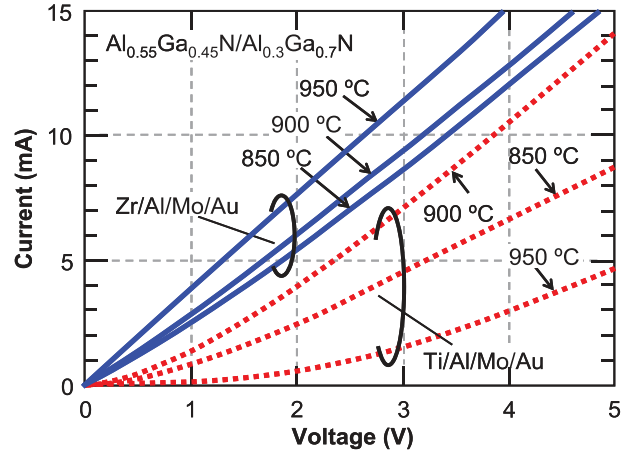


Fig. 6. I - V characteristics of $Al_{0.55}Ga_{0.45}N/Al_{0.3}Ga_{0.7}N$ HEMT with ohmic metals of $Zr/Al/Mo/Au$ and $Ti/Al/Mo/Au$ annealed at different temperatures.

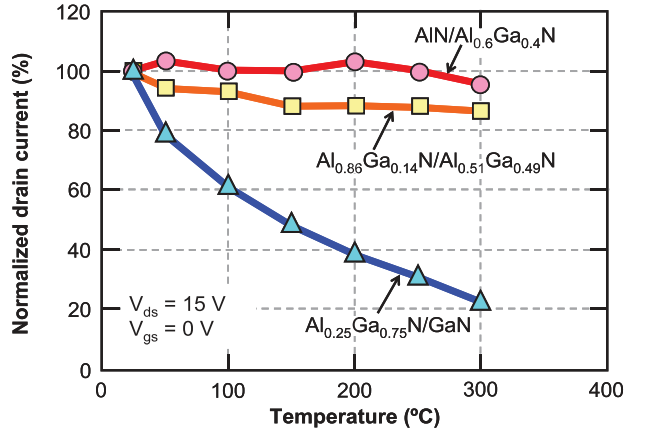


Fig. 7. Temperature dependences of drain current for AlGaIn-channel HEMTs ($AlN/Al_{0.6}Ga_{0.4}N$ and $Al_{0.86}Ga_{0.14}N/Al_{0.51}Ga_{0.49}N$) and conventional AlGaIn/GaN HEMT.

of >0.5 . For such HEMTs having a barrier layer with high-Al composition, Yafune *et al.* [57] have introduced a new metal stack of $Zr/Al/Mo/Au$ and reported a relatively low value in contact resistance. Fig. 6 shows current–voltage (I - V) characteristics measured between ohmic metals on $Al_{0.55}Ga_{0.45}N/Al_{0.3}Ga_{0.7}N$ HEMTs, in which ohmic metal stacks composed of $Zr/Al/Mo/Au$ and $Ti/Al/Mo/Au$ were annealed at 850 °C, 900 °C, and 950 °C. Almost linear I - V characteristics were observed for $Zr/Al/Mo/Au$ annealed at 950 °C, while only nonlinear characteristics were noticed for $Ti/Al/Mo/Au$. The reason for the better ohmic characteristics with $Zr/Al/Mo/Au$ is that metals (presumably Zr and Al) in $Zr/Al/Mo/Au$ penetrate more deeply into the $Al_{0.55}Ga_{0.45}N$ barrier layer than those in $Ti/Al/Mo/Au$, as reported in [57].

Two kinds of AlGaIn-channel HEMTs, composed of $AlN/Al_{0.6}Ga_{0.4}N$ and $Al_{0.86}Ga_{0.14}N/Al_{0.51}Ga_{0.49}N$, have been fabricated using $Zr/Al/Mo/Au$ as ohmic electrodes [58], [59]. Fig. 7 shows the temperature dependence of drain current measured at temperatures up to 300 °C. The drain current was normalized with the values measured at room temperature. It is evident that the AlGaIn-channel with a higher Al composition

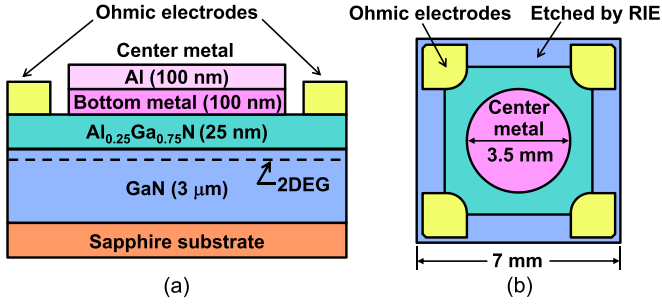


Fig. 8. (a) Cross-sectional and (b) top view of samples for measuring sheet electron density and mobility.

(>0.5) exhibits extremely stable drain current behaviors with respect to the device temperature up to 300 °C. Although further improvements, such as lowering the contact resistance are needed, the AlGaIn-channel HEMT is promising in terms of stable operation at elevated temperatures with sufficiently high-breakdown voltages [14].

Access resistances of HEMTs are determined by the product of 2-DEG density (n_s) and electron mobility (μ). Since those values measured at room temperature in AlGaIn/GaN heterostructures are typically in the range of $8\text{--}10 \times 10^{12} \text{ cm}^{-2}$ and $1500\text{--}2000 \text{ cm}^2/\text{Vs}$, respectively, the corresponding sheet resistance turns out to be $\sim 300\text{--}500 \text{ }\Omega/\text{sq}$. If one could increase both n_s and μ , the sheet resistance would be significantly reduced, leading to extremely low-access resistance. Regarding the 2-DEG mobility at low temperatures, theoretical and experimental studies have been extensively made to understand the dominant transport mechanisms at the AlGaIn/GaN interface. As a result, a mobility of $7500 \text{ cm}^2/\text{Vs}$ at 10 K was reported in [60], and it was increased year by year, reaching $167\,000 \text{ cm}^2/\text{Vs}$ at 0.3 K in [61]. Meanwhile, the room temperature mobility in AlGaIn/GaN heterostructure has not been increased over the last decade, even though significant progress has been achieved in epitaxial growth technologies. This is because, the room temperature electron mobility is mainly governed by polar-optical phonon scattering [62], [63], which is almost independent of the epitaxial layer quality. A room temperature electron mobility of $2019 \text{ cm}^2/\text{Vs}$ has been reported in [64], and it was increased to $2200 \text{ cm}^2/\text{Vs}$ in [65], corresponding to an increase by only 9%.

An interesting method to increase room temperature electron mobility is to introduce an additional tensile strain in an AlGaIn layer. Azize and Palacios [66] reported a mobility increase by etching a Si substrate from the backside for controlling the tensile strain in AlGaIn. Fehlberg *et al.* [67] reported a high mobility of $2380 \text{ cm}^2/\text{Vs}$ by depositing SiN films with varied deposition conditions, resulting in introduction of strain in an AlGaIn layer. In addition, Im *et al.* [68] controlled the tensile strain by varying the buffer layer thickness and fabricated an AlGaIn/GaN HEMT with increased n_s of $1 \times 10^{14} \text{ cm}^{-2}$, though the mobility was low.

Our group have reported significant increases in both n_s and μ by annealing an AlGaIn/GaN heterostructure in vacuum with deposited metals composed of Ti/Al or Ni/Al [69]–[71]. Fig. 8 shows cross-sectional and top views of the measured AlGaIn/GaN sample with

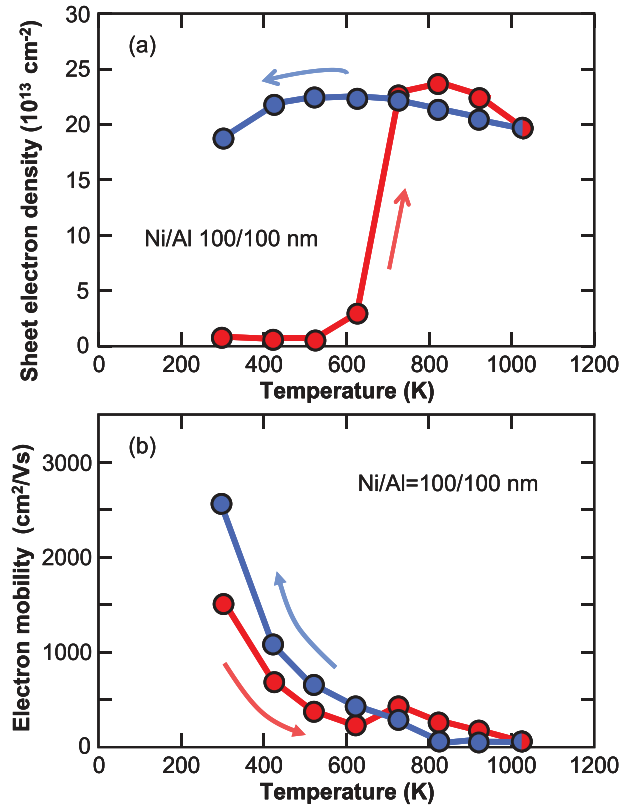


Fig. 9. Temperature dependences of (a) sheet electron density and (b) electron mobility of AlGaIn/GaN heterostructure deposited with Ni/Al. Arrows indicate temperature change directions.

van der Pauw configuration. Ohmic electrodes were formed in the four corners and an additional metal consisting of double layered metal stack, center metal, was evaporated. The sample was set to the Hall measurement system and the temperature dependence of n_s and μ was measured with increasing and decreasing sample temperatures ranging from 300 (room temperature) to 1020 K. All the measurements were carried out in vacuum ($\sim 1 \times 10^{-3}$ torr). An example of the measured temperature dependences in n_s and μ is shown in Fig. 9(a) and (b). With increasing temperature from 300 K, n_s showed a sudden increase at ~ 650 K and had a peak at 820 K followed by a slight decrease. By decreasing temperature from 1020 K, n_s was slightly increased and then decreased. Meanwhile, μ showed a hump at ~ 650 K with increasing temperature from 300 K. It is to be noted that both n_s and μ did not take the same paths in the cycle of heating and cooling, resulting in an increase in n_s by one order of magnitude and an increase in μ by 70% at room temperature. The amount of increased μ depended on the thickness of Ni/Al and a value of over $3000 \text{ cm}^2/\text{Vs}$ was achieved by optimizing the thickness [72].

The microstructure of the annealed metals with Ti/Al (Bottom: Ti, Top: Al) and with Al/Ti (Bottom: Al, Top: Ti) was investigated by Auger electron spectroscopy [70], where the increases in n_s and μ were only observed for Ti/Al and not for Al/Ti. The results indicate that the top Al layer diffuses into the bottom Ti layer for Ti/Al, whereas Ti does not diffuse into Al for Al/Ti, suggesting that the diffusion of top Al layer

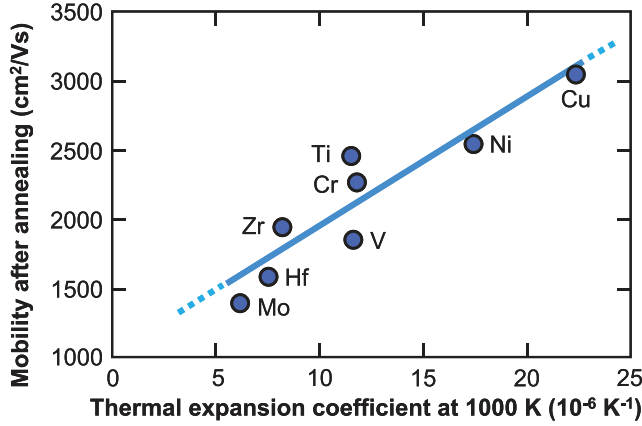


Fig. 10. Room temperature electron mobility of AlGaIn/GaN heterostructures after annealing as a function of thermal expansion coefficient of the bottom metal.

into the bottom metal plays a key role for the increases in n_s and μ . One might think that the measured n_s and μ did not correspond to those of 2-DEG, but those of center metal. However, such possibilities are excluded by the experimental results [69]–[71]. The increases in n_s and μ were not observed when the center metal was a monolayer of Ti or Al, and also the amount of increase in n_s was independent of the total thickness of Ni/Al, indicating that the measured n_s and μ were not dependent on the center metal resistivity. In addition, μ never exceeds $2000 \text{ cm}^2/\text{Vs}$ if electrons flow in the metallic layer.

Further experiments showed that μ also depended on the species of the bottom metal and had a close correlation with the thermal expansion coefficient of the bottom metal, as shown in Fig. 10, where the top metal was Al [73]. It was found that μ was more increased with a metal having larger thermal expansion coefficient and a value of over $3000 \text{ cm}^2/\text{Vs}$ was obtained with Cu/Al.

From a series of experiments, a model was proposed for the increases in n_s and μ by annealing as follows. The center metal, expanded with the temperature increase, induces a tensile strain in AlGaIn/GaN layer, resulting in the increase in n_s due to the increase in piezoelectric charge. The key of this process is that the induced tensile strain brings about inelastic deformation at high temperatures and maintains it after cooling the sample to room temperature. The increase in μ is accompanied by the increase in n_s , presumably due to the reduced effective mass [66], and/or increased electron screening [74]. The larger thermal expansion coefficient of Cu and Ni induces a stronger tensile strain at high temperatures, and thus, higher μ is observed with Cu and Ni. To confirm the validity of the model, Raman microprobe spectroscopy measurements have been performed. An Ar laser with a wavelength of 496.5 nm was used as an incident beam and the Raman shift was measured on the AlGaIn surface close to the Ni/Al metal ($1 \mu\text{m}$ apart from the metal edge). Fig. 11 shows the measured stress for three samples, i.e., bare AlGaIn surface without Ni/Al, as-deposited Ni/Al (100/100 nm), and annealed Ni/Al (100/100 nm) at 1020 K. Note that the stress

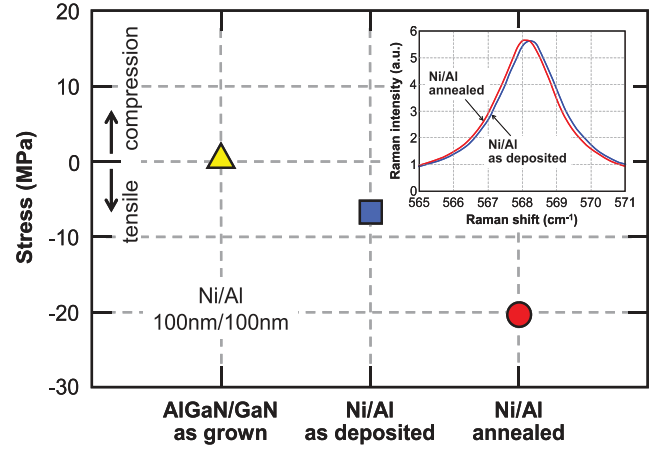


Fig. 11. Stress values measured by Raman spectroscopy for AlGaIn/GaN bare surface, as-deposited Ni/Al (100/100 nm), and annealed Ni/Al (100/100 nm). Inset: Raman signal.

value is plotted as the deviation from that for the bare surface. The increased tensile stress was 7 and 20.6 MPa for samples of as-deposited Ni/Al and annealed Ni/Al, respectively, indicating the evidence that an additional tensile stress was introduced by metal deposition and annealing.

The thermal annealing technique with a metal stack mentioned above is effective to increase both n_s and μ , and hence to reduce sheet resistance drastically. For example, the sheet resistance of AlGaIn/GaN heterostructure after annealing with Ni/Al was as low as $10 \Omega/\text{sq}$, which is fifty times lower than that before annealing. The decreased sheet resistance will be useful to reduce the access resistance between source and gate, leading to improvement in transconductance and reduction in ON-resistance in AlGaIn/GaN HEMTs.

IV. CURRENT COLLAPSE

It is widely known that applying high drain biases into GaN-based HEMTs frequently gives rise to the decrease in drain current accompanied by the increase in the dynamic ON-resistance. This undesirable phenomenon is known as current collapse and must be eliminated before pushing GaN-based HEMT devices into market. Since the pioneering work on the fabrication of AlGaIn/GaN HEMTs [75], extensive efforts have been made to suppress current collapse and to elucidate the origin of it. It is widely recognized that current collapse can be effectively mitigated by surface passivation, especially with a SiN film [76]. The proper surface treatment before depositing surface passivation films is also reported to be effective to reduce current collapse. One such example is the surface treatment by oxygen plasma exposure. The results are shown in Fig. 12, where the dynamic ON-resistance of AlGaIn/GaN HEMTs after oxygen plasma treatment with an RF power of 100 W for 60 s is dramatically reduced after passivation with either SiN, SiO₂, or AlN [77]. The oxygen termination of unbonded Ga and Al atoms near the AlGaIn surface may be responsible for the reduced current collapse, and the evidence of reduced trap levels after oxygen plasma exposure has been identified by analyzing transients

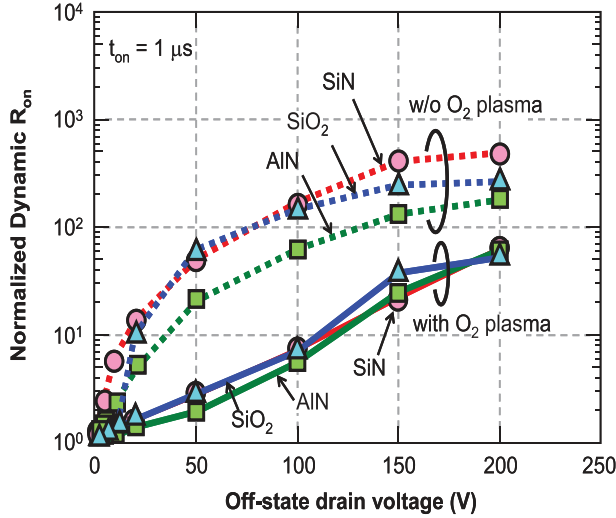


Fig. 12. Dynamic ON-resistance as a function of OFF-state drain voltage for devices with and without O₂ plasma treatment passivated with SiN, SiO₂, and AlN.

of ON-state drain current assuming Shockley-Read-Hall statistics [77].

The use of FP is another way for the suppression of current collapse. The comprehensive study on the effect of FP on the dynamic ON-resistance has been performed by preparing a series of AlGaIn/GaN HEMTs with different FP lengths and gate-to-drain distance [78]. The results indicated that current collapse was dramatically improved by the introduction of gate FP, and the improvement was more enhanced using a longer FP length. It was also found that the current collapse reduction was more pronounced when the gate bias during ON-state was chosen at more positive values, suggesting that the gate FP is capable of instantly supplying additional electrons into the channel access region [78].

Although current collapse is recognized as closely related to carrier trapping phenomenon, its origin is not yet fully understood. Since current collapse was found to be affected by surface passivation, surface treatment, and the use of FP, it must be at least ascribed to trapping and detrapping of carriers at the semiconductor surface. However, several reports also suggested that the collapse was likely to be brought about from bulk traps in the buffer layer [79], [80]. Therefore, more work is definitely needed for the elimination of current collapse, including further optimization of surface processing and the reduction of bulk traps in the epitaxial layers.

V. FREE-STANDING GaN SUBSTRATE

AlGaIn/GaN heterostructures have been commonly grown on various substrates, such as sapphire, SiC, and Si. However, a large lattice mismatch between GaN and a foreign substrate leads to generation of high density of threading dislocations on the orders of 10^8 – 10^{10} cm⁻² in the epitaxial layer. Hence, the development of free-standing GaN substrates with low-dislocation densities becomes extremely important. Several groups have reported that a free-standing GaN substrate grown by hydride vapor phase epitaxy (HVPE) is effective to obtain ideal Schottky or p-n diode characteristics [81], [82].

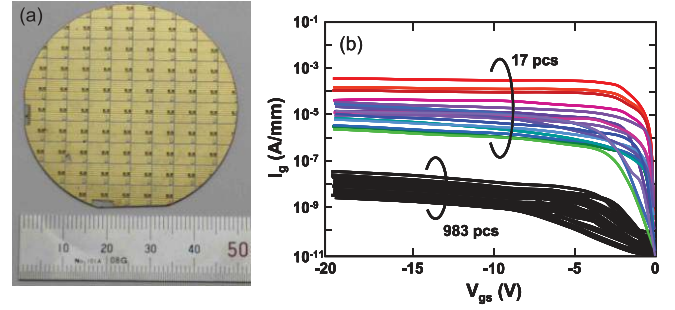


Fig. 13. (a) Top view image of AlGaIn/GaN HEMTs fabricated on 2-in diameter Na-flux bulk GaN substrate. (b) Gate leakage characteristics for 1000 HEMTs.

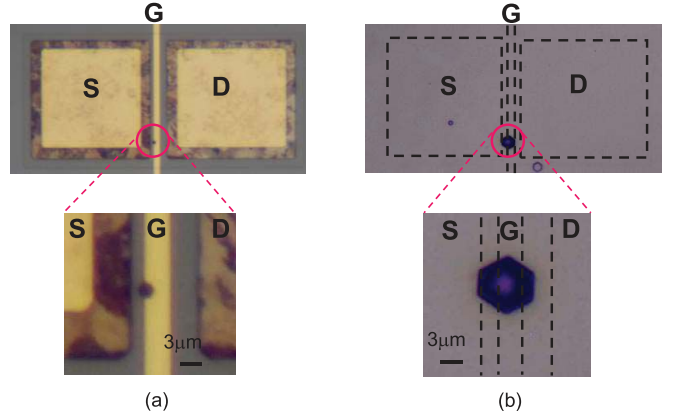


Fig. 14. Top view images of (a) device with metal electrodes and (b) device after etching off metal electrodes.

Although the dislocation density in the epitaxial layers can be significantly reduced by the use of HVPE-grown GaN substrate, the current dislocation density still stays as high as 10^6 cm⁻², and further reduction is required. A melt-grown method, such as Na-flux liquid phase epitaxy, has recently attracted significant attention to further reduce the dislocation density in a bulk GaN substrate [83], [84].

Using a free-standing GaN substrate with a dislocation density of less than 10^6 cm⁻² grown by Na-flux method, our group have studied the detailed correlation between the gate leakage current in AlGaIn/GaN HEMTs and point defects in the starting GaN substrate [85]. The epitaxial layers consisted of an undoped GaN channel layer and a 25-nm undoped AlGaIn barrier layer with an Al composition of 0.25. Fig. 13(a) shows a top view of the whole 2-in wafer after device fabrication.

The statistical data of the reverse gate leakage current were taken from randomly selected 1000 devices fabricated on a same wafer with a gate length of 3 μm. As shown in Fig. 13(b), there observed a large scatter by more than five orders of magnitude in the reverse gate leakage characteristics measured at a gate voltage of -20 V. The leakage characteristics were divided into two groups showing low- and high-gate leakage currents, i.e., almost all the devices (983 pcs.) showed a leakage current of $< 10^{-7}$ A/mm, whereas only 17 devices exhibited rather high leakage of $> 10^{-6}$ A/mm.

After electrical measurements, device areas were immersed in hot H_3PO_4 solution for 60 min to expose the surface of the GaN substrate. As shown in Fig. 14(a) and (b), there was a clear correspondence between locations of the gate electrode and the exposed hexagonal etch-pit defect with a size of several micrometers, indicating that the gate electrode of the leaky HEMT was formed in direct contact with the location of the etch-pit defect. These results suggest that the reverse gate leakage characteristics of GaN-based HEMTs are strongly influenced by the density of threading dislocations propagating through the substrate material, and thus continued efforts are of particular importance to reduce the density of point defects, and/or threading dislocations in free-standing GaN substrates.

VI. CONCLUSION

Technological advances for increasing the breakdown voltage, reducing the ON-resistance, suppressing current collapse, and reducing the gate leakage current in AlGaIn/GaN HEMTs have been described. Comprehensive characterization on breakdown voltages in AlGaIn/GaN HEMTs suggested that balancing of net positive and negative charges in the AlGaIn barrier was likely to account for the experimentally observed linear dependence of breakdown voltage on the gate-to-drain distance. However, the measured lateral critical electric field was still much lower than the theoretical limit. Attempts to achieve lower contact resistances for an AlGaIn-channel HEMT indicated that a new metal stack of Zr/Al/Mo/Au was effective and resulted in improved device stability at temperatures up to 300 °C. Extremely, low sheet resistance as well as high room temperature mobility exceeding 2500 $\text{cm}^2/\text{V}\cdot\text{s}$ were achieved for an AlGaIn/GaN heterostructure deposited with a metal stack, such as Ti/Al, Ni/Al, and Cu/Al. A model was proposed in which inelastic tensile strain plays a key role. Effects of oxygen plasma treatment and providing FP were presented to reduce current collapse in AlGaIn/GaN HEMTs and their mechanisms were discussed. The results strongly suggested the need for essential reduction in trapping states located on the surface or the bulk of the semiconductors. Finally, issues associated with the high-defect density in currently available AlGaIn/GaN epitaxial layers are presented and the strong demand for the development of low-defect density bulk GaN substrates were pointed out.

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