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Exploring fusion at extreme sub barrier energies with weakly bound nuclei

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Results of measurement of residues formed in fusion of 6Li with 198Pt in the energy range of 0.68<E/Vb<1.3 using a new sensitive off-beam technique are reported. The fusion excitation function and the derived average angular momenta do not indicate a change of slope at deep sub-barrier energies, contrary to recent observations. The present results for a system with weakly bound projectile confront the current understanding of the fusion hindrance at these low energies, underlying the role of internal re-organization on the dynamical path towards fusion.

Nuclear reactions around the Coulomb barrier are a vast reservoir for signatures of various aspects of basic quantum mechanics. The tunneling phenomena, in particular, can be probed under diverse conditions ranging from the effect of dissipation in heavy systems to the role of pairing. These unique features, in addition to the effects of inter-connectivity of intrinsic properties in the entrance channel on different processes, provide insights into various quantum mechanical effects [1]. In the last few years improved sensitivity in the measurements have challenged the understanding of the mechanism of tunneling through multidimensional barriers. The latest addition to this artillery is the study of the isospin degree of freedom and the effect of weak binding, which can be probed using recently available radioactive-ion beams and loosely bound stable projectiles.

Recent measurements with medium-heavy nuclei highlighted the change of slope of the fusion excitation function at deep sub-barrier energies compared to coupled channels calculations [2]. The energy where these deviations begin, referred to as the threshold energy for observing fusion hindrance, has been parametrized and its implications on the fusion with light nuclei of astrophysical relevance have been discussed [3]. Dasso and Pollarolo [4] pointed out that the cross-sections at deep sub-barrier energies could be used as an unique tool to obtain the value of the nuclear potential at small distances (see also [5]). More recently Ichikawa et al. showed that the potential energy at the touching point strongly correlates with this threshold energy [6]. Mišicu and Esbensen proposed a potential with a shallow pocket (as compared to that obtained from Woods-Saxon parametrization) based on a sudden approximation, where the reaction takes place so rapidly that the colliding nuclei overlap with each other without changing their density [7]. A repulsive core included to take into account the nuclear compressibility arising due to Pauli exclusion principle, modifies the depth and the shape of the minima of the inter-nuclear potential at small distances. They also showed that, depending on the choice of the couplings used in the calculations, there were surprising structures in the calculated average angular momentum at these low energies [7]. The nucleus-nucleus interaction potentials extracted from the microscopic time-dependent Hartree-Fock theory indicate that at low energies the frozen density approximation breaks down implying re-organization of the internal degrees of freedom [8]. Based on an adiabatic picture, a dynamical two-step model was proposed by Ichikawa et al. to explain the deep sub-barrier fusion data [9]. It should be noted that the above two approaches based on the sudden and adiabatic models predict different angular momentum distributions [10]. The measurement of the average angular momentum could also discriminate between the two approaches mentioned above [7, 9] that describe the fusion data equally well. In the sudden approach, using a shallow potential [7], the average angular momentum of the compound nucleus is always smaller than in the two-step adiabatic model [9] at low energies.

The fusion of weakly bound nuclei, which is a subject of current interest, has yet not been investigated at energies far below the barrier. For exotic weakly bound projectiles, a fully quantum mechanical time dependent wave-packet approach using a three body model also predicts a suppression of total fusion compared to corresponding stable nuclei over the entire energy range [11]. Experimental studies at deep sub-barrier energies have been restricted mainly to the measurement of fusion cross-sections of symmetric systems with the exception of 16O+204,208Pb systems, spanning a range of “stiffness”, reduced mass and Q-values [2, 3, 12–14]. Hence measurements of fusion cross-sections at low energies for a
were self-supporting rolled foils of $^{198}\text{Pt}$ (95.7% enriched, $\sim 1.3 \text{ mg/cm}^2$ thick) followed by an Al catcher foil of thickness $\sim 1 \text{ mg/cm}^2$. Fresh targets were used and back-ground data was collected before each irradiation. Two efficiency calibrated HPGe detectors were placed face to face for performing KX-γ-ray coincidence of the decay radiations from the irradiated sample. The sample was positioned symmetrically between the two detectors in a close geometry (1.5 mm from the face of each detector). The measurements were performed in a low background setup with a graded shielding. The reaction products were uniquely identified by means of their characteristic γ-ray energies and half-lives which in the case of fusion lead to $^{199-202}\text{Tl}$ residues. The γ-ray yields of the daughter nuclei were extracted by gating on their KX-ray transitions [15]. The resulting cross-sections of the residues are plotted in Fig. 1(a). Due to the increased sensitivity of the KX-γ-coincidence method, cross-sections down to a few nano-barns could be measured. The estimation of errors for low counting rates was made assuming Poisson statistics and using the method of maximum likelihood [17]. Statistical model calculations for the com-

FIG. 1: (color online) Excitation functions for $^6\text{Li}+^{198}\text{Pt}$ system (a) Evaporation residues from compound nuclear fusion. Dashed curves are results of statistical model calculations (see text). (b) Cross-sections of residues arising from d-capture ($^{198-200}\text{Au}$), 1n-pickup ($^{197}\text{Pt}$) and 1n-stripping ($^{199}\text{Pt}$) reactions. The dashed lines are to guide the eye.

FIG. 2: (color online) Fusion excitation function and derived observables for $^6\text{Li}+^{198}\text{Pt}$ system (a) Cross-sections for compound-nucleus formation and direct processes obtained from a sum of the partial cross-sections shown in Fig. 1 (a) and (b) respectively. The arrow indicates value of the Coulomb barrier ($V_B$). (b) Average angular momentum and (c) Logarithmic derivative of the fusion excitation function. The results of the coupled-channels calculations using the WS potential (solid line) along with single channel calculations using the WS potential (dashed line) and the M3Y potential with a repulsive core (dot-dashed lines) are shown in panels (a)-(c) (see text).
The fusion cross-sections, obtained from the sum of the measured evaporation residue cross-sections, are plotted in Fig. 2(a) for $^6\text{Li}+^{198}\text{Pt}$. Corrections for $^{196}\text{Pt}$ impurity in the target ($\sim 1\%$ even at highest energy). The cross-sections for the sum of deuteron-capture and neutron-transfer (plotted as open squares) are larger than those for fusion by orders of magnitude at deep sub-barrier energies. In the present work the average angular momenta ($\langle l \rangle$) have been derived from the fusion excitation function as suggested in Refs. [19, 20] and are plotted in Fig. 2(b).

Calculations using the coupled-channels (CC) code CCFULL [21] were performed with the ingoing-wave-boundary condition. Two sets of calculations, one using a standard Woods-Saxon potential (WS) ($V_0=110$ MeV, $r_0=1.1$ fm and $a=0.63$ fm) and the other based on the M3Y folded potential are presented. The potentials are plotted in Fig. 3. The calculations using the WS potential included the quadrupole excitation in $^{198}\text{Pt}$, considering coupling in the vibrational model. For $^6\text{Li}$ the $1^+$ (ground state) and the unbound $3^+$ states were assumed to be from a $\text{K}^* = 1^+$ rotational band. The results of the calculation with and without the inclusion of the couplings are shown in Fig. 2(a). At energies above the barrier the calculations overestimate the data, as expected from earlier studies involving weakly bound nuclei [22]. As can be seen in the figure, the CC calculations reproduce the data for energies around and well below the barrier. Plotted in Fig. 2(c) is the logarithmic derivative of the fusion cross-section ($L(E)=d[\text{ln}(\sigma E)]/dE$) obtained using three point numerical derivative. This representation provides an alternate way to illustrate any deviations in the slope of the fusion excitation function independent of the weight of the lowest barrier. The CC calculations reproduce well both the experimental slope $L(E)$ and the $\langle l \rangle$ values (Fig. 2(b)) over the entire range of energy. Thus for $^6\text{Li} + ^{198}\text{Pt}$, the CC calculations successfully explain the fusion excitation function along with the average angular momentum consistently, implying absence of the fusion hindrance at deep sub-barrier energies.

The lack of the fusion hindrance observed in the present case from the above calculations is also possible if the threshold value for the onset of fusion hindrance was not reached. This does not appear to be the case, as shown below. The threshold energy was computed following two independent approaches. The M3Y potential with repulsive core [7] was calculated taking the density distributions of $^6\text{Li}$ and $^{198}\text{Pt}$ from Ref. [23] and for the repulsive core, $V_{\text{rep}}=570$ MeV and $a_{\text{rep}}=0.35$ fm (yielding a value of $K=234$ MeV) as a representative choice for the parameters. The resulting potential (Fig. 3) has a minimum at $21.3$ MeV and as discussed in Refs. [4, 7] the threshold energy is larger than this value. Adopting a smaller value of $a_{\text{rep}} (=0.3$ fm) lowers the potential minimum ($=15.6$ MeV), but such a small value of $a_{\text{rep}}$ is inconsistent with that for other systems [7]. Alternatively following the two-step adiabatic model of Ichikawa with Krapple-Nix-Sierk potential [9], the energy at the touching configuration, related to the threshold energy is calculated to be $22.3$ MeV (Fig. 3). The present measurements extend down to $E_{\text{cm}}=19.8$ MeV, which is well below the threshold energy computed from both the approaches, although there may be some ambiguity for the definition of the touching point for a weakly bound nucleus.

Single-channel calculations using the above M3Y potential with a repulsive core were also performed as suggested in Ref. [7] and the results are shown in Fig. 2. The calculated fusion cross-sections, for energies lower than $22$ MeV, fall off steeply and are orders of magnitude lower than the corresponding single channel calculations using the WS potential (Fig. 2(a)). The effect of coupling on the calculated fusion cross-sections are found to be small from the CC calculations as seen in the same figure. A similar behavior was observed in Ref. [22]. Hence at these energies, even after including the effect of coupling the calculated fusion cross-sections using the M3Y+repulsive core potential will be much lower than the measured fusion cross-sections. The calculated $L(E)$ values also do
A shallow potential obtained using the M3Y interaction with a repulsive core successfully describes the fusion cross-sections at deep sub-barrier energies for symmetric, asymmetric and positive reaction Q-valued systems [7, 13, 16]. But for the present system with a weakly bound projectile this potential does not reproduce the trend of the fusion excitation function, \( L(E) \) and \( \langle l \rangle \). The present results suggest that the inner part of the interaction potential becomes deeper, going from a symmetric to a weakly bound asymmetric system, implying reduced contribution of the repulsive core. A plausible reason for this could be as follows. As the nuclei start overlapping, due to the weak binding of one partner, the Fermi energies of the two interacting nuclei are very different and will tend to equilibrate rather fast. Thus Pauli blocking is expected to be less effective for asymmetric systems involving weakly bound nuclei as compared to symmetric systems [24]. The actual form of the repulsive core is expected to depend also on the extent of the adiabatic nature of the collision [7]. At energies well below the barrier the adiabatic approximation is expected to be more appropriate where nuclear reactions take place following the minimum energy path allowing for the readjustment of the densities as a function of collective variables [8]. The predictions based on the adiabatic model of Ichikawa et al. [10] already appear to give the correct behavior for the average angular momentum in the medium-mass symmetric systems though currently such calculations are not possible for asymmetric systems.

In summary, we have presented the fusion excitation function for a very asymmetric system involving a weakly bound projectile at energies well below the barrier. This study shows the absence of fusion hindrance, pointing to the limitation of the sudden approximation for modeling reactions in such systems. It would be of interest to see whether this arises solely from the effect of weakly bound cluster structure [25] or also due to difference in transition from the sudden to the adiabatic potential. In order to address this question, both the sudden and adiabatic approaches would require extensions by taking into account the weakly bound nature of the projectile nucleus. An independent way to probe this conjuncture would be to analyze data of alpha-induced fusion at deep sub-barrier energies, for which measurements presently do not exist. Such data due to both experimental and theoretical simplicity could also provide an ideal testing ground for studying the effects of irreversible environmental couplings on the collision of nuclei [26].

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