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Resolution Enhancement in Multidimensional Solid-State NMR of $^{13}$C-Labeled Proteins Using Spin-State Selection

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Much progress has recently been made in the field of solid-state NMR of isotope ($^{13}$C, $^{15}$N) enriched bio-molecules, leading to the first protein structure solved by solid-state NMR. One of the principal factors limiting the study of larger systems remains spectral resolution. In uniformly $^{13}$C-labeled compounds such as proteins, the $^{13}$C-$^{13}$C J-couplings constitute a significant contribution to the linewidth in magic-angle-spinning (MAS) spectra. J-decoupling techniques for solid-state NMR using semi-selective pulses were first presented in 1996 by Straus et al.$^1$ to enhance the spectral resolution in indirectly detected spectral dimensions. Here, we present the application of spin-state selection and transition-selective polarization transfer to multi-dimensional solid-state NMR correlation experiments of $^{13}$C-labeled proteins. We show that single-transition selection removes the line broadening due to the $J_{CO}$ spin coupling in both direct and indirect dimensions of a two-dimensional CO-CO correlation experiment. The experiment is demonstrated on a sample of microcrystalline Crh, a 85-residue protein involved in carbon catabolite repression in Bacillus subtilis and leads to nearly a factor two improvement in spectral resolution.

Spin-state-selective NMR techniques have been developed in liquid-state NMR spectroscopy for the measurement of small spin-spin coupling constants$^2$ and for transverse-relaxation-optimized spectroscopy (TROSY)$^3$. Recently, we have demonstrated the feasibility of homonuclear spin-state selection in the solid state,$^4$ even for J-couplings that are not resolved, using an IPAP-type selection filter.$^4$ In order to perform a multi-dimensional experiment, the selected transition of the first spin, $I$, evolving during $t_1$ should be transferred to a single-transition of a second spin, $S$, by means of an appropriate mixing sequence. Generally, spin-state-selective coherence transfer is obtained through zero-quantum (ZQ) or double-quantum (DQ) rotations.$^2$ A ZQ rotation “conserves” the spin state ($I, S^x \rightarrow I^x S^x$) and a DQ rotation “reverses” the spin state ($I, S^x \rightarrow I^x S^x$). In liquid-state NMR, spin-state-selective coherence transfer is obtained using planar mixing or S'C'T building blocks.$^5$ Similar $J$-based spin-state-selective transfer techniques could be envisaged in the solid state. However, schemes using the dipolar coupling may be more appropriate for solid-state applications since most of the (many) existing solid-state dipolar coherence techniques under MAS have a ZQ or DQ average Hamiltonian, and may therefore be used directly for spin-state-selective polarization or coherence transfer. Common examples of such sequences are proton-driven spin diffusion (PDS)$^6$ and RFDR$^7$ with a ZQ-average Hamiltonian, and C7, POST-C7, and SPC5 with a DQ-average Hamiltonian.$^8$

The pulse sequence shown in Figure 1 yields a spin-state-selective CO-CO correlation experiment using PDS. After cross-polarization, both $I$ (CO) transitions are separated into different sub-spectra with respect to the $S$ (CO) spin by an IPAP sequence. The $t_1$ period is built into the IPAP sequence providing a constant time (CT) experiment.$^1$ PDS then yields spin-state-selective CO->C0 polarization transfer. During the PDS mixing period, the selected transition of the first spin, $I$, is transferred to $C^0$ polarization ($S^0$) by the dipolar coupling, but the two-spin order $2JS$, present before $t_1$ evolution, is not affected by the dipolar coupling. Since the build-up of $S$ and the decay of $2JS$, are dependent on the mixing time $t_{mix}$ and in order to obtain proper spin-state selection independent of $t_{mix}$, the two polarization-transfer pathways are separated into different sub-spectra by an appropriate 2-step phase cycle, as detailed in the caption of Figure 1. By separating the in-phase and anti-phase components of the S-spin coherence during detection ($t_2$), their relative amplitude can be adjusted by an appropriate scaling factor $k$, to yield proper spin-state selection.

Figure 1. Pulse sequence suitable for an I-S (e.g. CO-CO) correlation experiment. 90° and 180° rf pulses are represented by filled and open bars, respectively. The semi-selective CO pulses are applied with the shape of the center lobe of a sinc2 function, whereas the C0 pulses were rectangular with $\omega_0 = \Delta/\sqrt{15}$ (90°) and $\Delta/\sqrt{3}$ (180°). All pulses are applied along the $x$ axis unless indicated. An 8-step phase cycle was applied with $\phi_1 = \pi - \phi$; $\phi_2 = x y - x y$; $\phi_3 = x = x - x - x - x$ (180°). The constant time delay $2T_2$ is adjusted to $(2\tau_{CO})^3$. Four data sets (A1), (A2), (B1) and (B2) are recorded. A- and B-type experiments are used to separate the I-spin transitions using IPAP with the following settings: (A) $\alpha = T, \phi_1 = -\pi$ and (B) $\alpha = 0, \phi_2 = y$. Each of the experiments A and B are recorded twice by setting alternatively the phase $\phi_3$ to $x$ and $-x$. Addition of the two data sets yields in-phase selection (A1, B1) of the S-spin coherence, subtraction yields anti-phase selection (A2, B2). Linear combination (A1 + $\delta \phi B1) \pm k (A2 + \delta \phi B2)$, yields the four different single-transition correlation spectra with $\phi = +1$ or $-1$ and $k$ a scaling factor taking into account the different PDS transfer functions described in the text. The pulse sequence and phase cycle used here are available from our website, or upon request to the authors.

Figure 2 shows the (alpha)-sub-spectrum recorded for a microcrystalline sample of uniformly $^{13}$C, $^{15}$N labeled Crh using the sequence of Figure 1, compared to a standard PDS correlation spectrum. Spin-state selection provides a remarkable increase in resolution in both dimensions. Specifically, the resolution of the G49 cross peak is enhanced by 44% and 17% for C0 and CO dimensions, respectively. The resolution enhancement is more pronounced in the C0 dimension because of the longer acquisition time used for direct detection. We have limited here...
the spectral resolution in the CO dimension to the maximal CT of 9.2 ms. It can be further increased by lengthening $t_1$ in an additional non-constant acquisition time or by processing using stronger apodization functions and linear prediction (LP). Despite the CT frequency encoding, advanced prediction algorithms like mirror-image LP can not be used here. Indeed, in solid-state NMR, the $^1$C linewidth is dominated by refocusing interactions and the signal therefore still decays during the CT evolution period.

Figure 2. CO-C$^\alpha$ region of the standard (a) and the (αα)-spin-state selective (b) PDSD spectrum of microcrystalline Crh. The second spectrum was obtained from the linear combination $A1 + B1 + k(xA2 + B2)$, with $k = 0.7$. Both experiments were performed on a Bruker Avance spectrometer operating at a $^1$H frequency of 500.13 MHz with a 4 mm double-resonance CP/MAS probe. The temperature was set to 269 K, and the MAS frequency was 11 kHz. SPINAL-$^1$H decoupling was used with $\omega_0 = 78$ kHz. $\tau_{\text{mix}}$ was 30 ms, and $\tau_2^\text{CO}$ was 25 ms. For the standard and the spin-state-selective spectra, a total of 64 and 4x24 transients respectively were accumulated for each of 240 complex points in $t_1$. Cosine apodization was applied in both dimensions prior to Fourier transformation. For both spectra, the first contour level was set to 15% of the intensity of the G49 resonance, with a factor of 1.4 between levels.

The efficiency of the spin-state-selective experiment compared to the standard PDSD experiment was estimated from measured cross peak intensities to be about 35%. This loss is mainly due to the $2\tau \approx 9$ ms IPAP filter. As we have shown recently, this signal loss is greatly reduced by using high-power $^1$H decoupling during the filter delay and faster spinning of the sample. This is possible for protein samples if a better cooling system than we have is available. In addition, the four sub-spectra can be added after appropriate shifting of the spectrum by $\pm \Delta_d/2 = 27$ Hz along the CO and/or C$^\alpha$ dimensions, resulting in an additional gain of a factor of two in signal to noise. Therefore we expect that for an optimized experimental setup the spin-state-selective experiment not only increases spectral resolution but also sensitivity.

Figure 3 illustrates how the recorded 2D data sets are combined. The linear combinations $(A1 + B1) + k(xA2 + B2)$ yield all four single-transition-to-single-transition correlation spectra (αα), (αβ), (βα), and (ββ). For $\tau_{\text{mix}} = 30$ ms, the scaling factor was found to be $k = 0.7$. For spectra with $\tau_{\text{mix}} = 15$ ms, a scaling factor $k = 0.5$ was found.

At longer mixing times some polarization is transferred by PDSD from the CO to other side chain carbons (C$^\beta$, C$^\gamma$, C$^\delta$). Since there is no direct scalar coupling between the CO and the side chain carbons, there is no frequency shift along the detection dimension ($\omega_0$) between the (αα) and (αβ), or (βα) and (ββ) sub-spectra. Thus the new experiment provides a simple way of distinguishing C$^\alpha$ from other side chain carbons (e.g. C$^\beta$ of Thr and Ser residues) by comparison of sub-spectra.

In conclusion, we have introduced a new experiment providing significant resolution enhancement in multi-dimensional solid-state NMR correlation experiments. Resolution enhancement is achieved by using transition-selective excitation and transfer techniques. Spin-state-selective polarization transfer is obtained using standard ZQ solid-state NMR mixing sequences. Similar results are expected for transfer sequences based on DQ rotations. The new experiment can be easily extended to higher-dimensional experiments. In addition, spin-state-selective correlation experiments allow the distinction of ‘direct’ transfer peaks, involving covalently-bound nuclei, and ‘relayed’ transfer peaks.

The new experiment is expected to be very useful for the assignment of solid-state NMR spectra of proteins.

(7) Meissner, A.; Duus, J. O.; Sørensen, O. W. J. Biomol. NMR 1997, 10, 89-94.
(8) Sørensen, M. D.; Meissner, A.; Sørensen, O. W. J. Biomol. NMR 1997, 10, 181-186.
(9) Bloembergen, N. Physica 1949, 15, 386-426.
(13) http://www.ens-lyon.fr/STIM/NMR