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## Evidence from scanning tunneling microscopy in support of a structural model for the InSb(001)- $c(8 \times 2)$ surface

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In this letter we present evidence from scanning tunneling microscopy studies in support of a recently proposed structural model for the indium-terminated  $c(8 \times 2)$  surface of InSb(001). This structural model, by Norris and co-workers, is based on a surface x-ray diffraction study [Surf. Sci. **409**, 27 (1998)], and represents a significant departure from previously suggested models for the  $c(8 \times 2)$  reconstruction on any (001) surface of the common III–V semiconductor materials. Although filled state images of the InSb(001)- $c(8 \times 2)$  surface have previously been published, empty states image of sufficient quality to extract any meaningful information have not previously been reported. The observations are in excellent agreement with the recently proposed model for this surface reconstruction. © 1999 American Institute of Physics. [S0003-6951(99)03138-1]

In contrast to most III–V semiconductor surface reconstructions that are group V terminated, and whose structures have been determined, there has been considerable debate over the exact composition and atomic structure of the outermost layers of the group III terminated reconstructions.<sup>1,2</sup> This has clearly been demonstrated by the controversy surrounding the InSb(001)- $c(8 \times 2)$  surface reconstruction.<sup>3–5</sup>

The initial model for the InSb(001)- $c(8 \times 2)$  surface was proposed by John *et al.*,<sup>6</sup> based on high energy electron diffraction, soft x-ray photoemission data and comparison with the same surface reconstruction previously observed by Biegleson *et al.* for GaAs(001) using scanning tunneling microscopy (STM).<sup>7</sup> This model had the surface terminated with 0.75 monolayer (ML) of indium [where 1 ML is the In concentration in a  $(1 \times 1)$  unreconstructed In-terminated surface], made up of blocks of three In dimers and a missing dimer forming a  $(4 \times 2)$  block with the dimer bonds parallel to the  $[110]$  direction, on top of a complete Sb layer. The  $c(8 \times 2)$  unit mesh was then formed by staggering the rows of  $(4 \times 2)$  blocks in the  $[110]$  direction by  $\times 1$ .

A previous STM study of the  $c(8 \times 2)$  surface was performed by Schweitzer *et al.*<sup>3</sup> on an InSb(001) sample prepared *in situ* by cycles of low energy ion bombardment and annealing. Filled state images indicated double rows of bright dots aligned along the  $[110]$  direction, separated by darker regions in which no structure could be resolved. No empty state images were recorded. These authors could only identify a  $(4 \times 1)$  unit mesh in their STM images, which conflicted with the  $c(8 \times 2)$  translational symmetry observed in the low energy electron diffraction (LEED) pattern. This apparent conflict was resolved by attributing the pairs of bright rows in the STM image to tunneling from filled lone-

pair orbitals of the Sb atoms in the second layer at the location of the missing In dimers in the top layer of the structure. Although agreeing with the available data at the time, this assignment was later shown to be incorrect.

Varekamp *et al.*<sup>5</sup> also studied the InSb(001)- $c(8 \times 2)$  surface by STM, again prepared by low energy ion bombardment (500 eV) followed by annealing to 350–400 °C. These authors were able to determine a true  $c(8 \times 2)$  symmetry in their filled states images. They also observed that the spot separation in the double rows was not exactly one surface lattice constant in the  $[110]$  direction, but oscillated about this value. This alternating compression and expansion resulted in an observed two-fold periodicity, creating a  $(4 \times 2)$  unit mesh instead of the previously observed  $(4 \times 1)$  unit mesh.<sup>3</sup> They also observed a  $(4 \times 2)$  unit mesh in a relatively low quality empty states image, but there was insufficient data from which to derive any detailed structural information.<sup>5</sup> The observation of the true  $c(8 \times 2)$  unit mesh does, however, indicate that the double rows of bright spots cannot be due to second layer Sb atoms. Although noting this fact, Varekamp *et al.* were unable to suggest an alternative structural model for the surface that agreed with their STM images.

The most recent structural model for this InSb(001)- $c(8 \times 2)$  surface, proposed by Norris and co-workers, is based on a surface x-ray diffraction study.<sup>8</sup> This model is characterized by chains of metallic bonded indium atoms extending along the  $[110]$  direction, separated by pairs of Sb dimers, with both the metallic indium and the Sb dimers residing on top of an Sb terminated bulk structure. In this letter we present evidence from high resolution filled and empty state STM images which strongly support this structural model for the InSb(001)- $c(8 \times 2)$  surface.

The experimental work was carried out at Hamburg University, Germany, using a large ultrahigh vacuum system

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equipped with comprehensive sample preparation and analysis facilities (including LEED). Images were recorded with a commercial STM I (Omicron, GmbH). The *n*-type doped ( $n \sim 2 \times 10^{16} \text{ cm}^{-3}$ ) InSb(001) samples (Wafer Technology, UK) were chemically etched prior to insertion in vacuum and cleaned *in situ* by cycles of low energy  $\text{Ar}^+$  ion bombardment (500 eV, with the ion beam at  $45^\circ$  to the surface normal) followed by annealing to  $350^\circ\text{C}$  for up to 30 min. The resulting surface, which displayed a clear  $c(8 \times 2)$  LEED pattern, was checked for lack of contaminants using Auger electron spectroscopy prior to transfer to the STM, where all the imaging occurred at room temperature. Typical imaging conditions for both filled and empty states used bias voltages between  $\pm 2$  V with tunneling currents of 2–3 nA.

The STM image shown in Fig. 1 is a filled states image ( $70 \text{ \AA} \times 70 \text{ \AA}$ ) of the InSb(001)- $c(8 \times 2)$  surface following three cycles of  $\text{Ar}^+$  bombardment and annealing. Rows of bright features (marked A) are observed along the  $[110]$  direction [with a separation of  $4.5 \pm 0.2 \text{ \AA}$  ( $\times 1$ ) between features and  $19.6 \pm 0.2 \text{ \AA}$  ( $\times 4$ ) between adjacent rows]. Between these bright rows are two more rows of features (marked B) that are somewhat less intense and evenly distributed along the  $[110]$  direction. This image, recorded with a bias voltage of  $-2.0$  V and a tunnel current of 1.67 nA, differs from the previously published work of Schweitzer *et al.*,<sup>3</sup> but is similar to the image produced by Varekamp *et al.*<sup>5</sup> Marked on Fig. 1 are a series of  $(4 \times 2)$  blocks which have been arranged to form the overall  $c(8 \times 2)$  unit mesh seen in LEED. Within each  $(4 \times 2)$  block pairs of B features have been joined with a short line.

Figure 2 shows an empty states image ( $90 \text{ \AA} \times 90 \text{ \AA}$ ) of the same InSb(001)- $c(8 \times 2)$  surface. The image was recorded with a sample bias of  $+1.1$  V and a tunnel current of 3.2 nA, although similar images were also recorded at slightly lower tunnel currents. As in the filled states image, a series of bright features are seen (again marked A) aligned along the  $[110]$  direction with the same  $4.5 \pm 0.2 \text{ \AA}$  ( $\times 1$ ) separation. However, the structure observed between adjacent rows of bright dots is now very different. Bright, elongated features, marked C and aligned parallel to the  $[\bar{1}\bar{1}0]$  direction are seen linking adjacent rows of A features (rather like rungs on a ladder), with a separation of  $9.0 \pm 0.2 \text{ \AA}$  ( $\times 2$ ), and with each set of rungs shifted by  $\times 1$  ( $4.5 \text{ \AA}$ ) with respect to the adjacent rows. In addition between each rung, less well resolved features appear (marked D) the overall appearance being of a faint row of diffuse features halfway between and halfway along each rung. This image clearly shows the full  $c(8 \times 2)$  unit mesh expected from the LEED pattern, but not observed directly in previous STM studies, as empty states image of sufficient quality to extract any meaningful information have not previously been reported for this surface reconstruction.

When both the filled and empty state images are considered in terms of the structural model proposed by Norris and co-workers<sup>8</sup> (shown in Fig. 3), it is possible to assign the observed features in both images. The dominant features in both the filled and empty state images are the rows of bright features marked A, aligned along the  $[110]$  direction. Because of their prominence at both positive and negative bias voltages, these features are attributed to states associated

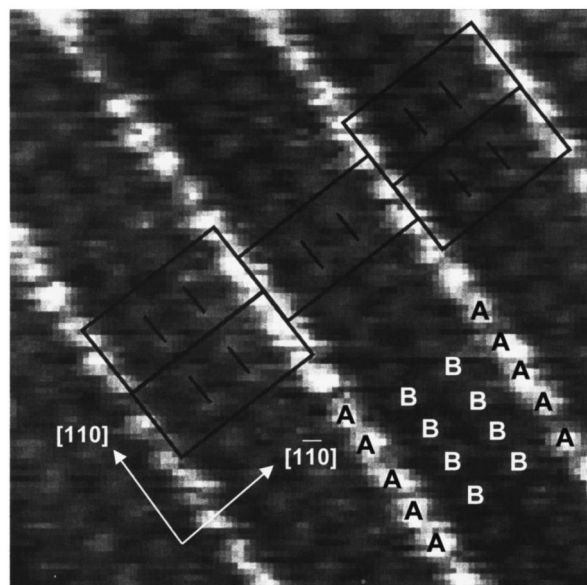


FIG. 1. Filled states STM image ( $70 \text{ \AA} \times 70 \text{ \AA}$ ) of the InSb(001)- $c(8 \times 2)$  surface showing the  $(4 \times 2)$  unit mesh arranged to form the overall  $c(8 \times 2)$  translational symmetry. Each unit mesh spans the adjacent rows of indiums (marked as A) aligned along the  $[110]$  direction and contains two Sb-Sb dimer pairs (marked by the dash, B-B) with the dimer direction also along the  $[110]$  direction. The image was recorded with a sample bias of  $-2.0$  V and a tunnel current of 1.67 nA.

with metallic bonding in the In rows.<sup>8</sup> Metallic bonding generally requires high coordination numbers and the arrangement of what is effectively a triple row of In atoms along the  $[110]$  direction, separated by the same interatomic distance observed in indium metal, satisfies this condition. We therefore identify the A features in Figs. 1 and 2 as due to the rows of indium atoms shown in Fig. 3.

If one now considers the filled states image in Fig. 1, the structure appears to be  $(4 \times 1)$ . However, closer scrutiny of

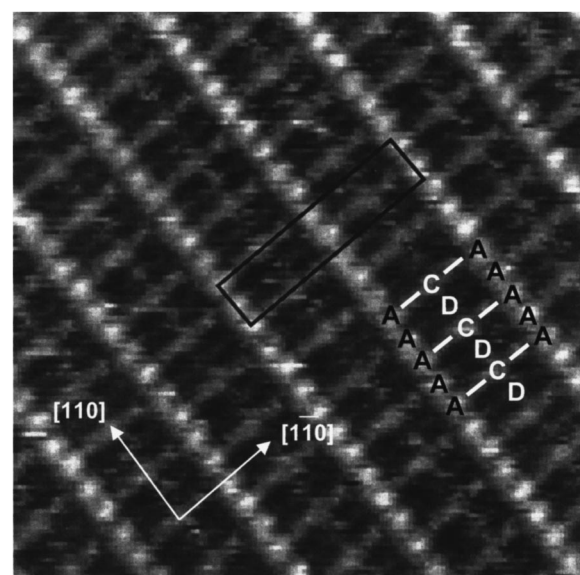


FIG. 2. Empty states STM image ( $90 \text{ \AA} \times 90 \text{ \AA}$ ) of the InSb(001)- $c(8 \times 2)$  surface. The rectangle shows the  $c(8 \times 2)$  unit mesh. Note the additional intensity on the right-hand side of the  $c(8 \times 2)$  unit mesh (marked as C), and the rows of bright features (marked as A) aligned along the  $[110]$  direction. Diffuse features (marked as D) are also seen between the C features also aligned along the  $[110]$  direction. The image was recorded with a sample bias of  $+1.1$  V and a tunnel current of 3.2 nA.

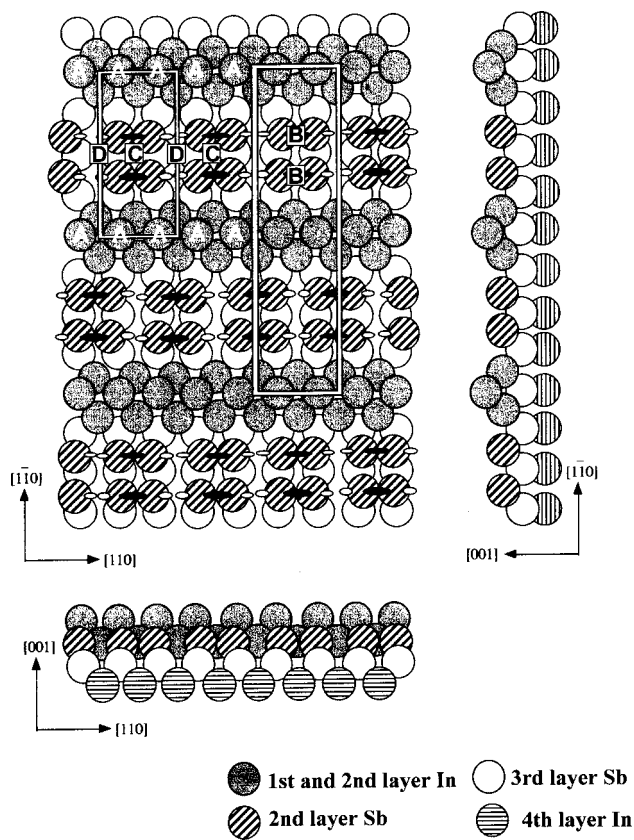


FIG. 3. Plan and side views of the structural model proposed for the In-terminated InSb(001)- $c(8 \times 2)$  surface based on surface x-ray diffraction measurements. Both the  $(4 \times 2)$  and  $c(8 \times 2)$  unit mesh are also shown and the model is marked with the A, B, C and D features that relate to the filled and empty states images.

the B features indicates that they tend to form pairs by moving slightly toward each other, allowing individual  $(4 \times 2)$  blocks to be identified. Within each  $(4 \times 2)$  block, four B features are seen, which we have divided into pairs marked by a short line. Schweitzer *et al.*<sup>3</sup> had originally attributed these features to electron lone pairs associated with second layer Sb atoms; however, it is now proposed that these features are due to lone pair orbitals associated with the Sb–Sb dimers in the second layer. The two Sb dimer pairs contained within each  $(4 \times 2)$  unit in the structural model (marked B) are shown in Fig. 3. The rather even spacing of the B features in Fig. 1 can be attributed to the STM probing the wave functions associated with the Sb dimers, rather than their nuclear positions.

With the exception of the bright rows of A features along the  $[110]$  direction in both STM images, the empty states image in Fig. 2 is markedly different from the filled states image in Fig. 1. The  $c(8 \times 2)$  unit mesh has already been identified, but it is somewhat more difficult to identify the individual features in terms of the structural model shown in Fig. 3. We know from the filled state image that there are two Sb dimer pairs present in each  $(4 \times 2)$  block. It is therefore proposed that the “rungs,” marked C in Fig. 2,

result from empty states associated with these two Sb dimer pairs. Comparison with empty states images from chemisorbed Sb–Sb dimers on the  $c(4 \times 4)$ -InSb(001) surface clearly indicate the presence of these states.<sup>4</sup> It should also be noted that the empty state image shows the A-C-A features aligned along the  $[1\bar{1}0]$  direction, whereas the structural model in Fig. 3 shows the C features lying halfway between two A features in a given row. We attribute this to the fact that the bright rows seen in Fig. 2 are in fact due to a combination of the states associated with In atoms in both the first and second layers. If one considers adjacent A's and the two second layer In atoms between them, and assigns the feature observed in the STM image to a combination of states from that cluster of four metallic bonded In atoms, then repeating this along the In rows indicates that the gaps between the bright A features seen in the image will lie on a line bisecting the Sb–Sb dimers and perpendicular to the dimer direction (as in the model). In addition, the formation of Sb–Sb dimers in the second layer (B features) exposes small sections of the third layer in this part of the unit cell (marked D) in both Figs. 2 and 3. This is relevant, since the structural model of Norris and co-workers<sup>8</sup> has the third layer as a complete Sb layer. These diffuse features are therefore assigned to empty states associated with the complete Sb third layer.

In summary, we have presented high resolution filled and empty states STM images of the In-terminated InSb(001)- $c(8 \times 2)$  surface. These STM images are in excellent agreement with a recently proposed structural model for this surface that represents a significant departure from previously suggested models, e.g., in comparison with the  $c(8 \times 2)$ -GaAs(001) surface reconstruction. Why the  $c(8 \times 2)$ -InSb(001) surface should be so different is not known; however, one possibility could be the increased degree of metallicity (i.e., the propensity to form metal–metal bonds) down a given group in the periodic table (hence  $\text{In} > \text{Ga}$ ). This could explain why the group III terminated reconstructions for these two surfaces with the same periodicity are so different.

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