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Improved Lattice Spectroscopy of Minimal Walking Technicolor

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We present an improved study of spectroscopic observables in the $SU(2)$ Yang-Mills theory with two adjoint fermions. We make an improvement on the precision of previous results which clarify the scale of finite volume effects present. This analysis adds to the evidence for near-conformal dynamics of this theory, while indicating a preference for a low anomalous mass dimension of the massless theory.

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1. Minimal Walking Technicolor

The mechanism of dynamical electro-weak symmetry breaking (often referred to as technicolor) [1] remains a possible explanation for the breaking of the electro-weak symmetry observed in nature. Gauge theories which possess an approximate infra-red fixed point have been proposed as preferred candidates for the technicolor sector in models of extended technicolor [2] as it has been argued that they would allow suppression of flavour changing neutral currents while permitting generation of large fermion masses. This is achieved through the non-trivial renormalisation dynamics between the technicolor and extended technicolor scales. This property of the theory has been coined “walking” in reference to the slow running of the coupling between the two scales ¹.

We can hope to generate an infra-red fixed point in a gauge theory, while minimising contributions to electroweak precision constraints, by adding a low number of fermion flavours in higher gauge representations to a gauge theory with a low number of colours. For these reasons the gauge theory theory with symmetry group $SU(2)$ and two flavours of Dirac fermion in the adjoint representation has been termed minimal walking technicolor (MWT) [4].

Conformal or near-conformal dynamics have not been conclusively identified in any theory as yet, although there are encouraging hints from a number of sources. Due to the non-perturbative nature of the problem, much of the investigation of this problem has arisen from lattice simulations, and this activity is growing. Some recent lattice studies of MWT [5, 6, 7] have attempted to identify a near-conformal behaviour directly from the behaviour of the coupling and anomalous dimensions of the theory under renormalisation flow. Others [8, 9, 10, 11, 12, 13], including this work, perform measurements of physical observables in the theory and from their behaviour attempt to identify signals of near-conformal dynamics.

This study builds on previous work in [9, 13]. In particular we seek to support these results and establish their reliability by investigating the effect of performing measurements and analysis using alternative methods. In addition, through this we expect to ascertain the scale of systematic uncertainties present, which have as yet been largely unexplored.

2. Signals of Conformality

The question of whether MWT in the chiral limit possesses an actual infra-red fixed point, i.e. it lies within the conformal window, or an approximate fixed point, has not been conclusively answered. Nevertheless, it is clear that MWT with a non-zero fermion mass and defined in a finite volume, as simulated on the lattice, cannot be conformal. If the chiral continuum theory possesses an infra-red fixed point, the lattice results will be described by a mass-deformed conformal gauge theory. Recent discussions of scaling laws in such theories [13, 14, 15] have derived scaling relations for hadronic masses and amplitudes in terms of the mass deformation.

In a theory with strong chiral symmetry breaking (like QCD), the mesonic spectrum departs significantly from degeneracy. The pseudoscalar mesons become massless in the chiral limit, while the rest of the spectrum is expected to retain a finite mass. In a mass deformed conformal theory, the outlook is different. In approaching a conformal limit, the theory respects the hyperscaling property, whereby all masses M in the theory scale identically. They must vanish in the limit of

¹For a recent technicolor review see [3].

vanishing quark mass m . From an analysis of the renormalisation of mesonic two-point functions we can deduce that all meson masses M in the theory vanish as $M \sim m^\rho$ where the critical exponent ρ is given by $\rho = \frac{1}{1+\gamma_*}$ where γ_* denotes the anomalous dimension of the chiral condensate in the conformal theory at the infra-red fixed point. This quantity is phenomenologically very interesting as it determines the extent to which the chiral condensate is enhanced at the extended technicolor scale with respect to the technicolor scale when a near-conformal gauge theory is included as the new strongly coupled gauge group in a model of extended technicolor. We will attempt to fit our lattice data with this hyperscaling relation in order to determine whether the massless limit of MWT is indeed a conformal theory, and also to deduce a range of γ_* preferred by the data.

In [13] it is also shown that in a theory confined to a finite box of size L , analysing the renormalisation of the free energy density and treating the box size L as a relevant parameter, there exist universal finite-size scaling laws for physical observables. These can be summarised as $LM \sim \Upsilon(Lm^\rho)$ for any observable M of mass dimension one. The function Υ will differ for each M , but this relation allows us to compare data across different lattice volumes in order to draw conclusions on the range of γ_* which are preferred by the data.

3. Systematic Spectroscopy

This study builds on the work described in [9, 13] where spectroscopic observables of MWT were measured through lattice simulations. The computation was performed using the *HiRep* code, designed to simulate theories of general number of colours N_c and number of flavours N_f of fermions in a generic representation R of the gauge group. The simulations used the Wilson gauge action, and the Wilson fermion formulation along with the RHMC algorithm. A number of lattice volumes have been analysed, from 16×8^3 to 64×24^3 with a range of bare quark masses. We describe analysis of ensembles at a single lattice spacing, with $\beta = 2.25$.

For this study we have performed some alternative analyses to those in [13]. We have modified the *Chroma* suite of lattice software [16] to operate with $N_c \neq 3$ and a number of fermion representations R other than the fundamental. Using the resulting code we have utilised the in-built smearing routines found in *Chroma* to perform measurements on the gauge configurations generated with *HiRep* with a number of different quark smearings.

The signals for the observables obtained from correlators using a wall-smear source operator are the cleanest available. This could be expected due to the enhancement of the overlap of the wall-smear quark bilinear with the mesonic ground states due to its projection onto zero momentum. As a result it was decided to produce a complete set of wall smeared correlators on our data, complementing the local correlators already analysed [13]. We present full results for all observables of interest in Sec. 4.

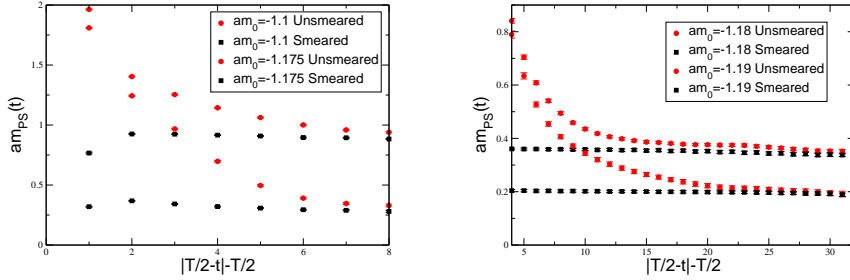
In addition to this investigation of the effect of smearing on the observables, we have also investigated possible systematic errors arising from our analysis methods. Firstly, we considered the algorithm for extracting meson masses from the correlators. The Prony method [13] using only the ground state mass was preferred to the effective mass method [17], and so the Prony method was implemented in all analyses. We then examined whether the effective observable for the PCAC quark mass and decay amplitudes was affected by the use of the effective meson mass or its fitted value. Finding little variation, it was decided to use the effective mass in all definitions.

4. Results

In the following we present results from the wall-smearred inversions on all available ensembles generated with HiRep. The analysis used every tenth configuration in the Monte Carlo chain. Meson masses were extracted from the correlators using the Prony method using only the ground state mass. The definition of effective observables for the PCAC mass and decay constants are those used in [13] and can be found explicitly in [17].

4.1 Systematics

As stated in Sec. 3 the wall-smearred correlators result in significantly improved signals for our mesonic observables. This is illustrated in Fig. 1 for a choice of bare masses, on our smallest and largest lattices. It can be seen that on the smallest lattice, the result from unsmeared correlators is contaminated with excited states and that their effects are not diluted even at the centre of the lattice. This will lead to a discrepancy between the results from the smeared correlators and those from the unsmeared. In contrast however, we see that on the largest lattice, the temporal extent is sufficient to suppress the effect of excited states in the local correlators towards the centre of the lattice, and the two results are in good agreement. As such we expect the smeared results to systematically differ from the local results on smaller lattices while we expect agreement on the larger lattices. Below we present the smeared results, for the full local results see [13].



(a) Effective pseudoscalar masses on the 16×8^3 lattice. (b) Effective pseudoscalar masses on the 64×24^3 lattice.

Figure 1: Effective pseudoscalar masses from both smeared and unsmeared correlators.

4.2 PCAC mass

In Fig. 2a results for the PCAC mass on all ensembles are presented. The inset illustrates a close up of the approach to the chiral limit, with a linear extrapolation to zero quark mass. Using this we find the critical bare quark mass to be $am_c = -1.2025(5)$, which compares very well to the result obtained from the local data [13].

4.3 Meson masses

Fig. 2b shows the results obtained for the pseudoscalar mass m_{PS} as a function of the PCAC quark mass m . In Figs. 2c, 2d we can see the ratio of m_{PS} and m_{PS}^2 , respectively, to m . We notice

that $\frac{m_{PS}}{m}$ appears to increase in the chiral limit, while $\frac{m_{PS}^2}{m}$ appears to vanish at zero quark mass. This is at odds with the behaviour expected in a theory described by a chiral effective field theory (like QCD) in which the pseudoscalar mass is expected to scale as $m_{PS} \sim \sqrt{m}$ in the chiral limit.

In Fig. 2e we see the mass of the vector meson m_V as a function of the quark mass m . We can see that it appears to vanish in the chiral limit, indicating that chiral symmetry is unbroken in this theory. Analysing the ratio of m_V to m_{PS} , Fig. 2f, we see that this does not diverge in the chiral limit, in fact it deviates little from unity, again in conflict with the expectations from a chirally broken theory where the pseudoscalar is expected to become massless in the chiral limit, with the rest of the mesonic spectrum retaining a finite mass. The behaviour we observe is more akin to that expected in a mass deformed conformal gauge theory where the property of hyperscaling insists that all masses must scale to zero with the same critical exponent in the chiral limit.

4.4 Decay constants

In Fig. 2g we show the measured pseudoscalar decay constant as a function of quark mass. Clearly, large finite-volume effects are present, indicated by the large discrepancies between results on different volumes. Using the Gell-Mann–Oakes–Renner relation we can analyse the scaling of the chiral condensate with the quark mass. Although again our data, shown in Fig. 2h, suffer from significant finite-size effects, there is no indication of a finite chiral condensate in the chiral limit.

4.5 Scaling and anomalous dimension

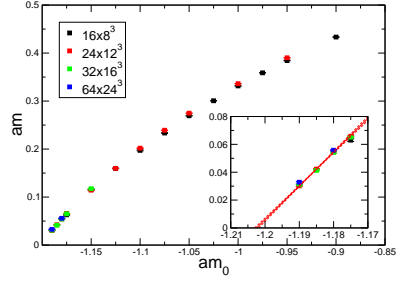
In Sec. 2 we described how the meson masses are expected to scale to zero with the same critical exponent $\rho = \frac{1}{1+\gamma_*}$ in the chiral limit. We have attempted to fit the behaviour of m_{PS} and m_V to such a power law and extract from these fits a value of γ_* . The χ -squared value of these fits is generally rather poor, except in a few cases. However the results indicate that no particular non-zero value of γ_* is preferred by our data, and we see this as consistent with a low value $\gamma_* \ll 1$. We have also attempted to determine γ_* from the finite-size scaling relations described in Sec. 2. This analysis also indicates a preference for a low γ_* from the data. For details, see [18].

5. Conclusions

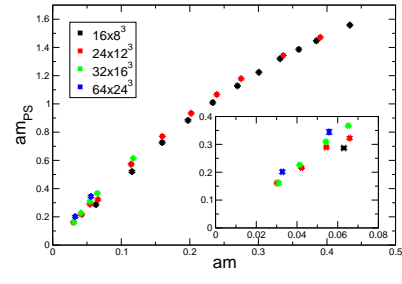
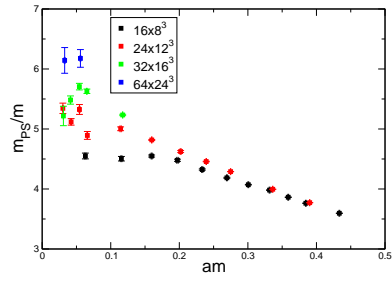
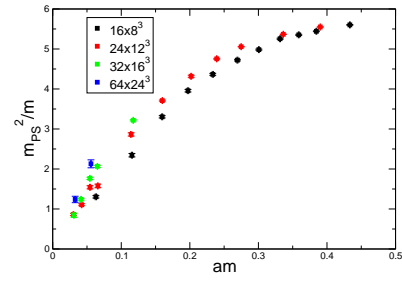
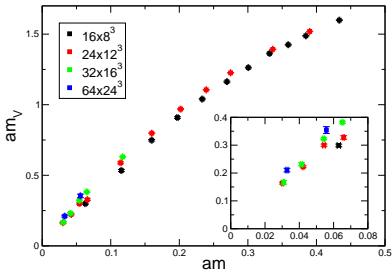
Our investigation of the systematic effects present in our analysis of mesonic observables in MWT indicates that they are largely under control, and finite volume effects are small on our larger lattices. We find further supporting evidence for the case that this theory is near-conformal, with indications that the anomalous dimension at the fixed point in the massless theory is small. This could be a problem for certain models which depend on a large γ_* to satisfy current experimental data, however it should be remembered that MWT must be embedded in a larger theory of extended technicolor in order to construct a model of electro-weak symmetry breaking and as such γ_* could be affected by other sectors of the theory, and so our result is nothing like the final answer.

Acknowledgements

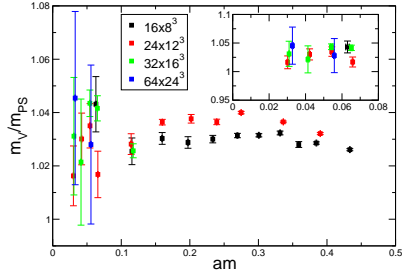
The numerical calculations presented in this work have been performed on the Horseshoe6 cluster at the University of Southern Denmark (SDU) funded by the Danish Centre for Scientific



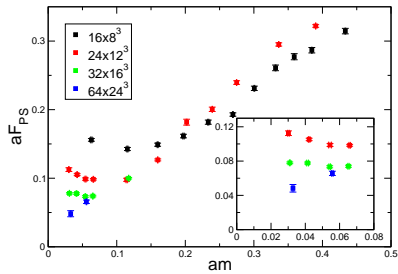
(a) PCAC quark mass.

(b) Pseudoscalar mass m_{PS} .(c) Ratio of m_{PS} to m .(d) Ratio of m_{PS}^2 to m .

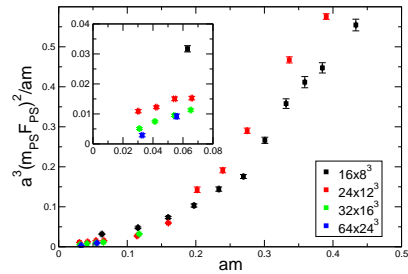
(e) Vector mass



(f) Ratio of vector mass to pseudoscalar mass.



(g) Pseudoscalar decay constant.



(h) Chiral condensate as given by the GMOR relation.

Figure 2: Fermionic observables in the chiral region.

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